## Trends in Dark Matter Physics

Paolo Gondolo<br>University of Utah

## Dark matter theory

- Fifty shades of dark
- The forbidden fruit
- Confusion of the mind
- That which does not kill us makes us stronger

Fifty shades of dark

## Evidence for cold dark matter



## Evidence for cold dark matter

Large Scale Structure

Galaxies spin faster or are hotter than gravity of visible mass can support (rotation curves, velocity dispersion)

## Evidence for cold dark matter

Empirical correlations found from thousands of spiral galaxy rotation curves


Salucci et al 2007

## Evidence for cold dark matter



Velocity dispersion measurements reveal dark matter in elliptical galaxies

$$
\begin{aligned}
& \sigma^{2} \propto \frac{G M}{r} \\
& M_{\mathrm{dyn}} \sim 10^{15} M_{\odot}
\end{aligned}
$$

## Evidence for cold dark matter

Dwarf galaxies are dominated by dark matter.



Adams et al 20/4

## Evidence for cold dark matter

## Large Scale Structure

Galaxy clusters are mostly invisible mass (motion of galaxies, gas density and temperature, gravitational lensing)

A. Riess

Dwarf Galaxies

Galaxy Clusters


Abell 1689 (HST/ACS, Benitezz et al: 2003).


## Evidence for cold dark matter

Large Scale Structure


Cosmic Microwave Background


Supernovae

An invisible mass makes the Cosmic Microwave Background fluctuations grow into galaxies (CMB and matter power spectra, or correlation functions)

## Evidence for cold dark matter

## Cosmic Microwave

## Background fluctuations

| Parameter | Planck+WP+highL+BAO |  |
| :---: | :---: | :---: |
|  | Best fit | 68\% limits |
| $\Omega_{\mathrm{b}} h^{2}$ | 0.022161 | $0.02214 \pm 0.00024$ |
| $\Omega_{\mathrm{c}} h^{2}$ | 0.11889 | $0.1187 \pm 0.0017$ |
| $100 \theta_{\text {MC }}$ | 1.04148 | $1.04147 \pm 0.00056$ |
| $\tau$ | 0.0952 | $0.092 \pm 0.013$ |
| $n_{\text {s }}$ | 0.9611 | $0.9608 \pm 0.0054$ |
| $\ln \left(10^{10} A_{\mathrm{s}}\right)$ | 3.0973 | $3.091 \pm 0.025$ |
| $\Omega_{\Lambda}$ | 0.6914 | $0.692 \pm 0.010$ |
| $\sigma_{8}$ | 0.8288 | $0.826 \pm 0.012$ |
| $z_{\text {re }}$ | 11.52 | $11.3 \pm 1.1$ |
| $H_{0}$ | 67.77 | $67.80 \pm 0.77$ |
| Age/Gyr | 13.7965 | $13.798 \pm 0.037$ |
| $100 \theta_{*}$ | 1.04163 | $1.04162 \pm 0.00056$ |
| $r_{\text {drag }}$ | 147.611 | $147.68 \pm 0.45$ |

## Evidence for cold dark matter

## $535 \pm 7 \mathrm{pJ} / \mathrm{m}^{3}$ dark energy



$$
\begin{array}{|l|}
201 \pm 2 \mathrm{pJ} / \mathrm{m}^{3} \\
\text { cold dark matter }
\end{array} \quad \begin{gathered}
\text { Cold Dark } \\
\text { Matter }
\end{gathered}
$$

1 to $4 \mathrm{pJ} / \mathrm{m}^{3}$ neutrinos
matter $p \ll \rho$ radiation $p=\rho / 3$ vacuum $p=-\rho$

Planck (20I5)
TT,TE,EE+lowP+lensing+ext
$1 \mathrm{pJ}=10^{-12} \mathrm{~J}$
$\rho_{\text {crit }}=1.68829 h^{2} \mathrm{pJ} / \mathrm{m}^{3}$

## Evidence for nonbaryonic cold dark matter

## GALAXY FORMATION

Matter fluctuations uncoupled to the plasma can gravitationally grow into galaxies in the given I3 Gyr


Dark matter is non-baryonic More than 80\% of all matter does not couple to the primordial plasma!


## Evidence for nonbaryonic cold dark matter

## BIG BANG NUCLEOSYNTHESIS

- The baryon-to-photon ratio has been the same since $\sim$ I minute after the Big Bang.
- Baryons are $\leqslant 15.7 \%$ of the mass in matter.



## Is dark matter an elementary particle?

## ELEMENTARY PARTICLES



Q couples to the plasma
Q disappears too quickly
Q is hot dark matter

No known particle can be nonbaryonic cold dark matter!

