

# Falsifying High-Scale Baryogenesis with $0\nu\beta\beta$ , LFV and the LHC

Frank F. Deppisch, JH, Martin Hirsch,  
Phys. Rev. Lett. 112, 221601 (2014), arXiv: 1312.4447 [hep-ph]

Frank F. Deppisch, JH, Martin Hirsch, Wei-Chih Huang, Heinrich Päs,  
arXiv: 1503.04825 [hep-ph]

Julia Harz

University College London

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WIN 2015, MPIK Heidelberg

- Observation of a baryon asymmetry of the Universe (BAU)

$$\eta_B^{\text{obs}} = \frac{n_B - n_{\bar{B}}}{n_\gamma} = (6.09 \pm 0.06) \times 10^{-10}$$

P. A. R. Ade et al. [Planck Collaboration], arXiv:1502.01589 [astro-ph.CO]

- Theoretical requirements: 3 Sakharov conditions

A. D. Sakharov, JETP Lett. 5, 24 (1967)

- $CP$  violation
- departure from thermal equilibrium
- (B-L)-violation



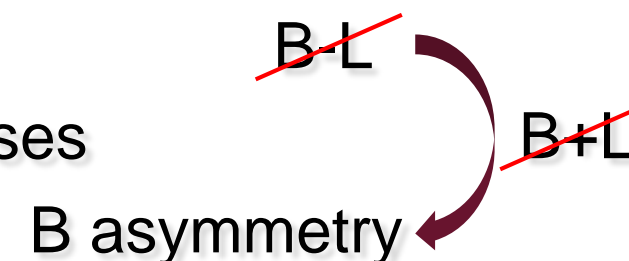
not fulfilled within the Standard Model



Physics beyond the Standard Model?

- Baryogenesis via Leptogenesis:

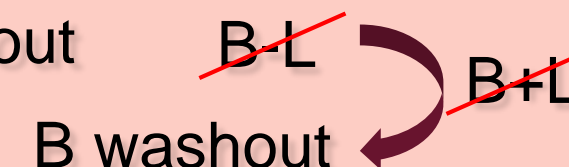
- generation via heavy neutrino decays
- competition with lepton number violating (LNV) washout processes
- conversion to baryon asymmetry via sphaleron processes

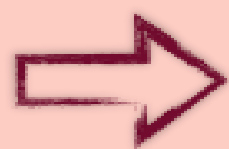
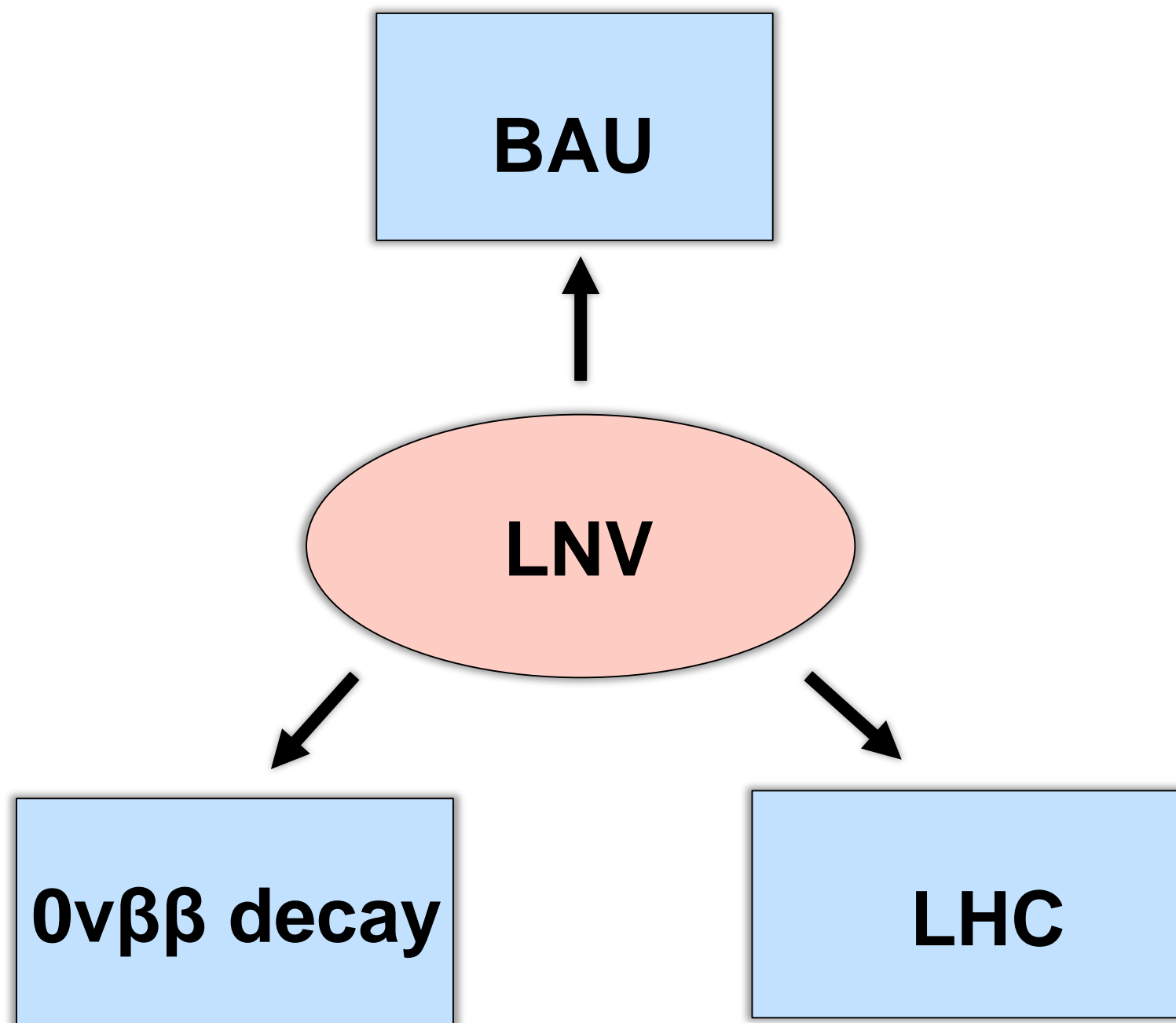


- In reverse:



- LNV processes as a probe of the strength of the washout
- probe if observable BAU still reachable

