

# Neutrino Decoherence

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# Outline

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- ✿ Quantum Decoherence and Neutrino Oscillation
- ✿ The Layer Structure
- ✿ Neutrino Decoherence
- ✿ Experimental Potentials

# Neutrino Oscillation (Quantum Coherence)

- ✿ Quantum coherence: Superposition of quantum states
- ✿ Neutrino oscillation:  
superposition of mass eigenstates in the measurement (flavor) basis

Density Matrix formalism (QM)

$$|\nu_\alpha\rangle = \sum_j U_{\alpha j} |\nu_j\rangle \quad \text{U: PMNS matrix}$$

$$\rho(t) = e^{iHt} \rho(0) e^{-iHt} = \sum_{jk} U_{j\delta} U_{k\delta}^* e^{-i(H_j - H_k)t} |\nu_j\rangle\langle\nu_k|$$

$$P_{\alpha\beta} = \text{Tr} \left( |\nu_\alpha\rangle\langle\nu_\beta| \rho(t) \right)$$

# Coherent

## S-Matrix formalism (QFT)

$$|\nu_\alpha\rangle = \sum_j U_{\alpha j} |\nu_j\rangle$$

$$iA_{\alpha\beta} = \text{Diagram} \quad \supset \sum_j U_{\alpha j}^* U_{\beta j} \int \frac{d^4 p}{(2\pi)^4} \frac{\not{p} + m_j}{p^2 - m_j^2 + i\epsilon} e^{-ipx}$$

The diagram shows a vertex labeled  $\nu$  with two outgoing lines.

$$P_{\alpha\beta} = |A_{\alpha\beta}|^2$$

## Density Matrix formalism (QM)

$$|\nu_\alpha\rangle = \sum_j U_{\alpha j} |\nu_j\rangle$$

$$\rho(t) = e^{iHt} \rho(0) e^{-iHt} = \sum_{jk} U_{j\delta} U_{k\delta}^* e^{-i(H_j - H_k)t} |\nu_j\rangle \langle \nu_k|$$

$$P_{\alpha\beta} = \text{Tr} \left( |\nu_\alpha\rangle \langle \nu_\beta| \rho(t) \right)$$

# Decoherent

## Consider Wavepackets

→ Not completely coherent to start with

- ❖ Some possible causes:

- ❖ Wavepacket separation ( $v_j \neq v_k$ ) [1001.4815], [2104.05806], ...
- ❖ Mass state identification [0905.1903]

## Consider an Open Quantum System

→ Lose information to the environment

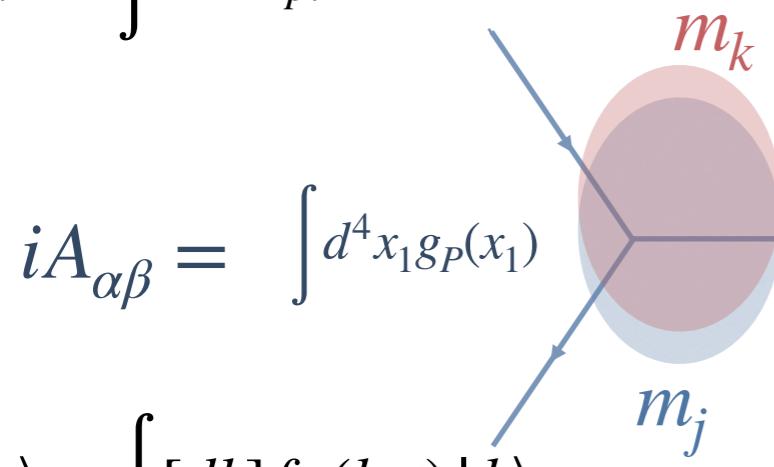
- ❖ Some possible causes:

- ❖ Interactions (matter effect, neutrino decay, ...) [1803.04438], ...
- ❖ Quantum gravity [hep-ph/0002053], [2007.00068], ...
- ❖ Cosmological [1704.06139]

# Decoherent

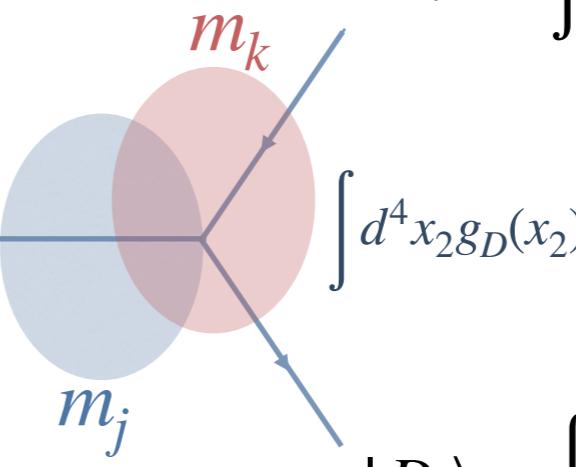
S-Matrix formalism (QFT) — Consider Wavepackets

$$|P_i\rangle = \int [dq] f_{pi}(q, t) |q\rangle$$



$$|P_f\rangle = \int [dk] f_{pf}(k, t) |k\rangle$$

$$|D_i\rangle = \int [dq'] f_{di}(q', t) |q'\rangle$$



$$|D_f\rangle = \int [dk'] f_{df}(k', t) |k'\rangle$$

Density Matrix formalism (QM) — Open Quantum System

Lindblad Equation:  $\frac{\partial \rho}{\partial t} = -i[H, \rho] - \mathcal{D}[\rho]$

$$\mathcal{D}[\rho] = -\frac{1}{2} \sum_k^{N^2-1} \left( [V_k, \rho V_k^\dagger] + [V_k \rho, V_k^\dagger] \right) \equiv \left( D_{\mu\nu} \rho^\nu \right) \lambda^\mu$$

