

Models for neutrino mass

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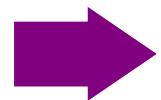
Outline

- lepton mixing
- tri-bimaximal ansatz (TBM)
- neutrino mass hierarchy: absolute scale & Onubb
- Majorana mass and seesaw mechanisms
- neutrino mass matrix and TBM
- discrete groups
- a prototype A4 model
- A4 breaking & vacuum alignment problem
- an A4 model with accidental TBM
- conclusions

Lepton mixing

$$\mathcal{L} = ig 2^{-1/2} W_\mu^- \sum_{a=1}^n \bar{E}_{aL} \gamma_\mu \rho_{aL} + \text{H.c.}$$

$$= ig 2^{-1/2} W_\mu^- \sum_{a,b,\alpha} \bar{e}_{bL} \gamma_\mu \Omega_{ab}^* U_{a\alpha} \nu_{\alpha L} + \text{H.c.}$$



$$K_{b\alpha} = \sum_{c=1}^n (\Omega^\dagger)_{bc} U_{c\alpha}$$

$$\begin{matrix} \nu_e & c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ \nu_\mu & -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ \nu_\tau & s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{matrix} \times \text{diag}(e^{i\alpha_1/2}, e^{i\alpha_2/2}, 1)$$

Lepton mixing

parameter	Ref. [1]		Ref. [2] (MINOS updated)	
	best fit $\pm 1\sigma$	3σ interval	best fit $\pm 1\sigma$	3σ interval
Δm_{21}^2 [10^{-5} eV 2]	$7.65^{+0.23}_{-0.20}$	7.05–8.34	$7.67^{+0.22}_{-0.21}$	7.07–8.34
Δm_{31}^2 [10^{-3} eV 2]	$\pm 2.40^{+0.12}_{-0.11}$	$\pm (2.07\text{--}2.75)$	-2.39 ± 0.12 $+2.49 \pm 0.12$	$-(2.02\text{--}2.79)$ $+(2.13\text{--}2.88)$
$\sin^2 \theta_{12}$	$0.304^{+0.022}_{-0.016}$	0.25–0.37	$0.321^{+0.023}_{-0.022}$	0.26–0.40
$\sin^2 \theta_{23}$	$0.50^{+0.07}_{-0.06}$	0.36–0.67	$0.47^{+0.07}_{-0.06}$	0.33–0.64
$\sin^2 \theta_{13}$	$0.01^{+0.016}_{-0.011}$	≤ 0.056	0.003 ± 0.015	≤ 0.049

Schwetz, Tortola, Valle,NJP10 (08')

Gonzalez,Maltoni, PR460 (08')

Table 1: Global 3ν oscillation analysis (2008): best-fit values and allowed n_σ ranges, from Ref. [4].

Parameter	$\delta m^2/10^{-5}$ eV 2	$\sin^2 \theta_{12}$	$\sin^2 \theta_{13}$	$\sin^2 \theta_{23}$	$\Delta m^2/10^{-3}$ eV 2
Best fit	7.67	0.312	0.016	0.466	2.39
1σ range	7.48 – 7.83	0.294 – 0.331	0.006 – 0.026	0.408 – 0.539	2.31 – 2.50
2σ range	7.31 – 8.01	0.278 – 0.352	< 0.036	0.366 – 0.602	2.19 – 2.66
3σ range	7.14 – 8.19	0.263 – 0.375	< 0.046	0.331 – 0.644	2.06 – 2.81

Fogli,Lisi,Marrone,Palazzo, PPNP57(06')

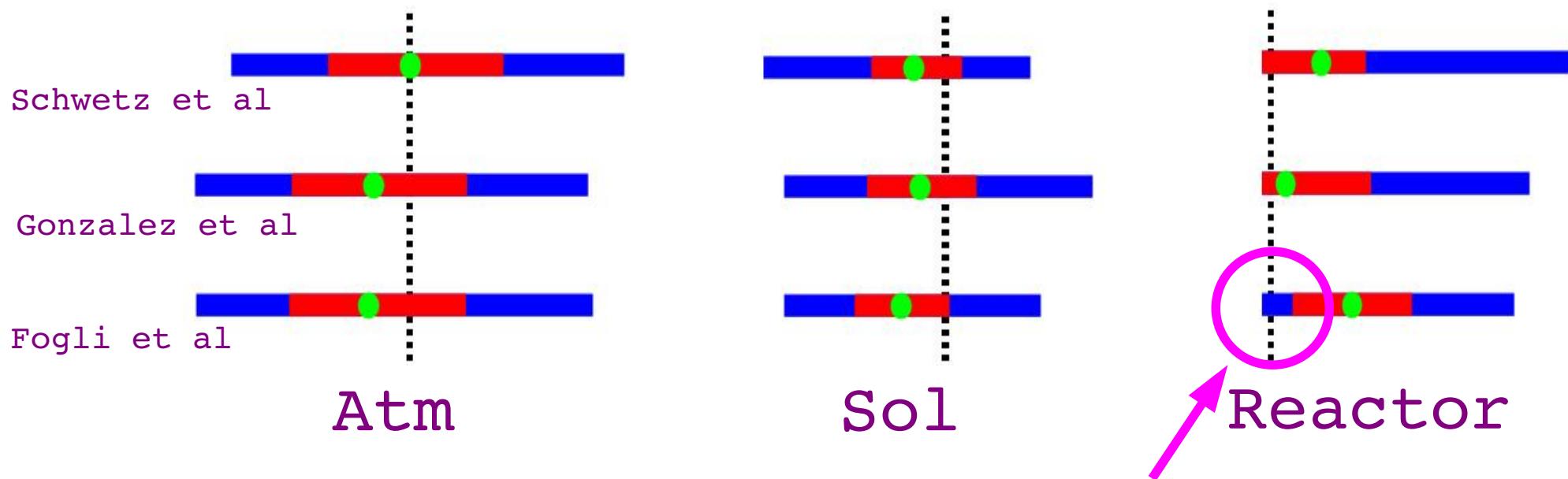
Tri-Bimaximal ansatz

Harrison, Perkins, Scott, 2002

$$\sin^2 \theta_{23} = 0.5$$

$$\sin^2 \theta_{12} = 1/3$$

$$\sin^2 \theta_{13} = 0$$



with this exception
tri-bimaximal is in good agreement within one sigma

Tri-Bimaximal mixing

$$U_{\text{HPS}} = \begin{pmatrix} \sqrt{2/3} & 1/\sqrt{3} & 0 \\ -1/\sqrt{6} & 1/\sqrt{3} & -1/\sqrt{2} \\ -1/\sqrt{6} & 1/\sqrt{3} & 1/\sqrt{2} \end{pmatrix}$$



trimaximal

$$\nu_2 = \frac{1}{\sqrt{3}}(\nu_e + \nu_\mu + \nu_\tau)$$

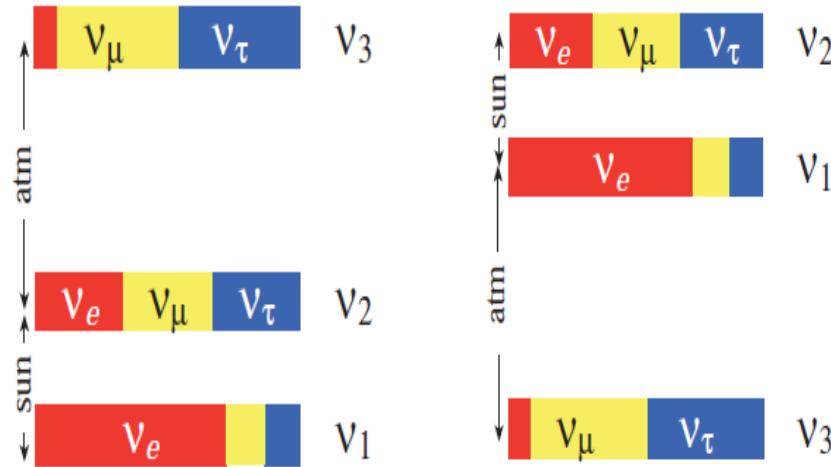


Bimaximal

$$\nu_3 = \frac{1}{\sqrt{2}}(-\nu_\mu + \nu_\tau)$$

- ◆ mu-tau
- ◆ trimaximal
- ◆ tetramaximal
- ◆ symmetric mixing
- ◆ bimaximal
- ◆ hexagon mixing
- ◆ golden
- ◆ quark-lepton complementarity

mass hierarchies



parameter	best fit	2σ	3σ
$\Delta m_{21}^2 [10^{-5}\text{eV}^2]$	$7.59^{+0.23}_{-0.18}$	7.22–8.03	7.03–8.27
$ \Delta m_{31}^2 [10^{-3}\text{eV}^2]$	$2.40^{+0.12}_{-0.11}$	2.18–2.64	2.07–2.75

