

WIN 2015

Antonio  
Palazzo

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Munich

Adrien  
Laura  
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Schwarz

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LBLs:  
a new  
window on  
sterile vs

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# Outline

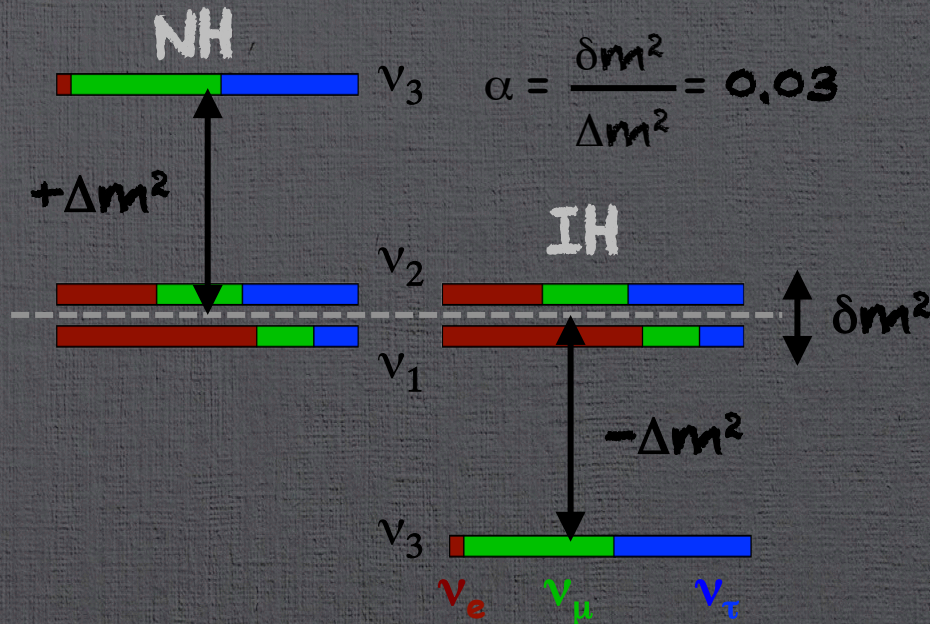
## Introduction

- Imprints of the CP violation effects induced by sterile neutrinos in T2K
- Impact of the  $4\nu$  interference effects on the interpretation of ICARUS and OPERA

## Conclusions



# The 3-flavor scheme



$$\alpha = \frac{\delta m^2}{\Delta m^2} = 0,03$$

unknowns:

CP-phase  $\delta$   
(Hints of  $\delta \neq 0, \pi$ )

NMH

(Hints of NH)

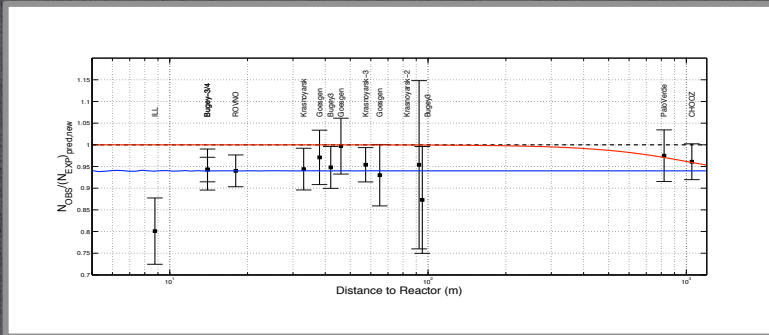
$$U = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$\theta_{23} \sim 41^\circ \quad \theta_{13} \sim 9^\circ \quad \theta_{12} \sim 34^\circ$$

