Neutrinoless Double Beta Decay and Particle Physics

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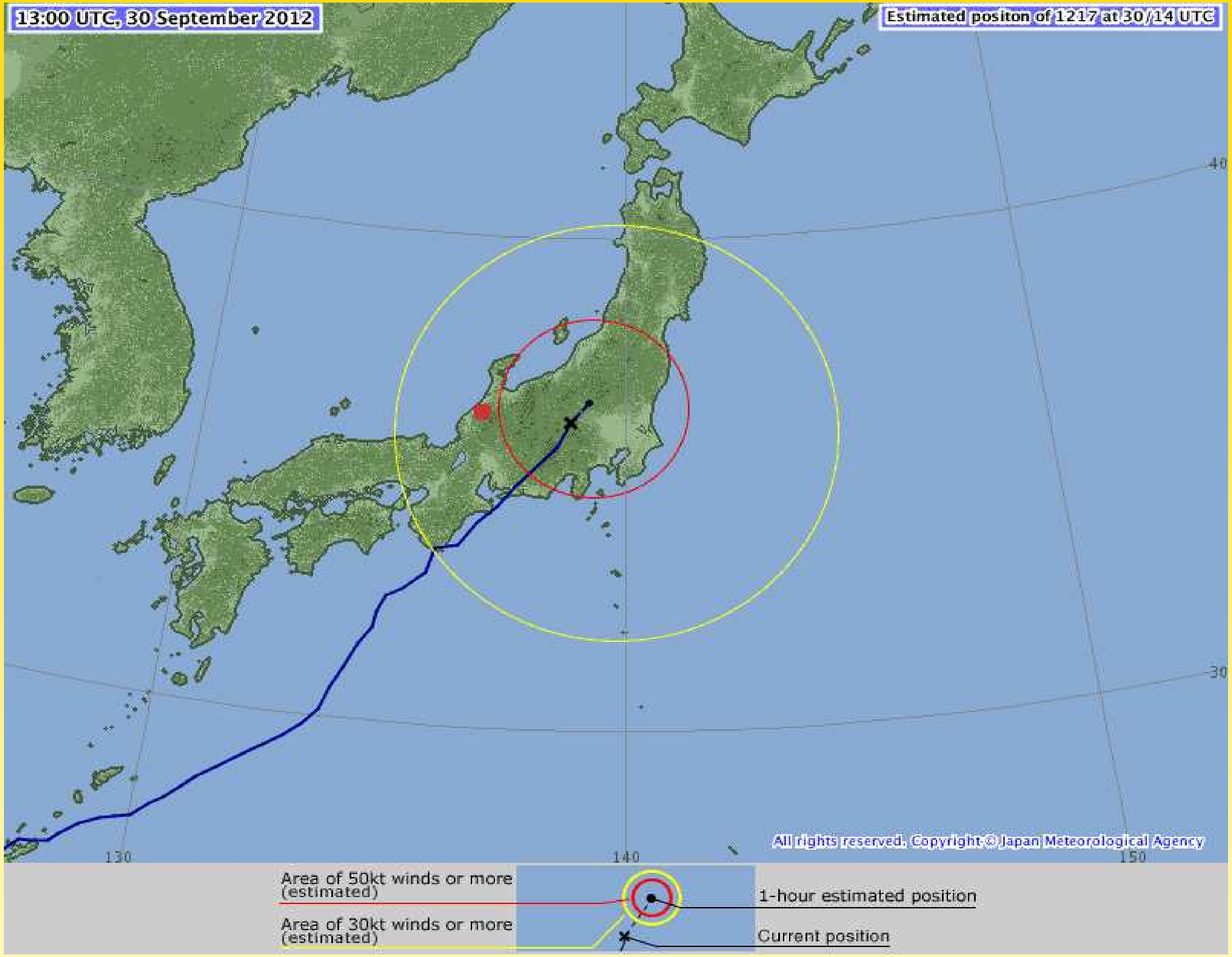
01/10/2012

MANITOP
Massive Neutrinos: Investigating their Theoretical Origin and Phenomenology


13:00 UTC, 30 September 2012
Estimated positon of 1217 at 30/14 UTC
40
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130
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150
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Area of 50kt winds or more (estimated)
Area of 30kt winds or more (estimated)
1-hour estimated position
Current position

Outline

$(A, Z) \rightarrow (A, Z+2) + 2\,e^- \quad (0\nu\beta\beta) \Rightarrow$ Lepton Number Violation

- **Standard Interpretation:**

  Neutrinoless Double Beta Decay is mediated by light and massive Majorana neutrinos (the ones which oscillate) and all other mechanisms potentially leading to $0\nu\beta\beta$ give negligible or no contribution

- **Non-Standard Interpretations:**

  There is at least one other mechanism leading to Neutrinoless Double Beta Decay and its contribution is at least of the same order as the light neutrino exchange mechanism

reviews on $0\nu\beta\beta$:

Int. J. Mod. Phys. **E20**, 1833 (2011); Focus issue J. Phys. **G** [1206.2560]

Interpretation of Experiments

Master formula:

$$\Gamma^{0\nu} = G_x(Q,Z)\,|\mathcal{M}_x(A,Z)\,\eta_x|^2$$

- $G_x(Q,Z)$: phase space factor
- $\mathcal{M}_x(A,Z)$: nuclear physics
- η_x: particle physics

Interpretation of Experiments

Master formula:

$$\Gamma^{0\nu} = G_x(Q,Z)\,|\mathcal{M}_x(A,Z)\,\eta_x|^2$$

- $G_x(Q,Z)$: phase space factor; calculable
- $\mathcal{M}_x(A,Z)$: nuclear physics; problematic
- η_x: particle physics; interesting

Upcoming/running experiments: exciting time!!

best limit was from 2001...

Name	Isotope	source = detector; calorimetric with			source ≠ detector
		high energy res.	low energy res.	event topology	event topology
AMoRE	^{100}Mo	✓	–	–	–
CANDLES	^{48}Ca	–	✓	–	–
COBRA	^{116}Cd (and ^{130}Te)	–	–	✓	–
CUORE	^{130}Te	✓	–	–	–
DCBA	^{82}Se or ^{150}Nd	–	–	–	✓
EXO	^{136}Xe	–	–	✓	–
GERDA	^{76}Ge	✓	–	–	–
KamLAND-Zen	^{136}Xe	–	✓	–	–
LUCIFER	^{82}Se or ^{100}Mo or ^{116}Cd	✓	–	–	–
MAJORANA	^{76}Ge	✓	–	–	–
MOON	^{82}Se or ^{100}Mo or ^{150}Nd	–	–	–	✓
NEXT	^{136}Xe	–	–	✓	–
SNO+	^{150}Nd(?)	–	✓	–	–
SuperNEMO	^{82}Se or ^{150}Nd	–	–	–	✓
XMASS	^{136}Xe	–	✓	–	–

multi-isotope determination good for 3 reasons

3 Reasons for Multi-isotope determination

1.) credibility

2.) test NME calculation

$$\frac{T_{1/2}^{0\nu}(A_1, Z_1)}{T_{1/2}^{0\nu}(A_2, Z_2)} = \frac{G(Q_2, Z_2)\,|\mathcal{M}(A_2, Z_2)|^2}{G(Q_1, Z_1)\,|\mathcal{M}(A_1, Z_1)|^2}$$

systematic errors drop out, ratio sensitive to NME model

3.) test mechanism

$$\frac{T_{1/2}^{0\nu}(A_1, Z_1)}{T_{1/2}^{0\nu}(A_2, Z_2)} = \frac{G_x(Q_2, Z_2)\,|\mathcal{M}_x(A_2, Z_2)|^2}{G_x(Q_1, Z_1)\,|\mathcal{M}_x(A_1, Z_1)|^2}$$

particle physics drops out, ratio of NMEs sensitive to mechanism

Experimental Aspects

particle physics:

$$(T^{0\nu}_{1/2})^{-1} \propto \text{ (particle physics)}^2$$

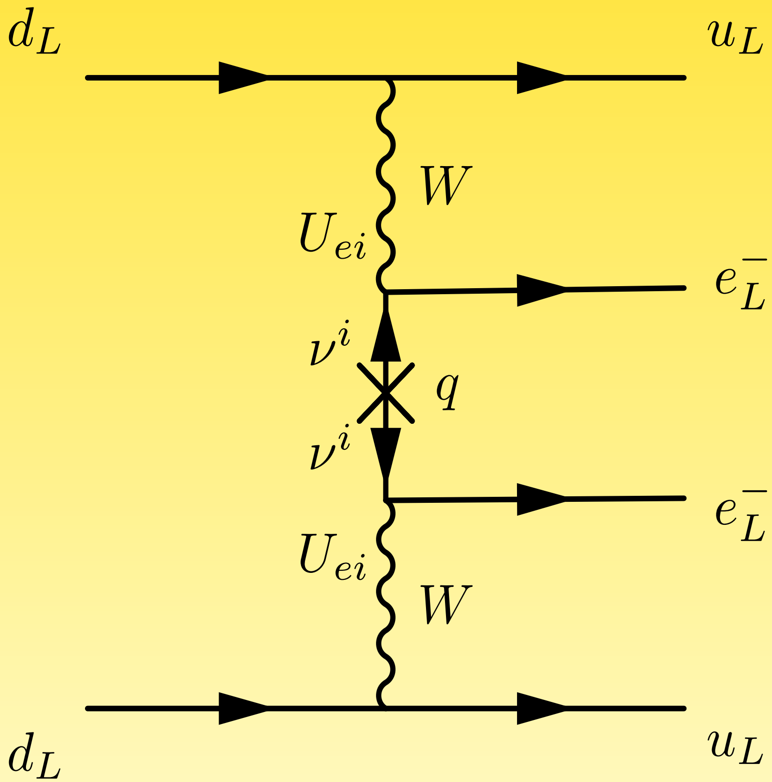
experimentally:

$$(T^{0\nu}_{1/2})^{-1} \propto \begin{cases} a\, M\, \varepsilon\, t & \text{without background} \\ a\, \varepsilon \sqrt{\dfrac{M\, t}{B\, \Delta E}} & \text{background-dominated} \end{cases}$$

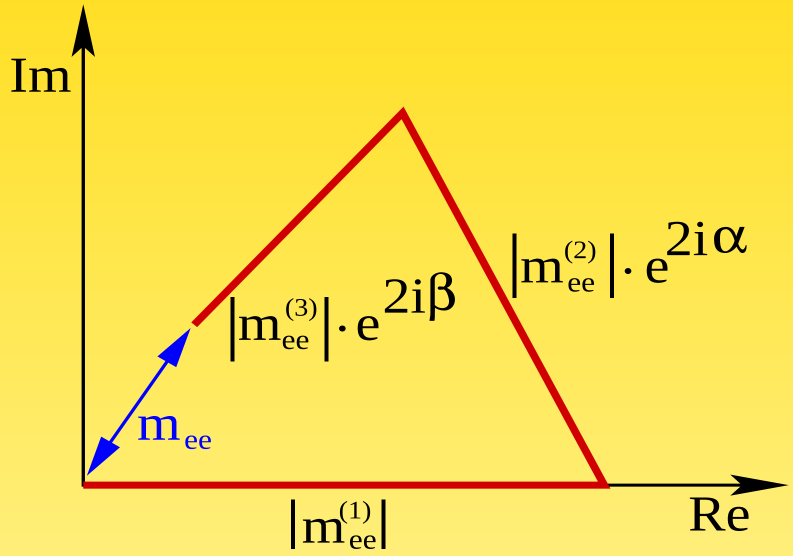
Note: factor 2 in particle physics is combined factor of 16 in $M \times t \times B \times \Delta E$

Standard Interpretation

Neutrinoless Double Beta Decay is mediated by light and massive Majorana neutrinos (the ones which oscillate) and all other mechanisms potentially leading to $0\nu\beta\beta$ give negligible or no contribution



$\Delta L \neq 0$: Neutrinoless Double Beta Decay

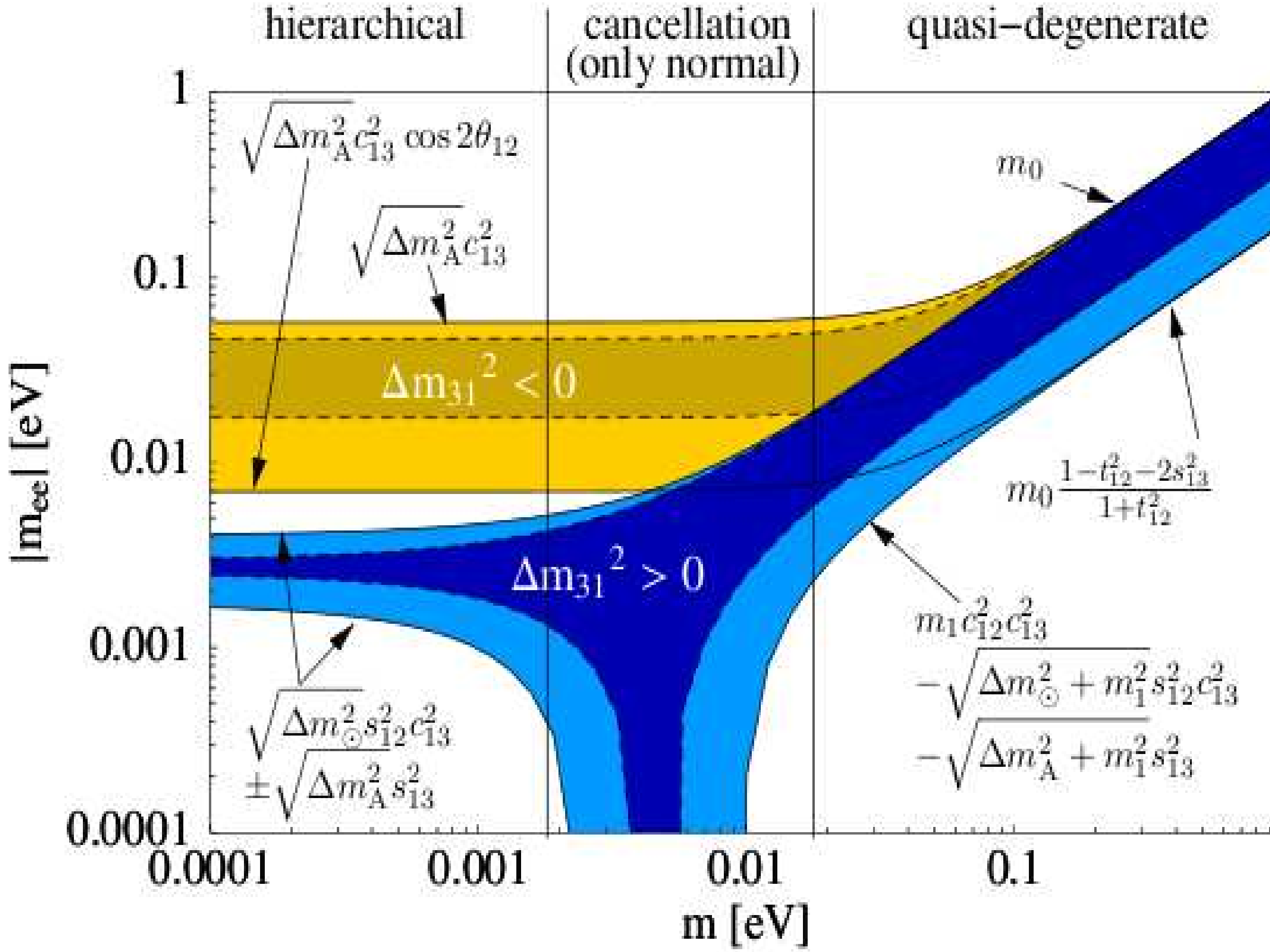


Amplitude proportional to coherent sum ("effective mass"):

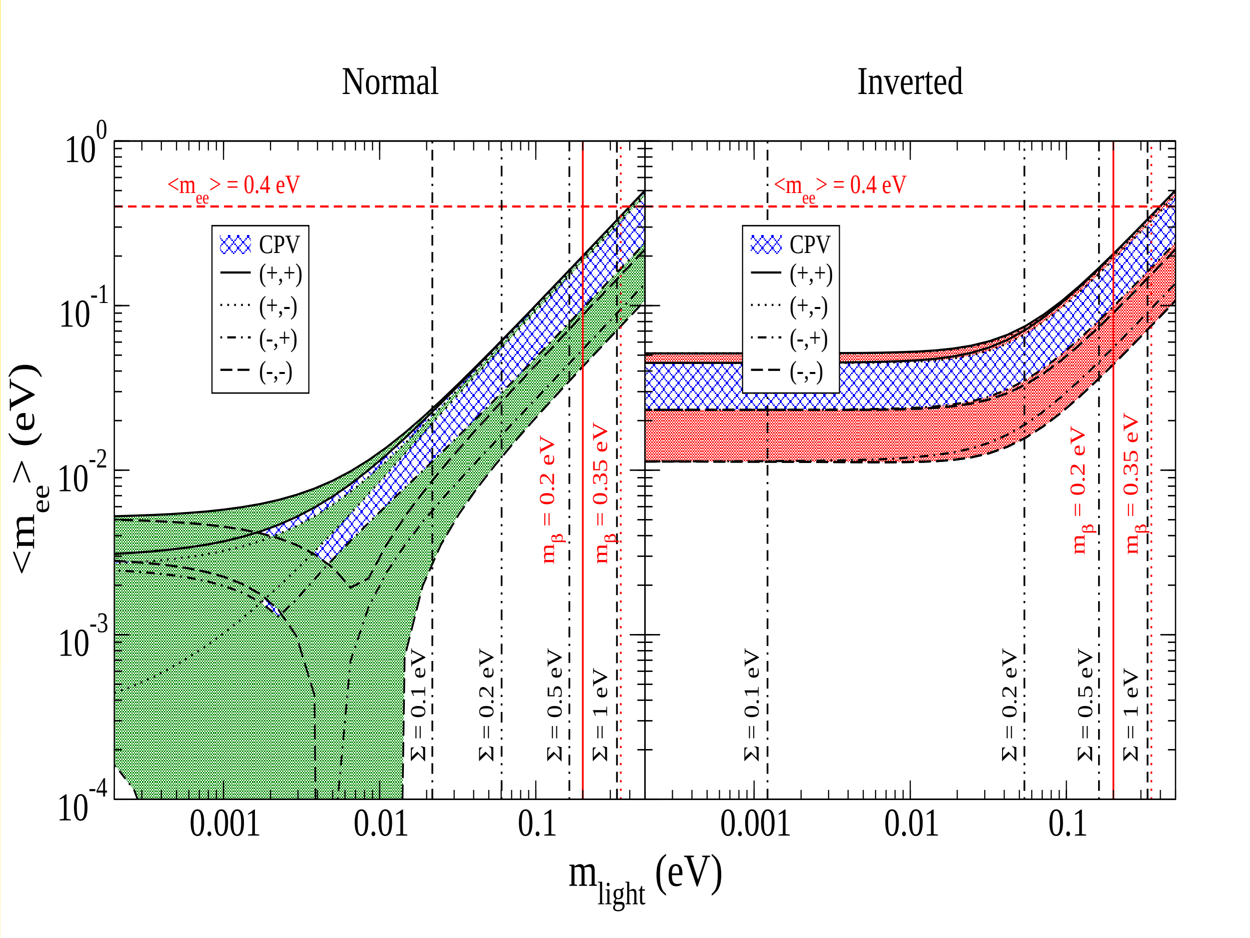
$$|m_{ee}| \equiv \left|\sum U_{ei}^2\, m_i\right| = \left||U_{e1}|^2\, m_1 + |U_{e2}|^2\, m_2\, e^{2i\alpha} + |U_{e3}|^2\, m_3\, e^{2i\beta}\right|$$

7 out of 9 parameters of m_ν!

