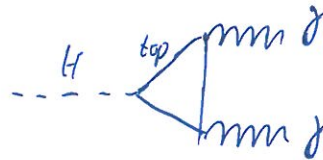
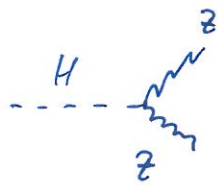


Recap: •) $\Phi \sim Z_{L,1}$ with mass \lesssim few 100 GeV ³⁻⁷²⁻⁷⁴
 seems to make sense; plus: needed for consistency

•) measured / detected @ LHC (July 2012)

most important channels: 



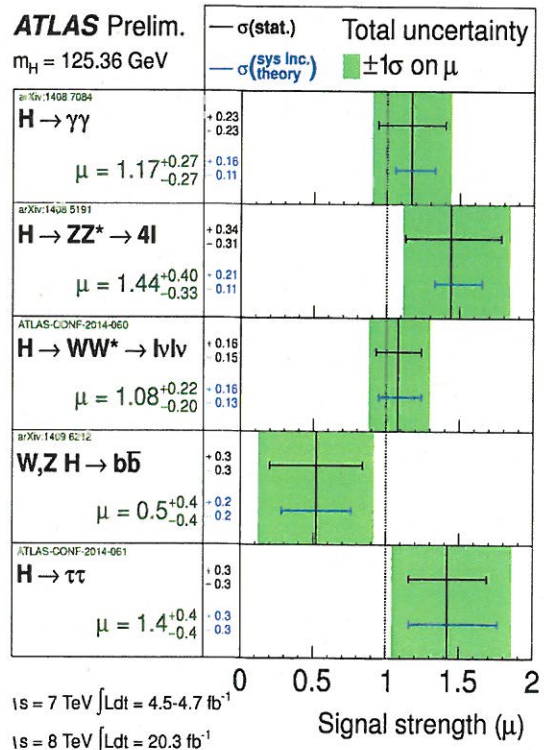
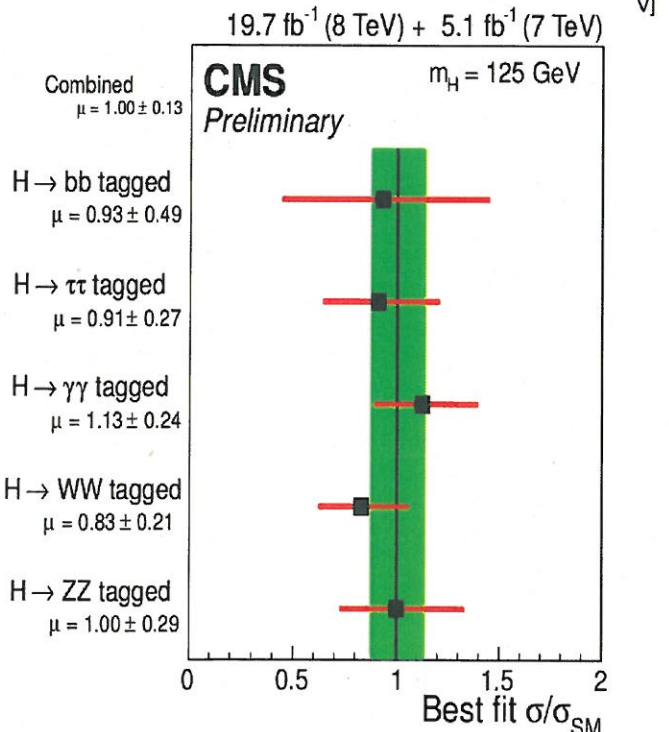
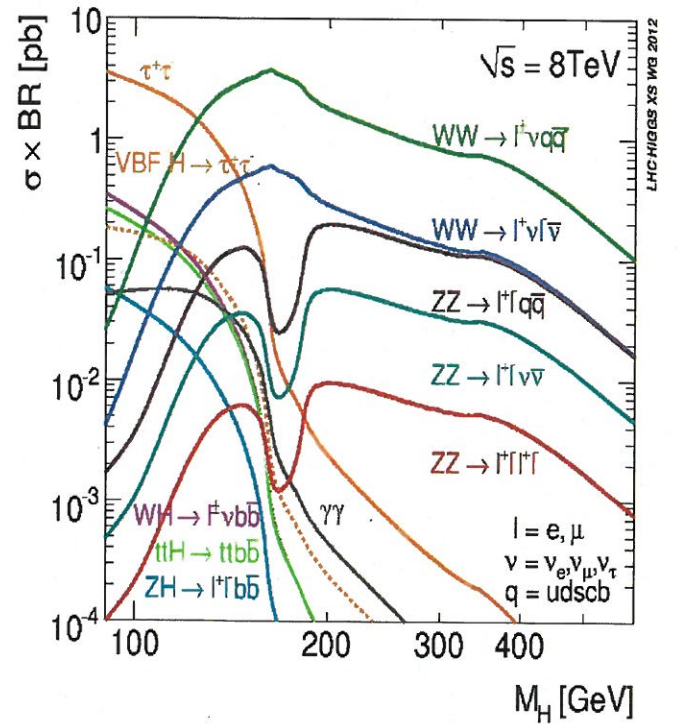
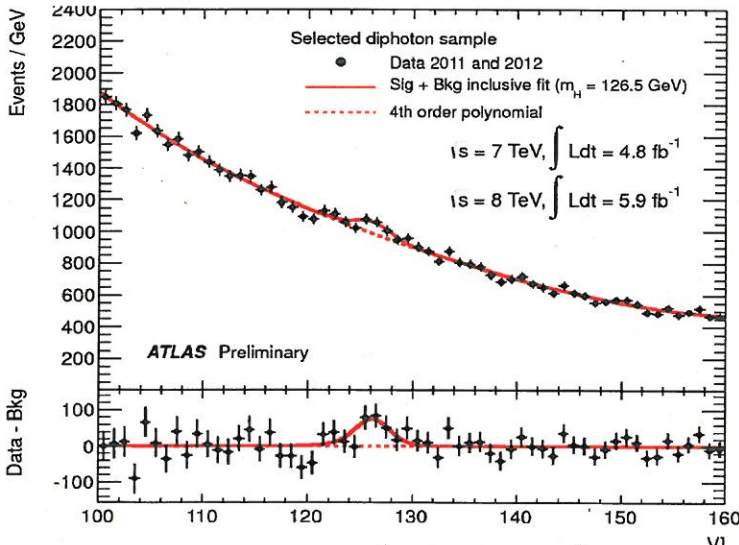
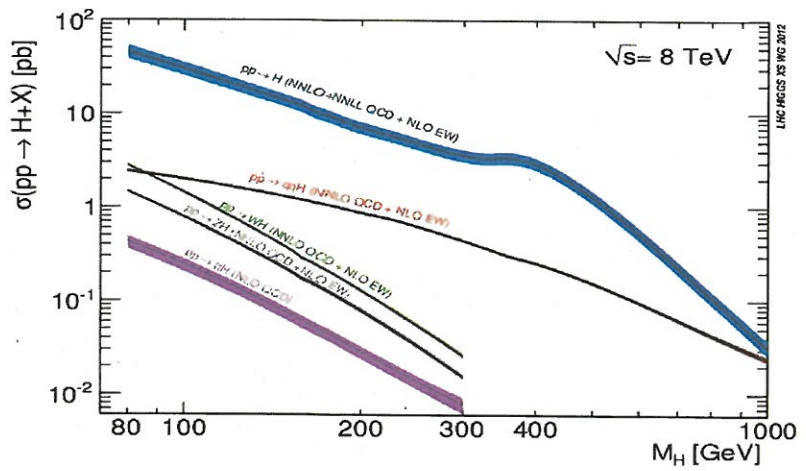
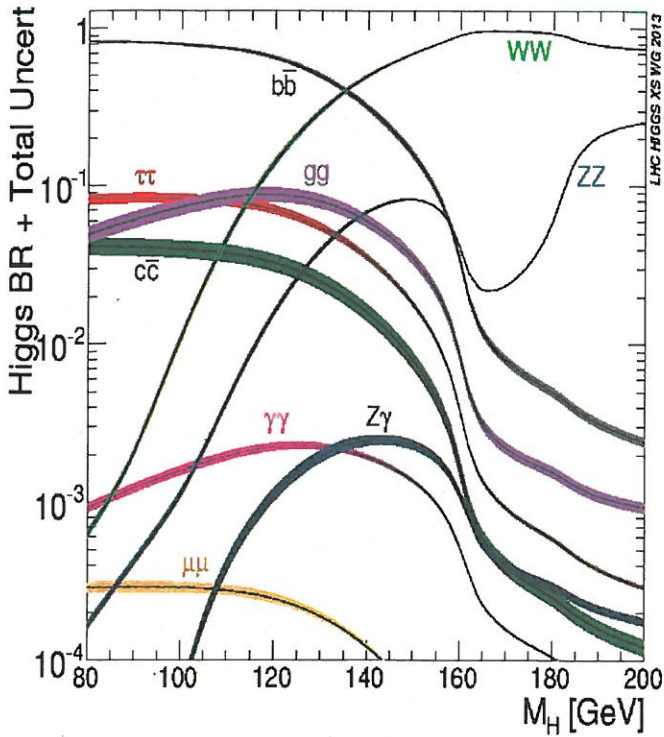
$$m_H \approx 125.36 \pm 0.37 \pm 0.78 \text{ GeV}$$

