

# Exercises to “Standard Model of Particle Physics II”

Winter 2022/23

Prof. Dr. Manfred Lindner and Dr. Werner Rodejohann

Sheet 02 - October 26, 2022

---

**Tutor:** Sophie Klett   **e-mail:** [sophie.klett@mpi-hd.mpg.de](mailto:sophie.klett@mpi-hd.mpg.de)

**Lecture webpage:** <https://www.mpi-hd.mpg.de/manitop/StandardModel2/index.html>

**Hand-in of solutions:**

November 2, 2022 - 09:15, Phil. 12, KHS

**Discussion of solutions:**

November 2, 2022 - 11:15, Phil. 12, R105

---

## Problem 1: $SU(N)$ [10 Points]

Let  $U \in SU(N)$ , i.e.  $\det U = 1$  and  $U^\dagger U = 1$ . Any element of  $SU(N)$  can be written as  $U = \exp(-i\theta^a T^a)$ , where the  $T^a$  are generators of the group with normalization  $\text{Tr}(T^a T^b) = \frac{1}{2}\delta^{ab}$ .

- a) Show that the  $T^a$  are traceless hermitian matrices.
- b) How many linear independent generators are there?
- c) The structure constants,  $d_{abc}$  and  $f_{abc}$ , are defined through

$$[T^a, T^b] = if_{abc} T^c, \quad \{T^a, T^b\} = \frac{1}{N}\delta_{ab} + d_{abc} T^c.$$

Show that

$$\begin{aligned} \text{Tr}(T^a T^b T^c) &= \frac{1}{4}(d_{abc} + if_{abc}), \\ \left[ \sum_a T^a T^a, T^b \right] &= 0, \\ [T^a, [T^b, T^c]] + [T^c, [T^a, T^b]] + [T^b, [T^c, T^a]] &= 0. \end{aligned}$$

- d) Show that the structure constants form a representation of  $SU(N)$ , i.e. take  $(T^a)_{bc} = -if_{abc}$  as a generator. This is the so-called adjoint representation.
- e) Calculate the  $f_{abc}$  for

$$T^a = \frac{\sigma^a}{2},$$

where  $a$  runs from 1 to 3, and  $\sigma^a$  are the Pauli matrices:

$$\sigma^1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma^2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma^3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}.$$

**Problem 2: From QED to QCD [10 points]**

In **Problem 2** on the last sheet we showed that in QED  $F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$  is *invariant* under the U(1) gauge transformation  $A_\mu \rightarrow A_\mu + \frac{1}{e}\partial_\mu\alpha(x)$ . For QCD, this gauge transformation is generalised to  $A_\mu \rightarrow U(x) \left( A_\mu + \frac{i}{g}\partial_\mu \right) U(x)^\dagger$ , while the quarks obey  $\psi \rightarrow U(x)\psi$ , where  $\psi$  carries one SU( $N$ ) index,  $A_\mu \equiv A_\mu^a T^a$  is now matrix-valued, and  $U(x) \in \text{SU}(N)$ .

- a) Using the U(1) covariant derivative  $D_\mu = \partial_\mu - ieA_\mu$ , show that

$$[D_\mu, D_\nu] \psi = -ie F_{\mu\nu} \psi,$$

where  $\psi$  is the electron field.

- b) For a set of parameters  $\alpha^a$ , the transformation matrix of QCD is  $U(x)_{ij} = \exp(i\alpha^a T^a)_{ij}$ . Using the infinitesimal version of the gauge transformation of the gluon field,  $A_\mu \rightarrow A_\mu + \frac{1}{g}(\partial_\mu\alpha^a)T^a + i[\alpha^a T^a, A_\mu^b T^b]$ , show that

$$D_\mu \psi \rightarrow (1 + i\alpha^a T^a) D_\mu \psi,$$

with the QCD covariant derivative  $D_\mu = \partial_\mu - igA_\mu^a T^a$ .

- c) In analogy to QED, we can define the QCD field strength matrix  $F_{\mu\nu} = F_{\mu\nu}^a T^a$  via  $[D_\mu, D_\nu] \psi = -igF_{\mu\nu}^a T^a \psi$ , which is no longer invariant. Compute  $F_{\mu\nu}^a$ .
- d) Show that the QCD Lagrangian is gauge-invariant.

$$\mathcal{L}_{\text{QCD}} = \bar{\psi}(i\not{D} - m)\psi - \frac{1}{4}F_{\mu\nu}^a F^{a\mu\nu}.$$

*Hint:* This is a one-line proof if you use b) and c).