

Precision calculations for BSM physics at the LHC

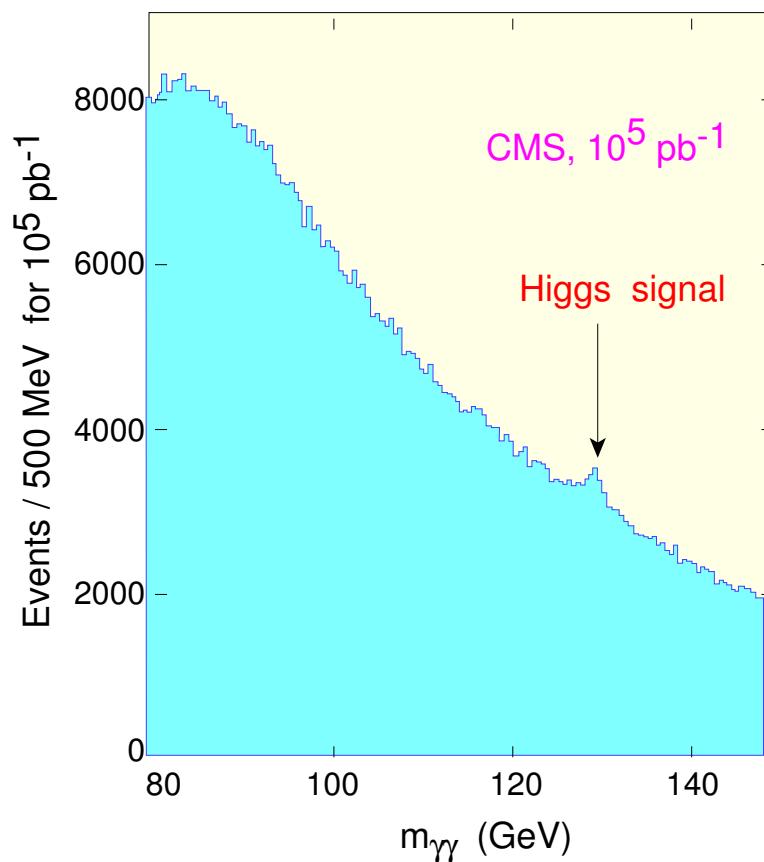
Michael Krämer (RWTH Aachen)

Sponsored by:

SFB/TR 9 “Computational Particle Physics”, Helmholtz Alliance “Physics at the Terascale”, BMBF-Theorie-Verbund, Marie Curie Research Training Network “Heptools”, Graduiertenkolleg “Elementarteilchenphysik an der TeV-Skala”

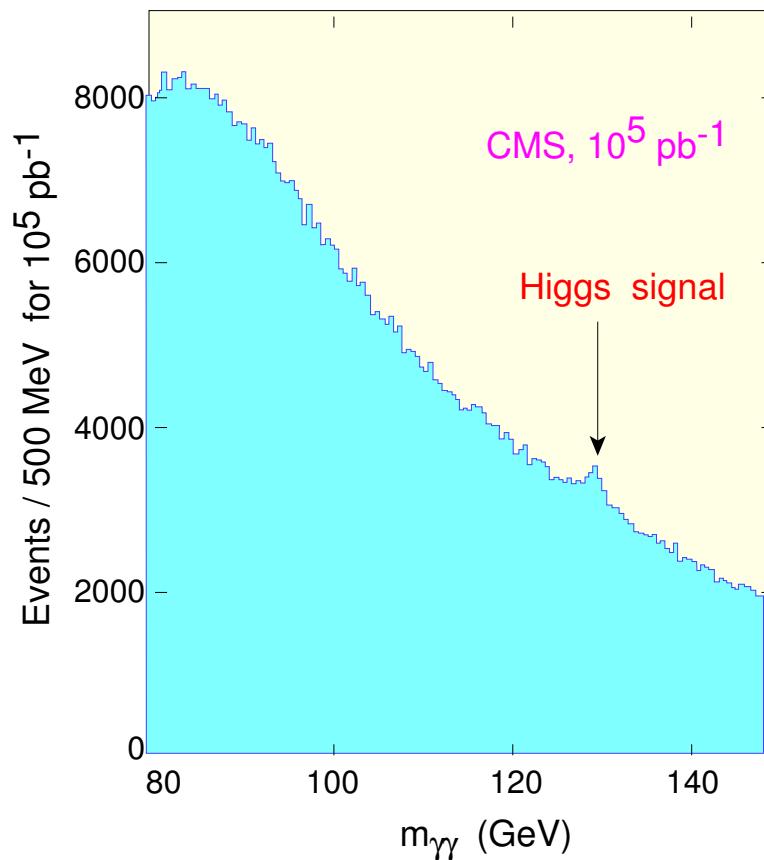
Why precision calculations for LHC physics?

- We do not need theory input for LHC discoveries



Why precision calculations for LHC physics?

- We do not need theory input for LHC discoveries



but only if we see a resonance...

Contents

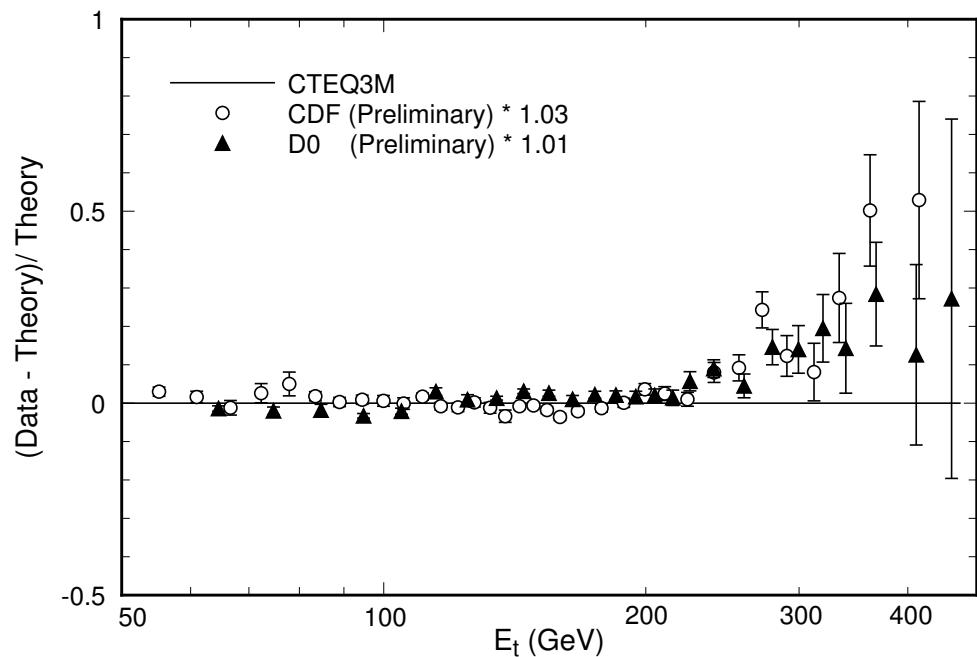
- Why precision calculations for LHC physics?
 - lessons from the past
 - prospects: SUSY searches at the LHC

