

Lecture:

Standard Model of Particle Physics

Heidelberg SS 2013

Tests of the Standard Model II

Contents

- Intro: Statistics and limit setting
- Higgs Mass Predictions
- Higgs Searches at LEP
- Higgs Searches at Hadron Colliders
- WW-Scattering Amplitude

Statistics and Limit Setting

chi² fit:

$$\chi^2 = \sum_i \frac{(y_i - \mu_i)^2}{\sigma_i^2}$$

y_i measurement

μ_i model prediction (nuisance parameter)

σ_i uncertainty (statistical and systematical)

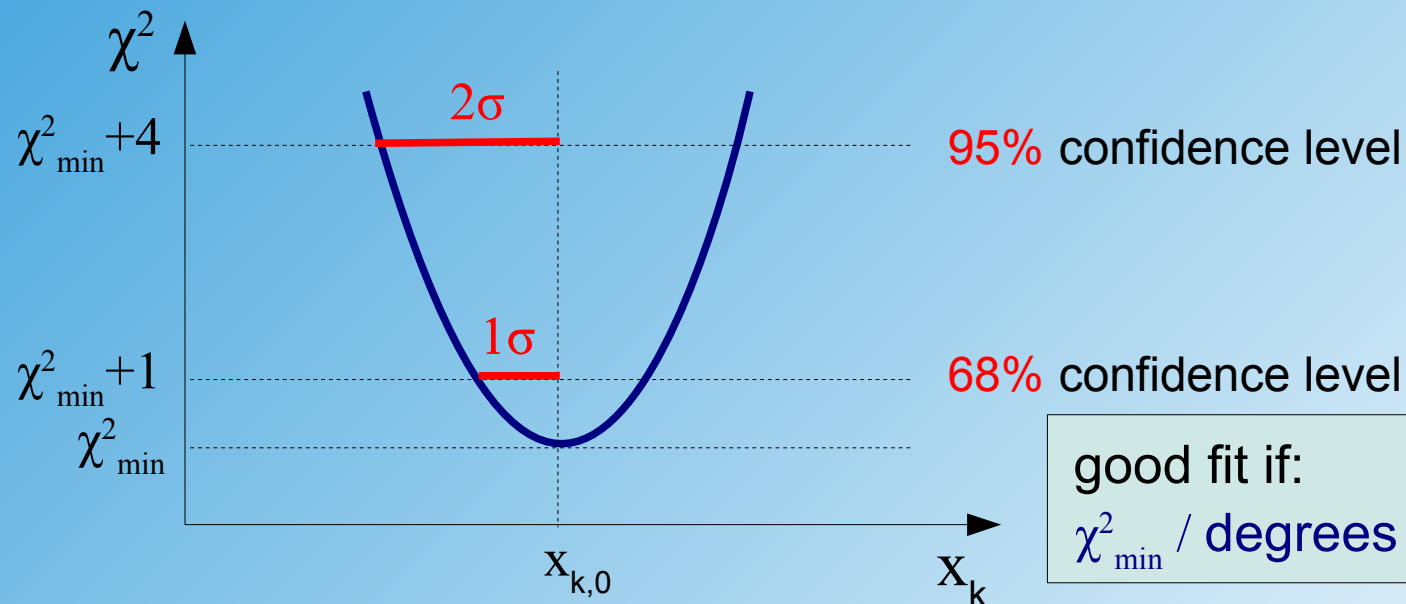
chi² fit with correlated errors:

$$\chi^2 = \sum_i \sum_j (y_i - \mu_i) \text{cov}_{ij}^{-1} (y_j - \mu_j) \quad \text{cov}_{ij} \text{ covariance (error) matrix}$$

Parameter Fit:

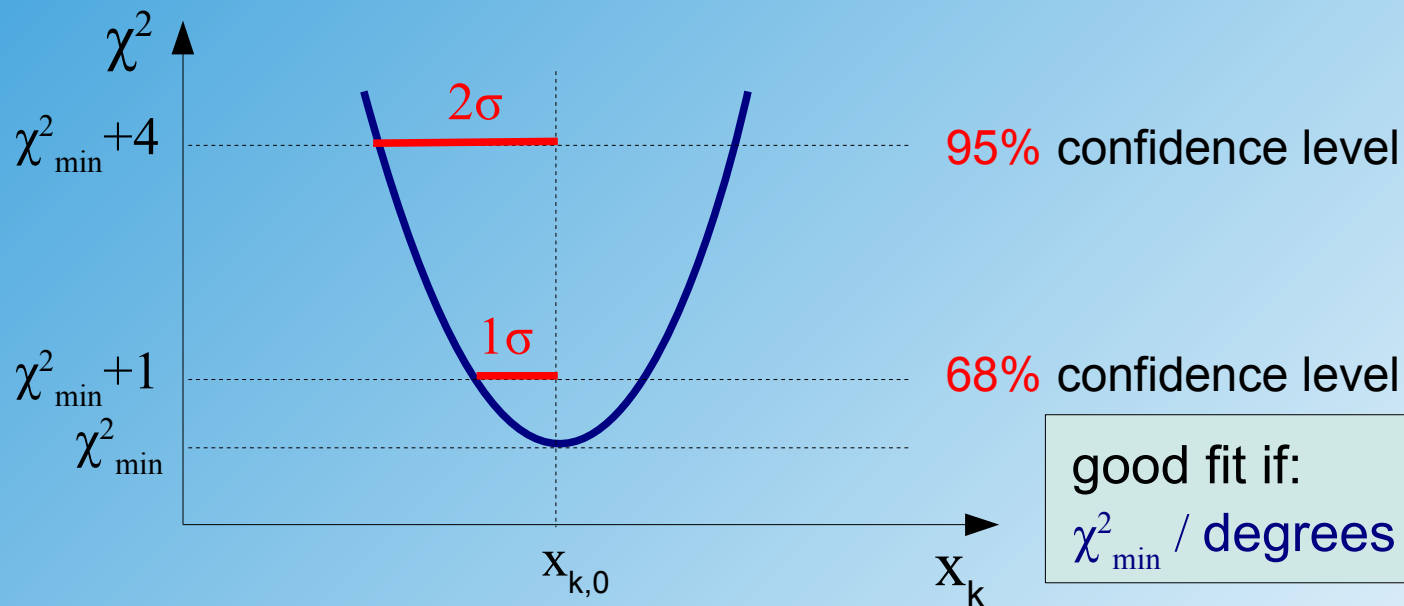
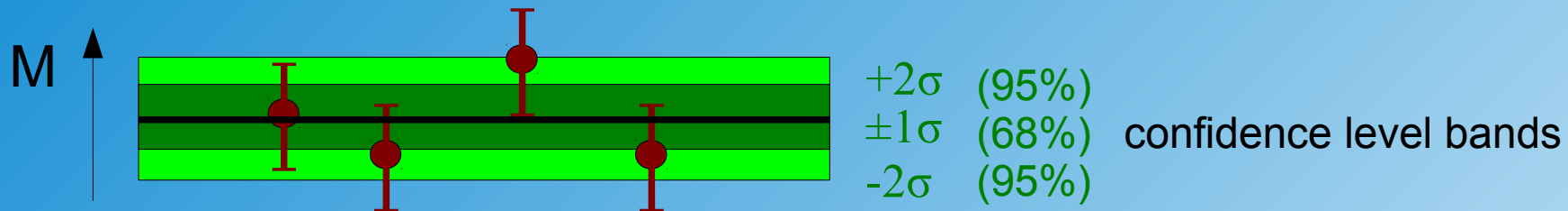
$$\mu_i = \mu_i(x_1, x_2, \dots, x_n)$$

x_k model parameter



Example Fit

Measurement of some mass (1-Parameter fit) from 4 experiments:



Probability Densities

for the above example a gaussian probability density was used

Gaussian (normal) distribution:

$$f(x) = \frac{1}{\sqrt{2\pi}\sigma} e^{-\frac{(x-\mu)^2}{2\sigma^2}} \rightarrow P(a < x < b) = \int_a^b f(x) dx$$

→ used for systematic uncertainties (symmetric)

Poissonian distribution:

$$P(N) = \frac{e^{-\mu} \mu^N}{N!}$$

→ used for statistical uncertainties

Note, Poisson distribution approaches Gaussian distribution for large μ

Limit Setting Philosophies

Baysian Method:

- based on the experiment posterior, exclusion limits are calculated
- low probability models are excluded
- probabilities are assigned to models using a prior

“natural method” but choice of prior is arbitrary

Frequentist Method:

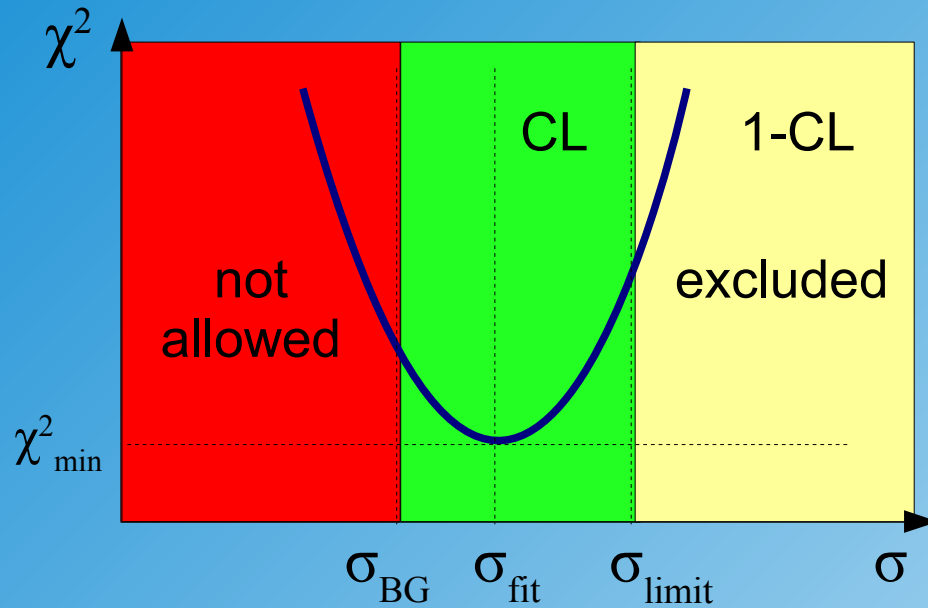
- based on Monte Carlo toy experiments probabilities are assigned to all possible experimental outcomes
- exclusion limit is set by that model which excludes this experimental outcome with certain confidence interval

computationally expensive and might give unphysical results
(e.g. negative cross sections)

Bayesian Method

Model: $N = N_{BG} + N_{Sig}$ (background + signal)

$\sigma = \sigma_{BG} + \sigma_{Sig} = N/L$ (choose cross section as "prior")



Probability:

$$P \propto e^{-\frac{1}{2}(\chi^2(\sigma) - \chi_{fit}^2)}$$

additional constraint:

$$\sigma > \sigma_{BG} \quad \text{because} \quad \sigma_{Sig} > 0$$

$$CL = \frac{\int_{\sigma > \sigma_{BG}}^{\sigma_{CL}} P(\sigma) d\sigma}{\int_{\sigma > \sigma_{BG}}^{\infty} P(\sigma) d\sigma}$$

CL = confidence level

Choice of Prior

- Cross sections depend on couplings
- Choose alternatively coupling as prior

$$\sigma_{Sig} = \sigma_0 \alpha^2$$

$$\frac{d\sigma}{d\alpha} = \frac{d\sigma_{Sig}}{d\alpha} = 2\sigma_0 \alpha$$

$$CL = \frac{\int_{\alpha>0}^{\alpha_{CL}} P(\alpha) d\alpha}{\int_{\alpha>0}^{\infty} P(\alpha) d\alpha} = \frac{\int_{\alpha>0}^{\alpha_{CL}} P(\alpha)/\alpha d\sigma}{\int_{\alpha>0}^{\infty} P(\alpha)/\alpha d\sigma} \neq \frac{\int_{\sigma_{Sig}>0}^{\sigma_{CL}} P(\sigma) d\sigma}{\int_{\sigma_{Sig}>0}^{\infty} P(\sigma) d\sigma}$$

Result depends on choice of prior!

Frequentist Method

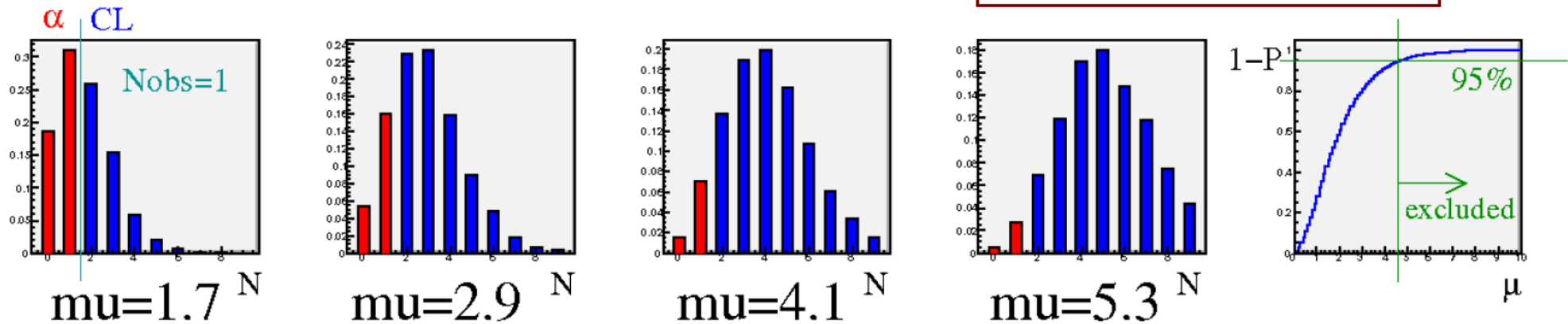
S.Schmitt

- Frequentist limit: exclude all theories which produce the data at small probability α less than $1-CL$ (typically: $CL=0.95$)

$$\alpha = P_{\mu}(N \leq N_{\text{obs}}) < 1 - CL$$

α : also called p-value

Frequentist limit:
sum (integrate) over observations up to N_{obs}
Repeat for each model



- set limit with 95% confidence level for $\mu=4.6$
- experiment has a 5% probability to happen

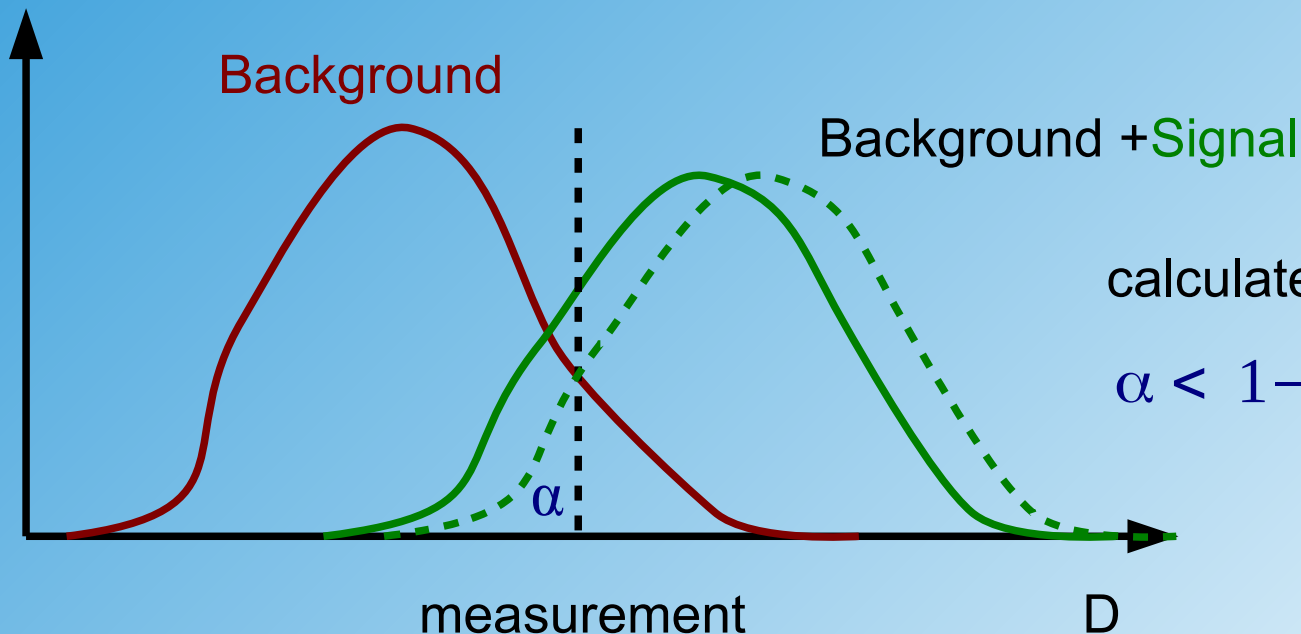
Frequentist Method

In case of many observables x_k a combined discriminator variable is often defined:

$$D = D(x_1, x_2, x_3, \dots, X_k)$$

- large discriminator means high probability
- small discriminator means low probability

Often, the output from artificial neural nets or other multivariate methods is used as discriminator variable

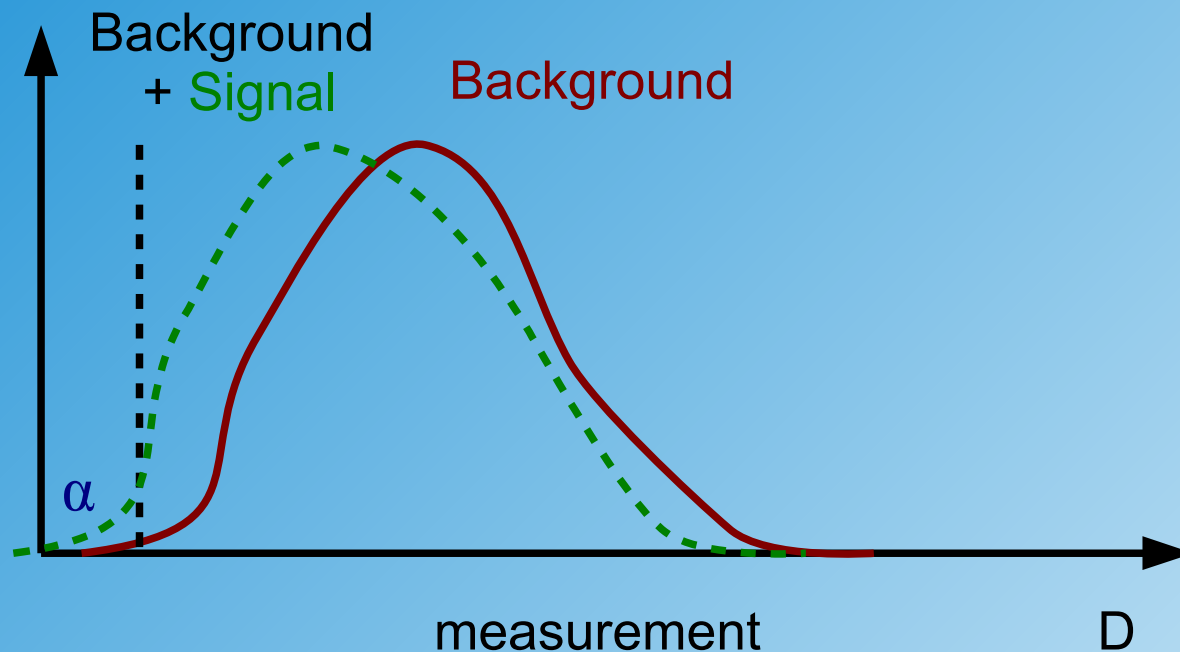


Problem with Frequentist Method

(same problem applies to Bayesian Method as well)

Problem in case of a very small measurement value with $P(\text{BG}) < (1 - \text{CL})$

- would require a negative signal cross section:



unphysical solution!

