

Lecture:

# Standard Model of Particle Physics

Heidelberg SS 2012

## Experimental Tests of QED Part 1

# Overview

## PART I

- Cross Sections and QED tests
- Accelerator Facilities + Experimental Results

## PART II

- Tests of QED in Particle Decays and Resonances
- QED Radiative Effects

# Measurement of Cross Sections

$$e^+e^- \rightarrow X$$

$$(e^-e^- \rightarrow X)$$

→ test predictions of QED

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Breakdown at higher energies?

# Measurement of Cross Sections

$$e^+e^- \rightarrow X$$
$$(e^-e^- \rightarrow X)$$

→ test predictions of QED

Why is  $\alpha_{em} = 1/137$  so small?

Breakdown at higher energies?

Reactions depend on center of mass → many different accelerators

Synchrotron Radiation Law::

$$P \propto \frac{E^4}{R^2}$$

→ large accelerators required for high energies

# List of ee-Accelerators

Accelerator	Location	Years of operation	Shape and circumference	Electron energy	Positron energy	Experiments	Notable Discoveries
AdA	Frascati, Italy; Orsay, France	1961–1964	Circular, 3 meters	250 MeV	250 MeV		Touschek effect (1963); first $e^-e^-$ interactions recorded (1964)
Princeton-Stanford ( $e^-e^-$ )	Stanford, California	1962–1967	Two-ring, 12 m	300 MeV	300 MeV		$e^-e^-$ interactions
VEP-1 ( $e^-e^-$ )	INP, Novosibirsk, Soviet Union	1964–1968	Two-ring, 2.70 m	130 MeV	130 MeV		$e^-e^-$ scattering; QED radiative effects confirmed
VEPP-2	INP, Novosibirsk, Soviet Union	1965–1974	Circular, 11.5 m	700 MeV	700 MeV	OLYA, CMD	multihadron production (1966), $e^+e^- \rightarrow \varphi$ (1966), $e^+e^- \rightarrow \gamma\gamma$ (1971)
SPEAR	SLAC	1972-1990(?)				Mark I, Mark II, Mark III	Discovery of Charmonium states
VEPP-2M	BINP, Novosibirsk	1974–2000	Circular, 17.88 m	700 MeV	700 MeV	ND, SND, CMD-2	$e^+e^-$ cross sections, radiative decays of $\rho$ , $\omega$ , and $\varphi$ mesons
DORIS	DESY	1974–1993	Circular, 300m	5 GeV	5 GeV	ARGUS, Crystal Ball, DASP, PLUTO	Oscillation in neutral B mesons
PETRA	DESY	1978–1986	Circular, 2 km	20 GeV	20 GeV	JADE, MARK-J, PLUTO, TASSO	Discovery of the gluon in three jet events
CESR	Cornell University	1979–2002	Circular, 768m	6 GeV	6 GeV	CUSB, CHES, CLEO, CLEO-2, CLEO-2.5, CLEO-3	First observation of B decay, charmless and "radiative penguin" B decays
PEP	SLAC	1980-1990(?)				Mark II	
SLC	SLAC	1988-1998(?)	Addition to SLAC Linac	45 GeV	45 GeV	SLD, Mark II	First linear collider
LEP	CERN	1989–2000	Circular, 27 km	104 GeV	104 GeV	Aleph, Delphi, Opal, L3	Only 3 light ( $m \leq m_Z/2$ ) weakly interacting neutrinos exist, implying only three generations of quarks and leptons
BEPC	China	1989–2004	Circular, 240m	2.2 GeV	2.2 GeV	Beijing Spectrometer (I and II)	
VEPP-4M	BINP, Novosibirsk	1994-	Circular, 366m	6.0 GeV	6.0 GeV	KEDR	Precise measurement of Y-meson masses
PEP-II	SLAC	1998–2008	Circular, 2.2 km	9 GeV	3.1 GeV	BaBar	Discovery of CP violation in B meson system
KEKB	KEK	1999–2009	Circular, 3 km	8.0 GeV	3.5 GeV	Belle	Discovery of CP violation in B meson system
DAΦNE	Frascati, Italy	1999-	Circular, 98m	0.7 GeV	0.7 GeV	KLOE	Crab-waist collisions (2007)
CESR-c	Cornell University	2002–2008	Circular, 768m	6 GeV	6 GeV	CHES, CLEO-c	
VEPP-2000	BINP, Novosibirsk	2006-	Circular, 24.4m	1.0 GeV	1.0 GeV	SND, CMD-3	Round beams (2007)
BEPC II	China	2008-	Circular, 240m	3.7 GeV	3.7 GeV	Beijing Spectrometer III	

# AdA Accelerator

- First  $e^+ e^-$  collider ever
- AdA = Anello di Accumulazione (Frascati/Orsay, 1961-64)
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## Motivation:

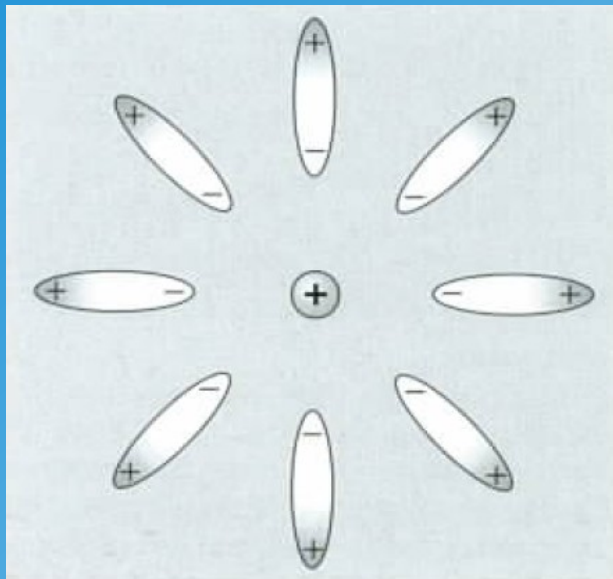
- **Bruno Touschek**: excite the dielectric vacuum to create vector mesons (e.g. rho meson predicted to be light!)

**Note: at that time all new particles had been discovered in hadronic interactions (ie. proton beams)!**



# Dielectric Vacuum

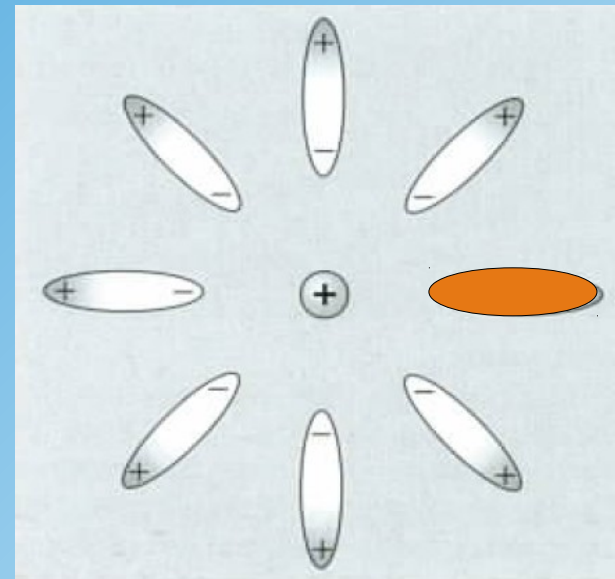
classical dielectric



bare electrical charge shielded  
by induced dipoles

high energy  
→

“excited dielectric”



bare charge shielded  
by vacuum polarisation

rho-meson

# AdA Accelerator

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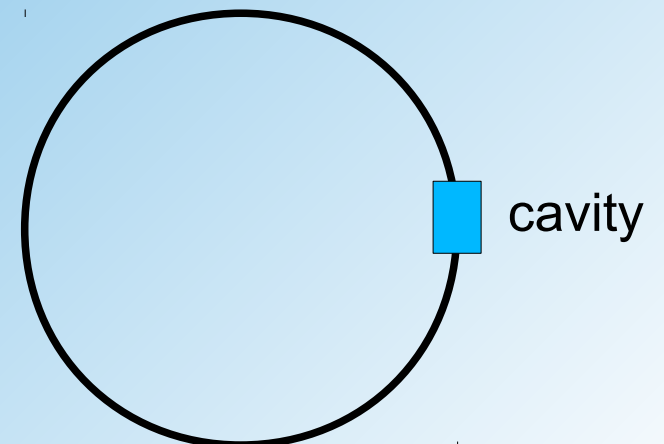
**Note: at that time all new particles had been discovered in hadronic interactions (ie. proton beams)!**

**“Revolutionary” concept as the rho-meson is electrically neutral and was predicted to explain (as carrier) strong interactions**

Remark: Indeed, Touschek was right. The strong force can be tested in  $e^+ e^-$  collisions. But not in AdA (too low luminosity, too low energy)

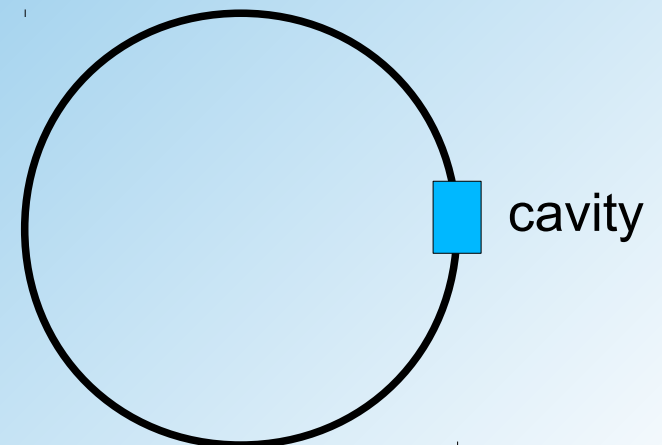
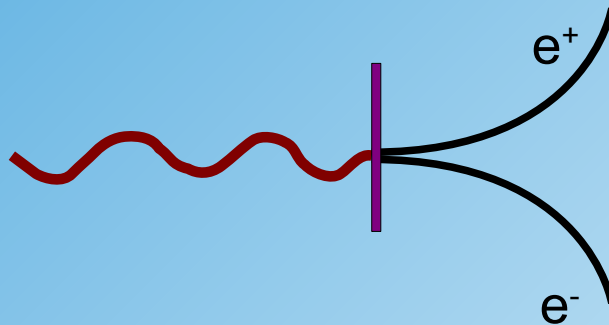
# AdA Challenges I

- How to store electrons and positrons?
  - magneto-optical storage ring (→ known at this time, synchrotron radiation facilities)



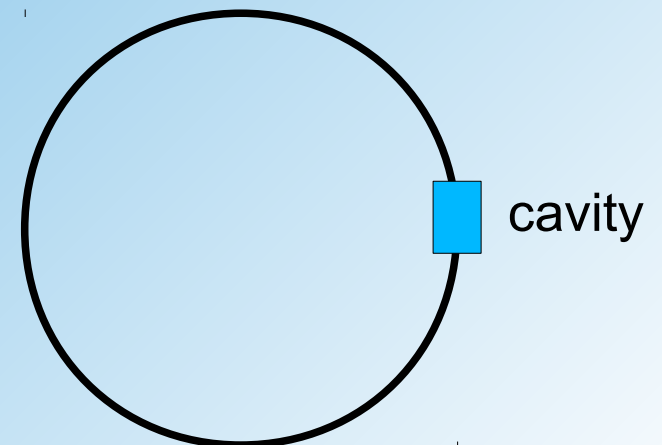
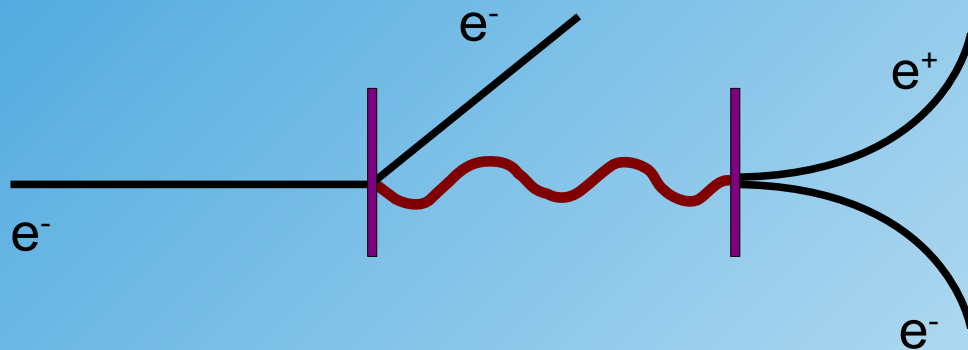
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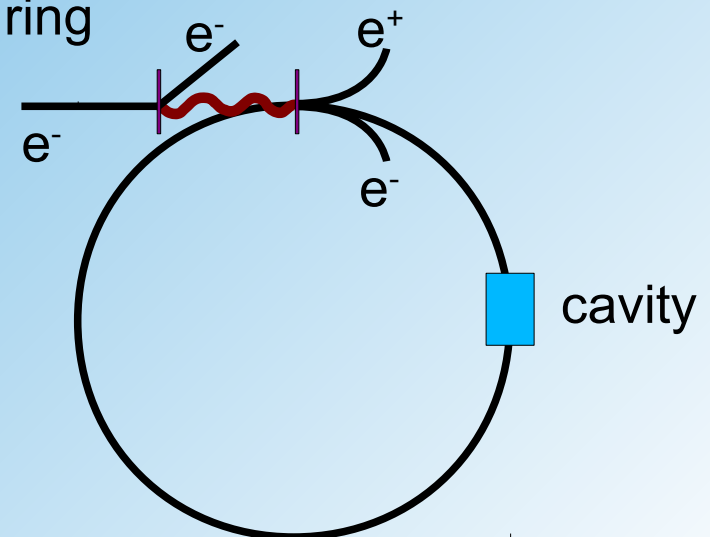
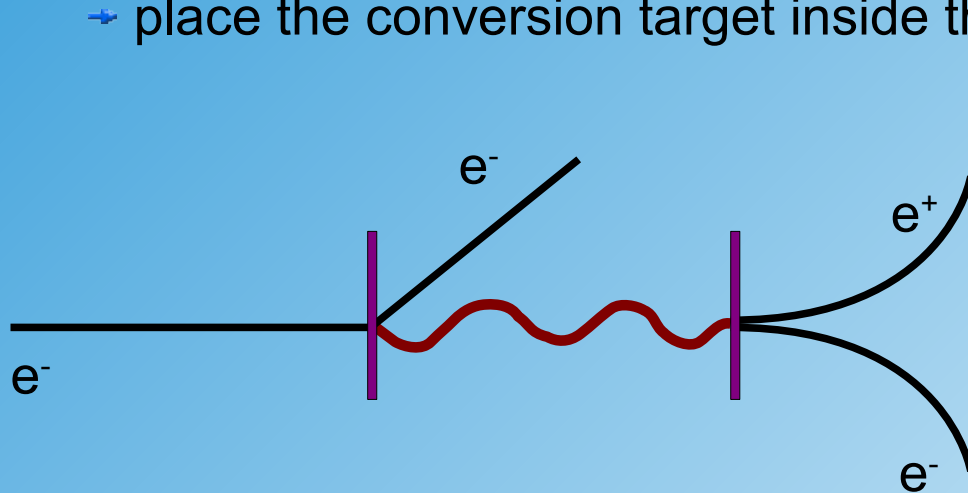
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  - Bremsstrahlung from high energetic electrons at target  
 $e^- N \rightarrow \gamma e^- N$  using a linear electron accelerator (→ also known)



