Higgs width and coupling constraints from the high-mass $H \to VV$ process

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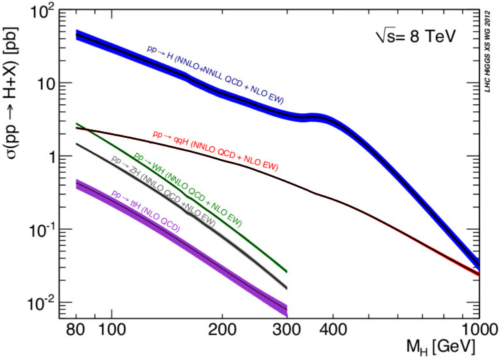
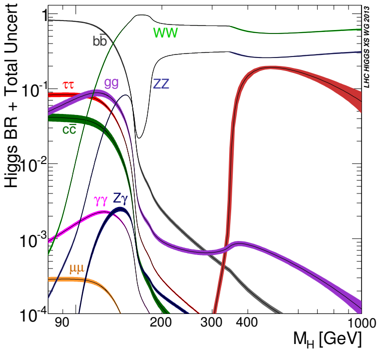
Outline

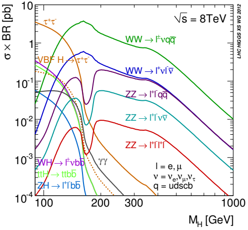
- $H \to ZZ, WW$ in ggF & VBF: sizeable off-shell Higgs signal contribution with large signal-backg. interference
- $H \to WW/ZZ \to \ell\bar{\nu}_\ell\bar{\ell}\nu_\ell$ and $ZZ \to 4\ell$ interference in ggF
- Interference effects for semileptonic decay modes
- Interference for $pp \to H \to ZZ$ + jet
- Precision predictions for $gg\ (\to H) \to VV$ interference/bkg.
- Higgs width measurement in a nutshell
- Off-shell Higgs boson signal strength & novel Higgs width bound
- BSM/EFT: exploiting the off-shell $H \to VV$ region
- BSM: model dependence of the off-shell Higgs width bound
- Higgs width constraints from $gg \to H \to \gamma\gamma$
- BSM: heavy Higgs-light Higgs-bkg. interference effects
- Higgs width/coupling constraints including LEP EW PO
- High-mass $H \to VV$ signal at a linear collider
- Summary

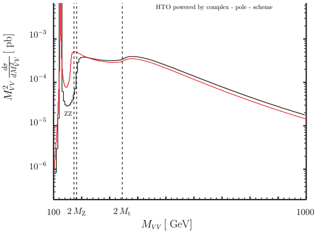
SM Higgs boson production and decay at the Large Hadron Collider (LHC)

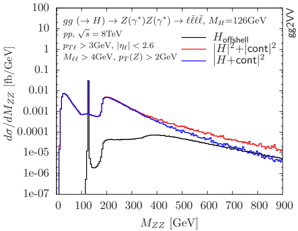




$gg \to H \to ZZ, WW$: sizeable off-shell Higgs signal with large signal-background interference



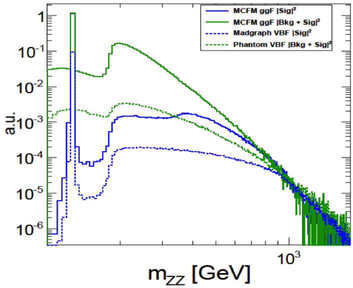


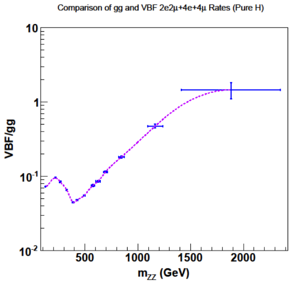


- $gg \to H \to VV \to 4\ell$ and $2\ell 2\nu$ signal-background interference very well studied at LO: Glover, van der Bij (1989); Kao, Dicus (1991); Binoth, Ciccolini, NK, Krämer (2006) (gg2WW); Campbell, Ellis, Williams (2011) (MCFM); NK (2012) (gg2VV); NK, Passarino (2012); Campanario, Li, Rauch, Spira (2012); Bonvini, Caola, Forte, Melnikov, Ridolfi (2013); Caola, Melnikov (2013); NK (2013) (gg2VV); Campbell, Ellis, Williams (2013) (MCFM); Campbell, Ellis, Williams (2014) (MCFM); Campbell, Ellis, Furlan, Röntsch (2014); related interference effects: Bredenstein, Denner, Dittmaier, Weber (2006) (PROPHECY4f); YR3: Denner, Dittmaier, Mück (2013) and Anderson, Bolognesi, Caola, Gao, Gritsan, Martin, Melnikov, Schulze, Tran, Whitbeck, Zhou (2013); Chen, Cheng, Gainer, Korytov, Matchev, Milenovic, Mitselmakher, Park, Rinkevicius, Snowball (2013); Chen, Vega-Morales (2013)
- tools for ggF: MCFM-6.8, gg2VV-3.1.7 ($gg \to VV$ parton-calculators and LO event generators), MadGraph5, Sherpa+OpenLoops (allows for merging of $gg \to VV + \{0, 1\}$ jets)
- loop technology closing in on NLO calculation (see below) Karlsruhe, Zurich, FNAL-RWTH, ...
- gluon-fusion Higgs production and semileptonic decay: Dobrescu, Lykken (2010); Lykken, Martin, Winter (2012); Kao, Sayre (2012); ATLAS arXiv:1206.2443; ATLAS arXiv:1206.6074; CMS PAS HIG-13-008

Sizeable off-shell Higgs signal in vector boson fusion

- similar effect in VBF $H \to VV$ (NK, Passarino): $\mathcal{O}(10\%)$ of Higgs signal is off-shell
  note: no exp. sensitivity to off-shell $H \to VV$ tail in VH and $t\bar{t}H$ channels (see $\sigma_{\text{prod}}(M_H)$)
- total off-shell Higgs signal has $\sim 10\%$ VBF contribution





Covarelli, Anderson, Sarica

figures taken from Covarelli's talk at LHC HXSWG workshop (12 Jun 2014)

- tools for VBF: MadGraph5 Alwall et al., Phantom Ballestrero et al., VBFNLO Baglio et al., Sherpa+OpenLoops Cascioli et al.

