Lecture:

Standard Model of Particle Physics

Heidelberg SS 2012

Experimental Tests of QED
Part 2

Overview

<u>PART I</u>

- Cross Sections and QED tests
- Accelerator Facilities + Experimental Results and Tests

<u>PART II</u>

- Tests of QED in Partice Decays
- QED Radiative Effects (Bremsstrahlung, Higher Order Processes)

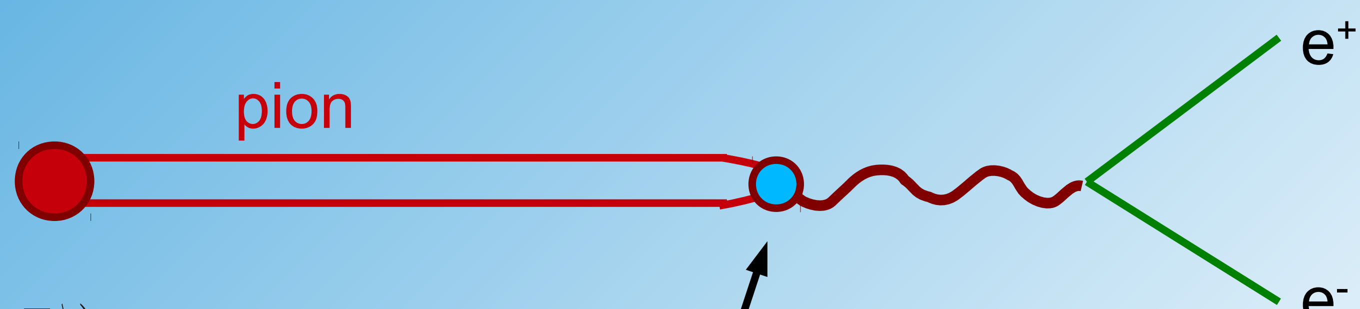
Electromagnetic Decay of Pion

- Pion is the lightest hadron (meson)
- consists of *u* and *d* quarks (isospin triplet)
- No strong or weak decay possible

Can the electromagnetic decay of the pion be described by QED?

<u>Complication</u>: quarks are involve large QCD corrections!

<u>Solution</u>: introduce pion form factor



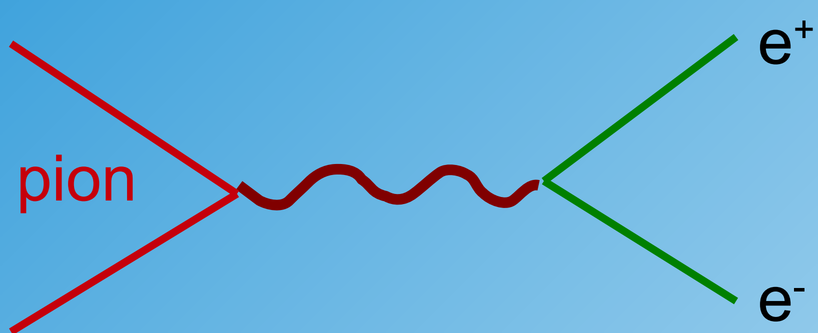
$$|\pi^0\rangle = \frac{1}{\sqrt{2}}\left(|u\bar{u}\rangle - |d\bar{d}\rangle\right)$$

Test Pion Branching Ratios

- dominant decay: $B(\pi^0 \to \gamma\gamma) = 98.823\ \%$
- radiative decay: $B(\pi^0 \to e^+e^-\gamma) = 1.174\ \%$
- 2-prong decay: $B(\pi^0 \to e^+e^-) = 6.46 \times 10^{-8}$

Can QED describe this surprisingly small branching ratio?

<u>Vector Currents</u>: $j^{\mu}_{elm} = \bar{v}(x)\gamma^{\mu}u(x)$ per photon/fermion vertex



- Vector currents conserve helicity.
- Resulting spin should be $J(\pi^0)=1$
- But pion is a Pseudo-scalar $J(\pi^0)=0$
  - ➔ contradiction → helicity suppression

Polarisation of helicity state is given by fermion velocity:

$$\langle\lambda\rangle = \pm\beta$$ for right/left chiralities

<u>Resulting suppression factor:</u>

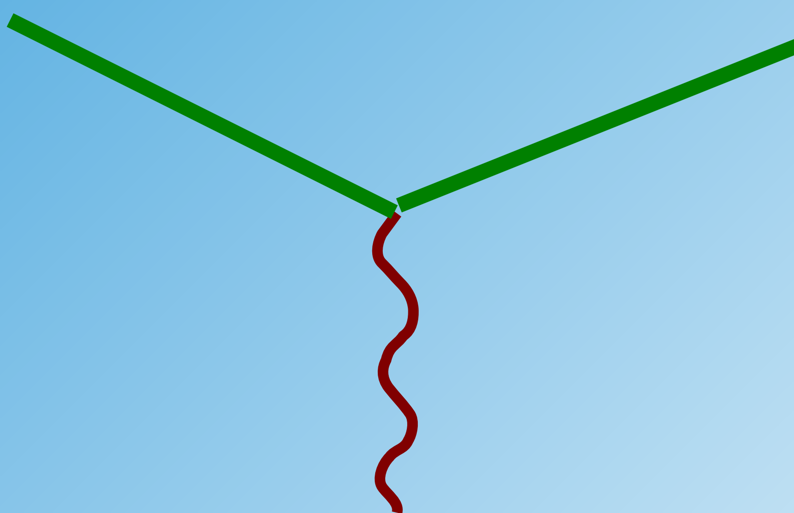
$$\frac{1}{16}\left(1-\beta_q^2\right)\left(1-\beta_e^2\right) = \frac{1}{16\,\gamma_q^2\gamma_q^2} = \frac{m_q^2 m_e^2}{m_\pi^4}$$

no scalar couplings!

$\sim 3 \cdot 10^8$

QED Radiative Effects

- **Bremsstrahlung**
- **Higher Order Processes**



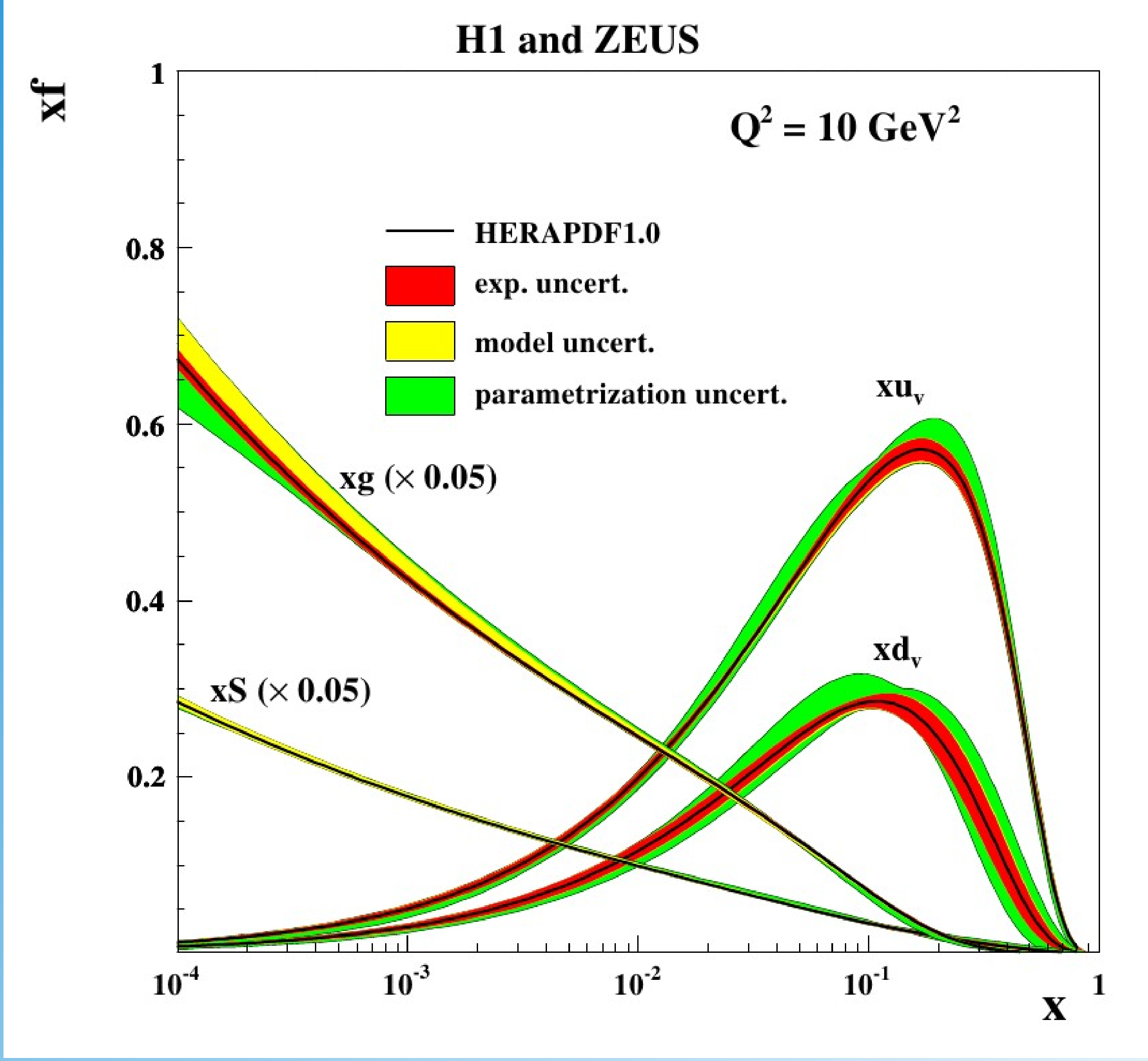
<u>Experimentally very important:</u>

- can be exploited for measurements (luminosity, radiative returns)
- but can also disturb measurements!

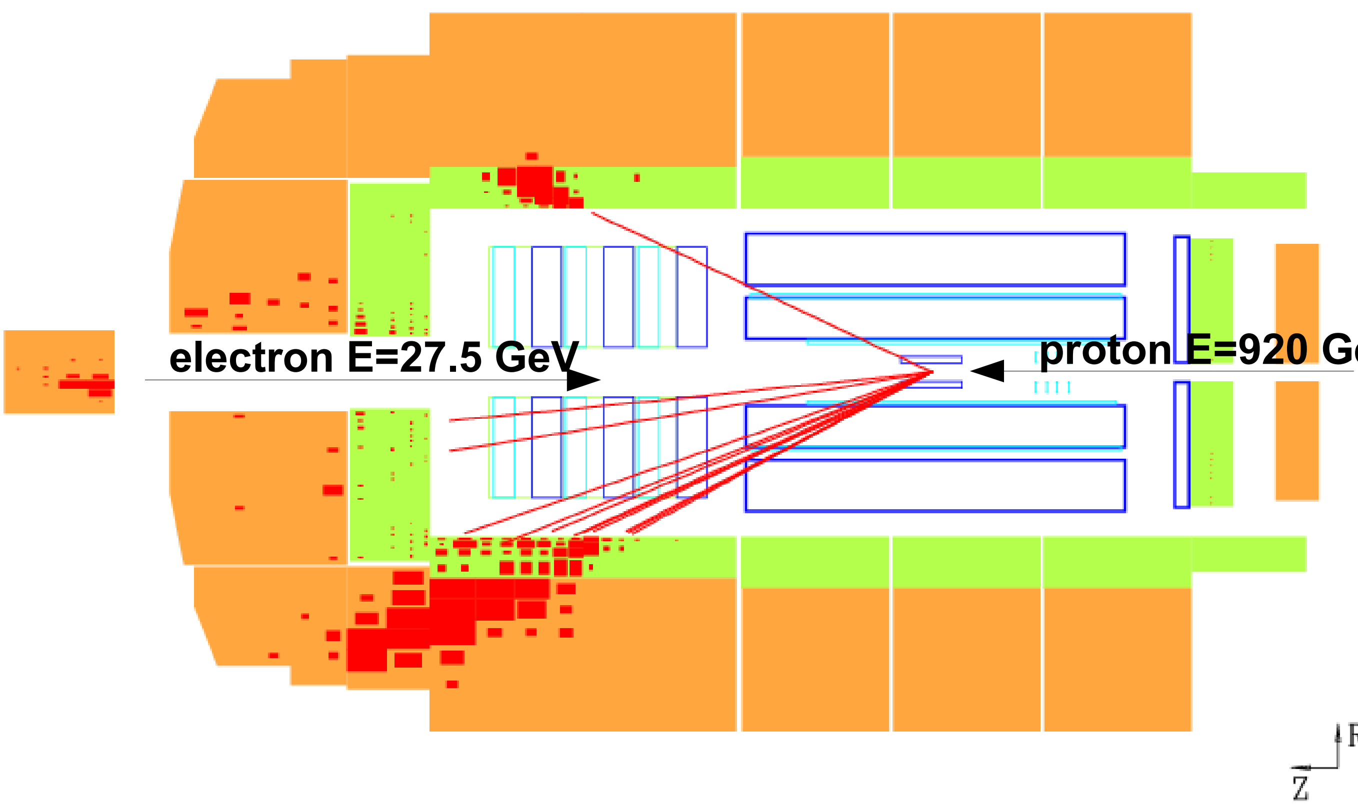
Electron-Proton Collider HERA

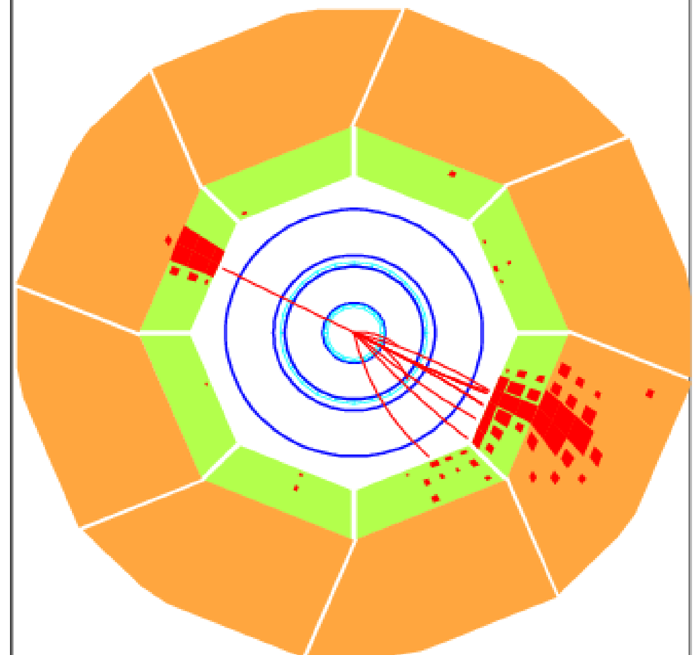
E_e = 26.7 GeV E_p = 920 GeV

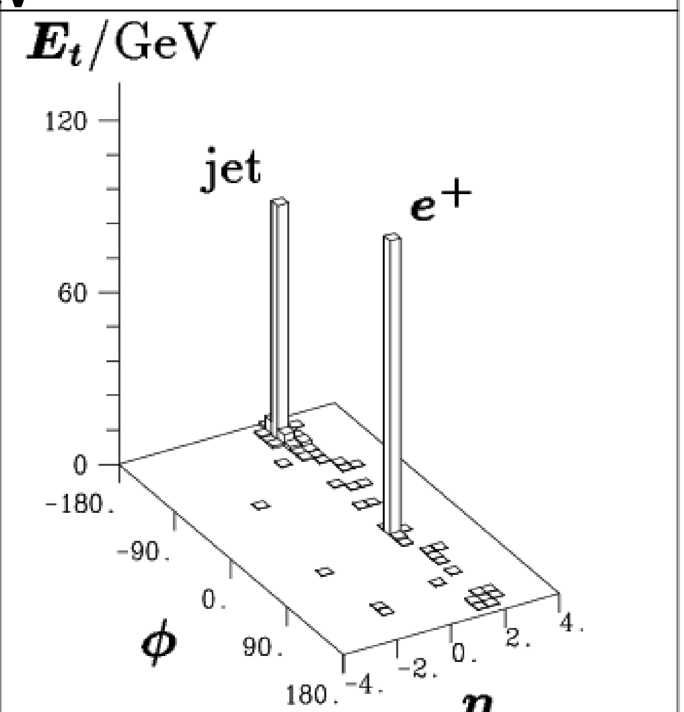
Proton Parton Densities



$$Q^2 = 25030\ \mathrm{GeV}^2, \quad y = 0.56, \quad M = 211\ \mathrm{GeV}$$







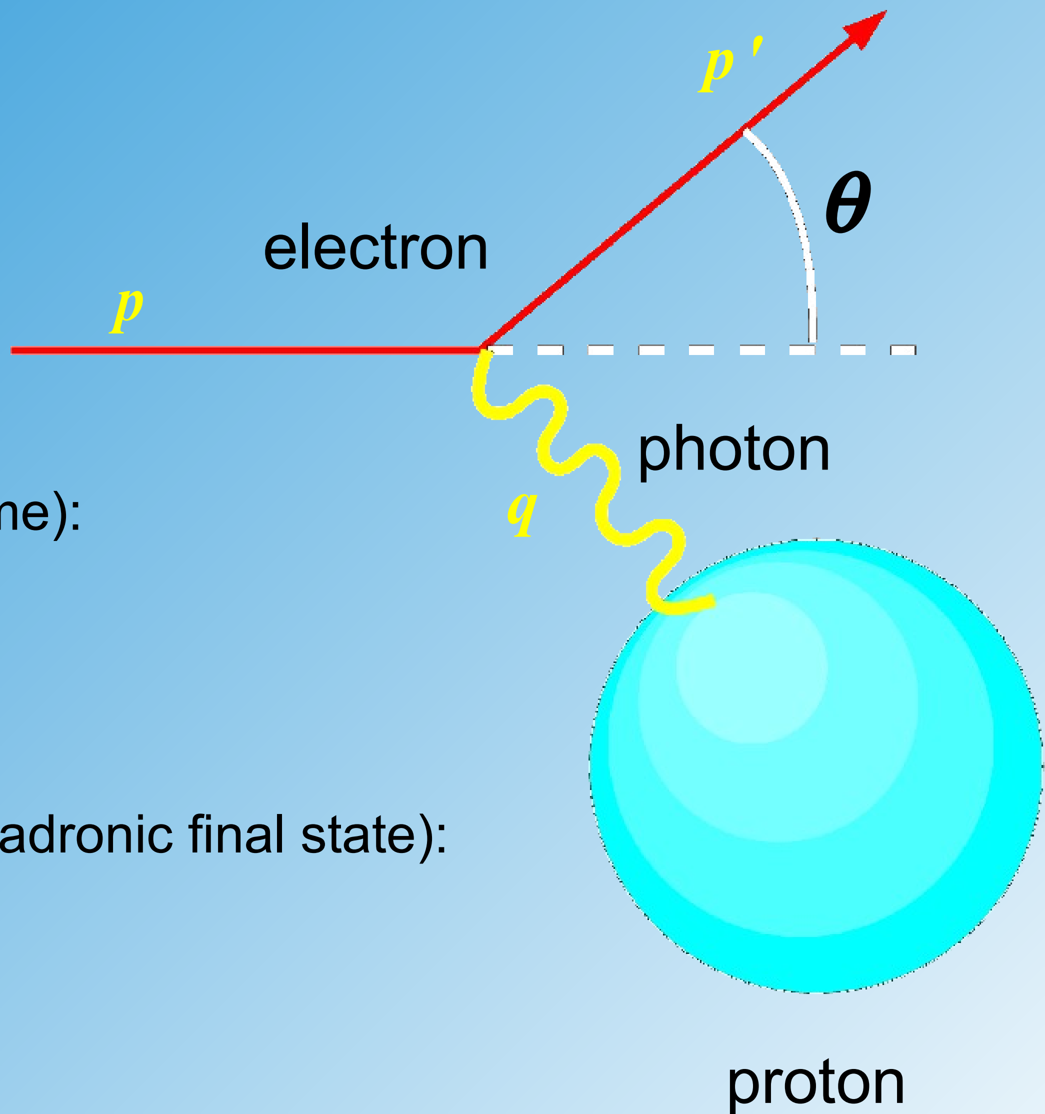
ep → e X (neutral current)

