

# Neutrino Masses, Interactions, and Asymmetries from Cosmology

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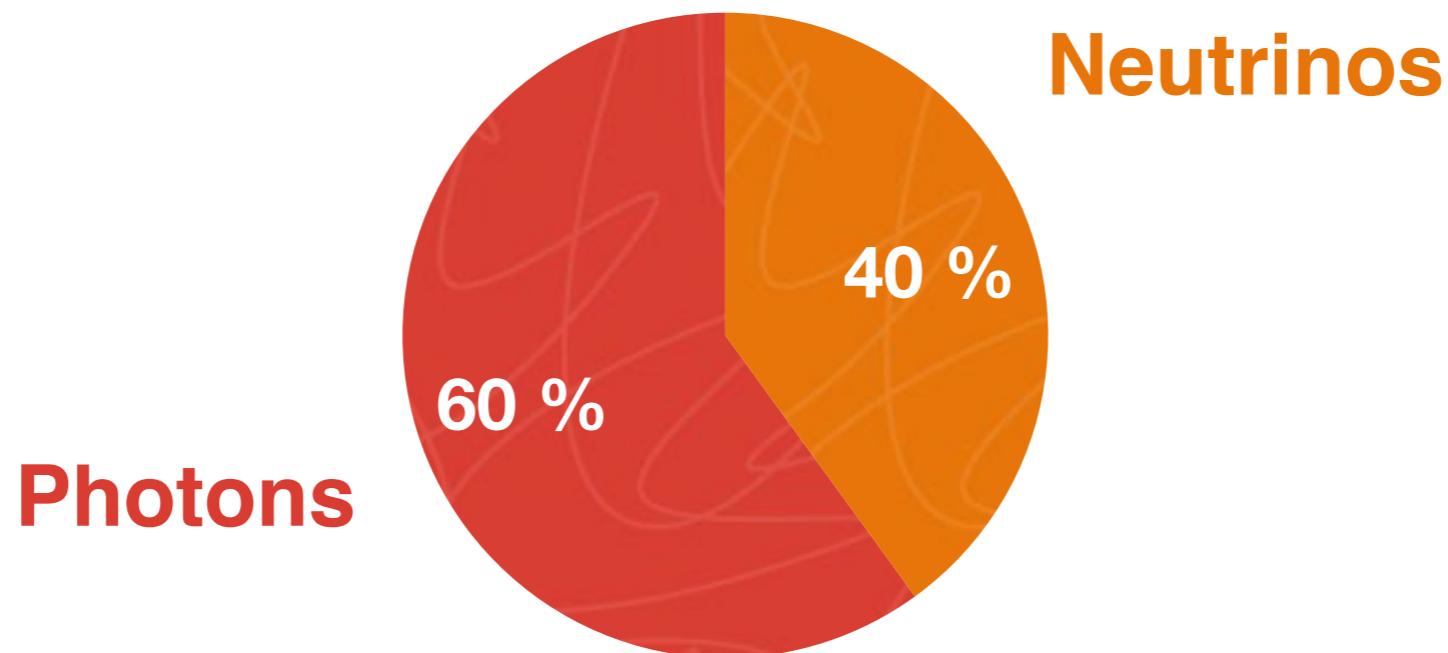
**Max Planck Institute for Kernphysik  
Heidelberg  
12-12-2022**

# Motivation

- Neutrino masses are the only laboratory evidence of physics beyond the Standard Model

**Use neutrinos to understand open problems in cosmology**

- Neutrinos are ubiquitous in Cosmology



**Use cosmological data to understand their properties**

# Topics Covered

## Topics covered:

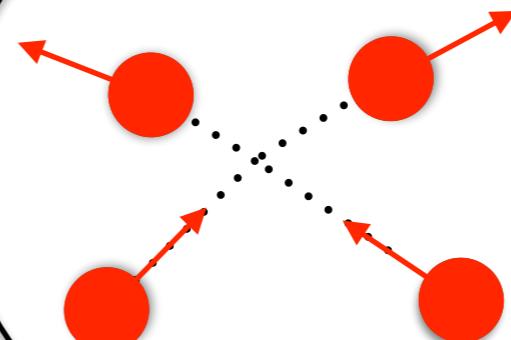
CNB

$$N_{\text{eff}}$$

Neutrino Masses

$$\sum m_\nu$$

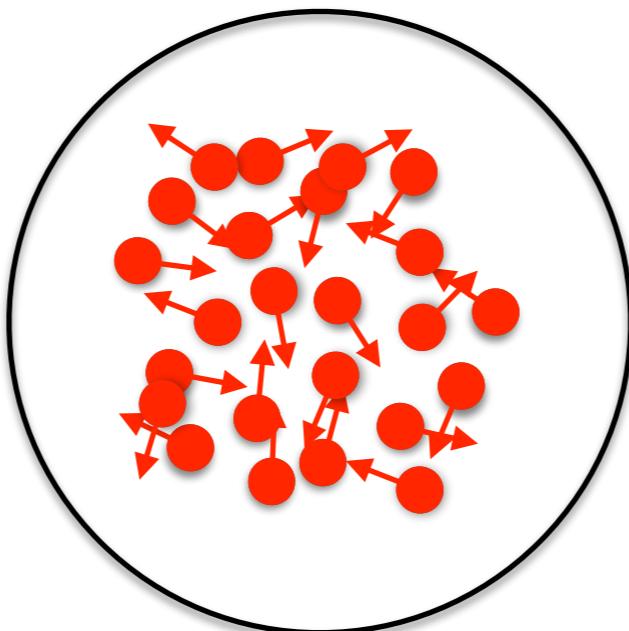
Neutrino Interactions



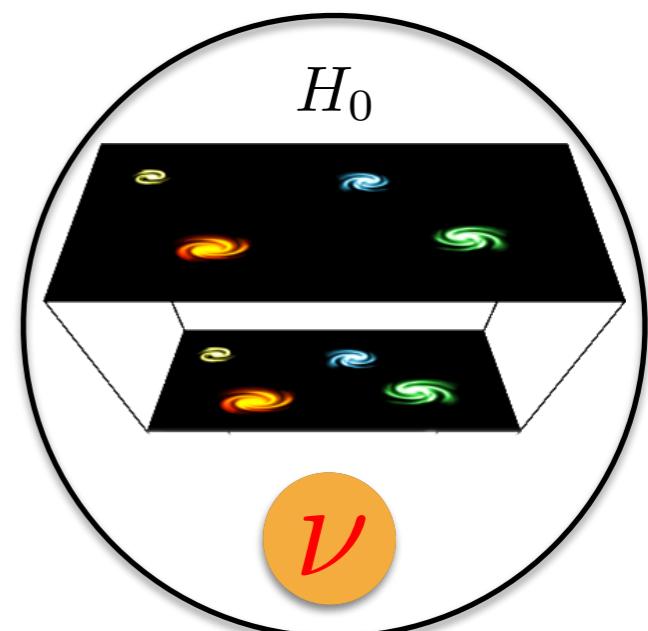
Neutrino Asymmetries

$$\eta_\nu = (n_\nu - n_{\bar{\nu}})/n_\gamma$$

CNB Detection



Neutrinos and H0



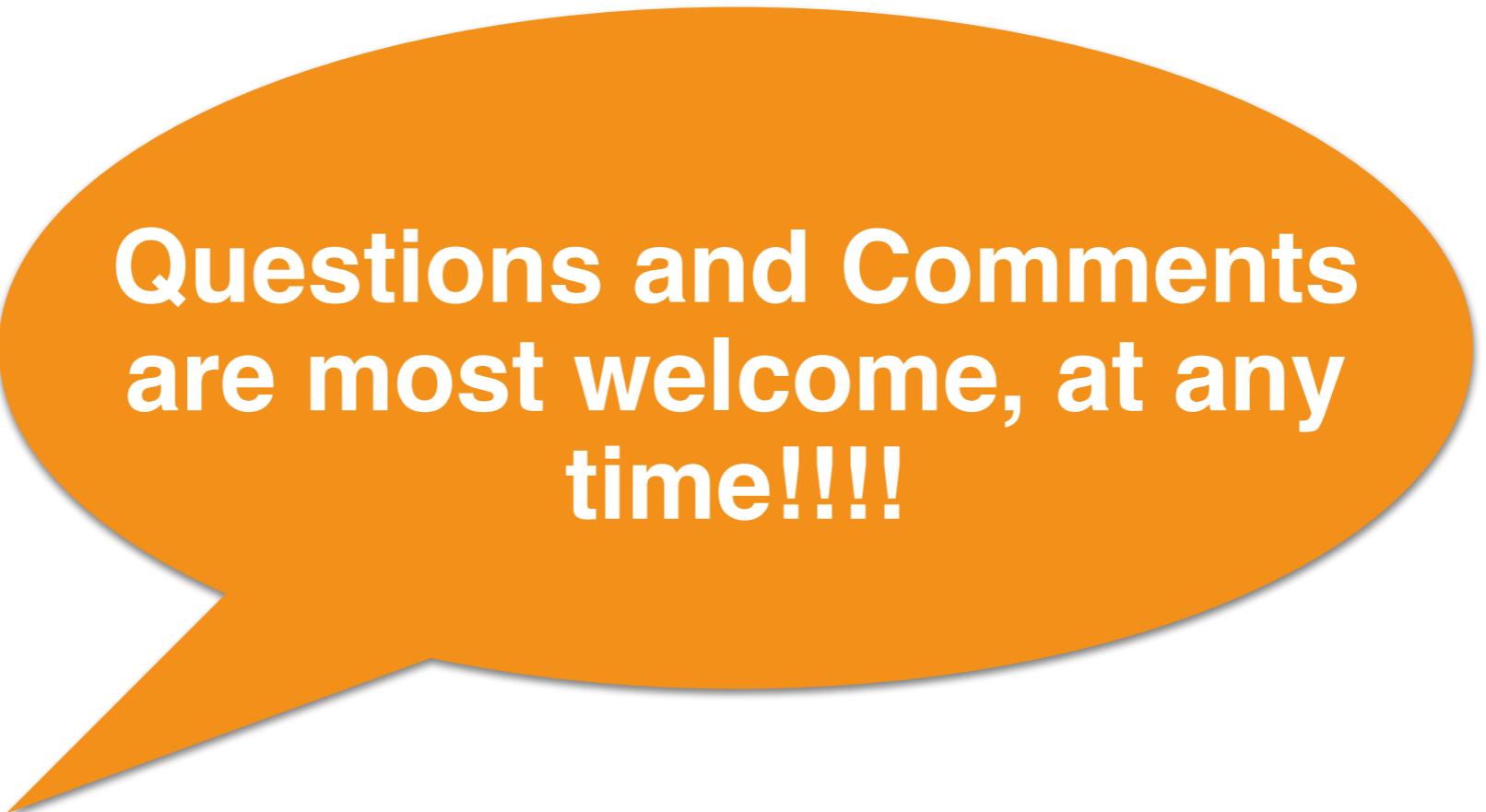
Topics not covered:  
(but happy to discuss offline)

# Outline

- 1) Neutrinos in  $\Lambda$ CDM**
- 2) Neutrino Masses in and beyond  $\Lambda$ CDM**
- 3) Neutrinos interactions and CMB observations**
- 4) A Hint for a Large Primordial Lepton Asymmetry?**
- 5) Conclusions and Outlook**

# Set Up

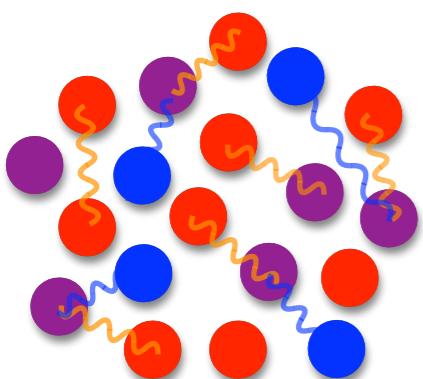
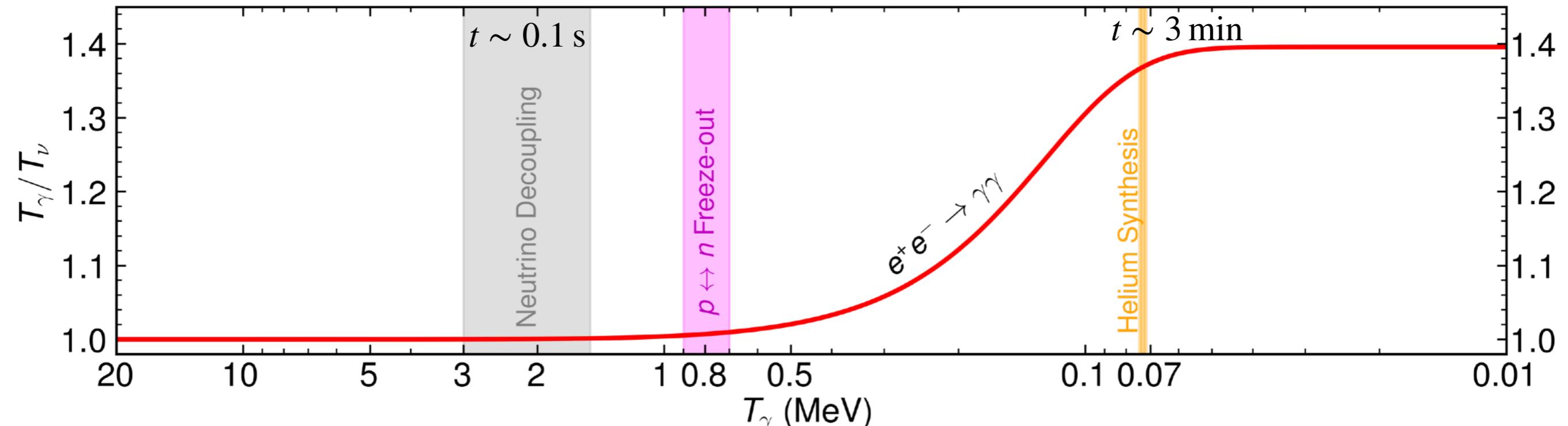
**Unlike neutrinos, I like to interact 😊**



**Questions and Comments  
are most welcome, at any  
time!!!!**

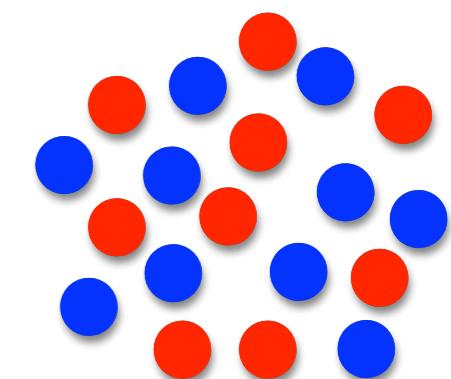
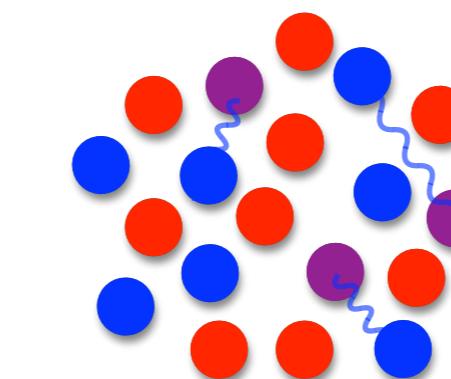
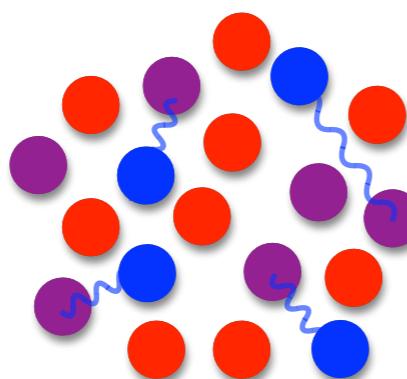
# Neutrino Decoupling

