

# Exercises to “Standard Model of Particle Physics II”

Winter 2018/19

Prof. Dr. Manfred Lindner and Dr. Werner Rodejohann      Sheet 10      December 19, 2018

---

## Problem 20: *Neutrino oscillations in matter* [20 Points]

The MSW effect describes oscillations of neutrinos in matter. Consider the case of two flavour oscillations in a realistic (non-constant) mass distribution.

The time evolution of the eigenfunctions of the Hamiltonian in matter is described by the following expression:

$$i \frac{d}{dt} \begin{pmatrix} |\nu_A\rangle \\ |\nu_B\rangle \end{pmatrix} = \begin{pmatrix} E_A(t) & -i\dot{\theta}(t) \\ i\dot{\theta}(t) & E_B(t) \end{pmatrix} \begin{pmatrix} |\nu_A\rangle \\ |\nu_B\rangle \end{pmatrix},$$

where

$$-E_A(t) = E_B(t) = \Delta m^2 \sin(2\theta_0)/4E \sin(2\theta)$$

with the (constant) vacuum mixing angle  $\theta_0$  and electron number density  $N_e(t)$  in matter. The eigenstates in matter are  $|\nu\rangle = (|\nu_A\rangle, |\nu_B\rangle)^T$  and at any given time they are connected to the flavor states via:  $|\tilde{\nu}\rangle = \tilde{U}(t)|\nu\rangle$  with

$$|\tilde{\nu}\rangle = \begin{pmatrix} |\tilde{\nu}_e\rangle \\ |\tilde{\nu}_\mu\rangle \end{pmatrix}, \quad \tilde{U} = \begin{pmatrix} \cos\theta & \sin\theta \\ -\sin\theta & \cos\theta \end{pmatrix},$$

and the effective (time-dependent) mixing angle

$$\tan(2\theta) = \frac{\frac{\Delta m^2}{2E} \sin(2\theta_0)}{\frac{\Delta m^2}{2E} \cos(2\theta_0) - \sqrt{2} G_F N_e(t)}. \quad (1)$$

The effective Hamiltonian is not diagonal, since matter eigenstates and thus the mixing angle  $\theta$  are not constant. If, however,  $\dot{\theta}$  is small, i.e. the change in  $\theta$  occurs slowly, we obtain the adiabatic approximation, where we can neglect  $\dot{\theta}$ . Suppose that at the time  $t = t_i$  an electron neutrino is produced,

$$|\tilde{\nu}_\alpha(t_i)\rangle \equiv \delta_{\alpha e} |\tilde{\nu}_\alpha\rangle$$

- a) What is the expression for  $|\tilde{\nu}_\alpha(t_f)\rangle$  at some later time  $t_f$  in the adiabatic approximation?
- b) Show that in the adiabatic approximation the result for the transition probability is given by

$$P(\nu_e \rightarrow \nu_\mu) = \frac{1}{2} - \frac{1}{2} \cos(2\theta_i) \cos(2\theta_f) - \frac{1}{2} \sin(2\theta_i) \sin(2\theta_f) \cos\Phi_{AB}, \quad (2)$$

with  $\theta_i \equiv \theta(t = t_i)$ ,  $\theta_f \equiv \theta(t = t_f)$ , and

$$\Phi_{IJ} \equiv \int_{t_i}^{t_f} [E_I(t) - E_J(t)] dt,$$

where  $I, J \in \{A, B\}$ .

- c) Show that in the case of constant electron density, Eq. (2) reduces to the standard formula of two-flavor oscillations

$$P(\nu_e \rightarrow \nu_\mu)(t) = \sin^2(2\theta) \sin^2\left(\frac{\pi t}{L_M^{\text{osc}}}\right),$$

and determine  $L_M^{\text{osc}}$ . What happens in the limit  $N_e \rightarrow 0$ ?

- d) Use Eq. (1) to find an expression for  $\sin^2(2\theta)$ . The formula for  $\sin^2(2\theta)$  describes a Breit-Wigner distribution. We define the resonance width at half height,  $\Delta r$ , as the width at which the curve of the distribution is larger than one half [ $\sin^2(2\theta) > 1/2$ ]. Find an expression for  $\Delta r$  in terms of  $\theta_0$ .
- e) The *adiabaticity parameter* is defined as

$$\gamma_r \equiv \frac{|E_A|}{|\dot{\theta}|},$$

which becomes large if the adiabatic approximation is good. To interpret this, first show that

$$\dot{\theta} = \frac{1}{2} \frac{\sin^2(2\theta)}{\Delta m^2 \sin(2\theta_0)} 2\sqrt{2}EG_F\dot{N}_e.$$

Now assume that the change in matter density is approximately linear,  $\dot{N}_e \approx \text{const}$ , and that the density passes exactly through the resonance, i.e.  $\Delta N_e = N_e(t_f) - N_e(t_i) \propto \Delta r$  to formulate  $\gamma_r$  in terms of  $L_M^{\text{osc}}$  and  $\Delta t = t_f - t_i$ . Discuss the adiabaticity condition  $\gamma_r > 3$ .

**Problem 21: The Koide relation [0 Points]**

The Koide formula reads

$$Q = \frac{m_e + m_\mu + m_\tau}{\left(\sqrt{m_e} + \sqrt{m_\mu} + \sqrt{m_\tau}\right)^2}.$$

- a) Assuming only that the masses are positive, determine the range in which  $Q$  may take values.
- b) Calculate  $Q$  for the current best-fit values according to PDG<sup>1</sup>,  $m_e = 0.510\,998\,946\,1$  MeV,  $m_\mu = 105.658\,374\,5$  MeV,  $m_\tau = 1776.86$  MeV.
- c) The empirical value of  $Q$  lies very close to a rational number which has a special position in the range of possible  $Q$ -values. Can you find an explanation for this astonishing apparent coincidence?

**Happy Holidays!**

**Tutor:**

Ingolf Bischer

e-mail: bischer@mpi-hd.mpg.de

Lecture webpage: <https://www.mpi-hd.mpg.de/manitop/StandardModel2/index.html>

**Hand-in and discussion of sheet:**

January 9, 2019, 15:45, Philosophenweg 12, R106

<sup>1</sup><http://pdg.lbl.gov/2018/tables/rpp2018-sum-leptons.pdf>