## Exercises to "Standard Model of Particle Physics II"

Winter 2024/25

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Hand-in of solutions:	Discussion of solutions:
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## Problem 1: The Rho-parameter [10 Points]

We encountered the Rho-parameter defined by

$$\rho \equiv \frac{M_W^2}{M_Z^2 \cos^2(\theta_W)}$$

whose value predicted by the SM is 1. Now we generalize our findings while keeping the above definition.

- a) From representation theory we know that in a *n*-dimensional representation (i.e. *n*-plet, or spin-(n-1)/2 representation), an eigenbasis to the generators  $t^3$  of SU(2) can be found, such that  $t_3$  acts like  $t^3 = \text{diag}(j, j-1, \ldots, -j)$ , where j = (n-1)/2. Show that if we expand a Higgs *n*-plet  $\Phi$  with a hypercharge of *y* in this basis its components will be eigenstates of  $Q = t^3 + \hat{Y}$ , where  $\hat{Y}\Phi = y\mathbf{1}_{n\times n}\Phi = y\Phi$ . What is the condition on *y* to ensure that there is a neutral component?
- b) In this basis the raising and lowering operators defined from  $t^{\pm} = (t^1 \pm it^2)$  can be described by their action on the orthonormal basis vectors  $e_m$  (labelled by their  $t^3$  eigenvalues m):

$$t^+e_m = \sqrt{j(j+1) - m(m+1)}e_{m+1}, \qquad t^-e_m = \sqrt{j(j+1) - m(m-1)}e_{m-1}.$$

Assume that the hypercharge is such that  $\phi_m$  is neutral and that only this neutral component of  $\Phi$  acquires a vev,  $\vec{v} = ve_m$ . Recall that the covariant derivative reads

$$D_{\mu}\Phi = (\partial_{\mu} - igt^{i}W^{i}_{\mu} - ig'\hat{Y}B_{\mu})\Phi$$

and, show that

$$m_W^2 = \frac{g^2}{2} \vec{v}^{\dagger} (t^+ t^- + t^- t^+) \vec{v} \,.$$

c) Show that this leads to

$$m_W^2 = g^2 v^2 (j(j+1) - m^2).$$

d) With the standard mixing angle  $\tan \theta_W = g'/g$ , derive the mass of Z and use it to show that

$$\rho = \frac{v^2(j(j+1) - y^2)}{2v^2y^2}$$

e) Check that the Higgs doublet with hypercharge 1/2 satisfies the condition that  $\rho = 1$ . What would be the next combination of weak isospin and (rational) hypercharge that can satisfy this bound?

## Problem 2: Integrating out the W-boson [10 Points]

Heavy particles decouple from a given theory when the relevant energy scale of a process is well below their mass. An important example are the  $W^{\pm}$  bosons of the SM, which we will consider in this exercise. The relevant part of the SM Lagrangian reads

$$\mathcal{L}_W = -\frac{1}{2}W^-{}_{\mu\nu}W^{+\mu\nu} + m_W^2 W^+{}_{\mu}W^{-\mu} + \frac{g}{\sqrt{2}}W^+{}_{\mu}J^{-\mu} + \frac{g}{\sqrt{2}}W^-{}_{\mu}J^{+\mu},$$

with the electroweak charged currents acting as source terms,

$$J^+{}_{\mu} = \overline{e}_L \gamma_{\mu} \nu_L, \quad J^-{}_{\mu} = \overline{\nu}_L \gamma_{\mu} e_L.$$

- a) Derive the classical equation of motion for the  $W^{\pm}$  bosons and expand to lowest order in  $\frac{p}{m_W}$ , where p is the momentum of the W (off-shell), i.e.  $\partial_{\mu}W^{\pm}{}_{\nu} = -ip_{\mu}W^{\pm}{}_{\nu}$ .
- b) Solve the equation to lowest order in  $\frac{p}{m_W}$  ignoring non-linear interactions among the W's.
- c) Show that this solution can be used to obtain the Lagrangian of Fermi's effective theory of weak interactions,

$$\mathcal{L}_{\rm Fermi} = -2\sqrt{2}G_F J^+{}_{\mu}J^{-\mu}$$

and determine  $G_F$  in terms of the W mass,  $m_W$ , and the weak gauge coupling, g.