

# Nonequilibrium QFT approach to leptogenesis

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PT  
FS FOR PRECISION TESTS  
OF FUNDAMENTAL  
SYMMETRIES

Based on [arXiv:1211.2140](https://arxiv.org/abs/1211.2140) in collaboration with M. Garny, A. Hohenegger, A. Kartavtsev, and D. Mitrouskas

- Baryogenesis
- Leptogenesis
- Boltzmann approach to leptogenesis
- Nonequilibrium QFT approach to leptogenesis
- Conclusion

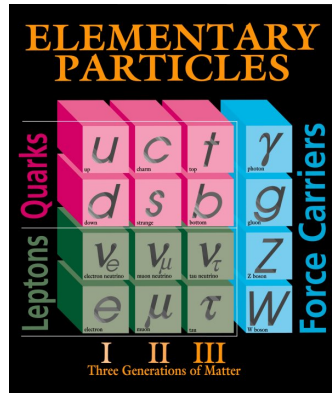
# Baryogenesis

- The Earth is made of matter
- The Sun is made of matter
- The Milky Way is made of matter
- The Universe is made of matter



The Universe contains almost only matter, and no antimatter!

The Laws of Nature are *almost* the same for matter and for antimatter!

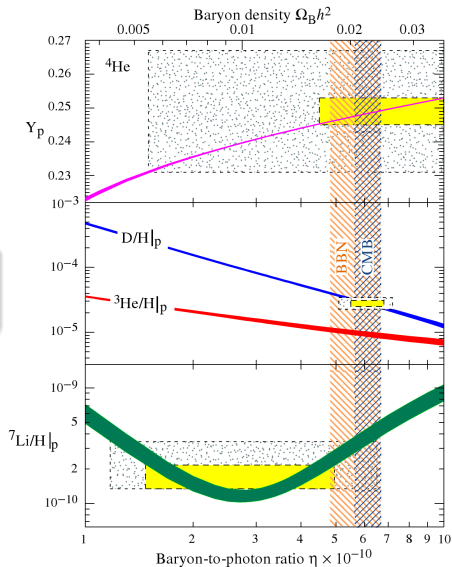


# Baryon-to-photon ratio

$$\eta = \frac{n_B - n_{\bar{B}}}{n_\gamma} \sim 6 \times 10^{-10}$$

$t_{\text{BBN}} \sim 1$  minute

$t_{\text{CMB}} \sim 380\,000$  years



[astro-ph/0601514]

# Sakharov conditions

Two different attitudes:

- The Universe contains initially more particles than antiparticles  
→ no need to produce the asymmetry
- The Universe was initially matter-antimatter **symmetric**, and the observed asymmetry must be dynamically produced

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## 3 Sakharov conditions

- Baryon number violation
- $C$  and  $CP$  violation
- out of equilibrium dynamics

## Standard Model

- ↔ sphalerons
- ↔ chirality, CKM matrix
- ↔ expanding Universe

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**Beyond SM physics:**

- Electroweak baryogenesis, Affleck-Dine baryogenesis, **leptogenesis**, ...

# Leptogenesis

## Leptogenesis: particular model of baryogenesis

- **very simple model**: SM particles +  $n$  (at least 2) heavy, gauge-singlet right-handed neutrinos  $N_i$
- RH neutrinos are unstable
- main decay channels:
  - lepton-Higgs pair
  - antilepton-antiHiggs pair
- $CP$  violation  $\rightarrow$  decay rates are unequal
- in an expanding Universe  $\rightarrow$  production of **lepton asymmetry**
- EW sphaleron processes transfer part of the lepton asymmetry to the baryon sector

$\Rightarrow$  Production of a net baryon asymmetry!