Gell-Mann matrices  
 $9 \times 9$  entries

# Decoherent from the Layer Structure

Consider Wavepackets + QFT description

→ Not completely coherent to start with

Layer 1

State Decoherence

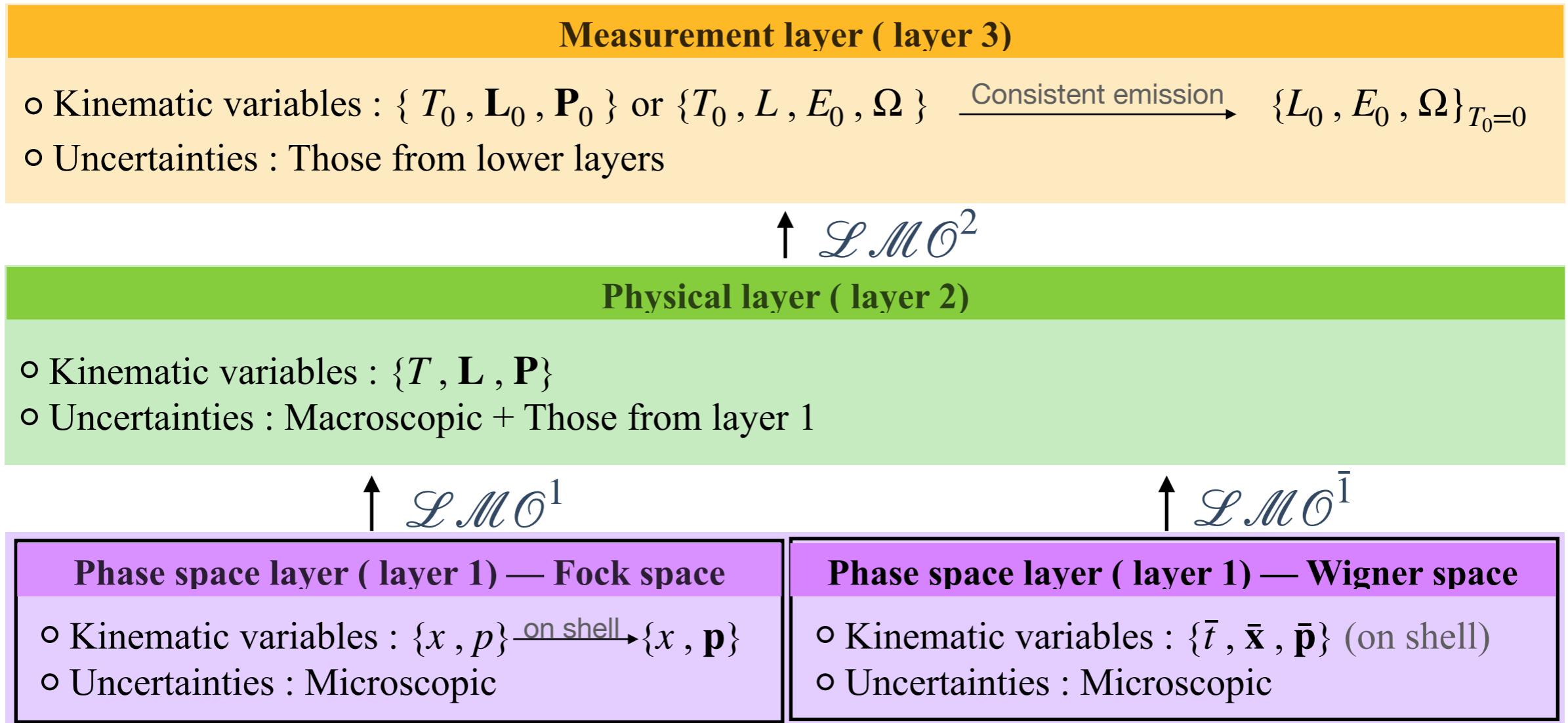
Consider an Open (Quantum) System

→ Lose information to the environment

Layer 1 + Layer 2

Phase Decoherence

# The Layer Structure



# The Layer Structure

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- ❖ The layer moving operator:

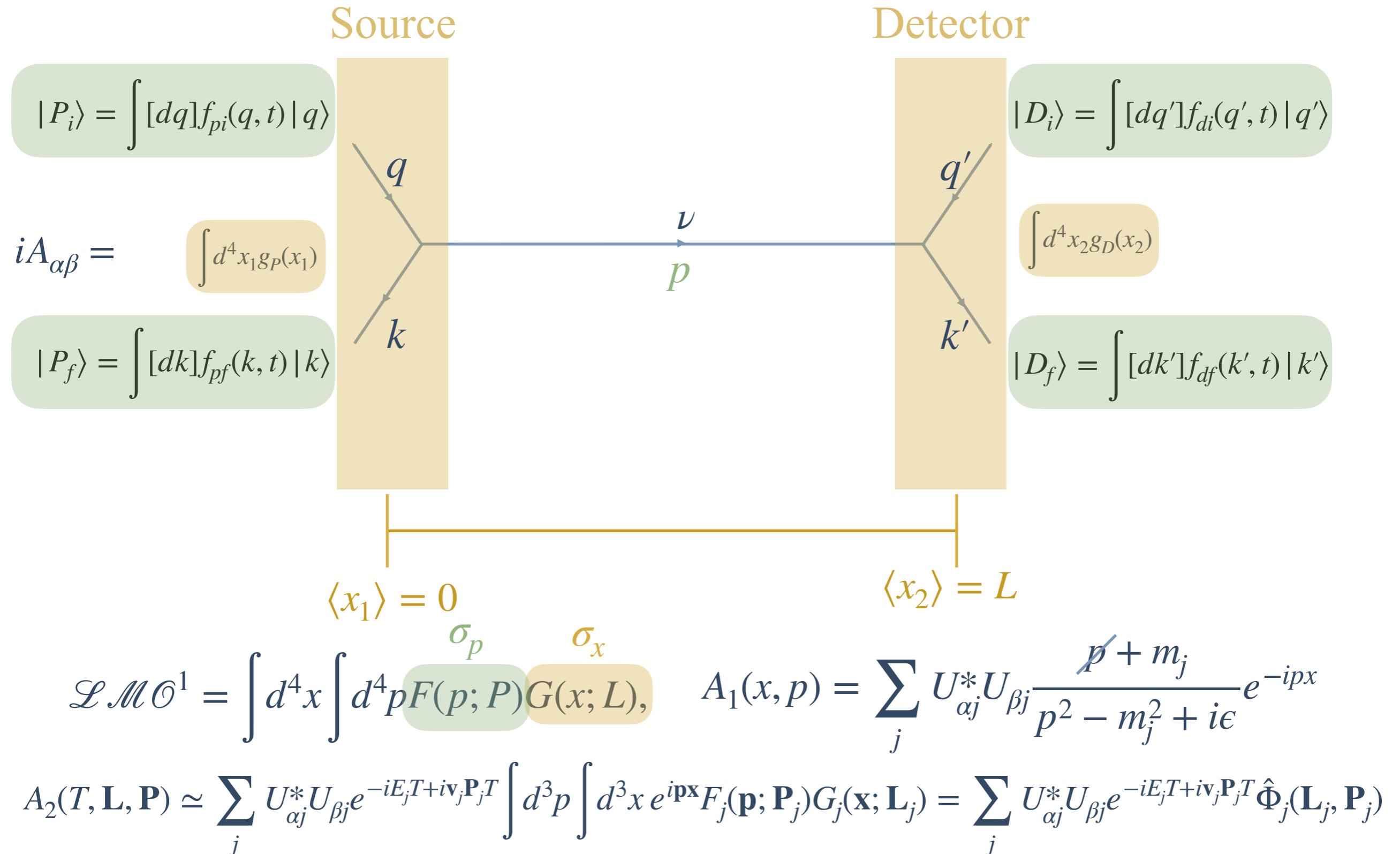
$$\mathcal{LMO}^i = \int d^n x_i \int d^m p_i W_i(x_i, p_i; x_{i+1}, p_{i+1}),$$

such that  $B_{i+1}(x_{i+1}, p_{i+1}) = \mathcal{LMO}^i B_i(x_i, p_i)$

- ❖ Flavor transition probability:

$$\begin{aligned} P_3(T_0, \mathbf{L}_0, \mathbf{P}_0) &= \mathcal{LMO}^2 P_2(T, \mathbf{L}, \mathbf{P}) = \mathcal{LMO}^2 \{ A_2^*(T, \mathbf{L}, \mathbf{P}) A_2(T, \mathbf{L}, \mathbf{P}) \} \\ &= \mathcal{LMO}^2 \{ \mathcal{LMO}^1 A_1^*(x', p') \mathcal{LMO}^1 A_1(x, p) \} \\ &= \mathcal{LMO}^2 \{ \mathcal{LMO}^{\bar{1}} P_{\bar{1}}(\bar{x}, \bar{p}) \}. \end{aligned}$$

# Layer 1: QFT Transition Amplitude $\sigma_p$ $\sigma_x$



# Layer $\bar{1}$ : Wigner Transition Probability $\sigma_p$ $\sigma_x$

Bridge of QM to statistical phase space

$$\mathcal{LMO}^1 A_1^*(x', p') \mathcal{LMO}^1 A_1(x, p) = \mathcal{LMO}^{\bar{1}} P_{\bar{1}}(\bar{x}, \bar{p}).$$

$$(x, x') \rightarrow (\bar{x} = \frac{1}{2}(x + x'), \Delta x = x - x')$$

$$\mathcal{LMO}^{\bar{1}} = \int d^3 \bar{x} \int d^3 \bar{p}, \quad P_{\bar{1}jk} = \tilde{W}_{jk}^G(\bar{x}, \bar{p}) \tilde{W}_{jk}^F(\bar{x}, \bar{p}),$$

$$\tilde{W}_{jk}^G(\bar{x}, \bar{p}) = \int d^3(\Delta x) e^{i\Delta x \bar{p}} G_j^*(\bar{x} - \frac{1}{2}\Delta x; \mathbf{L}_j) G_k(\bar{x} + \frac{1}{2}\Delta x; \mathbf{L}_k)$$

$$\tilde{W}_{jk}^F(\bar{x}, \bar{p}) = \int d^3(\Delta p) e^{i\Delta p \bar{x}} F_j^*(\bar{p} - \frac{1}{2}\Delta p; \mathbf{P}_j) F_k(\bar{p} + \frac{1}{2}\Delta p; \mathbf{P}_k)$$

