A history of some recent attempts to go beyond the standard model

Francesco Vissani INFN, Gran Sasso & GSSI

introduction

The standard model (SM) plays a role that resembles the one of thermodynamics, of Newton's, Maxwell's Einstein's theories or alike. I mean, it is a **principle-based theory**, which allows us to perform useful deductions concerning microphysics and to obtain expectations obeyed in most circumstances.

As usual, however, Ist rule is **to save the phenomena**, that in the specific field of competence of SM means the available experimental & observational facts of microphysics – the facts.

introduction

Certain phenomena contradict SM's predictions, thereby calling particle physicists to maximum attention.

Other facts go beyond the direct competences of the standard model, but put it into play along with other theories, especially in astrophysical/cosmological contexts.

Some structural aspects of this theory raise interesting questions, which – even without amounting to actual problems - deserve consideration.

roots of standard model

the first W.I.N. revolution



Absohrift/15.12.

Offener Brief an die Gruppe der Radicektiven bei der Gauvereins-Tagung zu Tübingen.

Abschrift

Physikalisches Institut der Eidg. Technischen Hochschule Zurich

Zürich, 4. Des. 1930 Cloriastrasse

Liebe Radioaktive Damen und Herren,

Wie der Ueberbringer dieser Zeilen, den ich huldvollst anzuhören bitte, Ihnen des näheren auseinandersetzen wird, bin ich angesichts der "falschen" Statistik der N- und Li-6 Kerne, sowie des kontinuierlichen beta-Spektrums auf einen versweifelten Ausweg verfallen um den "Wechselsats" (1) der Statistik und den Energiesats su retten. Nämlich die Möglichkeit, es könnten elektrisch neutrale Teilchen, die ich Neutronen nennen will, in den Kernen existieren, welche den Spin 1/2 haben und das Ausschliessungsprinzip befolgen und den von Lichtquanten musserden noch dadurch unterscheiden, dass sie det mit Lichtgeschwindigkeit laufen. Die Masse der Neutronen inste von derselben Grossenordnung wie die Elektronenwasse sein und Seignfalls nicht grösser als 0.01 Protonenmasse.- Das kontinuierliche beta- Spektrum wäre dann verständlich unter der Annahme, dass beim beta-Zerfall mit dem Elektron jeweils noch ein Neutron emittiert wird, derart, dass die Summe der Energien von Neutron und Elektron konstant ist.

Also, liebe Radioaktive, prüfst, und richtst.- Leider kann ich nicht personlich in Täbingen erscheinen, da sch infolge eines in der Macht vom 6. sum 7 Des. in Zürich stattfindenden Balles hier unabkömmlich bin.- Mit vielen Grügsen an Euch, sowie an Herrn Back, Buer untertanigster Diener

ges. W. Pauli

Asia Pacific Physics Newsletter | Vol. 01, No. 01, pp. 27-30 (2012) Fermi's β-decay Theory

Chen Ning Yang

https://doi.org/10.1142/S2251158X12000045 | Cited by: 0

Ι.

Throughout his lifetime Enrico Fermi (1901–1954) had considered his 1934 β -decay theory as his most important contribution to theoretical physics. E. Segrè (1905–1989) had vividly written about an episode at the inception of that paper:¹

Fermi gave the first account of this theory to several of his Roman friends while we were spending the Christmas vacation of 1993 in the Alys. It was in the evening after a full day of skiing; we were all sitting on one bed in a hotel room, and I could hardly keep still in that position, bruised as I was after several falls on Icy snow. Fermi was fully aware of the importance of his accomplishment and said that he would be remembered for this paper, his best so far. He sent a letter to *Nature* advancing his theory, but the editor refused it because he thought it contained speculations that were too remote from physical reality; and instead the paper ("Tentative Theory of Beta Rays" (FP 76)) was published in Italian and in the *Zeitschrift Far Physik*^b

In 2001 there was a centennial celebration of Fermi's 100th birthday. I contributed a paper to that celebration. One passage of my paper read:⁹

- One day in the 1970's, I had the following conversation with Eugene Wigner in the cafeteria of Rockefeller University:
- Y: What do you think was Fermi's most important contribution to theoretical physics?
- W: β -decay theory.
- Y: How could that be? It is being replaced by more fundamental ideas. Of course it was a very important contribution which had sustained the whole field for

some forty years: Fermi had characteristically swept what was unknowable at that time under the rug, and focused on what can be calculated. It was beautiful and agreed with experiment. But it was not permanent. In contrast the Fermi distribution is permanent.

- W: No, no, you do not understand the impact it produced at the time. Von Neumann and I had been thinking about β -decay for a long time, as did everybody else. We simply did not know how to create an electron in a nucleus.
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What this passage reflected was not just the very different evaluations, by Wigner and me, of Fermi's β -decay theory, but in fact the very different evaluations of this theory by Fermi himself and his generation of physicists and by my generation. Recently I looked into the old literature and was able to understand better the reason for this difference.

II.