Evolution in the Standard Model



$$e^+ e^- \leftrightarrow \bar{\nu}_i \nu_i$$

$$e^\pm \nu_i \leftrightarrow e^\pm \nu_i$$



**Neutrinos**



**Electrons**



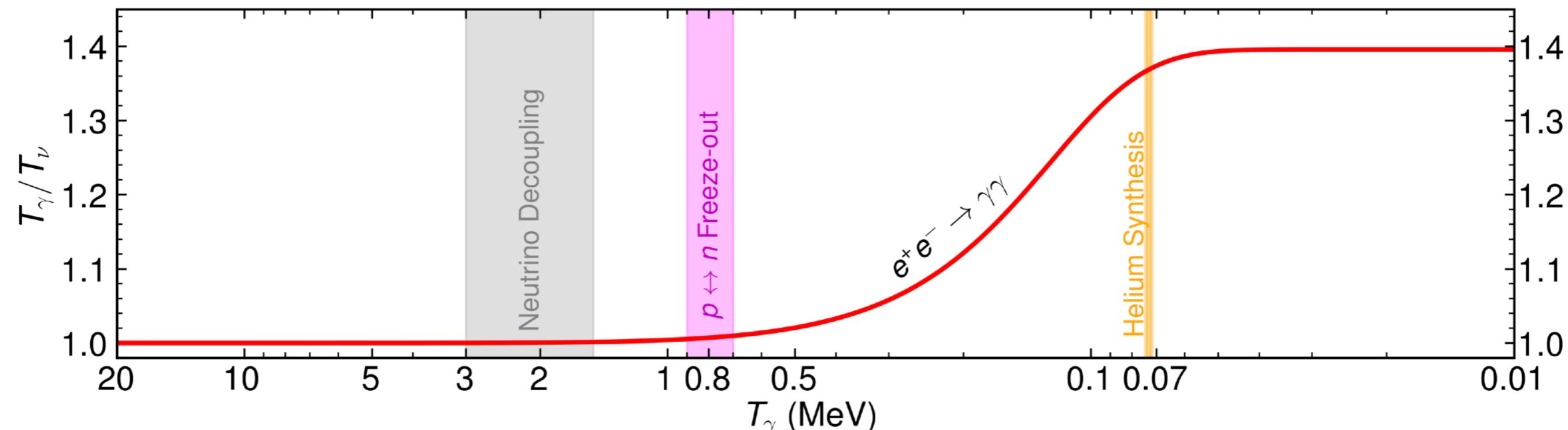
**Photons**



**Z-W (off-shell)**

# Cosmic Neutrino Background

Evolution in the Standard Model



- $N_{\text{eff}} \equiv \frac{8}{7} \left( \frac{11}{4} \right)^{4/3} \left( \frac{\rho_{\text{rad}} - \rho_{\gamma}}{\rho_{\gamma}} \right)$

$$N_{\text{eff}} = 3 \left( \frac{1.4 T_{\nu}}{T_{\gamma}} \right)^4$$

- $N_{\text{eff}}^{\text{SM}} = 3.044(1)$

de Salas & Pastor 1606.06986  
 Bennett, Buldgen, Drewes & Wong 1911.04504  
 Escudero 2001.04466

Akita & Yamaguchi 2005.07047  
 Froustey, Pitrou & Volpe 2008.01074  
 Gariazzo, de Salas, Pastor et al. 2012.02726  
 Hansen, Shalgar & Tamborra 2012.03948

## Relic Neutrino Decoupling

$$t \sim 0.1 \text{ s}$$

$$T_{\nu} \sim 2 \text{ MeV}$$

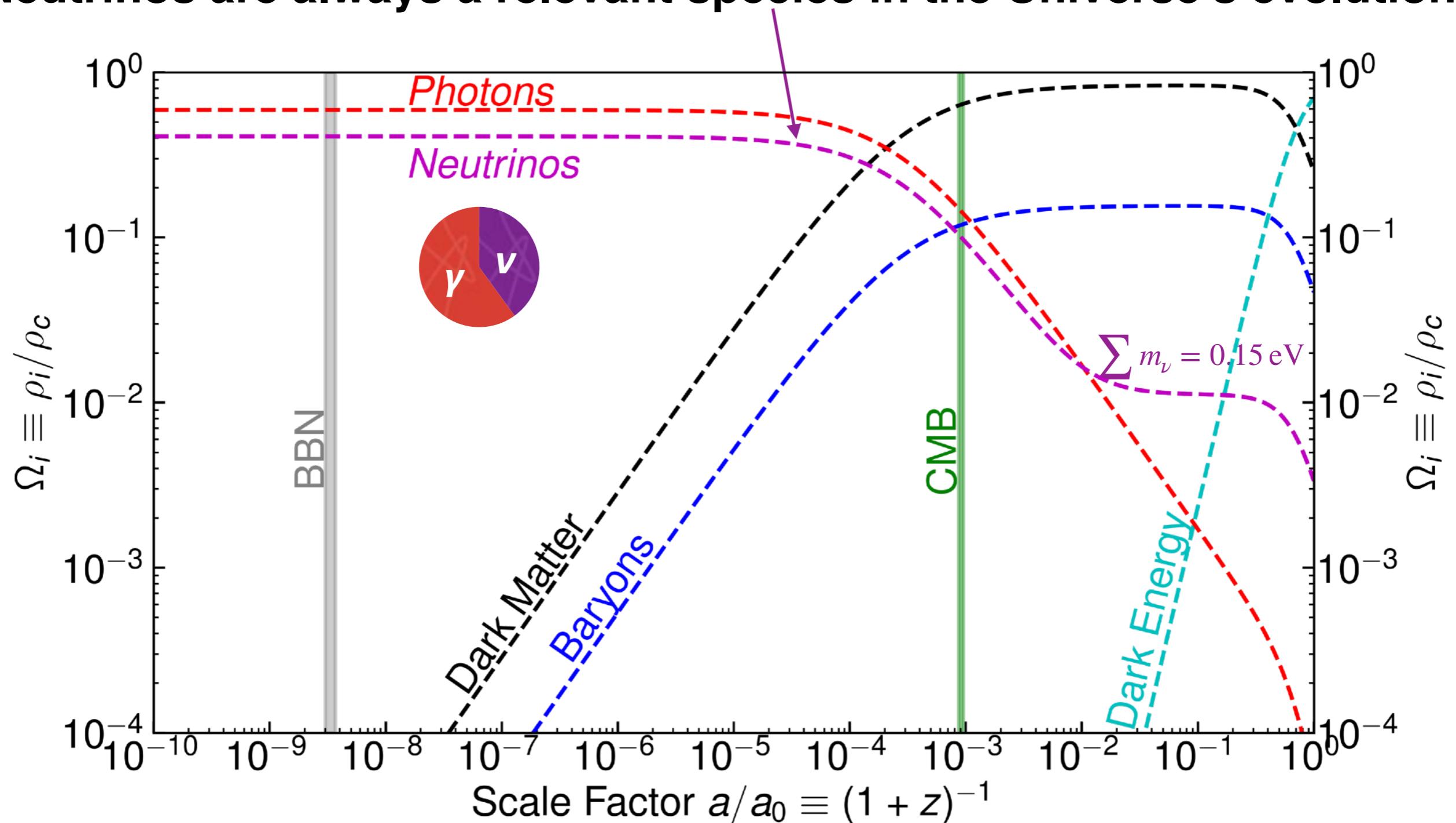
## Why is it not 3?

- Some  $e^+e^-$  heating
- Non-instantaneous decoupling
- QED thermal corrections
- Neutrino Oscillations

Excellent review  
 by Dolgov hep-ph/0202122

# Neutrino Evolution

Neutrinos are always a relevant species in the Universe's evolution



$$\text{Non-Rel: } z_\nu^{\text{non-rel}} \approx 110 \frac{m_\nu}{0.06 \text{ eV}}$$

$$\text{Hot DM: } \Omega_\nu h^2 = \sum m_\nu / (93.14 \text{ eV})$$

# Evidence for Cosmic Neutrinos

## Big Bang Nucleosynthesis

Current measurements are consistent with the SM picture\*

● H ~ 75%



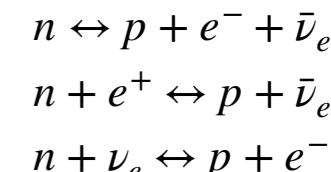
${}^4\text{He} \sim 25\%$



D ~ 0.005%

This implies that neutrinos should have been present:

1) It is impossible to have successful BBN without neutrinos.  
They participate in  $p \leftrightarrow n$  conversions up to  $T \gtrsim 0.7 \text{ MeV}$



2) Neutrinos contribute to the expansion rate  $H \propto \sqrt{\rho}$

By comparing predictions against observations, we know:

$$N_{\text{eff}}^{\text{BBN}} = 2.86 \pm 0.28$$

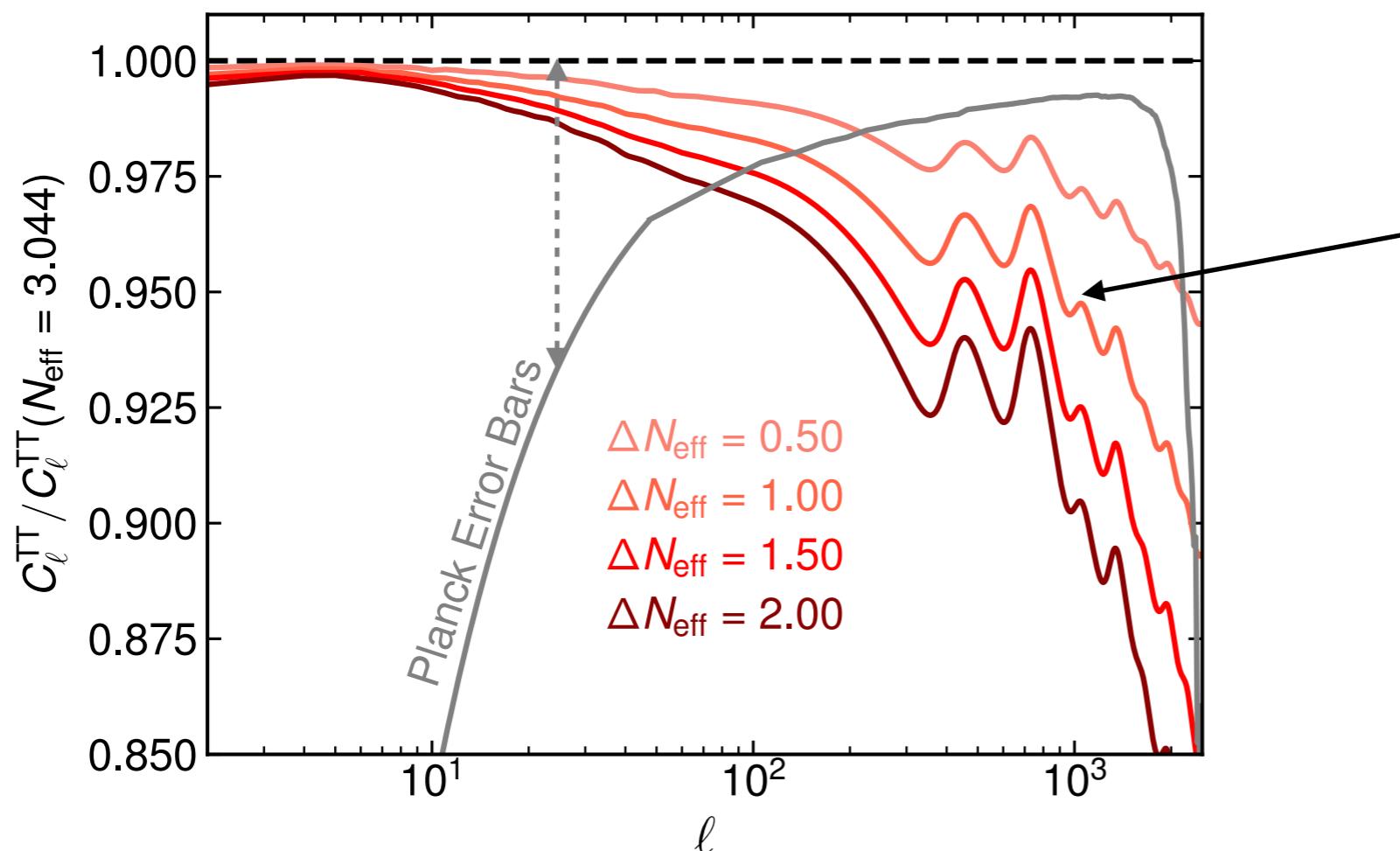
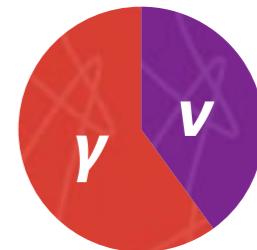
see e.g. Pisanti et al. 2011.11537

\*A very recent measurement of the primordial Helium abundance  $3\sigma$  smaller than the SM expectation could indicate the presence of a large lepton asymmetry:  $L_{\nu_e} \sim 10^{-3}$  (but which would not alter the  $N_{\text{eff}}^{\text{BBN}}$  bound significantly)

# Evidence for Cosmic Neutrinos

## Cosmic Microwave Background Why?

Ultra-relativistic neutrinos represent a large fraction of the energy density of the Universe,  $H \propto \sqrt{\rho}$



$N_{\text{eff}}$  is constrained by the high- $\ell$  multipoles,  
i.e. Silk damping

$$N_{\text{eff}}^{\text{CMB+BAO}} = 2.99 \pm 0.17$$

Planck 2018 1807.06209

# Evidence for Cosmic Neutrinos

- Current constraints

BBN

$$N_{\text{eff}}^{\text{BBN}} = 2.86 \pm 0.28$$

Pisanti et al. 2011.11537

Planck+BAO

$$N_{\text{eff}}^{\text{CMB}} = 2.99 \pm 0.17$$

Planck 2018, 1807.06209

- Standard Model prediction:  $N_{\text{eff}}^{\text{SM}} = 3.044(1)$

- Data is in excellent agreement with the Standard Model prediction
- This provides strong (albeit indirect) evidence for the Cosmic Neutrino Background

## Implications:

1) Stringent constraint on many BSM scenarios

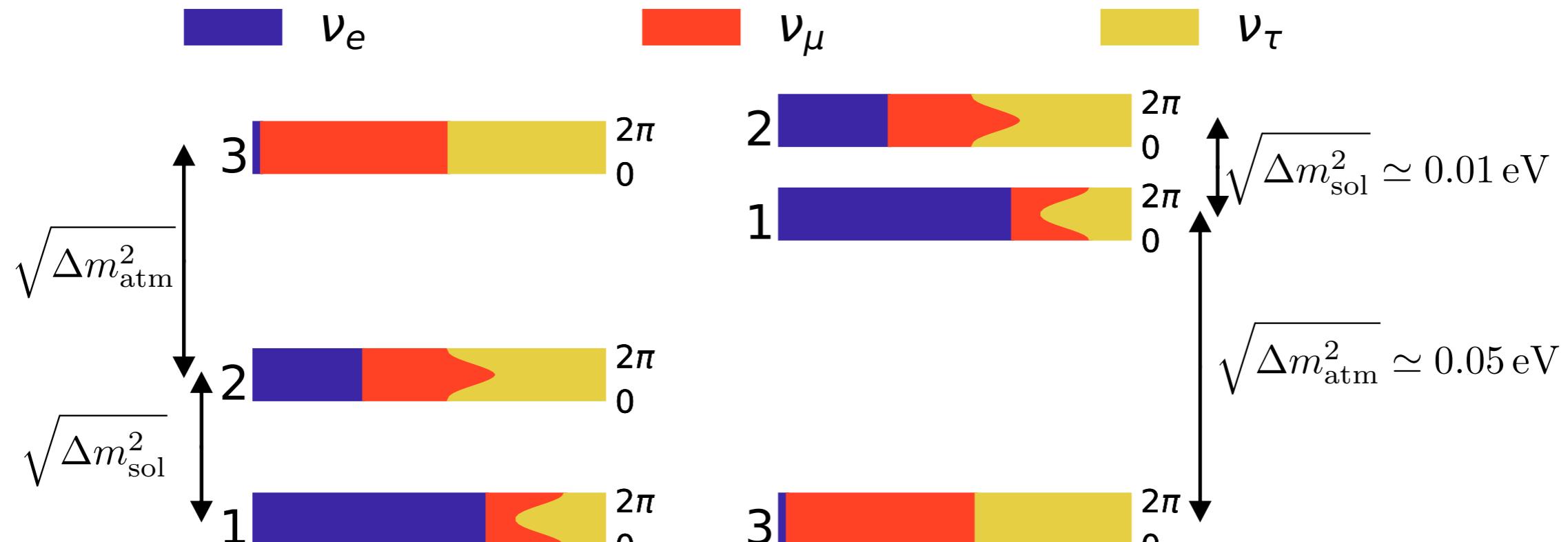
Sterile neutrinos, Goldstone bosons, hidden sector particles, GW ...