Wigner quasi-probability distribution function

# Layer 2: Physical Transition Probability $\sigma_L$ $\sigma_E$

$$P_2(T, \mathbf{L}, \mathbf{P}) = A_2^*(T, \mathbf{L}, \mathbf{P})A_2(T, \mathbf{L}, \mathbf{P}) = \mathcal{LCM}^{\bar{1}} P_{\bar{1}}(\bar{\mathbf{x}}, \bar{\mathbf{p}}).$$

- Recall the first layer weighting function:  $W_1^j(\mathbf{x}, \mathbf{p}; T, \mathbf{L}, \mathbf{P}) = F_j(\mathbf{p}; \mathbf{P}_j)G_j(\mathbf{x}; \mathbf{L}_j)$
- Relation between  $\{T, \mathbf{L}, \mathbf{P}\}$  and  $\{\mathbf{P}_j, \mathbf{L}_j\}$ :  $|\mathbf{P}_j| \equiv E - \delta E_j$  and  $\mathbf{L}_j = \mathbf{L} - \mathbf{v}_j T$   
Energy budget from the external particles , i.e. relativistic neutrinos
- $\delta E_j \neq 0$  and  $\mathbf{v}_j \neq 1$  results from neutrinos carrying mass:

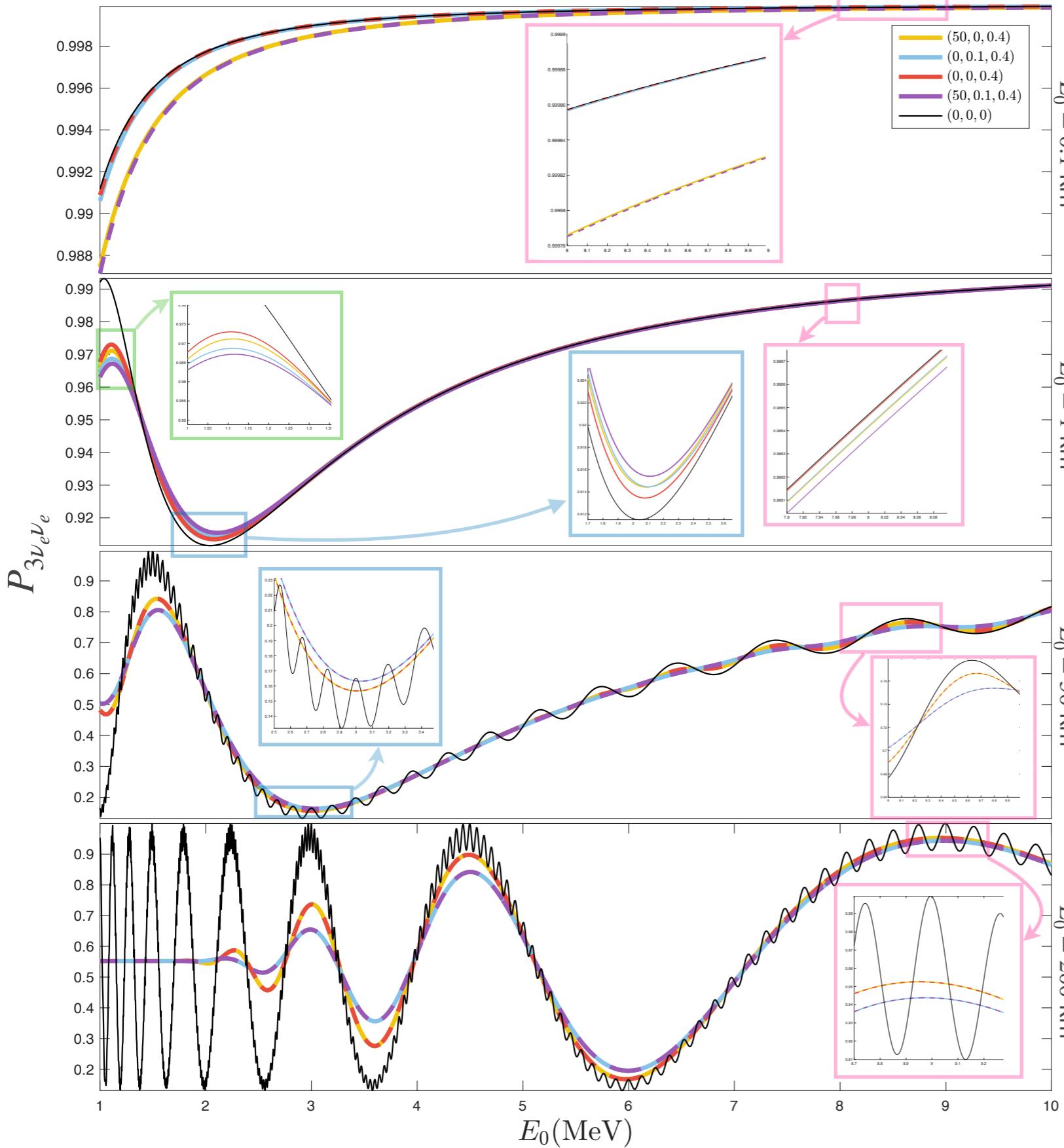
$$\mathbf{P}_j \simeq \mathbf{P} - \vec{\xi}_j \frac{m_j^2}{2E}, = \vec{\xi}_p E - \vec{\xi}_j \frac{m_j^2}{2E}, \text{ and } \mathbf{v}_j = \frac{\mathbf{P}_j}{E_j} \simeq \vec{\xi}_p \left( 1 - \frac{m_j^2}{2E^2} \right)$$

- Physical transition probability :  $P_2(T, \mathbf{L}, \mathbf{P}) \simeq \sum_{jk} U_{\alpha j}^* U_{\beta j} U_{\alpha k}^* U_{\beta k} e^{i \frac{\Delta m_{jk} L}{2E}} e^{-S_x^2 - S_p^2}$

With (state) decoherence terms:  $S_x = \frac{\Delta m_{jk} \sigma_x}{2\sqrt{2\Delta E}}$  and  $S_p = \frac{\Delta m_{jk} \sigma_p L}{2\sqrt{2\Delta E^2}} = \frac{L}{L_{jk}^{coh}}$ ,  $\Delta^2 = 1 + \sigma_p^2 \sigma_x^2$

- $\mathcal{LCM}^2 = \int dL \int dE \int dT \int d\Omega H_L(L; L_0) \overset{\sigma_L}{H_E(E; E_0)} \overset{\sigma_E}{H_T(T; T_0)} H_\Omega(\Omega; \Omega_0)$

# Layer 3: Measurement Transition Probability



$$P_3(T_0, \mathbf{L}_0, \mathbf{P}_0) = \mathcal{LMO}^2 P_2(T, \mathbf{L}, \mathbf{P})$$