In 1932 there were two discoveries that greatly shocked the world of physicists:

On February 17th, J. Chadwick (1891–1974) sent a short article to Nature⁴ with the title "Possible Existence of a Neutron". It had great immediate impact: One realized that nuclei were made of protons and neutrons rather than protons and electrons, thus understanding the many regularities of the composition of light nuclei. But at the same time it created a new difficult puzzle: Since there were

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the first W.I.N. revolution

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Nobel to Reines 65 years later

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the second W.I.N. revolution

A turning point in the understanding of weak interactions is the hypothesis that they violate parity, due to Lee and Yang (1956) a fact confirmed by the experiment of Wu (1957) and recognized by the Nobel committee in 1957. This was the key to understand the structure of weak interactions and it allowed Landau, Lee & Yang and Salam to conclude that, for neutrinos, the spin and the momentum have opposite directions while, for antineutrinos, the direction is the same one. One talks also of negative helicity of neutrinos and positive helicity of antineutrinos. The final proof of this picture was obtained by the impressive experiment of Goldhaber et al. (1958). Eventually, the theoretical picture was completed arguing for an universal vector-minus-axial (V–A) nature of the charged-current weak interactions (Sudarshan and Marshak, 1958; Feynman and Gell-Mann, 1958).

Introduction to the Formalism of Neutrino Oscillations Guido Fantini, Andrea Gallo Rosso, Francesco Vissani, Vanessa Zema Published in Adv.Ser.Direct.High Energy Phys. 28 (2018) 37-119

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Nobel to Lee & Yang 1957

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J. C. Ward

from 'Memoirs of a Theoretical Physicist'

"Quite soon after this triumph, the experiment of Mrs. C. S. Wu *et al.* at Columbia, acting upon the suggestion of Yang and Lee, definitely established the nonconservation of parity in weak interactions, surprising everyone.

I wrote a note to Abdus, telling him of the result, adding that Einstein must be spinning in his grave, clockwise presumably."



under the spell of gauge principle

- Yang+Mills 1954, non abelian gauge theories
- 2 Bludman, Leite Lopes 1958, neutral currents interactions
- (3) Glashow 1961, SU(2)_L × U(1)_Y and the Z^0 boson
- Anderson, Higgs, Englert+Brout,
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- (7) 't Hooft+Veltman, Lee, 1971, renormalizability
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Nobel prizes in physics 1979, 1984, 1999, 2008, 2013

SM: ingredients and recipe

- \blacktriangleright SU(3)_c × SU(2)_L × U(1)_Y gauge group
- > three families of Weyl spinors: q_L , l_L , u_R , d_R , e_R (6+2+3+3+1=15)
- ➤ one scalar doublet: H, whose negative mass squared breaks spontaneously $SU(2)_L \times U(1)_Y \rightarrow U(1)_{em}$

write all possible renormalizable terms and fix the value of the parameters by as many measurements as the parameters

SM: major implications

- one single mass scale (or two, I mean also the energy density) 17 dimensionless parameters (or 18, I mean also θ_{CP})
- ✓ baryon & leptons numbers, B, L_e, L_µ, L₊ conserved perturbatively.
 B-L, L_e-L_µ, L_e-L₊ conserved exactly (L=L_e+L_µ+L₊).
- \square calculable theory to any perturbation order.
- no flavor-changing neutral current at tree level.
 CP only mildly violated.
- non perturbative regime for SU(3)_c at low energies (glueballs?). lightest hadrons - proton, nuclei - are stable as the electron.
- ☑ neutrinos are massless.

proton decay

the age of gauge (theories)





Unified interactions of leptons and hadrons *

Harald Fritzsch, Peter Minkowski

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Abstract

It is suggested that a unifying description of leptons and hadrons can be obtained within a nonabelian gauge theory where the gauge group is a symmetry group of a set of massless elementary fermions (leptons, quarks). We investigate the consequences of such an approach for the strong, electromagnetic, and weak interactions. We study both gauge theories with and without fermion number conservation, e.g., theories based on the groups $SU_n \times SU_n$ (n = 8, 12, 16) and SO_n (n = 10, 14).

from 1979 Nobel lectures

Salam:

That summer [1973, ed] Jogesh Pati and I had predicted proton decay within the context of what is now called GUT.

Glashow:

GUT - perhaps along the lines of the original SU(5) theory of Georgi and me - must be essentially correct. This implies that the proton, and indeed all nuclear matter, must be inherently unstable.

Weinberg:

If effects of a tiny non-conservation of baryon or lepton number such as proton decay or neutrino masses are discovered experimentally, we will then be left with gauge symmetries as the only true internal symmetries of nature, a conclusion that I would regard as most satisfactory.

round table "Einstein and the physics of the future" published in Some Strangeness in the Proportion, ed. H. Woolf, 1980

Weinberg:

[...] the lifetime of the proton (this has been worked on by a number of people now) comes out to be of the order 10^{30} to 10^{32} years. The present experimental lower bound is 10^{29} years. Thus the time is ripe for an assault on the next few orders of magnitude in the proton lifetime.

Dyson:

[...] the modern view of particle theory, with the sub-nuclear world a playground of interlocking broken and unbroken symmetries, had its roots in Felix Klein's Erlanger Program of 1872 [...] I predict that in the next 25 years we shall see the emergence of unified physical theories in which general relativity, group theory, and field theory are tied together with bonds of rigorous maths.

