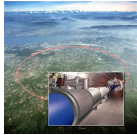


Developments in leptogenesis

Mathias Garny (CERN)

WIN MPIK June 9, 2015

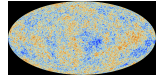
Physics beyond the Standard Model



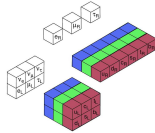
Collider exp.



Baryon
asymmetry

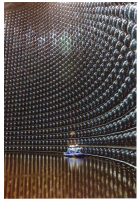


⋮



⋮

+ ?



Neutrino exp.



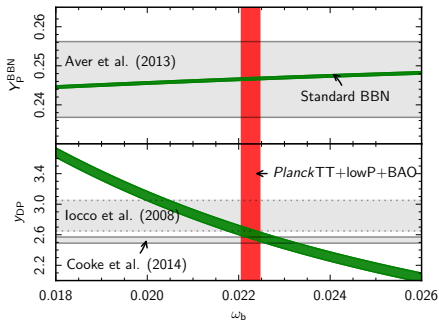
Dark matter



Baryon asymmetry

$$Y_P = 4n_{\text{He}}/n_b$$

$$y_{DP} = n_D/n_H \cdot 10^5$$



$$\omega_b = \Omega_b h^2 = \eta / (2.74 \cdot 10^{-8})$$

Consistent value inferred from Big Bang Nucleosynthesis ($T \sim \text{keV}$) and the cosmic microwave background ($T \sim \text{eV}$)

$$\eta = \frac{n_b - n_{\bar{b}}}{n_\gamma} = \begin{cases} (6.15 \pm 0.15) \cdot 10^{-10} & \text{WMAP9} \\ (6.10 \pm 0.08) \cdot 10^{-10} & \text{Planck} \end{cases}$$

Baryogenesis

- ▶ Electroweak baryogenesis (not in the SM, EWPT cross-over for $m_H \gtrsim 72$ GeV; MSSM: light RH stop would show up in direct search + Higgs properties)
- ▶ Leptogenesis
- ▶ Affleck-Dine
- ▶ GUT baryogenesis
- ▶ ...

Baryogenesis via leptogenesis

Fukugita, Yanagida 86

Light neutrino masses via seesaw mechanism (focus on type-I)*

$$\mathcal{L} = \mathcal{L}_{SM} + \frac{1}{2} \bar{N} i \not{\partial} N - \frac{1}{2} \bar{N} \hat{M} N - \bar{\ell} \tilde{\phi} h P_R N - \bar{N} P_L h^\dagger \tilde{\phi}^\dagger \ell$$

$$m_\nu = -v_{EW}^2 h \hat{M}^{-1} h^T$$

$$0.06\text{eV} \lesssim \sum m_\nu \lesssim 0.14..0.59\text{eV} \quad \rightarrow \quad \text{TeV} \lesssim M_i \lesssim M_{GUT} \text{ for } m_e/v_{EW} < h_{ij} < 1$$

Ly α BOSS 1410.7244 .. Planck 1502.01589 (0.23eV w BAO)

* recent analysis in type-II: Lavignac, Schmauch 1503.00629

Light neutrino masses via seesaw mechanism (focus on type-I)*

$$\mathcal{L} = \mathcal{L}_{SM} + \frac{1}{2} \bar{N}_i \not{\partial} N - \frac{1}{2} \bar{N} \hat{M} N - \bar{\ell} \tilde{\phi} h P_R N - \bar{N} P_L h^\dagger \tilde{\phi}^\dagger \ell$$

$$m_\nu = -v_{EW}^2 h \hat{M}^{-1} h^T$$

$$0.06\text{eV} \lesssim \sum m_\nu \lesssim 0.14..0.59\text{eV} \quad \rightarrow \quad \text{TeV} \lesssim M_i \lesssim M_{GUT} \text{ for } m_e/v_{EW} < h_{ij} < 1$$

Ly α BOSS 1410.7244 .. Planck 1502.01589 (0.23eV w BAO)

- ▶ B-violation via L-violating Majorana masses M_i
- ▶ CP-violation via Yukawa couplings $\text{Im}[(h^\dagger h)_{ij}] \neq 0$
- ▶ Out-of-equilibrium (inverse) decay $N_i \leftrightarrow \ell \phi^\dagger$ and $N_i \leftrightarrow \ell^c \phi$

$$(\Gamma_i/H)|_{T=M_i} \simeq \frac{(h^\dagger h)_{ii} M_i / 8\pi}{1.66 g_* T^2 / M_{pl}} \Big|_{T=M_i}$$

Light neutrino masses via seesaw mechanism (focus on type-I)*

$$\mathcal{L} = \mathcal{L}_{SM} + \frac{1}{2} \bar{N}_i \not{\partial} N - \frac{1}{2} \bar{N} \hat{M} N - \tilde{\ell} \phi h P_R N - \bar{N} P_L h^\dagger \tilde{\phi}^\dagger \ell$$

$$m_\nu = -v_{EW}^2 h \hat{M}^{-1} h^T$$

$$0.06\text{eV} \lesssim \sum m_\nu \lesssim 0.14..0.59\text{eV} \quad \rightarrow \quad \text{TeV} \lesssim M_i \lesssim M_{GUT} \text{ for } m_e/v_{EW} < h_{ij} < 1$$

Ly α BOSS 1410.7244 .. Planck 1502.01589 (0.23eV w BAO)

- ▶ B-violation via L-violating Majorana masses M_i
- ▶ CP-violation via Yukawa couplings $\text{Im}[(h^\dagger h)_{ij}] \neq 0$
- ▶ Out-of-equilibrium (inverse) decay $N_i \leftrightarrow \ell \phi^\dagger$ and $N_i \leftrightarrow \ell^c \phi$

$$(\Gamma_i/H)|_{T=M_i} \simeq \frac{\tilde{m}_i/8\pi}{1.66g_*v_{EW}^2/M_{pl}} \quad \text{where } \tilde{m}_i = v_{EW}^2(h^\dagger h)_{ii}/M_i$$