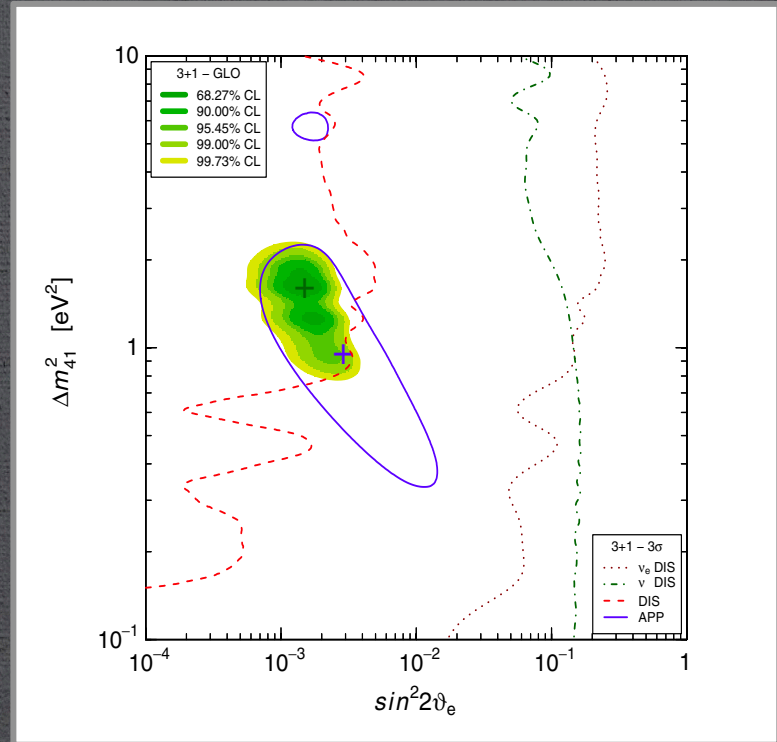


# SBL anomalies point to a 4<sup>th</sup> neutrino

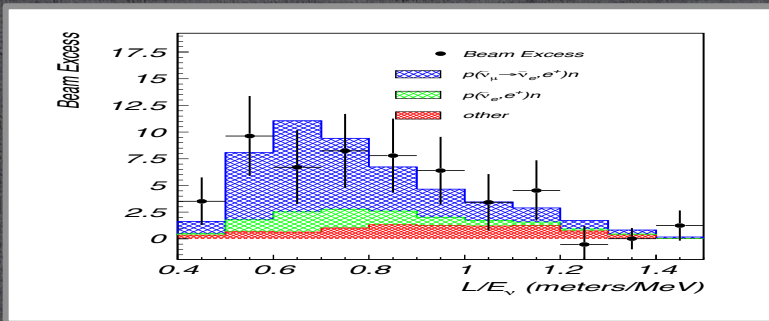
Reactor & Gallium:  $P_{ee} < 1$



Giunti et al., PRD 2013



Accelerators:  $P_{\mu e} > 0$



$$\frac{L}{E} \sim \frac{m}{\text{MeV}}$$

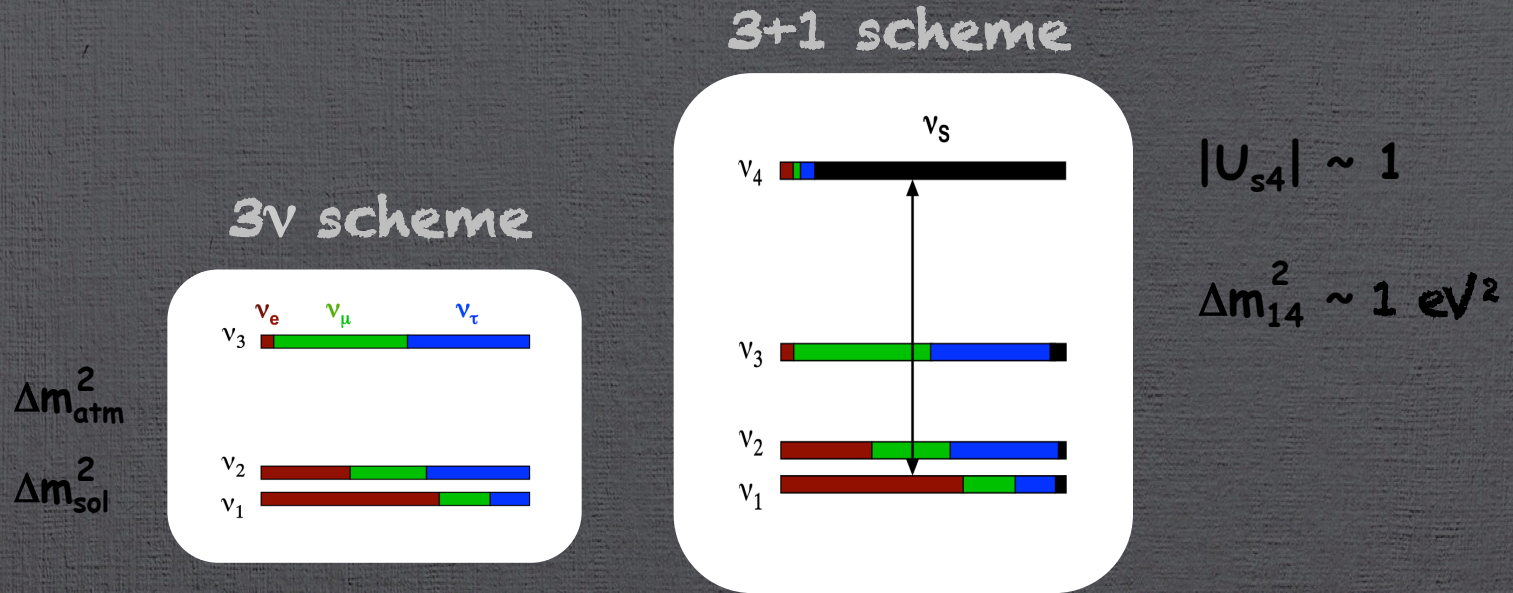


$$\begin{aligned} \Delta_{12} &\simeq 0 \\ \Delta_{13} &\simeq 0 \end{aligned}$$

Need of a new larger  $\Delta m^2$   
 $\sim 1 \text{ eV}^2$



# Introducing a sterile neutrino



Only small perturbations to the 3v framework

However, 3v CP-violation effects are very small!

Can new 4v CPV effects compete with the 3v ones?



# Mixing matrix in 3+1 scheme

$$U = \tilde{R}_{34} R_{24} \tilde{R}_{14} R_{23} \underbrace{\tilde{R}_{13} R_{12}}_{3\nu}$$

$$R_{ij} = \begin{bmatrix} c_{ij} & s_{ij} \\ -s_{ij} & c_{ij} \end{bmatrix}$$

$$\tilde{R}_{ij} = \begin{bmatrix} c_{ij} & \tilde{s}_{ij} \\ -\tilde{s}_{ij}^* & c_{ij} \end{bmatrix}$$

$$\begin{aligned} s_{ij} &= \sin \theta_{ij} \\ c_{ij} &= \cos \theta_{ij} \\ \tilde{s}_{ij} &= s_{ij} e^{-i\delta_{ij}} \end{aligned}$$

$$3\nu \begin{cases} 3 \text{ mixing angles} \\ 1 \text{ Dirac CP-phases} \\ 2 \text{ Majorana phases} \end{cases}$$

$$3+1 \begin{cases} 6 \\ 3 \\ 3 \end{cases}$$

$$3+N \begin{cases} 3+3N \\ 1+2N \\ 2+N \end{cases}$$

$$\theta_{14} = \theta_{24} = \theta_{34} = 0 \rightarrow 3\text{-flavor case}$$



# An important remark

$$A_{\alpha\beta}^{\text{CP}} \equiv P(\nu_\alpha \rightarrow \nu_\beta) - P(\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta)$$

$$A_{\alpha\beta}^{\text{CP}} = -16 J_{\alpha\beta}^{12} \sin \Delta_{21} \sin \Delta_{13} \sin \Delta_{32}$$

if  $\Delta \equiv \Delta_{13} \simeq \Delta_{23} \gg 1$   $\rightarrow$   $\langle \sin^2 \Delta \rangle = 1/2$   
Osc. averaged out by finite E resol.

It can be:

$$A_{\alpha\beta}^{\text{CP}} \neq 0$$

(if  $\sin \delta \neq 0$ )

The bottom line is that if one of the three  $\nu_i$  is  $\infty$  far from the other two ones this does not erase CPV

(relevant for the 4v case)