Yang:

beautiful mathematics is the language of fundamental physics [...]

Maybe it is my prejudice - maybe it is my ignorance - but I do not believe that any of these graded Lie algebras has the intrinsic and fundamental beauty of Lie algebras and Lie groups, not as yet!

KGF, IMB, NUSEX, KAMIOKANDE, ICARUS...



... JUNO, DUNE, HYPER-KAMIOKANDE



... JUNO, DUNE, HYPER-KAMIOKANDE



on expectations

- we have hints of gauge coupling unification
- but a prediction requires a complete theory
- significant uncertainties from choice of the scalar field, intermediate scales, threshold effects, nuclear matrix elements
- fermion masses are one of the few constraint

on expectations

- we have hints of gauge coupling unification
- but a prediction requires a complete theory
- significant uncertainties from choice of the scalar field, intermediate scales, threshold effects, nuclear matrix elements
- fermion masses are one of the few constraint
- (new particles? dark matter? supersymmetry? baryogenesis?)

What to some seemed a blitzkrieg has become a trench warfare, with little room for progress: can we still fight the fight? Surely, we still have nothing like a "GUT standard model".

neutrino masses

the idea that worked – neutrino oscillations was proposed earlier (late fifties / early sixties)





full recognition 1/2 a century later

Nobelpriset i fysik 2015

Nobelpriset i fysik 2015



Takaaki Kajita Super-Kamiokande Collaboration University of Tokyo, Kashiwa, Japan

Arthur B. McDonald Sudbury Neutrino Observatory Collaboration Queen's University, Kingston, Canada

"för upptäckten av neutrinooscillationer, som visar att neutriner har massa" "for the discovery of neutrino oscillations, which shows that neutrinos have mass"



The Nobel Prize in Physics 2015

solar neutrinos odyssey

- * lot of resistance to the idea of solar neutrino experiments, see Bahcall
- moreover, the solar neutrino anomaly has long been ignored and it took 20 years before a test was conducted
 - > in part it was distrust of astrophysics and nuclear physics
 - > in part this was due to widespread prejudice against neutrino mass
 - > acceptance begun when it was realized that small mixing could work
- * today, oscillation of solar neutrinos is considered an obvious thing
- (even if there are still doubts and things to do)

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where we are now

tests with terrestrial experiments, from reactor to accelerator

overall consistency of the indications

important role of "global fits" – i.e., of taking seriously the hypothesis of massive neutrinos

also appearance was seen (more on that later)

first hints of "normal spectrum" - aka ordering, aka hierarchy

now, we should see true neutrino mass – we hope in KATRIN



Progress in Particle and Nuclear Physics Volume 102, September 2018, Pages 48-72



Review Current unknowns in the three-neutrino framework

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F. Capozzi <sup>a</sup>, E. Lisi <sup>b</sup> ≈ , A. Marrone <sup>c, b</sup>, A. Palazzo <sup>c, b</sup>
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https://doi.org/10.1016/j.ppnp.2018.05.005

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Abstract

We present an up-to-date global analysis of data coming from neutrino oscillation and non-oscillation experiments, as available in April 2018, within the standard framework including three massive and mixed neutrinos. We discuss in detail the status of the three-neutrino (3ν) mass-mixing parameters, both known and unknown. Concerning the latter, we find that: normal ordering (NO) is favored over inverted ordering (IO) at 3σ level; the Dirac CP phase is constrained within $\sim 15\%$ ($\sim 9\%$) uncertainty in NO (IO) around nearly-maximal CP-violating values; the octant of the largest mixing angle and the absolute neutrino masses remain undetermined. We briefly comment on other unknowns related to theoretical and experimental uncertainties (within 3ν) or possible new states and interactions (beyond 3ν).

SM seen as an effective theory

Weinberg 79

$$\delta \mathcal{L} = \frac{(\ell H)^2}{M} + \frac{\ell q q q}{M'^2} + \frac{(\ell q d^c)^2}{M''^5} \text{ with } \begin{cases} M < 10^{11} \text{ TeV} & \text{for dim.5} \\ M' > 10^{12} \text{ TeV} & \text{for dim.6} \\ M'' > 5 \text{ TeV} & \text{for dim.9} \end{cases}$$

- the 1st is the SM-invariant formulation of Majorana neutrino masses
- the 2nd is one of the operators that cause the instability of the proton
- the 3rd violates lepton number

Der Springer Link



Journal of High Energy Physics

Updated global analysis of neutrino oscillations in the presence of eV-scale sterile neutrinos

Authors Authors and affiliations

Mona Dentler⊡, Álvaro Hernández-Cabezudo, Joachim Kopp, Pedro Machado, Michele Maltoni, Ivan Martinez-Soler, Thomas Schwetz