## Physicists have many ideas



## Particle dark matter

- neutrinos
- sterile neutrinos, gravitinos
- lightest supersymmetric particle
- lightest Kaluza-Klein particle
- Bose-Einstein condensates, axions, axion clusters
- solitons (Q-balls, B-balls, ...)
- supermassive wimpzillas

| Mass range |
| :---: |
| $10^{-22} \mathrm{eV}\left(10^{-56} \mathrm{~g}\right)$ B.E.C.s |
| $10^{-8} \mathrm{M}\left(10^{+25} \mathrm{~g}\right)$ axion clusters |

(hot)
(warm)
(cold)
(cold)


Interaction strength range
Only gravitational: wimpzillas
Strongly interacting: B-balls

## Particle dark matter

Hot dark matter

- relativistic at kinetic decoupling (start of free streaming)
- big structures form first, then fragment
light neutrinos
Cold dark matter
- non-relativistic at kinetic decoupling
- small structures form first, then merge neutralinos, axions,WIMPZILLAs, solitons

Warm dark matter

- semi-relativistic at kinetic decoupling
- smallest structures are erased
sterile neutrinos, gravitinos


## Particle dark matter

Thermal relics
in thermal equilibrium in the early universe
neutrinos, neutralinos, other WIMPs, ....
Non-thermal relics
not in thermal equilibrium in the early universe
axions, WIMPZILLAs, solitons, ....

Axions

## Axions as dark matter

## Hot

Produced thermally in early universe
Important for $m_{a}>0.1 e V\left(f_{a}<10^{8}\right)$, mostly excluded by astrophysics

## Cold

Produced by coherent field oscillations around mimimum of $V(\theta)$
(Vacuum realignment)
Produced by decay of topological defects
(Axionic string decays)

## Axion cold dark matter parameter space

PQ symmetry breaking scale


Expansion rate at end of inflation
Visinelli, Gondolo 2009 + updates

## Neutrinos

## Heavy active neutrinos

## PHYSICAL REVIEW LETTERS

25 JULY 1977

Cosmological Lower Bound on Heavy-Neutrino Masses
Deajamin W, Lee ${ }^{\text {Wo }}$

and
Steves Weinberg ${ }^{t 0}$




 seder of it Cev.
$2 \mathrm{GeV} / \mathrm{c}^{2}$ for $\Omega_{\mathrm{c}}=1$
Now $4 \mathrm{GeV} / \mathrm{c}^{2}$ for $\Omega_{\mathrm{c}}=0.25$

## Cosmic density of massive neutrinos

Fourth-generation Standard Model neutrin

$$
\sim \text { few } \mathrm{GeV}
$$ preferred cosmological mass

Excluded as dark matter (199I)
Lee \& Weinberg 1977


## Sterile neutrino dark matter

Standard model + right-handed neutrinos
Active and sterile neutrinos oscillate into each other.

Sterile neutrinos can be warm dark matter (mass > 0.3 keV ) Dodelson, Widrow I994; Shi, Fuller 1999; Laine, Shaposhnikov 2008

## Supersymmetric particles

## Supersymmetric dark matter

Neutralinos (the most fashionable/studied WIMP)
Goldberg I983; Ellis, Hagelin, Nanopoulos, Olive, Srednicki I984; etc.
Sneutrinos (also WIMPs)
Falk, Olive, Srednicki 1994;Asaka, Ishiwata, Moroi 2006; McDonald 2007; Lee, Matchev, Nasri 2007; Deppisch, Pilaftsis 2008; Cerdeno, Munoz, Seto 2009; Cerdeno, Seto 2009; etc.

Gravitinos (SuperWIMPs)
Feng, Rajaraman, Takayama 2003; Ellis, Olive, Santoso, Spanos 2004; Feng, Su, Takayama, 2004; etc.

Axinos (SuperWIMPs)
Tamvakis, Wyler I982; Nilles, Raby I982; Goto, Yamaguchi I992; Covi, Kim, Kim, Roszkowski 200 I; Covi, Roszkowski, Ruiz de Austri, Small 2004; etc.

## Neutralino dark matter: impact of LHC

- The CMSSM is in dire straights

Constrained Minimal
Superssymetric Standard Model
"a Higgs mass of $\sim 125 \mathrm{GeV}$ excludes the least fine-tuned CMSSM points; remaining viable models may be difficult to probe with dark matter searches"

- But there are many supersymmetric models



## Neutralino dark matter: impact of LHC

Cahill-Rowell et al I 305.692 I
"the only pMSSM models remaining [with neutralino being 100\% of CDM] are those with bino coannihilation"
pMSSM (phenomenological MSSM)

$$
\begin{aligned}
& \mu, m_{A}, \tan \beta, A_{b}, A_{t}, A_{\tau}, M_{1}, M_{2}, M_{3} \\
& m_{Q_{1}}, m_{Q_{3}}, m_{u_{1}}, m_{d_{1}}, m_{u_{3}}, m_{d_{3}} \\
& m_{L_{1}}, m_{L_{3}}, m_{e_{1}}, m_{e_{3}} \\
& \text { (l9 parameters) }
\end{aligned}
$$



## Neutralino dark matter: impact of LHC

"Supersymmetry cannot be experimentally ruled out: it can either be discovered or abandoned."