CL_S Method

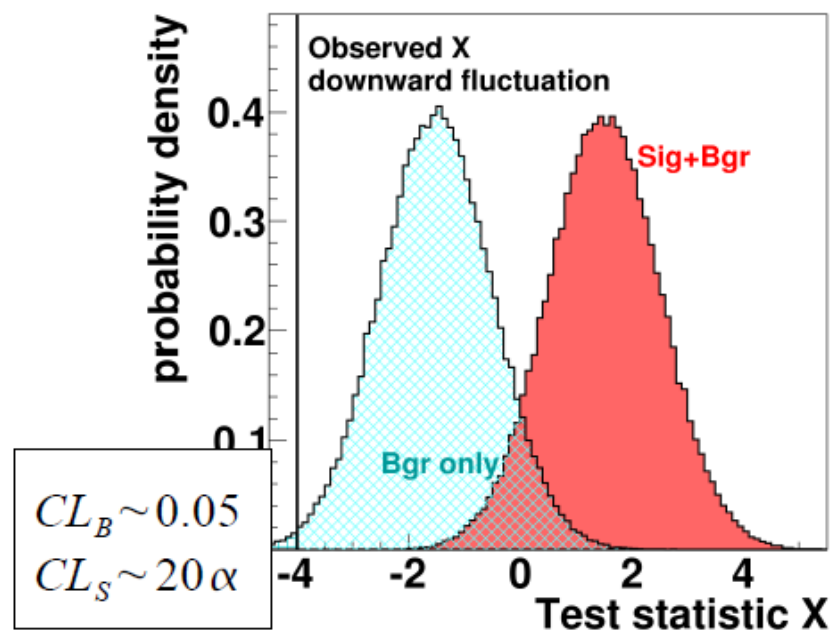
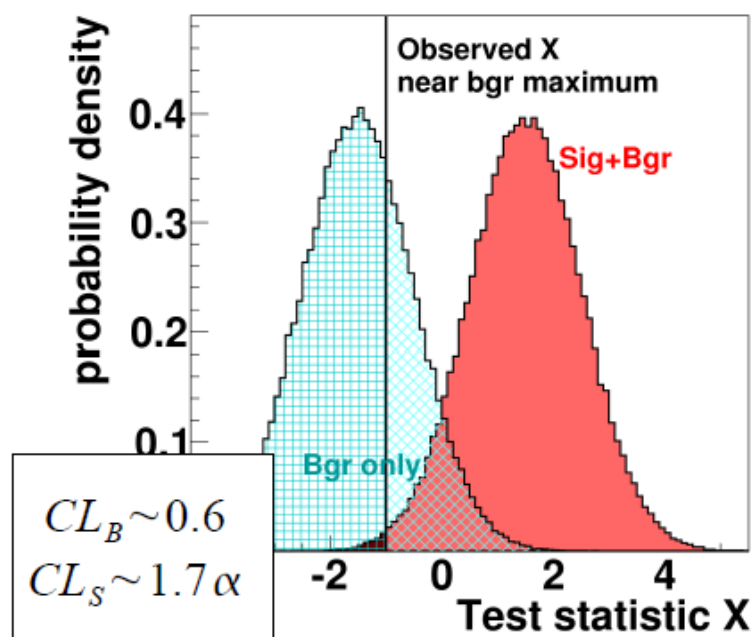
- Use ratio of two probabilities CL_S instead of α to test against CL

$$CL_{SB} = \alpha = \int_{X < X_{obs}} P(X | \text{signal} + \text{bgr}) dX$$

$$CL_B = \int_{X < X_{obs}} P(X | \text{bgr}) dX$$

$$CL_S = \frac{CL_{SB}}{CL_B}$$

- Standard model has CL_S=1 and is never excluded



S.Schmitt

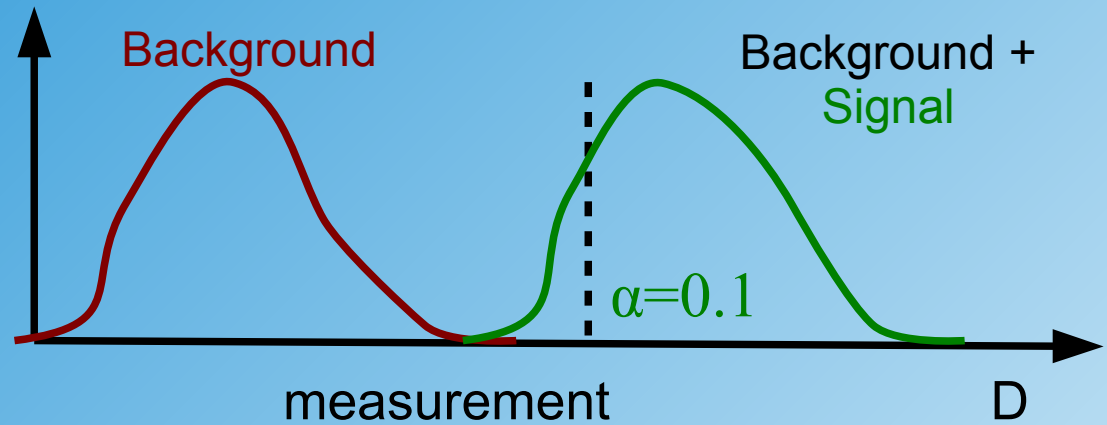
CL_S > CL_{SB} by definition!

Another Example CL_S Method

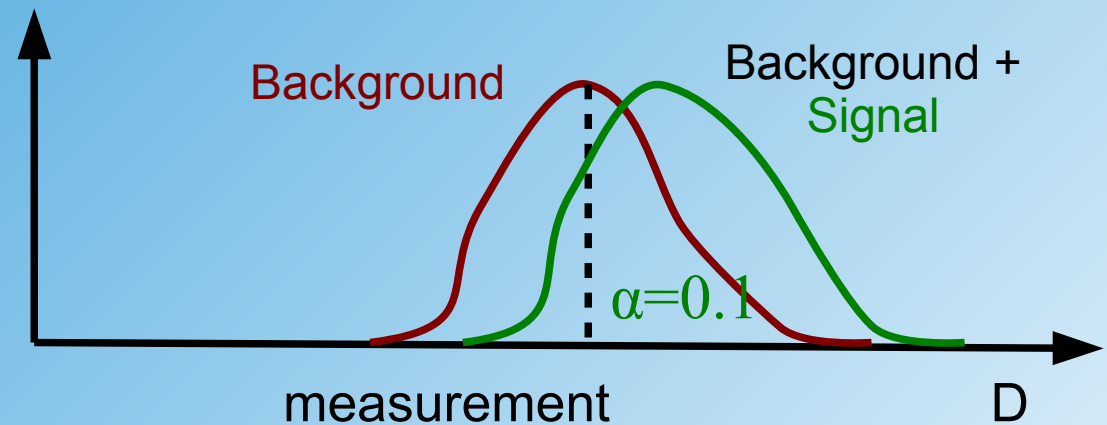
Definition:

$$CL_S = \frac{CL_{SB}}{CL_B}$$

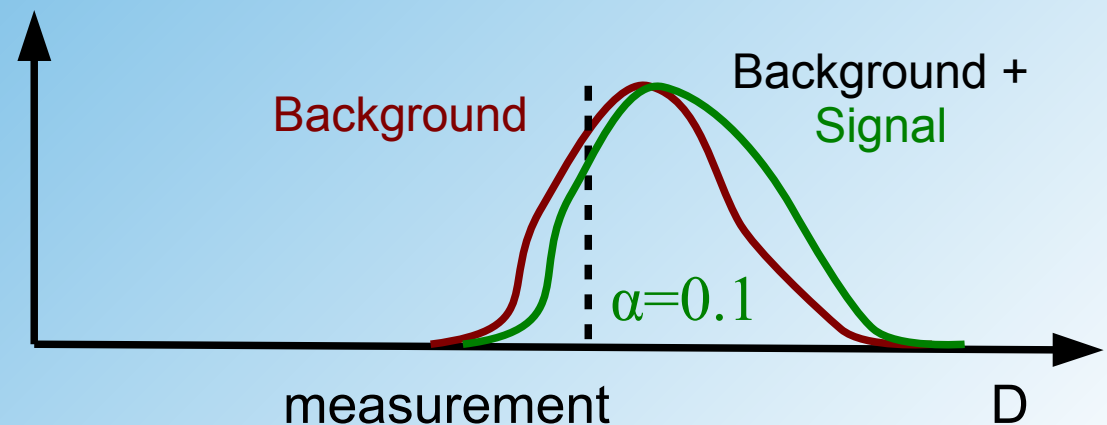
$$0.1 = \frac{0.1}{1.0}$$



$$0.2 = \frac{0.1}{0.5}$$



$$0.5 = \frac{0.1}{0.2}$$

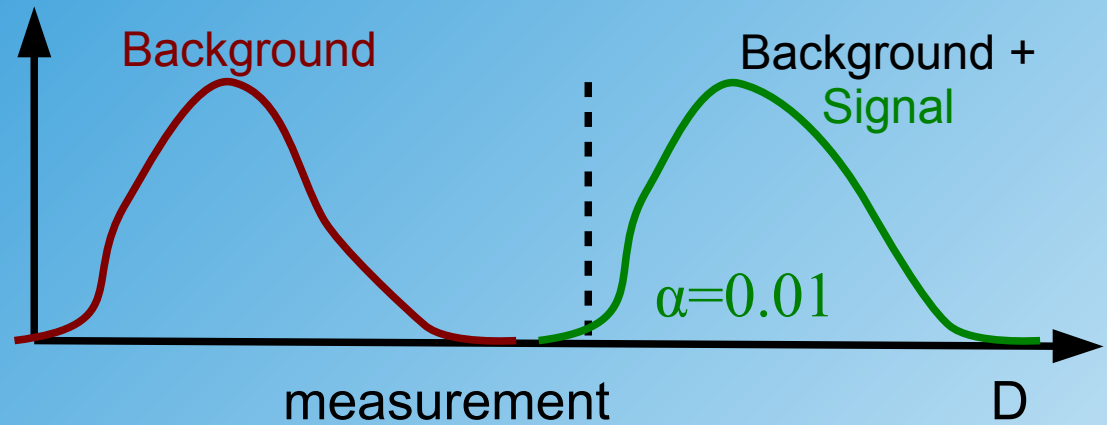


Another Example CL_S Method

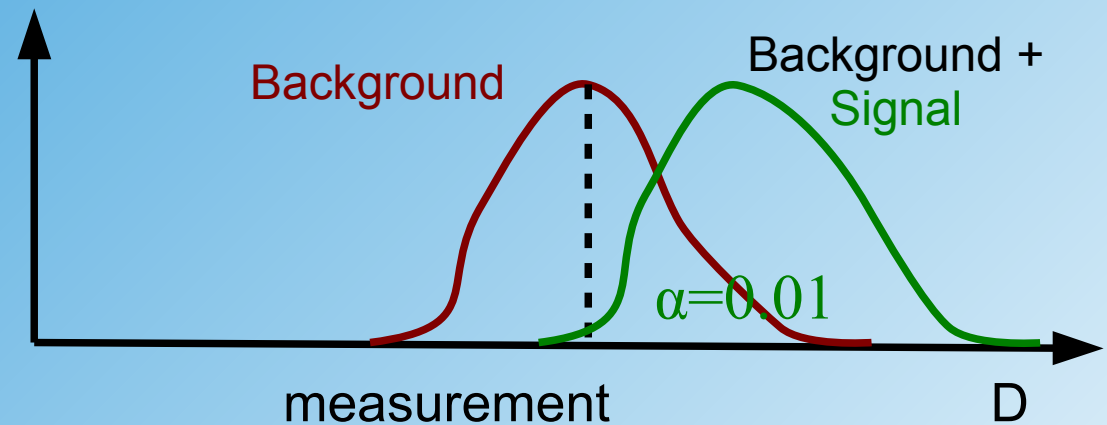
Definition:

$$CL_S = \frac{CL_{SB}}{CL_B}$$

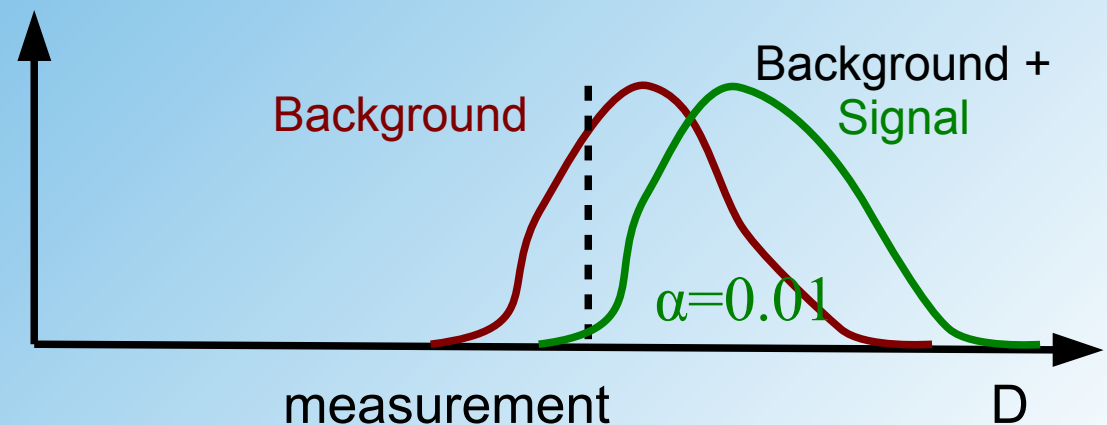
$$0.01 = \frac{0.01}{1.0}$$



$$0.02 = \frac{0.01}{0.5}$$



$$0.05 = \frac{0.01}{0.2}$$



excluded at 95% CL

Higgs Mass Constraints

Radiative Corrections and Indirect Higgs Constraints

$$\sin^2 \theta_w = 1 - \frac{M_W^2}{M_Z^2} \quad \sin \theta_w = \frac{e}{g}$$

$$\rho = \frac{m_W^2}{m_Z^2 \cos^2 \theta_W} = 1$$

$$\sin^2 \theta_W = 1 - \frac{m_W^2}{m_Z^2}$$

$$m_W^2 = \frac{\pi \alpha}{\sqrt{2} \sin^2 \theta_W G_F}$$

$\alpha(0)$

Lowest order SM predictions

Including radiative corrections

$$\bar{\rho} = 1 + \Delta\rho$$

$$\sin^2 \theta_{\text{eff}} = (1 + \Delta\kappa) \sin^2 \theta_W$$

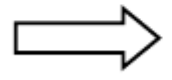
$$m_W^2 = \frac{\pi \alpha}{\sqrt{2} \sin^2 \theta_W G_F} (1 + \Delta r)$$

$$\alpha(m_Z^2) = \frac{\alpha(0)}{1 - \Delta\alpha}$$

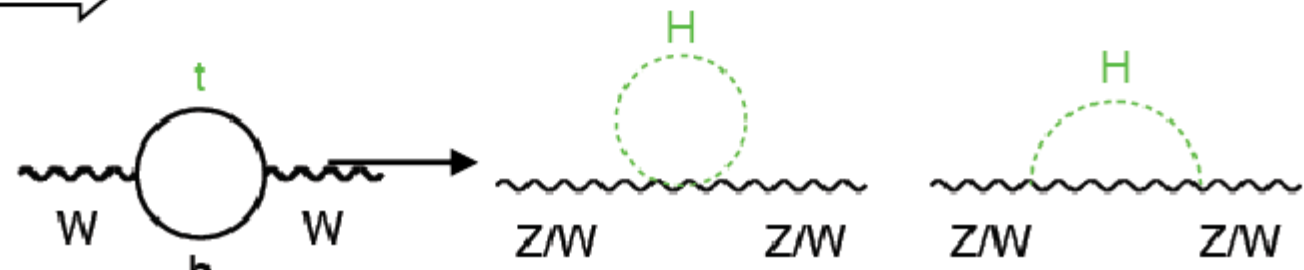
with : $\Delta\alpha = \Delta\alpha_{\text{lept}} + \Delta\alpha_{\text{top}} + \Delta\alpha_{\text{had}}^{(5)}$

$$\sin^2 \theta_w$$

$$g_A, g_V$$



$$\Delta\rho, \Delta\kappa, \Delta r = f(m_t^2, \log(m_H), \dots)$$

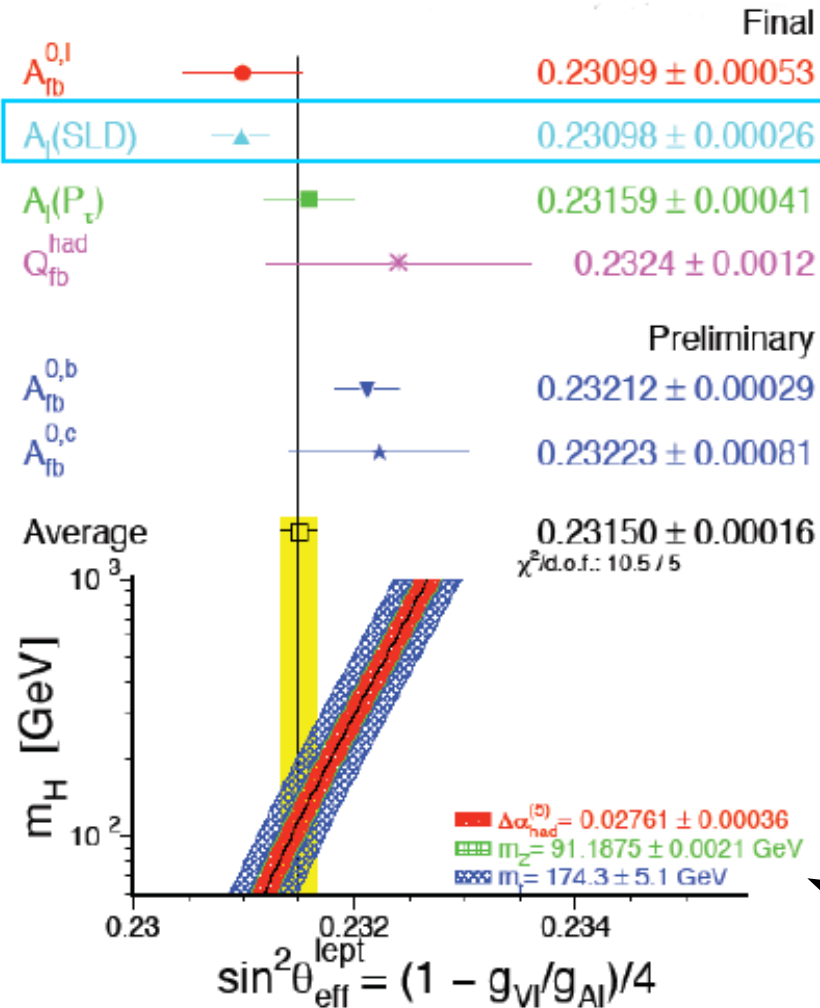


$$\sin^2 \theta_{\text{eff}}$$

$$\bar{g}_A, \bar{g}_V$$

$$\bar{g}_A = \sqrt{\bar{\rho}} T^3 \quad \bar{g}_V = \sqrt{\bar{\rho}} (T^3 - 2Q \sin^2 \theta_{\text{eff}})$$