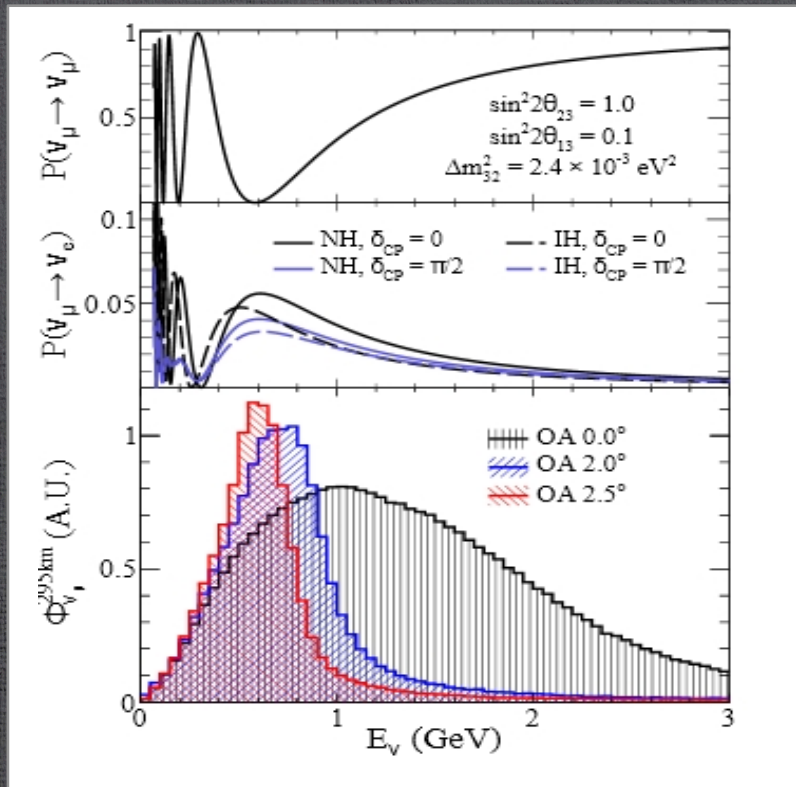
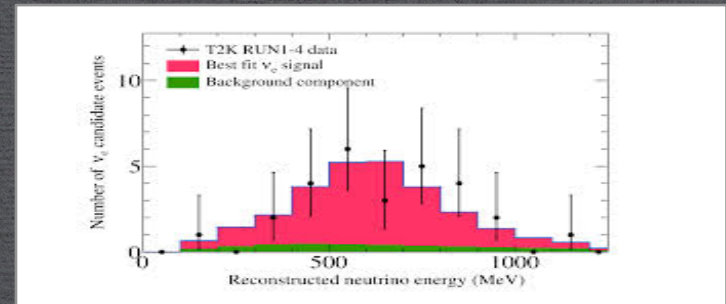
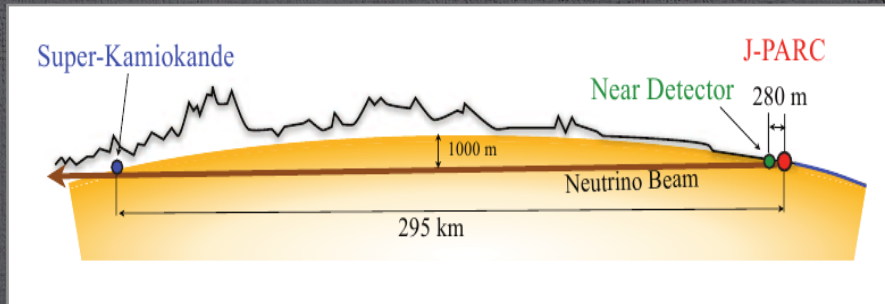


# Hints on the new CP-phases from T2K (and $\theta_{13}$ -reactor experiments)

N. Klop and A.P., PRD 91 073017 (2015)



# Outline of the T2K experiment



$E = 0.6 \text{ GeV}$   
 $L = 295 \text{ km}$   
 $\Delta m_{13}^2 = 2.4 \times 10^{-3}$

$$\Delta = \frac{\Delta m_{13}^2 L}{4E} \simeq \frac{\pi}{2}$$

First oscillation maximum



# T2K: 3-flavor transition probability

$$P_{\nu_{\mu} \rightarrow \nu_e}^{3\nu} = P^{\text{ATM}} + P^{\text{SOL}} + P^{\text{INT}}$$

In vacuum:

$$P^{\text{ATM}} = 4s_{23}^2 s_{13}^2 \sin^2 \Delta$$

$$P^{\text{SOL}} = 4c_{12}^2 c_{23}^2 s_{12}^2 (\alpha \Delta)^2$$

$$P^{\text{INT}} = 8s_{23}s_{13}c_{12}c_{23}s_{12}(\alpha \Delta) \sin \Delta \cos(\Delta + \delta_{CP})$$

$$\Delta = \frac{\Delta m_{31}^2 L}{4E}, \quad \alpha = \frac{\Delta m_{21}^2}{\Delta m_{31}^2}$$

$$\Delta \sim \pi/2$$

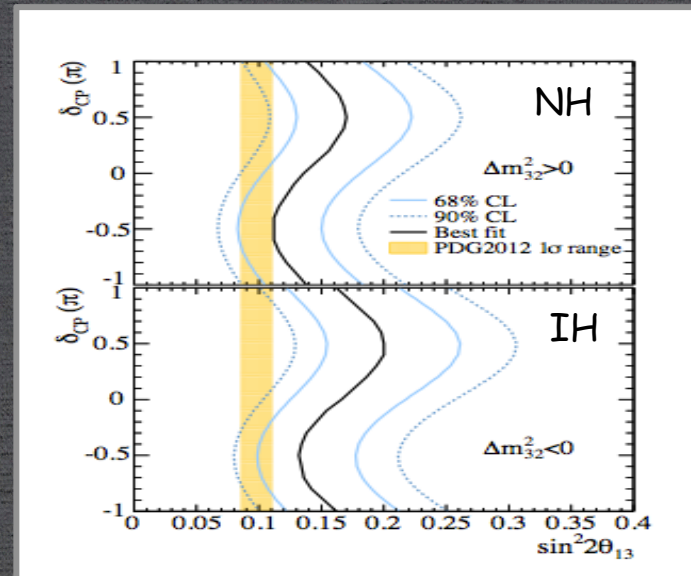
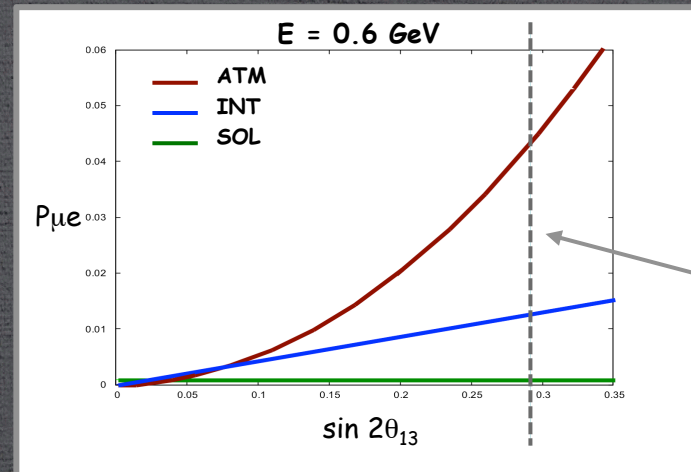
$$\alpha \sim 0.03$$