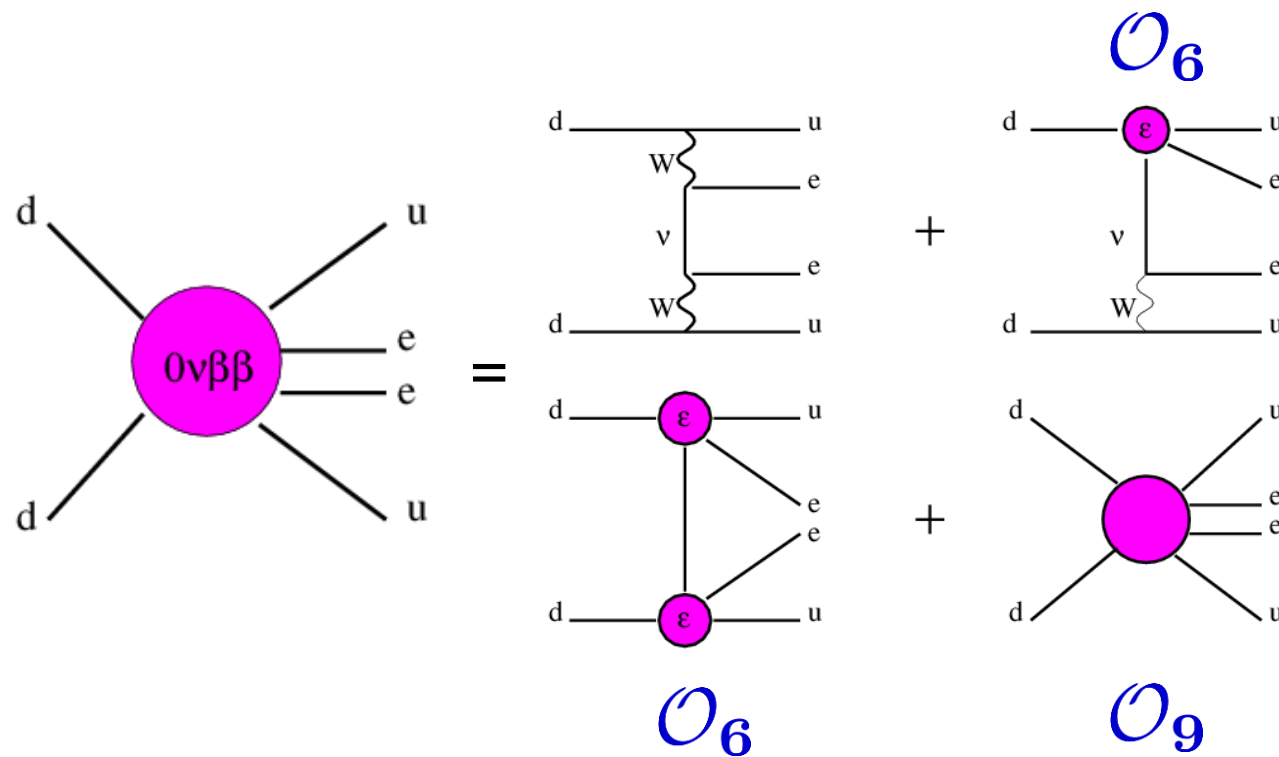




Observation of low energy LNV will have far-reaching consequences on mechanisms of baryogenesis

# Neutrinoless Double Beta Decay (0νββ)

- current limits on 0νββ:  $T_{1/2}^{76\text{Ge}} > (1.1 - 1.9) \times 10^{25}$  y EXO-200, KamLAND-Zen
- $T_{1/2}^{136\text{Xe}} > 2.1 \times 10^{25}$  y GERDA



long range contributions:

$$\mathcal{L} = \frac{G_F}{\sqrt{2}} \left\{ j_{V-A}^\mu J_{V-A,\mu}^\dagger + \sum_{\alpha,\beta} \epsilon_\alpha^\beta j_\beta J_\alpha^\dagger \right\}$$

$$J_\alpha^\dagger = \bar{u} \mathcal{O}_\alpha d \quad j_\beta = \bar{e} \mathcal{O}_\beta \nu$$

$$T_{1/2}^{-1} = |\epsilon_\alpha^\beta|^2 G_i |M_i|^2$$

similar treatment for short range contribution ...

Isotope	$ \epsilon_{V-A}^{V+A} $	$ \epsilon_{V+A}^{V+A} $	$ \epsilon_{S-P}^{S+P} $	$ \epsilon_{S+P}^{S+P} $	$ \epsilon_{TL}^{TR} $	$ \epsilon_{TR}^{TR} $
$^{76}\text{Ge}$	$3.3 \cdot 10^{-9}$	$5.9 \cdot 10^{-7}$	$1.0 \cdot 10^{-8}$	$1.0 \cdot 10^{-8}$	$6.4 \cdot 10^{-10}$	$1.0 \cdot 10^{-9}$
$^{136}\text{Xe}$	$2.6 \cdot 10^{-9}$	$5.1 \cdot 10^{-7}$	$6.2 \cdot 10^{-9}$	$6.2 \cdot 10^{-9}$	$4.4 \cdot 10^{-10}$	$7.4 \cdot 10^{-10}$

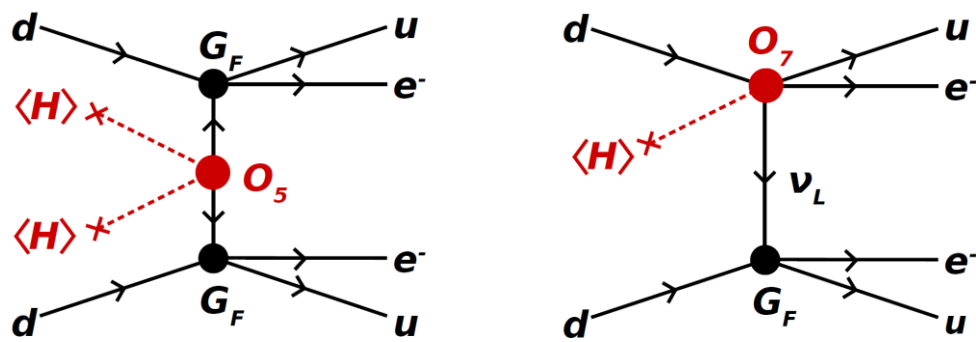
F. Deppisch, M. Hirsch, H. Päs, J. Phys. G 39 (2012) 124007, arXiv:1208.0727 [hep-ph], updated



half life sets constraints on effective coupling parameters

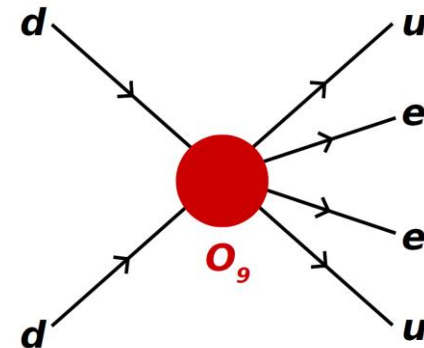
- Complete list of all LNV  $\Delta L = 2$  effective operators

K. S. Babu, C. N. Leung, Nucl. Phys. B 619 (2001), arxiv:0106054 [hep-ph]  
 A. de Gouvea, J. Jenkins, PRD 77 (2008), arXiv:0708.1344 [hep-ph]

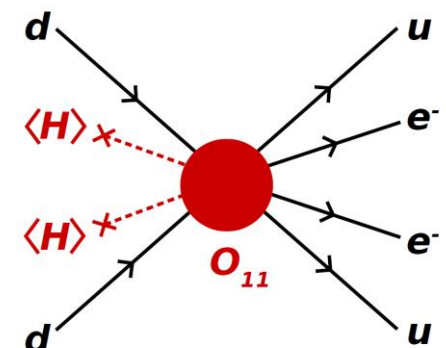


$$\mathcal{O}_5 = (L^i L^j) H^k H^l \epsilon_{ik} \epsilon_{jl}$$

$$\mathcal{O}_7 = (L^i d^c) (\bar{e}^c \bar{u}^c) H^j \epsilon_{ij}$$

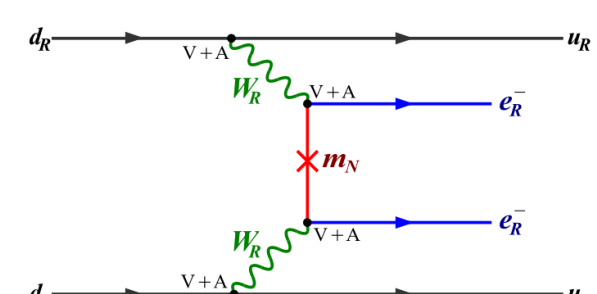
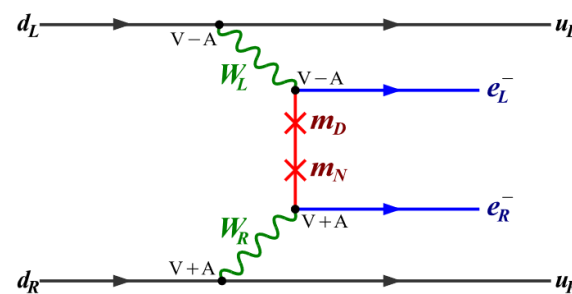
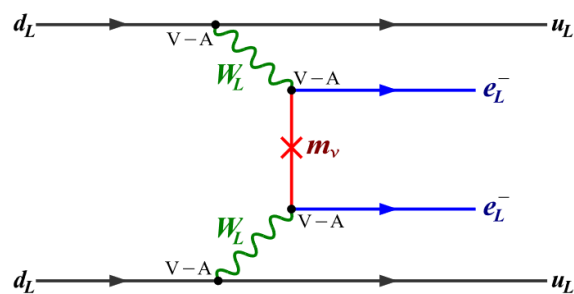