Indirect Higgs Mass Prediction



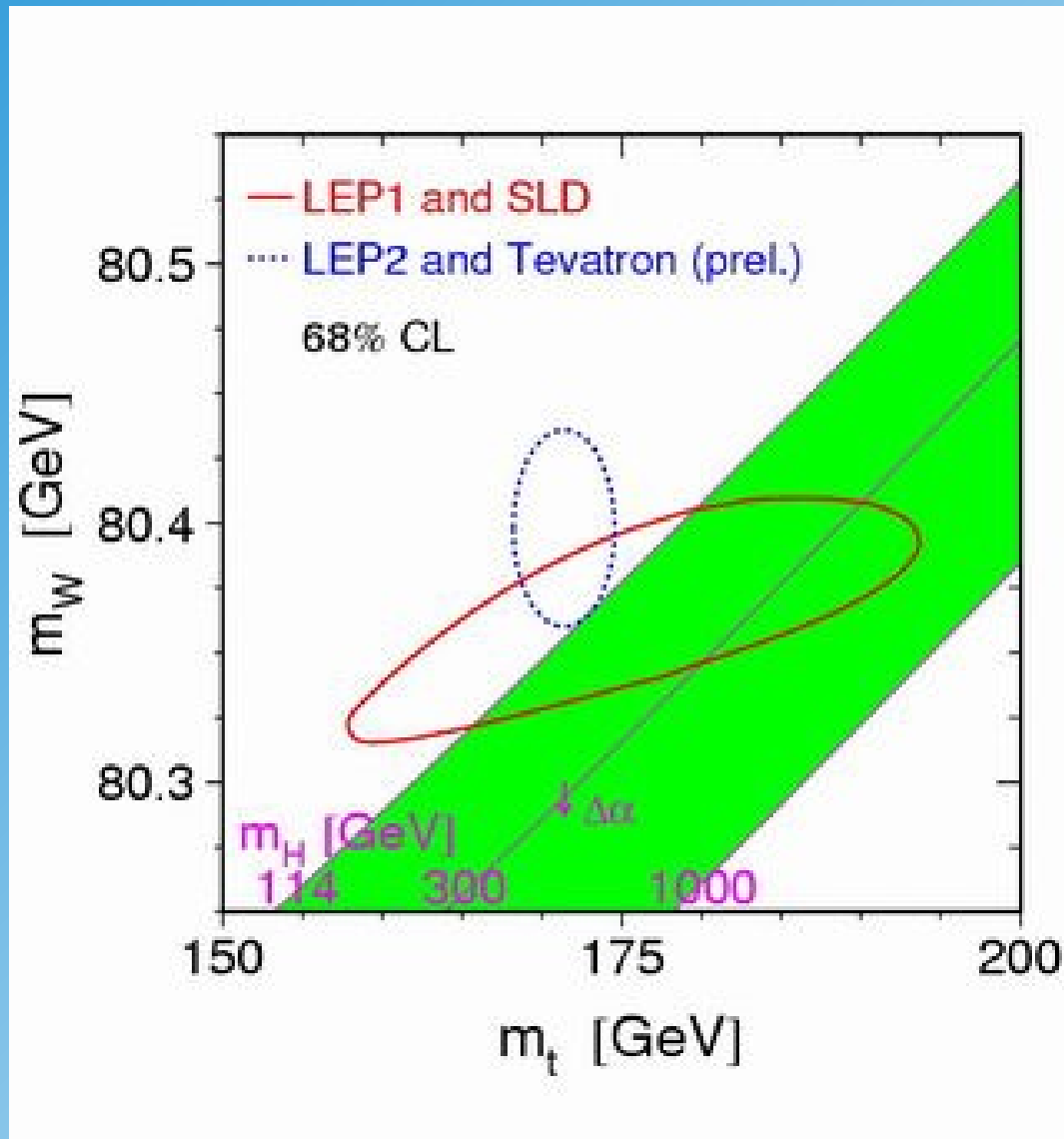
Take the top mass from direct measurements and use the radiative corrections to determine the Higgs mass.

$$\Delta r(m_t, M_H) = -\frac{3\alpha \cos^2 \theta_w}{16\pi \sin^4 \theta_w} \frac{m_t^2}{M_W^2} - \frac{11\alpha}{48\pi \sin^2 \theta_w} \ln \frac{M_H^2}{M_W^2} \dots$$

outdated

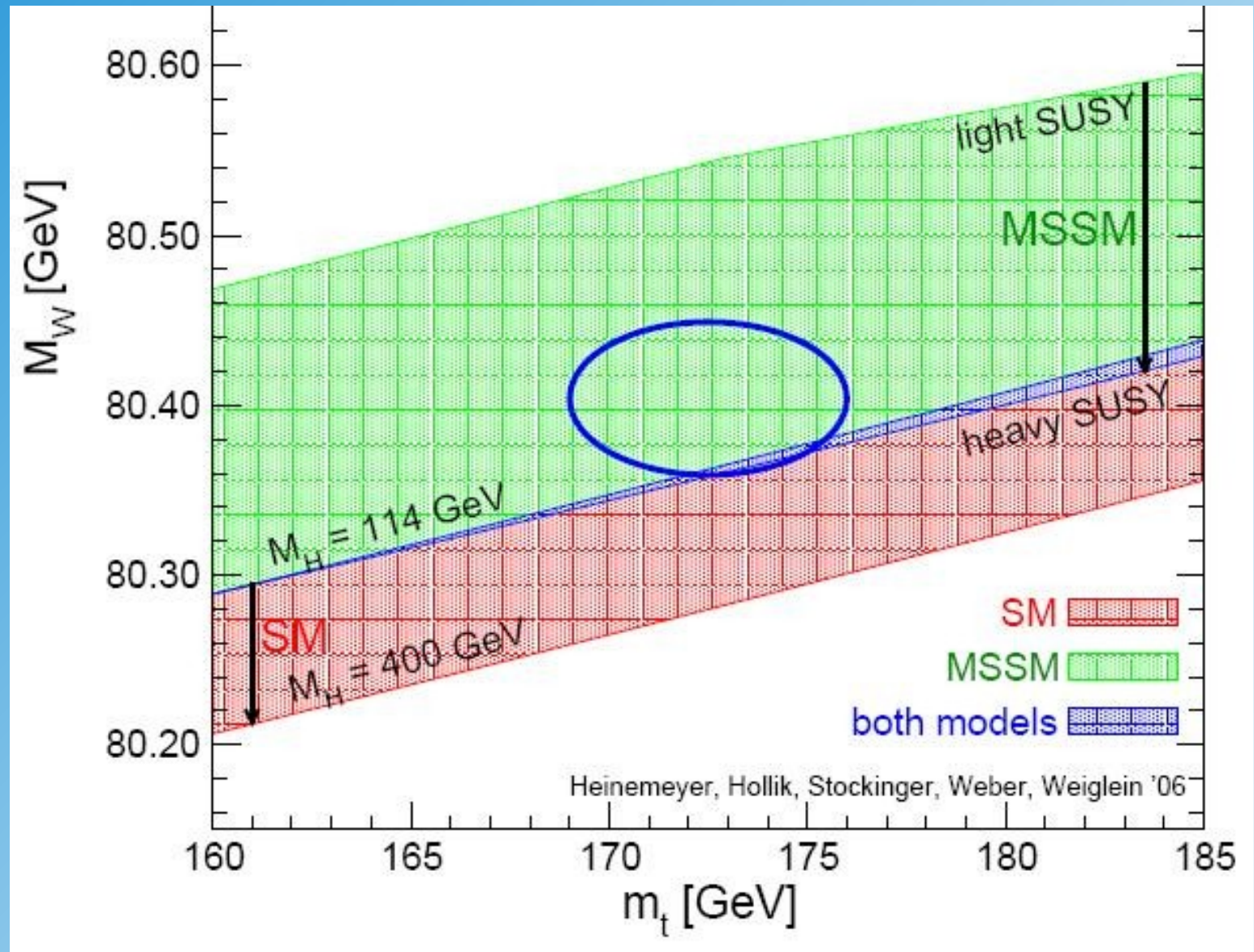
Theoretical prediction of $\sin^2 \theta_{\text{eff}}$ as function of the Higgs mass.

SM Higgs Mass Predictions



W-Top-Higgs Mass Relation and SUSY

SUSY Higgs is light!



Higgs Direct Searches at LEP

Higgs-Fermion Coupling:

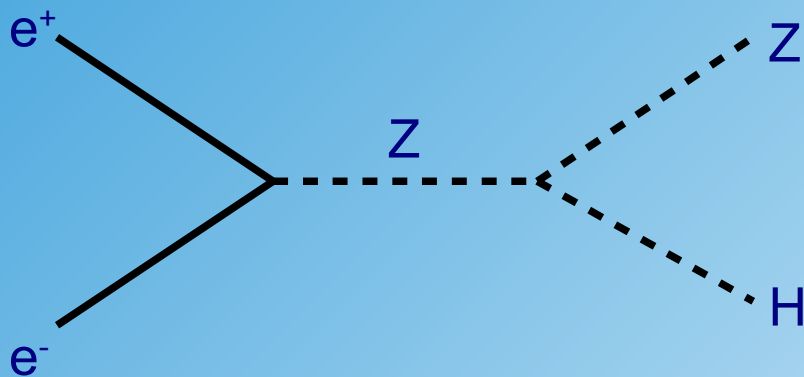
$$L_Y = -g_b \bar{L} \Phi b_R - g_t \bar{L} \tilde{\Phi} t_R \quad \text{with} \quad L = \begin{pmatrix} t \\ b \end{pmatrix}, \quad \Phi = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v+h \end{pmatrix}, \quad \tilde{\Phi} = i\tau_2 \Phi^*$$

Higgs couples to masses:

$$m_{b,t} = \frac{g_{b,t} v}{\sqrt{2}}$$

Electron-Positron annihilation has tiny Higgs coupling!

LEP Process (Higgs-Strahlung): $e^+ e^- \rightarrow ZH$



ZZH coupling is large!

ZH Signature at LEP

Higgs decays dominantly into heaviest fermions:

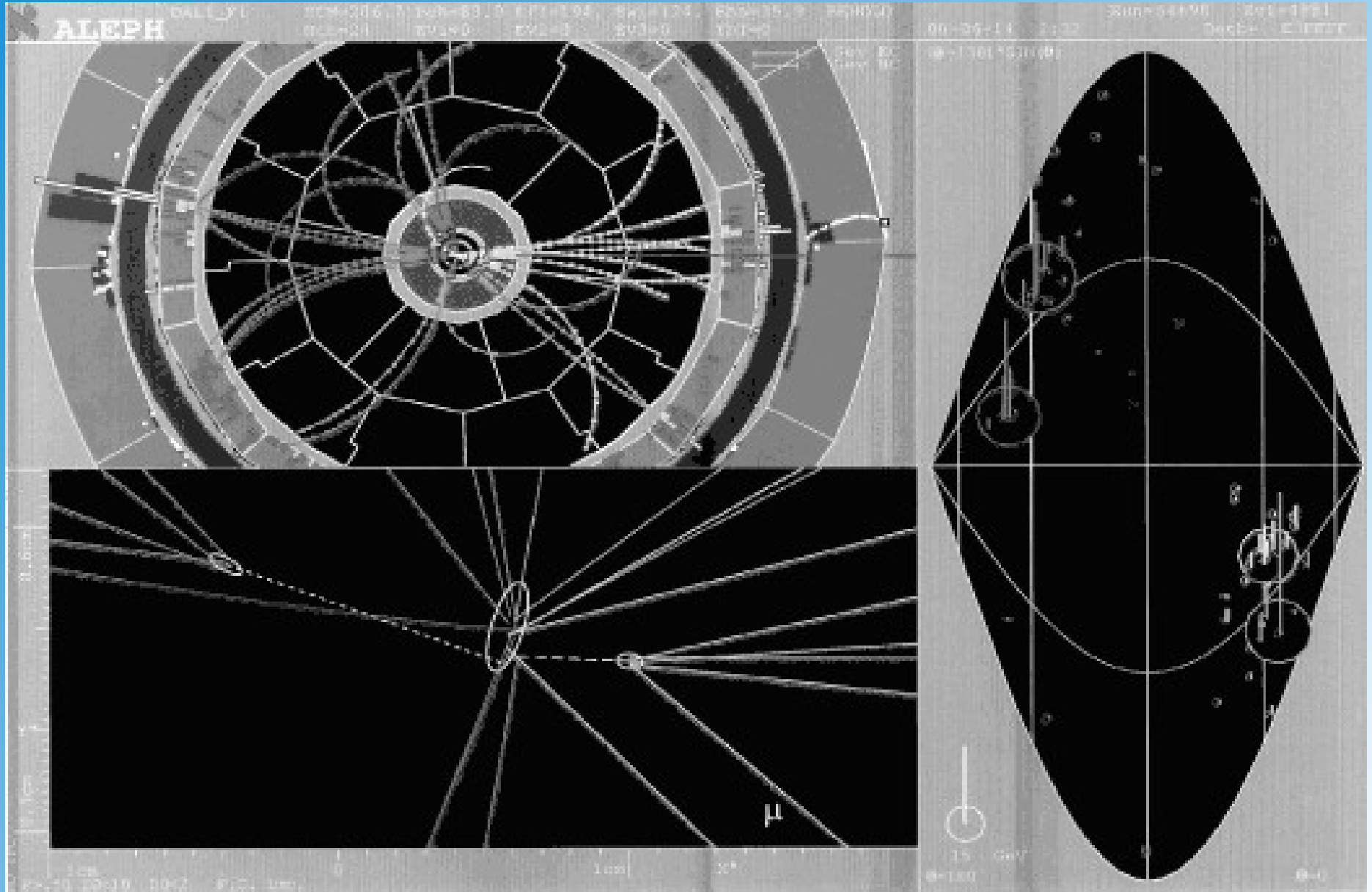
$$ZH \rightarrow Z bb \rightarrow ll bb, jj bb, bbbb, \nu\nu bb$$

All decay channels require (double) b-tag (lifetime)

LEP2 with $E_{\text{cms}} = 205 \text{ GeV}$:

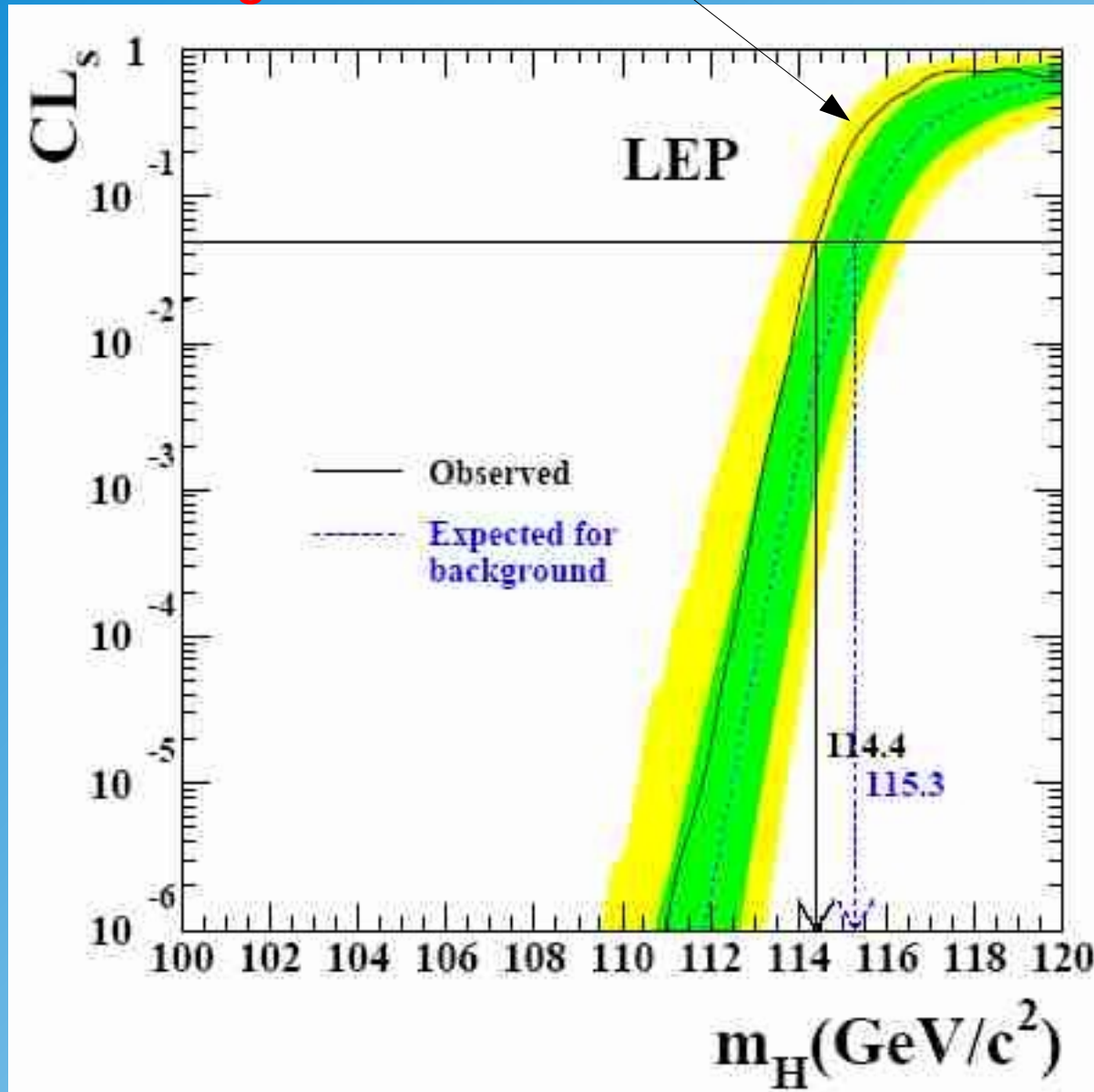
- installation of vertex detectors
- sensitivity up to $m_H = 114 \text{ GeV}$ ($E_{\text{cms}} = m_H + m_Z$)

bbj-candidate at ALEPH



Combined LEP2 Higgs Limit

no significant excess



Higgs excluded with:

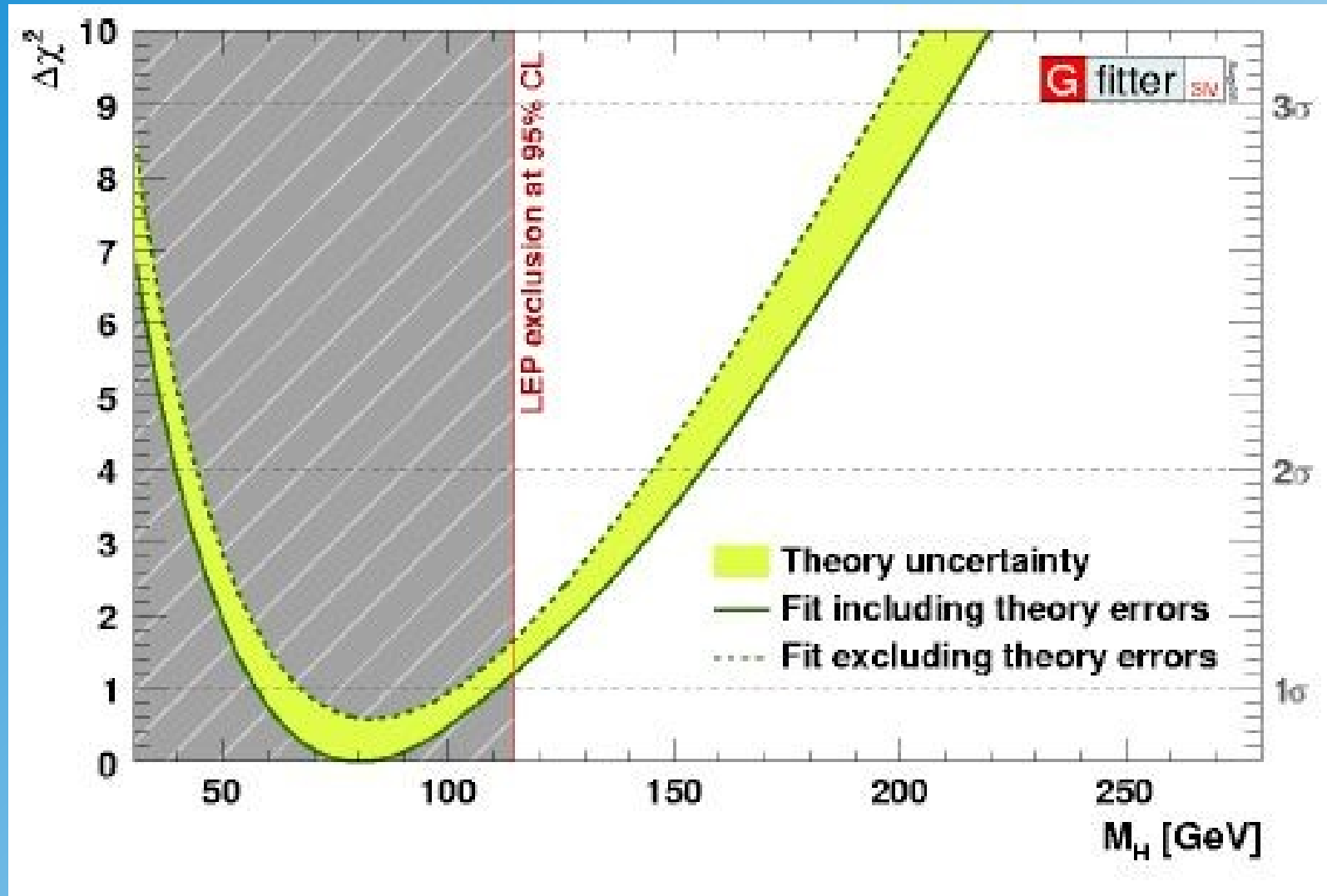
$$m_H < 114.4 \text{ GeV}$$

(expected 115.3)

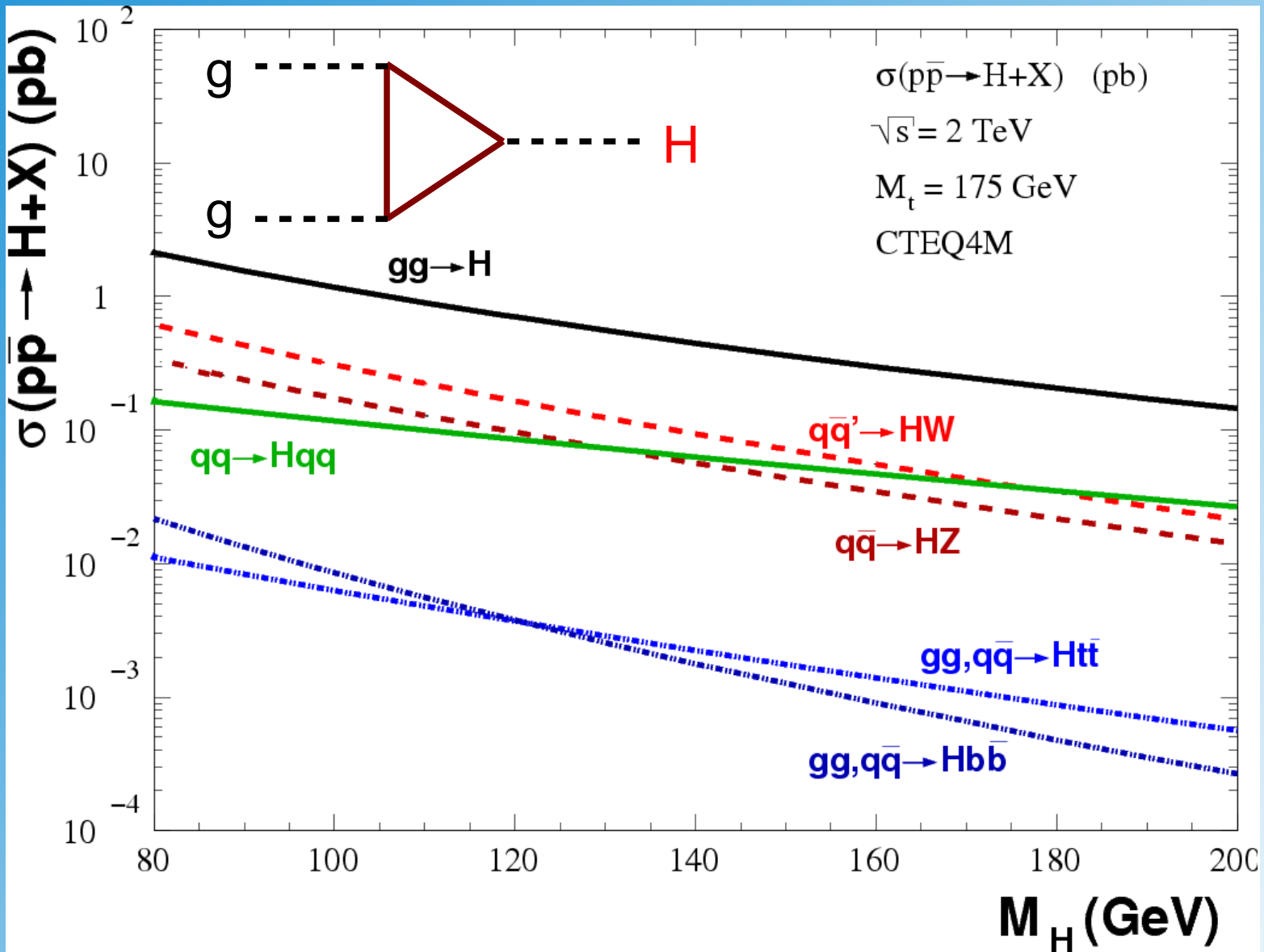
at 95% CL

The bands shows the 1 sigma and 2 sigma contours of the expected limit

Higgs LEP2 Direct Limits

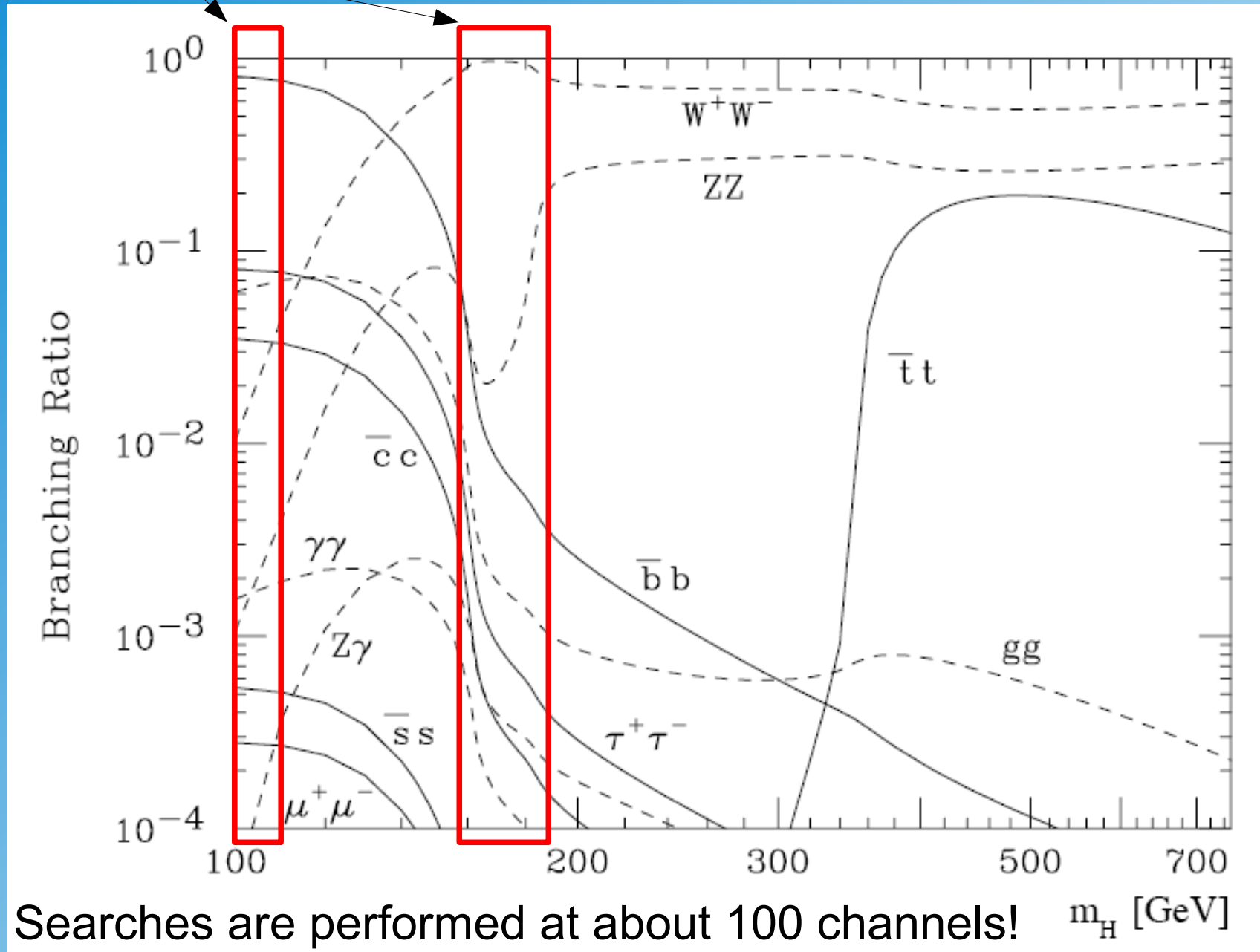


Higgs Production at Tevatron

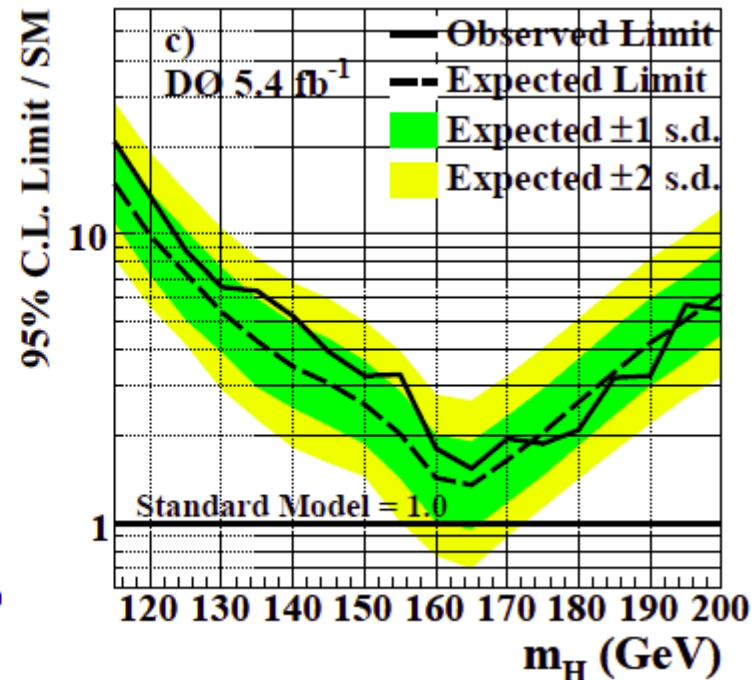
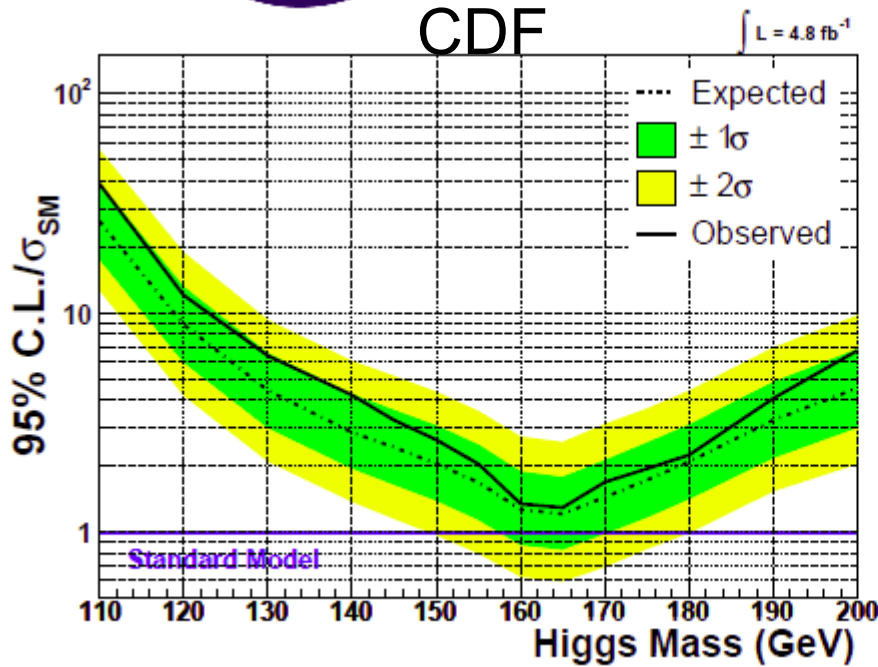
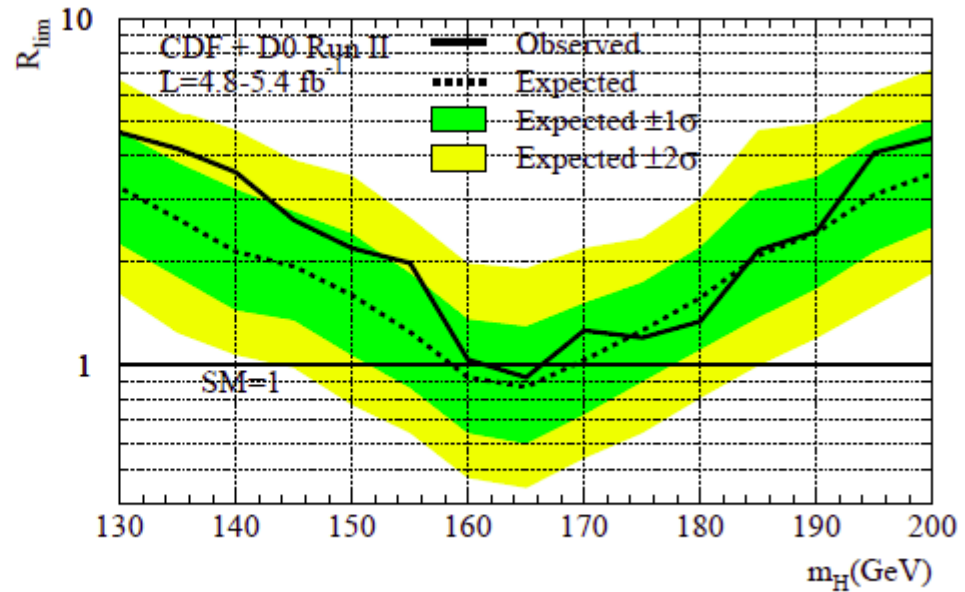
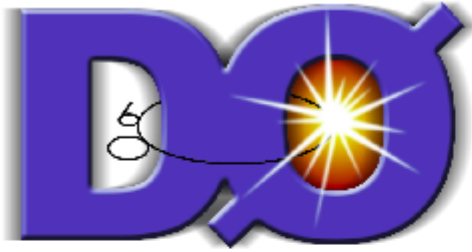


Higgs Decays

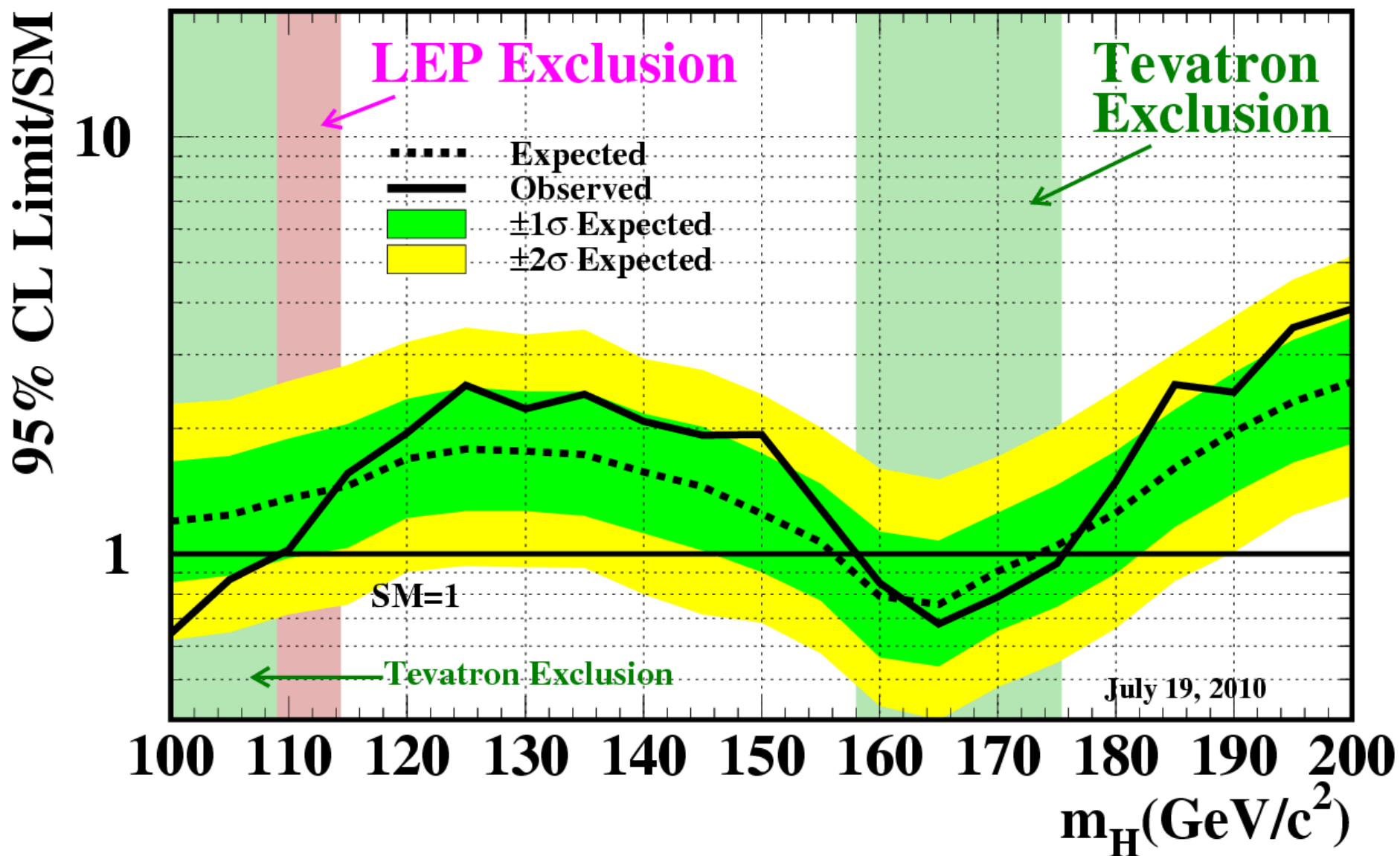
“easy” regions



Combination Tevatron Searches



Tevatron Run II Preliminary, $\langle L \rangle = 5.9 \text{ fb}^{-1}$



Situation before LHC

Fits to electro-weak data:

$$m_H = 89^{+35}_{-26} \text{ GeV}$$

$$m_H < 158 \text{ GeV (95\% CL)}$$

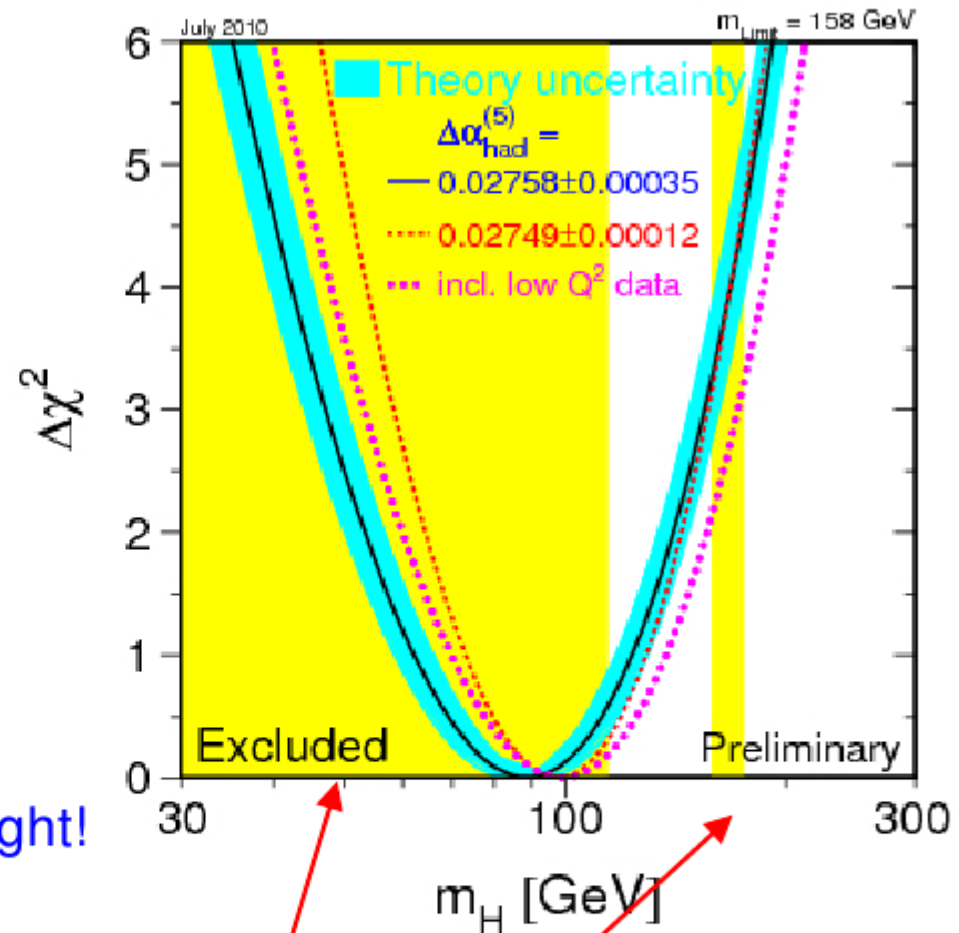
Assumption for fit:

- SM including Higgs

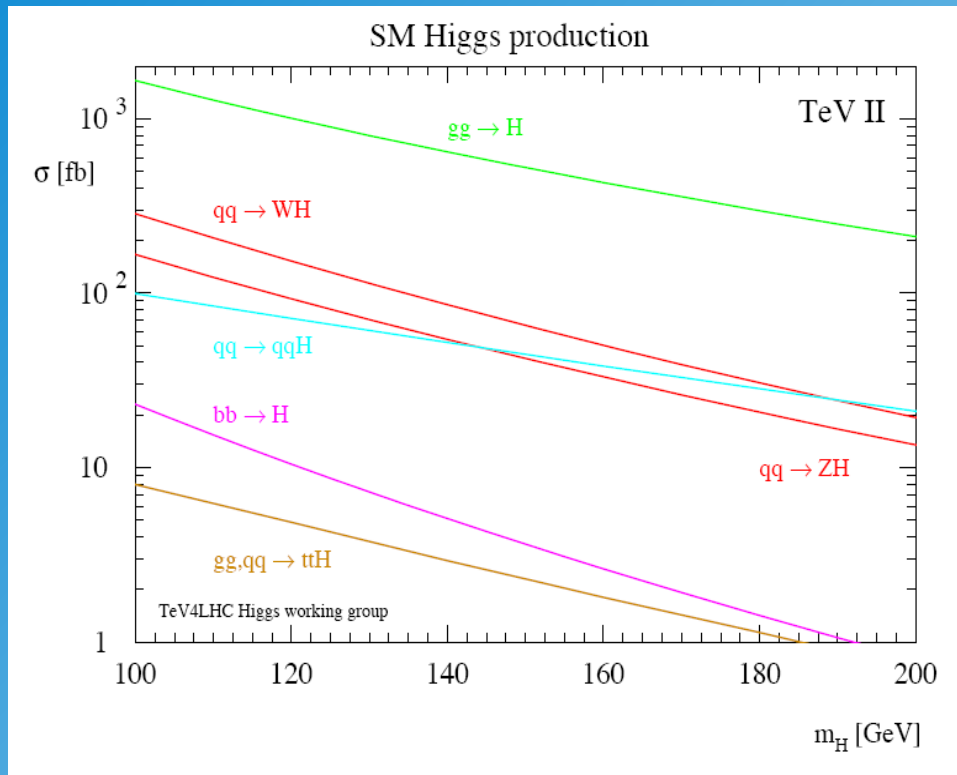
If existing, Higgs seems to be light!

- Direct searches at LEP:
 $m_H > 114.4 \text{ GeV @ 95\% CL}$
- Direct searches at Tevatron

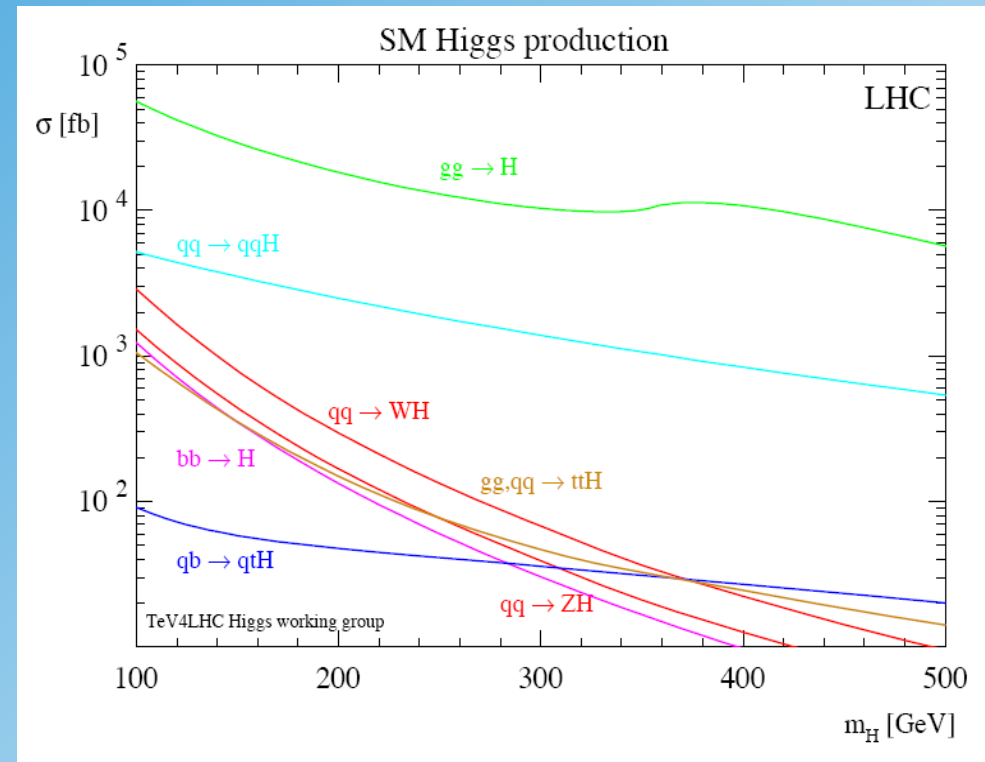
<http://lepewwg.web.cern.ch/LEPEWWG/>



Higgs Production at LHC



proton-antiproton at Tevatron $s^{1/2}=2$ TeV

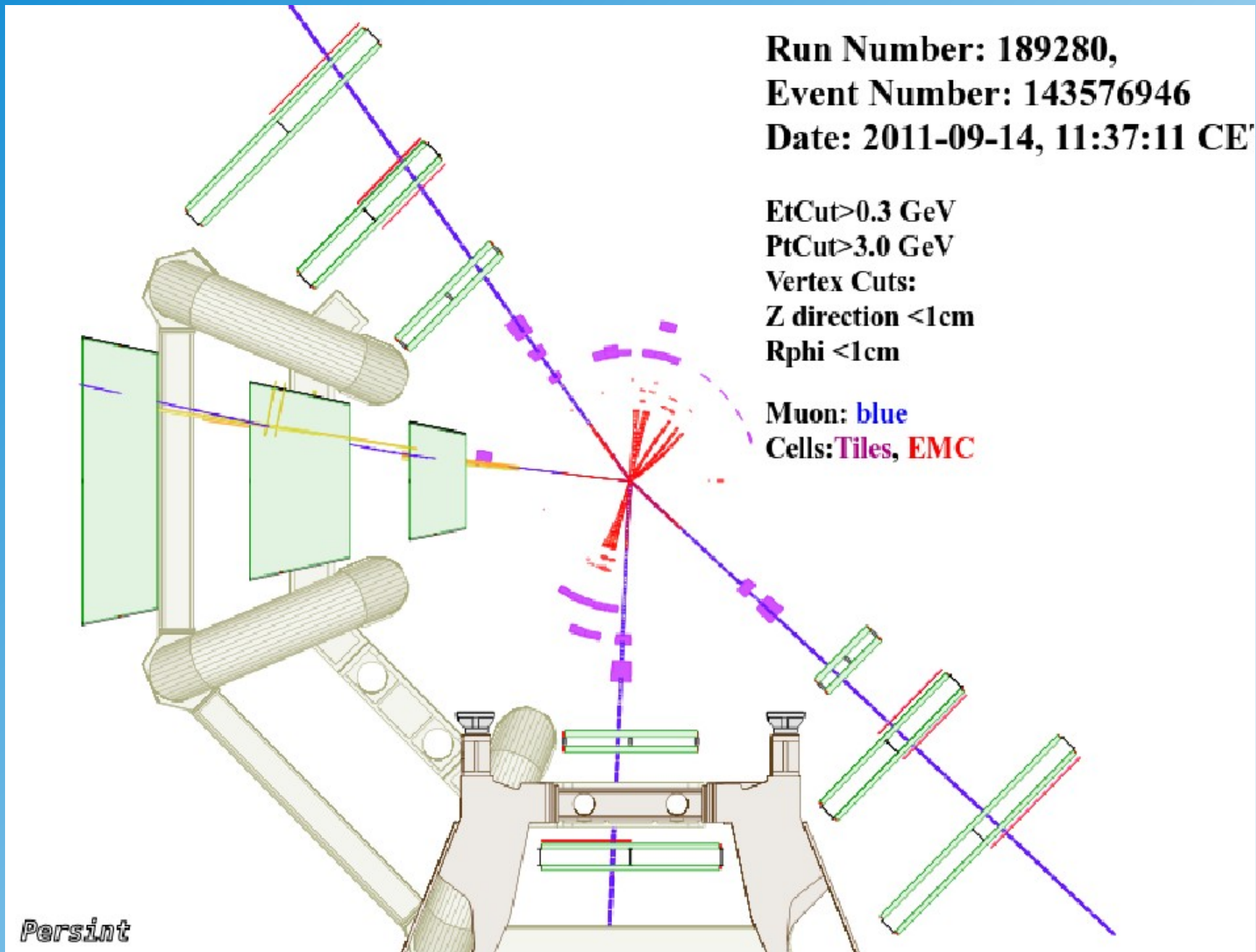


proton-proton at LHC $s^{1/2}=14$ TeV

much larger cross sections!

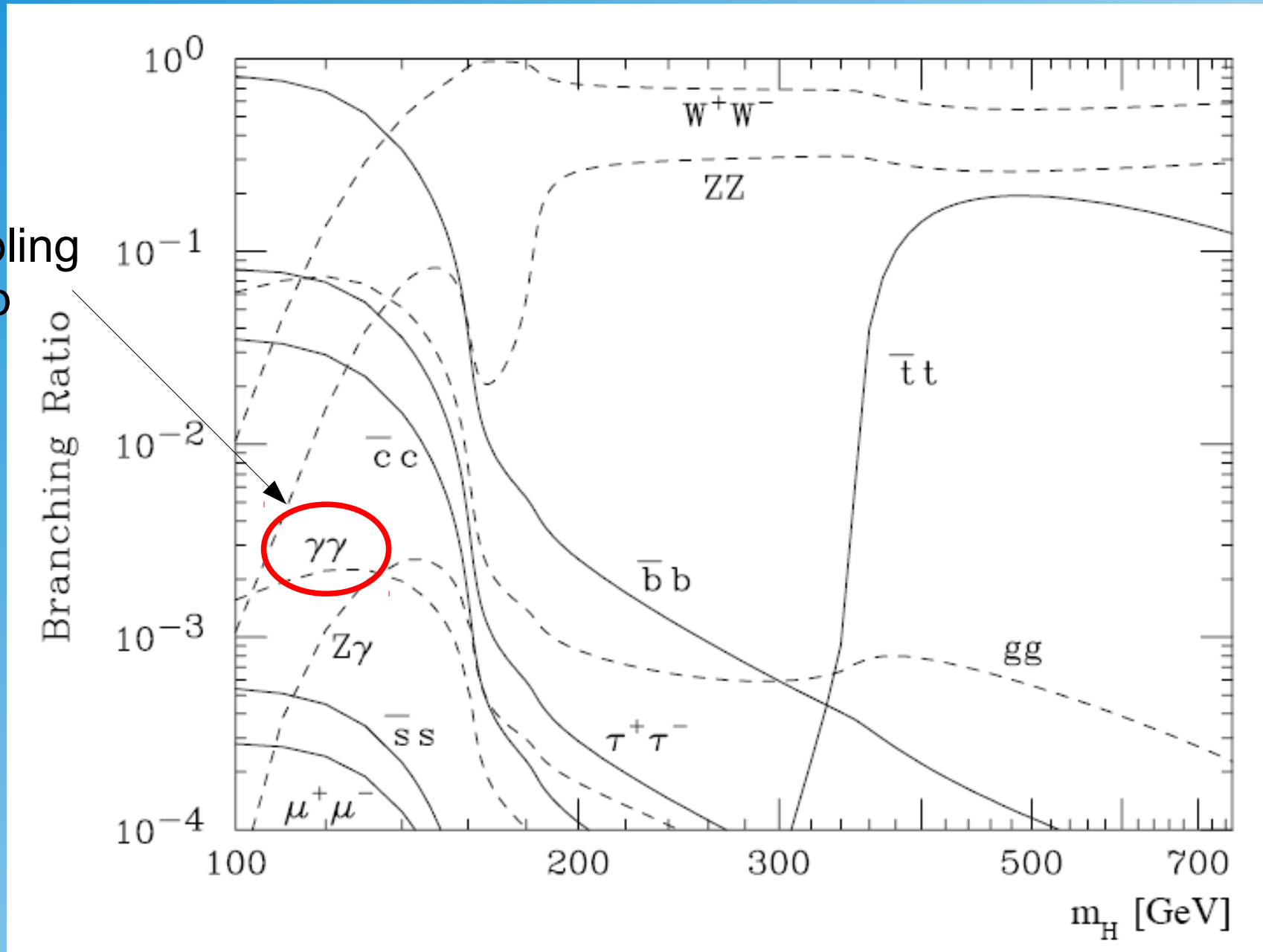
also higher luminosity!