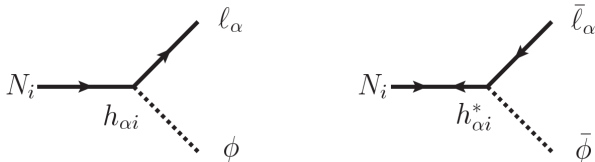


# Type-I seesaw Lagrangian

SM +  $n$  heavy right-handed gauge-singlet Majorana neutrinos:

$$\mathcal{L} = \mathcal{L}_{SM} + \bar{N}_i i \not{\partial} N_i - \frac{1}{2} M_i (\bar{N}_i^c N_i + \bar{N}_i N_i^c) - h_{\alpha i} \bar{\ell}_\alpha \tilde{\phi} N_i - h_{i\alpha}^\dagger \bar{N}_i \tilde{\phi}^\dagger \ell_\alpha$$

- the mass matrix violates lepton number
- Yukawa couplings  $h$  generate the vertices

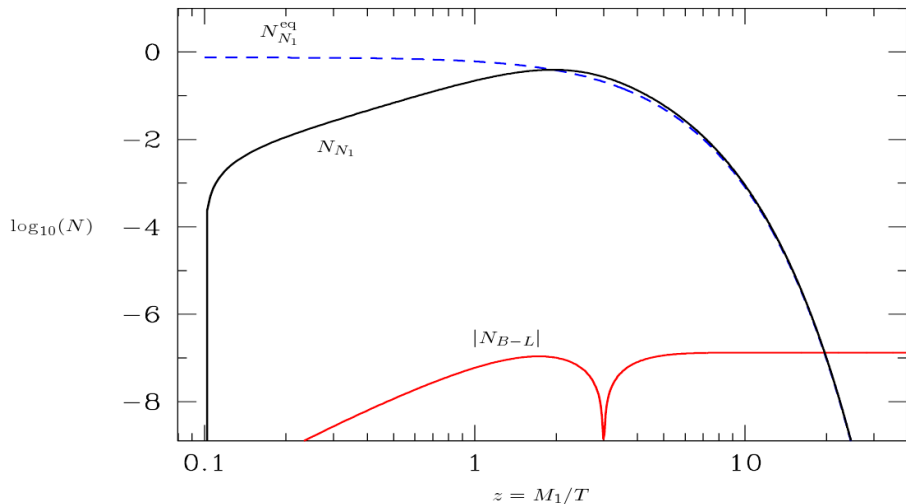


- gives **mass to active neutrinos** through the seesaw formula

$$m_\nu = -m_D M^{-1} m_D, \quad \text{where} \quad m_D = v h$$

- “heavy” means  $\mathcal{O}(10^{10} - 10^{13})$  GeV  
 $\Rightarrow$  no hope to see these particles in future experiments

# Evolution of the asymmetry



[hep-ph/0502169]

# Transition amplitudes

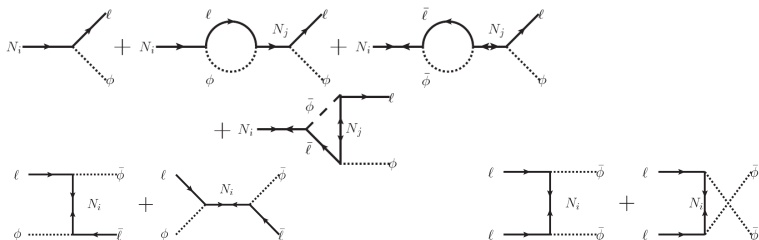
Generation of the asymmetry:

	production	washout
$N_i$ (inverse)decay	$\mathcal{O}(h^4)$	$\mathcal{O}(h^2)$
$\Delta L = 2$ scattering	$\mathcal{O}(h^6)$	$\mathcal{O}(h^4)$
top scattering	$\mathcal{O}(\lambda_t^2 h^4)$	$\mathcal{O}(\lambda_t^2 h^2)$
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$\Rightarrow$  into the **Boltzmann** equation for the leptons and antileptons!