Leszek Roszkowski



## The forbidden fruit

## Searches for particle dark matter



Indirect detection


Annihilation

## The power of the WIMP



Direct detection

Large scale structure


## Dark matter creation with particle accelerators

Searching for the conversion protons $\rightarrow$ energy $\rightarrow$ dark matter


The ATLAS detector
Particle production at the Large Hadron Collider

## Indirect detection of particle dark matter

## The principle

Dark matter particles transform into ordinary particles, which are then detected or inferred

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Dark matter particles transform into ordinary particles, which are then detected or inferred

Gunn, Lee, Lerche, Schramm, Steigman 1978; Stecker 1978



## Indirect detection of particle dark matter

## The principle

Dark matter particles transform into ordinary particles, which are then detected or inferred

The first stars to form in the universe may have been powered by dark matter instead of nuclear fusion.

They were dark-matter powered stars or for short
Dark Stars

- Explain chemical elements in old halo stars
- Explain origin of supermassive black holes in early quasars

Spolyar, Freese, Gondolo 2007-2008

## The principle of direct detection

## Dark matter particles that arrive on Earth scatter off nuclei in a detector

Goodman,

Witten
1985


Low-background underground detector

## Expected event rate is small

## Expected <br> WIMP spectrum


~I event/kg/year (nuclear recoils)

## Expected event rate is small

## Expected <br> WIMP spectrum


~I event/kg/year (nuclear recoils)

Measured
banana spectrum

~100 events/kg/second (electron recoils)

## Expected event rate is small

Expected
WIMP spectrum
Measured
banana spectrum


## Confusion of the mind

## Evidence for cold dark matter particles?

## $\mathrm{GeV} \gamma$-rays

Hooper et al
$2009-14$
3.5 keV X-ray line


Bulbul et al 2014
$135 \mathrm{GeV} \gamma$-ray line


Weniger 2012

Annual modulation

## Positron excess



Adriani et al 2009;Ackerman et al 201 I; Aguilar et al 2013

Drukier, Freese, Spergel 1986
 |997-20| 2

## Gamma-rays from dark matter?

## 1 GeV gamma-ray excess?

Goodenough, Hooper 2009; Hooper, Goodenough; Boyarsky, Malyshev, Ruchayskiy; Hooper, Linden 20 I I;Abazajian, Kaplinghat 20I 2; Gordon, Macias 20 I3;Abazajian, Canac, Horiuchi, Kaplinghat; Daylan et al 2014

Fermi-LAT

Fit diffuse + Fermi-bubble, find residual



## Gamma-rays from dark matter (2015)

(similar for $\tau^{+} \tau^{-}, W^{+} W^{-}, \ldots$ )


## Gamma-rays from dark matter (2015)

(similar for $\tau^{+} \tau^{-}$)


Positrons from dark matter?

## Excess in cosmic ray positrons

Ackernmann et al [Fermi-LAT] 201 I

High energy cosmic ray positrons are more than expected


Adriani et al. [PAMELA ,2008


Accardo et al [AMS-02] 2014

## Excess in cosmic ray positrons

Grasso et al [Fermi-LAT] 2009

## Dark matter?

Pulsars?
Secondaries from extra primaries?


Bergstrom, Edsjo, Zaharijas 2009



## Dynamical dark matter

Dienes, Thomas 20 I I, 2012
Dienes, Kumar, Thomas 20I2, 20I3
A vast ensemble of fields decaying one into another
Example: Kaluza-Klein tower of axions in extra-dimensions


Phenomenology obtained through scaling laws

$$
\begin{aligned}
& m_{n}=m_{0}+n^{\delta} \Delta m \\
& \rho_{n} \sim m_{n}^{\alpha}, \tau_{n} \sim m_{n}^{-\gamma}
\end{aligned}
$$

X-rays from dark matter?

## Sterile neutrino dark matter

The main decay mode of keV sterile neutrinos $\left(v_{s} \rightarrow 3 v\right)$ is undetectable Radiative decay of sterile neutrinos $\nu_{s} \rightarrow \gamma \nu_{a}$

$$
\begin{aligned}
& \text { X-ray line } \\
& E_{\gamma}=\frac{1}{2} m_{s}
\end{aligned}
$$



Figure from Kusenko 0906.2968

$$
\begin{aligned}
\Gamma_{\nu_{s} \rightarrow \gamma \nu_{a}} & =\frac{9}{256 \pi^{4}} \alpha_{\mathrm{EM}} \mathrm{G}_{\mathrm{F}}^{2} \sin ^{2} \theta m_{\mathrm{s}}^{5} \\
& =\frac{1}{1.8 \times 10^{21} \mathrm{~S}} \sin ^{2} \theta\left(\frac{m_{\mathrm{s}}}{\mathrm{keV}}\right)^{5}
\end{aligned}
$$

