

# Supernova Neutrinos

John Beacom

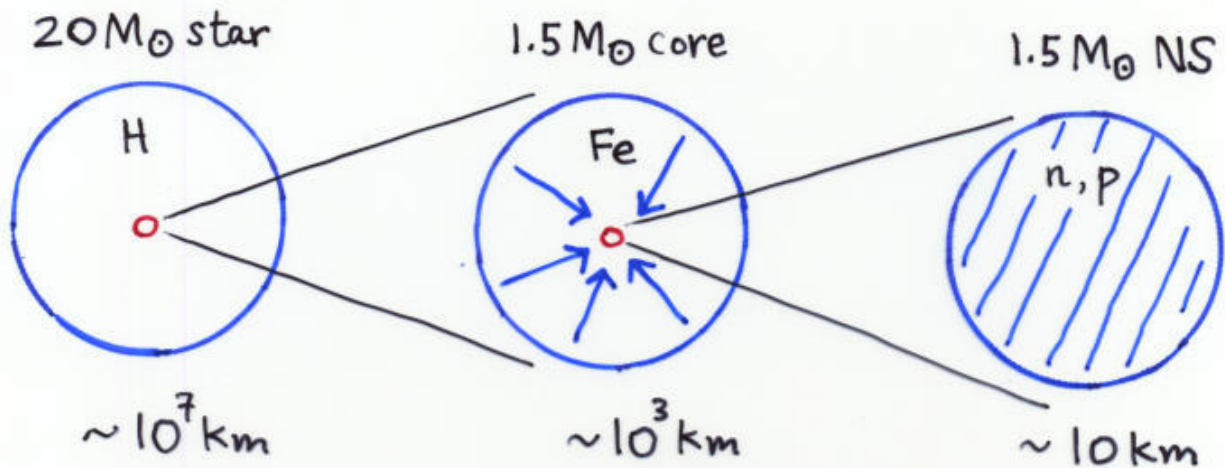
Theoretical Astrophysics Group

Fermilab

## Executive Summary:

No supernova neutrinos  
detected recently.

## Supernova: Core-Collapse



type-II SN : core collapse of an  $M > 8 M_{\odot}$  star

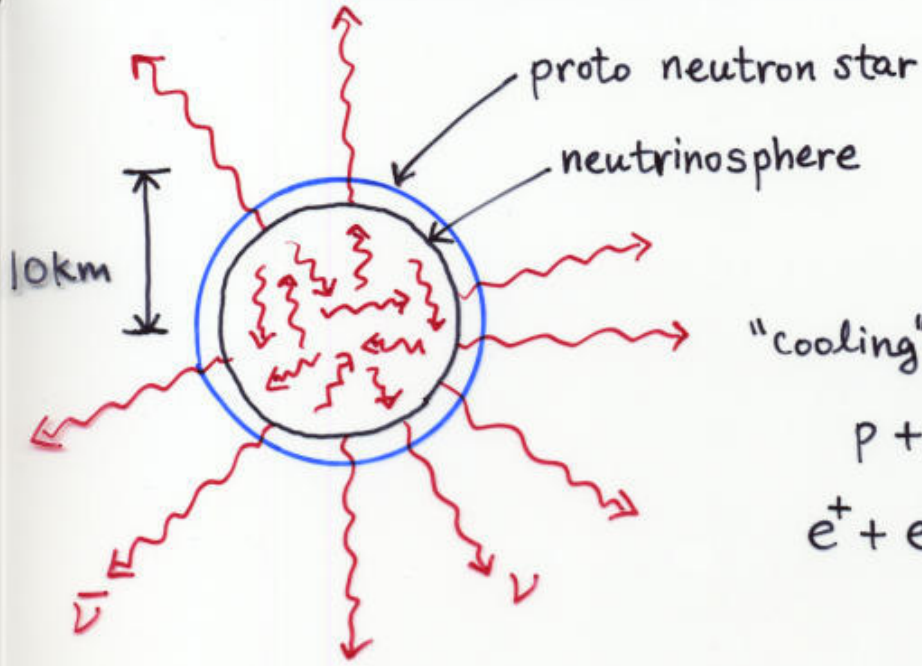
$$\Delta E_B \approx \frac{GM_{\odot}^2}{R_{NS}} - \frac{GM_{\odot}^2}{R_{core}} \approx \boxed{3 \times 10^{53} \text{ ergs}}$$
$$\approx 2 \times 10^{59} \text{ MeV}$$

observations :

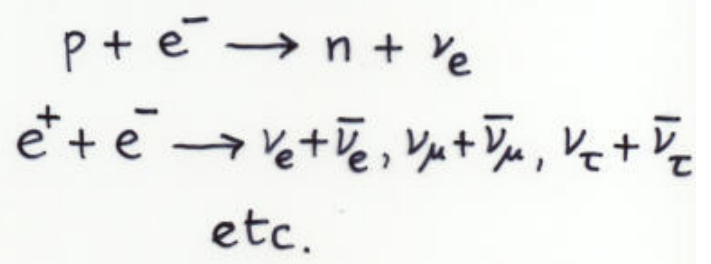
kinetic energy of explosion  $\approx 10^{-2} \cdot \Delta E_B$

electromagnetic radiation  $\approx 10^{-4} \cdot \Delta E_B$

# Supernova: Energy Release

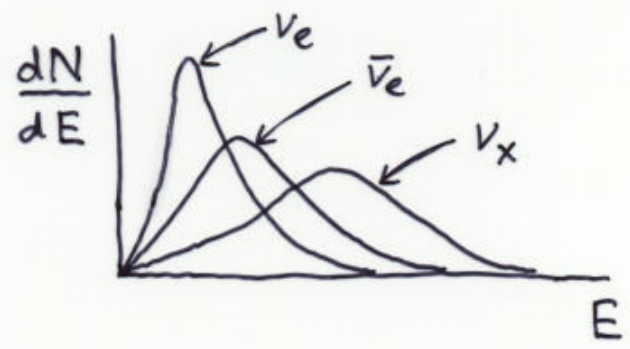


"cooling" by neutrino emission:

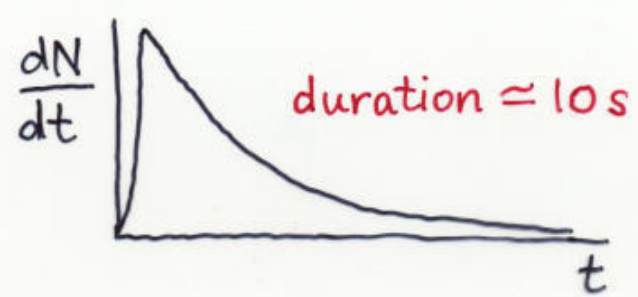


diffusion until  $\lambda = 1/\rho\sigma$  from surface, then escape

- $\langle E_{\nu_e} \rangle \approx 11 \text{ MeV}$
- $\langle E_{\bar{\nu}_e} \rangle \approx 16 \text{ MeV}$
- $\langle E_{\nu_x} \rangle \approx 25 \text{ MeV}$



$$L_{\nu_e}(t) \approx L_{\bar{\nu}_e}(t) \approx L_{\nu_x}(t)$$



## Supernova Models :

Bad news : Supernovae explode in nature,  
but not in computers.

Good news : Lots of progress.

1d models with full Boltzmann  
transport and Newtonian or full  
GR gravity don't explode.

[Janka et al.; Mezzacappa et al.]

Next : 1d  $\rightarrow$  2d  $\rightarrow$  3d.

Lots of work on emissivities  
and opacities.

[Janka ; Raffelt ; Horowitz ; others]

This is a very important problem.



## Detectors:

water	Super-Kamiokande	32 kton $H_2O$
	SNO	1 kton $D_2O$
	AMANDA	1.4 kton $H_2O$
		lots of ice

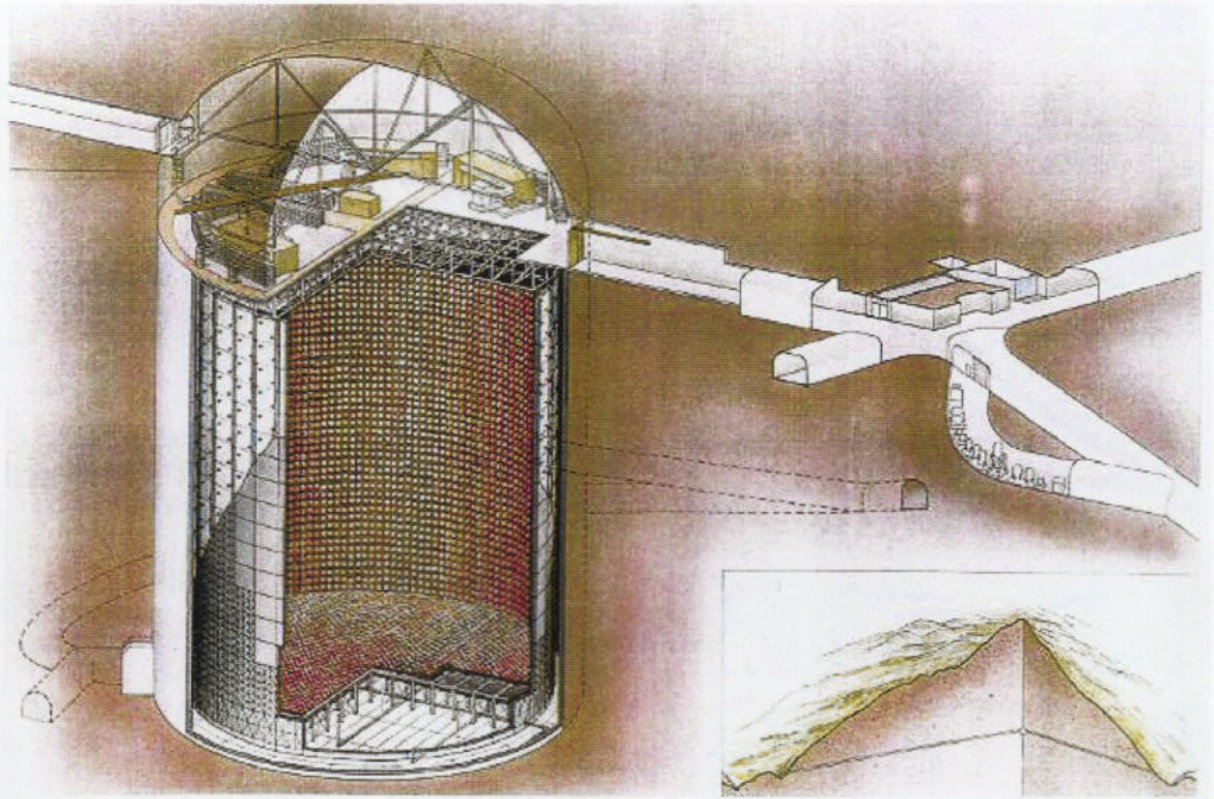
oil	LVD	1 kton
	KamLAND	1 kton
	MiniBooNE	0.6 kton
	Borexino	0.3 kton
	Baksan	0.2 kton

Biggest yield:  $\bar{\nu}_e p \rightarrow e^+ n$ ,  $\sim 300$  events/kton

Some other proposed targets:

Pb (OMNIS), Ar (ICARUS), Mo (MOON),

Cl (Homestake Hybrid),  $\sim 1$  Mton  $H_2O$



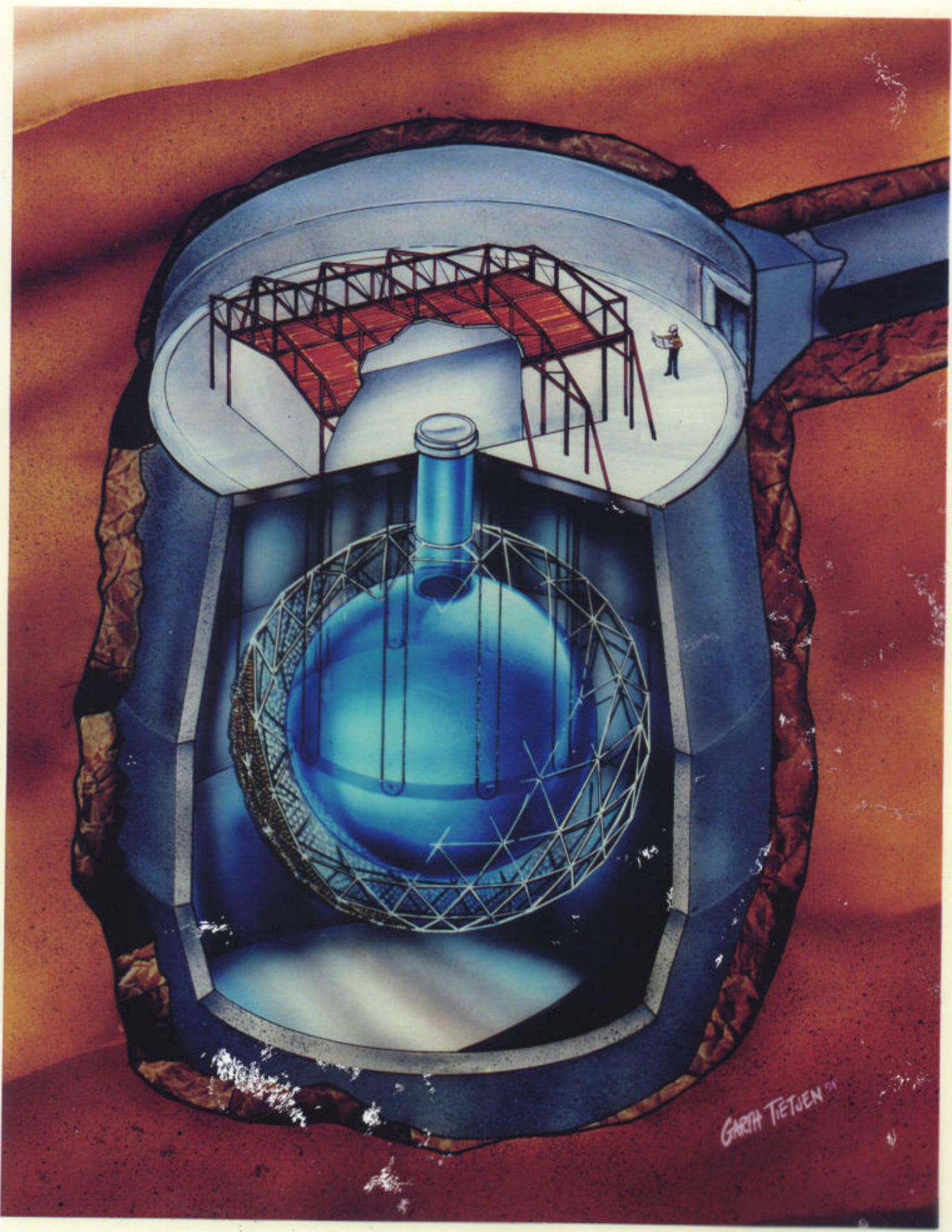
SUPERKAMIOKANDE INSTITUTE FOR COSMIC RAY RESEARCH UNIVERSITY OF TOKYO

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TABLE I. Calculated numbers of events expected in SK with a 5 MeV threshold and a supernova at 10 kpc. The other parameters (e.g., neutrino spectrum temperatures) are given in the text. In rows with two reactions listed, the number of events is the total for both. The second row is a subset of the first row that is an irreducible background to the reactions in the third and fourth rows.

Reaction	No. of events
$\bar{\nu}_e + p \rightarrow e^+ + n$	<i>detected particle: <math>e^+</math></i> 8300
$\bar{\nu}_e + p \rightarrow e^+ + n$ ( $E_{e^+} \leq 10$ MeV)	$e^+$ 530
$\nu_\mu + {}^{16}\text{O} \rightarrow \nu_\mu + \gamma + X$ $\bar{\nu}_\mu + {}^{16}\text{O} \rightarrow \bar{\nu}_\mu + \gamma + X$	$\gamma$ 355
$\nu_\tau + {}^{16}\text{O} \rightarrow \nu_\tau + \gamma + X$ $\bar{\nu}_\tau + {}^{16}\text{O} \rightarrow \bar{\nu}_\tau + \gamma + X$	$\gamma$ 355
$\nu_e + e^- \rightarrow \nu_e + e^-$ $\bar{\nu}_e + e^- \rightarrow \bar{\nu}_e + e^-$	$e^-$ 200
$\nu_\mu + e^- \rightarrow \nu_\mu + e^-$ $\bar{\nu}_\mu + e^- \rightarrow \bar{\nu}_\mu + e^-$	$e^-$ 60
$\nu_\tau + e^- \rightarrow \nu_\tau + e^-$ $\bar{\nu}_\tau + e^- \rightarrow \bar{\nu}_\tau + e^-$	$e^-$ 60





GARTH TETJEN '91

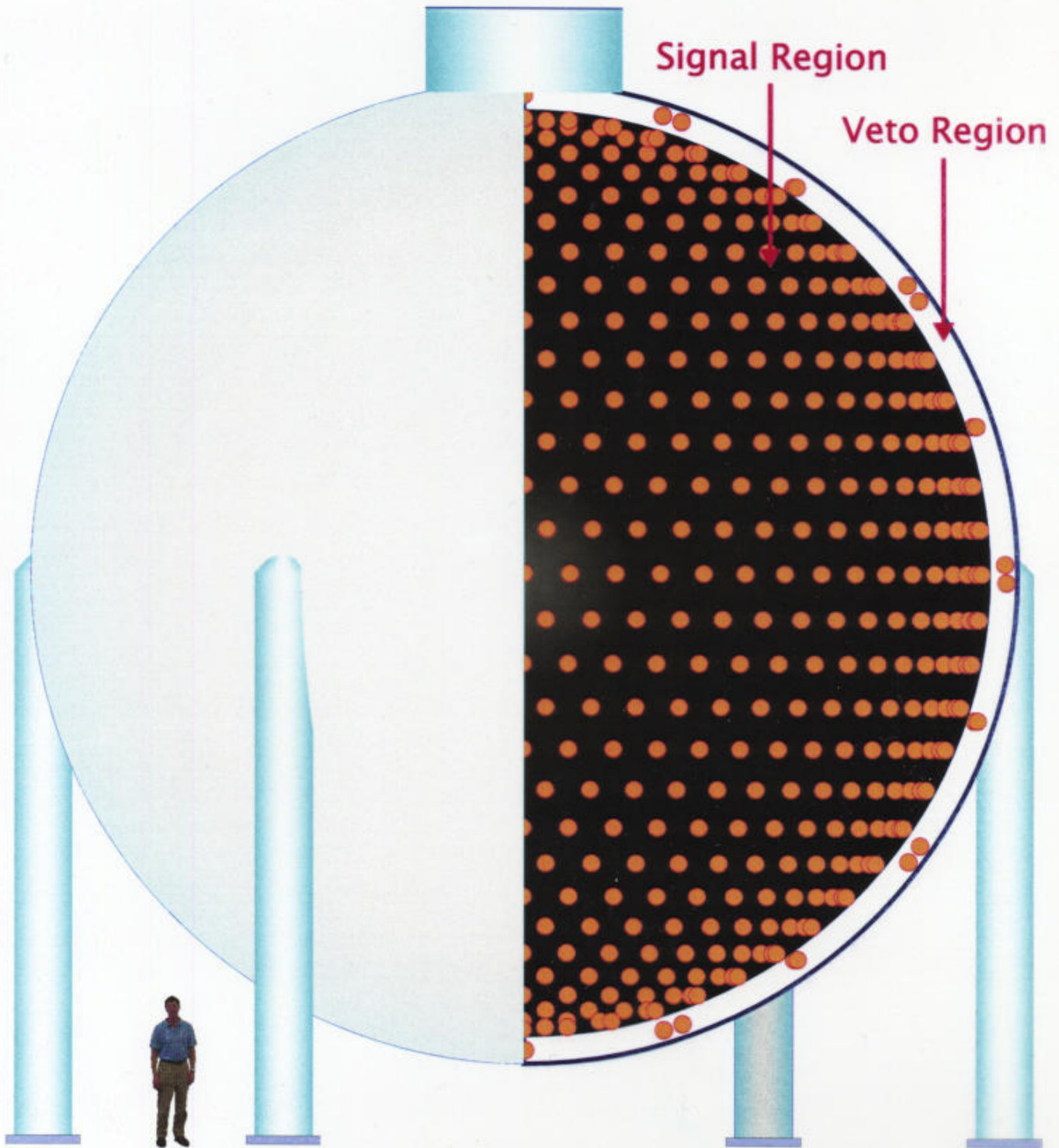


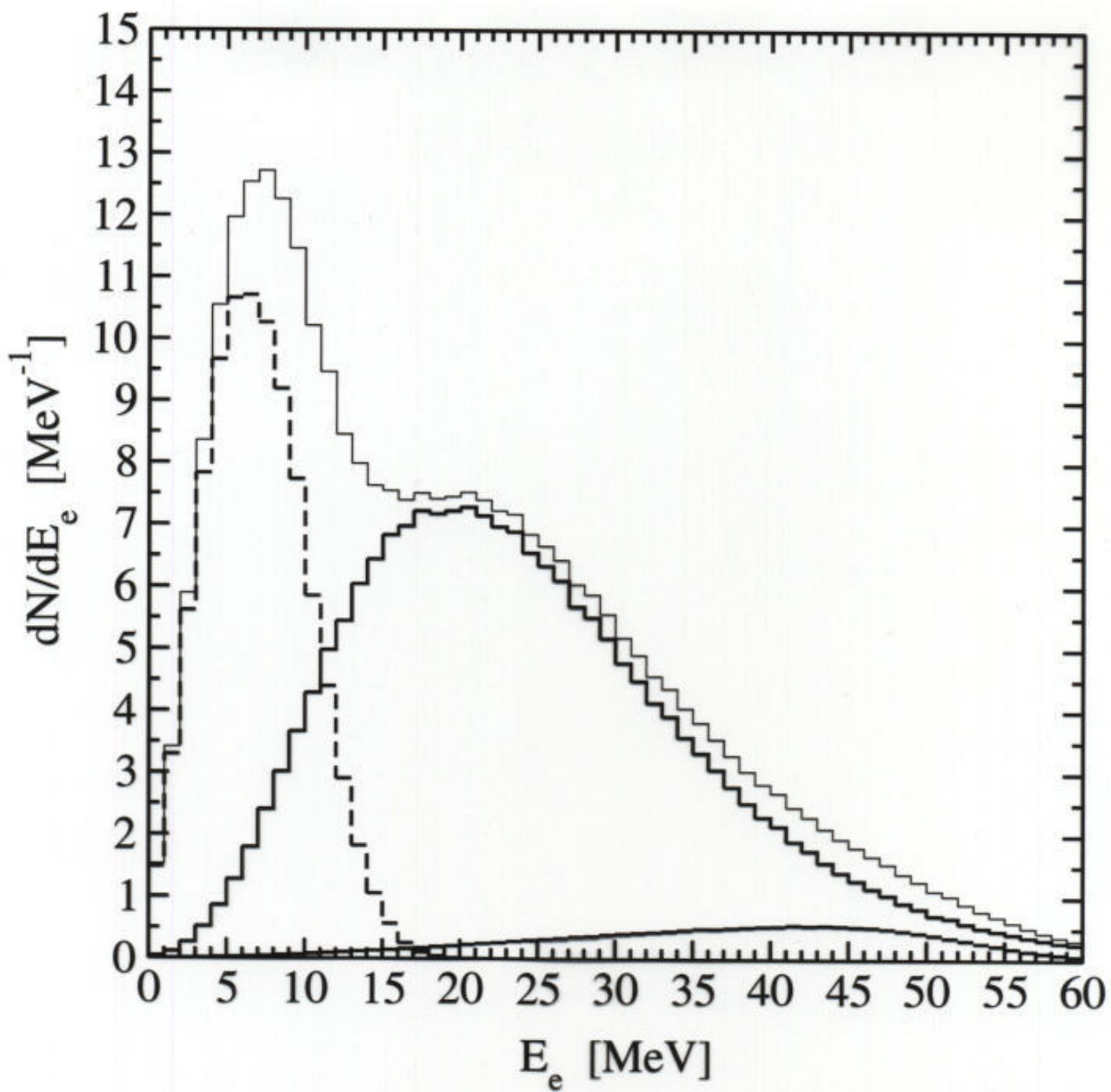
TABLE I. Calculated numbers of events expected in SNO for a supernova at 10 kpc. The other parameters (e.g., neutrino spectrum temperatures) are given in the text. In rows with two reactions listed, the number of events is the total for both. The notation  $\nu$  indicates the sum of  $\nu_e$ ,  $\nu_\mu$ , and  $\nu_\tau$ , though they do not contribute equally to a given reaction, and  $X$  indicates either  $n+^{15}\text{O}$  or  $p+^{15}\text{N}$ .

