

#) Neutrino mass?

$$- \mathcal{L} = \bar{\ell}'_L M^{(R)} \ell'_R \quad \text{with} \quad \ell'_{L,R} = \begin{pmatrix} \ell' \\ \mu' \\ \tau' \end{pmatrix}_{L,R} \quad \nu'_L = \begin{pmatrix} \nu'_\mu \\ \nu'_\tau \\ \nu'_\tau \end{pmatrix}_L$$

$$= \underbrace{\bar{\ell}'_L}_{\equiv \bar{\ell}_L} \underbrace{u_L u_L^\dagger M^{(R)} V_L V_L^\dagger}_{D^R} \underbrace{\ell'_R}_{\equiv \ell_R}$$

$$\Rightarrow \mathcal{L}_{cc} = \frac{g}{\sqrt{2}} W_\mu^+ \underbrace{\bar{\ell}'_L}_{\bar{\ell}_L} u_L u_L^\dagger \gamma^\mu \underbrace{u_L u_L^\dagger}_{\nu_L} \nu'_L$$

if no mass term for neutrinos: U_s arbitrary

\Rightarrow choose $U_s = U_R \Rightarrow$ no mixing in

lepton sector \Rightarrow no analogue of CKM

If there is mass for neutrinos, the PMNS

(Ponteconvo-Maki-Nakagawa-Sakata) matrix

exists

III 2) Neutrino Oscillations

- only physics beyond the Standard Model that can be directly tested in lab
- connected to new representations, and (very possibly) with new energy scales and new concepts
- ν interact very weakly \leftrightarrow challenging expts, but option to do source physics as well (solar interior, earth interior, cosmic rays...)

A) Oscillations in Vacuum

suppose neutrinos have mass:

$$\nu_\alpha = U_{\alpha i}^* \nu_i$$

↓
flavor

↓
mass

PMNS-matrix

as usual, $\frac{1}{2} N(N-1)$ angles and $\frac{1}{2} (N-2)(N-1)$ phases

(note: this assumes mass term $\bar{\nu}_L \nu_R$; also possible $\bar{\nu}_L \nu_L^c$ & $\bar{\nu}_L \nu_L^*$ \Rightarrow rephasing does not work, can only rephase N charged lepton fields $\Rightarrow N-1$ additional phases in U , but oscillations are not affected.)

#) How to obtain the oscillation probability

$$|\nu(0)\rangle = |\nu_\alpha\rangle = U_{\alpha j}^* |\nu_j\rangle$$

flavor state produced in CC reaction (e.g.) at time 0

$$|\nu(t)\rangle = U_{\alpha j}^* e^{-iE_j t} |\nu_j\rangle \quad \text{is time evolution}$$

amplitude to find state $|\nu_\beta\rangle$:

$$\begin{aligned} A(\nu_\alpha \rightarrow \nu_\beta, t) &= \langle \nu_\beta | \nu(t) \rangle = U_{\beta i} U_{\alpha j}^* e^{-iE_j t} \langle \nu_i | \nu_j \rangle \\ &= U_{\beta i}^* U_{\alpha j} e^{-iE_i t} \end{aligned}$$

probability:

$$P(\nu_\alpha \rightarrow \nu_\beta, t) = P_{\alpha\beta} = |\mathcal{A}(\nu_\alpha \rightarrow \nu_\beta, t)|^2$$

$$= \sum_{ij} \underbrace{U_{\alpha i}^* U_{\beta i} U_{\beta j}^* U_{\alpha j}}_{\mathcal{J}_{ij}^{\alpha\beta}} \underbrace{e^{-i(E_i - E_j)t}}_{e^{-i\Delta_{ij}}}$$

$$= \dots = \delta_{\alpha\beta} - 4 \sum_{j>i} \text{Re}\{\mathcal{J}_{ij}^{\alpha\beta}\} \sin^2 \frac{\Delta_{ij}}{2} + 2 \sum_{j>i} \text{Im}\{\mathcal{J}_{ij}^{\alpha\beta}\} \sin \Delta_{ij}$$

