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### Constraining Dipolar Gravitational Radiation through Pulsar Timing

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- Pulsar Timing
- Tensor-Scalar Theories: J1141-6545
- Generic Tests for Dipole Radiation
  - Introduction
  - Why not 1141?
  - J1012+5307 & J1738+0333
- Anticipated Results

## **Pulsar Timing**





### PSR J1012+5307: 15 years of observations with EPTA

 $P = 5.255749014115410 \pm 0.000000000000015 \text{ ms}$  $P_{\text{orb}} = 52243.722450 \pm 0.000003 \text{ s}$ 

[Lazaridis et al. 2009]



Champion et al., 2010

## The PSR J1141-6545 System

- Pulse period
   P ~ 400 ms
- Binary period
   P<sub>b</sub> ~ 4.7 hr
- Ecc ~ 0.17
- Dist > 3.7 kpc
- WD-NS binary (Antoniadis et al., 2010)



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Right Ascension (J2000)

### **Observable** Relativistic Effects

Periastron Advance:  $\dot{\omega} \propto \frac{M_{Tot}^{2/3}}{(P_b^{5/3})(1-e^2)}$ 

**Shapiro Delay (shape):**  $s \propto \frac{M_{Tot}^{2/3} a \sin i}{P_b^{2/3} M_c}$ 

Orbital Period Decay:  $\dot{P}_b \propto \frac{M_{psr}M_c}{M_{Tot}^{1/3}P_b^{5/3}}(1+\frac{73}{24}e^2+\frac{37}{96}e^4)$ Gravitational Redshift:  $\gamma \propto P_b^{1/3}e^2 \frac{M_c(M_{Tot}+M_c)}{M_{Tot}^{4/3}}$ 

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Determined through scintillation studies (Ord et al., 2002)

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### Mass-Mass Diagram



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### **Tensor-Scalar Theories**

- Amongst the most well-founded alternatives to GR
- Coupling strength btw. scalar field & matter, e.g. quadratic:  $a(\psi) = \alpha_0 \psi + 0.5 \beta_0 \psi^2$ ; GR:  $(\alpha_0; \beta_0) = (0;0)$ JFBD:  $(\alpha_0; \beta_0) = (\alpha_0;0)$
- Predict emission of dipolar GWs, depending on difference in compactness ( $\Delta\epsilon$ ) of objects
- J1141-6545: WD (ε~10-4) NS (ε~0.2)



# Esposito-Farèse (2009) arXiv:0905.2575; Data: Bhat, Bailes & Verbiest (PRD, 2008)

# Figure:





Esposito-Farèse (2009) arXiv:0905.2575; Data: Bhat, Bailes & Verbiest (PRD, 2008)

Figure:

**Forecas** SR J141-6545



### Problems with J1141-6545

- No theory-invariant mass estimates
- Not applicable in generic frameworks
- Practical problems: geodetic precession, glitch, unknown proper motion, ...

# **Generic Limits on dipolar GW Emission**

- Dipolar GW emission is predicted by most alternative gravitational theories
- Strongly dependent on difference in effective scalar coupling:

$$\dot{P}_{\rm b}^{\rm dipole} = -\frac{4\pi^2}{P_{\rm b}} \frac{G\mu}{c^3} \frac{1+e^2/2}{(1-e^2)^{5/2}} \left(\alpha_{\rm p} - \alpha_{\rm c}\right)^2$$
Negligible if NS-NS Substantial if NS-WI

### Generic Test with PSR J1012+5307



### Combined Limits for a Variation in G and **Dipole Radiation**

$$\frac{\dot{P}_{b}^{exc}}{P_{b}} = -2\frac{\dot{G}}{G} \left[ 1 - \left( 1 + \frac{m_{c}}{2M} \right) s_{p} \right] - 4\pi^{2} \frac{T_{\odot} \mu}{P_{b}^{2}} \kappa_{D} s_{p}^{2} \right]$$

$$\frac{PSR J1012 + 5307}{Orbital \text{ period}} = 0.605 \text{ days}$$

$$\frac{PSR J0437 - 4715}{Orbital \text{ period}} = 5.75 \text{ days}$$

[Lazaridis et al. 2009]

 $\dot{G}/G \; [10^{-12} \, {\rm yr}^{-1}]$ 

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### PSR J1738+0333

### Timing observations of the pulsar

[Freire et al. 2011]





 $dP_{b}/dt = (-3 \pm 1) \times 10^{-14}$ 

### **Optical observations of the WD companion**

[Antoniadis et al. 2011]







 $m_{\rm c} = 0.178 + 0.015 / -0.009 \ {\rm M}_{\rm sun}$ 

 $m_{\rm p}/m_{\rm c} = 8.20 \pm 0.19$ 

### PSR J1738+0333 - GW Limit



## TeVeS Limits from J1738+0333



### Summary & Outlook

 Generic Strong-Field Tests of Gravity with Binary Pulsar Systems

 PSR J1141-6545: Quadratic Tensor-Vector Theories

 PSR J1738+0333: TeVeS Theories — Must be "unnatural"

# **Additional Slides**

### "Natural" MOND and Binary Pulsars



If  $\alpha_0^2 < 0.002$  MOND effects should be apparent in the solar system dynamics

 $\alpha_0^2 > 0.002$  would lead to a significant emission of dipolar gravitational radiation in pulsar-white dwarf binaries with ~1 day orbital period:

$$\dot{P}_{\rm b}^{\rm dipole} = -\frac{4\pi^2}{P_{\rm b}} \frac{G\mu}{c^3} \frac{1+e^2/2}{\left(1-e^2\right)^{5/2}} \left(\alpha_0 s_{\rm p}\right)^2$$

### **TeVeS Becomes Unnatural**

 $\alpha_0^2 = 10^{-4}$  requires an extremely unnatural behaviour of f(x):



[Bruneton & Esposito-Farese 2007]

If  $\alpha_0^2 \sim 10^{-5}$  stability and consistency conditions violated.

The vector field  $U_{\mu}$  has to be parallel to the matter proper time, in order to prevent preferred frame effects. This leads to additional gravitational wave damping. Therefore constraints on TeVeS are most likely even tighter.

### TeVeS Theories and Binary Pulsar Experiments

MOND (MOdified Newtonian Dynamics) is not a relativistic theory of gravity (cannot explain lensing, cosmology, etc.)

To be an alternative to GR, a relativistic theory is needed that

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- Has MOND as a weak-field, slow motion limit
- Can explain existing experimental and observational results

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Most promising attempt so far: tensor-vector-scalar (TeVeS) theories (Bekenstein). With sufficient fine-tuning (free function, free parameters) TeVeS theories can pass solar system tests

- Post-Newtonian dynamics of the planets
- Light-deflection and light-propagation effects
- Preferred-frame effect (PFE) tests

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However, general considerations by Bruneton & Esposito-Farese (2007) already show, that the scalar coupling used to explain solar system tests ( $\alpha^2 = 2 \times 10^{-3}$ ) is in stark contrast to gravitational wave tests with binary pulsars: ( $\alpha^2 < 2 \times 10^{-4}$ ).

Further, there is an issue with PFEs due to the vector field. If the vector field is adjusted to the matter frame to prevent PFEs, it would lead to <u>additional GW damping</u> in binary pulsars.