see [2006.16182](#) Vaskonen & Tenkanen et al. (Escudero & Poulin)

2) We can use cosmological data to test neutrino properties

# Neutrino Properties

Figure from de Salas et al. 1806.11051



**Normal**

$$\sum m_\nu \gtrsim 0.06 \text{ eV}$$

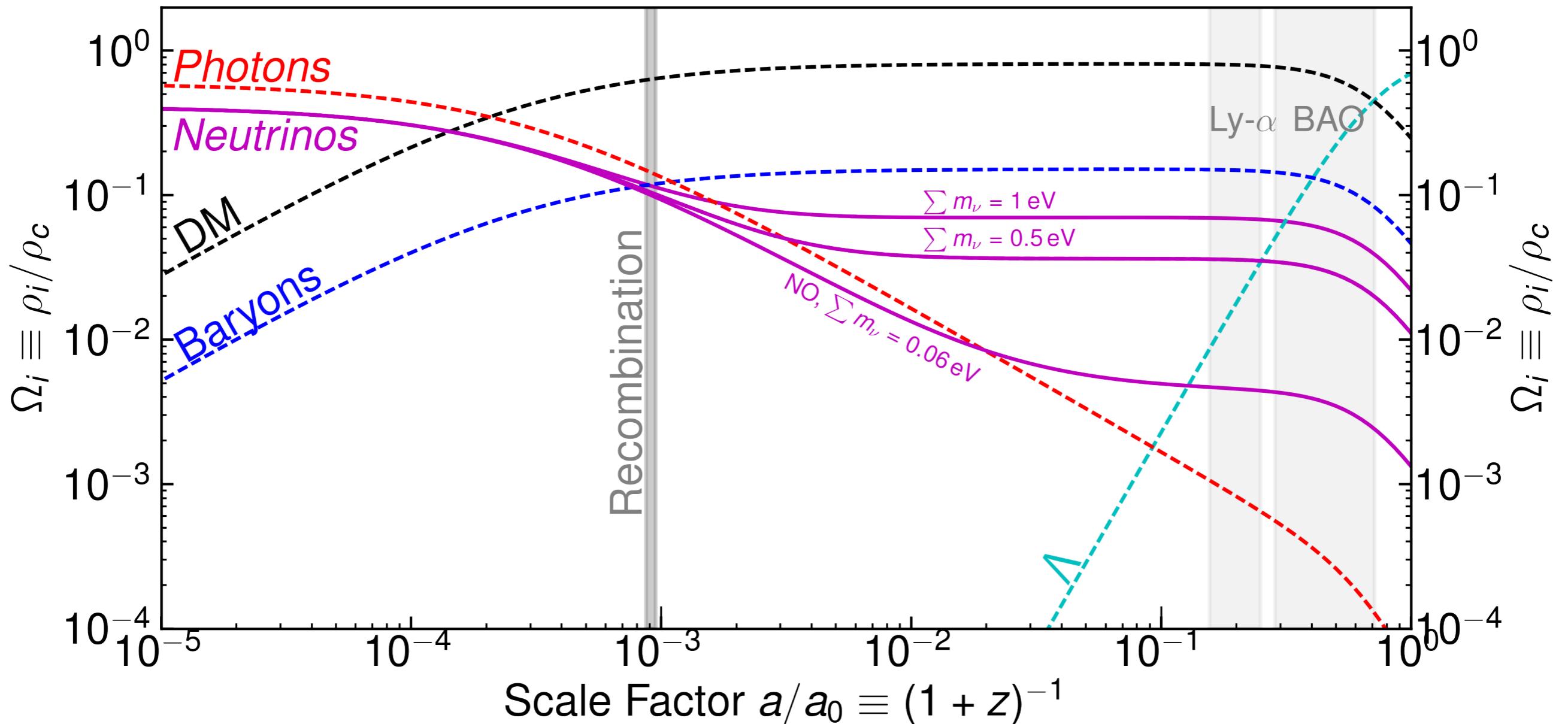
**Inverted**

$$\sum m_\nu \gtrsim 0.10 \text{ eV}$$

- Mass differences and mixings measured with high precision
- What is  $\delta_{\text{CP}}$  and what is the mass ordering? **Neutrino Oscillations**
- Are Neutrinos Dirac or Majorana Particles? **0v2 $\beta$  Experiments**
- What is the neutrino mass scale? i.e.  $\sum m_\nu$ ? i.e.  $m_{\text{lightest}}$ ? **Cosmology**

# Neutrino Masses in Cosmology

- 1) Massive neutrinos modify the expansion history



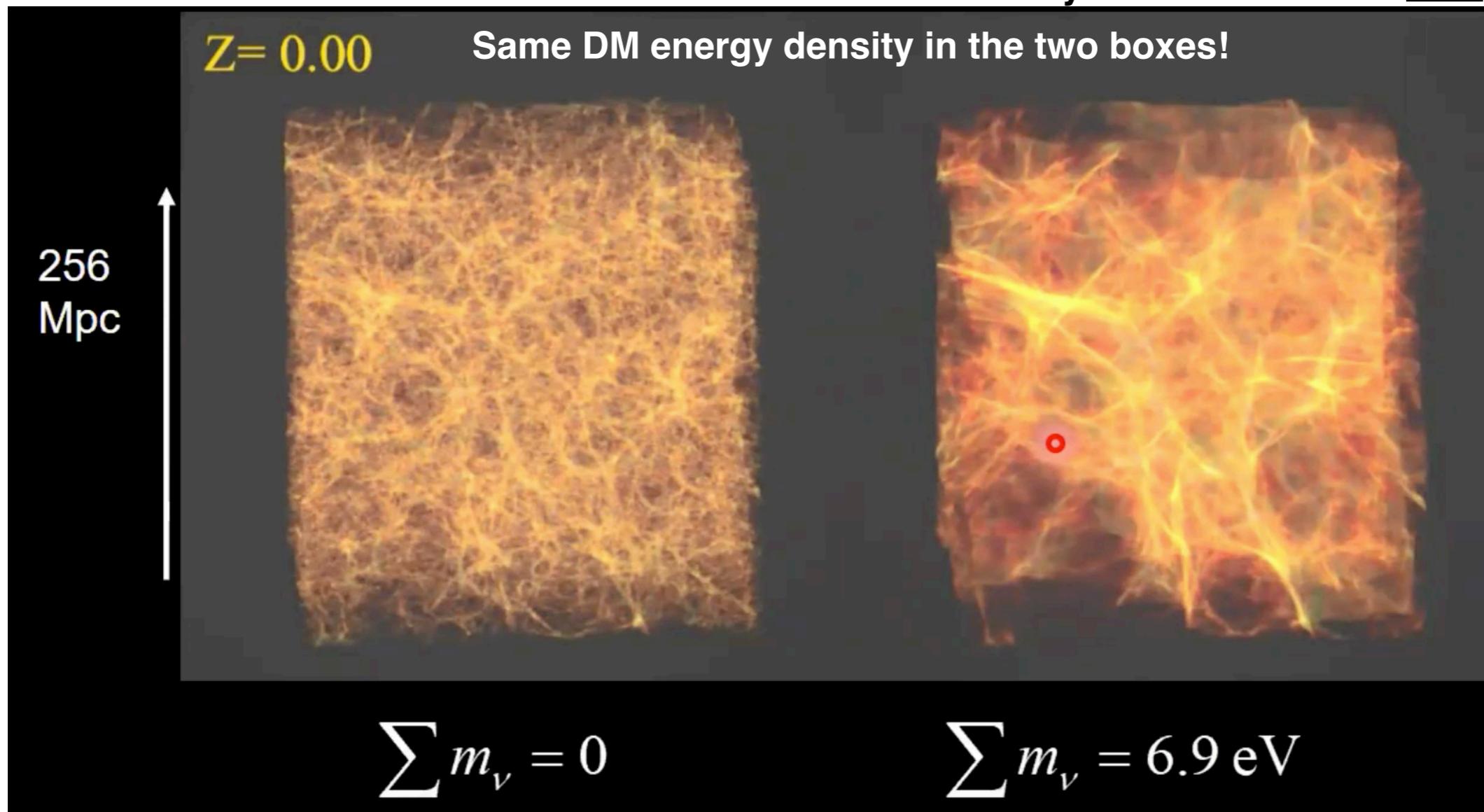
**Non-Rel:**  $z_{\nu}^{\text{non-rel}} \simeq 200 \frac{m_{\nu}}{0.1 \text{ eV}}$

**Hot DM:**  $\Omega_{\nu} h^2 = \sum m_{\nu} / (93.14 \text{ eV})$

# Neutrino Masses in Cosmology

- 2) Massive neutrinos suppress the growth of structure

Taken from a talk by Steen Hannestad [Link](#).



This happens because neutrinos travel very fast and therefore cannot fall in gravitational potentials. The effect of this smoothing is proportional to  $\Omega_\nu$

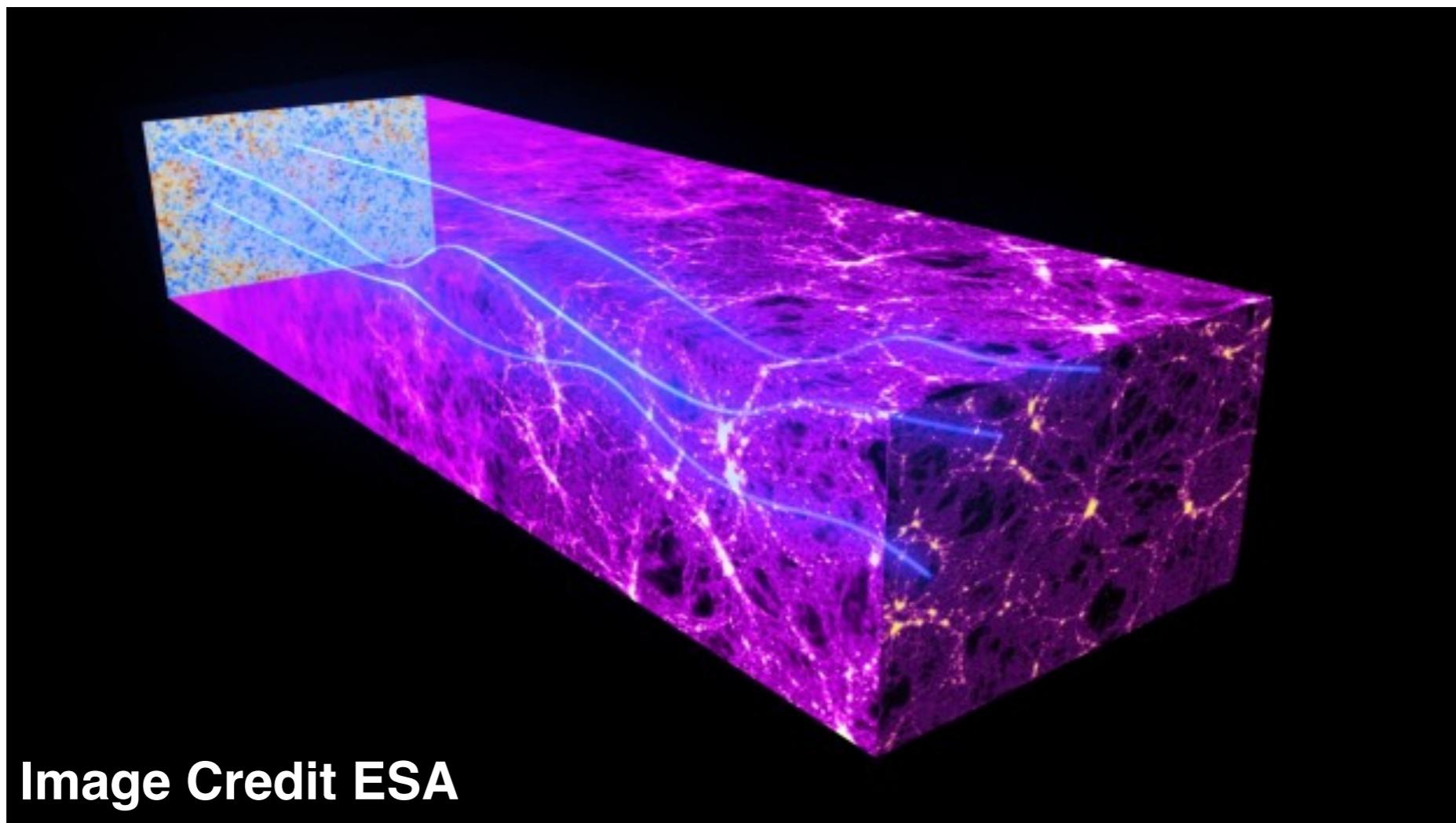
# Neutrino Masses in Cosmology

## Cosmic Microwave Background Anisotropies

Neutrinos of  $m_\nu < 0.5 \text{ eV}$  become non-relativistic after recombination.

That means that their effect on the anisotropies is somewhat small!

The most relevant impact is through the effect of gravitational lensing:

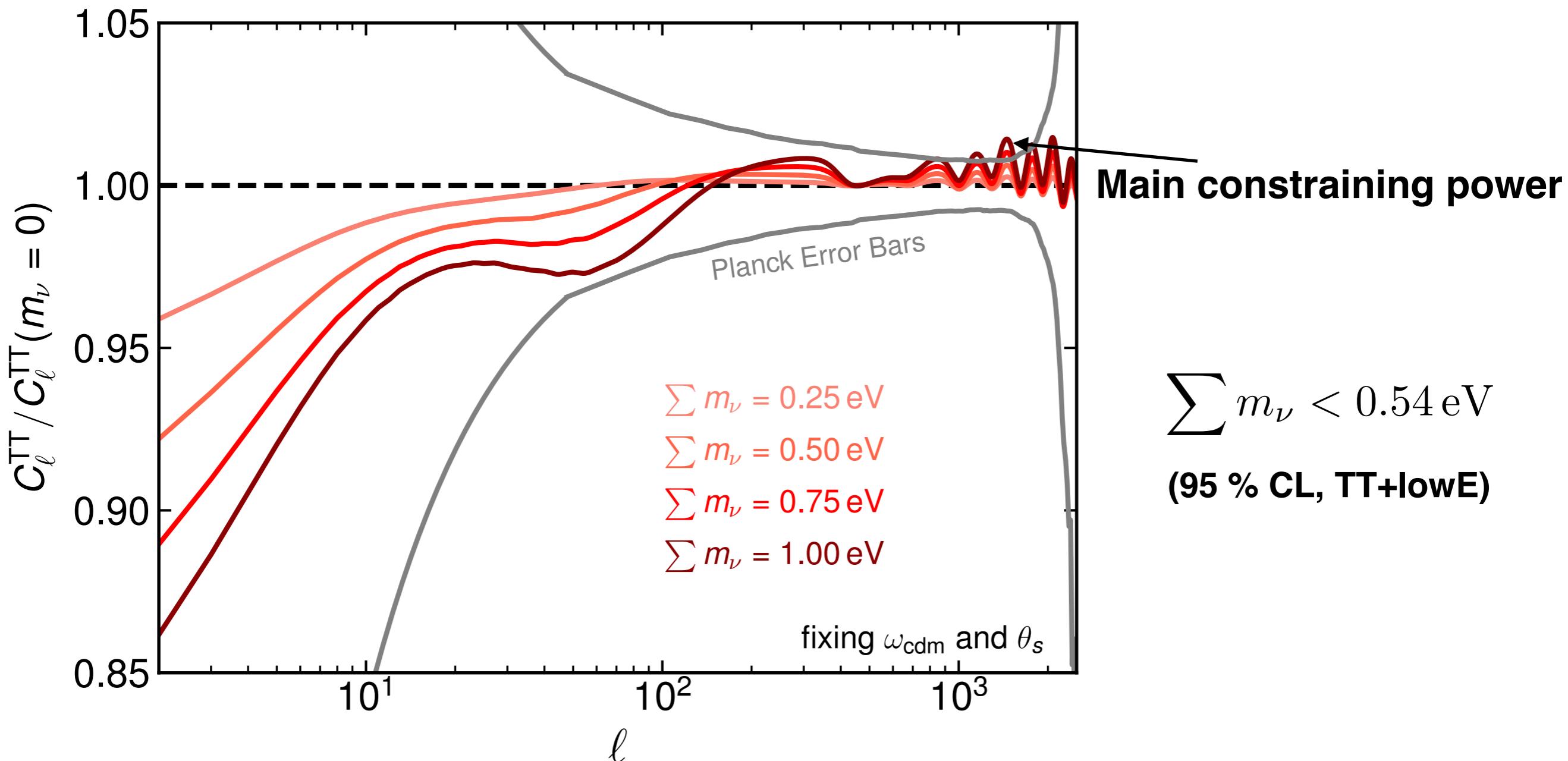


**The larger the neutrino mass the less is the CMB light lensed!**

# Neutrino Masses in Cosmology

## Cosmic Microwave Background Anisotropies

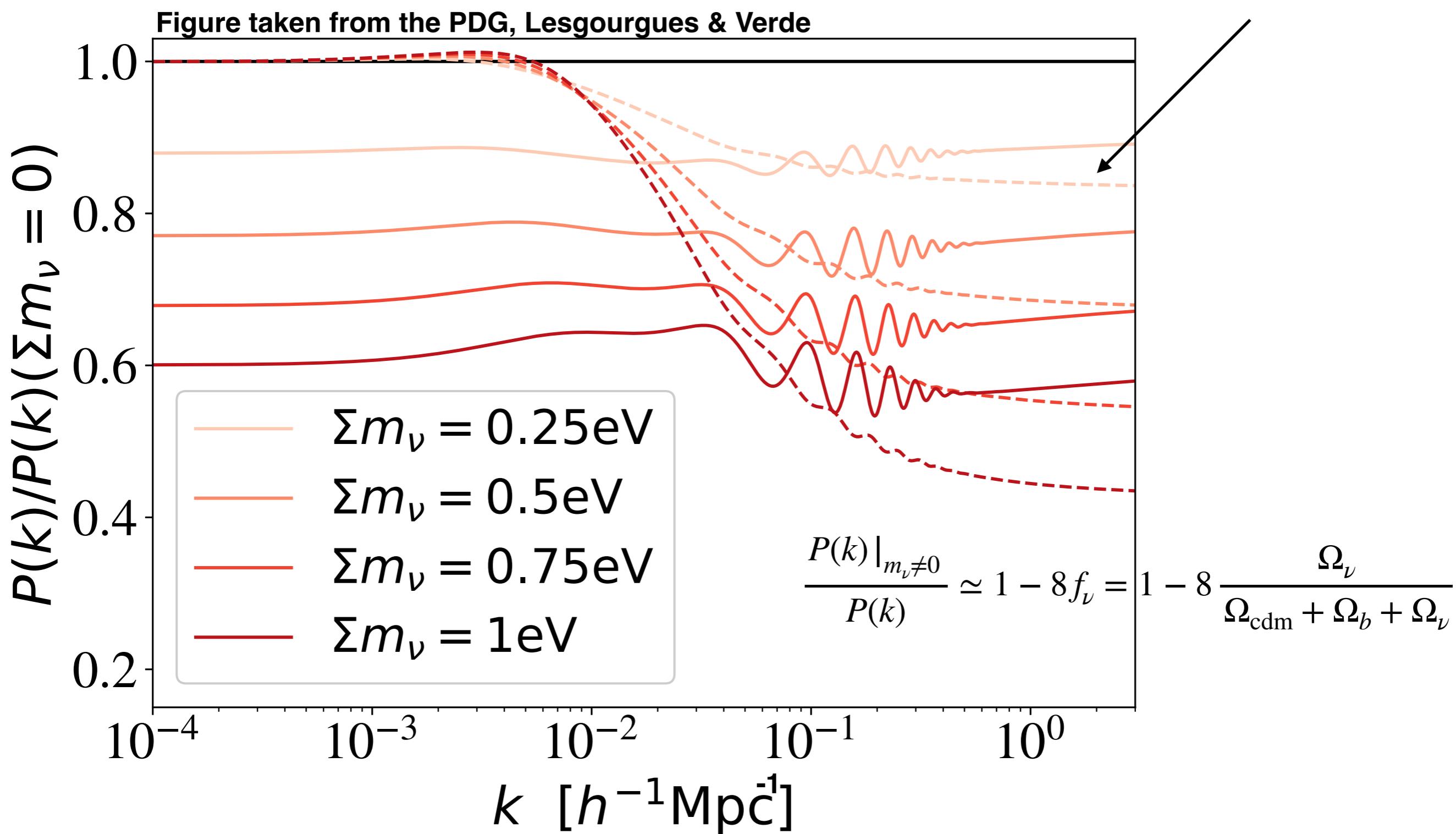
The effect of neutrino masses in the CMB:



# Neutrino Masses in Cosmology

## Galaxy Surveys

Suppression from  $\Omega_\nu h^2$



# Neutrino Masses from Cosmology

## Planck 2018 for $\Lambda$ CDM (1807.06209)

$$\sum m_\nu < 0.54 \text{ eV} \quad (95\% \text{ CL, TT+lowE})$$

$$\sum m_\nu < 0.26 \text{ eV} \quad (95\% \text{ CL, TTTEEE+lowE})$$

$$\sum m_\nu < 0.24 \text{ eV} \quad (95\% \text{ CL, TTTEEE+lowE+lensing})$$

$$\sum m_\nu < 0.12 \text{ eV} \quad (95\% \text{ CL, TTTEEE+lowE+lensing+BAO})$$

To be compared to the KATRIN bound:  $\sum m_\nu < 2.4 \text{ eV}$

**Very robust bounds from linear Cosmology  $\Delta T/T \sim 10^{-5}$**

**What about other non-linear cosmological data?**

**Importantly, all cosmological bounds are cosmological model dependent**

**What is the dependence upon the assumed Cosmological Model?**

# Neutrino Masses from Cosmology

## Data beyond Planck and BAO within $\Lambda$ CDM

$\sum m_\nu < 0.26 \text{ eV}$	<b>Planck</b>	Planck 1807.06209
$\sum m_\nu < 0.12 \text{ eV}$	<b>Planck+BAO</b>	Planck 1807.06209
$\sum m_\nu < 0.86 \text{ eV}$	<b>BOSS P(k)</b>	Ivanov et al. 1909.05277
$\sum m_\nu < 0.16 \text{ eV}$	<b>Planck+BOSS P(k)</b>	Ivanov et al. 1912.08208
$\sum m_\nu < 0.58 \text{ eV}$	<b>Lyman-}\alpha\text{+H}_0\text{prior}</b>	Palanque-Delabrouille et al. 1911.09073
$\sum m_\nu < 0.10 \text{ eV}$	<b>Planck+Lyman-}\alpha</b>	Choudhury & Hannestad 1907.12598
$\sum m_\nu < 0.08 \text{ eV}$	<b>Planck+BAO+H}_0</b>	di Valentino, Gariazzo & Mena 2106.15267
$\sum m_\nu < 0.09 \text{ eV}$	<b>Planck+BAO+SN+RSD</b>	

- Planck is driving current cosmological constraints
- Non-linear or mildly non-linear data sets break degeneracies in the fit
- The larger  $H_0$  is, the stronger the constraint on  $\sum m_\nu$  is (However, this comes from combining two data sets in strong tension!)

# Neutrino Masses from Cosmology

## Cosmological Model Dependence

Planck+BAO and 3 degenerate neutrinos

$$\sum m_\nu < 0.12 \text{ eV}$$

**Standard Case**

$\Lambda\text{CDM}+m_\nu$

Planck 1807.06209

$$\sum m_\nu < 0.25 \text{ eV}$$

**Dark Energy dynamics**

$\text{CDM}+m_\nu+\omega_a+\omega$

Choudhury & Hannestad 19'

$$\sum m_\nu < 0.15 \text{ eV}$$

**Varying Curvature**

$\Lambda\text{CDM}+m_\nu+\Omega_k$

Choudhury & Hannestad 19'

$$\sum m_\nu < 0.13 \text{ eV}$$

**Varying  $N_{\text{eff}}$**

$\Lambda\text{CDM}+m_\nu+N_{\text{eff}}$

Planck 1807.06209

$$\sum m_\nu < 0.17 \text{ eV}$$

**Varying  $N_{\text{eff}}+\omega+a_s+m_\nu$**

$\text{CDM}+m_\nu+N_{\text{eff}}+\omega+a_s+m_\nu$

di Valentino et al. 1908.01391

- Constraints are robust upon standard modifications of  $\Lambda\text{CDM}$

# Neutrino Masses from Cosmology

## Cosmological Model Dependence

### Non-standard Neutrino Cosmologies:

#### Invisible Neutrino Decay

$$\nu_i \rightarrow \nu_j \phi$$

$$\sum m_\nu \lesssim 0.2 \text{ eV}$$

Oldengott, Wong et al. 2203.09075 & 2011.01502

Escudero & Fairbairn 1907.05425

$$\nu_i \rightarrow \nu_4 \phi$$

$$\sum m_\nu \lesssim 0.42 \text{ eV}$$

Abellán, Poulin et al. 1909.05275, 2112.13862

Escudero, López-Pavón, Rius & Sandner 2007.04994

#### Time Dependent Neutrino Masses

Late phase transition

$$\sum m_\nu < 1.4 \text{ eV}$$

Dvali & Funcke 1602.03191

Lorenz et al. 1811.01991 & 2102.13618

#### Ultralight scalar field screening

$$\sum m_\nu < 3 \text{ eV}$$

Esteban & Salvadó 2101.05804

Wetterich et al. 1009.2461

#### Non-standard Neutrino Populations

$$T_\nu < T_\nu^{\text{SM}} + \text{DR}$$

$$\sum m_\nu < 3 \text{ eV}$$

Farzan & Hannestad 1510.02201

Escudero, Schwetz & Terol-Calvo 2211.01729

$$\langle p_\nu \rangle > 3.15 T_\nu^{\text{SM}}$$

$$\sum m_\nu < 3 \text{ eV}$$

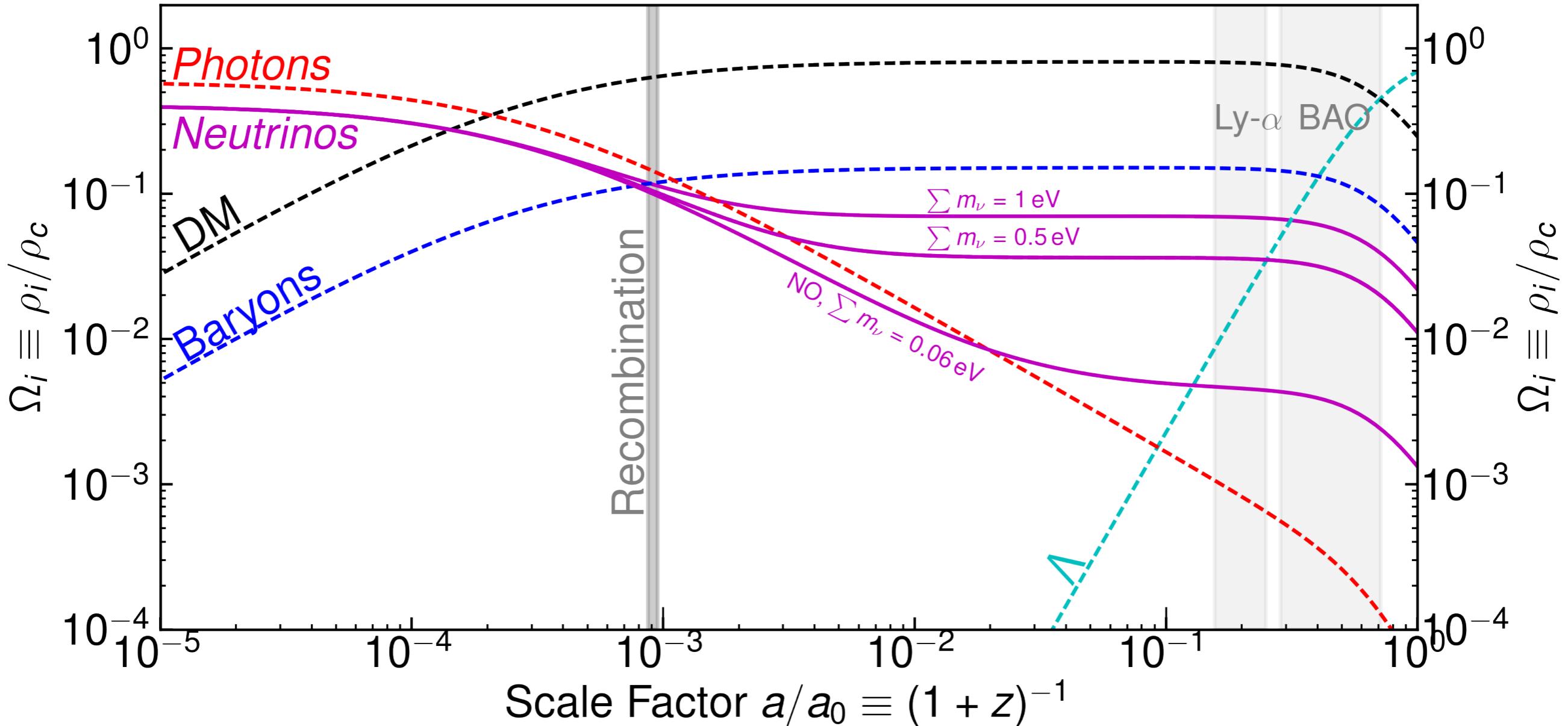
Oldengott et al. 1901.04352

Alvey, Escudero & Sabti 2111.14870

- Bounds can be significantly relaxed in some extensions of  $\Lambda$ CDM. They require modifications to the neutrino sector.

## But Why? and How?

# Neutrino Masses from Cosmology



**CMB peaks fix:**

$$\theta_s \equiv r_s / D_M(z_*)$$

**Massive neutrinos**

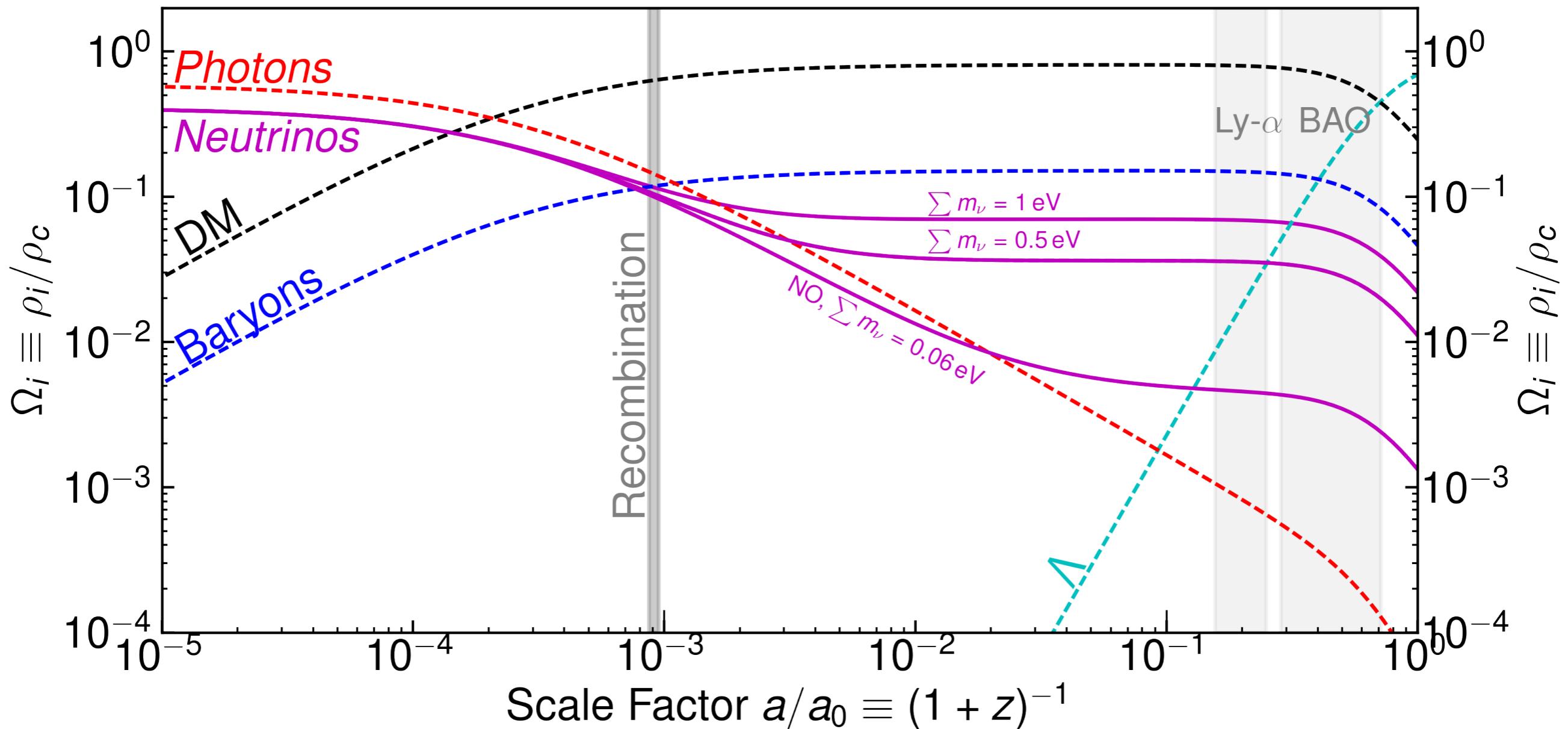
$$r_s = \int_{z_*}^{\infty} \frac{c_s}{H(z')} dz'$$