phase: $\frac{1}{2} \Delta_{ij} = \frac{1}{2} (E_i - E_j)t = \frac{1}{2} (\sqrt{p_i^2 + m_i^2} - \sqrt{p_j^2 + m_j^2}) L$

$$\approx \frac{1}{2} \left[p_i \left(1 + \frac{m_i^2}{2p_i^2}\right) - p_j \left(1 + \frac{m_j^2}{2p_j^2}\right) \right] L$$

$$\approx \frac{m_i^2 - m_j^2}{4E} L$$

$$= 1.27 \left(\frac{\Delta m_{ij}^2}{\text{eV}^2} \right) \left(\frac{L}{\text{km}} \right) \left(\frac{\text{GeV}}{E} \right)$$

$\alpha = \beta$: survival probability

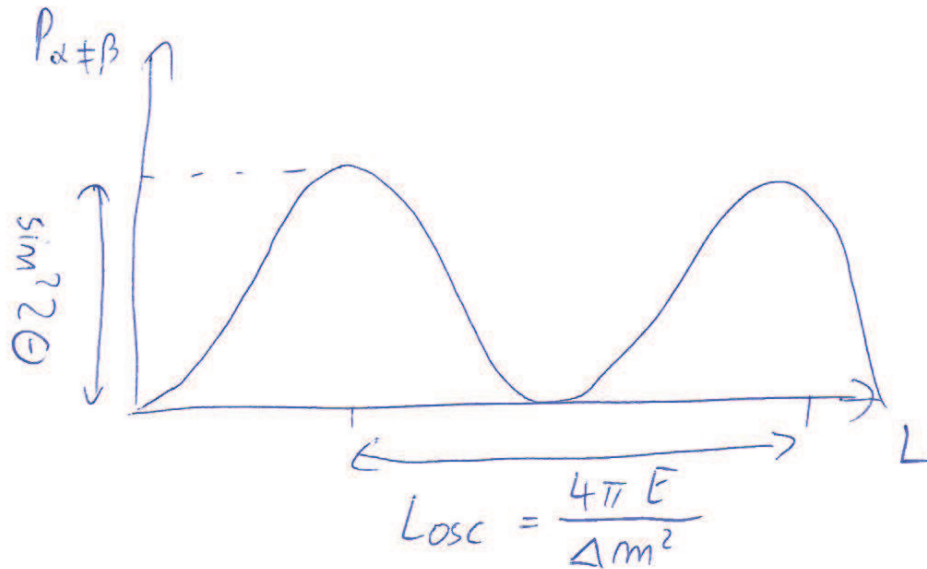
$\alpha \neq \beta$: transition "

requires $U \neq \mathbb{1}$ and $\Delta m_{ij}^2 \neq 0$

(67)

2 flavor case: $U = \begin{pmatrix} c & s \\ -s & c \end{pmatrix} \Rightarrow \begin{aligned} \omega_e &= c\omega_1 + s\omega_2 \\ \omega_\mu &= c\omega_2 - s\omega_1 \end{aligned}$

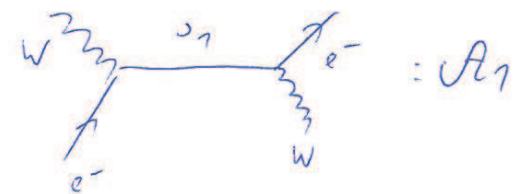
$$P_{\alpha \neq \beta} = \sin^2 2\Theta \sin^2 \frac{\Delta m^2 L}{4E}$$



oscillation length

e.g. $E = \text{GeV}, \Delta m^2 = 10^{-3} \text{ eV}^2 \Rightarrow L_{osc} \sim 1000 \text{ km}$

quantum mechanical interference on macroscopic scales!