$P^{\text{ATM}}$  leading  $\rightarrow \theta_{13} > 0$

$P^{\text{INT}}$  subleading  $\rightarrow \delta$  dependence

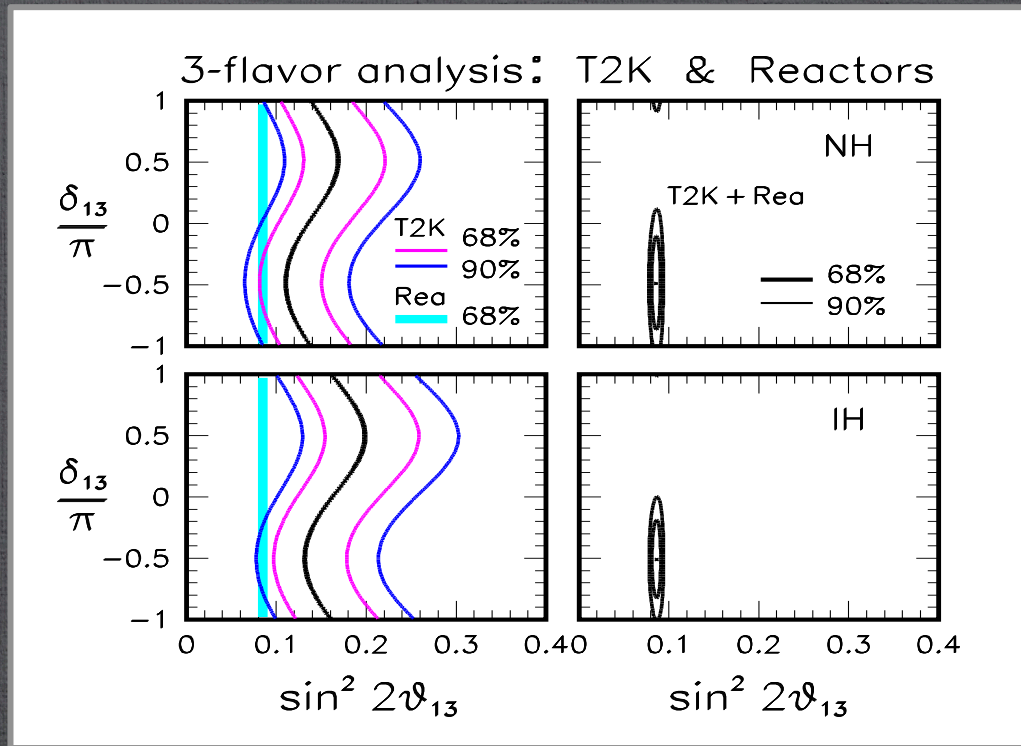
$P^{\text{SOL}}$  negligible

Matter effects induce some difference among NH and IH





# Present data have some sensitivity to $\delta$



Slight  $\theta_{13}$  mismatch  
T2K vs Reactors

No CPV ( $\delta = 0, \pi$ )  
disfavored at  
 $\sim 90\%$  C.L.

Best fit  $\delta \sim -\pi/2$

NH slightly  
favored  $\Delta\chi^2 \sim -1$   
(similar finding in  
SK atmospheric vs)

Note that  $\delta$  is not extracted from observation of manifest CPV

Combination of  $\left\{ \begin{array}{l} P_{ee} \text{ (}\delta\text{-independent), LBL Reactors} \\ P_{\mu e} \text{ (}\delta\text{-dependent), LBL Accelerators (T2K)} \end{array} \right.$



# T2K: 4-flavor transition probability

- $\Delta m^2_{14} \gg \Delta m^2_{13}$  : fast oscillations induced by  $\Delta m^2_{14}$  are averaged out
- Phase information (value of  $\Delta m^2_{14}$ ) gets lost (in contrast to SBL)
- Unlike SBL, interf. of  $\Delta m^2_{14} \not\equiv \Delta m^2_{13,12}$  observable: sensitivity to CP-phases

In vacuum, for  $\Delta m^2_{14} \rightarrow \infty$

$$P_{\nu_\mu \rightarrow \nu_e}^{4\nu} = 4|U_{\mu 3}|^2|U_{e 3}|^2 \sin^2 \Delta + 4|U_{\mu 2}|^2|U_{e 2}|^2(\alpha\Delta)^2 + 8|U_{\mu 3}^*||U_{e 3}||U_{\mu 2}||U_{e 2}|(\alpha\Delta) \sin \Delta \cos(\Delta + \delta_{13}) + 4|U_{\mu 3}^*||U_{e 3}||U_{\mu 4}||U_{e 4}^*| \sin \Delta \sin(\Delta + \delta_{13} - \delta_{14}) - 4|U_{\mu 2}^*||U_{e 2}||U_{\mu 4}||U_{e 4}^*|(\alpha\Delta) \sin \delta_{14} + 2|U_{\mu 4}|^2|U_{e 4}|^2$$

$$P_{\nu_\mu \rightarrow \nu_e}^{4\nu} \sim (1 - |U_{e 4}|^2 - |U_{\mu 4}|^2)P_{\mu e}^{3\nu} + P_{II}^{INT} + P_{III}^{INT} + P^{STR}$$

$$P_{II}^{INT} = 2 \sin 2\theta_{\mu e} s_{13} s_{23} \sin \Delta \sin(\Delta + \delta_{13} - \delta_{14})$$

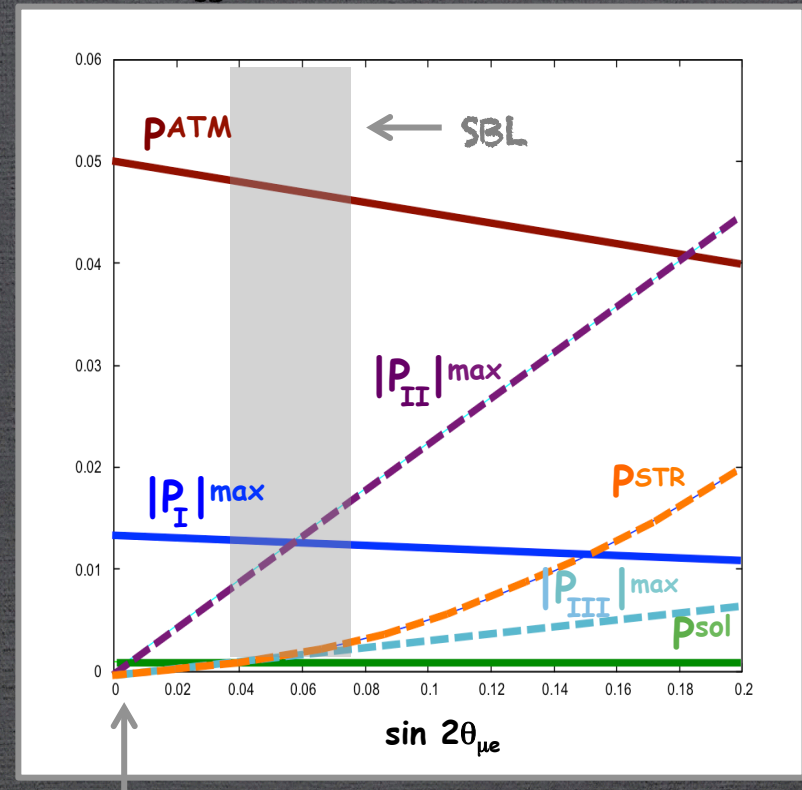
$$P_{III}^{INT} = -2 \sin 2\theta_{\mu e} c_{23} s_{12} c_{12} (\alpha\Delta) \sin \delta_{14}$$