$$\mathcal{O}_9 = (L^i L^j) (\bar{Q}_i \bar{u}^c) (\bar{Q}_j \bar{u}^c)$$



$$\mathcal{O}_{11} = (L^i L^j) (Q_k d^c) (Q_l d^c) H_m \bar{H}_i \epsilon_{jk} \epsilon_{lm}$$

- Example for an UV completion: Left-right symmetric model



- If  $0\nu\beta\beta$  was observed, the scale of the underlying operator can be determined

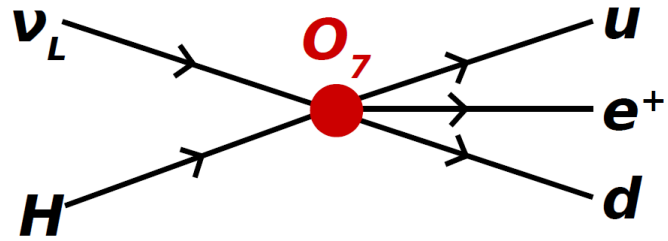
$$m_e \epsilon_5 = \frac{g^2 v^2}{\Lambda_5} \quad \frac{G_F \epsilon_7}{\sqrt{2}} = \frac{g^3 v}{2\Lambda_7^3} \quad \frac{G_F^2 \epsilon_{\{9,11\}}}{2m_p} = \left\{ \frac{g^4}{\Lambda_9^5}, \frac{g^6 v^2}{\Lambda_{11}^7} \right\}$$

$\mathcal{O}_D$	$\Lambda_D^0$ [GeV]
$\mathcal{O}_5$	$9.1 \times 10^{13}$
$\mathcal{O}_7$	$2.6 \times 10^4$
$\mathcal{O}_9$	$2.1 \times 10^3$
$\mathcal{O}_{11}$	$1.0 \times 10^3$

F. Deppisch, JH, W. Huang, M. Hirsch, H. Päs, arXiv:1503.07632 [hep-ph]



- Study washout of pre-existing net lepton asymmetry introduced by single D-dim operator



- 20 combinations of  $O_7$  to create  $2 \rightarrow 3$  and  $3 \rightarrow 2$  processes
- $1 \rightarrow 4$  phase space suppressed

$$O_7 = (L^i d^c)(\bar{e}^c \bar{u}^c) H^j \epsilon_{ij}$$

$$z H n_\gamma \frac{d\eta_N}{dz} = - \sum_{a,i,j,\dots} \left( \frac{n_N n_a \dots}{n_N^{\text{eq}} n_a^{\text{eq}} \dots} - \frac{n_i n_j \dots}{n_i^{\text{eq}} n_j^{\text{eq}} \dots} \right) \gamma^{\text{eq}} (Na \dots \leftrightarrow ij \dots)$$

$$n_\gamma H T \frac{d\eta_L}{dT} = c_D \frac{T^{2D-4}}{\Lambda_D^{2D-8}} \eta_L$$

$$\gamma^{\text{eq}} \propto \frac{T^{2D-4}}{\Lambda_D^{2D-8}}$$

$c_D$  operator specific factor

$\eta_L$  lepton density

- Washout effective if

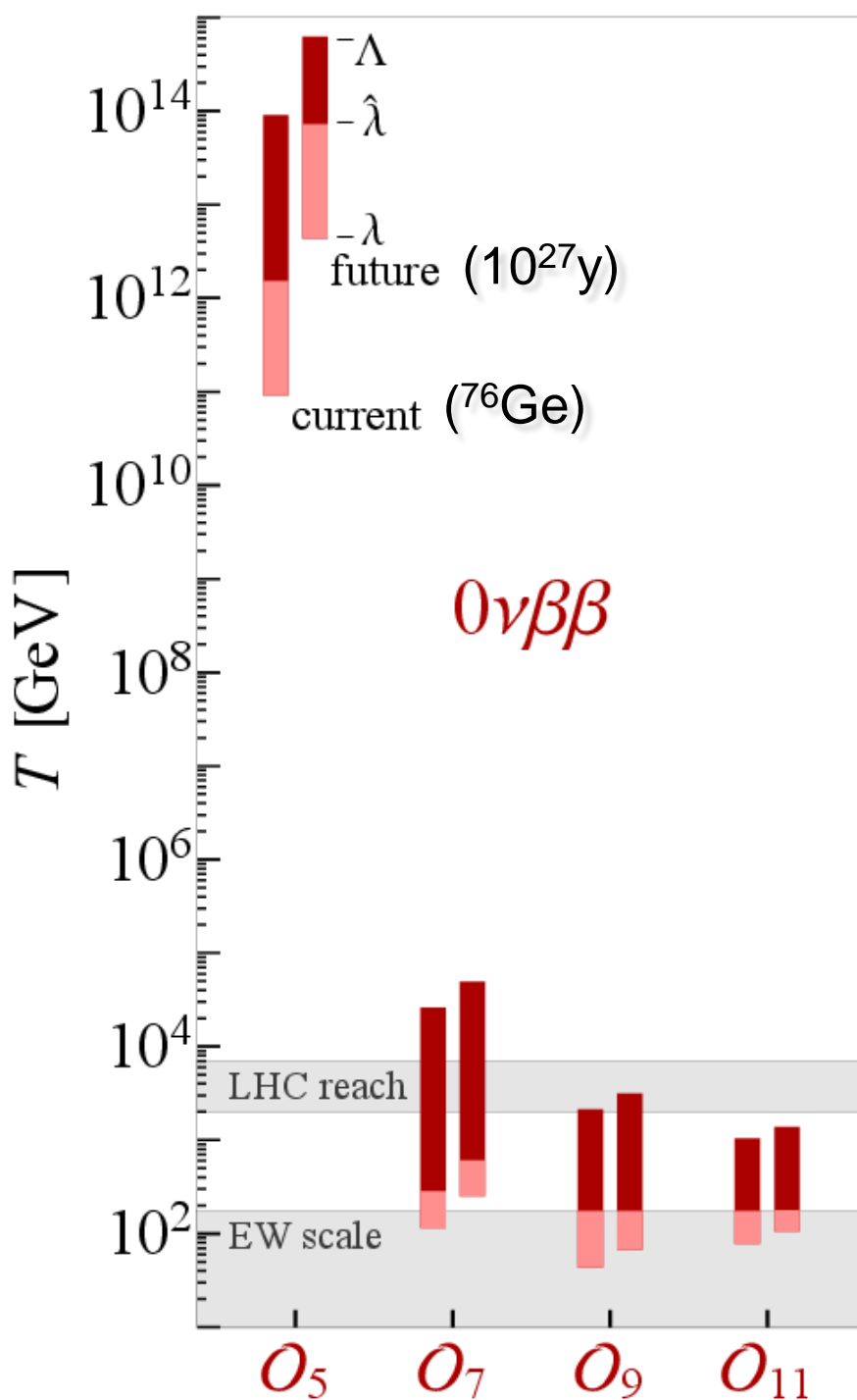
$$\frac{\Gamma_W}{H} \equiv \frac{c_D}{n_\gamma H} \frac{T^{2D-4}}{\Lambda_D^{2D-8}} = c'_D \frac{\Lambda_{\text{Pl}}}{\Lambda_D} \left( \frac{T}{\Lambda_D} \right)^{2D-9} > 1$$