## Sterile neutrino dark matter

An unidentified 3.5 keV X-ray line has been reported in stacked images of 73 galaxy clusters and in the Andromeda galaxy

Bulbul et al 2014
Boyarsky et al 2014



## Sterile neutrino dark matter

```
            vMSM
mv}=7.1 ke
    \mp@subsup{\operatorname{sin}}{}{2}(20)=7\times1\mp@subsup{0}{}{-11}
```



Laine, Shaposhnikov 2008

## Direct detection of dark matter?

## Annual modulation in direct detection

- DAMA observes more nuclei are "hit" in Summer, fewer in Winter


Bernabei et al 2003-2008

- This is exactly what is expected of dark matter WIMPs

Drukier, Freese, Spergel I986


## DAMA modulation

## Model Independent Annual Modulation Result

DAMA/Nal + DAMA/LIBRA-phase 1 Total exposure: $487526 \mathrm{~kg} \times$ day $=\mathbf{1 . 3 3}$ ton $\times \mathbf{y r}$
EPJC 56(2008)333, EPJC 67(2010)39, EPJC 73(2013)2648
The measured modulation amplitudes (A), period (T) and phase $\left(t_{0}\right)$ from the single-hit residual rate vs time

|  | A(cpd/kg/keV) | $\mathbf{T}=\mathbf{2} \pi / \omega$ ( $\mathbf{r r}$ ) | $t_{0}$ (day) | C.L. | $\operatorname{Acos}\left[\omega\left(t-t_{0}\right)\right]$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| DAMA/Nal+DAMA/LIBRA-phase1 |  |  |  |  |  |
| (2-4) keV | $0.0190 \pm 0.0020$ | $0.996 \pm 0.0002$ | $134 \pm 6$ | $9.5 \sigma$ |  |
| (2-5) keV | $0.0140 \pm 0.0015$ | $0.996 \pm 0.0002$ | $140 \pm 6$ | $9.3 \sigma$ |  |
| (2-6) keV | $0.0112 \pm 0.0012$ | $0.998 \pm 0.0002$ | $144 \pm 7$ |  |  |



[^0]Comparison between single hit residual rate (red points) and multiple hit residual rate (green points); Clear modulation in the single hit events; No modulation in the residual rate of the multiple hit events $A=-(0.0005 \pm 0.0004) \mathrm{cpd} / \mathrm{kg} / \mathrm{keV}$


This result offers an additional strong support for the presence of DM particles in the galactic halo further excluding any side effect either from hardware or from software procedures or from background

## DAMA modulation

## Model Independent Annual Modulation Result

DAMA/NaI + DAMA/LIBRA-phase 1 Total exposure: $487526 \mathrm{~kg} \times \mathrm{day}=\mathbf{1 . 3 3}$ ton $\times \mathbf{y r}$
EPJC 56(2008)333, EPJC 67(2010)39, EPJC 73(2013)2648

- No modulation above 6 keV
- No modulation in the whole energy spectrum
- No modulation in the $2-6 \mathrm{keV}$ multiple-hit events
$R(t)=S_{0}+S_{m} \cos \left[\omega\left(t-t_{0}\right)\right]$
here $T=2 \pi / \omega=1 \mathrm{yr}$ and $t_{0}=152.5$ day



No systematics or side processes able to quantitatively account for the measured modulation amplitude and to simultaneously satisfy the many peculiarities of the signature are available.

## DAMA modulation



## CoGeNT made their data public

CoGeNT decided to publish energy and time of their events
Independent groups reanalyzed the CoGeNT data
Pulse-shape discrimination of surface/bulk events

No significant annual modulation found

```
The CoGeNT region of interest results from a biased analysis, and has no statistical meaning.
```

Davis, McCabe, Boehm I 405.0495
The likelihood gets worse when including a WIMP component either as a standard halo or Sagittarius like stream

Bellis, Collar, Field, Kelso at IDM2OI4

## CoGeNT made their data public

## CoGeNT decided to publish energy and time of their events

Maximum Likelihood Signal Extraction Method Applied to 3.4 years of CoGeNT Data C.E. Aalseth, ${ }^{1}$ P.S. Barbeau, ${ }^{2, *}$ J. Colaresi, ${ }^{3}$ J.I. Collar, ${ }^{2}$ J. Diaz Leon, ${ }^{4}$ J.E. Fast, ${ }^{1}$ N.E. Fields, ${ }^{2}$ T.W. Hossbach, ${ }^{1}$
A. Knecht, ${ }^{4, \dagger}$ M.S. Kos, ${ }^{1, \ddagger}$ M.G. Marino,,${ }^{4}$ H.S. Miley, ${ }^{1}$ M.L. Miller, ${ }^{4}$, ${ }^{\text {® }}$ J.L. Orrell, ${ }^{1}$ and K.M. Yocum ${ }^{3}$ (CoGeNT Collaboration)

$$
\text { arXiv:1401.6234v1 } 24 \text { Jan } 2014
$$

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$$
\text { arXiv:1401.6234v2 } 27 \text { Jan } 2014
$$