Subdominant  
& absorbed into

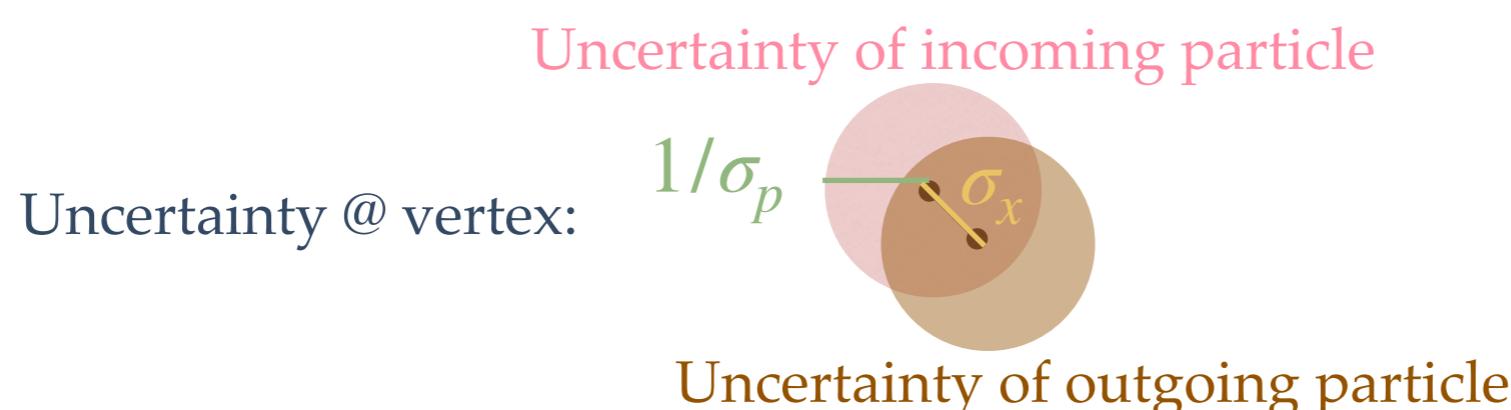
	$\sigma_x$	$\sigma_L$	$\sigma_E$	$\sigma_p$
—	X	0	X	X
—	0	X	X	X
—	0	0	0	X
—	X	X	X	X
—	0	0	0	0

- Dominance:  
Short Distance  $\rightarrow \sigma_L$   
Long Distance  $\rightarrow \sigma_E, \sigma_p$

# Uncertainties

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- ✿ Possible contributions to each uncertainties:
  - ✿  $\sigma_p$ : external particle WP size
  - ✿  $\sigma_x$ : non central collisions
  - ✿  $\sigma_E$ : energy resolution, reconstruction model
  - ✿  $\sigma_L$ : production profile, reactor core size, decay length, space fluctuation (e.g. quantum gravity)



# Neutrino Decoherence

- ❖ Operation definition of transition probability:

Neutrino oscillation experiment

$$P_{3jk}(X_3) = \frac{\int dX_2 H(X_2; X_3) \Gamma_{2jk}(X_2; X_3)}{\sqrt{\int dX_2 H(X_2; X_3) \Gamma_{2jj}(X_2; X_3)} \sqrt{\int dX_2 H(X_2; X_3) \Gamma_{2kk}(X_2; X_3)}},$$

Neutrino mass measuring experiment

$\approx$

$$\frac{\int dX_2 \Gamma_\alpha^{\text{pro.}}(X_2) P_{2\alpha \rightarrow \beta}(X_2; X_3) \sigma_\beta^{\text{det.}}(X_2) H(X_2; X_3)}{\int dX_2 \Gamma_\alpha^{\text{pro.}}(X_2) \sigma_\beta^{\text{det.}}(X_2) H(X_2; X_3)}$$

$$\Gamma_{2jk}(X_2; X_3) \rightarrow \Gamma_\alpha^{\text{pro.}}(X_2) \sigma_\beta^{\text{det.}}(X_2) P_{2jk}(X_2; X_3)$$

Automatic  $\sum_\alpha P_{3\alpha \rightarrow \beta} = 1$

# Neutrino Decoherence

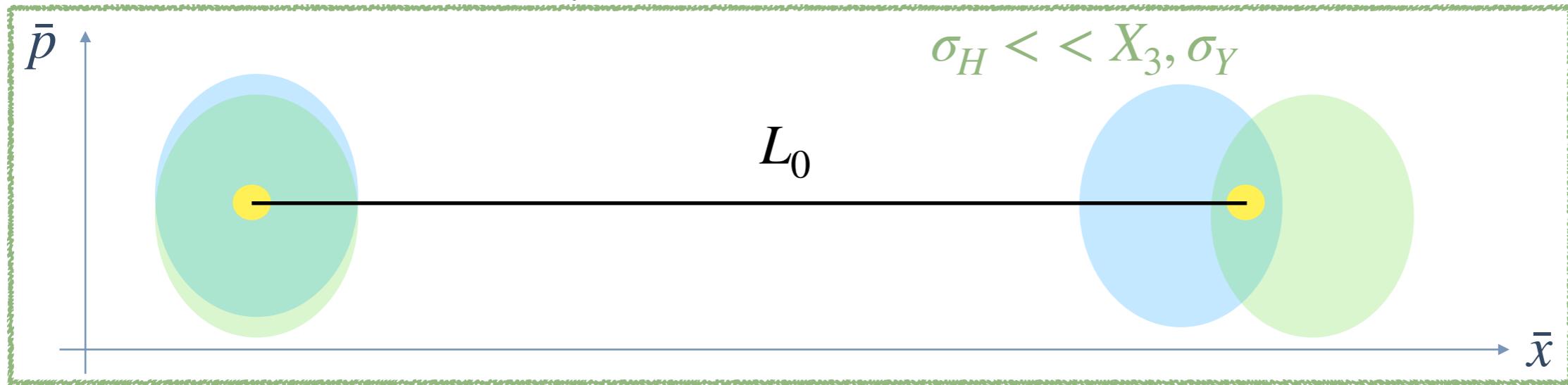
- ❖ State Decoherence and Phase Decoherence

$$P_{3jk}(X_3) = \frac{\int dX_2 H |\Gamma_{2jk}|}{\sqrt{\int dX_2 H \Gamma_{2jj}} \sqrt{\int dX_2 H \Gamma_{2kk}}} \times \frac{\int dX_2 H \Gamma_{2jk}}{\int dX_2 H |\Gamma_{2jk}|}$$

The equation is enclosed in a dashed box. Below the equation, there are two diagrams. The left diagram shows two overlapping circles, one blue and one red, with a horizontal line below them. The right diagram shows the same two overlapping circles, but with diagonal hatching lines (blue on top, red on bottom) and a horizontal line below them.

# State Decoherence and Phase Decoherence

- SD term:  $S_{3jk}(X_3) = \frac{\int dX_2 Y_{jk}}{\sqrt{\int dX_2 Y_{jj}} \sqrt{\int dX_2 Y_{kk}}} \underset{X_2=X_3}{\underset{\simeq}{\downarrow}} S_{2jk}(X_2)$



$S_{2jk}(X_2)$       Washout effect at the Wigner layer ( $\bar{1}$ )

$$Y_{jk} = H |\Gamma_{2jk}| = H \left[ e^{-i\theta_{jk}} \frac{\int d\bar{x} \int d\bar{p} \Gamma_{1jk}(\bar{x}, \bar{p}; X_2)}{\int d\bar{x} \int d\bar{p} |\Gamma_{1jk}(\bar{x}, \bar{p}; X_2)|} \right] \int d\bar{x} \int d\bar{p} |\Gamma_{1jk}(\bar{x}, \bar{p}; X_2)|$$