Candidate $H \rightarrow ZZ \rightarrow \mu\mu\mu\mu$

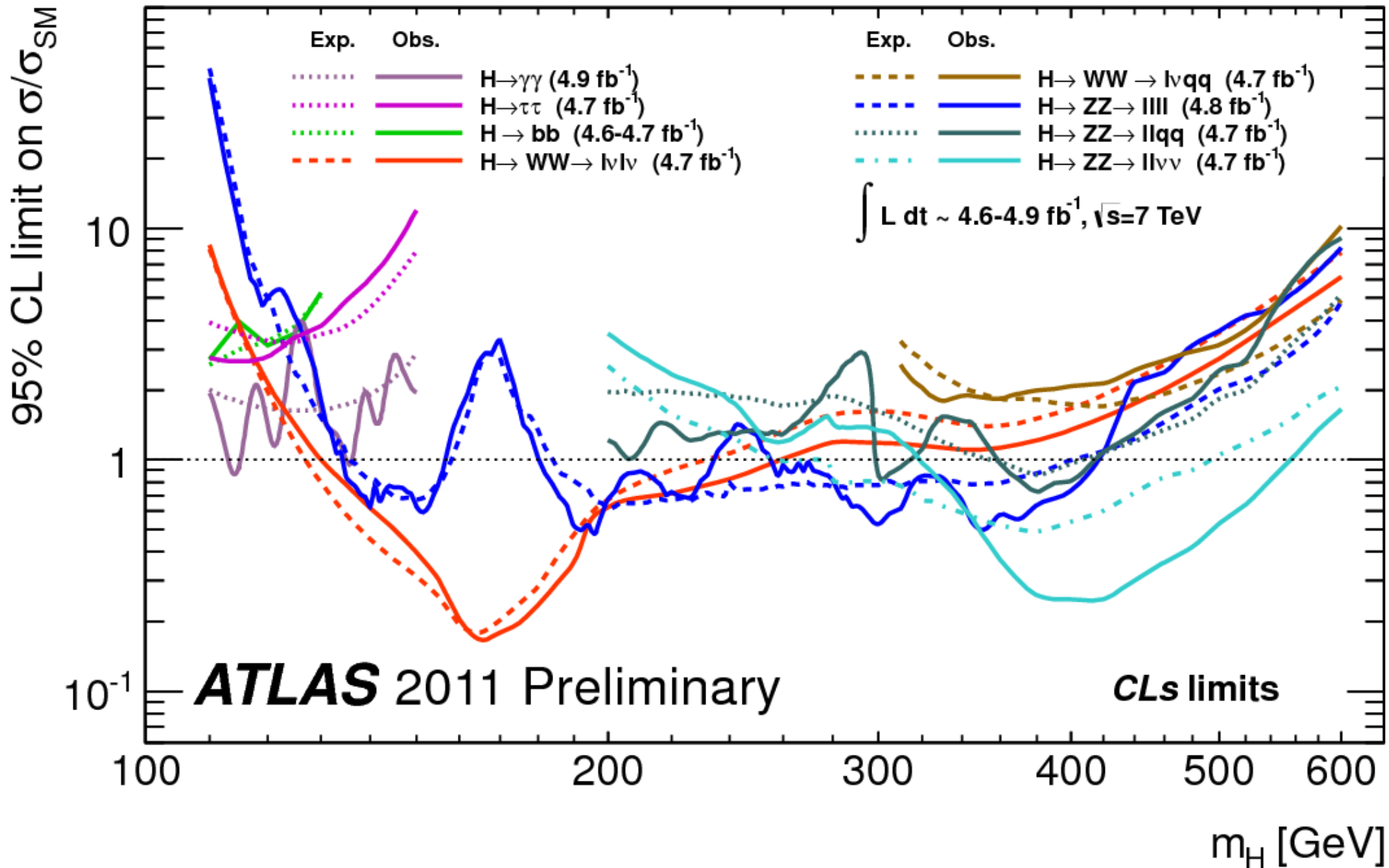


Special Importance $H \rightarrow \gamma\gamma$

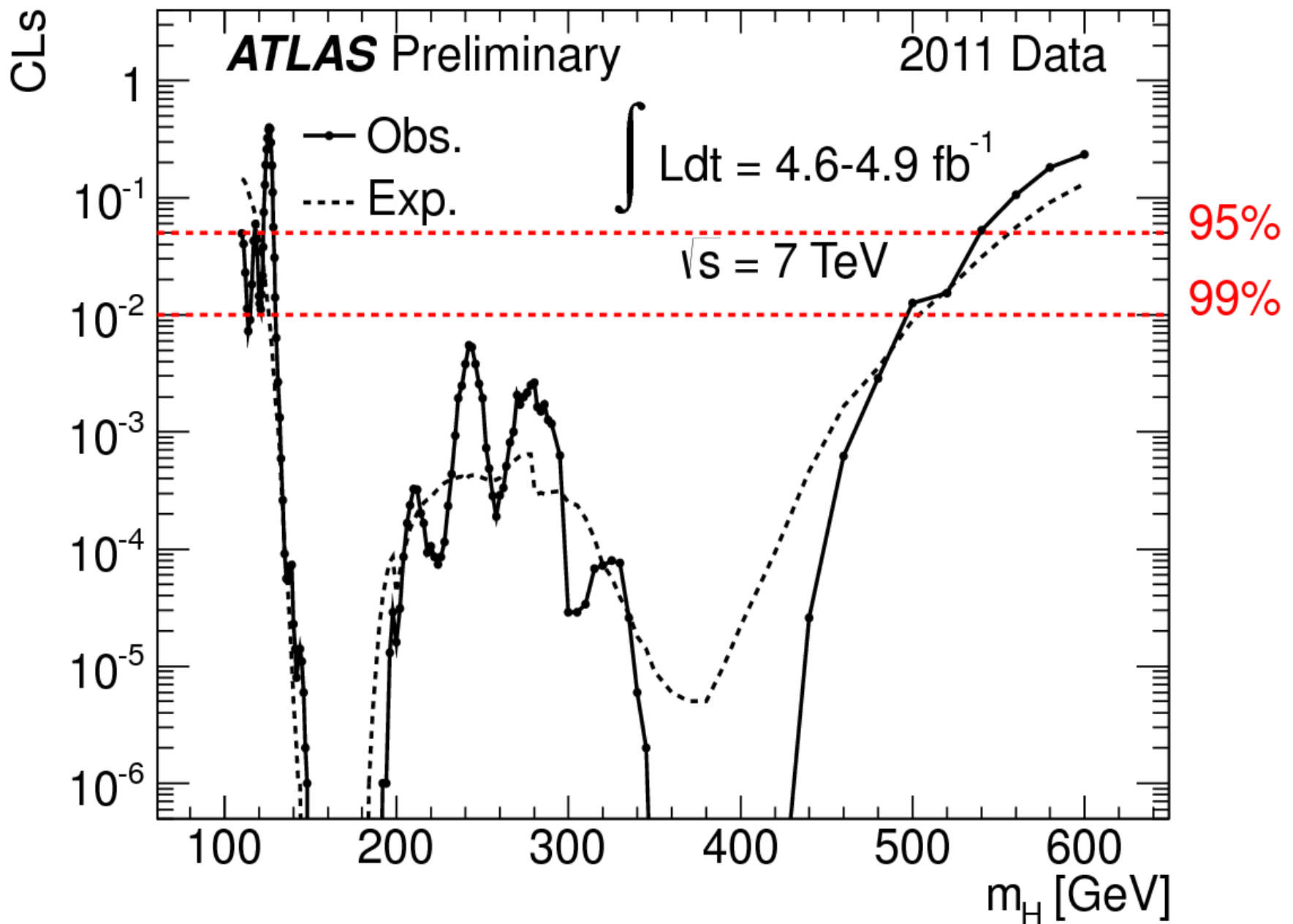
Higgs coupling
via top loop



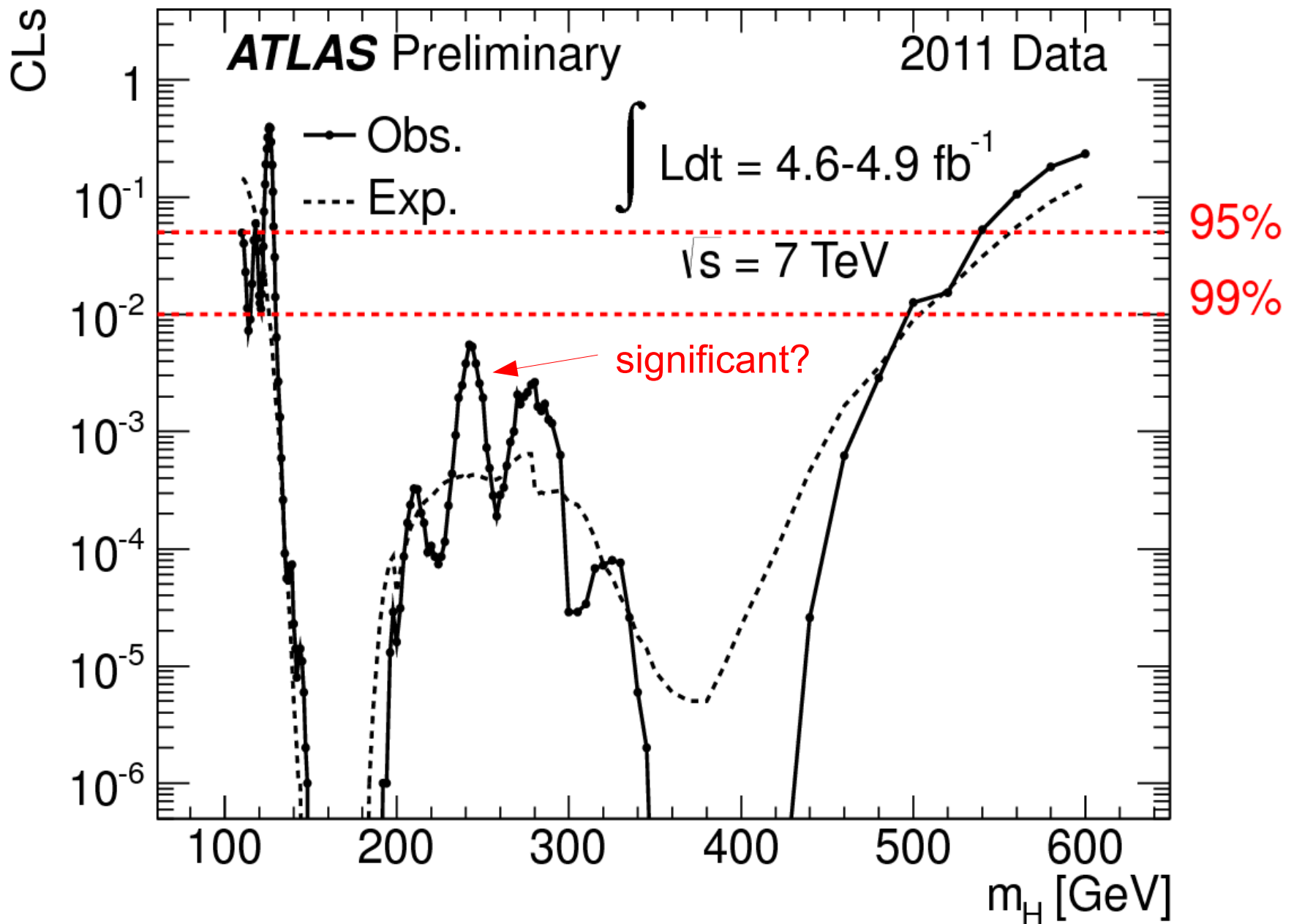
Summary ATLAS Searches



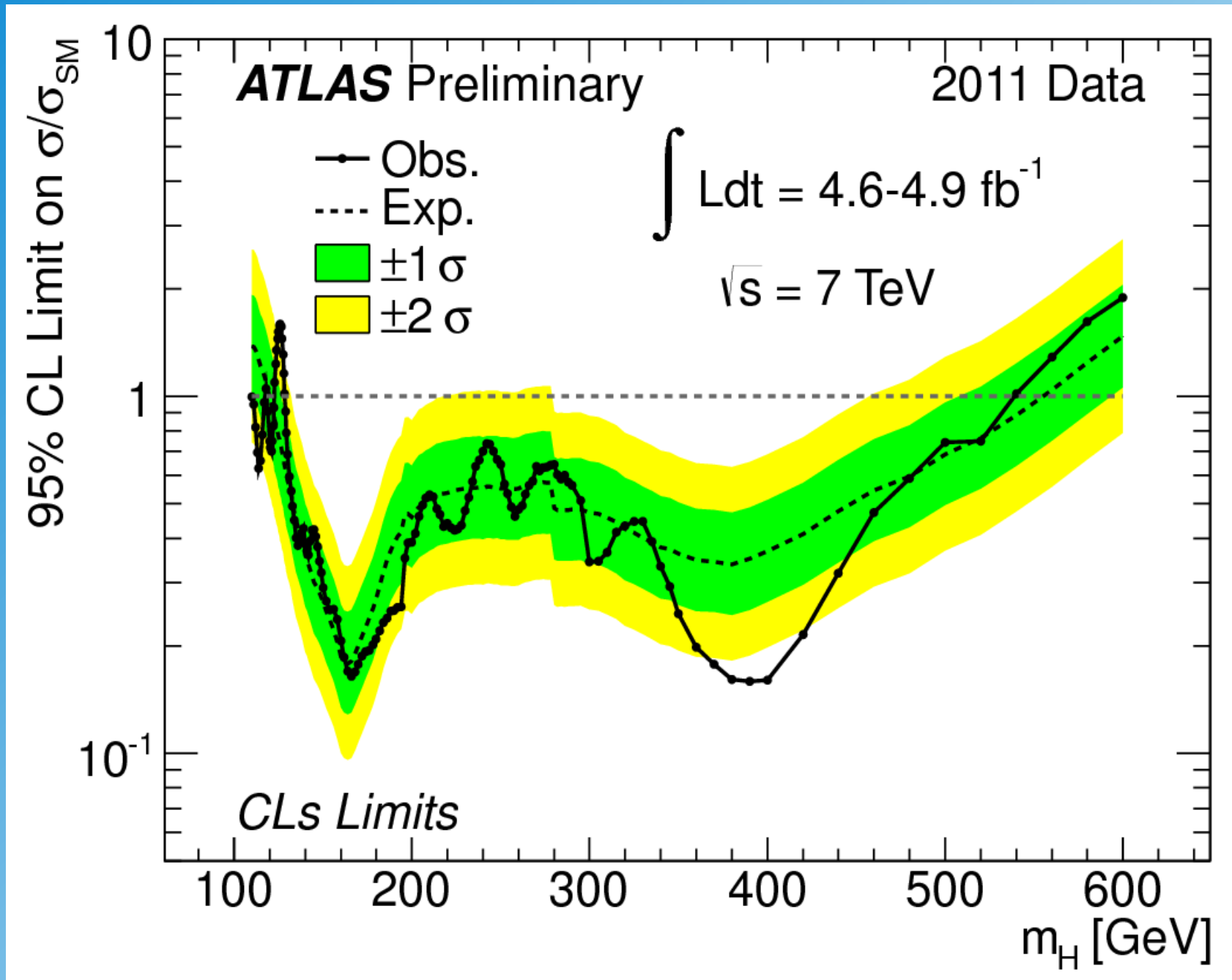
ATLAS Combined Result



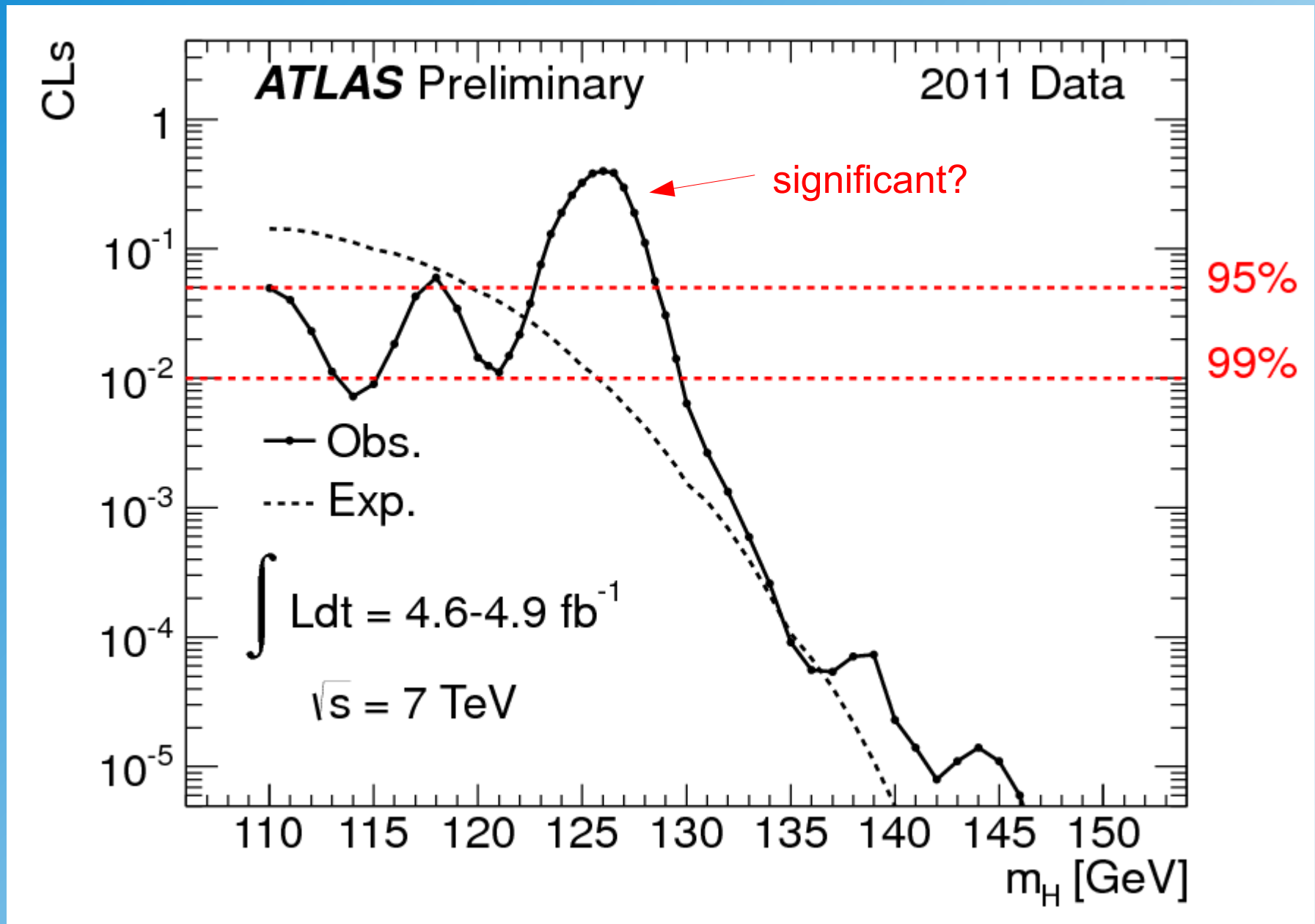
ATLAS Combined Result



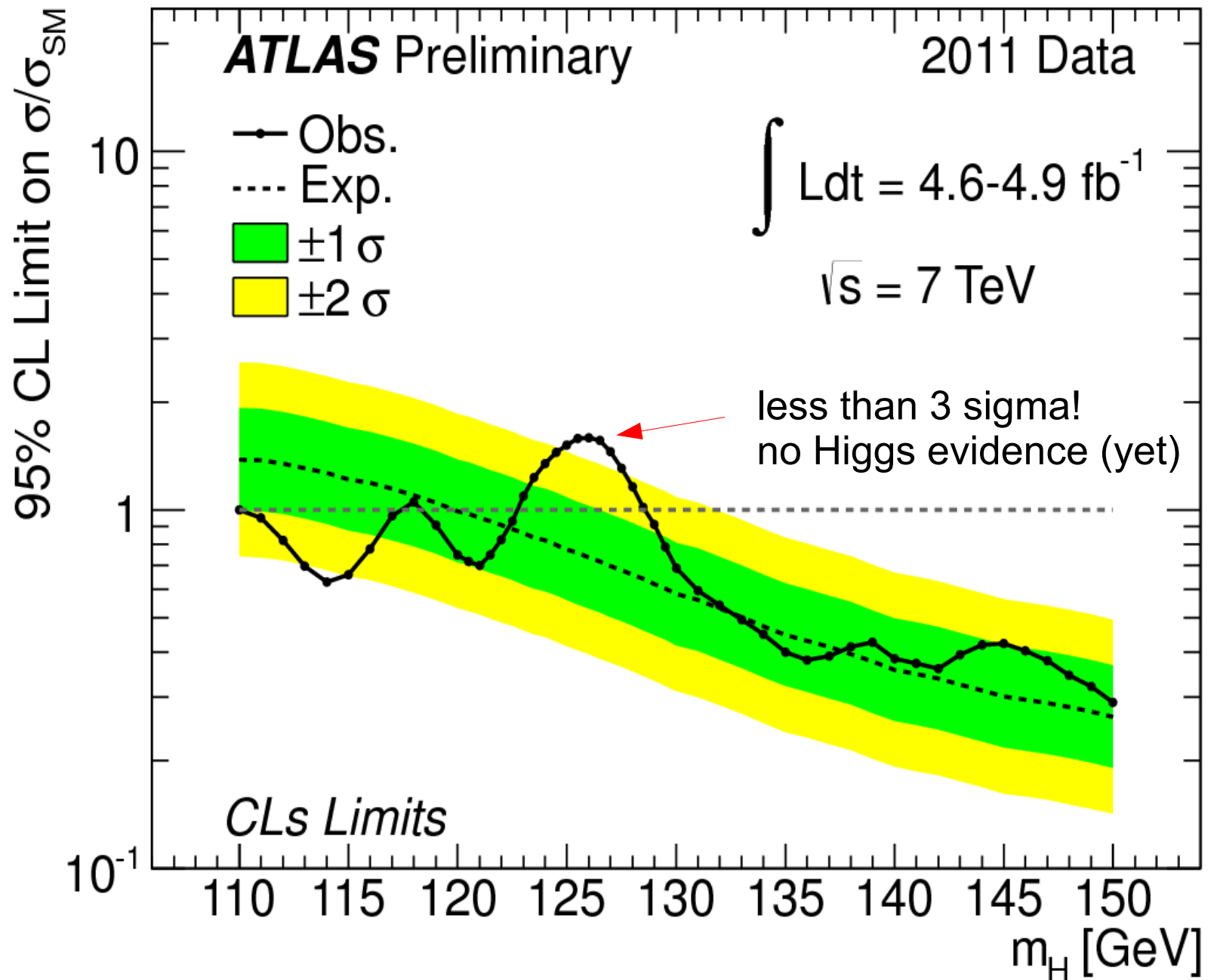
ATLAS Combined Result



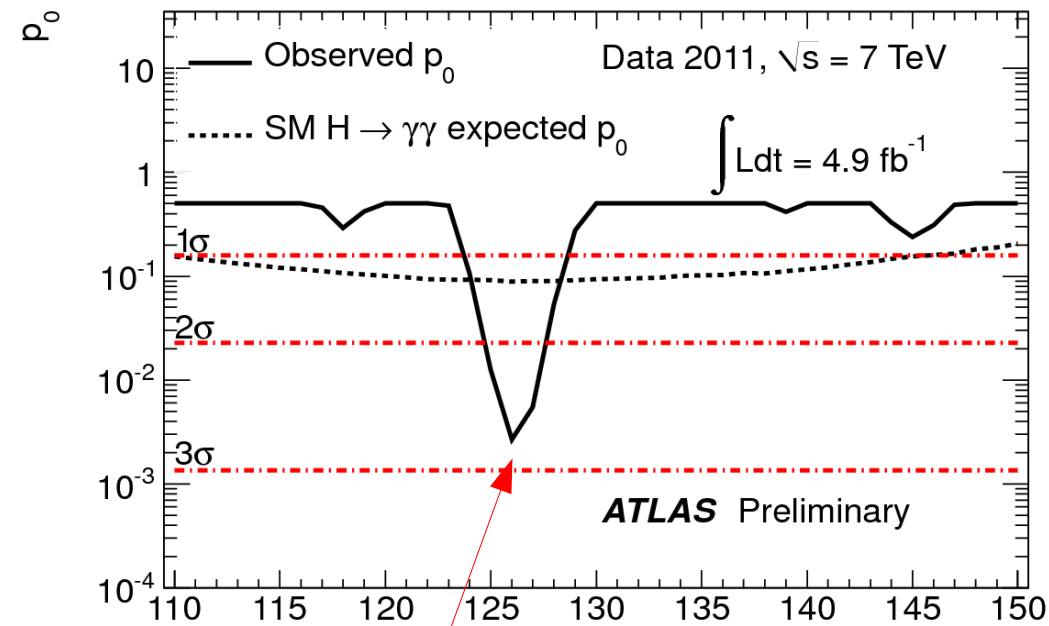
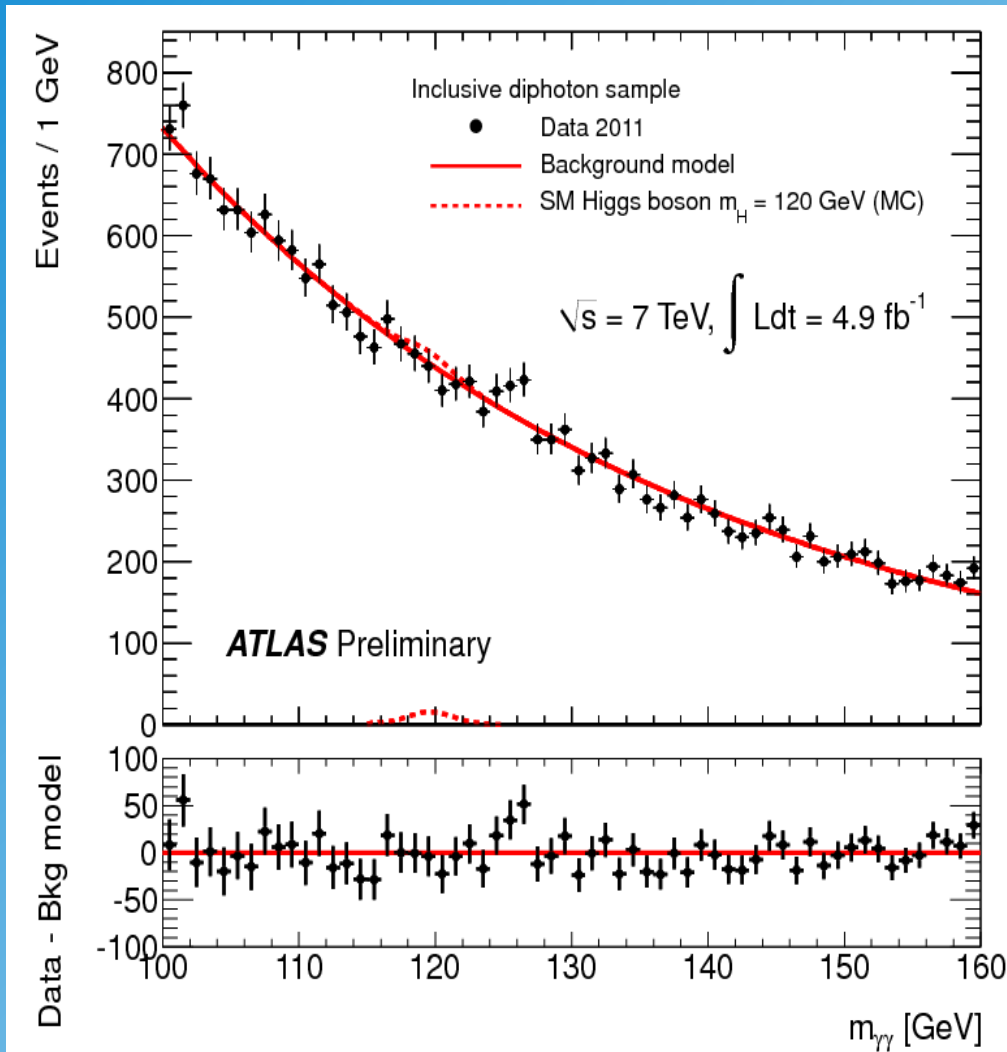
