

Chances for SUSY-GUT in the LHC Epoch

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based on paper [arXiv:1505.04950](https://arxiv.org/abs/1505.04950)

in collaboration with Z. Berezhiani, G. Miele and S. Morisi*

* at this workshop



Beyond Standard Model

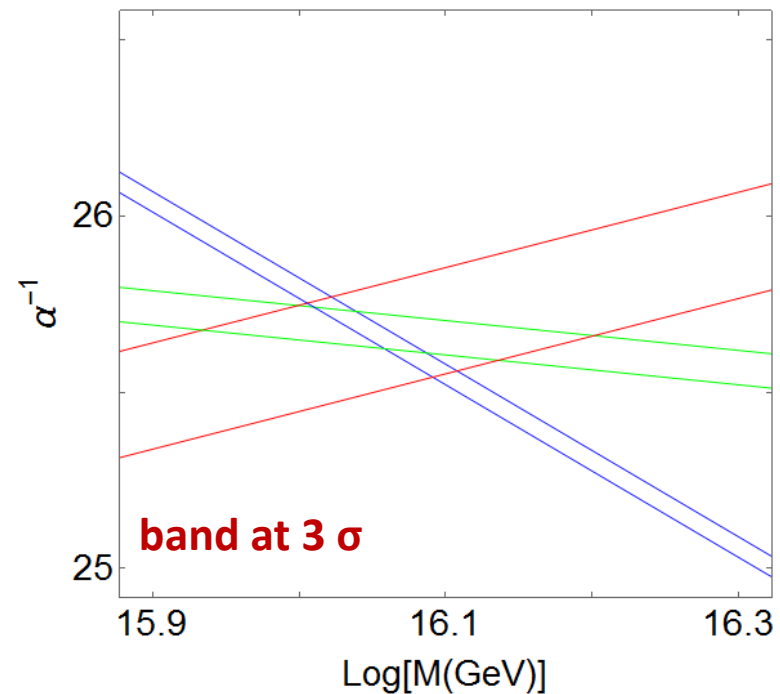
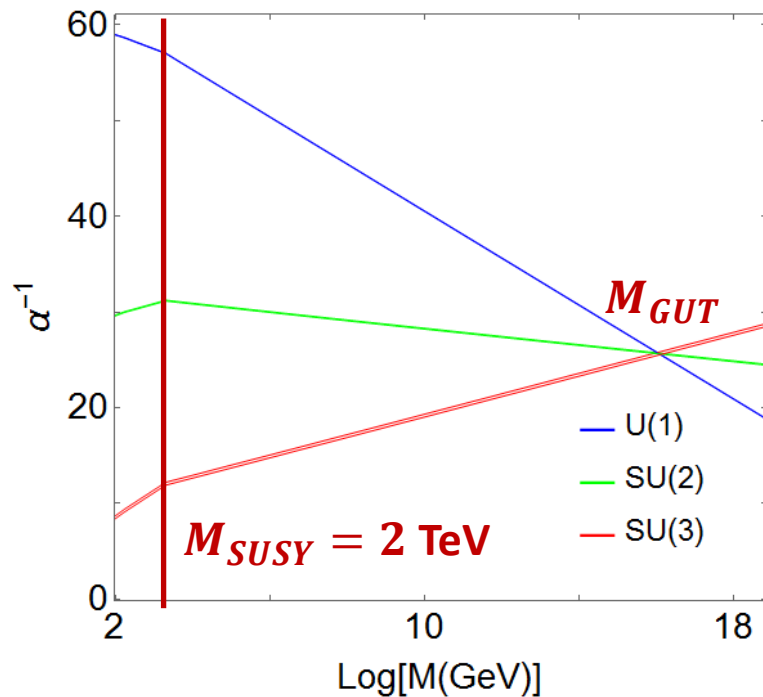
- Several phenomena or open problems suggest the presence of physics beyond Standard Model (SM):
 - hint of gauge unification;
 - structure of fermion masses;
 - hierarchy problem;
 - baryogenesis.
- The Supersymmetric Grand Unified Theories (SUSY-GUTs) are able to address part of these problems.
- After the 8 TeV LHC run I, we try:
 - to reanalyze the room still remaining for SUSY-GUT inspired models;
 - to determine the **upper bound M_{UB}** for the energy below which SUSY signatures have to show up.

