

Lecture:

Standard Model of Particle Physics

Heidelberg SS 2013

W- and Z-Bosons

Contents

- Discovery of “real” W- and Z-bosons
- Intermezzo: QCD at Hadron Colliders
- LEP + Detectors
- W- and Z- Physics at LEP
- W- and Z-Physics at Hadron Colliders (Tevatron+LHC)

Prediction of W and Z masses

SM predictions:

$$e = g \sin \theta_W = g' \cos \theta_W$$

Measurement of Weinberg angle:

$$\sin^2 \theta_W \approx 0.25 \quad \rightarrow \quad g \approx 0.6$$

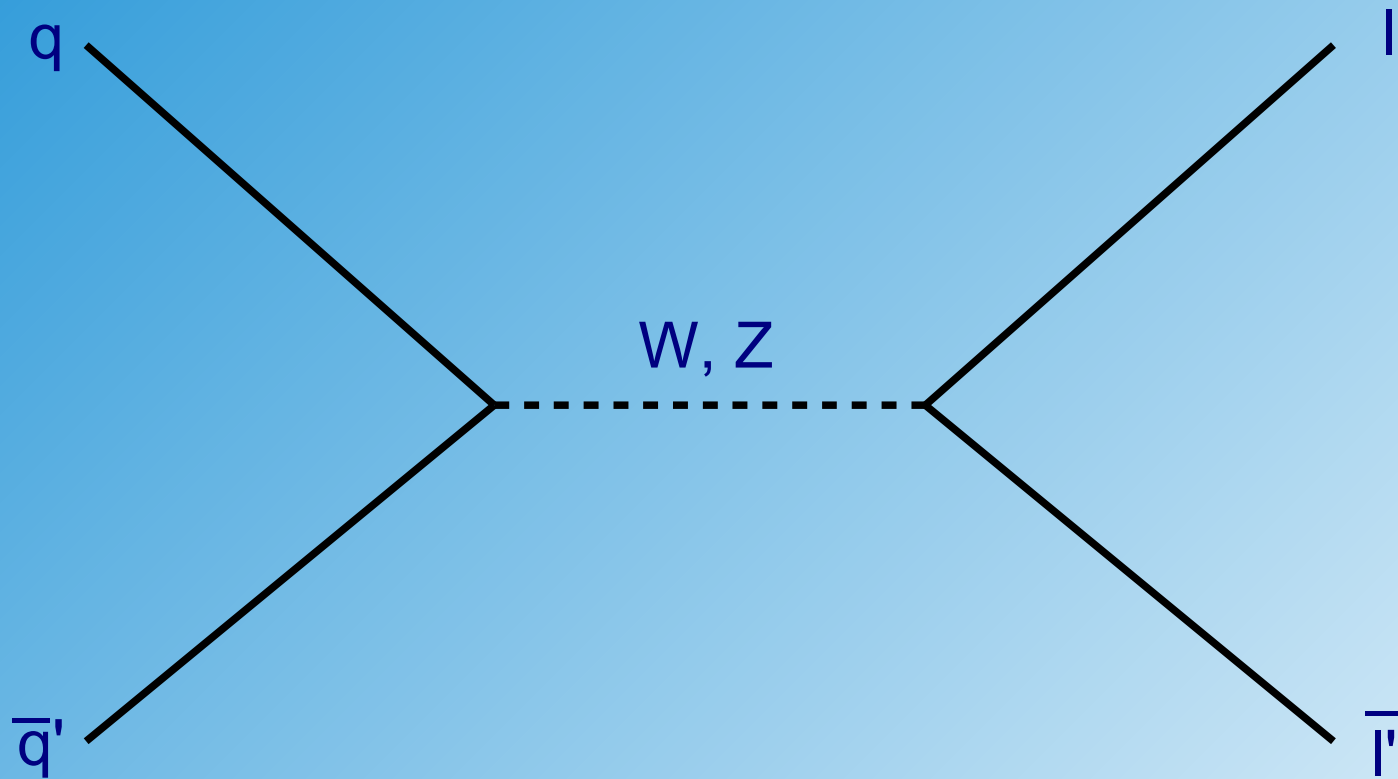
Low energy limit of W-propagator

$$G_F / \sqrt{2} = g^2 / 8 M_W^2 \quad \rightarrow \quad M_W \approx 80 \text{ GeV}$$

Relation from vector-boson mass matrix (Higgs mechanism)

$$\frac{M_W^2}{M_Z^2} = \frac{g^2}{g^2 + g'^2} = \cos^2 \theta_W \quad \rightarrow \quad M_Z \approx 90 \text{ GeV}$$

W,Z Physics at Hadron Colliders



Intermezzo QCD

QCD Lagrangian (physical fields)

$$L_{phys} = -\frac{1}{4}F^{\alpha}_{\mu\nu}(x)F_{\alpha}{}^{\mu\nu}(x) + \sum_k \frac{i}{2}(\bar{q}_k(x)\gamma^{\mu}\nabla_{\mu}q_k(x) - \nabla_{\mu}\bar{q}_k(x)\gamma^{\mu}q_k(x)) .$$

vector coupling

Covariant derivative:

$$\nabla_{\mu}q(x) = \partial_{\mu}q(x) - i g G_{\mu}^{\alpha}(x)\hat{t}_{\alpha}q(x) ;$$

SU(3) group generators

Gluon field: non-abelian coupling

$$F_{\mu\nu}^{\alpha}(x) = \partial_{\mu}G_{\nu}^{\alpha}(x) - \partial_{\nu}G_{\mu}^{\alpha}(x) + g f_{\beta\gamma}^{\alpha}G_{\mu}^{\beta}(x)G_{\nu}^{\gamma}(x) ;$$

SU(3) structure
constants

self coupling

SU(3) Group Representation

color states

$$r = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} \quad g = \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} \quad b = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

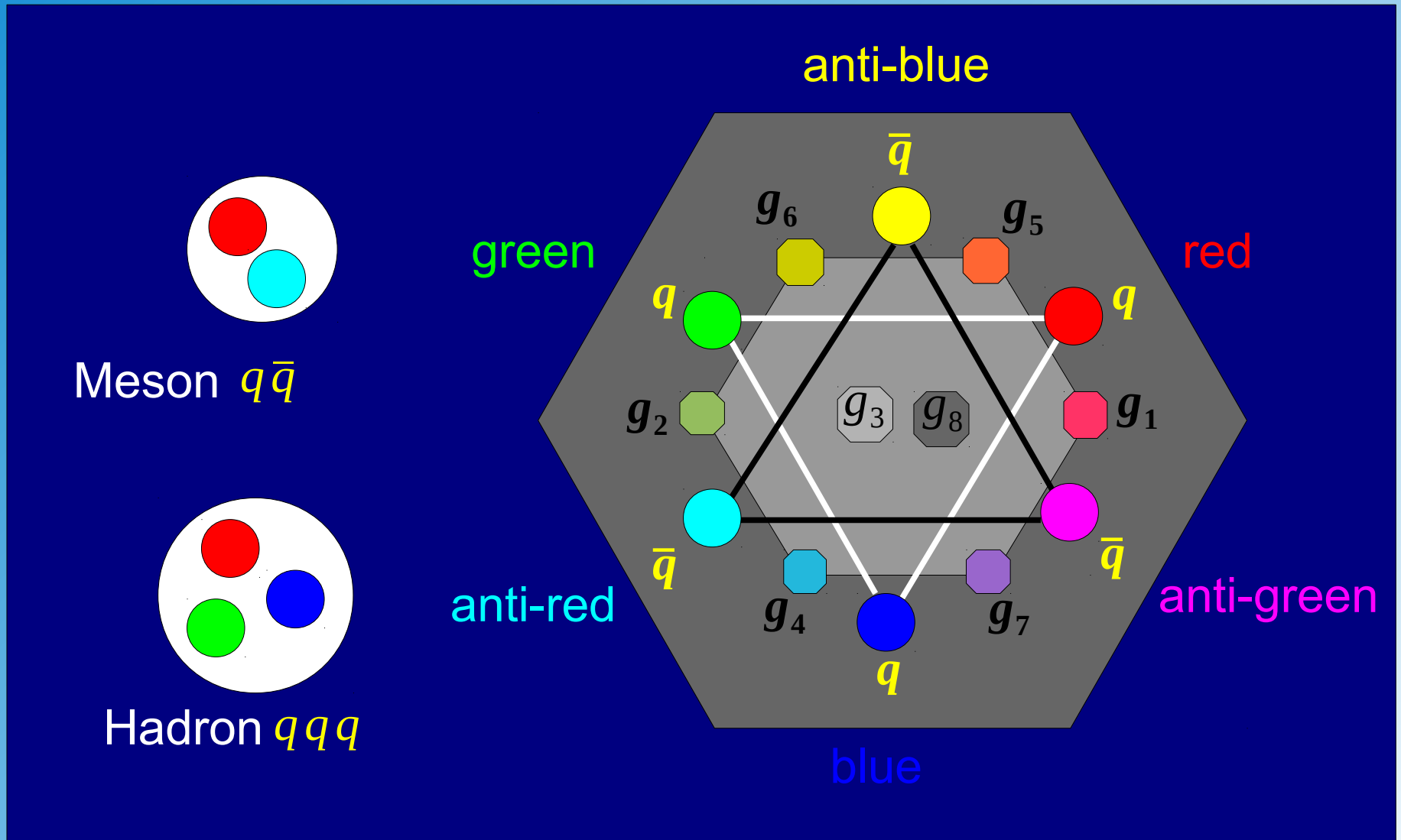
8 generators (N*N-1)

$$t_1 = \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad t_2 = \begin{pmatrix} 0 & -i & 0 \\ i & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad t_3 = \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

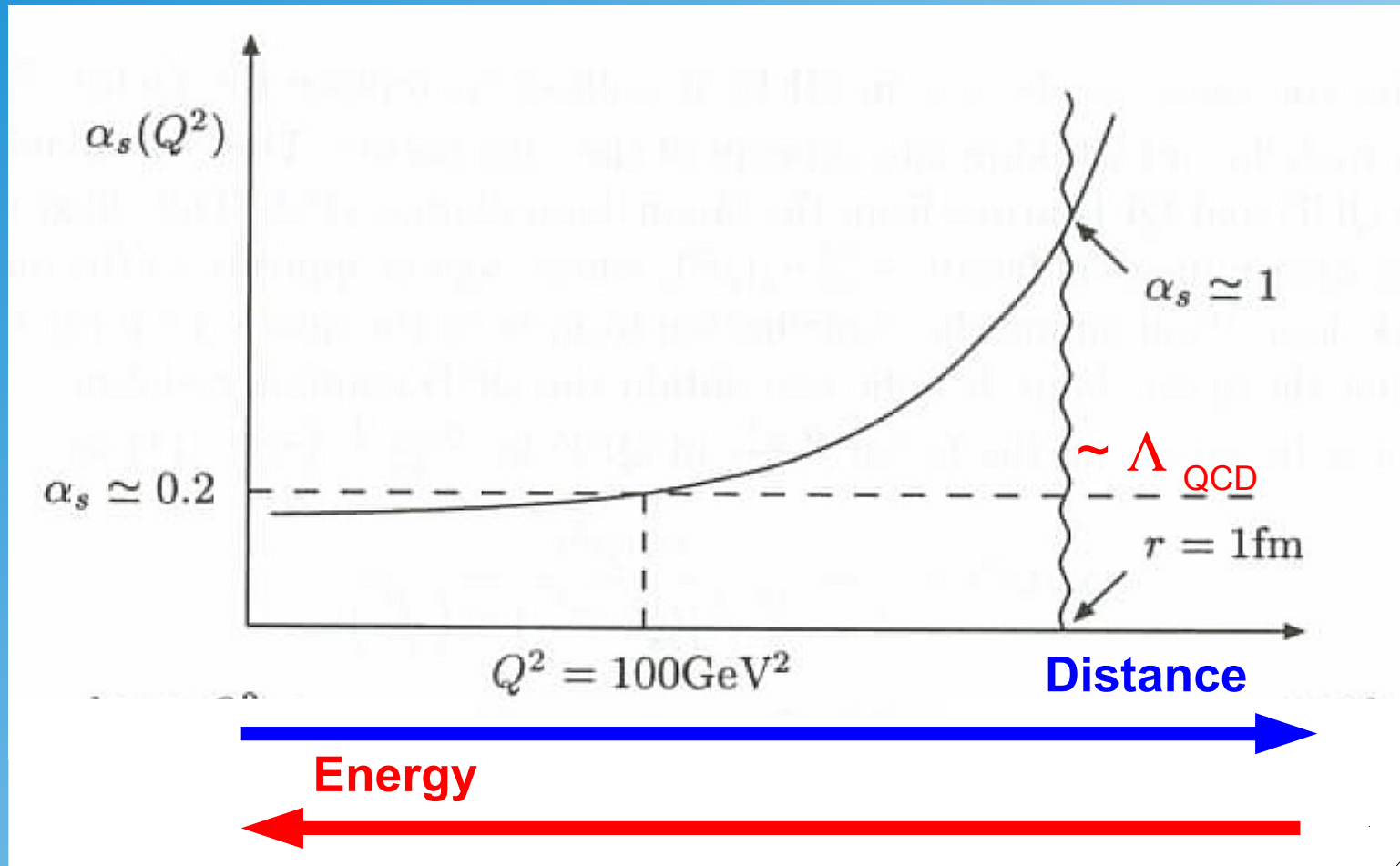
$$t_4 = \begin{pmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{pmatrix} \quad t_5 = \begin{pmatrix} 0 & 0 & -i \\ 0 & 0 & 0 \\ i & 0 & 0 \end{pmatrix} \quad t_6 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}$$

$$t_7 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & -i \\ 0 & i & 0 \end{pmatrix} \quad t_8 = \frac{1}{\sqrt{3}} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -2 \end{pmatrix}$$

Quantum Chromodynamics

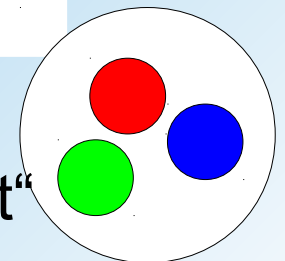


Running of α_s

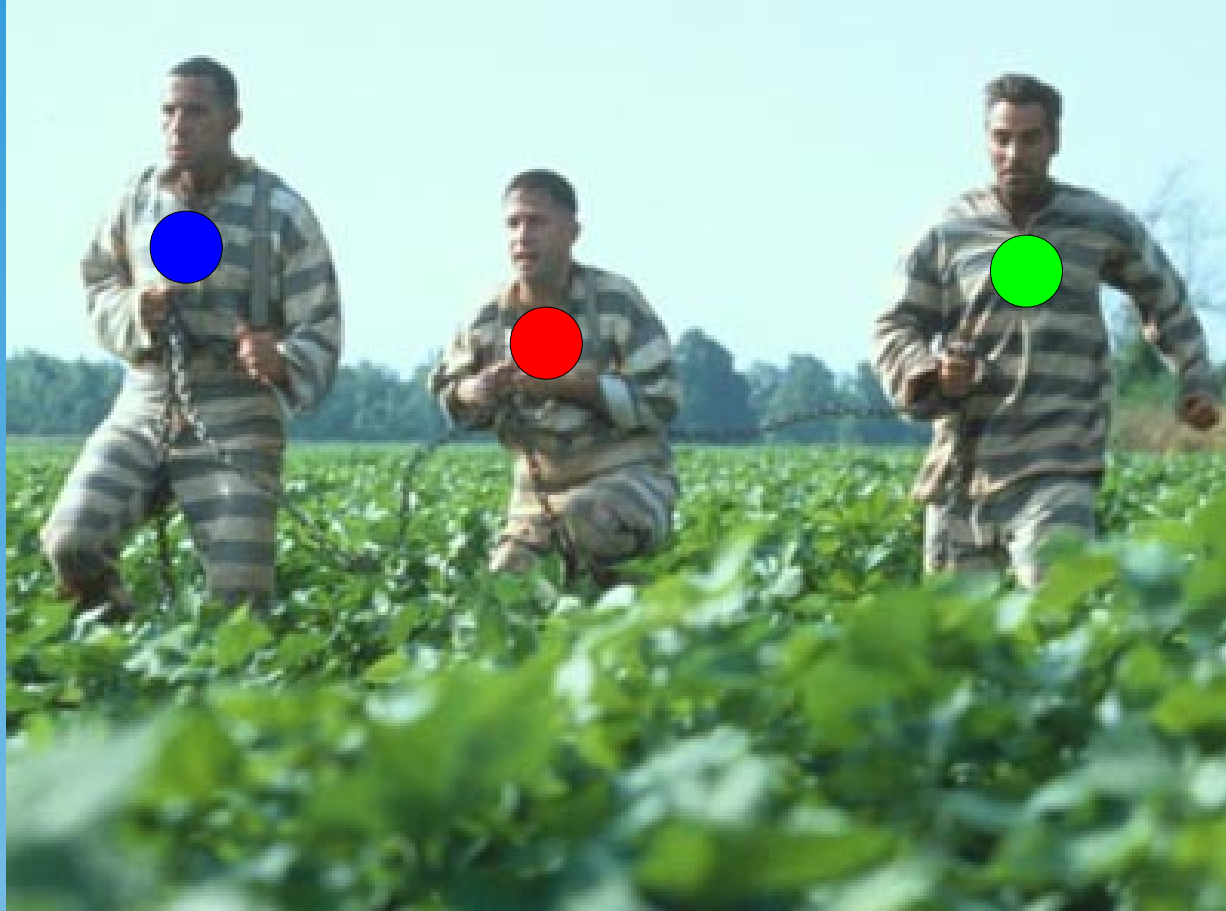


„Asymptotic Freedom“

„Confinement“



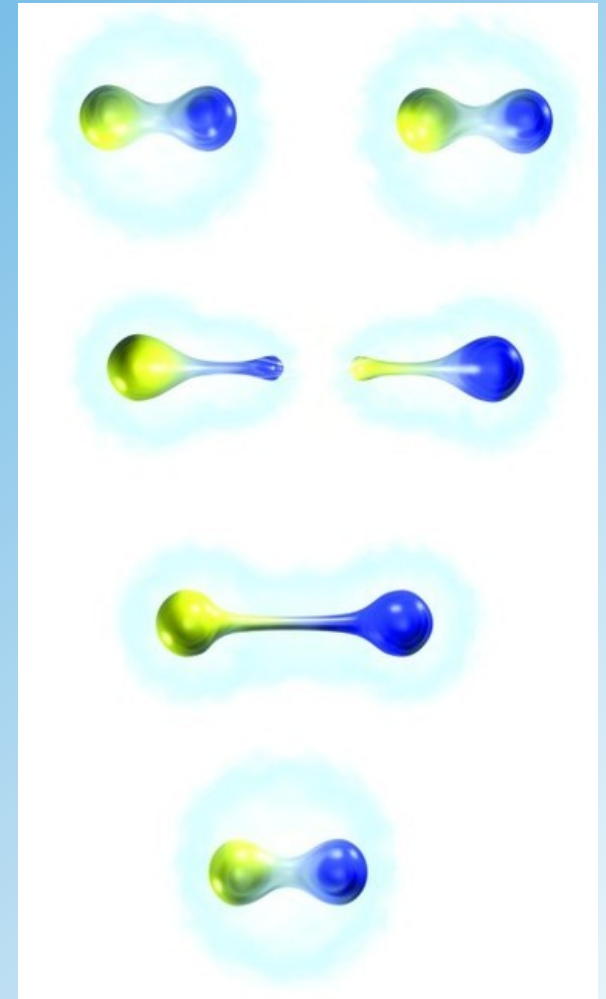
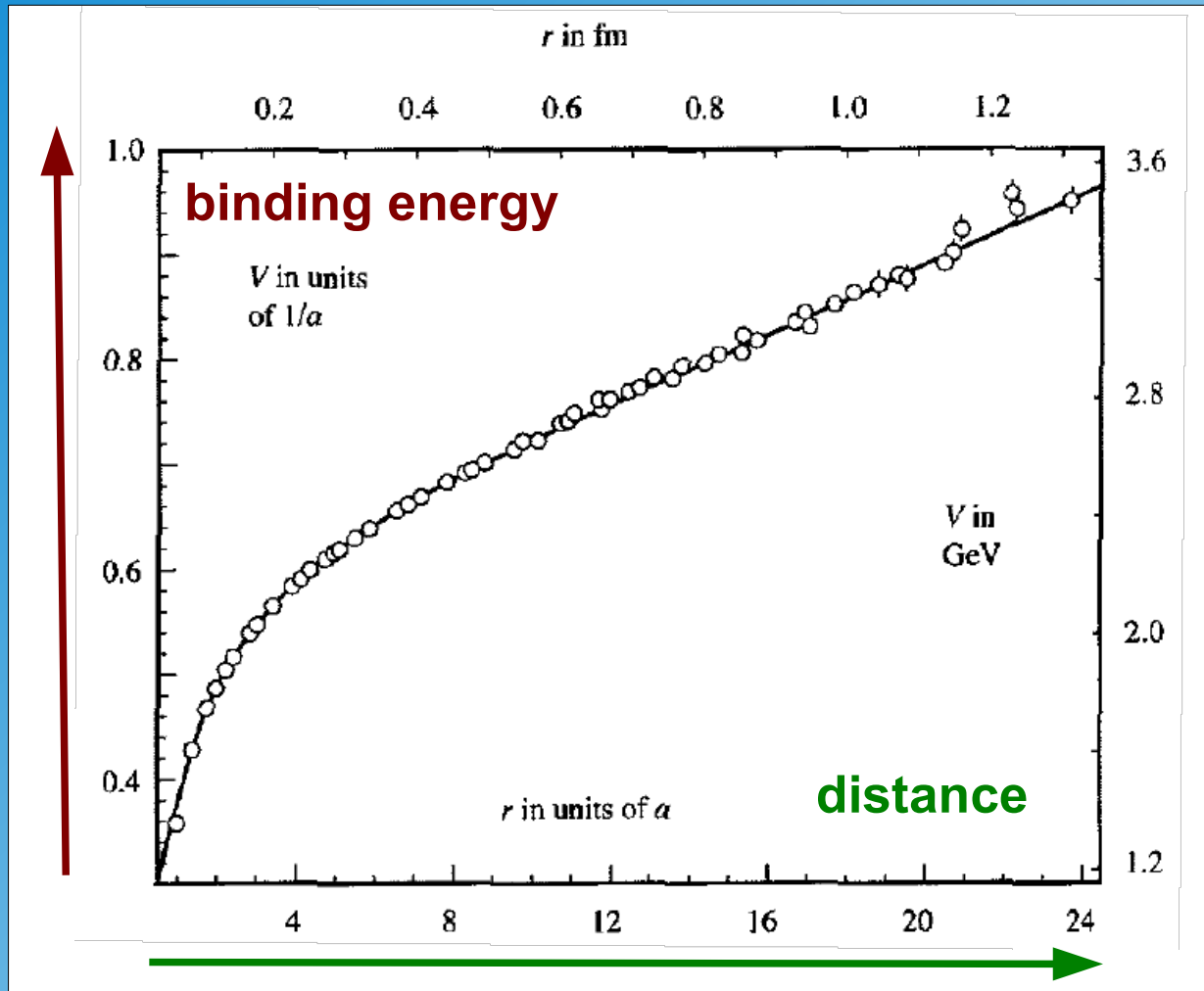
„Asymptotic Freedom“



„*Oh Brother, where art thou?*“ (2000)

Confinement

The force between two quarks is 50000 N !!!