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 $e^- N \rightarrow \gamma e^- N$  using a linear electron accelerator (→ also known)
- How to fill the storage ring with electrons and positrons???

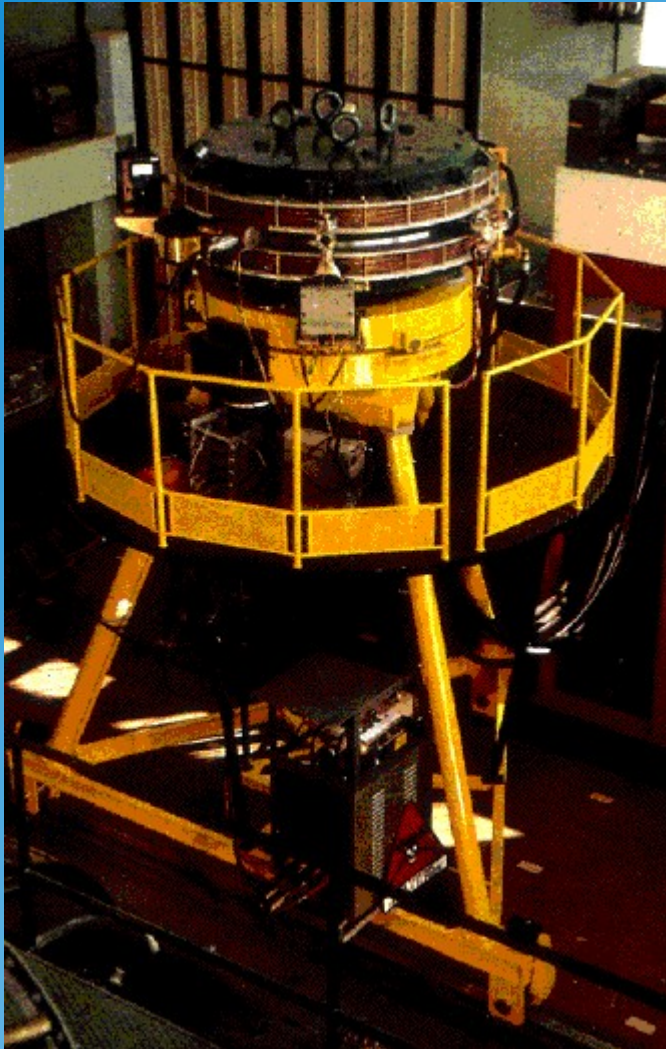




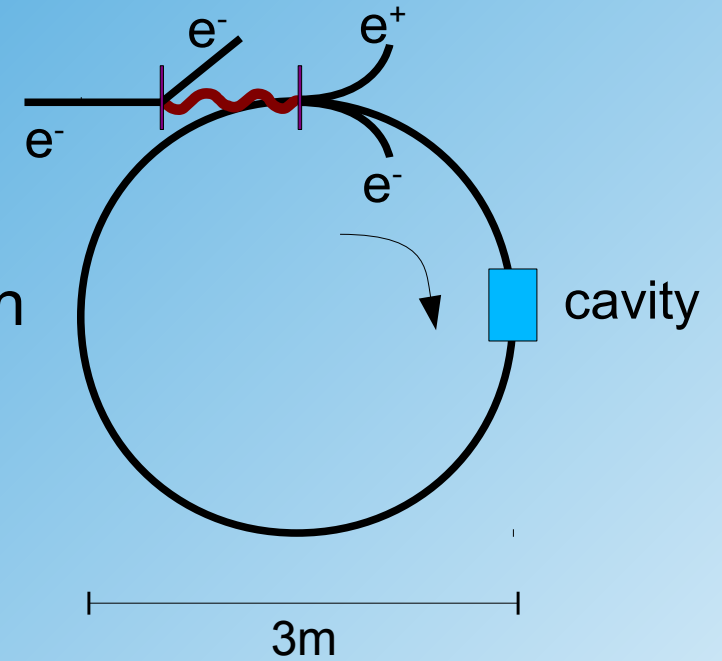
MAGNETIC DISCUSSION

*brunswick*

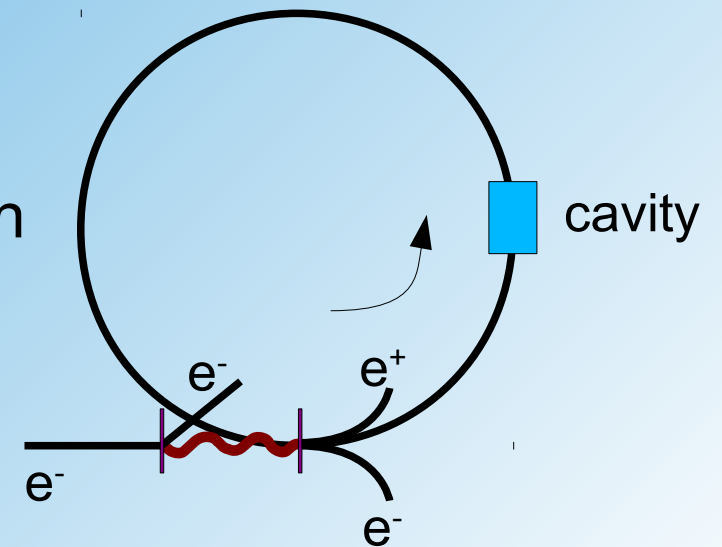
# AdA Concept



Electron Injection



Positron Injection





# AdA Challenges II

- How to make electrons and positrons collide?

Note: AdA is a single storage ring: electrons and positrons see same optics but in reverse direction

*B.Touschek: It is guaranteed that an electron and a positron necessarily meet in a single orbit because QED is CP (charge-parity)*



*Bruce D. Joseph*

# AdA Challenges II

- How to make electrons and positrons collide?

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B.Touschek: *It is guaranteed that an electron and a positron necessarily meet in a single orbit because QED is CP (charge-parity)*

If a ring collider works, then CP(T) invariance of QED is confirmed!!!

Note: CP(T) invariance says that a positron can be regarded as an electron traveling in reverse time direction.

**Touschek was right, in a very short time AdA was commissioned and electron-positron collisions were observed – much more than just a technical achievement!**

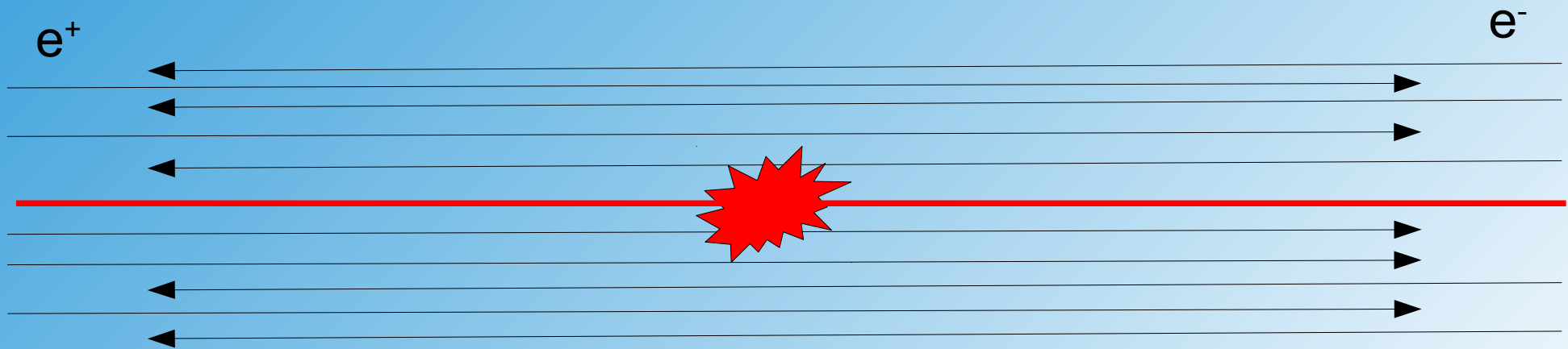
# AdA Challenges III

- How to measure that electron-positron collisions take place?
- How many collisions?

## Definition of “Luminosity” Measurement (source factor)

$$R = L \sigma$$

Relation between rate of events and cross section of process



# Luminosity Measurement in Ring

Accelerator Formula

$$L = \frac{N_1 N_2 f}{4 \pi A}$$

$N_1$  and  $N_2$  and beam cross section  $A$  are unknown and have to be precisely measured → difficult

More simple ansatz – use reference process(es):

$e^+ e^- \rightarrow e^+ e^-$

$$\frac{d\sigma}{d(\cos\theta)} = \frac{\pi\alpha^2}{s} \left( u^2 \left( \frac{1}{s} + \frac{1}{t} \right)^2 + \left( \frac{t}{s} \right)^2 + \left( \frac{s}{t} \right)^2 \right)$$

ultrarelativistic approx. (Bhabha 1936)

$e^+ e^- \rightarrow \gamma \gamma$

$$\frac{d\sigma}{d\Omega}(e^+ e^- \rightarrow \gamma \gamma) = \frac{\alpha^2}{2s} \frac{u^2 + t^2}{tu}$$

annihilation process (Compton-like)

**Both processes are forward peaked!**

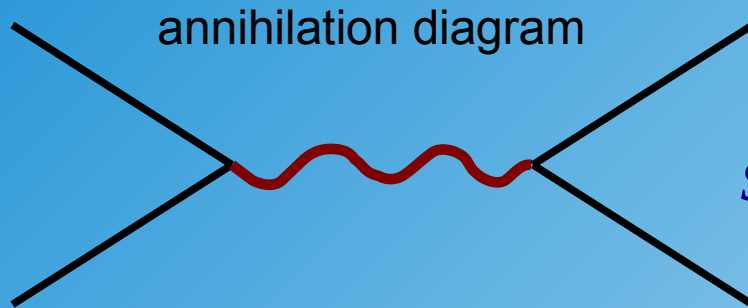
t-pole

$$t = -s \sin^2(\theta/2)$$

# Bhabha Scattering

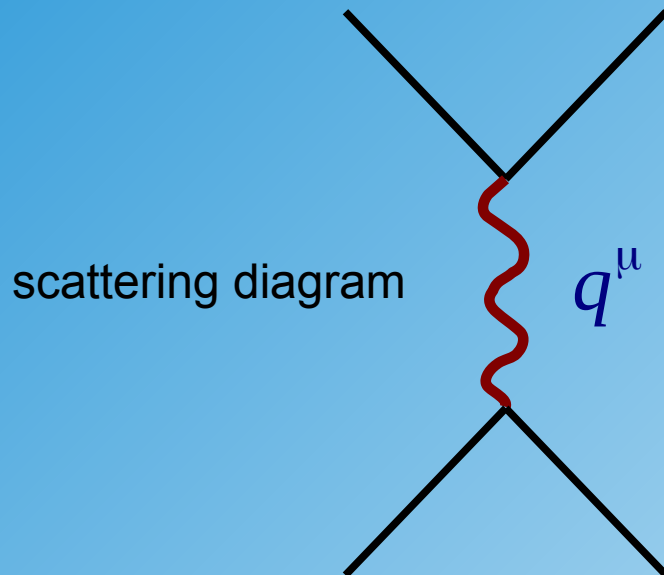
$$\frac{d\sigma}{d\Omega}(e^+ e^- \rightarrow e^+ e^-) = \frac{\alpha^2}{2s} \left( \frac{u^2 + t^2}{s^2} + \frac{s^2 + u^2}{t^2} + \frac{2u^2}{st} \right)$$

annihilation   scattering   interference



$$s = q^\mu q_\mu$$

s-pole from photon propagator

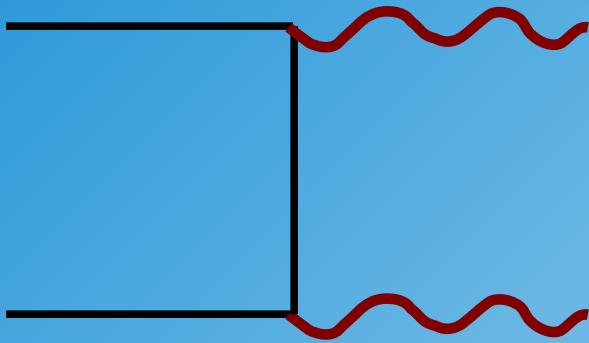


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t-pole from photon propagator