\uparrow \uparrow
 stat sys

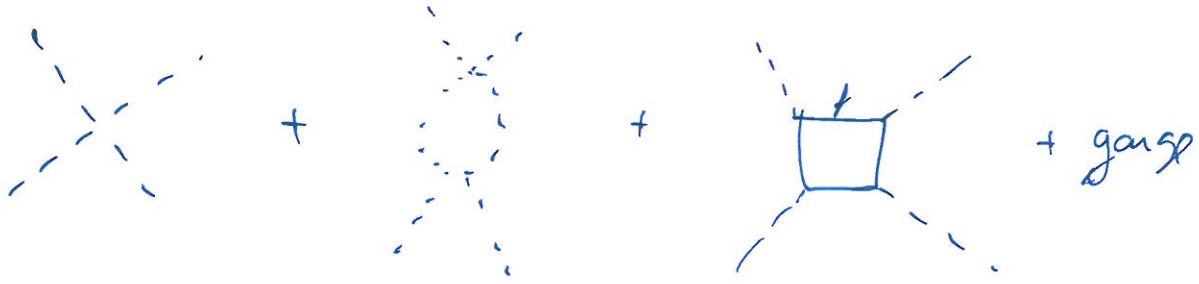
$$\Gamma_{\text{pred}} = 4.27 \times 10^{-3} \text{ GeV} \Rightarrow \tau = 1.56 \times 10^{-22} \text{ s}$$

(some) experimental aspects \rightarrow sheet with BR,
 σ , $\sigma \times \text{BR}$, etc.

today: extrapolation of Higgs to higher E



RG-Analysis of Higgs



$\Rightarrow d = d(Q^2) = f(\lambda, g, g', m_t/v)$ \Rightarrow top-quark gives largest fermion-contribution

$$16\pi^2 \frac{d\lambda}{d\log Q^2} = 12\lambda^2 + 6\lambda\gamma_t^2 - 3\gamma_t^4 - \frac{3}{2}\lambda(3g'^2 + g^2) + \frac{3}{16}(2g'^4 + (g' + g)^2)$$

a) large λ

$$\frac{d\lambda}{d\log Q^2} = \frac{3}{4\pi^2} \lambda \quad : \quad \lambda \nearrow \text{ for } Q^2 \nearrow$$

$$\Rightarrow d(Q^2) = \frac{\lambda(\mu^2)}{1 - \frac{3}{4\pi^2} \lambda(\mu^2) \log \frac{Q^2}{\mu^2}}$$

choose $\mu^2 = v^2$
as start

$$\text{denominator} = 0 \quad \text{for} \quad Q = v \exp\left\{ \frac{2\pi^2}{3\lambda(v^2)} \right\}$$

$$= v \exp\left\{ \frac{4\pi^2 v^2}{3m_\lambda^2} \right\}$$

Landau-pole only avoidable for $\lambda=0 \Rightarrow$ "triviality bound"

$$\Leftrightarrow \lambda(\text{low energy}) \stackrel{!}{<} \lambda_{\max} \Rightarrow m_\lambda < m_\lambda^{\max}$$

a) small λ

$$\frac{d\lambda}{d \log Q^2} = -\frac{3}{16\pi^2} \gamma_t^4 \quad : \quad \lambda \downarrow \text{ for } Q^2 \uparrow$$