- This is the case in the temperature interval

$$\Lambda_D \left( \frac{\Lambda_D}{c'_D \Lambda_{\text{Pl}}} \right)^{\frac{1}{2D-9}} \equiv \lambda_D < T < \Lambda_D$$

$O_D$	$\lambda_D^0$ [GeV]	$\Lambda_D^0$ [GeV]
$O_5$	$9.2 \times 10^{10}$	$9.1 \times 10^{13}$
$O_7$	$1.2 \times 10^2$	$2.6 \times 10^4$
$O_9$	$4.3 \times 10^1$	$2.1 \times 10^3$
$O_{11}$	$7.8 \times 10^1$	$1.0 \times 10^3$

F. Deppisch, JH, W. Huang, M. Hirsch, H. Päs, arXiv:1503.07632 [hep-ph]



- $\Lambda$  scale of operator
- $\lambda$  scale above which washout highly effective  $\frac{\Gamma_W}{H} > 1$
- $\hat{\lambda}$  scale above which a max. lepton asymmetry of 1 is washed out to  $\eta_B^{\text{obs}}$  or less

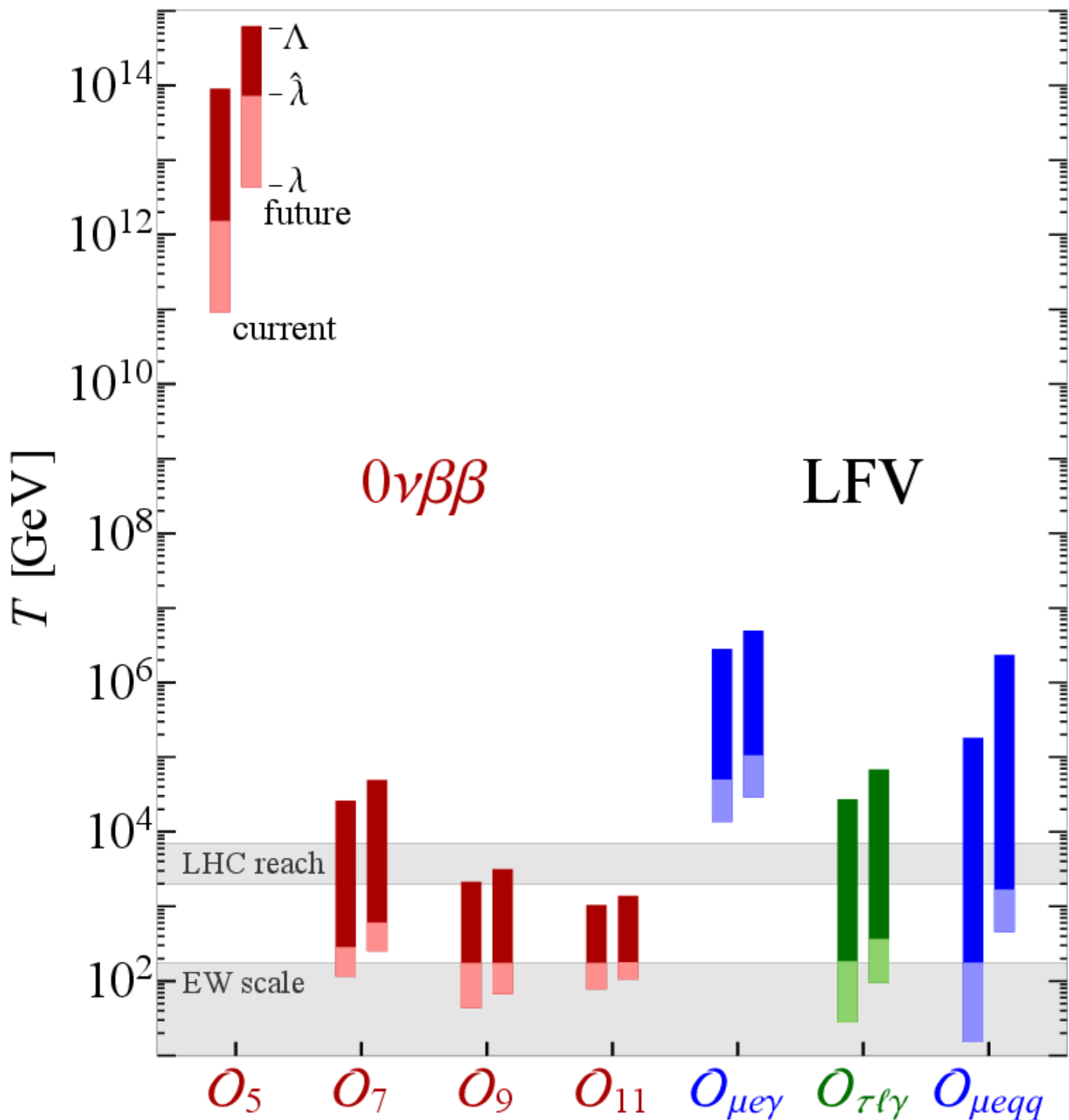
$$\hat{\lambda}_D \approx \left[ (2D - 9) \ln \left( \frac{10^{-2}}{\eta_B^{\text{obs}}} \right) \lambda_D^{2D-9} + v^{2D-9} \right]^{\frac{1}{2D-9}}$$

• IF  $0\nu\beta\beta$  was observed via non-standard mechanism, resulting washout would rule out baryogenesis mechanisms above  $\lambda$

• observation of  $0\nu\beta\beta$  via  $O_9$  and  $O_{11}$  will imply observation of LNV at LHC

•  $0\nu\beta\beta$  decay probes only electron-electron component of LNV operators

$$\frac{1}{\Lambda_9^5} \rightarrow \frac{c_{\alpha\beta}}{\Lambda_9^5}$$



- Most stringent limits on LFV set by 6-dim  $\Delta L = 0$  operators

$$\mathcal{O}_{ll\gamma} = C_{ll\gamma} \bar{L}_\ell \sigma^{\mu\nu} \bar{\ell}^c H F_{\mu\nu}$$

$$\mathcal{O}_{llqq} = C_{llqq} (\bar{\ell} \Pi_1 \ell) (\bar{q} \Pi_2 q)$$

$$C_{llqq} = \frac{g^2}{\Lambda_{llqq}^2} \quad C_{ll\gamma} = \frac{eg^3}{16\pi^2 \Lambda_{ll\gamma}^2}$$

- Current & future limits:

$$\text{Br}_{\mu \rightarrow e \gamma} < 5.7 \times 10^{-13} \quad (6.0 \times 10^{-14})$$

$$\text{Br}_{\tau \rightarrow l \gamma} < 4.0 \times 10^{-8} \quad (1.0 \times 10^{-9}), \quad l = e, \mu$$