Kinematics Scattering Process

The virtuality of the exchanged photon is given by:

$$Q^2 = -q^2 = -(p - p')^2$$

$$\propto \frac{1}{\sin^4 \theta / 2}$$



Energy transfer (in Proton-rest frame):

$$\nu = E_{Elektron} - E_{Elektron}'$$

Energy of photon-proton system (hadronic final state):

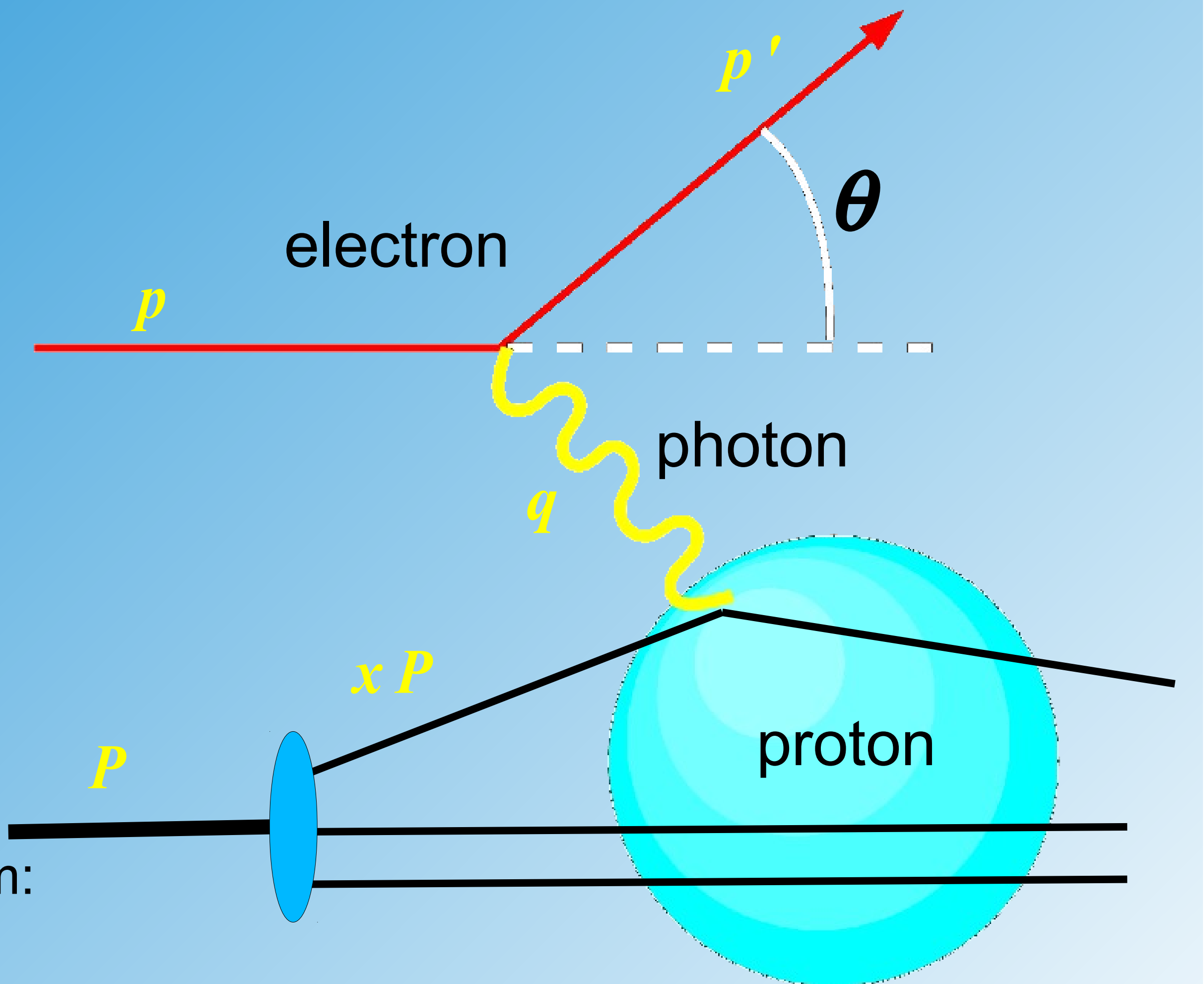
$$W = m_p + \nu$$

Lorentz Invariant Kinematics of Deep Inelastic Scattering Process

The virtuality of the exchanged photon is given by:

$$Q^2 = -q^2 = -(p - p')^2$$

$$\propto \frac{1}{\sin^4 \theta / 2}$$



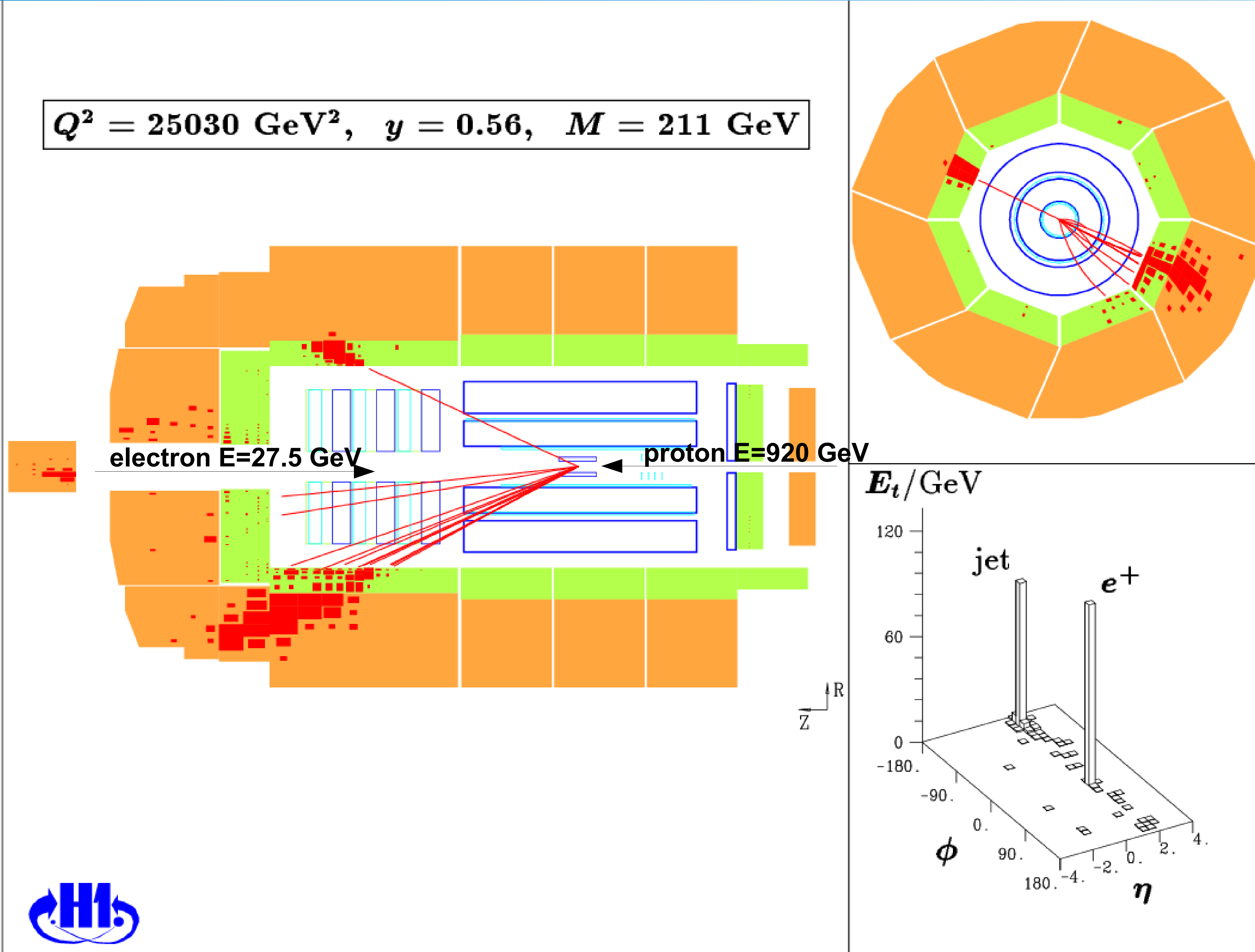
Relative energy loss (inelasticity):

$$y = \frac{\nu}{E_{Elektron}} = \frac{q\,P}{p\,P}$$

relative fraction of parton momentum:

$$x = \frac{q^2}{2\,q\,P} = \frac{Q^2}{s\,y}$$

with cms energy: $s = 2\,p\,P$

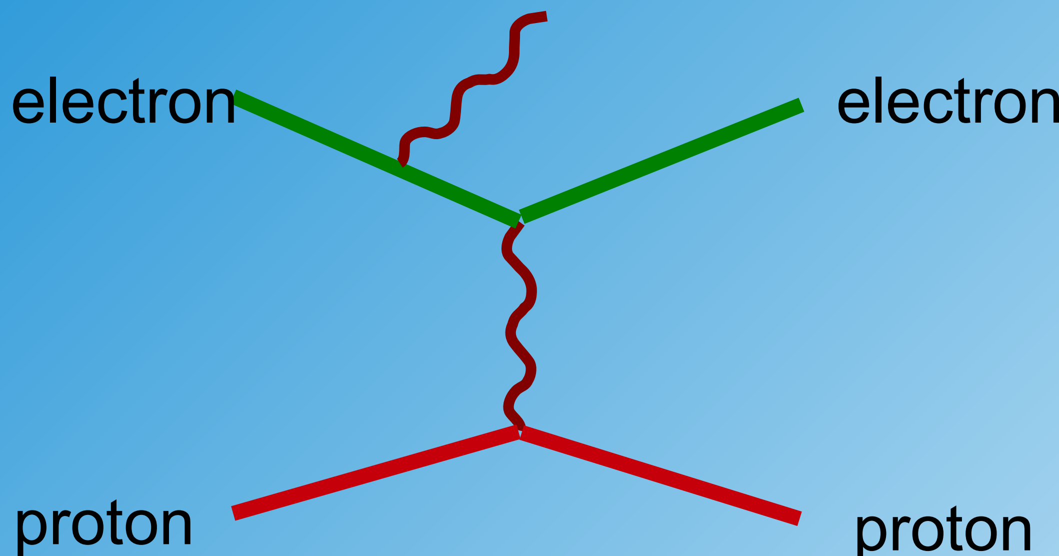

$Q^2 = 25030\ GeV^2,\ \ y = 0.56,\ \ M = 211\ GeV$
electron E=27.5 GeV
proton E=920 GeV
R
Z
E_t/GeV
120
60
0
jet
e^+
-180.
-90.
0.
90.
180.
ϕ
-4.
-2.
0.
2.
4.
η

Determination of Luminosity?

Reference Process

<u>Need process with large cross section</u>

- Bethe Heitler Process: e p → ep γ
- Bremsstrahlung-Process with large cross section (~ 1 barn)!



Note,

- the proton stays intact (elastic)!
- the cross section can only be given for a minimum photon energy! (infrared singularity)

Radiative Effects in Electon-Proton Scattering

from last lecture:

Bremsstrahlung proportional to

$$\frac{1}{(p-k)^2 - m^2}$$

Bremsstrahlung effects are large for particles with <u>low mass</u>

Electron mass: $m_e = 0.511$ MeV → heavy radiation!

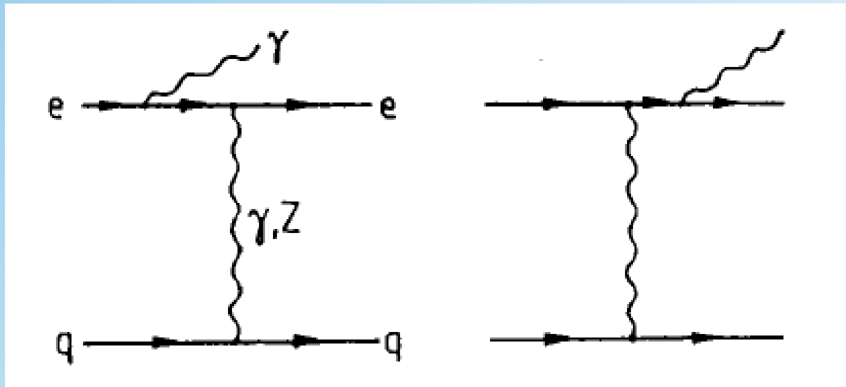
Proton mass: $m_p = 938$ GeV

Radiation is large if

1. photon is **soft**
2. photon is emitted **collinear**

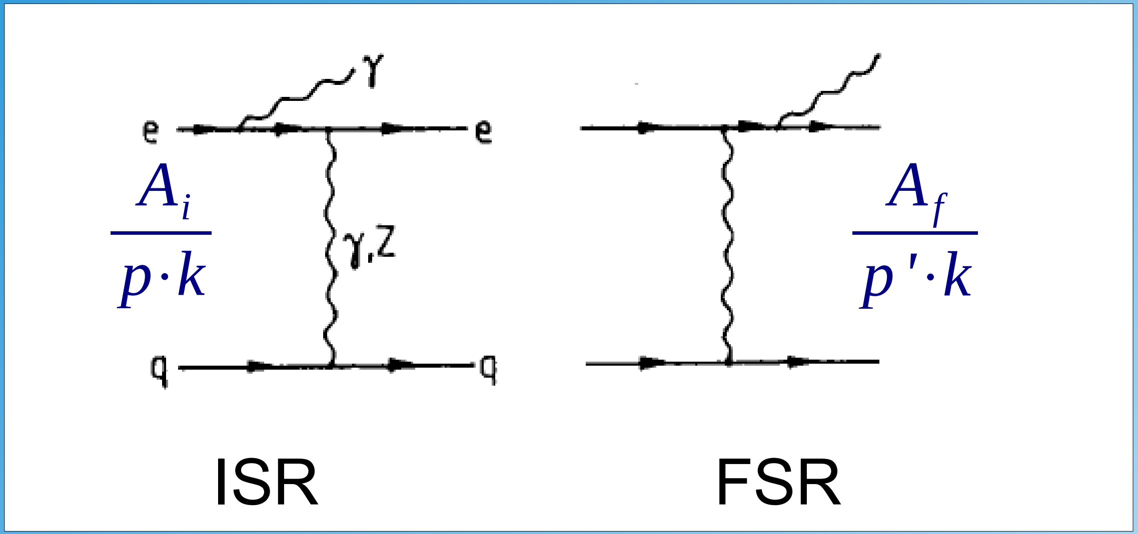
Photons can be emitted from

1. incoming electron (initial state radiation, **ISR**)
2. outgoing electron (final state radiation, **FSR**)

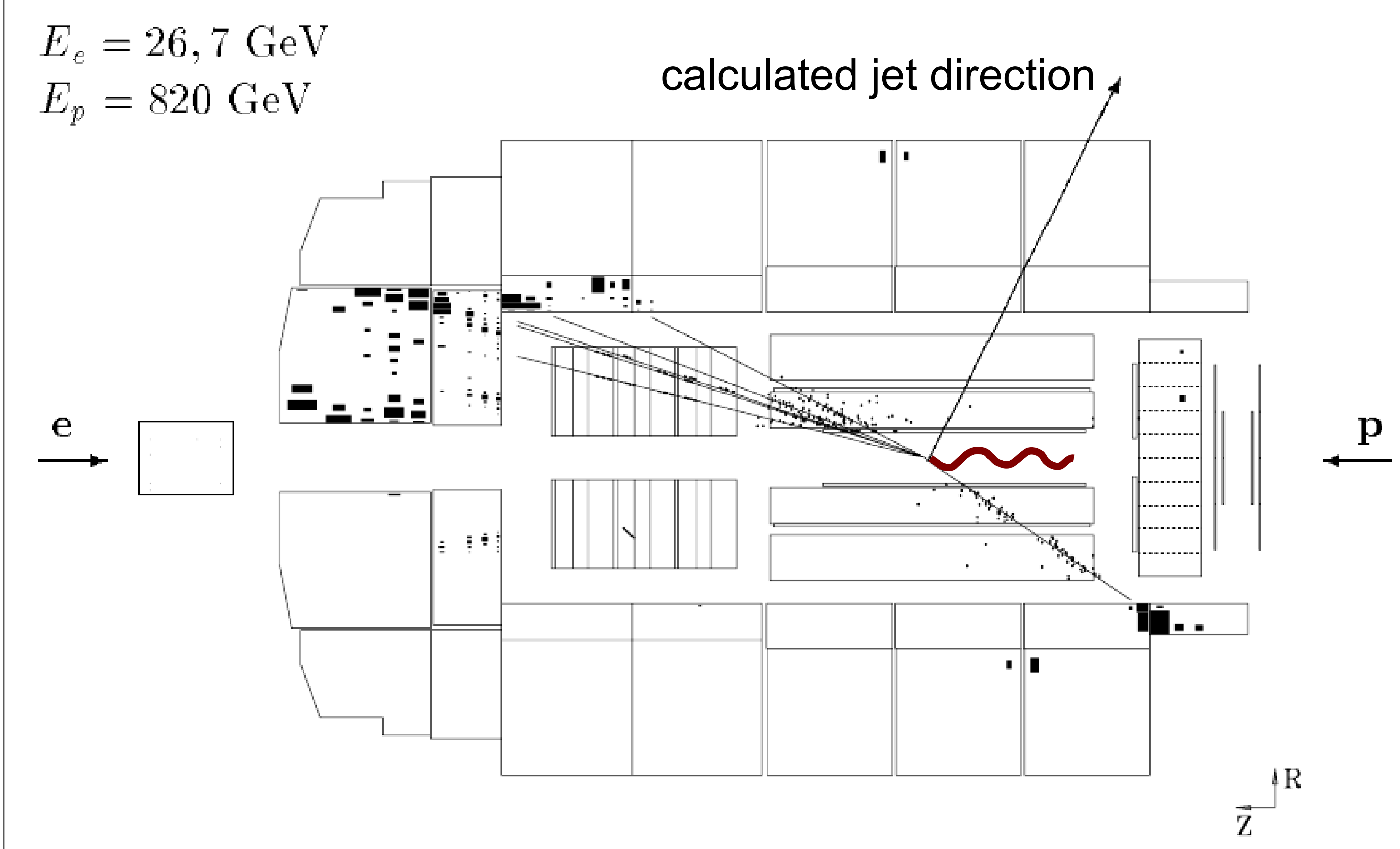


Radiative Poles

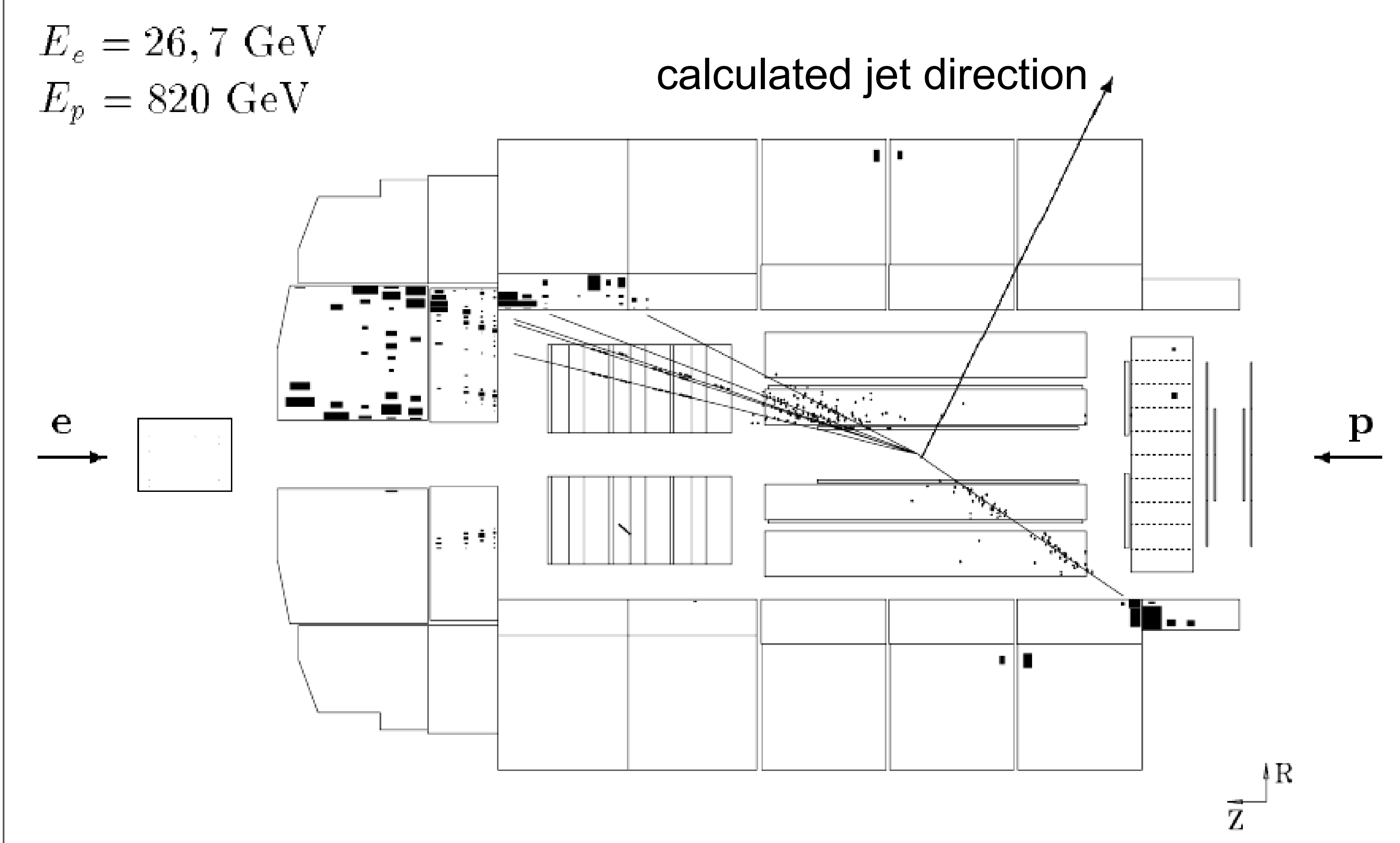
Neglecting the small electron mass, the poles become:



Note ISR reduces the available center of mass energy $s^{1/2}$



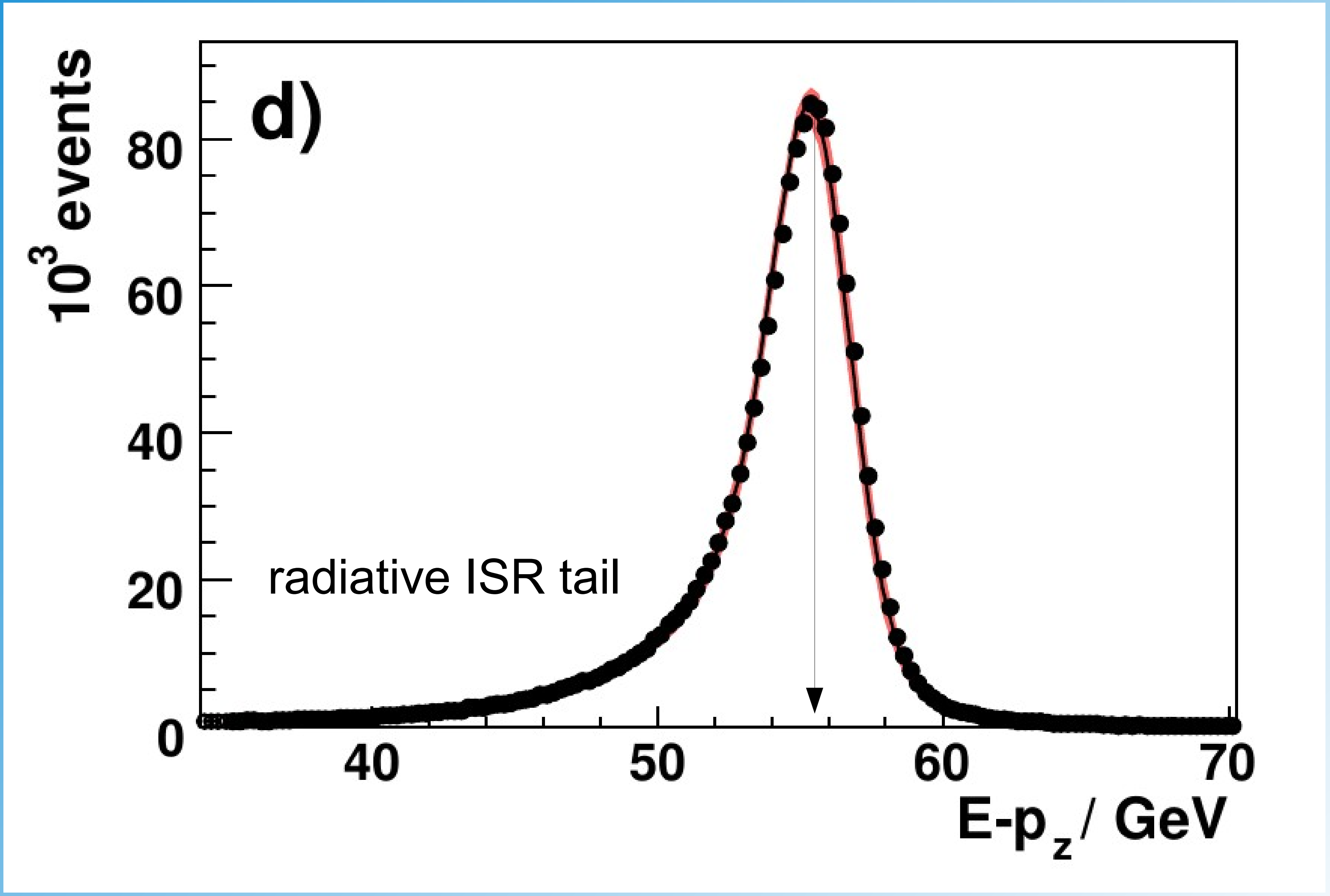
Initial State Radiation?



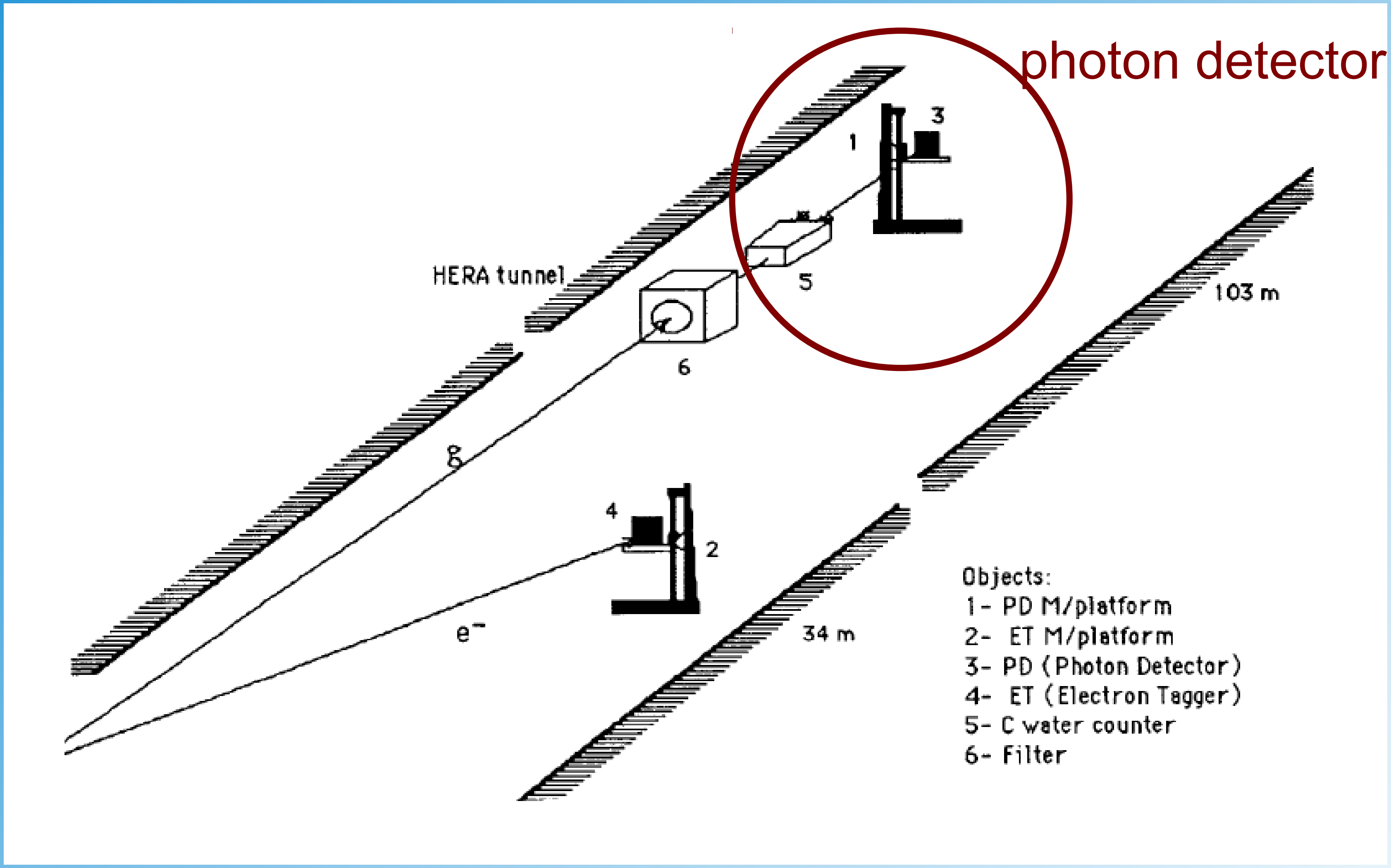
Energy *E* conserved and momentum *p* conserved *E*-p_z conserved

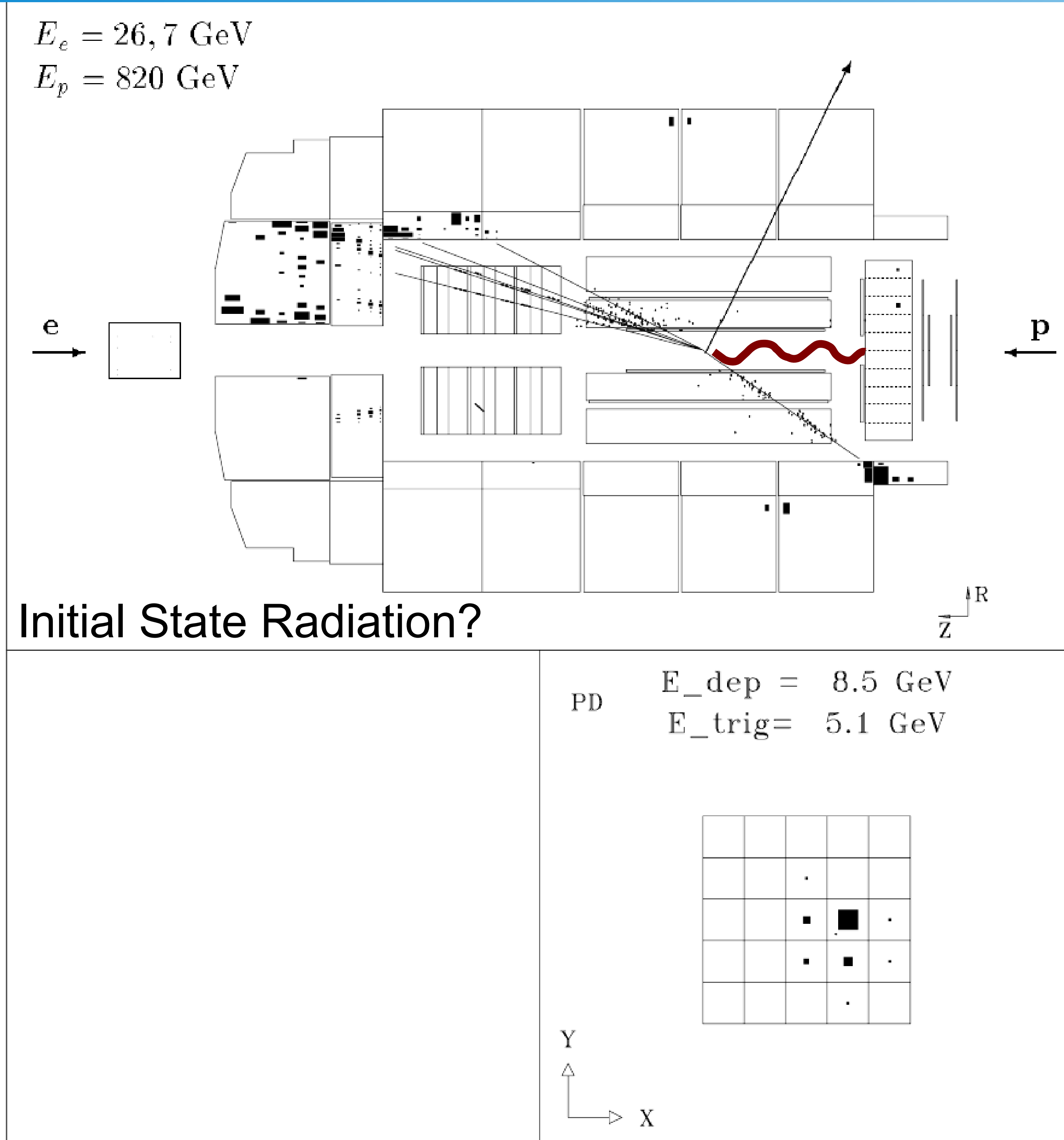
E-p_z (proton)=0   E-p_z (electron) = 2 x 26.7 GeV

Kinematic Reconstruction of ISR



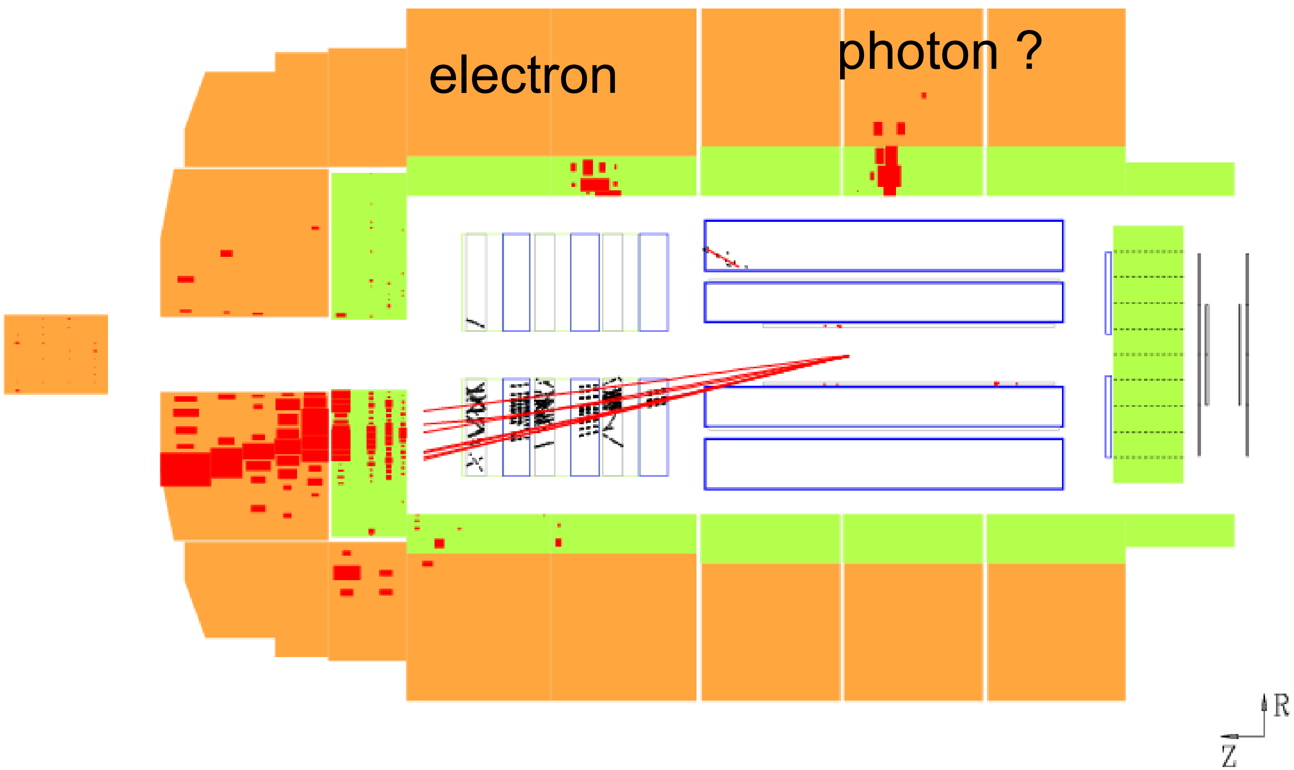
Photondetector at H1

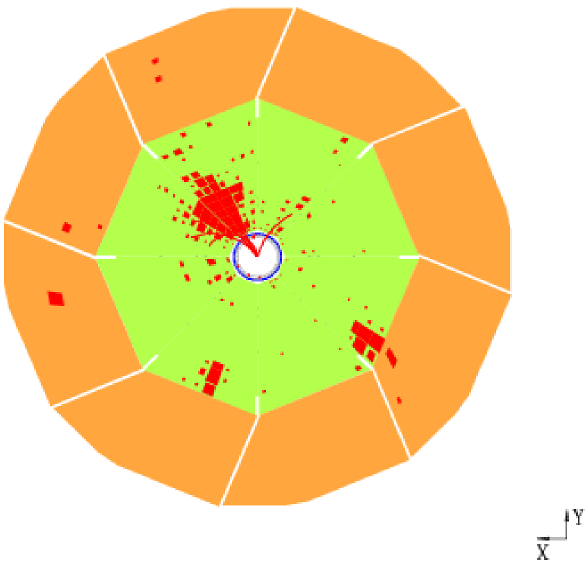


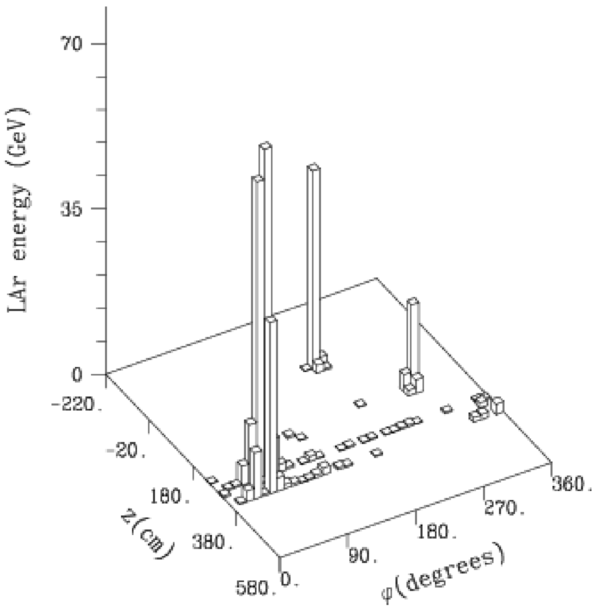


Initial State Radiation?