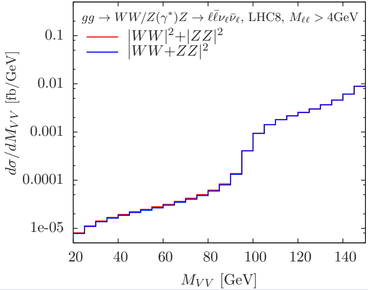
$gg \to H \to WW/ZZ \to \ell\bar{\nu}_\ell\bar{\ell}\nu_\ell$ interference (gg2VV)

Integrated cross sections (SM Higgs)

$gg\ (\to H) \to VV \to \ell\bar{\nu}_\ell\bar{\ell}\nu_\ell$, σ [fb], pp, $\sqrt{s} = 8$ TeV, $M_H = 126$ GeV, min. cuts, $\mu_R = \mu_F = M_{\ell\bar{\nu}_\ell\bar{\ell}\nu_\ell}/2$				interference	
VV	H	cont	$\lvert H$+cont$\rvert^2$	R_1=(S+B+I)/(S+B)	R_2=(S+I)/S
WW	17.318(4)	16.925(4)	32.803(8)	0.9580(3)	0.9169(6)
ZZ	0.8822(2)	2.1553(6)	2.872(1)	0.9455(4)	0.813(2)
WW/ZZ	17.402(3)	19.084(4)	34.884(7)	0.9561(3)	0.9079(5)
R_3	0.9562(3)	1.0002(3)	0.9778(3)	$\sigma(\lvert WW+ZZ\rvert^2)/\sigma(\lvert WW\rvert^2+\lvert ZZ\rvert^2)$	
R_4	0.9540(3)	1.0002(4)	0.9759(4)	$(\sigma(\lvert WW\rvert^2)+I_{WW/ZZ})/\sigma(\lvert WW\rvert^2)$	
R_6	0.05094(2)	0.12735(5)	0.08756(4)	$\sigma(\lvert ZZ\rvert^2)/\sigma(\lvert WW\rvert^2)$	

minimal cuts: WW/ZZ interference: Higgs signal: ≈5%, gg continuum: negligible NK arXiv:1310.7011

Differential cross sections



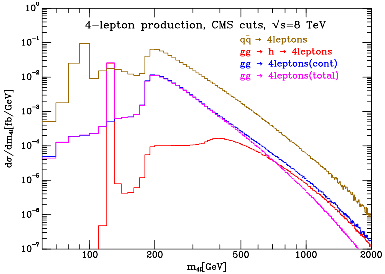
$gg \to WW/ZZ$ continuum: M_{VV} distribution

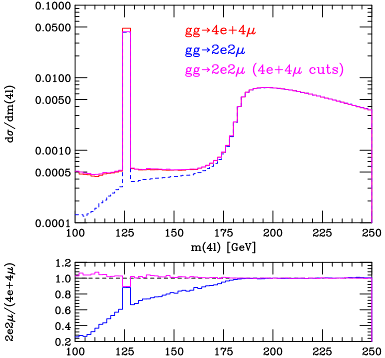
$M_{VV} > 95$ GeV: WW/ZZ interference negligible

$M_{VV} < 95$ GeV: WW/ZZ interference of ≈ 5%

see also: $H \to WW/ZZ \to \ell\bar{\nu}_\ell\bar{\ell}\nu_\ell$ interference at LO & NLO Mück, Bredenstein, Denner, Dittmaier, Weber YR3 arXiv:1307.1347, Sec. 2.2

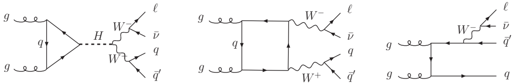
$gg \to H \to ZZ \to 4\ell$ interference (MCFM)





Campbell, R.K. Ellis, Williams   figures taken from arXiv:1408.1723

Interference for semileptonic H decay modes in ggF



$$\mathcal{M} = \mathcal{M}_{signal}\text{ (LO) } + \mathcal{M}_{background} = \mathcal{M}_{signal} + \mathcal{M}_{loop} + \mathcal{M}_{tree}$$

Notation for amplitude contributions to cross sections:

$$S \sim |\mathcal{M}_{signal}|^2$$
$$I_{tree} \sim 2\,\mathrm{Re}(\mathcal{M}^*_{signal}\,\mathcal{M}_{tree})$$
$$I_{loop} \sim 2\,\mathrm{Re}(\mathcal{M}^*_{signal}\,\mathcal{M}_{loop})$$
$$I_{full} \sim 2\,\mathrm{Re}(\mathcal{M}^*_{signal}\,\mathcal{M}_{background})$$

$\mathcal{M}_{loop}$ contains all quark loop graphs. (NLO EW corrections to I_{tree} not included.)

relative measure for interf. with bkg. i:

$$R_i = \frac{\sigma(|\mathcal{M}_{signal}|^2 + 2\,\mathrm{Re}(\mathcal{M}^*_{signal}\mathcal{M}_i))}{\sigma(|\mathcal{M}_{signal}|^2)}$$

Interference for semileptonic H decay modes in ggF

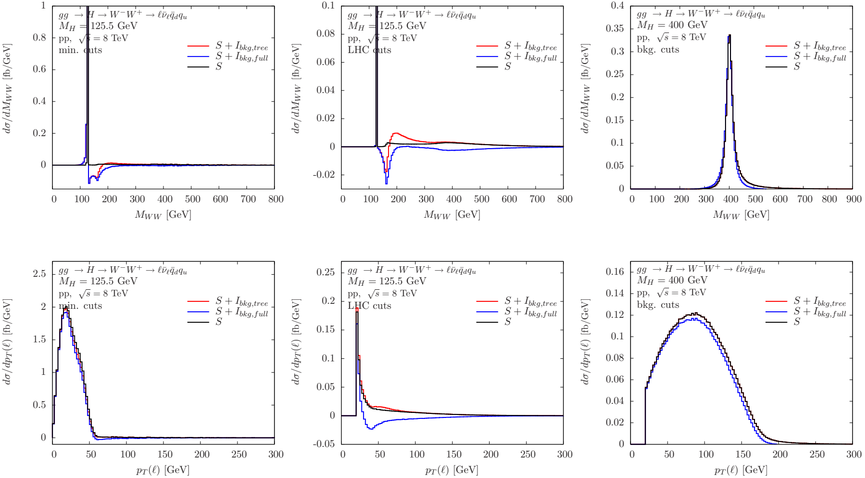
$gg \to H \to ZZ \to \ell\bar{\ell}q_u\bar{q}_u$ σ [fb], pp, $\sqrt{s} = 8$ TeV		interference			ratio		
cuts	S	I_{tree}	I_{loop}	I_{full}	R_{tree}	R_{loop}	R_{full}
min.	1.96(1)	–0.190(4)	–0.343(3)	–0.541(5)	0.903(7)	0.825(7)	0.724(7)
LHC	0.1166(6)	0.017(2)	–0.194(2)	–0.176(6)	1.15(2)	–0.67(2)	–0.51(5)
bkg.	1.342(7)	–0.0012(2)	–0.0882(9)	–0.0892(9)	0.999(7)	0.934(7)	0.934(7)