**Comoving sound horizon  
(Early Universe)**

$$D_M(z) = \int_0^z \frac{1}{H(z')} dz'$$

**Comoving angular diameter distance  
(Late Universe)**

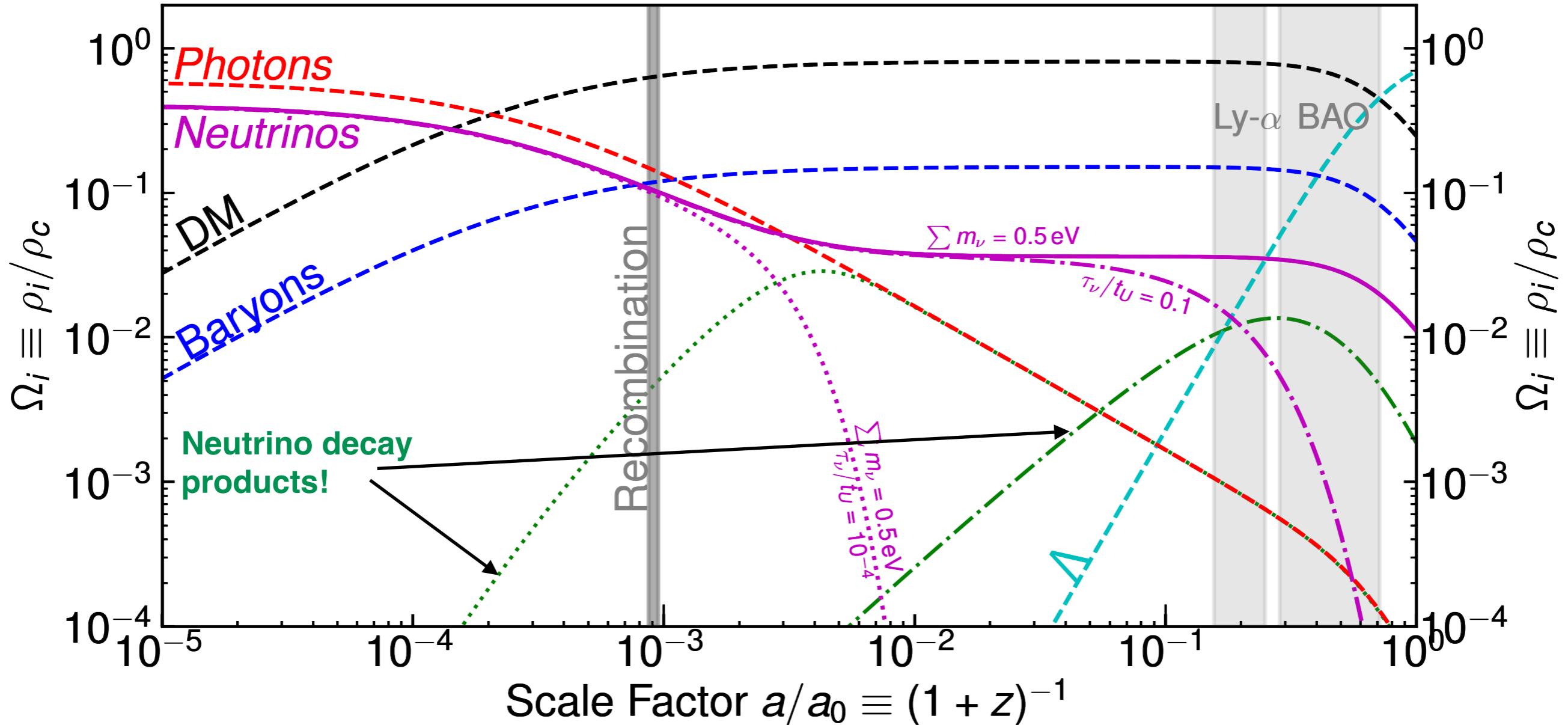
# Neutrino Masses from Cosmology



**Not only a background effect:**

**Massive neutrinos also affect CMB lensing  $\propto \Omega_\nu$**

# Neutrino Decays



**Neutrinos decaying with  $\tau_\nu \lesssim t_U/10$  do not impact  $D_M(z_{\text{CMB}})$**

**Effect of induced neutrino Lensing is substantially reduced**

**Unstable Neutrinos can relax the bounds on  $\Sigma m_\nu$ !**

# Neutrino Masses from Cosmology

## Non-standard Neutrino Cosmologies:

### Invisible Neutrino Decay

$$\nu_i \rightarrow \nu_j \phi$$
$$\sum m_\nu < 0.2 \text{ eV}$$

Oldengott, Wong et al. 2203.09075 & 2011.01502

Escudero & Fairbairn 1907.05425

Archidiacono & Hannestad 1311.3873

$$\nu_i \rightarrow \nu_4 \phi$$

at least:  $\sum m_\nu \lesssim 0.42 \text{ eV}$

Abellán, Poulin et al. 1909.05275, 2112.13862  
Escudero, López-Pavón, Rius & Sandner 2007.04994

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Esteban, Mena & Salvadó 2202.04656

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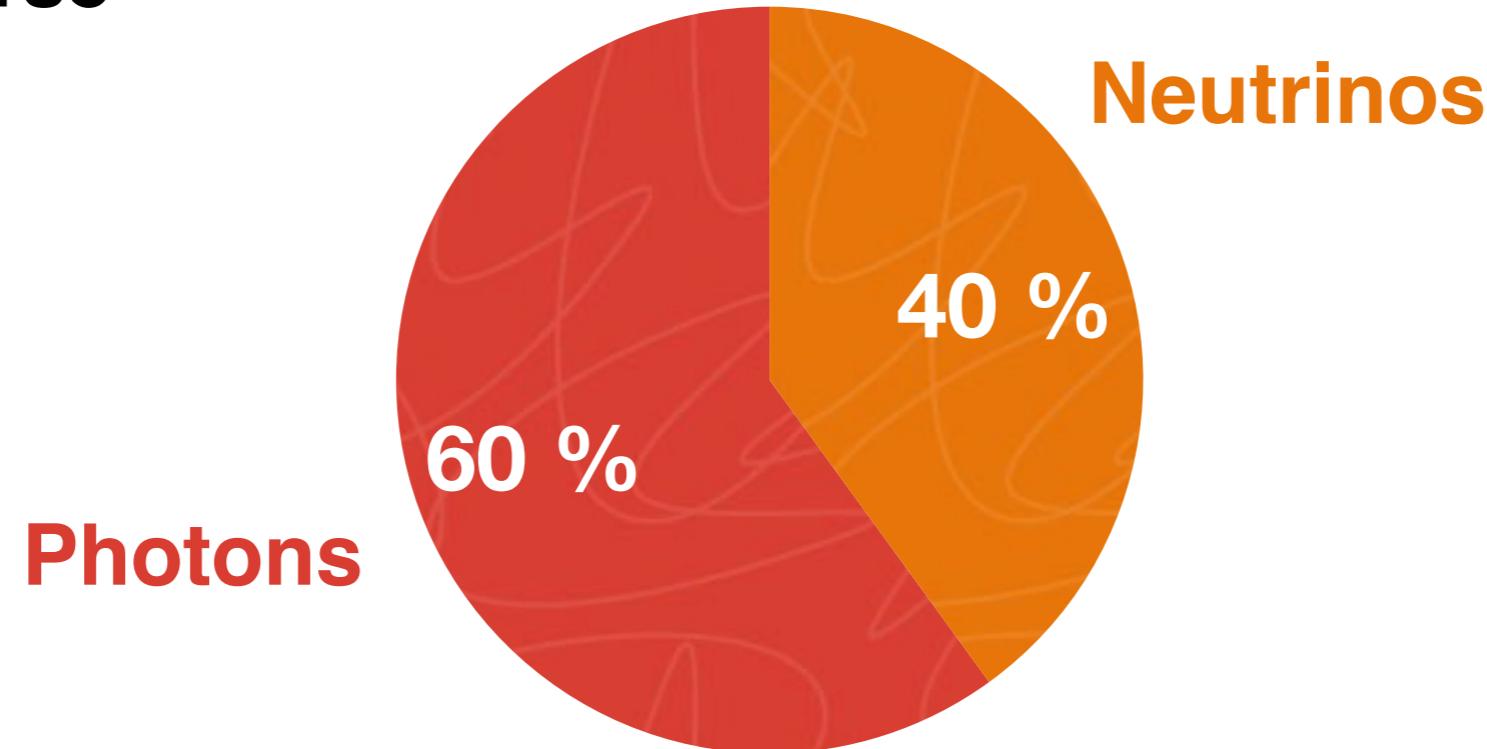
Oldengott et al. 1901.04352  
Alvey, Escudero & Sabti 2111.14870

## Take Away Messages:

- Cosmology can only constrain  $\Omega_\nu(z)$  and not directly  $m_\nu$
- Of course, in  $\Lambda$ CDM there is a direct link between  $\Omega_\nu(z)$  and  $m_\nu$
- All these models reduce  $\Omega_\nu(z)$  with respect to the one in  $\Lambda$ CDM and are in excellent agreement with all known cosmological data
- Importantly, they entail non-standard neutrino properties

# Neutrino Interactions

- Neutrinos represent a large component of the energy density of the Universe



- Neutrinos have very special cosmological perturbations
  - 1) They are ultrarelativistic until  $z \sim 200 m_\nu/0.1 \text{ eV}$
  - 2) In the SM: since  $t_U \sim 0.1 \text{ s}$  ( $T \sim 2 \text{ MeV}$ ), they are free streaming i.e. do not interact with anything

These together actually mean that CMB observations can probe potential neutrino interactions!

# Why?

First discussed by Bashinsky & Seljak in [astro-ph/0310198] and applied by Chacko, Hall, Okui & Oliver [hep-ph/0312267]