(top Yukawa)

$$\Rightarrow \lambda(Q^2) = \lambda(v^2) - \frac{3\gamma_t^4}{16\pi^2} \log \frac{Q^2}{v^2} \quad \left(\text{starting value } \lambda(v^2) = \frac{m_h^2}{2v^2} \right)$$

when is $\lambda(Q^2) = 0$?

$$\lambda(v^2) = \frac{3}{16\pi^2} \gamma_t^4 \log \frac{Q^2}{v^2} \quad ; \quad \gamma_t = \frac{\sqrt{2} m_t}{v}$$

$$\Rightarrow m_h = \frac{3m_t^4}{2\pi^2 v^2} \log \frac{Q^2}{v^2} \quad "$$

eg: $m_h = 126 \text{ GeV}$; $m_t = 173 \text{ GeV}$

$$\Rightarrow Q = \exp\left\{ \frac{m_h^2 \pi^2 v^2}{3m_t^4} \right\} v \approx 10^5 \text{ GeV}$$

$$\lambda(\text{low energy}) > \lambda_{\text{min}} \Rightarrow m_h \geq m_h^{\text{min}}$$

"vacuum stability bound"

[λ must be > 0 to have a minimum in V ;
 $\Leftrightarrow V$ bounded from below!]

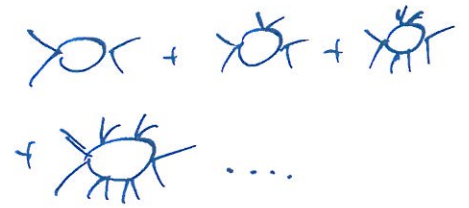
otherwise theory not defined!

the larger Φ , the more negative the Hamiltonian

•) the story doesn't end here... further corrections:

e.g. $V_{\text{eff}} \propto -\gamma_t^4 \Phi^4 \ln \frac{g^2 \Phi^2}{\mu^2}$ from summing up:

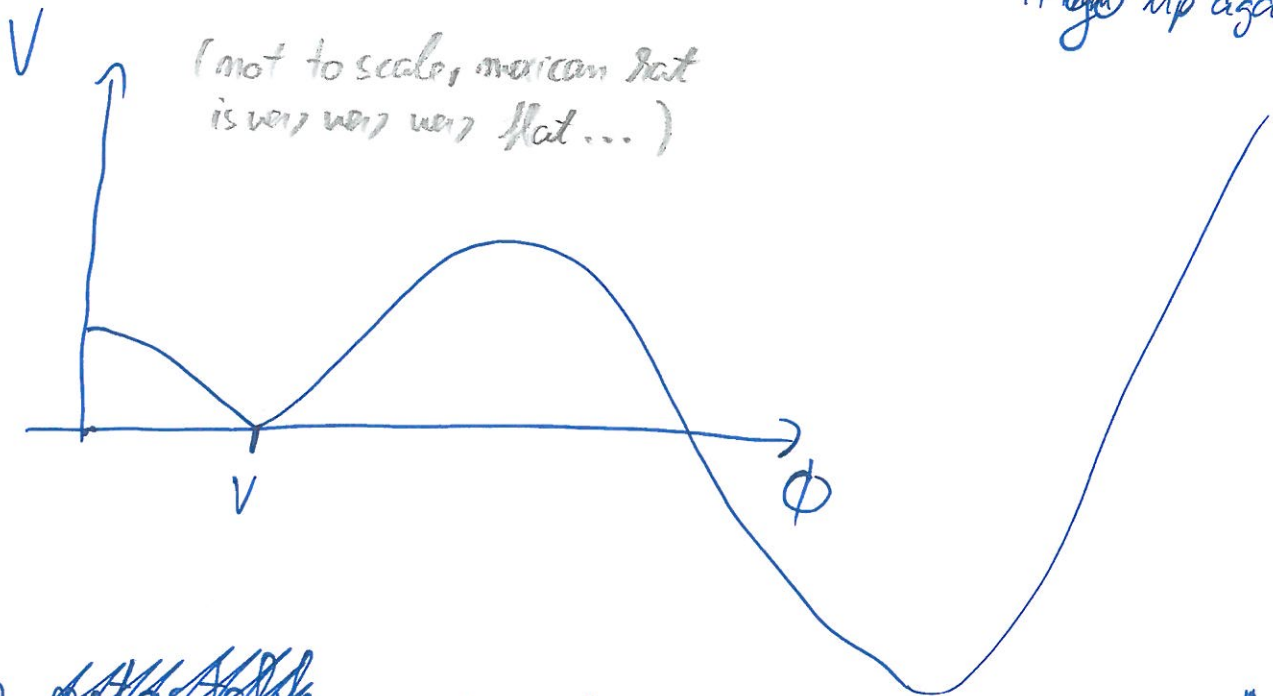
[if no quadratic term in V : conformal symmetry!
 \rightarrow minimum generated by V_{eff} !
 "dimensional transmutation" (\leftrightarrow for imCCD)
 gives too small Higgs mass in SM...]
 Coleman-Weinberg