Light neutrino masses via seesaw mechanism (focus on type-I)*

$$\mathcal{L} = \mathcal{L}_{SM} + \frac{1}{2} \bar{N}_i \hat{\phi} N - \frac{1}{2} \bar{N} \hat{M} N - \bar{\ell} \tilde{\phi} h P_R N - \bar{N} P_L h^\dagger \tilde{\phi}^\dagger \ell$$

$$m_\nu = -v_{EW}^2 h \hat{M}^{-1} h^T$$

$$0.06\text{eV} \lesssim \sum m_\nu \lesssim 0.14..0.59\text{eV} \quad \rightarrow \quad \text{TeV} \lesssim M_i \lesssim M_{GUT} \text{ for } m_e/v_{EW} < h_{ij} < 1$$

Ly α BOSS 1410.7244 .. Planck 1502.01589 (0.23eV w BAO)

- ▶ B-violation via L-violating Majorana masses M_i
- ▶ CP-violation via Yukawa couplings $\text{Im}[(h^\dagger h)_{ij}] \neq 0$
- ▶ Out-of-equilibrium (inverse) decay $N_i \leftrightarrow \ell \phi^\dagger$ and $N_i \leftrightarrow \ell^c \phi$

$$\begin{aligned}
 (\Gamma_i/H)|_{T=M_i} &\simeq \frac{\tilde{m}_i/8\pi}{1.66g_* v_{EW}^2/M_{pl}} \quad \text{where } \tilde{m}_i = v_{EW}^2 (h^\dagger h)_{ii}/M_i \\
 &\simeq \tilde{m}_i/\text{meV}
 \end{aligned}$$

Baryogenesis via leptogenesis

Light neutrino masses via seesaw mechanism (focus on type-I)*

$$\mathcal{L} = \mathcal{L}_{SM} + \frac{1}{2} \bar{N}_i \not{\partial} N - \frac{1}{2} \bar{N} \hat{M} N - \tilde{\ell} \phi h P_R N - \bar{N} P_L h^\dagger \tilde{\phi}^\dagger \ell$$

$$m_\nu = -v_{EW}^2 h \hat{M}^{-1} h^T$$

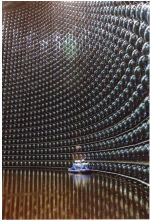
$$0.06\text{eV} \lesssim \sum m_\nu \lesssim 0.14..0.59\text{eV} \quad \rightarrow \quad \text{TeV} \lesssim M_i \lesssim M_{GUT} \quad \text{for } m_e/v_{EW} < h_{ij} < 1$$

Ly α BOSS 1410.7244 .. Planck 1502.01589 (0.23eV w BAO)

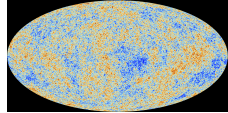
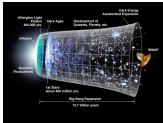
- ▶ B-violation via L-violating Majorana masses M_i
- ▶ CP-violation via Yukawa couplings $\text{Im}[(h^\dagger h)_{ij}] \neq 0$
- ▶ Out-of-equilibrium (inverse) decay $N_i \leftrightarrow \ell \phi^\dagger$ and $N_i \leftrightarrow \ell^c \phi$

$$\begin{aligned} (\Gamma_i/H)|_{T=M_i} &\simeq \frac{\tilde{m}_i/8\pi}{1.66g_* v_{EW}^2/M_{pl}} \quad \text{where } \tilde{m}_i = v_{EW}^2 (h^\dagger h)_{ii}/M_i \\ &\simeq \tilde{m}_i/\text{meV} \quad \sim \mathcal{O}(1-100) \leftrightarrow \begin{array}{l} \text{leptogenesis works well} \\ \text{for observed } \nu \text{ mass scale} \end{array} \end{aligned}$$

Leptogenesis



$m_\nu, \theta_{ij}, \delta_{CP}$



$$\eta = \frac{n_b - n_{\bar{b}}}{n_\gamma} = (6.10 \pm 0.08) \cdot 10^{-10}$$

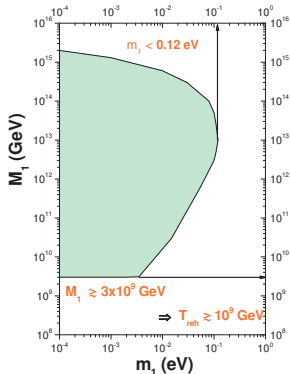
- ▶ Theory of leptogenesis
- ▶ Testing leptogenesis

Vanilla Leptogenesis

$$\frac{dN_{N_i}}{dt} = - \underbrace{\Gamma_{N_i}}_{\text{equilibration rate}} (N_{N_i} - N_{N_i}^{eq})$$

$$\frac{dN_{B-L}}{dt} = \underbrace{\sum_i \Gamma_{CP,i} (N_{N_i} - N_{N_i}^{eq})}_{\text{source term}} - \underbrace{\Gamma_W N_{B-L}}_{\text{washout term}}$$

$$\Gamma_{CP,i} \propto \text{[triangle diagram]} + \text{[circle diagram]}$$



Di Bari 1206.3168 (unflavoured)

- ▶ Lower bound $M_{N_1} \gtrsim 10^9 \text{ GeV}$ (unless $\Delta M_N / M_{N_1} \ll 1$)

Davidson, Ibarra 02; Hamaguchi, Murayama, Yanagida 01

- ▶ Upper bound $m_{\nu_1} \lesssim 0.12 \text{ eV}$ (unflavoured), $\sim 1 \text{ eV}$ (flavoured)

Developments: theory

Systematic approach based on CTP/Kadanoff-Baym

- ▶ Resonant enhancement
- ▶ Flavor effects
- ▶ Thermal corrections

- ▶ Partial spectator equilibration *Garbrecht, Schwaller 1404.2915*
- ▶ Non-relativistic expansion *Bodeker, Wormann 1311.2593*
- ▶ ...