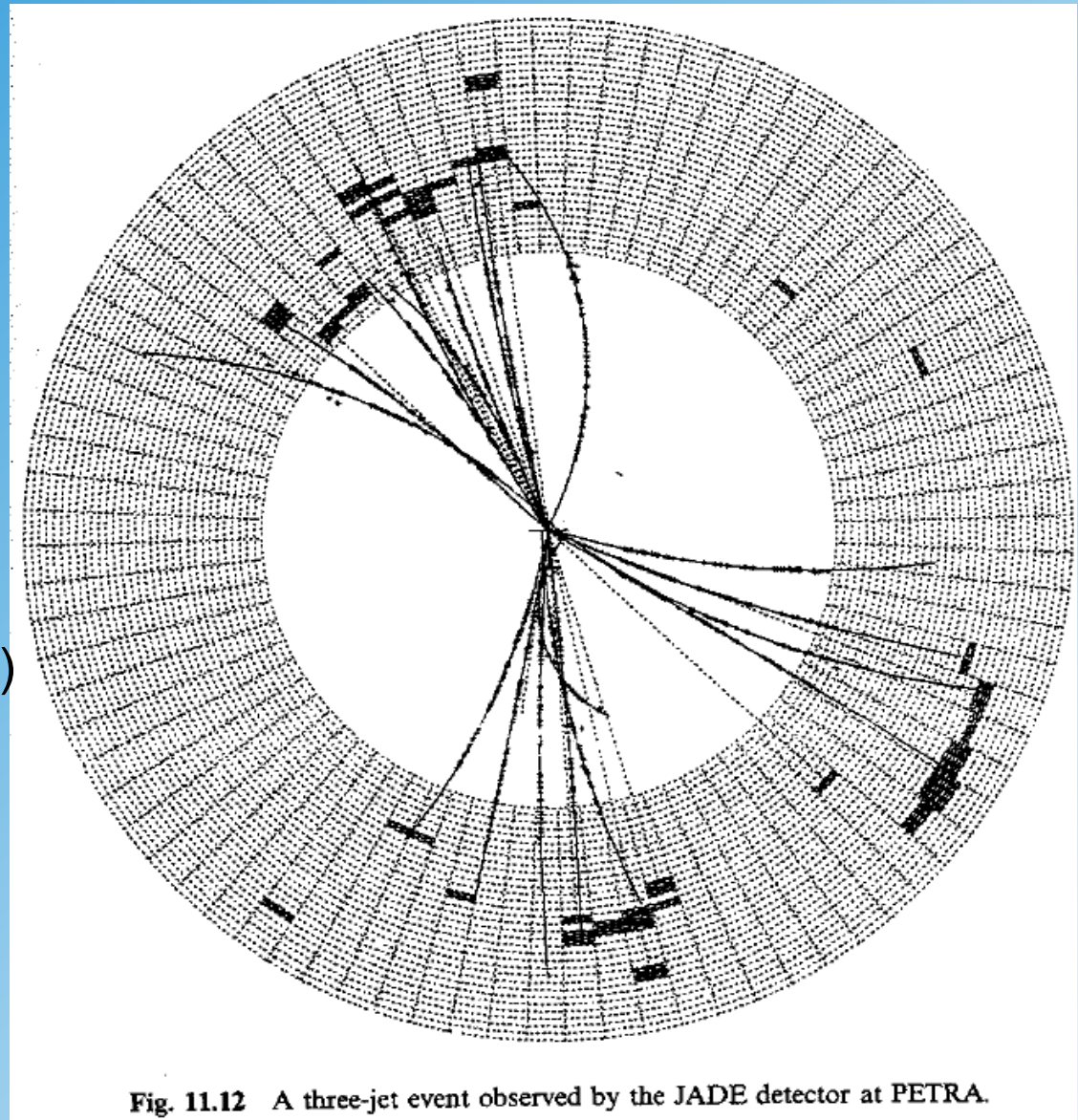
consequence: free quarks or gluons are not observable

Three-Jet Event at PETRA

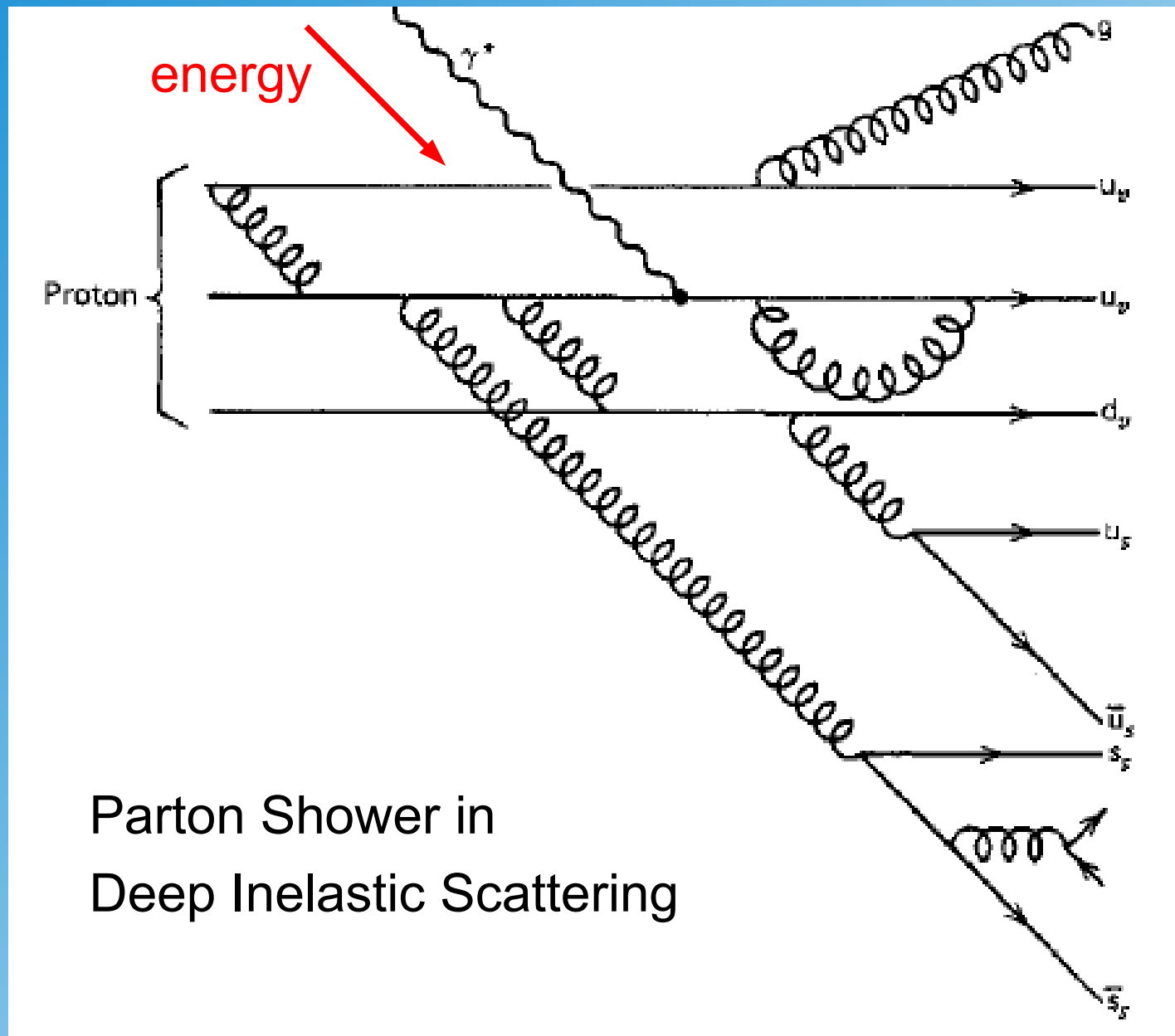
Reaction:



- Hard gluon emission
 - calculable in pQCD
 - event topology
- Soft gluon emissions
 - parton showers (non-pQCD)
 - high particle multiplicities
 - collinear emissions makes “jet” structure
- Hadronisation
 - long distance scale
 - formation of hadrons from quarks and gluons



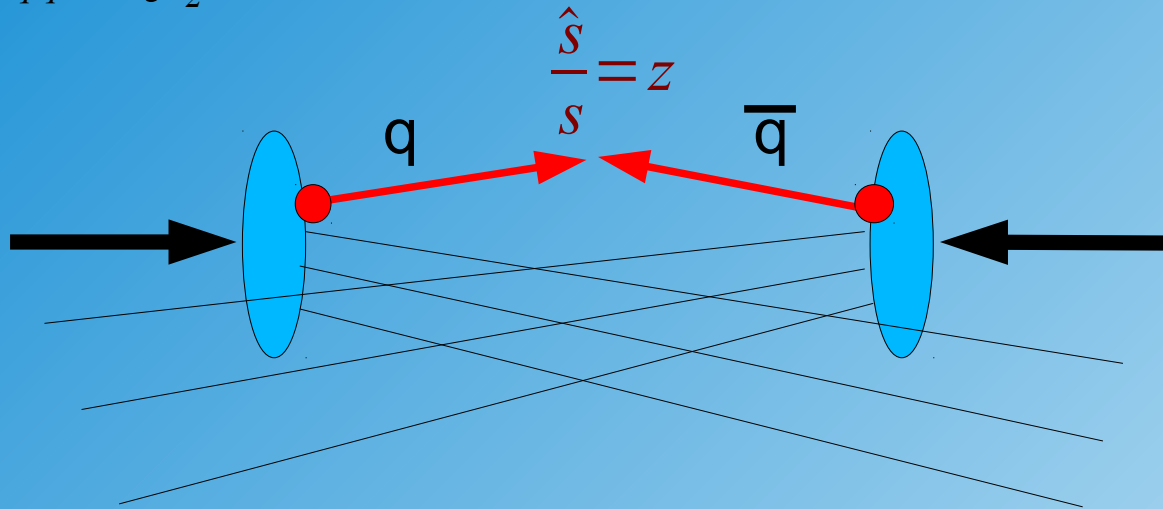
Parton Showers



Luminosity-Function

At Hadron Colliders: how to get from the proton to the parton?

$$L_{q\bar{q}} = \int_z^1 q(z/z_2) \bar{q}(z_2) dz_2/z_2 \quad \text{with} \quad z = z_1 z_2$$



s = total cms energy

\hat{s} = cms energy of
hard parton interaction

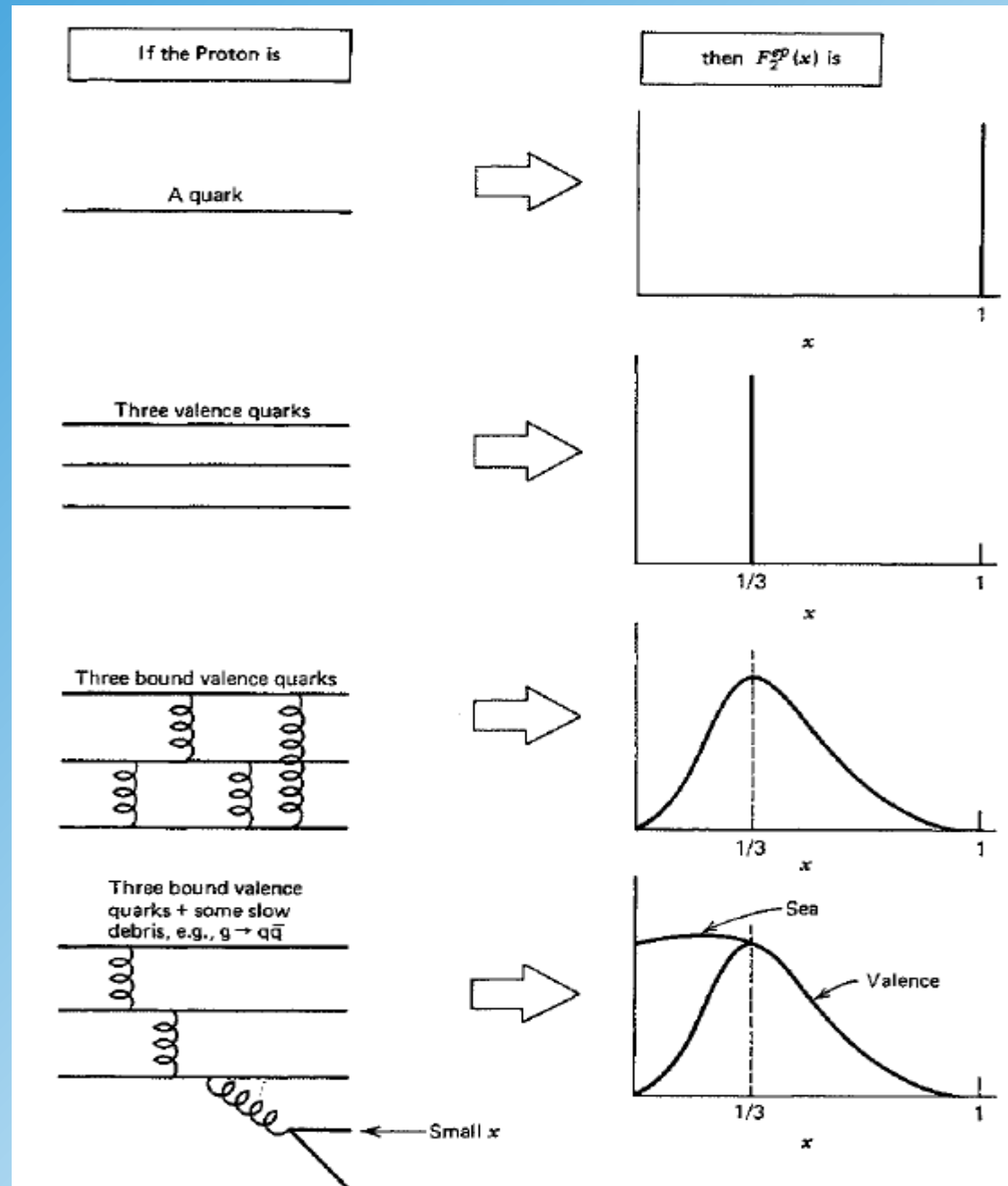
Parton density function $q = q(x, \mu^2)$

- In Lepton-Nucleon Scattering parton splitting (factorisation) scale $\mu = Q^2$
- **Question:** Which scale determines parton splitting in hadron-colliders?
- **Answer:** factorisation scale typically: $\mu_F = \hat{s}$

Input from lepton-nucleon scattering needed!

Parton Dynamics

- The x -dependence of $q(x, \mu)$ can not be calculated from first principles!
- Parton densities have to be measured by experiments
- Evolution of parton densities in Q^2 is described by DGLAP equations (splitting functions)



...coming back to W,Z Production

W,Z Production in Hadron Collisions

Reaction:

$$q \bar{q} \rightarrow W (Z) X$$

Collider energy:

$$s^{1/2} \sim 500 \text{ GeV}$$

Boson masses

$$M_{W,Z} \sim 100 \text{ GeV}$$

$$M_{W,Z} \sim \hat{s} = x_1 x_2 s$$

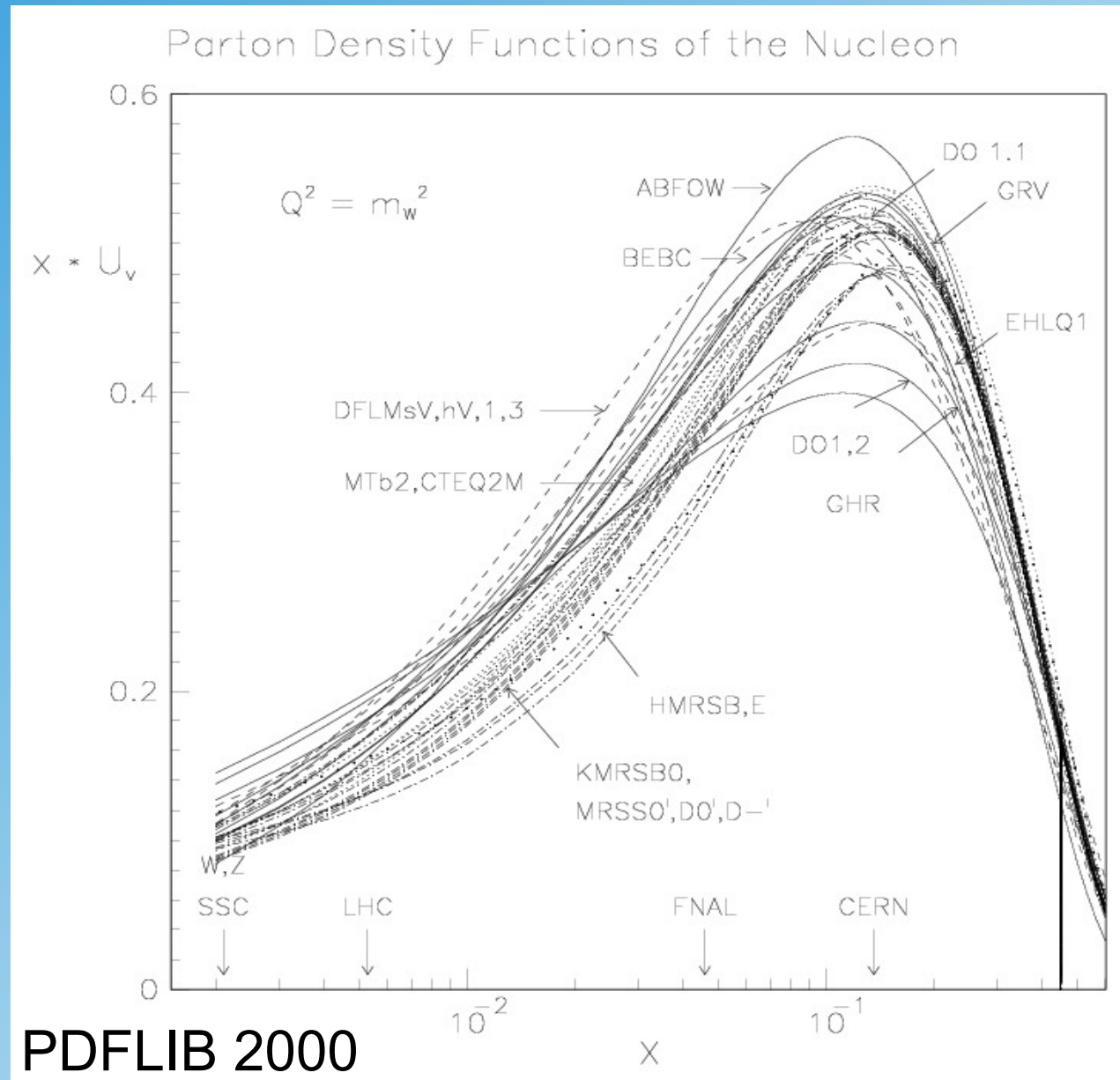
parton momentum fractions:

$$x_1, x_2 \sim 0.2$$

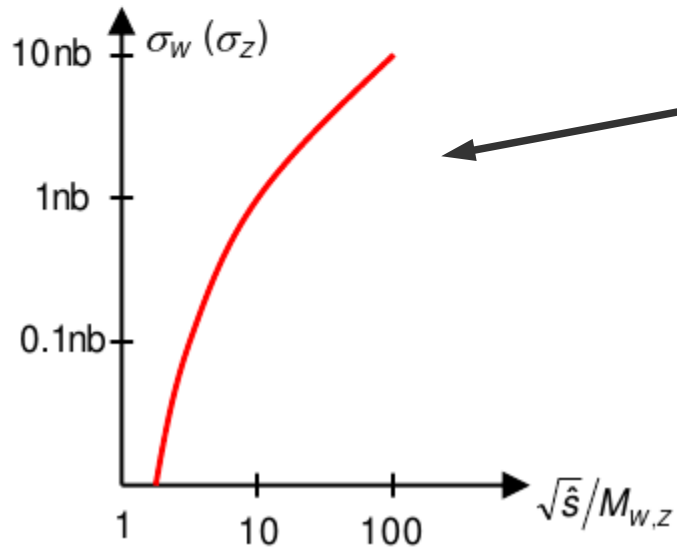
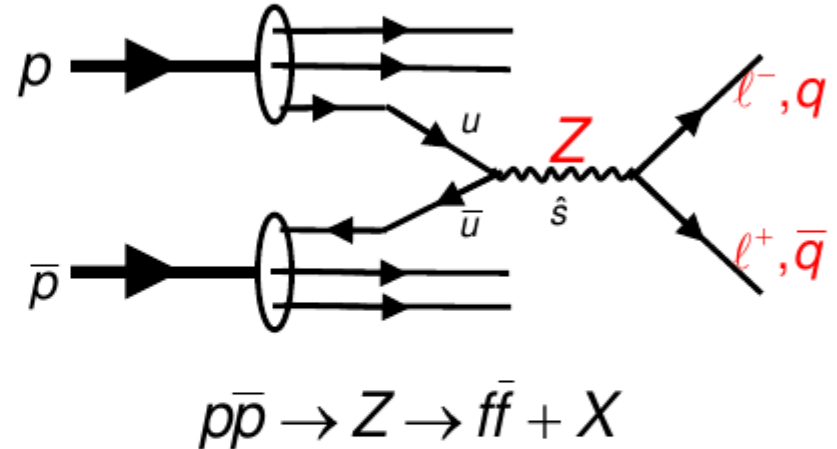
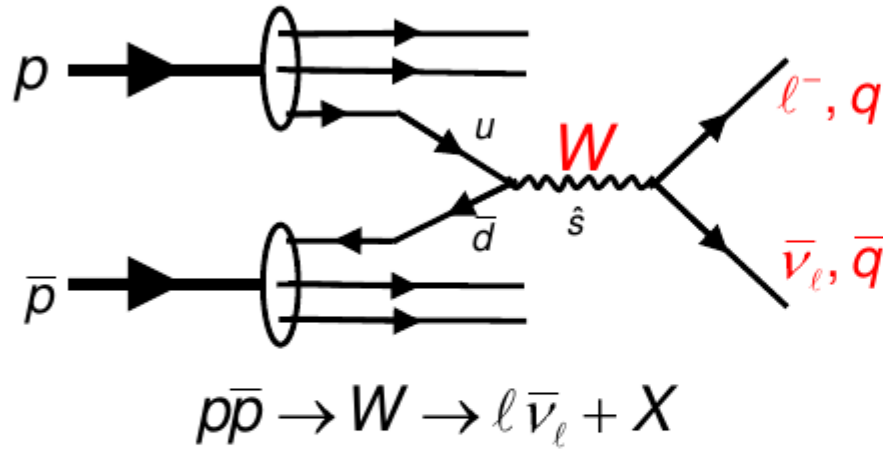
→ valence-quark region

$$p \bar{p} \rightarrow W (Z) X$$

anti-protons needed!