absolute value

NH

IH

$$\lambda_C < \frac{m_2^\nu}{m_3^\nu} < 1$$

$$0 < \frac{m_1^\nu}{m_2^\nu} < 1$$



Absolute neutrino mass scale
unknown

charged fermions

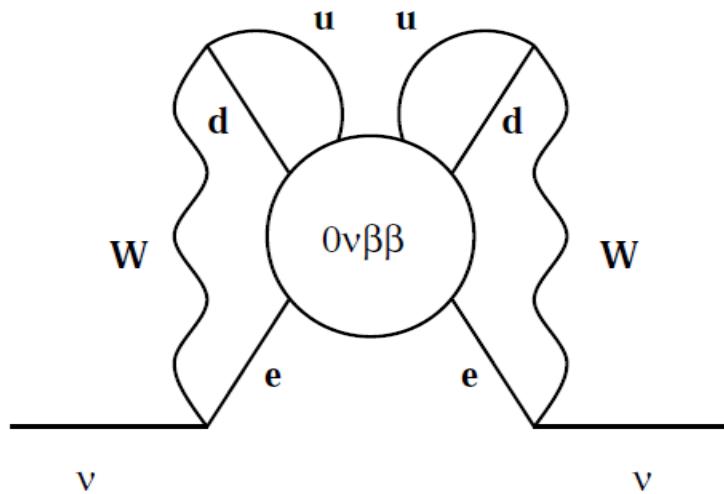
$$\frac{m_2^u}{m_3^u} \approx \lambda_C^4$$

$$\frac{m_1^u}{m_2^u} \approx \lambda_C^3$$

$$\frac{m_2^{d,l}}{m_3^{d,l}} \approx \lambda_C^2$$

$$\frac{m_1^{d,l}}{m_2^{d,l}} \approx \lambda_C^2$$

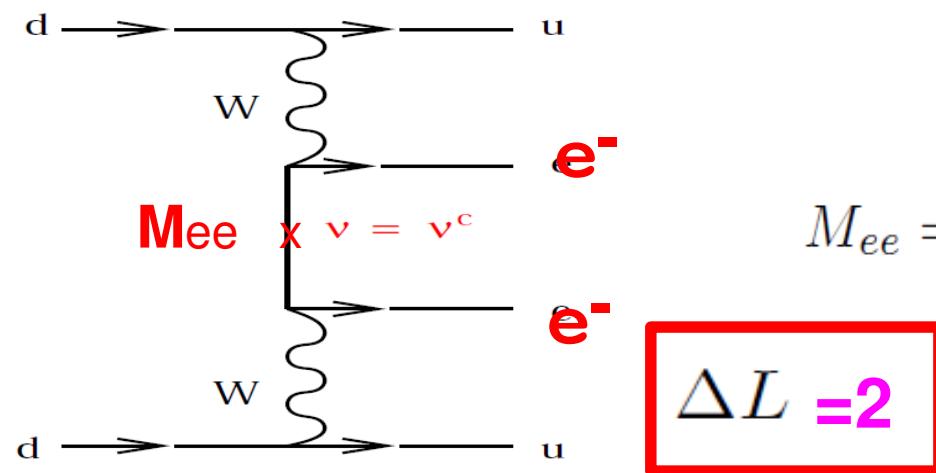
absolute neutrino mass scale



BLACK BOX THEOREM:

If neutrinoless double beta decay is observed, neutrino has Majorana mass

Schechter, Valle PRD25 (82')

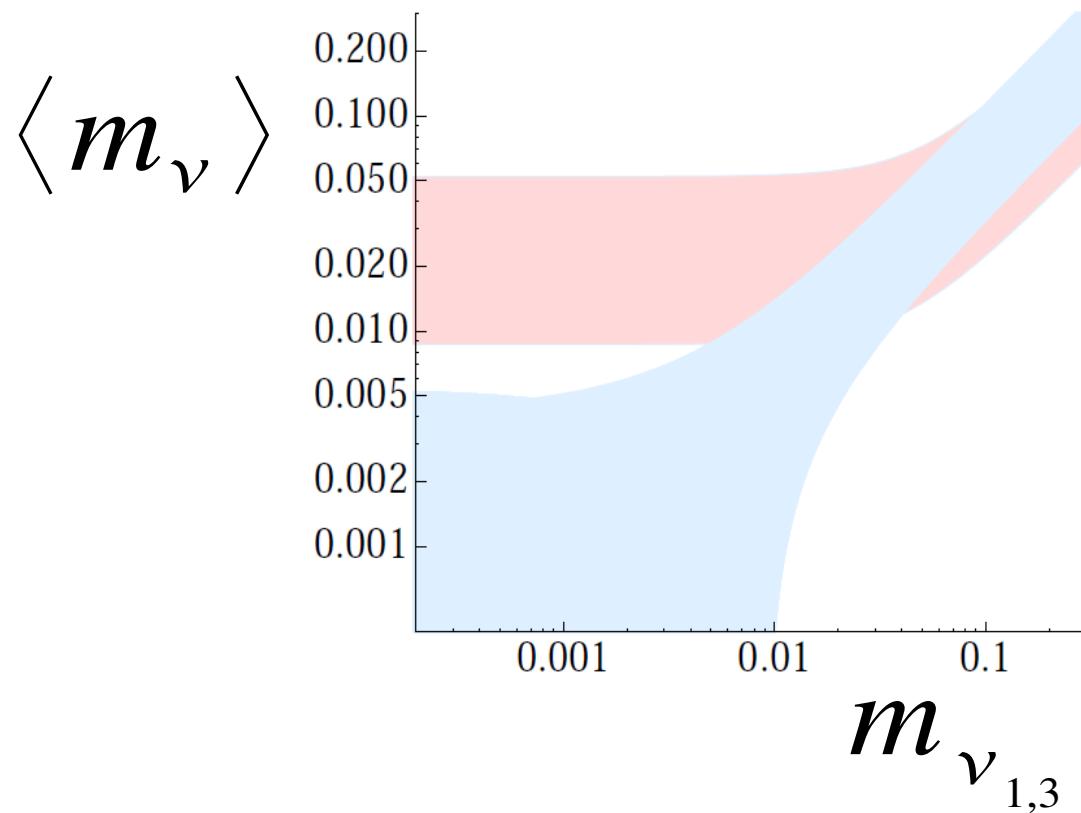


$$M_{ee} = |m_1 U_{e1}^2 + m_2 U_{e2}^2 + m_3 U_{e3}^2|$$

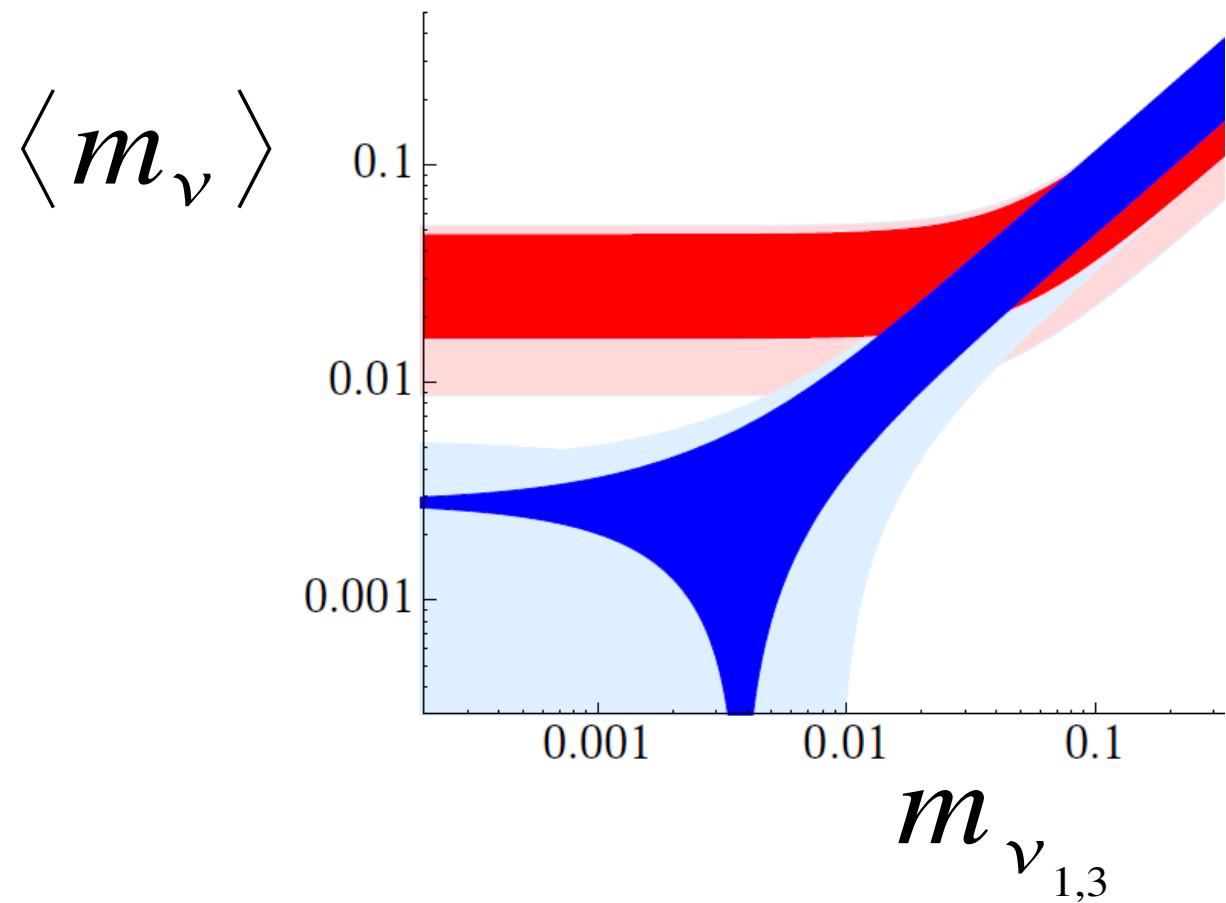
$$\Delta L = 2$$

$$\langle m_\nu \rangle = c_\odot^2 c_R^2 m_{\nu_1} + s_\odot^2 c_R^2 e^{i\alpha} \sqrt{m_{\nu_1}^2 + \Delta m_\odot^2} + s_R^2 e^{i\beta} \sqrt{m_{\nu_1}^2 + \Delta m_\odot^2 + \Delta m_{\text{Atm}}^2}, \quad \text{normal hierarchy}$$

$$\langle m_\nu \rangle = c_\odot^2 c_R^2 \sqrt{m_{\nu_3}^2 - \Delta m_\odot^2 + \Delta m_{\text{Atm}}^2} + s_\odot^2 c_R^2 e^{i\alpha} \sqrt{m_{\nu_3}^2 + \Delta m_{\text{Atm}}^2} + s_R^2 e^{i\beta} m_{\nu_3} \quad \text{inverse hierarchy}$$