Assumptions

- To perform our study, we require:
 - One step gauge unification at a single energy scale M_{GUT} , without intermediate symmetry scales;



SU(5) Bottleneck



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SU(5) Bottleneck

- Consistency of third family fermion masses and possible Yukawa b- τ unification;

$$y_b(M_{GUT}) = y_\tau(M_{GUT}) \left(1 + \mathcal{O} \left(\frac{y_\mu(M_{GUT})}{y_\tau(M_{GUT})} \right) \right)$$

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- Consistency of third family fermion masses and possible Yukawa b- τ unification;

$$y_b(M_{GUT}) = y_\tau(M_{GUT}) \left(1 + \mathcal{O} \left(\frac{y_\mu(M_{GUT})}{y_\tau(M_{GUT})} \right) \right)$$

- Consistency with the experimental limit on proton decay;
- Absence of special fine tunings among the parameters in the GUT scenario, implying couplings $\mathcal{O}(1)$ above M_{GUT} .



Naturalness

Renormalization Group Equations

- We have developed a *Mathematica* code, which resolves all the RGEs up to **2-loop order** with numerical iterative method.

$$\frac{d}{dt}X_i = \frac{1}{16\pi^2}\beta_{X_i}^{(1)}(X_j) + \frac{1}{(16\pi^2)^2}\beta_{X_i}^{(2)}(X_j)$$

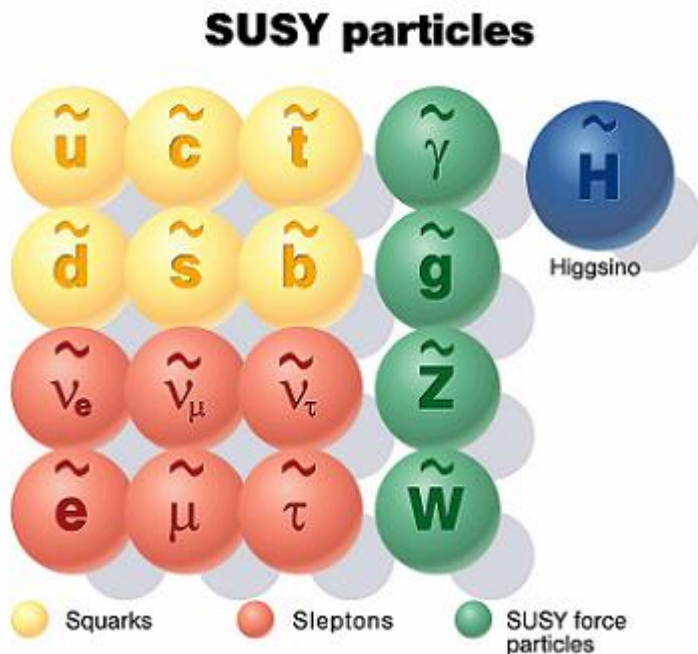
- The analysis takes into account all the matching and threshold relations at 1-loop level.
- We consider the general possibility of:
 - several SUSY thresholds (**multi-scale approach**);
 - GUT threshold.
- Only the third generation Yukawa couplings are relevant.

Input and output parameters

- The set of input parameters are

$$\{\tilde{m}_h, \tilde{m}_g, \tilde{m}_{sq}, \chi, \tan \beta\}$$

SUSY thresholds



$\tilde{m}_h \equiv$ higgsinos

$\tilde{m}_g \equiv$ gluinos

$\tilde{m}_{sq} \equiv$ squarks

SUSY μ -term

F-term

D-term

- Two other constrained masses.

At M_{GUT}

$$\text{SU}(5) \longrightarrow \frac{\tilde{m}_g}{\tilde{m}_W} = 1 \quad \frac{\tilde{m}_{sq}}{\tilde{m}_{sl}} = \frac{\tilde{m}_{10}}{\tilde{m}_5}$$

IRR: $\bar{5} + 10$

Higgs sector and SUSY thresholds

- The mass matrix of the Higgs scalars H_u and H_d involves mass parameters of different origin.

$$\mathcal{M}^2 = \begin{pmatrix} \tilde{M}_u^2 + \mu^2 & \mu B_\mu \\ \mu B_\mu & \tilde{M}_d^2 + \mu^2 \end{pmatrix}$$

SUSY μ -term

$$\mu^2 = \mathcal{O}(\tilde{m}_h)$$

F-term

$$\mu B_\mu = \mathcal{O}(\tilde{m}_g)$$

D-term

$$\tilde{M}_{u,d}^2 = \mathcal{O}(\tilde{m}_{sq})$$

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D-term

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- A fine-tuning condition has to be imposed in order to get the SM Higgs.