Systematic approach based on CTP/Kadanoff-Baym

Statistical propagator $S_F^{ij}(x, y) = \langle N_i(x)\bar{N}_j(y) - \bar{N}_j(y)N_i(x) \rangle / 2$

Spectral function $S_\rho^{ij}(x, y) = i \langle N_i(x)\bar{N}_j(y) + \bar{N}_j(y)N_i(x) \rangle$

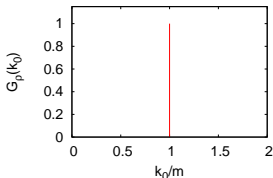
Systematic approach based on CTP/Kadanoff-Baym

Statistical propagator $S_F^{ij}(x, y) = \langle N_i(x)\bar{N}_j(y) - \bar{N}_j(y)N_i(x) \rangle / 2$

Spectral function $S_\rho^{ij}(x, y) = i \langle N_i(x)\bar{N}_j(y) + \bar{N}_j(y)N_i(x) \rangle$

Boltzmann limit

- ▶ on-shell quasi-stable particles



$$S_\rho^{ij}(k) \sim \delta^{ij} \delta(k^2 - m_i^2)$$

- ▶ equilibrium-like
fluctuation-dissipation relation

$$S_F^{ii}(t, k) = \left(\frac{1}{2} - f_k^i(t) \right) S_\rho^{ii}(k)$$

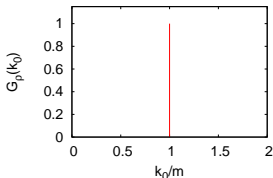
Systematic approach based on CTP/Kadanoff-Baym

Statistical propagator $S_F^{ij}(x, y) = \langle N_i(x)\bar{N}_j(y) - \bar{N}_j(y)N_i(x) \rangle / 2$

Spectral function $S_\rho^{ij}(x, y) = i \langle N_i(x)\bar{N}_j(y) + \bar{N}_j(y)N_i(x) \rangle$

Boltzmann limit

- ▶ on-shell quasi-stable particles



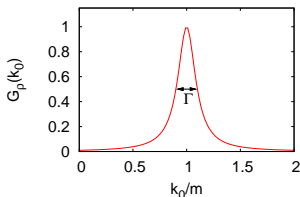
$$S_\rho^{ij}(k) \sim \delta^{ij} \delta(k^2 - m_i^2)$$

- ▶ equilibrium-like fluctuation-dissipation relation

$$S_F^{ii}(t, k) = \left(\frac{1}{2} - f_k^i(t) \right) S_\rho^{ii}(k)$$

CTP/Kadanoff-Baym

- ▶ spectrum with (thermal) width

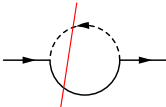
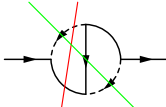
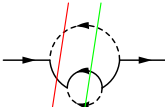


$$S_\rho^{ij}(t, k) \propto \frac{\delta^{ij} 2k_0 \Gamma_i(t)}{(k^2 - m_{th,i}^2(t))^2 + k_0^2 \Gamma_i(t)^2} + \dots$$

- ▶ coherent $N_1 - N_2$ transitions

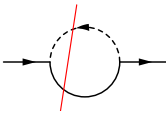
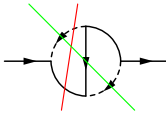
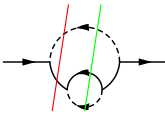
$$S_F^{ij}(t, k) = \begin{pmatrix} S_F^{11} & S_F^{12} \\ S_F^{21} & S_F^{22} \end{pmatrix}$$

CTP/Kadanoff-Baym approach to leptogenesis

			
$N \leftrightarrow l\phi^\dagger$ $N \leftrightarrow l^c\phi$	$ tree ^2$	tree \times vertex-corr.	tree \times wave-corr.
$l\phi^\dagger \leftrightarrow l^c\phi$		$s \times t$	$s \times s, t \times t$

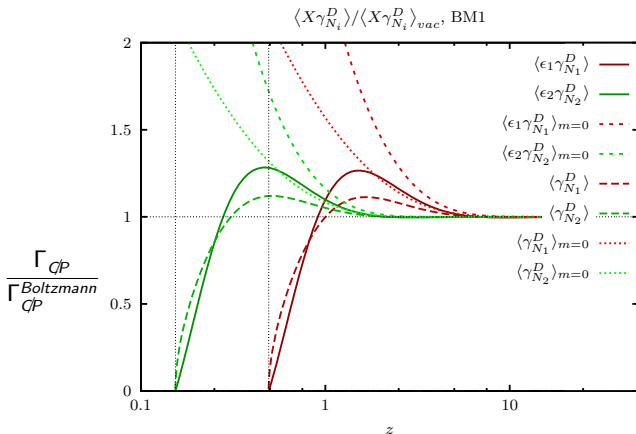
$$\begin{aligned}
 \frac{dn_L}{dt} = & i \int_0^t dt' \int \frac{d^3p}{(2\pi)^3} \text{tr} \left[\Sigma_{l\rho\mathbf{p}}^{\alpha\gamma}(t, t') S_{F\mathbf{p}}^{\gamma\beta}(t', t) - \Sigma_{lF\mathbf{p}}^{\alpha\gamma}(t, t') S_{l\rho\mathbf{p}}^{\gamma\beta}(t', t) \right. \\
 & \left. - S_{l\rho\mathbf{p}}^{\alpha\gamma}(t, t') \Sigma_{lF\mathbf{p}}^{\gamma\beta}(t', t) + S_{lF\mathbf{p}}^{\alpha\gamma}(t, t') \Sigma_{l\rho\mathbf{p}}^{\gamma\beta}(t', t) \right]
 \end{aligned}$$

CTP/Kadanoff-Baym approach to leptogenesis

			
$N \leftrightarrow l\phi^\dagger$ $N \leftrightarrow l^c\phi$	$ tree ^2$	tree \times vertex-corr.	tree \times wave-corr.
$l\phi^\dagger \leftrightarrow l^c\phi$		$s \times t$	$s \times s, t \times t$

- ▶ unified description of CP-violating **decay**, **inverse decay**, **scattering**
- ▶ no need for real intermediate state subtraction
- ▶ flavor-covariant by construction, (partial) flavor coherence
- ▶ no problems with unstable particles in asymptotic in/out states, finite-width (important for resonant enhancement)