Events in 1 kton D <sub>2</sub> O		
$\nu + d \rightarrow \nu + p + n$	detected particle(s) : n	485
$\bar{\nu} + d \rightarrow \bar{\nu} + p + n$		
$\nu_e + d \rightarrow e^- + p + p$	$e^-, e^+ nn$	160
$\bar{\nu}_e + d \rightarrow e^+ + n + n$		
$\nu + ^{16}\text{O} \rightarrow \nu + \gamma + X$	$\gamma, \gamma n$	20
$\bar{\nu} + ^{16}\text{O} \rightarrow \bar{\nu} + \gamma + X$		
$\nu + ^{16}\text{O} \rightarrow \nu + n + ^{15}\text{O}$	n	15
$\bar{\nu} + ^{16}\text{O} \rightarrow \bar{\nu} + n + ^{15}\text{O}$		
$\nu + e^- \rightarrow \nu + e^-$	$e^-$	10
$\bar{\nu} + e^- \rightarrow \bar{\nu} + e^-$		
Events in 1.4 kton H <sub>2</sub> O		
$\bar{\nu}_e + p \rightarrow e^+ + n$	$e^+$	365
$\nu + ^{16}\text{O} \rightarrow \nu + \gamma + X$	$\gamma$	30
$\bar{\nu} + ^{16}\text{O} \rightarrow \bar{\nu} + \gamma + X$		
$\nu + e^- \rightarrow \nu + e^-$	$e^-$	15
$\bar{\nu} + e^- \rightarrow \bar{\nu} + e^-$		

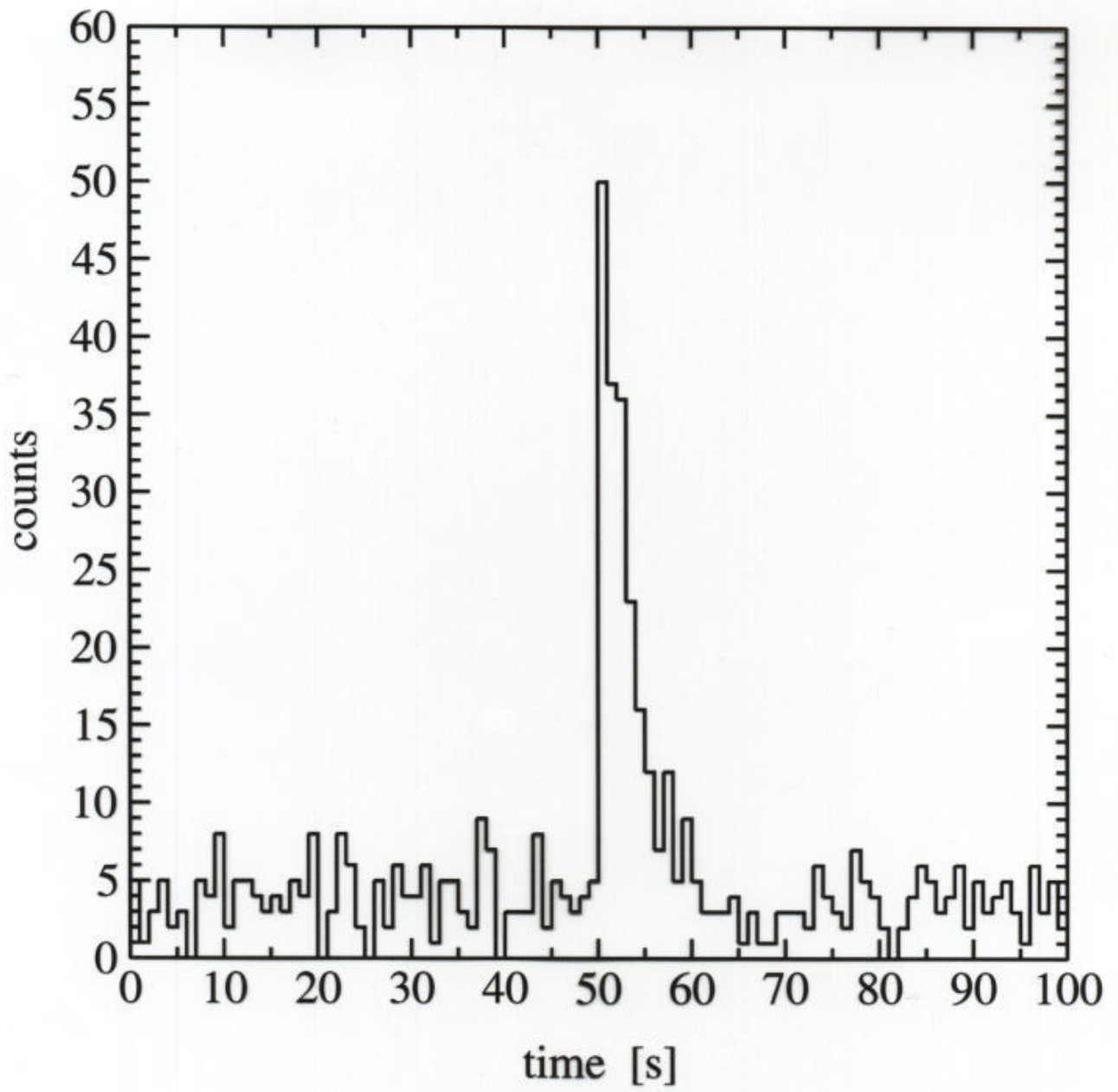
- detected particles
- NC : dominated by  $\nu_\mu, \nu_\tau$
- CC :  $\nu_e, \bar{\nu}_e$  only
- CC, NC separation easy

# MiniBooNE Detector









## CC Measurements:

At 10 kpc:

$$\text{SK: } \simeq 8000 \quad \bar{\nu}_e + p \rightarrow e^+ + n$$

$$\text{SNO: } \simeq 80 \quad \nu_e + d \rightarrow e^- + p + p$$

} "clean"  
reactions

spectral quality:

$$E_\nu = E_e + \Delta + \mathcal{O}(1/M_p)$$

$$\rightarrow \frac{dN_e}{dE_e} \sim f(E_\nu) \sigma(E_\nu)$$

shape:

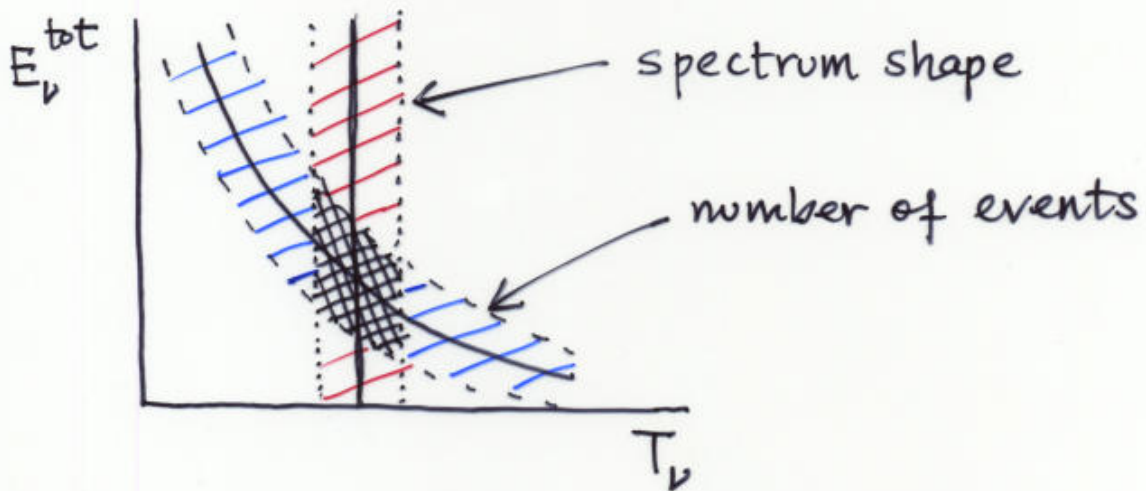
$$\frac{\delta T}{T} \sim \frac{\delta \langle E_\nu \rangle_{fo}}{\langle E_\nu \rangle_{fo}} \sim \frac{\delta \langle E_e \rangle}{\langle E_e \rangle} \sim \frac{1}{\sqrt{N}}$$

$$\rightarrow \begin{array}{l} \frac{\delta T_{\bar{\nu}_e}}{T_{\bar{\nu}_e}} \sim 1\% \quad \text{SK} \\ \frac{\delta T_{\nu_e}}{T_{\nu_e}} \sim 10\% \quad \text{SNO} \end{array}$$

normalization:

$$N \sim \frac{E_\nu^{\text{tot}} \langle \sigma \rangle}{T_\nu}$$

→  $N \sim E_\nu^{\text{tot}} T_\nu$  for  $\sigma \sim E_\nu^2$



These are the most basic observables (model comparisons,  $E_B$ , r-process oscillations).

Next step: time dependence

$$\frac{dN}{dt}(t) \sim L(t), \text{ modulo } T(t) \text{ effects}$$

## NC Measurements:

$\nu_\mu, \nu_\tau, \bar{\nu}_\mu, \bar{\nu}_\tau$  : No CC reactions  
: dominate the NC signals

Why are these so important?

- $\approx \frac{2}{3}$  of binding energy is radiated in these flavors
- not observed in SN1987A
- equipartition untested
- temperature controversial
- generally unaffected by oscillations
- needed for computing  $E_B \sim \frac{GM^2}{R}$



BUT .....  $E_\nu$  not measured in NC reactions!

at 10 kpc:

SK:  $\approx 120$   $\nu + e^- \rightarrow \nu + e^-$   
recoil spectrum  
very hard to separate events

SNO:  $\approx 500$   $\nu + d \rightarrow \nu + p + n$   
detect n capture

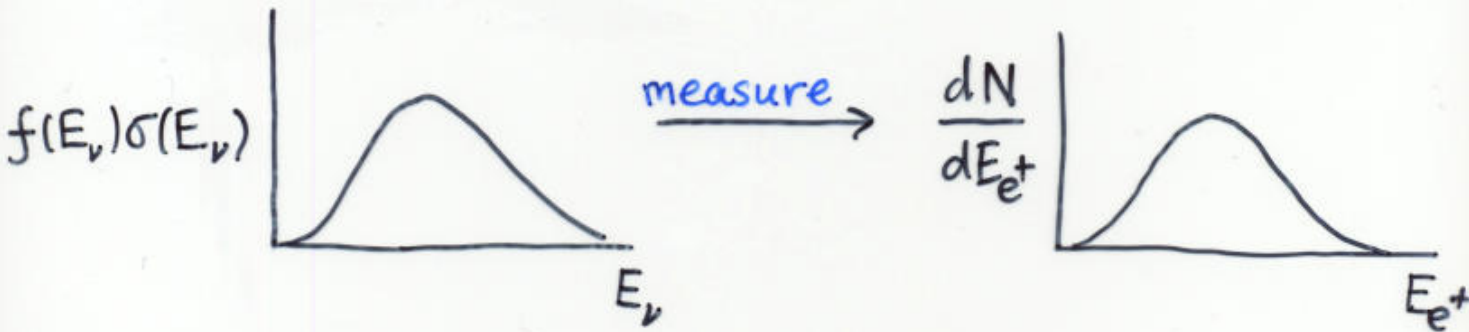
SK:  $\approx 700$   $\nu + {}^{16}\text{O} \rightarrow$   
 ${}^{15}\text{O} + n + \gamma$   
 ${}^{15}\text{N} + p + \gamma$   
 $E_\gamma \approx 5-10 \text{ MeV}$