$gg \to H \to ZZ \to \ell\bar{\ell}q_d\bar{q}_d$ σ [fb], pp, $\sqrt{s} = 8$ TeV		interference			ratio		
cuts	S	I_{tree}	I_{loop}	I_{full}	R_{tree}	R_{loop}	R_{full}
min.	2.51(2)	–0.248(3)	–0.439(6)	–0.680(7)	0.901(7)	0.825(7)	0.729(7)
LHC	0.1497(8)	0.0223(6)	–0.245(5)	–0.227(3)	1.149(9)	–0.64(3)	–0.52(2)
bkg.	1.720(9)	–0.00130(5)	–0.113(1)	–0.114(1)	0.999(7)	0.934(7)	0.934(7)

higher-order background contributions can induce leading interference effects

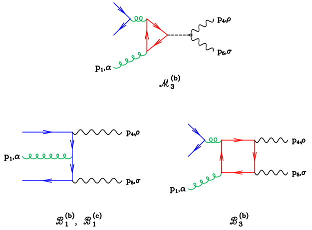
NK, C. O'Brien, E. Vryonidou arXiv:1506.01694 (gg2VV and MG5_aMC@NLO)

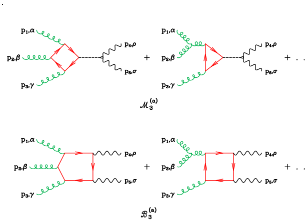
Interference for semileptonic H decay modes in ggF

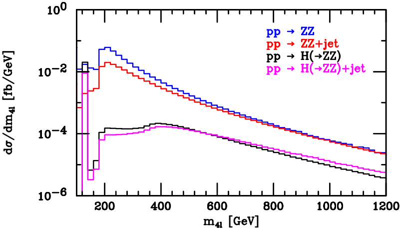


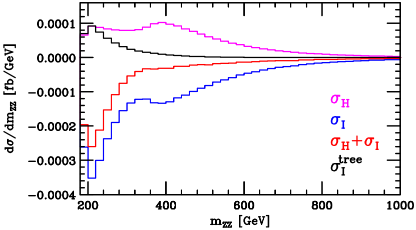
NK, C. O'Brien, E. Vryonidou arXiv:1506.01694

Interference for $pp \to H \to ZZ$ + jet









off-shell Higgs cross sections for ZZ and ZZ+jet comparable ($p_{Tj} > 30$ GeV)

Campbell, R.K. Ellis, Furlan, Röntsch figures taken from arXiv:1409.1897

Z bosons treated in zero-width approximation (validated for ZZ final state: excellent for $m_{4l} > 300$ GeV)

Precision predictions for $gg\ (\to H) \to VV$ signal-background interference

Signal: $gg \to H$ cross section at NLO QCD with finite t and b mass effects (important for off-shell Higgs with $M_{VV} \gtrsim 2M_t$: 5–10% correction) (scale uncertainty: 10–15%) Djouadi, Spira, Zerwas, Graudenz (1991-1995); N^{3}LO in soft expansion with $M_t \to \infty$ (scale uncertainty $\approx 3\%$) C. Anastasiou, C. Duhr, F. Dulat, F. Herzog, B. Mistlberger arXiv:1503.06056; NLO EW corrections important for off-shell Higgs (8% at $M_{VV} \sim 500$ GeV) A. Bredenstein, A. Denner, S. Dittmaier, M. Weber arXiv:hep-ph/0604011 (also arXiv:1111.6395)

Background: $pp \to ZZ$ and $pp \to WW$ at NNLO QCD with massless quarks (scale uncertainty $\approx 3\%$), F. Cascioli, T. Gehrmann, M. Grazzini, S. Kallweit, P. Maierhofer, A. von Manteuffel, S. Pozzorini, D. Rathlev, L. Tancredi, E. Weihs arXiv:1405.2219 and T. Gehrmann, M. Grazzini, S. Kallweit, P. Maierhofer, A. von Manteuffel, S. Pozzorini, D. Rathlev, L. Tancredi arXiv:1408.5243

$gg \to VV$ enters $pp \to VV$ at NNLO QCD $\to$ LO (loop-induced) with $\sim$ 20–25% scale uncertainty, unknown NLO K-factor, but expected to be similar to signal, i.e. ~ 1.6

11–17% (9–12%) NNLO correction to $pp \to ZZ$ (WW) for $\sqrt{s} = 7$–14 TeV

$gg \to VV$ contributes to full NNLO correction with 60% (35%) for $pp \to ZZ$ (WW)

$\to$ NLO $gg \to VV$ correction is of similar size or larger than residual $pp \to VV$ scale uncertainty $\Rightarrow$ calculation is important and by a similar argument the calculation of the NLO correction to signal-background interference

Precision predictions for $gg\ (\to H) \to VV$ signal-background interference

Work towards $gg\ (\to H) \to VV$ signal-background interference and $gg \to VV$ continuum background beyond leading order:

M. Bonvini, F. Caola, S. Forte, K. Melnikov, G. Ridolfi arXiv:1304.3053:

NLO and NNLO calculation for $gg\ (\to H) \to WW \to \ell\nu\ell\nu$ interference with $M_H = 600$ GeV in soft-gluon approximation (very good accuracy for inclusive signal cross section)
$\to$ interference K-factors are generally very similar to signal K-factors (also for kinematic distributions)

C. Li, H. Li, D. Shao, J. Wang arXiv:1504.02388:

Soft gluon resummation to all orders for $gg\ (\to H) \to ZZ \to \ell\ell\ell'\ell'$ interference, 100 GeV $< M_{ZZ} < 1000$ GeV, effects signal like

Technical bottleneck for unapproximated $gg \to VV$ calc. at NLO: two-loop virtual corrections