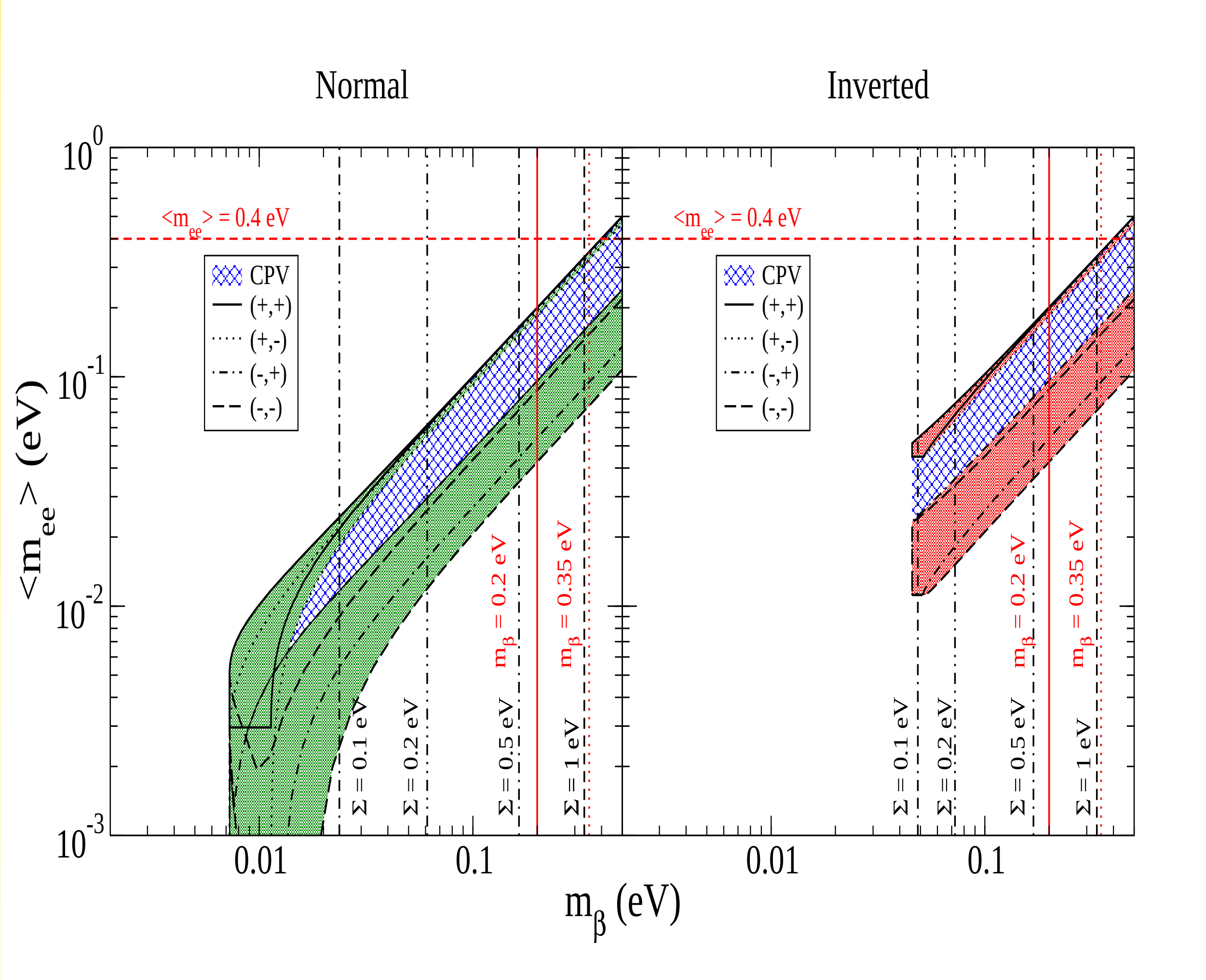
The usual plot



The usual plot

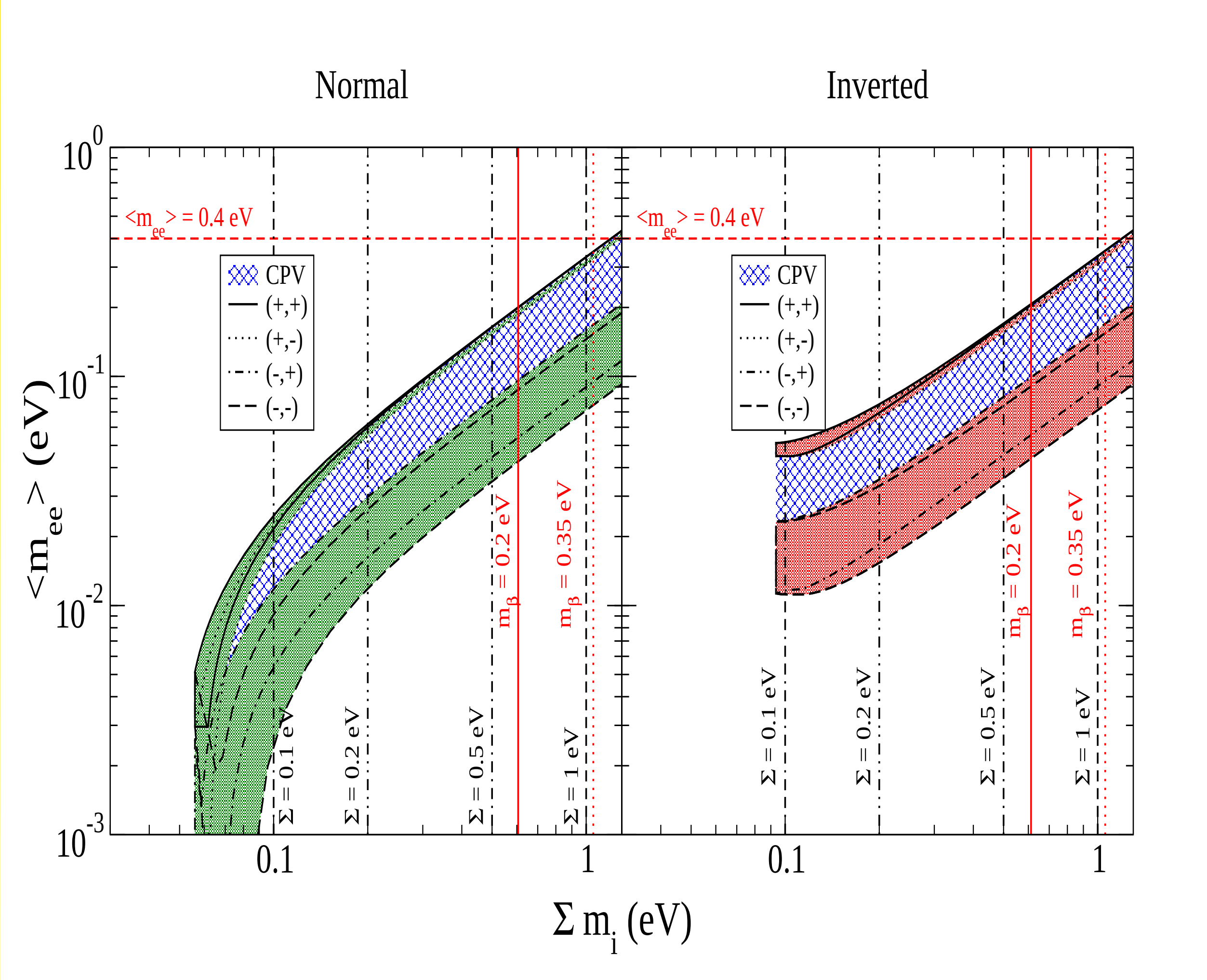


Plot against other observables



Complementarity of $|m_{ee}| = U_{ei}^2\, m_i$ and $m_\beta = \sqrt{|U_{ei}|^2\, m_i^2}$ and $\Sigma = \sum m_i$

Plot against other observables

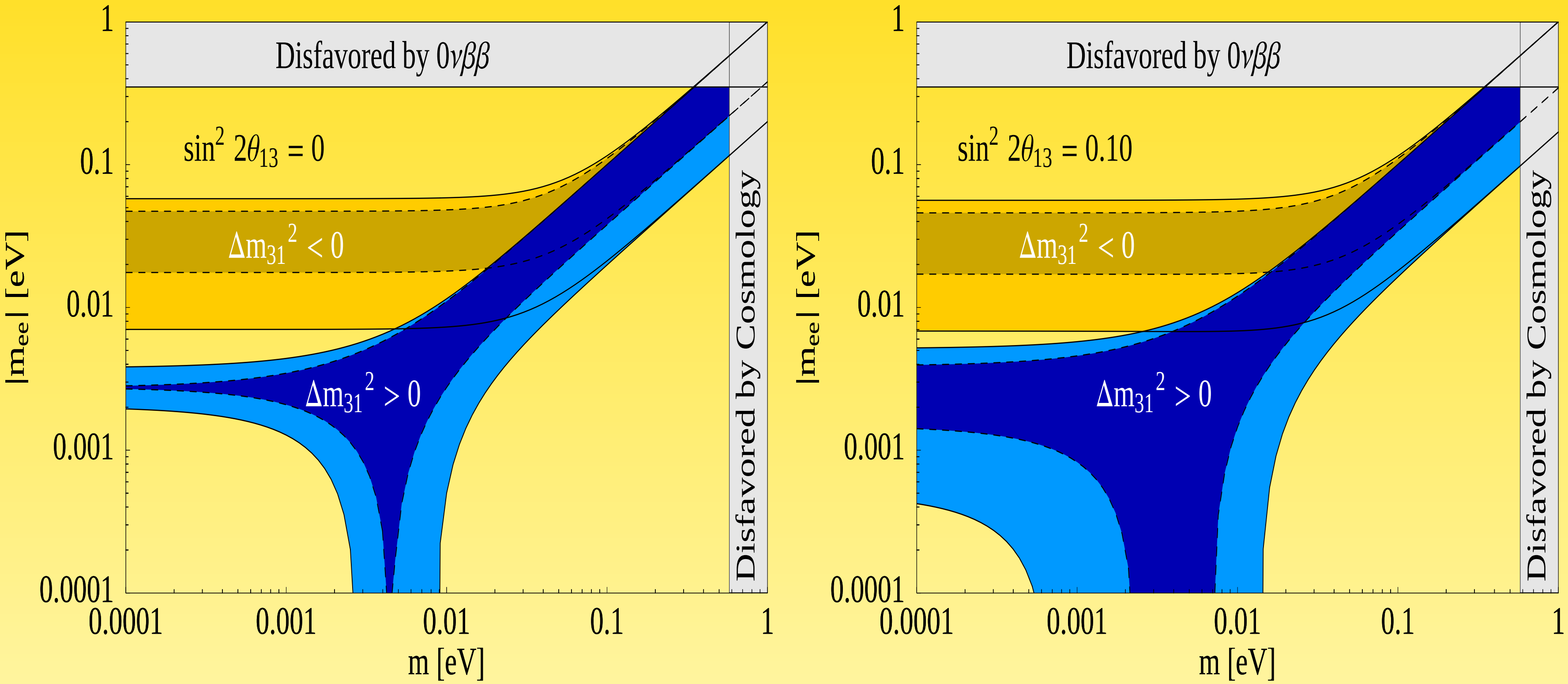


Complementarity of $|m_{ee}| = U_{ei}^2\, m_i$ and $m_\beta = \sqrt{|U_{ei}|^2\, m_i^2}$ and $\Sigma = \sum m_i$

Neutrino Mass Matrix

		KATRIN		$0\nu\beta\beta$		cosmology	
		yes	no	yes	no	yes	no
KATRIN	yes	–	–	QD + Majorana	QD + Dirac	QD	N-SC
	no	–	–	N-SI	low IH or NH or Dirac	$m_\nu \lesssim 0.1$ eV or N-SC	NH
$0\nu\beta\beta$	yes	•	•	–	–	(IH or QD) + Majorana	N-SC or N-SI
	no	•	•	–	–	low IH or (QD + Dirac)	NH
cosmology	yes	•	•	•	•	–	–
	no	•	•	•	•	–	–

$0\nu\beta\beta$ and U_{e3}



Lindner, Merle, W.R., hep-ph/0512143

Vanishing effective mass...

conspiracy of 7 parameters?

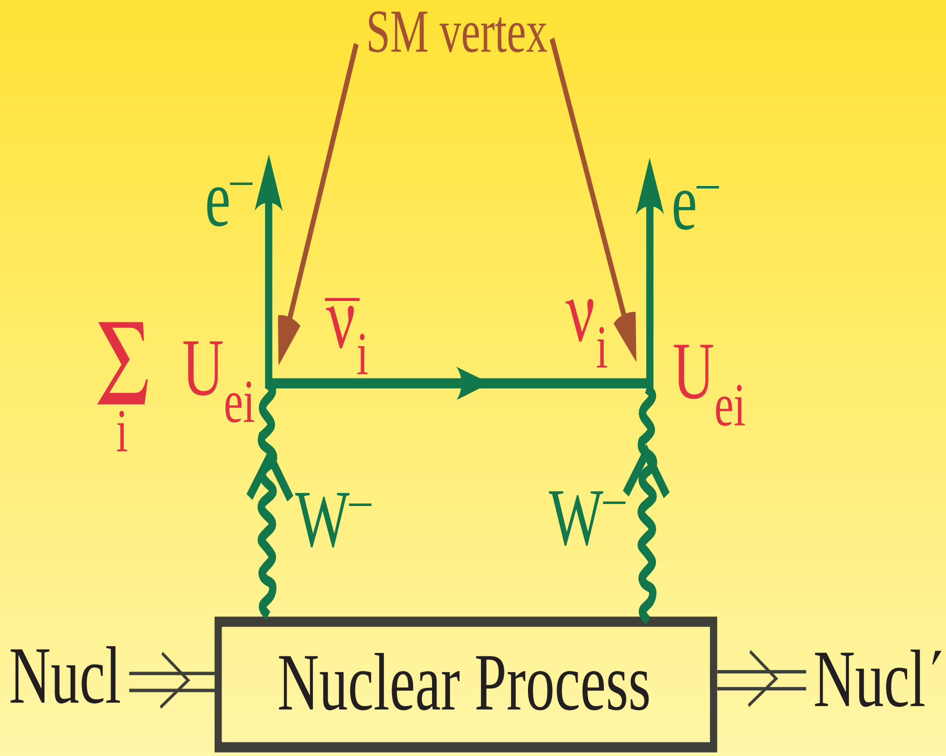
	$(\overline{L_1}, \overline{L_2})$	$\overline{L_3}$	(ν_{R_1}, ν_{R_2})	e_R	μ_R	τ_R	(ϕ_1, ϕ_2)	χ
S_3	2^*	1_S	2	1_S	1_S	1_S	2	1_S
Z_3	ω	ω	ω	ω	1	ω^2	ω	ω

leads to $|m_{ee}| = 0$ and

$$\tan^2\theta_{13} = \sqrt{\frac{\Delta m^2_\odot}{\Delta m^2_{\rm A}}}\,\sin^2\theta_{12}$$

W.R., Tanimoto, Watanabe, 1201.4936

From life-time to particle physics: Nuclear Matrix Elements



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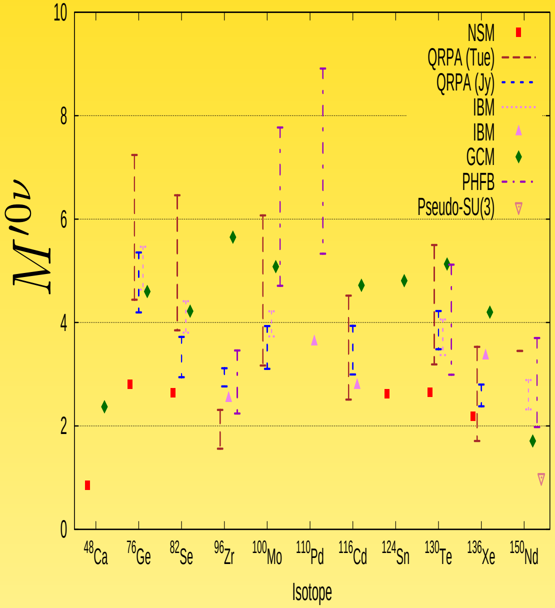
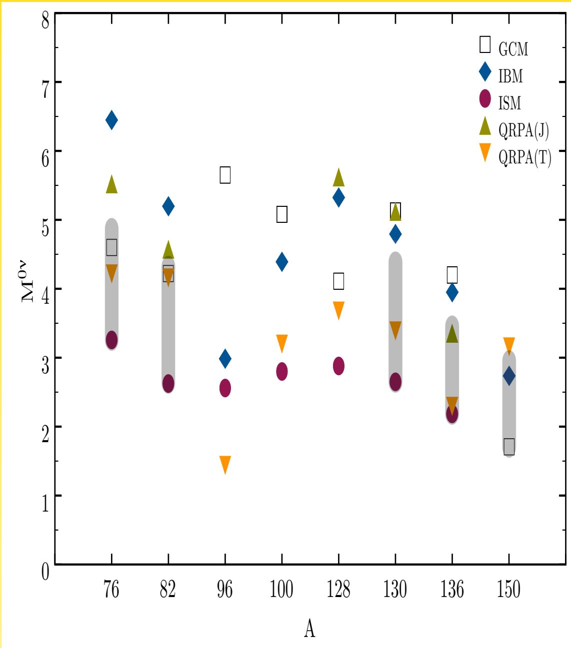
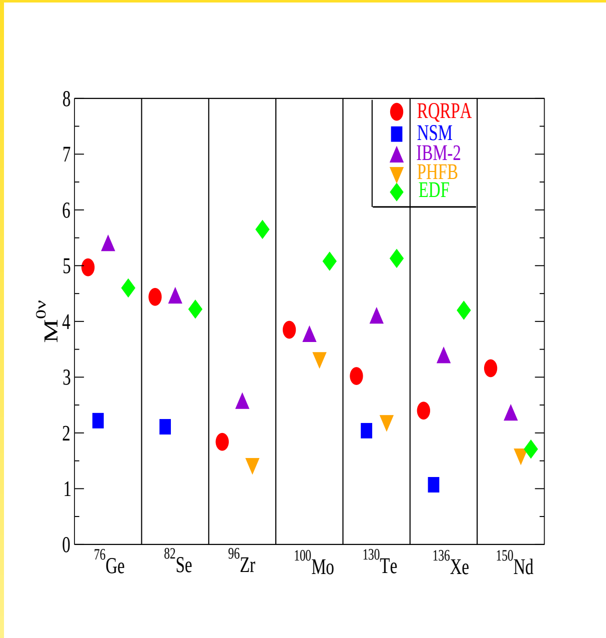






Dueck, W.R., Zuber, PRD **83** Gomez-Cadenas *et al.*, 1109.5515 Vogel, 1208.1992

(current) uncertainty of factor 2 to 3, directly translates into uncertainty on particle physics parameter

Isotope	$T^{0\nu}_{1/2}$ [yrs]	Experiment	$\lvert m_{ee}\rvert^{\rm lim}_{\rm min}$ [eV]	$\lvert m_{ee}\rvert^{\rm lim}_{\rm max}$ [eV]	
^{48}Ca	5.8×10^{22}	CANDLES	3.55	9.91	$\times 0.98$
^{76}Ge	1.9×10^{25}	HDM	0.21	0.53	$\times 1.04$
	1.6×10^{25}	IGEX	0.25	0.63	$\times 1.04$
^{82}Se	3.2×10^{23}	NEMO-3	0.85	2.08	$\times 1.04$
^{96}Zr	9.2×10^{21}	NEMO-3	3.97	14.39	$\times 1.06$
^{100}Mo	1.0×10^{24}	NEMO-3	0.31	0.79	$\times 1.06$
^{116}Cd	1.7×10^{23}	SOLOTVINO	1.22	2.30	$\times 1.06$
^{130}Te	2.8×10^{24}	CUORICINO	0.27	0.57	$\times 1.09$
^{136}Xe	1.6×10^{25}	EXO-200	0.15	0.36	$\times 1.10$
^{150}Nd	1.8×10^{22}	NEMO-3	2.35	5.08	$\times 1.12$

(recent reevaluation of phase space factors by Iachello+Kotila)

HDM limit reached/improved by EXO-200 !