ATLAS Combined Result



ATLAS Combined Result

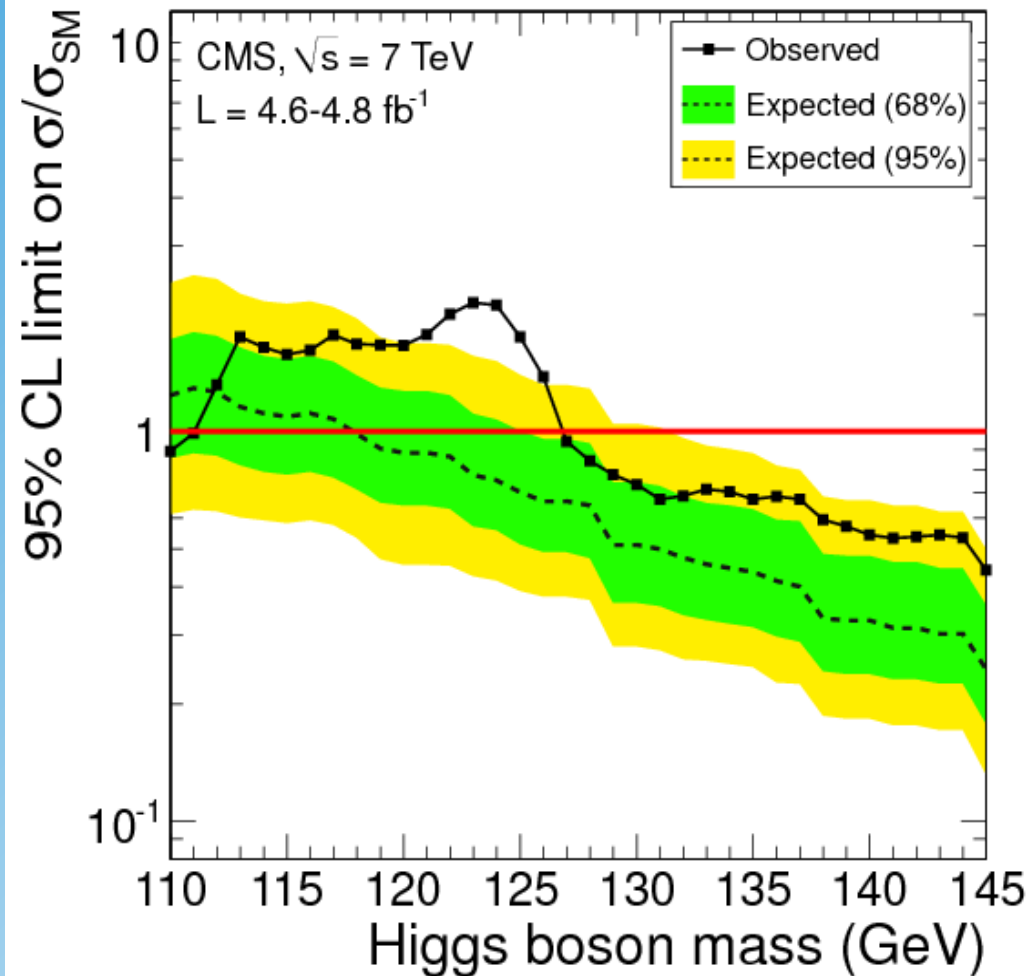
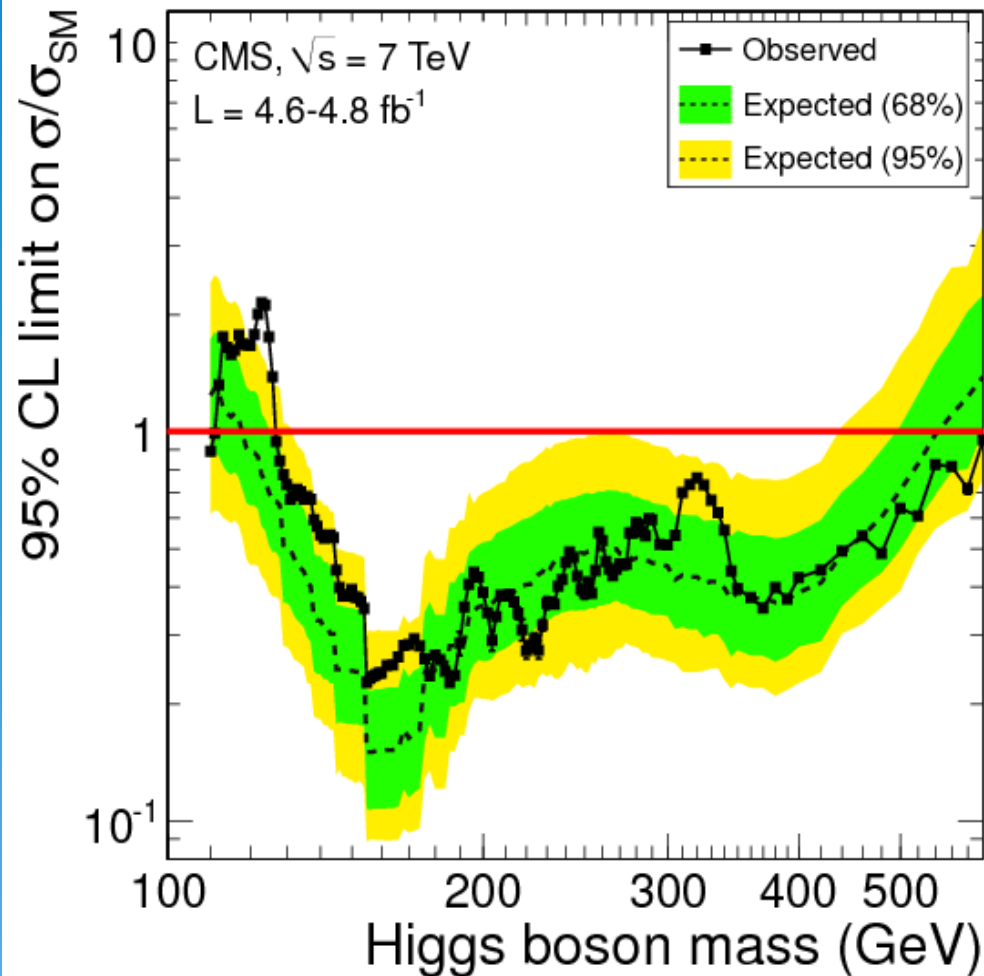


Higgs $\rightarrow \gamma\gamma$ Search



note, signal too strong for Higgs!

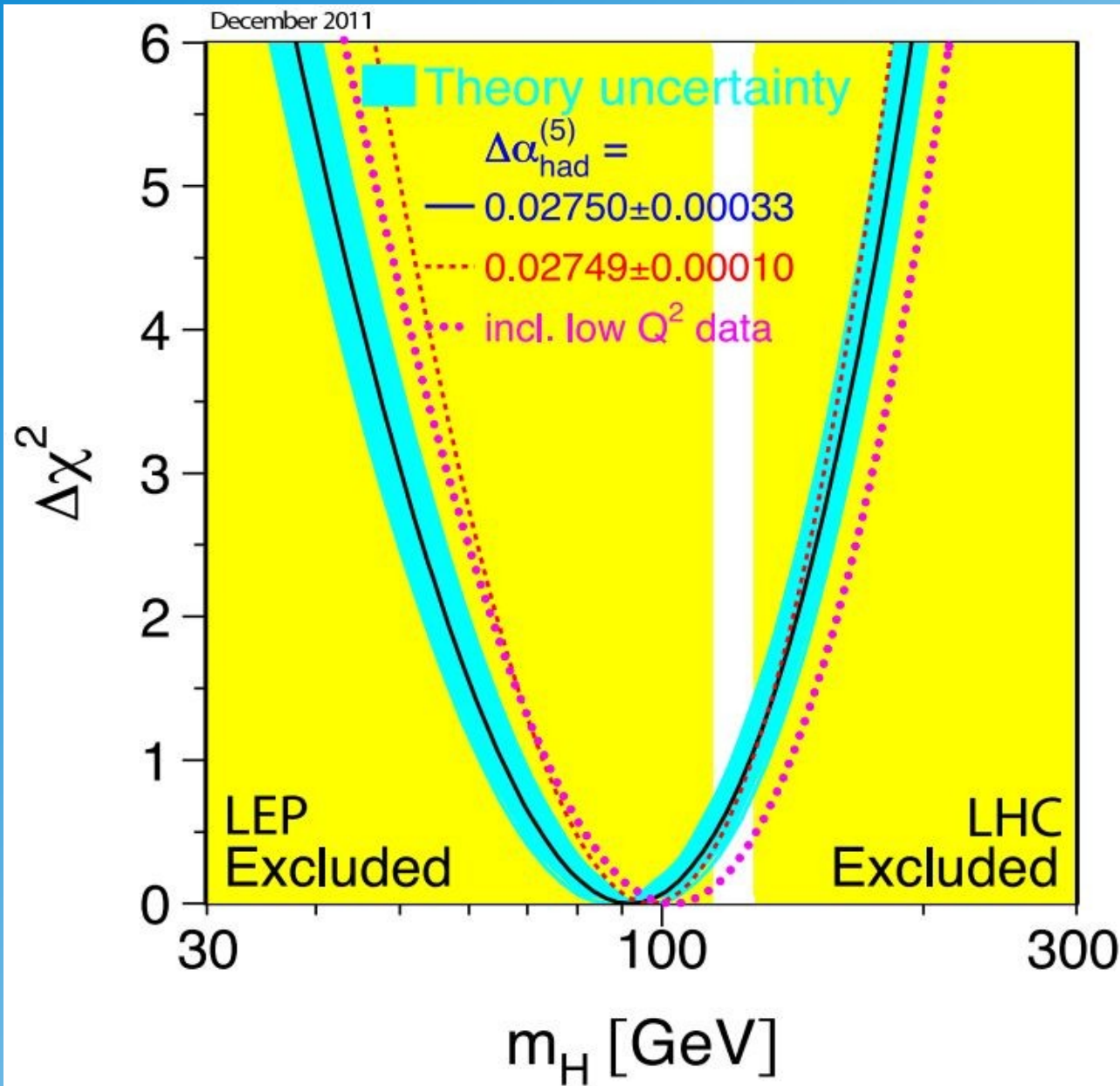
CMS Higgs Search (2011 data)



very consistent with ATLAS results, Higgs?

Situation in June 2012

(before summer conferences)

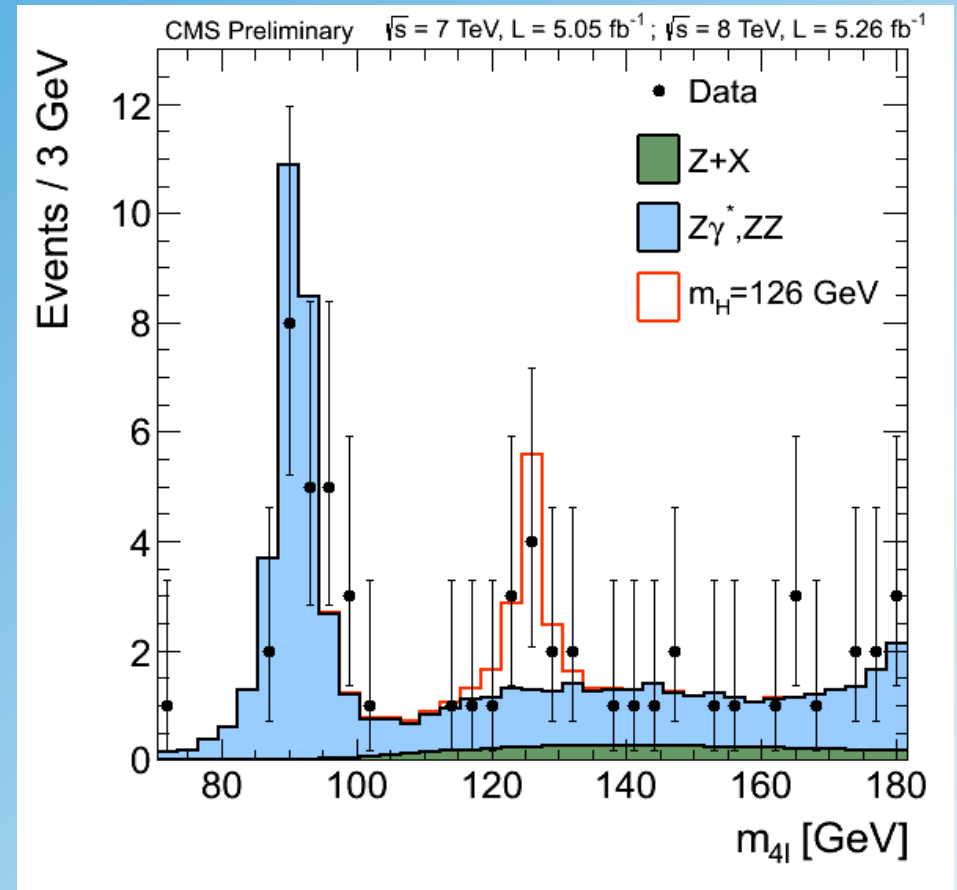
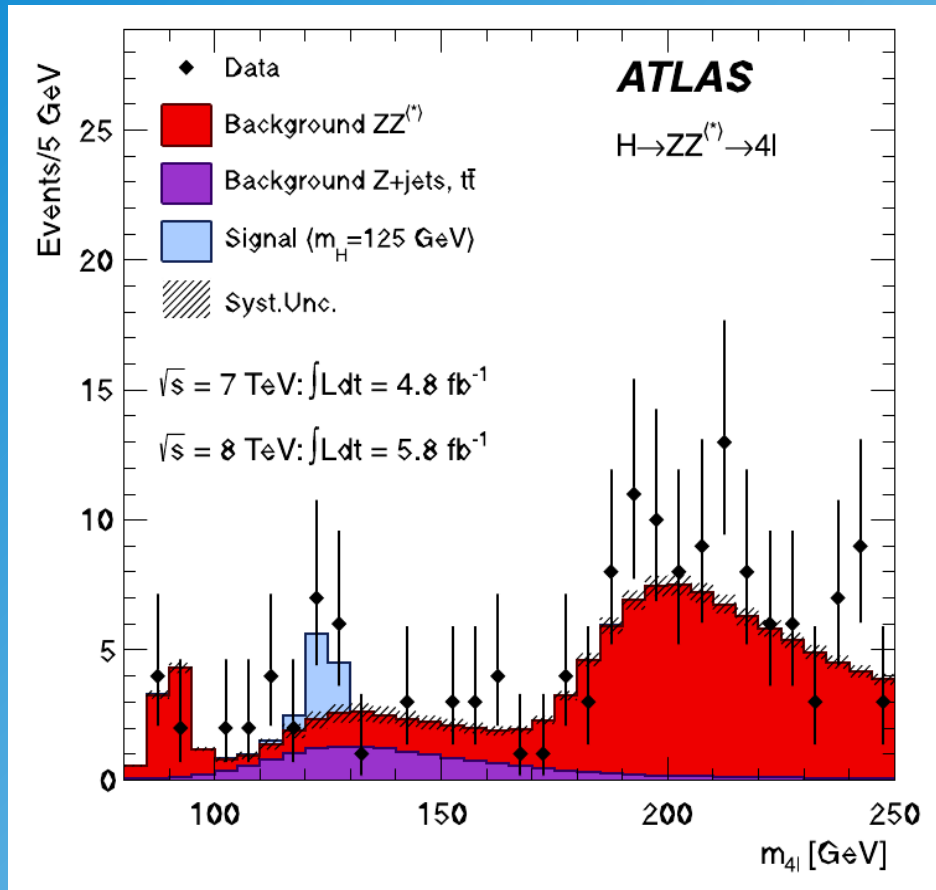


Non-excluded region
is the most difficult!

small window around
125 GeV not excluded
at 95% CL

→ Updates July 4th 2012!