# Boltzmann equation

$$\begin{aligned}
 \frac{df_\ell}{dt} = & \frac{1}{2E_\ell} \sum_{\{a\dots\},\{i\dots\}} \int \frac{d^3p_a}{2E_a} \dots \frac{d^3p_i}{2E_i} \dots (2\pi)^4 \delta(\sum p_a - \sum p_i - p_\ell) \\
 & \times f_a \dots (1 \pm f_i) \dots (1 - f_\ell) |\mathcal{M}_{a+\dots \rightarrow i+\dots+\ell}|^2 \\
 - \frac{1}{2E_\ell} \sum_{\{a\dots\},\{i\dots\}} & \int \frac{d^3p_a}{2E_a} \dots \frac{d^3p_i}{2E_i} \dots (2\pi)^4 \delta(\sum p_a + p_\ell - \sum p_i) \\
 & \times f_a \dots f_\ell (1 \pm f_i) \dots |\mathcal{M}_{a+\dots+\ell \rightarrow i+\dots}|^2
 \end{aligned}$$

+ similar equation for antilepton

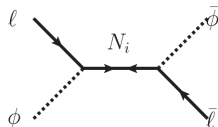
In **equilibrium**,  $\frac{d(f_\ell - f_{\bar{\ell}})}{dt}$  is non zero if one uses the Feynman amplitudes!

# Double counting problem

BE equations + *naive* Feynman amplitudes are **not consistent** with the third Sakharov condition

→ unstable particles as in/out-state

- some processes are counted twice



- can be solved if one neglects quantum statistical factors:

$$(1 \pm f_i)(1 \pm f_j) \dots \rightarrow 1$$

- subtraction of the real intermediate state from the s-channel

$$|\mathcal{M}_{\ell\phi \rightarrow \bar{\ell}\bar{\phi}}|^2:$$

→ so-called **RIS-subtraction**

# First principles approach

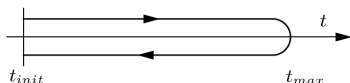
- RIS-subtraction is an **unsatisfactory** solution
- derive Boltzmann-like equations from **first principles**
  - free of double counting problem
  - include thermal masses
  - include thermal amplitudes
  - include thermal width
- keep track of the approximations done
  - range of validity of the resulting equations

⇒ Nonequilibrium quantum field theory approach!

# Basics of nonequilibrium QFT

Lepton propagator on CTP:

$$S_{\alpha\beta}(x, y) = \langle T_C l_\alpha(x) \bar{l}_\beta(y) \rangle$$

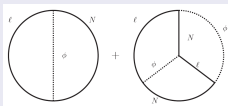


Closed Time Path

## Schwinger-Dyson equation

$$\hat{S}^{-1}(x, y) = \hat{S}_0^{-1}(x, y) - \hat{\Sigma}(x, y), \quad \Sigma_{\alpha\beta}(x, y) = -\frac{\delta i\Gamma_2}{\delta S_{\beta\alpha}(y, x)}$$

$\Gamma_2 =$  2PI effective action  $=$



$$\rightarrow \text{doubling of the d.o.f.: } \hat{S}(x, y) = \underbrace{\hat{S}_F(x, y)}_{\text{statistical}} - \frac{i}{2} \text{sign}_C(x^0 - y^0) \underbrace{\hat{S}_\rho(x, y)}_{\text{spectral}}$$



# Kadanoff-Baym equations

## Kadanoff-Baym equations

$$i\partial_x \hat{S}_F(x, y) = \int_0^{x^0} d^4 z \hat{\Sigma}_\rho(x, z) \hat{S}_F(z, y) - \int_0^{y^0} d^4 z \hat{\Sigma}_F(x, z) \hat{S}_\rho(z, y)$$