$$-m^2 = \frac{1}{2} \left(2\mu^2 + \tilde{M}_u^2 + \tilde{M}_d^2 - \sqrt{4\mu^2 B_\mu^2 + (\tilde{M}_u^2 - \tilde{M}_d^2)^2} \right) \sim -(100 \text{ GeV})^2$$



$$\tilde{m}_h \sim \tilde{m}_g \sim \tilde{m}_{sq}$$

**same order of
magnitude**

Input and output parameters

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GUT threshold

- The Higgs Σ , which breaks SU(5) down to the MSSM, can have mass M_Σ smaller than M_{GUT} .

$$V_\Sigma = \frac{M_\Sigma}{2} \Sigma^2 + \frac{\lambda_\Sigma}{2} \Sigma^3$$

$$\chi = \frac{M_{GUT}}{M_\Sigma}$$



$$\lambda_\Sigma = \frac{\sqrt{2\pi\alpha_{GUT}}}{\chi} = \mathcal{O}(1)$$

Naturalness

Input and output parameters

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Yukawa matching condition

- In the transition between SM and MSSM we have to impose

$$v_u = \langle H_u^0 \rangle$$

$$v_d = \langle H_d^0 \rangle$$

$$\tan \beta = \frac{v_u}{v_d}$$



$$m_t = y_t v \sin \beta$$

$$m_b = y_b v \cos \beta$$

$$m_\tau = y_\tau v \cos \beta$$

Input and output parameters

- The set of input parameters are

$$\{\tilde{m}_h, \tilde{m}_g, \tilde{m}_{sq}, \chi, \tan \beta\}$$

- The outputs are

$$\{\alpha_3(M_Z), M_{GUT}, \alpha_{GUT}, y_t(M_{GUT}), y_b(M_{GUT}), y_\tau(M_{GUT})\}$$

- These two values must be compatible with the experimental measurements.

EW measurement

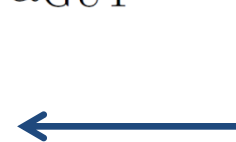
$$\alpha_3^{\text{exp}}(M_Z) = 0.1184 \pm 0.0007$$

proton decay

$$\frac{M_{GUT}}{\sqrt{\alpha_{GUT}}} \gtrsim 3 \cdot 10^{16} \text{ GeV}$$



compatibility
constraints



Proton decay

- We take into account the **6d** operators describing the proton decay mediated by leptoquarks.

$$\Gamma(p \rightarrow e^+ \pi^0) = \frac{\pi}{4} \frac{\alpha_{\text{GUT}}^2}{M_{\text{GUT}}^4} \frac{m_p}{f_\pi^2} \alpha_H^2 |1 + D + F|^2 \left(1 - \frac{m_\pi^2}{m_p^2}\right)^2 \left[\left(A_{\text{R}}^{(1)}\right) + \left(A_{\text{R}}^{(2)}\right) \left(1 + |V_{ud}|^2\right)^2 \right]$$

Hisano, Kobayashi, Nagata, PL B716 (2012)

$$\tau_p / \text{Br}(p \rightarrow e^+ \pi^0) > 1.29 \cdot 10^{34} \text{ yr}$$



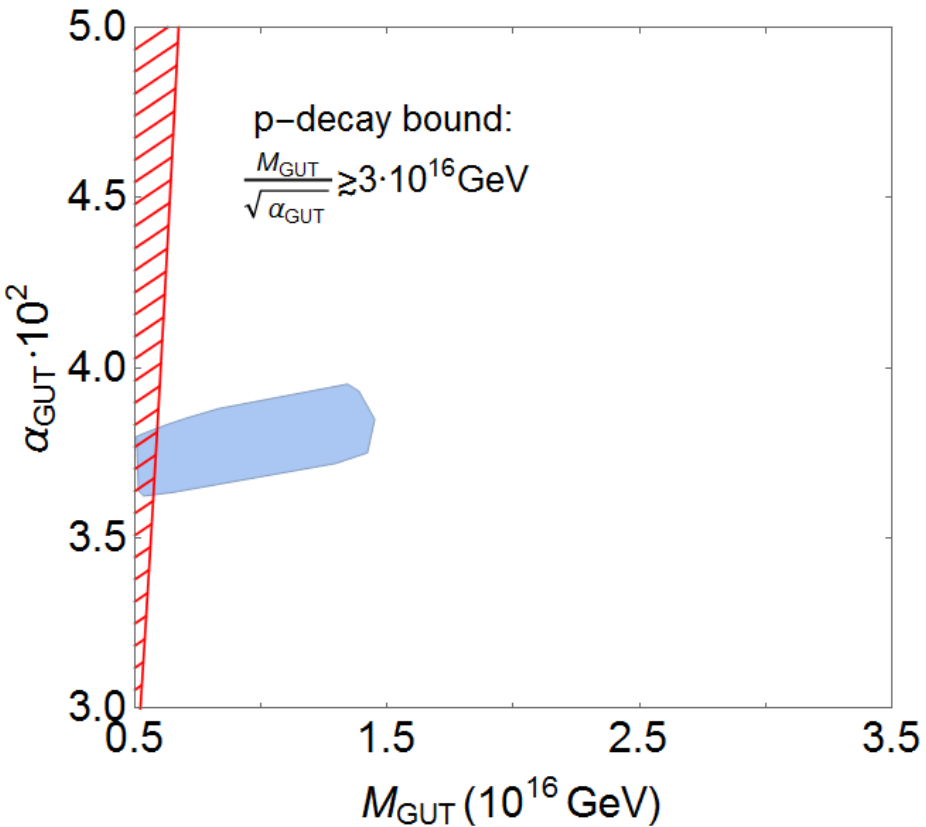
$$\frac{M_{\text{GUT}}}{\sqrt{\alpha_{\text{GUT}}}} \gtrsim 3 \cdot 10^{16} \text{ GeV}$$

