SU(3) Generators and Gauge Fields

$$Gell - Ram matrices:$$

$$\lambda_{1} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \quad \lambda_{2} = \begin{pmatrix} 0 & -i & 0 \\ 0 & 0 & 0 \end{pmatrix}, \quad \lambda_{3} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix},$$

$$\lambda_{4} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \quad \lambda_{5} = \begin{pmatrix} 0 & 0 & -i \\ 0 & 0 & 0 \end{pmatrix}, \quad \lambda_{6} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix},$$

$$\lambda_{7} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & -i \end{pmatrix}, \quad \lambda_{8} = \frac{1}{(3)} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & -2 \end{pmatrix}$$

$$gamgs fields in Su(3)_{W}:$$

$$\frac{1}{2} A_{1}^{0} \lambda_{1}^{0} = \frac{1}{2} \begin{pmatrix} A_{1}^{0} + A_{1}^{0} & A_{1}^{0} - i A_{1}^{0} & A_{1}^{0} - i A_{1}^{0} \\ A_{1}^{0} + i A_{1}^{0} & A_{1}^{0} - i A_{1}^{0} & A_{1}^{0} - i A_{1}^{0} \\ A_{1}^{0} + i A_{1}^{0} & A_{1}^{0} + i A_{1}^{0} & -2 A_{2} A_{1}^{0} \end{pmatrix}$$

Derivation of Effective Potential

Calculation of 1-loop effective polanhal [according to Antoniadia et al.
take

$$Veqt(\Phi) = \frac{1}{2} \sum_{n=1}^{\infty} (-1)^{\frac{n}{2}} \int_{(2\pi)^{n}}^{\frac{1}{2}} \log \left[p^{2} + \Pi_{2}^{+}(\Phi) \right]$$
Schwinger representation

$$Veqt(\Phi) = -\frac{1}{2} \sum_{n=1}^{\infty} (-1)^{\frac{n}{2}} \int_{0}^{\infty} \frac{dt}{dt} \int_{(2\pi)^{n}}^{\frac{1}{2}} e^{-\frac{1}{2} \left[p^{2} + \Pi_{2}^{+}(\Phi) \right]}$$

$$= \frac{1}{32\pi^{n}} \sum_{n=1}^{\infty} (-1)^{\frac{n}{2}} \int_{0}^{\infty} \frac{dt}{dt^{2}} e^{-\frac{1}{2} \left[n^{\frac{1}{2}} + \Pi_{2}^{+}(\Phi) \right]}$$

$$= \frac{1}{32\pi^{n}} \sum_{n=1}^{\infty} (-1)^{\frac{n}{2}} \int_{0}^{\infty} \frac{dt}{dt^{2}} e^{-\frac{1}{2} \left[n^{\frac{1}{2}} + \Pi_{2}^{+}(\Phi) \right]}$$

$$= \frac{1}{32\pi^{n}} \sum_{n=1}^{\infty} (-1)^{\frac{n}{2}} \int_{0}^{\infty} \frac{dt}{dt} R = \frac{1}{n^{\frac{n}{2}}} \left[n^{\frac{1}{2}} + \frac{1}{n^{\frac{n}{2}}} \int_{0}^{\infty} \frac{dt}{dt} R = \frac{1}{n^{\frac{n}{2}}} \left[n^{\frac{1}{2}} + \frac{1}{n^{\frac{n}{2}}} \int_{0}^{\infty} \frac{dt}{dt} R = \frac{1}{n^{\frac{n}{2}}} \left[n^{\frac{1}{2}} + \frac{1}{n^{\frac{n}{2}}} \int_{0}^{\infty} \frac{dt}{dt} R = \frac{1}{n^{\frac{n}{2}}} \left[n^{\frac{1}{2}} + \frac{1}{n^{\frac{n}{2}}} \int_{0}^{\infty} \frac{dt}{dt} R = \frac{1}{n^{\frac{n}{2}}} \left[n^{\frac{1}{2}} + \frac{1}{n^{\frac{n}{2}}} \int_{0}^{\infty} \frac{dt}{dt} R = \frac{1}{n^{\frac{n}{2}}} \left[n^{\frac{1}{2}} + \frac{1}{n^{\frac{n}{2}}} + \frac{1}{n^{\frac{n}{2}}} \int_{0}^{\infty} \frac{dt}{dt} R = \frac{1}{n^{\frac{n}{2}}} \left[n^{\frac{n}{2}} + \frac{1}{n^{\frac{n}{2}}} + \frac{1}{n^{\frac{n}{2}}} \int_{0}^{\infty} \frac{dt}{dt} R = \frac{1}{n^{\frac{n}{2}}} \left[n^{\frac{n}{2}} + \frac{1}{n^{\frac{n}{2}}} + \frac{1}{n^{\frac{n}{2}}} \int_{0}^{\infty} \frac{dt}{dt} R = \frac{1}{n^{\frac{n}{2}}} \left[n^{\frac{n}{2}} + \frac{1}{n^{\frac{n}{2}}} +$$