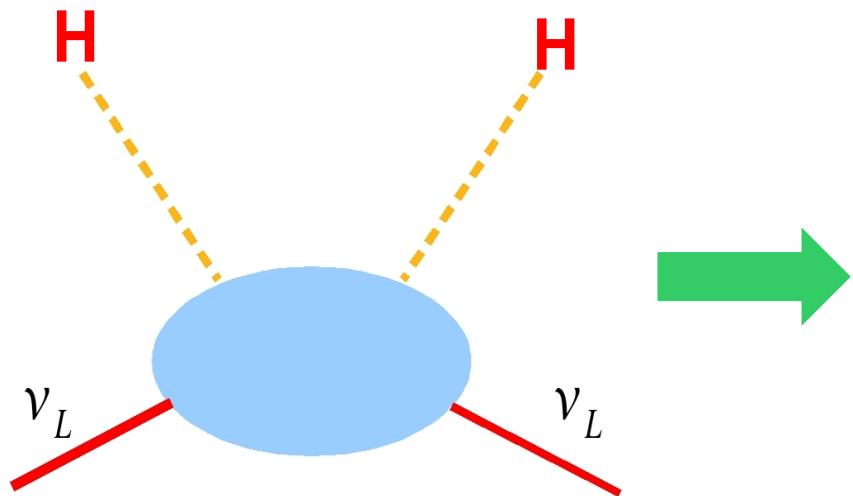


Onubb & TBM



Majorana mass

If lepton number is violated



$$LH LH / \mathcal{M}$$

Dim-5 operator, Weinberg (80)

There are many seesaw realizations:
type-I, II, III, linear, inverse seesaw,
R parity breaking, radiatively

Seesaw type-I & II

In the basis ν_L, ν_L^c

$$M_\nu = \begin{pmatrix} M_L & M_D \\ M_D^T & M_N \end{pmatrix}$$

Minkowski 77
Yanagita 79
Gell-Mann, Ramond, Slansky 79
Mohapatra, Senjanovic 80

$m_\nu = -M_D M_N^{-1} M_D^T$

Schechter, Valle 80, 82
Cheng, Li 80
Mohapatra, Senjanovic 81

type-III

Similar to type-I:
right-handed neutrino replaced
with a SU(2)-triplet $Y=0$

$$\Sigma = \begin{pmatrix} \Sigma^0/\sqrt{2} & \Sigma^+ \\ \Sigma^- & -\Sigma^0/\sqrt{2} \end{pmatrix}$$

Foot et al 88

Ma, 98

Ma,Roy, 02

Bajc,Senjanovic, 07

Abada et al, 07

$$M_\nu = \begin{pmatrix} 0 & M_D \\ M_D^T & M_\Sigma \end{pmatrix} \rightarrow m_\nu = -M_D^T M_\Sigma^{-1} M_D$$

Inverse seesaw

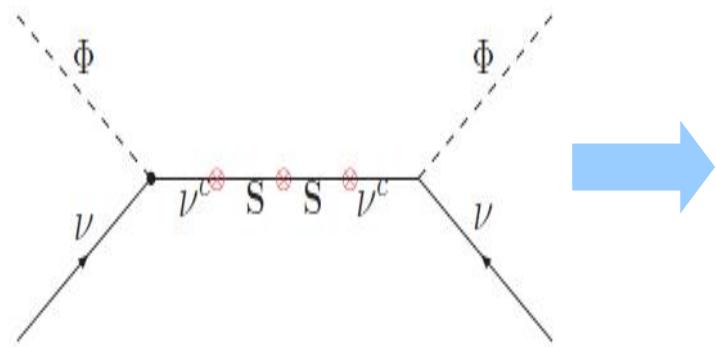
Mohapatra, Valle PRD34

Extra singlets **S**

$$M_\nu = \begin{pmatrix} \nu, \nu^c, S & \\ 0 & M_D & 0 \\ M_D^T & 0 & M \\ 0 & M^T & \mu \end{pmatrix}$$

μ Breaks lepton number

M at EW or TeV scale
 μ at KeV scale



$$m_\nu = M_D M^{T^{-1}} \mu M^{-1} M_D^T$$

$\mu \rightarrow 0$ Lepton number is conserved

Three massless neutrinos

Linear seesaw

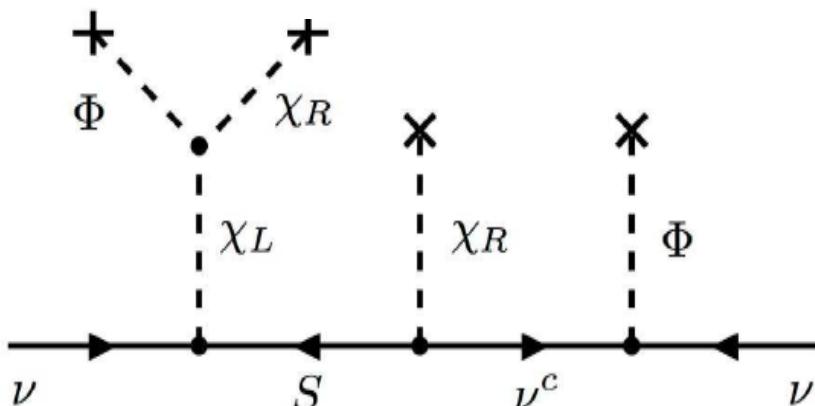
Malinsky, Romao, Valle, PRL95 (05)

$$SO(10) \xrightarrow{M_G} 3_c 2_L 2_R 1_{B-L} \xrightarrow{V_R} 3_c 2_L 1_{I_{3R}} 1_{B-L} \xrightarrow{v_R} 3_c 2_L 1_Y$$

Breaks B-L

$$M_\nu = \begin{pmatrix} 0 & Yv & Fv_L \\ Y^T v & 0 & \tilde{F}v_R \\ F^T v_L & \tilde{F}^T v_R & 0 \end{pmatrix}$$

$$v_L \sim \frac{v v_R}{M_G}$$



Linear in Y

$$M_\nu \simeq \frac{v^2}{M_G} \rho \left[Y(F\tilde{F}^{-1})^T + (F\tilde{F}^{-1})Y^T \right]$$

Neutrino mass is suppressed by M_G
irrespectively how low is $B-L$ breaking scale

Neutrino mass matrix and mu-tau

Grimus, Joshipura, Kaneko, Lavoura,Tanimoto JHEP0407

$$m = \begin{pmatrix} x & y & y \\ y & x + v & y - v \\ y & y - v & x + v \end{pmatrix}$$

$\longleftrightarrow \mu \leftrightarrow \tau$

$$O = \begin{pmatrix} -\cos \theta & \sin \theta & 0 \\ \frac{1}{\sqrt{2}} \sin \theta & \frac{1}{\sqrt{2}} \cos \theta & -\frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} \sin \theta & \frac{1}{\sqrt{2}} \cos \theta & \frac{1}{\sqrt{2}} \end{pmatrix}$$

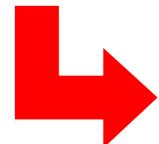
$\sin^2 \theta_{13} = 0$

$\sin^2 \theta_{12} = 0.5$

Neutrino mass matrix and TBM

$$m = \begin{pmatrix} x & y & y \\ y & x + v & y - v \\ y & y - v & x + v \end{pmatrix}$$

$$m_{\nu_{1,1}} + m_{\nu_{1,3}} = m_{\nu_{2,2}} + m_{\nu_{2,3}}$$



$$\sin^2 \theta_{23} = 1/3$$

The Flavor Problem:

- why fermion mass hierarchies?
 - why flavor eigenstates are mixed?
 - why quarks and leptons mixing are so different?
- Why tri-bimaximal lepton mixing?

$$SU_c(3) \times SU_L(2) \times U_Y(1) \times G_f$$

$G_f = SO(3)$ or $SU(3)$  discrete subgroups
spontaneously

Adulpravitchai, Blum, Lindner JHEP 0909
Berger, Grossman JHEP 1002

Discrete groups

A group **G={A,B,C,...}** which consist of a finite number of elements **g** is a *finite group if*

- ★ the set is close with respect to the composition law
- ★ associative
- ★ cancelling rule: $A X = B X$ and $Y A = Y B \rightarrow A = B$