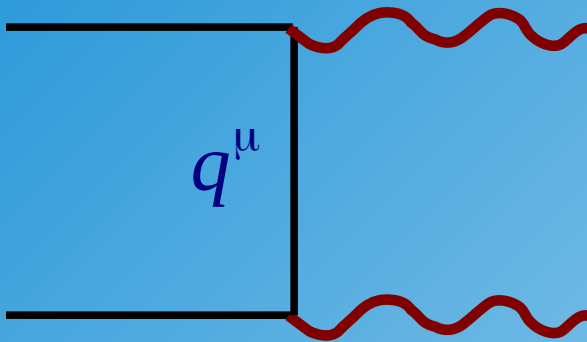
# Photon Pair Production

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$$\frac{d\sigma}{d\Omega}(e^+ e^- \rightarrow \gamma \gamma) = \frac{\alpha^2}{2s} \left( \frac{u}{t} + \frac{t}{u} \right)$$



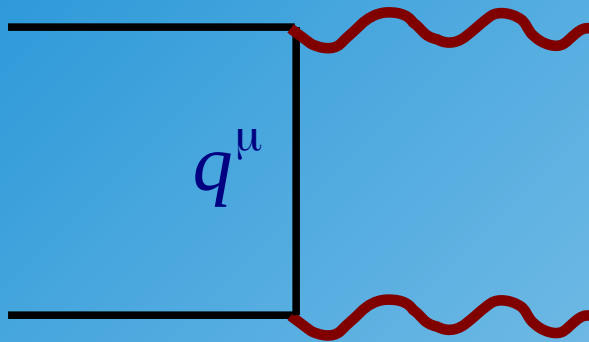
$$t = q^\mu q_\mu$$

t-pole from electron propagator



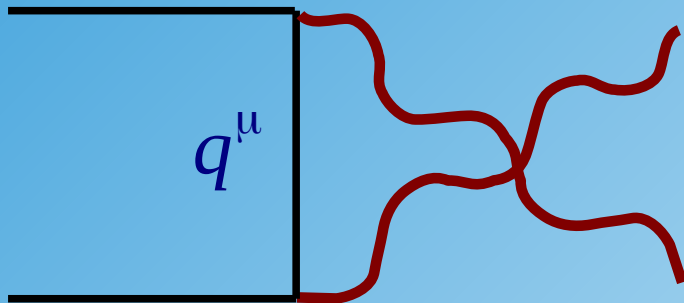
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$$t = q^\mu q_\mu$$

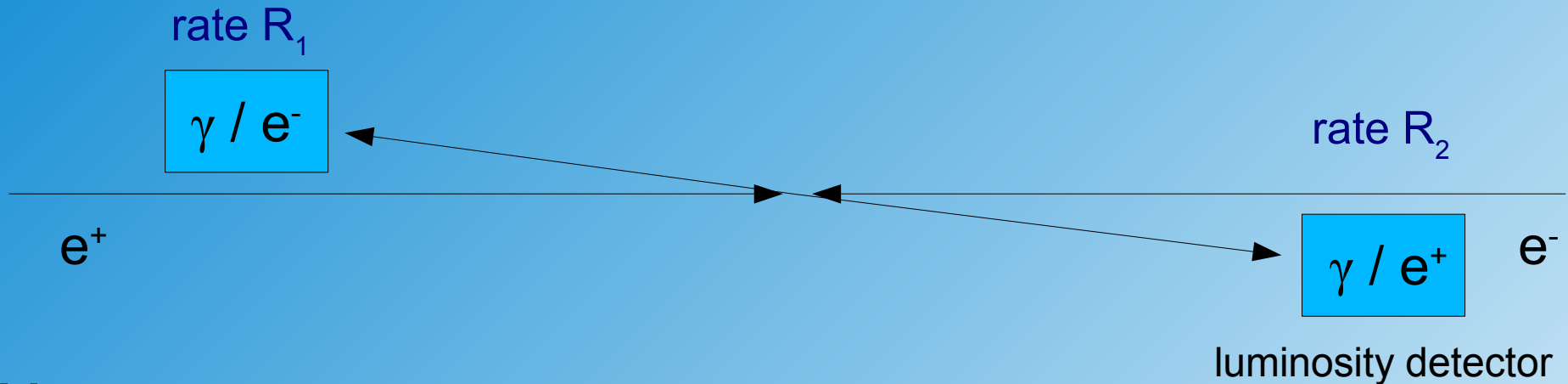
t-pole from electron propagator



$$u = q^\mu q_\mu$$

u-pole from electron propagator

# Sketch of Luminosity Measurement



Measurement: Rates  $R_1$  and  $R_2$  (counts/s)

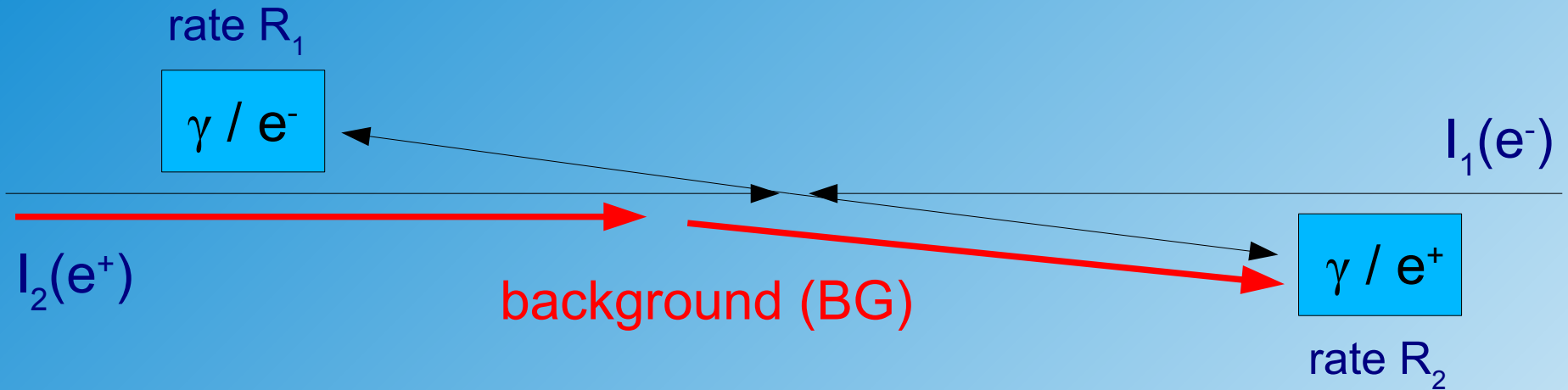
Use:  $R = L \sigma_{\text{Detector}} \leftrightarrow L = R / \sigma_{\text{Detector}}$

$$\sigma_{\text{Detector}} = \int_{\text{Detector}} \frac{d\sigma}{d\Omega} d\Omega$$

**acceptance calculation is an experimental task!**

$\sigma_{\text{Detector}}$  is the **observable** cross section  $\neq$  total cross section

# Background for Luminosity Measurement



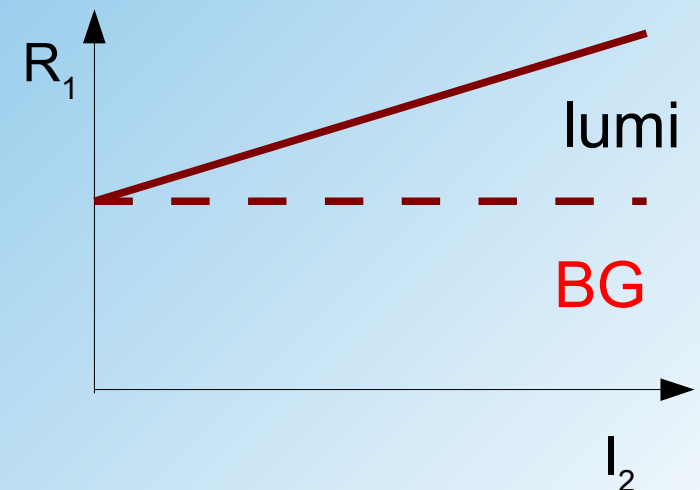
Problem: beam induced background, e.g. electron-rest gas scattering)

Ansatz:

$$R_1 = a_1 I_1 + b I_1 I_2 = I_1 (a_1 + b I_2)$$

$$R_2 = a_2 I_2 + b I_1 I_2 = I_2 (a_2 + b I_1)$$

**BG** lumi

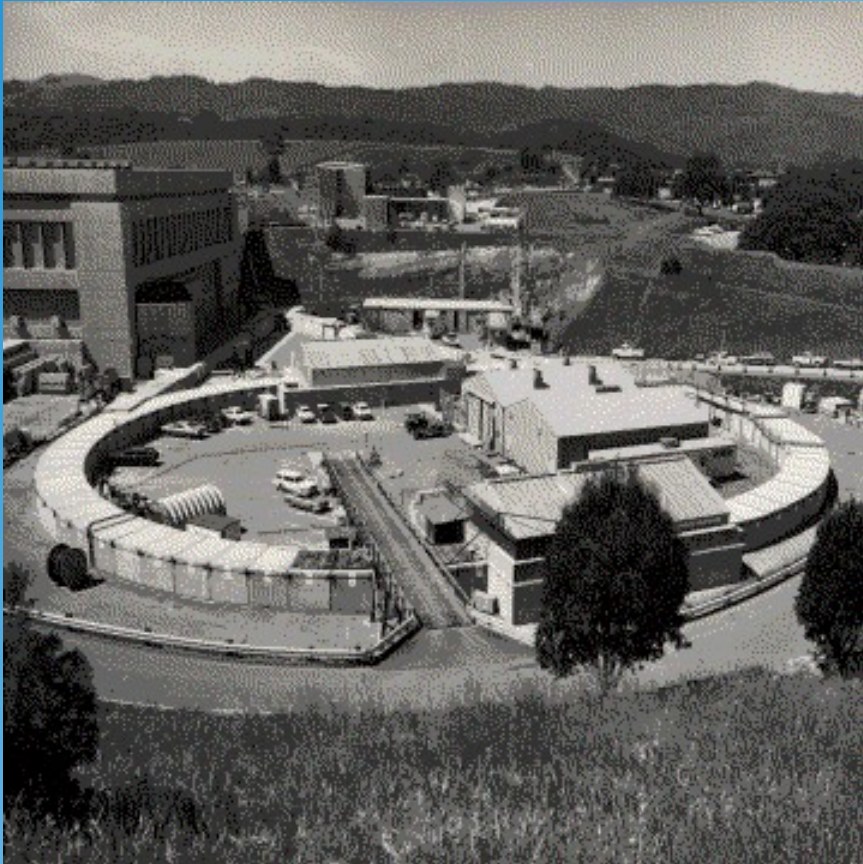


# The Big $e^+e^-$ Accelerators

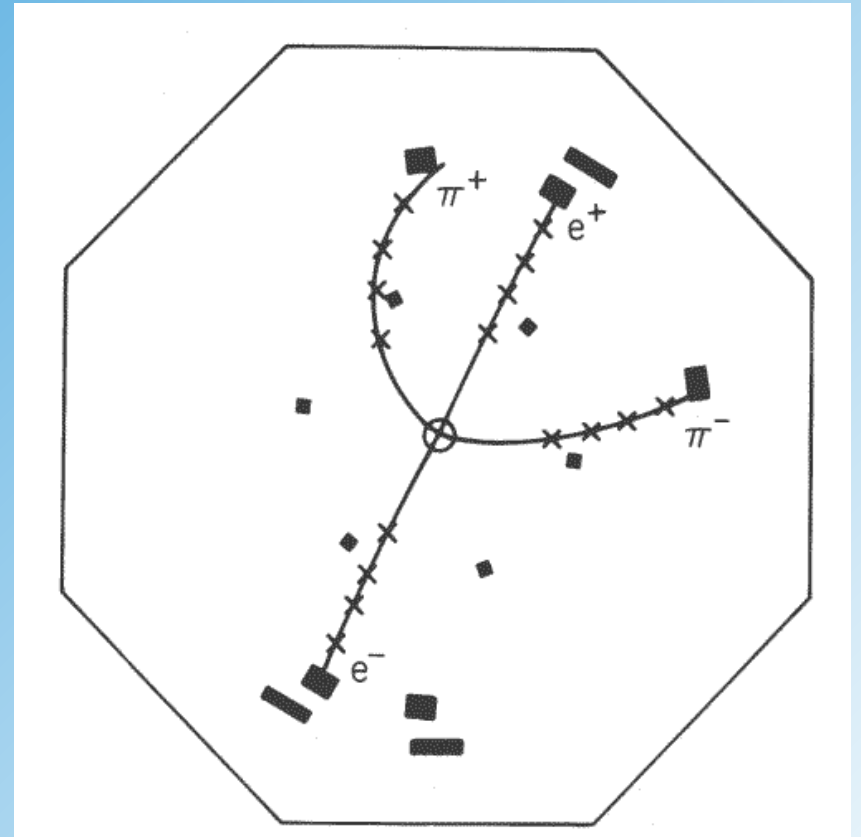
- **SPEAR** (Stanford Positron Electron Accelerator Ring) at SLAC (1974-1990),  $s^{1/2}=3-8$  GeV, Discovery of the Charm Quark
- **PETRA** (Positron Electron Tandem Ringanlage) at DESY (1978-1986),  $s^{1/2}=38$  GeV, Discovery of Gluon-Jets
- **TRISTAN** at KEK, Japan (1986-1989)  $s^{1/2}=50-64$  GeV (discovery of the “desert”)
- Large Electron-Positron Collider, Geneva (1988-2000):  
 $s^{1/2}=90$  GeV (**LEP I**, Z-factory),  $s^{1/2}=200$  GeV (**LEP II**, WW factory)
- Stanford Linear Accelerator at SLAC, Stanford (1991-1998)  
 $s^{1/2}=90$  GeV (**SLC**, Z-factory)