Contents

- Why precision calculations for LHC physics?
 - lessons from the past
 - prospects: SUSY searches at the LHC

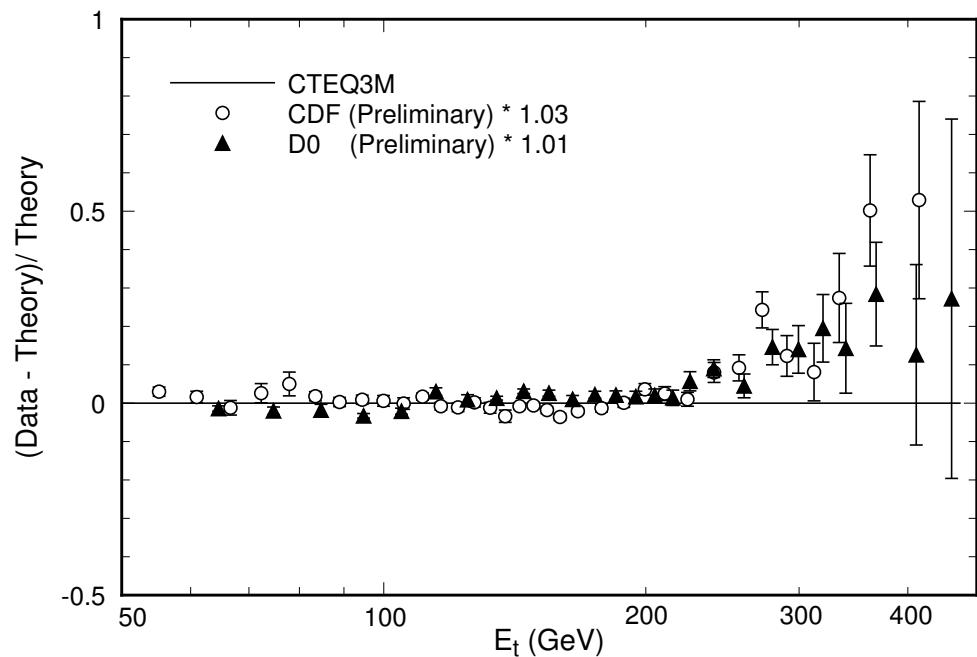
Why precision calculations for LHC physics?

- Tevatron jet cross section at large E_T

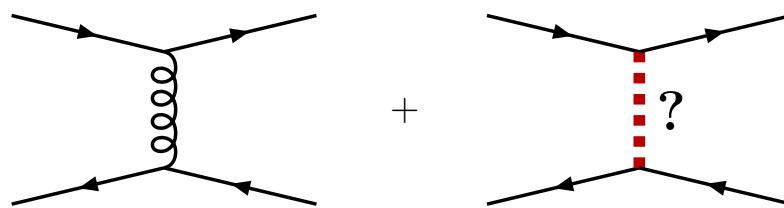


Why precision calculations for LHC physics?

- Tevatron jet cross section at large E_T

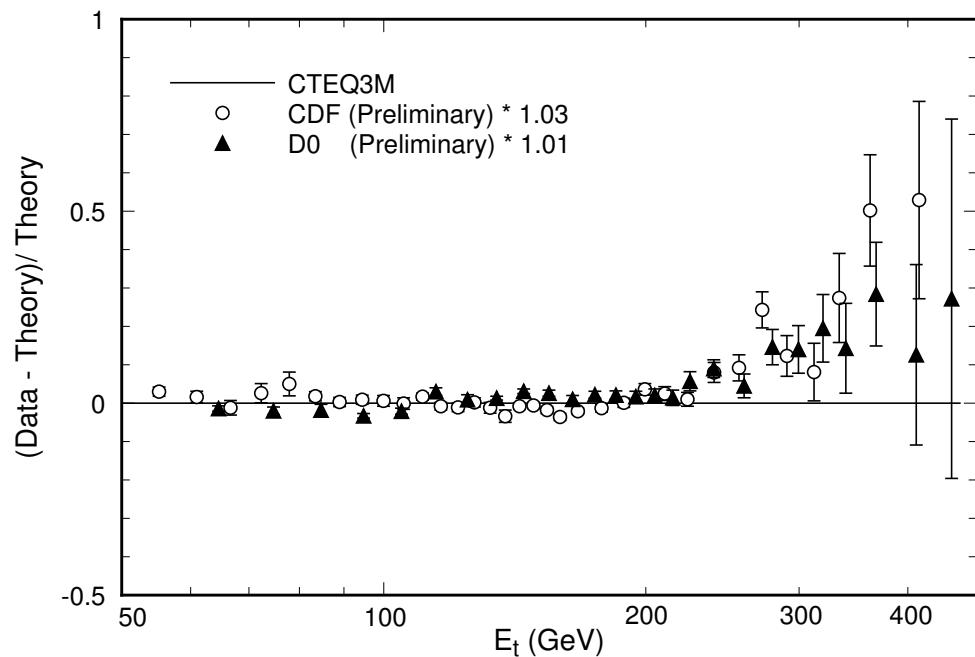


⇒ new physics?



Why precision calculations for LHC physics?

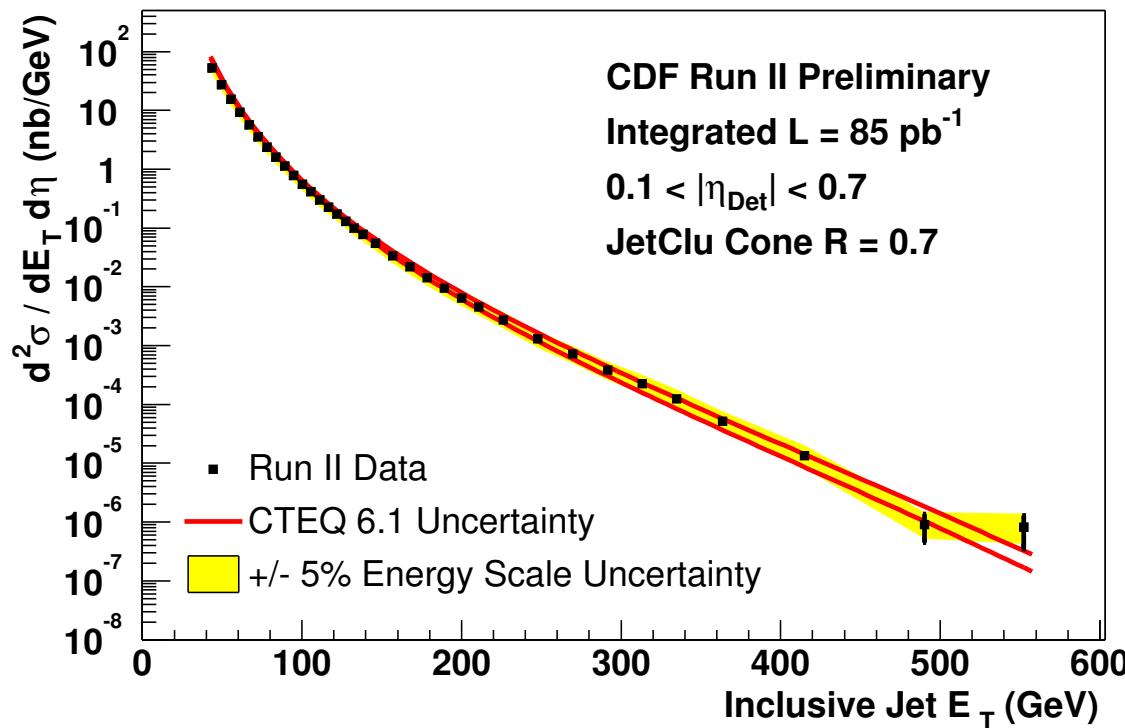
- Tevatron jet cross section at large E_T