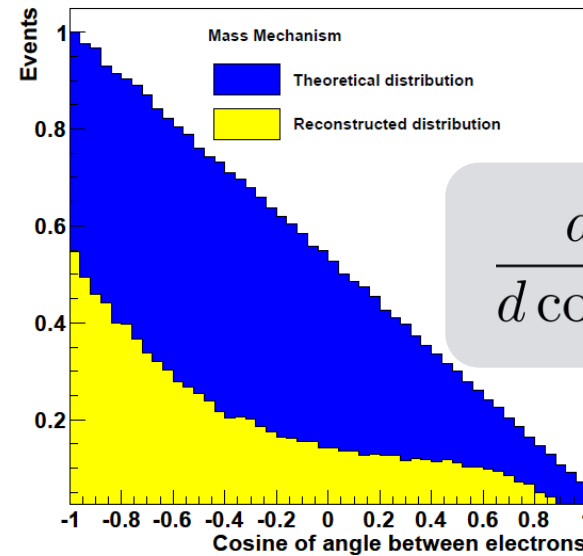
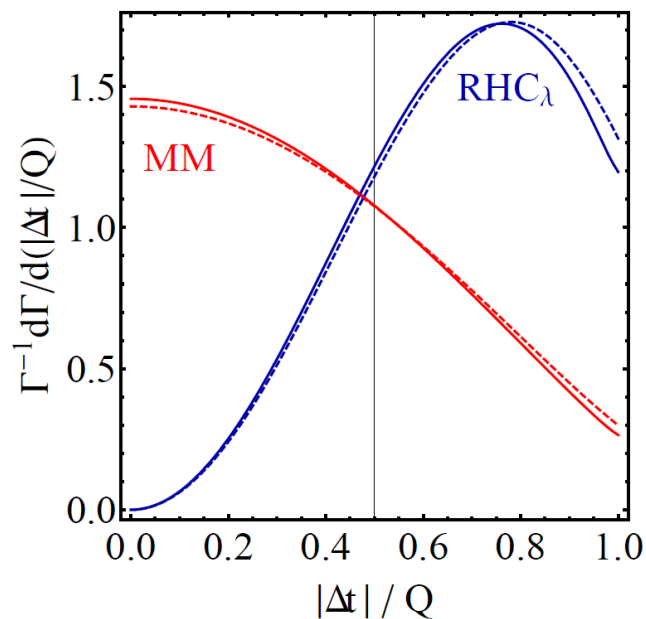
$$R_{\mu \rightarrow e}^{\text{Au}} < 7.0 \times 10^{-13} \quad (2.7 \times 10^{-17})$$

- determine temperature interval in which LFV process equilibrate pre-existing flavour asymmetry

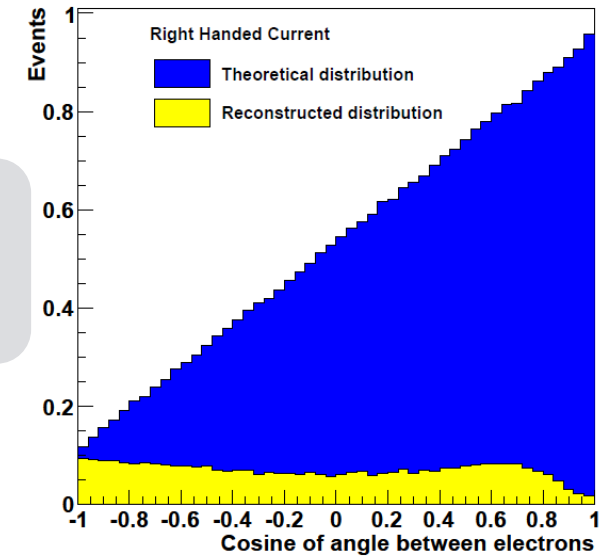
- **IF** LFV processes are observed as well, loophole of asymmetry being stored in another flavour sector is ruled out



- SuperNEMO can discriminate  $O_7$  from other mechanisms, due to  $e^-_R$  and  $e^-_L$  in final state

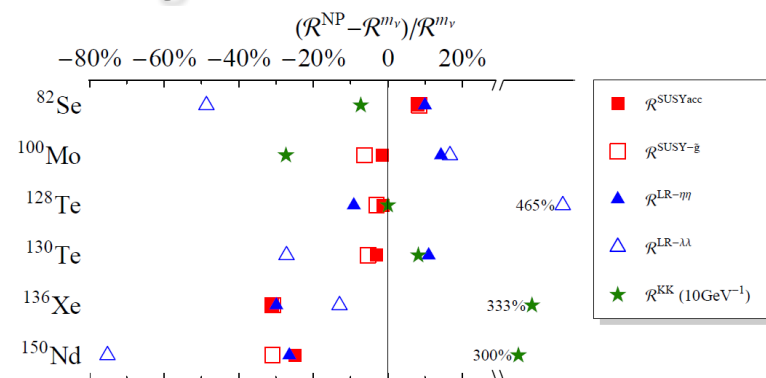


$$\frac{d\Gamma}{d \cos \theta_{12}} = \frac{\Gamma}{2} (1 - k_\theta \cos \theta_{12})$$



SuperNEMO collaboration, arXiv:1005.1241 [hep-ex]

- Potential discrepancy between neutrino mass (cosmology) and  $0\nu\beta\beta$  half live measurement could be an indication for  $0\nu\beta\beta$  being triggered by non-standard mass mechanism
- Distinguishing between different mechanisms via measurements in different isotopes

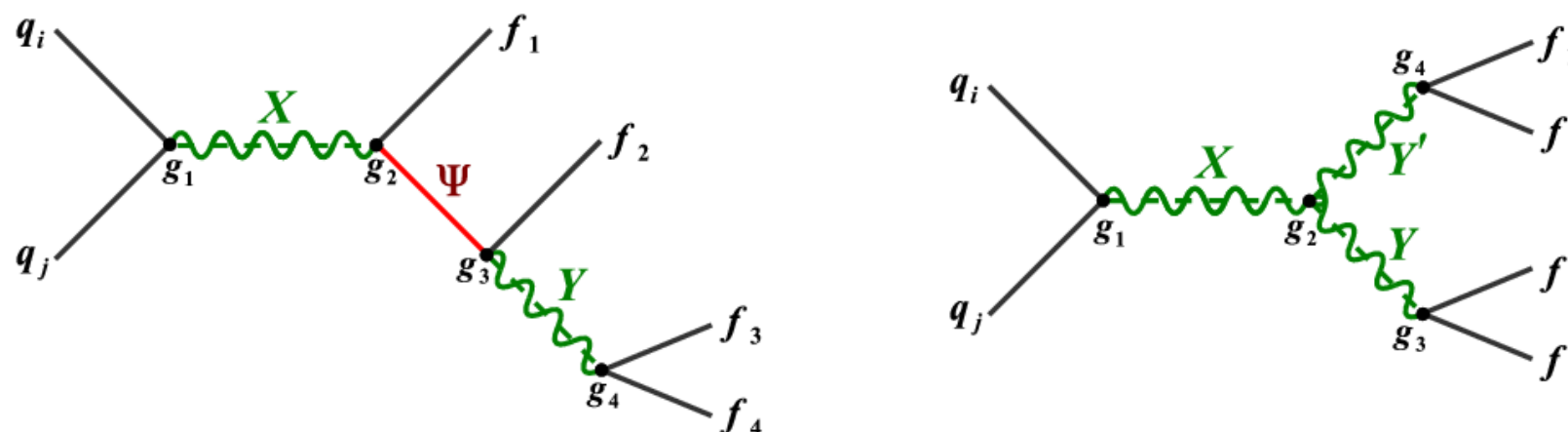


$$\frac{T_{1/2}(^A X)}{T_{1/2}(^A X)} = \frac{|\mathcal{M}(^{76}Ge)|^2 G(^{76}Ge)}{|\mathcal{M}(^A X)|^2 G(^A X)}$$