EXO-200 vs. Klapdor

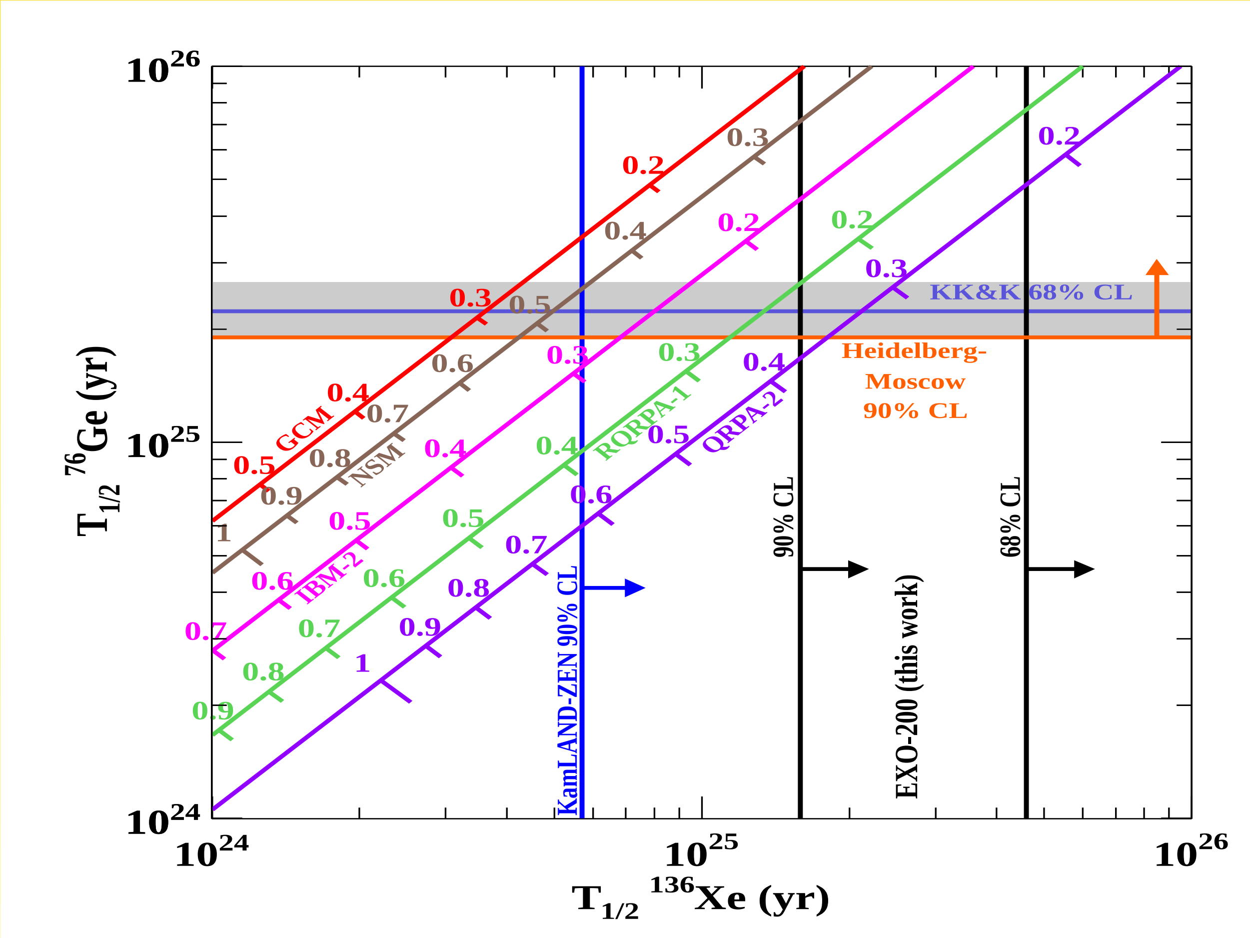
$$T_{\rm Ge}^{-1} = G_{\rm Ge}\,|\mathcal{M}_{\rm Ge}|^2\,|m_{ee}|^2 = \left(2 \times 10^{25}\,{\rm yrs}\right)^{-1}$$
$$T_{\rm Xe}^{-1} = G_{\rm Xe}\,|\mathcal{M}_{\rm Xe}|^2\,|m_{ee}|^2 \le \left(1.6 \times 10^{25}\,{\rm yrs}\right)^{-1}$$

Ge-claim is ruled out when

$$T_{\rm Xe} \ge 2.9 \times 10^{24} \left|\frac{\mathcal{M}_{\rm Ge}}{\mathcal{M}_{\rm Xe}}\right|^2 \;{\rm yrs}$$

With compilation from `Vogel, 1208.1992`:

$$\left|\frac{\mathcal{M}_{\rm Ge}}{\mathcal{M}_{\rm Xe}}\right|^2 \simeq \begin{cases} 4.0 & ({\rm RQRPA}) & \Rightarrow T_{\rm Xe} \ge 1.2 \times 10^{25}\,{\rm yrs} \\ 4.2 & ({\rm NSM}) & \Rightarrow T_{\rm Xe} \ge 1.2 \times 10^{25}\,{\rm yrs} \\ 2.5 & (\text{IBM-2}) & \Rightarrow T_{\rm Xe} \ge 7.3 \times 10^{24}\,{\rm yrs} \\ 1.2 & ({\rm EDF}) & \Rightarrow T_{\rm Xe} \ge 3.5 \times 10^{24}\,{\rm yrs} \end{cases}$$



EXO, 1205.5608

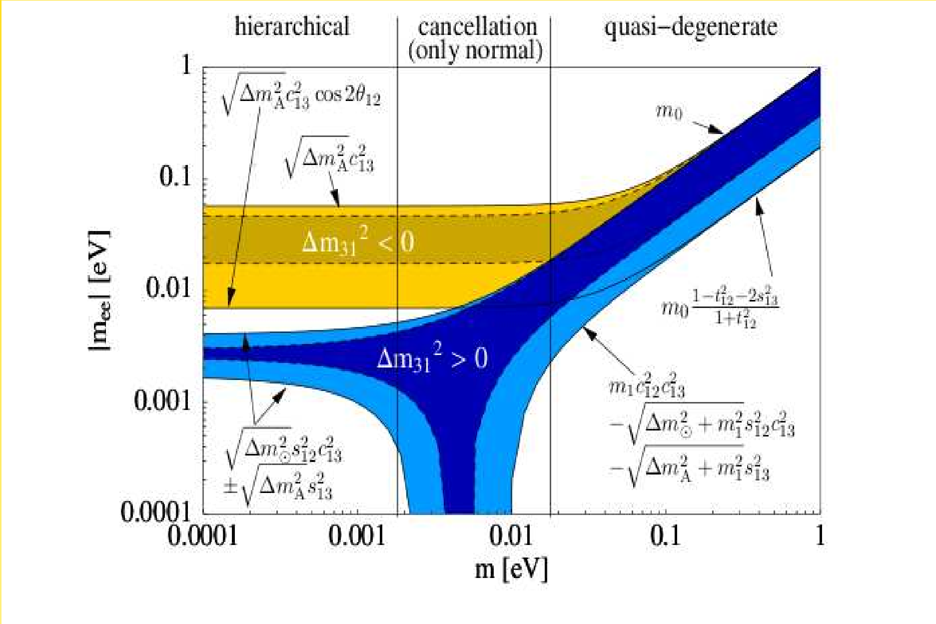
Experiment	Isotope	Mass of Isotope [kg]	Sensitivity $T_{1/2}^{0\nu}$ [yrs]	Status	Start of data-taking	Sensitivity $\langle m_\nu \rangle$ [eV]
GERDA	^{76}Ge	18	3×10^{25}	running	$\sim$ 2011	0.17-0.42
		40	2×10^{26}	in progress	$\sim$ 2012	0.06-0.16
		1000	6×10^{27}	R&D	$\sim$ 2015	0.012-0.030
CUORE	^{130}Te	200	$6.5 \times 10^{26*}$	in progress	$\sim$ 2013	0.018-0.037
			$2.1 \times 10^{26**}$			0.03-0.066
MAJORANA	^{76}Ge	30-60	$(1-2) \times 10^{26}$	in progress	$\sim$ 2013	0.06-0.16
		1000	6×10^{27}	R&D	$\sim$ 2015	0.012-0.030
EXO	^{136}Xe	200	6.4×10^{25}	running	$\sim$ 2011	0.073-0.18
		1000	8×10^{26}	R&D	$\sim$ 2015	0.02-0.05
SuperNEMO	^{82}Se	100-200	$(1-2) \times 10^{26}$	R&D	$\sim$ 2013-15	0.04-0.096
KamLAND-Zen	^{136}Xe	400	4×10^{26}	running	$\sim$ 2011	0.03-0.07
		1000	10^{27}	R&D	$\sim$ 2013-15	0.02-0.046
SNO+	^{150}Nd	132	1.8×10^{25}	in progress	$\sim$ 2014	0.09-0.18

(with *same* lifetime: ^{150}Nd and ^{100}Mo do best...)

With $0\nu\beta\beta$ one can

- test Majorana nature of neutrinos
- probe neutrino mass scale
- test inverted ordering
- extract Majorana phase
- test flavor symmetry models: neutrino mass “sum-rules”

Inverted Ordering



Nature provides 2 scales:

$$\langle m_\nu \rangle^{\text{IH}}_{\text{max}} \simeq c_{13}^2 \sqrt{\Delta m_A^2} \quad \text{and} \quad \langle m_\nu \rangle^{\text{IH}}_{\text{min}} \simeq c_{13}^2 \sqrt{\Delta m_A^2} \cos 2\theta_{12}$$

requires $\mathcal{O}(10^{26} \ldots 10^{27})$ yrs

Ruling out Inverted Hierarchy

$$|m_{ee}|^{\rm IH}_{\rm min} = (1 - |U_{e3}|^2)\sqrt{|\Delta m^2_{\rm A}|}\,(1 - 2\sin^2\theta_{12}) = \begin{cases} (0.016\ldots0.020)~{\rm eV} & 1\sigma \\ (0.013\ldots0.024)~{\rm eV} & 3\sigma \end{cases}$$

- small $|U_{e3}|$
- large $|\Delta m^2_{\rm A}|$
- small $\sin^2\theta_{12}$

Current 3σ range of $\sin^2\theta_{12}$ gives factor of 2 uncertainty for $|m_{ee}|^{\rm IH}_{\rm min}$

$\Rightarrow$ combined factor of 16 in $M \times t \times B \times \Delta E$

$\Rightarrow$ need precision determination of θ_{12}

Dueck, W.R., Zuber, PRD **83**

Sterile Neutrinos and $0\nu\beta\beta$

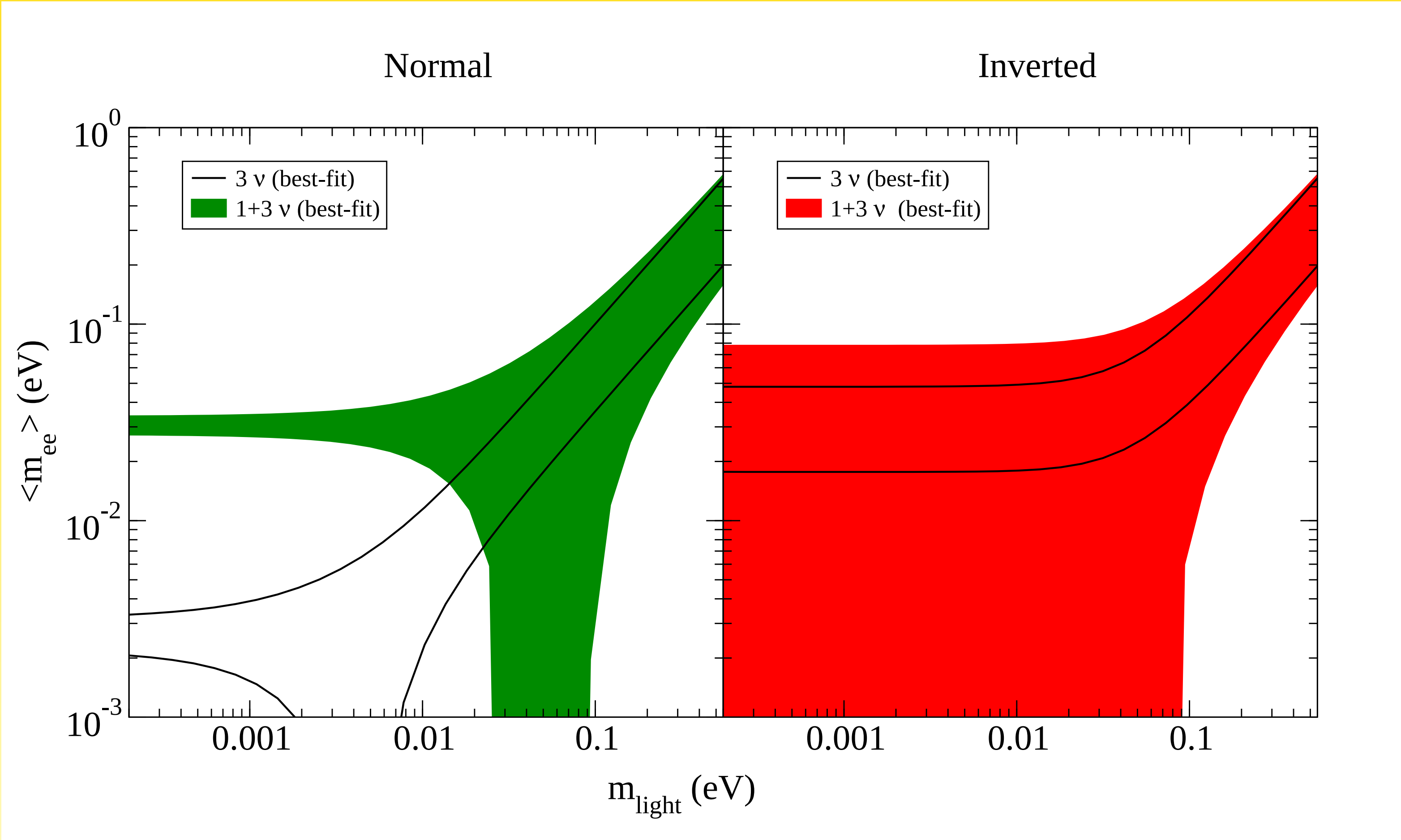
- recall: $|m_{ee}|^{\rm act}_{\rm NH}$ can vanish and $|m_{ee}|^{\rm act}_{\rm IH} \sim 0.03$ eV cannot vanish
- $|m_{ee}| = \Big| \underbrace{|U_{e1}|^2\, m_1 + |U_{e2}|^2\, m_2\, e^{2i\alpha} + |U_{e3}^2|\, m_3\, e^{2i\beta}}_{m_{ee}^{\rm act}} + \underbrace{|U_{e4}|^2\, m_4\, e^{2i\Phi_1}}_{m_{ee}^{\rm st}} \Big|$
- $\Delta m^2_{\rm st} \simeq 1.8$ eV^2 and $|U_{e4}| \simeq 0.13$
- sterile contribution to $0\nu\beta\beta$ (assuming 1+3):

$$|m_{ee}|^{\rm st} \simeq \sqrt{\Delta m^2_{\rm st}}\, |U_{e4}|^2 \simeq 0.03 \text{ eV} \begin{cases} \gg |m_{ee}|^{\rm act}_{\rm NH} \\ \simeq |m_{ee}|^{\rm act}_{\rm IH} \end{cases}$$

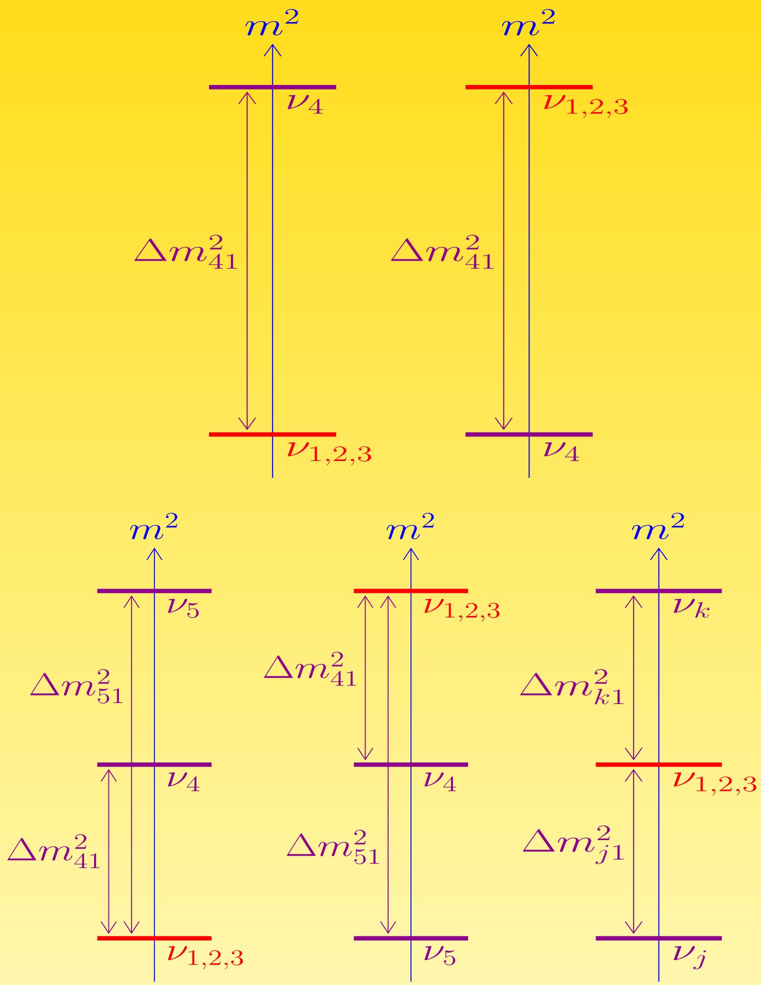
- $\Rightarrow |m_{ee}|_{\rm NH}$ cannot vanish and $|m_{ee}|_{\rm IH}$ can vanish!

usual phenomenology gets completely turned around!

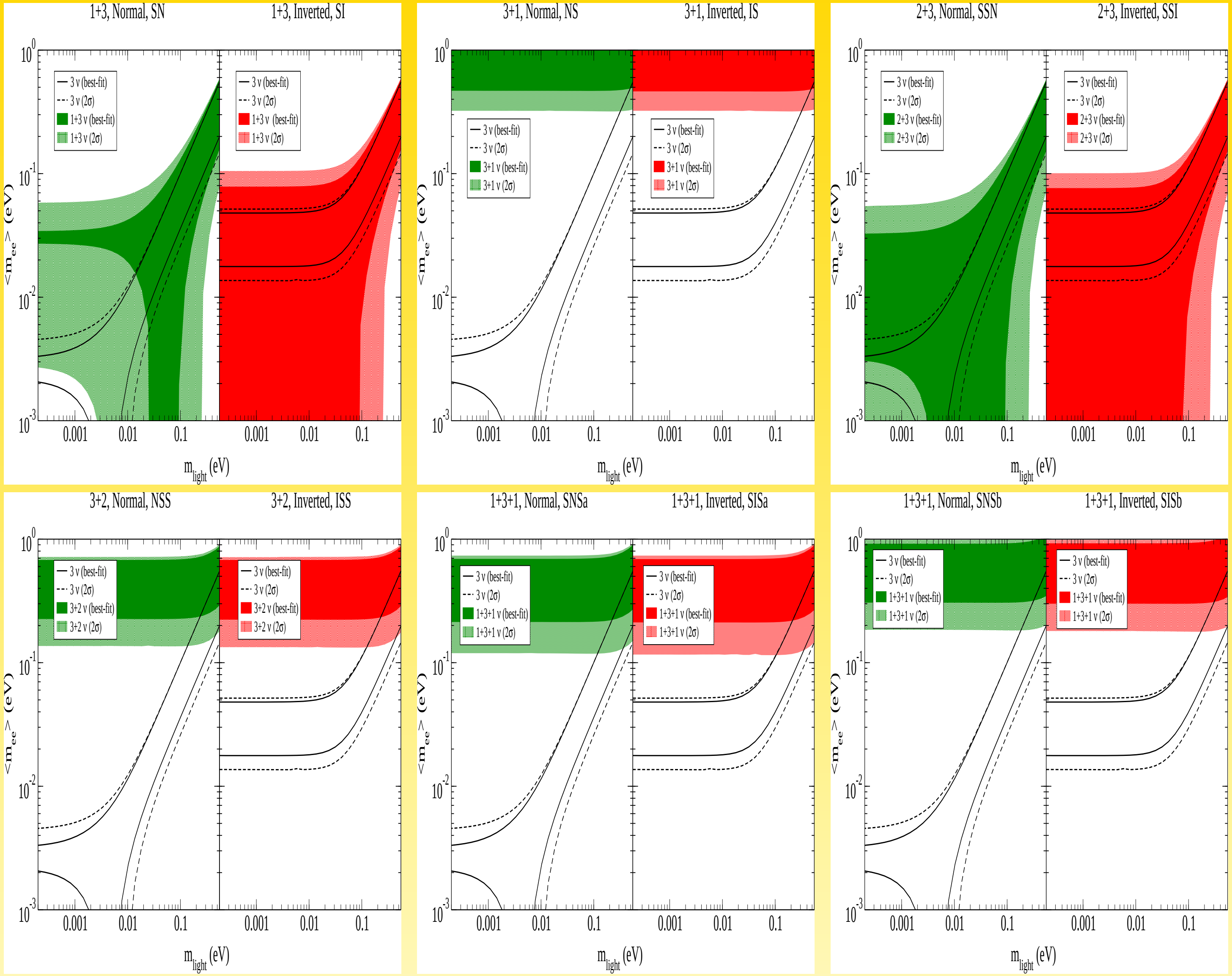
Usual plot gets completely turned around!