$$\mathcal{L}_{\text{BSM}} \supseteq \frac{\tilde{g}^2}{M^2} \bar{\psi} \gamma^\mu \psi \bar{\psi} \gamma_\mu \psi \quad \Rightarrow \quad \frac{\text{data} - \text{theory}}{\text{theory}} \propto \tilde{g}^2 \frac{E_T^2}{M^2}$$

Why precision calculations for LHC physics?

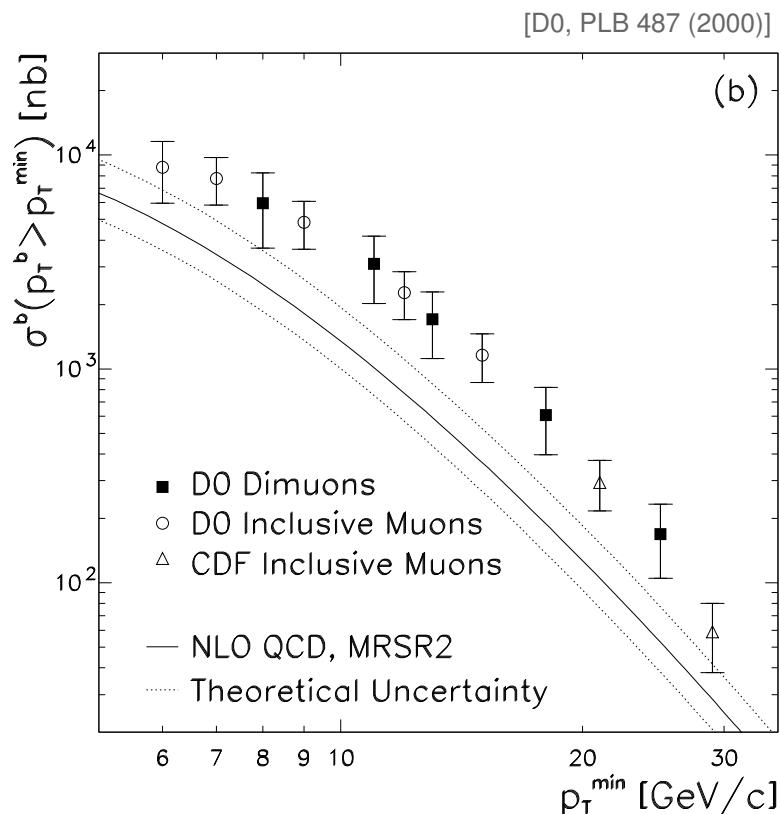
- Tevatron jet cross section at large E_T



⇒ just QCD... (uncertainty in gluon pdf at large x)

Why precision calculations for LHC physics?

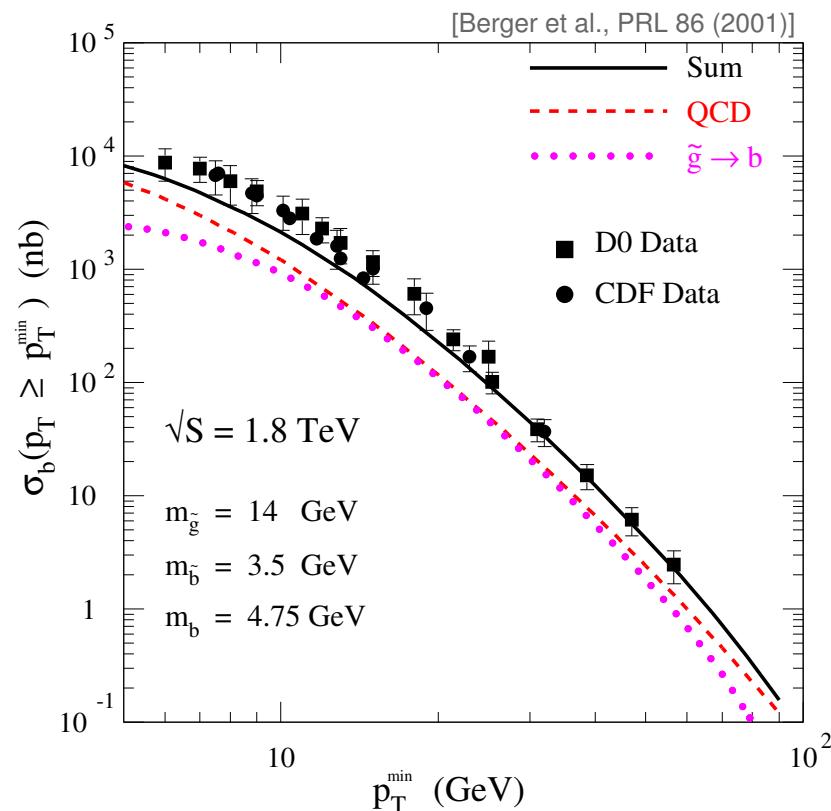
- Tevatron bottom quark cross section



⇒ new physics?

Why precision calculations for LHC physics?