Two-loop $gg \to VV \to 4$ leptons amplitudes with massless quarks calculated by two groups:
F. Caola, J. Henn, K. Melnikov, A. Smirnov, V. Smirnov arXiv:1503.08759
A. v. Manteuffel, L. Tancredi arXiv:1503.08835

Calculation of NLO $gg \to ZZ$ cross section in model where Z bosons only couple to t quarks in s/M_t^2 expansion (LO) yields K-factor of 1.5–2 for 180 GeV $< M_{ZZ} < 340$ GeV
(LO QCD comparison with exact M_t: $M_t \to \infty$ poor for $M_{ZZ} \gtrsim 300$ GeV)
K. Melnikov, M. Dowling arXiv:1503.01274

Higgs width measurement in a nutshell

- Total Higgs width Γ_H is not a fundamental parameter of the theory, but of great phenomenological interest (Higgs mechanism $\rightarrow$ overall coupling strength)
- Direct Higgs width measurement via resonance shape is limited at LHC by experimental mass resolution of $\mathcal{O}(1)$ GeV (CMS: $\Gamma_H < 2.4$ GeV, but note that $\Gamma_{H,SM} \approx 4$ MeV)
- All resonant Higgs cross sections depend on Γ_H, therefore Γ_H and couplings cannot be determined at the LHC (on-peak) without theoretical assumptions M. Duhrssen et al. (2004), LHC Higgs Cross Section WG (2012)
- For broad class of models, assuming upper limit for HWW or HZZ coupling (e.g. SM) $\rightarrow$ upper bound for Γ_H $(\Gamma_H = \mathcal{O}(\Gamma_{H,SM}))$ M. Peskin (2012); B. Dobrescu, J. Lykken (2013)
- Assuming no BSM Higgs decays, and Higgs coupling parameterisations, can fit Higgs width to data and agreement with SM Higgs width is found V. Barger, M. Ishida, W. Keung (2012); K. Cheung, J. Lee, P. Tseng (2013); J. Ellis, T. You (2013); A. Djouadi, G. Moreau (2013); P. Bechtle, S. Heinemeyer, O. Stal, T. Stefaniak, G. Weiglein (2014)
- $e^+e^- \rightarrow Z(H \rightarrow \text{all})$: construct recoil mass and measure HZZ coupling $\rightarrow \Gamma_H$ can be determined indirectly, ILC: 6%–11% accuracy M. Peskin (2013), T. Han et al. (2013)
- Direct threshold scan at muon collider: Γ_H accuracy 4%–9% T. Han, Z. Liu (2013)
- Higgs width determination could provide first evidence for BSM Higgs interactions

Off-shell Higgs signal and Higgs width determination

indirect Higgs width determination via on- and off-peak Higgs cross section

F. Caola, K. Melnikov (2013) arXiv:1307.4935

$$|\mathcal{M}_{i\to H\to f}|^2 = \frac{|\mathcal{M}_i|^2\ |\mathcal{M}_f|^2}{|p_H^2 - M_H^2 + iM_H\Gamma_H|^2}$$

resonance contribution to signal cross section ("on-peak"): $\sigma_{i\to H\to f} \overset{\text{NWA}}{\propto} \frac{g_i^2 g_f^2}{\Gamma_H}$

NWA scaling degeneracy: σ unchanged if $g_i \to \xi\, g_i,\ \ g_f \to \xi\, g_f,\ \ \Gamma_H \to \xi^4\, \Gamma_H$

cf. L. Dixon, Y. Li arXiv:1305.3854 (see below)

$$\sqrt{p_H^2} - M_H \gg \mathcal{O}(\Gamma_H) \ \to\ p_H^2 - M_H^2 \gg M_H\Gamma_H \ \to\ |\mathcal{M}_{i\to H\to f}|^2 \approx \frac{|\mathcal{M}_i|^2\ |\mathcal{M}_f|^2}{|p_H^2 - M_H^2|^2}$$

off-resonance contribution ("off-peak"): $\sigma_{i\to H\to f}\left(\sqrt{p_H^2} - M_H \gg \mathcal{O}(\Gamma_H)\right) \propto g_i^2 g_f^2$

sizeable off-resonance contribution to signal cross section is independent of Higgs width, and therefore "breaks" NWA scaling degeneracy: $\sigma_{\text{off-peak}}/\sigma_{\text{on-peak}} \propto \Gamma_H$

competitive constraints on Higgs width without assumptions(?) feasible with LHC data

large interference with cont. background (necessary to prevent unitarity violation) weakens bounds

MCFM analysis

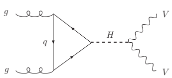
J. Campbell, K. Ellis, C. Williams (2013) (update of Caola-Melnikov analysis)

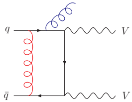
CMS Higgs selection cuts are applied:

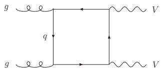
$p_{T\ell 1} > 20$ GeV, $p_{T\ell 2} > 10$ GeV, $p_{Te} > 7$ GeV, $p_{T\mu} > 5$ GeV, $|\eta_\mu| < 2.4$, $|\eta_e| < 2.5$, $M_{\ell\bar{\ell}} > 4$ GeV, $M_{4\ell} > 100$ GeV, 40 GeV $< M_{\ell\bar{\ell},\text{closest}} < 120$ GeV, 12 GeV $< M_{\ell\bar{\ell},\text{other}} < 120$ GeV (relative to M_Z)

Best prediction cross sections for $pp \to H \to ZZ \to e^- e^+ \mu^- \mu^+$ in fb, obtained using the running scale $m_{4\ell}/2$:

Energy	PDF	σ^H_{peak}	$m_{4\ell} > 130$ GeV σ^H_{off}	$m_{4\ell} > 130$ GeV σ^I_{off}	$m_{4\ell} > 300$ GeV σ^H_{off}	$m_{4\ell} > 300$ GeV σ^I_{off}
7 TeV	MSTW	0.203	0.025	-0.053	0.017	-0.025
	CTEQ	0.192	0.021	-0.047	0.015	-0.021
8 TeV	MSTW	0.255	0.034	-0.073	0.025	-0.036
	CTEQ	0.243	0.031	-0.065	0.022	-0.031
13 TeV	MSTW	0.554	0.108	-0.215	0.085	-0.122
	CTEQ	0.530	0.100	-0.199	0.077	-0.111







MCFM analysis

Higgs width bounds from matrix element method ($H \to ZZ$)

Matrix element method: optimize discrimination using fully differential information

Associate probabilistic weight with each event:

$$P(\phi) = \frac{1}{\sigma} \sum_{i,j} \int dx_1 dx_2\, \delta(x_1 x_2 s - Q^2) f_i(x_1) f_j(x_2) \hat{\sigma}_{ij}(x_1, x_2, \phi)$$

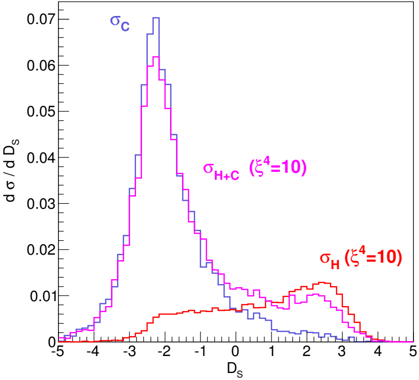
$P_{q\bar{q}}$: $q\bar{q}$ induced continuum background

P_{gg} : gg induced contributions
(incl. Higgs signal, cont. bkg. & interf.)

