Prof. Dr. Manfred Lindner and Dr. Werner Rodejohann Sheet 12 January 23, 2019

Problem 24: Left-right symmetric electroweak model [10 Points]

The left-right symmetric model can be introduced by assuming right-handed fermion doublets in analogy to the left-handed ones. The quark and lepton spectra consist of

$$Q_{\mathrm{L, R}}^{i} = \begin{pmatrix} U_{\mathrm{L, R}}^{i} \\ D_{\mathrm{L, R}}^{i} \end{pmatrix}; \qquad L_{\mathrm{L, R}}^{i} = \begin{pmatrix} \nu_{\mathrm{L, R}}^{i} \\ e_{\mathrm{L, R}}^{i} \end{pmatrix}$$

with the following $SU(2)_L$, $SU(2)_R$, $U(1)_{B-L}$ transformation properties:

$$Q_{L}: (2_{L}, 1_{R}, 1/3); \qquad L_{L}: (2_{L}, 1_{R}, -1); Q_{R}: (1_{L}, 2_{R}, 1/3); \qquad L_{R}: (1_{L}, 2_{R}, -1).$$

The Higgs sector contains a bi-doublet ϕ and two triplets Δ_L and Δ_R with transformation properties

 $\phi: (2_{\mathrm{L}}, 2_{\mathrm{R}}, 0); \qquad \Delta_{\mathrm{L}}: (3_{\mathrm{L}}, 1_{\mathrm{R}}, 2); \qquad \Delta_{\mathrm{R}}: (1_{\mathrm{L}}, 3_{\mathrm{R}}, 2).$

Note that the bi-doublet and the triplets can be expressed by the following 2×2 matrices:

$$\phi = \begin{pmatrix} \phi_1^0 & \phi_2^+ \\ \phi_1^- & \phi_2^0 \end{pmatrix}; \qquad \Delta_{\mathrm{L, R}} = \begin{pmatrix} \Delta^+/\sqrt{2} & \Delta^{++} \\ \Delta^0 & -\Delta^+/\sqrt{2} \end{pmatrix}$$

- a) Why does this model not work with only the bi-doublet?
- b) Construct the Lagrange density for the fermion-Higgs interactions \mathscr{L}_{Yukawa} (including all possible gauge singlets).
- c) Use the assumption that after spontaneous symmetry breaking the vacuum is electrically neutral to derive the fermion mass terms in the broken phase.
- d) We leave the quark and lepton sector unchanged, but modify the symmetry-breaking part of the model. In the Higgs sector we still have the bi-doublet ϕ , but instead of the triplets we introduce two scalar doublets $A_{L, R}$ and a fermionic singlet χ with the following transformation properties under SU(2)_L, SU(2)_R and U(1)_{B-L}:

$$A_{\rm L}: (2_{\rm L}, 1_{\rm R}, 1); \qquad A_{\rm R}: (1_{\rm L}, 2_{\rm R}, 1); \qquad \chi: (1_{\rm L}, 1_{\rm R}, 0).$$

Left-right symmetry implies the invariance of the Lagrange density under the following transformations (where Ψ denotes any fermion field):

$$\Psi_{\rm L} \leftrightarrow \Psi_{\rm R} \qquad \qquad A_{\rm L} \leftrightarrow A_{\rm R} \qquad \qquad \phi \leftrightarrow \phi^{\dagger}.$$

Construct the Lagrange density \mathscr{L}_{Yukawa} for the fermion masses in this model. (Again, you have to construct singlets under the whole gauge group.)

Problem 25: Z' physics [10 Points]

The general effective Lagrange density after breaking the $SU(3)_C \times SU(2)_L \times U(1)_Y \times U(1)'$ symmetry to $SU(3)_C \times U(1)_{em}$ can be written as

$$\mathscr{L} = \mathscr{L}_{\mathrm{SM}} + \mathscr{L}_{Z'} + \mathscr{L}_{\mathrm{mix}},$$

where the relevant part of the Standard Model Lagrangian is

$$\mathscr{L}_{\rm SM} = -\frac{1}{4}\hat{B}_{\mu\nu}\hat{B}^{\mu\nu} - \frac{1}{4}\hat{W}^a_{\mu\nu}\hat{W}^{\mu\nu,a} + \frac{1}{2}\hat{M}^2_Z\hat{Z}_\mu\hat{Z}^\mu - \frac{e}{c_W}j^\mu_B\hat{B}_\mu - \frac{e}{s_W}j^{\mu,a}_W\hat{W}^a_\mu$$

and the hats merely denote that the fields are not mass eigenstates. The Z' part reads

$$\mathscr{L}_{Z'} = -\frac{1}{4} \hat{Z'}_{\mu\nu} \hat{Z'}^{\mu\nu} + \frac{1}{2} \hat{M}_{Z'}^2 \hat{Z'}_{\mu} \hat{Z'}^{\mu} - g' j'_{\mu} \hat{Z}'_{\mu} ,$$

where g' denotes the U(1)' gauge coupling, and the kinetic- and mass-mixing terms can be parameterized as

$$\mathscr{L}_{\rm mix} = -\frac{\sin\chi}{2}\hat{Z}'_{\mu\nu}\hat{B}^{\mu\nu} + \delta\hat{M}^2\hat{Z}'_{\mu}\hat{Z}^{\mu}$$

- a) Determine the mass eigenstates Z_1^{μ} and Z_2^{μ} and determine the couplings of $Z_{1,2}$ to the currents j_B , j_W and j'. Set the kinetic mixing angle χ to zero for simplicity. Hint: Reexpress \hat{B}_{μ} and \hat{W}_{μ}^3 in terms of A_{μ} and Z_{μ} .
- b) Since the mass of the physical Z boson changes compared to the SM, the ρ parameter is no longer equal to one (at tree-level). Use the current value $\rho = 1.0008^{+0.0017}_{-0.0007}$ to constrain the Z-Z' mixing. You can assume $\hat{M}_{Z'} \gg \hat{M}_Z \gg \delta \hat{M}$.
- c) A well-motivated extension of the SM is a gauged B L symmetry (baryon minus leptonnumber). Write down explicitly the corresponding current for the SM fermions:

$$j'_{\mu} = \sum_{\psi} \bar{\psi} \gamma_{\mu} (B - L) \psi$$

where (B - L) denotes the B - L charge operator.

d) To cancel quantum anomalies in the B - L model, one also has to introduce three righthanded neutrinos N_i which are then part of the current: $\Delta j'_{\mu} = -\sum_i \bar{N}_i \gamma_{\mu} P_R N_i$. Due to Z - Z' mixing, the SM-like Z_1 will also couple to these new neutrinos. Can you give a constraint on the Z - Z' mixing from the well-measured (invisible) decay width of Z (for $M(N_i) \ll M_Z/2$)? See the PDG for numbers.

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Lecture webpage: https://www.mpi-hd.mpg.de/manitop/StandardModel2/index.html

Hand-in and discussion of sheet:

January 30, 2019, 15:45, Philosophenweg 12, R106