To each finite group correspond a multiplication table

g=6	I	A	B	C	D	E
I	I	A	B	C	D	E
A	A	B	I	E	C	D
B	B	I	A	D	E	C
C	C	D	E	I	A	B
D	D	E	C	B	I	A
E	E	C	D	A	B	I

Not all the product are independent:

$$A C = E, C B = E, B B = A, C B = E, A E = D \rightarrow C A = D$$

It exist a set of **elements** and a set of **independent relations** associated to each multiplication table

Generators of the group

Set of elements

A, C

Set of relations

$$A^3 = C^2 = (AC)^2 = I$$

$$I, A, A^2, C, AC, CA$$

	I	A	B	C	D	E	
$I =$	I		A	B	C	D	E
$A =$	A	A	B		E	C	D
$A^2 =$	B	B		A	D	E	C
$C =$	C	C	D	E		A	B
$CA =$	D	D	E	C	B		A
$AC =$	E	E	C	D	A	B	

Classification of the group of order < 32

Frampton and Kephart, PRD64 (01)

order	groups
6	$S_3 \equiv D_3$
8	$D_4, Q = Q_4$
10	D_5
12	$D_6, Q_6, T \equiv A_4$
14	D_7
16	$D_8, Q_8, Z_2 \times D_4, Z_2 \times Q$
18	$D_9, Z_3 \times D_3$
20	D_{10}, Q_{10}
22	D_{11}
24	$D_{12}, Q_{12}, Z_2 \times D_6, Z_2 \times Q_6, Z_2 \times T, Z_3 \times D_4, Z_3 \times Q, Z_4 \times D_3, S_4$
26	D_{13}
28	D_{14}, Q_{14}
30	$D_{15}, D_5 \times Z_3, D_3 \times Z_5$

For a review of discrete group see also
Ishimori et al 1003.3552

Classification of the group of order < 32

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26	D_{13}
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$$\langle A, B | A^2, B^3, (AB)^4 \rangle$$

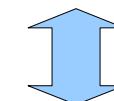
S_n Permutation group of n objects of order $n!$

Classification of the group of order < 32

Frampton and Kephart, PRD64 (01)

Hypercomplex number

order	groups	$a + bi + cj + dk$
6	$S_3 \equiv D_3$	
8	$D_4, Q = Q_4$	$i^4 = 1, i^2 = j^2, ji = i^3 j$
10	D_5	
12	$D_6, Q_6, T \equiv A_4$	
14	D_7	
16	$D_8, Q_8, Z_2 \times D_4, Z_2 \times Q$	$A^4 = I, A^2 = B^2, AB = B^3 A$
18	$D_9, Z_3 \times D_3$	
20	D_{10}, Q_{10}	
22	D_{11}	
24	$D_{12}, Q_{12}, Z_2 \times D_6, Z_2 \times Q_6, Z_2 \times T, Z_3 \times D_4, Z_3 \times Q, Z_4 \times D_3, S_4$	
26	D_{13}	
28	D_{14}, Q_{14}	
30	$D_{15}, D_5 \times Z_3, D_3 \times Z_5$	



Q_n Quaternion series of order 2^n

Classification of the group of order < 32

Frampton and Kephart, PRD64 (01)

order	groups
6	$S_3 \equiv D_3$
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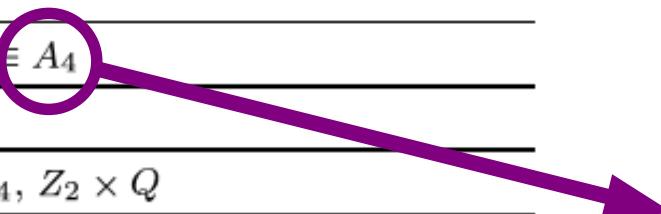
$$\langle A, B | A^n, B^2, (AB)^2 \rangle$$

D_n Dihedral series of order $2 n$

Classification of the group of order < 32

Frampton and Kephart, PRD64 (01)

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6	$S_3 \equiv D_3$
8	$D_4, Q = Q_4$
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24	$D_{12}, Q_{12}, Z_2 \times D_6, Z_2 \times Q_6, Z_2 \times T, Z_3 \times D_4, Z_3 \times Q, Z_4 \times D_3, S_4$
26	D_{13}
28	D_{14}, Q_{14}
30	$D_{15}, D_5 \times Z_3, D_3 \times Z_5$



**Smallest group
with triplet irrep 3**

$$\left\langle S, T \mid S^2, T^3, (ST)^3 \right\rangle$$

A_n Alternating series of order $n!/2$

A4 group

12 elem.

C1: I

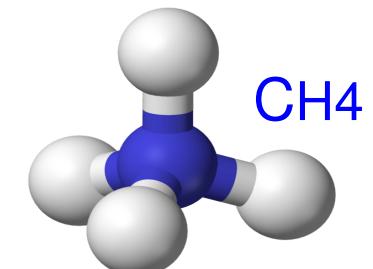
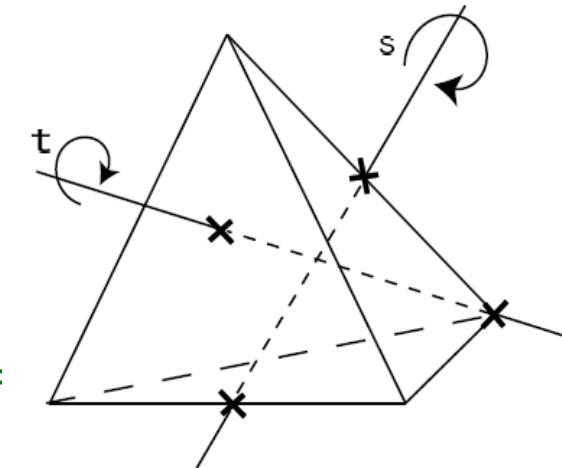
C2: T, ST, TS,STS

C3: TT, STT, TTS, TST

C4: S, TTST, TSTT

Isomorphic to group of
tetrahedron rotations

12 rotations



$$\langle S, T | S^2 = I, T^3 = I, (ST)^3 = I \rangle$$

Z_2 ;C2

Z_3 ;C4

Subgroups of A4

A4 product rules

Class	χ^1	$\chi^{1'}$	$\chi^{1''}$	χ^3
C_1	1	1	1	3
C_2	1	ω	ω^2	0
C_3	1	ω^2	ω	0
C_4	1	1	1	-1

$$1' \times 1' = 1'', 1' \times 1'' = 1, 1'' \times 1'' = 1' \text{ etc.}$$

$$3 \times 3 = 1 + 1' + 1'' + 3 + 3$$

$$\omega^3 = 1$$

$$S = \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{pmatrix}, T = \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 1 & 0 & 0 \end{pmatrix} \quad \rightarrow$$

S-diag basis

$$1 = a_1 b_1 + a_2 b_2 + a_3 b_3$$

$$1' = a_1 b_1 + \omega^2 a_2 b_2 + \omega a_3 b_3$$

$$1'' = a_1 b_1 + \omega a_2 b_2 + \omega^2 a_3 b_3$$

$$3 \sim (a_2 b_3, a_3 b_1, a_1 b_2)$$

$$3 \sim (a_3 b_2, a_1 b_3, a_2 b_1)$$

A4 product rules

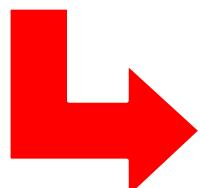
$$T' = VTV^\dagger = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \omega & 0 \\ 0 & 0 & \omega^2 \end{pmatrix},$$

$$S' = VSV^\dagger = \frac{1}{3} \begin{pmatrix} -1 & 2 & 2 \\ 2 & -1 & 2 \\ 2 & 2 & -1 \end{pmatrix}$$

$$1 = a_1 b_1 + a_2 b_3 + a_3 b_2$$

$$1' = a_3 b_3 + a_1 b_2 + a_2 b_1$$

$$1'' = a_2 b_2 + a_1 b_3 + a_3 b_3$$



$$3_{symm} \sim \frac{1}{3} (2a_1 b_1 - a_2 b_3 - a_3 b_2, 2a_3 b_3 - a_1 b_2 - a_2 b_1, 2a_2 b_2 - a_1 b_3 - a_3 b_1)$$

$$3_{antisymm} \sim \frac{1}{2} (a_2 b_3 - a_3 b_2, a_1 b_2 - a_2 b_1, a_1 b_3 - a_3 b_1)$$

Models with A4 symmetry: matter assignments

Type	L_i^e	ℓ_i^e	ν_i^e	Δ
A1	<u>3</u>	<u>1</u> , <u>1'</u> , <u>1''</u>	-	-
A2	<u>3</u>	<u>1</u> , <u>1'</u> , <u>1''</u> , <u>3</u>	-	<u>1</u> , <u>1'</u> , <u>1''</u> , <u>3</u>
B1	<u>3</u>	<u>1</u> , <u>1'</u> , <u>1''</u>	<u>3</u>	-
B2	<u>3</u>	<u>1</u> , <u>1'</u> , <u>1''</u>	-	<u>1</u> , <u>3</u>
C1	-	-	-	-
C2	<u>3</u>	<u>3</u>	-	<u>1</u>
C3	-	-	-	<u>1</u> , <u>3</u>
C4	-	-	-	<u>1</u> , <u>1'</u> , <u>1''</u> , <u>3</u>
D1	-	-	-	-
D2	<u>3</u>	<u>3</u>	<u>3</u>	<u>1</u>
D3	-	-	-	<u>1'</u>
D4	-	-	-	<u>1'</u> , <u>3</u>
E	<u>3</u>	<u>3</u>	<u>1</u> , <u>1'</u> , <u>1''</u>	-
F	<u>1</u> , <u>1'</u> , <u>1''</u>	<u>3</u>	<u>3</u>	<u>1</u> or <u>1'</u>
G	<u>3</u>	<u>1</u> , <u>1'</u> , <u>1''</u>	<u>1</u> , <u>1'</u> , <u>1''</u>	-
H	<u>3</u>	<u>1</u> , <u>1</u> , <u>1</u>	-	-
I	<u>3</u>	<u>1</u> , <u>1</u> , <u>1</u>	<u>1</u> , <u>1</u> , <u>1</u>	-
J	<u>3</u>	<u>1</u> , <u>1</u> , <u>1</u>	<u>3</u>	-