# SPEAR at SLAC

- Stanford Positron Electron Accelerator Ring (1974-1990),  $s^{1/2}=3-7$  GeV, Discovery the J/Psi

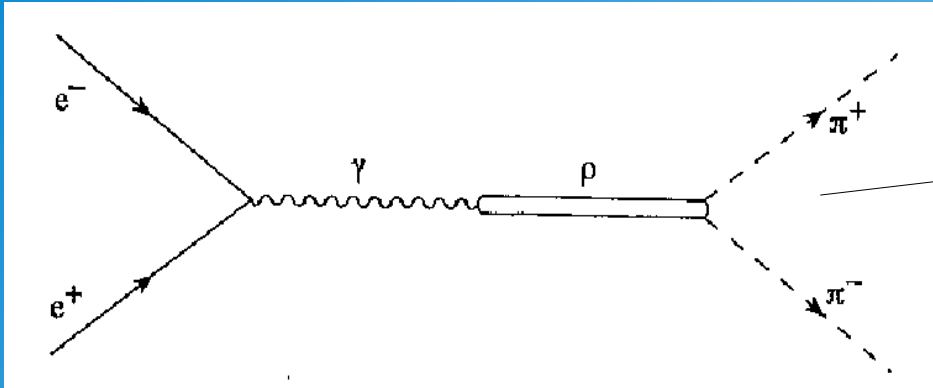


Discovery of the Charm Quark



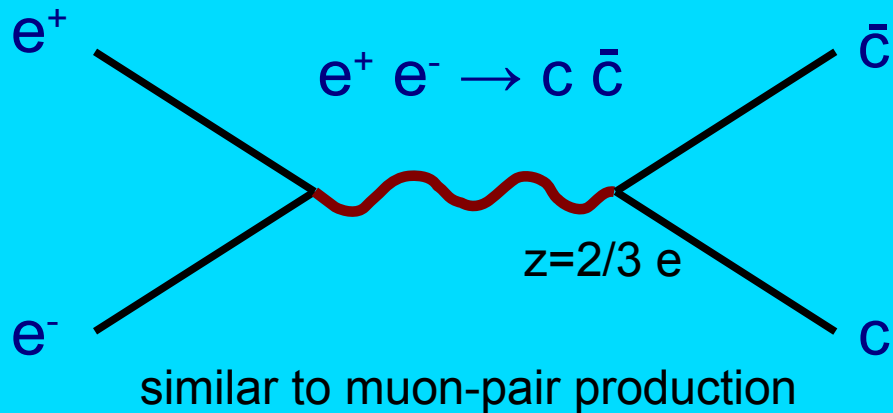
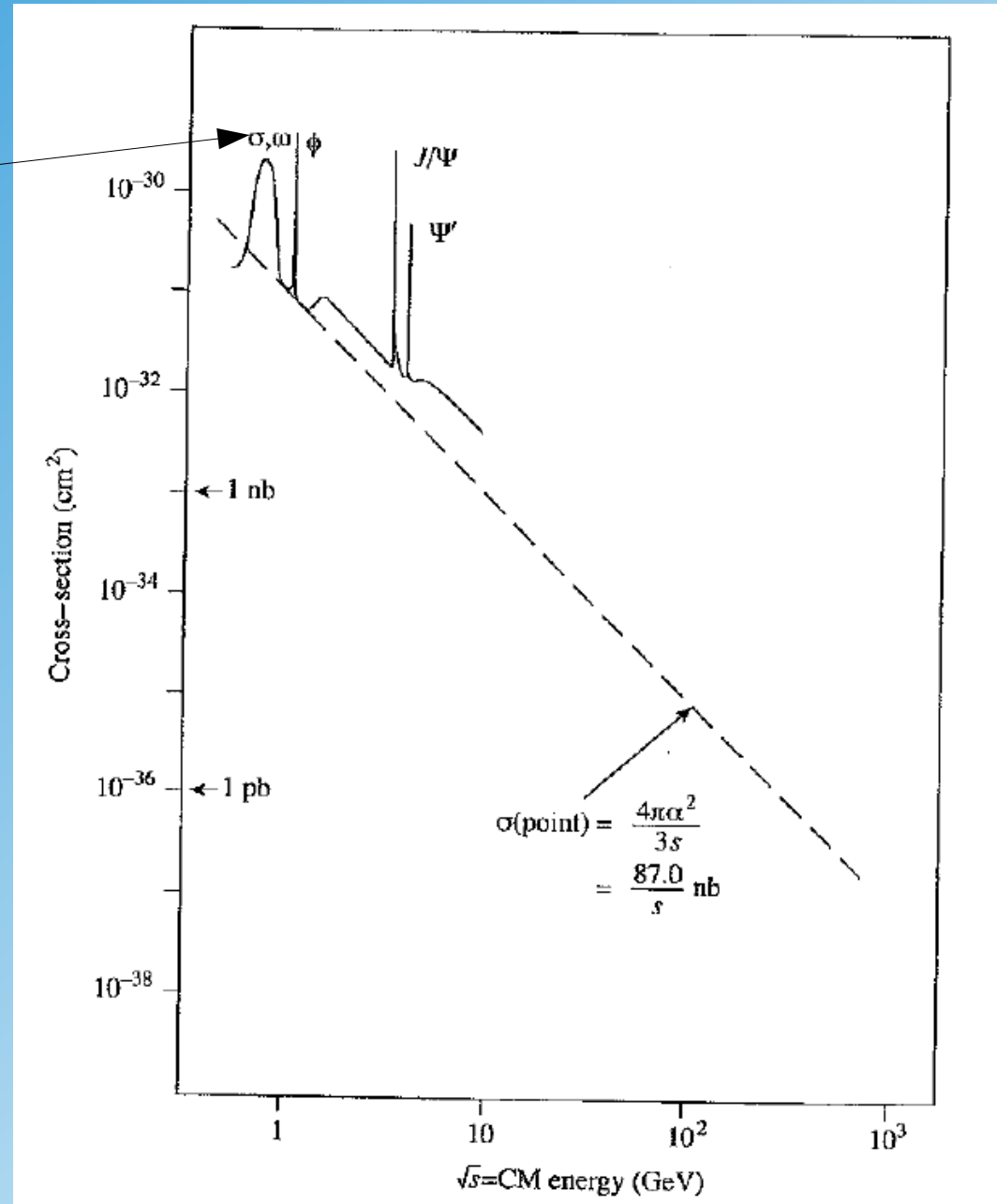
$$\Psi(2S) \rightarrow J/\Psi \pi^+ \pi^- \rightarrow e^+ e^- \pi^+ \pi^-$$

# Quark-Pair Production



Resonant Rho production

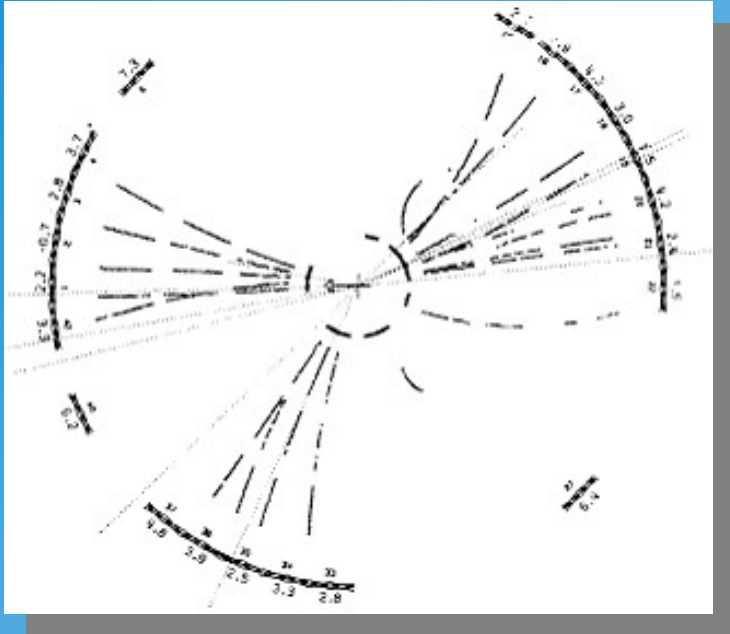
→ later



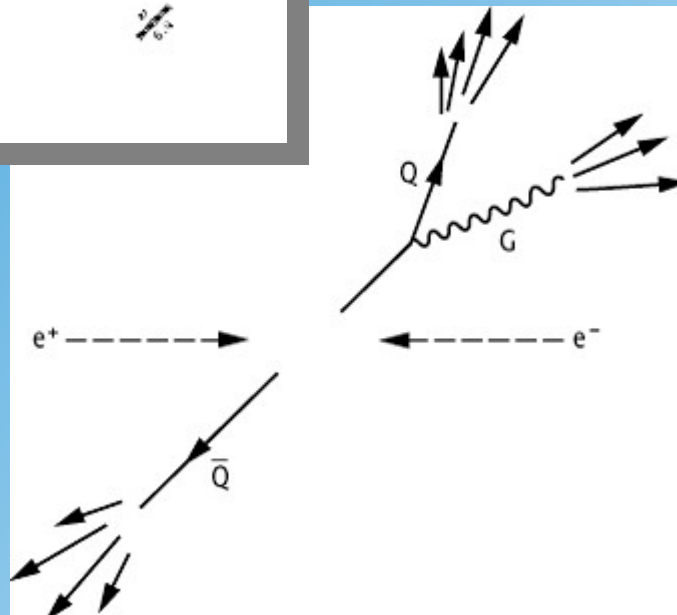
$$\frac{d\sigma}{d\Omega}(e^+ e^- \rightarrow c \bar{c}) = \frac{\alpha^2 q^2}{2s} \left( \frac{u^2 + t^2}{s^2} \right)$$

# PETRA at DESY

- Positron Electron Tandem Ring Anlage (1978-1986),  $\sqrt{s}=38$  GeV, Discovery of Gluon Jets



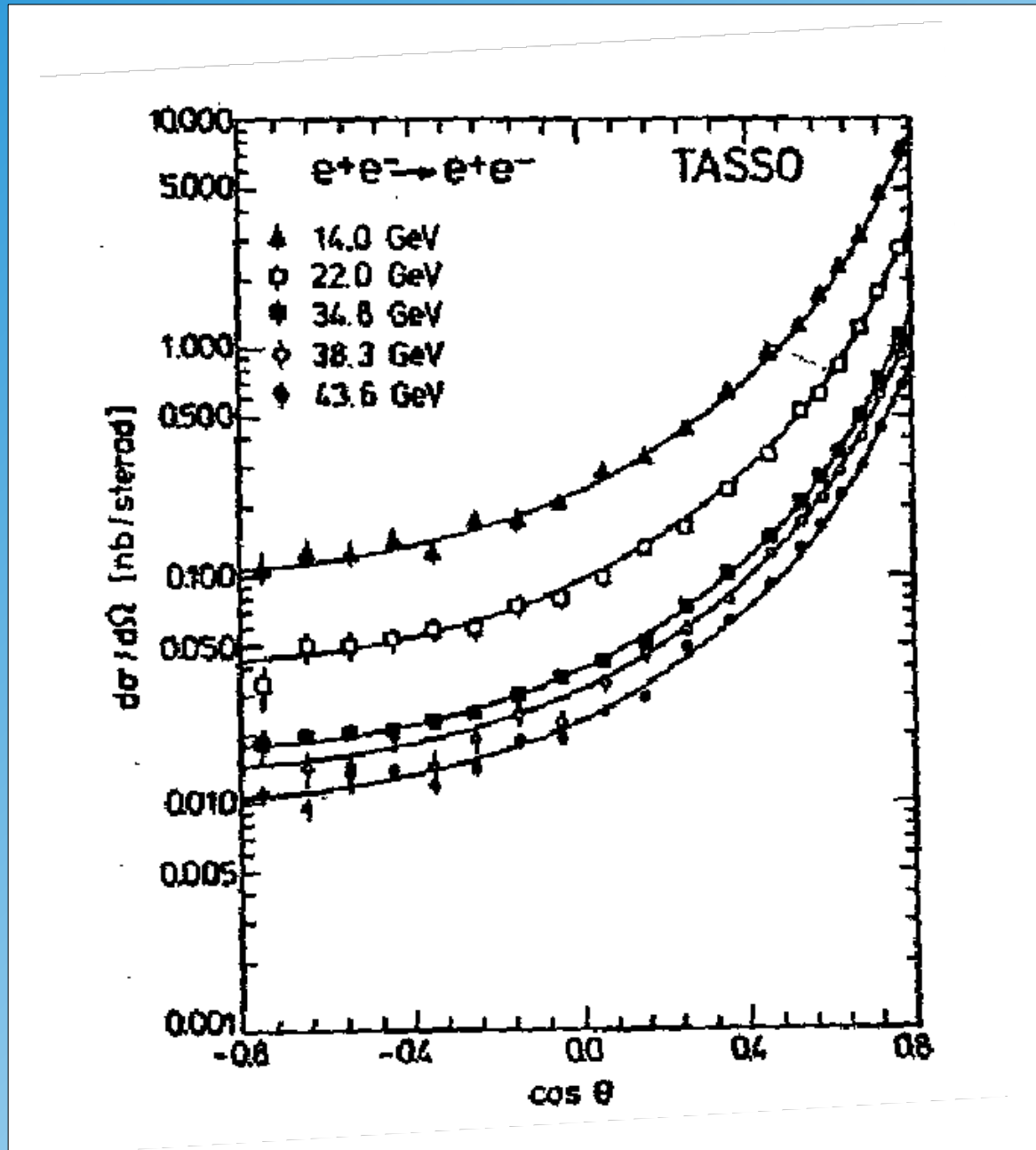
predicted by  
QCD!!!