The likelihood gets worse when including a WIMP component either as a standard halo or Sagittarius like stream

Bellis Collar, Field, Kelso at IDM2014
CoGeNT leader

## Direct WIMP searches (2015)



Billard et al 2013, Snowmass 2013, LUX 20I3, SuperCDMS 2014

## Direct WIMP searches (2015)

## Spin-dependent interactions




Best limits are from IceCube from the Sun)
(indirect detection of high-energy neutrindion) 20/3, Oberlack at IDM2014

## Evidence for light dark matter particles?



No significant modulation
Same target material
Ahmed et al (CDMS)
I203.I309

Not so many events
Akerib et al (LUX) 2013


## That which does not kill us makes us stronger

## Make no assumptions

## All particle physics models

- Consider all possible interactions between dark matter and standard model particles
- This program has been carried out in some limits (e.g., non-relativistic conditions, heavy mediators)


## All astrophysical models

- Halo-independent methods of analysis have been developed
- Ideally they require no assumption on the astrophysical density and velocity distributions of dark matter particles


## All particle physics models

Write down and analyze all possible WIMP interactions with ordinary matter

## Effective operators

if mediator mass > exchanged energy


Four-particle effective operator

There are many possible operators. Interference is important although often, but not always, neglected.

Long(ish) distance interactions are not included.

## Effective operators: LHC \& direct detection

| Name | Operator | Coefficient |
| :---: | :---: | :---: |
| D1 | $\bar{\chi} \chi \bar{q} q$ | $m_{q} / M_{*}^{3}$ |
| D2 | $\bar{\chi} \gamma^{5} \chi \bar{q} q$ | $i m_{q} / M_{*}^{3}$ |
| D3 | $\bar{\chi} \chi \bar{q} \gamma^{5} q$ | $i m_{q} / M_{*}^{3}$ |
| D4 | $\bar{\chi} \gamma^{5} \chi \bar{q} \gamma^{5} q$ | $m_{q} / M_{*}^{3}$ |
| D5 | $\bar{\chi} \gamma^{\mu} \chi \bar{q} \gamma_{\mu} q$ | $1 / M_{*}^{2}$ |
| D6 | $\bar{\chi} \gamma^{\mu} \gamma^{5} \chi \bar{q} \gamma_{\mu} q$ | $1 / M_{*}^{2}$ |
| D7 | $\bar{\chi} \gamma^{\mu} \chi \bar{q} \gamma_{\mu} \gamma^{5} q$ | $1 / M_{*}^{2}$ |
| D8 | $\bar{\chi} \gamma^{\mu} \gamma^{5} \chi \bar{q} \gamma_{\mu} \gamma^{5} q$ | $1 / M_{*}^{2}$ |
| D9 | $\bar{\chi} \sigma^{\mu \nu} \chi \bar{q} \sigma_{\mu \nu} q$ | $1 / M_{*}^{2}$ |
| D10 | $\bar{\chi} \sigma_{\mu \nu} \gamma^{5} \chi \bar{q} \sigma_{\alpha \beta} q$ | $i / M_{*}^{2}$ |
| D11 | $\bar{\chi} \chi G_{\mu \nu} G^{\mu \nu}$ | $\alpha_{s} / 4 M_{*}^{3}$ |
| D12 | $\bar{\chi} \gamma^{5} \chi G_{\mu \nu} G^{\mu \nu}$ | $i \alpha_{s} / 4 M_{*}^{3}$ |
| D13 | $\bar{\chi} \chi G_{\mu \nu} \tilde{G}^{\mu \nu}$ | $i \alpha_{s} / 4 M_{*}^{3}$ |
| D14 | $\bar{\chi} \gamma^{5} \chi G_{\mu \nu} \tilde{G}^{\mu \nu}$ | $\alpha_{s} / 4 M_{*}^{3}$ |