- The key is in Einstein's equations

$$G_{\mu\nu} = 8\pi G T_{\mu\nu}$$

Background expansion: Neff

$$\delta G_{\mu\nu} = 8\pi G \delta T_{\mu\nu}$$

Perturbations: can tell about interactions

Neutrino anisotropic stress

$$\sigma_\nu$$

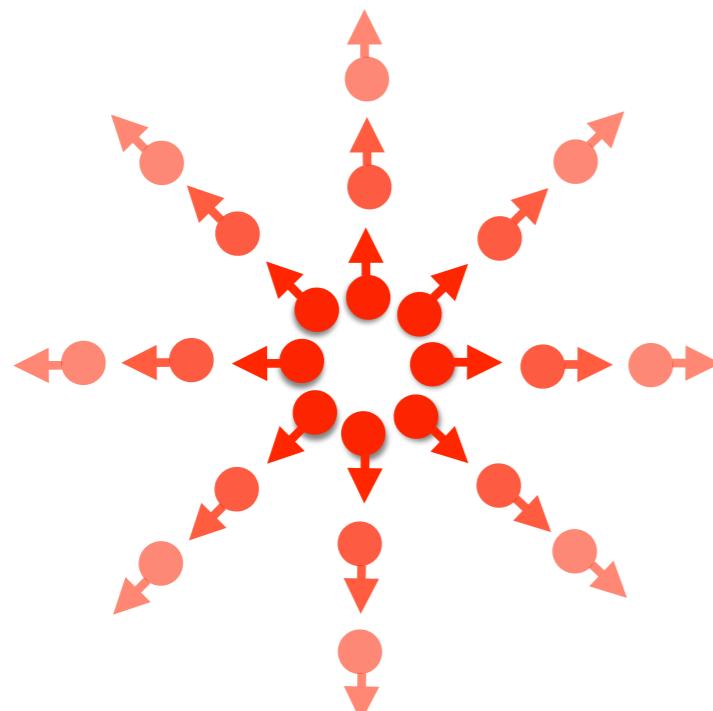
Metric

$$\delta g_{\mu\nu}$$

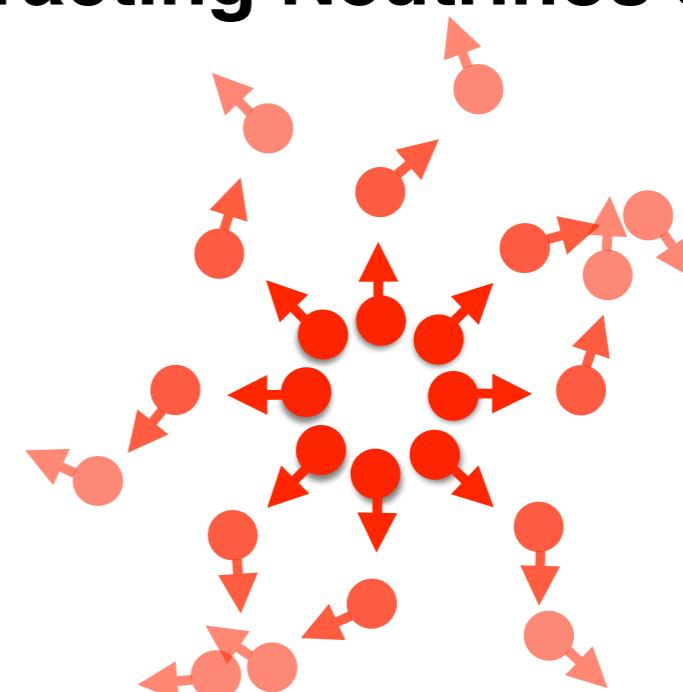
CMB spectra

$$\Delta T_\gamma$$

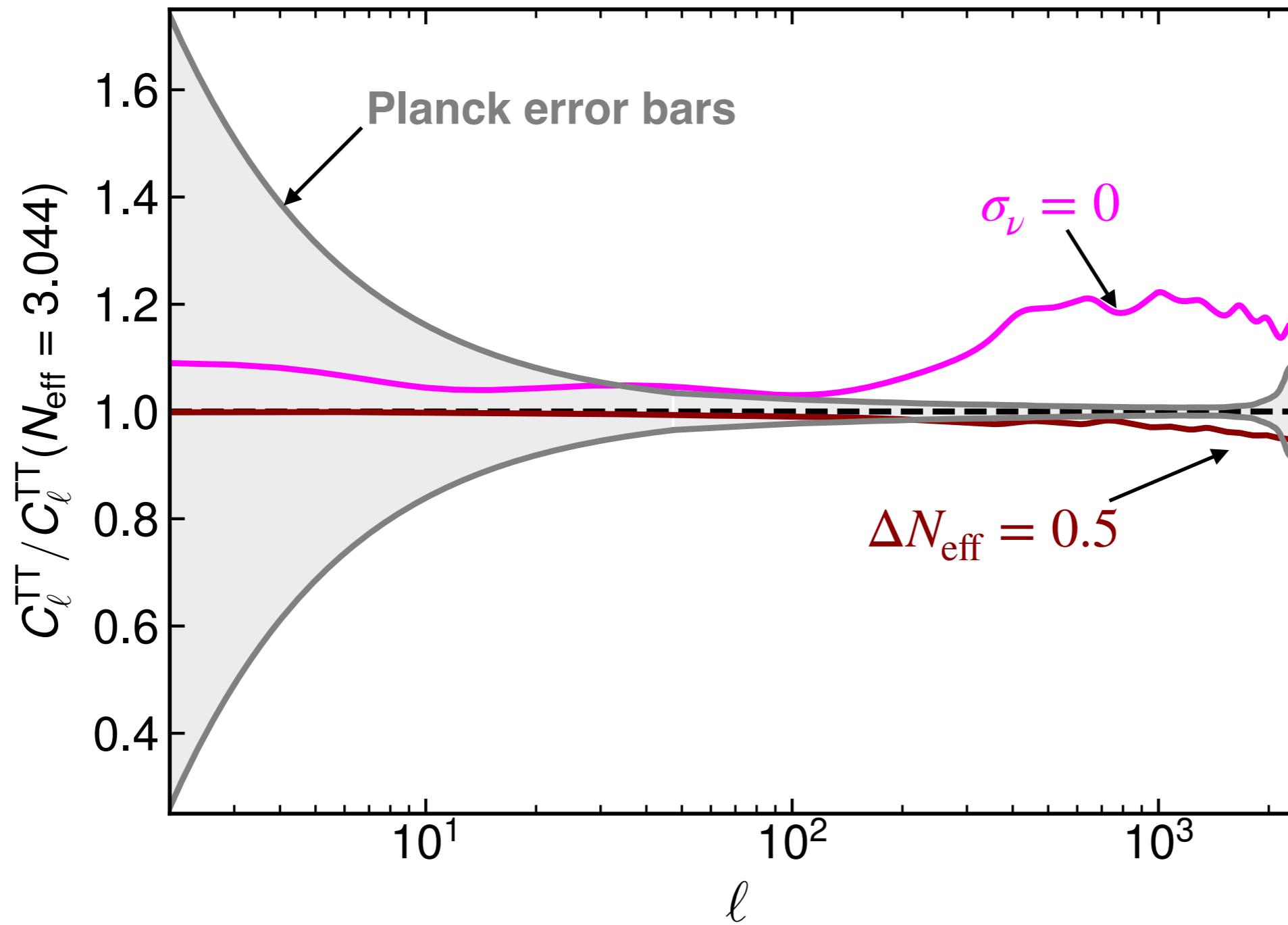
Free Streaming Neutrinos  $\sigma_\nu \neq 0$



Interacting Neutrinos  $\sigma_\nu \rightarrow 0$

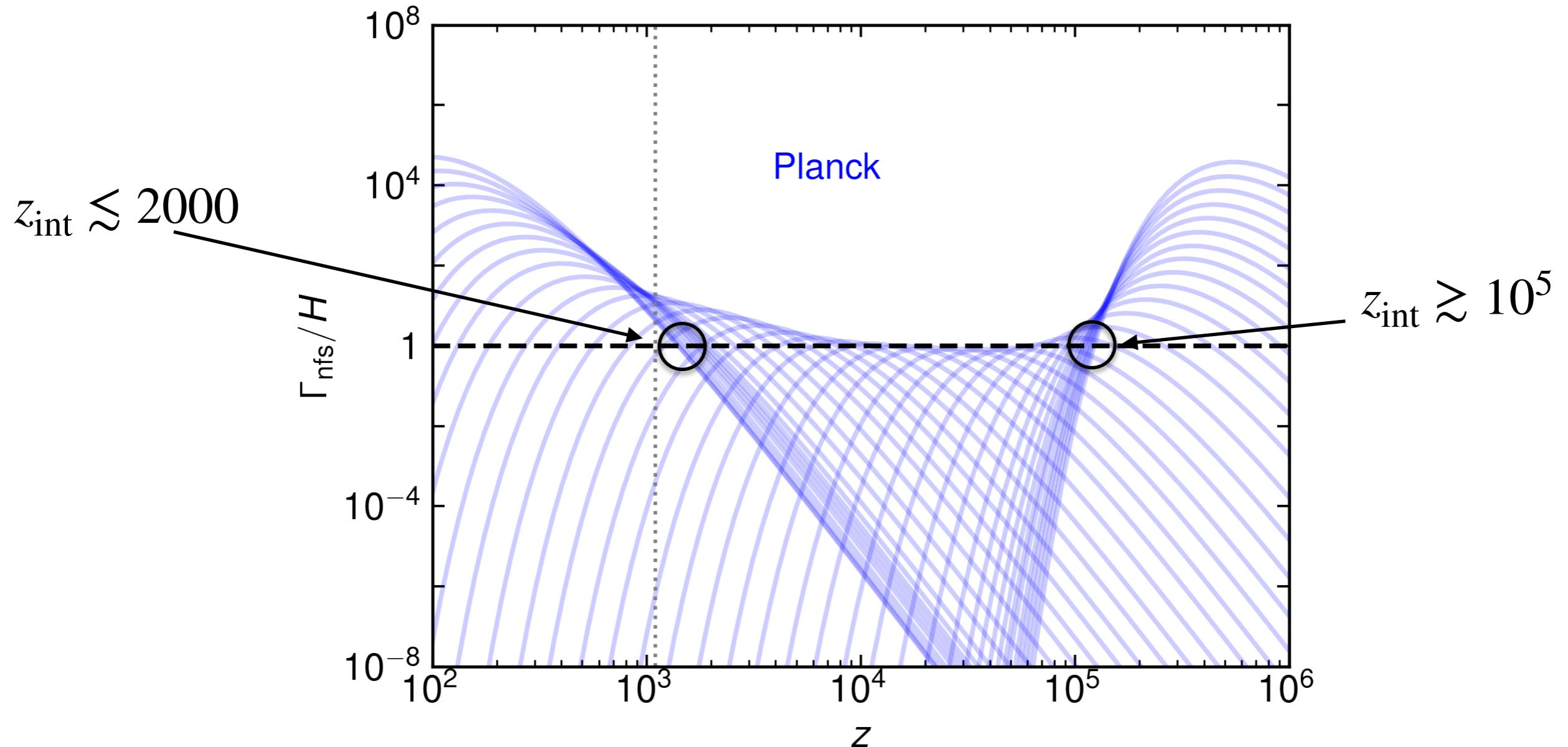


# Effect of Neutrino Free-streaming in the CMB



# The Neutrino Freestreaming window

Together with Petter Taule and Mathias Garny in [2207.04062](#) we have recently established the presence of a neutrino free streaming window.



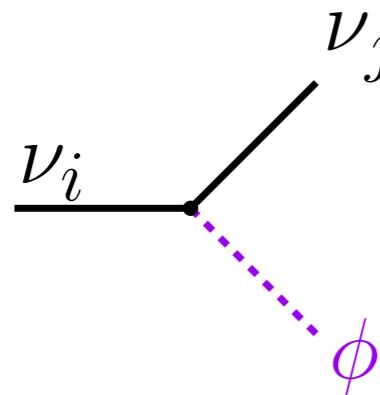
This analysis reinforces the fact that neutrino free streaming cannot be suppressed in

$$2000 \lesssim z_{\text{int}} \lesssim 10^5$$

$$0.3 \text{ eV} \lesssim T_\nu \lesssim 15 \text{ eV}$$

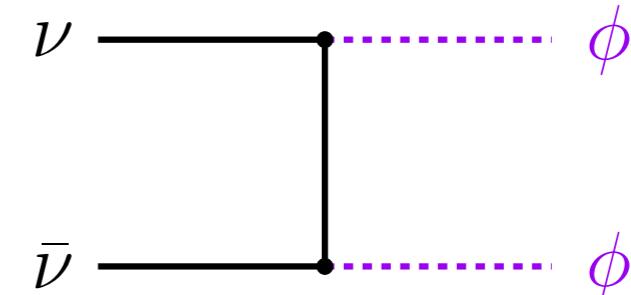
# Models

## Neutrino Decays



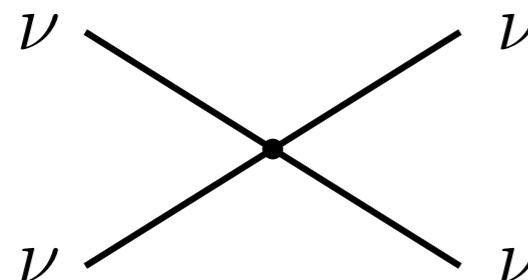
Hannestad & Raffelt [hep-ph/0509278]  
Basbøll, Bjaelde, Hannestad & Raffelt [0806.1735]  
Escudero & Fairbairn [1907.05425]  
Chacko, Dev, Du, V. Poulin and Y. Tsai [1909.05275]  
Barenboim, Chen, Hannestad, Oldengott, Tram & Wong [2011.01502]  
Abellán, Chacko, Dev, Du, Poulin & Tsai [2112.13862]  
Chen, Oldengott, Pierobon & Wong [2203.09075]

## Neutrino Annihilations



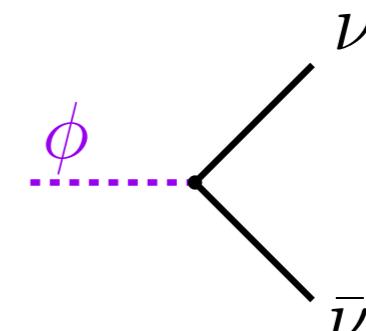
Beacom, Bell & Dodelson [astro-ph/0404585]  
Hannestad [astro-ph/0411475]  
Archidiacono & Hannestad [1311.3873]  
Forastieri, Lattanzi & Natoli [1904.07810]

## Neutrino Scatterings



Cyr-Racine & Sigurdson [1306.1536]  
Lancaster, Cyr-Racine, Knox & Pan [1704.06657]  
Oldengott, Tram, Rampf & Wong [1706.02123]  
Kreisch, Cyr-Racine & Doré [1902.00534]  
Das & Ghosh [2011.12315]  
Choudhury, Hannestad & Tram [2012.07519]  
Brinckmann, Chang & LoVerde [2012.11830]

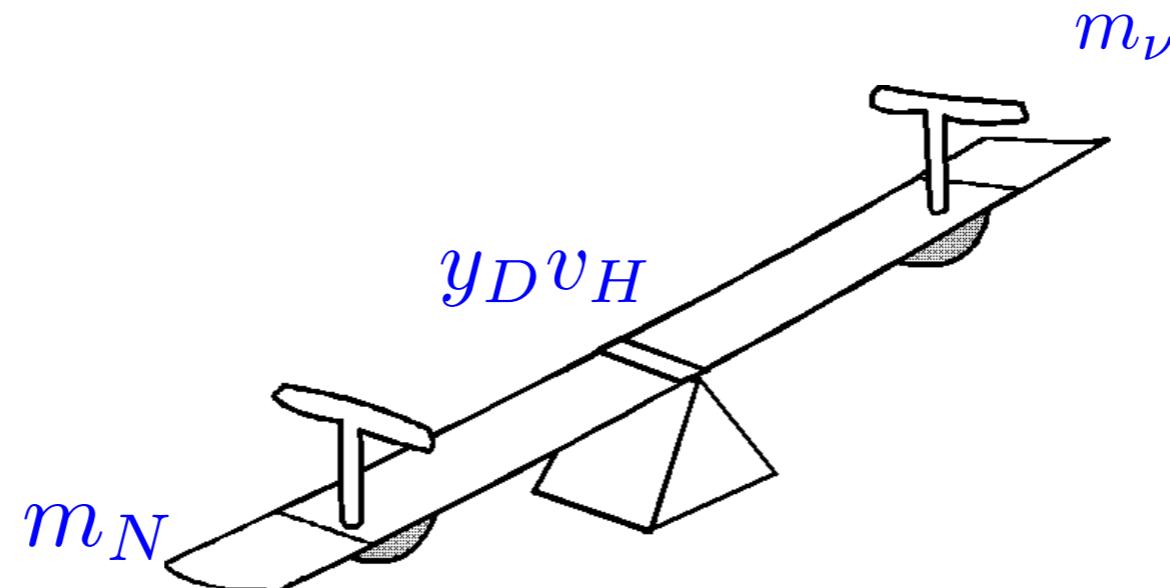
## eV-scale neutrinophilic bosons



Chacko, Hall, Okui & Oliver [hep-ph/0312267]  
Escudero & Witte [1909.04044]  
Escudero & Witte [2103.03249]

# The case of the Majoron

Type-I seesaw



**Neutrinos are very light Majorana particles:**  $m_\nu \simeq 0.03 \text{ eV} \left( \frac{y_D}{10^{-6}} \right)^2 \frac{\text{TeV}}{M_N}$

Are There Real Goldstone Bosons Associated with Broken Lepton Number? #3

Y. Chikashige (Munich, Max Planck Inst.), Rabindra N. Mohapatra (Munich, Max Planck Inst. and Munich U.), R.D. Peccei (Munich, Max Planck Inst.) (Sep, 1980)

Published in: *Phys.Lett.B* 98 (1981) 265-268

[DOI](#) [cite](#) [claim](#)

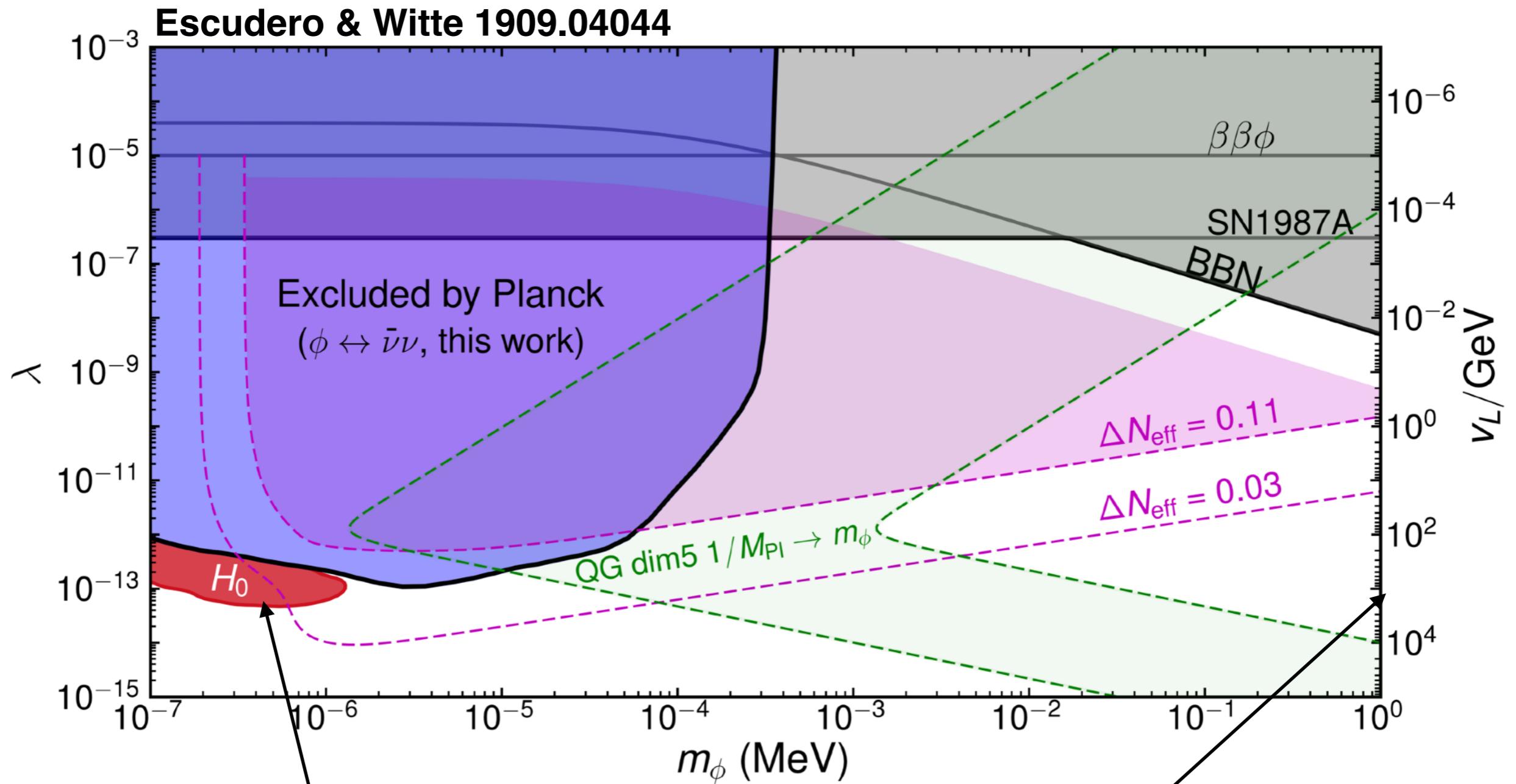
[reference search](#)