Deppisch, Paes, PRL 98 (2007)  
Gehmann, Elliott, J. Phys G 34 (2007)

- Comparison of  $0\nu\beta^-\beta^-$  with  $0\nu\beta^+\beta^+$  Hirsch, Muto, Oda, Klapdor-Kleingrothaus, Z. Phys A347 (1994)
- observation of  $0\nu\beta\beta$  via  $O_9$  and  $O_{11}$  will imply observation of LNV at LHC

- Signature:**  $\Delta L = 2$  LNV at LHC through resonant process  $pp \rightarrow l^\pm l^\pm + 2 \text{ jets}$  with two same-sign leptons and two jets without missing energy



$$\frac{\Gamma_W}{H} = \frac{1}{n_\gamma H} \frac{T}{32\pi^4} \int_0^\infty ds s^{3/2} \sigma(s) K_1\left(\frac{\sqrt{s}}{T}\right) \quad \sigma(s) = \frac{4 \cdot 9 \cdot s}{f_{q_1 q_2} (M_X / \sqrt{s})} \sigma_{\text{LHC}}$$

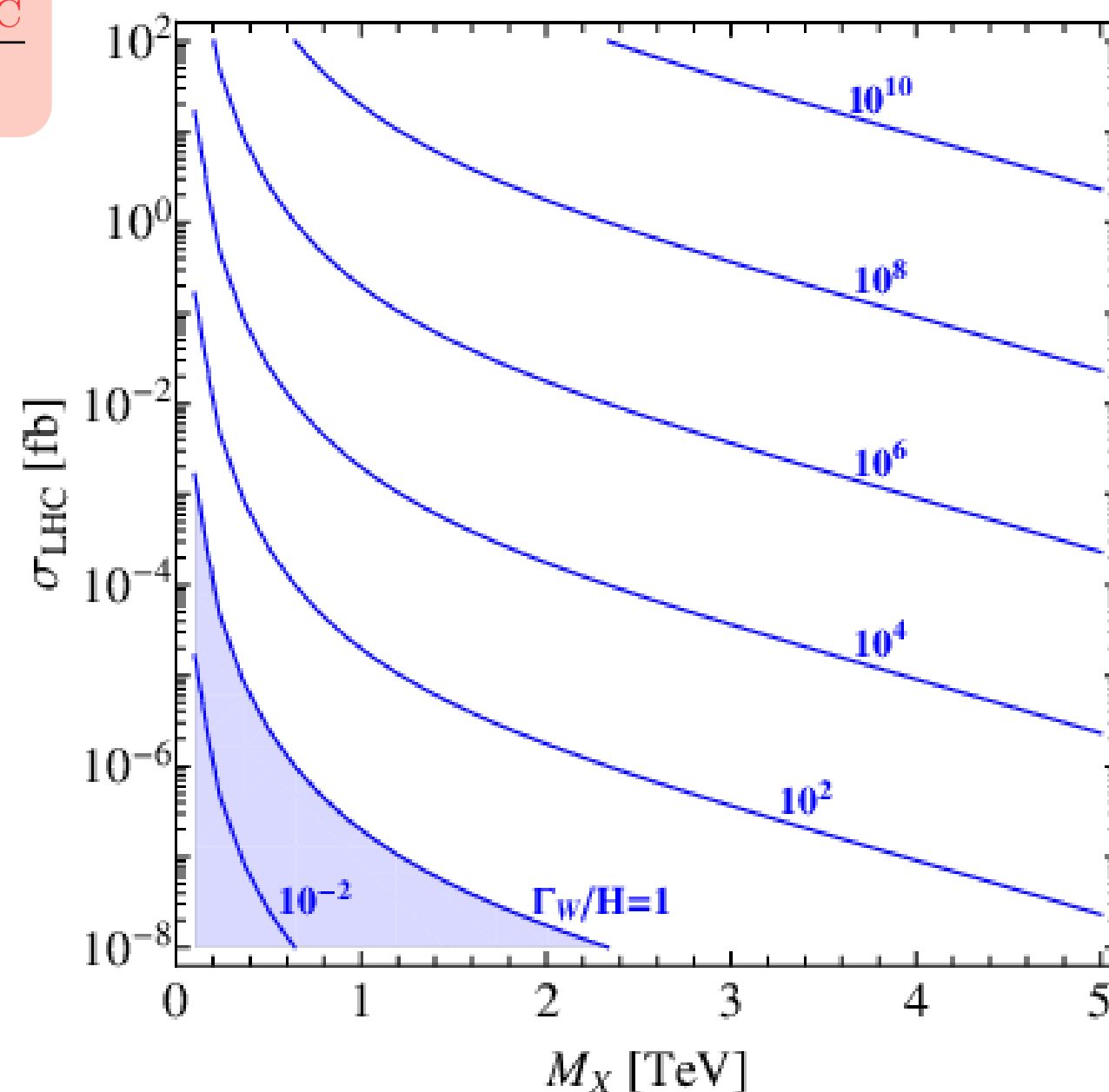
$$\frac{\Gamma_W}{H} = \frac{0.028}{\sqrt{g_*}} \frac{M_{\text{P}} M_X^3}{T^4} \frac{K_1(M_X/T)}{f_{q_1 q_2} (M_X / \sqrt{s})} \times (s \sigma_{\text{LHC}})$$

$$\log_{10} \frac{\Gamma_W}{H} > 6.9 + 0.6 \left( \frac{M_X}{\text{TeV}} - 1 \right) + \log_{10} \frac{\sigma_{\text{LHC}}}{\text{fb}}$$

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- For any realistic cross section at LHC with  $\sigma_{\text{LHC}} > 10^{-2} \text{ fb}$  washout highly effective