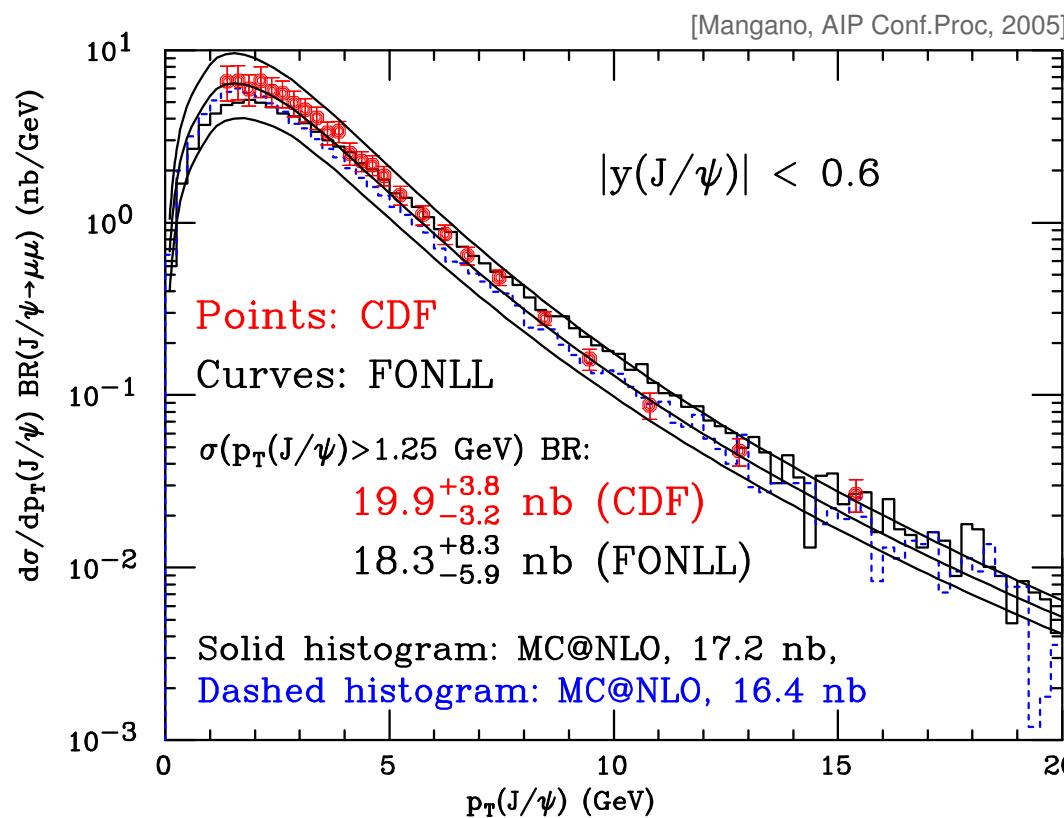
• Tevatron bottom quark cross section



⇒ light gluino/sbottom production?

Why precision calculations for LHC physics?

- Tevatron bottom quark cross section



⇒ just QCD... (α_s , low- x gluon pdf, $b \rightarrow B$ fragmentation & exp. analyses)

Why precision calculations for LHC physics?

- Signals of discovery include
 - mass peaks
 - anomalous shapes of kinematic distributions
 - excess of events after kinematic selection

Why precision calculations for LHC physics?

- Signals of discovery include
 - mass peaks
 - anomalous shapes of kinematic distributions
 - e.g. jet production at large E_T
 - excess of events after kinematic selection
 - e.g. bottom cross section

Why precision calculations for LHC physics?

- Signals of discovery include
 - mass peaks
e.g. Higgs searches in $H \rightarrow \gamma\gamma/ZZ^*$ final states
 - anomalous shapes of kinematic distributions
e.g. SUSY searches in multijet + $E_{T,\text{miss}}$ final states
 - excess of events after kinematic selection
e.g. Higgs searches in $H \rightarrow WW^*$ final states

Why precision calculations for LHC physics?

- Signals of discovery include
 - mass peaks
e.g. Higgs searches in $H \rightarrow \gamma\gamma/ZZ^*$ final states
 - anomalous shapes of kinematic distributions
e.g. SUSY searches in multijet + $E_{T,\text{miss}}$ final states
 - excess of events after kinematic selection
e.g. Higgs searches in $H \rightarrow WW^*$ final states
- Experiments measure multi-parton final states in restricted phase space region
 - Need flexible and realistic precision calculations
 - precise → including loops & many legs
 - flexible → allowing for exclusive cross sections & phase space cuts
 - realistic → matched with parton showers and hadronization

Contents

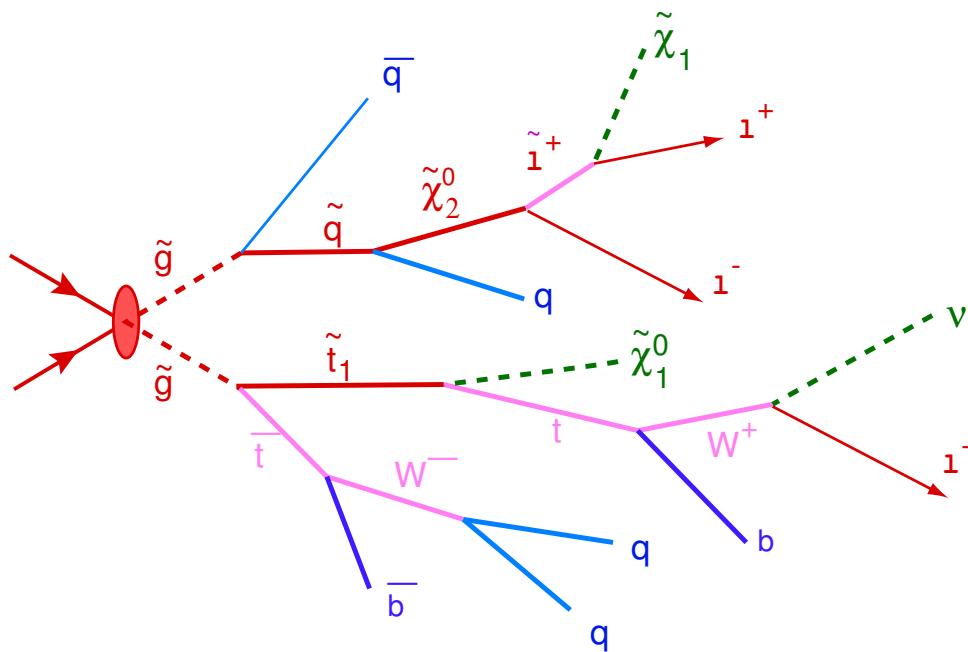
- Why precision calculations for LHC physics?
 - lessons from the past
 - prospects: SUSY searches at the LHC

BSM searches at the LHC

- Many model for new physics are addressing two problems:
 - the hierarchy/naturalness problem
 - the origin of dark matter
- spectrum of new particles at the TeV-scale with weakly interacting & stable particle

BSM searches at the LHC

- Many model for new physics are addressing two problems:
 - the hierarchy/naturalness problem
 - the origin of dark matter
- spectrum of new particles at the TeV-scale with weakly interacting & stable particle
- generic BSM signature at the LHC involves cascade decays with missing energy,
eg. supersymmetry