P_H : gg induced Higgs amplitude squared

Discriminant:

$$D_S = \log\left(\frac{P_H}{P_{gg} + P_{q\bar{q}}}\right)$$



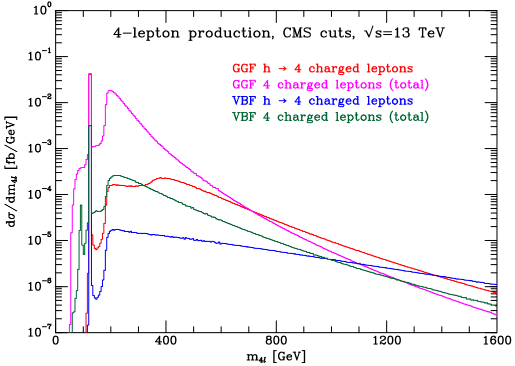
$\Gamma_H < \left(15.7^{-2.9}_{+3.9}\right)\, \Gamma_H^{SM}$ (95% CL), $(D_S > 1)$

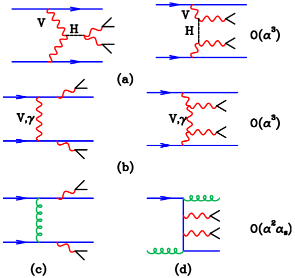
bound $1.6\times$ better than for $m_{4\ell} > 300$ GeV

Higgs (width) constraints from vector boson fusion

J.M. Campbell, R.K. Ellis, arXiv:1502.02990

(see also C. Englert, M. Spannowksy arXiv:1405.0285, C. Englert, M. McCullough, M. Spannowksy arXiv:1504.02458)





most sensitive off-sh. channel: $W^\pm W^\pm$ due to lower bkg. (t-channel Higgs!)

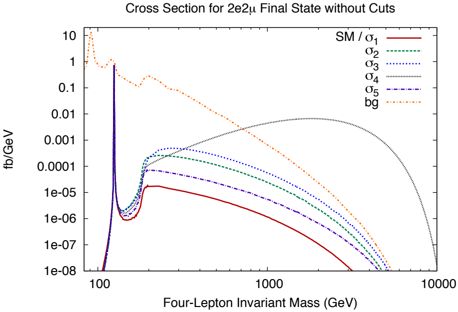
$\Gamma_H < 61\ \Gamma_H^{SM}$ (LHC Run 1)

$\Gamma_H < 4.4\ \Gamma_H^{SM}$ (LHC 100 fb^{-1} data)

$\Gamma_H < 3.2\ \Gamma_H^{SM}$ (LHC 300 fb^{-1} data)

BSM searches and model builder's considerations

Constraining higher dimensional operators with the off-shell Higgs (see below)
Disentangling New Physics with the off-shell Higgs boson
EFT studies including the off-shell Higgs boson
Limitations of model independence



$\mathcal{O}_1 = -\frac{M_Z^2}{v} H Z_\mu Z^\mu$ (SM), $\mathcal{O}_2 = -\frac{1}{2v} H Z_{\mu\nu} Z^{\mu\nu}$, $\mathcal{O}_3 = -\frac{1}{2v} H Z_{\mu\nu} \tilde{Z}^{\mu\nu}$, $\mathcal{O}_4 = \frac{M_Z^2}{M_H^2 v} Z_\mu Z^\mu \partial^2 H$,

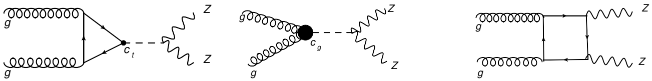
$\mathcal{O}_5 = \frac{2}{v} H Z_\mu \partial^2 Z^\mu$ J. Gainer, J. Lykken, K. Matchev, S. Mrenna, M. Park arXiv:1403.4951

Also: modification of lepton angular distributions $\longrightarrow$ good control with 300 fb^{-1} I. Anderson et al. arXiv:1309.4819

EFT analysis of on- and off-shell $H \to ZZ \to 4\ell$ data

A. Azatov, C. Grojean, A. Paul, E. Salvioni (2014)

(see also G. Cacciapaglia, A. Deandrea, G. Drieu La Rochelle, J. Flament (PRL 2014))

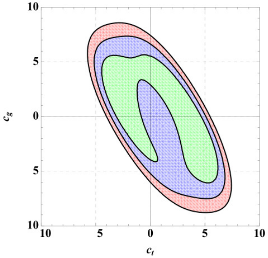


$$\mathcal{L} = -c_t \frac{m_t}{v} \bar{t}th + \frac{g_s^2}{48\pi^2} c_g \frac{h}{v} G_{\mu\nu}G^{\mu\nu}$$

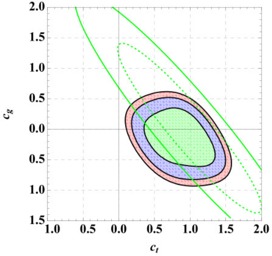
$$\mathcal{M}_{gg\to ZZ} = \mathcal{M}_h + \mathcal{M}_{bkg} = c_t\, \mathcal{M}_{c_t} + c_g\, \mathcal{M}_{c_g} + \mathcal{M}_{bkg}$$

$\sigma \sim |c_t + c_g|^2$: on-shell degeneracy $c_t + c_g = \mathrm{const}$ is broken by far-off-shell data

Constraints in (c_t, c_g) plane (68%, 95% and 99% probability contours): (not MELA improved!)