Latest Results ATLAS+CMS

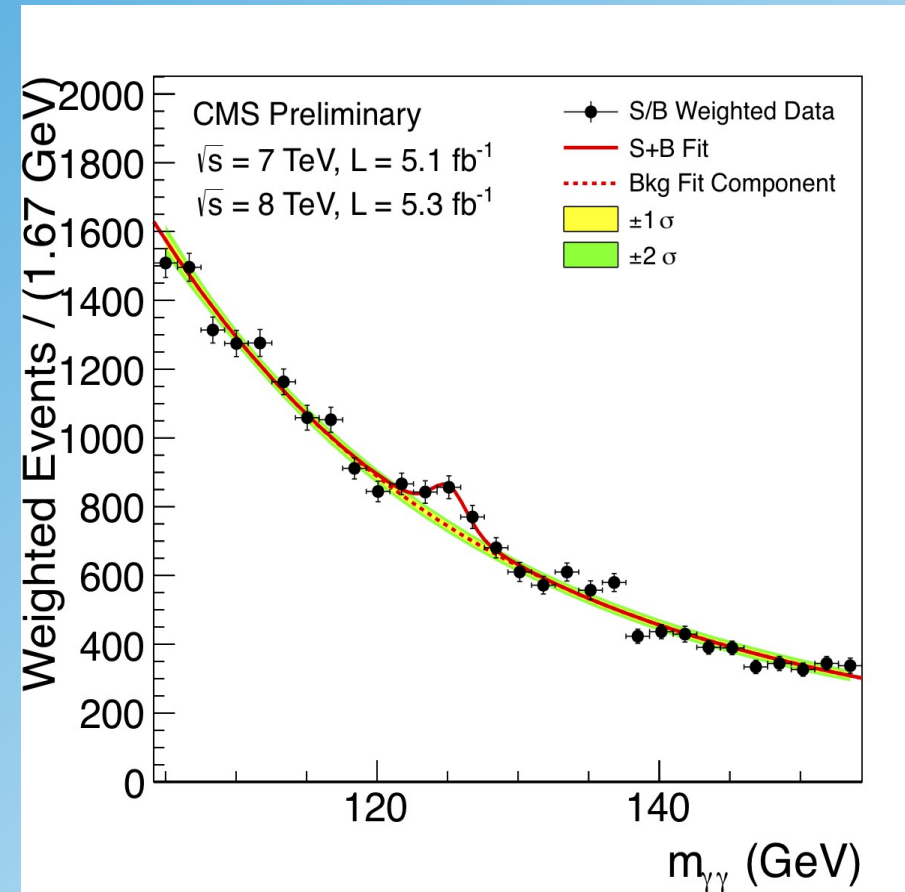
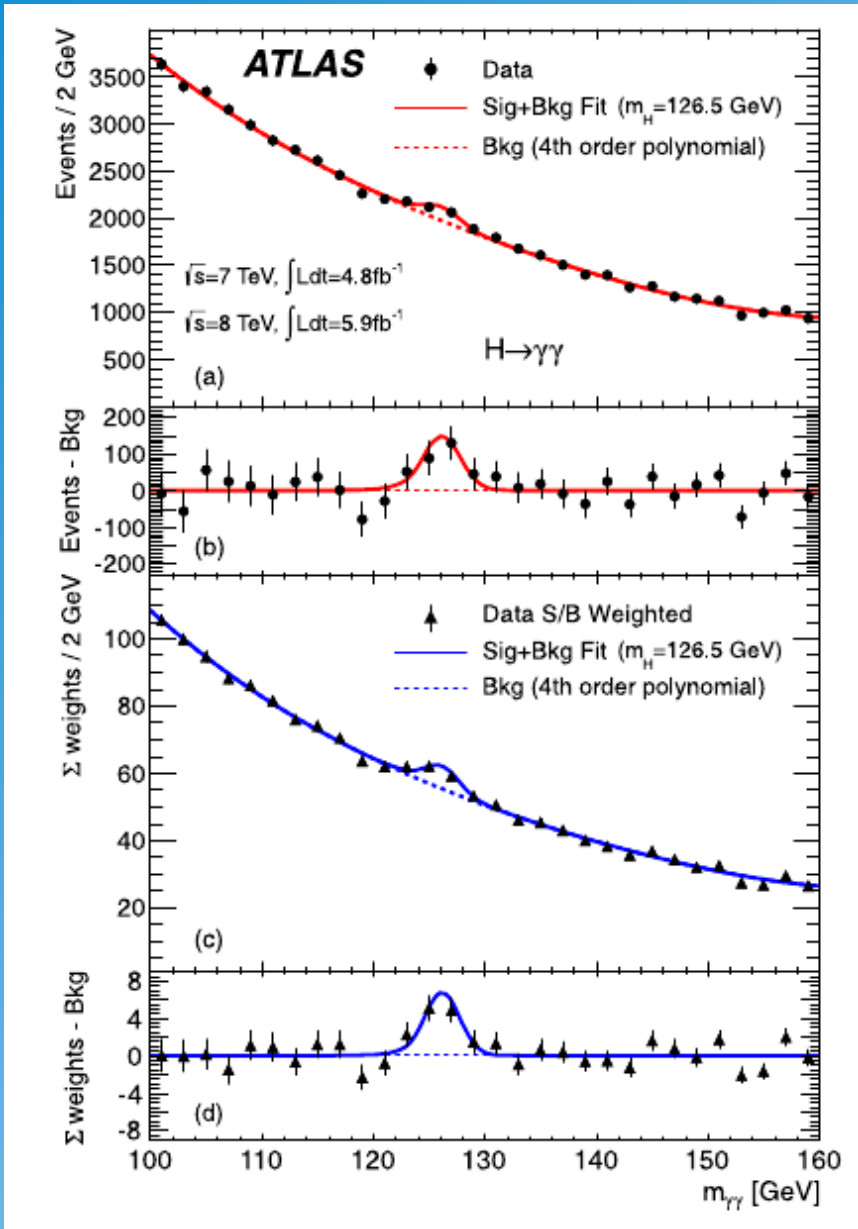


$H \rightarrow ZZ \rightarrow$ leptons

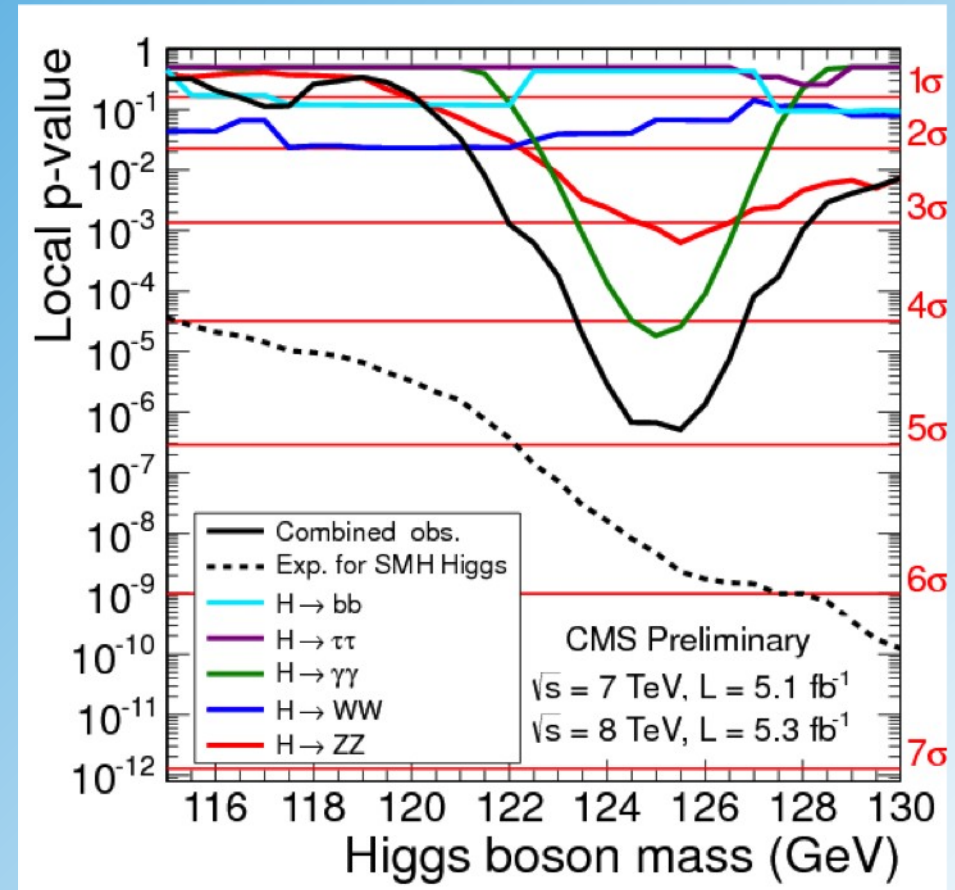
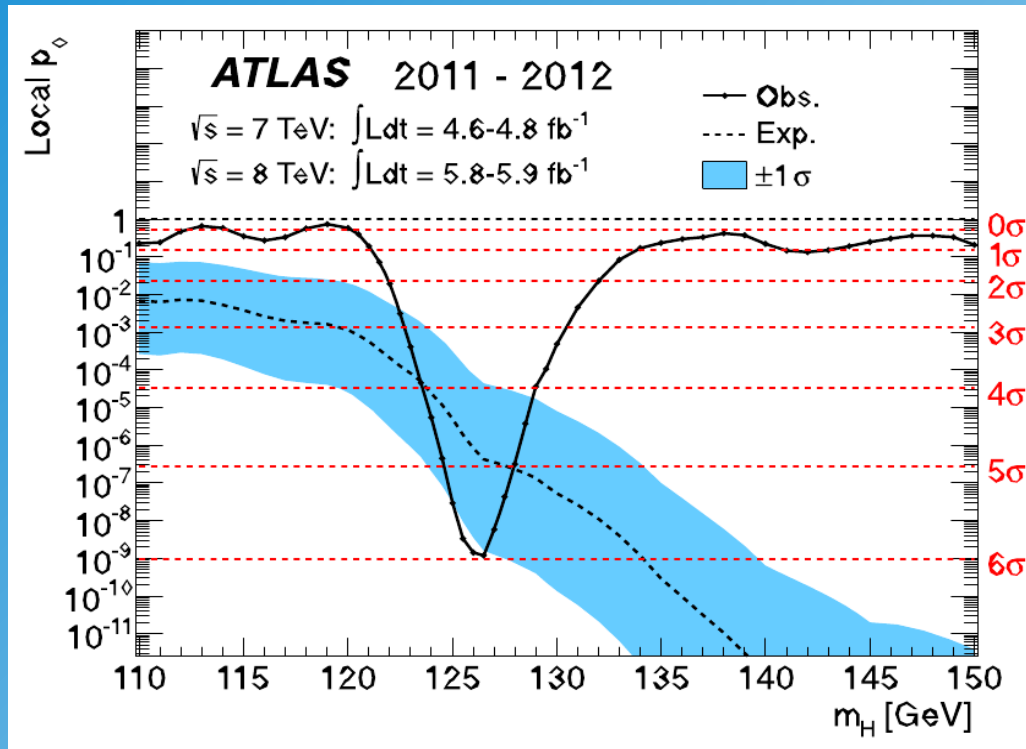
Growing Signal $H \rightarrow ZZ$

Latest Results ATLAS+CMS

H → gamma gamma

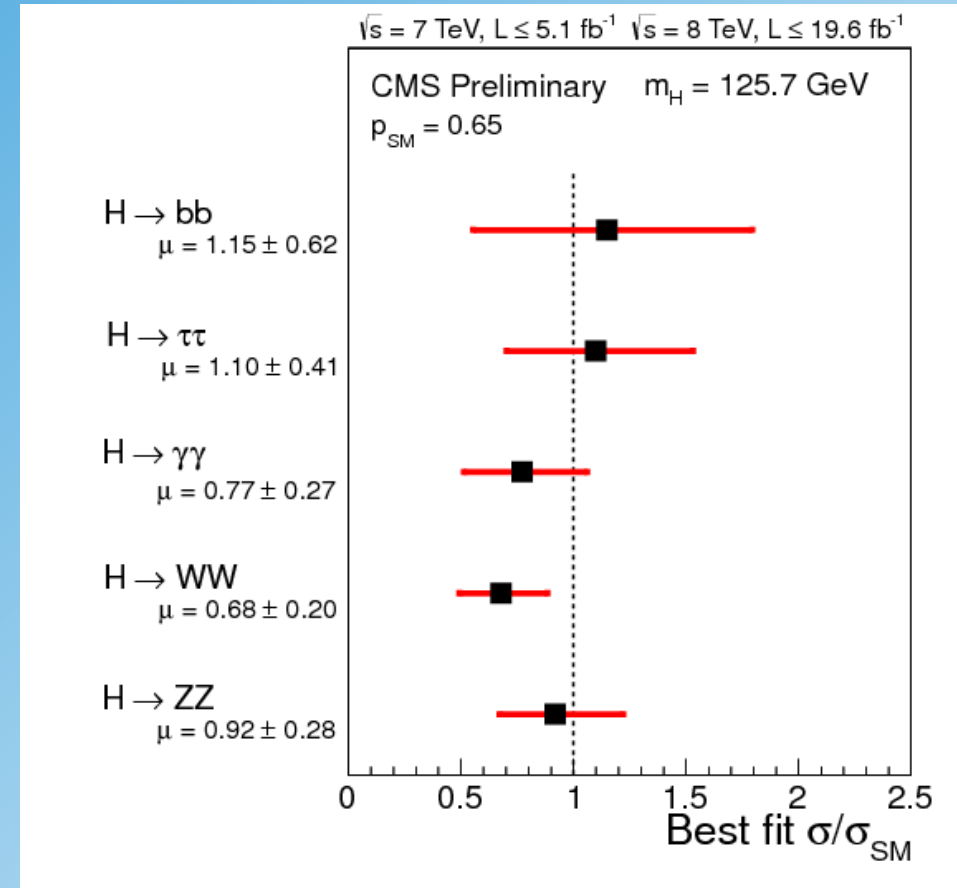
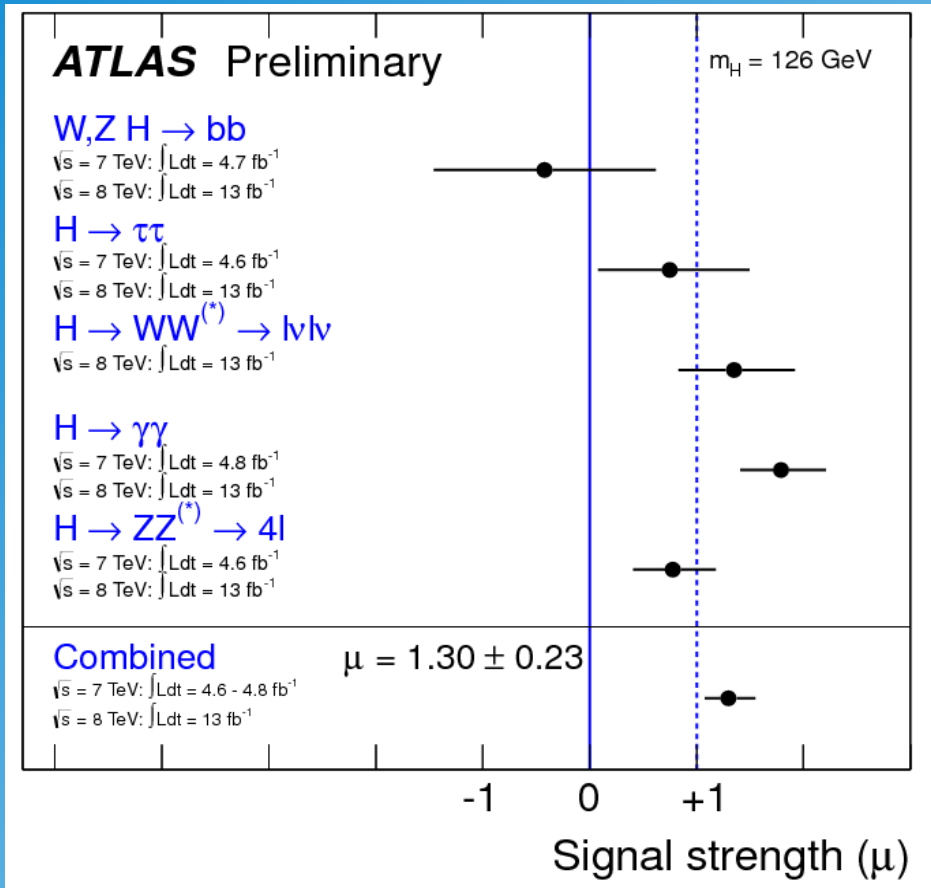


Latest Results ATLAS+CMS



Combined local probabilities

Latest Results ATLAS+CMS



Couplings strengths consistent with SM expectation

Summary

- Indirect constraints predict a light Higgs in the Standard Model
- Experimentally a Higgs with $m_H \sim 125$ GeV is most difficult to find
- Direct searches exclude (June 2012) Higgs except for a small region around $m_H \sim 125$ GeV
- Excess of Higgs candidate events at $m_H \sim 125$ by ATLAS and CMS. More events are seen than expected by SM Higgs model.
- Interestingly, $m_H \sim 125-130$ GeV is theoretically also favored by vacuum stability and triviality reasons (\rightarrow Werner Rodejohann). No new physics required up to very high mass scales!