W,Z Cross Section



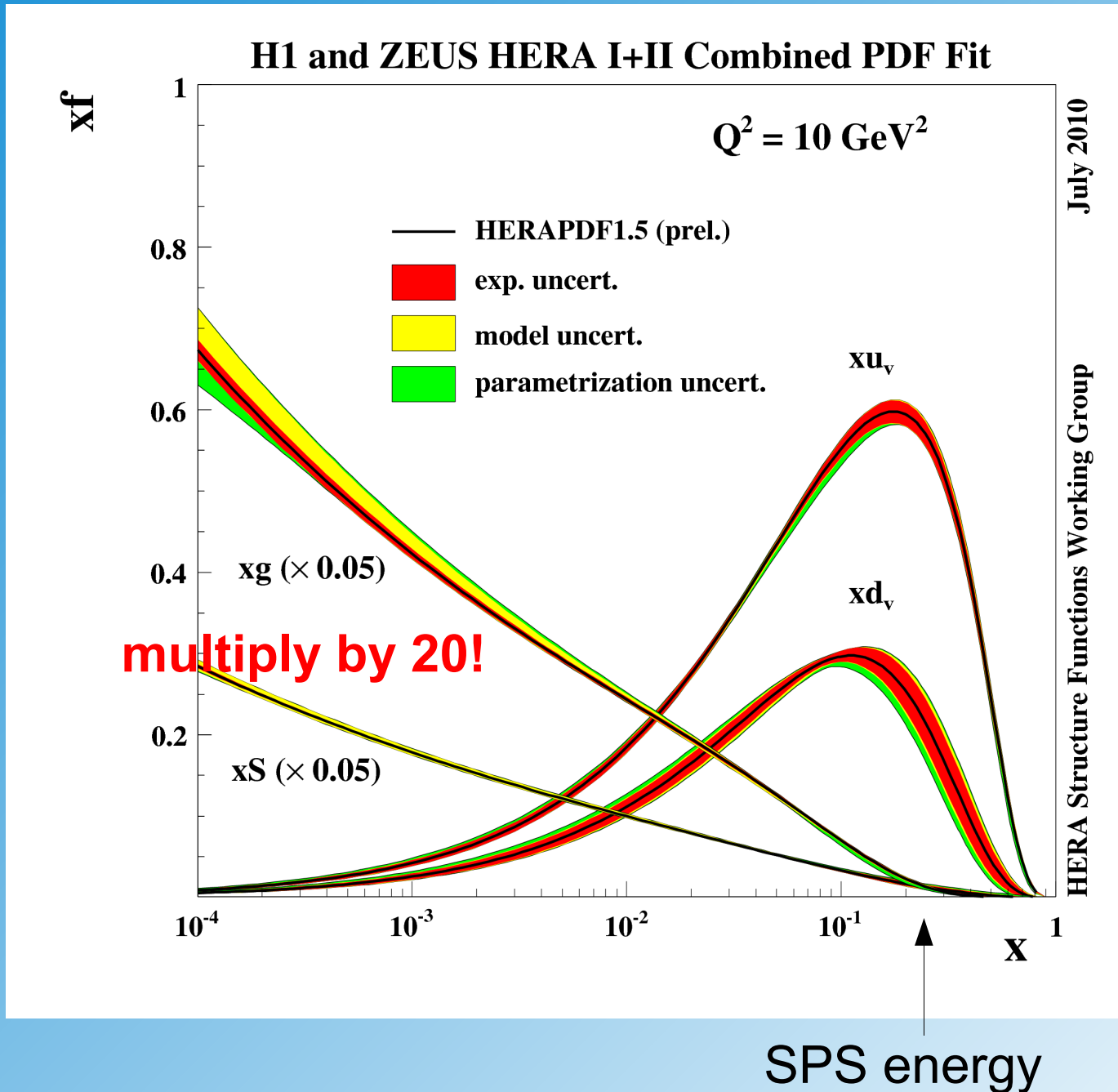
Reasonable cross section of 0.1 nb at
 $s^{1/2}/M_W \sim 2$

Typically:

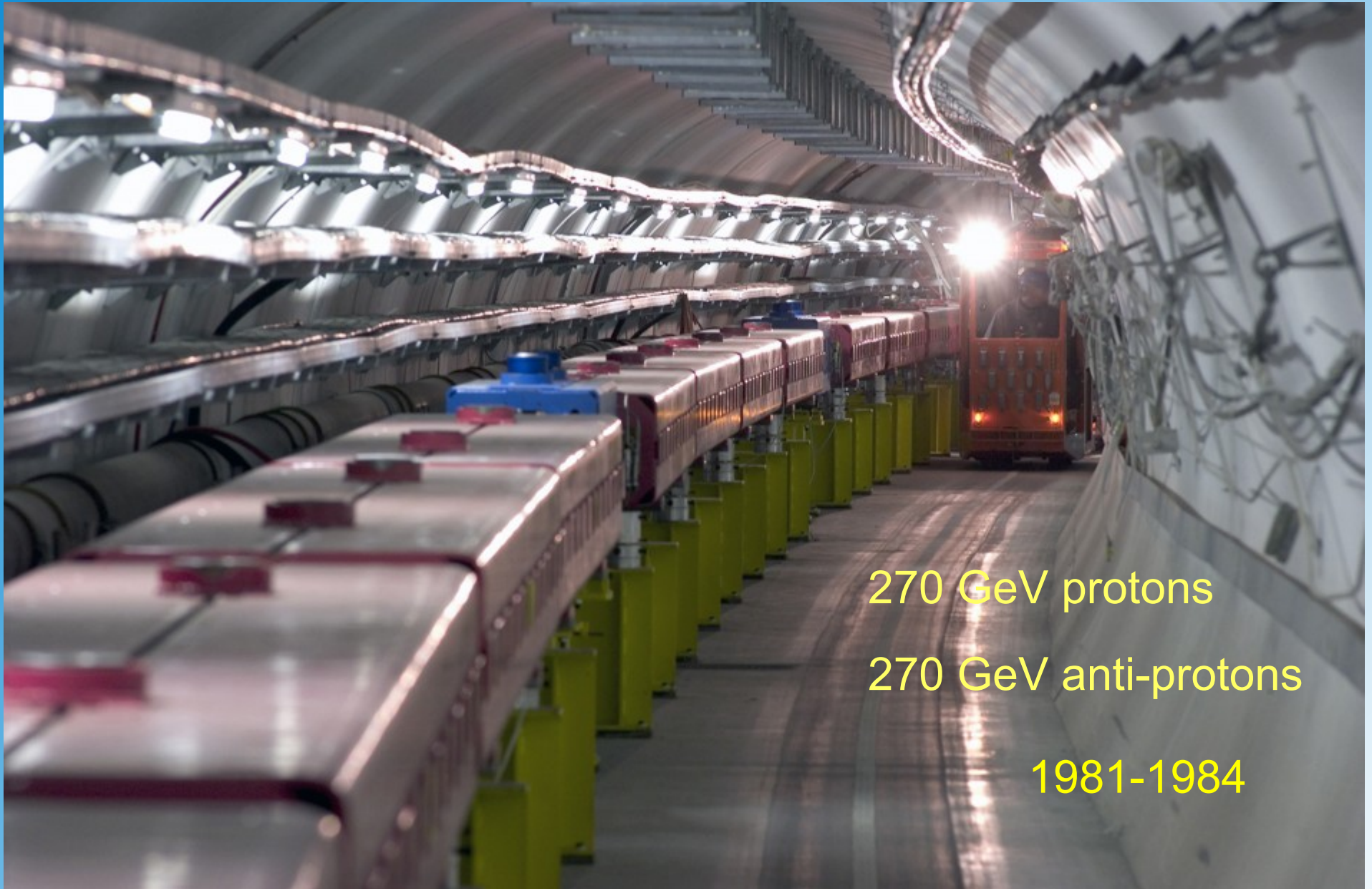
$$x_1, x_2 \sim 0.4$$

need high luminosity!

Proton Parton Densities



Super Proton (Antiproton) Synchrotron



270 GeV protons

270 GeV anti-protons

1981-1984

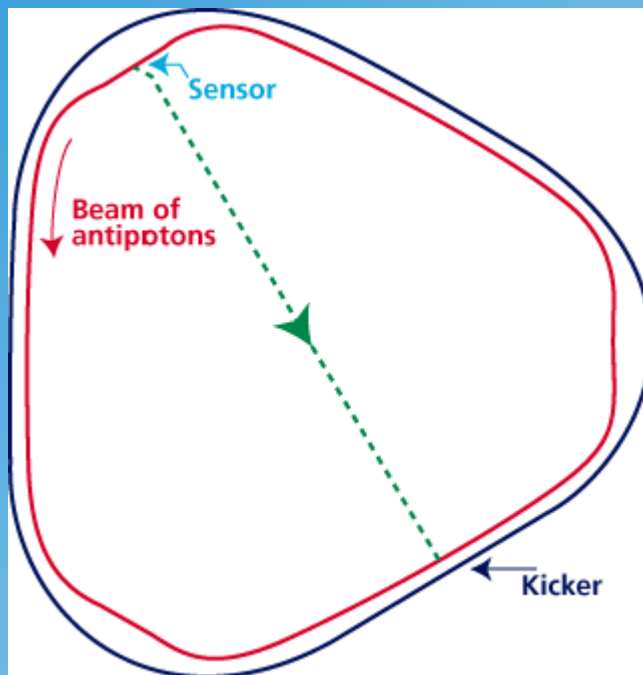
Cooling of Anti-protons

- High luminosities are obtained for small beam emittances !
- Antiprotons are hot after production!