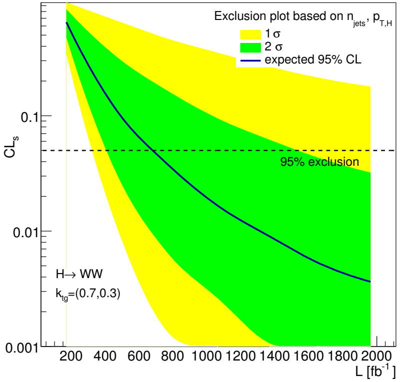


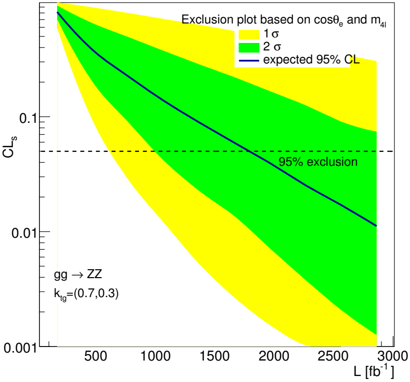
LHC 8 TeV CMS data



LHC 14 TeV 3 ab^{-1} data

Effective ggH coupling: boosted v. off-shell Higgs sensitivity

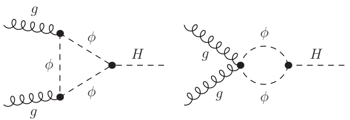


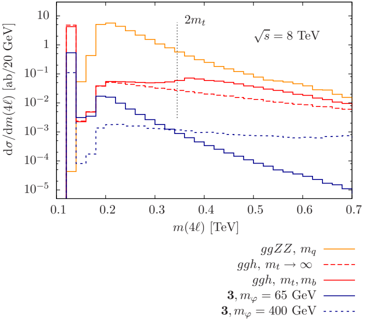


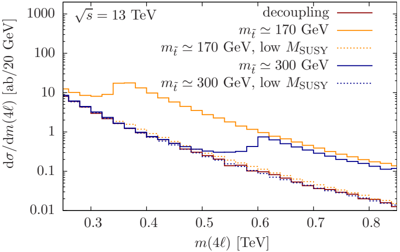
left: boosted analysis, right: off-shell analysis (not MELA improved)

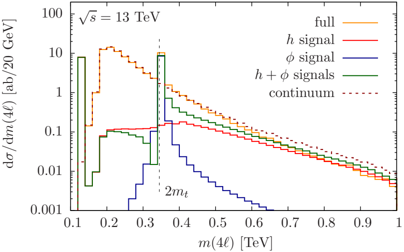
M. Buschmann, D. Goncalves, S. Kuttimalai, M. Schoenherr, F. Krauss, T. Plehn (2014) (1410.5806)

BSM benchmark scenario studies





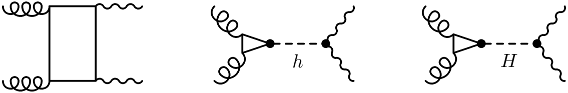




C. Englert, M. Spannowsky arXiv:1405.0285,1410.5440

Loophole: additional light scalar in the s-channel

[H.E. Logan, 1412.7577]



SM: h cancels growth $\propto E/v$ of $t\bar{t} \to ZZ$ amplitude.

Modified h couplings: cancellation imperfect; growth of amplitude with E provides LHC sensitivity at high m_{ZZ}!

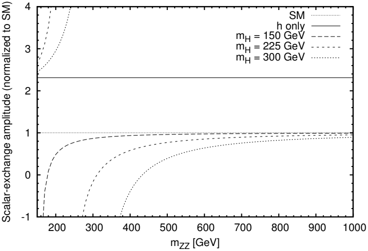
Extended Higgs sector: Require $\kappa_t^h \kappa_Z^h + \kappa_t^H \kappa_Z^H = 1$ for unitarity of $t\bar{t} \to ZZ$ (automatic in renormalizable models): $\kappa_t^h \kappa_Z^h = 1 + \Delta > 1, \quad \kappa_t^H \kappa_Z^H = -\Delta$

Amplitude relative to SM:

$$\frac{\mathcal{M}_h + \mathcal{M}_H}{\mathcal{M}_{h_{\rm SM}}} = (1+\Delta) - \Delta \frac{p^2 - m_h^2}{p^2 - m_H^2}$$

$$\simeq 1 - \Delta \frac{(m_H^2 - m_h^2)}{p^2}$$

$$\to 1 \text{ for } p^2 \gg m_{h,H}^2$$

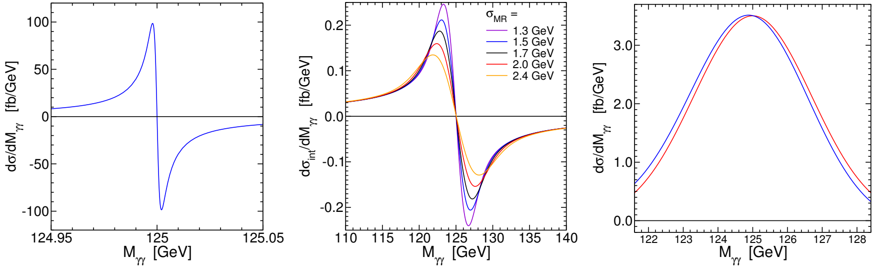


Presence of H at low mass (well below 350 GeV) causes $gg \to ZZ$ cross section to be SM-like at high m_{ZZ}, even if $\kappa_t^h \kappa_Z^h$ is strongly non-SM-like.

Higgs width via interferometry in $gg \to H \to \gamma\gamma$

S. Martin arXiv:1208.1533 (LO analysis of Higgs mass peak shift)

Higgs signal continuum background interference induces sizeable peak shift in $gg \to H \to \gamma\gamma$ (but negligible in $gg \to H \to ZZ^*$)



left fig.: interference contribution (real term) before detector resolution effects

center fig.: interference contribution (real term) for different mass resolutions (Gaussian, σ_{MR})

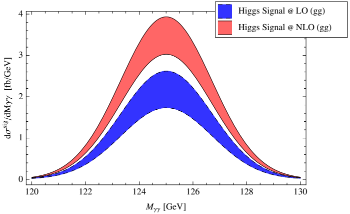
right fig.: peak shift of invariant mass distribution ($\sigma_{\text{MR}} = 1.7$ GeV): $\Delta M_{\gamma\gamma} = -120$ MeV at LO