Thermal correction to CP asymmetry



Bose enhancement at $M_N \sim T$, suppression for large T due to thermal masses

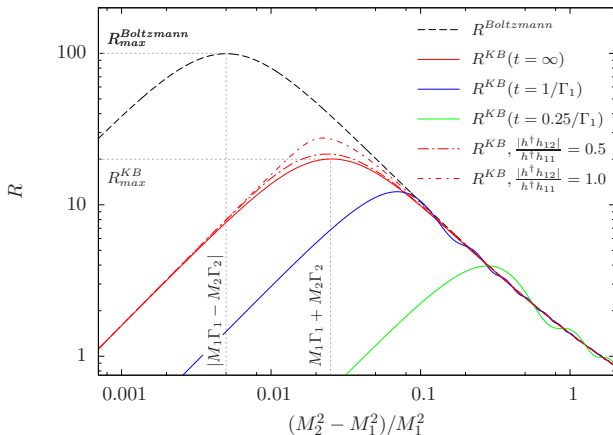
Frossard, MG, Hohenegger, Kartavtsev, Mitrouskas 12

Symmetry quantum statistics vs thermal loop corr. important for soft

leptogenesis ($\epsilon^{vac} = 0 \Rightarrow \epsilon^{th} = 0$)

Garbrecht, Ramsey-Musolf 13

Resonant enhancement



$$R_{max}^{Boltzmann} = M_1 M_2 / (2|\Gamma_1 M_1 - \Gamma_2 M_2|), \quad R_{max}^{KB} \simeq M_1 M_2 / (2(\Gamma_1 M_1 + \Gamma_2 M_2))$$

MG, Kartavtsev, Hohenegger *Annals Phys.* 328 (2013) 26

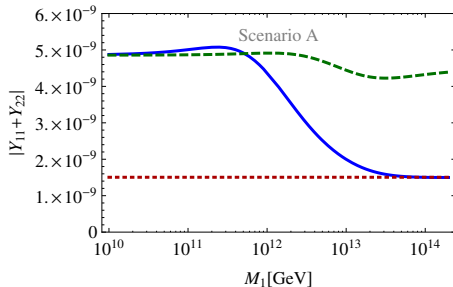
related works: Iso, Shimada, Yamanaka 1312.7680; Dev, Millington, Pilaftsis, Teresi 1403.1003; Iso Shimada 1404.4806; Garbrecht, Gautier, Klaric 1406.4190

Flavoured leptogenesis

- ▶ Effect of charged-lepton Yukawas (flavor effects)

$$\ell_\tau \leftrightarrow \tau_R \phi \quad \text{vs} \quad (h_{1e} \ell_e + h_{1\mu} \ell_\mu + h_{1\tau} \ell_\tau) \phi \leftrightarrow N_1$$

- ▶ Unflavoured: $\Gamma_{LR} \ll H \Rightarrow$ project on flavor that couples to N_1
- ▶ Flavoured: $\Gamma_{LR} \gg H \Rightarrow$ project on flavor that couples to τ (and \perp or e, μ)



- ▶ KB/CTP: matrix-valued, flavor covariant eqs.
- ▶ Interpolate between both regimes, off-diagonal comp. important

Flavoured leptogenesis

Source term $N_1 \rightarrow \ell_a \phi$ ($\mathcal{O}(h^4)$)

$$S_{ab} = \sum_j \left[\underbrace{(h_{a1} h_{1c}^T h_{cj}^* h_{jb} - h_{b1}^* h_{1c}^\dagger h_{cj} h_{ja}^*)}_{\text{Tr} \propto \epsilon_1} \underbrace{S^{LNV}}_{\propto M_1 M_j} + \underbrace{(h_{a1} h_{1c}^\dagger h_{cj} h_{jb} - h_{b1}^* h_{1c}^T h_{cj}^* h_{ja}^*)}_{\text{Tr} = 0} \underbrace{S^{LFV}}_{\propto T^2} \right]$$

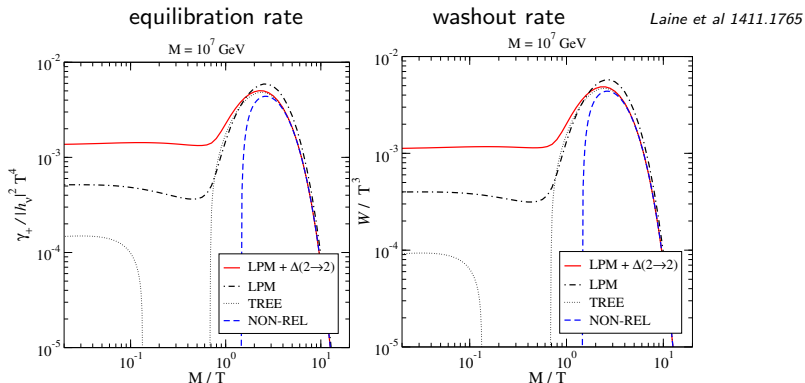
Garbrecht, Drewes 12

Washout $\ell_a \phi \rightarrow N_1$

$$W_{ab} \propto h_{a1} h_{1b}^\dagger$$

- ▶ Potentially large effects for ultrarelativistic N 's, which decay after sphaleron freeze-out *Garbrecht, Drewes 12, cf. also Pilaftsis et. al. 14*
- ▶ Use freedom in h_{ai} to 'hide' asymmetry from washout when $N_{1,2,3}$ have comparable masses \Rightarrow possibility of GeV-scale leptogenesis w/o need for resonant enhancement *Garbrecht, Drewes 14*

Thermal corrections from SM top/gauge int.