Super-Kamiokande, PR D85:112001 (2012)

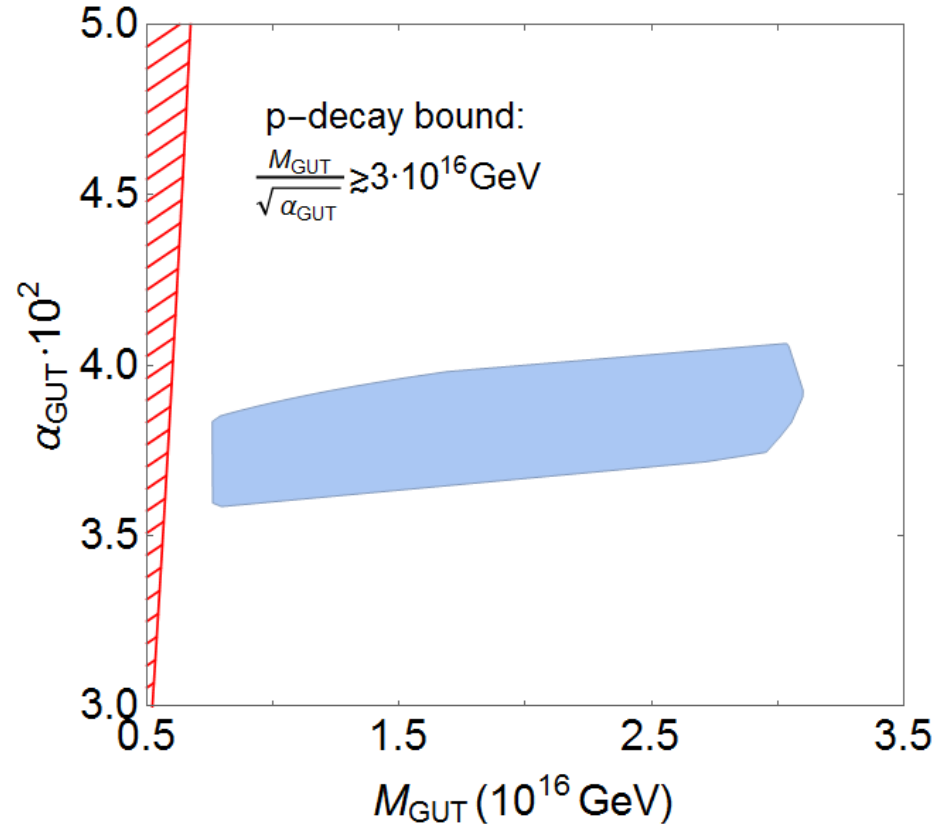
- We do not consider the **5d** operators since they are strongly model dependent.