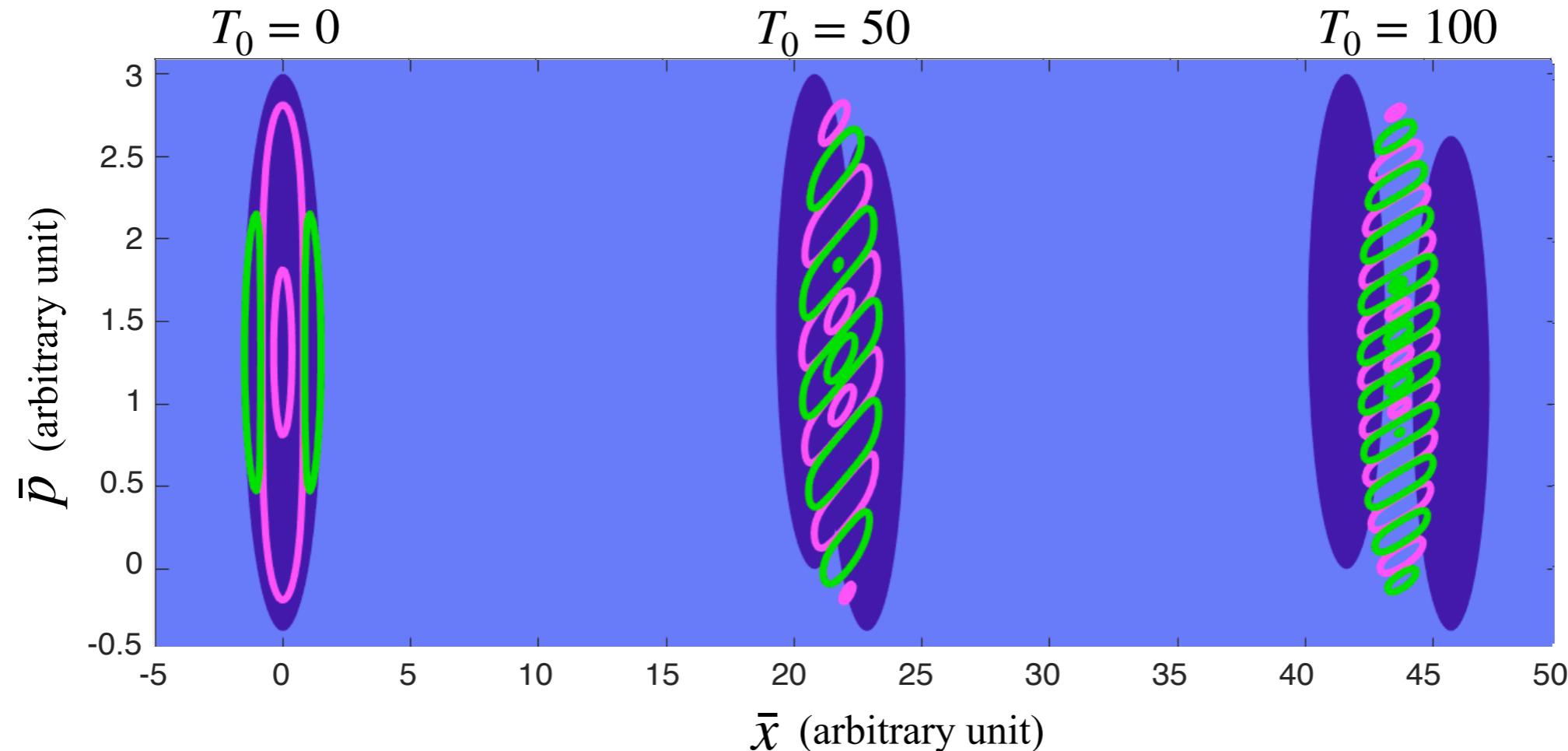
- PD term:  $\Phi_{jk}(X_3) = \frac{\int dX_2 Y_{jk}(X_2; X_3) e^{i\theta_{jk}(X_2)}}{\int dX_2 Y_{jk}(X_2; X_3)}$  Washout effect at the physical layer (2)

# State Decoherence and Phase Decoherence

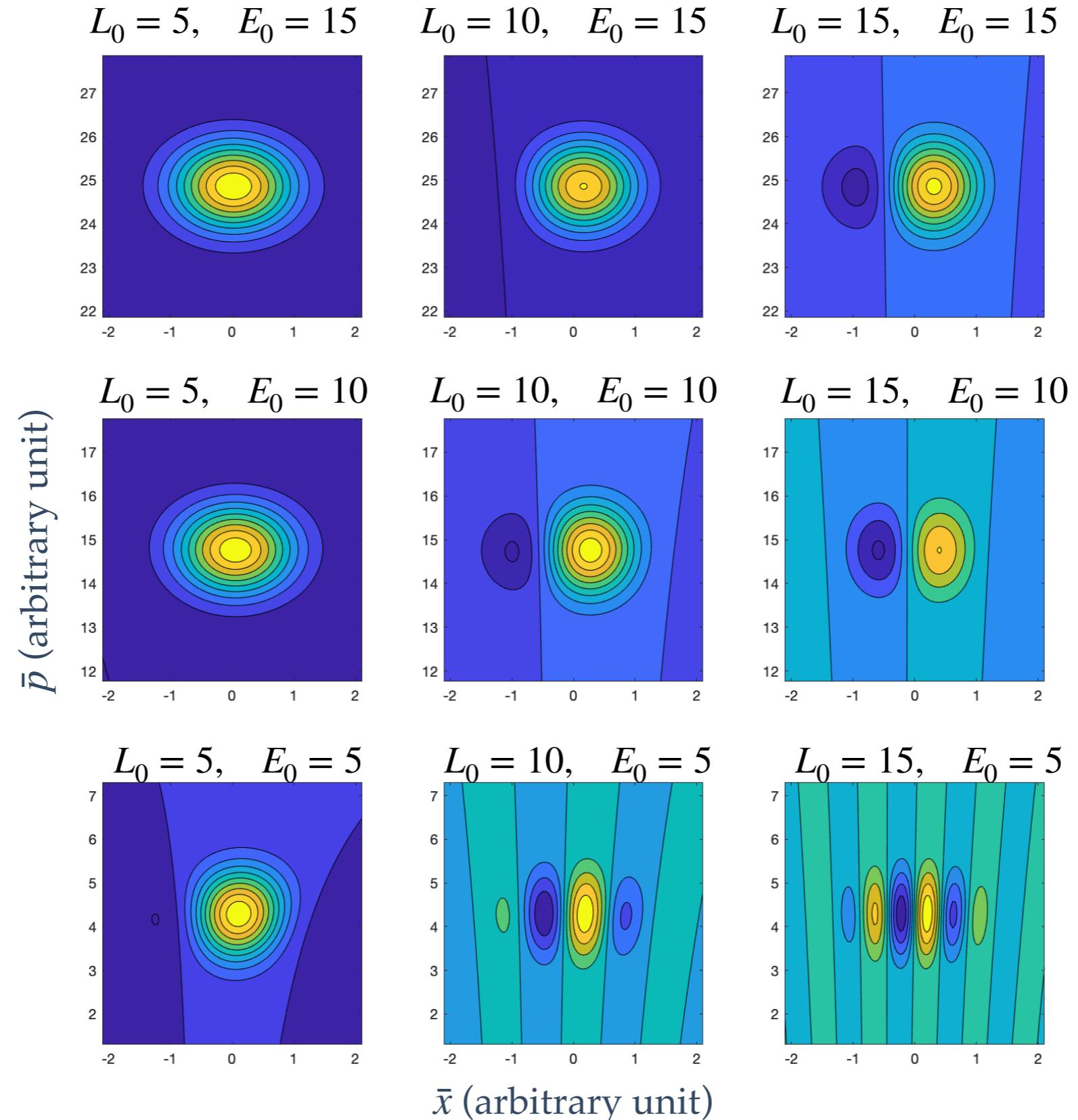
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- ✿ Phase structure on the Wigner phase space layer  $P_{\bar{1}}(\bar{x}, \bar{p}) = |P_{\bar{1}}(\bar{x}, \bar{p})| e^{i\eta_{jk}}$ 
  - ✿ Pulse neutrino emission :  $\eta_{jk}(\bar{x}, \bar{p}; T, E)$
  - ✿ Continuous neutrino emission:  $\eta_{jk}(\bar{x}, \bar{p}; L, E)$
- ✿ Phase structure on the Physical phase space layer  $P_2(x, p) = |P_2(x, p)| e^{i\theta_{jk}}$ 
  - ✿ Pulse neutrino emission :  $\theta_{jk}(T, L, E)$
  - ✿ Continuous neutrino emission:  $\theta_{jk}(L, E) \simeq -\frac{i\Delta m_{jk}^2 L}{2E}$