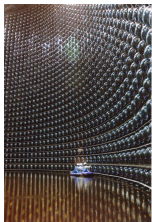
$T \ll M_N$: Important to add virtual and $1 \rightarrow 3$ to scatterings, $g^2 \log(g)$ cancel
 $T \gg M_N$: Rates enhanced by factor ~ 6 (resummation of soft gauge bosons).

cf. also Besak Bodeker 1208.1288; Laine 1307.4909; Garbrecht, Glowna, Herranen 1302.0743; Garbrecht, Glowna, Schwaller 1303.5498

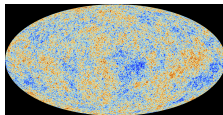
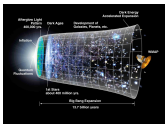
Developments: theory (summary)

- ▶ Systematic approach based on CTP/Kadanoff-Baym
 - ▶ Basic mechanism confirmed, quantitative corrections of $\mathcal{O}(1)$
 - ▶ Some models are not viable ($\epsilon^{vac} = 0$)
- ▶ Resonant enhancement:
 - ▶ saturation; decays vs oscillations
 - ▶ no extra enhancement in double degenerate limit
 $M_1 \rightarrow M_2, \Gamma_1 \rightarrow \Gamma_2$
- ▶ Flavor effects:
 - ▶ transition between flavored/unflavored $M_N = 10^{12\pm 1}, 10^{9\pm 1} \text{ GeV}$
 - ▶ finite-T enhancement of flavored asymmetries $\propto (T/M_N)^2$
- ▶ Thermal corrections: no thermal blocking
- ▶ ...

Leptogenesis



$m_\nu, \theta_{ij}, \delta_{CP}$



$$\eta = \frac{n_b - n_{\bar{b}}}{n_\gamma} = (6.10 \pm 0.08) \cdot 10^{-10}$$

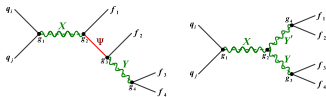
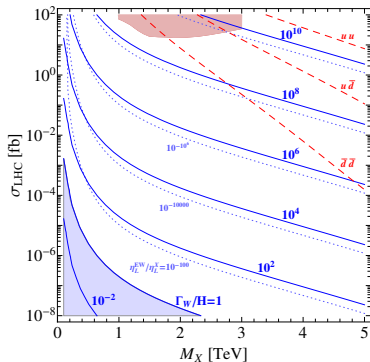
- ▶ Theory of leptogenesis
- ▶ Testing leptogenesis

Falsifying high-scale Baryogenesis

Prediction: no LNV beyond Weinberg operator $LLHH/\Lambda$. Observation of additional LNV would imply washout of lepton/baryon asymmetry at $T \gtrsim T_{ew}$.

Falsifying high-scale Baryogenesis

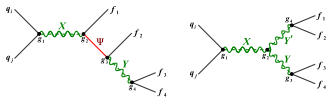
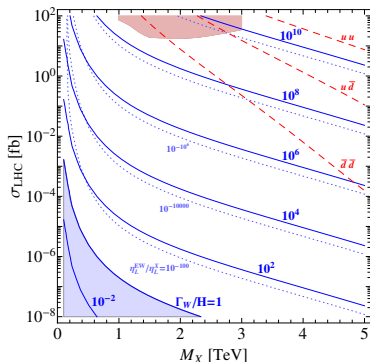
Prediction: no LNV beyond Weinberg operator $LLHH/\Lambda$. Observation of additional LNV would imply washout of lepton/baryon asymmetry at $T \gtrsim T_{ew}$.
 E.g. at LHC ($pp \rightarrow X \rightarrow \ell^\pm \ell^\pm jj$) *Deppisch, Harz 1408.5351, 1312.4447* → talk by J. Harz



$$\log_{10} \frac{\Gamma_W}{H} \gtrsim 6.9 + 0.6 \left(\frac{M_X}{\text{TeV}} - 1 \right) + \log_{10} \frac{\sigma_{\text{LHC}}}{\text{fb}}$$

Falsifying high-scale Baryogenesis

Prediction: no LNV beyond Weinberg operator $LLHH/\Lambda$. Observation of additional LNV would imply washout of lepton/baryon asymmetry at $T \gtrsim T_{ew}$.
 E.g. at LHC ($pp \rightarrow X \rightarrow \ell^\pm \ell^\pm jj$) *Deppisch, Harz 1408.5351, 1312.4447* → talk by *J. Harz*

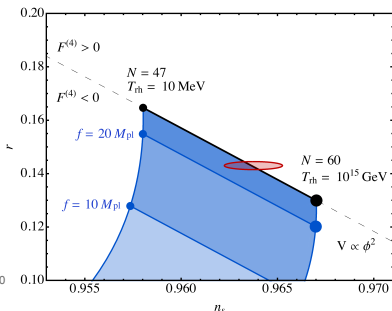
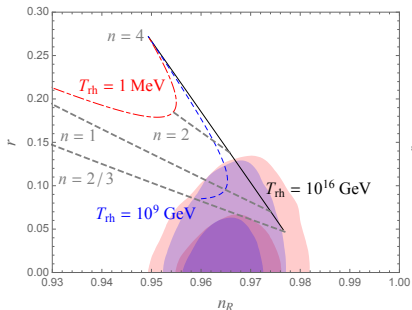


$$\log_{10} \frac{\Gamma_W}{H} \gtrsim 6.9 + 0.6 \left(\frac{M_X}{\text{TeV}} - 1 \right) + \log_{10} \frac{\sigma_{\text{LHC}}}{\text{fb}}$$

Similar argument for $0\nu\beta\beta$ *Deppisch, Harz, Huang, Hirsch, Päs 1503.04825* and W_R *Hambye+08*; does not apply to EW-scale scenarios ($W_R > 3 \text{ TeV}$ *Dev, Mohapatra, 1408.2820*);
 loop-hole τ -flavor if no LFV observed

High-scale baryogenesis and inflation

- Preference for larger n_s for higher T_R ($n_s - 1 \propto -1/N$) for generic slow-roll models, dependence on EoS during reheating and inflaton potential



Creminelli, Nacir, Simonovic, Trevisan, Zaldarriaga 14; Dai, Kamionkowski, Wang 14; Gong, Leung, Pi 15

- Non-thermal leptogenesis (via Higgs-oscillations after inflation *Kusenko+ 15*;
 $U(1)_{B-L}$ phase transition *Buchmüller, Domcke, Kamada, Schmitz 13*)