Mass Orderings



3 active neutrinos can be normally or inversely ordered



Barry, W.R., Zhang, JHEP **1107**

Sterile Neutrinos, Seesaw and $0\nu\beta\beta$

- if the eV-steriles are from seesaw: individual cancellations in flavor symmetry models, e.g.:

$$U_{e2}^2 \, m_2 + U_{e4}^2 \, m_4 = 0$$

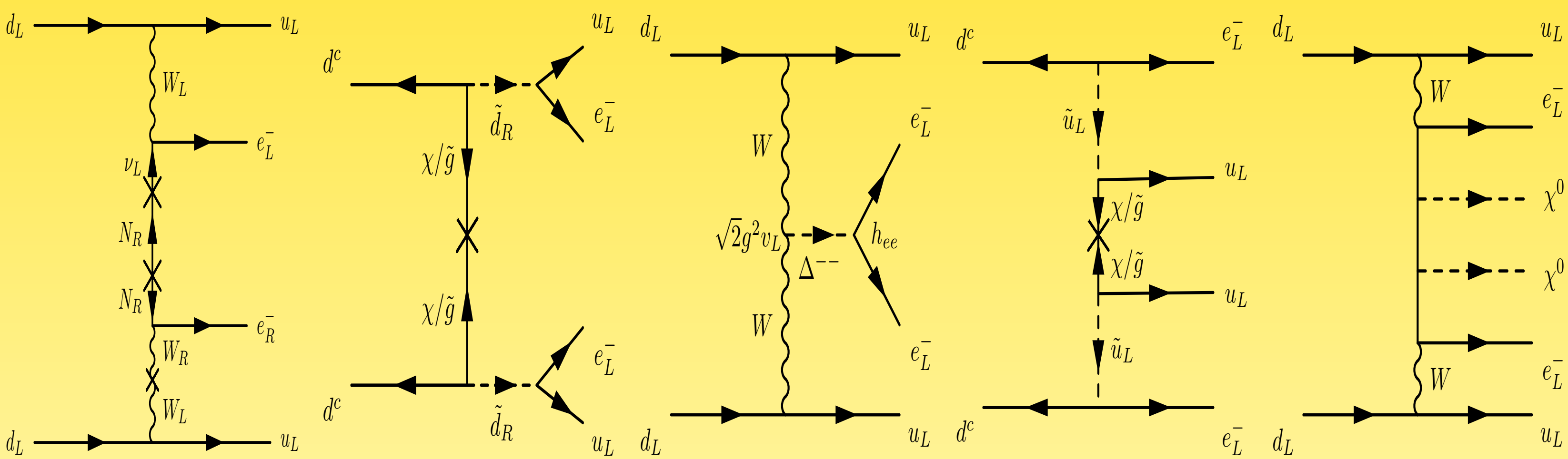
- if seesaw scale is below 100 MeV: No double beta decay!

$$\sum_{i=1}^{6} U_{ei}^2 \, m_i = 0 \text{ since } \mathcal{M} = \left(\begin{array}{cc} 0 & m_D \\ m_D^T & M_R \end{array} \right) = U \left(\begin{array}{cc} m_\nu^{\rm diag} & 0 \\ 0 & M_R^{\rm diag} \end{array} \right) U^T$$

Barry, W.R., Zhang, JCAP **1201**

Non-Standard Interpretations:

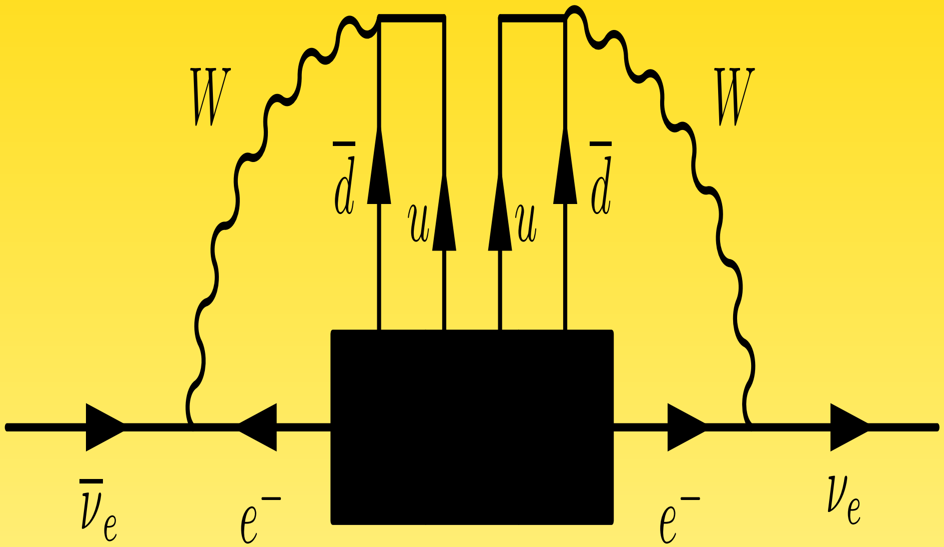
There is at least one other mechanism leading to Neutrinoless Double Beta Decay and its contribution is at least of the same order as the light neutrino exchange mechanism



Clear experimental signature:

KATRIN and/or cosmology see nothing but $0\nu\beta\beta$ does

Schechter-Valle theorem: no matter what process, neutrinos are Majorana:



is 4 loop diagram: $m_\nu \sim \frac{1}{(16\pi^2)^4} \frac{\mathrm{MeV}^5}{m_W^4} \lesssim 10^{-23}\,\mathrm{eV}$

explicit calculation: Duerr, Lindner, Merle, 1105.0901

note: often there are 1-loop diagrams leading to m_ν: direct vs. indirect contribution (Choubey, Duerr, Mitra, W.R., 1201.3031)

mechanism	physics parameter	current limit	test
light neutrino exchange	$\left\lvert U_{ei}^2\, m_i \right\rvert$	0.4 eV	oscillations, cosmology, neutrino mass
heavy neutrino exchange	$\left\lvert \frac{S_{ei}^2}{M_i} \right\rvert$	2×10^{-8} GeV^{-1}	LFV, collider
heavy neutrino and RHC	$\left\lvert \frac{V_{ei}^2}{M_i\, M_{W_R}^4} \right\rvert$	4×10^{-16} GeV^{-5}	flavor, collider
Higgs triplet and RHC	$\left\lvert \frac{(M_R)_{ee}}{m_{\Delta_R}^2 \quad M_{W_R}^4} \right\rvert$	10^{-15} GeV^{-5}	flavor, collider e^- distribution
λ-mechanism with RHC	$\left\lvert \frac{U_{ei}\, \tilde{S}_{ei}}{M_{W_R}^2} \right\rvert$	1.4×10^{-10} GeV^{-2}	flavor, collider, e^- distribution
η-mechanism with RHC	$\tan\zeta \left\lvert U_{ei}\, \tilde{S}_{ei} \right\rvert$	6×10^{-9}	flavor, collider, e^- distribution
short-range $\not{R}$	$\frac{\left\lvert \lambda'^2_{111} \right\rvert}{\Lambda_{\rm SUSY}^5}$ $\Lambda_{\rm SUSY} = f(m_{\tilde{g}}, m_{\tilde{u}_L}, m_{\tilde{d}_R}, m_{\chi_i})$	7×10^{-18} GeV^{-5}	collider, flavor
long-range $\not{R}$	$\left\lvert \sin 2\theta^b\, \lambda'_{131}\, \lambda'_{113} \left(\frac{1}{m_{\tilde{b}_1}^2} - \frac{1}{m_{\tilde{b}_2}^2} \right) \right\rvert$ $\sim \frac{G_F}{q} m_b \frac{\left\lvert \lambda'_{131}\, \lambda'_{113} \right\rvert}{\Lambda_{\rm SUSY}^3}$	2×10^{-13} GeV^{-2} 1×10^{-14} GeV^{-3}	flavor, collider
Majorons	$\lvert\langle g_\chi \rangle\rvert$ or $\lvert\langle g_\chi \rangle\rvert^2$	$10^{-4} \ldots 1$	spectrum, cosmology

Distinguishing Mechanisms

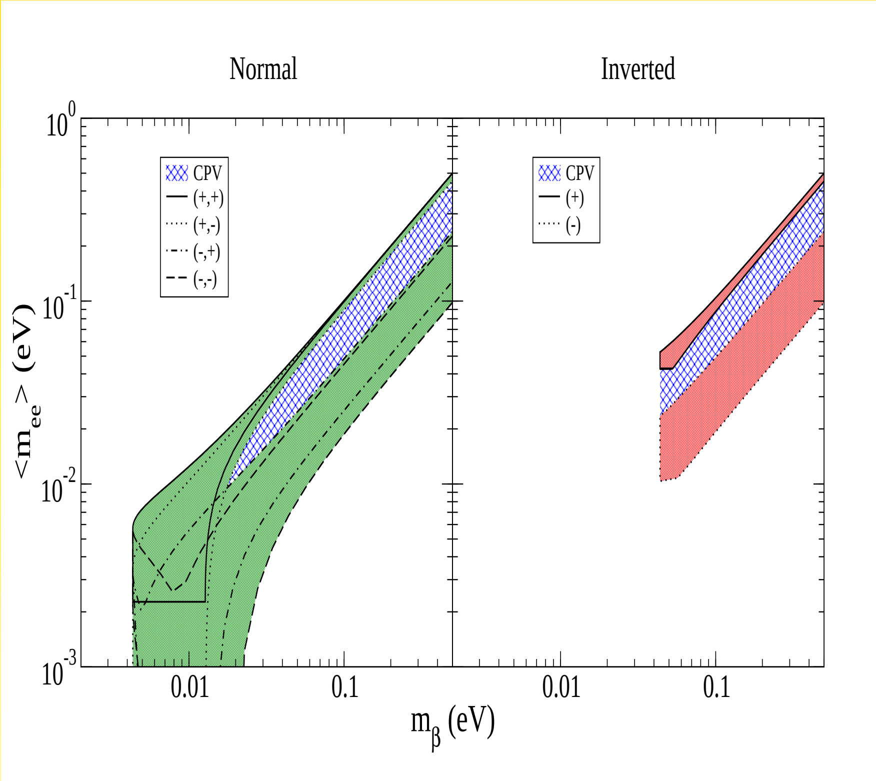
The inverse problem of 0$\nu\beta\beta$

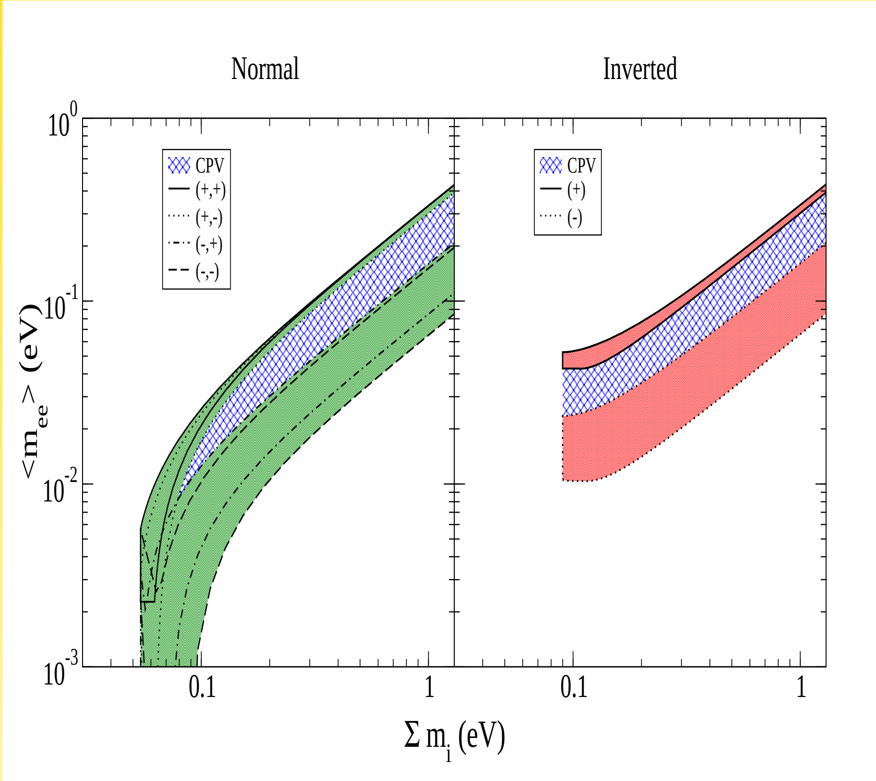
1.) Other observables (LHC, LFV, KATRIN, cosmology,...)

2.) Decay products (individual e^- energies, angular correlations, spectrum,...)

3.) Nuclear physics (multi-isotope, 0νECEC, $0\nu\beta^+\beta^+$,...)

1.) Distinguishing via other Observables





standard mechanism: KATRIN, cosmology

Energy Scale:

Note: *standard amplitude* for light Majorana neutrino exchange:

$$\mathcal{A}_{\rm l} \simeq G_F^2 \, \frac{|m_{ee}|}{q^2} \simeq 7 \times 10^{-18} \left(\frac{|m_{ee}|}{0.5 \text{ eV}} \right) \text{GeV}^{-5} \simeq 2.7 \text{ TeV}^{-5}$$

if new heavy particles are exchanged:

$$\mathcal{A}_{\rm h} \simeq \frac{c}{M^5}$$

$\Rightarrow$ for $0\nu\beta\beta$ holds:

$$\boxed{1 \text{ eV} = 1 \text{ TeV}}$$

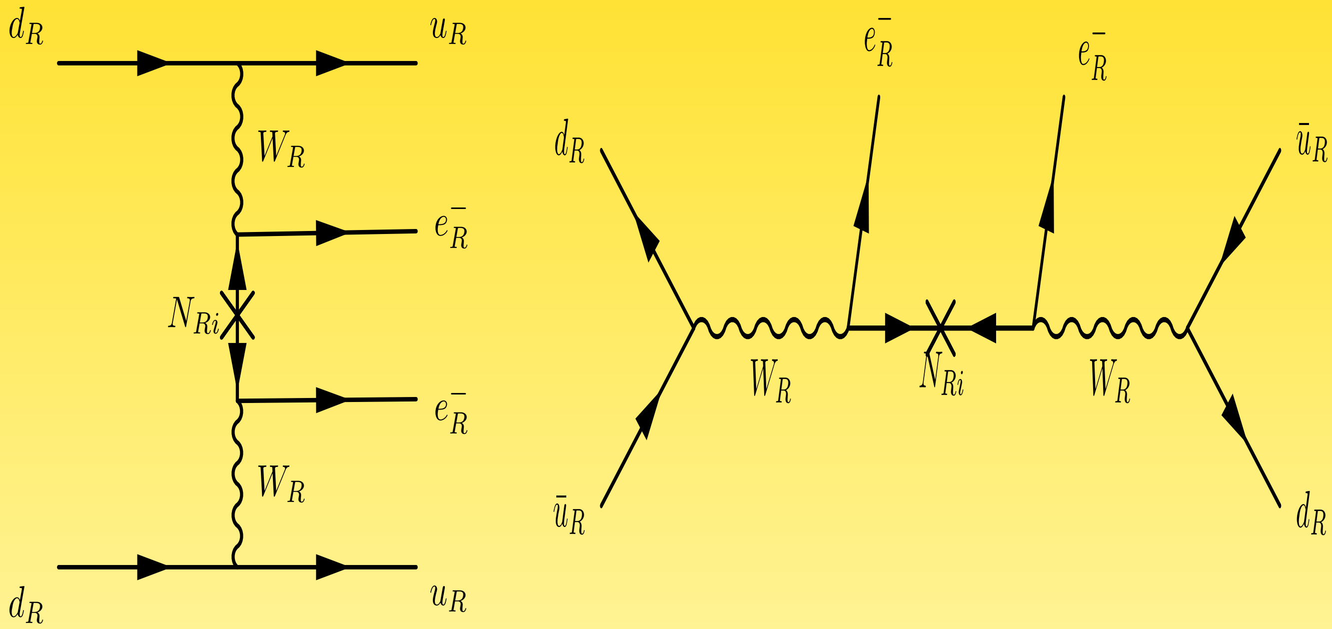
$\Rightarrow$ Phenomenology in colliders, LFV

Examples

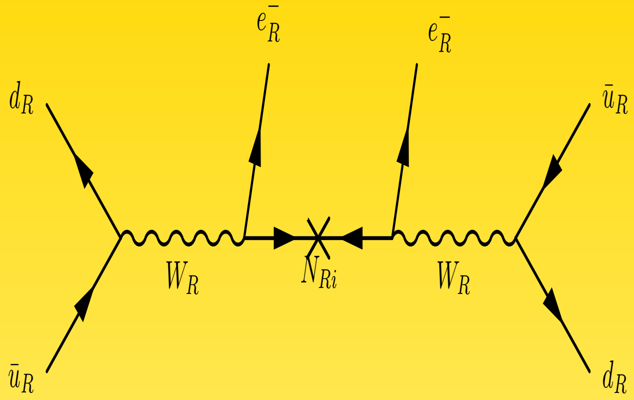
- R-parity violating supersymmetry (Allanach, Paes, Kom)
- TeV seesaw neutrinos (Ibarra, Petcov *et al.*; Mitra, Senjanovic, Vissani)
- Left-right symmetric theories (Senjanovic *et al.*; Goswami *et al.*)
- Color seesaw (Choubey, Duerr, Mitra, W.R.)

...focus only on one example here...

