

# Exercises to “Standard Model of Particle Physics”

Summer 2013

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Sheet 7

13.6.12

## Exercise 15: Goldstone theorem [20 Points]

The Goldstone–Theorem states: every generator of a continuous global symmetry, which does not annihilate the vacuum, corresponds to a massless scalar particle.

As “proof”, consider the Lagrangian  $\mathcal{L}(\Phi_i)$  with real scalar fields  $\Phi_i$ , invariant under the global transformation  $\vec{\Phi} \rightarrow \exp(i\theta^a T^a) \vec{\Phi}$ .

1. Which properties have the  $iT^a$ , if the  $\Phi_i$  are real and if  $\Phi^T \Phi$  is invariant under the transformation?
2. Show that from the conservation of the Noether current it follows that:

$$\frac{\partial \mathcal{L}}{\partial \Phi_i} (iT^a)_{ij} \Phi_j + \frac{\partial \mathcal{L}}{\partial (\partial^\mu \Phi_i)} (iT^a)_{ij} \partial^\mu \Phi_j = 0.$$

3. Let  $\mathcal{L} = \frac{1}{2}(\partial_\mu \vec{\Phi})^T (\partial^\mu \vec{\Phi}) - V(\vec{\Phi})$ . Show with the help of the first two results that

$$\frac{\partial V}{\partial \Phi_i} (iT^a)_{ij} \Phi_j = 0.$$

4. Let  $\vec{v}$  be the minimum of  $V(\vec{\Phi})$ . Show with the results from above that  $M^2 (iT^a) \vec{v} = 0$ , where  $M_{ij}^2$  is given by

$$(M^2)_{ij} := \left. \frac{\partial^2 V}{\partial \Phi_i \partial \Phi_j} \right|_{\vec{\Phi}=\vec{v}}$$

5. The matrix  $M^2$  is interpreted as a mass matrix. The broken vacuum  $\vec{v}$  is in general not invariant under the transformations  $\exp(i\theta^a T^a)$ . If  $T^a \vec{v} = 0$ , one has a *Wigner–Weyl realization* of the symmetry, if  $T^a \vec{v} \neq 0$  one has a realization à la *Nambu–Goldstone*. With the result from the last point it follows that for all  $a$  with  $T^a \vec{v} \neq 0$  the mass matrix has an eigenvector with eigenvalue 0, i.e. the scalar field  $\phi_i (iT^a)_{ij} v_j$  (with  $\phi_i = \Phi_i - v_i$ ) is massless.

Consider e.g. the Lagrangian  $\mathcal{L} = \frac{1}{2}(\partial_\mu \vec{\Phi})^T (\partial^\mu \vec{\Phi}) - V(\vec{\Phi}^T \vec{\Phi})$  with the potential  $V(\vec{\Phi}^T \vec{\Phi}) = \frac{1}{2}\mu^2 \vec{\Phi}^T \vec{\Phi} + \frac{1}{4}\lambda (\vec{\Phi}^T \vec{\Phi})^2$ , where  $\vec{\Phi}^T = (\phi_1, \phi_2, \phi_3)$  is a triplet of real and scalar particles. The coupling obeys  $\lambda > 0$ .

- (i) Show that  $\mathcal{L}$  is invariant under a  $SU(2)$  transformation  $\vec{\Phi} \rightarrow \exp(i\theta^a T^a) \vec{\Phi}$ . The 3 generators  $T^j$  of  $SU(2)$  are written here as  $(T^j)_{kl} = -i\epsilon_{jkl}$  (this is called the adjoint representation of  $SU(2)$ ).
- (ii) For  $\mu^2 < 0$  the potential has a minimum for  $\vec{v}^T = (0, 0, v)$ . How are  $v$ ,  $\mu^2$  and  $\lambda$  connected? Show that the vacuum state  $\vec{v}$  is invariant under the transformation with  $T^3$ , but not with  $T^1$  and  $T^2$ . That is, the vacuum state does *not* possess the symmetry of the Lagrangian (spontaneous symmetry breaking). Why do we speak in this case of spontaneous symmetry breaking from  $SU(2)$  to  $U(1)$ ?
- (iii) The scalar fields will be expanded around the minimum:  $\vec{\Phi}^T \equiv (\phi_1, \phi_2, v + \phi_3)$ . Show by inserting into the above Lagrangian that there is one massive scalar fields and 2 massless Goldstone bosons. What is the mass of  $\phi_3$ ? Convince yourself that the matrix  $M^2$  from part 4 is really the mass matrix.

**Tutors:**

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Tutorials homepage: <http://www.mpi-hd.mpg.de/manitop/StandardModel/exercise.html>

**Hand-in of sheet:**

during lecture on 17.6.

**Discussion of sheet:**

Thursday, 20.6. 2.15 pm, INF 227 SR 2.402

Friday, 21.6. 2.15 pm, INF 227 SR 1.403