NC - DIS with hard Photon



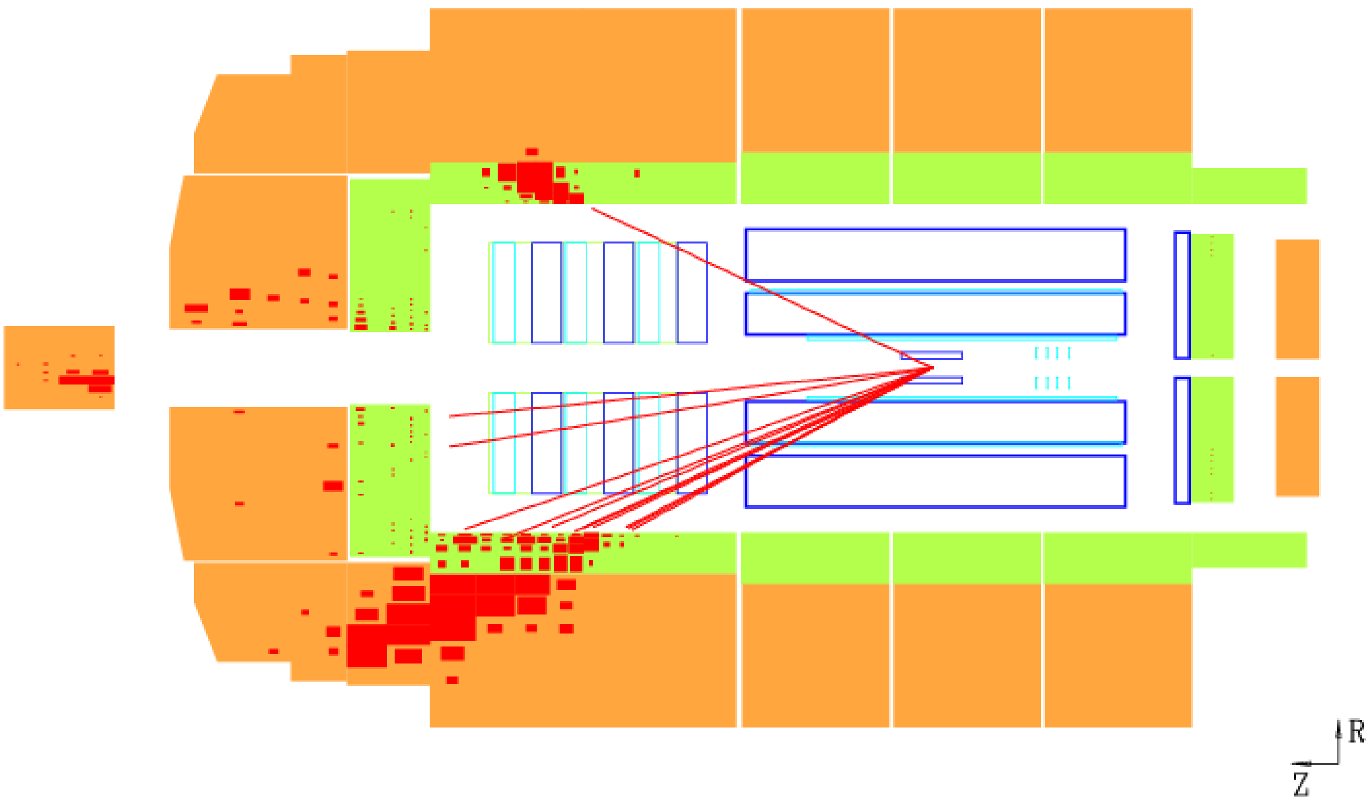


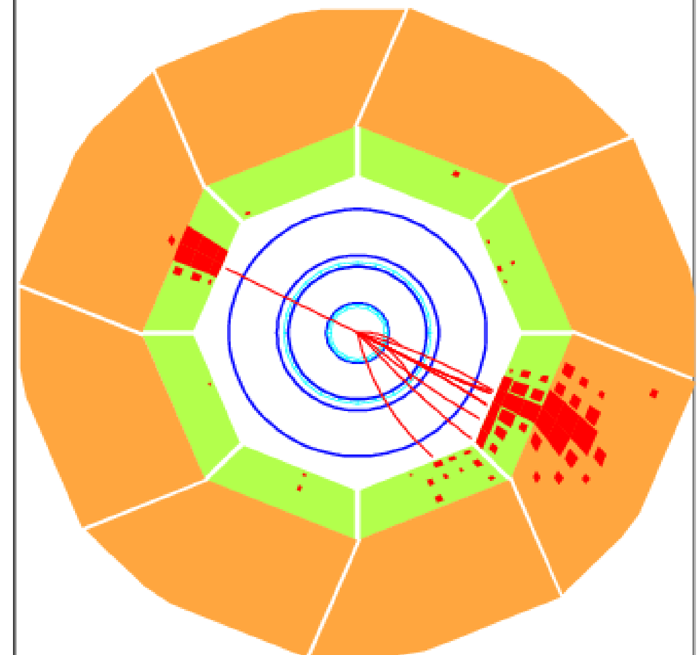


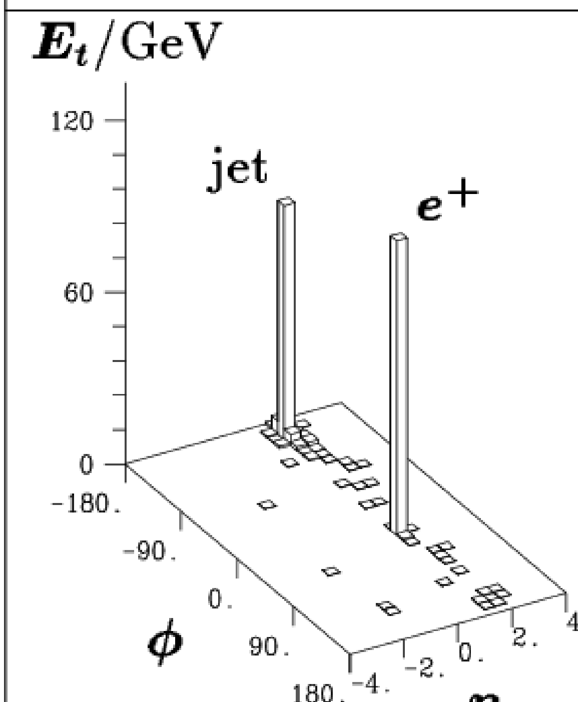
Inner Drift Chamber operational?

Final State Radiation?

$Q^2 = 25030\ \text{GeV}^2, \quad y = 0.56, \quad M = 211\ \text{GeV}$



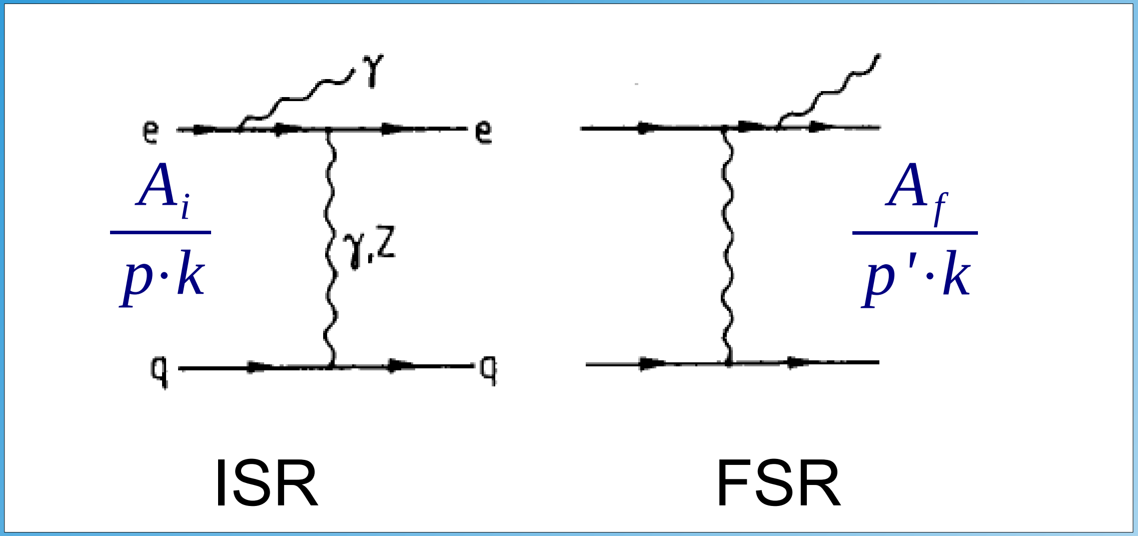




Question: Final State Radiation?

Radiative Poles

Neglecting the small electron mass, the poles become:



Note ISR reduces the available center of mass energy $s^{1/2}$

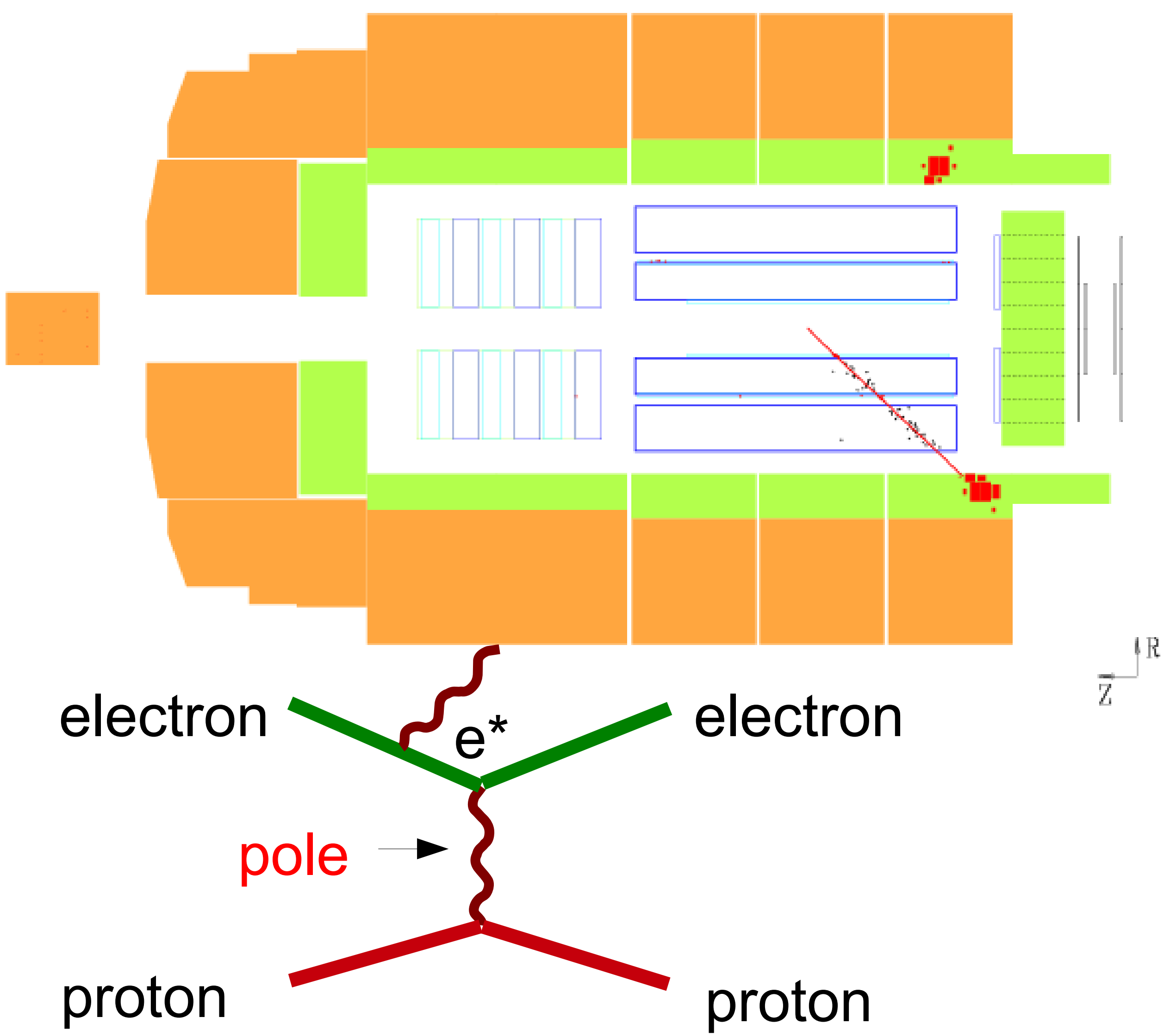
- There is a third pole (Compton events / Wide Angle Bremsstrahlung):

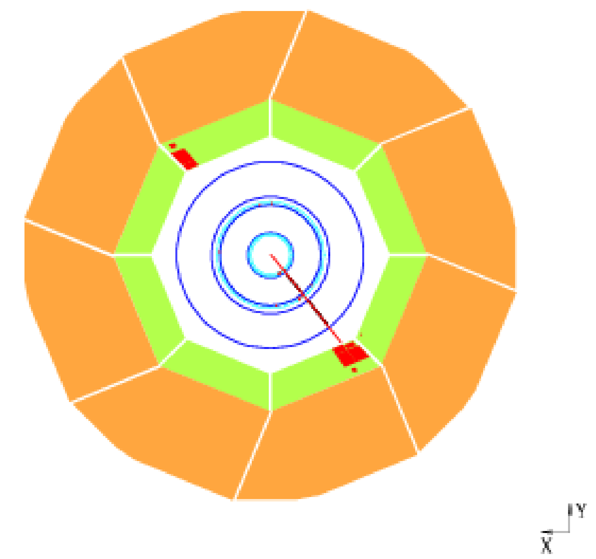
$$\frac{A_c}{(p-p'-k)^2}=\frac{1}{q^2}$$

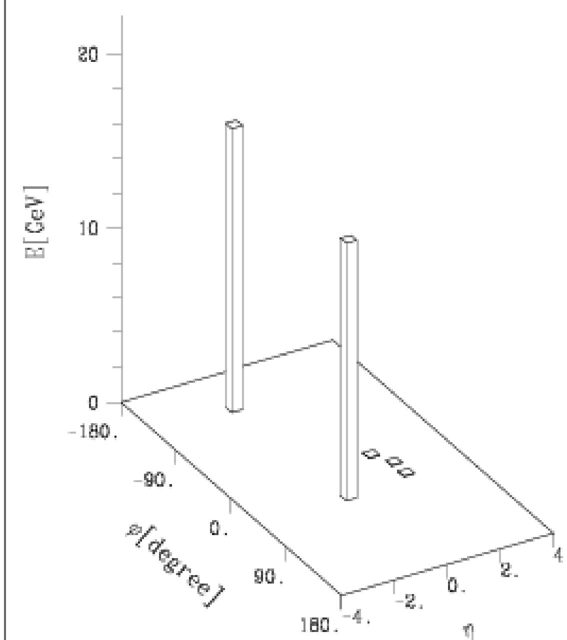
(energy is transferred from the proton to the electron line)

Final State electron and photon can have large opening angles and are p_T balanced

QED Compton

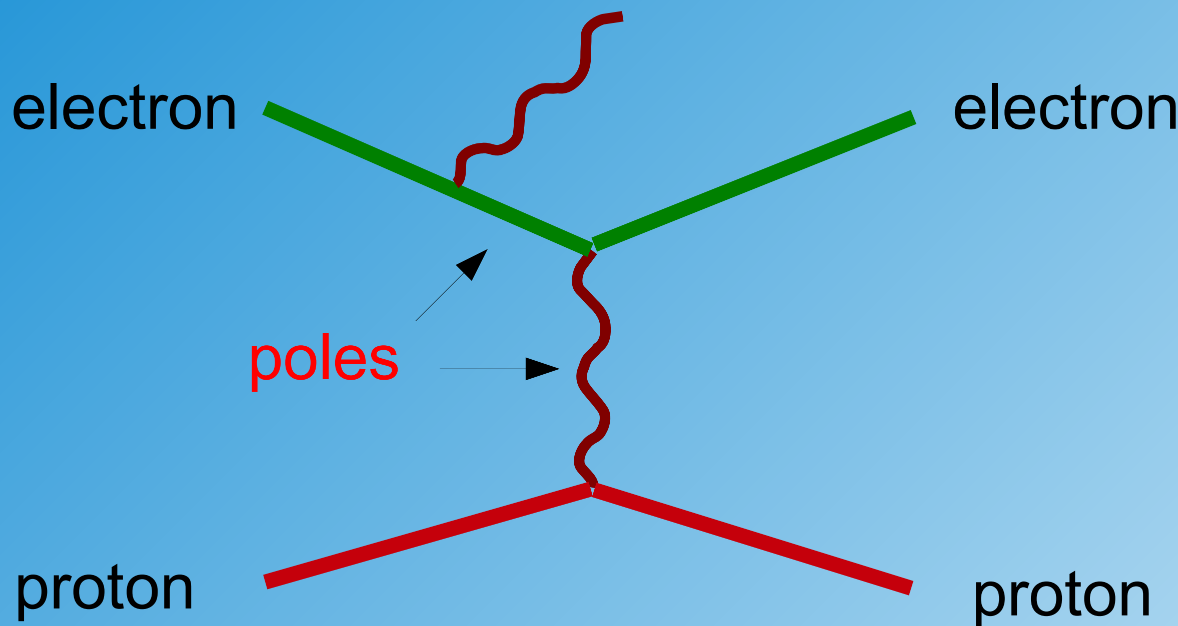

R
Z
electron
e*
electron
pole
proton
proton


Y
X


E[GeV]
20
10
0
-180.
-90.
0.
90.
180.
φ[degree]
-4.
-2.
0.
2.
4.
η

Bethe Heitler Process

diagram:

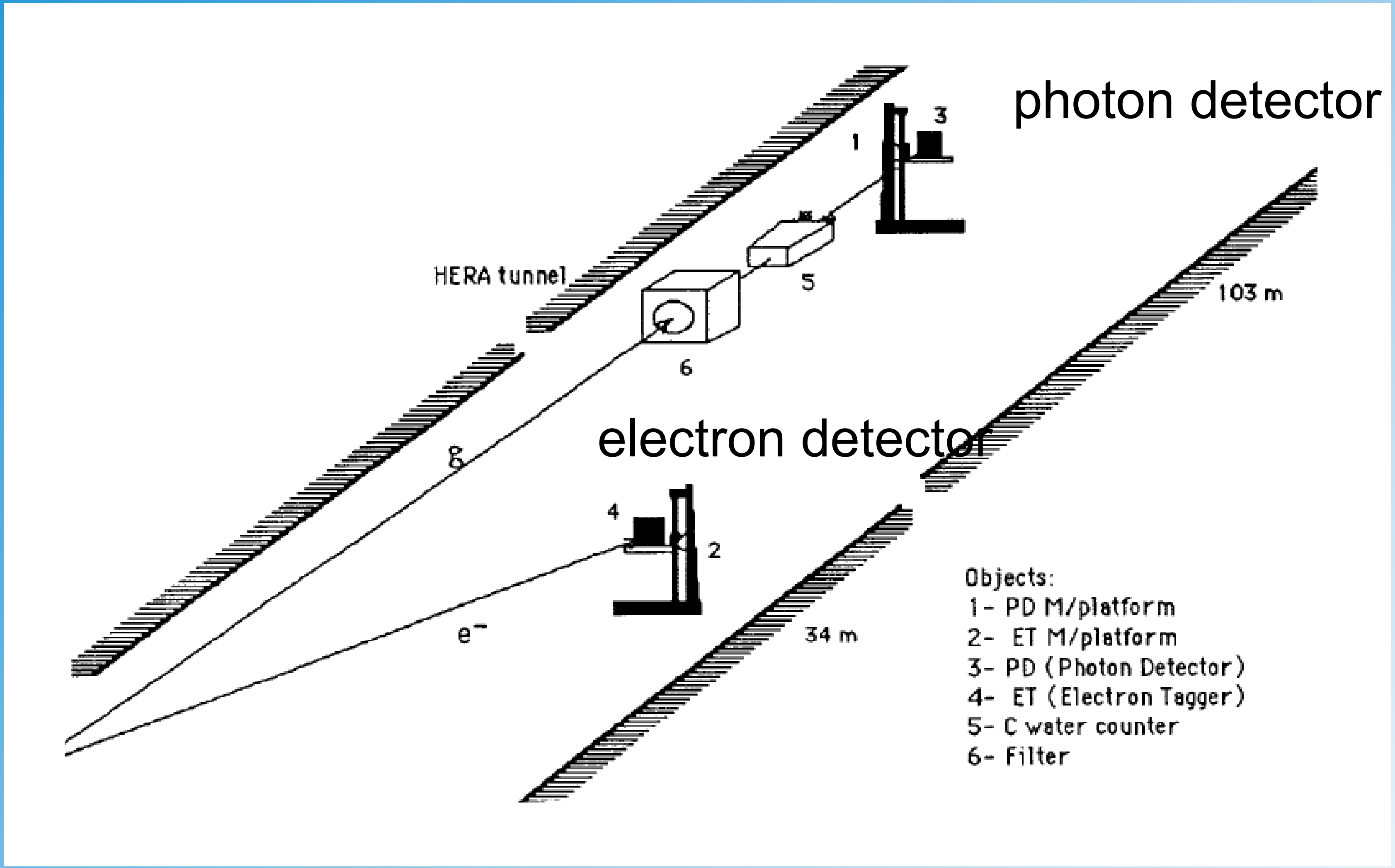


Cross section is largest if electron and photon are scattered (emitted) at small angle (Compton and ISR/FSR poles combined)