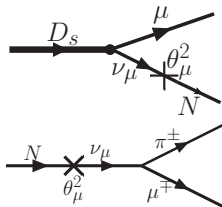
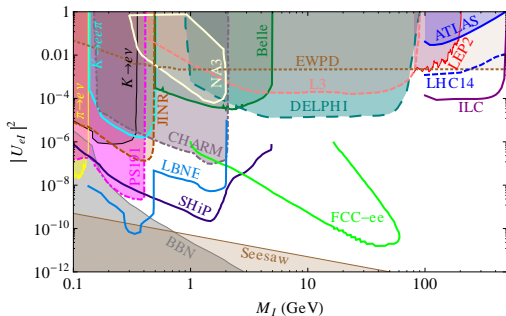
EW-scale Leptogenesis (resonant)

- ▶ Leptogenesis for $M_N < 10^9 \text{ GeV}$ with resonant enhancement
- ▶ Generically tiny Yukawa couplings \Rightarrow difficult
- ▶ RG-induced splitting is not enough *Dev, Milington, Pilaftsis, Teresi 1504.07640*
- ▶ Testable in specific regions of parameter space (hide asymmetry in L_τ , stronger coupling to e, μ compatible with signal in LFV and LNV)
Dev, Milington, Pilaftsis, Teresi 1404.1003

\rightarrow talk by B. Dev

Low-scale Leptogenesis

- ▶ Leptogenesis with GeV-scale RH neutrinos via resonant or thermal enhancement
- ▶ Search for active-sterile mixing $\theta^2 = U^2 \sim v^2 h^2 / M_N^2$
- ▶ Beam-dump (SHiP) $D \rightarrow \ell N$, $N \rightarrow \pi^\pm \ell^\mp$
- ▶ $Z \rightarrow N \nu$ with displaced vertex ($M_N < 45 \text{ GeV}$) or off-shell

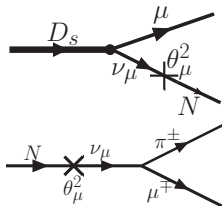
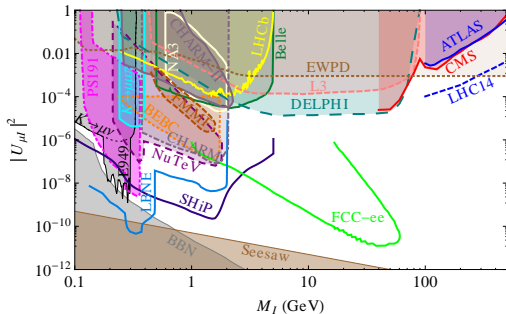


Seesaw: $U^2 \gtrsim \mathcal{O}(m_\nu / M_N)$

Leptogenesis: $U^2 \lesssim \begin{cases} 10^{-6..8} & N = 2, \text{ resonant enh. } (\nu\text{MSM}) \\ 10^{-4..5} & N = 3, \text{ thermal enh.} \end{cases}$

Low-scale Leptogenesis

- ▶ Leptogenesis with GeV-scale RH neutrinos via resonant or thermal enhancement
- ▶ Search for active-sterile mixing $\theta^2 = U^2 \sim v^2 h^2 / M_N^2$
- ▶ Beam-dump (SHiP) $D \rightarrow \ell N$, $N \rightarrow \pi^\pm \ell^\mp$
- ▶ $Z \rightarrow N \nu$ with displaced vertex ($M_N < 45\text{GeV}$) or off-shell

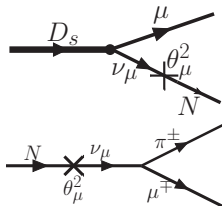
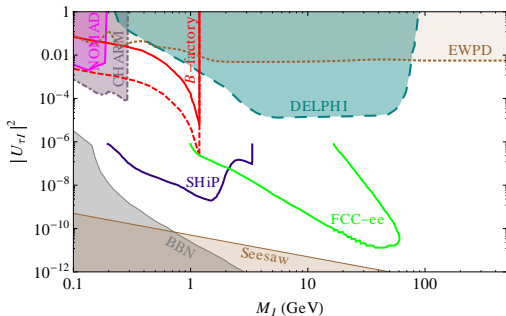


Seesaw: $U^2 \gtrsim \mathcal{O}(m_\nu / M_N)$

Leptogenesis: $U^2 \lesssim \begin{cases} 10^{-6..8} & N = 2, \text{ resonant enh. } (\nu\text{MSM}) \\ 10^{-4..5} & N = 3, \text{ thermal enh.} \end{cases}$

Low-scale Leptogenesis

- ▶ Leptogenesis with GeV-scale RH neutrinos via resonant or thermal enhancement
- ▶ Search for active-sterile mixing $\theta^2 = U^2 \sim v^2 h^2 / M_N^2$
- ▶ Beam-dump (SHiP) $D \rightarrow \ell N$, $N \rightarrow \pi^\pm \ell^\mp$
- ▶ $Z \rightarrow N \nu$ with displaced vertex ($M_N < 45 \text{ GeV}$) or off-shell



Seesaw: $U^2 \gtrsim \mathcal{O}(m_\nu / M_N)$

Leptogenesis: $U^2 \lesssim \begin{cases} 10^{-6..8} & N = 2, \text{ resonant enh. } (\nu\text{MSM}) \\ 10^{-4..5} & N = 3, \text{ thermal enh.} \end{cases}$

Conclusion

- ▶ Lots of progress in **theory of leptogenesis**, both on qualitative and quantitative level
- ▶ So far all measurements are **consistent with expectations for typical leptogenesis scenarios** (no low-energy LNV beyond Weinberg, sub-eV neutrino masses, large reheating temp.)
- ▶ Observation of $0\nu\beta\beta$ compatible with $\sum m_\nu$ [or primordial tensors r or $n_s \gtrsim 0.965$] would **strengthen the generic leptogenesis scenario**; observation of LNV at LHC or non-standard $0\nu\beta\beta$ would **falsify** it
- ▶ **Specific scenarios can be tested** (low-scale leptogenesis via resonant (SHiP) or thermal enhancement (LHCb))