| Name | Operator | Coefficient |
| :---: | :---: | :---: |
| C 1 | $\chi^{\dagger} \chi \bar{q} q$ | $m_{q} / M_{*}^{2}$ |
| C 2 | $\chi^{\dagger} \chi \bar{q} \gamma^{5} q$ | $i m_{q} / M_{*}^{2}$ |
| C 3 | $\chi^{\dagger} \partial_{\mu} \chi \bar{q} \gamma^{\mu} q$ | $1 / M_{*}^{2}$ |
| C 4 | $\chi^{\dagger} \partial_{\mu} \chi \bar{q} \gamma^{\mu} \gamma^{5} q$ | $1 / M_{*}^{2}$ |
| C 5 | $\chi^{\dagger} \chi G_{\mu \nu} G^{\mu \nu}$ | $\alpha_{s} / 4 M_{*}^{2}$ |
| C 6 | $\chi^{\dagger} \chi G_{\mu \nu} \tilde{G}^{\mu \nu}$ | $i \alpha_{s} / 4 M_{*}^{2}$ |
| R 1 | $\chi^{2} \bar{q} q$ | $m_{q} / 2 M_{*}^{2}$ |
| R 2 | $\chi^{2} \bar{q} \gamma^{5} q$ | $i m_{q} / 2 M_{*}^{2}$ |
| R 3 | $\chi^{2} G_{\mu \nu} G^{\mu \nu}$ | $\alpha_{s} / 8 M_{*}^{2}$ |
| R 4 | $\chi^{2} G_{\mu \nu} \tilde{G}^{\mu \nu}$ | $i \alpha_{s} / 8 M_{*}^{2}$ |

Table of effective operators relevant for the collider/direct detection connection

Goodman, Ibe, Rajaraman, Shepherd, Tait, Yu 2010

## Effective operators: LHC \& direct detection

LHC limits on WIMP-quark and WIMP-gluon interactions are competitive with direct searches

Beltran et al, Agrawal et al., Goodman et al., Bai et al., 20I0; Goodman et al., Rajaraman et al. Fox et al., 20II; Cheung et al., Fitzptrick et al., March-Russel et al., Fox et al., 20 I2.......


> These bounds do not apply to SUSY, etc.

Complete theories contain sums of operators (interference) and not-so-heavy mediators (Higgs)

Fox, Harnik, Primulando, Yu 2012

## Effective operators: direct detection

All short-distance operators classified
Fitzpatrick et al 2012

$$
\begin{array}{lllll}
\text { 1, } \quad \vec{S}_{\chi} \cdot \vec{S}_{N}, & v^{2}, & i\left(\vec{S}_{\chi} \times \vec{q}\right) \cdot \vec{v}, & i \vec{v} \cdot\left(\vec{S}_{N} \times \vec{q}\right), & \left(\vec{S}_{\chi} \cdot \vec{q}\right)\left(\vec{S}_{N} \cdot \vec{q}\right) \quad i \vec{S}_{N} \cdot \vec{q}, i \vec{S}_{\chi} \cdot \vec{q}, \\
\vec{v}^{\perp} \cdot \vec{S}_{\chi}, & \vec{v}^{\perp} \cdot \vec{S}_{N}, & i \vec{S}_{\chi} \cdot\left(\vec{S}_{N} \times \vec{q}\right) . & \left(i \vec{S}_{N} \cdot \vec{q}\right)\left(\vec{v}^{\perp} \cdot \vec{S}_{\chi}\right), & \left(i \vec{S}_{\chi} \cdot \vec{q}\right)\left(\vec{v}^{\perp} \cdot \vec{S}_{N}\right) .
\end{array}
$$