at large Φ -values: V dominated by $\lambda_{\text{eff}} \Phi^4$

(not to scale)

\downarrow
 becomes > 0 again:
 $\gamma_t \downarrow$, gauge couplings make
 it go up again



\Rightarrow ~~metastable~~ \Rightarrow tunneling in true minimum!

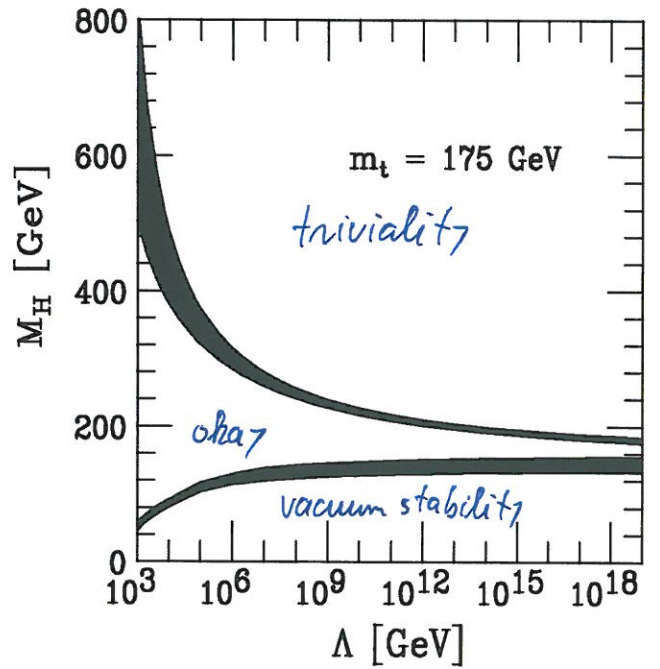
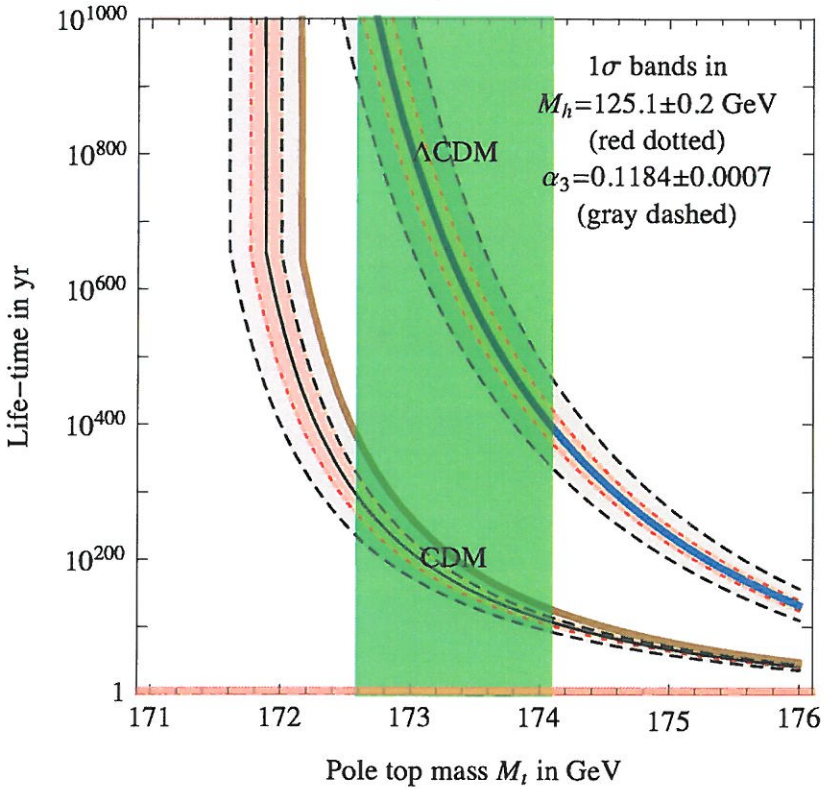
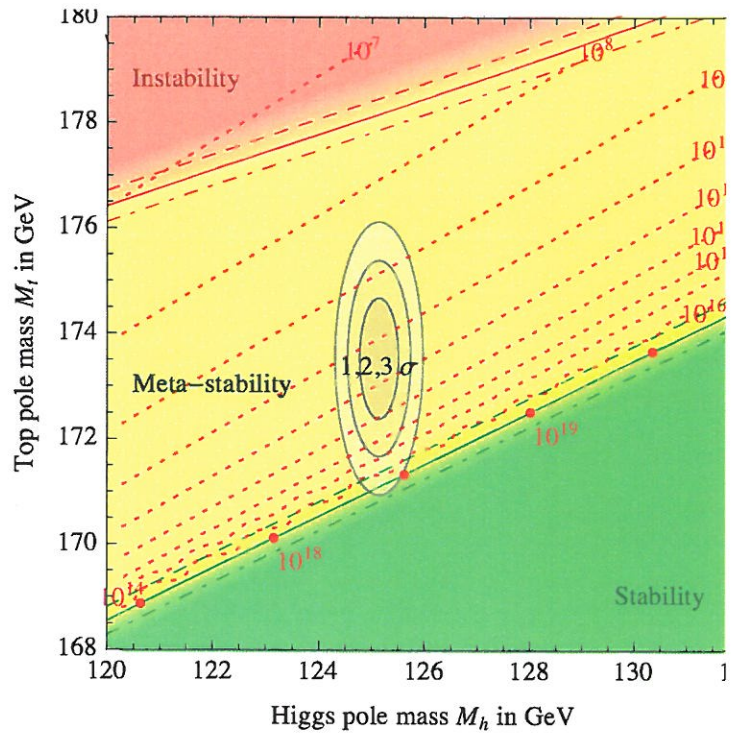
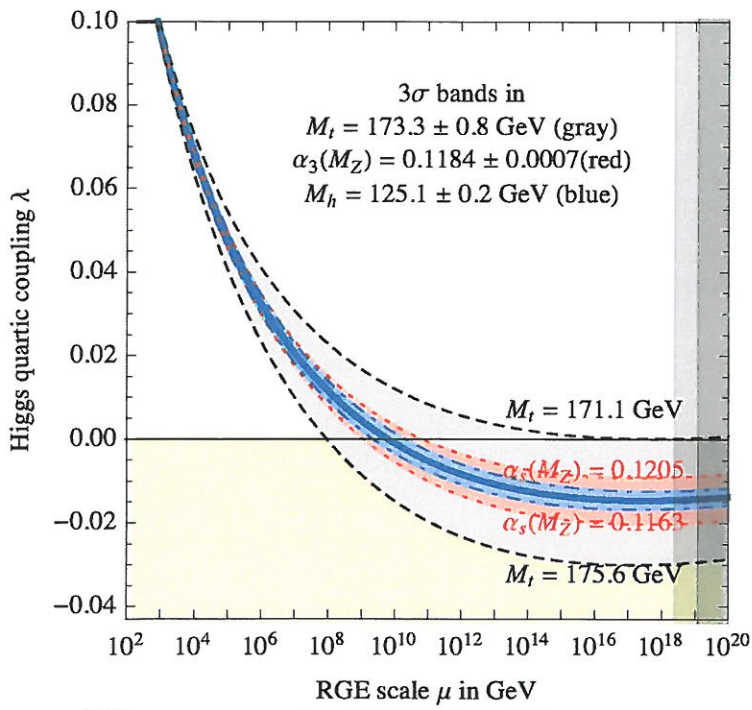
if $\tau > T_{\text{universe}}$: metastable

$\tau < T_{\text{universe}}$: unstable (bad!)

stable for $m_H > 129.6 \text{ GeV} + 2(m_t - 173.35 \text{ GeV}) = \frac{\alpha_s(m_Z) - 0.1184}{0.0014} \pm 0.36$

- « Vacuum decay » :
-) quantum tunneling (depends on cosmology - past and future...)
 -) « classical transition » via bubble formation ; radius R ;
surface costs energy $\propto R^2$
volume produces energy $\propto R^3$

-) experimental question (value of m_t, α_s)
-) new physics coupling to SM modifies all calculations !
-) stretching all parameters to corners of parameter space : $\lambda(M_{pl}) = 0$ is possible
(boundary conditions from e.g. Quantum Gravity?)