Hierarchical limit

$$\begin{aligned}
 \partial_t n_L^{KB} &= 16\pi\epsilon_1 \int_{\mathbf{p}, \mathbf{q}, \mathbf{q}', \mathbf{k}, \mathbf{k}'} \frac{\mathbf{k} \cdot \mathbf{k}'}{M_1} (2\pi)^3 \delta(\mathbf{p} - \mathbf{k} - \mathbf{q}) (2\pi)^3 \delta(\mathbf{p} - \mathbf{k}' - \mathbf{q}') \\
 &\quad \times \operatorname{Re} \frac{\Gamma_{\ell\phi}}{(\omega_p - k - q)^2 + \Gamma_{\ell\phi}^2/4} \frac{\Gamma_{\ell\phi}}{(\omega_p - k' - q')^2 + \Gamma_{\ell\phi}^2/4} \\
 &\quad \times (f_{N_1}(\mathbf{p}) - f_{N_1}^{\text{eq}}(\mathbf{p})) (1 - f_\ell(k) + f_\phi(q)) (1 - f_\ell(k') + f_\phi(q')) \\
 \partial_t n_L^{\text{Boltzmann}} &= 16\pi\epsilon_1 \int_{\mathbf{p}, \mathbf{q}, \mathbf{q}', \mathbf{k}, \mathbf{k}'} \frac{\mathbf{k} \cdot \mathbf{k}'}{M_1} (2\pi)^4 \delta(\mathbf{p} - \mathbf{k} - \mathbf{q}) (2\pi)^4 \delta(\mathbf{p} - \mathbf{k}' - \mathbf{q}') \\
 &\quad \times 2\pi\delta(\omega_p - k_0 - q_0) 2\pi\delta(\omega_p - k'_0 - q'_0) \\
 &\quad \times (f_{N_1}(\mathbf{p}) - f_{N_1}^{\text{eq}}(\mathbf{p})) (1 - f_\ell(k) + f_\phi(q))
 \end{aligned}$$

The 'width' of lepton and Higgs $\Gamma_{\ell\phi} = \Gamma_\ell + \Gamma_\phi$ leads to a replacement of the on-shell delta function in the Boltzmann equations by a Breit-Wigner curve

MG, Kartavtsev, Hohenegger, Lindner 09; Beneke, Garbrecht, Herranen, Schwaller 10; Anisimov, Buchmüller, Drewes, Mendizabal 11

Hierarchical limit

$$\begin{aligned} \partial_t n_L^{KB} &= 16\pi\epsilon_1 \int_{\mathbf{p}, \mathbf{q}, \mathbf{q}', \mathbf{k}, \mathbf{k}'} \frac{\mathbf{k} \cdot \mathbf{k}'}{M_1} (2\pi)^3 \delta(\mathbf{p} - \mathbf{k} - \mathbf{q}) (2\pi)^3 \delta(\mathbf{p} - \mathbf{k}' - \mathbf{q}') \\ &\times \operatorname{Re} \frac{\Gamma_{\ell\phi}}{(\omega_p - k - q)^2 + \Gamma_{\ell\phi}^2/4} \frac{\Gamma_{\ell\phi}}{(\omega_p - k' - q')^2 + \Gamma_{\ell\phi}^2/4} \\ &\times (f_{N_1}(\mathbf{p}) - f_{N_1}^{eq}(\mathbf{p})) (1 - f_\ell(k) + f_\phi(q)) (1 - f_\ell(k') + f_\phi(q')) \end{aligned}$$

$$\begin{aligned} \partial_t n_L^{\text{Boltzmann}} &= 16\pi\epsilon_1 \int_{\mathbf{p}, \mathbf{q}, \mathbf{q}', \mathbf{k}, \mathbf{k}'} \frac{\mathbf{k} \cdot \mathbf{k}'}{M_1} (2\pi)^4 \delta(\mathbf{p} - \mathbf{k} - \mathbf{q}) (2\pi)^4 \delta(\mathbf{p} - \mathbf{k}' - \mathbf{q}') \\ &\times 2\pi\delta(\omega_p - k_0 - q_0) 2\pi\delta(\omega_p - k'_0 - q'_0) \\ &\times (f_{N_1}(\mathbf{p}) - f_{N_1}^{eq}(\mathbf{p})) (1 - f_\ell(k) + f_\phi(q)) \end{aligned}$$

Thermal correction to loop in CP asymmetry, comparable to quantum statistical corrections

MG, Kartavtsev, Hohenegger, Lindner 09; Beneke, Garbrecht, Herranen, Schwaller 10; Anisimov, Buchmüller, Drewes, Mendizabal 11

Degenerate case

Analytical result for $\text{Re}(h^\dagger h)_{12} \ll (h^\dagger h)_{ii}$ (mass basis \sim 'int. basis')

$$\begin{aligned}
 n_L(t) &= \frac{\text{Im}[(h^\dagger h)_{12}^2]}{8\pi} \frac{M_1 M_2 (M_2^2 - M_1^2)}{(M_2^2 - M_1^2)^2 + (M_1 \Gamma_1 - M_2 \Gamma_2)^2} \int \frac{d^3 p d^3 q d^3 k}{(2\pi)^9 \omega_p 2q 2k} 4k \cdot p \\
 &\times (2\pi)^3 \delta(\mathbf{p} - \mathbf{k} - \mathbf{q}) \frac{\Gamma_{\ell\phi}}{(\omega_p - k - q)^2 + \Gamma_{\ell\phi}^2/4} (1 + f_\phi - f_\ell) f_{FD}(\omega_p) \\
 &\times \left[\sum_{l=1,2} \frac{1 - e^{-\Gamma_{pl} t}}{\Gamma_{pl}} - 4 \text{Re} \left(\frac{1 - e^{-i(\omega_{p1} - \omega_{p2})t} e^{-(\Gamma_{p1} + \Gamma_{p2})t/2}}{\Gamma_{p1} + \Gamma_{p2} + 2i(\omega_{p1} - \omega_{p2})} \right) \right]
 \end{aligned}$$

$$\begin{aligned}
 n_L^{\text{Boltzmann}}(t) &= \frac{\text{Im}[(h^\dagger h)_{12}^2]}{8\pi} \frac{M_1 M_2 (M_2^2 - M_1^2)}{(M_2^2 - M_1^2)^2 + (M_1 \Gamma_1 - M_2 \Gamma_2)^2} \int \frac{d^3 p d^3 q d^3 k}{(2\pi)^9 \omega_p 2q 2k} 4k \cdot p \\
 &\times (2\pi)^3 \delta(\mathbf{p} - \mathbf{k} - \mathbf{q}) 2\pi \delta(\omega_p - k - q) (1 + f_\phi - f_\ell) f_{FD}(\omega_p) \\
 &\times \left[\sum_{l=1,2} \frac{1 - e^{-\Gamma_{pl} t}}{\Gamma_{pl}} \right]
 \end{aligned}$$