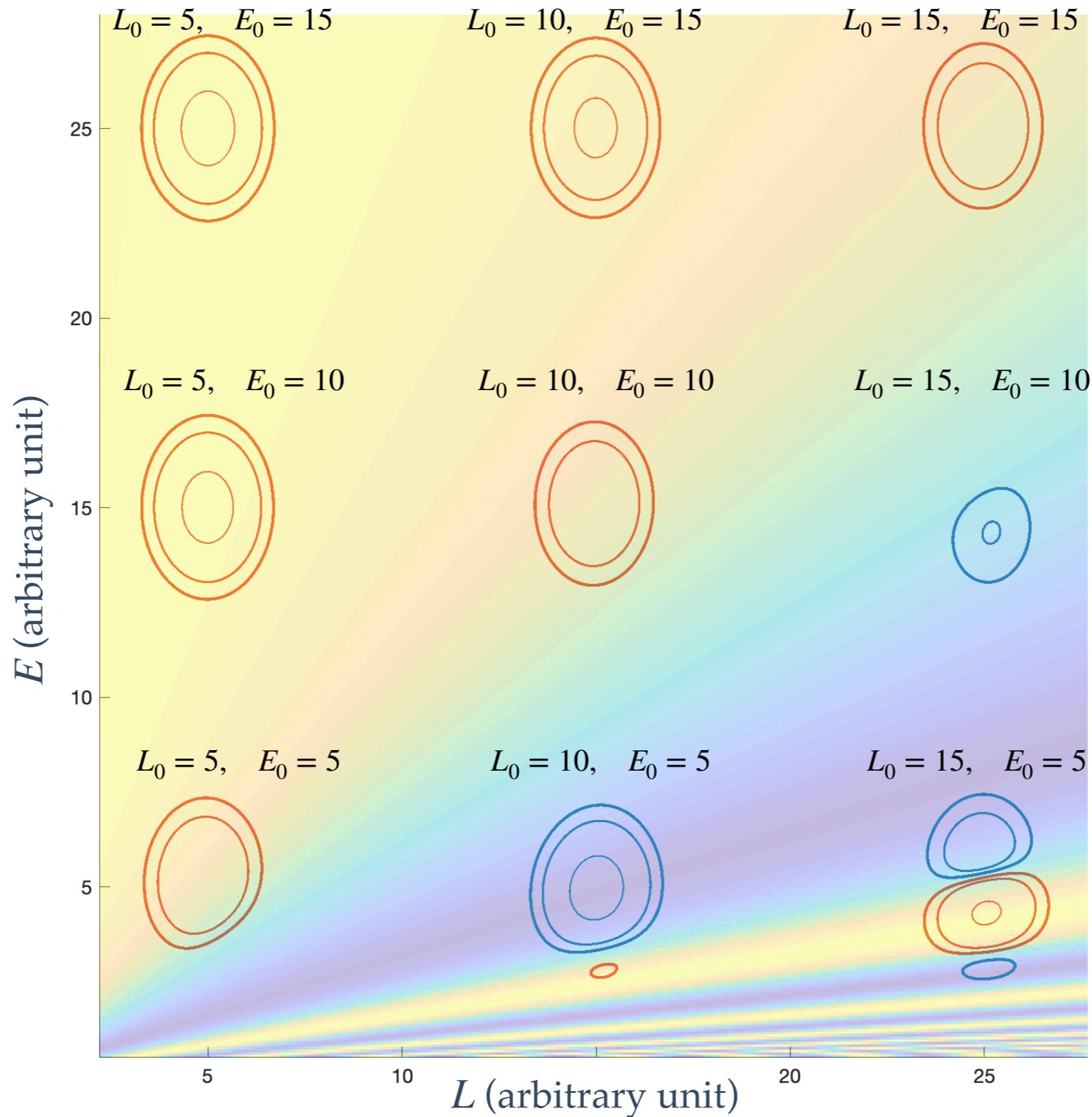
# State Decoherence = Pulse neutrino emission



# State Decoherence — Continuous neutrino emission



# Phase Decoherence — Continuous neutrino emission



# Experimental Potential

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- ❖ Phase shift:
  - ❖ Asymmetry of the PDF w.r.t. the phase structure  
(If all PDFs are Gaussian, only  $\sigma_E$  is not symmetric)
  - ❖ Independent of other (smearing) decoherence effect
- ❖ Long distance: Effects of  $\frac{\sigma_p}{\sqrt{1 + \sigma_p \sigma_x}}$  and  $\sigma_E$  grows with distance
- ❖ Short distance:  $\sigma_L$  isn't affected by distance

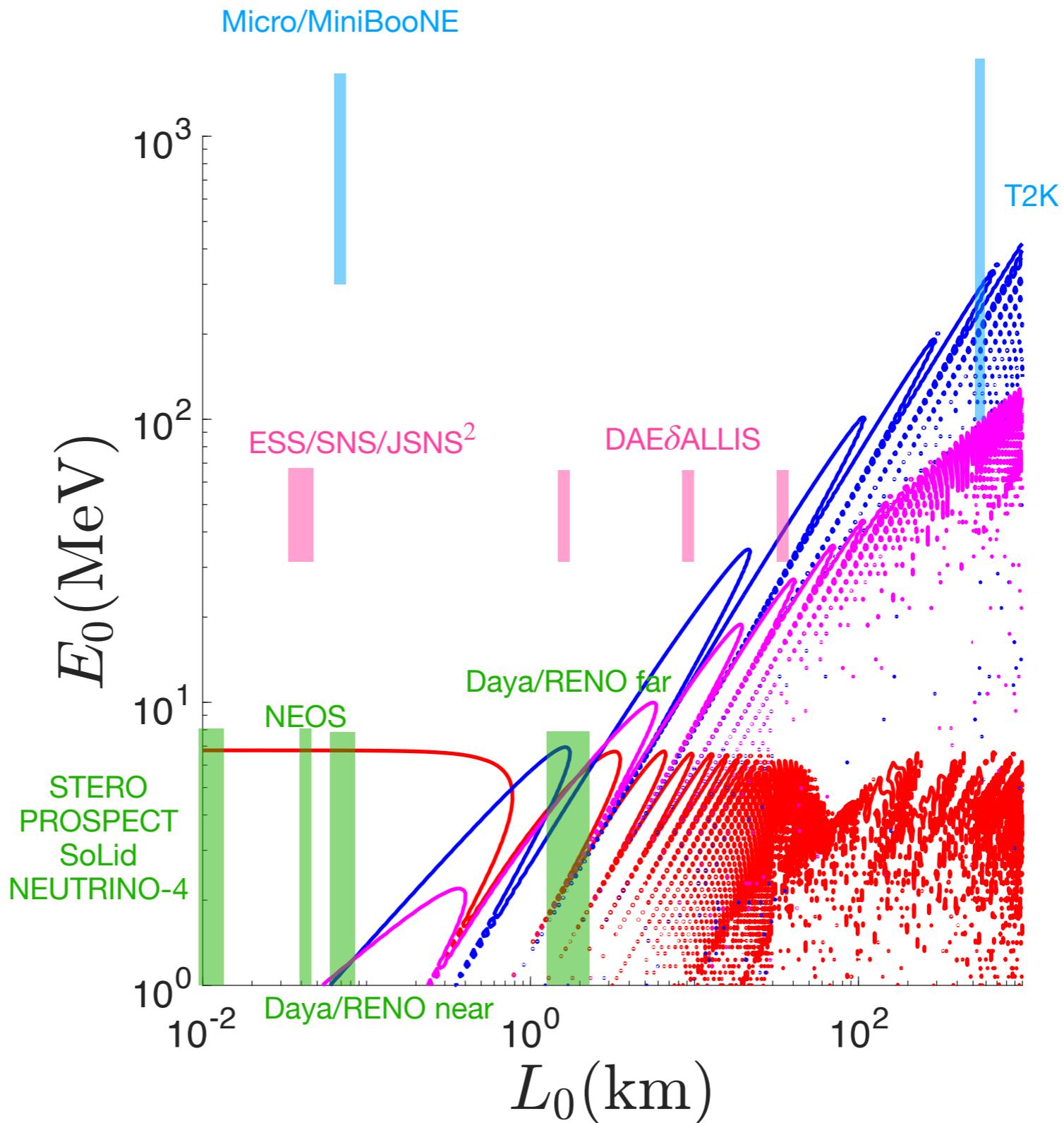
With these differences, one can better identify the uncertainties through combined analysis of different experiments