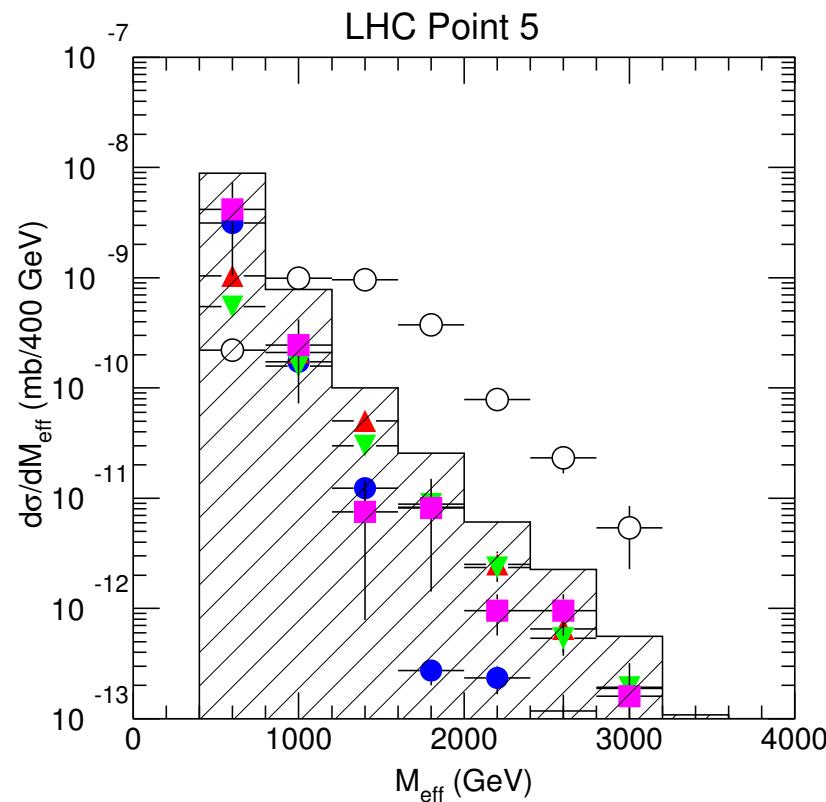


BSM searches at the LHC

- simple discriminant for SUSY searches: effective mass

$$M_{\text{eff}} = E_{\text{T,miss}} + \sum_{i=1}^4 p_{\text{T}}^{\text{jet}}$$

- excess of events with large M_{eff}
→ initial discovery of SUSY

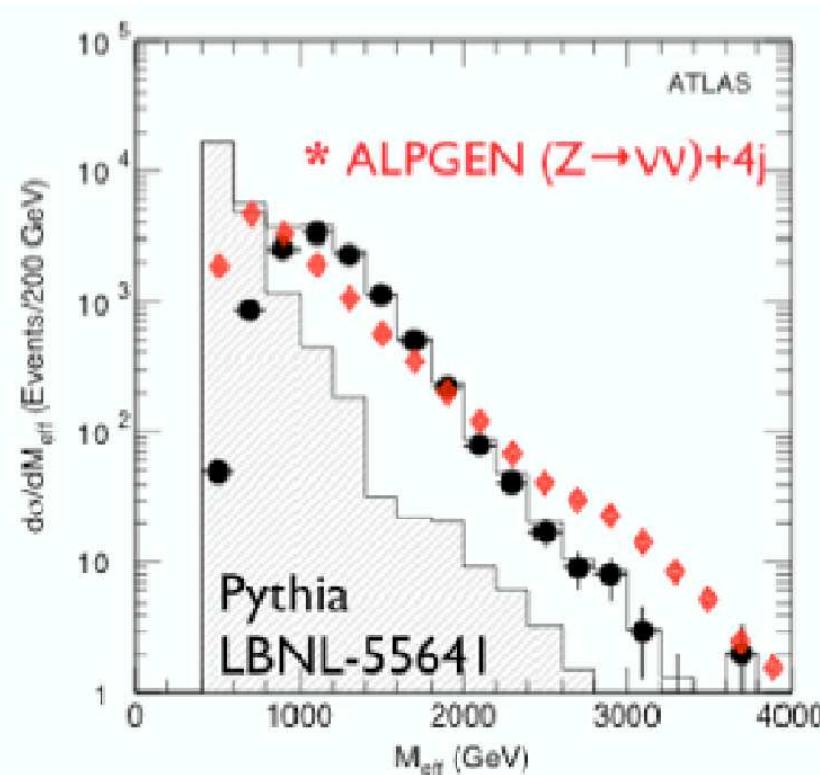


BSM searches at the LHC

- simple discriminant for SUSY searches: effective mass

$$M_{\text{eff}} = E_{\text{T,miss}} + \sum_{i=1}^4 p_{\text{T}}^{\text{jet}}$$

- excess of events with large M_{eff}
→ initial discovery of SUSY



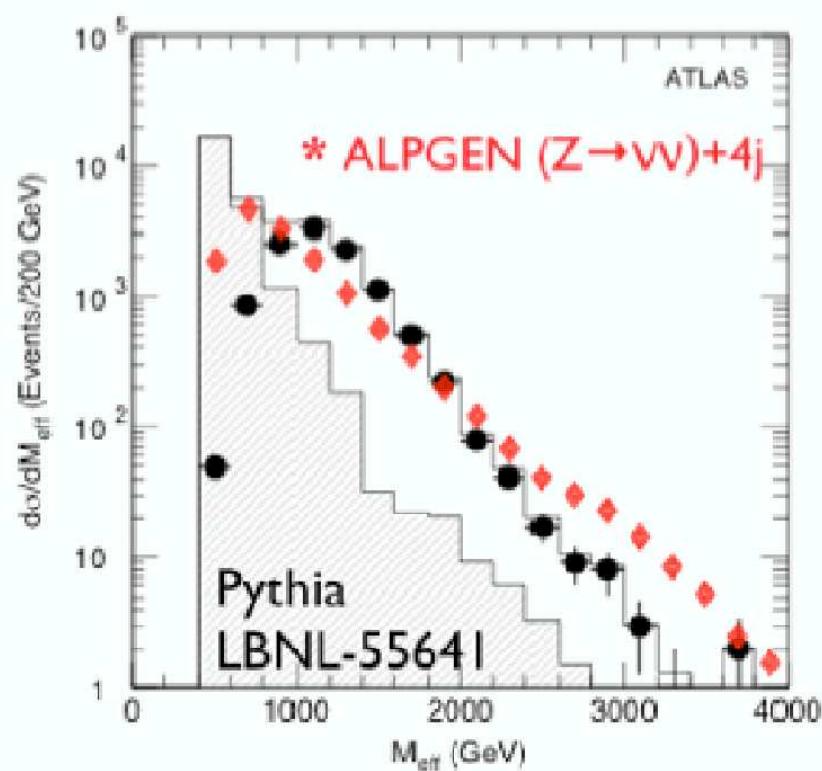
- Modern LO Monte Carlo tools predict much larger multi-jet background