electron cooling of anti-protons

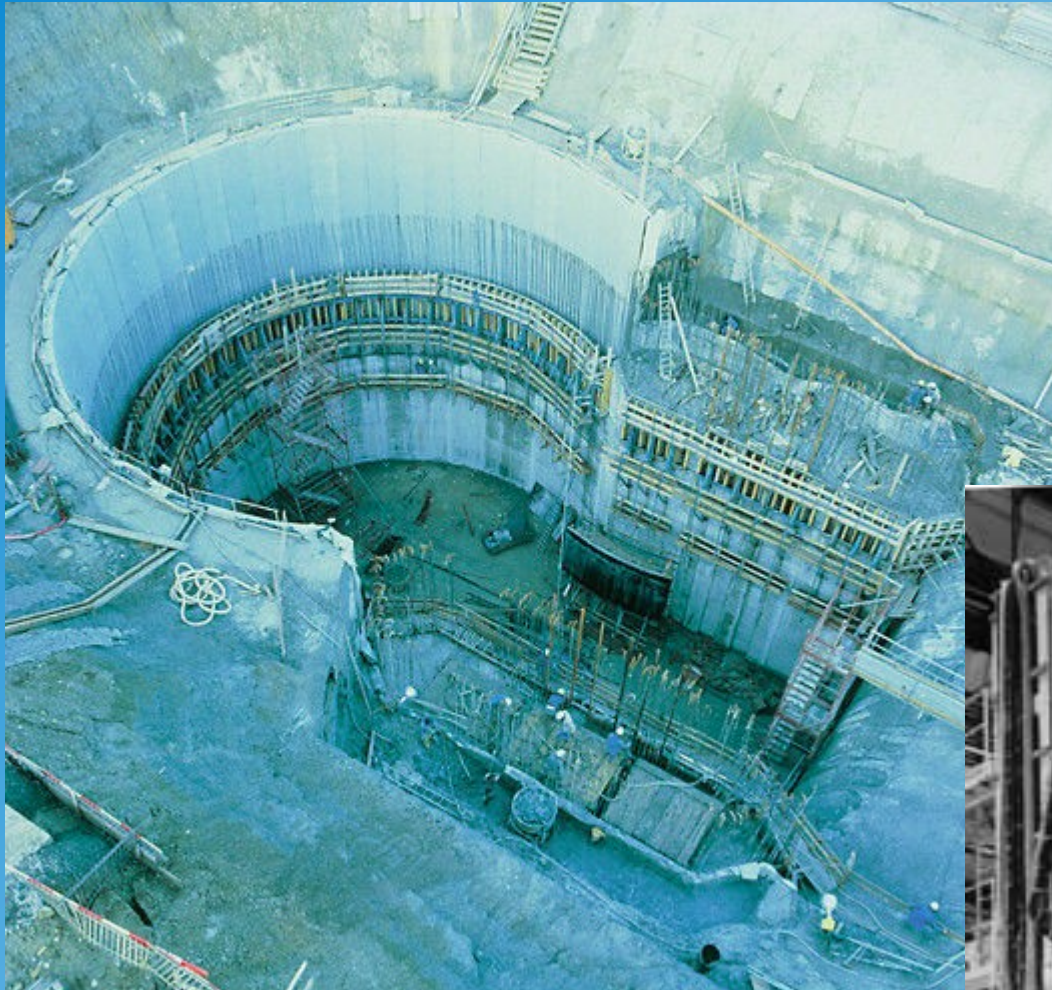


Stochastic cooling of anti-protons

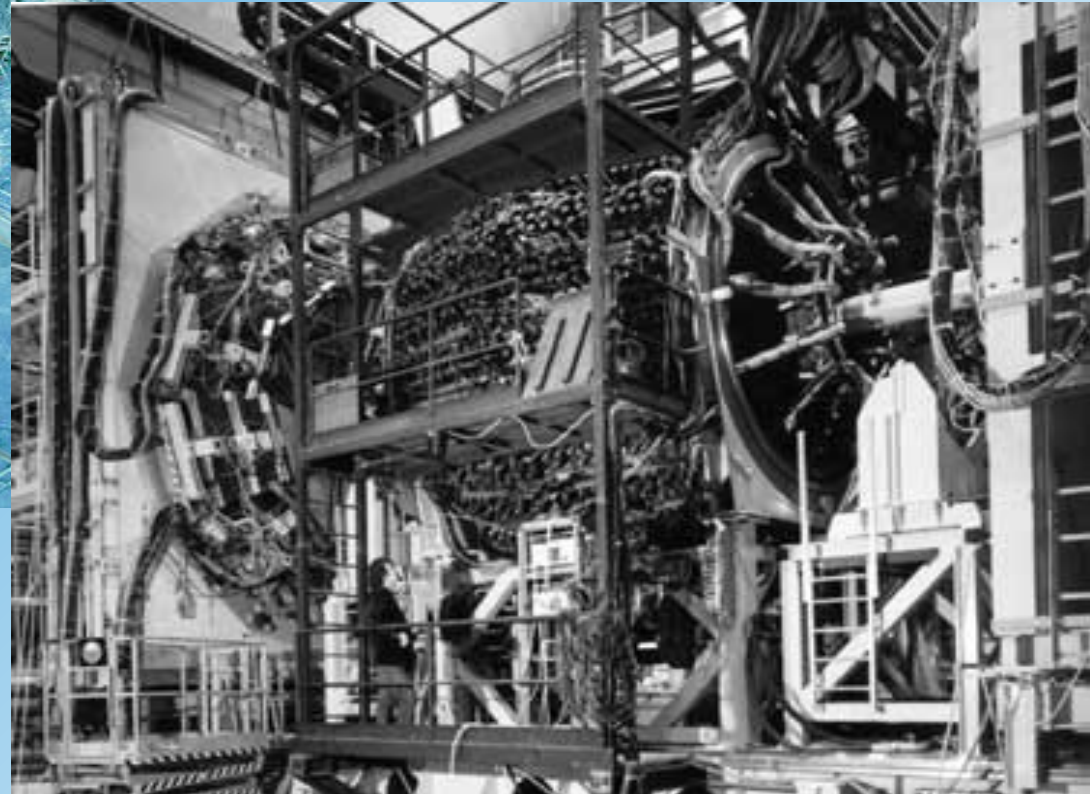


Simon
van de Meer

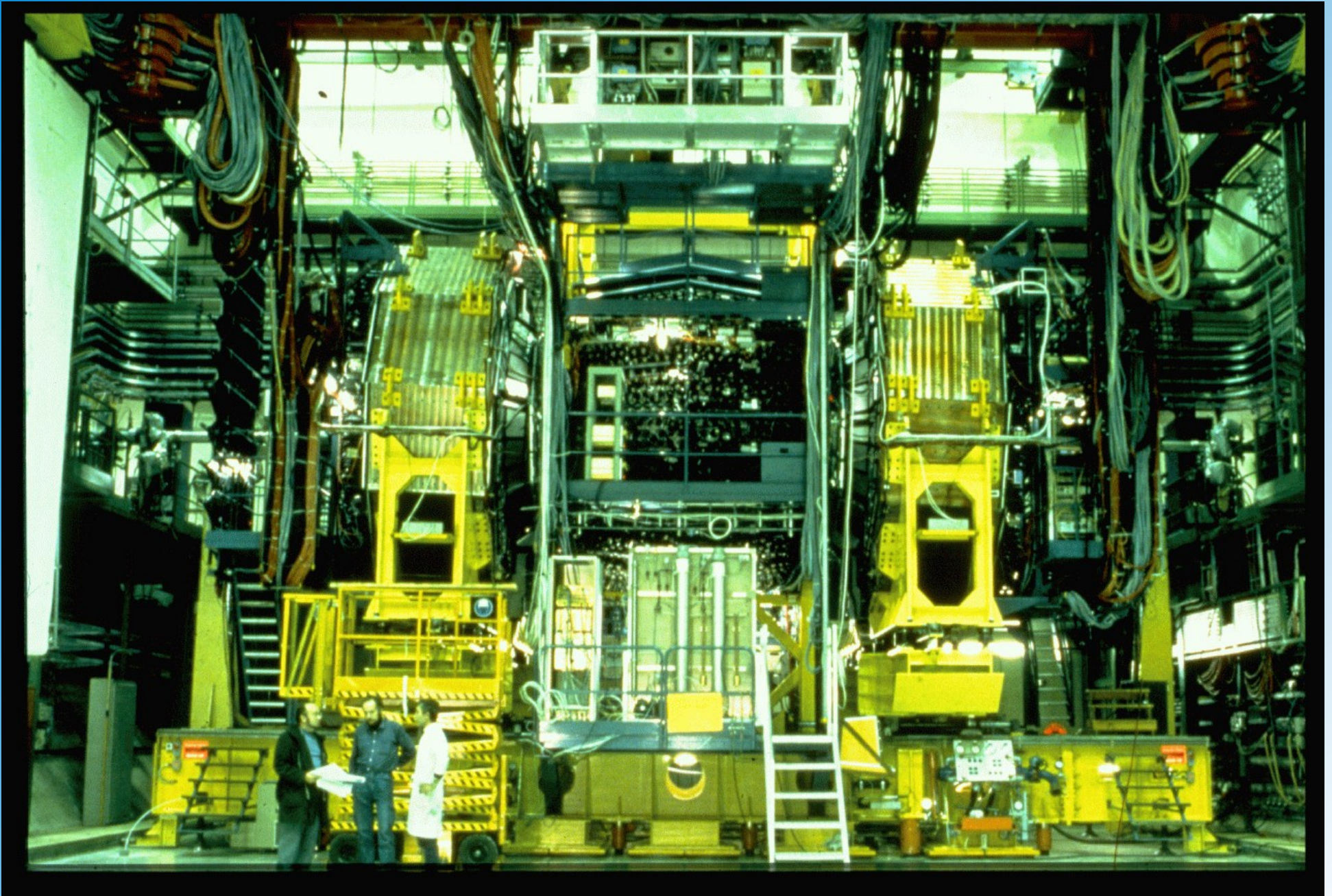
UA1 Experiment



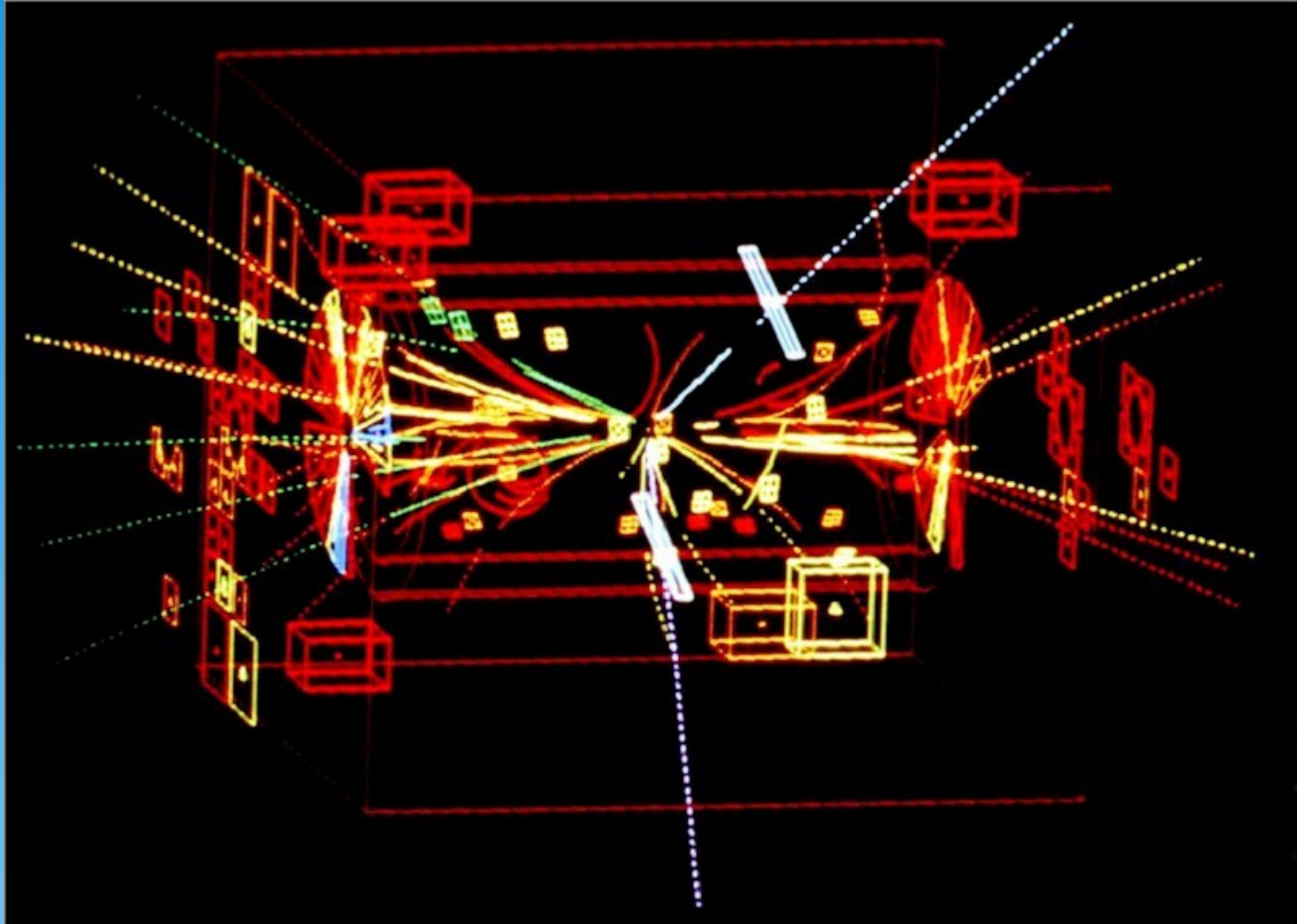
“modern” high energy
collider experiment able
to run at high collision rates
(fast electronics)



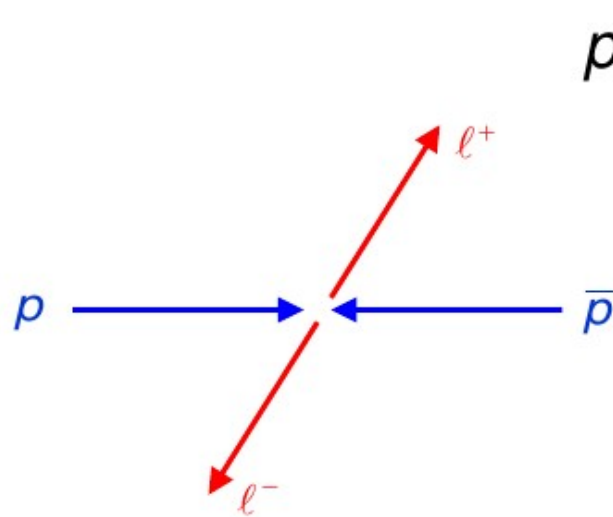
UA2 experiment



Candidate $Z \rightarrow ee$



Z-candidate Event Signature



$$p\bar{p} \rightarrow Z \rightarrow f\bar{f} + X$$

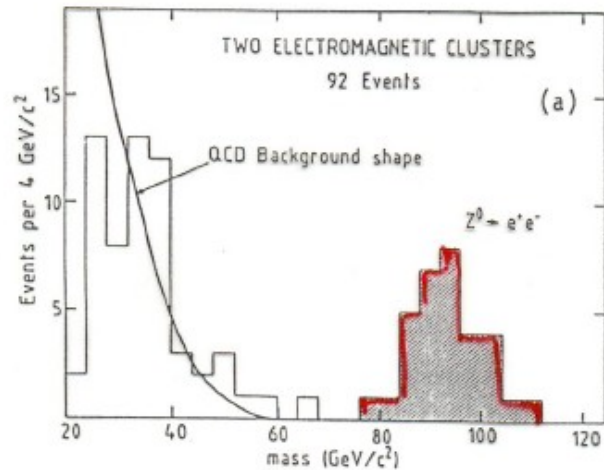
$$p + \bar{p} \rightarrow Z^0 + X$$

$$\downarrow$$

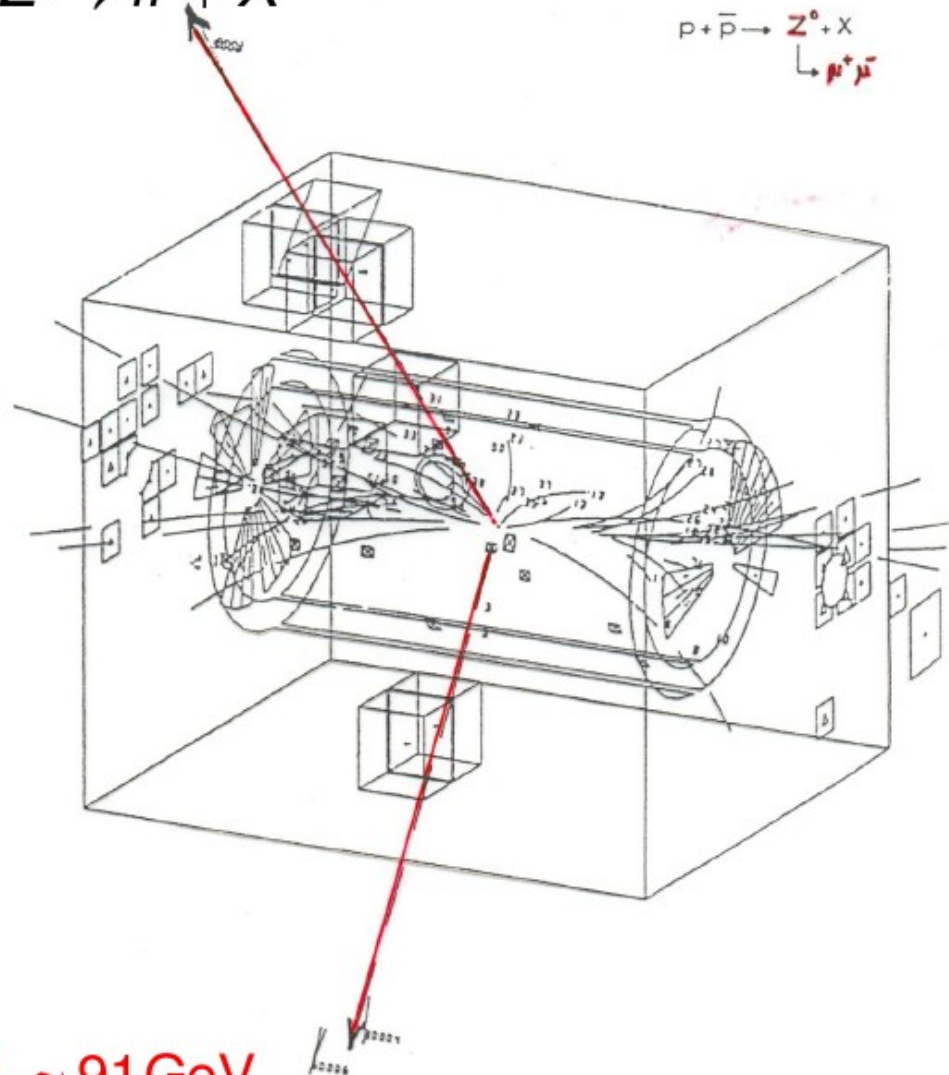
$$\mu^+ \mu^-$$

High-energy lepton pair:

$$m_{\ell\ell}^2 = (p_{e^+} + p_{e^-})^2 = M_Z^2$$



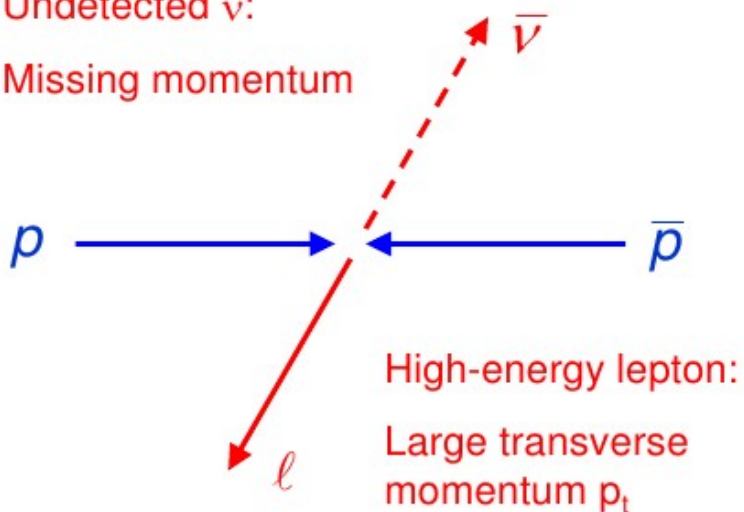
$$M_Z \approx 91 \text{ GeV}$$



W-candidates

$$p\bar{p} \rightarrow W \rightarrow l\bar{\nu}_l + X \quad W^- \rightarrow e\bar{\nu}$$

Undetected ν :
Missing momentum



High-energy lepton:
Large transverse momentum p_T

How can the W mass be reconstructed ?

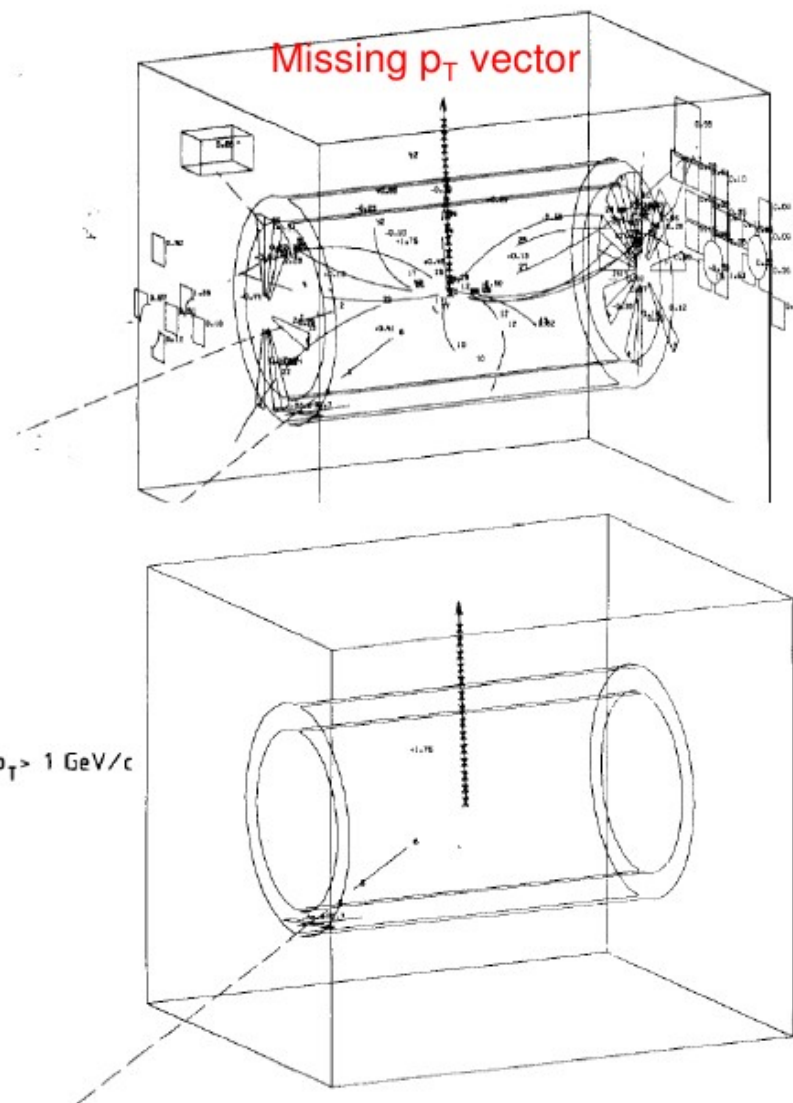
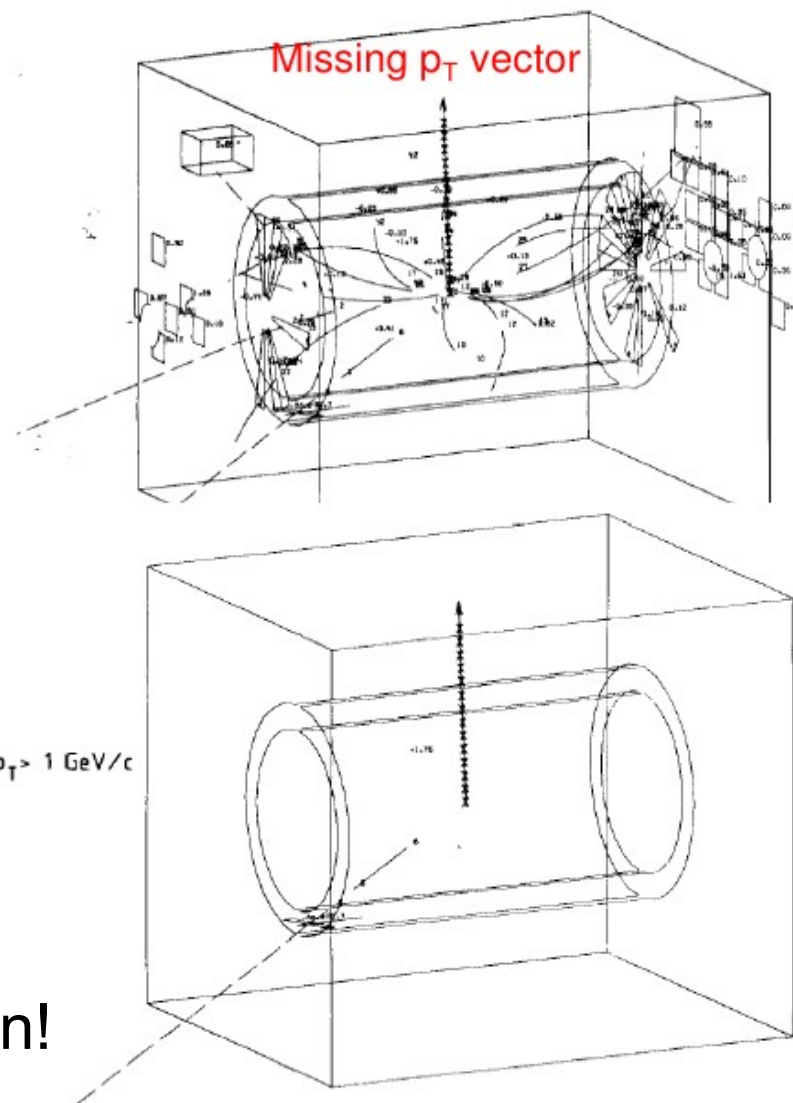
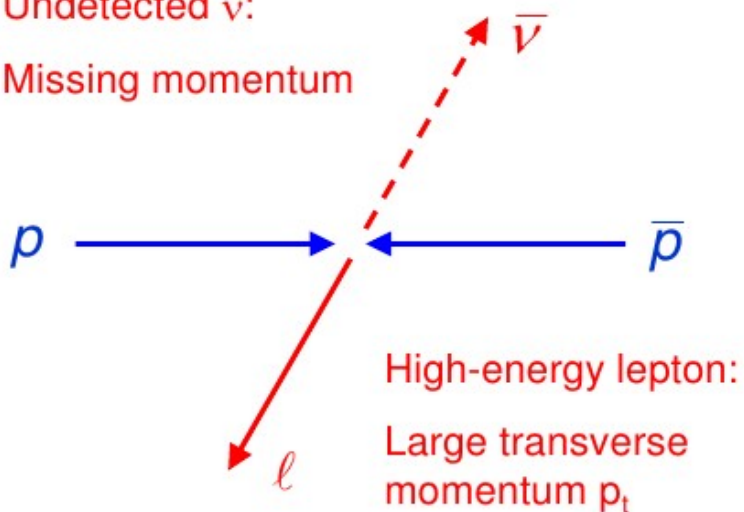


Fig. 1.6b. The same as picture (a), except that now only particles with $p_T > 1$ GeV/c and calorimeters with $E_c > 1$ GeV are shown.

W-candidates

$$p\bar{p} \rightarrow W \rightarrow l\bar{\nu}_l + X \quad W^- \rightarrow e\bar{\nu}$$

Undetected ν :
Missing momentum

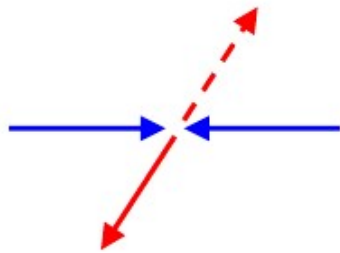


How can the W mass be reconstructed ?

exploit momentum conservation!

Kinematic Reconstruction of W-bosons

W mass measurement



In the W rest frame:

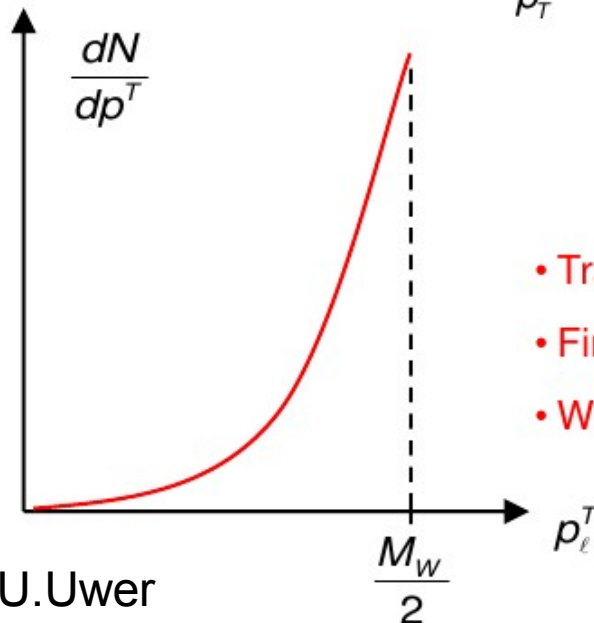
- $|\vec{p}_\ell| = |\vec{p}_\nu| = \frac{M_W}{2}$
- $|p_\ell^T| \leq \frac{M_W}{2}$

In the lab system:

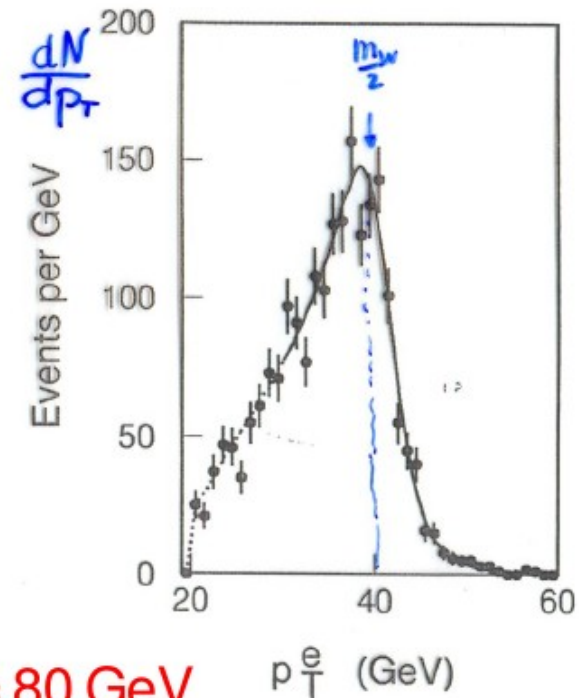
- W system boosted only along z axis
- p_T distribution is conserved: maximum $p_T = M_W / 2$

Jacobian Peak:

$$\frac{dN}{p_T} \sim \frac{2p_T}{M_W} \cdot \left(\frac{M_W^2}{4} - p_T^2 \right)^{-1/2}$$



- Trans. Movement of the W
- Finite W decay width
- W decay not isotropic

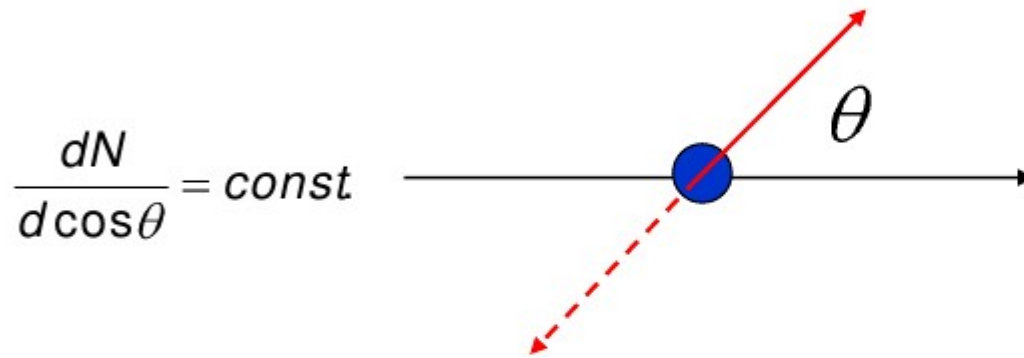


$M_W \approx 80 \text{ GeV}$

Jacobian Peak

Assume isotropic decay of the W boson in its CM system:

(Not really correct: W boson has spin=1 → decay is not isotropic!)



$$\frac{dN}{d\cos\theta} = \text{const}$$

$$\sin\theta = \frac{p_T}{p} = \frac{p_T}{M_W/2}$$

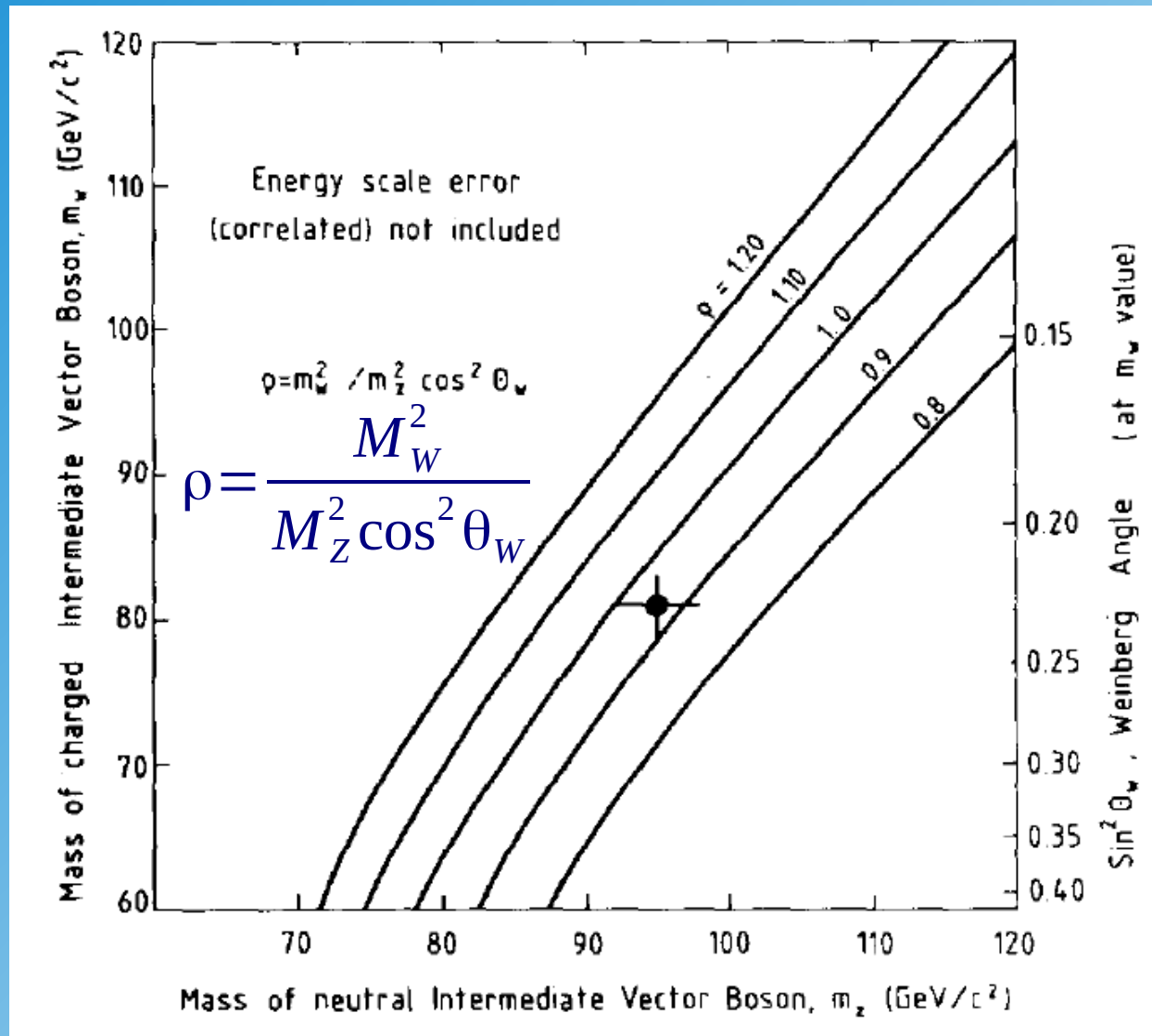
$$1 - \cos^2\theta = \left(\frac{p_T}{M_W/2}\right)^2$$

$$d\cos\theta \sim \frac{p_T}{(M_W/2)^2} \frac{dp_T}{\cos\theta}$$

$$\frac{dN}{dp_T} = \left(\frac{dN}{d\cos\theta}\right) \cdot \boxed{\left(\frac{d\cos\theta}{dp_T}\right)} \sim \frac{2p_T}{M_W} \cdot \left(\frac{M_W^2}{4} - p_T^2\right)^{-1/2}$$

Jacobian

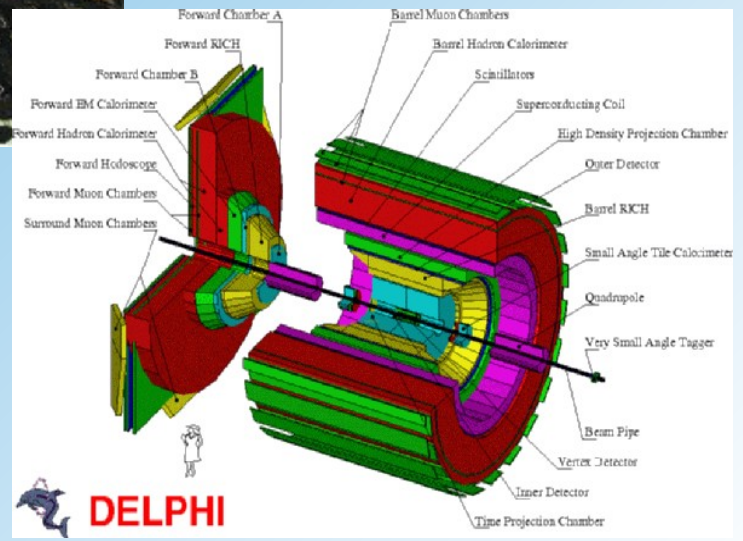
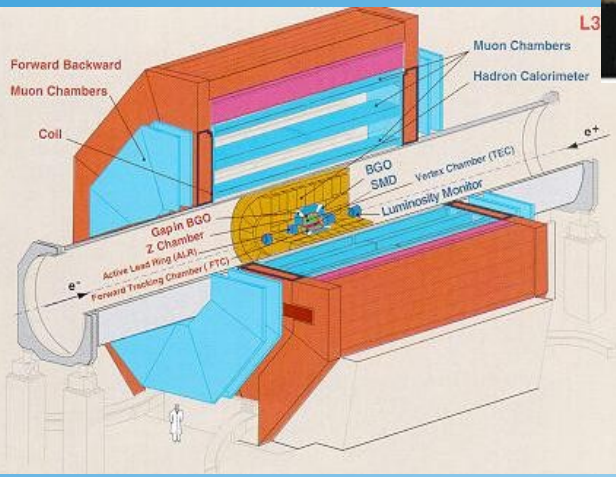
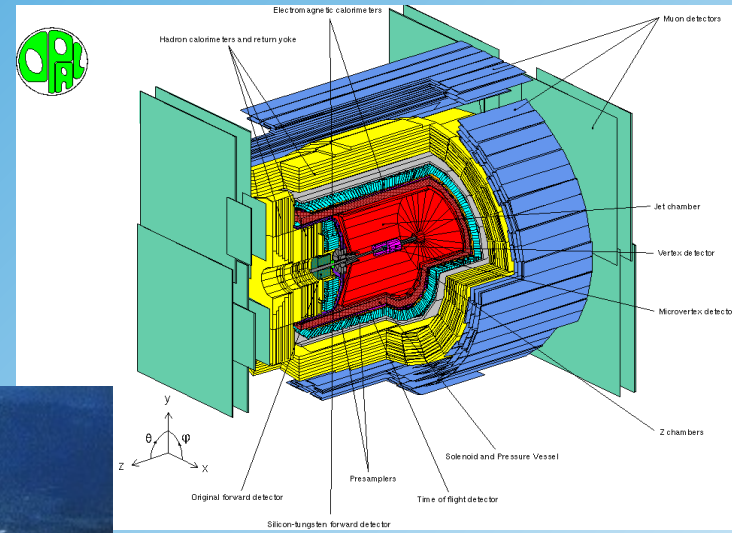
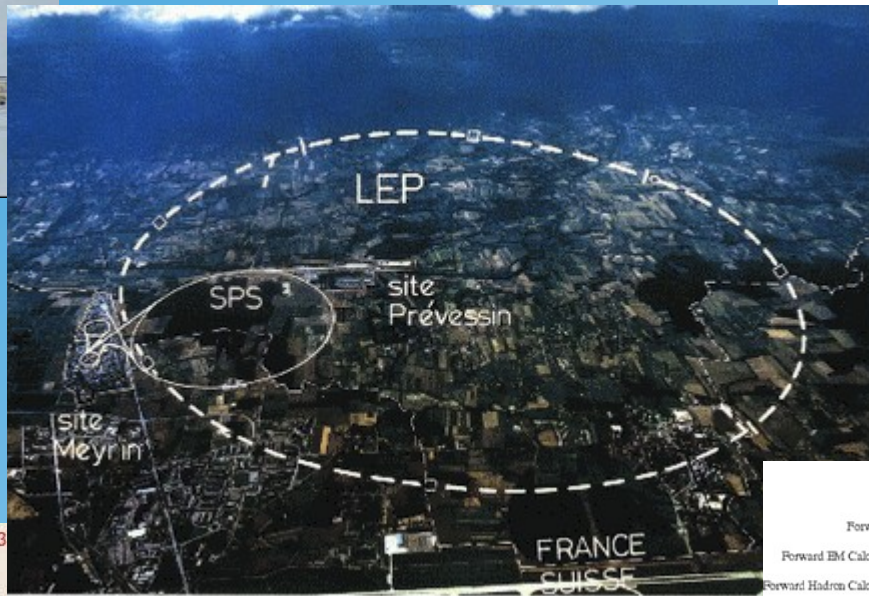
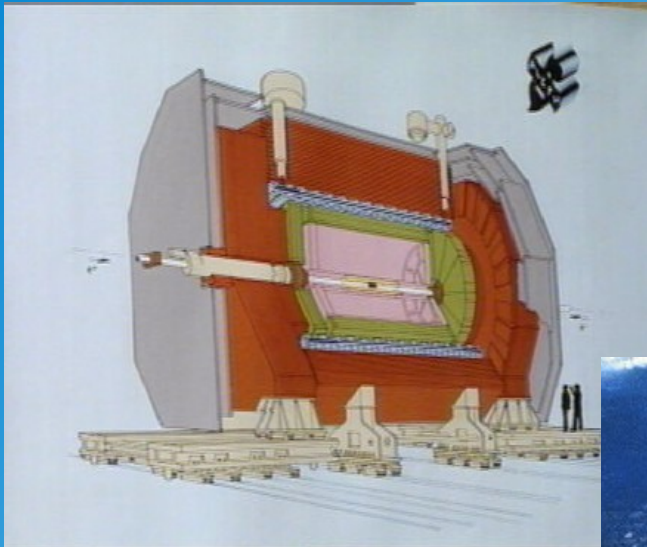
Final Result



Rho parameter consistent with 1 \rightarrow confirmation of the SM

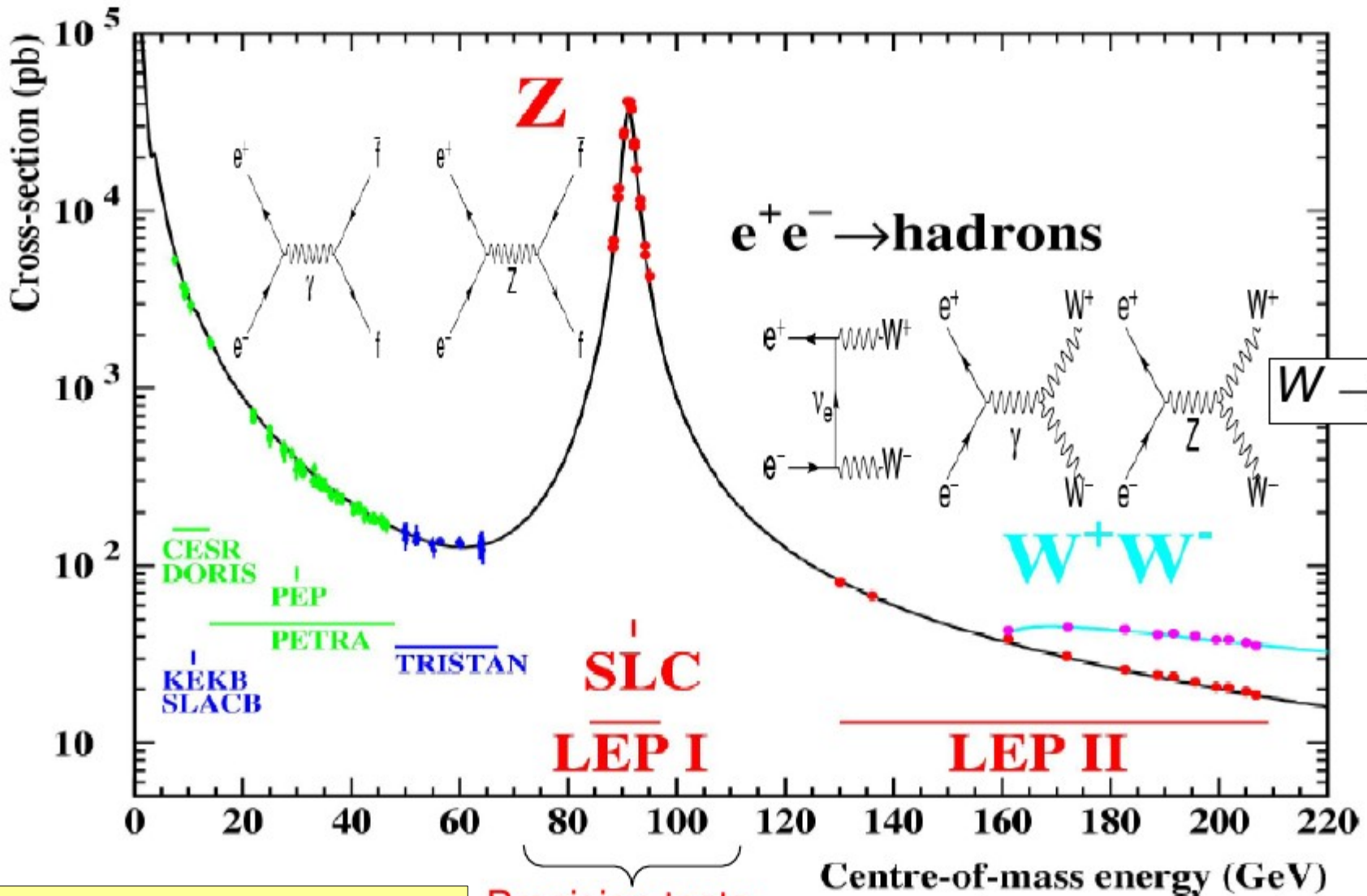
Nobel Prize for Physics 1984: C. Rubbia and S. van de Meer

Large Electron Positron Collider



e^+e^- collider
 $s^{1/2} = 90-200 \text{ GeV}$

Hadron Production in e^+e^-



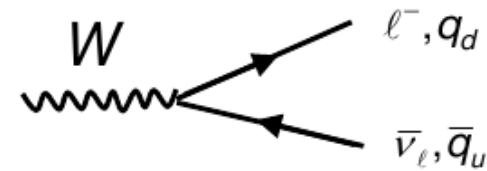
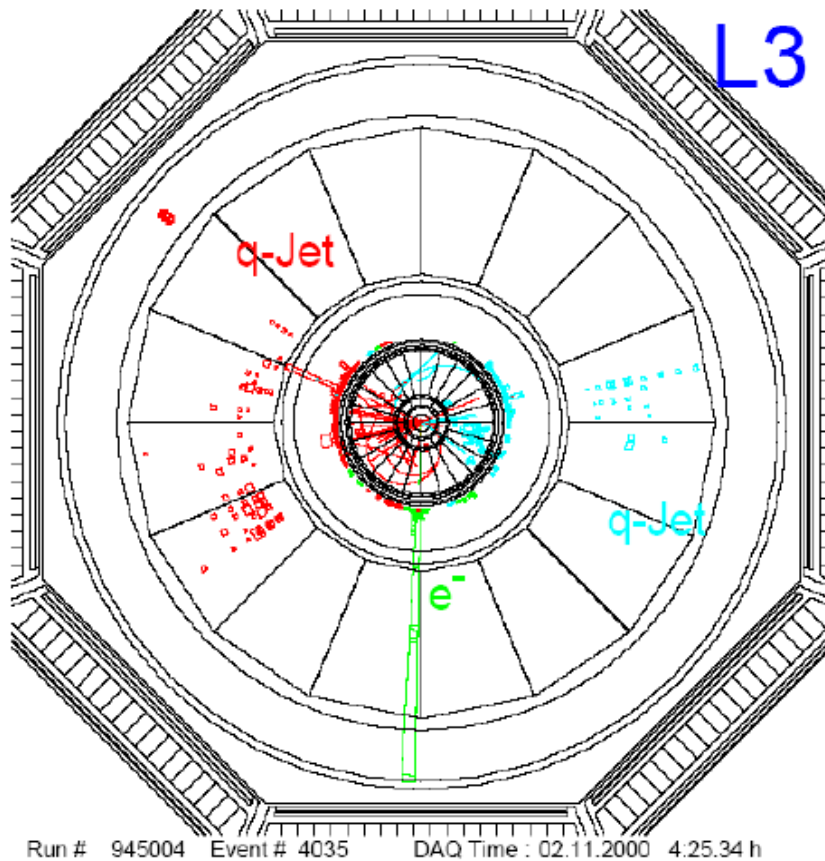
$M_Z = 91.1876 \pm 0.0021 \text{ GeV}$

Precision tests of the Z sector

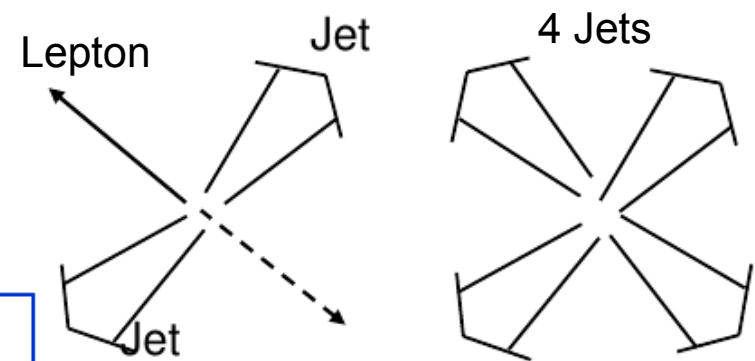
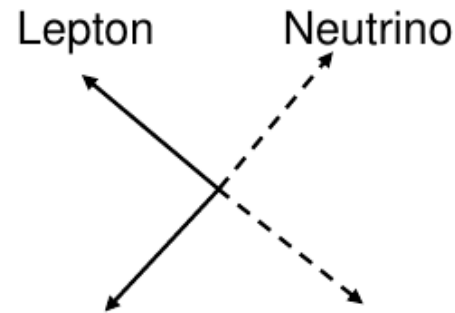
Tests of the W sector

WW Pair Production at LEP

W decays



WW → $\left\{ \begin{array}{l} qq\ell\nu \text{ 44\%} \\ qq\bar{q}q \text{ 45\%} \\ \ell\nu\ell\nu \text{ 11\%} \end{array} \right.$