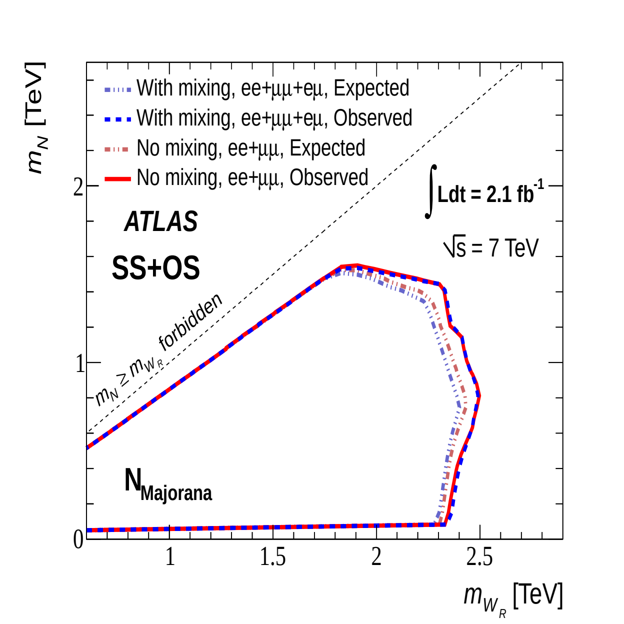
Left-right symmetry



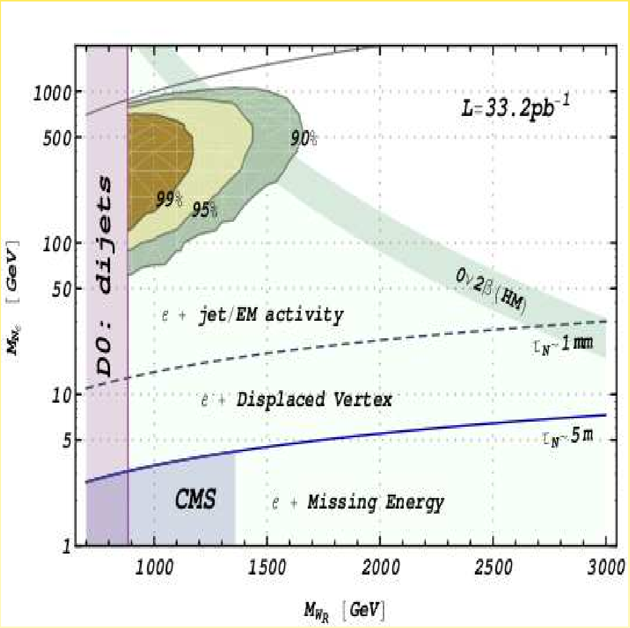
Senjanovic *et al.*, 1011.3522; 1103.1627



d_R
e_R^-
e_R^-
$\bar{u}_R$
W_R
N_{Ri}
W_R
$\bar{u}_R$
d_R

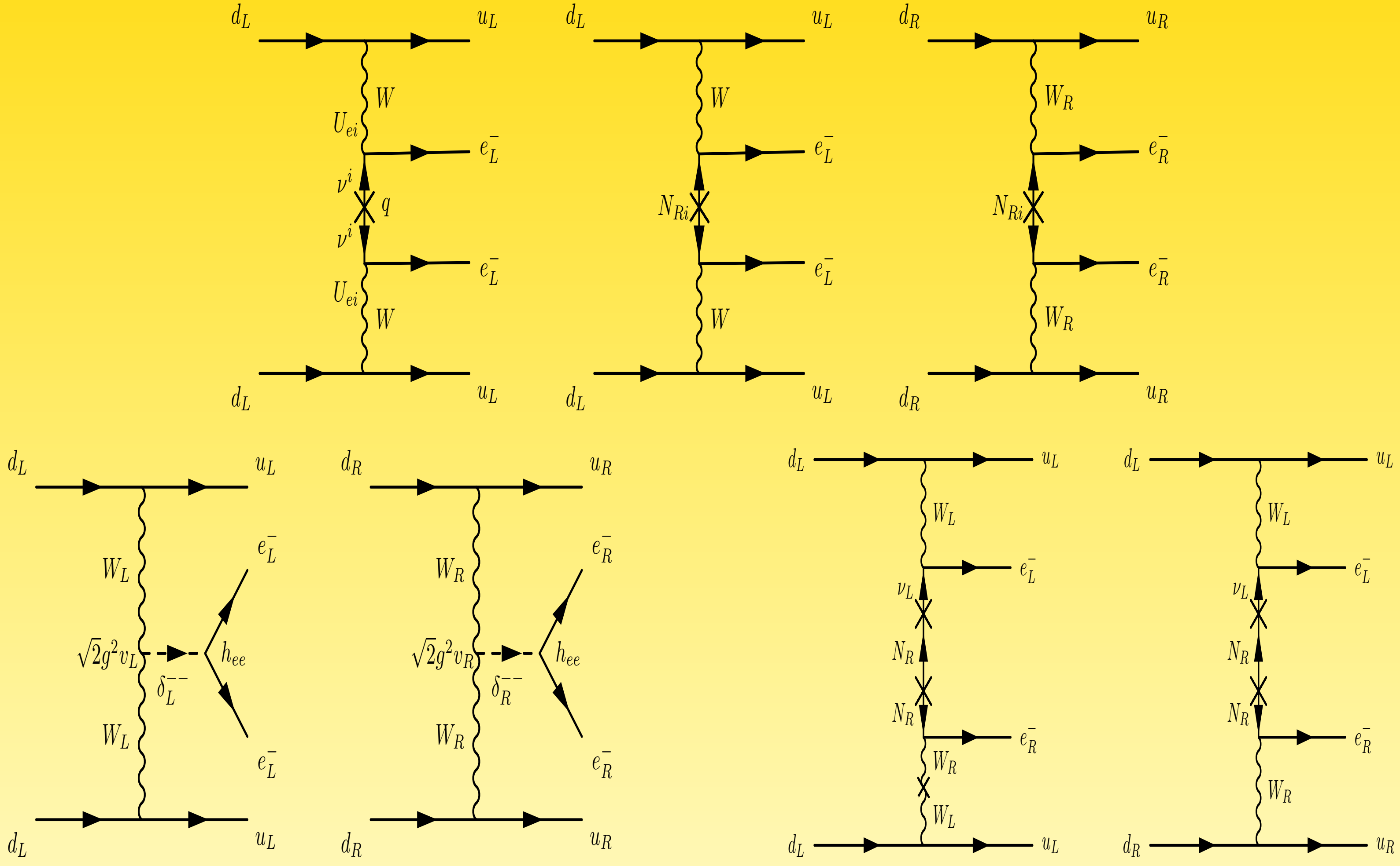


With mixing, ee+μμ+eμ, Expected
With mixing, ee+μμ+eμ, Observed
No mixing, ee+μμ, Expected
No mixing, ee+μμ, Observed
ATLAS
SS+OS
$\int$Ldt = 2.1 fb^{-1}
$\sqrt{s}$ = 7 TeV
$m_N \geq m_{W_R}$ forbidden
$N_{Majorana}$
m_N [TeV]
m_{W_R} [TeV]

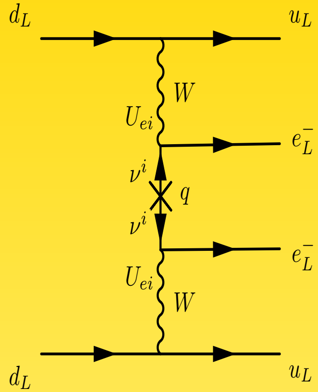


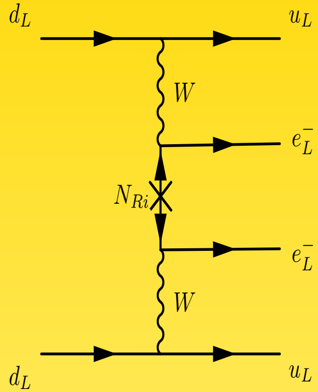
L=33.2pb^{-1}
90%
99%
95%
D0: dijets
$0\nu2\beta$ (HM)
e + jet/EM activity
τ_N~1 mm
e + Displaced Vertex
τ_N~5 m
CMS
e + Missing Energy
M_{N_e} [GeV]
M_{W_R} [GeV]

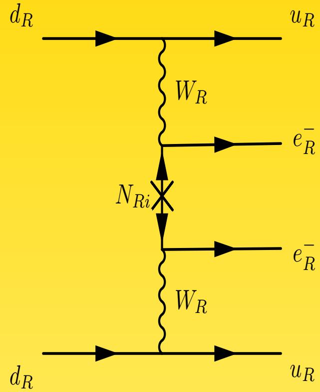
Left-right symmetry



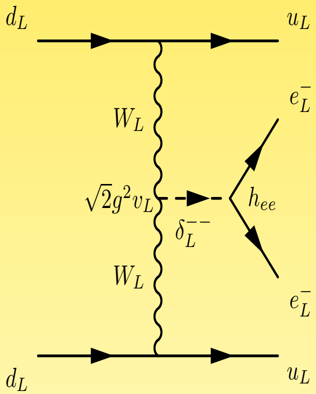
Left-right symmetry

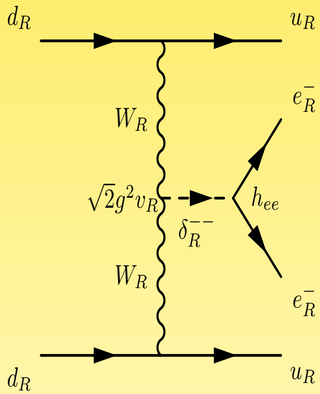


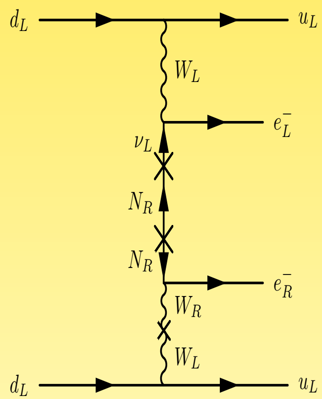


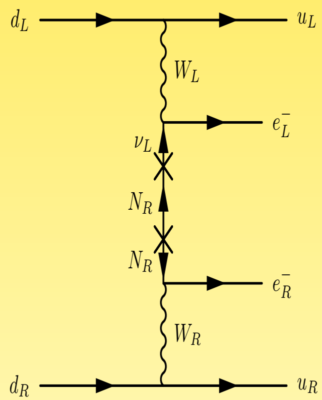


$U_{ei}^2\, m_i$ $\quad$ $\dfrac{S_{ei}^2}{M_i}$ $\quad$ $\dfrac{V_{ei}^2}{M_{W_R}^4\, M_i}$



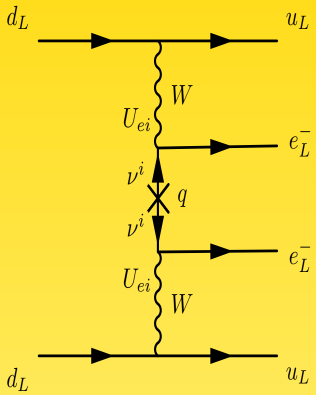




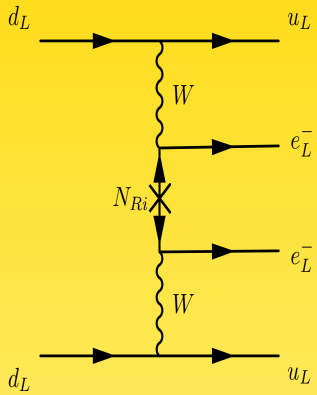


$\le \dfrac{U_{ei}^2\, m_i}{M_{\Delta_L}^2}$ $\quad$ $\dfrac{V_{ei}^2\, M_i}{M_{W_R}^4\, M_{\Delta_R}^2}$ $\quad$ $U_{ei}\, T_{ei} \tan\zeta$ $\quad$ $\dfrac{U_{ei}\, T_{ei}}{M_{W_R}^2}$

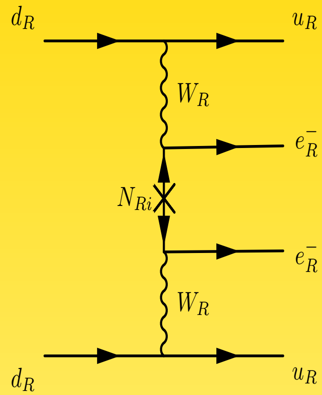
Left-right symmetry



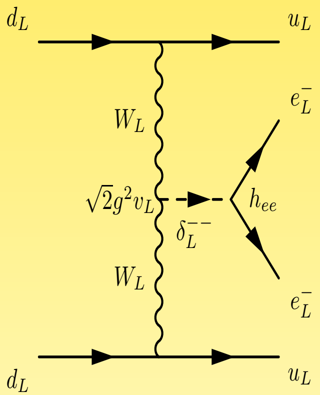
0.4 eV



2×10^{-8} GeV^{-1}



4×10^{-16} GeV^{-5}



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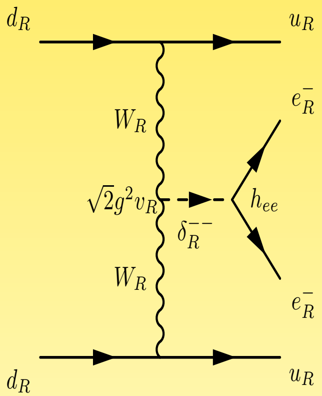
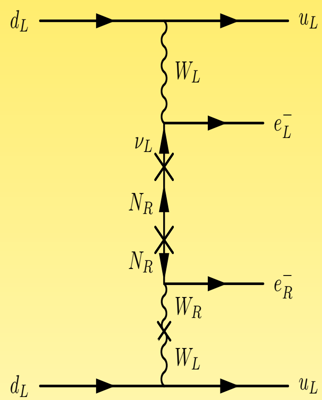
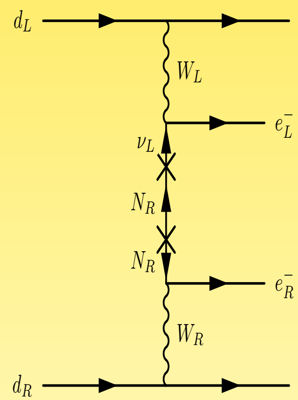


10^{-15} GeV^{-5}



6×10^{-9}



1.4×10^{-10} GeV^{-2}

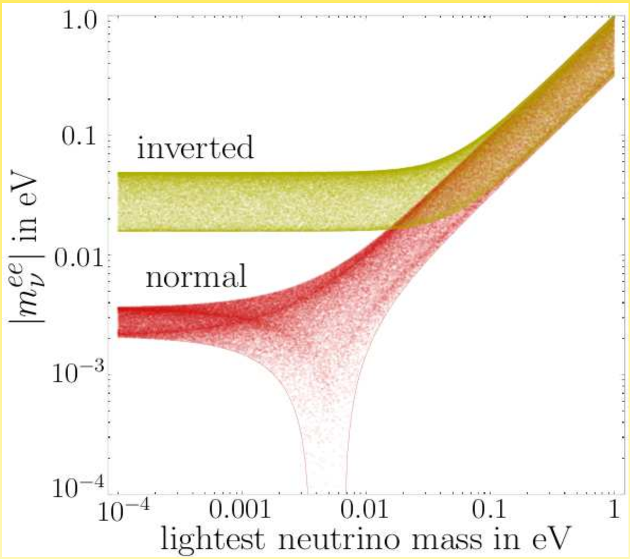
Simple and interesting scenario

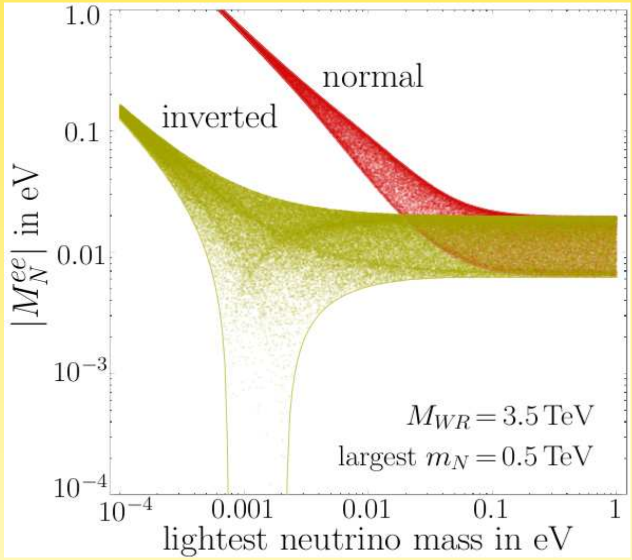
$$m_\nu = M_L - m_D\, M_R^{-1}\, m_D^T = v_L\, h - m_D\,(v_R\, f)^{-1}\, m_D^T$$

suppose M_L dominates in m_ν and $h = f$: $\Rightarrow M_R \propto m_\nu$

Triplet can mediate $\mu \to 3e$ at tree-level: $m_\Delta \gg M_i$

$$\Rightarrow \mathcal{A}_{N_R} \simeq G_F^2 \left(\frac{m_W}{M_{W_R}}\right)^4 \sum \frac{V_{ei}^2}{M_i} \propto \sum \frac{U_{ei}^2}{m_i}$$

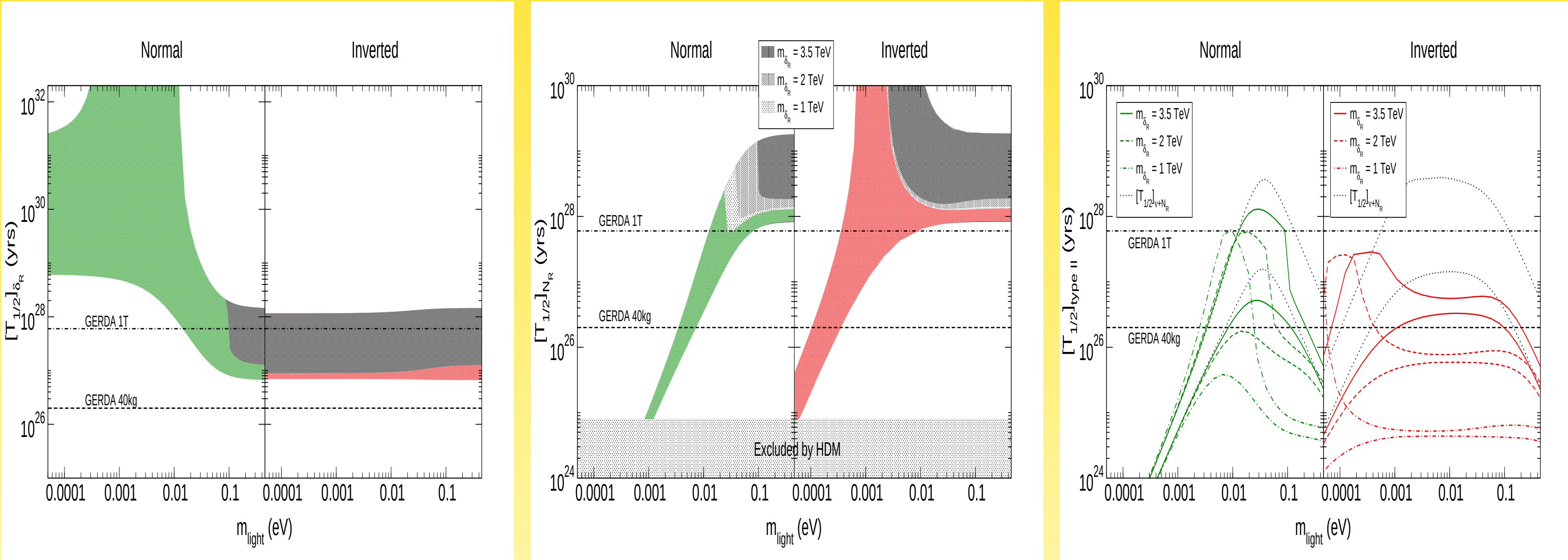




Tello *et al.*, 1011.3522

Interplay of diagrams in left-right symmetry

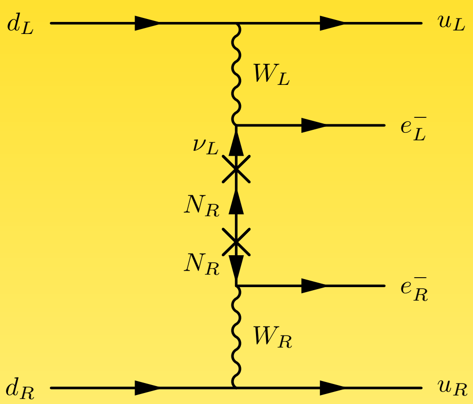
Interference of diagrams, constraints from LFV, neutrino data,...