Kolbe  
Langmuir  
Vogel

KamLAND:  $\approx 60$   $\nu + {}^{12}\text{C} \rightarrow \nu + {}^{12}\text{C} + \gamma$   
 $E_\gamma = 15.11 \text{ MeV}$

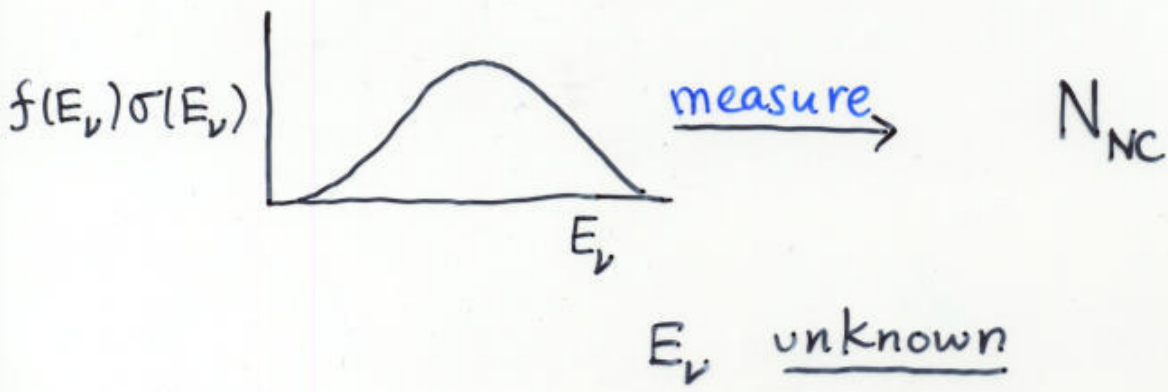
$$N \sim \int dE_\nu f(E_\nu) \sigma(E_\nu)$$

- charged-current, e.g.,  $\bar{\nu}_e p \rightarrow e^+ n$



since  $E_{e^+} \approx E_\nu - 1.8 \text{ MeV}$

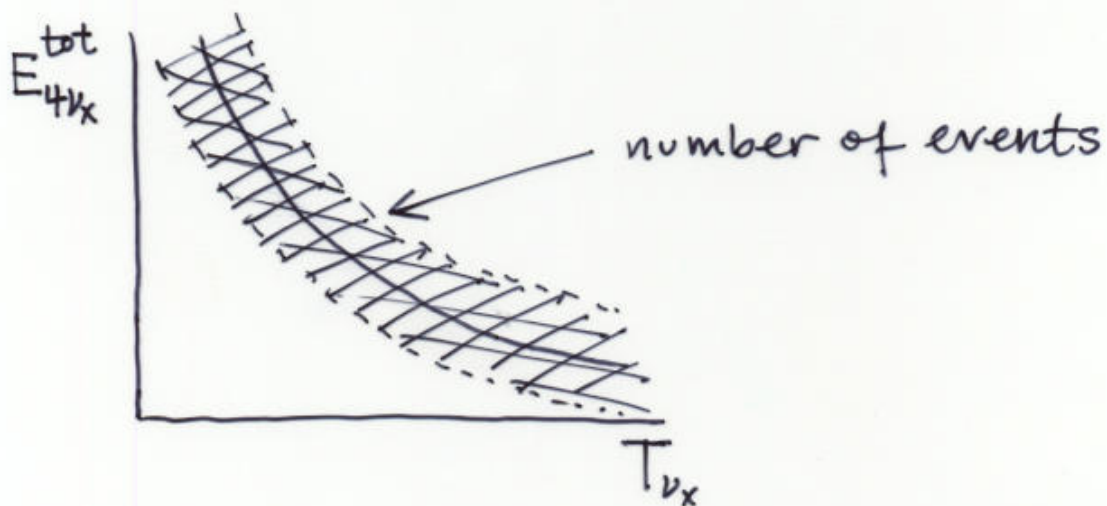
- neutral-current, e.g.,  $\nu d \rightarrow \nu pn$



So, can only measure

$$N \sim E_{4\nu_x}^{\text{tot}} \frac{\langle \sigma \rangle}{T_{\nu_x}}$$

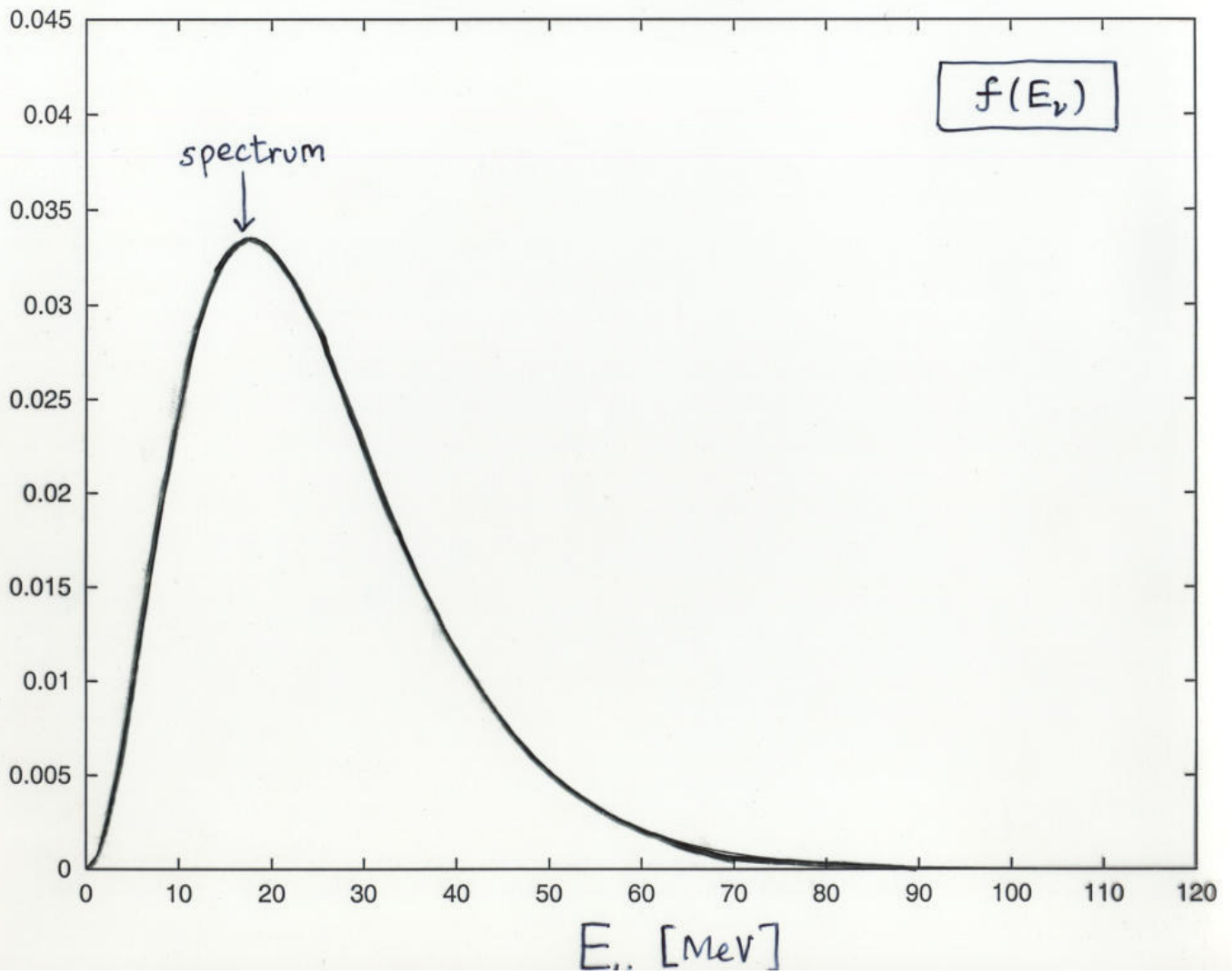
→  $N \sim E_{4\nu_x}^{\text{tot}} T_{\nu_x}$  for  $\sigma \sim E_\nu^2$



How can we break the degeneracy?

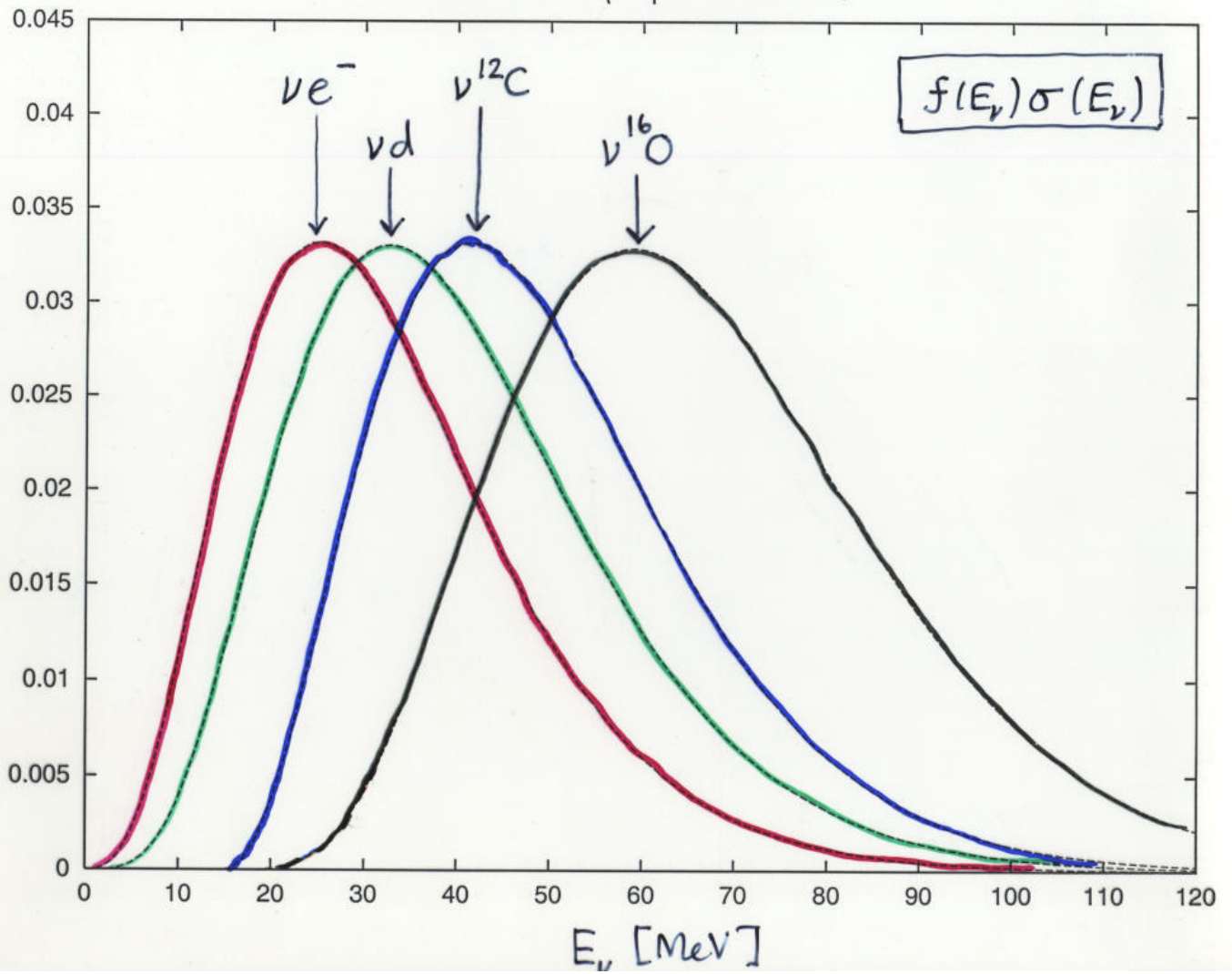
Option A: exploit different spectral responses