# Tasso at PETRA

QED Test:

Bhabha  
scattering





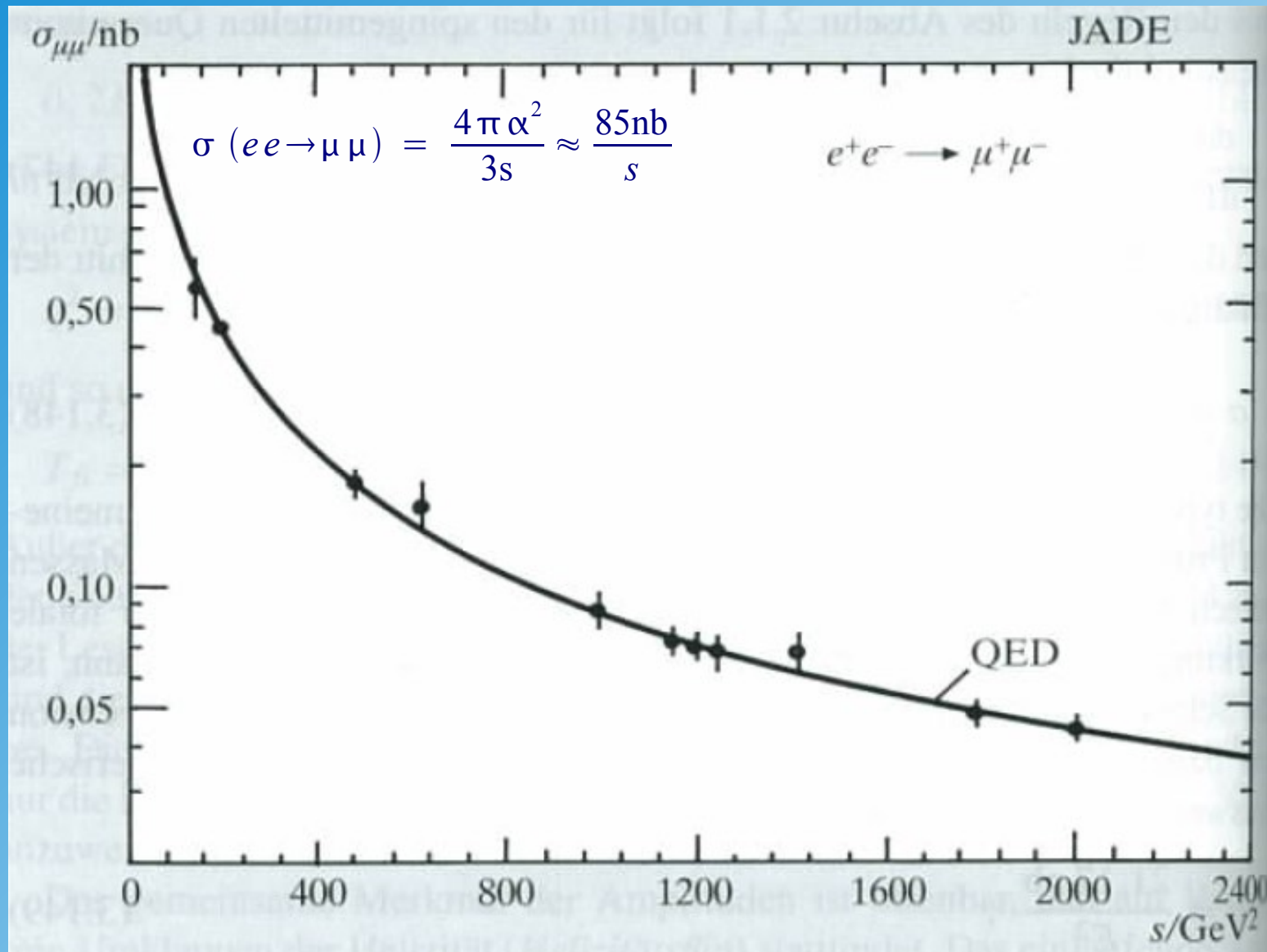
# Total Muon Pair Production C.S.

$$\frac{d\sigma}{dt} = -\frac{2\pi\alpha^2}{s^2} \frac{t^2 + u^2}{s^2} \quad s+t+u = \sum m_i^2 \approx 0 \quad (\text{only two independent})$$
$$\rightarrow u^2 = t^2 + s^2 + 2ts$$

$$\frac{d\sigma}{dt} = -2\pi\alpha^2 \frac{2t^2 + s^2 + 2ts}{s^4}$$

$$\sigma = -\int_{-s}^0 2\pi\alpha^2 \frac{2t^2 + s^2 + 2ts}{s^4} dt = -2\pi\alpha^2 \frac{-2/3s^3 - s^3 + s^3}{s^4} = \frac{4\pi\alpha^2}{3s}$$

# Myon Pair Production



PETRA accelerator (DESY)

# Quark-Pair Production

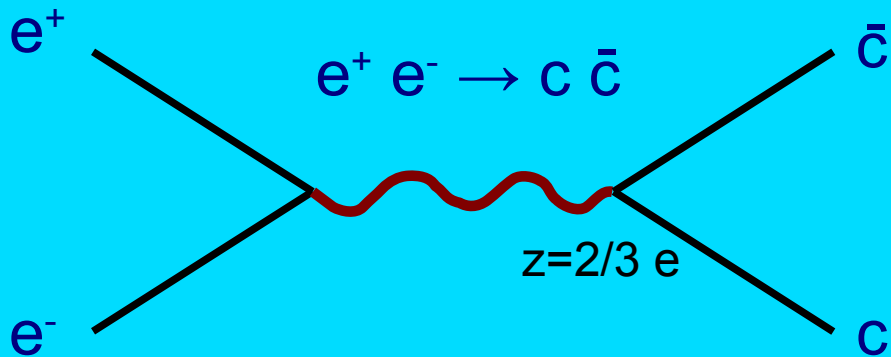
## Difficulty:

quarks and antiquarks are exp. difficult to distinguish

$$t = -\frac{s}{2} (1 \mp \cos \theta) \quad \text{different signs for quarks and antiquarks}$$

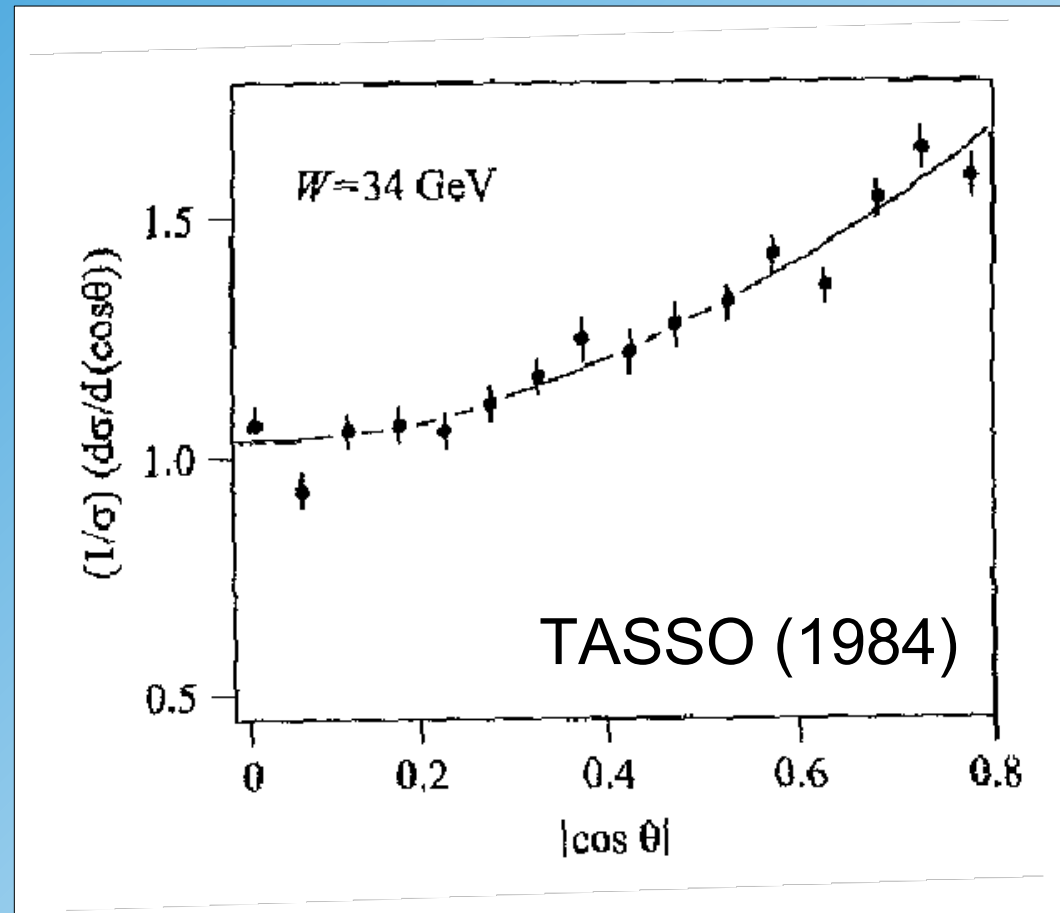
$$u = -\frac{s}{2} (1 \pm \cos \theta)$$

$$t^2 + u^2 = -\frac{s}{2} (1 + \cos^2 \theta) \quad \text{quarks and antiquarks averaged!}$$