GUT region

$$\chi = 1$$

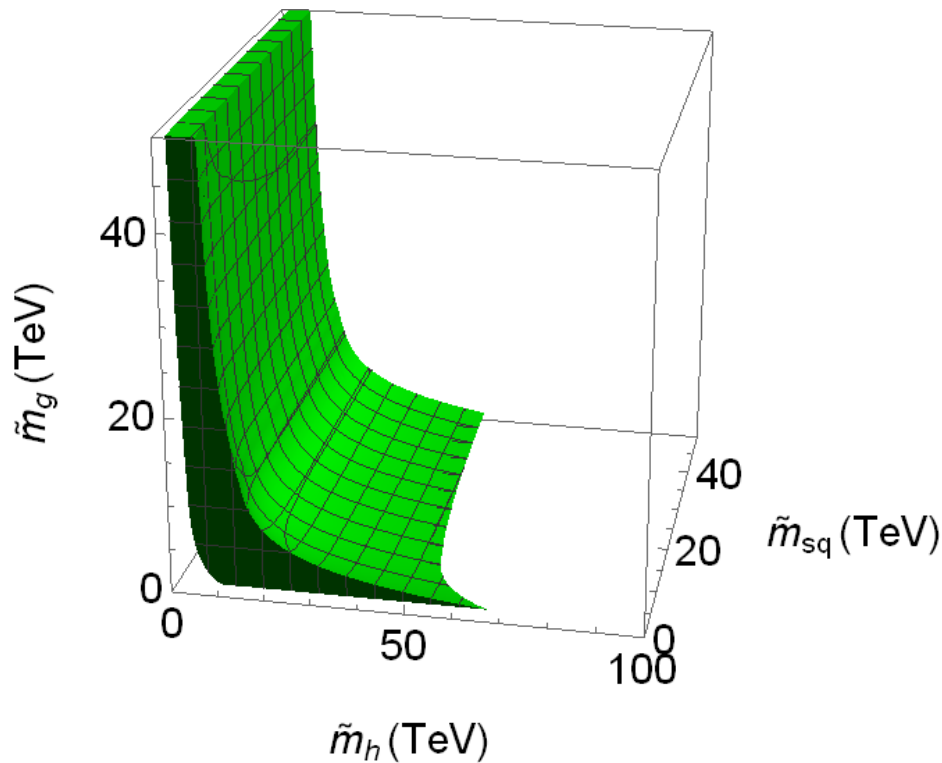


$$\chi = 10$$

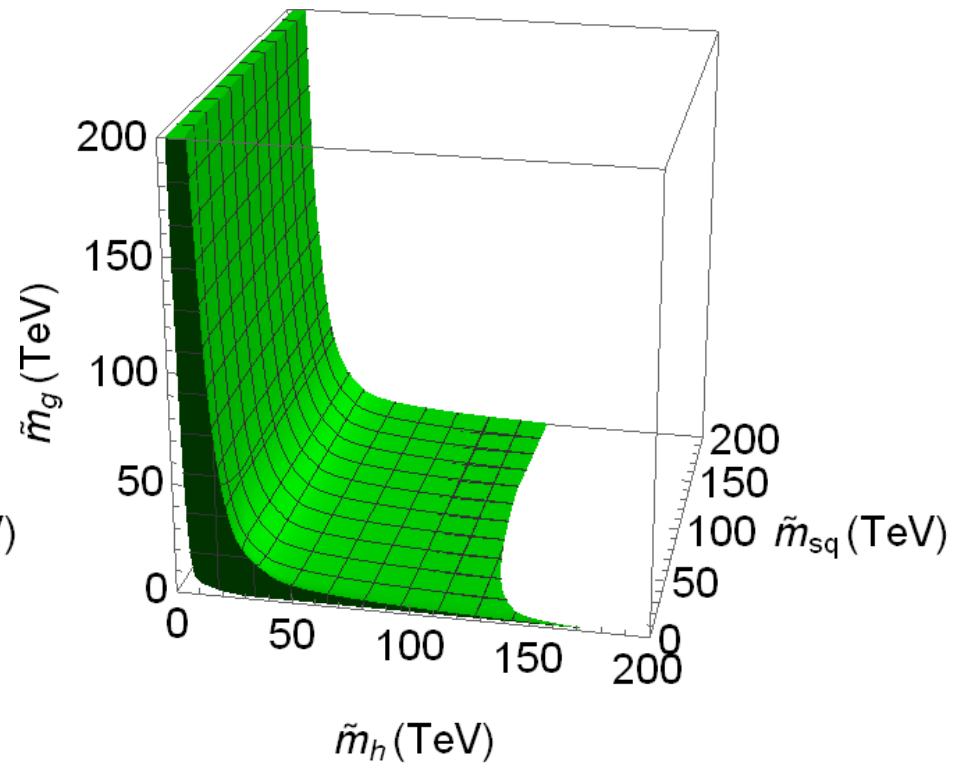


SUSY mass spectrum