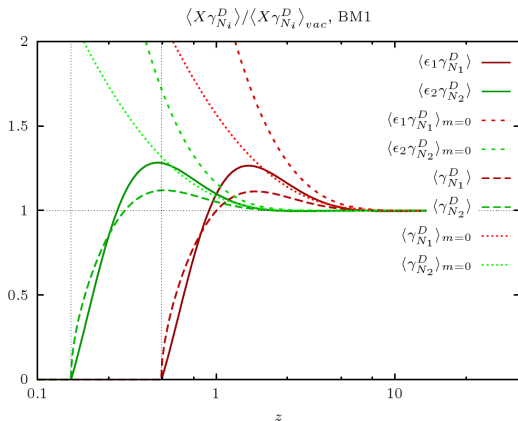
$$i\partial_x \hat{S}_\rho(x, y) = \int_{y^0}^{x^0} d^4 z \hat{\Sigma}_\rho(x, z) \hat{S}_\rho(z, y)$$

- equivalent to the Schwinger-Dyson equation
- **exact** equations
- only Gaussian initial conditions
- self-energies are functional of the propagators
  - need a loop expansion
- extremely difficult to solve numerically
  - **approximations** needed!



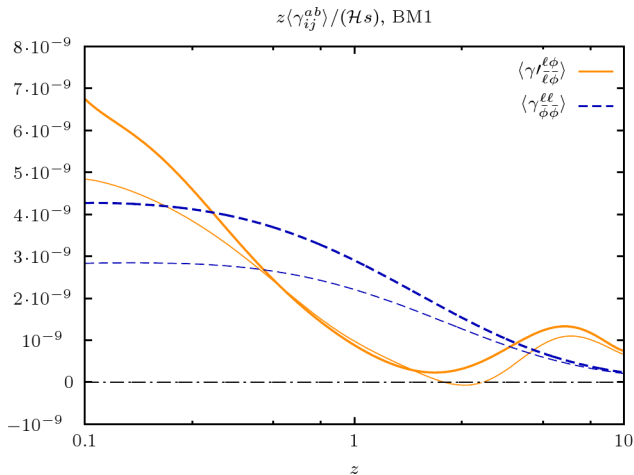
## Results: RH neutrino decay

- $z = M_1/T$
- $\epsilon_i = \frac{|\mathcal{M}_{N_i \rightarrow \ell \phi}|^2 - |\mathcal{M}_{N_i \rightarrow \bar{\ell} \bar{\phi}}|^2}{|\mathcal{M}_{N_i \rightarrow \ell \phi}|^2 + |\mathcal{M}_{N_i \rightarrow \bar{\ell} \bar{\phi}}|^2}$
- $\langle \dots \gamma_{N_i}^D \rangle$ : reaction density  
 $\sim$  thermally averaged amplitudes



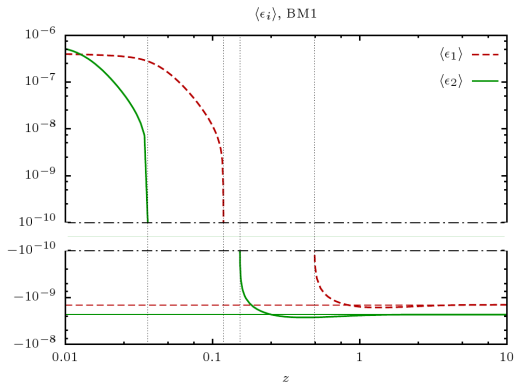
$\Rightarrow$  thermal enhancement partially compensated by thermal masses!

## Results: scattering process



# Results: Higgs decay

- due to thermal masses the **phase space** for heavy neutrino **shrinks** at high temperature
- at even higher temperature the Higgs decay becomes **kinematically allowed**



# Conclusion

- Nonequilibrium QFT is the appropriate tool to study leptogenesis
- lead to consistent set of equations
  - free of double counting problem
  - include thermal masses
  - include thermal corrections to the amplitudes
- total amplitudes are not very sensitive to thermal corrections
  - tree level
- CP-violating parameter is more sensitive to thermal corrections
  - loop corrections
- inclusion of top scattering (work in progress)
- inclusion of gauge scattering (more complicated ...)

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Thank you for your attention!