All nuclear form factors classified

| Response $\times\left[\frac{4 \pi}{2 J_{i}+1}\right]^{-1}$ | Leading <br> Multipole | Long-wavelength Limit | Response <br> Type |
| :---: | :---: | :---: | :---: |
| $\sum_{=}^{\infty}\left\|\left\langle J_{i}\left\\|M_{J M}\right\\| J_{i}\right\rangle\right\|^{2}$ | $M_{00}\left(q \vec{x}_{i}\right)$ | $\frac{1}{\sqrt{4 \pi}} 1(i)$ | $M_{J M}$ : Charge |
| $\left.\sum_{J=1,3, \ldots}^{\infty}\left\|\left\langle J_{i}\right\|\right\| \Sigma_{J M}^{\prime \prime} \\| J_{i}\right\rangle\left.\right\|^{2}$ | $\Sigma_{1 M}^{\prime \prime}\left(q \vec{x}_{i}\right)$ | $\frac{1}{2 \sqrt{3 \pi}} \sigma_{1 M}(i)$ | $L_{J M}^{5}$ : Axial <br> Longitudinal |
| $\left.\sum_{J=1,3, \ldots}^{\infty}\left\|\left\langle J_{i}\right\|\right\| \Sigma_{J M}^{\prime} \\| J_{i}\right\rangle\left.\right\|^{2}$ | $\Sigma_{1 M}^{\prime}\left(q \vec{x}_{i}\right)$ | $\frac{1}{\sqrt{6 \pi}} \sigma_{1}$ | $T_{J M}^{\mathrm{el}}$ : Axial Transverse Electric |
| $\sum_{J=1,3, \ldots}^{\infty}\left\|\left\langle J_{i}\left\\|\frac{q}{m_{N}} \Delta_{J M}\right\\| J_{i}\right\rangle\right\|^{2}$ | $\frac{q}{m_{N}} \Delta_{1 M}\left(q \vec{x}_{i}\right)$ | $-\frac{q}{2 m_{N} \sqrt{6 \pi}} \ell_{1 M}(i$ | $T_{J M}^{\text {mag }}:$ Transverse Magnetic |
| $\sum_{J=0,2, \ldots}^{\infty}\left\|\left\langle J_{i}\left\\|\frac{q}{m_{N}} \Phi_{J M}^{\prime \prime}\right\\| J_{i}\right\rangle\right\|^{2}$ | $\frac{q}{m_{N}} \Phi_{00}^{\prime \prime}\left(q \vec{x}_{i}\right)$ | $-\frac{q}{3 m_{N} \sqrt{4 \pi}} \vec{\sigma}(i) .$ | $L_{J M}$ : <br> Longitudinal |
| $\sum_{J=2,4, \ldots}^{\infty}\left\|\left\langle J_{i}\left\\|\frac{q}{m_{N}} \tilde{\Phi}_{J M}^{\prime}\right\\| J_{i}\right\rangle\right\|^{2}$ | $\begin{aligned} & \frac{q}{m_{N}} \Phi_{2 M}^{\prime \prime}\left(q \vec{x}_{i}\right) \\ & \frac{q}{m_{N}} \tilde{\Phi}_{2 M}^{\prime}\left(q \vec{x}_{i}\right) \end{aligned}$ | $\begin{aligned} & -\frac{q}{m_{N} \sqrt{30 \pi}}\left[x_{i} \otimes\left(\vec{\sigma}(i) \times \frac{1}{i} \vec{\nabla}\right)_{1}\right]_{2 M} \\ & -\frac{q}{m_{N} \sqrt{20 \pi}}\left[x_{i} \otimes\left(\vec{\sigma}(i) \times \frac{1}{i} \vec{\nabla}\right)_{1}\right]_{2 M} \end{aligned}$ | $T_{J M}^{\mathrm{el}}$ : <br> Transverse Electric |

nuclear oscillator model
Fitzpatrick et al 2012

## Effective operators: direct detection

Experimental limits on single operators...
Schneck et al (SuperCDMS) 2015


| Operater confficieat | SuperCDMES Soudan |
| :---: | :---: |
| $\left(\mathrm{Cl}^{(1)}\right)^{2} \cdot \mathrm{~m}_{\text {+ }}^{4}$ | $8.98 \times 10^{-5}(-)$ |
| $(\mathrm{el})^{2} * \mathrm{~m}^{4}$ | $3.14 \times 10^{4}(-)$ |
| $\left.(4)^{2}\right)^{2} \cdot \mathrm{~m}^{4}$ | $8.77 \times 10^{1}(-)$ |
| $(\mathrm{c})^{2} * \mathrm{~m}^{4}+$ | $6.34 \times 10^{5}(-)$ |
| $\left(\epsilon_{5}\right)^{2} \cdot \mathrm{~m}^{4}+\mathrm{t}$ | $4.54 \times 10^{8}(-)$ |
| $(\mathrm{d})^{2} \cdot \mathrm{~m}^{4}+$ | $8.4 .4 \times 10^{7}(-)$ |
| $\left(e^{(2)}\right)^{2} \cdot m^{4}$ cou | $4.30 \times 10^{2}(-)$ |
| $\left(c^{9}\right)^{2}=\mathrm{m}_{\text {+ }}^{4}$ | $1.95 \times 10^{3}(-)$ |
| $\left(\mathrm{cos}_{0}\right)^{2} \times \mathrm{m}^{4}+$ | $9.22 \times 10^{4}(-)$ |
|  | $5.13 \times 10^{-1}(-)$ |
| $\left(\mathrm{cfin}^{2}\right)^{2}+\mathrm{m}_{-1-1}^{4}$ | $1.00 \times 10^{3}(-)$ |
| $\left(\mathrm{c}_{2}\right)^{7}+\mathrm{m}^{4}$ - | $4.28 \times 10^{8}(-)$ |
| $\left(\mathrm{c}_{4}\right)^{2}+\mathrm{m}^{4}$ - | $5.00 \times 10^{11}(-)$ |
| $\left(c_{i n}\right)^{2}+\mathrm{m}_{\text {ctat }}^{4}$ | $1.32 \times 10^{2}(-)$ |

## Effective operators: direct detection

## Combined analysis of short-distance operators



## Effective operators: direct detection

## Combined analysis of short-distance operators



## All astrophysics models

Do not assume any particular
WIMP density or velocity distribution

## DM-nucleus elastic scattering

$\binom{$ event }{ rate }$=\binom{$ detector }{ response }$\times\binom{$ particle }{ physics }$\times($ astrophysics $)$

Dark matter particle


## Detector response model

$$
\binom{\text { event }}{\text { rate }}=\binom{\text { detector }}{\text { response }} \times\binom{\text { particle }}{\text { physics }} \times(\text { astrophysics })
$$

## Is a nuclear recoil detectable?

Counting efficiency, energy resolution, scintillation response, etc.

$$
\binom{\text { detector }}{\text { response }}=\mathcal{G}\left(E, E_{R}\right)
$$

Probability of detecting an event with energy (or number of photoelectrons) $E$, given an event occurred with recoil energy $E_{R}$.