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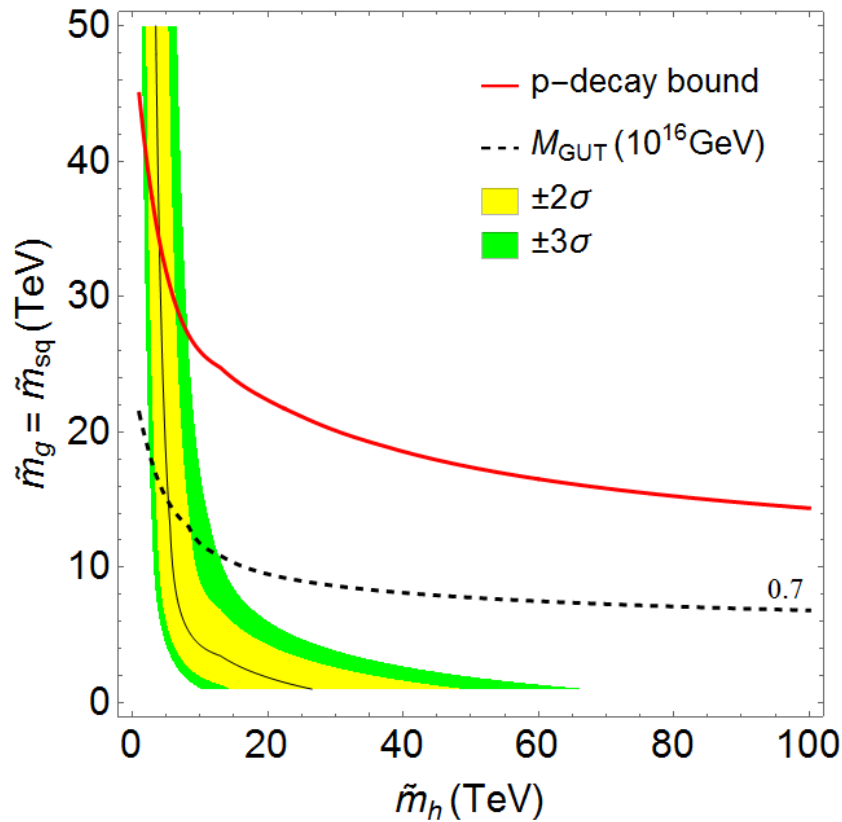


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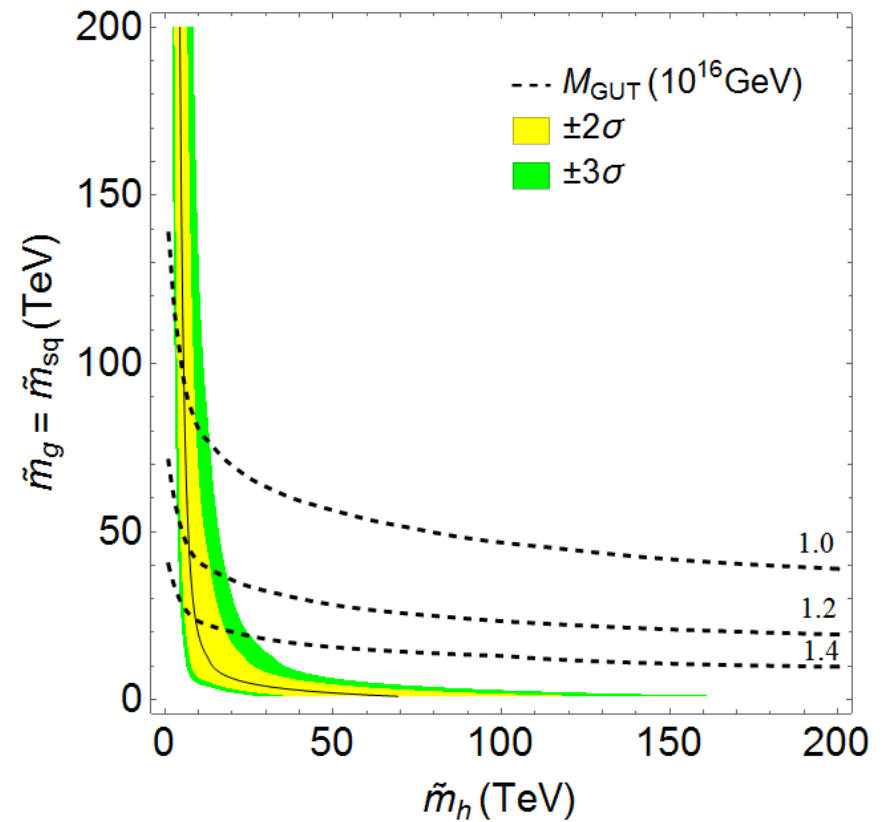