)

Effective Potential



Figure 1: Different contributions to the effective potential (in arbitrary units): the bulk and boundary fermion contributions (upper left) and the full potential (upper right) for $\lambda = 1.57$, $\epsilon_1 = 3.1$, $\epsilon_2 = 0.7$ and $\delta = 0$; the bulk and boundary fermion contributions (lower left) and the full potential (lower right) for $\lambda = 1.83$, $\epsilon_1 = 6.4$, $\epsilon_2 = 6.1$ and $\delta = 1$. Taken from *C.A. Scrucca et al.*.



Figure 2: The full effective potential (in arbitrary units) in the presence of high-rank bulk fermions. Left: r = 8, $\lambda = 3.47$, $\epsilon_1 = \epsilon_2 = 9$ and $\delta = 0$, resulting in $m_H = 112$ GeV and 1/R = 830 GeV. Right: r = 6, $\lambda = 2.23$, $\epsilon_1 = 7$, $\epsilon_2 = 1$ and $\delta = 1$, resulting in $m_H = 104$ GeV and 1/R = 600 GeV. Taken from *C.A. Scrucca et al.*.

... with localized gauge kinetic terms



Figure 3: Gauge contribution to the effective potential (in arbitrary units) in the presence of localized gauge kinetic terms, with $c_2 = 0$ and increasing values of c_1 . Taken from *C.A. Scrucca et al.*.





Figure 4: Plot of the Higgs potential (in arbitrary units) from the gauge bosons (dashed-red), tau and top (blue), twisted fermions (dashed-green), and the total (thick black), for one light generation with $\kappa_l = 3$. The other parameters are like described in the text. Taken from *G. Cacciapaglia et al.*.



Figure 5: Plot of the Higgs potential (in arbitrary units) from the gauge bosons (reddashed), top (blue), bottom (3) and tau (10) (green-dashed), and the total (thick black), for $\epsilon = 1.25$.

Taken from G. Cacciapaglia et al..

Effective Potential with SO(4,1) Lorentz violation (and antiperiodic fermions)



Figure 6: The Higgs potential in the $\delta = 0$ set-up, obtained with input parameters $\lambda^t = 0.99$, $\lambda^b = 6.9$, $\lambda^A = 0.24$, $k_t = 2.42$, $k_b = 2.26$, $k_A = 3.14$, $\varepsilon_1^t = 1.9$, $\varepsilon_2^t = 1.6$, $\varepsilon_1^b = 2.9$, $\varepsilon_2^b = 3.4$. Taken from *G. Panico et al.*.

χ^2 fit on the EWPT for $SU(3)_w$ model with Z_2 symmetry

Taken from G. Panico et al..



Figure 7: Constraints coming from a χ^2 fit on the EWPT. The contours represent the allowed regions in the $(1/R, m_H)$ plane at 68%, 90% and 99% confidence level (2 d.o.f.). The shaded band shows the experimentally excluded values for the Higgs mass ($m_H < 115$ GeV). The blue dots represent the predictions of our model for different values of the microscopic parameters (only points with the correct top and bottom masses are plotted).

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