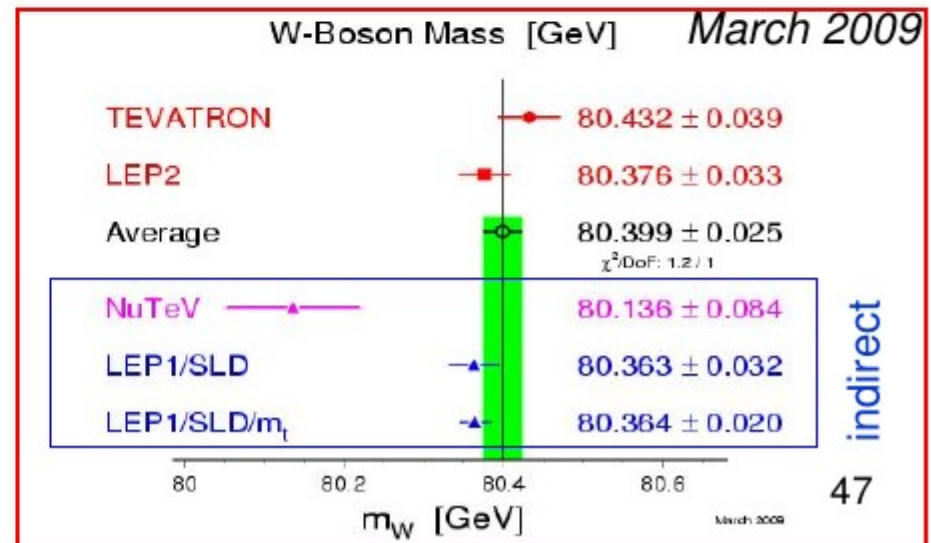
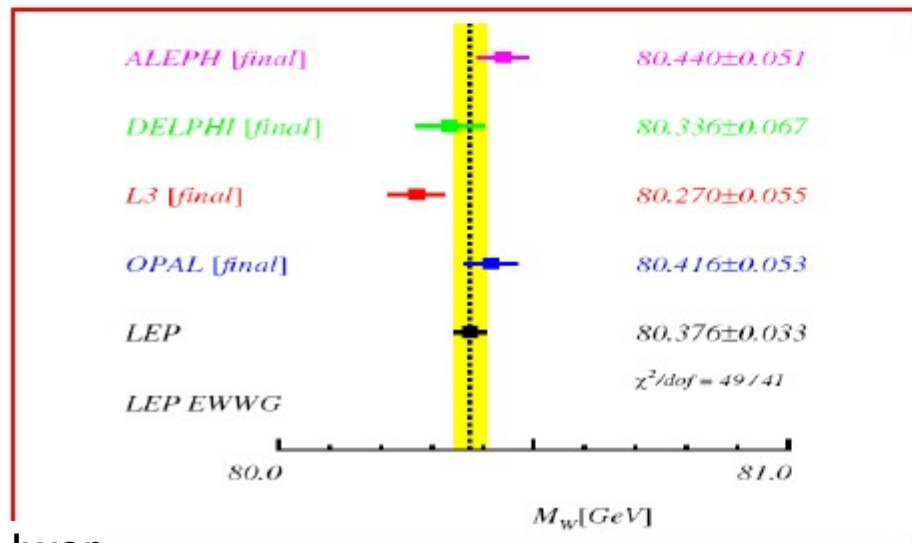
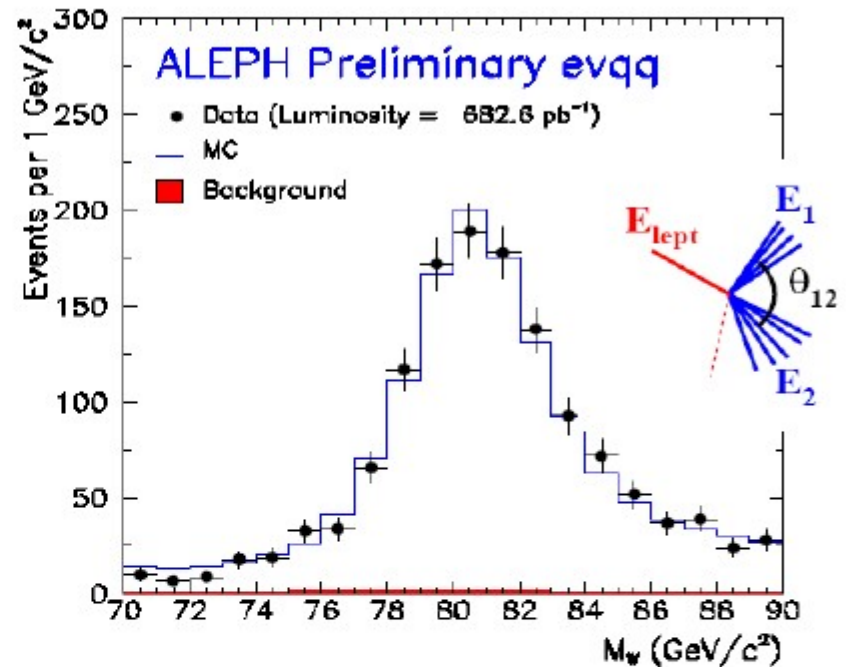
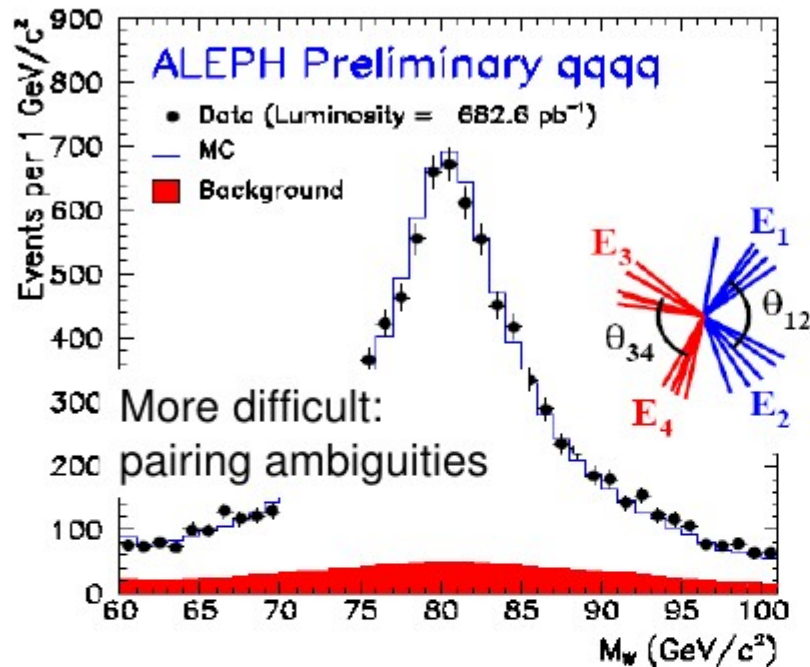
Easiest signature for a mass measurement:

$W_1 \rightarrow \ell\nu$ $W_2 \rightarrow \text{JetJet}$: use JetJet invariant mass

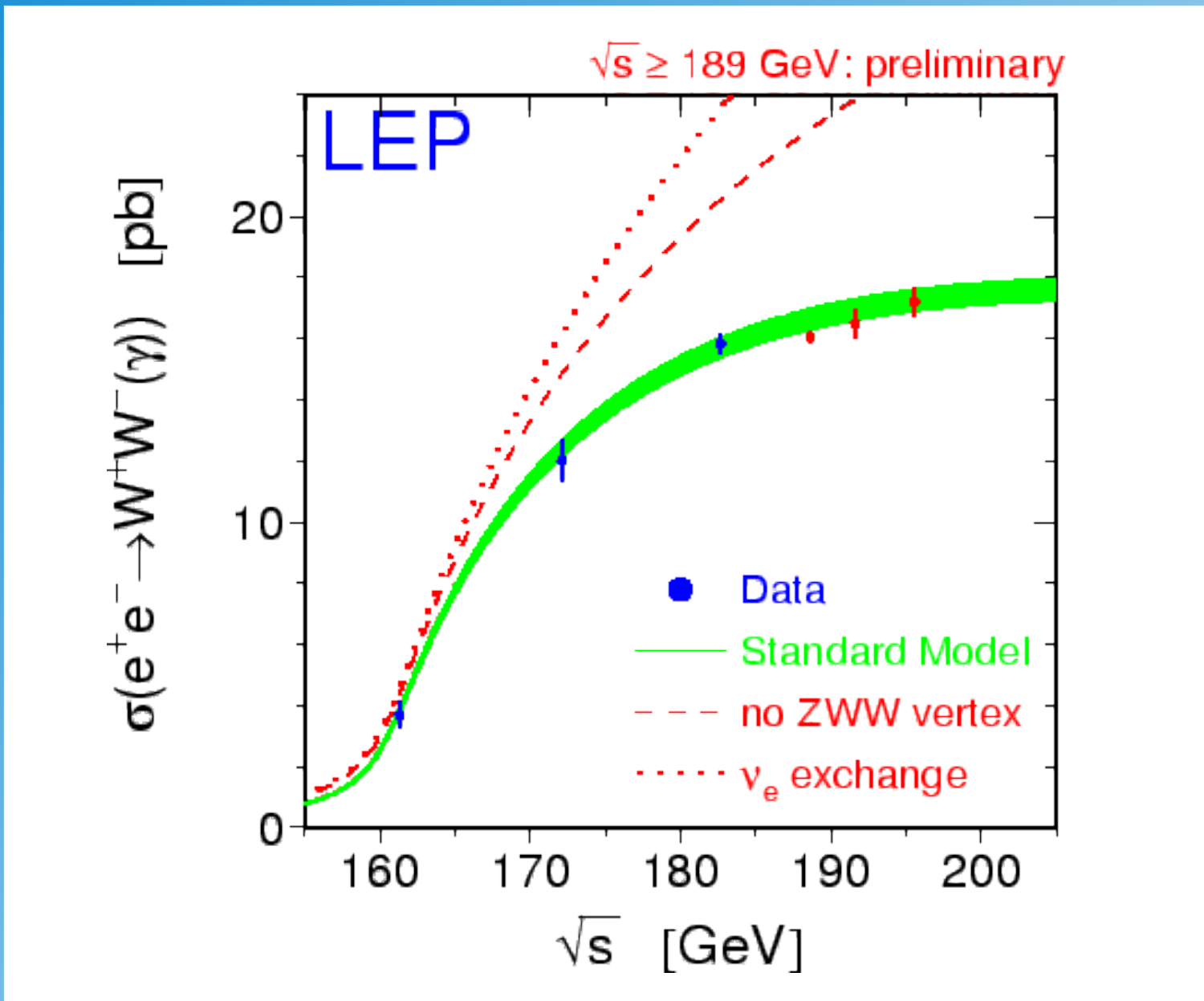
45

U.Uwer

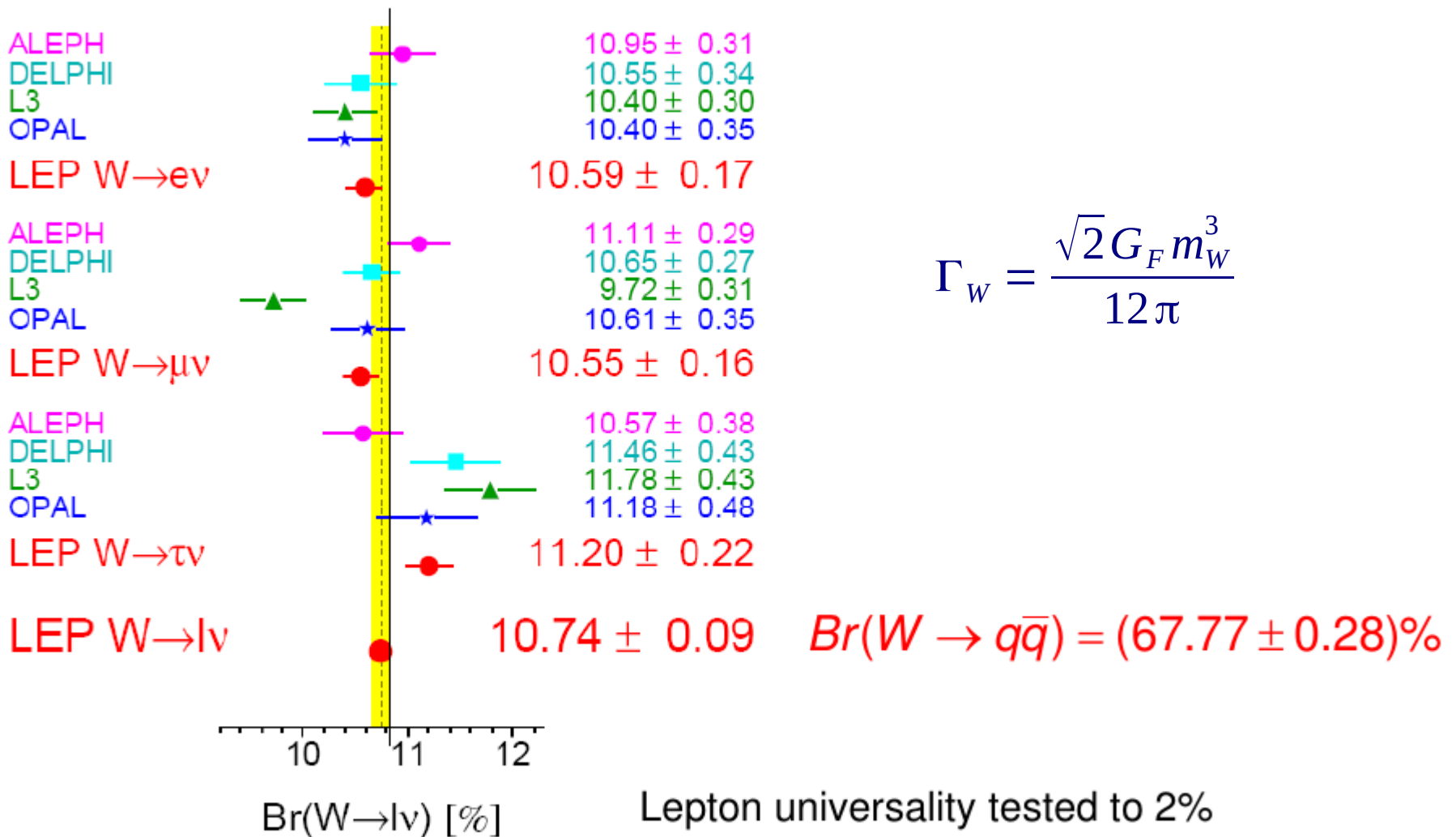
Invariant W mass reconstruction



W-Pair Production at LEP2



W leptonic branching fractions



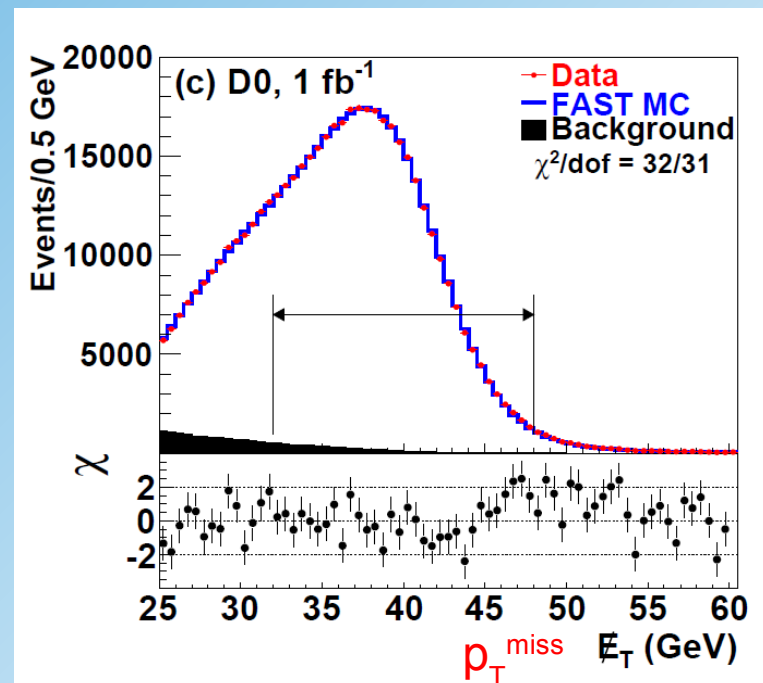
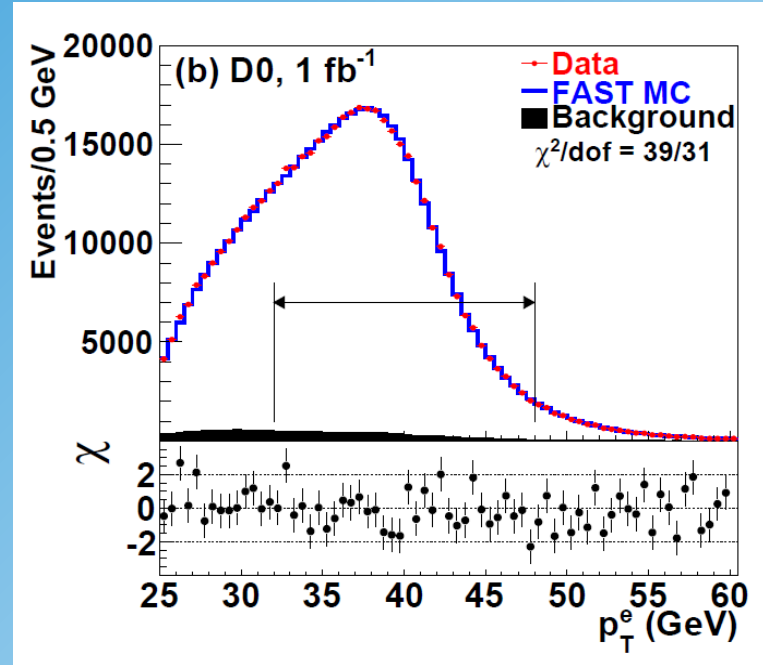
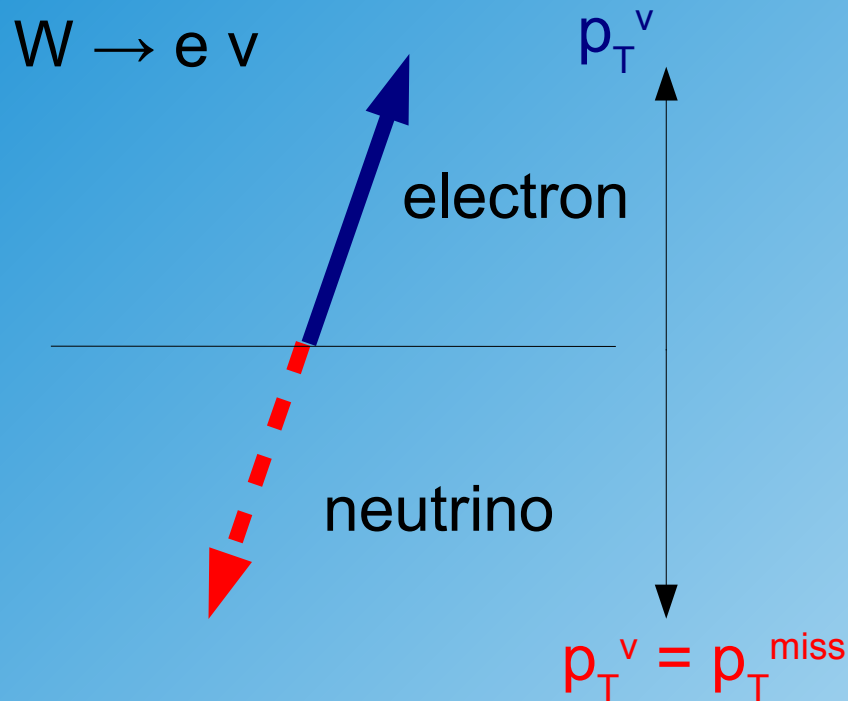
Tevatron at Fermilab



Proton – Antiproton Collider at $\sqrt{s} = 2 \text{ TeV}$

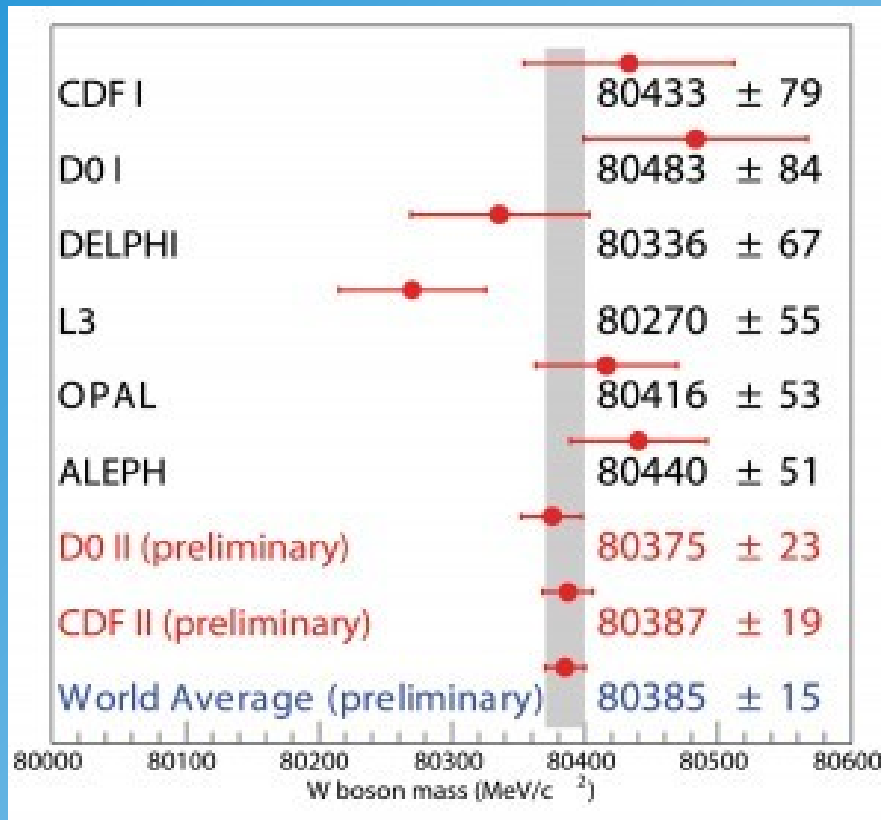
Missing Transverse Momentum

Jacobian peak at D0:



Latest Results W-mass

Method: normalise W-mass measurement to Z-mass measurement and take input (precise Z-mass) from LEP



} Tevatron Run 1

LEP2

} Tevatron Run 2

World Average

March 2012

W-mass measurement important for Top and Higgs Mass predictions

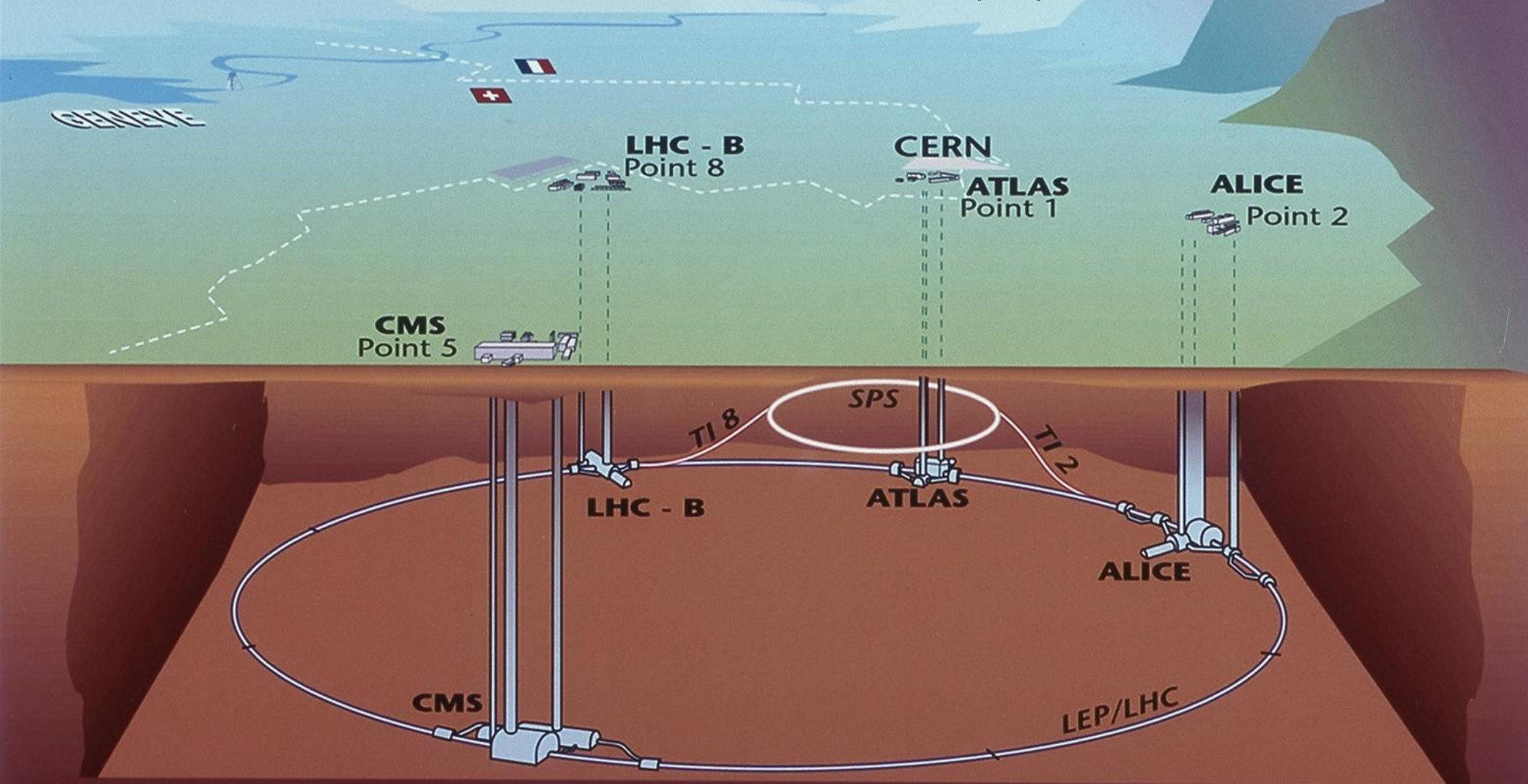
Overall view of the LHC experiments.

proton-proton collisions!