Surface $\tilde{m}_g = \tilde{m}_{sq}$

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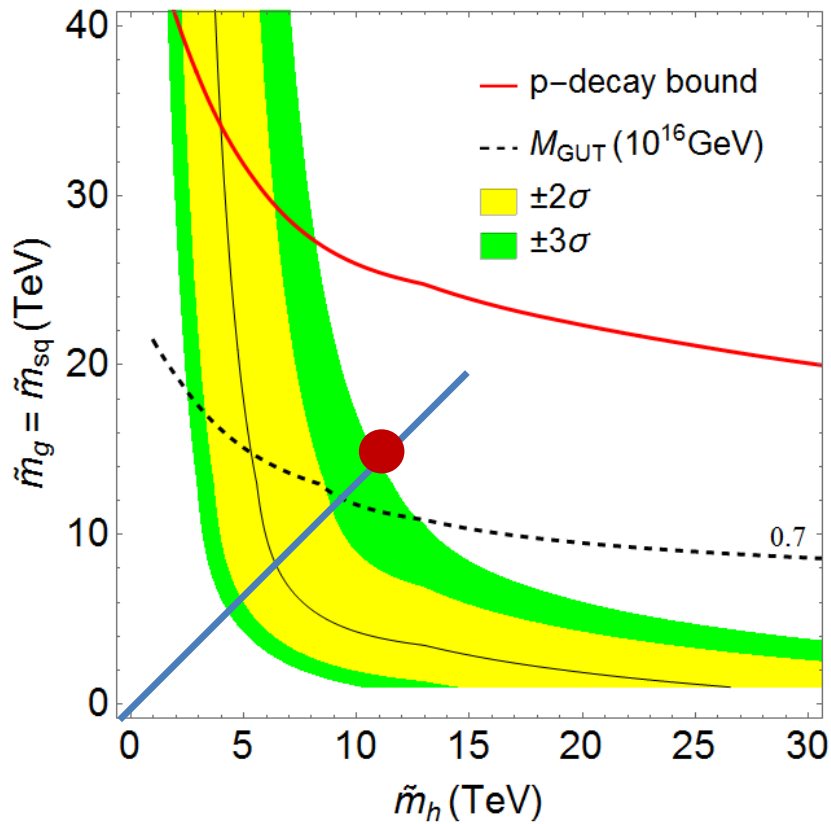


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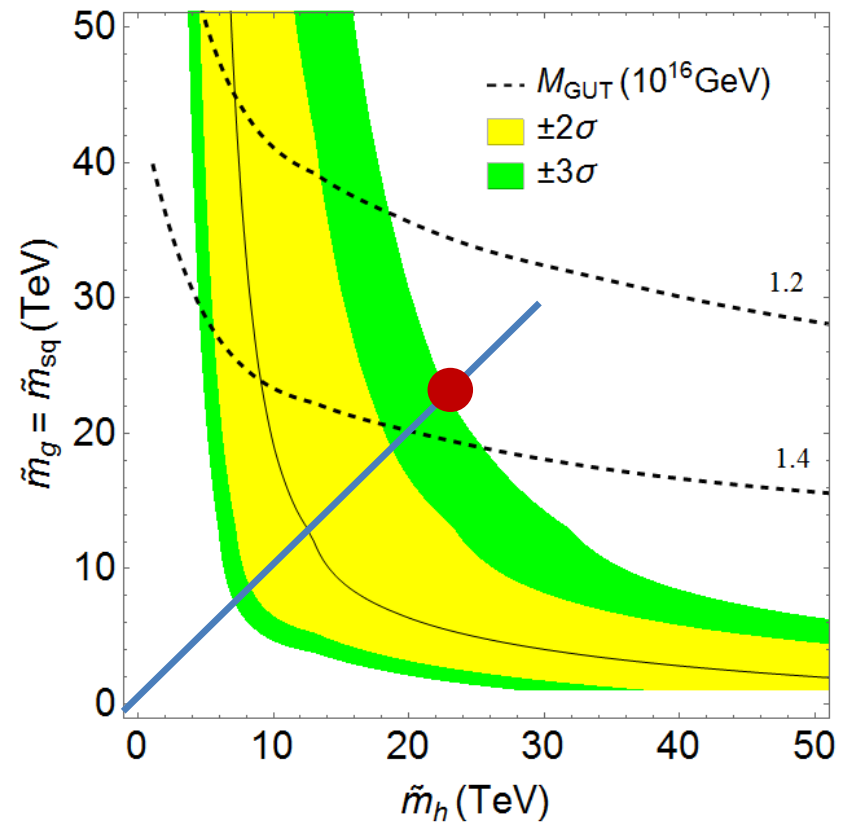