Models with A4 symmetry

Ma, Rajasekaran PRD64(01')
 Babu, Ma, Valle PLB552(02')

	L	l_1^c	l_2^c	l_3^c	ν_1^c	H	ϕ	ϕ'
$SU(2)$	2	1	1	1	1	2	1	1
A_4	3	1	$1''$	$1'$	3	1	3	3

- Altarelli, Feruglio NPB720 (05')

$$\mathcal{L} = y_e(L\phi')l_1^c + y_\mu(L\phi')'l_2^c + y_\tau(L\phi'')l_3^c + y_D(L\nu^c)H + y_a(\nu^c\nu^c) + y_b(\nu^c\nu^c\phi)$$

$$\phi \longleftrightarrow \phi'$$

We assume extra Abelian symmetries Z_3, Z_4, \dots
 to forbid unwanted terms

Models with A4 symmetry

	L	l_1^c	l_2^c	l_3^c	ν_1^c	H	ϕ	ϕ'
$SU(2)$	2	1	1	1	1	2	1	1
A_4	3	1	$1''$	$1'$	3	1	3	3

$$\mathcal{L} = y_e(L\phi')l_1^c + y_\mu(L\phi')'l_2^c + y_\tau(L\phi'')l_3^c + y_D(L\nu^c)H + y_a(\nu^c\nu^c) + y_b(\nu^c\nu^c\phi)$$

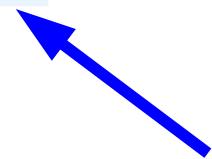
$$(L\phi') = L_1\phi'_1 + L_2\phi'_3 + L_3\phi'_2$$

Models with A4 symmetry

	L	l_1^c	l_2^c	l_3^c	ν_1^c	H	ϕ	ϕ'
$SU(2)$	2	1	1	1	1	2	1	1
A_4	3	1	$1''$	$1'$	3	1	3	3

$$\mathcal{L} = y_e(L\phi')l_1^c + y_\mu(L\phi')'l_2^c + y_\tau(L\phi'')l_3^c + y_D(L\nu^c)H + y_a(\nu^c\nu^c) + y_b(\nu^c\nu^c\phi)$$

$$(L\phi') = L_1\phi'_1 + L_2\cancel{\phi'_3} + L_3\cancel{\phi'_2}$$



$$\langle \phi' \rangle \sim (1, 0, 0)$$

Models with A4 symmetry

	L	l_1^c	l_2^c	l_3^c	ν_1^c	H	ϕ	ϕ'
$SU(2)$	2	1	1	1	1	2	1	1
A_4	3	1	$1''$	$1'$	3	1	3	3

$$\mathcal{L} = y_e(L\phi')l_1^c + y_\mu(L\phi')'l_2^c + y_\tau(L\phi'')l_3^c + y_D(L\nu^c)H + y_a(\nu^c\nu^c) + y_b(\nu^c\nu^c\phi)$$

$$M_l = \begin{pmatrix} y_e & 0 & 0 \\ 0 & y_\mu & 0 \\ 0 & 0 & y_\tau \end{pmatrix}$$

$$(L\phi') = L_1\phi'_1 + L_2\cancel{\phi'_3} + L_3\cancel{\phi'_2}$$

$$\langle \phi' \rangle \sim (1, 0, 0)$$

Models with A4 symmetry

	L	l_1^c	l_2^c	l_3^c	ν_1^c	H	ϕ	ϕ'
$SU(2)$	2	1	1	1	1	2	1	1
A_4	3	1	$1''$	$1'$	3	1	3	3

$$\mathcal{L} = y_e(L\phi')l_1^c + y_\mu(L\phi')'l_2^c + y_\tau(L\phi'')l_3^c + y_D(L\nu^c)H + y_a(\nu^c\nu^c) + y_b(\nu^c\nu^c\phi)$$

$$(L\phi')' = L_3\phi'_3 + L_1\phi'_2 + L_2\phi'_1$$

$$\langle \phi' \rangle \sim (1, 0, 0)$$

$$M_l = \begin{pmatrix} y_e & 0 & 0 \\ 0 & y_\mu & 0 \\ 0 & 0 & y_\tau \end{pmatrix}$$

Models with A4 symmetry

	L	l_1^c	l_2^c	l_3^c	ν_1^c	H	ϕ	ϕ'
$SU(2)$	2	1	1	1	1	2	1	1
A_4	3	1	$1''$	$1'$	3	1	3	3

$$\mathcal{L} = y_e(L\phi')l_1^c + y_\mu(L\phi')'l_2^c + \boxed{y_\tau(L\phi')''l_3^c} + y_D(L\nu^c)H + y_a(\nu^c\nu^c) + y_b(\nu^c\nu^c\phi)$$

$$(L\phi')'' = L_2\phi'_2 + L_1\phi'_3 + \boxed{L_3\phi'_1}$$

$$\langle \phi' \rangle \sim (1, 0, 0)$$

$$M_l = \begin{pmatrix} y_e & 0 & 0 \\ 0 & y_\mu & 0 \\ 0 & 0 & y_\tau \end{pmatrix}$$



Models with A4 symmetry

	L	l_1^c	l_2^c	l_3^c	ν_1^c	H	ϕ	ϕ'
$SU(2)$	2	1	1	1	1	2	1	1
A_4	3	1	$1''$	$1'$	3	1	3	3

$$\mathcal{L} = y_e(L\phi')l_1^c + y_\mu(L\phi')'l_2^c + y_\tau(L\phi'')l_3^c + \boxed{y_D(L\nu^c)H + y_a(\nu^c\nu^c) + y_b(\nu^c\nu^c\phi)}$$

$$\psi_1\varphi_1 + \psi_2\varphi_3 + \psi_3\varphi_2 \sim 1 \quad \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}$$

Models with A4 symmetry

	L	l_1^c	l_2^c	l_3^c	ν_1^c	H	ϕ	ϕ'
$SU(2)$	2	1	1	1	1	2	1	1
A_4	3	1	$1''$	$1'$	3	1	3	3

$$\mathcal{L} = y_e(L\phi')l_1^c + y_\mu(L\phi')'l_2^c + y_\tau(L\phi'')l_3^c + y_D(L\nu^c)H + y_a(\nu^c\nu^c) + y_b(\nu^c\nu^c\phi)$$

$$(\nu^c\nu^c)_3 = \begin{pmatrix} 2\nu_1^c\nu_1^c - \nu_2^c\nu_3^c - \nu_3^c\nu_2^c \\ 2\nu_3^c\nu_3^c - \nu_1^c\nu_2^c - \nu_2^c\nu_1^c \\ 2\nu_2^c\nu_2^c - \nu_1^c\nu_3^c - \nu_3^c\nu_1^c \end{pmatrix}$$

Models with A4 symmetry

	L	l_1^c	l_2^c	l_3^c	ν_1^c	H	ϕ	ϕ'
$SU(2)$	2	1	1	1	1	2	1	1
A_4	3	1	$1''$	$1'$	3	1	3	3

$$\mathcal{L} = y_e(L\phi')l_1^c + y_\mu(L\phi')'l_2^c + y_\tau(L\phi'')l_3^c + y_D(L\nu^c)H + y_a(\nu^c\nu^c) + y_b(\nu^c\nu^c\phi)$$

$$(\nu^c\nu^c)_3 = \begin{pmatrix} 2\nu_1^c\nu_1^c - \nu_2^c\nu_3^c - \nu_3^c\nu_2^c \\ 2\nu_3^c\nu_3^c - \nu_1^c\nu_2^c - \nu_2^c\nu_1^c \\ 2\nu_2^c\nu_2^c - \nu_1^c\nu_3^c - \nu_3^c\nu_1^c \end{pmatrix} \quad \langle \phi \rangle \sim (1, 1, 1)$$

$$M_R =$$

$$+ b \begin{pmatrix} 2 & -1 & -1 \\ -1 & 2 & -1 \\ -1 & -1 & 2 \end{pmatrix}$$



Models with A4 symmetry

	L	l_1^c	l_2^c	l_3^c	ν_1^c	H	ϕ	ϕ'
$SU(2)$	2	1	1	1	1	2	1	1
A_4	3	1	$1''$	$1'$	3	1	3	3

$$\mathcal{L} = y_e(L\phi')l_1^c + y_\mu(L\phi')'l_2^c + y_\tau(L\phi'')l_3^c + y_D(L\nu^c)H + y_a(\nu^c\nu^c) + y_b(\nu^c\nu^c\phi)$$

$$(\nu^c\nu^c)_1 = (\nu_1^c\nu_1^c + \nu_2^c\nu_3^c + \nu_3^c\nu_2^c)$$

$$M_R = a \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} + b \begin{pmatrix} 2 & -1 & -1 \\ -1 & 2 & -1 \\ -1 & -1 & 2 \end{pmatrix}$$