$$\frac{\Gamma_W}{H} \gg 1$$



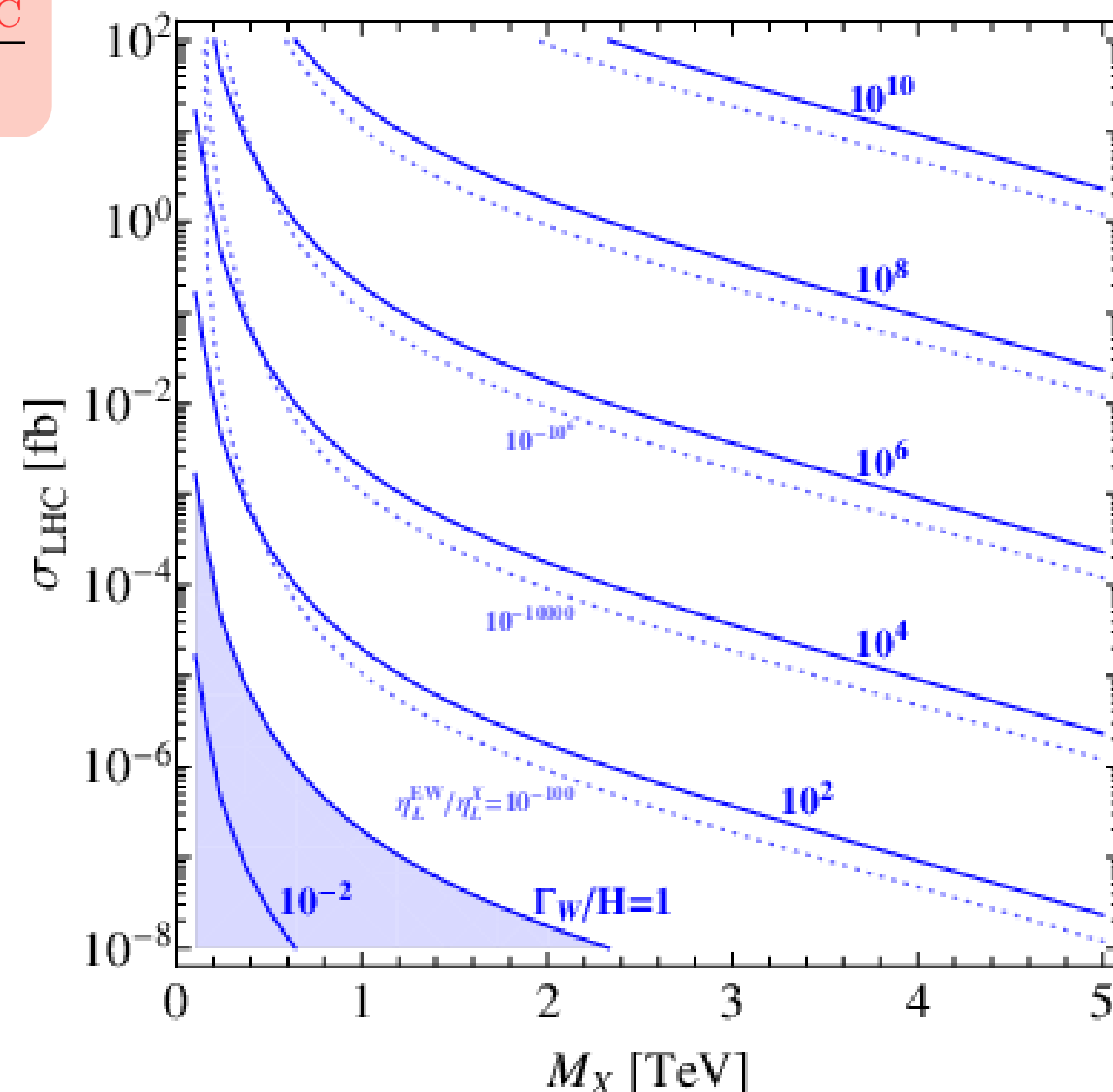
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- enormous washout of any pre-existing lepton asymmetry

$$\eta_L^{\text{EW}} / \eta_L^X \approx \exp(-\Gamma_W / H)$$



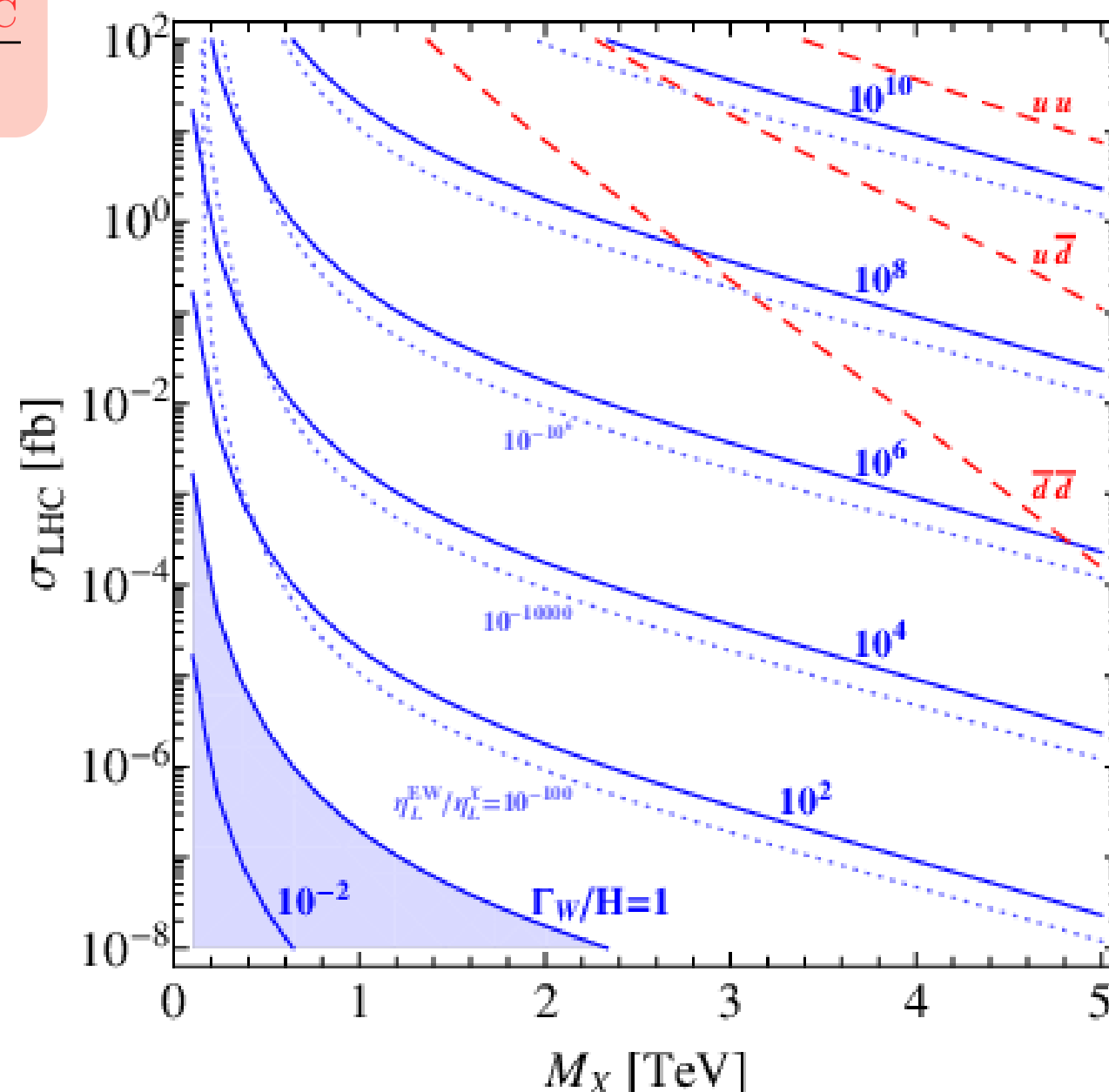
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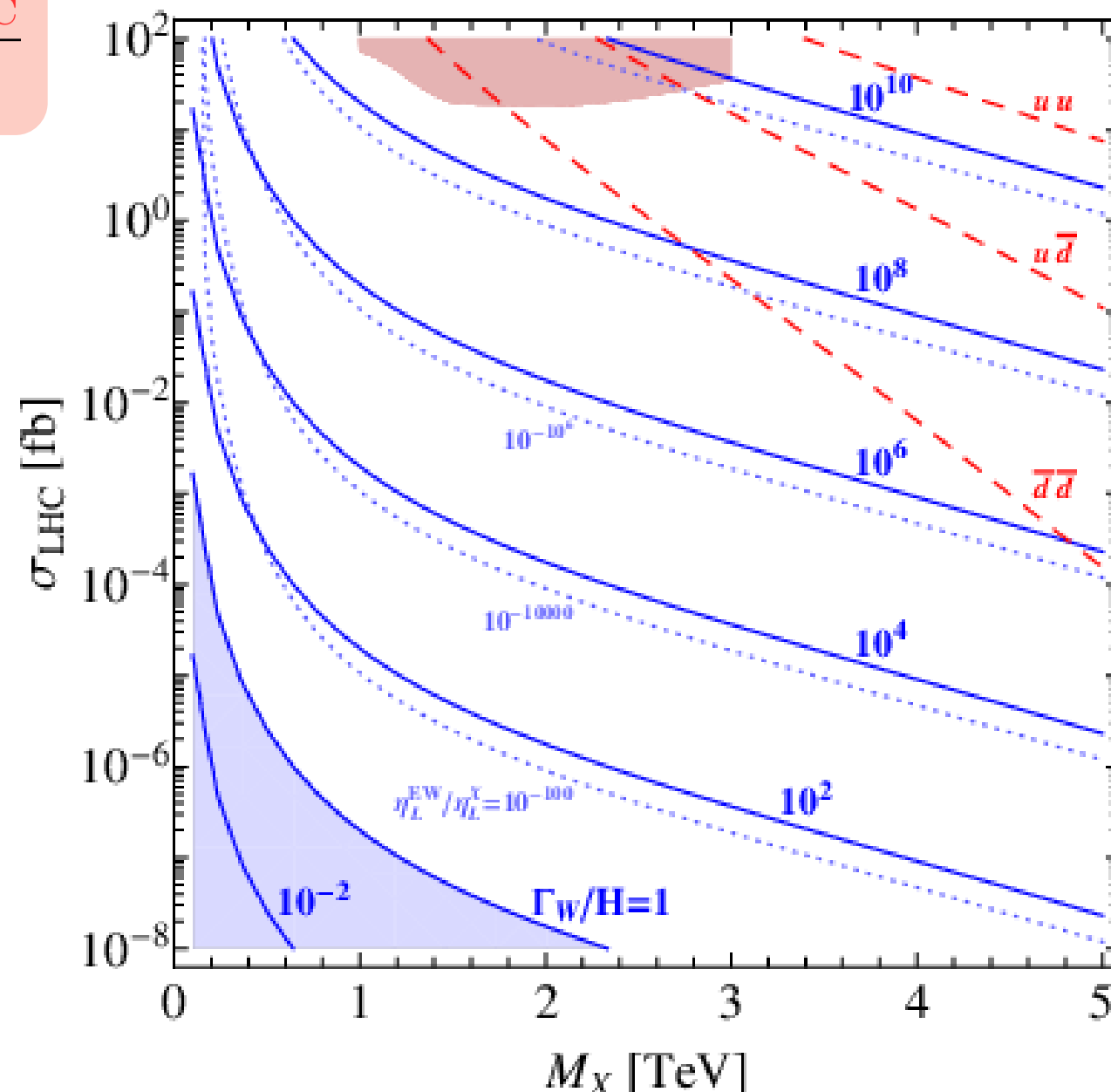
- enormous washout of any pre-existing lepton asymmetry