Degenerate case

Analytical result for $\text{Re}(h^\dagger h)_{12} \ll (h^\dagger h)_{ii}$ (mass basis \sim 'int. basis')

$$\begin{aligned}
 n_L(t) = & \frac{\text{Im}[(h^\dagger h)_{12}^2]}{8\pi} \frac{M_1 M_2 (M_2^2 - M_1^2)}{(M_2^2 - M_1^2)^2 + (M_1 \Gamma_1 - M_2 \Gamma_2)^2} \int \frac{d^3 p d^3 q d^3 k}{(2\pi)^9 \omega_p 2q 2k} 4k \cdot p \\
 & \times (2\pi)^3 \delta(\mathbf{p} - \mathbf{k} - \mathbf{q}) \frac{\Gamma_{\ell\phi}}{(\omega_p - k - q)^2 + \Gamma_{\ell\phi}^2/4} (1 + f_\phi - f_\ell) f_{FD}(\omega_p) \\
 & \times \left[\sum_{l=1,2} \frac{1 - e^{-\Gamma_{pl}t}}{\Gamma_{pl}} - 4 \text{Re} \left(\frac{1 - e^{-i(\omega_{p1} - \omega_{p2})t} e^{-(\Gamma_{p1} + \Gamma_{p2})t/2}}{\Gamma_{p1} + \Gamma_{p2} + 2i(\omega_{p1} - \omega_{p2})} \right) \right]
 \end{aligned}$$

$$\begin{aligned}
 n_L^{\text{Boltzmann}}(t) = & \frac{\text{Im}[(h^\dagger h)_{12}^2]}{8\pi} \frac{M_1 M_2 (M_2^2 - M_1^2)}{(M_2^2 - M_1^2)^2 + (M_1 \Gamma_1 - M_2 \Gamma_2)^2} \int \frac{d^3 p d^3 q d^3 k}{(2\pi)^9 \omega_p 2q 2k} 4k \cdot p \\
 & \times (2\pi)^3 \delta(\mathbf{p} - \mathbf{k} - \mathbf{q}) 2\pi \delta(\omega_p - k - q) (1 + f_\phi - f_\ell) f_{FD}(\omega_p) \\
 & \times \left[\sum_{l=1,2} \frac{1 - e^{-\Gamma_{pl}t}}{\Gamma_{pl}} \right]
 \end{aligned}$$

► Regulator $M_1 \Gamma_1 - M_2 \Gamma_2$ is confirmed

Degenerate case

Analytical result for $\text{Re}(h^\dagger h)_{12} \ll (h^\dagger h)_{ii}$ (mass basis \sim 'int. basis')

$$n_L(t) = \frac{\text{Im}[(h^\dagger h)_{12}^2]}{8\pi} \frac{M_1 M_2 (M_2^2 - M_1^2)}{(M_2^2 - M_1^2)^2 + (M_1 \Gamma_1 - M_2 \Gamma_2)^2} \int \frac{d^3 p d^3 q d^3 k}{(2\pi)^9 \omega_p 2q 2k} 4k \cdot p$$

$$\times (2\pi)^3 \delta(\mathbf{p} - \mathbf{k} - \mathbf{q}) \frac{\Gamma_{\ell\phi}}{(\omega_p - k - q)^2 + \Gamma_{\ell\phi}^2/4} (1 + f_\phi - f_\ell) f_{FD}(\omega_p)$$

$$\times \left[\sum_{l=1,2} \frac{1 - e^{-\Gamma_{pl}t}}{\Gamma_{pl}} - 4 \text{Re} \left(\frac{1 - e^{-i(\omega_{p1} - \omega_{p2})t} e^{-(\Gamma_{p1} + \Gamma_{p2})t/2}}{\Gamma_{p1} + \Gamma_{p2} + 2i(\omega_{p1} - \omega_{p2})} \right) \right]$$

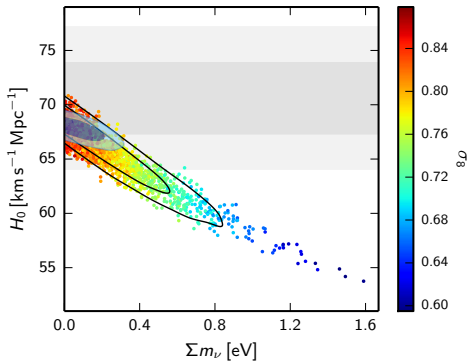
$$n_L^{\text{Boltzmann}}(t) = \frac{\text{Im}[(h^\dagger h)_{12}^2]}{8\pi} \frac{M_1 M_2 (M_2^2 - M_1^2)}{(M_2^2 - M_1^2)^2 + (M_1 \Gamma_1 - M_2 \Gamma_2)^2} \int \frac{d^3 p d^3 q d^3 k}{(2\pi)^9 \omega_p 2q 2k} 4k \cdot p$$

$$\times (2\pi)^3 \delta(\mathbf{p} - \mathbf{k} - \mathbf{q}) 2\pi \delta(\omega_p - k - q) (1 + f_\phi - f_\ell) f_{FD}(\omega_p)$$

$$\times \left[\sum_{l=1,2} \frac{1 - e^{-\Gamma_{pl}t}}{\Gamma_{pl}} \right]$$

- ▶ Regulator $M_1 \Gamma_1 - M_2 \Gamma_2$ is confirmed
- ▶ Additional oscillating contribution due to **coherent** $N_1 - N_2$ transitions

Neutrino mass



Planck 1502.01589