BSM searches at the LHC

- simple discriminant for SUSY searches: effective mass

$$M_{\text{eff}} = E_{\text{T,miss}} + \sum_{i=1}^4 p_{\text{T}}^{\text{jet}}$$

- excess of events with large M_{eff}
→ initial discovery of SUSY



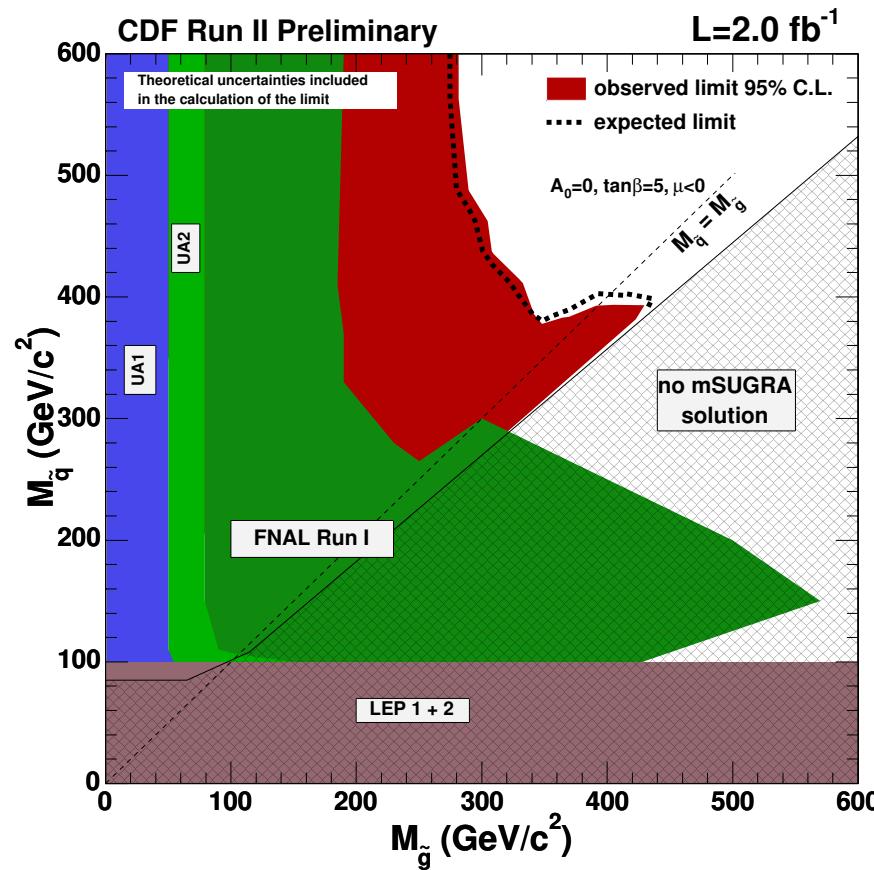
- Modern LO Monte Carlo tools predict much larger multi-jet background
- What will disagreement between Monte Carlo and data mean?
→ Need precision NLO calculations to tell!

Precision calculations for BSM physics at the LHC

- no excess over SM expectations
 - exclude models / set limits on model parameters
- excess in inclusive jets + $E_{T,\text{miss}}$ signal (early LHC phase)
 - discriminate BSM models
- exploring BSM models (later LHC phase)
 - determine masses & spins

Precision calculations for BSM physics at the LHC

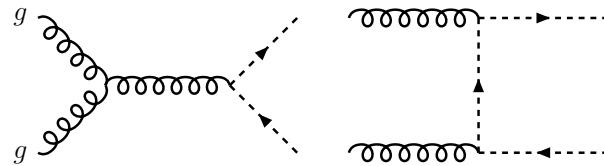
- no excess over SM expectations
→ exclude models / set limits on model parameters



→ needs accurate theoretical predictions

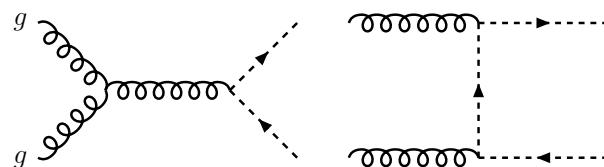
SUSY particle production at hadron colliders

SUSY particles would be produced copiously at hadron colliders via QCD processes, eg.



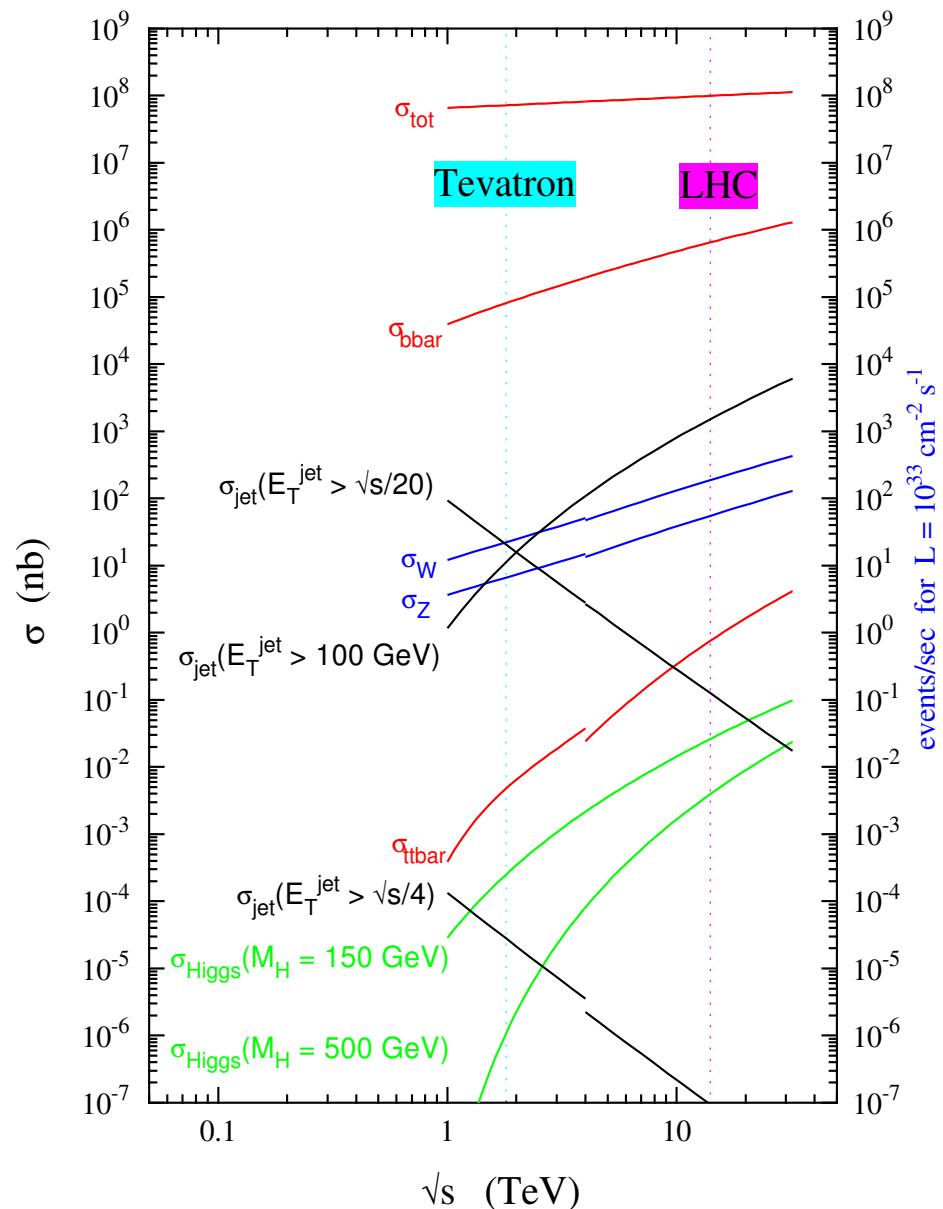
SUSY particle production at hadron colliders