$$P^{STR} = \frac{1}{2} \sin^2 2\theta_{\mu e}$$

$$\sin^2 2\theta_{\mu e} = 4|U_{e 4}|^2|U_{\mu 4}|^2$$

$P_{II}^{INT}$  can be as large as  $P_I^{INT}$

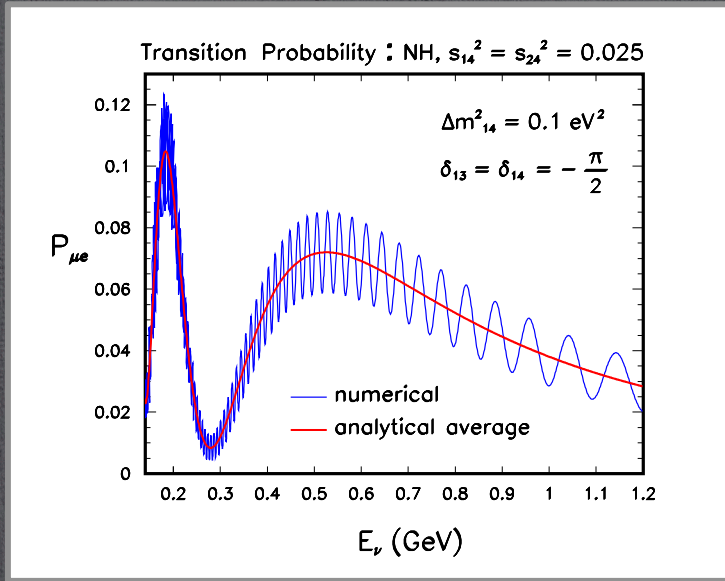
$\theta_{13} = 9^\circ$   $E = 0.6$  GeV



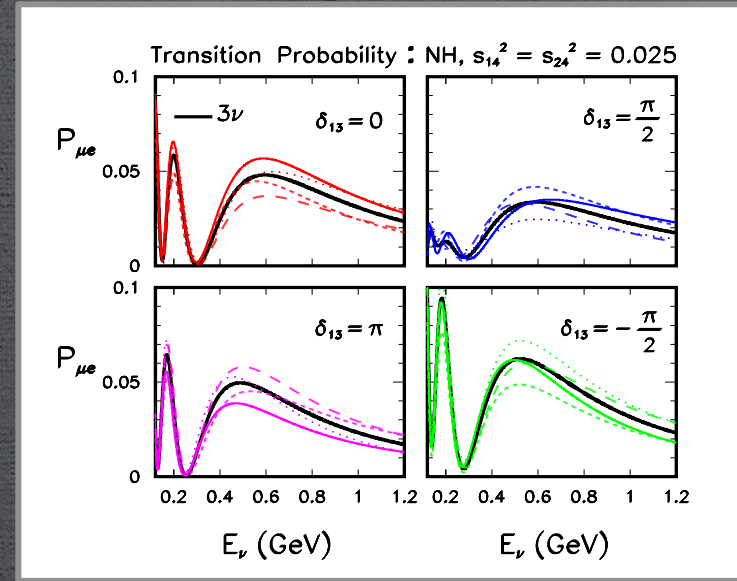
3ν limit



# Numerical examples of $4\nu$ probability



The fast oscillations get averaged out due to the finite energy resolution

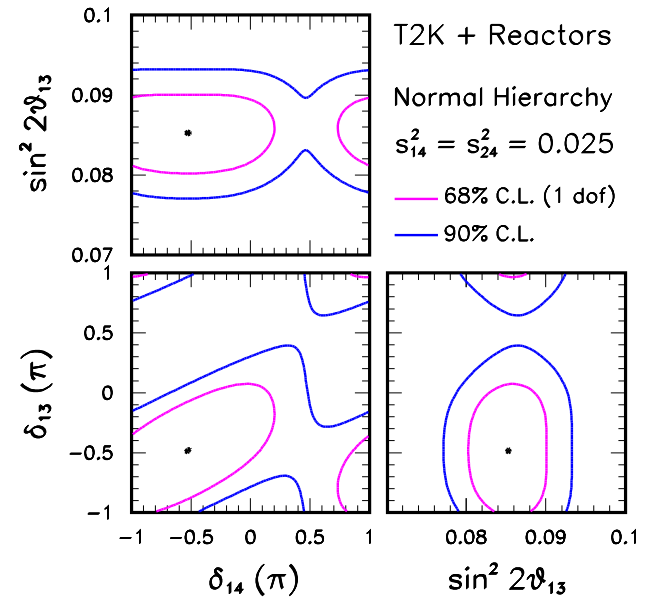
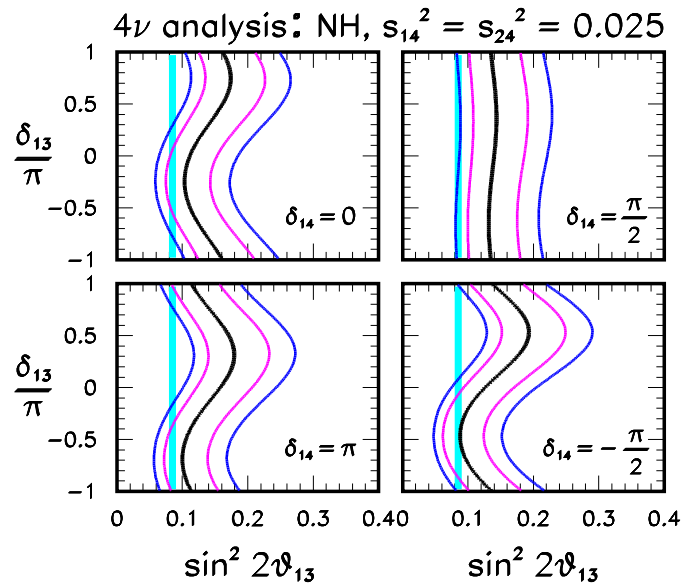


Different line styles  
 $\Leftrightarrow$   
 Different values of  $\delta_{14}$

The modifications induced by  $\delta_{14}$  are as large as those induced by the standard CP-phase  $\delta_{13}$



# Results of the 4ν analysis (NH)



Similar findings in IH

- Big impact on T2K "wiggles"
- Comparable sensitivity to  $\delta_{13} \nsubseteq \delta_{14}$
- Best fit values:  $\delta_{13} \sim \delta_{14} \sim -\pi/2$
- 4ν gives better agreement of T2K & Reactors