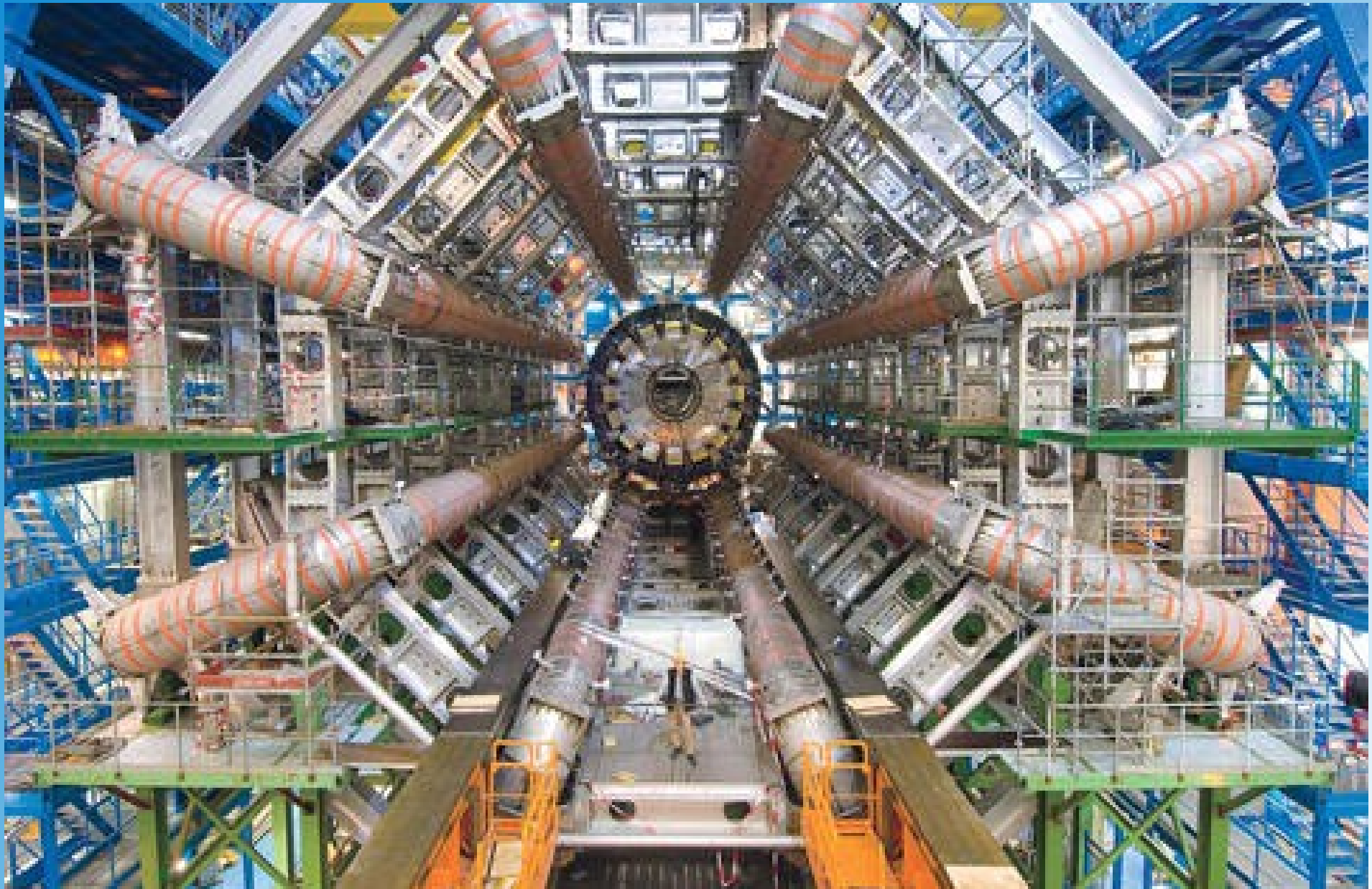
2011: $s^{1/2} = 7$ TeV

2012: $s^{1/2} = 8$ TeV

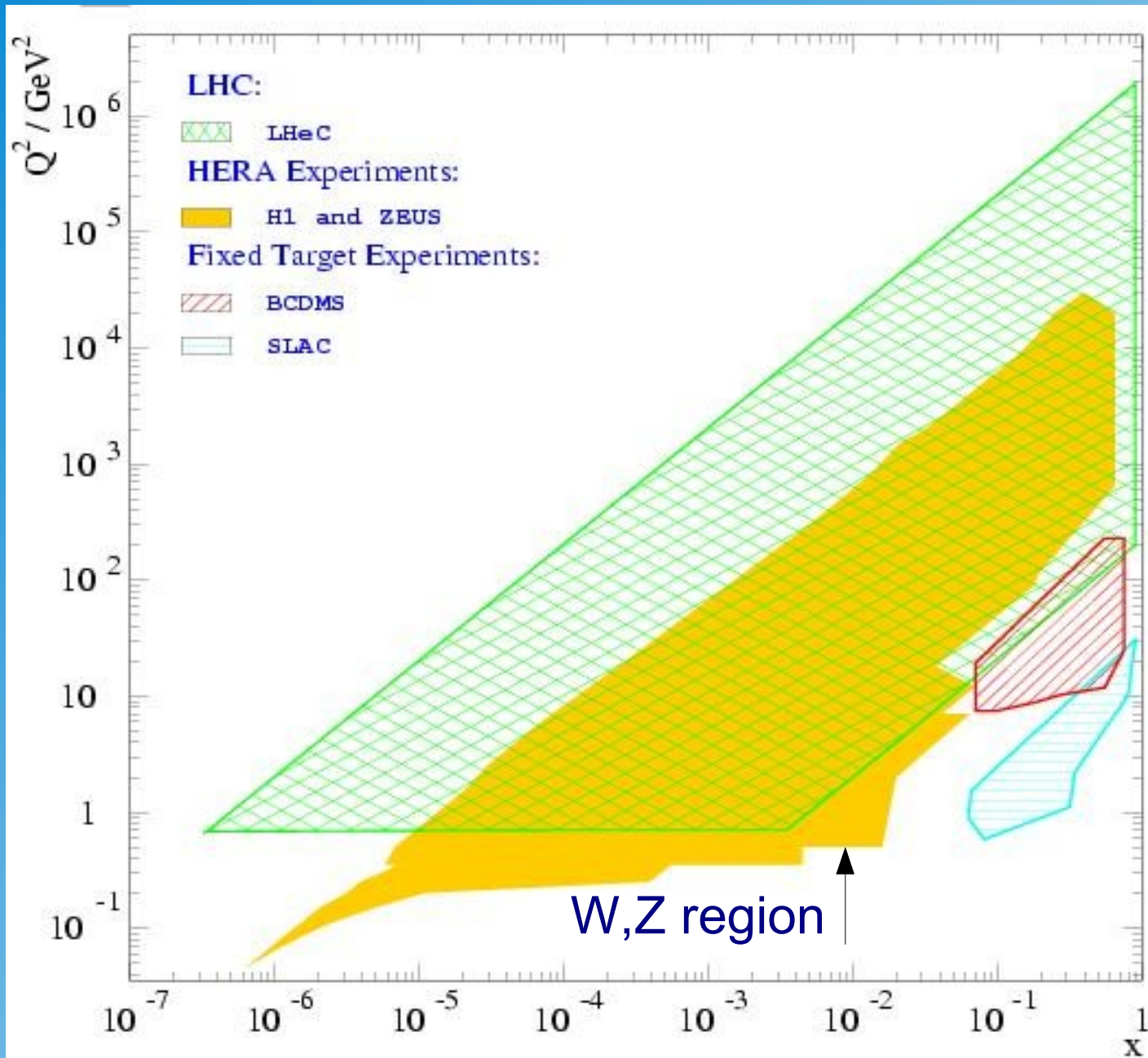
>2015: $s^{1/2} = 13$ (14) TeV



ATLAS Detector



LHC Kinematic Plane

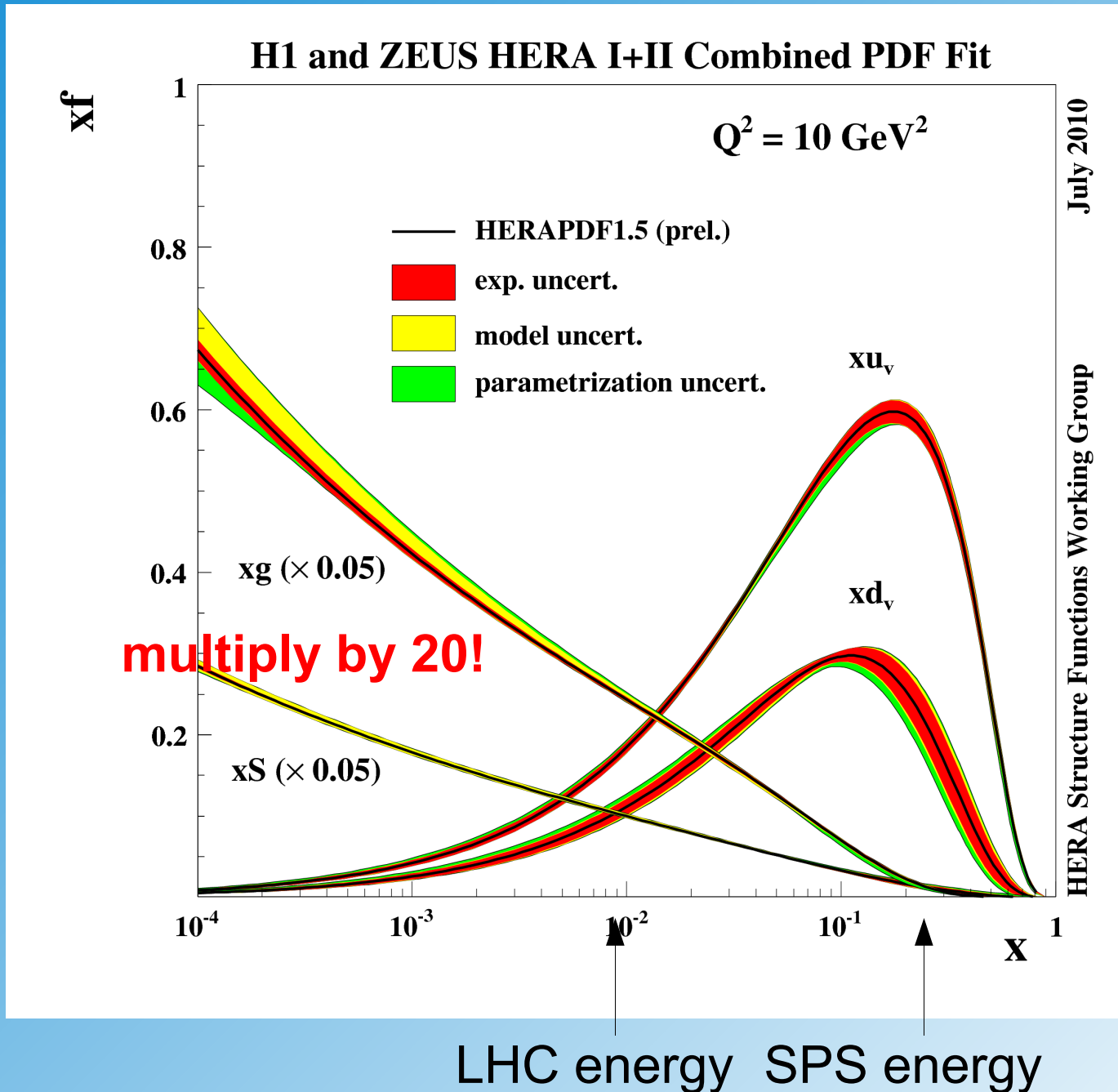


- W,Z production dominated by sea quarks

- low x-region very well constrained by HERA

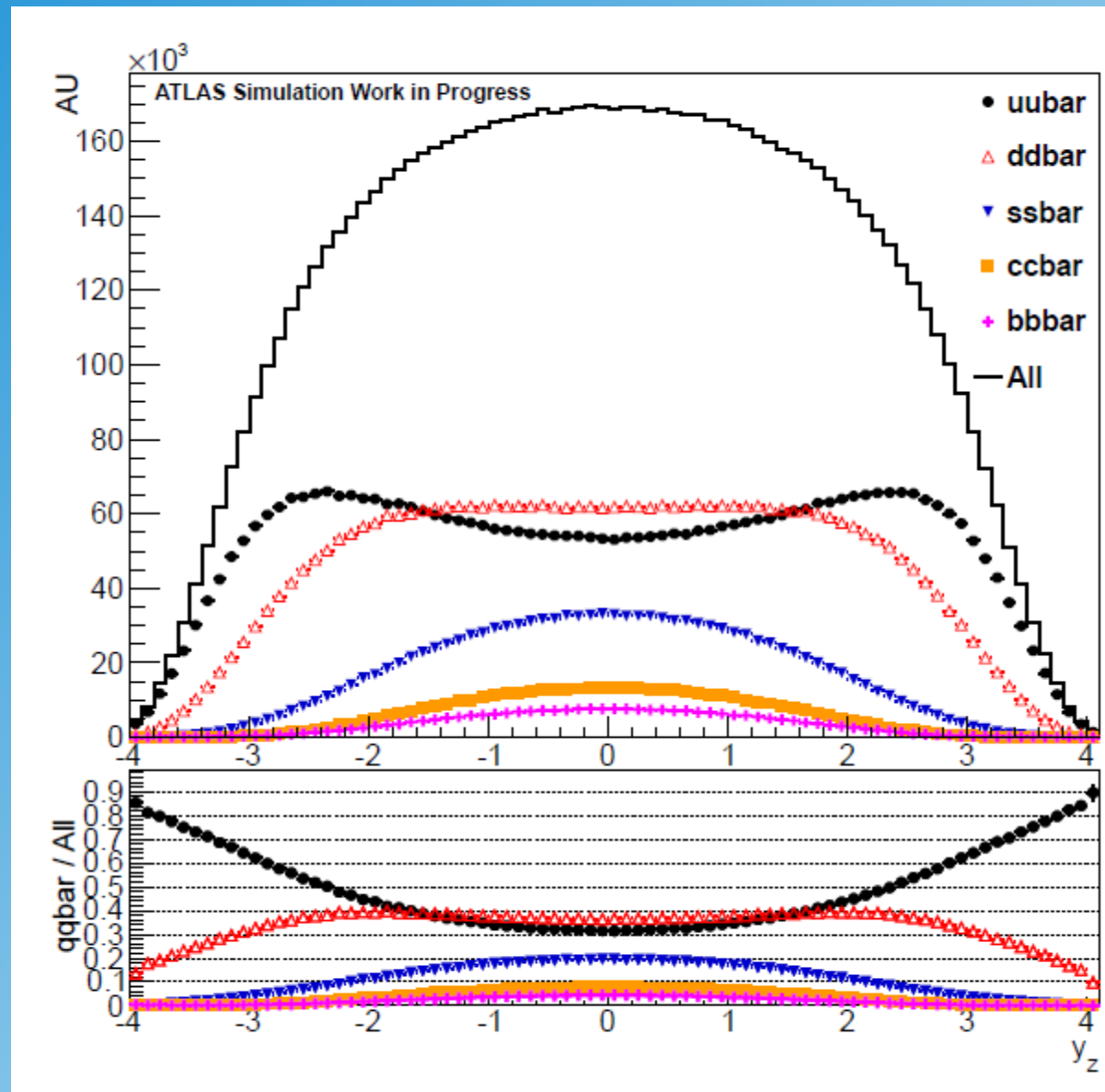
- W,Z production can be used to measure proton-PDFs and LHC luminosity!

Proton Parton Densities



Quark Flavors in Z Production

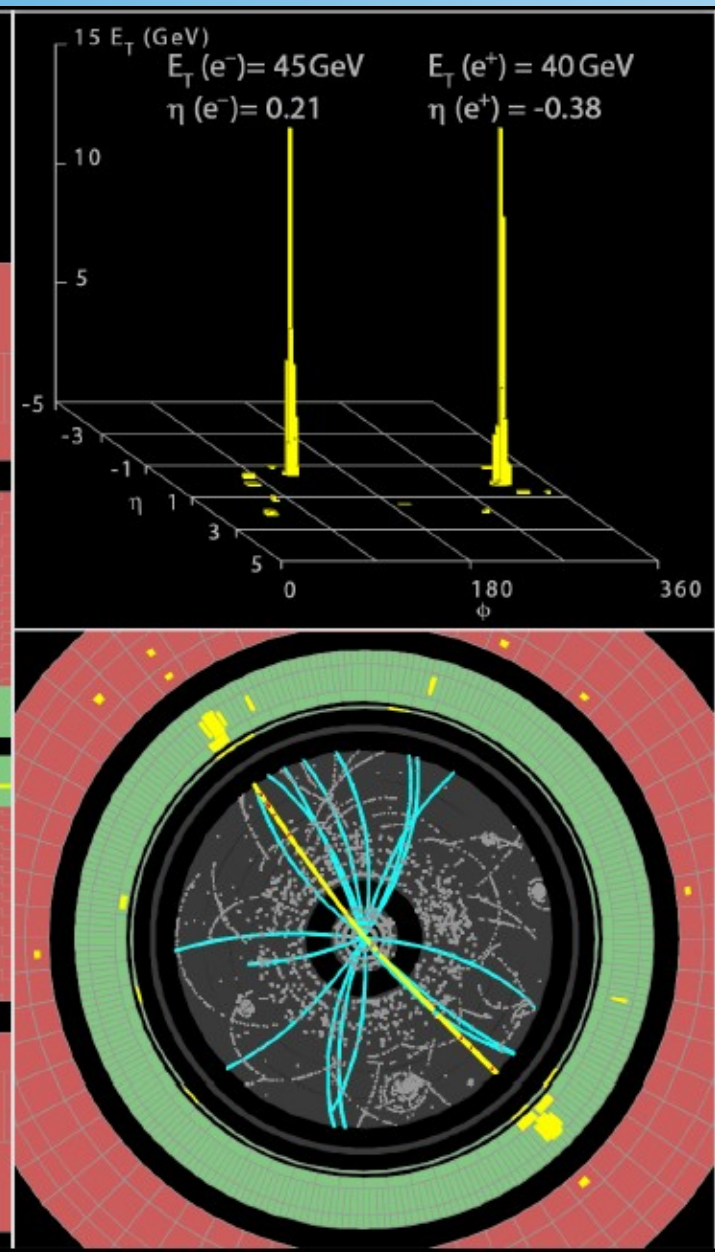
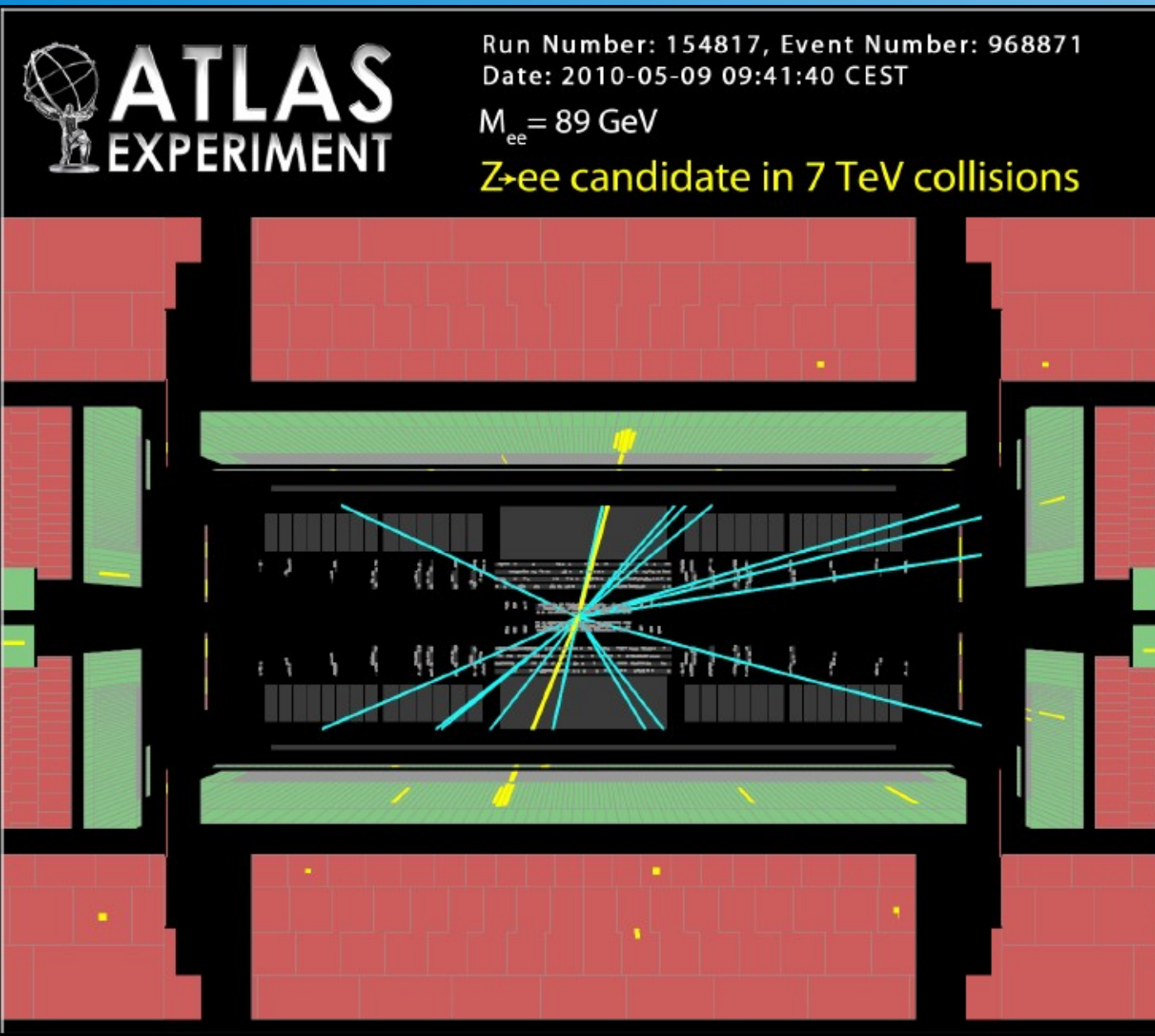
$q \bar{q} \rightarrow Z$