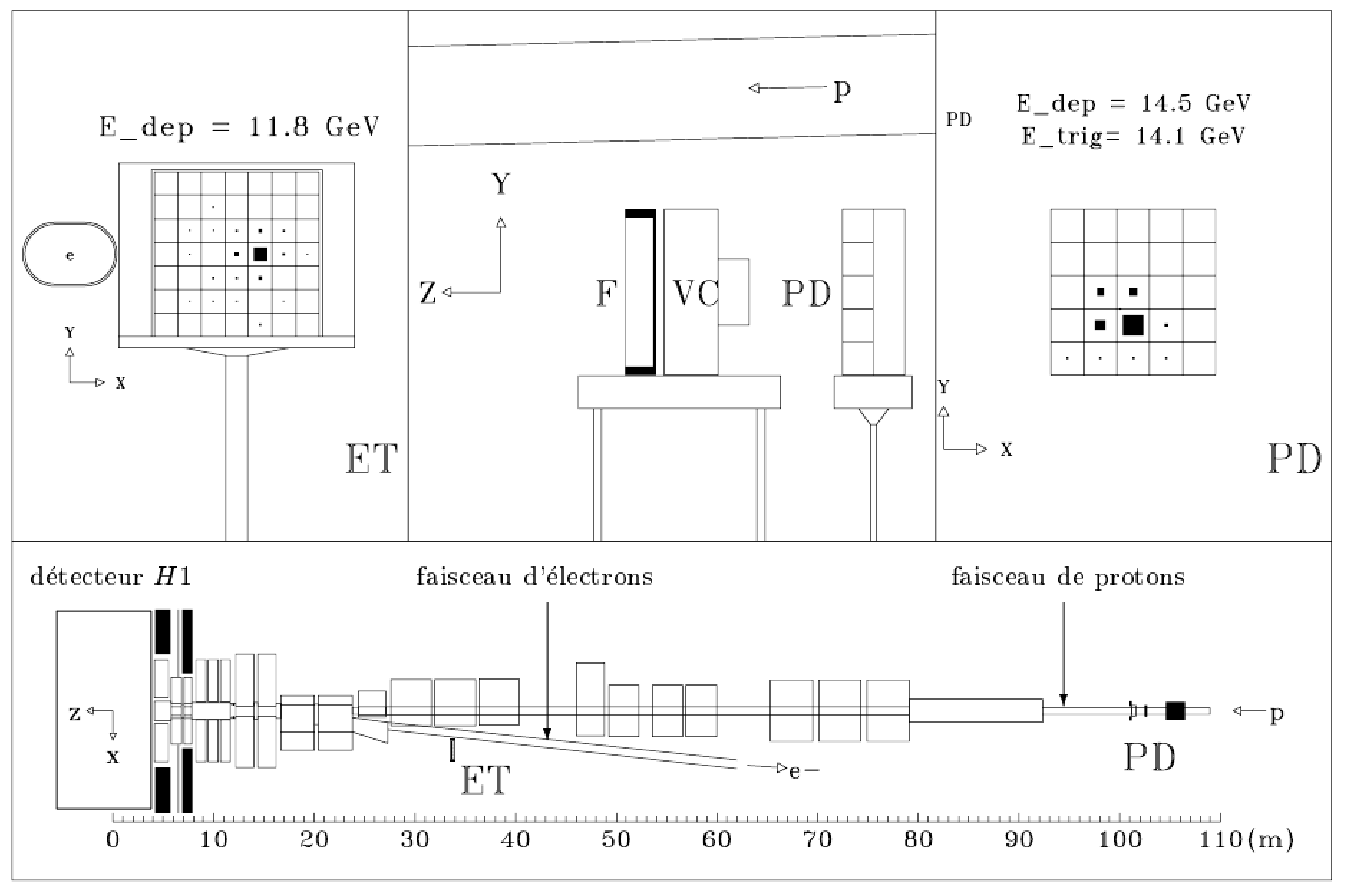
Electron and photon go down the beampipe and are not registered in the central detector

Luminosity Measurement at HERA

dedicated electron and photon detectors



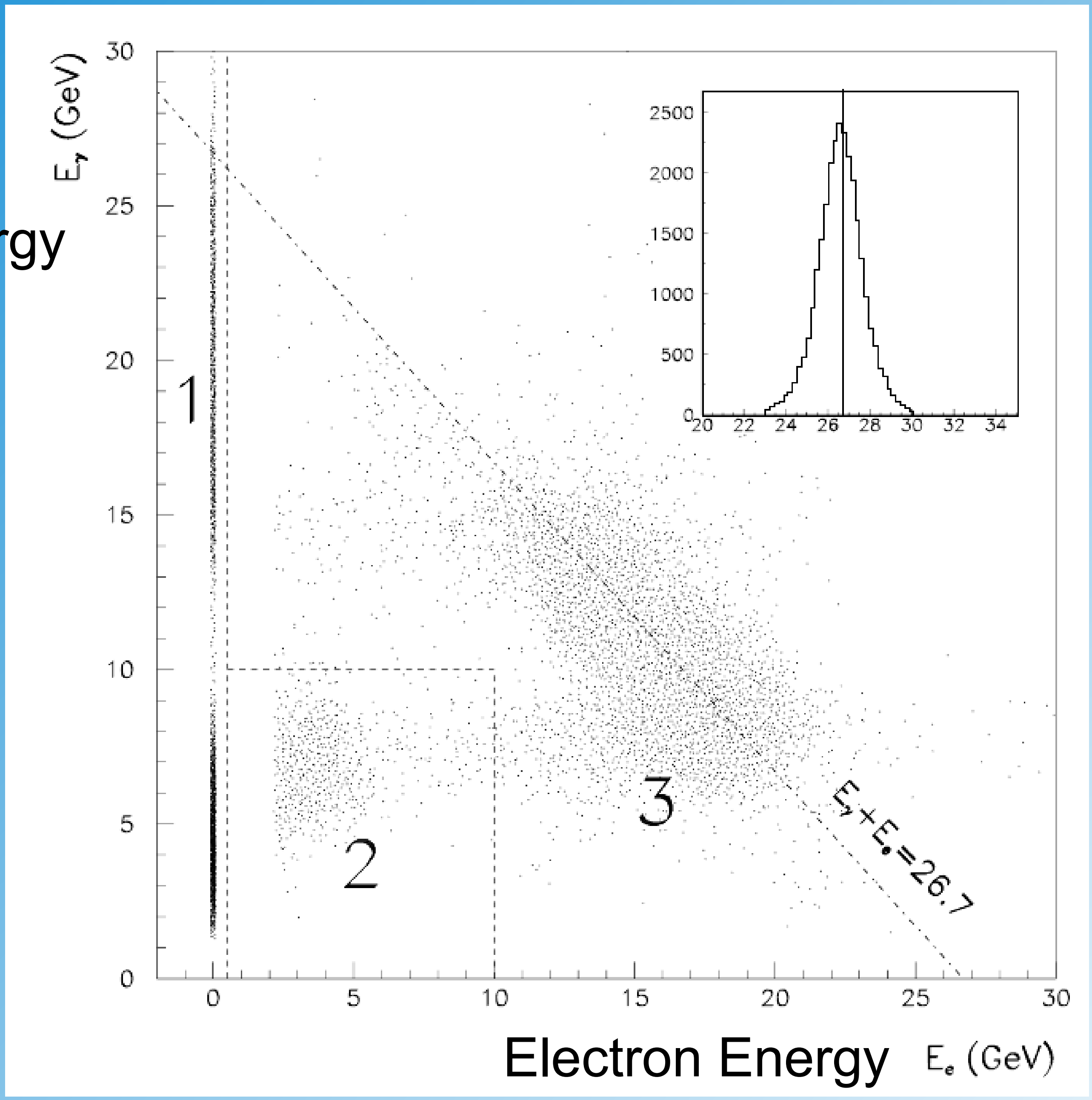
H1 Luminosity Detectors



Method: measure coincidence signal

Coincidence Technique

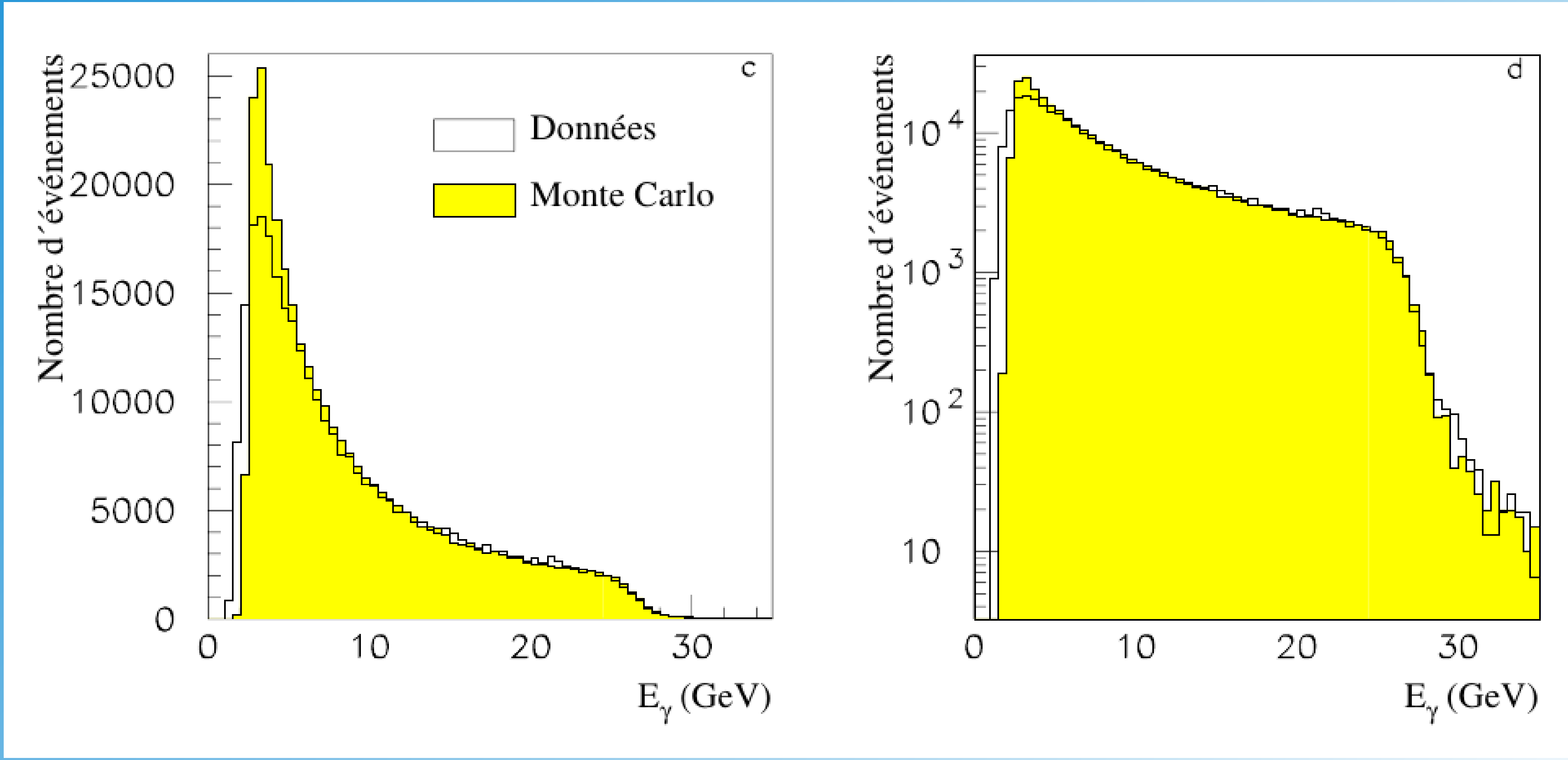
Photon Energy



Electron Energy

Data Monte Carlo Comparison

Photon Detector Energy



Principle of Luminosity Measurement

Integration of the instantaneous luminosity:

$$\mathcal{L}_\mathrm{b} = \int_0^{\Delta T} L_\mathrm{b}(t)\mathrm{d}t$$

Relation between Photon counts, Bethe-Heitler cross section and integrated Luminosity

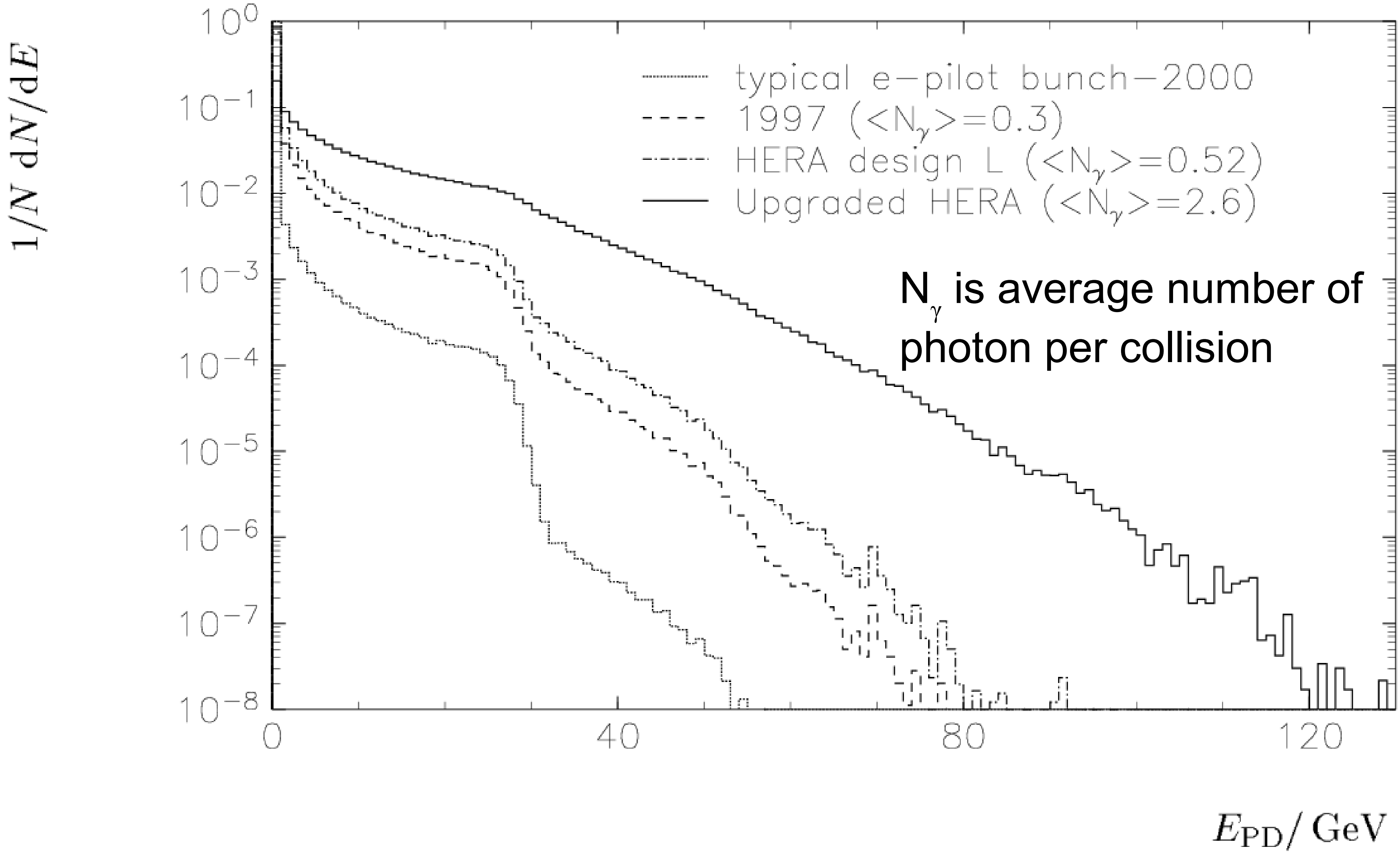
$$N_\gamma(E > \epsilon) = \mathcal{L}_\mathrm{b} A \int_\epsilon^\infty \frac{\mathrm{d}\sigma_\mathrm{BH}}{\mathrm{d}E}\mathrm{d}E$$

Breakdown of Systematic Errors

Item	Online Actual	Offline Actual	Online Upgrade
Theoretical uncertainty	0.5	0.5	0.5
Trigger efficiency	0.1	0.1	–
Background subtraction	0.3	0.2	0.2
Pile up corrections	0.3	0.2	–
Photon detector acceptance	1.0	0.6	0.5
E-scale, resolution	2.0	0.7 – 1.2	0.6
p-bunch satellites	2.0	0.5 – 1.0	

Pileup Correction →

Effect of Pileup



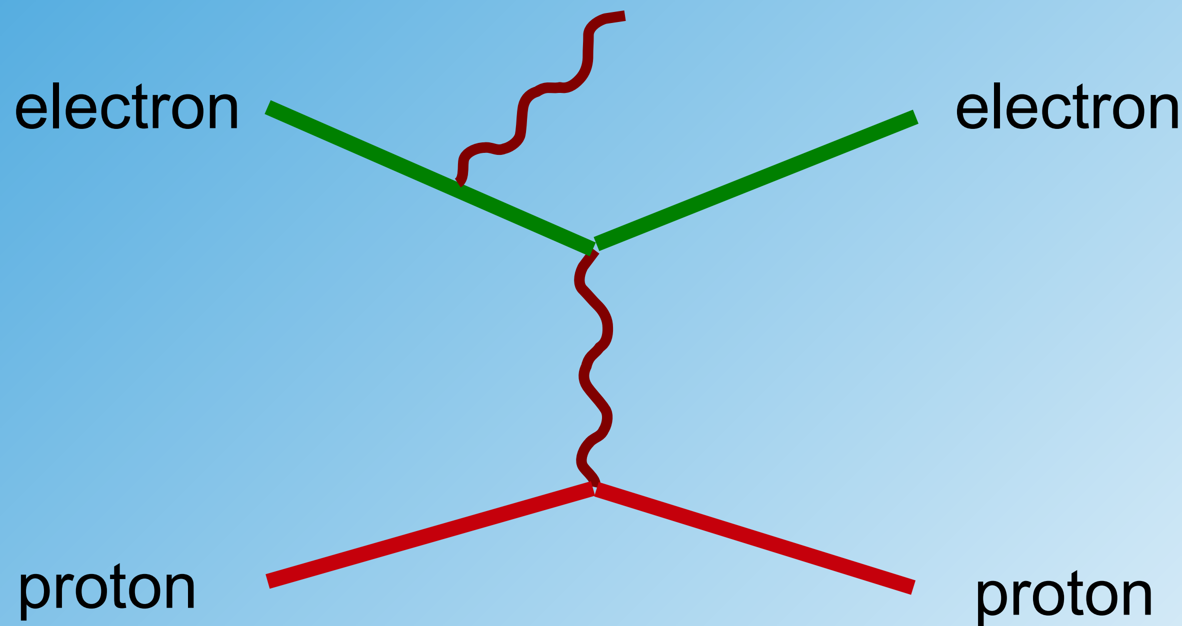
N_γ is average number of photon per collision