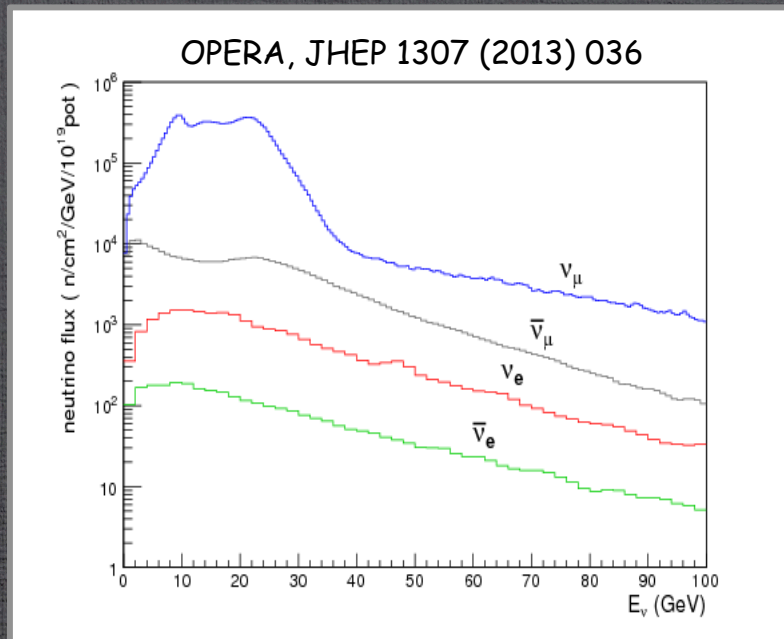
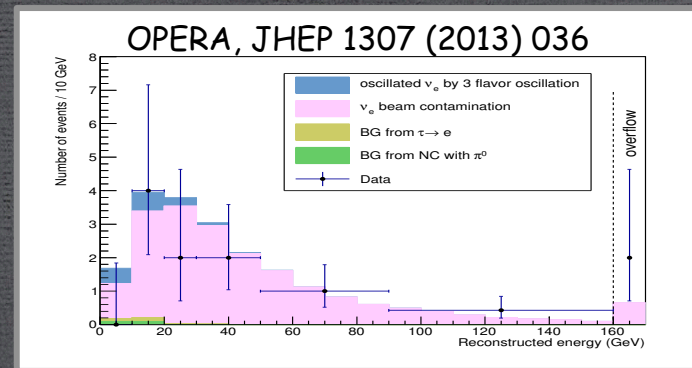
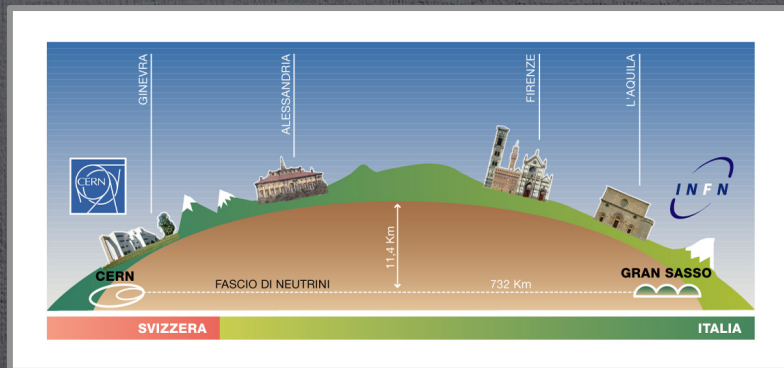


Impact of the new CP-phases  
on the interpretation  
of the  $\nu_\mu \rightarrow \nu_e$  sterile  $\nu$  searches  
of ICARUS & OPERA

A.P., PRD 91 091301 (2015) Rapid Communication



# Outline of the CNGS experiments



$$\langle E \rangle = 17 \text{ GeV}$$

$$L = 732 \text{ km}$$

$$\Delta m_{13}^2 = 2.4 \times 10^{-3}$$

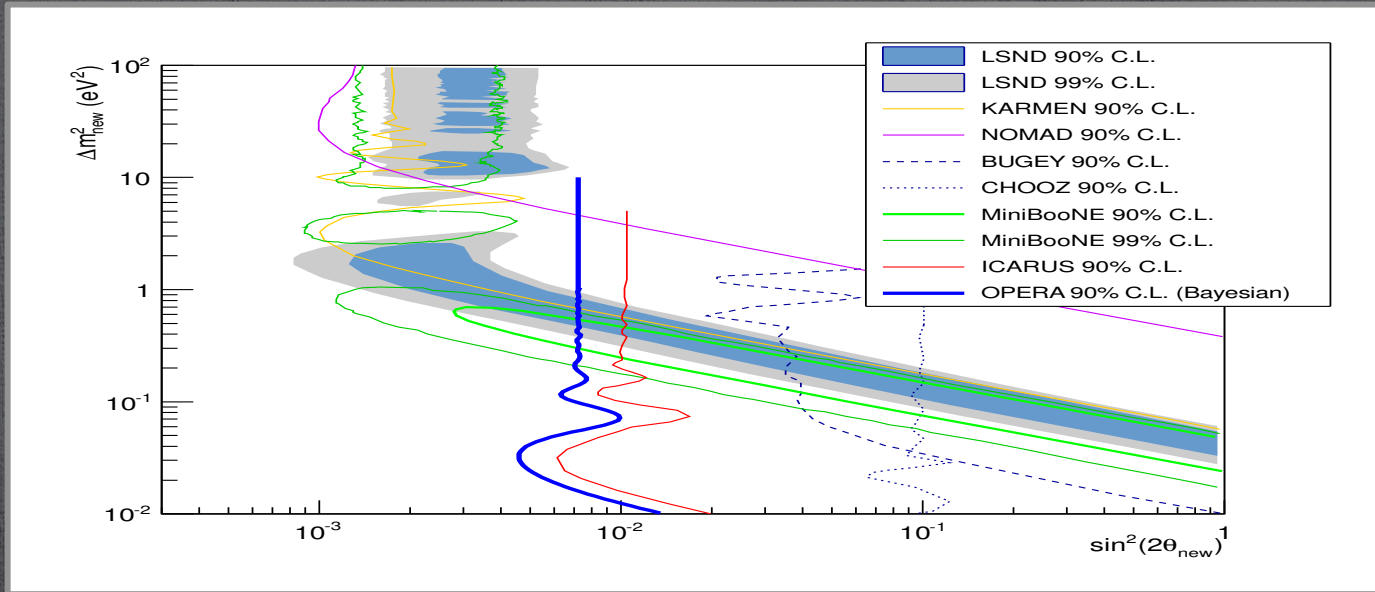
$$\Delta = \frac{\Delta m_{13}^2 L}{4E} \simeq 0.13$$