Open Access Regular Article - Theoretical Physics First Online: 03 August 2018



Abstract

We discuss the possibility to explain the anomalies in short-baseline neutrino oscillation experiments in terms of sterile neutrinos. We work in a 3 + 1 framework and pay special attention to recent new data from reactor experiments, IceCube and MINOS+. We find that results from the DANSS and NEOS reactor experiments support the sterile neutrino explanation of the reactor anomaly, based on an analysis that relies solely on the relative comparison of measured reactor spectra. Global data from the v_e disappearance channel favour sterile neutrino oscillations at the 3 σ level with Δm $_{41}$ 2 \approx 1.3 eV 2 and $|U_{e4}| \approx$ 0.1, even without any assumptions on predicted reactor fluxes. In contrast, the anomalies in the v_{ρ} appearance channel (dominated by LSND) are in strong tension with improved bounds on v_{μ} disappearance, mostly driven by MINOS+ and IceCube. Under the sterile neutrino oscillation hypothesis, the p-value for those data sets being consistent is less than 2.6×10^{-6} . Therefore, an explanation of the LSND anomaly in terms of sterile neutrino oscillations in the 3 + 1 scenario is excluded at the 4.7 σ level. This result is robust with respect to variations in the analysis and used data, in particular it depends neither on the theoretically predicted reactor neutrino fluxes, nor on constraints from any single experiment. Irrespective of the anomalies, we provide updated constraints on the allowed mixing strengths $|U_{\alpha 4}|$ ($\alpha = e, \mu, \tau$) of active neutrinos with a fourth neutrino mass state in the eV range.

more light neutrinos?

various neutrino experiments found anomalous results

individually, many of them could be explained invoking new light neutrinos

however, when one takes this hypothesis seriously, contradictions emerge and this interpretation is not supported by the experiments

the current situation with cosmology is also contradictory

(on top of that, while new light neutrinos might exist, there is no strong theoretical argument in their favor)

residual global symmetries

there is only one basic type of lepton

(=at the scrutiny of T2K, NOvA, OPERA, SK, DeepCore, only total lepton number L survived)

	ΔL _e	ΔL _μ	۵L	ΔL
v _µ →v _e	+1	-1	0	0
ν _μ →ν _τ	0	-1	+1	0

Appearance experiments proved that all global symmetries of SM are violated, except L (with B). Conversion among families is possible, we have only two types of **matter particles**: leptons and quarks

In "standard model" B and L are not conserved individually.

Only B-L is conserved exactly, thus: light & heavy matter types are connected





	Δ(L_e-L_μ)	Δ(L _μ -L _τ)	Δ(L _τ -L _e)	Δ(B-L)
v _µ →v _e	+2	-1	-1	0
ν _μ →ν _τ	+1	-2	+1	0

Appearance experiments proved that **all anomaly free symmetries** of SM are violated, except one. B+L is not a conserved number in the Standard Model --- leptons and baryons conversion is possible

probe lepton number by neutrinoless double beta decay, i.e, by creation of electrons




usually we see ultrarelativistic (anti) neutrinos



if we stopped them, we would see the spin states



SM seen as an effective theory

Weinberg 79

$$\delta \mathcal{L} = \frac{(\ell H)^2}{M} + \frac{\ell q q q}{M'^2} + \frac{(\ell q d^c)^2}{M''^5} \text{ with } \begin{cases} M < 10^{11} \text{ TeV} & \text{for dim.5} \\ M' > 10^{12} \text{ TeV} & \text{for dim.6} \\ M'' > 5 \text{ TeV} & \text{for dim.9} \end{cases}$$

- the Ist is the SM-invariant formulation of Majorana neutrino masses which violates B-L
- the 2nd is one of the operators that cause the instability of the proton **but** conserves B-L
- the 3rd violates lepton number **and contributes to** $0\nu 2\beta$.
- At dim.9 also **B violation** appears

a systematic, updated study of 0v2β



sensitivity to Majorana mass



Munich, 04/11/19





cosmology & SM

$$\rho_{\rm crit.} \equiv \frac{3(H_0)^2}{8\pi G_{\rm N}} \approx 5.0 \ \frac{m_{\rm p}}{{\rm m}^3}$$

$$\rho_{\Lambda} \equiv \frac{\Lambda c^2}{8\pi G_{\rm N}} \approx 3.5 \, \frac{m_{\rm p}}{{\rm m}^3}$$

$$\rho_{\text{CDM}} + \rho_{\text{bary.}} \approx 1.5 \, \frac{m_{\text{p}}}{\text{m}^3}$$

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we have $1/\sqrt{\Lambda} \approx 3$ Gpc, similar in size to visible universe

$$\rho_{\rm crit.} \equiv \frac{3(H_0)^2}{8\pi G_{\rm N}} \approx 5.0 \ \frac{m_{\rm p}}{{\rm m}^3}$$

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setting $\rho_{\Lambda}c^2 \equiv \mathcal{E}^4/(\hbar c)^3$ we get $\mathcal{E} \approx 2.2 \text{ meV}$

$$\rho_{\rm crit.} \equiv \frac{3(H_0)^2}{8\pi G_{\rm N}} \approx 5.0 \ \frac{m_{\rm p}}{{\rm m}^3}$$

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$$ho_{
m CDM}+
ho_{
m bary}pprox 1.5\,rac{m_{
m p}}{{
m m}^3}$$

most matter is non-baryonic $\rho_{\text{CDM}} \sim 5 \times \rho_{\text{bary}}$

a few of SM-related questions

* 4.9% of baryons

Why anti-baryons are absent? Can we have a theory of that? No way to explain this with SM during big-bang

* 26.4% of non-relativistic matter / DM

What's this? Why it is similar to previous number? Can we see it in lab? No other massive particle in SM (except primordial BH or neutrinos)

* 68.7% of dark energy

How to know more about that? No problem for SM to provide this, maybe the opposite problem



troubles with determination of H₀?