# Sensitivity on the measurement layer

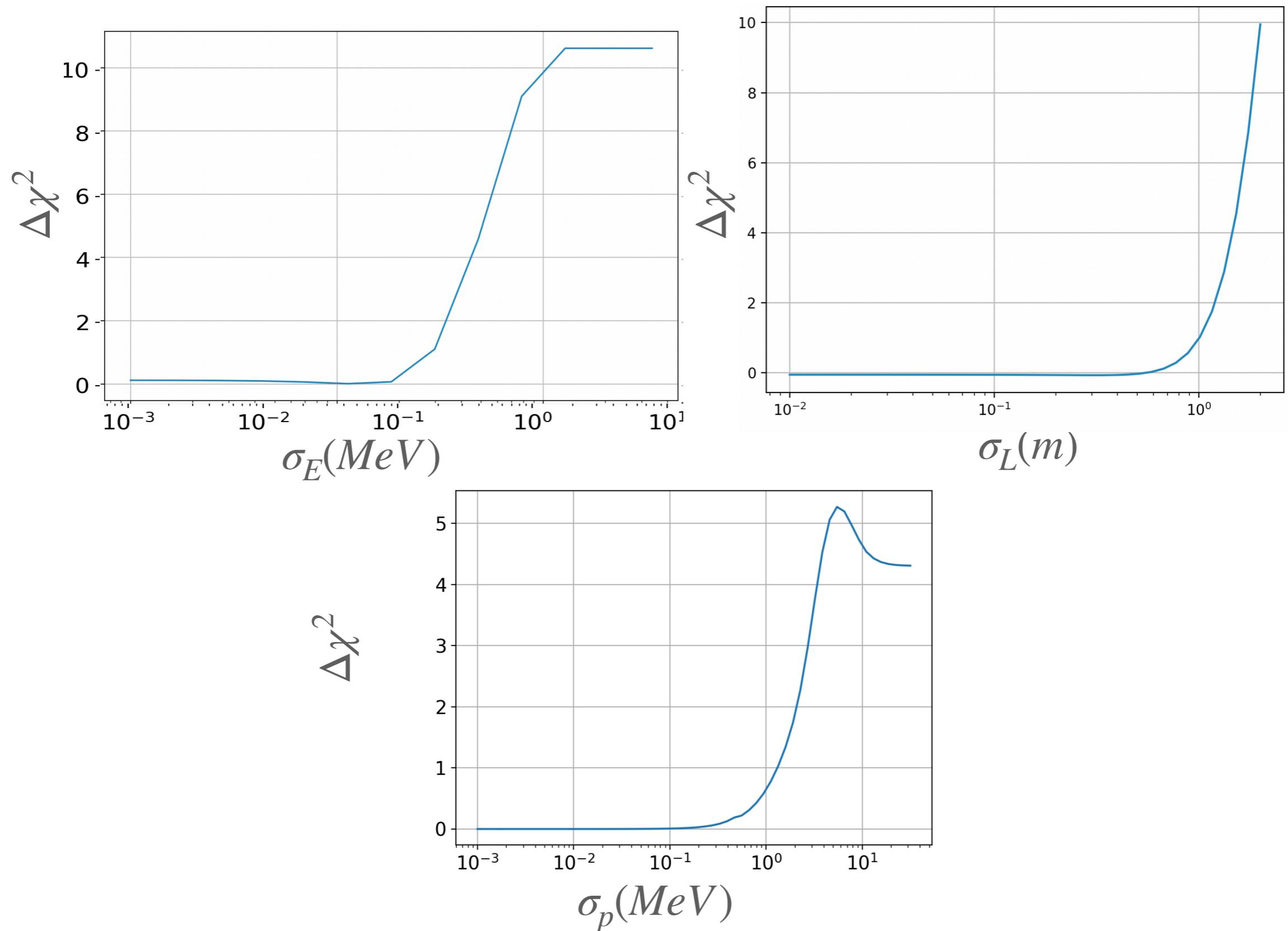
$\nu_e \rightarrow \nu_e$

- $\sigma_L = 50 \text{ m}$
- $\sigma_E = 0.1\sqrt{E_0} \text{ MeV}$
- $\sigma_p = 0.4 \text{ MeV}$

Contour:  
 $|P_3(\sigma \neq 0) - P_3(\sigma = 0)| = 10^{-4}$



# An example: RENO experiment



# Future prospects

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- ✿ Estimation of the sensitivity for each uncertainty for future experiments
- ✿ Develop (analysis, experimental) schemes to identify the uncertainties
- ✿ Consider specific mechanisms to contribute to the classified uncertainties

# Conclusion

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- ❖ Developed the layer structures which classifies uncertainties from microscopic to macroscopic, building on QFT calculation
- ❖ Classify decoherence into state decoherence and phase decoherence based on their causes and investigate their phenomenological properties
- ❖ State decoherence and phase decoherence and wash-out effects on the Wigner phase space and physical phase space respectively
- ❖ Showed sensitivity of decoherence effect for reactor neutrino experiments (not yet complete)