[1,118 citations](#)

**The Majoron is the pseudo-Goldstone boson associated with the spontaneous breaking of  $U(1)_L$**

$$\mathcal{L} = \lambda \phi \bar{\nu} \gamma_5 \nu \quad \lambda = m_\nu / v_\phi$$

# The case of the Majoron



- CMB observations can test a well motivated neutrino mass model up to TeV scales!
- The model features an enhanced expansion history and together with the majoron-neutrino interactions can ameliorate the Hubble tension

# Recent Anomalous Helium Measurement

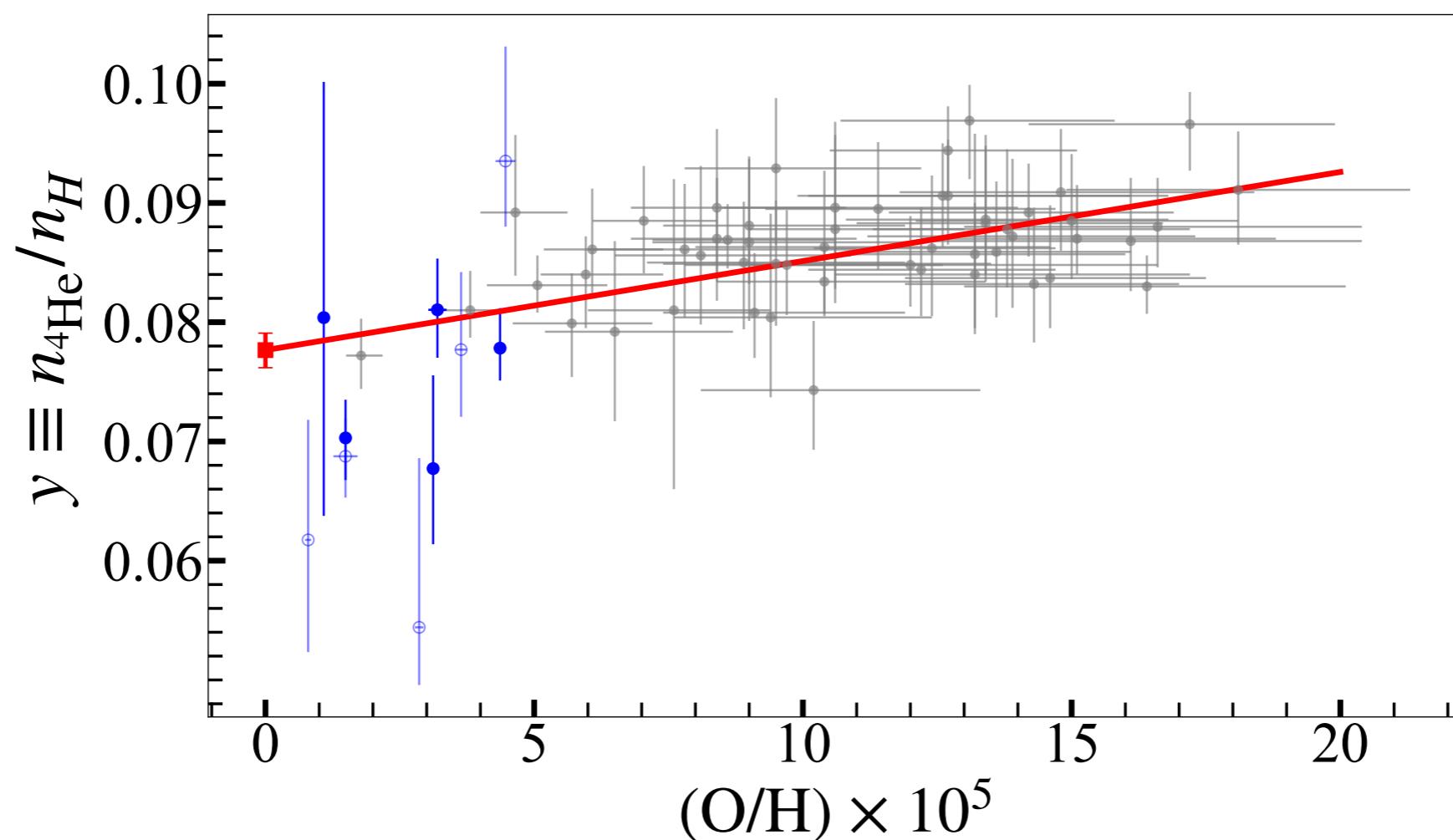
The EMPRESS collaboration reported a Primordial Helium abundance smaller than the SM value by 3 sigma

$$Y_P \equiv \frac{\rho_{^4\text{He}}}{\rho_b}$$

$$Y_P = 0.2370(34) \quad \textbf{Matsumoto et al. [2203.09617]}$$

$$Y_P^{\text{SBBN}} = 0.2471(02) \quad \textbf{Pitrou et al. [1801.08023]}$$

There are several other recent  $Y_p$  determinations that are in agreement with the SM but their new measurement is interesting because they have 5 new extremely metal poor systems:



# Primordial Lepton Asymmetries

The simplest BSM interpretation of a lower value of the primordial helium abundance is the presence of a non-zero lepton asymmetry in electron neutrinos at the time of BBN

[Matsumoto et al. \[2203.09617\]](#) [Burns, Tait & Valli \[2206.00693\]](#) [Escudero, Ibarra & Maura \[2208.03201\]](#)

## Some basics about lepton asymmetries:

The lepton asymmetries are parametrized with a comoving neutrino chemical potential:

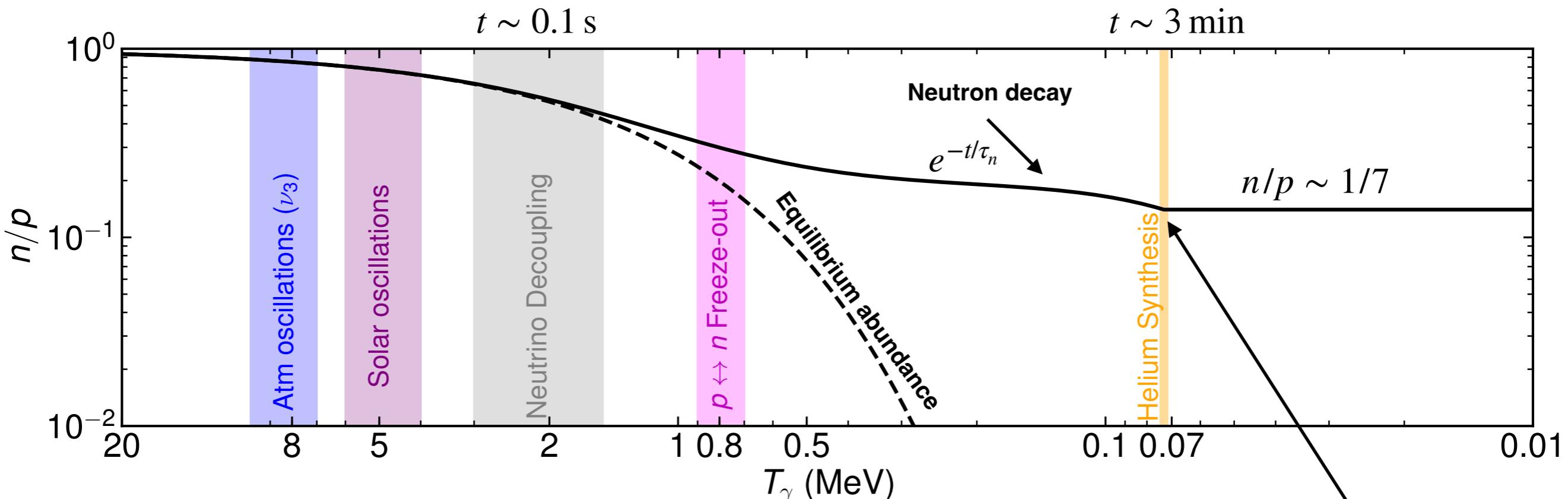
$$\xi_\nu \equiv \mu_\nu / T_\nu \quad \frac{n_{\nu_e} - n_{\nu_{\bar{e}}}}{n_\gamma} \simeq 0.25 \times \xi_{\nu_e}$$

Neutrinos start to oscillate at  $T \sim 8 \text{ MeV}$

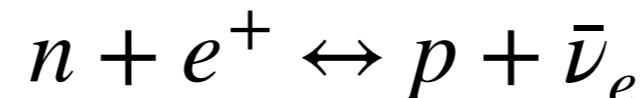
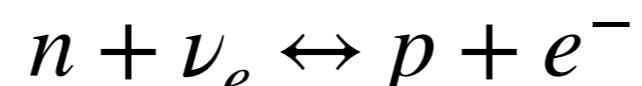
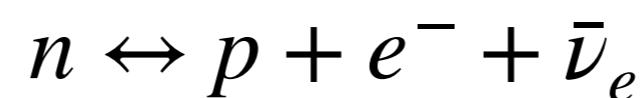
This means that  $|\xi_{\nu_e}| \simeq |\xi_{\nu_\mu}| \simeq |\xi_{\nu_\tau}|$  by the time of BBN  
see e.g.: 0809.0631 locco, Mangano, Miele, Pisanti & Serpico

This implies that a constraint on  $\xi_{\nu_e}$  generically applies to  $\xi_{\nu_{\mu/\tau}}$

# Primordial Helium and $\nu_e$ -Asymmetry

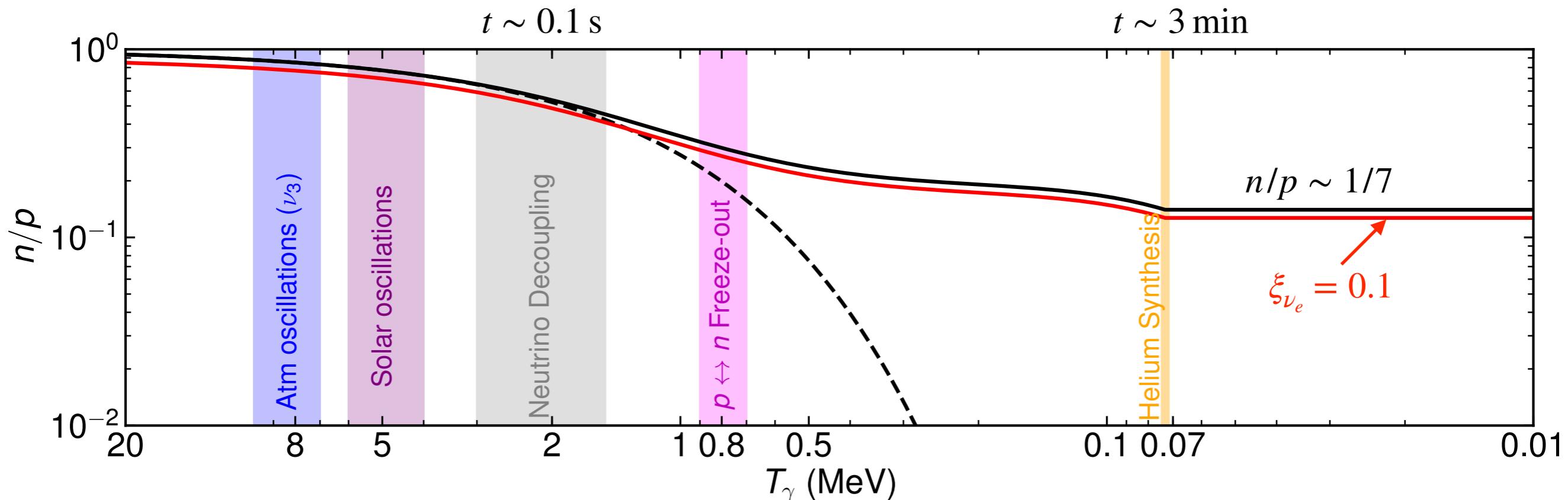


The key processes

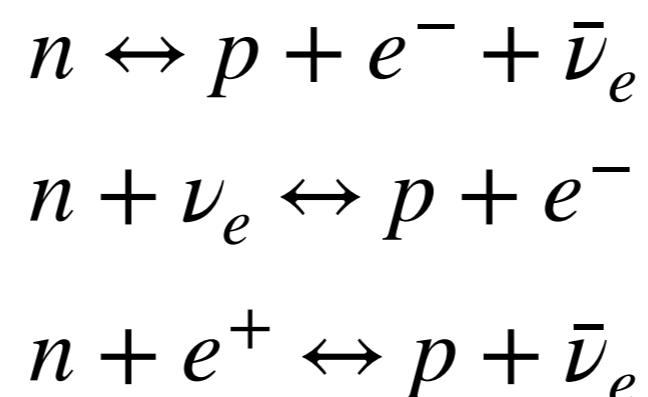


$$Y_P \approx \frac{2n/p}{1 + 2n/p} \Big|_{T=T_{\text{BBN}}}$$

# Primordial Helium and $\nu_e$ -Asymmetry



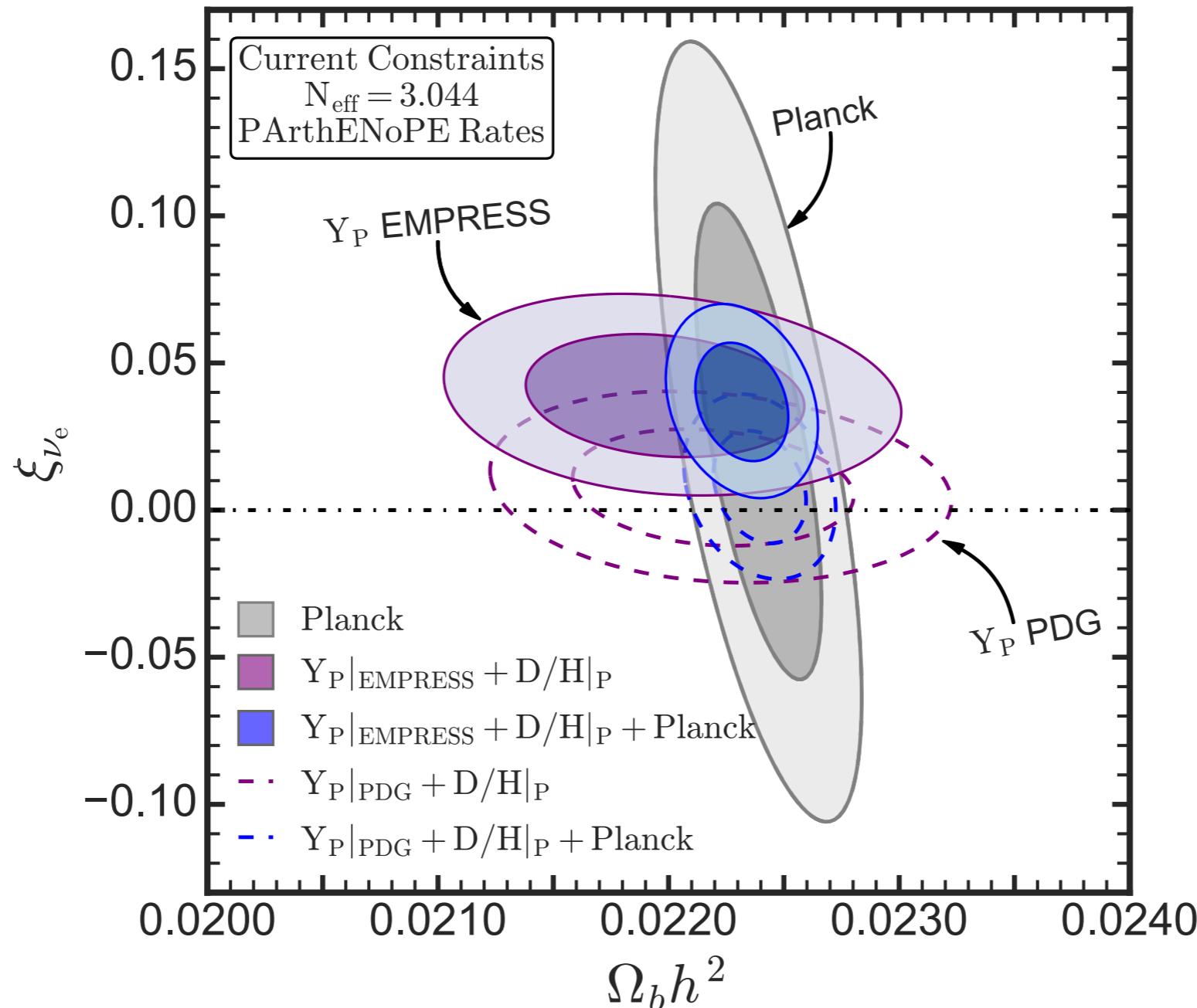
The key processes



$$Y_P \approx \left. \frac{2n/p}{1 + 2n/p} \right|_{T=T_{\text{BBN}}}$$

# $\nu_e$ -Asymmetry: Current Status

Escudero, Ibarra & Maura [2208.03201] (to appear in PRD)