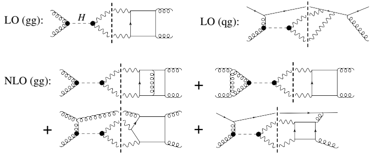
$(H \to \gamma\gamma)$+jet at LO: negligible mass peak shift (< 20 MeV for $p_{Tj} > 25$ GeV)

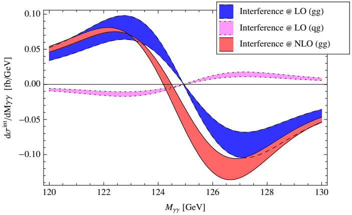
Daniel de Florian, et al. arXiv:1303.1397; S. Martin arXiv:1303.3342

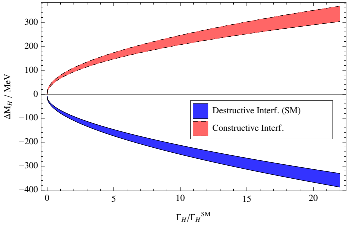
Higgs width via interferometry in $gg \to H \to \gamma\gamma$

L. Dixon, Y. Li arXiv:1305.3854 (NLO analysis and Higgs width constraint)









SM mass shift: $\Delta M_{\gamma\gamma} = -70$ MeV at NLO

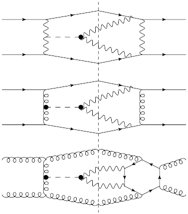
Vary Higgs width and couplings (maintaining on-peak SM signal strenghts):

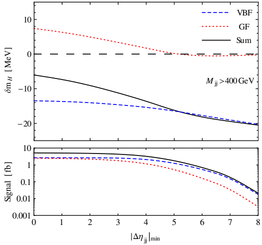
$$\Gamma_H < 15\,\Gamma_H^{SM} \text{ (14 TeV, 3 ab}^{-1}\text{, 95\% CL)}$$

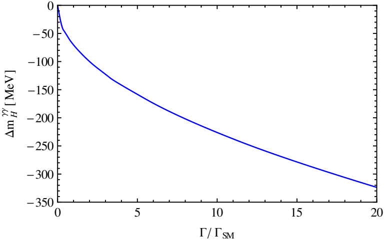
Higgs width via interferometry in $gg \to H \to \gamma\gamma$ with VBF $H \to \gamma\gamma$ mass peak as reference

Calculation of $pp \to H\ (\to \gamma\gamma) + 2$ jets signal (VBF and GF) and interference with background (LO)

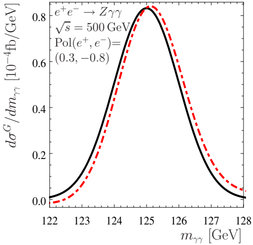
F. Coradeschi, D. de Florian, L. Dixon, N. Fidanza, S. Hoeche, H. Ita, Y. Li, J. Mazzitelli arXiv:1504.05215







$\Delta m_H^{\gamma\gamma} \equiv \delta m_{H+X,\text{NLO,incl}}^{\gamma\gamma} - \delta m_{H+2j,\text{LO,VBF cuts}}^{\gamma\gamma}$

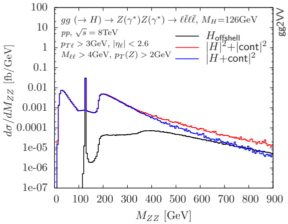


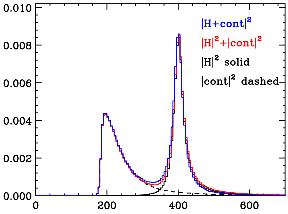
LC: $H \to \gamma\gamma$ mass peak shift S. Liebler arXiv:1503.07830

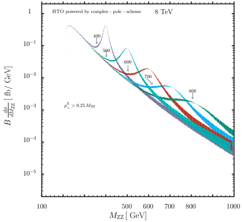
Heavy Higgs - light Higgs - continuum VV interference

consider a heavy Higgs h_2 (signal) in addition to a light Higgs h_1 at 125 GeV (background)

Two-Higgs model: SM & real EW singlet scalar, as defined in YR3 arXiv:1307.1347, Sec. 13.3







right fig.: G. Passarino (arXiv:1206.3824)

What is the impact of interference with the offshell tail of the 125 GeV Higgs for a heavy Higgs of 300, 600 or 900 GeV?

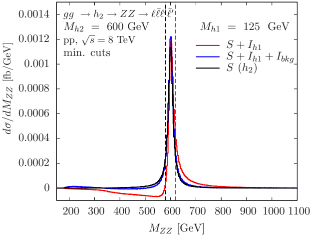
$$S \sim |\mathcal{M}_{h2}|^2$$

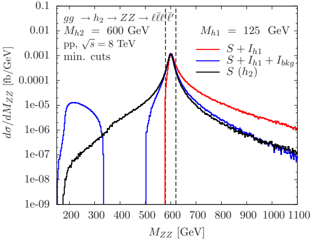
$$I_{h1} \sim 2\,\mathrm{Re}(\mathcal{M}^*_{h2}\,\mathcal{M}_{h1})$$

$$I_{bkg} \sim 2\,\mathrm{Re}(\mathcal{M}^*_{h2}\,\mathcal{M}_{bkg})$$

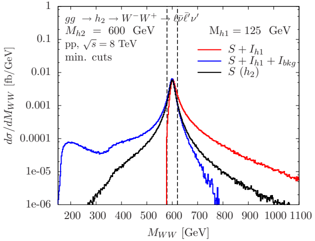
$$I_{full} \sim 2\,\mathrm{Re}(\mathcal{M}^*_{h2}\,(\mathcal{M}_{h1} + \mathcal{M}_{bkg}))$$

Heavy Higgs - light Higgs - continuum VV interference

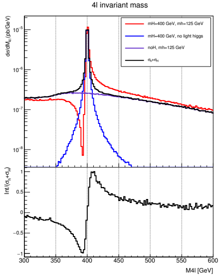


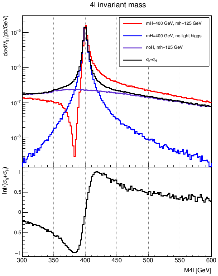


NK, C. O'Brien arXiv:1502.04113



NK, C. O'Brien



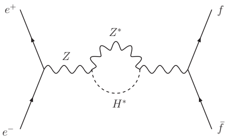


E. Maina arXiv:1501.02139

see also C. Englert, I. Low, M. Spannowsky arXiv:1502.04678 and C. Englert, Y. Soreq, M. Spannowsky arXiv:1410.5440