Option B: see below





J.F. Beacom, hep-ph/9909231



## Detector

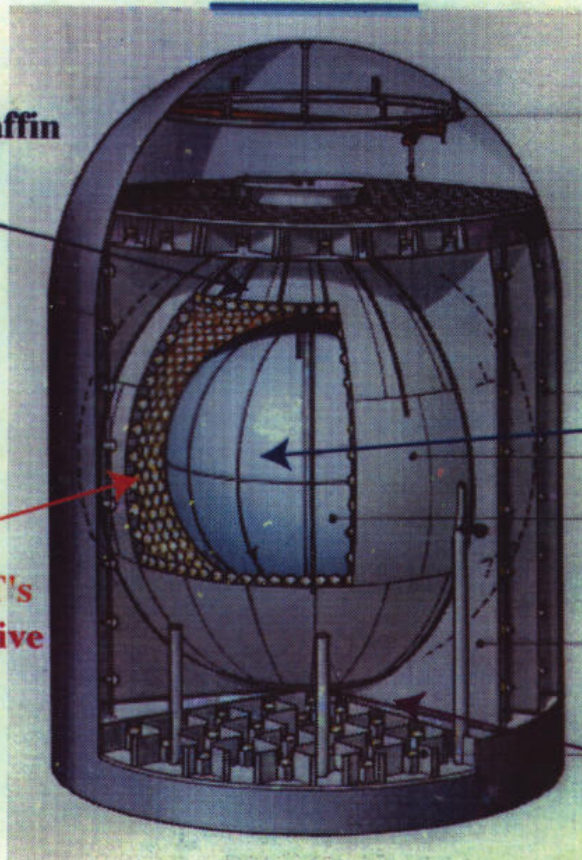
3,000 m<sup>3</sup> Isoparaffin  
Buffer Tank

1,200 m<sup>3</sup> Scintillator  
Balloon

1,280 17 inch-PMT's  
(~ 22% photosensitive  
coverage)

+  
630 PMTs  
(U.S. plan)

Water Cherenkov  
Anti-Counter

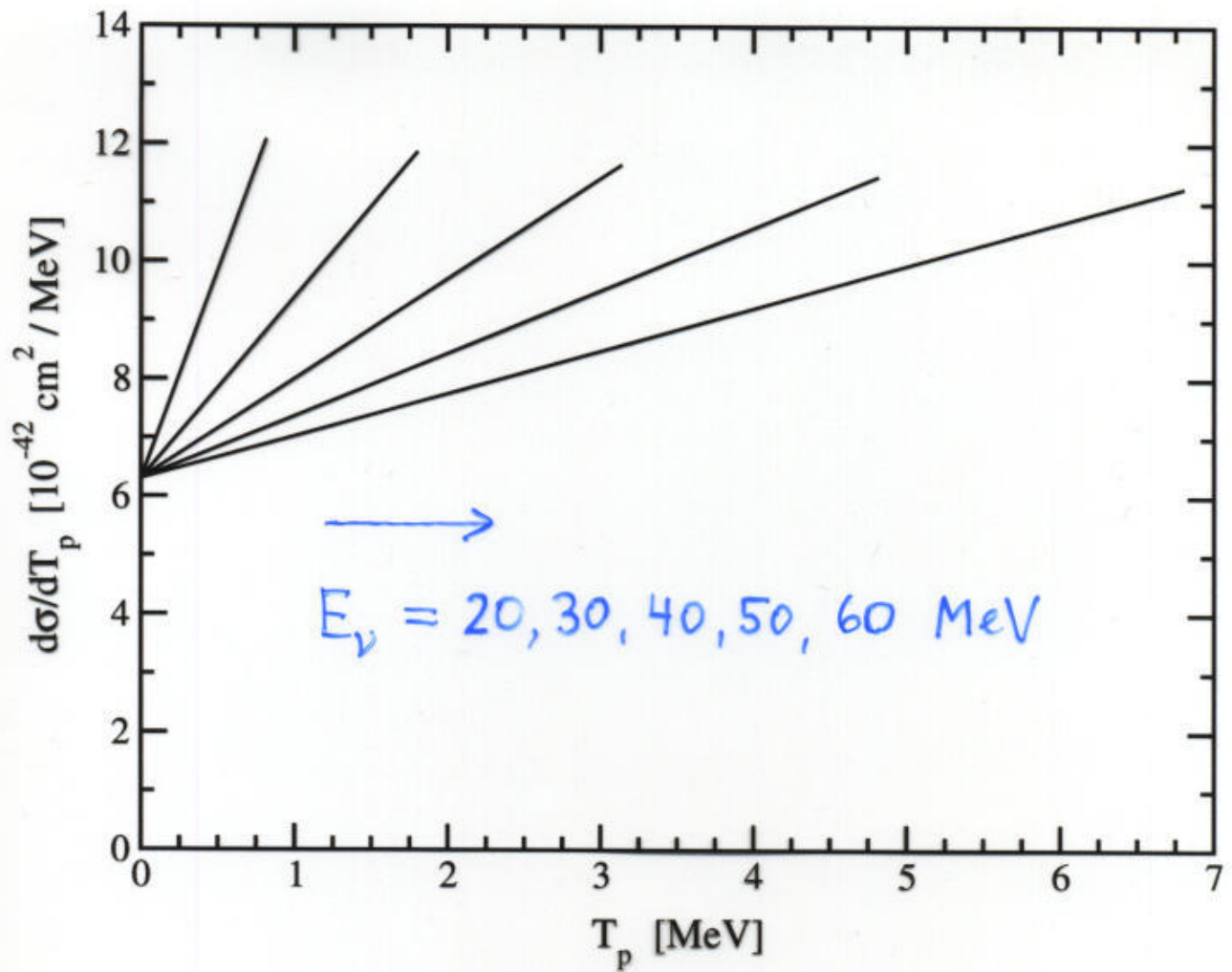


## $\nu + p \rightarrow \nu + p$ in KamLAND:

Why this seems crazy .....but..... why it isn't

- |  |   |
|--|---|
| 1. $T_p \sim \frac{E_\nu^2}{M_p} \sim \text{MeV}$  | 1'. scintillation detector  |
| 2. light quenching<br>( $\approx 14$ for alphas)   | 2'. quenching $\approx 4$ for protons,<br>threshold $\approx 0.2 \text{ MeV}$ |
| 3. NC vector coupling<br>$\sim 1 - 4 \sin^2 \theta_w \approx 0$  | 3'. NC axial coupling ok,<br>favors large $T_p$                               |
| 4. NC cross section<br>$\approx \frac{1}{4}$ CC $\bar{\nu}_e + p$ cross section  | 4'. 4 flavors contribute,<br>$E_\nu$ is higher                                |
| 5. $\nu_e, \bar{\nu}_e$ NC contributions   | 5'. $\nu_\mu, \nu_\tau, \bar{\nu}_\mu, \bar{\nu}_\tau$ dominate               |
| 6. confusion with<br>$\bar{\nu}_e + p \rightarrow e^+ + n$<br>$\nu + e^- \rightarrow \nu + e^-$<br>$\nu + {}^{12}\text{C} \rightarrow \nu + {}^{12}\text{C} + \gamma$<br>etc | 6'. easy to separate<br>in practice   |





$\nu p \rightarrow \nu p$  (and  $\bar{\nu} p \rightarrow \bar{\nu} p$ )

$$\frac{d\sigma}{dT_p} \simeq \frac{G_F^2 M_p}{\pi} \left( 1 + \frac{T_p}{T_p^{\text{max}}} \right) C_A^2$$

$$T_p^{\text{max}} \simeq \frac{2E_\nu^2}{M_p}$$

$$\sigma^{\text{tot}} \simeq \frac{G_F^2 E_\nu^2}{\pi} \cdot 3C_A^2$$



## Comparison to $\nu A \rightarrow \nu A$ :

(Freedman 1974, Drukier & Stodolsky 1984)

$$E_\nu \lesssim 50 \text{ MeV} \longrightarrow \lambda_\nu \gtrsim 4 \text{ fm}$$

coherent scattering from whole  $^{12}\text{C}$  nucleus

$$\sigma_{\nu^{12}\text{C}}^{\text{tot}} \approx \frac{G_F^2 N^2 E_\nu^2}{\pi} \sim \sigma_{\nu p}^{\text{tot}}$$

but

$$(i) T_{\nu^{12}\text{C}}^{\text{max}} = \frac{1}{12} \cdot T_{\nu p}^{\text{max}}$$

(ii)  $\frac{d\sigma}{dT}$  strongly favors  $T=0$

(iii) light output from recoil  $^{12}\text{C}$   
is heavily quenched

## Comparison to $\nu e^- \rightarrow \nu e^-$ :

Easy to detect an event.

Simple two-body kinematics.

but

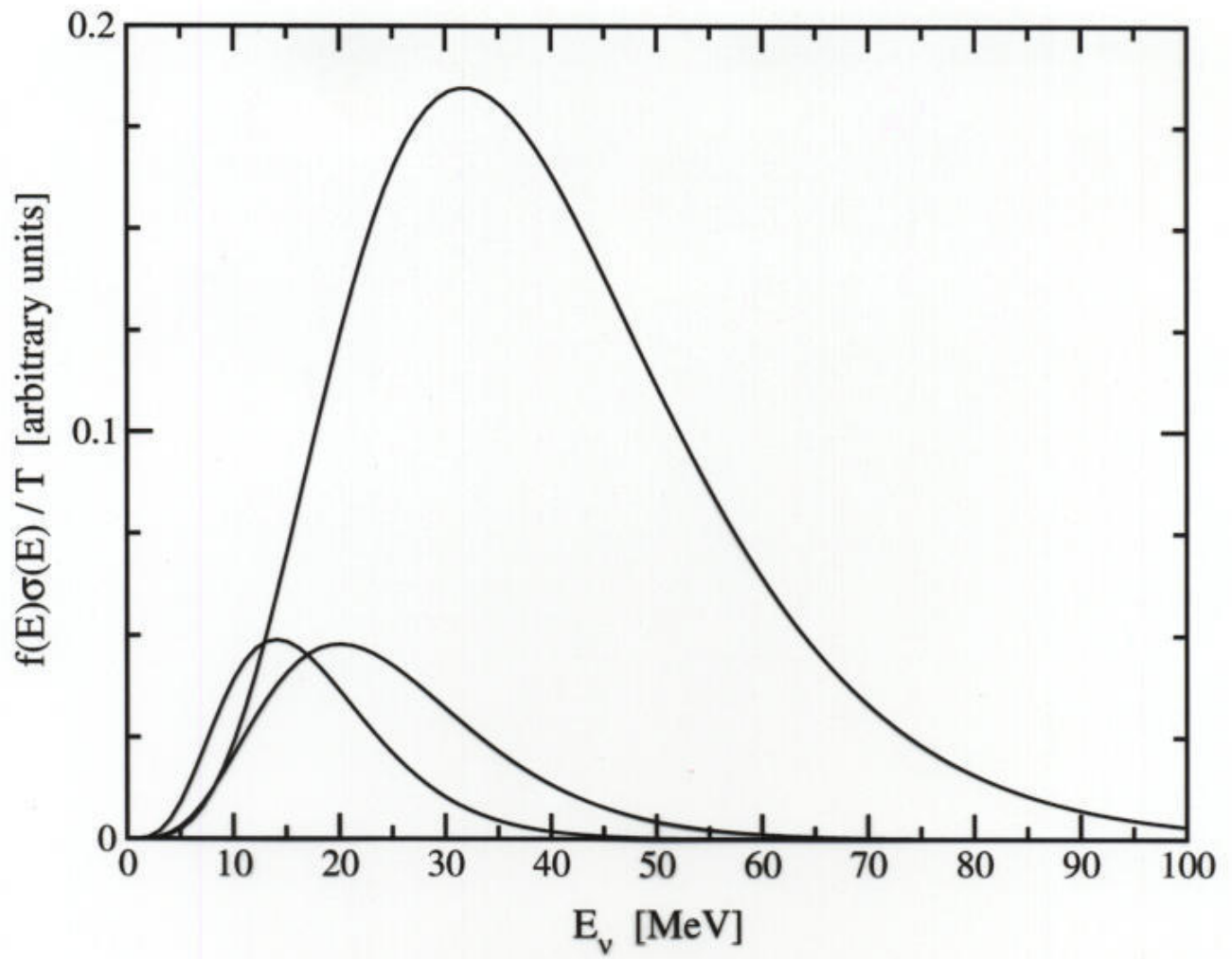
(i) Can't measure  $E_\nu$  for each event.