Adding some additional relativistic species help relieve the tension. In principle, new neutrinos could help, but their mass contribution should be negligible. Not a painless modification.

CMB is sensitive to $\Sigma = m_1 + m_2 + m_3$



PHYSICAL REVIEW D 100, 073003 (2019)



FIG. 2. Differential distributions of $m_{\beta\beta}$ obtained by projecting the density plot of Fig. 1 (red line) and as a result of the analytic procedure discussed in the text (blue line). The 1σ , 2σ and 3σ C. L. intervals are reported.

a few remarks



on right handed neutrinos

- it is plausible that right handed neutrinos exist
- with them, B-L can be promoted to a gauge symmetry
- if lighter <10¹⁴⁻¹⁵ GeV, light neutrino masses can be explained
- they can also give reason of baryon asymmetry
- if heavier $> 10^{7-8}$ GeV, they mean "hierarchy problem"
- (but plenty of similar problems with cosmological constant)

SUPERSYMMETRY, SUPERGRAVITY AND PARTICLE PHYSICS

H.P. NILLES*

Département de Physique Théorique, Université de Genève, 1211 Genève 4, Switzerland

and

CERN, Genève, Switzerland

Received 16 February 1984

future. Experiments within the next five to ten years will enable us to decide whether supersymmetry as a solution of the naturalness problem of the weak interactions scale is a myth or reality. This far for the 100 GeV region. The extension of the models to higher energies is of course only

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2019-1984=35 y and 35/(5 to 10)≥3.5, thus the CL of the 2nd option seems > 3.5σ (0.05%) ☺

quosque tandem, gravity?

consider the mass that realizes

$$\frac{e^2}{r} \equiv G_{\rm N} \, \frac{m_{\rm s}^2}{r}$$

this defines the Stoney mass (1881)

$$\Rightarrow m_{\rm s} \equiv \sqrt{\frac{e^2}{G_{\rm N}}} \approx 2\mu g$$

energy at which gravity equates electric force

$$\Rightarrow E_{\rm s} \equiv m_{\rm s} \times c^2 \approx 10^{18} \, {\rm GeV}$$

not the only pending issue

- o B-physics, 3σ
- o (g-2)_μ, 4 σ
- o p-radius / μ -atoms, 6 σ (?)
- o 8 Be decay (Atomki), $6.8\,\sigma$

- □ nature of dark energy [doubt]
- □ nature of inflation [doubt]
- $\square \quad value of H_0 \qquad [cosm. problem]$
- □ lithium abundance [problem]

- deficit of v_{e} (gallium), 3 σ
- deficit of reactor $\overline{\mathbf{v}}_{\mathrm{e}}$, $3\,\sigma$
- $\overline{\mathbf{v}}_{e}$ appearance (LSND + MiniBoone), 6 σ

- shape of high energy cosmic neutrino spectra, 2σ
- shape of solar neutrino spectra, 3σ
- nuclear recoils (DAMA), II σ

Workshop on "Gravity, Information and Fundamental Symmetries" Munich, MPQ, 04.-06.11.2019

Begin	End	Sunday		Monday		Tuesday		Wednesday
08:55	09:00			Welcome				
09:00	10:00		Chair: Klaus Blaum	Francesco Vissani, LNGS+Gran Sasso Inst. Aspects, limitations and cracks of the standard model of elementary particles	Chair: Ignacio Cirac	Jun Ye, Univ. of Colorado, Boulder Quantum matter, clocks, and fundamental physics	Chair: Dieter Lüst	Cumrun Vafa, Harvard Quantum Gravity and String Theory Landscape
10:00	10:45			David DeMille, Yale Probing PeV scale physics via particle electric dipole moments		Piet Schmidt, PTB Testing Fundamental Physics with Quantum Logic Spectroscopy of Highly Charged Ions		Xiaoliang Qi, Stanford University Operator growth in the SYK mode
10:45	11:15			Coffee + Tea		Coffee + Tea		Coffee + Tea
11:15	12:00			Stefan Ulmer, RIEKEN+CERN Antimatter under the Microscope - High Precision Comparisons of the Fundamental Properties of Matter/Antimatter Conjugates		Tanya Zelevinsky, Columbia Molecular clocks for precision metrology		Silke Weinfurtner, Univ. of Nottingham Quantum simulators of gravitational effects
12:00	12:45			Martin Fertl, JGU Mainz Precision low-energy physics with muons		Ronald Walsworth, Harvard Quantum Tools to Explore the Universe		Kyriakos Papadodimas, ICTP Trieste The black hole information paradox and space-time behind the horizon
12:45	14:00			Lunch		Lunch		Lunch
14:00	14:45	Arrival	Chair: Manfred Lindner	Gustaaf Brooijmans, Columbia Neutrons, Anti-Neutrons and Us	Chair: Allen Caldwell	Javier Redondo, Zaragoza Dark matter and axions	Chair: Ignacio Grac	Andrew Geraci, Northwestern University Searching for gravitational waves, ultra-light dark matter, and "fifth" forces with levitated opto-mechanics
14:45	15:30			Kate Scholberg, Duke Interactions of Neutrinos with Matter		Babette Döbrich, CERN Axion searches across different mass scales		Fernando Pastawski, PsiQuantum Holography, locality and quantum error correction
15:30	16:00			Coffee + Tea		Coffee + Tea		Coffee + Tea
16:00	16:45			Susanne Mertens, MPP+ TUM Neutrino mass measurements		Teresa Marrodan Undagoitia, MPIK Dark Matter and neutrino physics with liquid xenon detectors		Tadashi Takayanagi, Kyoto University Entanglement Wedges and AdS/CFT
16:45	17:30			Xianguo Lu, Oxford Neutrino Physics with Accelerators		Regina Caputo, NASA Recent dark matter searches with Fermi-LAT		Bryan Swingle, Univ. of Maryland Black Hole Information Scrambling and Spacetime Locality
17:30	18:15			Veronika Hubeny, UC Davis Entanglement structure for holography		Discussion		Discussion
18:15	19.15			Discussion				