$P = |A_1 + A_2|^2$ coherence!

This derivation was wrong:

- 1) $E_i - E_j$ not Lorentz-invariant
 - 2) different p_i , same E : violates energy-momentum conservation
 - 3) define $p = \gamma m e^{iPx}$: no localization
- \Rightarrow do in quantum mechanics with wave packets: same result... (62)

o) same E , different p ?

$E_j \neq \sqrt{E_j^2 - m_j^2} = p_j$, do Taylor-expansion:

$$\Rightarrow p_j = E + m_j^2 \frac{\partial p_j}{\partial m_j^2} \Big|_{m_j=0} \equiv E - \xi \frac{m_j^2}{2E} \text{ with } \xi = -2E \frac{\partial p_j}{\partial m_j^2} \Big|_{m_j=0}$$

$$E_j = p_j + m_j^2 \frac{\partial E_j}{\partial m_j^2} \Big|_{m_j=0} = p_j + \frac{m_j^2}{2p_j} = E + \frac{m_j^2}{2E} (1 - \xi)$$

consider π -decay: $E_j = \frac{m_\pi^2}{2} \left(1 - \frac{m_\mu^2}{m_\pi^2} \right) + \frac{m_j^2}{2m_\pi^2}$

$$\Rightarrow \xi = \frac{1}{2} \left(1 + \frac{m_\mu^2}{m_\pi^2} \right) \approx 0.8$$

and $E_i - E_j \approx (1 - \xi) \frac{\Delta m^2}{2E}$

phase difference would depend on neutrino source ...

On the correct way to get $P_{\alpha\beta}$:

$$\Psi_i \propto \exp\left\{-i(E_i t - p_i x) - \frac{(x - v_i t)^2}{4\sigma_x^2}\right\}$$

with $\sigma_x \approx 1/\sigma_p$ and group velocity $v_i = \frac{\partial E_i}{\partial p_i} = p_i/E_i$

2 conditions to see oscillations:

1) wave packet separation should be smaller than σ_x !

$$\Rightarrow L \Delta v \stackrel{!}{\leq} \sigma_x \Rightarrow \frac{L}{L_{osc}} < \frac{p}{\sigma_p}$$

(loss of coherence: interference impossible)

2) m_s^2 should not be known too precisely !

if known too well: $\Delta m^2 \Rightarrow \delta m_s^2 = \frac{\partial m_s^2}{\partial p_s} \delta p_s$

$$\Rightarrow \delta x_s \gg \frac{2p}{\Delta m^2} = \frac{L_{osc}}{2\pi}$$

() know which state ν_i is exchanged }

proper calculation gives:

$$P \propto \exp \left\{ -i \frac{\Delta m_{ij}^2}{2E} L - \left(\frac{L}{L_{ij}^{\text{coh}}} \right)^2 - (2\pi^2) (1 - \dots)^2 \left(\frac{\sigma_x}{L_{ij}^{\text{osc}}} \right)^2 \right\}$$

$$\text{with } L_{ij}^{\text{coh}} = \frac{4\sqrt{2}E}{|\Delta m_{ij}^2|} \sigma_x \quad \text{and} \quad L_{ij}^{\text{osc}} = \frac{4\pi E}{|\Delta m_{ij}^2|}$$

~~$$P_j = \dots$$~~

$$P_j \equiv E - \sum m_j^2 / 2E$$

the last 2 terms suppress oscillation if neutrinos still mix;

$$P_{\alpha\beta} = \sum |U_{\alpha i}|^2 |U_{\beta i}|^2$$

Example of such neutrinos losing coherence:

IceCube: neutrinos as cosmic rays

[the same happens when $L \gg L_{\text{osc}}$ "fast oscillations"

$$\langle \sin^2 \frac{\Delta m^2 L}{4E} \rangle \Rightarrow = \langle \sin^2 \frac{\pi L}{L_{\text{osc}}} \rangle = 1/2$$

