
Exercises for Neutrino Physics: Theory and Experiment

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Sheet 12: Theory

1. Majorana Neutrinos [10 Points]

The Lagrangian for the coupling of a fermion pair f with the Z -Boson is given by

$$\mathcal{L} = \frac{g}{2 \cos \theta_W} \bar{f} \gamma^\mu (v_f - a_f \gamma_5) f Z_\mu \quad (1)$$

For neutrinos we have $a_\nu = v_\nu = 1/2$

- On sheet 9 we discussed the decay of the Z -boson. Calculate the decay width for $Z \rightarrow \bar{\nu} \nu$ in the Standard Model but keeping a possible neutrino mass in the expression.
- Neutrinos could be Majorana particles, which obey the relation $\nu^C = \nu$. The superscript C denotes charge conjugation

$$\nu^C = C \bar{\nu}^T, \quad (2)$$

with $C = i\gamma_2\gamma_0$ in the Dirac basis. Show the following properties:

$$-C = C^T = C^{-1} = -C^* = C^\dagger, \quad (3)$$

$$C^{-1} \gamma_\mu C = -\gamma_\mu^T, \quad (4)$$

$$C^{-1} \gamma_5 C = \gamma_5^T, \quad (5)$$

$$\overline{\Psi^C} = -\Psi^T C^{-1}, \quad (6)$$

$$(\Psi_L)^C = (\Psi^C)_R, \quad (7)$$

where $(\Psi^C)_L = P_L (\Psi^C)$.

- Show that for Majorana neutrinos the vector current $\bar{\nu} \gamma_\mu \nu$ vanishes. What happens with $\bar{\nu} \gamma_5 \nu$, $\bar{\nu} \gamma_\mu \gamma_5 \nu$ and $\bar{\nu} [\gamma_\mu, \gamma_\nu] \nu$?
- Using the previous result calculate the decay width $Z \rightarrow \nu \nu$ for Majorana neutrinos and compare with a)
(Hint 1: Note that due to the Majorana properties the final state particles are indistinguishable from each other resulting in a factor of $1/2$.
Hint 2: Compared to the Dirac case the Majorana axial-vector contribution gains a factor of 2)

2. Seesaw [10 Points]

The Higgs mechanism generates Dirac masses for the active neutrinos when right-handed sterile neutrinos N_R are introduced through the following term:

$$\mathcal{L}_{\text{Dirac}} = -\bar{\nu}_L M_D N_R + \text{h.c.} \quad (8)$$

where $(\nu_L = \nu_L^1, \nu_L^2, \nu_L^3)^T$ is the column vector of the active neutrinos and N_R the corresponding vector for the sterile neutrinos. The matrix M_D is (in general) a complex 3×3 matrix. Additionally, the right-handed sterile neutrinos can have a Majorana mass with the Lagrangian density

$$\mathcal{L}_{\text{Majorana}} = -\frac{1}{2} \overline{(N_R)^C} M_R N_R + \text{h.c.} \quad (9)$$

where M_R is a symmetric 3×3 matrix and $\Psi^C = C\bar{\Psi}^T$ for a general Dirac spinor Ψ and the charge conjugation matrix $C = i\gamma^2\gamma^0$.

Let m_D and m_R denote the mass scale of M_D and M_R , respectively. Suppose the entries of M_R are much larger than the ones of M_D ($m_R \gg m_D$)

- (a) Show that it is possible to rewrite the whole mass matrix in the flavour basis in the following way:

$$\mathcal{L}_{\text{mass}} \equiv \mathcal{L}_{\text{Dirac}} + \mathcal{L}_{\text{Majorana}} = -\frac{1}{2} \overline{\Psi^C} M \Psi + \text{h.c.} \quad (10)$$

with

$$\Psi = \begin{pmatrix} (\nu_L)^C \\ N_R \end{pmatrix} \quad \text{and} \quad M \equiv \begin{pmatrix} 0 & M_D \\ M_D^T & M_R \end{pmatrix} \quad (11)$$

For the solution prove and use the identity

$$\overline{\nu_L} M_D N_R = \overline{(N_R)^C} M_D^T (\nu_L)^C . \quad (12)$$

- (b) Using the (unitary) transformation $\Psi = U\xi$ with $U = \begin{pmatrix} 1 & \rho \\ -\rho^\dagger & 1 \end{pmatrix}$ (change of basis) it is possible to convert the 6×6 matrix M into a block diagonal form, i.e. that it takes on the following form

$$U^T M U \simeq \begin{pmatrix} M_1 & 0 \\ 0 & M_2 \end{pmatrix}, \quad (13)$$

with symmetric 3×3 matrices M_1, M_2 . The matrix ρ in the transformation matrix U is assumed to be proportional to the scale m_R^{-1} and terms of order m_R^{-2} (and smaller) can be neglected in the calculation.

Determine ρ, M_1 and M_2 from Eq.(13). What is the connection between the fields χ_1, χ_2 and the original fields ν_L, N_R ? (Hint: You may assume that M_R is invertible.)

- (c) Where does the name "seesaw" come from?