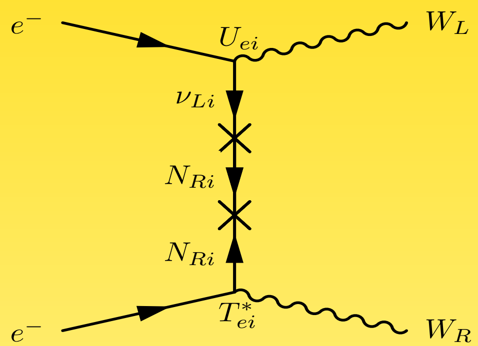


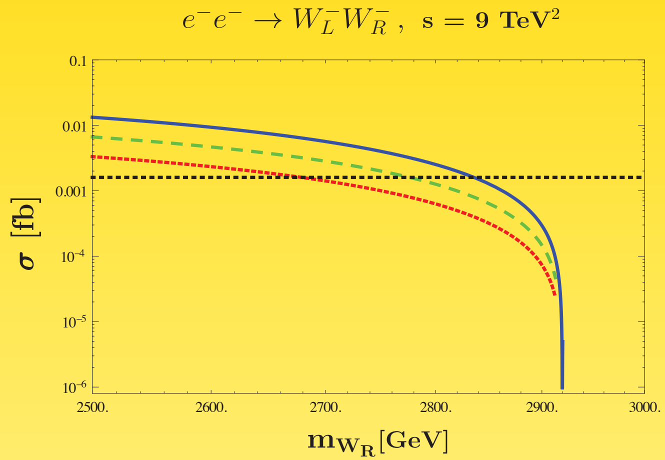
Barry, W.R., to appear; see also Goswami *et al.*, 1204.2527

Cleanest Probe: $e^- e^-$ collisions: "inverse $0\nu\beta\beta$"

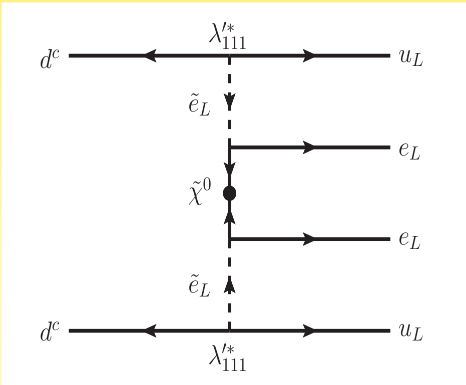
- LR-symmetry: (Barry, Dorame, W.R., 1204.3365)

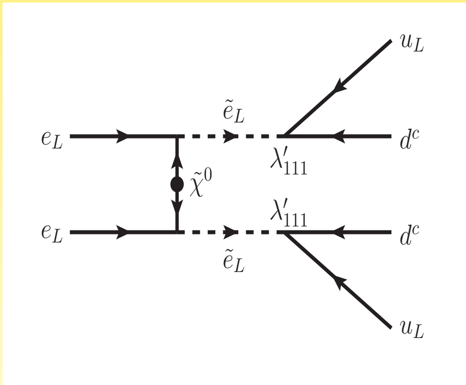


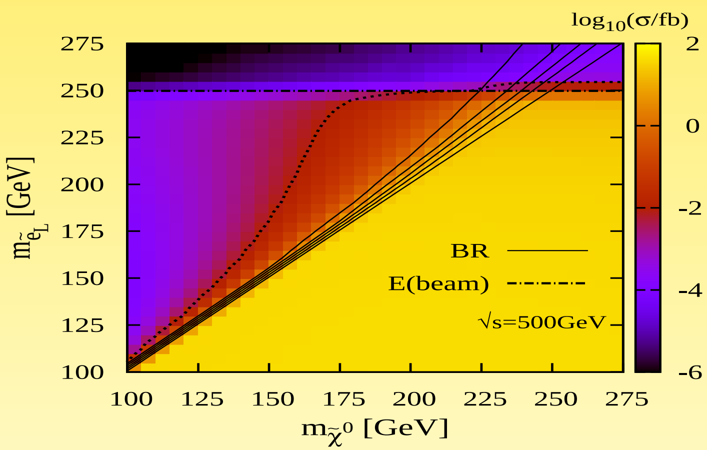




- SUSY: (Kom, W.R., 1110.3220)







"Inverse $0\nu\beta\beta$"

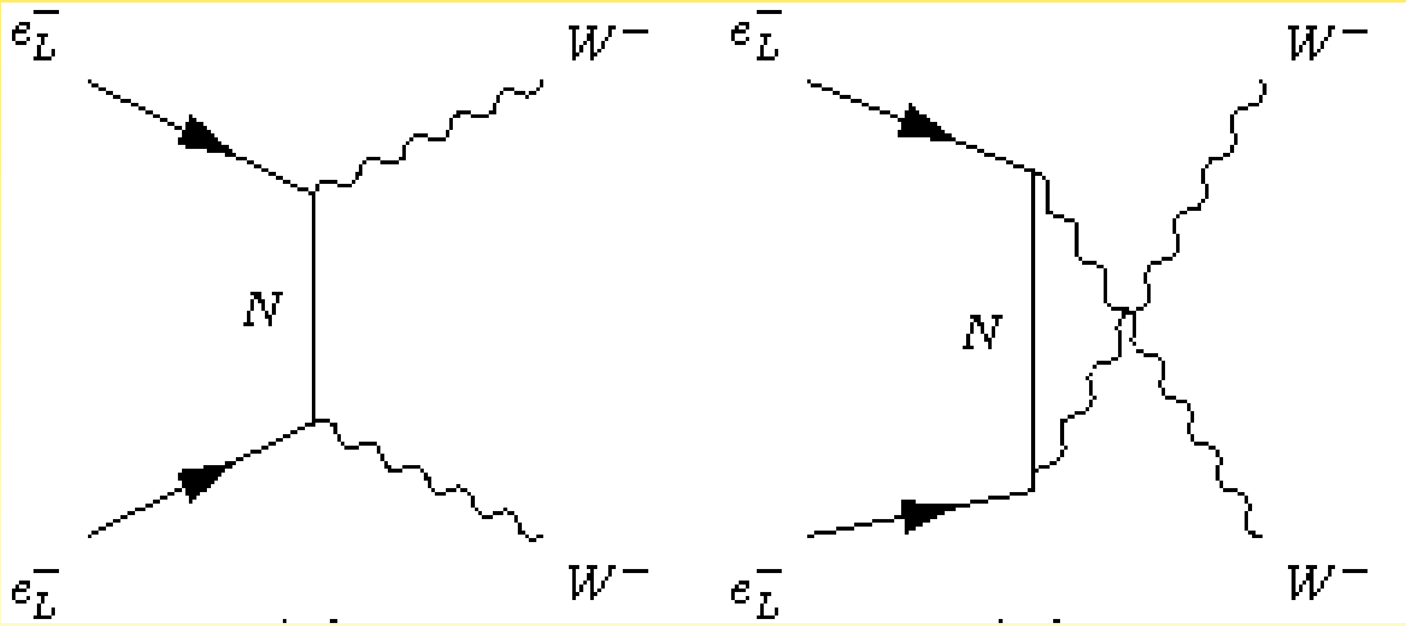
this is not

$$^{76}\mathrm{Se}^{++} + e^- + e^- \rightarrow \ ^{76}\mathrm{Ge}$$

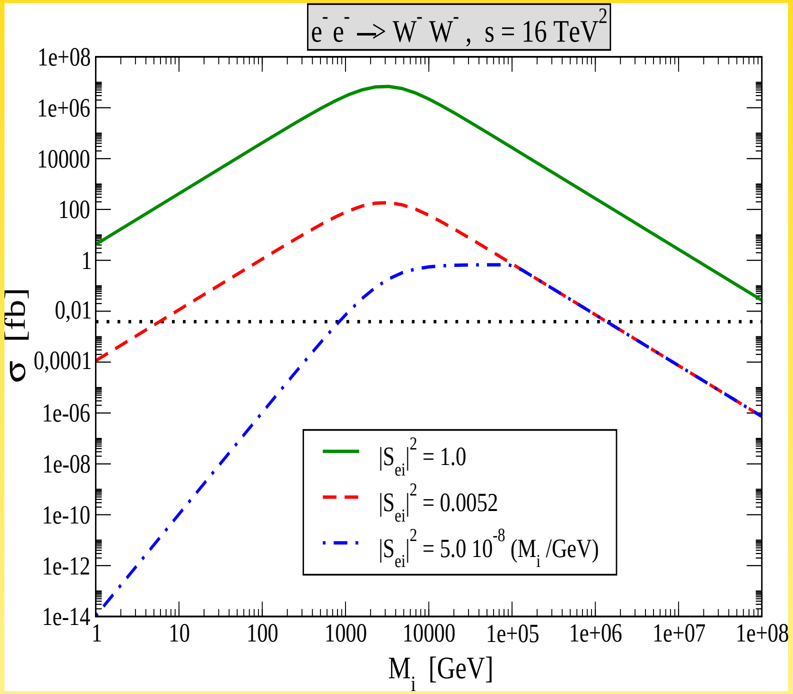
but rather

$$e^- + e^- \rightarrow W^- + W^-$$

Rizzo; Heusch, Minkowski; Gluza, Zralek; Cuypers, Raidal;...



Inverse Neutrinoless Double Beta Decay



W.R., PRD **81**

$$\frac{d\sigma}{d\cos\theta} = \frac{G_F^2}{32\,\pi}\left\{\sum (m_\nu)_i\,\mathcal{U}_{ei}^2\left(\frac{t}{t-(m_\nu)_i}+\frac{u}{u-(m_\nu)_i}\right)\right\}^2$$

Inverse Neutrinoless Double Beta Decay

Extreme limits:

- light neutrinos:

$$\sigma(e^- e^- \to W^- W^-) = \frac{G_F^2}{4\,\pi}\,|m_{ee}|^2 \leq 4.2 \cdot 10^{-18} \left(\frac{|m_{ee}|}{1\,\text{eV}}\right)^2 \text{fb}$$

$\Rightarrow$ way too small

- heavy neutrinos:

$$\sigma(e^- e^- \to W^- W^-) = 2.6 \cdot 10^{-3} \left(\frac{\sqrt{s}}{\text{TeV}}\right)^4 \left(\frac{S_{ei}^2/M_i}{5 \cdot 10^{-8}\,\text{GeV}^{-1}}\right)^2 \text{fb}$$

$\Rightarrow$ too small

- $\sqrt{s} \to \infty$:

$$\sigma(e^- e^- \to W^- W^-) = \frac{G_F^2}{4\pi} \left(\sum \mathcal{U}_{ei}^2\,(m_\nu)_i\right)^2$$

$\Rightarrow$ amplitude grows with $\sqrt{s}$? Unitarity??

Unitarity

high energy limit $\sqrt{s} \to \infty$:

$$\sigma(e^- e^- \to W^- W^-) = \frac{G_F^2}{4\pi} \left(\sum \mathcal{U}_{ei}^2 \, (m_\nu)_i \right)^2$$

$\leftrightarrow$ amplitude grows with $\sqrt{s}$?

Answer: exact see-saw relation $\mathcal{U}_{ei}^2 \, (m_\nu)_i = 0$

$$\mathcal{M} = \left(\begin{array}{cc} 0 & m_D \\ m_D^T & M_R \end{array} \right) = \mathcal{U} \left(\begin{array}{cc} m_\nu^{\rm diag} & 0 \\ 0 & M_R^{\rm diag} \end{array} \right) \mathcal{U}^T$$

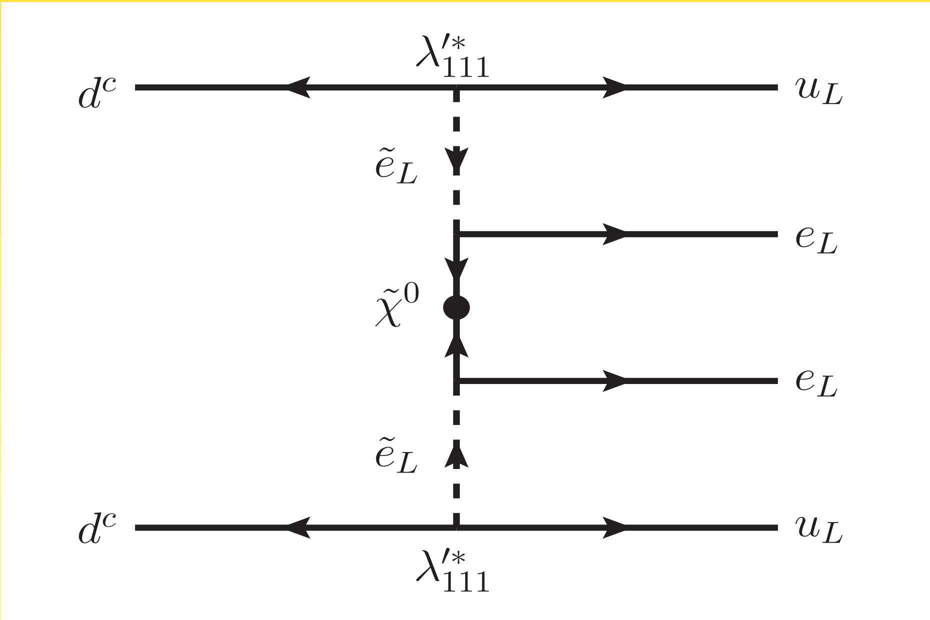
if Higgs triplet is present: unitarity also conserved

$$\sigma(e^- e^- \to W^- W^-) = \frac{G_F^2}{4\pi} \left((\mathcal{U}_{ei}^2 \, (m_\nu)_i - (m_L)_{ee} \right)^2 = 0$$

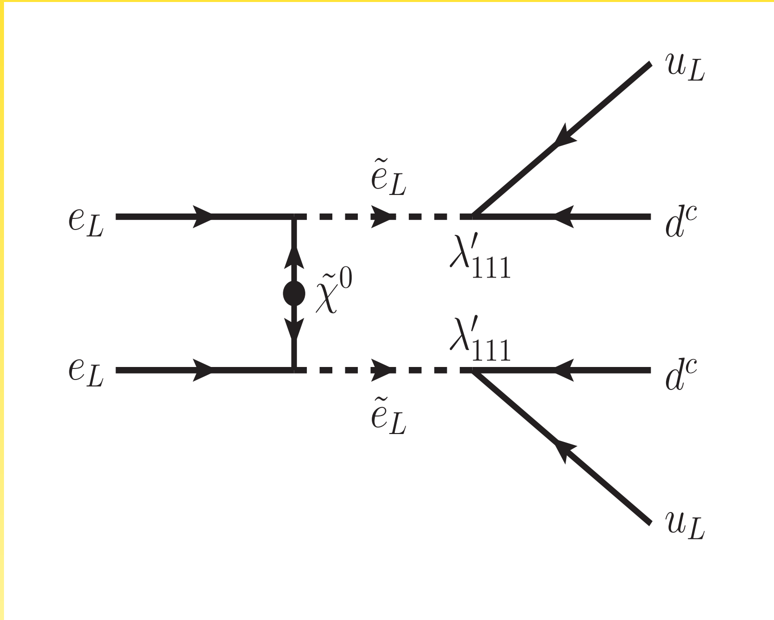
W.R., PRD **81**

Inverse $0\nu\beta\beta$ and RPV SUSY

$$\mathcal{W} = \lambda'_{111} L_1 Q_1 D_1^c \;\Rightarrow\; e^- e^- \to 4 \text{ jets}$$



$0\nu\beta\beta$



resonant selectron production

via *gauge* interactions

Cross section

$$\sigma(e_L^- e_L^- \to \tilde{e}_L^- \tilde{e}_L^-) = \frac{\pi\alpha^2|g_L|^4}{s} \frac{2m_{\tilde{\chi}^0}^2}{s + 2m_{\tilde{\chi}^0}^2 - 2m_{\tilde{e}_L}^2} \left[L + \frac{2\lambda}{(s + 2m_{\tilde{\chi}^0}^2 - 2m_{\tilde{e}_L}^2)^2 - \lambda^2} \right]$$

where

$$L = \ln \frac{s + 2m_{\tilde{\chi}^0}^2 - 2m_{\tilde{e}_L}^2 + \lambda}{s + 2m_{\tilde{\chi}^0}^2 - 2m_{\tilde{e}_L}^2 - \lambda}$$

$$\lambda = \lambda(s, m_{\tilde{e}_L}^2, m_{\tilde{e}_L}^2) = \sqrt{s^2 - 4sm_{\tilde{e}_L}^2}$$

Keung, Littenberg, 1983

adjustable parameters

$$m_{\tilde{\chi}^0}, \quad m_{\tilde{g}}, \quad m_{\tilde{e}_L}, \quad m_{\tilde{u}_L}, \quad m_{\tilde{d}_R}, \quad \lambda'_{111}$$

squarks and gluinos decoupled;

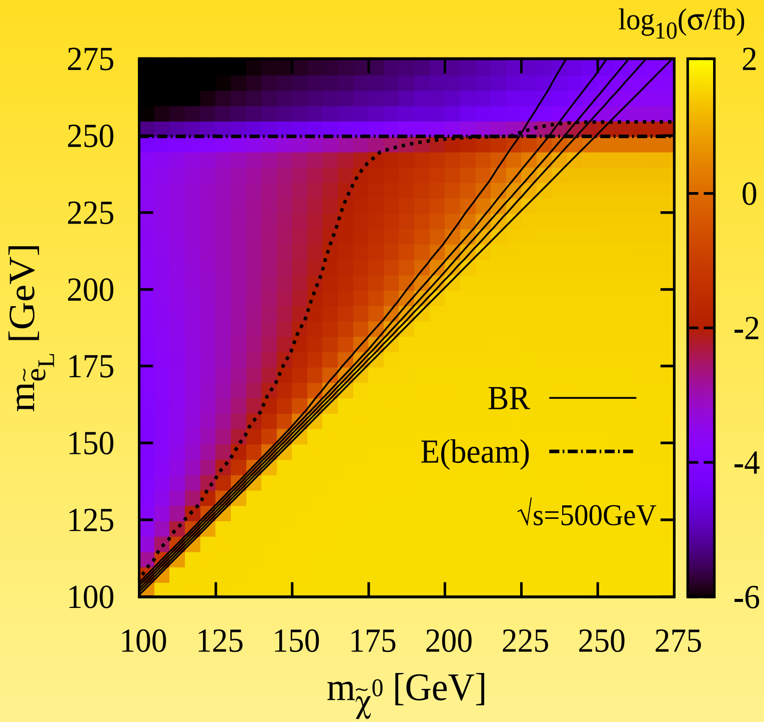
competing decays $\tilde{e}_L \to e\,\tilde{\chi}^0$ and $\tilde{e}_L \to jj$

competing decays $\tilde{e}_L \to e\,\tilde{\chi}^0$ and $\tilde{e}_L \to jj$:

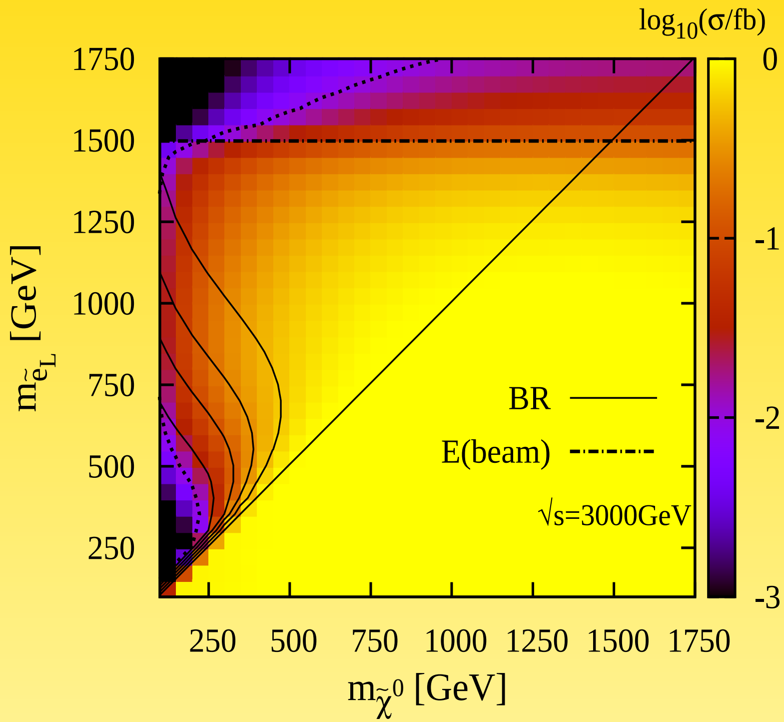
- $0\nu\beta\beta$-limit goes with $\Lambda^5_{\rm SUSY} \Rightarrow \lambda'_{111}$ can be $\mathcal{O}(1)$ and thus BR$(\tilde{e}_L \to jj) >$ BR$(\tilde{e}_L \to e\,\tilde{\chi}^0)$
- even for low masses, large BR$(\tilde{e}_L \to jj)$ possible for narrow band around $m_{\tilde{e}_L} - m_{\tilde{\chi}^0} \ll m_{\tilde{e}_L}$

reconstruction:

- mass and width of $\tilde{e}_L$: dijet invariant mass distribution
- BR$(\tilde{e}_L \to jj)$ and thus λ'_{111}: $\tilde{e}_L$ decays
- mass of $\tilde{\chi}^0$: rate of $e^-_L e^-_L \to \tilde{e}^-_L \tilde{e}^-_L$



1.9×10^{25} yrs



1.0×10^{27} yrs

Kom, W.R., 1110.3220

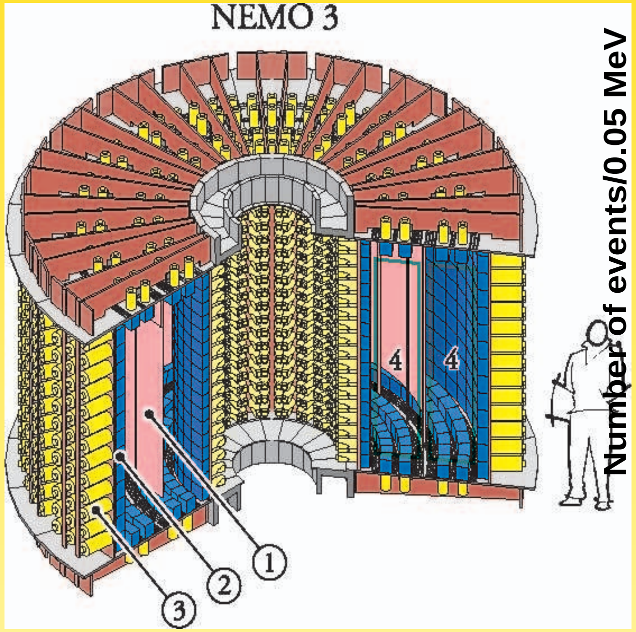
Summary

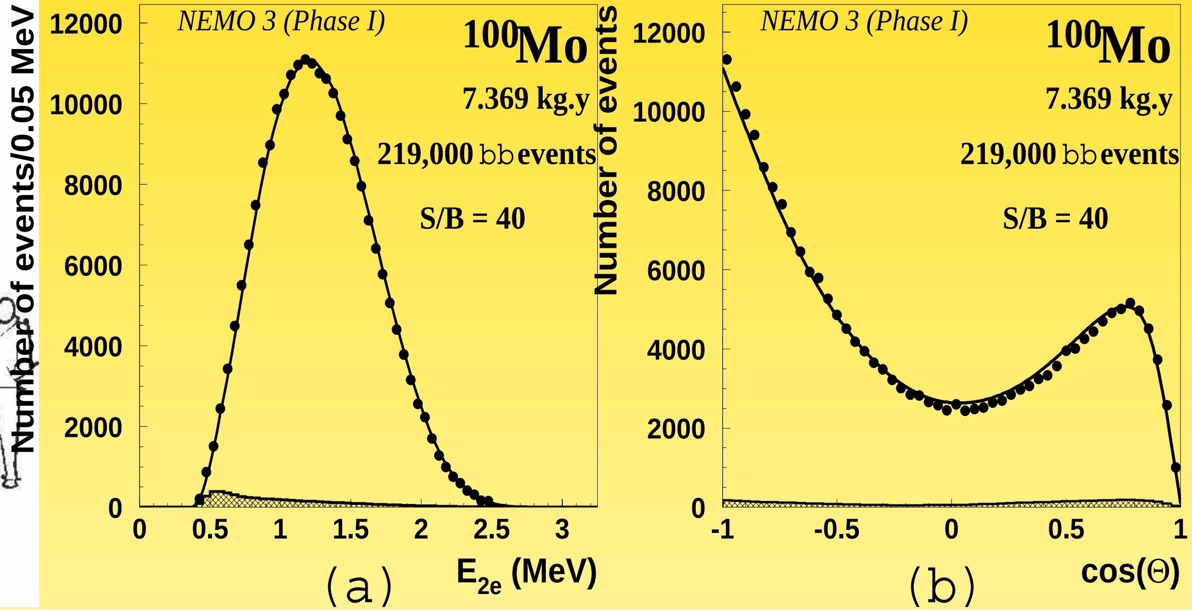
Chi l'ha visto?