## Particle physics model

## $\binom{$ event }{ rate }$=\binom{$ detector }{ response }$\times\binom{$ particle }{ physics }$\times$ (astrophysics)

## What force couples dark matter to nuclei?

Coupling to nucleon number density, nucleon spin density, ...


## Astrophysics model

$$
\binom{\text { event }}{\text { rate }}=\binom{\text { detector }}{\text { response }} \times\binom{\text { particle }}{\text { physics }} \times \text { (astrophysics) }
$$

## How much dark matter comes to Earth?

Velocity distribution
Local halo density
$($ astrophysics $)=\eta\left(v_{\min }, t\right) \equiv \rho_{\chi} \int_{v>v_{\min }} \frac{f(\mathbf{v}, t)}{v} \mathrm{~d}^{3} v$
Minimum WIMP speed to impart recoil energ
$v_{\min }=\left(M E_{R} / \mu+\delta\right) / \sqrt{2 M E_{R}}$

## Astrophysics model: velocity distribution

Standard Halo Model

| truncated |
| :--- | :--- |
| Maxwellian |\(f(\mathbf{v})= \begin{cases}\frac{1}{N_{esc} \pi^{3 / 2} \overline{\widetilde{v}}_{0}^{3}} e^{-\left|\mathbf{v}+\mathbf{v}_{obs}\right| / \bar{v}_{0}^{2}} \& |\mathbf{v}|<v_{esc} <br>

0 \& otherwise\end{cases}\)

The spherical cow of direct WIMP searches

## Astrophysics model: local density

The dark matter density near the Solar System is known reasonably well


## Astrophysics model: velocity distribution

We know very little about the dark matter velocity distribution near the Sun


Cosmological N-Body simulations including baryons are challenging but underway


Odenkirchen et al 2002 (SDSS) Streams of stars have been observed in the galactic halo SDSS, 2MASS, SEGUE, .......

## Astrophysics model: velocity distribution

## $\binom{$ event }{ rate }$=\binom{$ detector }{ response }$\times \underbrace{\binom{\text { particle }}{\text { physics }}}_{\text {FIXED }} \times \underbrace{(\text { astrophysics })}_{\text {FIXED }}$



Agnese et al (SuperCDMS) 2014

## Astrophysics-independent approach



## Astrophysics-independent approach

Gondolo Gelmini 2012

- The measured rate is a "weighted average" of the astrophysical factor.


## Measured rate

$$
R=\int_{0}^{\infty} d v \mathcal{R}(v) \tilde{\eta}(v)
$$

Response function

- Every experiment is sensitive to a "window in velocity space."



## Spin-independent isoscalar interactions

$$
\sigma_{\chi A}=A^{2} \sigma_{\chi \mathrm{p}} \mu_{\chi A}^{2} / \mu_{\chi \mathrm{p}}^{2}
$$

## Astrophysics-independent

 approach

CDMS-Si event rate is similar to yearly modulated rates

Still depends on particle model

## In the next episodes

## In the next episodes..... Revenge



## In the next episodes..... Giant detectors

## SuperCDMS, XENON1T, XENONnT, Darwin, .......



[^1]
## In the next episodes..... Precision cosmic rays

## AMS (Alpha Magnetic Spectrometer)

AMS-02 can measure isotopic ratios to $\sim 1 \%$ precision up to Fe and $\sim 100 \mathrm{GeV} /$ nucleon, and much better at lower energies.


## In the next episodes..... WIMP astronomy

- Directional direct detection
- measure direction of nuclear recoil
- Several R\&D efforts
- DRIFT
- Dark Matter TPC
- NEWAGE
- MIMAC
- D3
- Emulsion Dark Matter Search
- Columnar recombination


DMTPC

Only ~10 events needed to confirm extraterrestrial signal

## In the next episodes..... WIMP astronomy



## Synopsis

- Fifty shades of dark
- There is evidence for nonbaryonic cold dark matter.
- There are many candidates for nonbaryonic dark matter particles.
- The forbidden fruit
- WIMP interaction rates in direct searches are very small.
- No bananas in the lab.
- Confusion of the mind
- Some experiments claim dark matter detection while others exclude it.
- That which does not kill us makes us stronger
- Move to consider all possible WIMP-SM currents.
- Do not assume any specific dark halo model.


[^0]:    No systematics or side reaction able to account for the measured modulation amplitude and to satisty all the peculiarities of the signature

[^1]:    Oberlack, IDM20/4