*3ν oscillations  
play a minor role*

*Good place where  
to look for sterile vs*



# Official bounds from OPERA & ICARUS

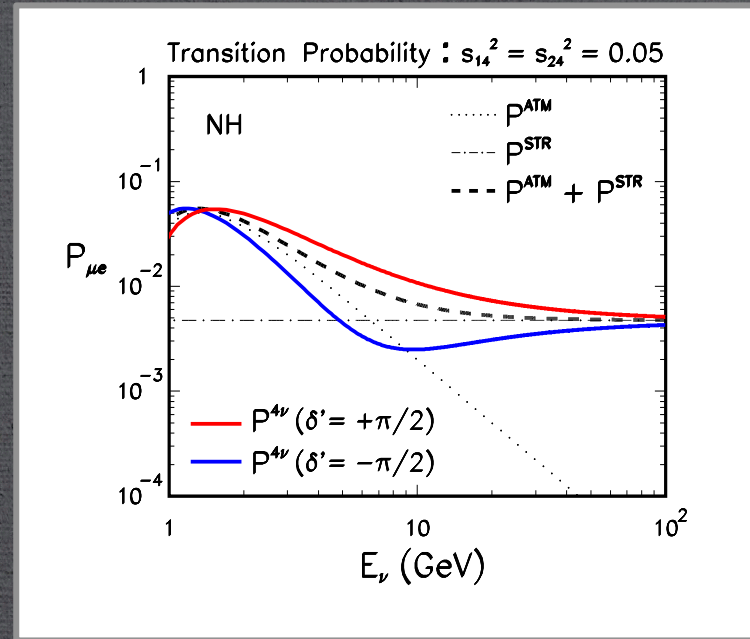
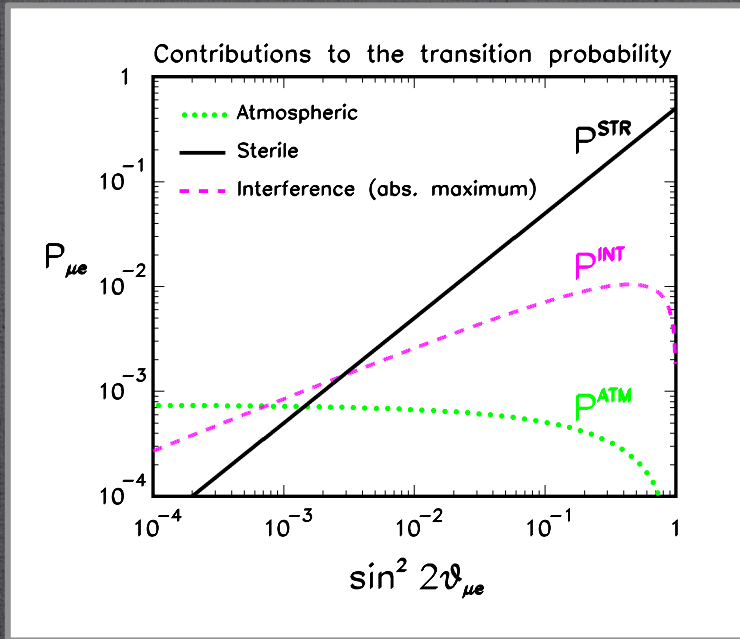


2-flavor  
treatment  
adopted by both  
collaborations

$$\begin{cases} \mathcal{P}(\nu_\mu \rightarrow \nu_e) = 4 \sin^2 2\theta_{\mu e} \sin^2 \Delta_{14} \\ \quad + \text{small Atm. term} \\ \mathcal{P}(\nu_e \rightarrow \nu_e) = 1 \quad (\nu_e \text{ bck fixed}) \end{cases}$$



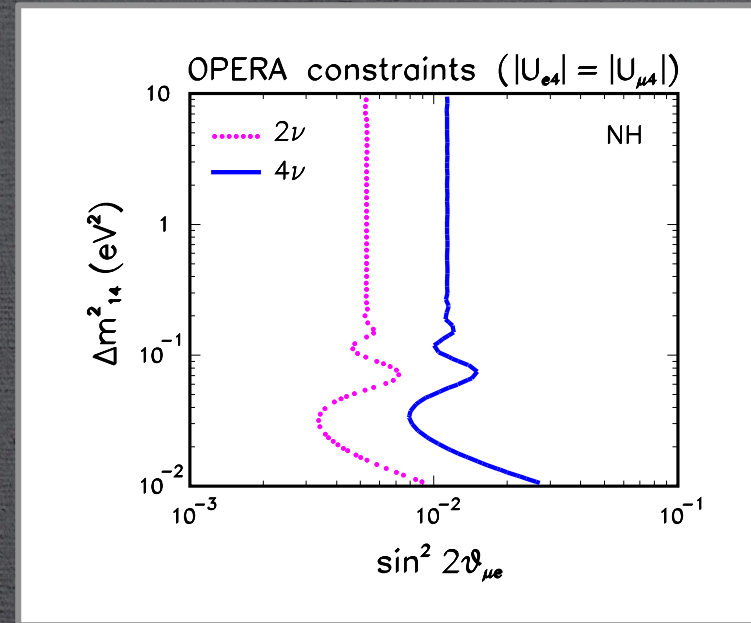
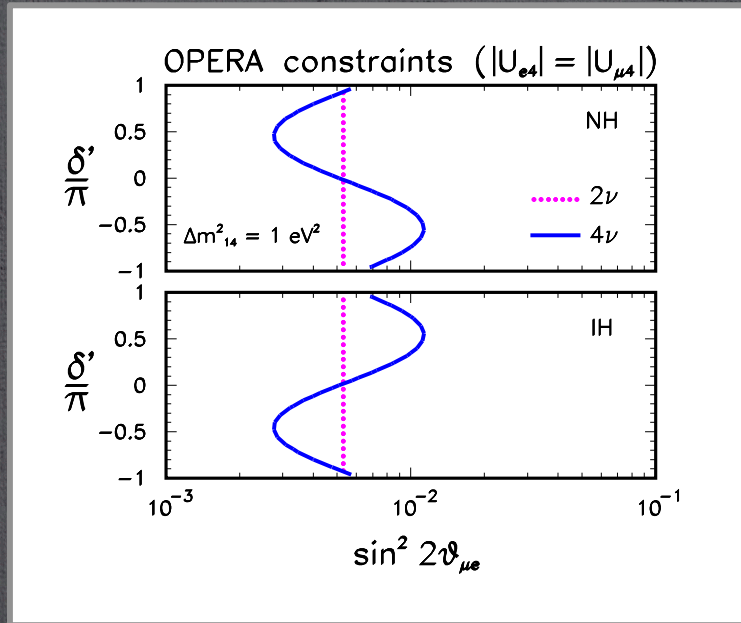
# 4ν effects at the CNGS beam



- Interference has substantial impact on  $P(\nu_\mu \rightarrow \nu_e)$
- The official analyses neglect the interference term
- Proper inclusion of such effects is necessary



# Impact of the 4ν interference term



Upper bound depends on the (unknown) CP-phase  $\delta'$

After marginalization of the CP-phase...