Ettore Majorana, ordinario di fisica teorica all' Università di Napoli, è misteriosamente scomparso dagli ultimi di marzo. Di anni 31, alto metri 1,70, snello, con capelli neri, occhi scuri, una lunga cicatrice sul dorso di una mano. Chi ne sapesse qualcosa è pregato di scrivere al R. P. E. Marianecci, Viale Regina Margherita 66 - Roma.

2.) Distinguishing via decay products

SuperNEMO

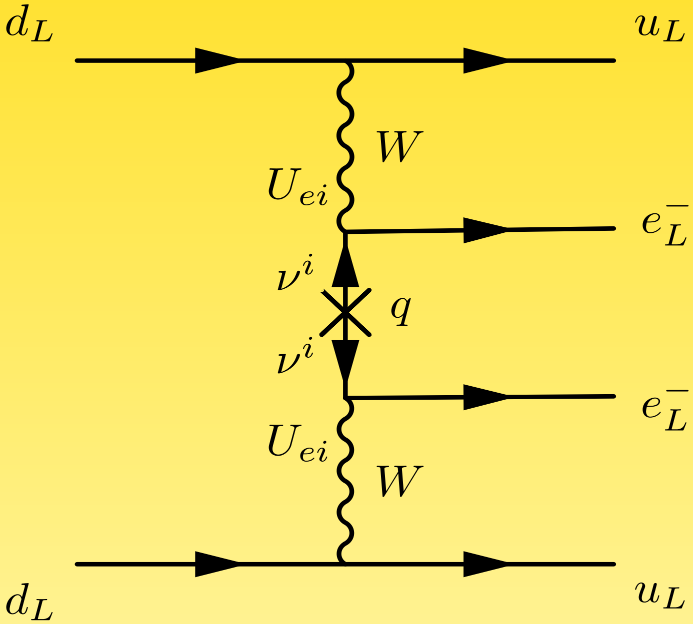


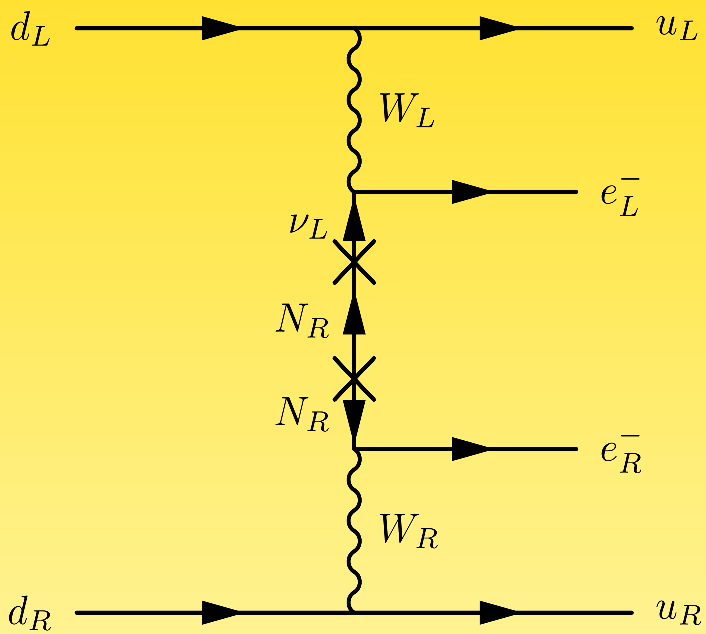


- source foils in between plastic scintillators
- individual electron energy, and their relative angle!

2.) Distinguishing via decay products

Consider standard plus λ-mechanism





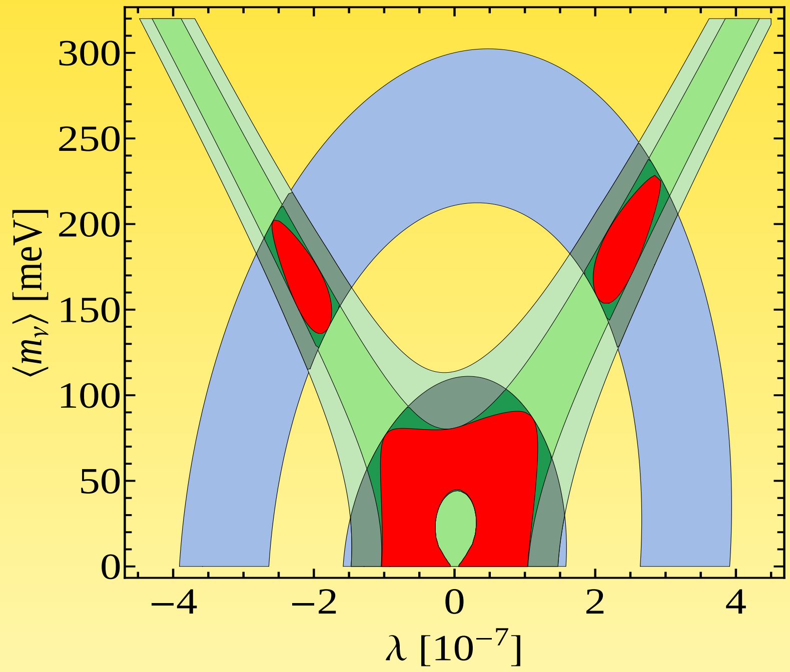
$$\frac{d\Gamma}{dE_1\, dE_2\, d\cos\theta} \propto (1 - \beta_1\, \beta_2\, \cos\theta) \qquad \frac{d\Gamma}{dE_1\, dE_2\, d\cos\theta} \propto (E_1 - E_2)^2\, (1 + \beta_1\, \beta_2\, \cos\theta)$$

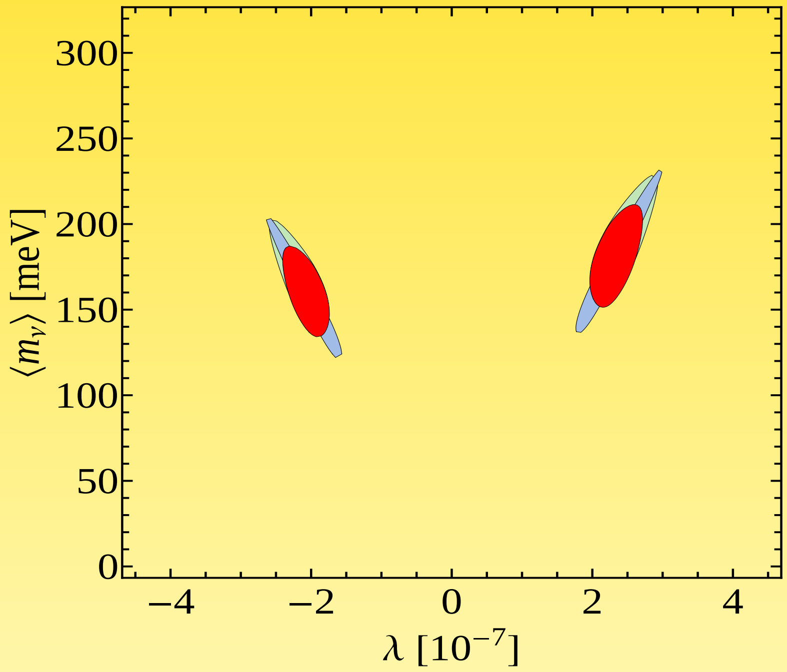
Arnold *et al.*, 1005.1241

2.) Distinguishing via decay products

Defining asymmetries

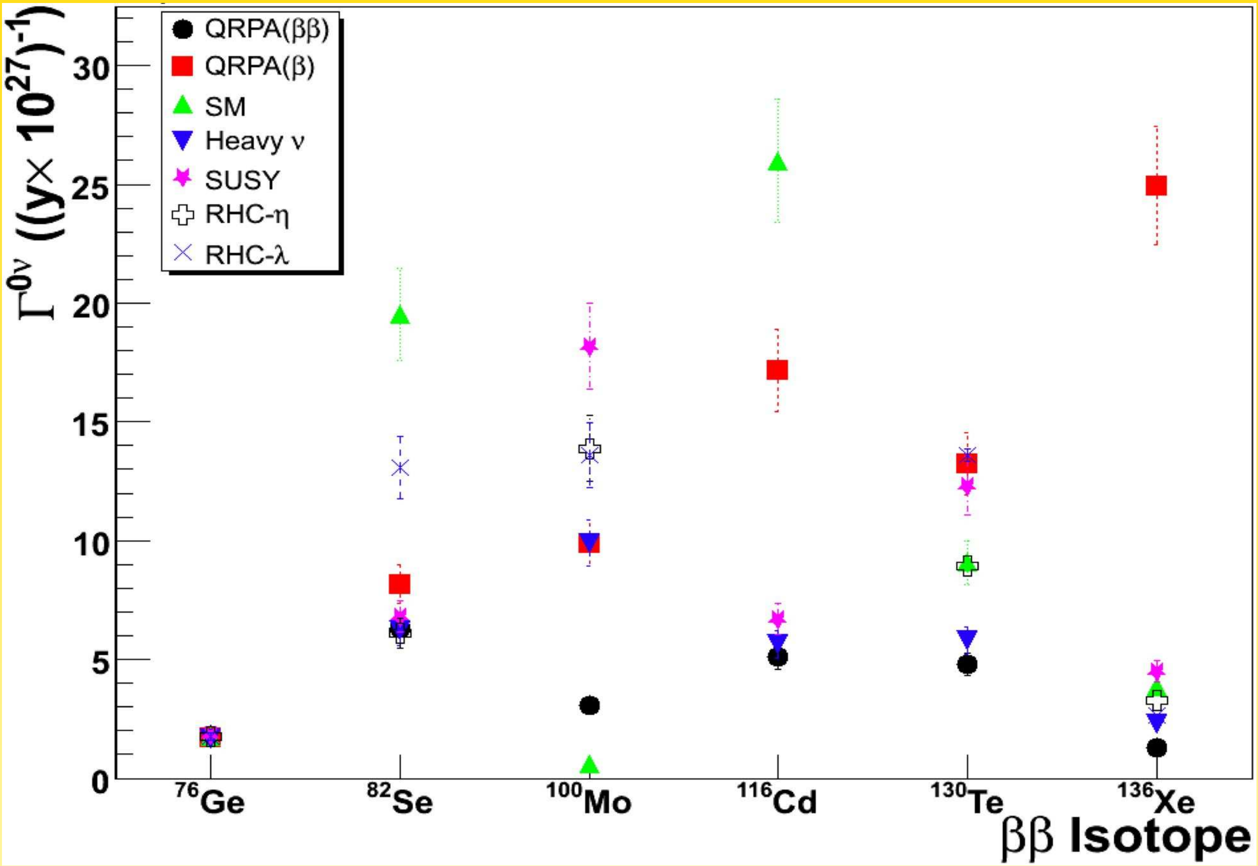
$A_\theta = (N_+ - N_-)/(N_+ + N_-)$ and $A_E = (N_> - N_<)/(N_> + N_<)$





SuperNEMO: Arnold *et al.*, 1005.1241

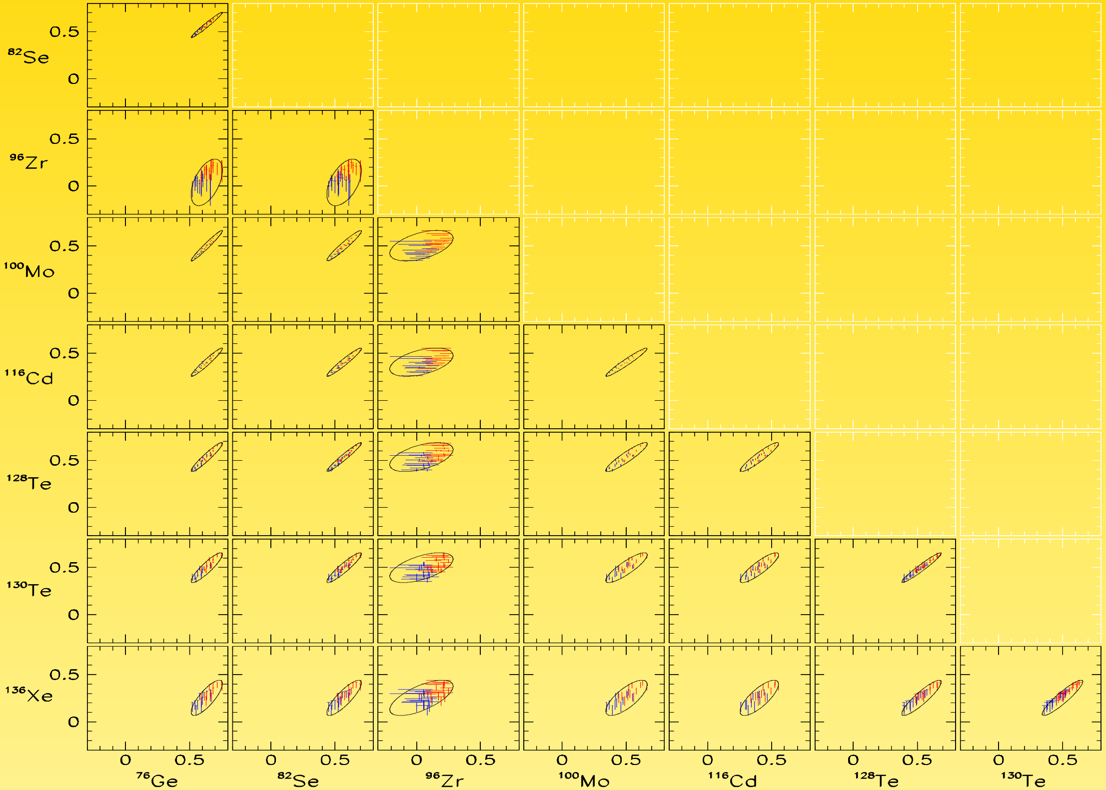
3.) Distinguishing via nuclear physics



Gehman, Elliott, hep-ph/0701099

3 to 4 isotopes necessary to disentangle mechanism

to better estimate error range: correlations need to be understood:



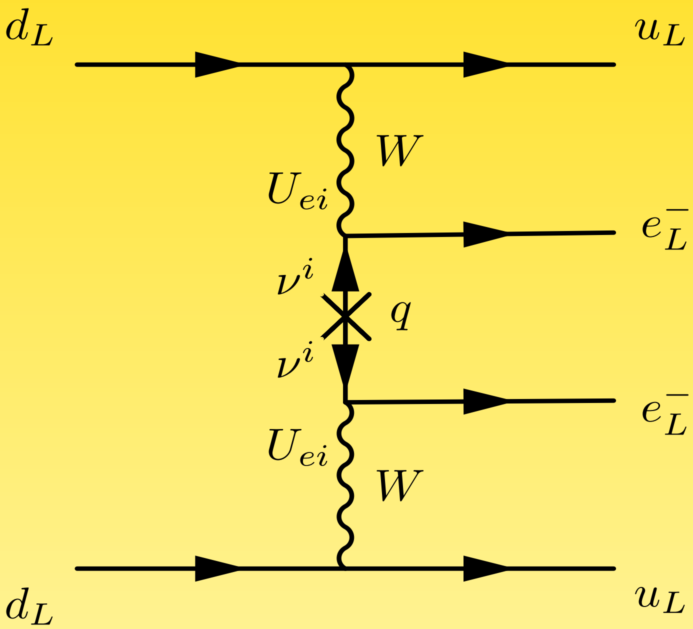
Faessler, Fogli *et al.*, PRD **79**

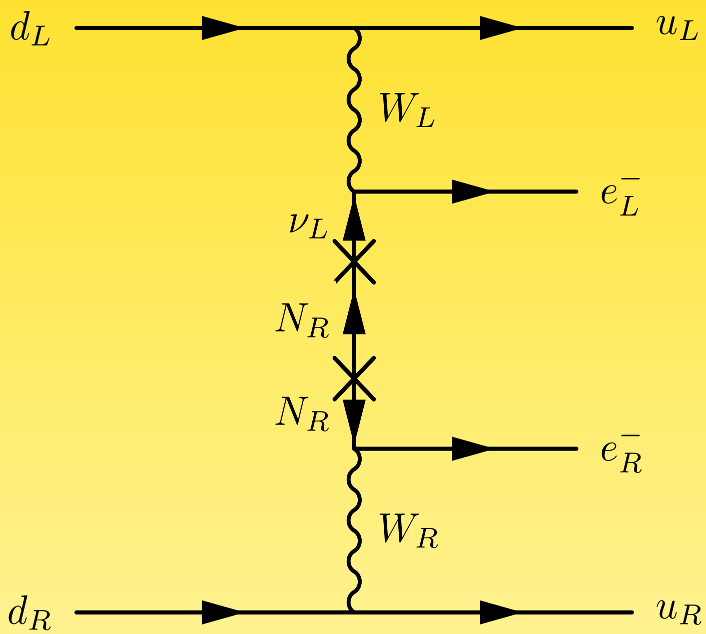
ellipse major axis: SRC (blue, red) and g_A

ellipse minor axis: g_{pp}

Distinguishing via decay products

Consider standard plus λ-mechanism

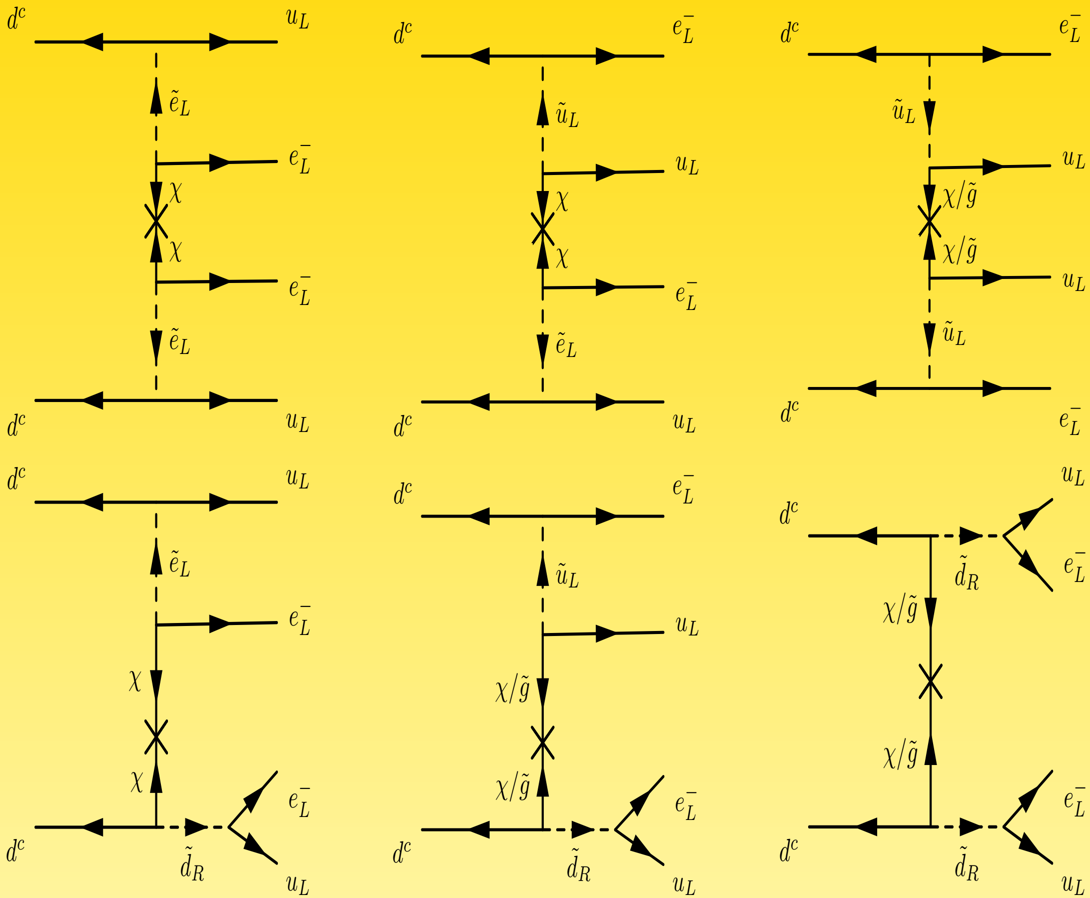




$$\frac{d\Gamma}{dE_1\, dE_2\, d\cos\theta} \propto (1 - \beta_1\, \beta_2\, \cos\theta) \qquad \frac{d\Gamma}{dE_1\, dE_2\, d\cos\theta} \propto (E_1 - E_2)^2\, (1 + \beta_1\, \beta_2\, \cos\theta)$$