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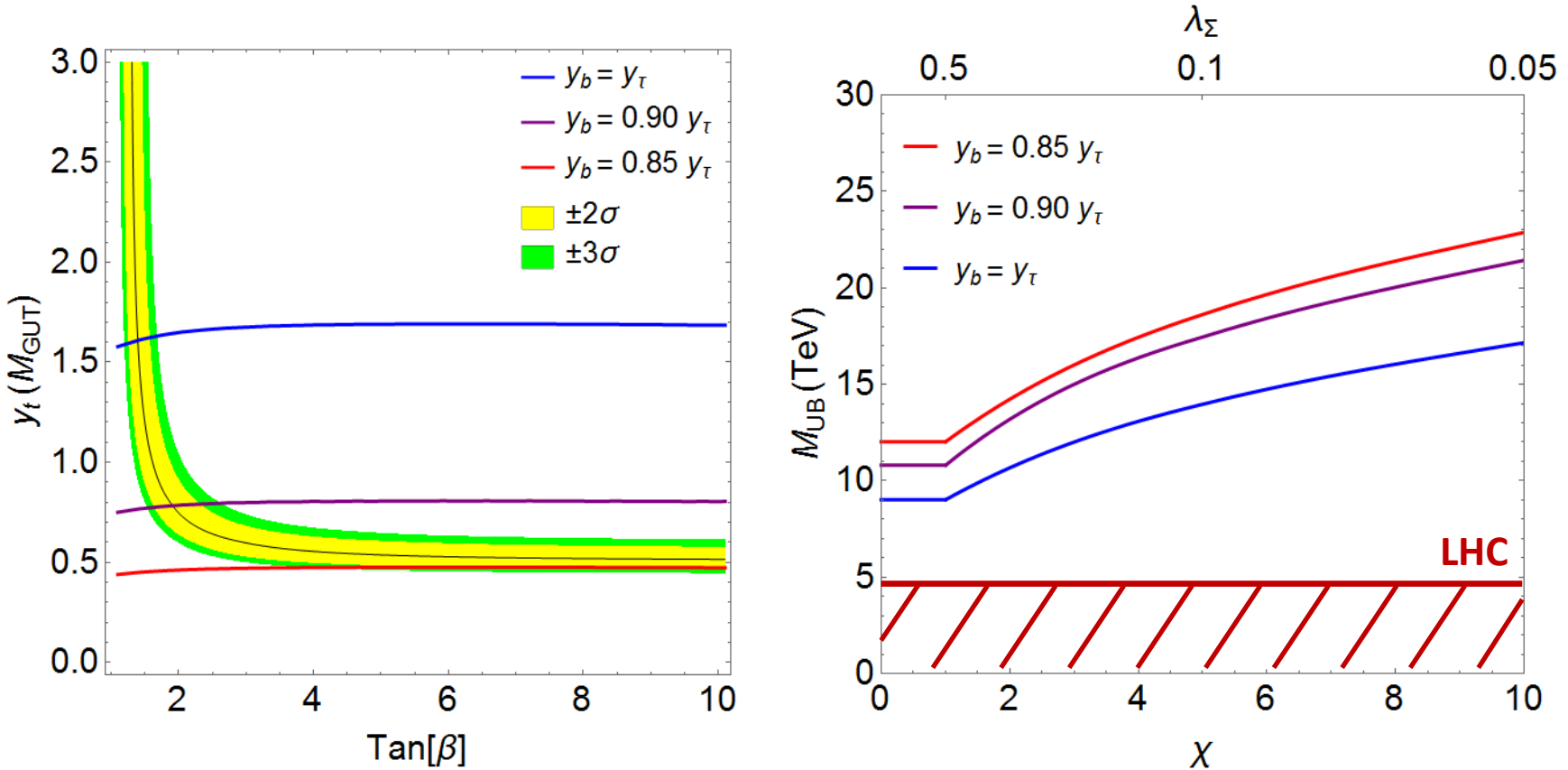
$\chi = 1$



$\chi = 10$



Upper bound for SUSY physics



- Larger values for GUT threshold χ are not allowed since:
 - naturalness requirement implies $\lambda_\Sigma \sim \mathcal{O}(1)$;
 - M_{GUT} unnaturally approaches M_{Plack} .



$M_{UB} \sim 20$ TeV

Conclusions

- After LHC run I (8 TeV), for planning new colliders it is of interest:
 - to reanalyze the room still remaining for SUSY-GUT inspired models;
 - to determine the **upper bound M_{UB}** for the energy below which SUSY signatures have to show up.
- Assuming one step unification (SU(5) bottleneck), under natural assumptions we have obtained general bounds on SUSY mass spectrum.
- We claim that if a SUSY-GUT model is the proper way to describe physics beyond the SM, the lightest gluino or higgsino cannot have a mass larger than

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Thanks for your attention