similar to muon-pair production

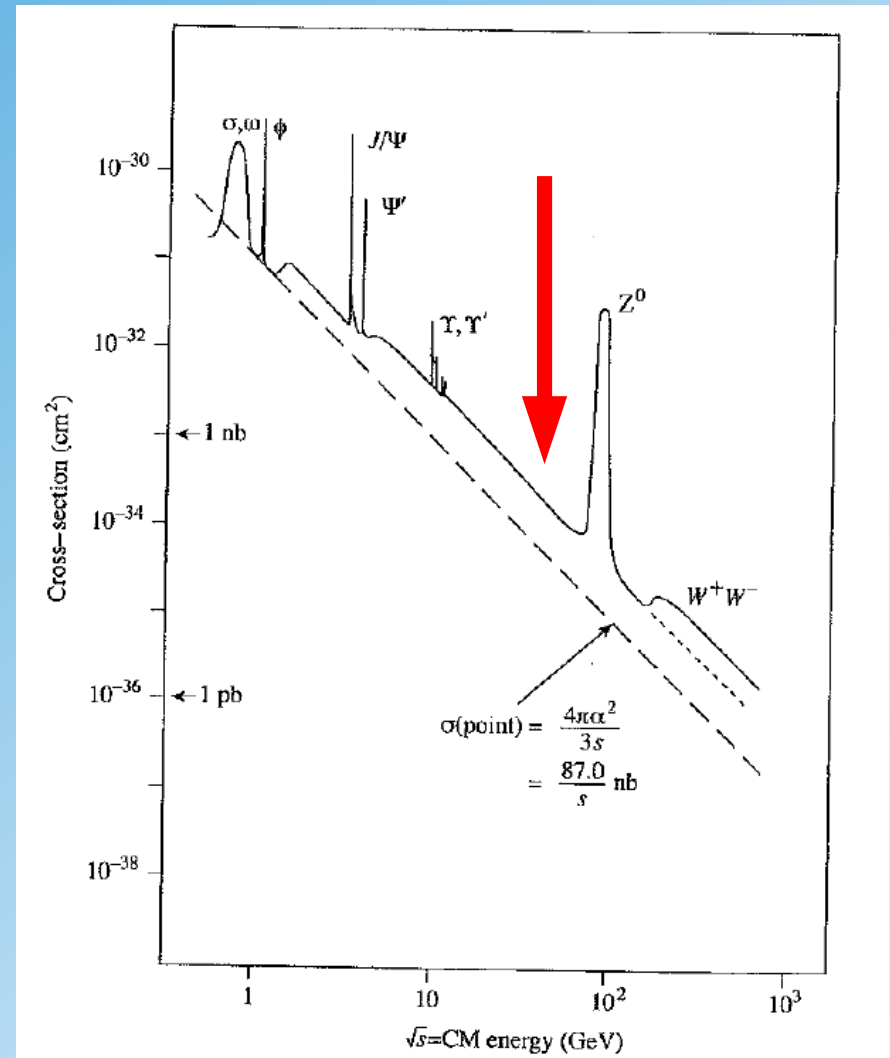
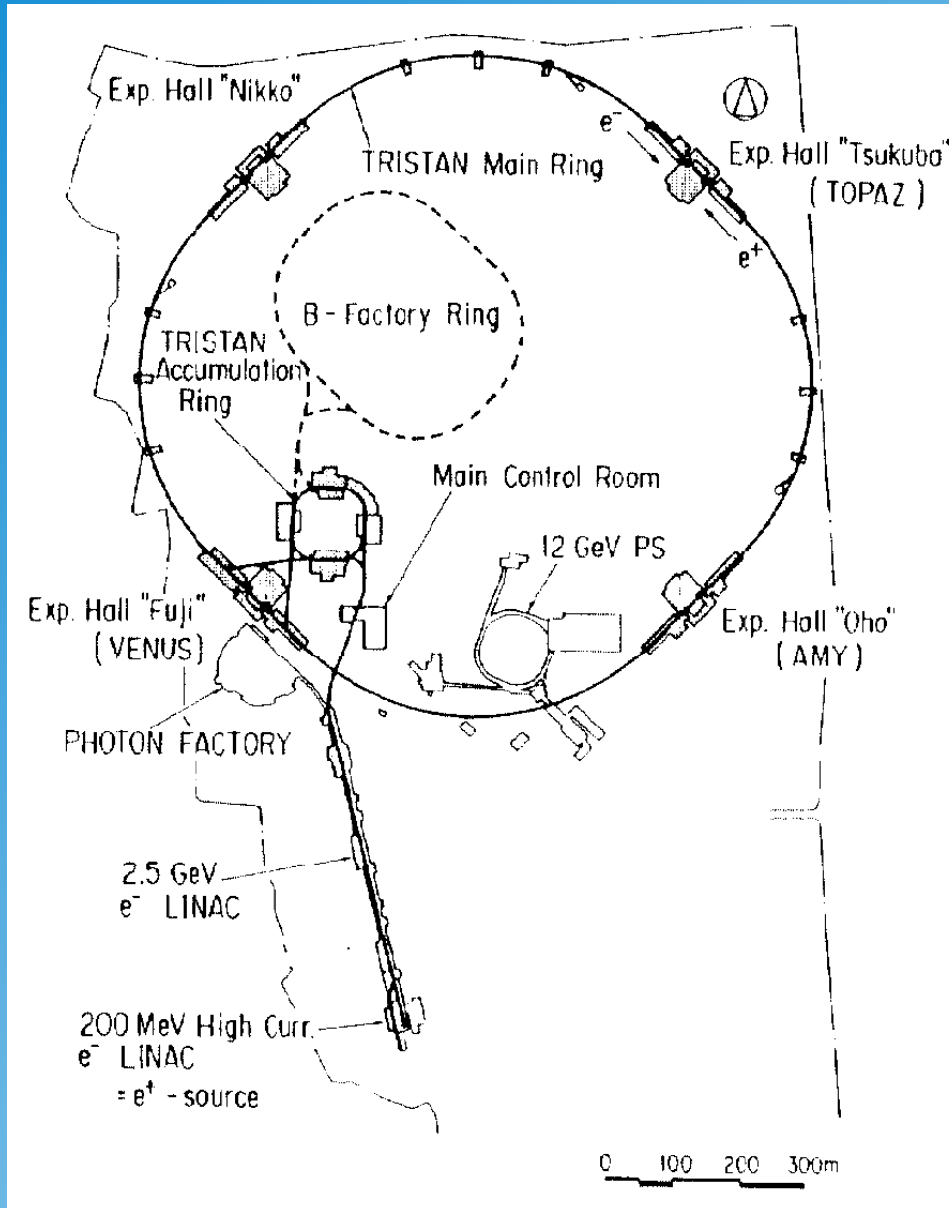
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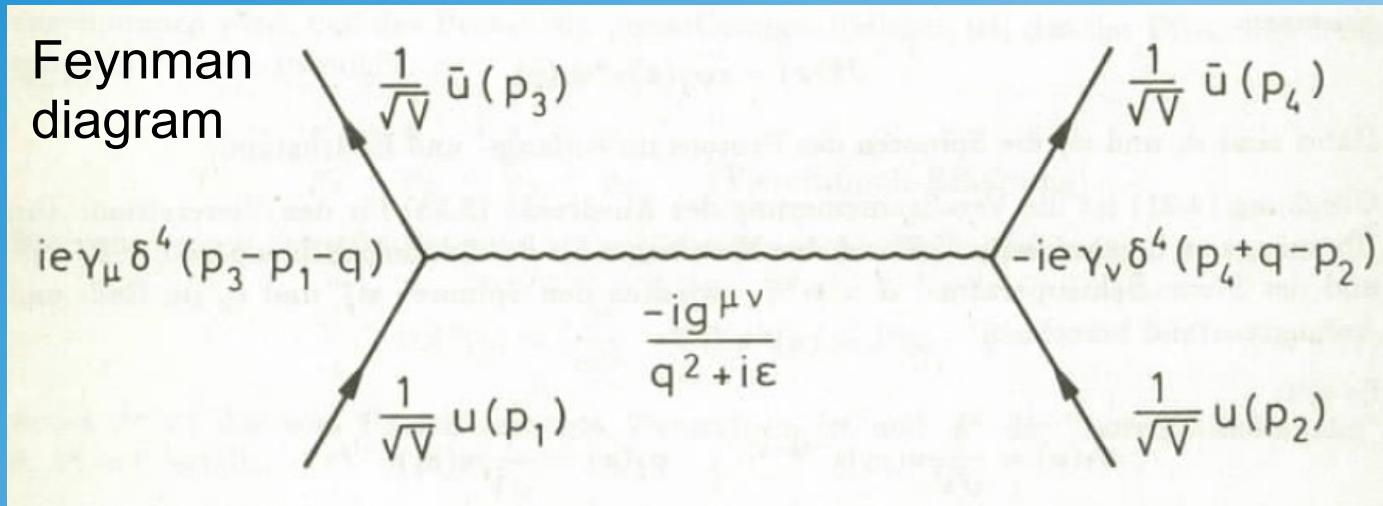
# Tristan Collider at KEK

1986-1989:  $s^{1/2}=50-64$  GeV

Search for the top in the "desert"



# QED Tests in $e^+e^-$ collisions



## Possible tests:

- universality of charges (leptons, quarks, ...)
- energy dependence of coupling (“running”)
- test of perturbation theory
- Lorentz structure of coupling
- propagator effect  $\rightarrow$  new physics
- test crossing symmetries ( $\rightarrow$  gauge invariance)

# Measurement of $R_{\text{had}}$

## Test of Quark Charges

$$R = \frac{e^+ e^- \rightarrow \text{hadrons}}{e^+ e^- \rightarrow \mu^+ \mu^-}$$

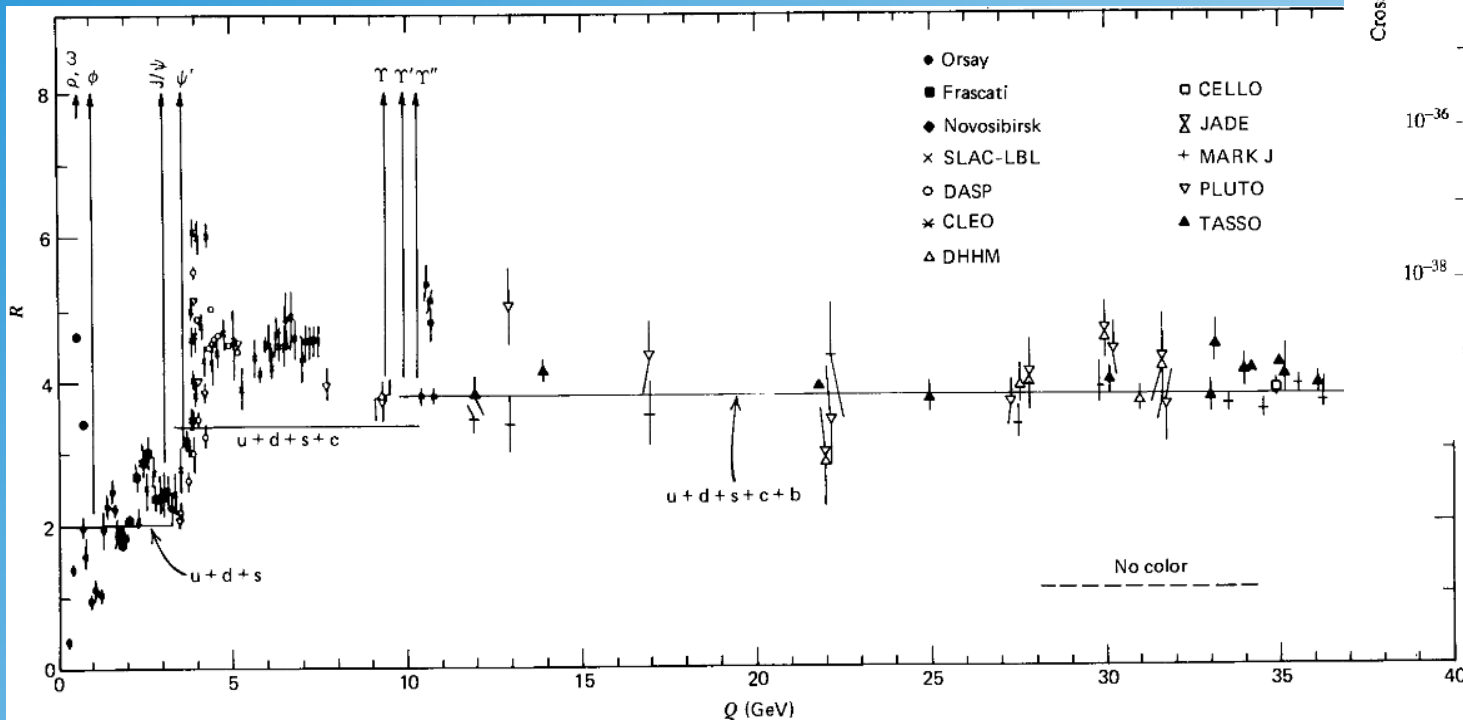
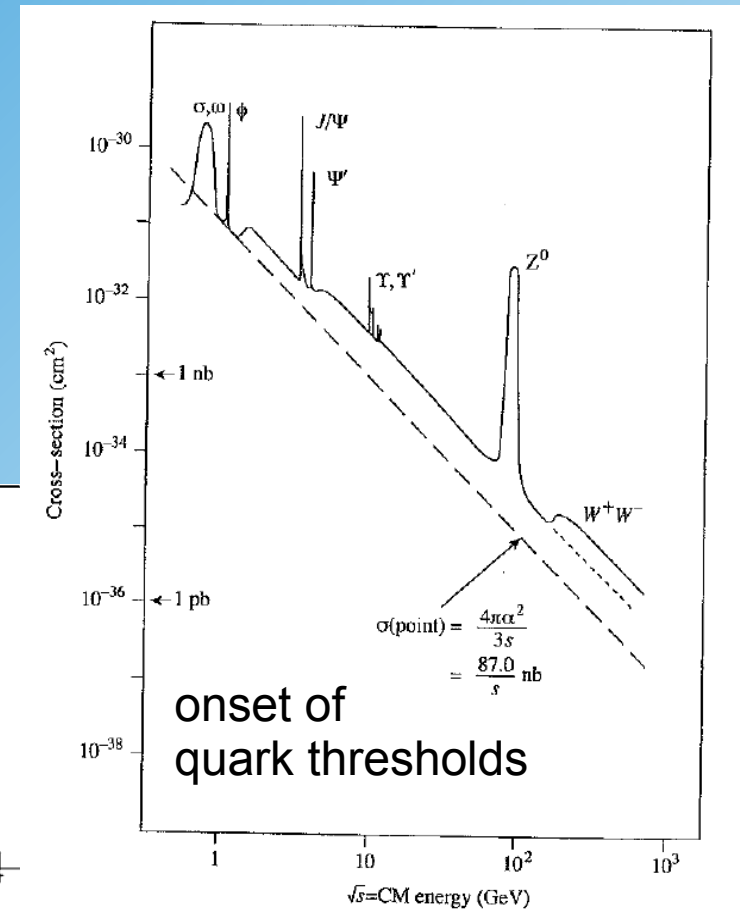
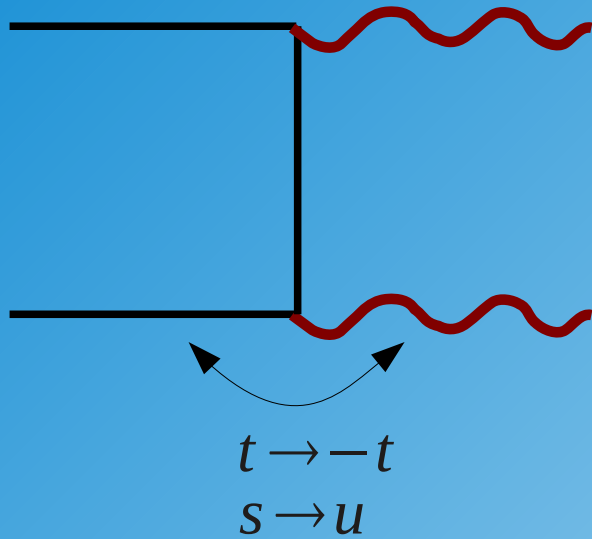


Fig. 11.3 Ratio  $R$  of (11.6) as a function of the total  $e^-e^+$  center-of-mass energy. (The sharp peaks correspond to the production of narrow  $1^-$  resonances just below or near the flavor thresholds.)