QED Compton Analysis

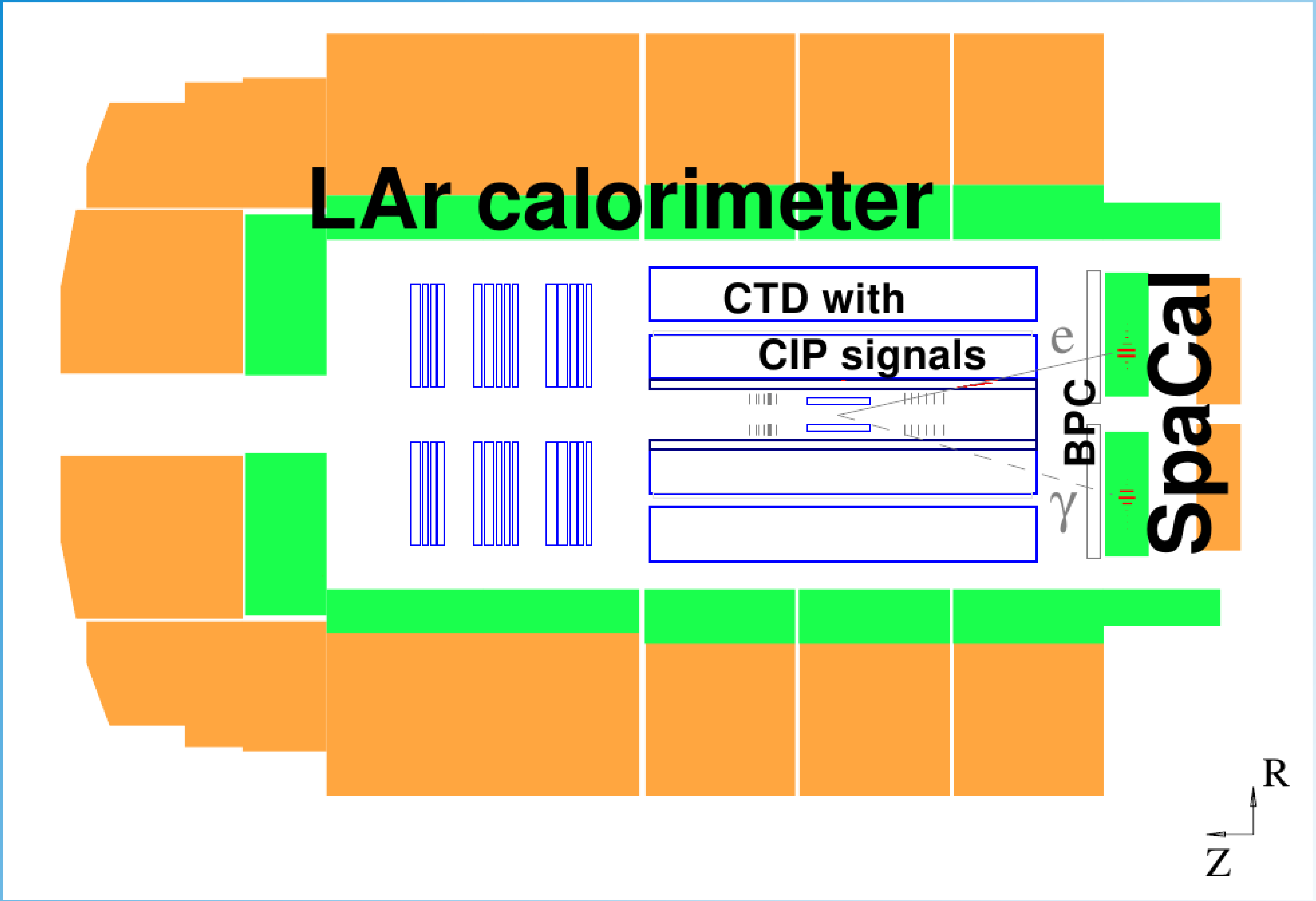
Determination of the Integrated Luminosity in the H1 Experiment at HERA using Elastic QED Compton Events (version of April 18, 2012)

to be published in the next days

H1 Collaboration

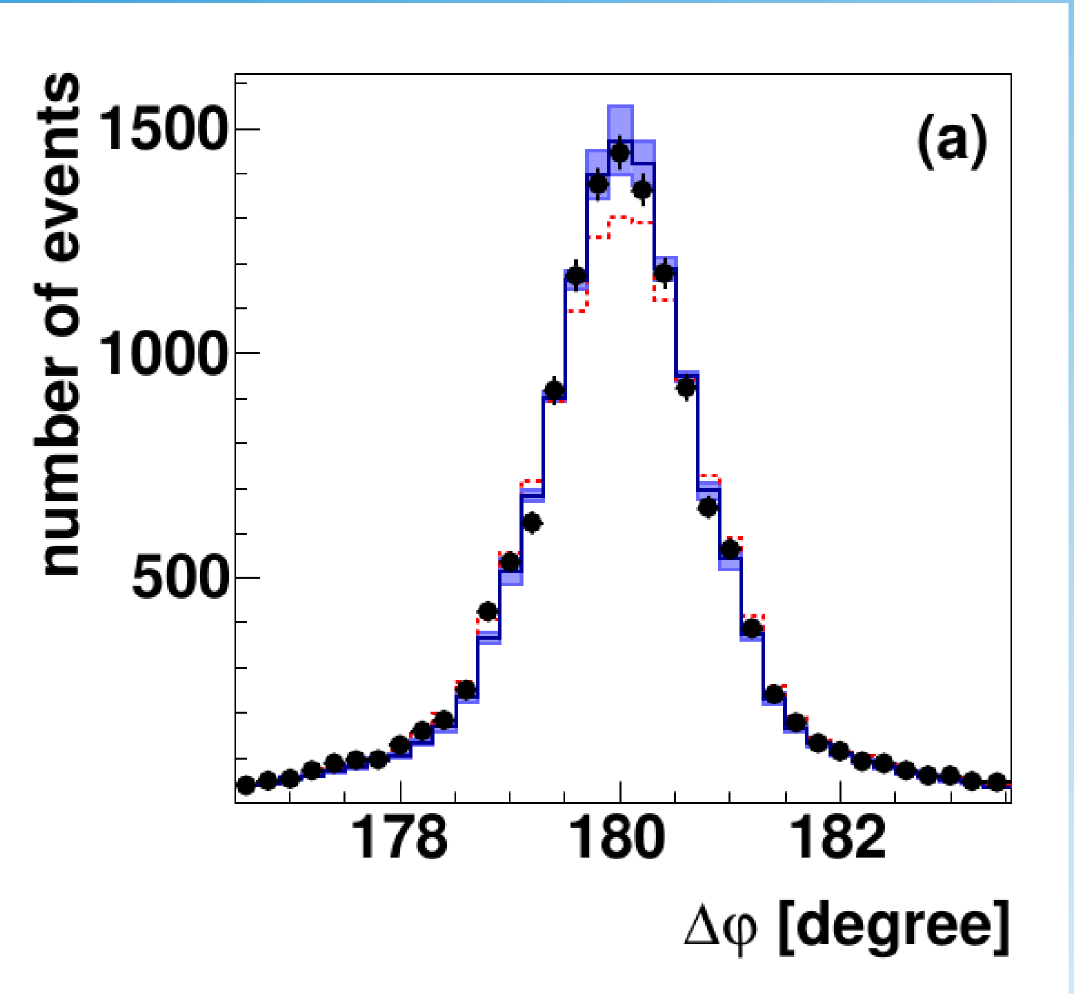


QED Compton Event

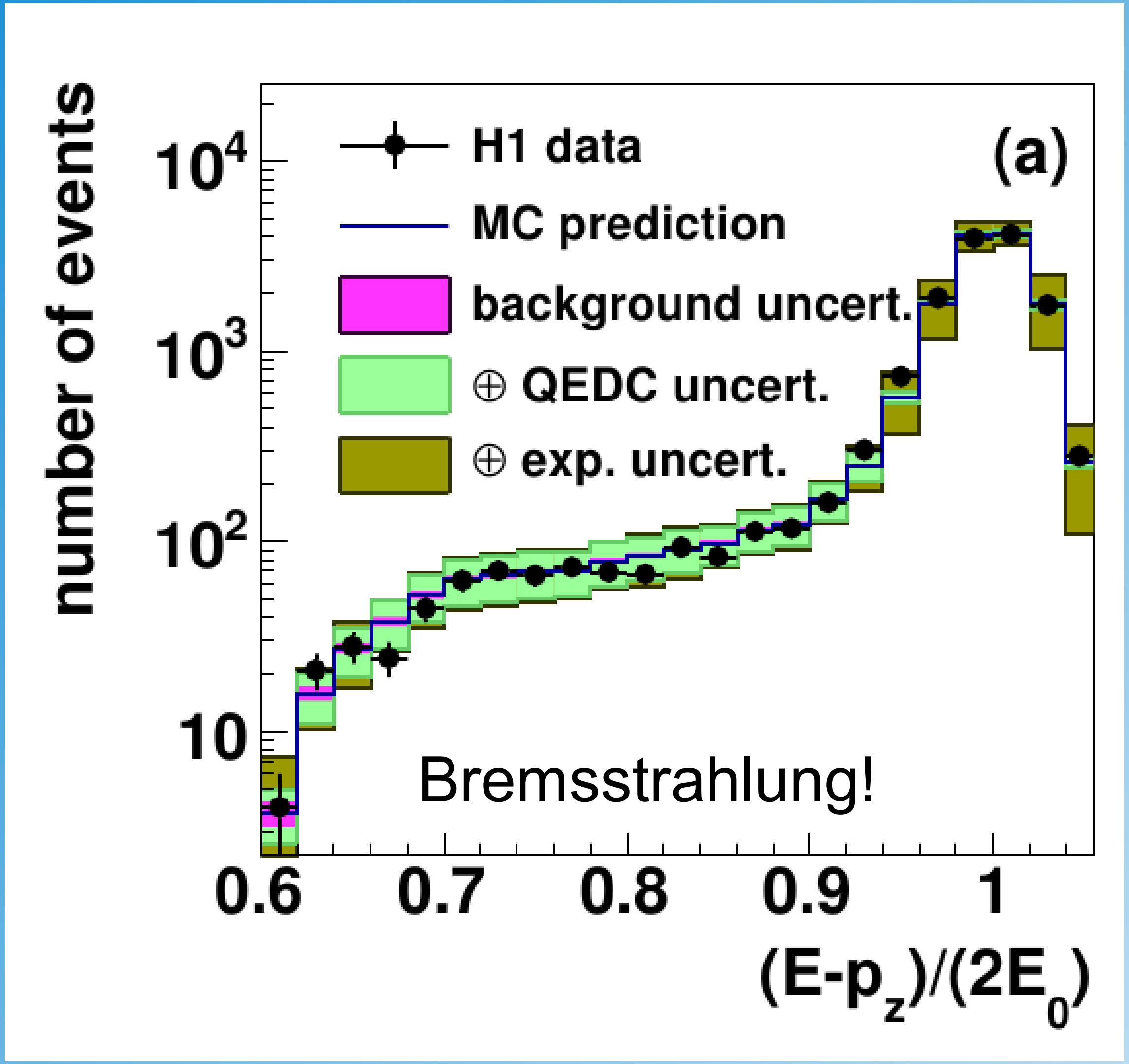


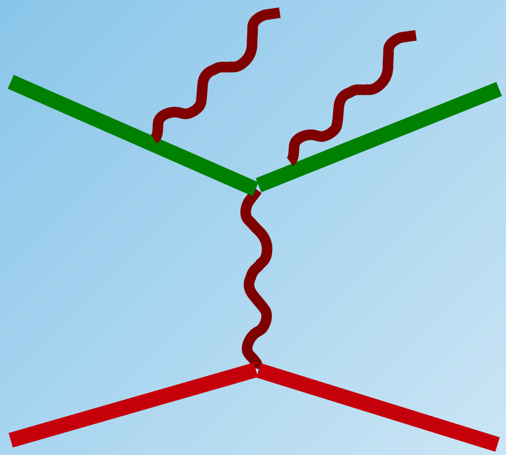
QED Compton Events

Back-to-Back Topology:



Radiative Corrections to Compton Events!





Uncertainties to Compton Analysis

Trigger uncertainties	
Trigger efficiency	0.02%
Veto inefficiency	0.22%
	0.22%

Reconstruction uncertainties	
SpaCal energy scale	0.56%
SpaCal energy resolution	1.10%
SpaCal position resolution	0.34%
CIP efficiency	0.18%
CIP resolution	$< 0.01\%$
Conversion probability	0.33%
Alignment	0.39%
z-vertex distribution	0.24%
SpaCal cluster finder	0.04%
CTD efficiency	0.03%
LAr energy veto	0.05%
	1.41%

Background uncertainties	
non-elastic QEDC	1.11%
elastic DVCS	0.25%
p-dissociative DVCS	0.08%
diffractive ρ	0.03%
diffractive ϕ	$< 0.01\%$
diffractive ω	0.03%
diffractive J/ψ	0.03%
diffractive ψ'	0.08%
diffractive Υ	0.01%
non-resonant diffractive DIS	0.16%
$ep \to epe^+e^-$	0.13%
	1.17%

QEDC theory uncertainties	
Higher order corrections	0.99%
Proton form factor	0.35%
	1.05%

total systematic error [in percent]: $0.22^2 + 1.41^2 + 1.17^2 + 1.05^2 = \mathbf{2.12}^2$

Bremsstrahlung in Charged Currents

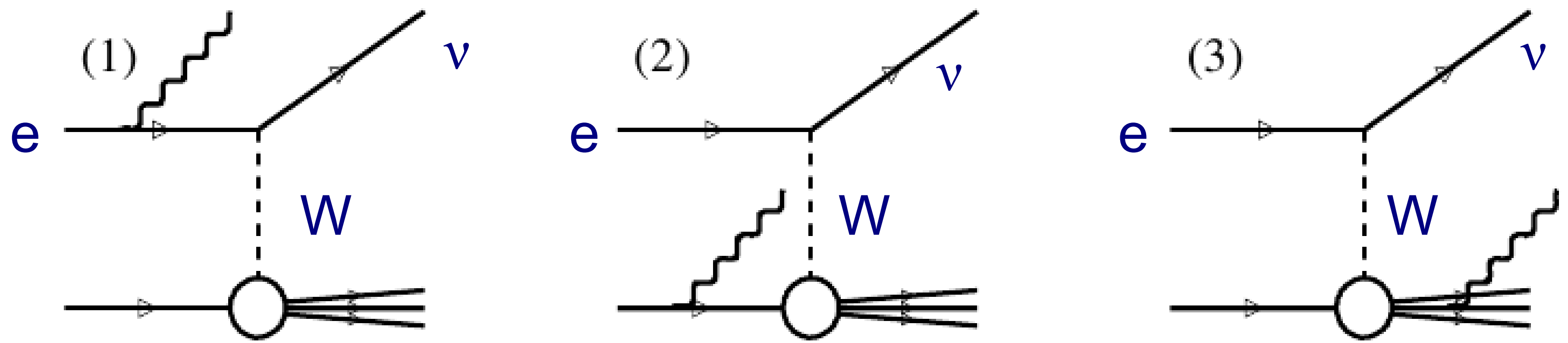
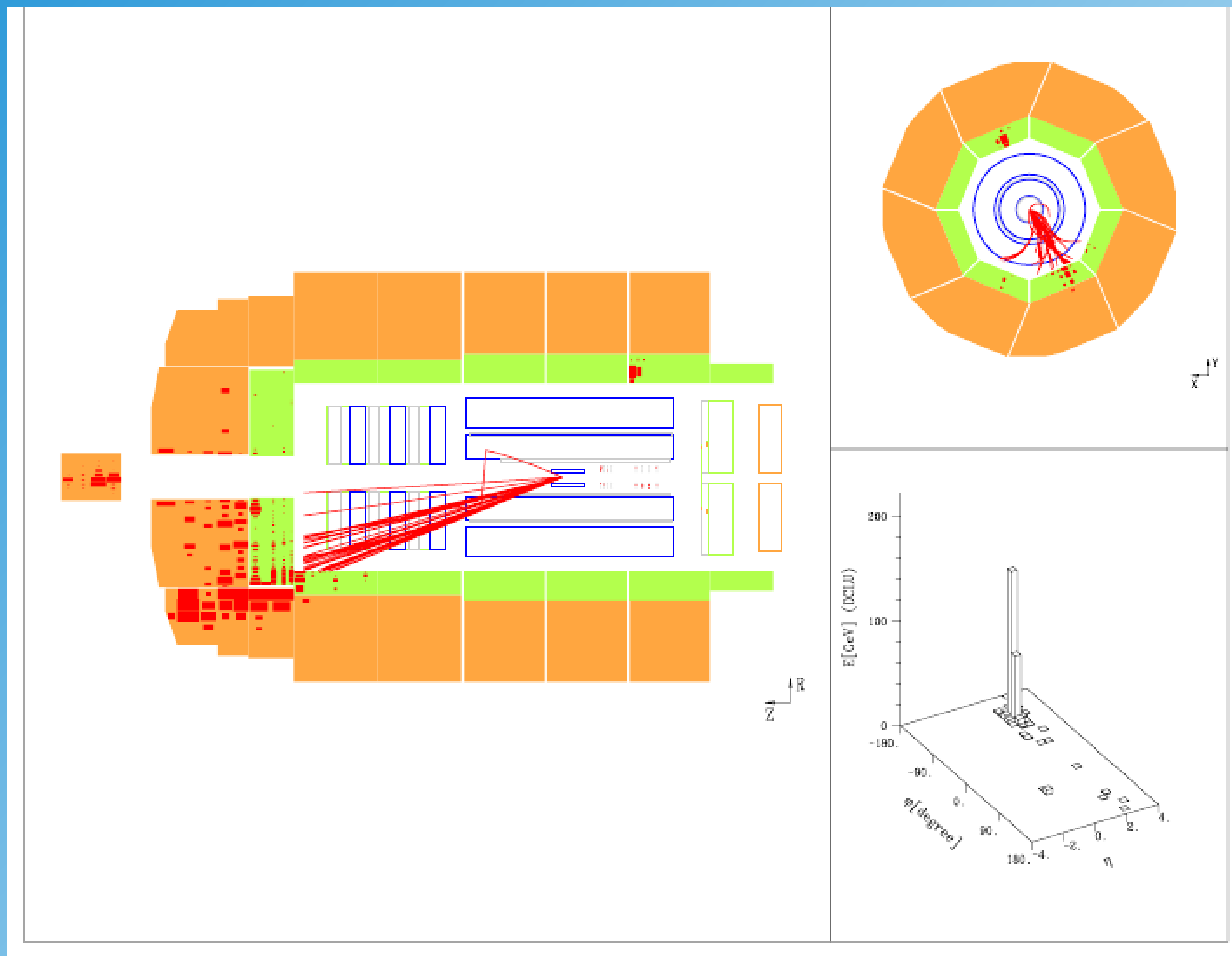


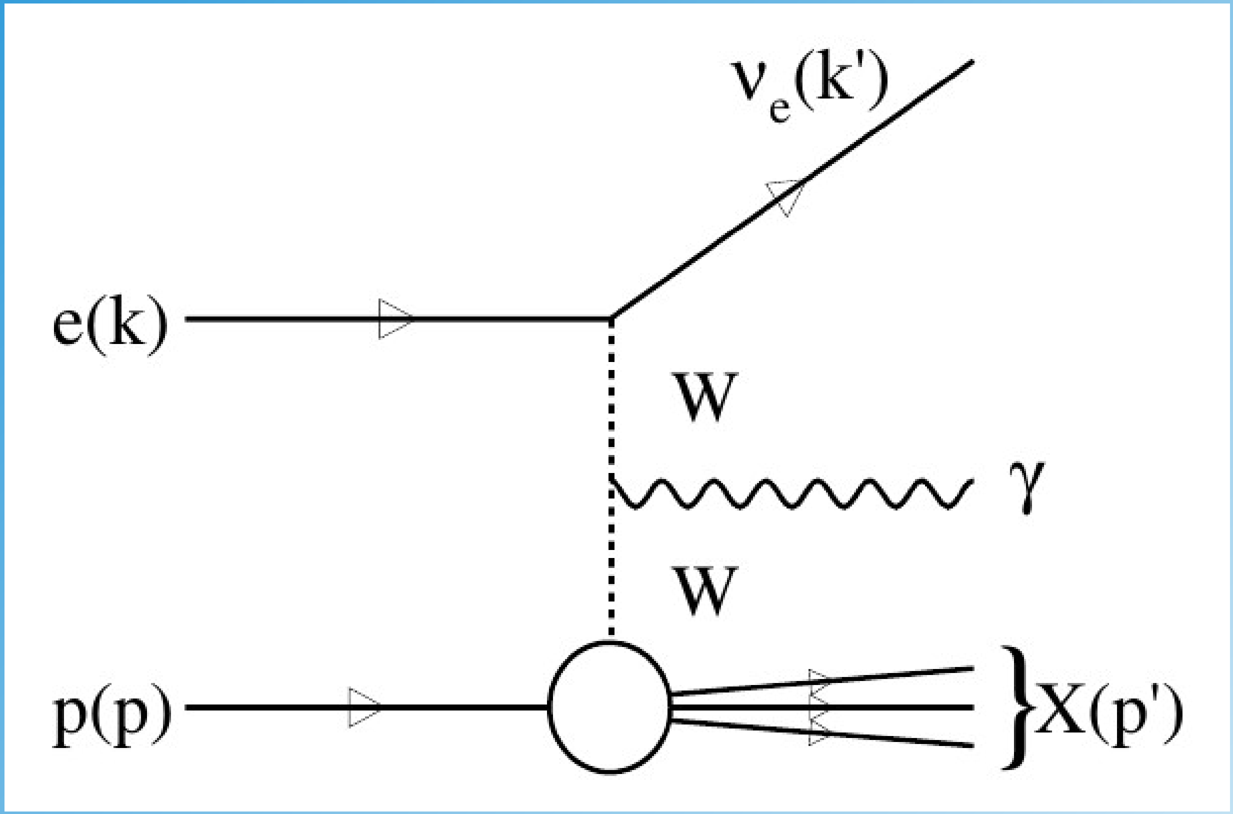
Figure 1.7: Feynman diagrams for scattering with radiation.