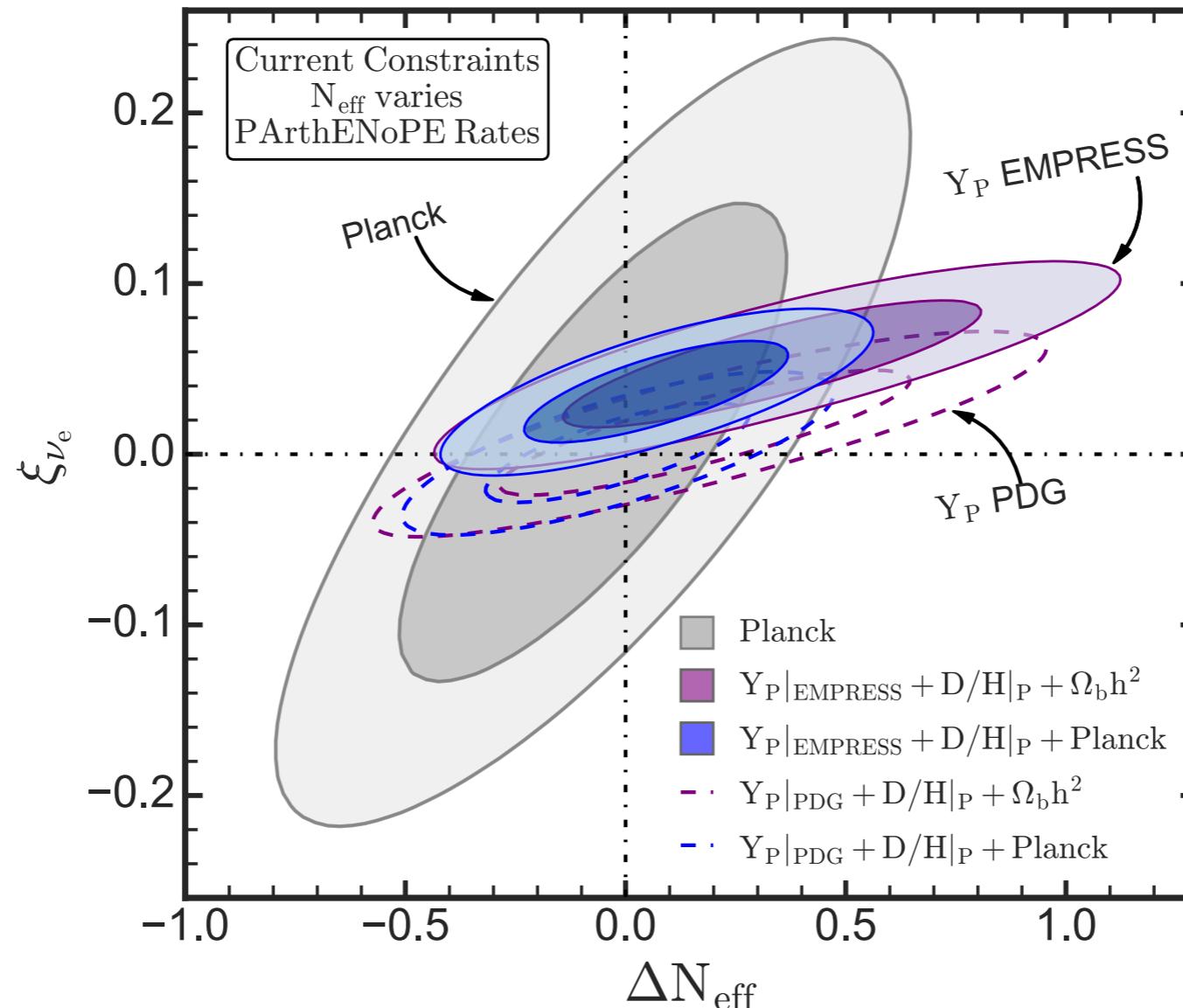


**Planck is compatible with a zero lepton asymmetry**

**The preference for a non-zero asymmetry is driven only by EMPRESS. Taking the PDG-21 recommended value points to  $\xi_{\nu_e} = 0$**

# $\nu_e$ -Asymmetry: Current Status

Escudero, Ibarra & Maura [2208.03201] (to appear in PRD)



The presence of a non-zero  $\Delta N_{\text{eff}}$  changes the inference on  $\xi_{\nu_e}$

Taking EMPRESS: Preference for  $\xi_{\nu_e} > 0$ :

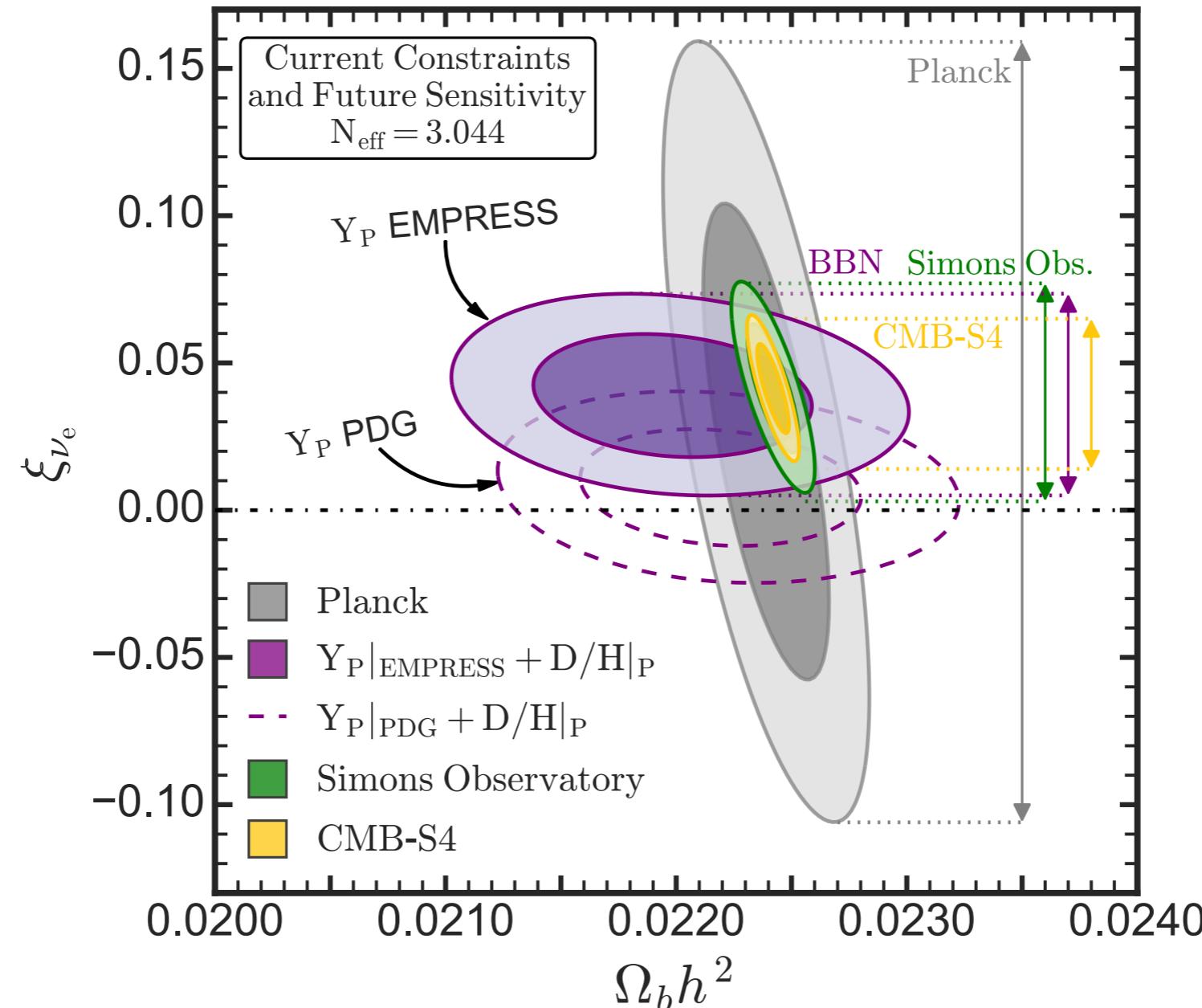
$2.9\sigma$  **Fixed Neff = 3.044**

$1.9\sigma$  **Varying Neff**

$2.9\sigma$  **Varying Neff but with  $\Delta N_{\text{eff}} > 0$**

# $\nu_e$ -Asymmetry: Future Sensitivity

Escudero, Ibarra & Maura [2208.03201] (to appear in PRD)



**Simons will reach a similar sensitivity than current BBN analyses. CMB-S4 will do even better:**

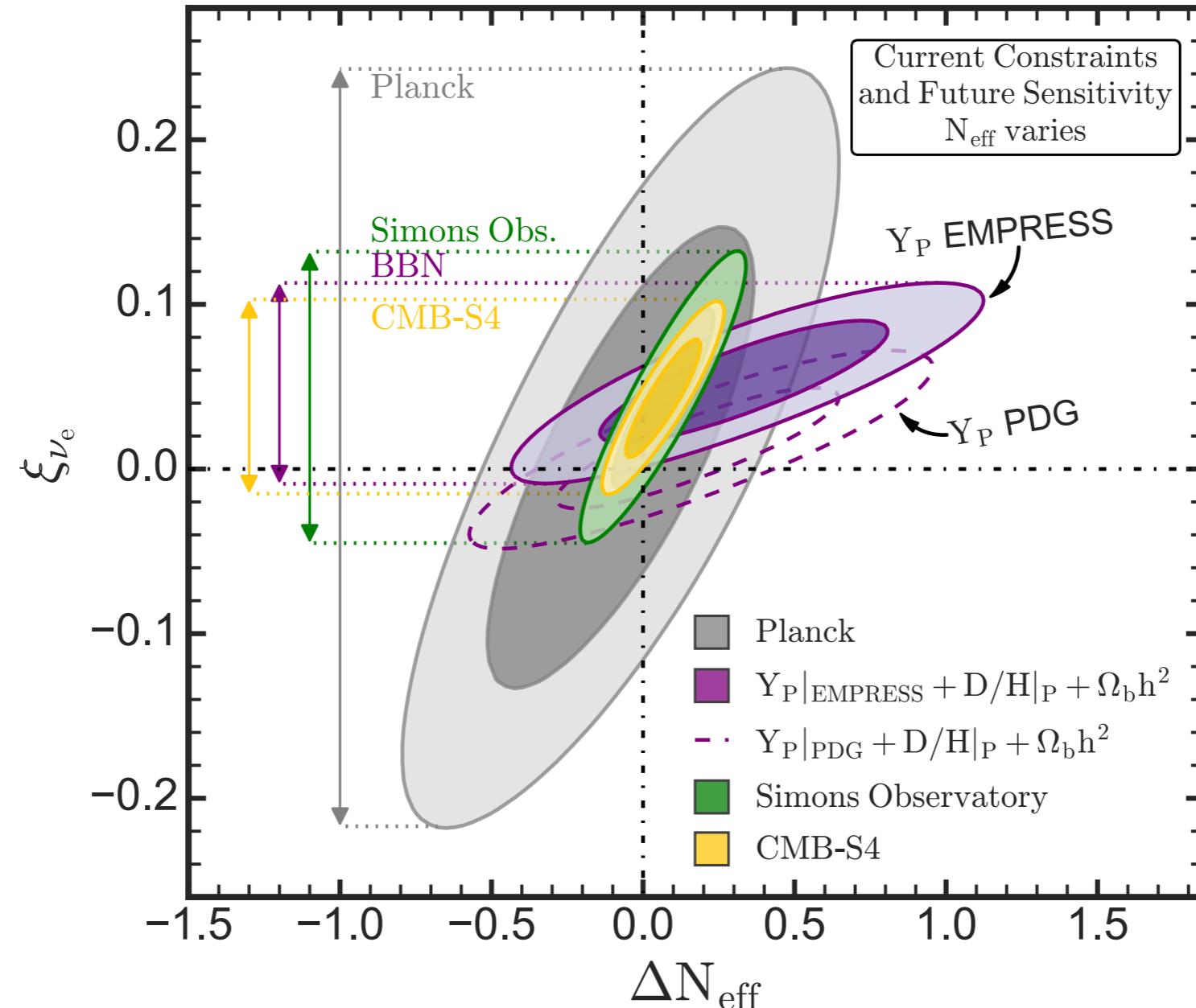
**EMPRESS:  $2.9\sigma$**

**EMPRESS+Simons:  $4.3\sigma$**

**EMPRESS+CMB-S4:  $5.2\sigma$**

# $\nu_e$ -Asymmetry: Future Sensitivity

Escudero, Ibarra & Maura [2208.03201] (to appear in PRD)



With  $\Delta N_{\text{eff}}$  it will be harder to break the degeneracies, even with Simons or CMB-S4. The tension will have low significance

EMPRESS:  $1.9\sigma$

EMPRESS+Simons:  $2.8\sigma$

EMPRESS+CMB-S4:  $3.1\sigma$

# Conclusions

## Neutrino Masses:

Cosmological bounds are very stringent in  $\Lambda$ CDM:  $\sum m_\nu < 0.12 \text{ eV}$

There are several non-standard cosmologies where this bound can be evaded

Unstable neutrinos are a plausible particle physics possibility

## Neutrino Interactions:

The CMB is a powerful probe of neutrino interactions

We have shown that there is a well defined redshift region where neutrinos must free stream

$$2 \times 10^3 \lesssim z \lesssim 10^5$$

These bounds are relevant for many particle physics scenarios

## Neutrino Asymmetries:

The EMPRESS survey has recently reported an anomalously low Helium abundance

We have performed a global BBN and CMB analysis showing:

$$\frac{n_{\nu_e} - n_{\bar{\nu}_e}}{n_\gamma} \sim 1 \%$$

# Outlook

## Neutrino Masses:

**KATRIN reach:**

$$\sum m_\nu < 0.6 \text{ eV} \quad (90\% \text{ CL})$$

**Next generation of  $0\nu 2\beta$  experiments, e.g. LEGEND:**  $m_{\beta\beta} \sim 0.02 - 0.04 \text{ eV}$

**JUNO/DUNE expected to determine the neutrino ordering**

**Next Galaxy Surveys+CMB should detect neutrino masses**

e.g.: 1308.4164 Font-Ribera et al., 1408.7052 Kitching et al.

**DESI/EUCLID+Planck:**  $\sigma \left( \sum m_\nu \right) \simeq 0.02 \text{ eV} \quad (1\sigma)$

## Lepton Asymmetries:

**Simons Observatory in  $\sim 6$  years:**  $\sigma(\xi_{\nu_e}) = 0.014$ ,  $\sigma(N_{\text{eff}}) = 0.06$

**CMB-S4 in  $>10$  years**  $\sigma(\xi_{\nu_e}) = 0.010$ ,  $\sigma(N_{\text{eff}}) = 0.03$

**Should be able to clarify the situation independently on local measurements.**

**Of course, further scrutiny is needed in the local measurements that are currently dominated by systematic effects**

**What about BSM models?**

# Time for Questions and Comments

**Upcoming years are going to be exciting!**



**Thank you for your attention!**

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