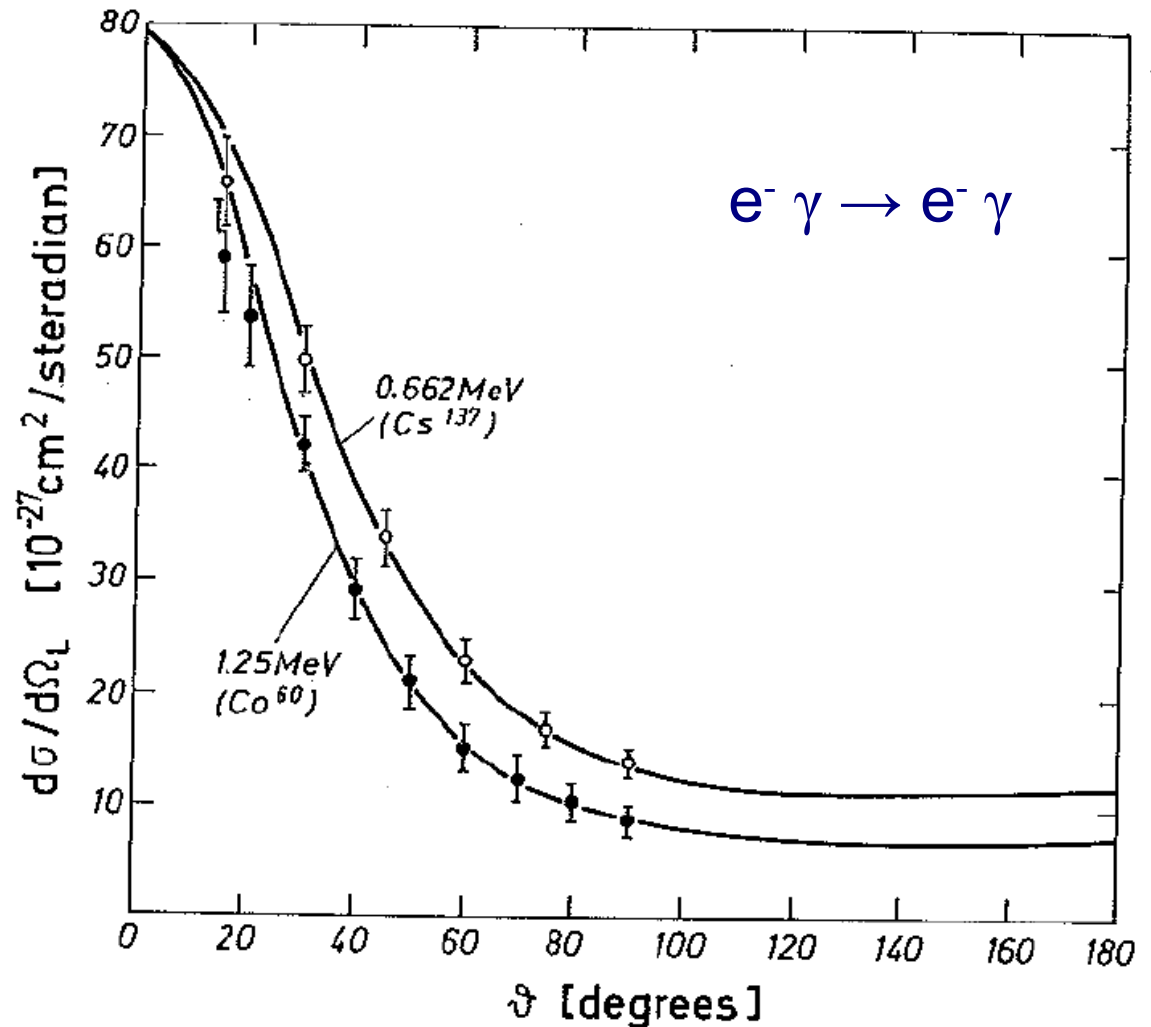
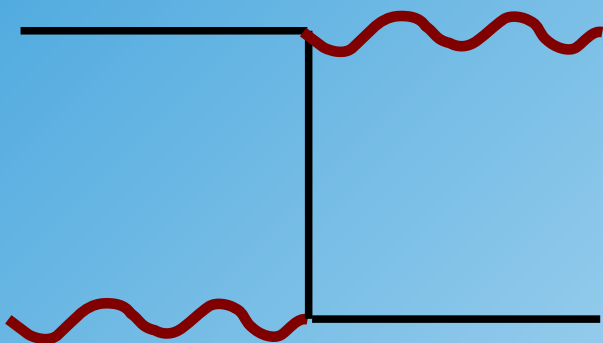


# Crossing Symmetries

$$\frac{d\sigma}{d\Omega}(e^+e^- \rightarrow \gamma\gamma) = \frac{\alpha^2}{2s} \left( \frac{u}{t} + \frac{t}{u} \right)$$



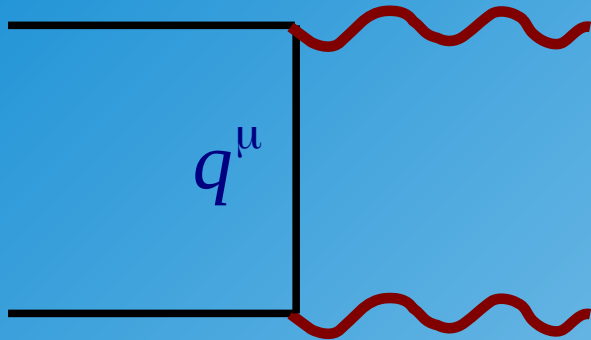
$$\frac{d\sigma}{d\Omega}(\gamma e^- \rightarrow \gamma e^-) = \frac{\alpha^2}{2s} \left( \frac{-s}{u} + \frac{-u}{s} \right)$$



**Figure 10.7** Experimental results for Compton scattering. The curves correspond to the Klein–Nishina formula (10.41) for photon energies  $\omega = 0.662$  MeV and  $\omega = 1.25$  MeV. The experimental data are from Hofstadter (1949) and Bernstein (1956) (after Evans 1958).

# The Low Energy Limit

Electromagnetic coupling at low energy:



$$Q^2 = -t = -(q^\mu q_\mu) \rightarrow 0$$

$Q^2 = m_e^2$  is kinematically forbidden

Thompson scattering cross section

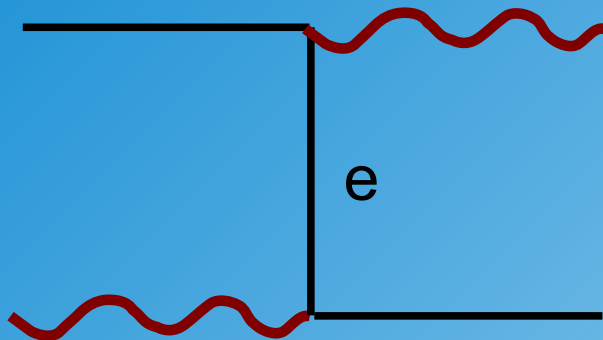
$$\sigma_t = \frac{8\pi}{3r_e^2} = \frac{8\pi}{3} \left( \frac{\alpha \lambda_c}{2\pi} \right)^2$$

used to determine  $\alpha$



# The Low Energy Limit

Electromagnetic coupling at low energy:



$$Q^2 = -t = -(q^\mu q_\mu) \rightarrow m_e^2$$

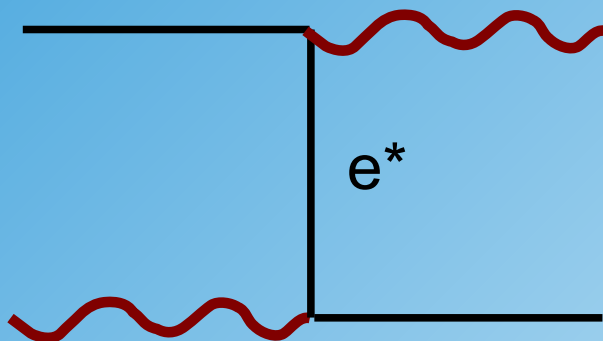
$Q^2 = m_e^2$  is kinematically forbidden

Thompson scattering cross section

$$\sigma_t = \frac{8\pi}{3r_e^2} = \frac{8\pi}{3} \left( \frac{\alpha \lambda_c}{2\pi} \right)^2$$

not dependent on energy! Used to determine  $\alpha$

General cross section:

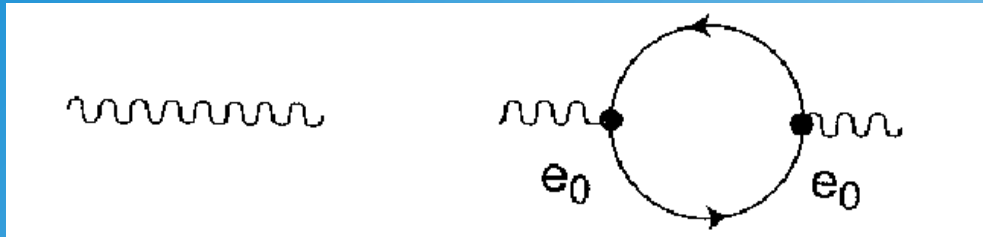


$$\frac{d\sigma}{d\Omega} (\gamma e^- \rightarrow \gamma e^-) = \frac{\alpha^2}{2s} \left( \frac{-s}{u} + \frac{-u}{s} \right)$$

# Running of $\alpha_{em}$

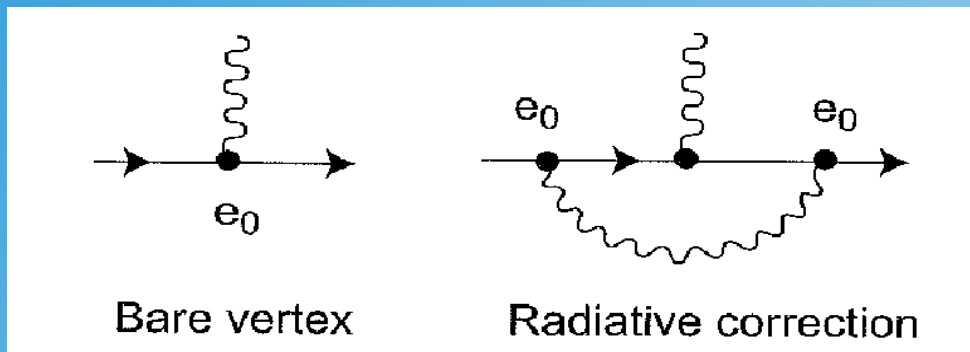
$$\alpha(Q=0)=1/137 \longrightarrow \alpha(Q=90 \text{ GeV})=1/128$$

self-energy corrections



alpha is not a constant!

vertex corrections



dressed charge!

- measure em. coupling for different (high) energies
- search for new physics effects at mass scale  $\Lambda$

$$\frac{d\sigma}{d\Omega} = \frac{d\sigma(QED)}{d\Omega} \left( 1 + \frac{s}{\Lambda^2} \right)$$

# Lorentz-Structure of Electromagnetic Interaction

From Maxwell Equations:

$$\partial_\nu \partial^\nu A^\mu(x) = e J^\mu(x) \quad \text{in QED: } \partial_\mu j_V^\mu = 0 \quad (\text{conservation of currents})$$

electromagnetic interaction described by vector currents!

Also true at high energies?

Vector Current:

$$j_V^\mu = \bar{\psi} \gamma^\mu \psi$$

scalar coupling:

$$\lambda = \bar{\psi} \psi$$

Axial-vector Current:

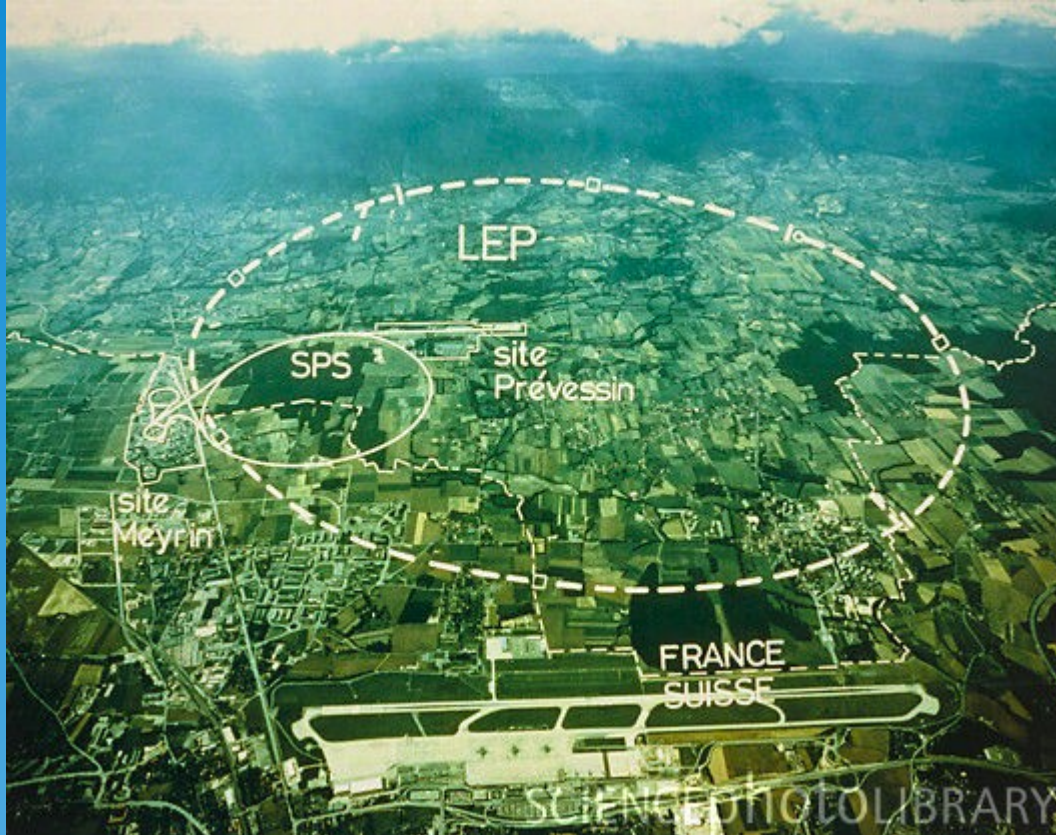
$$j_A^\mu = \bar{\psi} \gamma^\mu \gamma^5 \psi$$

pseudoscalar coupling:

$$\lambda = \bar{\psi} \gamma^5 \psi$$

→ lead in general to different angular distributions!