SUSY particles would be produced copiously at hadron colliders via QCD processes, eg.



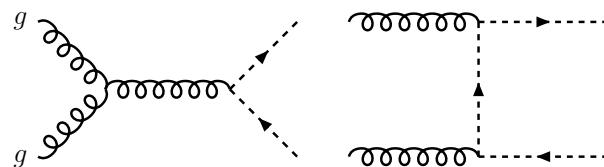
SUSY signal LHC →

$$\sigma(\tilde{q}\tilde{q} + \tilde{g}\tilde{g} + \tilde{g}\tilde{q}) \approx 2 \text{ nb} \quad (M_{\tilde{q}, \tilde{g}} \approx 300 \text{ GeV})$$



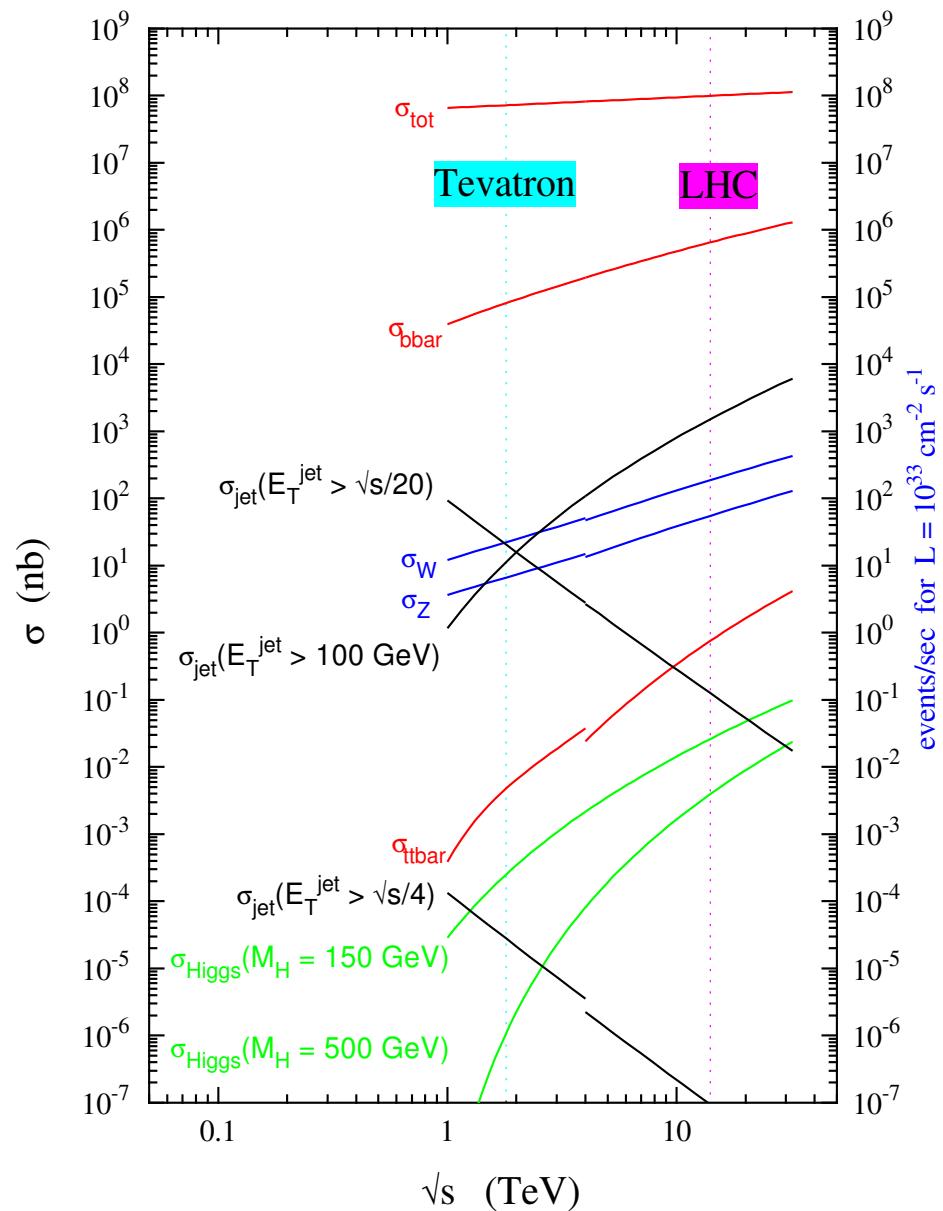
SUSY particle production at hadron colliders

SUSY particles would be produced copiously at hadron colliders via QCD processes, eg.



SUSY signal Tevatron →

$$\sigma(\tilde{q}\tilde{q} + \tilde{g}\tilde{g} + \tilde{g}\tilde{q}) \approx 2 \text{ nb} \quad (M_{\tilde{q}, \tilde{g}} \approx 300 \text{ GeV})$$



SUSY particle production at hadron colliders

● MSSM sparticle pair production at NLO (Beenakker, Höpker, MK, Plehn, Spira, Zerwas)

— squarks & gluinos $pp/p\bar{p} \rightarrow \tilde{q}\bar{\tilde{q}}, \tilde{g}\bar{\tilde{g}}, \tilde{q}\bar{g}$

— stops $pp/p\bar{p} \rightarrow \tilde{t}\bar{\tilde{t}}$

— gauginos $pp/p\bar{p} \rightarrow \tilde{\chi}^0\tilde{\chi}^0, \tilde{\chi}^\pm\tilde{\chi}^0, \tilde{\chi}^+\tilde{\chi}^-$

— sleptons $pp/p\bar{p} \rightarrow \tilde{l}\bar{\tilde{l}}$

— associated production $pp/p\bar{p} \rightarrow \tilde{q}\tilde{\chi}, \tilde{g}\tilde{\chi}$

References of LO calculations:

Kane, Leveillé, '82; Harrison, Llewellyn Smith, '83; Reya, Roy, '85; Dawson, Eichten, Quigg, '85; Baer, Tata, '85, ...