See: xxx.lanl.gov/abs/1307.3536

•) Hierarchy problem

Embedding of SM in larger theory, cutoff Λ

$$m_{\text{fermion}} = m_{\text{fermion}}^0 \left(1 + c \log \frac{\Lambda}{\mu}\right)$$

-) multiplicative
-) $m_f \rightarrow 0$ for $m_f^0 = 0$
-) logarithmic

$$m_h^2 = (m_h^0)^2 + \frac{3}{8\pi^2 v^2} (4m_t^2 - 2m_W^2 - 4m_Z^2 - m_h^2) \Lambda^2$$

-) quadratic!
-) want $\Delta m_h \lesssim m_h^0$ "natural theory"

\Rightarrow fine-tuning: if $\Lambda = M_{\text{Pl}} = \sqrt{\frac{1}{6N}}$

$$m_h^2 = (m_h^0)^2 + 10^{37} \text{ GeV}^2$$

\downarrow
 $(10^2 \text{ GeV})^2$

\Rightarrow 10^{-33} fine-tuning



Solutions:

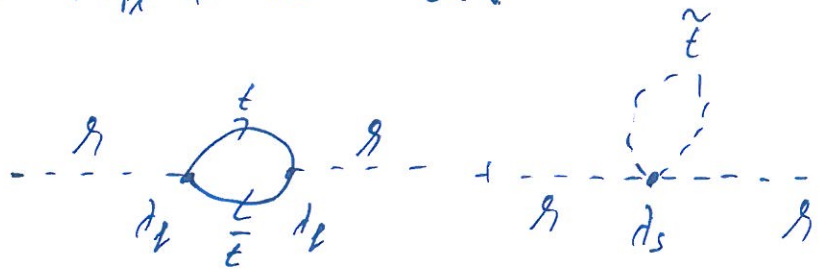
•) $4 m_t^2 - 2 m_W^2 - 4 m_Z^2 - m_h^2 = 0$ (Veltman condition)

→ at electroweak scale : $m_h = 374 \text{ GeV}$ ↯

(→ works at Planck scale...)

•) Extra dimensions : $M_{Pl} \approx 10^{19} \text{ GeV}$

•) Supersymmetry



scalar partner of top = "stop"

$$\Delta m_h^2 \propto (-d_t^2 + d_s) \Lambda^2$$

$$= 0 \text{ for exact SUSY } m_t = m_{\tilde{t}}$$

must be broken : $\Delta m_h^2 \approx \frac{1}{16\pi^2} m_{\text{SUSY}}^2 \log \frac{\Lambda}{m_{\text{SUSY}}}$

(must be zero for $m_{\text{SUSY}} = 0$
and should be $\propto \Lambda^2$ or $\ln \frac{\Lambda}{m_{\text{SUSY}}}$)

to have $\Delta m_h^2 \lesssim 10^2 \text{ GeV}$:

new particles at 10^3 GeV ...

•) avoid scalar particles!

Recall: pions (scalar...) formed by condensates of
u,d quarks; dynamical symmetry breaking
($d_s \nearrow$, confinement)

$\rightarrow \pi^\pm, \pi^0$ have quantum numbers of Goldstone-
bosons that are eaten by gauge bosons.

$\rightarrow m_{W, Z} \simeq g f_\pi$ with $f_\pi \simeq 100 \text{ GeV}$

too small (but included in calculations of EW
precision observables)

\rightarrow introduce new interaction and fermions
with $f_\pi = v$

"Technicolor" SU(N_{TC})

\Rightarrow Higgs Mechanism, but no Higgs particle

\Rightarrow Fermion masses?

for all this Higgs-stuff and more, see the write-up
of Tilman Plehn's LHC lectures from his homepage