The upper bounds get relaxed by a factor of two

$$(2\nu) \sin^2 2\theta_{\mu e} < 5 \times 10^{-3} \quad \rightarrow \quad (4\nu) \sin^2 2\theta_{\mu e} < 1.2 \times 10^{-2}$$



# A further remark on 4v effects

In a 4v scheme:

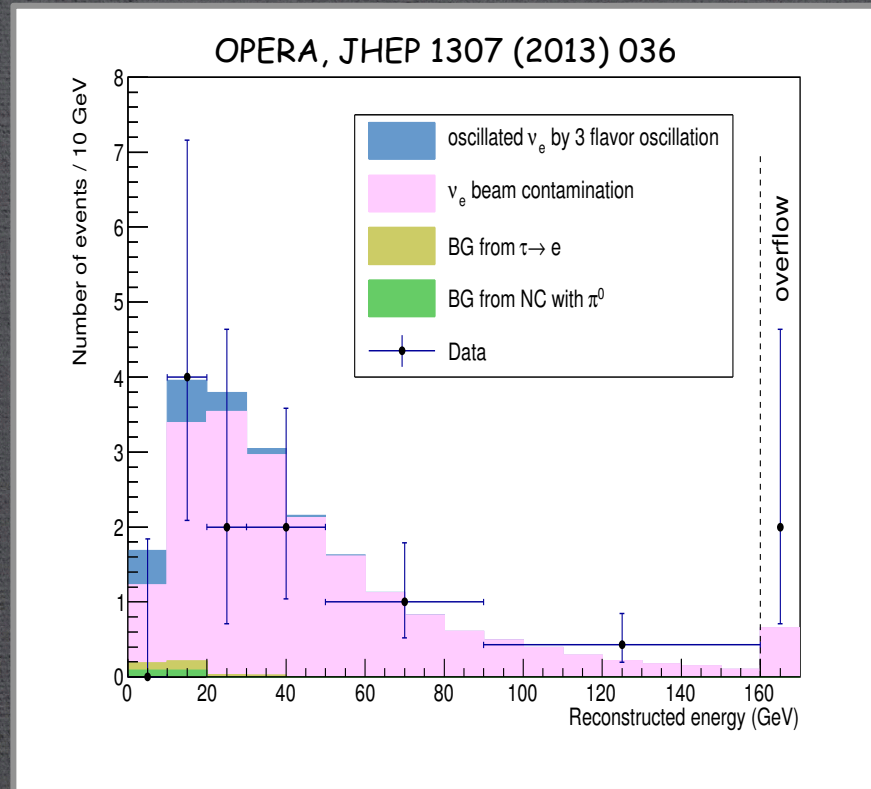
$$P_{ee} \sim 1 - 2 U_{e4}^2 < 1$$

$\nu_e$  bkg is not fixed!

Relevant because  
ICARUS & OPERA  
are bkg-dominated

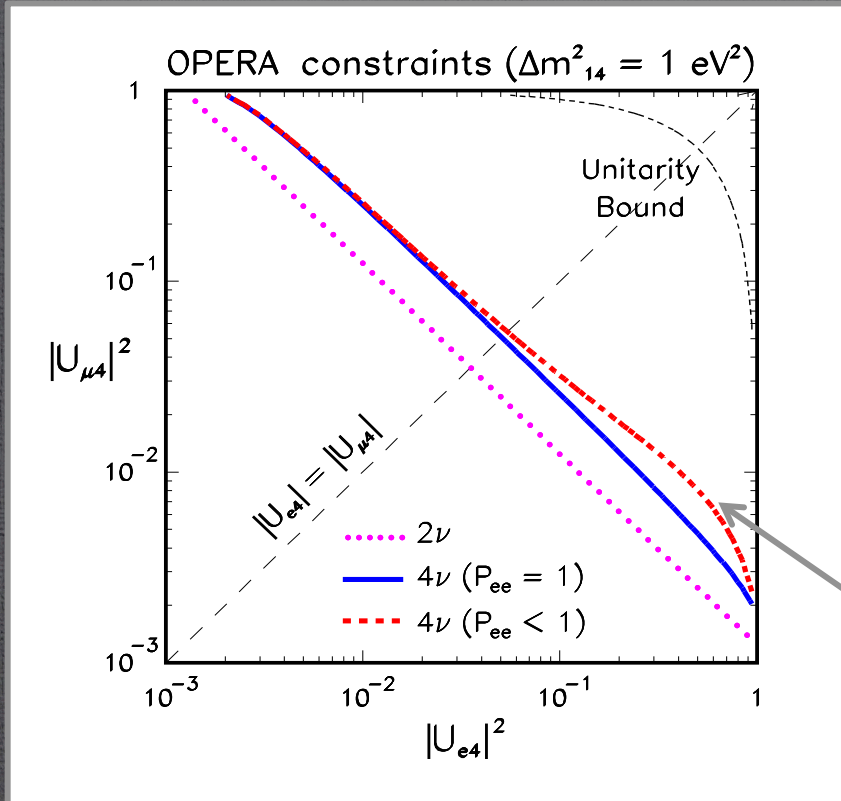
Measured # of events  
smaller than bkg

Expected bkg tends to be lower for  
 $U_{e4} \neq 0$  allowing for a larger signal





# General analysis with $(U_{e4}, U_{\mu4})$ free



Fit prefers big values of  $|U_{e4}|^2$

Larger values of  $\sin^2 2\theta_{\mu e}$  tolerated

$\sin^2 2\theta_{\mu e} < 1.7 \times 10^{-2}$   
at the 90% C.L.

Overall, bounds relaxed by a factor of 3 with respect to the 2-flavor case ( $\sin^2 2\theta_{\mu e} < 5 \times 10^{-3}$ )



# Summary

- Several indications of light sterile  $\nu$  species
- Sterile neutrinos are sources of additional CPV
- LBL expts. can give info on the new CP-phases
- The experiment T2K has already some sensitivity
- Accurate treatment of  $4\nu$  effects is important for a correct interpretation of the LBL results

Investigation of sterile  $\nu$ s and related CPV at LBL experiments is a unique opportunity



Back up slides



# CPV is a genuine 3-flavor effect

$$\Delta_{ij} = \frac{\Delta m_{ij}^2 L}{4E}$$

$$A_{\alpha\beta}^{\text{CP}} \equiv P(\nu_\alpha \rightarrow \nu_\beta) - P(\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta)$$

$$A_{\alpha\beta}^{\text{CP}} = -16 J_{\alpha\beta}^{12} \sin \Delta_{21} \sin \Delta_{13} \sin \Delta_{32}$$

$$J_{\alpha\beta}^{ij} \equiv \text{Im} [U_{\alpha i} U_{\beta j} U_{\alpha j}^* U_{\beta i}^*] \equiv J \sum_{\gamma=e,\mu,\tau} \epsilon_{\alpha\beta\gamma} \sum_{k=1,2,3} \epsilon_{ijk}$$

$J$  is parameterization independent (Jarlskog invariant)

In the standard parameterization:

$$J = \frac{1}{8} \sin 2\theta_{12} \sin 2\theta_{23} \sin 2\theta_{13} \cos \theta_{13} \sin \delta$$

Conditions for CPV:

- No degenerate  $(\nu_i, \nu_j)$  ✓
- No  $\theta_{ij} = (0, \pi/2)$  ✓
- $\delta \neq (0, \pi)$  ?



# Results of the T2K 4ν analysis (IH)