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- LHC starts to exclude top of parameter plane

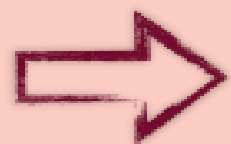
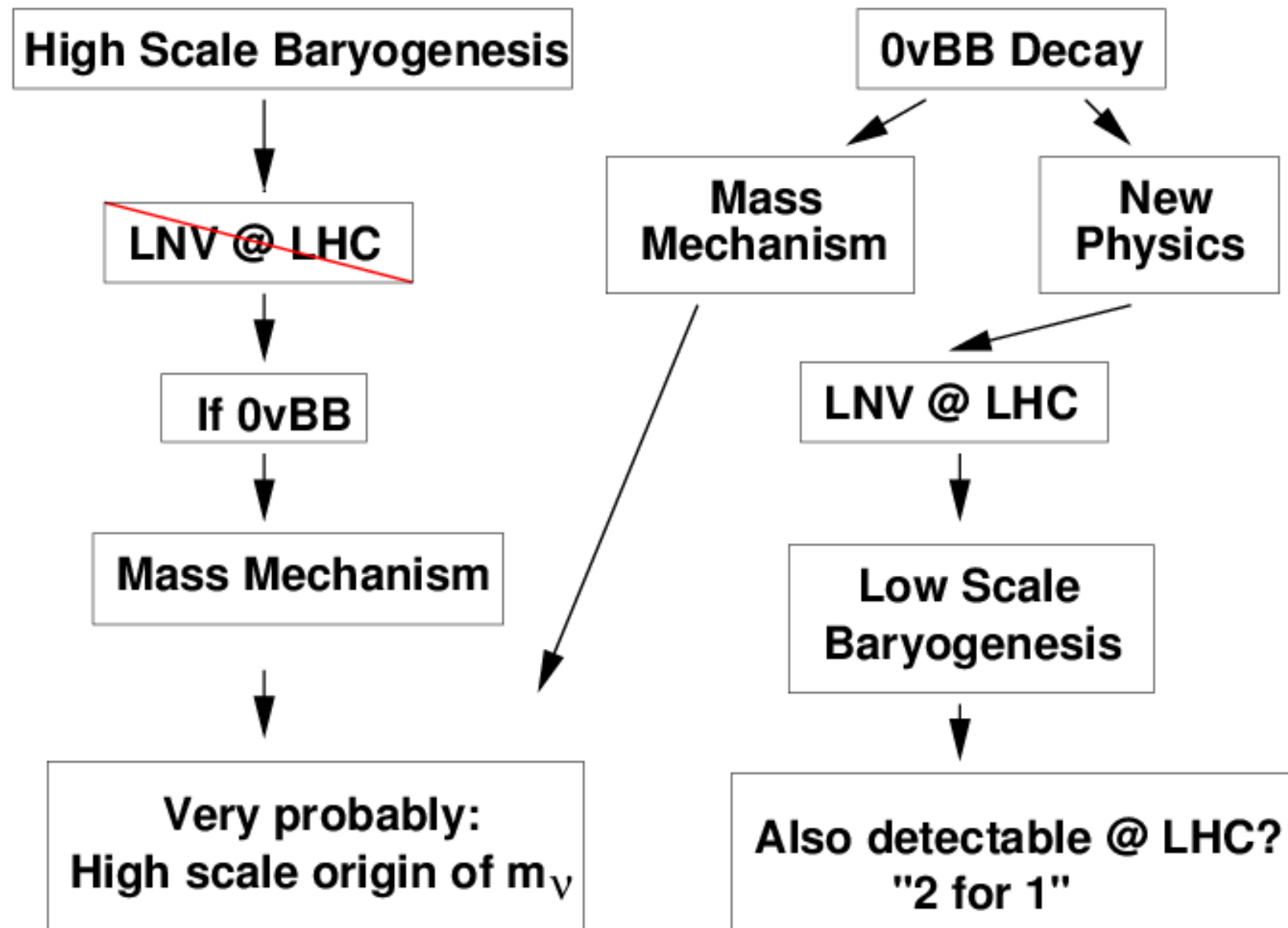


- observation of LNV processes sets serious bounds on washout
- excludes LG models which generate asymmetry above



F. Deppisch, JH, M. Hirsch, PRL 112 (2014) 221601, arXiv:1312.4447 [hep-ph]

- LNV process at LHC involves right-handed leptons, but SM sphaleron processes only affect EW fermion doublets
  - left- and right-handed fermions are in thermal equilibrium around EW-scale
- Possible generation of LNV only in one flavour family
  - observation of same-sign signatures in different flavours
  - observation of LFV processes
- LNV models with new conserved quantum numbers or hidden sectors may be exempt
  - S. Weinberg, PRD 22 (1980)
  - A. Antaramian, L. Hall, A. Rasin, PRD 49 (1994), arXiv:9311279 [hep-ph]
- Baryon asymmetry could be generated below the EW scale



observation of low energy LNV processes (e.g. in  $0\nu\beta\beta$  or LHC) can washout any pre-existing baryon asymmetry irrespective of the baryogenesis mechanism

THANK YOU!