Models with A4 symmetry

$$M_R = a \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} + b \begin{pmatrix} 2 & -1 & -1 \\ -1 & 2 & -1 \\ -1 & -1 & 2 \end{pmatrix}$$

$$m_D \frac{1}{M_R} m_D^T \sim \begin{pmatrix} -a^2 + 2ab + 3b^2 & b(-a + 3b) & b(-a + 3b) \\ b(-a + 3b) & b(2a + 3b) & -a^2 - ab + 3b^2 \\ b(-a + 3b) & -a^2 - ab + 3b^2 & b(2a + 3b) \end{pmatrix}$$

$\mu \leftrightarrow \tau$



$$m_{\nu_{1,1}} + m_{\nu_{1,3}} = m_{\nu_{2,2}} + m_{\nu_{2,3}}$$

TBM

Models with A4 symmetry

$$M_R = a \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} + b \begin{pmatrix} 2 & -1 & -1 \\ -1 & 2 & -1 \\ -1 & -1 & 2 \end{pmatrix}$$

$$m_D \frac{1}{M_R} m_D^T \sim \begin{pmatrix} -a^2 + 2ab + 3b^2 & b(-a + 3b) & b(-a + 3b) \\ b(-a + 3b) & b(2a + 3b) & -a^2 - ab + 3b^2 \\ b(-a + 3b) & -a^2 - ab + 3b^2 & b(2a + 3b) \end{pmatrix}$$

Only three free parameters, predicts three angles and absolute neutrino mass scale

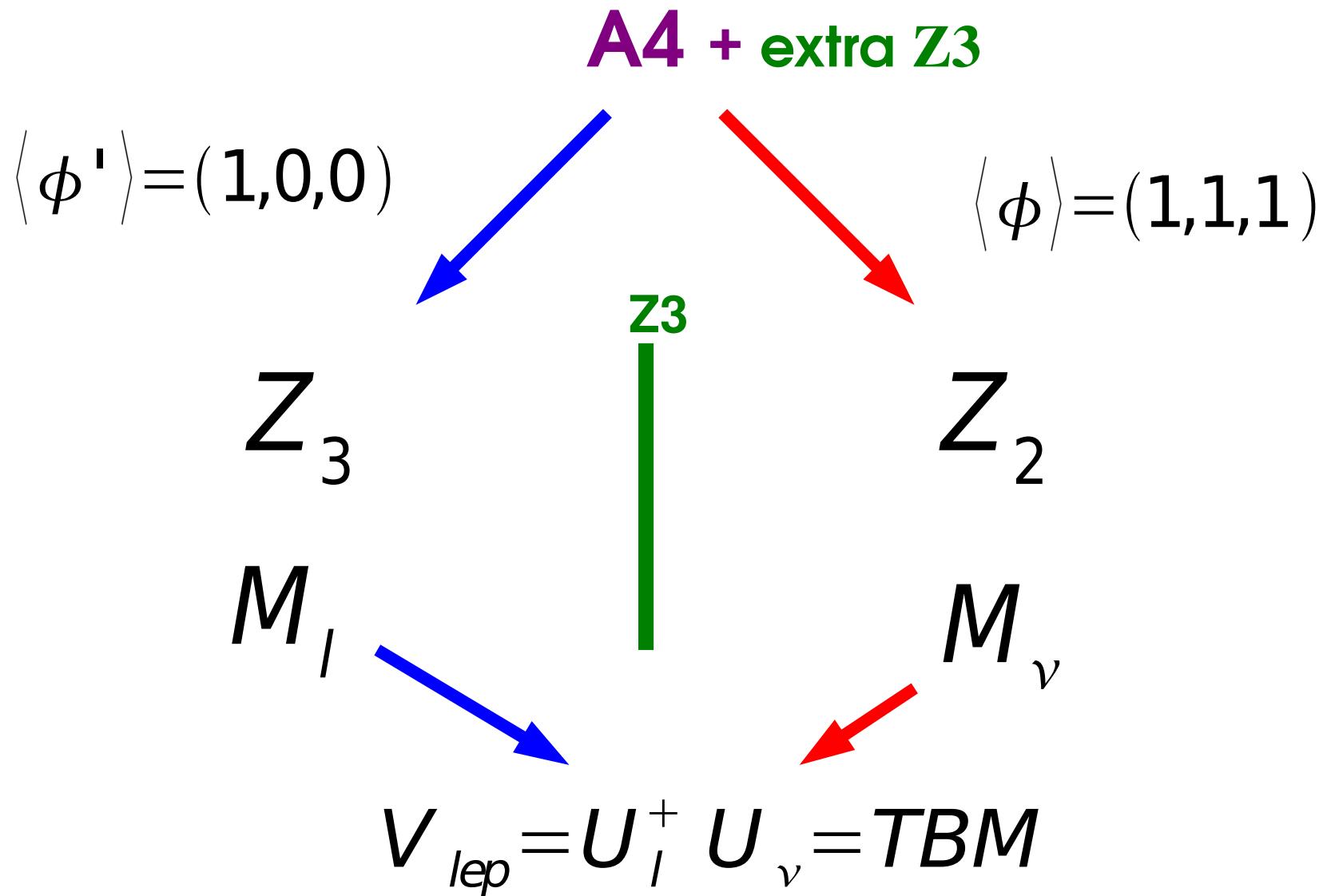
A4 breaking

$$T' = VTV^\dagger = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \omega & 0 \\ 0 & 0 & \omega^2 \end{pmatrix}, \quad \xrightarrow{\text{invariant}} \quad \langle \phi' \rangle \sim (1, 0, 0)$$

$$S' = VSV^\dagger = \frac{1}{3} \begin{pmatrix} -1 & 2 & 2 \\ 2 & -1 & 2 \\ 2 & 2 & -1 \end{pmatrix} \quad \xrightarrow{\text{invariant}} \quad \langle \phi \rangle \sim (1, 1, 1)$$

$$T T T = I \quad A4 \xrightarrow{\phi'} Z3$$

$$S S = I \quad A4 \xrightarrow{\phi} Z2$$



Large neutrino
mixing

$\phi \neq \phi'$
Misalignment

Vacuum alignment problem

Altarelli, Feruglio, NPB 720 05'

$$B_1 = \varphi_1^2 + \varphi_2^2 + \varphi_3^2$$

$$B_2 = \varphi_1'^2 + \varphi_2'^2 + \varphi_3'^2$$

$$T_1 = \varphi_1 \varphi_2 \varphi_3$$

$$T_2 = \varphi_1 \varphi_2' \varphi_3' + \varphi_2 \varphi_3' \varphi_1' + \varphi_3 \varphi_1' \varphi_2'$$

$$Q_1 = \varphi_1^2 \varphi_2^2 + \varphi_2^2 \varphi_3^2 + \varphi_3^2 \varphi_1^2$$

$$Q_2 = |\varphi_1^2 + \omega^2 \varphi_2^2 + \omega \varphi_3^2|^2$$

$$Q_3 = \varphi_1'^2 \varphi_2'^2 + \varphi_2'^2 \varphi_3'^2 + \varphi_3'^2 \varphi_1'^2$$

$$Q_4 = |\varphi_1'^2 + \omega^2 \varphi_2'^2 + \omega \varphi_3'^2|^2$$

$$Q_5 = \varphi_1 \varphi_2 \varphi_1' \varphi_2' + \varphi_2 \varphi_3 \varphi_2' \varphi_3' + \varphi_3 \varphi_1 \varphi_3' \varphi_1'$$

$$Q_6 = (\varphi_1^2 + \varphi_2^2 + \varphi_3^2)(\varphi_1'^2 + \varphi_2'^2 + \varphi_3'^2)$$

$$Q_7 = (\varphi_1^2 + \omega^2 \varphi_2^2 + \omega \varphi_3^2)(\varphi_1'^2 + \omega \varphi_2'^2 + \omega^2 \varphi_3'^2)$$

$$\begin{aligned} V &= \frac{M_1^2}{2} B_1^2 + \frac{M_2^2}{2} B_2^2 + \mu_1 T_1 + \mu_2 T_2 \\ &+ c_1 Q_1 + c_2 Q_2 + c_3 Q_3 + c_4 Q_4 \\ &+ c_5 Q_5 + c_6 Q_6 + (c_7 Q_7 + c.c.) \quad , \end{aligned}$$