Neutrinos are note seen in detector and create missing p_T

Charged Current Event with Bremsstrahlung



Radiative Charged Current Events and W-W-γ Coupling



The W-boson is electrically charged and can also radiate photons!

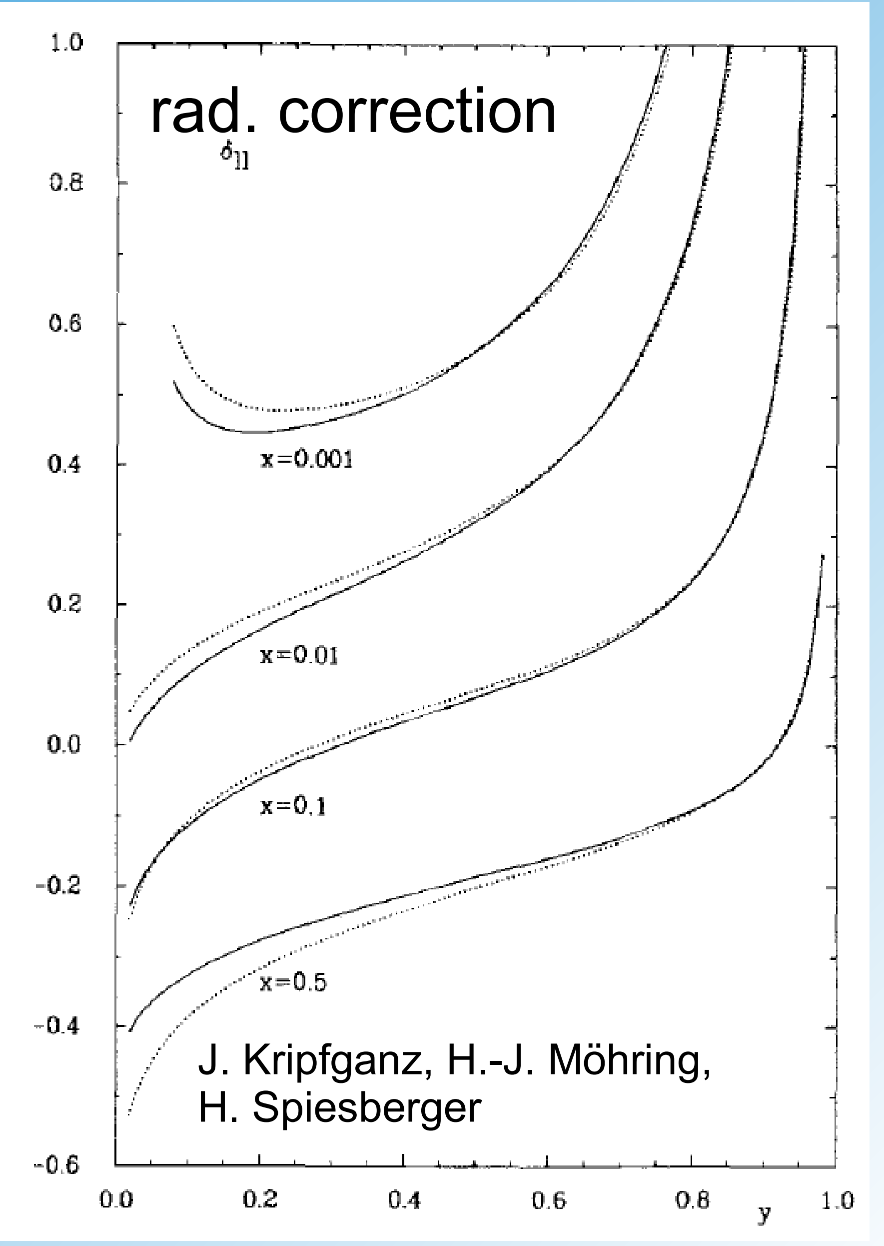
It is possible to test anomalous couplings

Radiative Corrections

Higher order leading logarithmic QED corrections to deep inelastic ep scattering at very high energies have significant impact on the kinematic reconstruction

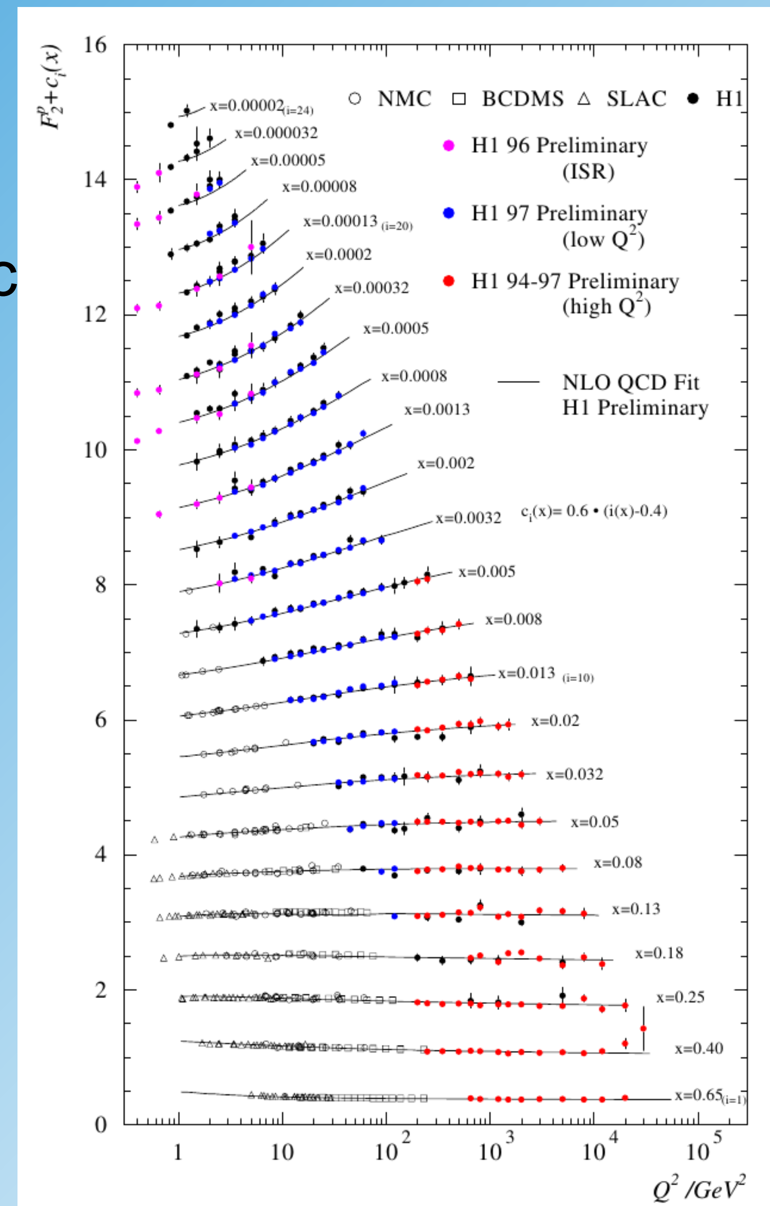
These logarithmic QED corrections come from multiple photon emissions

➜ systematic error of precision measurements



Exploitation of ISR

- ISR reduces the center of mass energy of the actual hard interaction
- This can be used to extend the kinematic phase space to lower energies



$$Q^2 = -q^2 = -(p - p')^2$$

allows to access smaller values of Q^2 which could not be reached otherwise!

Exploitation of ISR

- ISR
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Q

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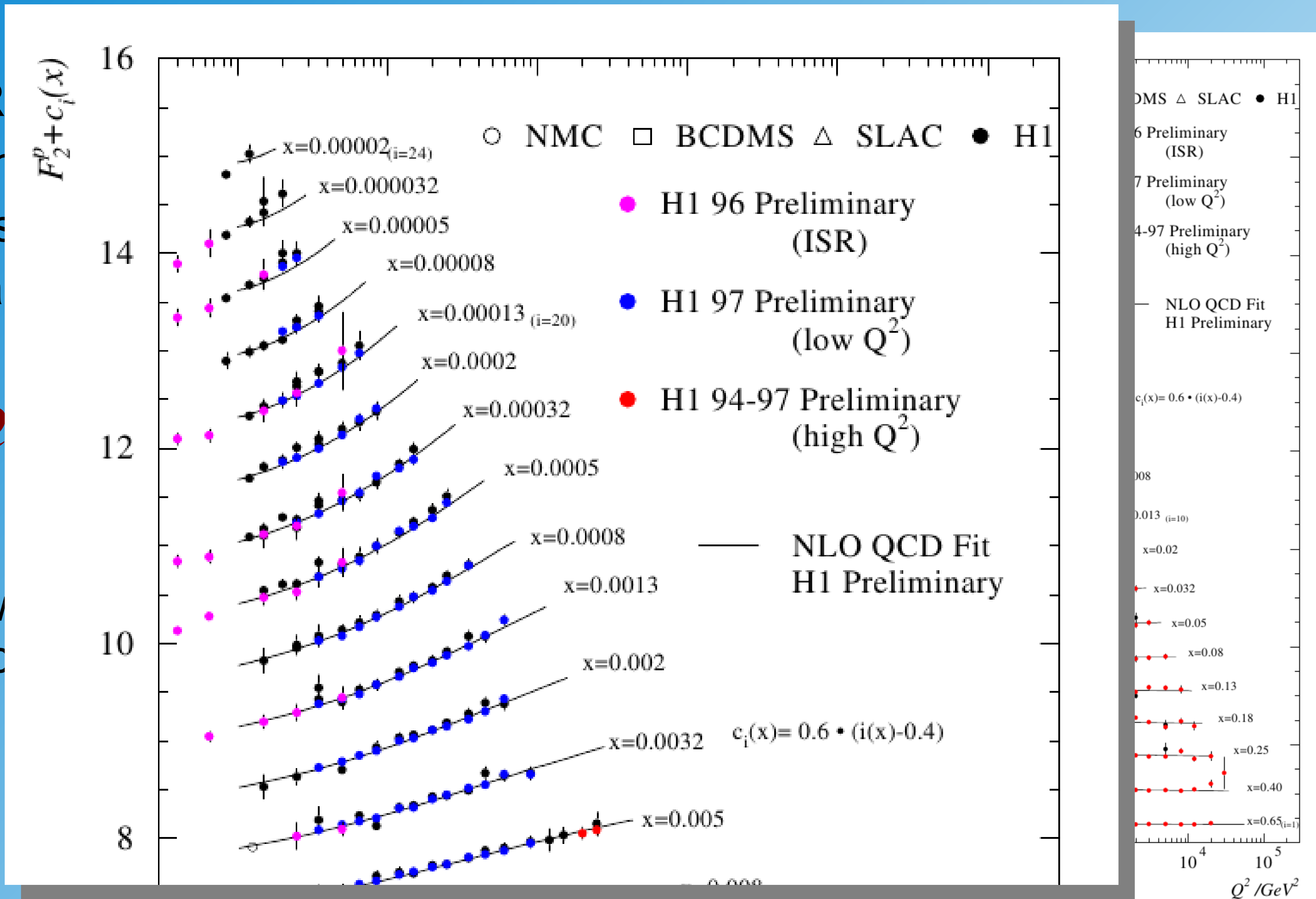
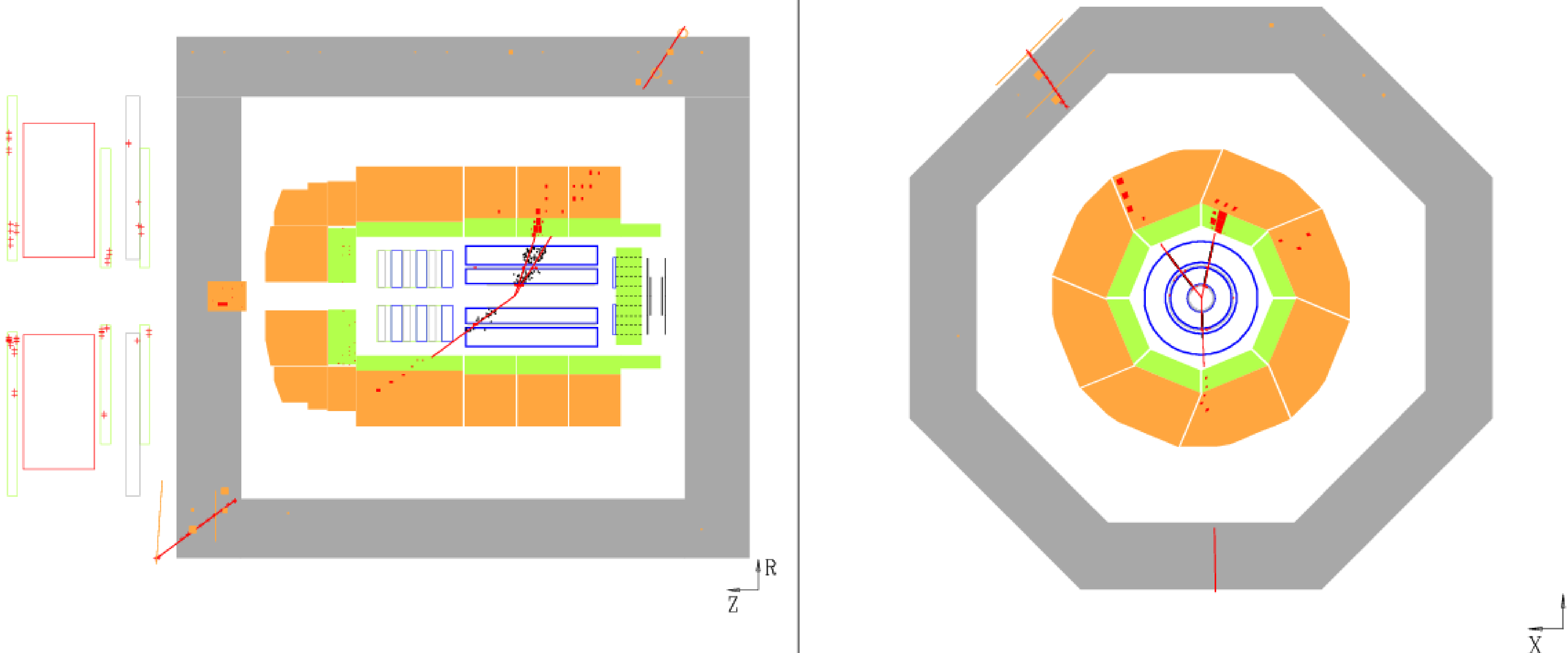


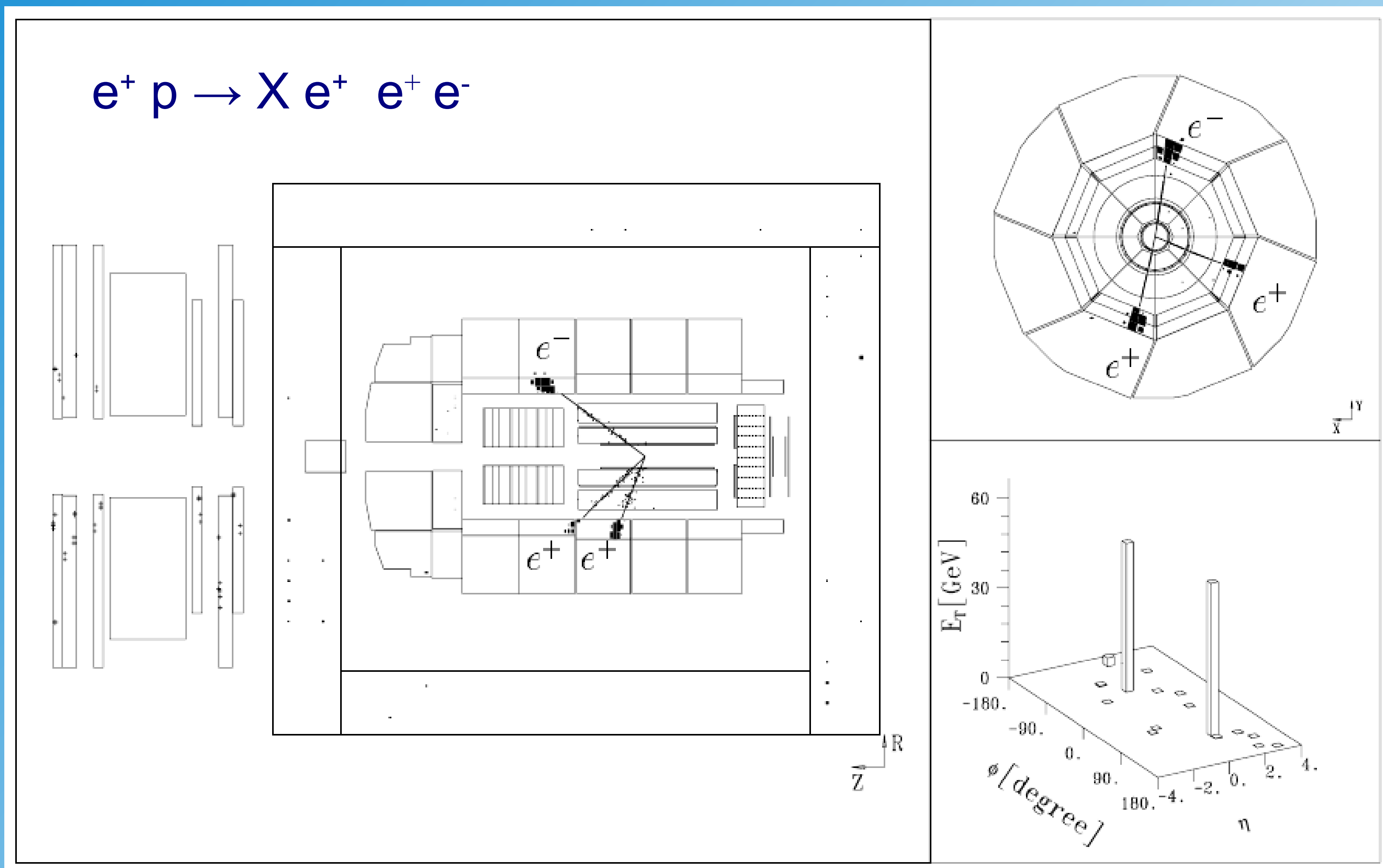
Higher Order Processes at HERA

Alpha[4] Processes at HERA



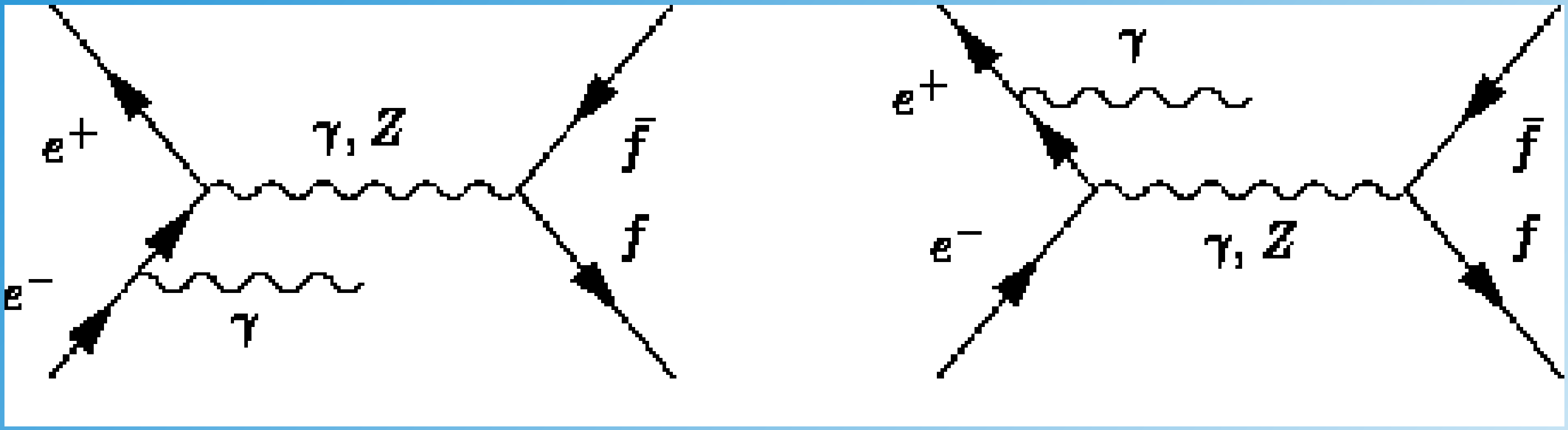
$e^{+} p \rightarrow e^{+} X \mu^{+} \mu^{-}$