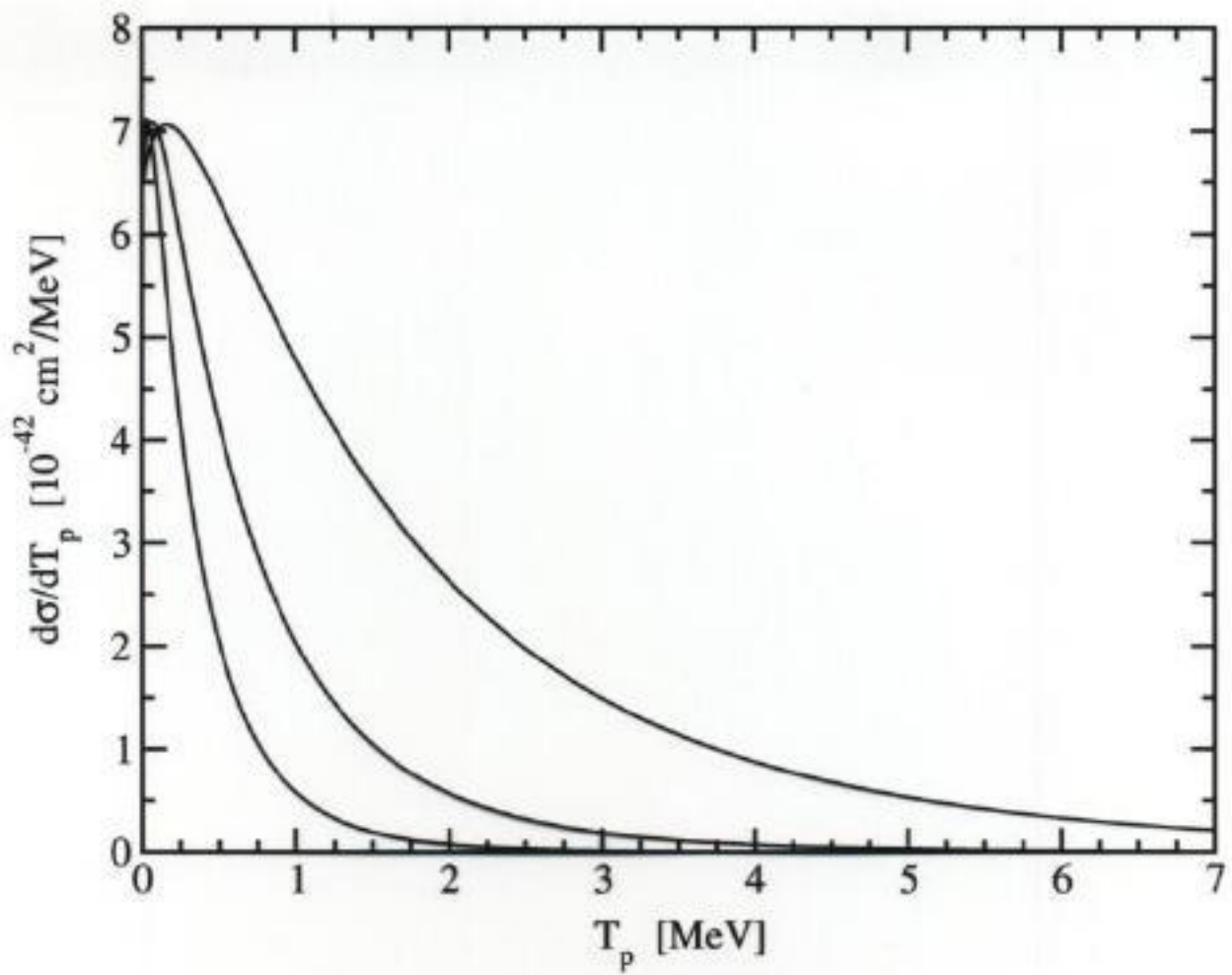
(ii) Hard to measure spectrum well.

(iii)  $\sigma_{\nu e^-}^{\text{tot}} \sim \frac{m_e}{E_\nu} \cdot \sigma_{\nu p}^{\text{tot}}$

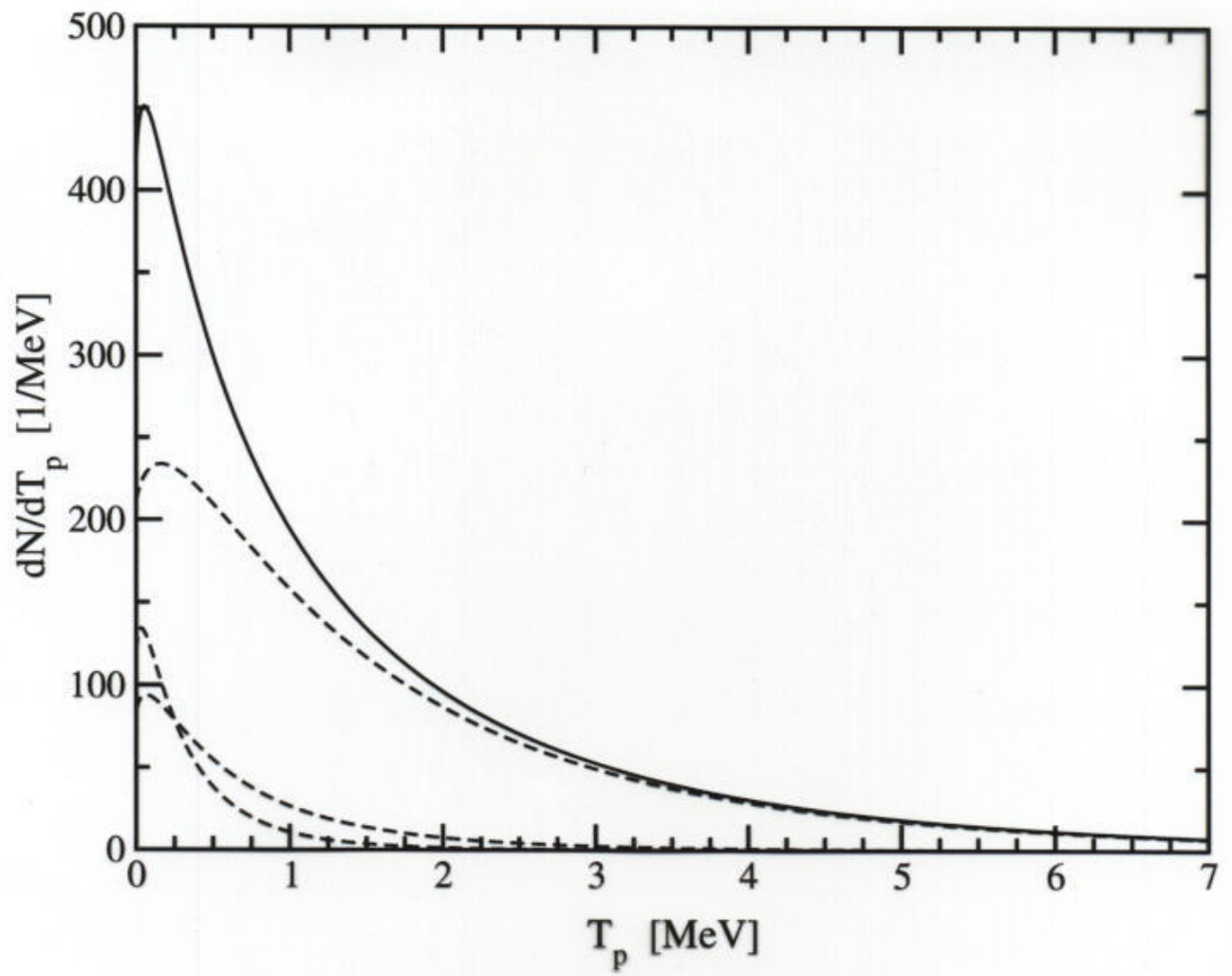
(iv)  $T_{\nu e^-}^{\text{max}} \simeq E_\nu$  vs.  $T_{\nu p}^{\text{max}} \simeq \frac{2E_\nu^2}{M_p}$

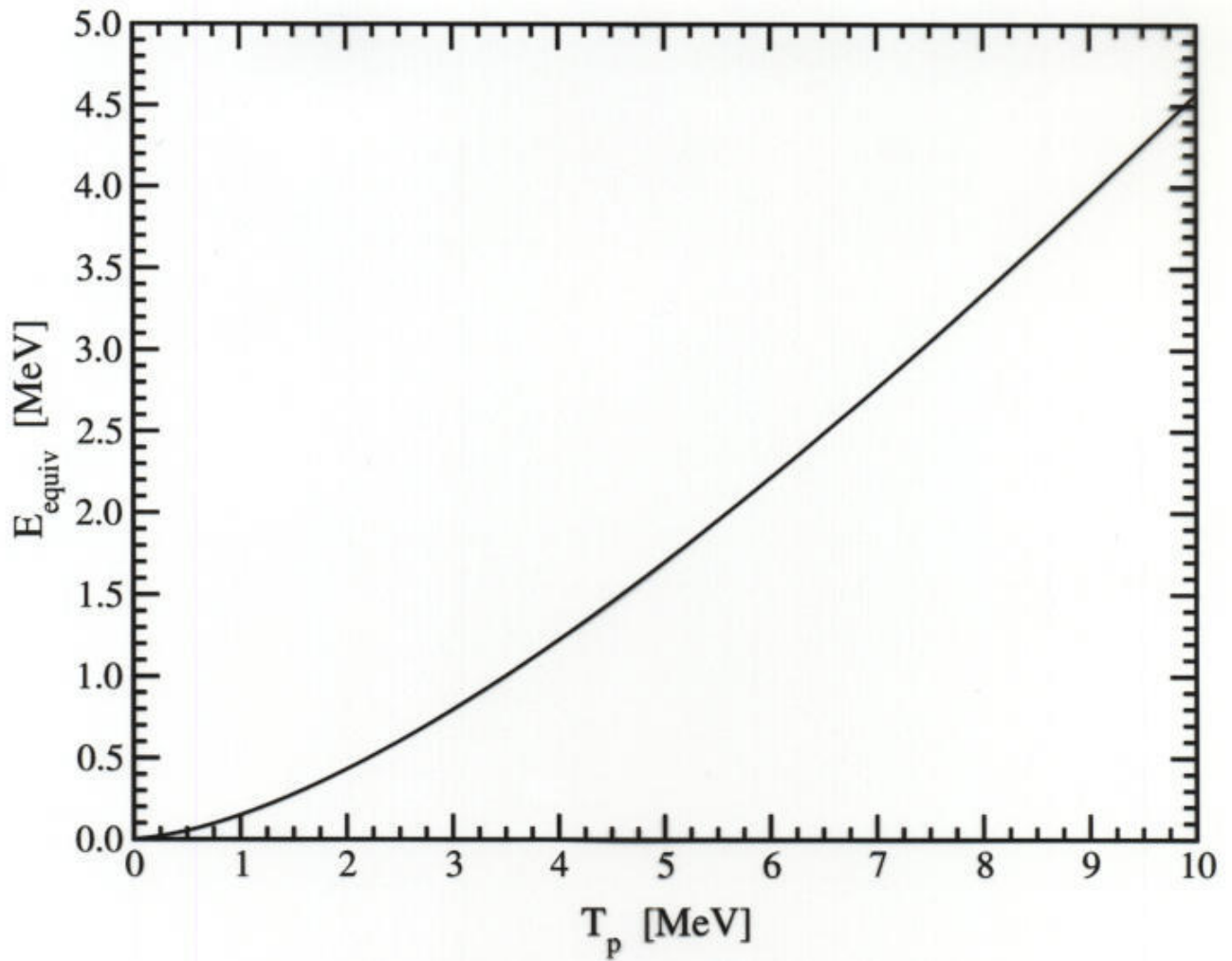
(v)  $\frac{d\sigma}{dT}$  weakly favors  $T=0$