Arnold *et al.*, 1005.1241

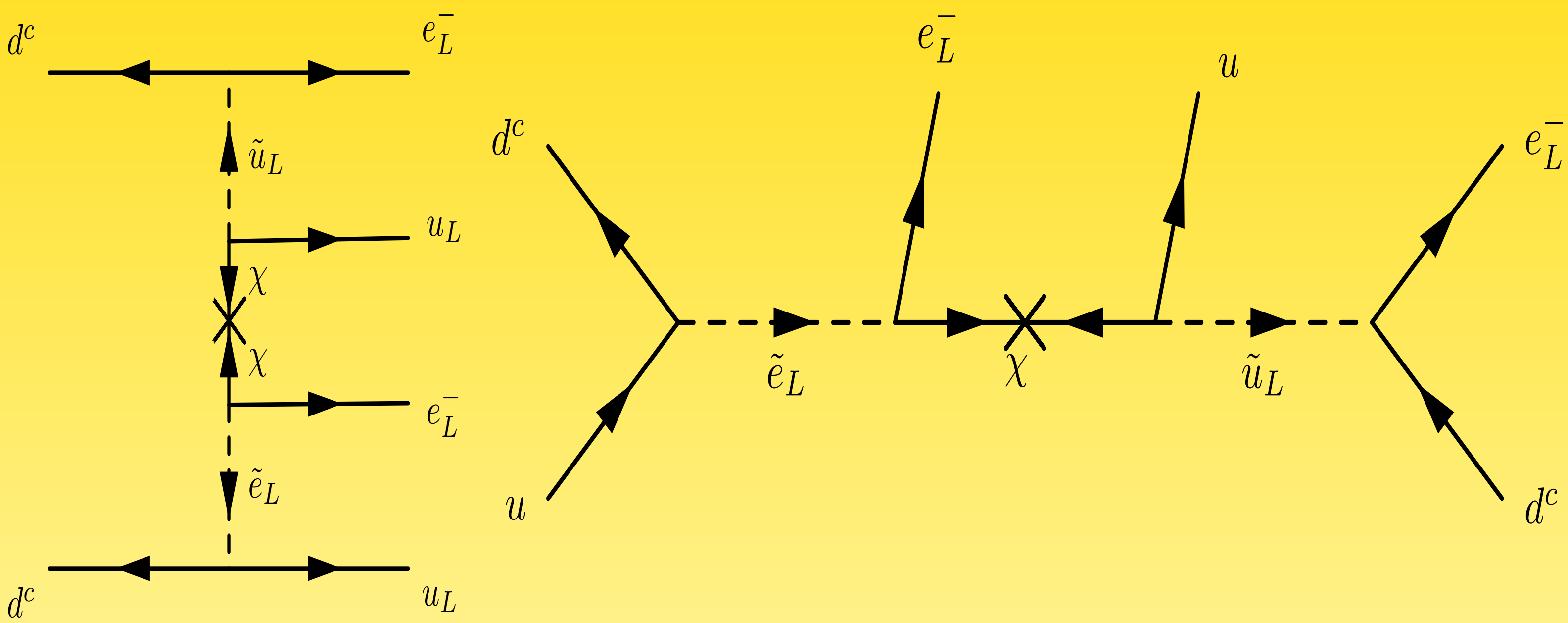
Supersymmetry: short range



$$\mathcal{A}_{\not{R}_1} \simeq \frac{\lambda'^2_{111}}{\Lambda^5_{\mathrm{SUSY}}}$$

Supersymmetry: short range

interplay with LHC:

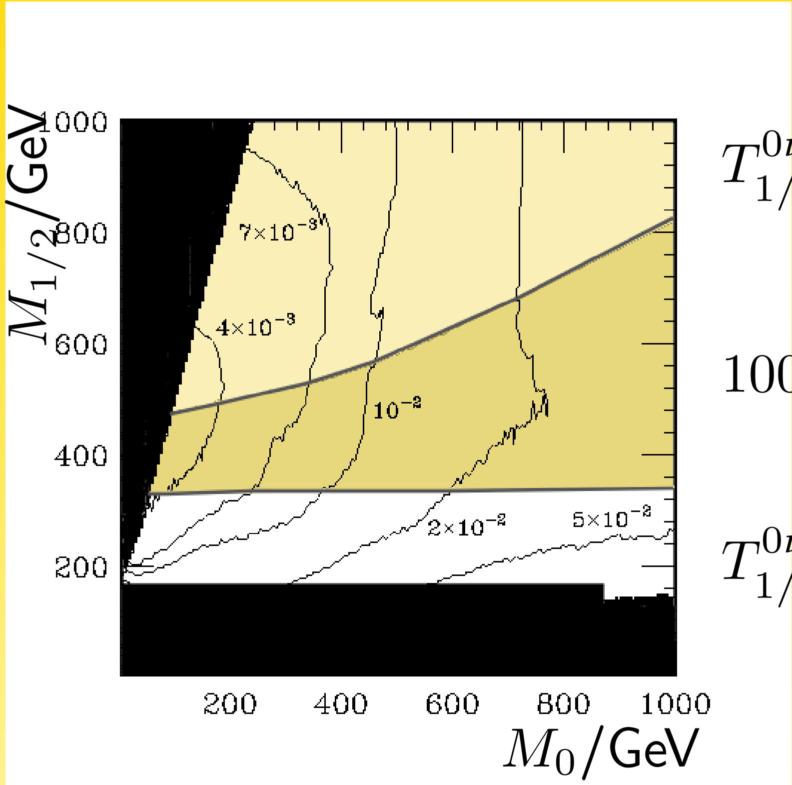


"resonant selectron production"

$$\hat{\sigma} \propto \frac{\lambda'^2_{111}}{\hat{s}}$$

Allanach, Kom, Paes, 0903.0347

$\tan\beta = 10,\ A_0 = 0,\ 10\ \text{fb}^{-1}$



$T_{1/2}^{0\nu\beta\beta}(\text{Ge}) > 1 \times 10^{27}$ yrs

$100 > T_{1/2}^{0\nu\beta\beta}(\text{Ge})/10^{25}$ yrs > 1.9

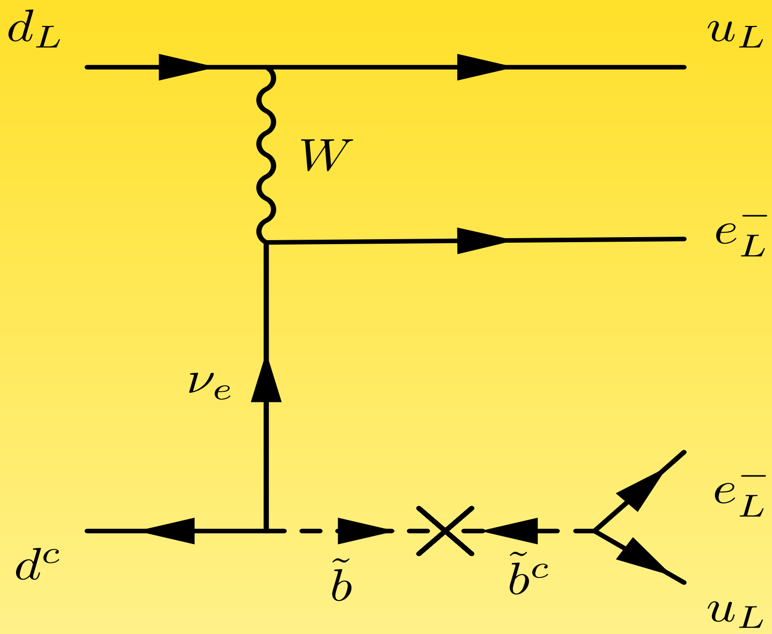
$T_{1/2}^{0\nu\beta\beta}(\text{Ge}) < 1.9 \times 10^{25}$ yrs

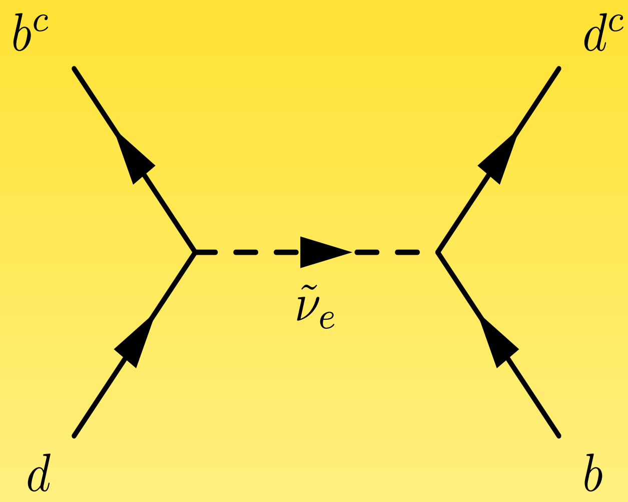
$\rightarrow$ observation in white region in conflict with $0\nu\beta\beta$

$\rightarrow$ if $0\nu\beta\beta$ observed: dark yellow region tests $R\!\!\!/$ SUSY mechanism

$\rightarrow$ light yellow region: no significant $R\!\!\!/$ contribution to $0\nu\beta\beta$

Supersymmetry: long range





$$\mathcal{A}^b_{\not{R}_2} \simeq G_F \, \frac{1}{q} \, U_{ei} \, m_b \, \frac{\lambda'_{131} \, \lambda'_{113}}{\Lambda^3_{\mathrm{SUSY}}}$$

$0\nu\beta\beta$

$$\frac{\lambda'_{131} \, \lambda'_{113}}{\Lambda^2_{\mathrm{SUSY}}}$$

B^0-$\bar{B}^0$ mixing

Flavor Symmetry Models

suppose your model predicts TBM:

$$(m_\nu)_{\text{TBM}} = \begin{pmatrix} x & y & y \\ \cdot & z+x & y-z \\ \cdot & \cdot & z+x \end{pmatrix}$$

$$m_1 = x - y \ , \quad m_2 = x + 2y \ , \quad m_3 = x - y + 2z$$

if $z = y + x/2$, then:

$$m_1 = x - y \ , \quad m_2 = x + 2y \ , \quad m_3 = 2x + y$$

and one has a neutrino mass sum-rule

$$m_1 + m_2 = m_3$$

The Zoo (of A_4 models)

Type	L_i	ℓ_i^c	ν_i^c	Δ	References
A1				-	[1–14] [15]$^{\#}$
A2	$\underline{3}$	$\underline{1}, \underline{1}', \underline{1}''$	-	$\underline{1}, \underline{1}', \underline{1}'', \underline{3}$	[16–18]
A3				$\underline{1}, \underline{3}$	[19]
B1	$\underline{3}$	$\underline{1}, \underline{1}', \underline{1}''$	$\underline{3}$	-	[4, 20–27]$^{\#}$ [28–30]* [31–45]
B2				$\underline{1}, \underline{3}$	[46]$^{\#}$
C1				-	[2, 47, 48]
C2	$\underline{3}$	$\underline{3}$	-	$\underline{1}$	[49, 50] [51]$^{\#}$
C3				$\underline{1}, \underline{3}$	[52]
C4				$\underline{1}, \underline{1}', \underline{1}'', \underline{3}$	[53]
D1				-	[54, 55]$^{\#}$ [56, 57]* [58]
D2	$\underline{3}$	$\underline{3}$	$\underline{3}$	$\underline{1}$	[59] [60]*
D3				$\underline{1}'$	[61]*
D4				$\underline{1}', \underline{3}$	[62]*
E	$\underline{3}$	$\underline{3}$	$\underline{1}, \underline{1}', \underline{1}''$	-	[63, 64]
F	$\underline{1}, \underline{1}', \underline{1}''$	$\underline{3}$	$\underline{3}$	$\underline{1}$ or $\underline{1}'$	[65]
G	$\underline{3}$	$\underline{1}, \underline{1}', \underline{1}''$	$\underline{1}, \underline{1}', \underline{1}''$	-	[66]
H	$\underline{3}$	$\underline{1}, \underline{1}, \underline{1}$	-	-	[67]
I	$\underline{3}$	$\underline{1}, \underline{1}, \underline{1}$	$\underline{1}, \underline{1}, \underline{1}$	-	[68]*
J	$\underline{3}$	$\underline{1}, \underline{1}, \underline{1}$	$\underline{3}$	-	[12, 39, 69, 70]
K	$\underline{3}$	$\underline{1}, \underline{1}, \underline{1}$	$\underline{1}, \underline{1}$	$\underline{1}$	[71]*
L	$\underline{3}$	$\underline{1}, \underline{1}, \underline{1}$	$\underline{1}$	-	[72]*
M	$\underline{1}, \underline{1}', \underline{1}''$	$\underline{1}, \underline{1}'', \underline{1}'$	$\underline{3}, \underline{1}$	-	[73, 74]
N	$\underline{1}, \underline{1}', \underline{1}''$	$\underline{1}, \underline{1}'', \underline{1}'$	$\underline{3}, \underline{1}', \underline{1}''$	-	[75]

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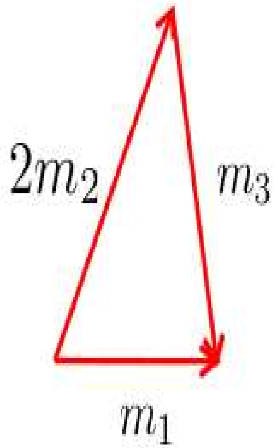
http://www.mpi-hd.mpg.de/personalhomes/jamesb/Table_A4.pdf

Sum-rules in Models and $0\nu\beta\beta$



Sum-rule	Flavour symmetry
$2m_2 + m_3 = m_1$	$A_4, T', (S_4)$
$m_1 + m_2 = m_3$	$S_4, (A_4)$
$\frac{2}{m_2} + \frac{1}{m_3} = \frac{1}{m_1}$	A_4, T'
$\frac{1}{m_1} + \frac{1}{m_2} = \frac{1}{m_3}$	S_4

constrains masses and Majorana phases

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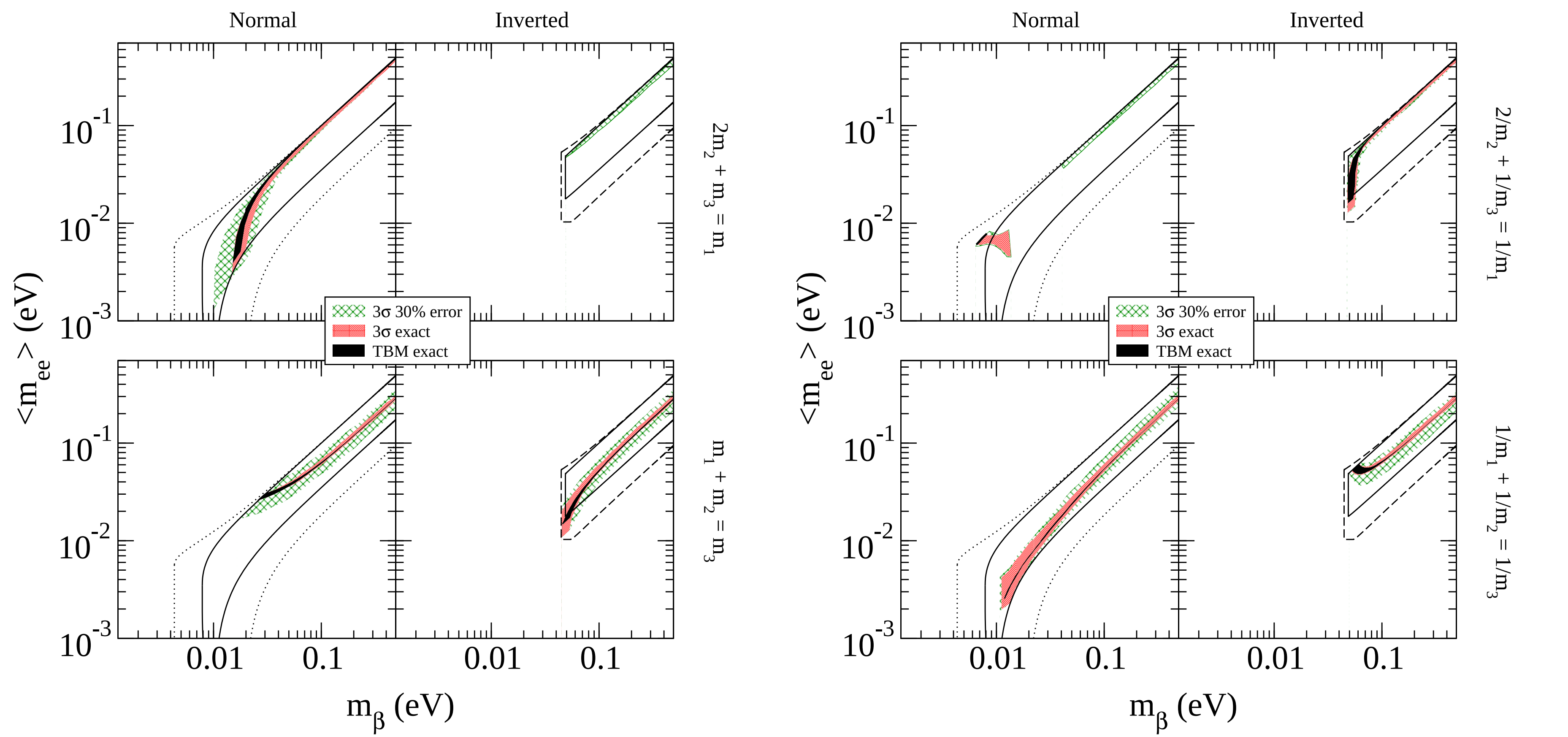

$$m_1 + m_2 - m_3 = \epsilon\, m_{\rm max}$$

stable: new solutions not before $\epsilon \simeq 0.2$