# LEP Collider



biggest electron-positron  
collider with up to 200 GeV  
centre of mass

4 experiments:  
ALEPH, DELPHI, L3, OPAL

LEP1: “Z-factory”

LEP2: “WW factory”



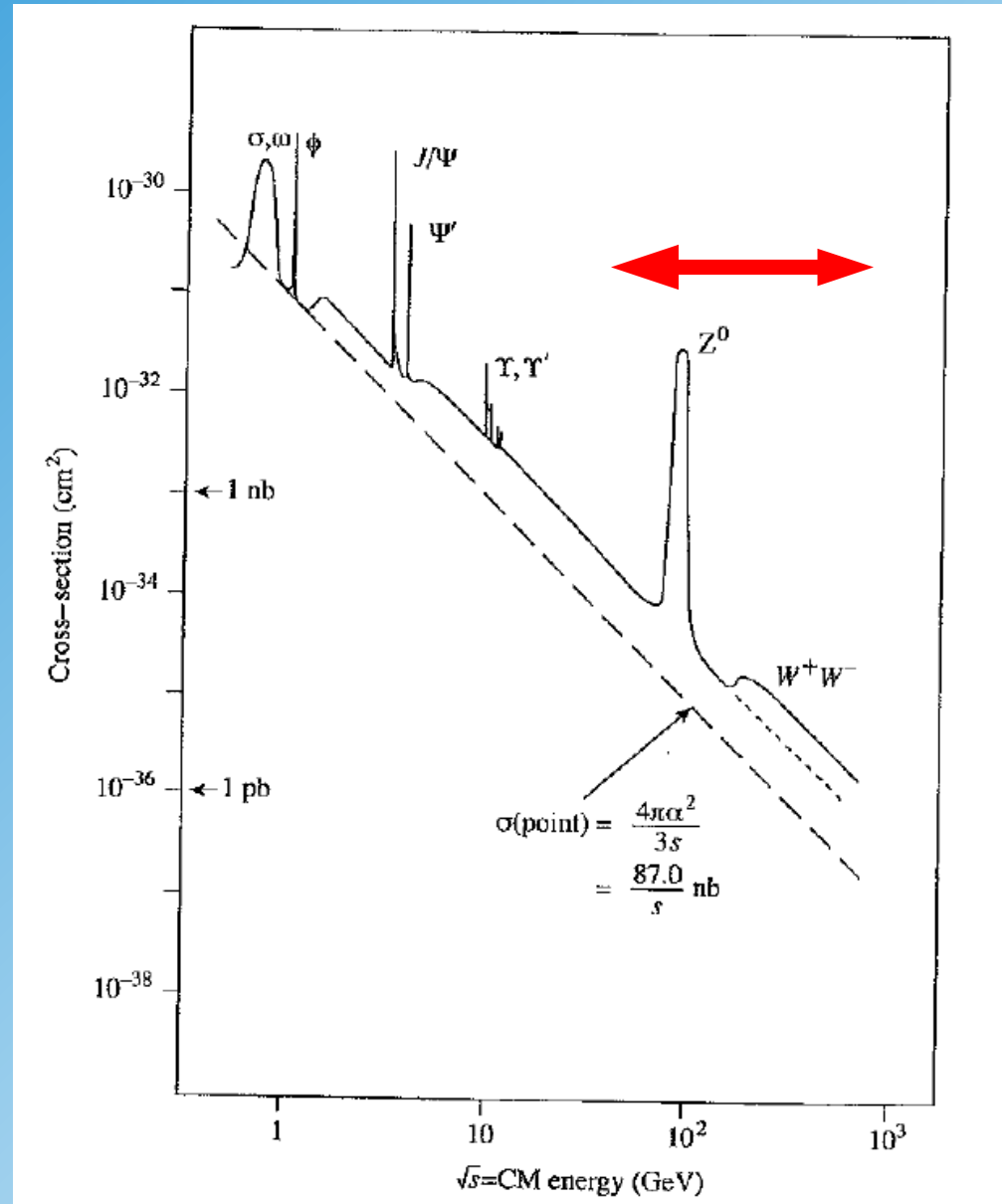
# LEP Regime

Processes:

$$e^+ e^- \rightarrow Z$$

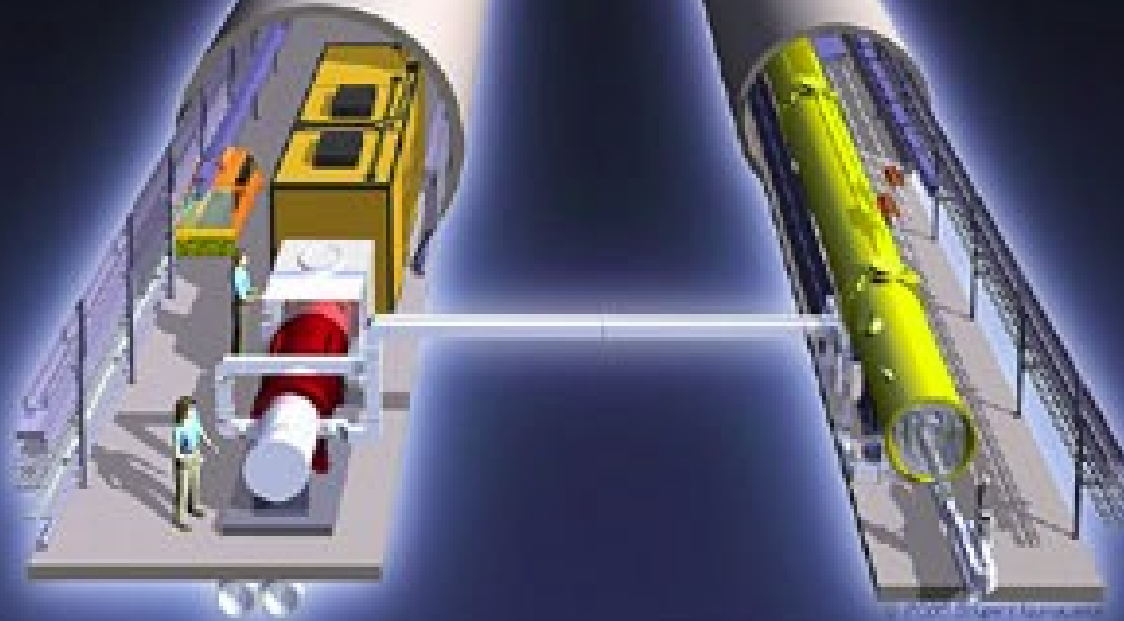
$$e^+ e^- \rightarrow W^+ W^-$$

at LEP energies  
radiative effects  
play an important role!



# International Linear Collider

500 GeV electrons x 500 GeV positrons



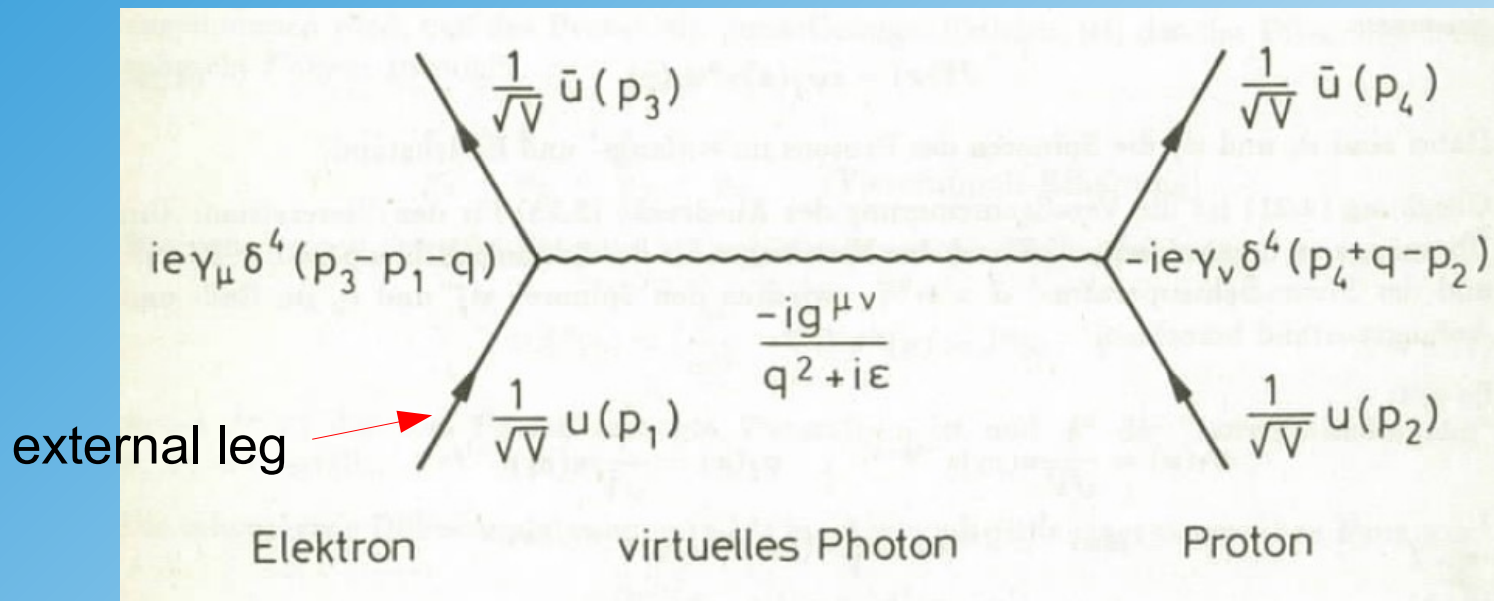
# Backup

# Fermion-Fermion Scattering

$$S_{fi}^{(1)} = ie^2 \int \frac{d^4 q}{(2\pi)^4} \delta^4(p_3 - p_1 - q) \delta^4(p_4 + q - p_2) (2\pi)^8$$

$$\cdot \frac{1}{V^2} \bar{u}(p_3) \gamma_\mu u(p_1) \cdot \frac{-g^{\mu\nu}}{q^2 + i\epsilon} \cdot \bar{u}(p_4) \gamma_\nu u(p_2)$$

particle-particle (here electron-proton) scattering



lowest order perturbation theory: leading order graph (Born)



# Gamma Matrices I

Gamma matrices  $\gamma^\mu$  are chosen such that  $\gamma^0$  is hermitian while  $\gamma^k$  ( $k=1,2,3$ ) are anti-hermitian

$$\begin{aligned}(\gamma^0)^\dagger &= \gamma^0, & (\gamma^0)^\mu &= 1, \\ (\gamma^k)^\dagger &= -(\gamma^k), & (\gamma^k)^\mu &= -1 \quad (k=1,2,3)\end{aligned}$$

We define  $\gamma^5$  as the hermitian matrix:  $\gamma^5 = i \gamma^0 \gamma^1 \gamma^2 \gamma^3$ ,  $(\gamma^5)^2 = 1$ ,

The 4x4 gamma matrices can be represented by (representation where  $\gamma^0$  is diagonal):

$$\gamma^k = \begin{pmatrix} 0 & \underline{\sigma}^k \\ -\underline{\sigma}^k & 0 \end{pmatrix}, \quad \gamma^0 \stackrel{\text{def}}{=} \beta = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \quad \gamma^5 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \quad (\text{note several representations exist})$$

With the 2x2 Pauli matrices:

$$\underline{\sigma}^1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \underline{\sigma}^2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \underline{\sigma}^3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}.$$

Gamma matrices anti-commute:  $\gamma^i \gamma^k + \gamma^k \gamma^i = 0$  for  $i \neq k$

# Gamma Matrices II

$$\gamma^0 = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}$$

$$\gamma^1 = \begin{pmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & -1 & 0 & 0 \\ -1 & 0 & 0 & 0 \end{pmatrix}$$

$$\gamma^2 = \begin{pmatrix} 0 & 0 & 0 & -i \\ 0 & 0 & i & 0 \\ 0 & i & 0 & 0 \\ -i & 0 & 0 & 0 \end{pmatrix}$$

$$\gamma^3 = \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \\ -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{pmatrix}$$

$$\gamma^5 = \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{pmatrix}$$

(in other representations  
 $\mathfrak{g}$  is diagonal )