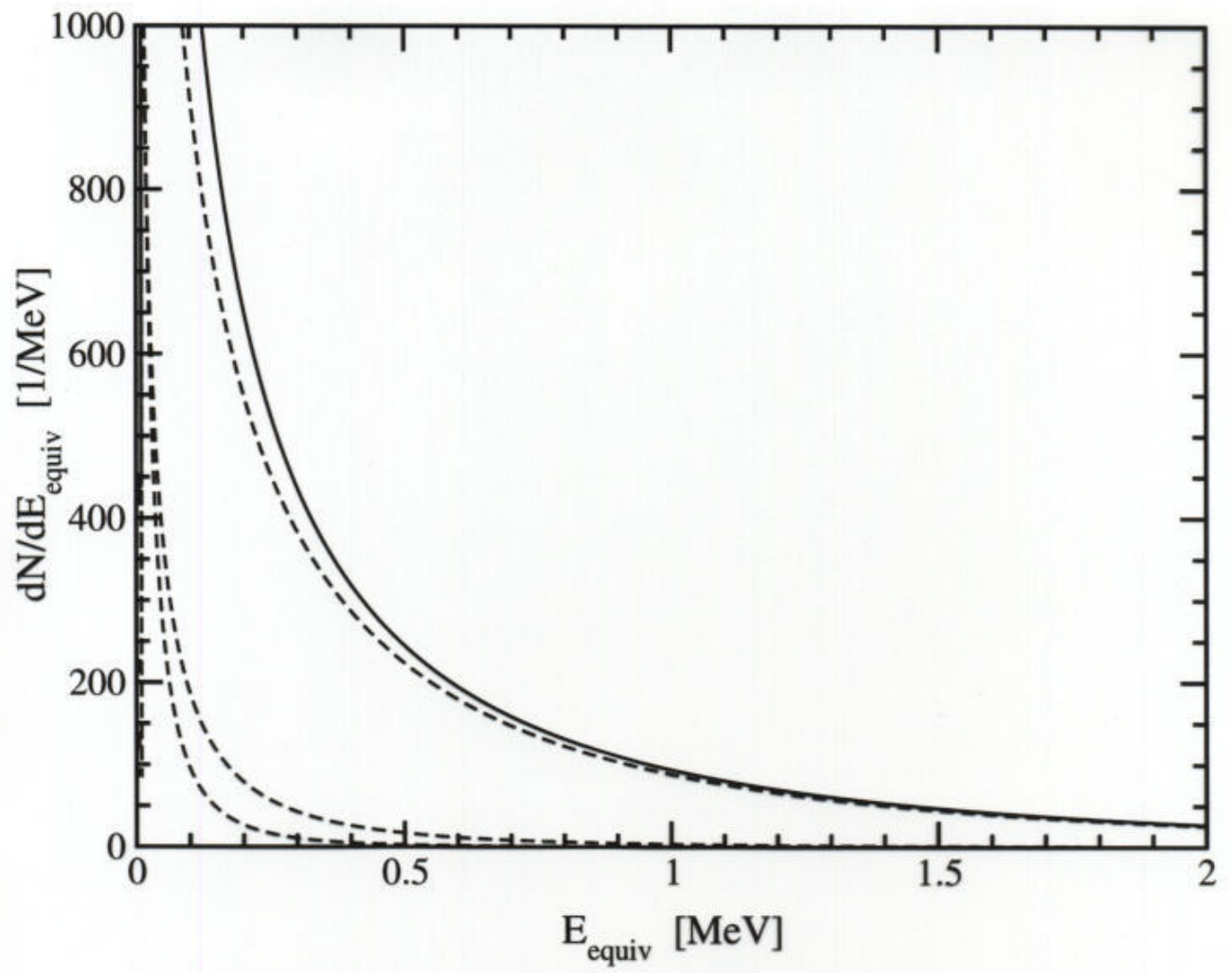


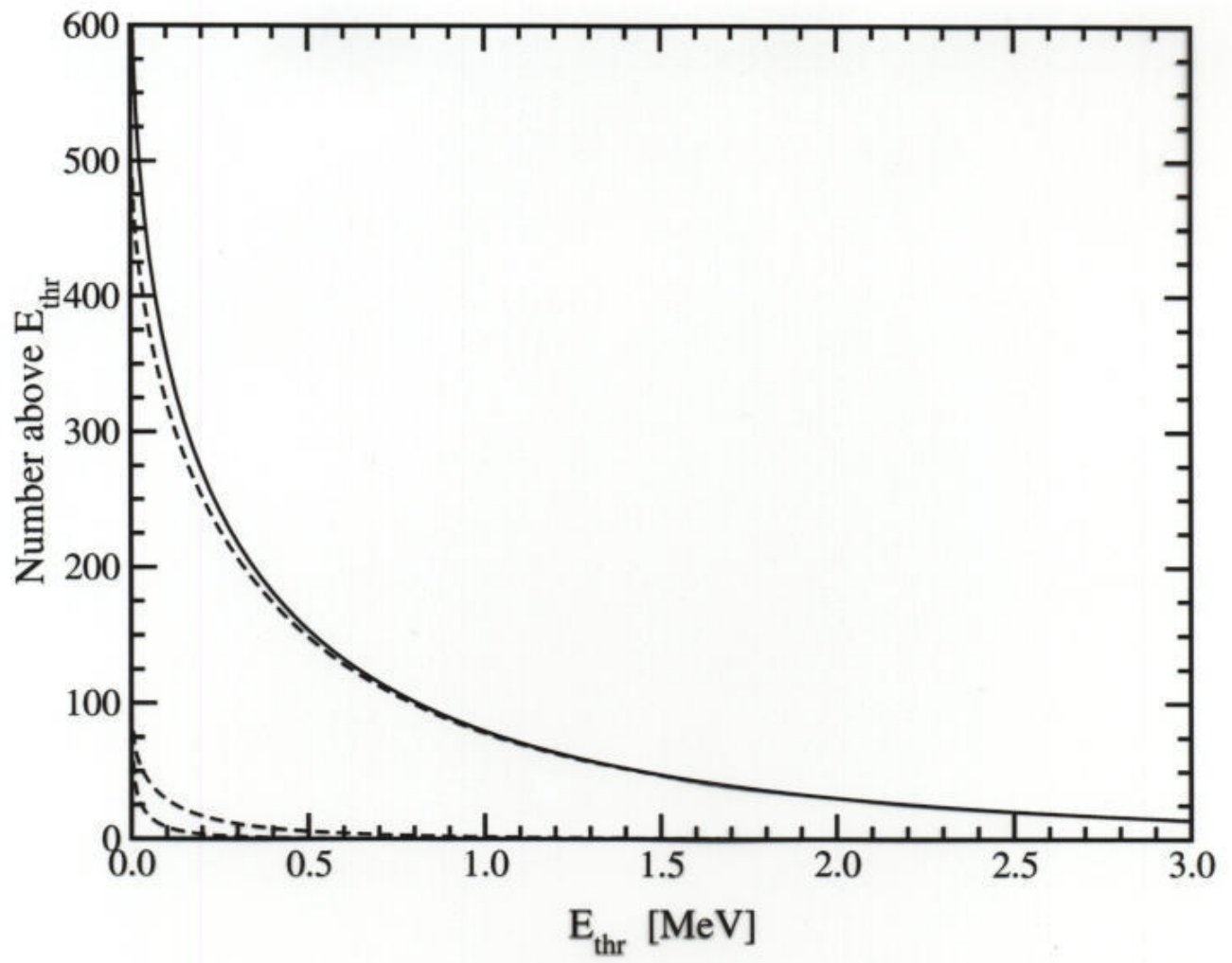






$$\frac{E_{\text{equiv}}}{T_p} = \text{"quenching factor"}$$





<u>flavor</u>	<u><math>E_{thr} = 0</math></u>	<u><math>E_{thr} = 0.2 \text{ MeV}</math></u>
$\nu_e$	60	5
$\bar{\nu}_e$	80	20
$\nu_\mu + \nu_\tau + \bar{\nu}_\mu + \bar{\nu}_\tau$	490	250
ALL	630	275

bottom line:

- large NC sample
- totally clean
- proton recoil spectrum reflects the incoming neutrino spectrum

→ Probably the best way to measure  $E_{4\nu_x}^{\text{tot}}$  and  $T_{\nu_x}$ .



## Two Sanity Checks:

- $\nu + p \rightarrow \nu + p$

$$E_\nu \approx 45 \text{ MeV} \longrightarrow T_p \approx 4 \text{ MeV}$$

or

$$n + p \rightarrow n + p$$

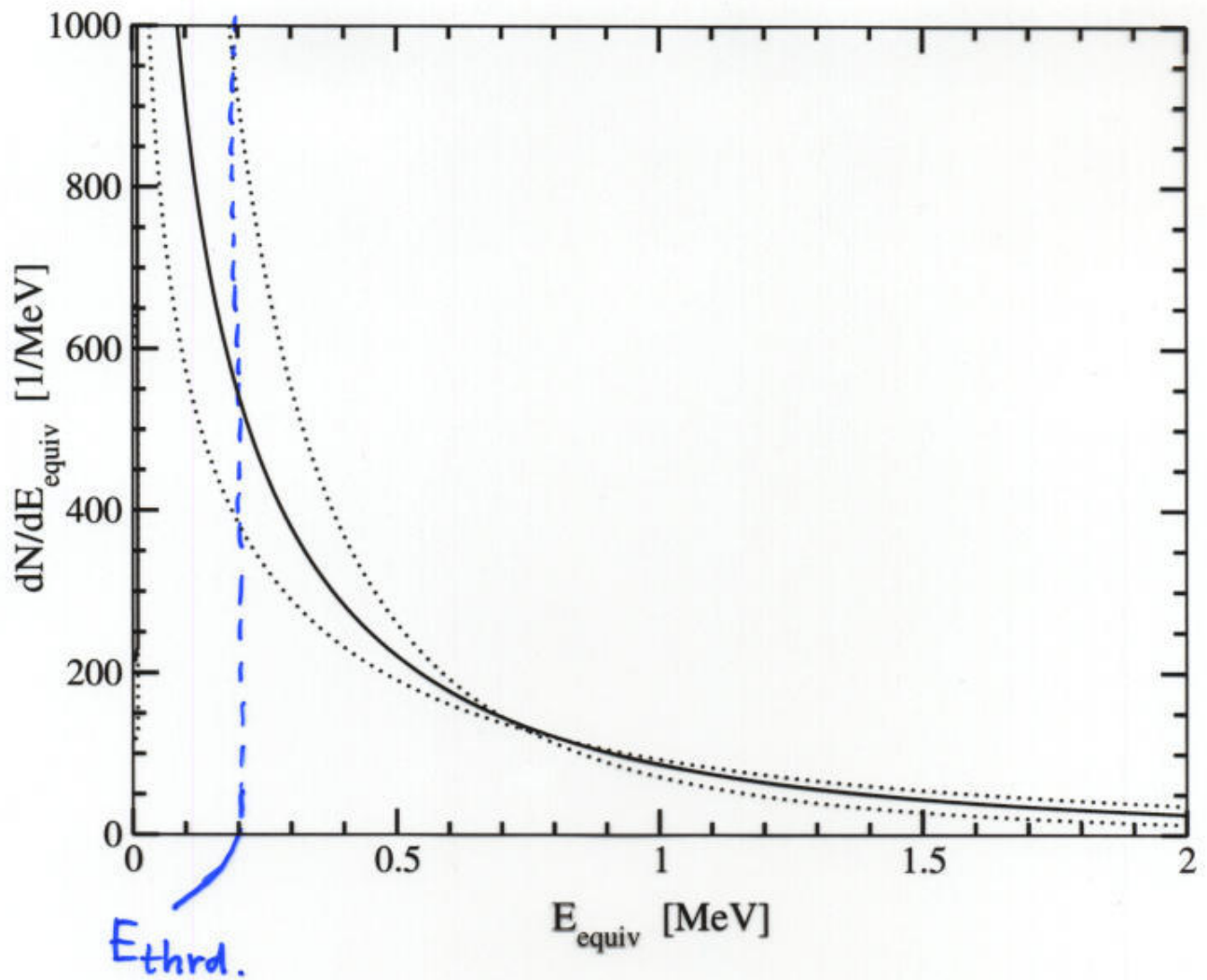
$$T_n \approx 8 \text{ MeV} \longrightarrow T_p \approx 4 \text{ MeV}$$

All aspects of detection are identical

- KamLAND solar background requirement above a similar threshold is

$$\text{Rate} \lesssim 10^{-3} \text{ Hz.}$$

We only need Rate  $\lesssim 1 \text{ Hz}$ .