summary and discussion

- The SM is a theory, based on principles, that describes a huge amount of facts about microphysics. The hope that the gauge principle could take us much further than the SM, however, has not been realized.
- Neutrinos & weak interactions have been of key relevance for SM foundation and for demonstrating its incompleteness. Their study could provide further intellectual access on what is matter.
- There are several interesting experiments ongoing, there are many things still to understand, there is room for big surprises. It seems important to continue to work on observational cosmology. Perhaps we should put more efforts in building truly predictive theories.
- Is it so correct to decide a priori which are the most fundamental scientific goals and which are the more earthly ones instead, to be put aside? History has shown that this attitude is not w/o dangers.

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Matter

From Wikipedia, the free encyclopedia

This article is about the concept in the physical sciences. For other uses, see Matter (disambiguation).

In classical physics and general chemistry, **matter** is any substance that has mass and takes up space by haultimately composed of atoms, which are made up of interacting subatomic particles, and in everyday as we made up of them, and any particles (or combination of particles) that act as if they have both rest mass and photons, or other energy phenomena or waves such as light or sound.^{[1][2]} Matter exists in various states (al as solid, liquid, and gas – for example water exists as ice, liquid water, and gaseous steam – but other states fermionic condensates, and quark–gluon plasma.^[3]

Usually atoms can be imagined as a nucleus of protons and neutrons, and a surrounding "cloud" of orbiting correct, because subatomic particles and their properties are governed by their quantum nature, which mear like waves as well as particles and they do not have well-defined sizes or positions. In the Standard Model o elementary constituents of atoms are quantum entities which do not have an inherent "size" or "volume" in a other fundamental interactions, some "point particles" known as fermions (quarks, leptons), and many comp particles under everyday conditions; this creates the property of matter which appears to us as matter taking

For much of the history of the natural sciences people have contemplated the exact nature of matter. The ide particulate theory of matter, was first put forward by the Greek philosophers Leucippus (~490 BC) and Demo

 Contents [hide]

 1 Comparison with mass

 2 Definition

 2.1 Based on atoms

 2.2 Based on protons, neutrons and electrons

 2.3 Based on quarks and leptons

 2.4 Based on elementary fermions (mass, volume, and space)



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manson with mass

2 Definition

2.1 Based on atoms

- 2.2 Based on protons, neutrons and electrons
- 2.3 Based on quarks and leptons
- 2.4 Based on elementary fermions (mass, volume, and space)

- In ordinary language we say: matter does not disappear
- In this panel, the main reaction that allows the Sun to work, that means that *neutrinos are matter particles* (leptons)
- Shown physics says: numbers of baryons and leptons do not change
- Only B-L exactly conserved (SM)



64

what makes "matter"?

elementary components	identifying feature	time period theory (exp.)	reason of inadequacy
atoms	mass, type	till 1838 (1909)	"atoms of electricity"
nuclei & electrons	mass, charge	till 1933 (1956)	neutrons & neutrinos
p, n, e, v _e	B, L_{e,}, spin	till 1961 (1968)	quarks / SM
quarks, leptons	Β, L_e, L_µ, L_τ, spin	till 1967 (2010)	neutrino appearance
quarks-antileptons	B-L , spin	till 1937!!! (??)	Majorana mass/0v2β
fermions	spin	t!ll šššš (ššššš)	\$\$\$\$\$\$\$\$\$\$

C N Yang at the Centennial of MIT,1961

our present knowledge is sufficient to enable us to say with some certainty that great clarification will come in the field of weak interactions in the next few years. With luck on our side we might even hope to see some integration of the various manifestations of the weak interactions.

Beyond that we are on very uncertain grounds.

- ★ Is the continuum concept of space time extrapolatable to regions of space 10^{-14} cm to 10^{-17} cm, and to regions smaller than 10^{-17} cm?
- ★ What is the unifying basis of the strong, the electromagnetic and the weak interactions?
- \star What is the role of the gravitational field relative to all these?

If it is difficult to locate singularities of functions by extrapolation, it is as difficult to predict revolutionary changes in physical concepts by forecasting.