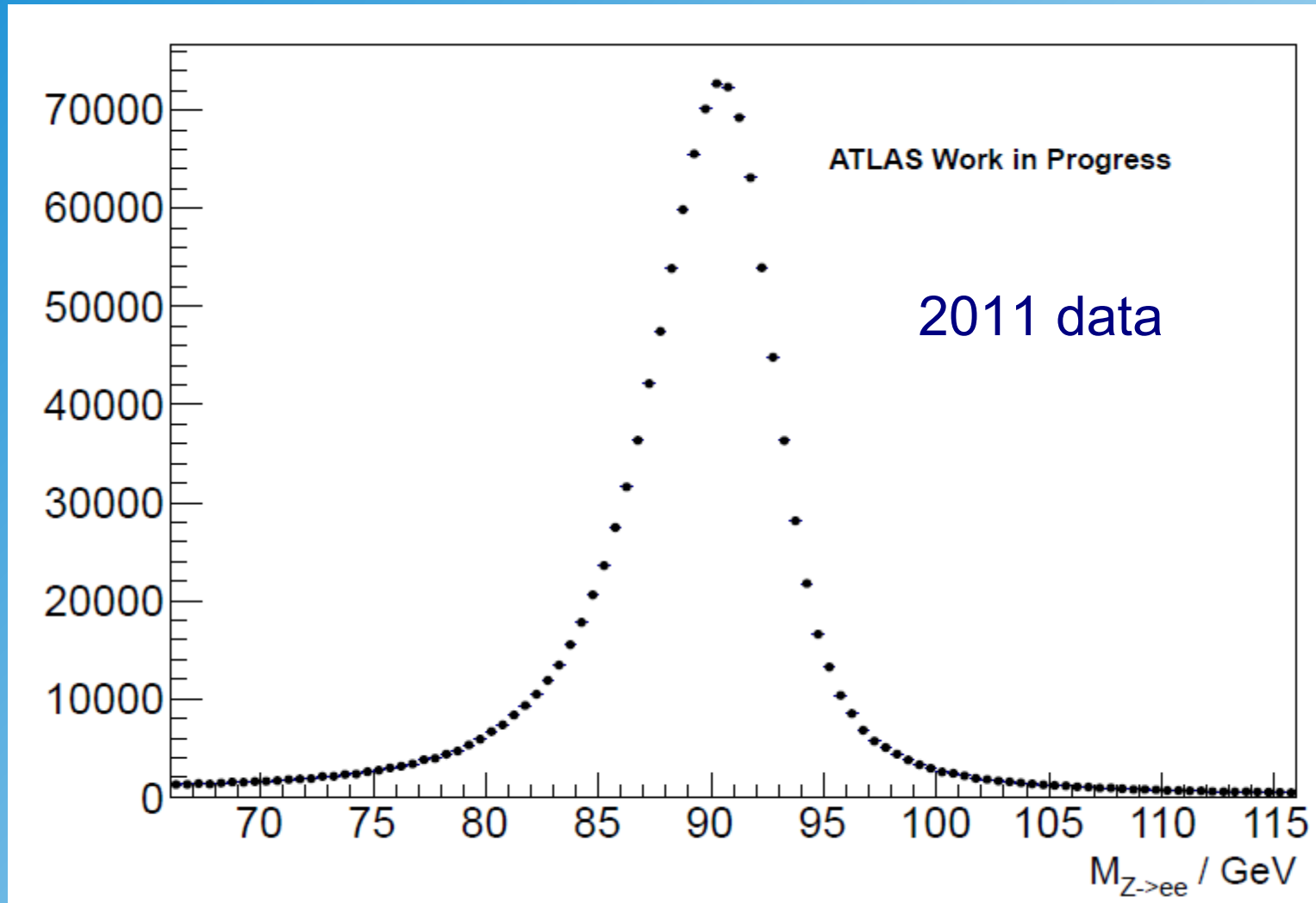
$y_z =$ pseudorapidity of Z-boson:

$$y = -\ln \tan \theta/2$$

$Z \rightarrow ee$ candidate at ATLAS



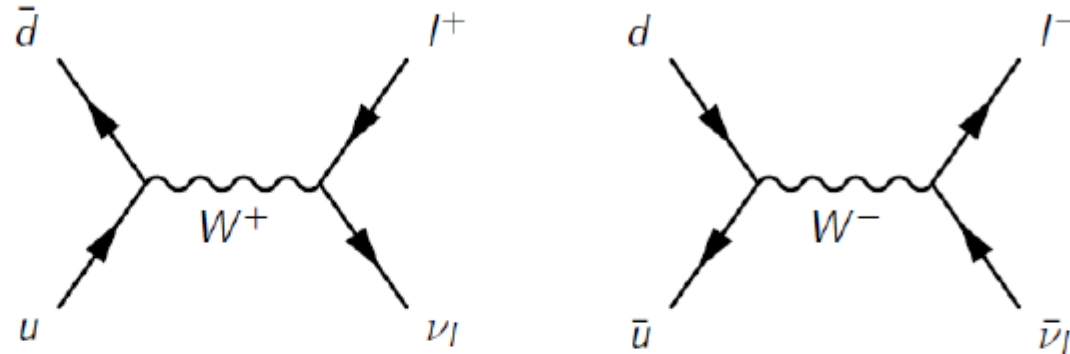
Z-Peak at ATLAS



LHC is a Vector-Boson factory!

W-Production at LHC

Valence quark +
sea quark



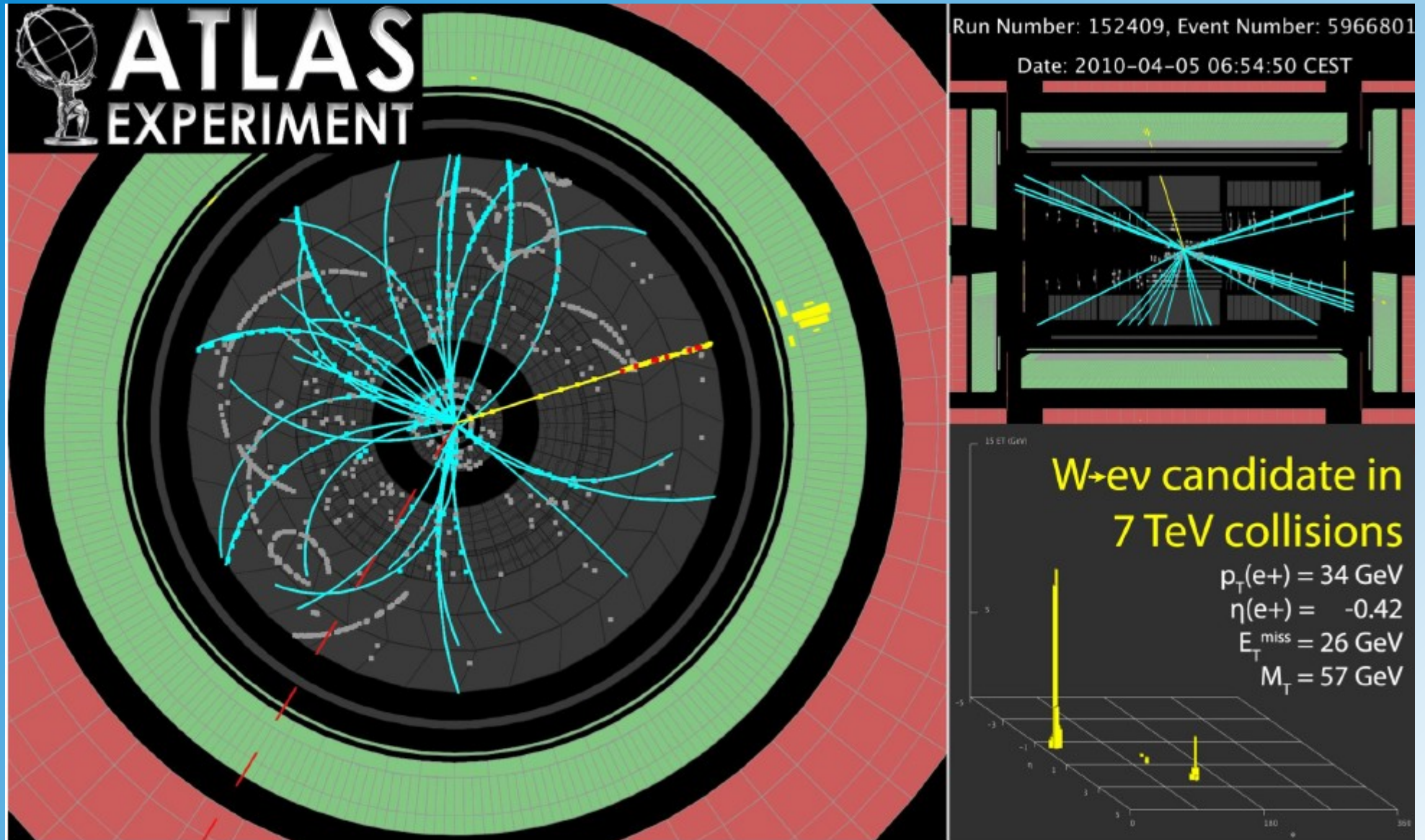
valence quark ratio $u/d = 2 \Rightarrow$ more W^+ than W^-

ATLAS 2010:

[nb]	Data
W^+	$6.257 \pm 0.017(\text{sta}) \pm 0.152(\text{sys}) \pm 0.213(\text{lum}) \pm 0.188(\text{acc})$
W^-	$4.149 \pm 0.014(\text{sta}) \pm 0.102(\text{sys}) \pm 0.141(\text{lum}) \pm 0.124(\text{acc})$
W	$10.391 \pm 0.022(\text{sta}) \pm 0.238(\text{sys}) \pm 0.353(\text{lum}) \pm 0.312(\text{acc})$

Charge Asymmetric! Tool to disentangle d and u valence quarks

W-boson Production at LHC

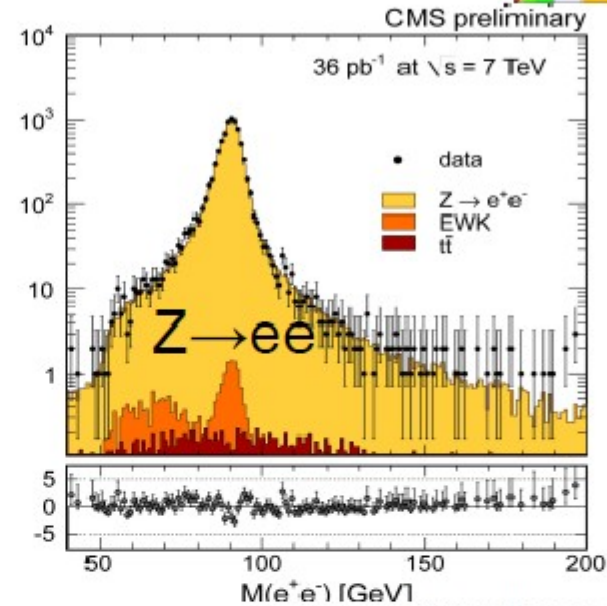
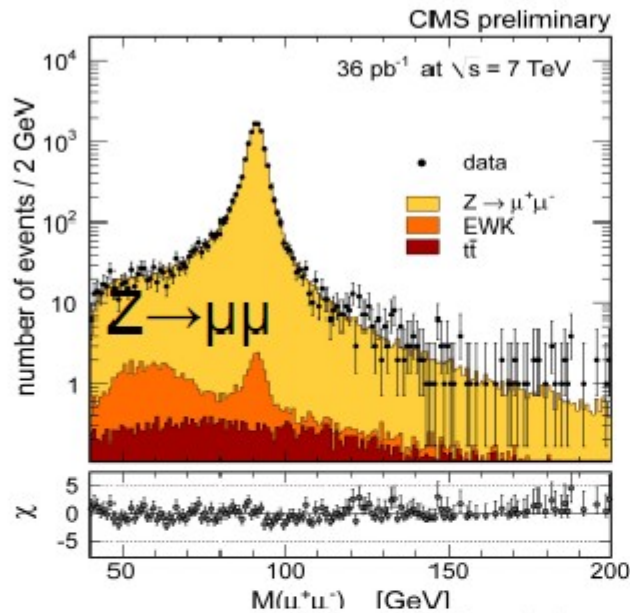


$$u \bar{d} \rightarrow W^+ \rightarrow e^+ \nu_e$$

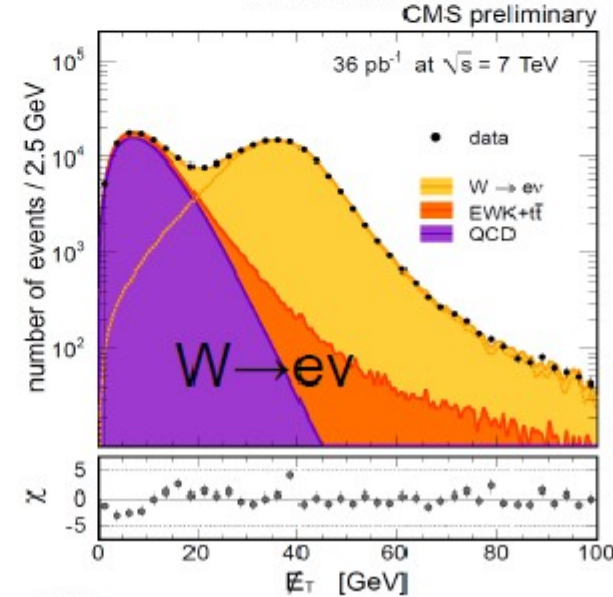
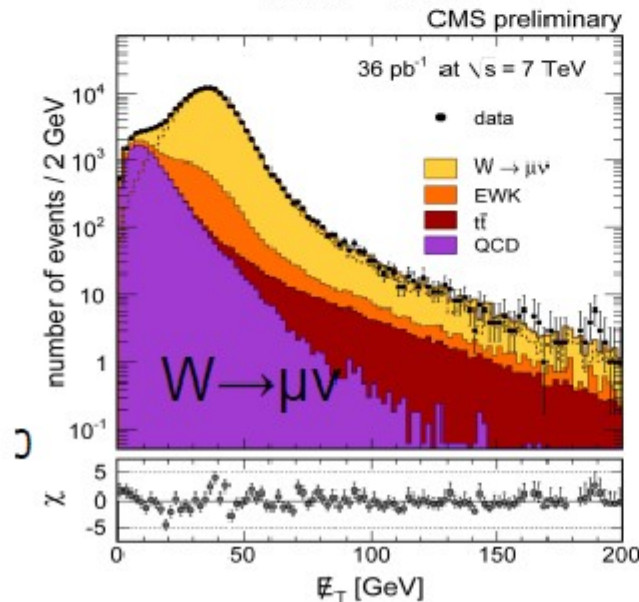
Anti-quarks from the sea!

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Z and W production at LHC

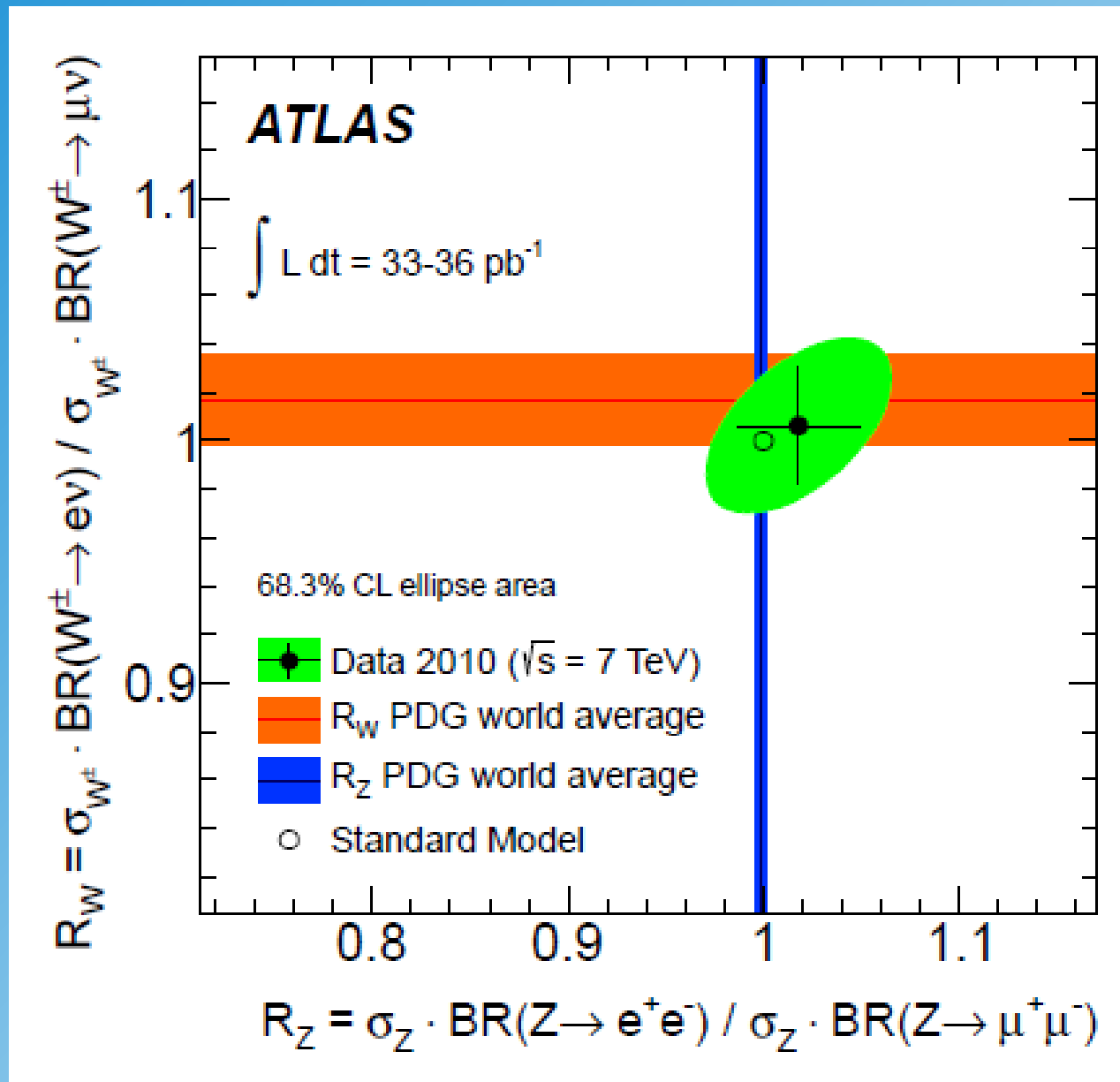


*P.Harris,
Moriond 2011*



Instead of E_{eT}
use E_T (i.e. $E_{\nu T}$)

Lepton Universality Check at LHC



Summary

- W , Z boson discovered in 1983
- W , Z masses consistent with SM predictions
- Ratio of W and Z mass consistent with Weinberg angle measured in Neutral Currents
- Lepton universality tested in W , Z Decays
- $W^+ W^-$ pair production cross section measured. Confirmation of triple gauge couplings (WWZ)
- W and Z mass relevant for Higgs mass predictions