Vacuum alignment problem

Altarelli, Feruglio, NPB 720 05'

$$\langle \varphi \rangle = (v, v, v)$$

$$\langle \varphi' \rangle = (v', 0, 0)$$

$$\frac{\partial V}{\partial \varphi_1} = M_1^2 v + \mu_1 v^2 + 4c_1 v^3 + 2c_6 v v'^2 + 2(c_7 + \bar{c}_7) v v'^2 = 0$$

$$\frac{\partial V}{\partial \varphi_2} = M_1^2 v + \mu_1 v^2 + 4c_1 v^3 + 2c_6 v v'^2 + 2(\omega^2 c_7 + \omega \bar{c}_7) v v'^2 = 0$$

$$\frac{\partial V}{\partial \varphi_3} = M_1^2 v + \mu_1 v^2 + 4c_1 v^3 + 2c_6 v v'^2 + 2(\omega c_7 + \omega^2 \bar{c}_7) v v'^2 = 0$$

$$\frac{\partial V}{\partial \varphi'_1} = M_2^2 v' + 4c_4 v'^3 + 6c_6 v^2 v' = 0$$

$$\frac{\partial V}{\partial \varphi'_2} = \mu_2 v v' + c_5 v^2 v' = 0$$

$$\frac{\partial V}{\partial \varphi'_3} = \mu_2 v v' + c_5 v^2 v' = 0 .$$

Incompatible unless $c_7=0$



$v, v' = 0$

Vacuum alignment problem

Altarelli, Feruglio, NPB 720 05'

$$\langle \varphi \rangle = (v, v, v)$$

$$\langle \varphi' \rangle = (v', 0, 0)$$

$$\frac{\partial V}{\partial \varphi_1} = M_1^2 v + \mu_1 v^2 + 4c_1 v^3 + 2c_6 v v'^2 .$$

$$\frac{\partial V}{\partial \varphi_2} = M_1^2 v + \mu_1 v^2 + 4c_1 v^3 + 2c_6 v v'^2 . \quad \text{c7=0}$$

$$\frac{\partial V}{\partial \varphi_3} = M_1^2 v + \mu_1 v^2 + 4c_1 v^3 + 2c_6 v v'^2 .$$

$$\frac{\partial V}{\partial \varphi'_1} = M_2^2 v' + 4c_4 v'^3 + 6c_6 v^2 v' = 0$$

$$\frac{\partial V}{\partial \varphi'_2} = \mu_2 v v' + c_5 v^2 v' = 0$$

$$\frac{\partial V}{\partial \varphi'_3} = \mu_2 v v' + c_5 v^2 v' = 0 .$$

v, v' different from zero if c6 = 0, mu2 = 0, c5 = 0

Vacuum alignment problem

Altarelli, Feruglio, NPB 720 05'

$$\begin{aligned}
 V = & \frac{M_1^2}{2} B_1^2 + \frac{M_2^2}{2} B_2^2 + \mu_1 T_1 + \mu_2 T_2 \\
 & + c_1 Q_1 + c_2 Q_2 + c_3 Q_3 + c_4 Q_4 \\
 & + c_5 Q_5 + c_6 Q_6 + (c_7 Q_7 + c.c) ,
 \end{aligned}$$

$c_7 = 0, c_6 = 0, \mu_2 = 0, c_5 = 0$

$$B_1 = \varphi_1^2 + \varphi_2^2 + \varphi_3^2$$

$$B_2 = \varphi'_1^2 + \varphi'_2^2 + \varphi'_3^2$$

$$T_1 = \varphi_1 \varphi_2 \varphi_3$$

$$\rightarrow T_2 = \varphi_1 \varphi'_2 \varphi'_3 + \varphi_2 \varphi'_3 \varphi'_1 + \varphi_3 \varphi'_1 \varphi'_2$$

$$Q_1 = \varphi_1^2 \varphi_2^2 + \varphi_2^2 \varphi_3^2 + \varphi_3^2 \varphi_1^2$$

$$Q_2 = |\varphi_1^2 + \omega^2 \varphi_2^2 + \omega \varphi_3^2|^2$$

$$Q_3 = \varphi'_1^2 \varphi'_2^2 + \varphi'_2^2 \varphi'_3^2 + \varphi'_3^2 \varphi'_1^2$$

$$Q_4 = |\varphi'_1^2 + \omega^2 \varphi'_2^2 + \omega \varphi'_3^2|^2$$

we don't want terms mixing

$$\varphi \quad \varphi'$$

Abelian symmetries do not forbid such a terms

$$\rightarrow Q_5 = \varphi_1 \varphi_2 \varphi'_1 \varphi'_2 + \varphi_2 \varphi_3 \varphi'_2 \varphi'_3 + \varphi_3 \varphi_1 \varphi'_3 \varphi'_1$$

$$\rightarrow Q_6 = (\varphi_1^2 + \varphi_2^2 + \varphi_3^2)(\varphi'_1^2 + \varphi'_2^2 + \varphi'_3^2)$$

$$\rightarrow Q_7 = (\varphi_1^2 + \omega^2 \varphi_2^2 + \omega \varphi_3^2)(\varphi'_1^2 + \omega \varphi'_2^2 + \omega^2 \varphi'_3^2)$$

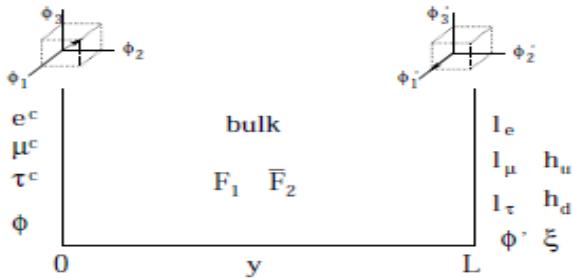
Vacuum alignment problem

- SUSY:
the renormalizable superpotential only contains cubic terms like

$$\varphi\varphi\varphi' \text{ not invariant under some abelian symm}$$

Altarelli, Feruglio NPB741

- Extra Dimension



Altarelli, Feruglio NPB720

- Enlarge the flavor group

$A_4 \ltimes (Z_2)^3$ each component of the scalar triplet is glued to a corresponding right-handed field

$$(L_1\varphi_1 + L_2\varphi_3 + L_3\varphi_2)l_1^c$$



Grimus, Lavoura JHEP0904
Mohapatra, Nasri PLB639
Morisi PRD79

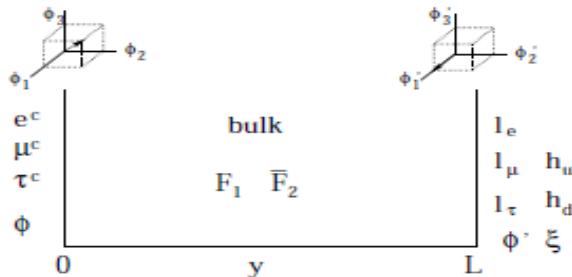
Vacuum alignment problem

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the renormalizable superpotential only contains cubic terms like

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Altarelli, Feruglio NPB741

- Extra Dimension



Altarelli, Feruglio NPB720

- Enlarge the flavor group

$A_4 \ltimes (Z_2)^3$ each component of the scalar triplet is glued to a corresponding right-handed field

$$(L_1\varphi_1 + L_2\varphi_3 + L_3\varphi_2)l_1^c$$

charged leptons diagonal

also with $\langle\varphi\rangle = (v, v, v)$

Grimus, Lavoura JHEP0904
Mohapatra, Nasri PLB639

Morisi PRD79

TBM as accidental mixing

Abbs, Smirnov 1004.0099

Morisi, Peinado PRD80 (09')

fields	L_i	l_i^c	H_i
$SU(2)_L$	2	1	2
A_4	3	3	3

- no flavons
- three Higgs doublets
- no extra abelian symmetry
- neutrino mass from dim 5 oper

Assuming the vevs complex, the minimization of the potential gives

$$\langle \phi_1 \rangle = v_1, \quad \langle \phi_2 \rangle = v e^{i\alpha/2}, \quad \langle \phi_3 \rangle = v e^{-i\alpha/2}$$

Lavoura,kuhbock EPJC55

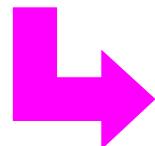
$$L_{\text{Yukawa}} = y_1 (\bar{L}_1 \phi_3 l_2^c + \bar{L}_2 \phi_1 l_3^c + \bar{L}_3 \phi_2 l_1^c) + \\ + y_2 (\bar{L}_1 \phi_2 l_3^c + \bar{L}_2 \phi_3 l_1^c + \bar{L}_3 \phi_1 l_2^c)$$



$$M_l = \begin{pmatrix} 0 & ae^{i\alpha} & be^{-i\alpha} \\ be^{i\alpha} & 0 & ar \\ ae^{-i\alpha} & br & 0 \end{pmatrix}$$

charged leptons

$$M_l M_l^T = \begin{pmatrix} a^2 + b^2 & abr & abr \\ abr & b^2 + a^2 r^2 & ab \\ abr & ab & a^2 + b^2 r^2 \end{pmatrix}$$



$$\begin{aligned} r &\approx \frac{m_\tau}{\sqrt{m_e m_\mu}} \sqrt{1 - \frac{m_e^2 m_\mu^2}{m_\tau^4}}, \\ a &\approx \frac{m_\mu}{m_\tau} \sqrt{m_e m_\mu} \left[1 + \frac{1}{2} \frac{m_\mu^2}{m_\tau^2} \right], \\ b &\approx \sqrt{m_e m_\mu} \left[1 - \frac{1}{2} \frac{m_\mu^2}{m_\tau^2} \right]. \end{aligned}$$



$$a < b \ll r$$

$$O_{l12} \approx \frac{b}{a} r^{-1}, \quad O_{l13} \approx \frac{a}{b} r^{-1}, \quad O_{l23} \approx \frac{a}{b} r^{-2}$$

$$O_l = \begin{pmatrix} 0.997 & 0.069 & 2.44 \times 10^{-4} \\ -0.069 & 0.997 & 1.075 \times 10^{-6} \\ -2.439 \times 10^{-4} & -1.800 \times 10^{-5} & 0.999 \end{pmatrix}$$