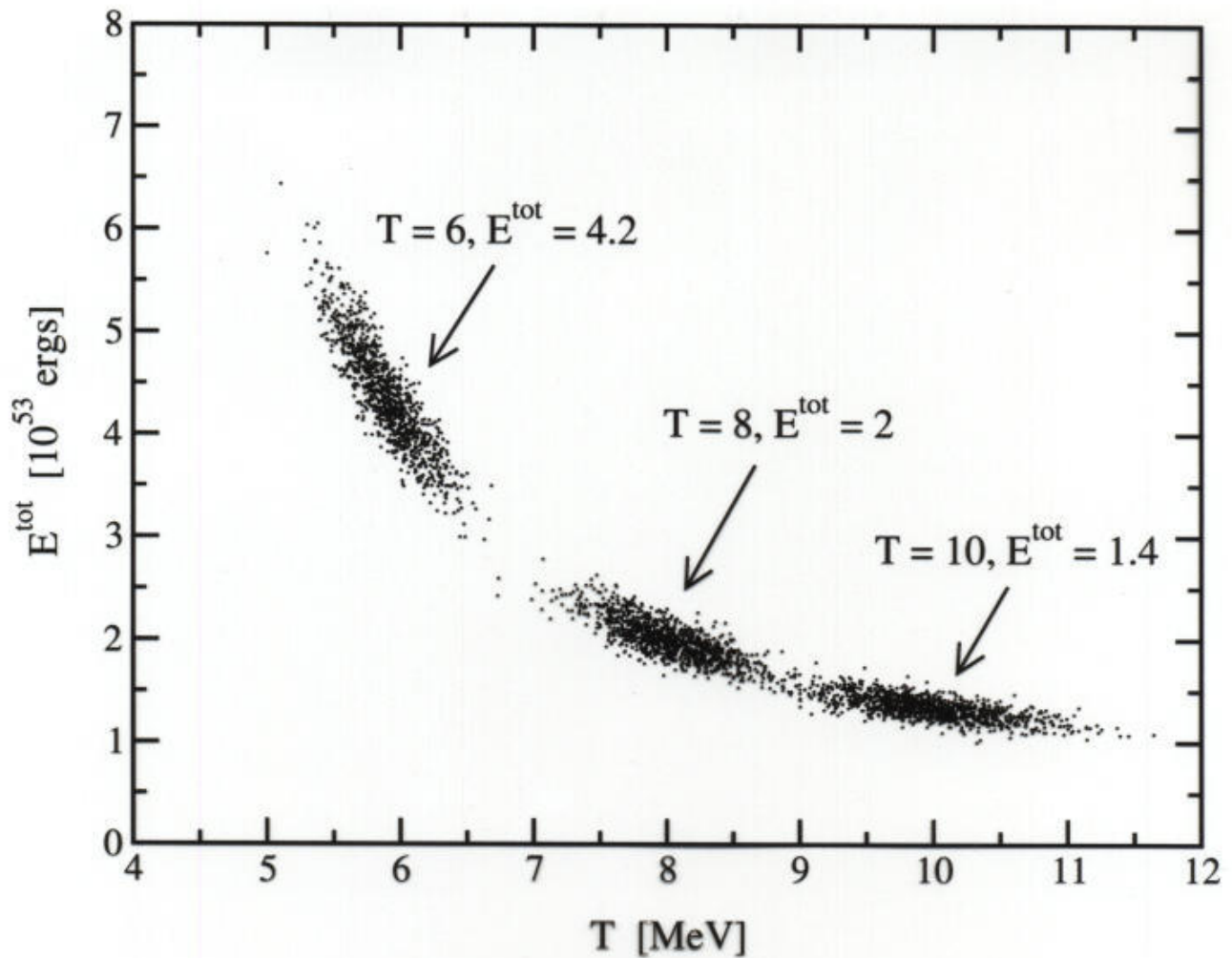


Solid line:  $E^{\text{tot}} = 2 \times 10^{53}$  ergs,  $T = 8$  MeV

upper line:  $E^{\text{tot}} = 1.4 \times 10^{53}$  ergs,  $T = 10$  MeV

lower line:  $E^{\text{tot}} = 4.2 \times 10^{53}$  ergs,  $T = 6$  MeV

All with the same number of events in  $E_{\text{equiv}}$  in  $[0.2 \text{ MeV}, \infty)$ .



If supernova distance is unknown,  
must marginalize over  $E^{\text{tot}}$ .

$$\delta E^{\text{tot}}, \delta T \sim 10\%$$

$10^3$  events used in each Monte Carlo.

## Conclusions

- $E^{\text{tot}}$  and  $T$  for  $\nu_\mu, \bar{\nu}_\mu, \nu_\tau, \bar{\nu}_\tau$  are presently quite uncertain. A lot of present work depends on assumptions.
- $\nu + p \rightarrow \nu + p$  and  $\bar{\nu} + p \rightarrow \bar{\nu} + p$  can be measured in KamLAND and Borexino, making them premier SN neutrino observatories.

$E^{\text{tot}}, T$  can be measured to  $\sim 10\%$ .

- Crucial for

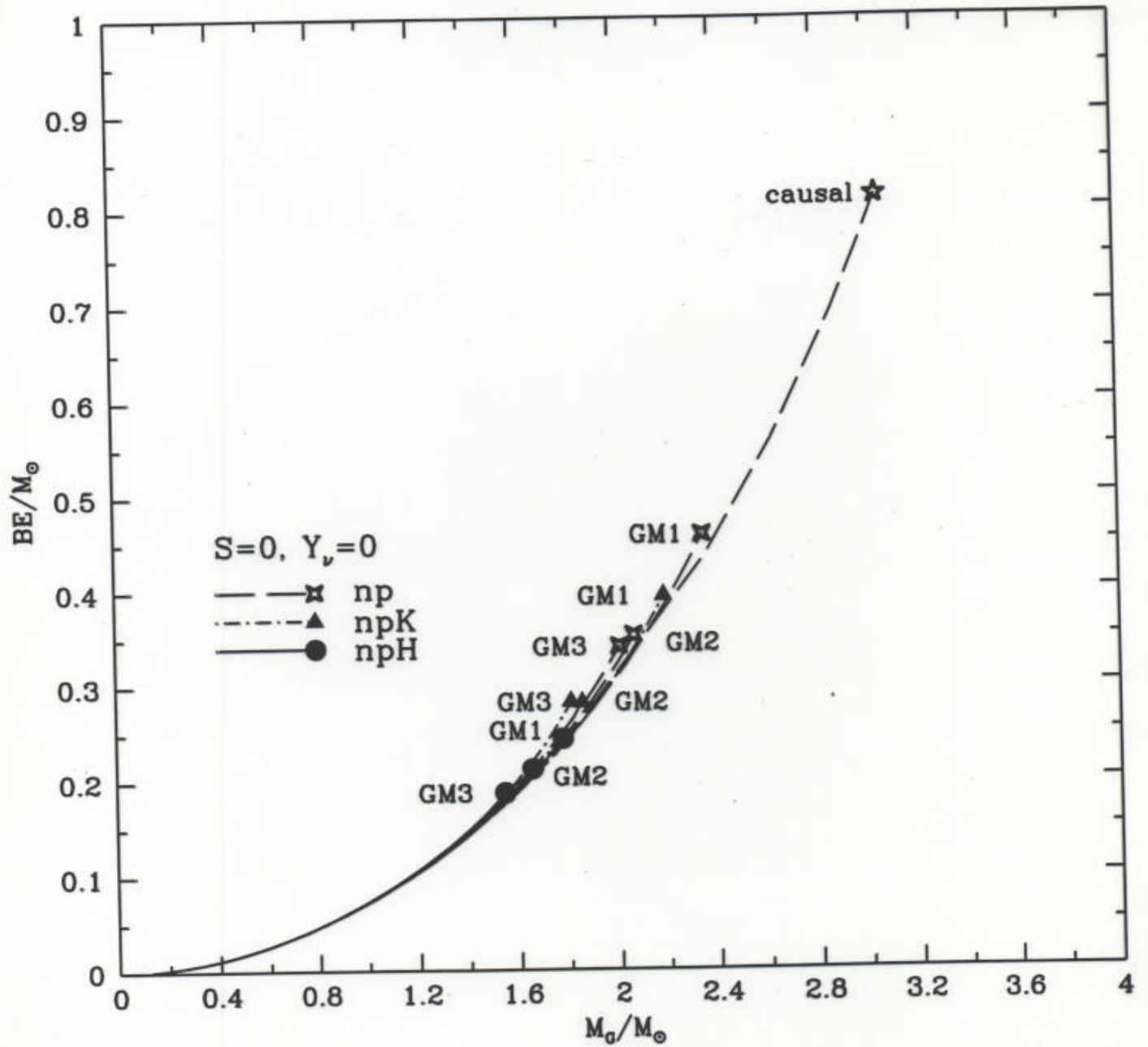
Comparison to SN models

Neutrino oscillations

Total binding energy

Beacom, Farr, and Vogel, hep-ph/0205220





Lattimer & Prakash

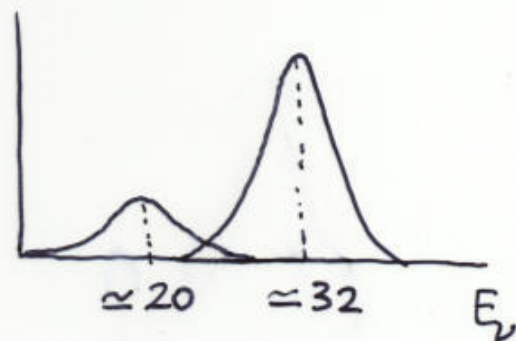
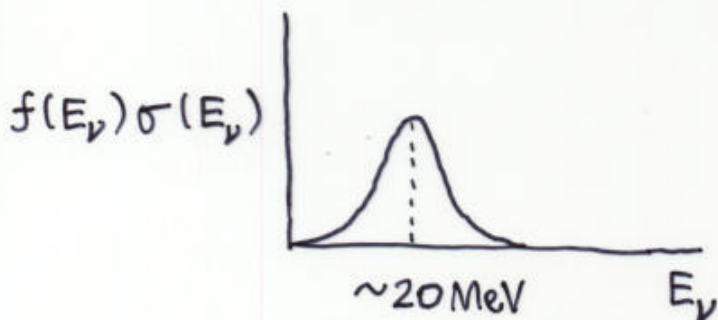
$$BE \approx \frac{3}{5} \frac{GM_{NS}^2}{R_{NS}}$$

## SN Neutrino Oscillations:

- Oscillations to steriles reduce numbers of events
- Oscillations among  $\nu_\mu, \nu_\tau, \bar{\nu}_\mu, \bar{\nu}_\tau$  irrelevant
- Oscillations  $\nu_e \leftrightarrow \nu_\mu, \nu_\tau$  or  $\bar{\nu}_e \leftrightarrow \bar{\nu}_\mu, \bar{\nu}_\tau$ :  
Main effect is to swap in a "hot" spectrum

Example:  $\bar{\nu}_e p \rightarrow e^+ n$

$$\sigma(E_\nu) \sim E_\nu^2$$



peak at  $\approx 4T$   
integral  $\approx \frac{\langle \sigma \rangle}{T} \approx T$

$$T_{\bar{\nu}_e} \approx 5 \text{ MeV}$$

$$T_{\nu_x} \approx 8 \text{ MeV}$$

- Similar for  $\nu + d, {}^{12}\text{C}, {}^{16}\text{O}, {}^{208}\text{Pb}, \dots$

## SN Neutrino Oscillations (cont.)

See recent papers by

Barger, Marfatia, Wood

Minakata, Nunokawa, Tomas, Valle

Akmedhov, Lunderdini, Smirnov

Sato et al.

Balantekin et al.

Fogli et al.

Schirato and Fuller

and many others.....