1st elementary matter particle (Thomson)

Hertz showed, however, that cathode particles possess another property which seemed inconsistent with the idea that they are particles of matter, for he found that they were able to penetrate very thin sheets of metal, e.g.

to the conclusion that the mass of the corpuscle is only about 1/1,700 of that of the hydrogen atom. Thus the atom is not the ultimate limit to the subdivision of matter; we may go further and get to the corpuscle, and at this stage the corpuscle is the same from whatever source it may be derived.

The corpuscle appear to form a part of all kinds of matter under the most diverse conditions; it seems natural therefore to regard it as one of the bricks of which atoms are built up.

universal V-A

Ruderman+Finkelstein 1949

Predictions of $R(pi \rightarrow enu)/R(pi \rightarrow munu)$ in various hypotheses

Durbin+Loar+Steinberger 1951

pion parity from deuterium photodissociation

Lokathan+Steinberger 1955 & Anderson+Lattes 1957

apparently $R(pi \rightarrow enu)$ is just absent, ruling out standard model (!)

Sudarshan+Marshak 1957 & Feynman+Gell-Mann 1958

V-A theory argues that previous result is inaccurate

Fazzini et al.1958 R(pi→enu)/R(pi→munu) confirmsV-A

 $(A, Z) \rightarrow (A, Z+1)+e^{-}+\bar{v}_{e}$ All OK $(A, Z) \rightarrow (A, Z+2)+2 e^{-1}$ L changes $n \rightarrow e^{-} + p$ L and spin change $p \rightarrow e^+ + \pi^0$ B and L change $n \rightarrow e^{-} + K^{+}$ B and L change $e^{-} \rightarrow \gamma + \nu_{e}$ **Charge changes** $n \rightarrow \bar{n}$ **B** changes



Figure 1: Sensitivity to $p \to K\bar{\nu}$; syst. not included. Water, $\epsilon = 14.6\%$ and b = 14/(Mton y) (2 methods, summed); Argon, $\epsilon = 97\%$ and b = 1/(Mton y). Impact of stat. fluctuations ≈ 2 .

how oscillations work

Fach neutrino is produced and observed as a mixture of several waves, which describe particles with different masses. Particles with different masses have different speeds, thus neutrinos transform as they propagate.



what we know on masses and mixings




Normal hierarchy → Normal ordering

Pee=0.7, 0.5, 0.3 through the Earth (La Thuile 2003)





Normal ordering → **Yearningly Expected Spectrum**

the "electron neutrino" mass

If the mass of the light v leads the transition, e.g. if new physics is at ultra-HE scale, the parameter that counts for $0v2\beta$ is,

$$m_{\beta\beta} \equiv |(M_{\nu})_{ee}| = \left| \sum_{i=1}^{3} |U_{ei}^{2}| e^{i\xi_{i}} m_{i} \right|$$

Symbols: first is the traditional one; second, ee-element of the ν mass matrix

The absolute mass scale and the (Majorana) **phases** ξ_i are not probed by oscillations: Only mass differences and electronic mixing $|U_{ei}|^2$ are measured.

STERILE NEUTRINOS?

◎LSND-MiniBooNe anomaly is with us since 1995; in latest data is mostly/only at lowest energies. 10 year later sterile neutrino were assessed by a global analysis w/o finding them

◎Ga- & reactor anomalies further tested with movable detector close to reactors. Not supported by DANSS, Stereo, NEOS

Adding sterile neutrinos does not help: the ensuing theory is predictive, leading to inconsistencies between these and other data



Available online at www.sciencedirect.com

Nuclear Physics B 708 (2005) 215-267



Probing oscillations into sterile neutrinos with cosmology, astrophysics and experiments

M. Cirelli^a, G. Marandella^b, A. Strumia^c, F. Vissani^d



Fig. 13. The LSND anomaly interpreted as oscillations of 3 + 1 neutrinos. Shaded region: suggested at 99% C.L. by LSND. Black dotted line: 99% C.L. global constraint from other neutrino experiments (mainly Karmen, Bugey, SK, CDHS). Continuos red line: $N_{\nu} = 3.8$ thermalized neutrinos. Dot-dashed orange line: $\Omega_{\nu}h^2 = 0.01$.

III. Sterile neutrino models and ν_{μ} disappearance

(3+1): tension among data samples

- Limits on ν_e → ν_e and ν_μ → ν_μ disappearance imply a bound on the ν_μ → ν_e appearance probability;
- such bound is stronger than what is required to explain the LSND and MiniBooNe excesses [A];
- hence, severe tension arises between **APP** and **DIS** data: $\chi^2_{PG}/dof = 29.6/2 \Rightarrow PG = 3.7 \times 10^{-7}$ [17];
- a similar result is visible when comparing " v_e -data" ($v_e \rightarrow v_e$ and $v_\mu \rightarrow v_e$) and " v_μ -data" ($v_\mu \rightarrow v_\mu$) [B];
- note: tension between APP and DIS data first pointed out in 1999 [34]. Full global fit in 2001 [35] cornered (3+1) models. No conceptual change since then...