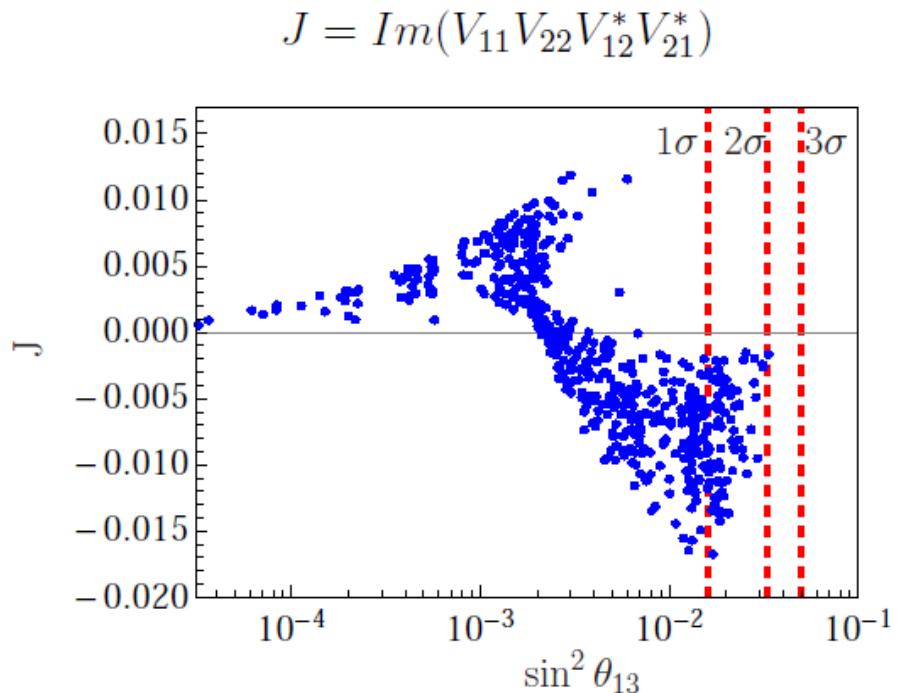
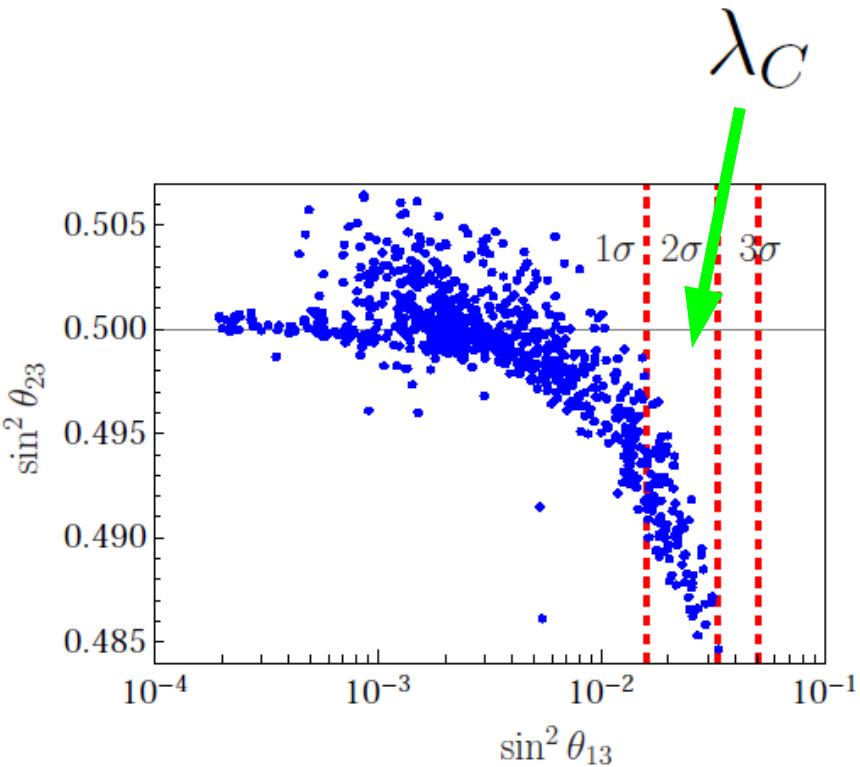
Neutrino mass matrix

$$\begin{aligned}\mathcal{L}_{5d} = & \beta (LL)_3 (HH)_3 + k (LL)_1 (HH)_1 + \alpha' (LL)_{1'} (HH)_{1''} + \alpha'' (LL)_{1''} (HH)_{1'} + \\ & + [a (LH)_{3^a} (LH)_{3^a} + b (LH)_{3^a} (LH)_{3^b} + c (LH)_{3^b} (LH)_{3^a} + d (LH)_{3^b} (LH)_{3^b}] + \\ & + l (LH)_1 (LH)_1 + l' [(LH)_{1'} (LH)_{1''} + (LH)_{1''} (LH)_{1'}],\end{aligned}$$

$$M_\nu = \begin{pmatrix} xr^2 & \kappa r e^{-i\alpha} & \kappa r e^{i\alpha} \\ \kappa r e^{-i\alpha} & zr^2 & \kappa \\ \kappa r e^{i\alpha} & \kappa & yr^2 \end{pmatrix}$$

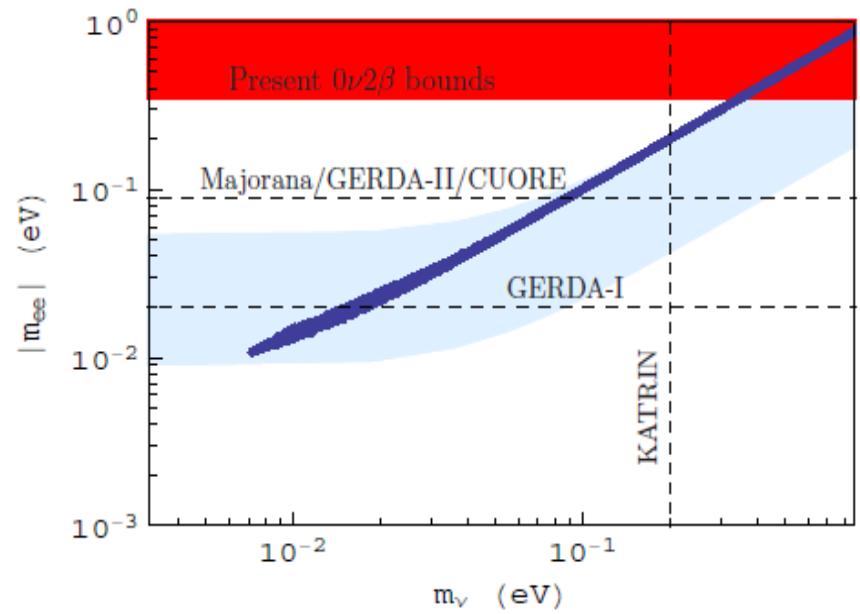
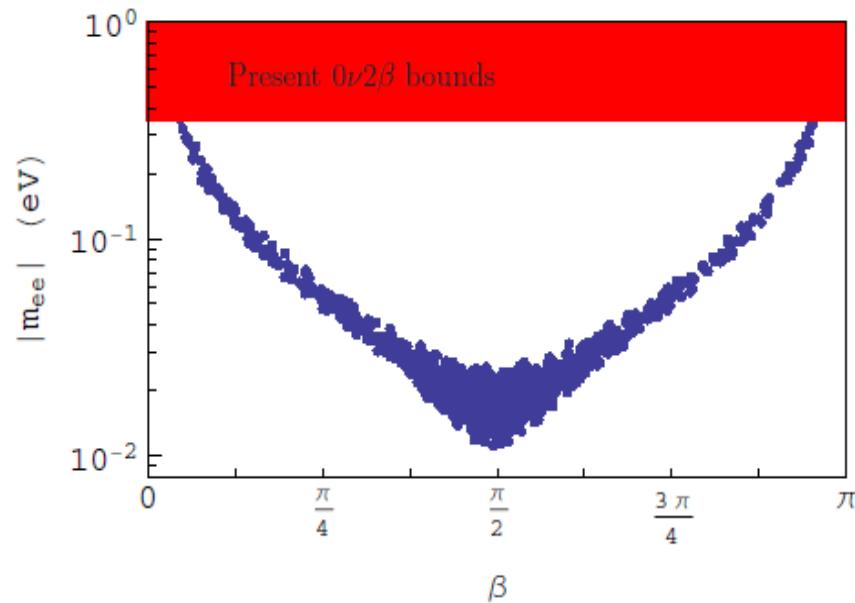
we have six free parameters, \mathbf{r} is already fixed

Neutrino phenomenology



- in general all mixing angle receive corrections of the same order
- allowed departures of the solar angle from the best fit $\mathcal{O}(\lambda_C^2)$
- then TBM models give corrections for the reactor angle of the same order

Neutrino phenomenology



$$m_{\text{light}} > 0.008$$

conclusions

Which is the flavor symmetry?

If tri-bimaximal will be confirmed from future experiments

may be a non-Abelian discrete group of $SU(2)$, $SU(3), \dots$

however...

conclusions

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...S3, T', Delta(27), Dn,..

conclusions

If tri-bimaximal will be ruled out

non-Abelian discrete groups are useful to have
small reactor and large atmospheric angles

in general models with discrete group have few
free parameters and are predictive