

# Exercises to “Standard Model of Particle Physics II”

Winter 2023/24

Prof. Dr. Manfred Lindner and PD Dr. Werner Rodejohann  
Sheet 03 - November 8, 2023

---

**Tutor:** Juan Pablo Garces   **e-mail:** [juan.garces@mpi-hd.mpg.de](mailto:juan.garces@mpi-hd.mpg.de)

**Lecture webpage:** <https://www.mpi-hd.mpg.de/manitop/StandardModel2/index.html>

**Hand-in of solutions:**

November 15, 2023 - 09:15, Phil. 12, kHS

**Discussion of solutions:**

November 15, 2023 - 11:15, Phil. 12, R105

---

## Problem 1: Goldstone’s theorem [15 Points]

Goldstone’s theorem states: *Every generator of a continuous global symmetry, which does not annihilate the vacuum, corresponds to a massless scalar particle.*

To demonstrate this, consider the Lagrangian  $\mathcal{L}(\phi_i, \partial_\mu \phi_i)$  with  $n$  real scalar fields  $\phi_i$ , invariant under the global transformation  $\phi_i \rightarrow \{\exp(i\theta^a T^a)\}_{ij} \phi_j$ . Defining  $\Phi \equiv (\phi_1, \dots, \phi_n)^T$ , we can write this transformation as  $\Phi \rightarrow \exp(i\theta^a T^a) \Phi$ .

- What are the properties of the  $iT^a$ , if the  $\phi_i$  are real, and if  $\Phi^T \Phi$  is invariant under the global transformation?
- Show that from the conservation of the Noether current, it follows that:

$$\frac{\partial \mathcal{L}}{\partial \phi_i} (iT^a)_{ij} \phi_j + \frac{\partial \mathcal{L}}{\partial (\partial^\mu \phi_i)} (iT^a)_{ij} \partial^\mu \phi_j = 0.$$

- Let  $\mathcal{L} = \frac{1}{2}(\partial_\mu \Phi)^T (\partial^\mu \Phi) - V(\Phi)$ . Show with the help of the first two results that

$$\frac{\partial V}{\partial \phi_i} (iT^a)_{ij} \phi_j = 0.$$

- Let  $\vec{v}$  be the minimum of  $V(\Phi)$ . Show with the results from above that  $M^2 (iT^a) \vec{v} = 0$ , where  $M_{ij}^2$  is given by

$$(M^2)_{ij} \equiv \left. \frac{\partial^2 V}{\partial \phi_i \partial \phi_j} \right|_{\Phi=\vec{v}}.$$

- The matrix  $M^2$  is interpreted as a mass matrix. The broken vacuum  $\vec{v}$  is in general not invariant under the transformations  $\exp(i\theta^a T^a)$ . If  $T^a \vec{v} = 0$ , one has a *Wigner–Weyl realization* of the symmetry, if  $T^a \vec{v} \neq 0$  one has a realization à la *Nambu–Goldstone*. With the result from d) it follows that for all  $a$  with  $T^a \vec{v} \neq 0$  the mass matrix has an eigenvector with eigenvalue 0.

As an explicit example, consider the Lagrangian

$$\mathcal{L} = \frac{1}{2}(\partial_\mu \Phi)^T (\partial^\mu \Phi) - V(\Phi^T \Phi)$$

with the potential

$$V(\Phi^T \Phi) = \frac{1}{2} \mu^2 \Phi^T \Phi + \frac{1}{4} \lambda (\Phi^T \Phi)^2,$$

where  $\Phi = (\phi_1, \phi_2, \phi_3)^T$  is a triplet of real and scalar particles. The coupling obeys  $\lambda > 0$ .

- (i) Show that  $\mathcal{L}$  is invariant under an SU(2) transformation  $\Phi \rightarrow \exp(i\theta^a T^a) \Phi$ . The three generators  $T^a$  of SU(2) are written here as  $(T^a)_{ij} = -i\epsilon_{aij}$  (this is called the adjoint representation of SU(2)).
- (ii) For  $\mu^2 < 0$  the potential has a minimum  $\vec{v}^T = (0, 0, v)$  with  $v \neq 0$ . Derive a relation between  $v$ ,  $\mu^2$ , and  $\lambda$ . Show that the vacuum state  $\vec{v}$  is invariant under a transformation generated by  $T^3$ , but not  $T^1$  and  $T^2$ . That is, the vacuum state does *not* possess the symmetry of the Lagrangian (spontaneous symmetry breaking). Why do we speak in this case of spontaneous symmetry breaking from SU(2) to U(1)?
- (iii) The scalar fields will be expanded around the minimum,  $\Phi^T \equiv (\phi_1, \phi_2, v + \phi_3)$ , i.e. we treat  $\phi_3$  as a field fluctuating around  $v$  instead of 0. Show by inserting into the above Lagrangian that there is one massive scalar field and two massless Goldstone bosons. What is the mass of  $\phi_3$ ? Convince yourself that the matrix  $M^2$  defined in d) is really the mass matrix.

**Problem 2:** *Noether's theorem revised* [5 Points]

Consider a Lagrange density of the general form  $\mathcal{L} = \mathcal{L}_0(u, \partial_\mu u)$ , invariant under some symmetry in such a way that, according to the Noether's theorem, there is a conserved current

$$j^\mu = \frac{\partial \mathcal{L}}{\partial(\partial_\mu u)} \delta u.$$

Now consider to extend the theory,  $\mathcal{L}_0 \rightarrow \mathcal{L} = \mathcal{L}_0 + \mathcal{L}_1$  to include  $\mathcal{L}_1(u, \partial_\mu u)$  that depends on the same fields as  $\mathcal{L}$  but is not invariant under the previously considered symmetry. The total current will not be conserved in this new case. Show that also the previously conserved Noether's current corresponding to the invariance of  $\mathcal{L}$  will not be conserved anymore.