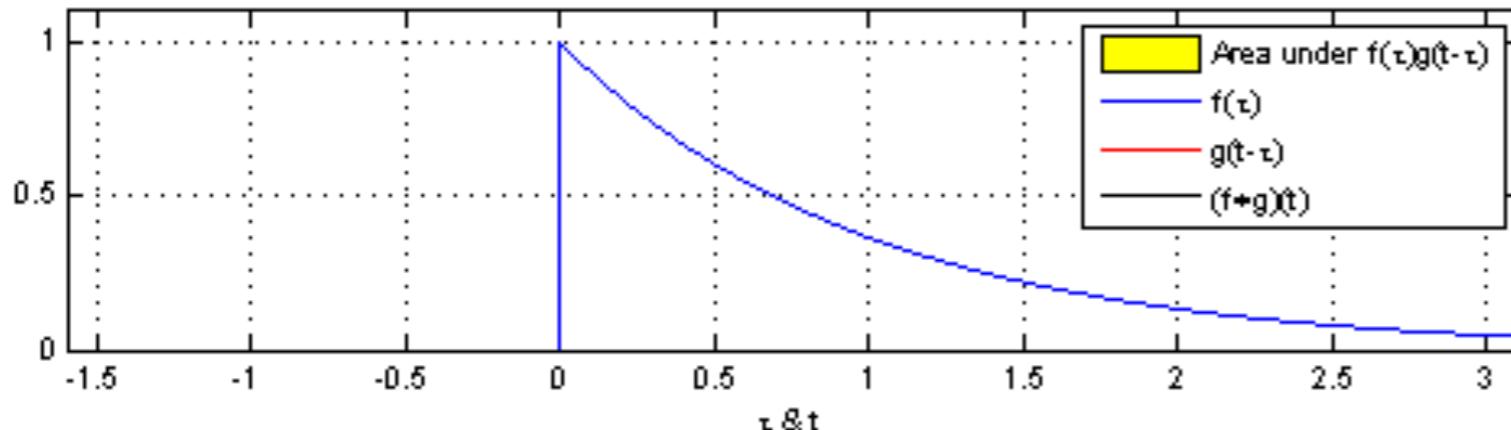
# Back Up Slides

# State Decoherence and Phase Decoherence

- ✿ Phase structure on the Wigner phase space layer  $P_{\bar{1}}(\bar{x}, \bar{p}) = |P_{\bar{1}}(\bar{x}, \bar{p})| e^{i\eta_{jk}}$ 
  - ✿ Pulse neutrino emission :  $\eta_{jk}(\bar{x}, \bar{p}; T, E) = -i\bar{p}T(\mathbf{v}_j - \mathbf{v}_k) - i\bar{x}(\mathbf{P}_j - \mathbf{P}_k)|_{T=T_0, E=E_0}$
  - ✿ Continuous neutrino emission:  $\eta_{jk}(\bar{x}, \bar{p}; \mathbf{L}, E) \simeq -\frac{i\Delta m_{jk}^2}{2E}((\mathbf{L} - \bar{x})\frac{\bar{p}}{E} + \bar{x})|_{\mathbf{L}=\mathbf{L}_0, E=E_0}$
- ✿ Phase structure on the Physical phase space layer  $P_2(\mathbf{x}, \mathbf{p}) = |P_2(\mathbf{x}, \mathbf{p})| e^{i\theta_{jk}}$ 
  - ✿ Pulse neutrino emission :  $\theta_{jk}(T, \mathbf{L}, E) = -iT(E_j - E_k) + i\mathbf{L}(\mathbf{P}_j - \mathbf{P}_k)$
  - ✿ Continuous neutrino emission:  $\theta_{jk}(\mathbf{L}, E) \simeq -\frac{i\Delta m_{jk}^2 L}{2E}$

# Uncertainties

- ❖ Convolution property:



$$G(L) = \int dx f(x)g(x - L)$$
$$\sigma_G \geq \sigma_f \text{ and } \sigma_g$$

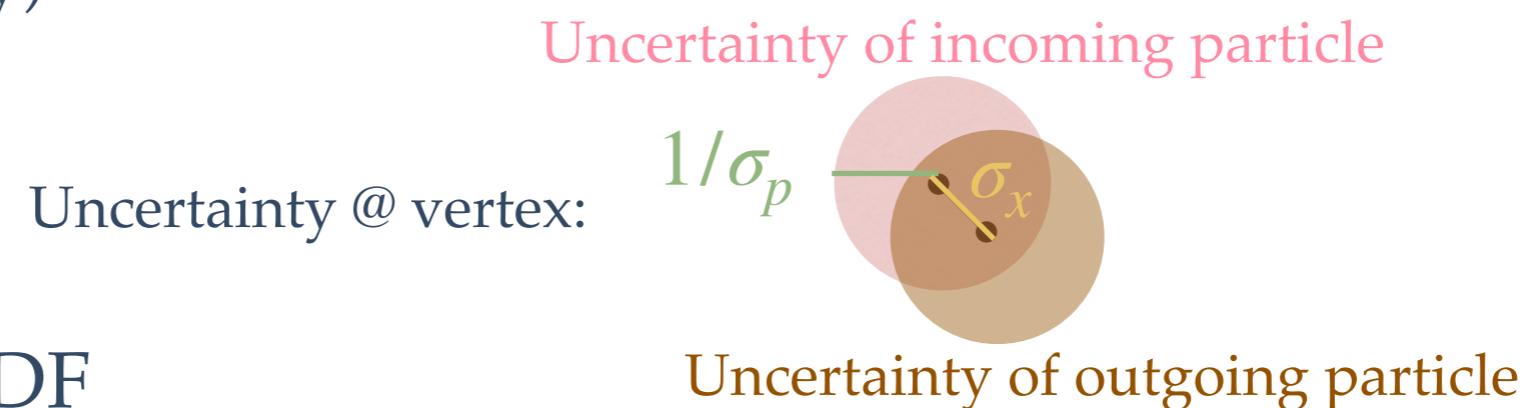
- ❖ Fourier Transformation property:  $F(p) = \int dx e^{ipx} f(x)$        $\sigma_{xp} \uparrow \Leftrightarrow \sigma_x \downarrow$

- ❖ Product of two PDF:  $G'(x) = f(x) g(x)$        $\sigma_{fg} \uparrow \leq \sigma_f, \sigma_g \uparrow$

e.g. for two Gaussians:  $\sigma_{fg}^2 = \frac{\sigma_f^2 \sigma_g^2}{\sigma_f^2 + \sigma_g^2} = \frac{\sigma_f^2}{\sigma_f^2/\sigma_g^2 + 1}$

# Uncertainties

- ✿ Possible contributions to each uncertainties:
  - ✿  $\sigma_p$ : external particle WP size
  - ✿  $\sigma_x$ : non central collisions
  - ✿  $\sigma_E$ : energy resolution, reconstruction model
  - ✿  $\sigma_L$ : production profile, reactor core size, decay length, space fluctuation (e.g. quantum gravity)



## ✿ Properties of PDF

- ✿ Convolution:  $\sigma_{f \otimes g} \geq \sigma_f$  and  $\sigma_g$
- ✿ Fourier Transformation:  $\sigma_{xp} \uparrow \Leftrightarrow \sigma_x \downarrow$
- ✿ Product of two PDF:  $\sigma_{fg} \uparrow \leq \sigma_f, \sigma_g \uparrow$

# Uncertainties

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- \* Layer 1 → 2 structure :

$$\int dp \int dx e^{i\mathbf{px}} F_j(\mathbf{p}; \mathbf{P}_j) G_j(\mathbf{x}; \mathbf{L}_j) = \int dp e^{i\mathbf{p}\mathbf{L}_j} F_j(\mathbf{p}; \mathbf{P}_j) \tilde{G}_j(\mathbf{p}; 0; \mathbf{L}_j) = e^{i\mathbf{p}\mathbf{L}_j} \hat{\Phi}(\mathbf{L}_j, \mathbf{P}_j)$$

$\sigma_p$        $\sigma_x \uparrow$                            $\sigma_p^{eff} \downarrow \leq \sigma_p, \sigma_{xp} \downarrow$        $\sigma_p^{(2)} \uparrow$

$$F_j(\mathbf{p}; \mathbf{P}_j) = \left( f_{pi}(q) M_q(q) \circledast f_{pf}(k) M_k(k) \right) \times \left( f_{di}(q') M_{q'}(q') \circledast f_{df}(k') M_{k'}(k') \right)$$

Production of **convolution** of incoming & outgoing particles @ production/detection site

$$G_j(\mathbf{x}; \mathbf{L}_j) = g_1(x_1) \circledast g_2(x_2)$$

**Convolution** of production and detection vertex uncertainties

- \* Layer 2 → 3 structure :

$$\int dL e^{i\frac{\Delta m_{jk} L}{2E_0}} H_L(L; L_0) \tilde{H}_E(L) \hat{\Phi}(\mathbf{L}_j, \mathbf{P}_j) = e^{i\frac{\Delta m_{jk} L_0}{2E_0}} D_{jk}(L_0)$$

$\sigma_L^{eff} \leq \sigma_L, \sigma_{EL}, \sigma_x^{(2)}$        $\sigma_{L_0} \uparrow$