Alpha^4 Processes at HERA



Radiative Returns

Z-boson are resonantly produced in $e^+ e^-$ collisions at resonance (m_Z ~91 GeV)

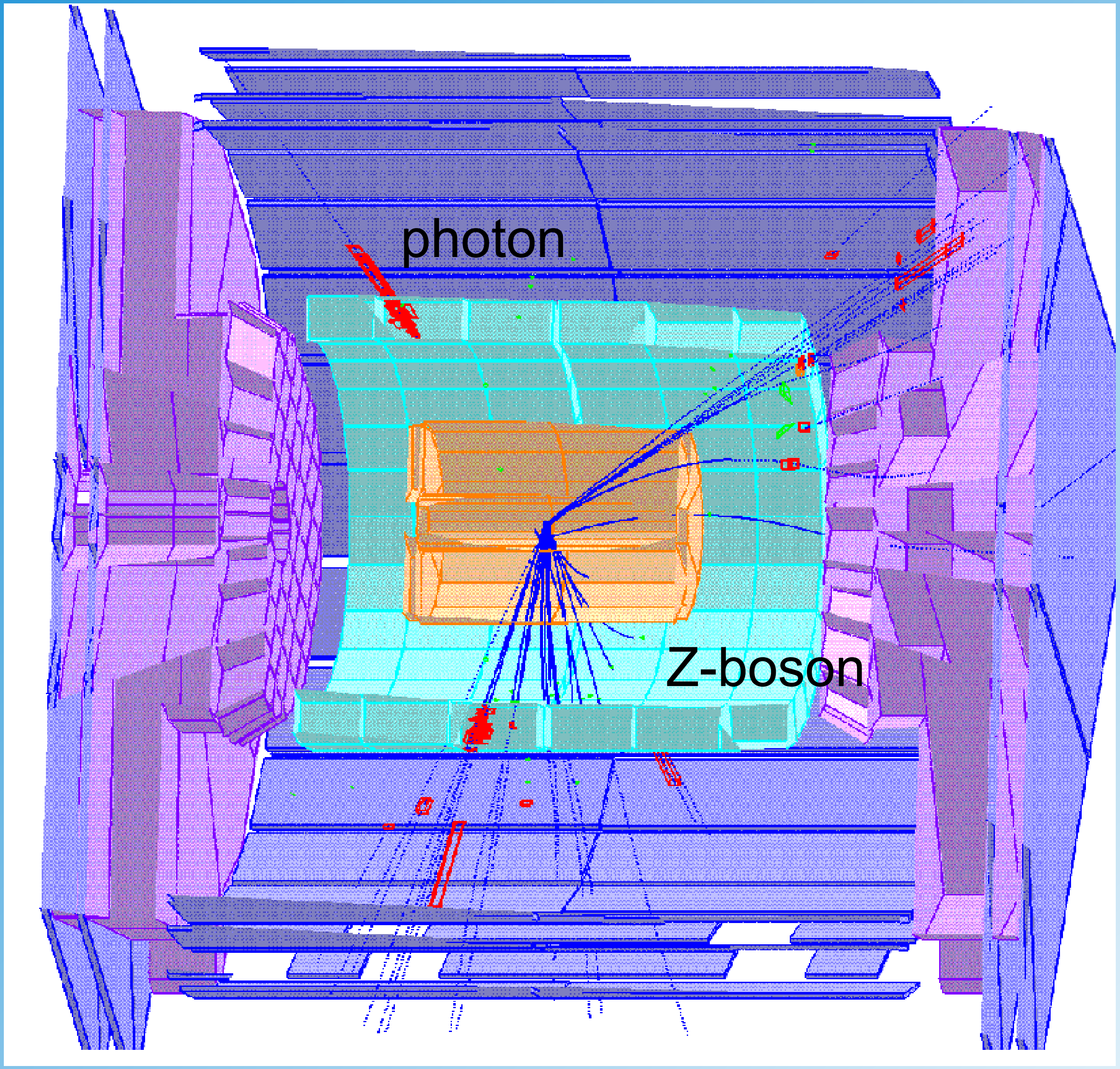


For $s^{1/2} > m_Z$ the radiative return allows a return to the resonance:

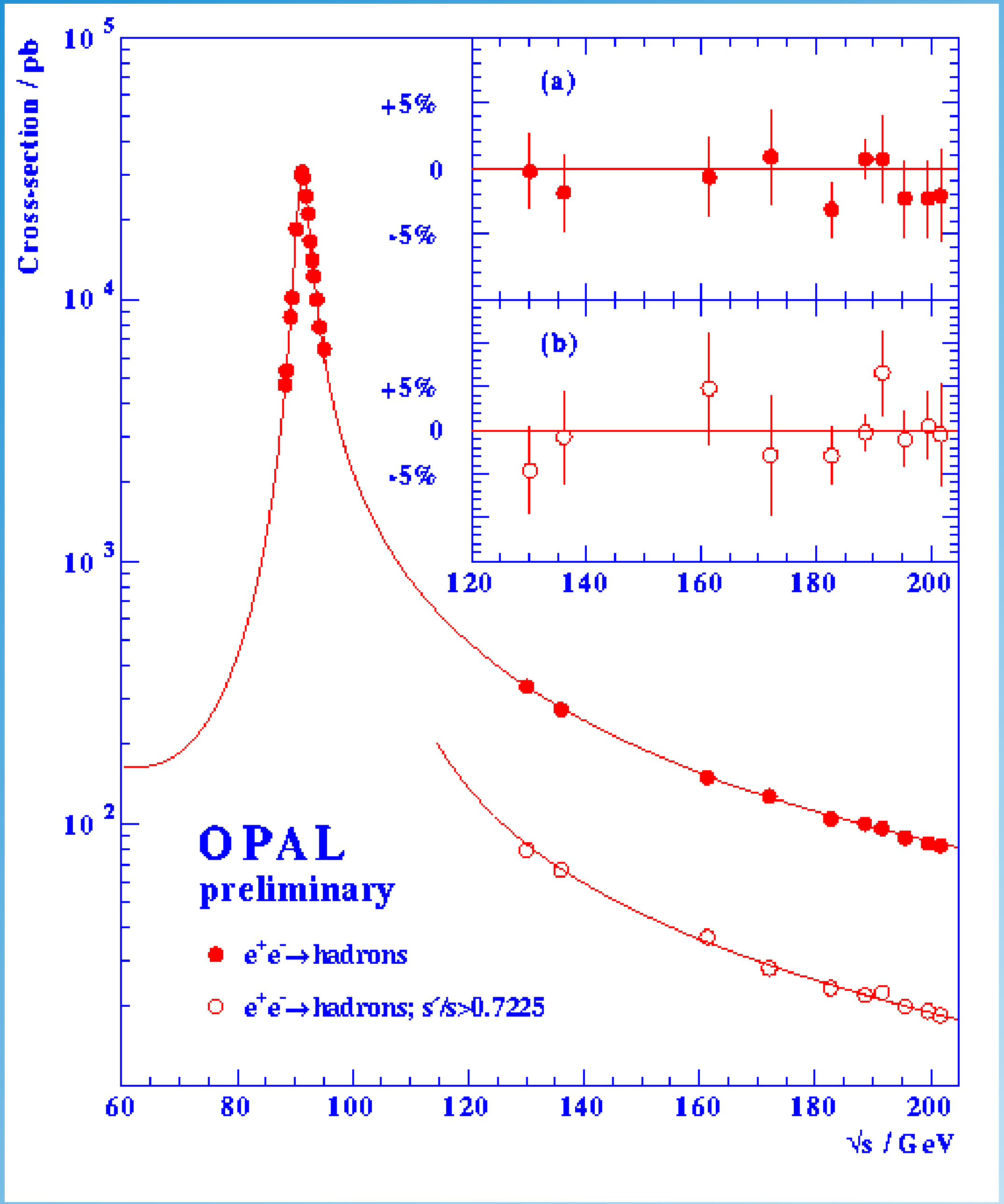
$$e^+ e^- \rightarrow Z \gamma$$

This effect can be large, e.g. at LEP

DELPHI Radiative Return Event at LEP



Measurement of Radiative Return

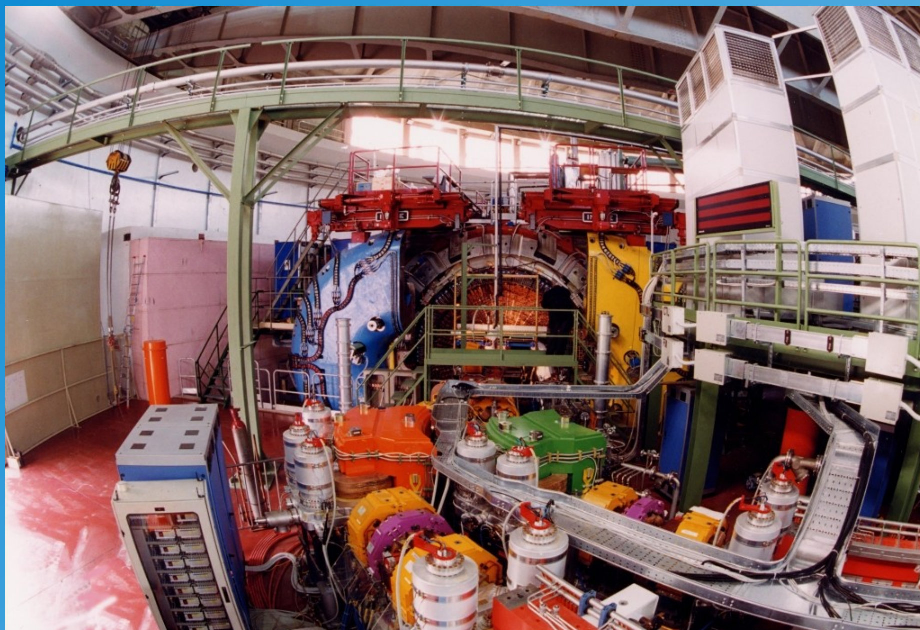


radiative returns
are extemely important!

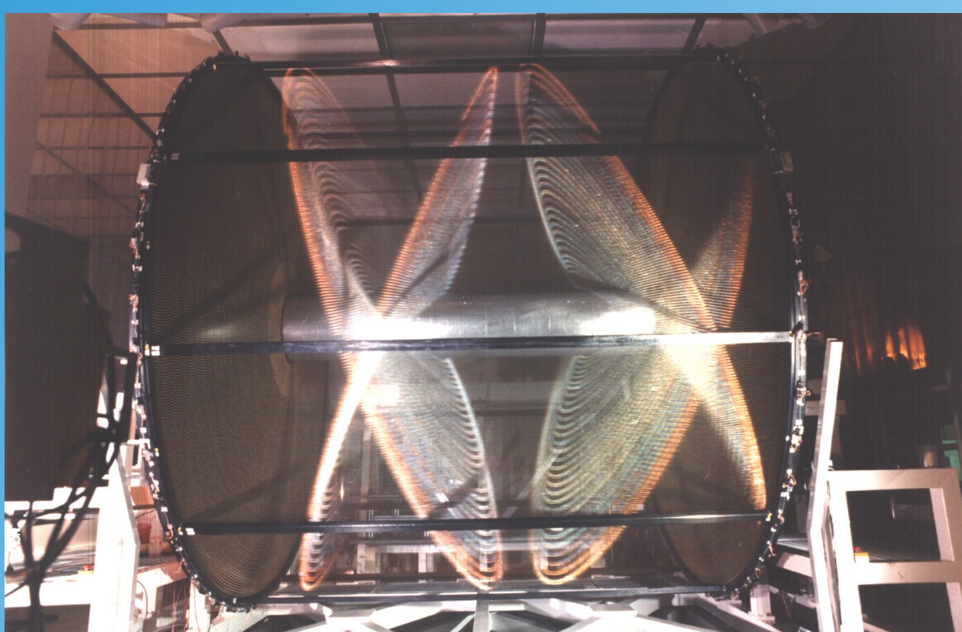
incl. radiative return

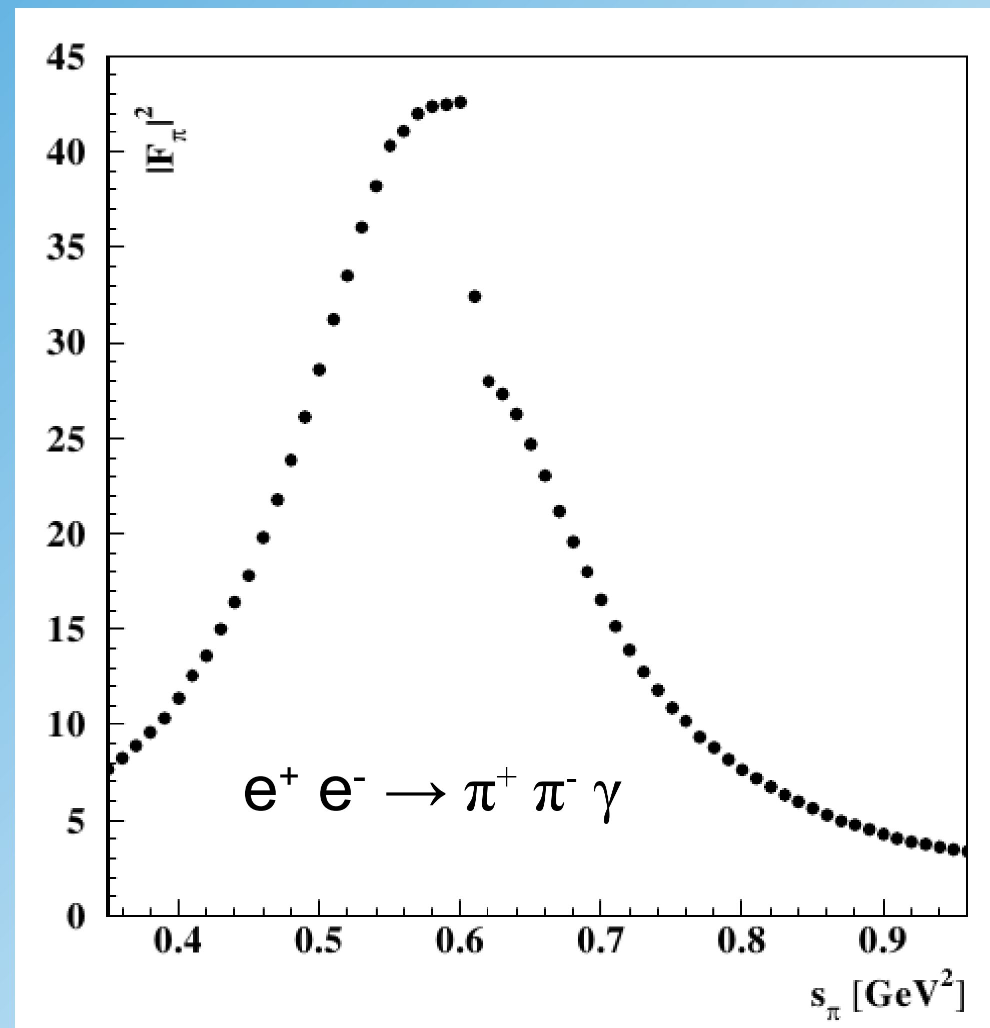
no radiative return

Radiative Return at KLEO (Daphne)

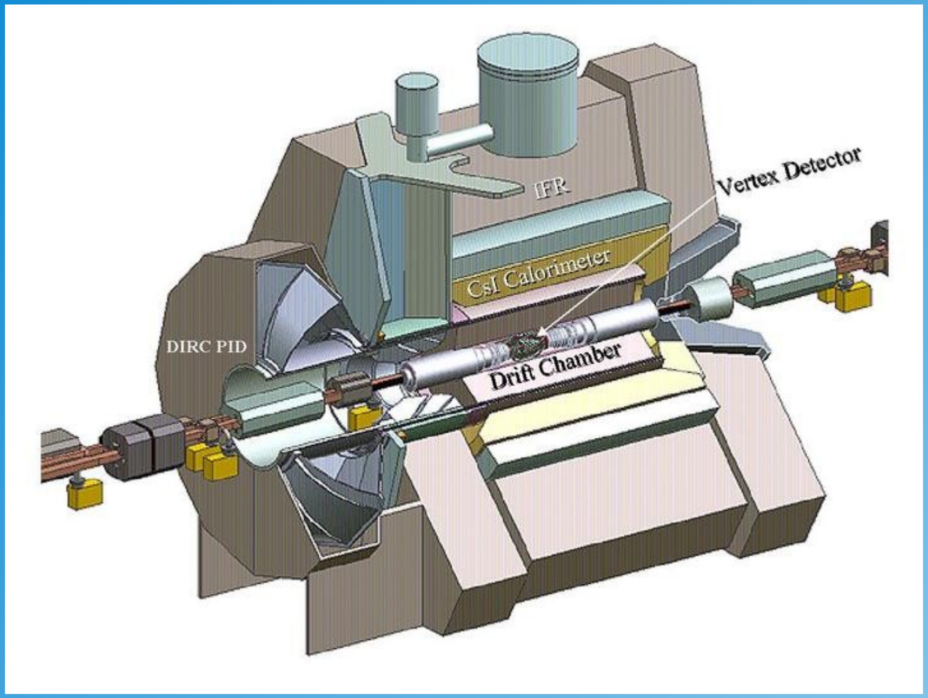


$e^+ e^-$ collider at $s^{1/2} = 1$ GeV





Radiative Returns at Barbar



$e^+ e^-$ collider at $s^{1/2}$ = 10 GeV

$e^+ e^- \rightarrow p \bar{p} \gamma$