Exercises to "Standard Model of Particle Physics II"

Winter 2018/19

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Sheet 8

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Problem 16: Running couplings [10 Points]

a) In QED the coupling $\alpha = e^2/4\pi$ is a running coupling. For low energies one obtains $\alpha(\mu = m_e) = 1/137$. The running of the QED coupling with only one lepton is given by

$$\alpha_{\text{QED}}(\mu) = \frac{\alpha(m_e)}{1 - \frac{2\alpha(m_e)}{3\pi} \ln(\frac{\mu}{m_e})} \ . \tag{1}$$

At which scale does $\alpha_{\text{QED}}(\mu)$ become infinite? (If further charged leptons and quarks are considered, the coupling increases even faster.)

b) In QCD, with $\alpha_s = g_3^2/4\pi$, the running coupling

$$\alpha_s(\mu) = \frac{\alpha_s(\mu_0)}{1 + \frac{(33 - 2f)\alpha_s(\mu_0)}{6\pi} \ln(\frac{\mu}{\mu_0})}$$
(2)

is obtained, where f denotes the respective number of quark flavours with mass $2m_q \leq \mu$ in the considered energy range.

The experimental boundary conditions are $\alpha_s(\mu_0 = m_Z = 91 \text{ GeV}) = 0.12$, $m_t = 175 \text{ GeV}$, $m_b = 4.8 \text{ GeV}$, $m_c = 1.4 \text{ GeV}$ and $m_u \simeq m_d \simeq m_s \simeq 0$.

- i) Determine the pole $\mu = \Lambda_{\rm OCD}$ of eq. (2).
- ii) For which μ does the coupling $\alpha_s(\mu)$ become very small (asymptotic freedom), and where does perturbation theory break down?
- iii) Determine the value of $\alpha_s(\mu)$ in the different energy ranges $(2m_q \leq \mu \leq 2m_t \text{ etc.})$ at the thresholds.
- iv) $\alpha_s(\mu)$ should be continuous. However, from the threshold values calculated in the last part we saw that this is not yet true. To make $\alpha_s(\mu)$ continuous, promote μ_0 to a function $\mu_0(f)$ with $\mu_0(f=5)=91$ GeV. Find a relation between $\mu_0(f)$ and $\mu_0(f-1)$ assuming that $\alpha_s(\mu_0(f))=0.12$ for all f. With this relation determine the values of $\mu_0(4)$ and $\mu_0(6)$.
- c) Draw $\alpha_s^{-1}(\mu)$ and $\alpha_{\text{QED}}^{-1}(\mu)$ as functions of $\ln(\mu)$. What could the intersection of the curves indicate?

Problem 17: Integrating out the W-boson [10 Points]

Heavy particles decouple from a given theory when the relevant energy scale of a process is well below their mass. An important example are the W^{\pm} bosons of the SM, which we will consider in this exercise. The relevant part of the SM Lagrangian reads

$$\mathcal{L}_W = -\frac{1}{2} W^-_{\mu\nu} W^{+\mu\nu} + m_W^2 W^+_{\mu} W^{-\mu} + \frac{g}{\sqrt{2}} W^+_{\mu} J^{-\mu} + \frac{g}{\sqrt{2}} W^-_{\mu} J^{+\mu} ,$$

with the electroweak charged currents acting as source terms,

$$J^{+}_{\mu} = \overline{e}_L \gamma_{\mu} \nu_L, \quad J^{-}_{\mu} = \overline{\nu}_L \gamma_{\mu} e_L.$$

Derive the classical equation of motion for the W^\pm bosons and expand to lowest order in $\frac{p}{m_W}$, where p is the momentum of the W, i.e. $\partial_\mu W^\pm_{\ \nu} = -ip_\mu W^\pm_{\ \nu}$. Solve the equation to lowest order in $\frac{p}{m_W}$ ignoring non-linear interactions among the W's. Show that this solution can be used to obtain the Lagrangian of Fermi's effective theory of weak interactions,

$$\mathcal{L}_{\text{Fermi}} = \frac{G_F}{\sqrt{2}} J^+_{\ \mu} J^{-\mu}$$

and determine G_F in terms of the W mass, m_W , and the weak gauge coupling, g.

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Lecture webpage: https://www.mpi-hd.mpg.de/manitop/StandardModel2/index.html

Hand-in and discussion of sheet:

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