Higgs width/coupling constraints including LEP EW PO

C. Englert, M. McCullough, M. Spannowksy arXiv:1504.02458



Higgs contribution to Z boson self-energy is Higgs width independent (at LO), but depends on HZZ coupling

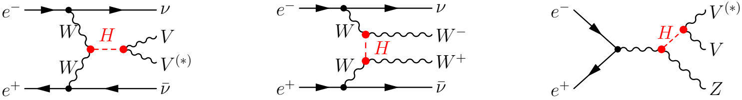
$\rightarrow$ similar characteristics as (tree-level) off-shell Higgs signal

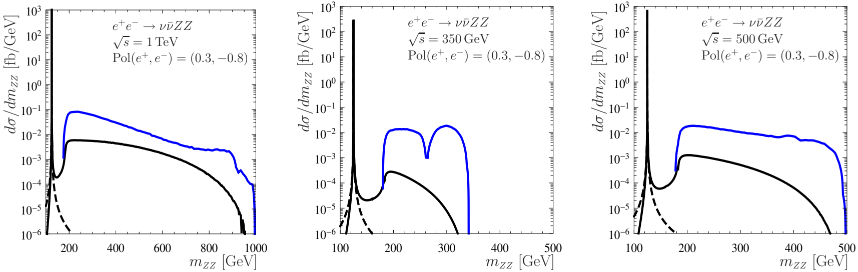
Consider Peskin-Takeuchi parameters S, T, U with rescaled HVV couplings ($g = c_V\, g_{\mathrm{SM}}$):

$$
\begin{aligned}
s_H(m_H, c_V) &= \frac{c_V^2}{\pi M_Z^2}\left(B^{OO}_{M_Z^2,Z} - B^{OO}_{0,Z} - M_Z^2(B^{O}_{M_Z^2,Z} - B^{O}_{0,Z})\right) \\
t_H(m_H, c_V) &= \frac{c_V^2}{4\pi M_W^2 S_W^2}\left(B^{OO}_{0,W} - B^{OO}_{0,Z} + M_Z^2 B^{O}_{0,Z} - M_W^2 B^{O}_{0,W}\right) \\
u_H(m_H, c_V) &= \frac{c_V^2}{\pi}\Big((B^{O}_{M_Z^2,Z} - B^{O}_{0,Z}) - (B^{O}_{M_W^2,W} - B^{O}_{0,W}) - \frac{1}{M_Z^2}(B^{OO}_{M_Z^2,Z} - B^{OO}_{0,Z}) \\
&\quad + \frac{1}{M_W^2}(B^{OO}_{M_W^2,W} - B^{OO}_{0,W})\Big)
\end{aligned}
$$

High-mass $H \to VV$ signal at a linear collider

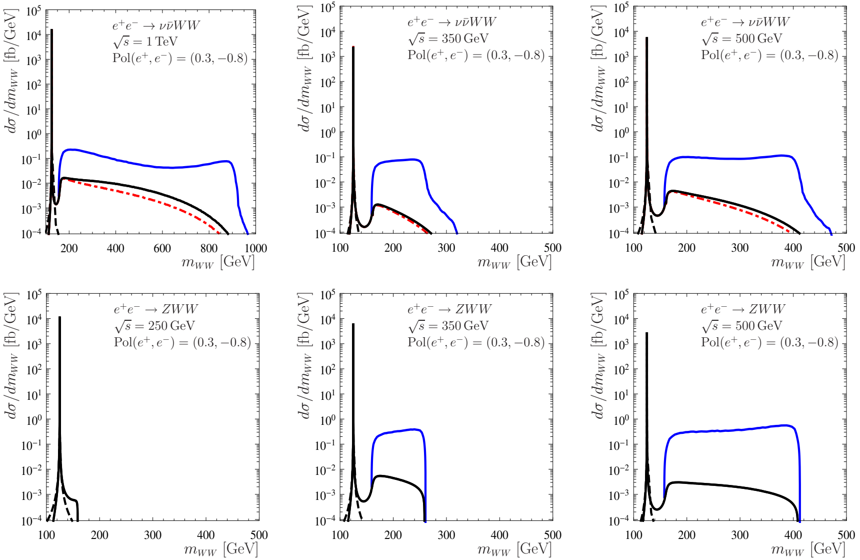
S. Liebler, G. Moortgat-Pick, G. Weiglein (2014), arXiv:1502.07970





High-mass $H \to VV$ signal at a linear collider

S. Liebler, G. Moortgat-Pick, G. Weiglein (2014), arXiv:1502.07970



Summary

- $H \to ZZ, WW$ in ggF & VBF @ LHC: $\mathcal{O}(10\%)$ off-shell high-mass Higgs signal contribution with large Higgs(-Higgs)-continuum interference: now taken into account, provides complementary physics information (similar at high-energy linear collider)
- $gg \to H \to ZZ, WW \to 2\ell 2\ell,\ 4\ell,\ 2\ell 2\nu$: interference studied in great detail, tools & events available (caveat: LO); NLO calculation: technically hard, impressive progress
- Semileptonic channels $gg \to H \to WW \to \ell\nu qq'$ and $gg \to H \to ZZ \to \ell\bar{\ell}q\bar{q}$ contribute to ongoing (heavy) Higgs analyses; first analysis of interference effects in semileptonic channels, new feature: interfering tree-level background
- First analysis of interference (& Higgs width bounds) for $pp \to H \to ZZ$ + jet
- First analysis of heavy Higgs-light Higgs-bkg. interference effects in $gg \to H \to VV$, complementary studies for VBF and linear collider
- Direct Higgs width measurement at LHC limited by mass resolution: $\Gamma_H < 600\,\Gamma_H^{SM}$
- high-mass Higgs tail not Higgs width dependent $\longrightarrow$ provides complementary constraints on Higgs couplings and Higgs width Γ_H (when combined with on-peak data)
- Assuming no E-dependence of relevant Higgs couplings, a bound on Γ_H can be obtained; optimise bound with fully differential discriminant (Matrix Element Method)
- LHC Run 1: ATLAS and CMS Higgs width bounds $\longrightarrow$ Roberto Covarelli's talk
- $H \to \gamma\gamma$: interference-facilitated bound $\Gamma_H < 15\,\Gamma_H^{SM}$ (14 TeV, 3 ab^{-1}, 95% CL)
- LHC Run 2: improved bounds (ggF & VBF), high-mass $H \to VV$ EFT and BSM benchmark studies