[17] M. Dentler *et al.*, arXiv:1803.10661.
[34] S.M. Bilenky *et al.*, PRD **60** (1999) 073007 [hep-ph/9903454].
[35] MM, Schwetz, Valle, PLB **518** (2001) 252 [hep-ph/0107150].





direct search of big-bang neutrinos

big-bang neutrinos produce 3 **neutrinocapture** lines for a radioactive target

their positions depend on m_{i} ; their intensity on $|\,U_{ei}^{2}|$

lightest neutrino gives the most intense line for normal hierarchy

Needs

- great energy resolution
- ➢ big target mass, ≥100g of tritium



Glancing beyond SM

- High dim. operators, invariant under SM symmetry, summarize new physics at ultra-high scales
- > (They play exactly the same role of Fermi interactions)
- > The one with lowest dimension describes Majorana neutrino masses
- Oscillations are matched by a huge mass, say, of GUT

$$m_{\rm overall}^{\nu} \sim \frac{M_W^2}{M_{\rm gut}} = 65~{\rm meV} \times \frac{10^{14}~{\rm GeV}}{M_{\rm gut}}$$

an explanation of small of neutrino masses



an explanation of small of neutrino masses



this is called "seesaw"









Bajc et al 2005; Bertolini et al 2009-2011; Joshipura et al 2011; Buccella et al 2012; Dueck et al 2013; Altarelli et al 2013; Ohlsson et al 2019

minimal SO(10)

(principled model)

- ★ 16-plet coupled to 10 and 126 higgs: heavy right-handed neutrinos
- ★ (Peccei Quinn symmetry to address strong CP and dark matter)



Figure 2: Evolution of the gauge coupling constants in a GUT model with intermediate scale. Here, $M_{\rm interm.} \approx 5 \times 10^{13}$ GeV.

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Be

201 201

al 2009-Buccella 3; Altarelli



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- \star normal mass hierarchy; m_{$\beta\beta$} in the few meV range



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- ★ 16-plet coupled to 10 and 126 higgs: heavy right-handed neutrinos
- ★ (Peccei Quinn symmetry to address strong CP and dark matter)
- ★ neutrinos are massive and fermion masses constrained
- ★ normal mass hierarchy; m_{BB} in the few meV range
- (potentially interesting expectations for proton decay)











June 07, 2019







Figure 1: Diagrams contributing to the vertex (Fig. 1a) and wave function (Fig. 1b) CP violation in the heavy singlet neutrino decay.

Covi et al. '96

by the SM effects described previously

Do experiments suggest a hierarchy problem?

Francesco Vissani International Centre for Theoretical Physics, Strada Costiera 11, I-34013 Trieste, Italy (Received 18 September 1997; published 14 April 1998)



FIG. 1. The Feynman diagram originating the corrections in Eq. (1); ν_R denotes the right-handed neutrino of mass M_R , $\ell_L = (\nu_L, e_L)$ the leptonic and H the Higgs doublets.



the standard model. So (if a joke is permitted) we present a prediction for LHC:

$$m_H = 135 \pm 5 \text{ GeV}$$
 (14)

the reason is that this value will increase the entropy in the minds of several theorists. Note, however, that the decay of a standard and supersymmetric HIGGS particles with the same mass (or also the production rate–"cross-sections") could be rather different; thence, these measurements would offer a possibility to distinguish between the SM and its supersymmetric extension even in this tricky case.

3 (Not quite a) conclusion

We would like to close this pages by spending few words of caution, to remind that failures of the standard model have been often claimed in the past years (today, several of them are considered dubious or simply wrong tracks). Here is an arbitrary selection:

THEORETICAL INTERPRETATION

EXPERIMENTAL ANOMALY

leptoquark	 High x and Q^2 events at HERA
compositeness	 Excess of 4-jet events at ALEPH
light gravitino	 $ee\gamma\gamma \not\!$
17 keV neutrino	 bump in β spectra (SIMPSON,)
monopole	 induced currents (CABRERA)
proton decay	 contained multitrack events at KGF

Is there any moral behind these stories? Maybe not; however:

 they suggest to go slowly and carefully from data to theories and back (because of possible pitfalls of interpretation, of suggestion, etc.);

2) they witness how hard is to reach the frontiers of standard model; and, also, how strong is the desire of particle physicists to find them! ЕЗОЛТЯАТ УЯАТИЗМАЛА ТО ЛЭДОМ ОЯАОИАТЕ ЗНТ 21 ТАНW ВАЛОГАРАТ 21 ТАНW И ОИА ПАЛОГА 11. В 10 МОНТ 11 ТО МОЛОГА 11 ГО МОЛОГА 11.

$$R_{\rm s} \equiv \frac{2G_{\rm N}M}{c^2}, \ M \equiv \frac{4\pi}{3}R_{\rm s}^3 \times \rho \Rightarrow$$

$$\rho \equiv \frac{3 c^2}{8\pi G_{\rm N}R_{\rm s}^2}; \ \text{imposing } \rho < \rho_{\rm crit.} \Rightarrow$$

$$R_{\rm s} < \frac{c}{H_0} \equiv R_{\rm H} = 4.4 \text{ Gpc}$$



gut health - latest news, breaking stories a. independent.co.uk

How gut bacteria may help to spot and add... medicalnewstoday.com