Summer 2013

Prof. Dr. Andre Schöning, Dr. Werner Rodejohann	Sheet 8	17.6.13
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Exercise 16: Stückelberg mechanism [20 Points]

For a gauged abelian symmetry U(1)' (it does not extend to non-abelian symmetries) there exists an interesting mechanism to generate a massive gauge boson, while retaining renormalizability. The method contains a real scalar field σ together with the Z'-boson associated to the U(1)'.

a) Consider the Lagrangian

$$\mathcal{L} = -\frac{1}{4}Z'^{\mu\nu}Z'_{\mu\nu} + \frac{1}{2}(M_{Z'}Z'_{\mu} + \partial_{\mu}\sigma)(M_{Z'}Z'^{\mu} + \partial^{\mu}\sigma) + i\overline{\psi}\gamma^{\mu}(\partial_{\mu} - ig'Y'Z'_{\mu})\psi - m\overline{\psi}\psi.$$

The gauge transformations for the Dirac fermion (with U(1)' charge Y') and gauge boson are given by

$$\psi \to e^{-ig'Y'\theta(x)}\psi, \qquad Z'_{\mu} \to Z'_{\mu} - \partial_{\mu}\theta(x).$$

Calculate the gauge transformation of the real scalar σ that makes the Lagrangian invariant, and show the invariance of the other terms, too. Can you fix a gauge to eliminate σ from the theory? Count degrees of freedom in both gauges.

b) Consider an extension of the SM gauge group $G_{\rm SM} \equiv SU(3) \times SU(2) \times U(1)_Y \rightarrow G_{\rm SM} \times U(1)'$ with Lagrangian

$$\mathcal{L} \supset -\frac{1}{4}\hat{B}^{\mu\nu}\hat{B}_{\mu\nu} - \frac{1}{4}\hat{W}_{3}^{\mu\nu}\hat{W}_{3\mu\nu} + \frac{1}{2}\hat{M}_{Z}^{2}\left(\hat{c}_{W}\hat{W}_{3}^{\mu} - \hat{s}_{W}\hat{B}^{\mu}\right)^{2} \\ -\frac{1}{4}\hat{Z}'^{\mu\nu}\hat{Z}'_{\mu\nu} - \frac{\sin\chi}{2}\hat{B}^{\mu\nu}\hat{Z}'_{\mu\nu} + \frac{1}{2}\left(\partial_{\mu}\sigma + \hat{M}_{Z'}\hat{Z}'_{\mu} + \hat{M}_{B}\hat{B}_{\mu}\right)^{2},$$

where σ is charged under both U(1)' (with gauge boson \hat{Z}') and $U(1)_Y$ (with gauge boson \hat{B}'). Find again how σ has to transform under $U(1)_Y \times U(1)'$ for the Lagrangian to be invariant and show that the sin χ term is gauge invariant. In the SM, you find the Z boson $Z = \hat{c}_W \hat{W}_3^{\mu} - \hat{s}_W \hat{B}^{\mu}$ and the massless photon $\gamma = \hat{s}_W \hat{W}_3^{\mu} + \hat{c}_W \hat{B}^{\mu}$, but the extra neutral gauge boson \hat{Z}' changes these results. To find the gauge boson mass eigenstates, we need to diagonalize the kinetic terms first (if the "kinetic mixing angle" χ is nonzero). To do this, find a (non-unitary) transformation matrix R that transforms $(\hat{B}, \hat{W}_3, \hat{Z}')$ into $(\tilde{B}, \tilde{W}_3, \tilde{Z}')$ with the properly normalized "diagonal" kinetic terms

$$\mathcal{L} \supset -\frac{1}{4} \tilde{B}^{\mu\nu} \tilde{B}_{\mu\nu} - \frac{1}{4} \tilde{W}_3^{\mu\nu} \tilde{W}_{3\mu\nu} - \frac{1}{4} \tilde{Z}'^{\mu\nu} \tilde{Z}'_{\mu\nu}$$

Hint: $\hat{W}_3 = \tilde{W}_3$.

c) The fields \tilde{B} , \tilde{W}_3 and \tilde{Z}' share a non-diagonal mass matrix. Write it down and show that it has at least one zero eigenvalue, i.e. there is still a massless photon in the model. Calculate the other two masses for the case $\chi = 0$ and expand in $\hat{M}_B \ll \hat{M}_Z < \hat{M}_{Z'}$.

Exercise 17: Effects of Z-Z' mixing [10 Points]

The general effective Lagrange density after breaking the $SU(3)_C \times SU(2)_L \times U(1)_Y \times U(1)'$ symmetry to $SU(3)_C \times U(1)_{\rm EM}$ can be written as

$$\mathcal{L} = \mathcal{L}_{\mathrm{SM}} + \mathcal{L}_{Z'} + \mathcal{L}_{\mathrm{mix}} \,,$$

where the relevant part of the Standard Model Lagrangian is

$$\mathcal{L}_{\rm SM} = -\frac{1}{4}\hat{B}_{\mu\nu}\hat{B}^{\mu\nu} - \frac{1}{4}\hat{W}^a_{\mu\nu}\hat{W}^{a\mu\nu} + \frac{1}{2}\hat{M}^2_Z\hat{Z}_\mu\hat{Z}^\mu - \frac{\hat{e}}{\hat{c}_W}j^\mu_B\hat{B}_\mu - \frac{\hat{e}}{\hat{s}_W}j^{a\mu}_W\hat{W}^a_\mu,$$

and the hats merely denote that the fields are not mass eigenstates. The Z' part reads

$$\mathcal{L}_{Z'} = -\frac{1}{4} \hat{Z}'_{\mu\nu} \hat{Z}'^{\mu\nu} + \frac{1}{2} \hat{M}'^2_Z \hat{Z}'_\mu \hat{Z}'^\mu - \hat{g}' j'^\mu Z'_\mu,$$

and the kinetic- and mass-mixing terms can be parameterized as

$$\mathcal{L}_{\rm mix} = -\frac{\sin\chi}{2} \hat{Z}^{\prime\mu\nu} \hat{B}_{\mu\nu} + \delta \hat{M}^2 \hat{Z}^{\prime}_{\mu} \hat{Z}^{\mu} \,.$$

- a) Determine the mass eigenstates Z_1^{μ} and Z_2^{μ} and determine the couplings of $Z_{1,2}$ to the currents j_B , j_W and j'. Set the kinetic mixing angle χ to zero for simplicity.
- b) Since the mass of the physical Z boson changes compared to the SM, the so-called ρ parameter $\rho = M_W^2/M_Z^2 c_W^2$ is no longer equal to one. Use the current value $\rho = 1.0008^{+0.0017}_{-0.0007}$ to constrain the Z–Z' mixing.

Tutors:

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Tutorials homepage: http://www.mpi-hd.mpg.de/manitop/StandardModel/exercise.html

Hand-in of sheet:

during lecture on 24.6.

Discussion of sheet:

Thursday, 27.6. 2.15 pm, INF 227 SR 2.402 Friday, 28.6. 2.15 pm, INF 227 SR 1.403