SUSY particle production at hadron colliders

● MSSM sparticle pair production at NLO (Beenakker, Höpker, MK, Plehn, Spira, Zerwas)

— squarks & gluinos $pp/p\bar{p} \rightarrow \tilde{q}\bar{\tilde{q}}, \tilde{g}\bar{\tilde{g}}, \tilde{q}\bar{g}$

— stops $pp/p\bar{p} \rightarrow \tilde{t}\bar{\tilde{t}}$

— gauginos $pp/p\bar{p} \rightarrow \tilde{\chi}^0\tilde{\chi}^0, \tilde{\chi}^\pm\tilde{\chi}^0, \tilde{\chi}^+\tilde{\chi}^-$

— sleptons $pp/p\bar{p} \rightarrow \tilde{l}\bar{\tilde{l}}$

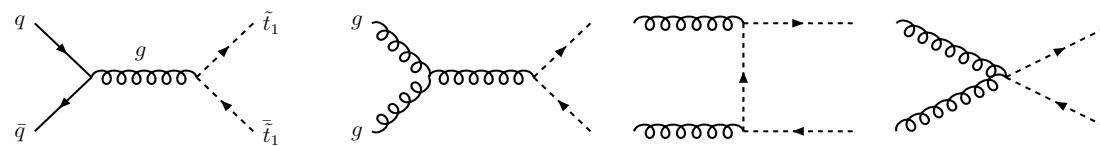
— associated production $pp/p\bar{p} \rightarrow \tilde{q}\tilde{\chi}, \tilde{g}\tilde{\chi}$

References of LO calculations:

Kane, Leveillé, '82; Harrison, Llewellyn Smith, '83; Reya, Roy, '85; Dawson, Eichten, Quigg, '85; Baer, Tata, '85, ...

Example: Top-squark production

● LO graphs

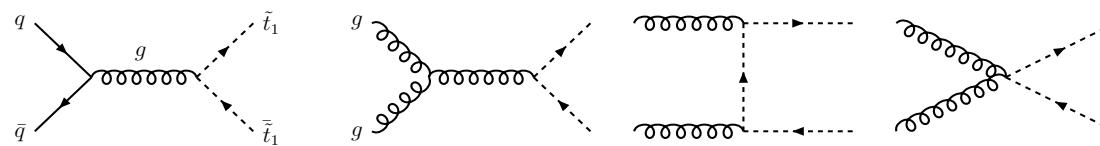


$$\hat{\sigma}_{\text{LO}}[gg] = \frac{\alpha_s^2 \pi}{s} \left[\beta \left(\frac{5}{48} + \frac{31m^2}{24s} \right) + \left(\frac{2m^2}{3s} + \frac{m^4}{6s^2} \right) \log \frac{1-\beta}{1+\beta} \right] \quad (\beta^2 = 1 - 4m^2/s)$$

⇒ no MSSM parameter dependence

Example: Top-squark production

- LO graphs



$$\hat{\sigma}_{\text{LO}}[gg] = \frac{\alpha_s^2 \pi}{s} \left[\beta \left(\frac{5}{48} + \frac{31m^2}{24s} \right) + \left(\frac{2m^2}{3s} + \frac{m^4}{6s^2} \right) \log \frac{1-\beta}{1+\beta} \right] \quad (\beta^2 = 1 - 4m^2/s)$$

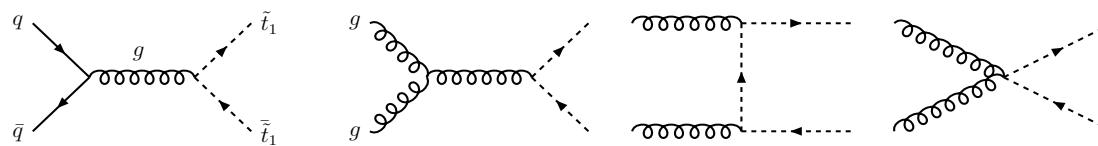
⇒ no MSSM parameter dependence

- pp cross section: $\sigma = F_{\text{non-pert.}}(\mu) \otimes (C_0 \alpha_s^B(\mu) + C_1(\mu) \alpha_s^{B+1}(\mu) + \dots)_{\text{pert.}}$

Scale dependence through factorization of IR contributions and renormalization of UV contributions

Example: Top-squark production

- LO graphs



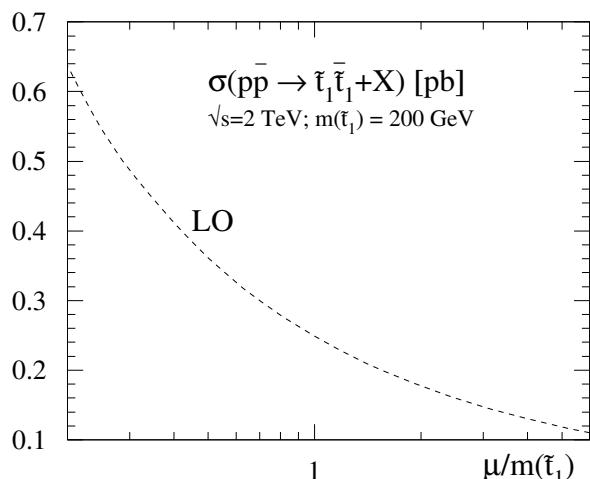
$$\hat{\sigma}_{\text{LO}}[gg] = \frac{\alpha_s^2 \pi}{s} \left[\beta \left(\frac{5}{48} + \frac{31m^2}{24s} \right) + \left(\frac{2m^2}{3s} + \frac{m^4}{6s^2} \right) \log \frac{1-\beta}{1+\beta} \right] \quad (\beta^2 = 1 - 4m^2/s)$$

⇒ no MSSM parameter dependence

- pp cross section: $\sigma = F_{\text{non-pert.}}(\mu) \otimes (C_0 \alpha_s^B(\mu) + C_1(\mu) \alpha_s^{B+1}(\mu) + \dots)_{\text{pert.}}$

Scale dependence through factorization of IR contributions and renormalization of UV contributions

- LO scale dependence

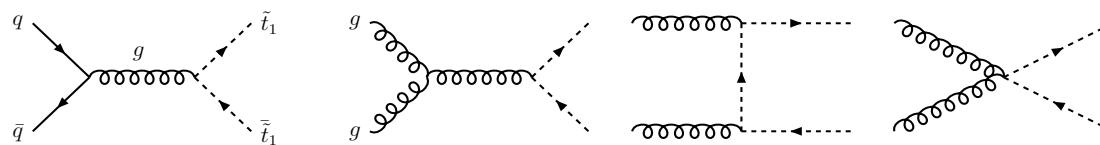


→ theoretical uncertainty $\approx \pm 100\%$ at LO

⇒ must include NLO corrections

Example: Top-squark production

- LO graphs



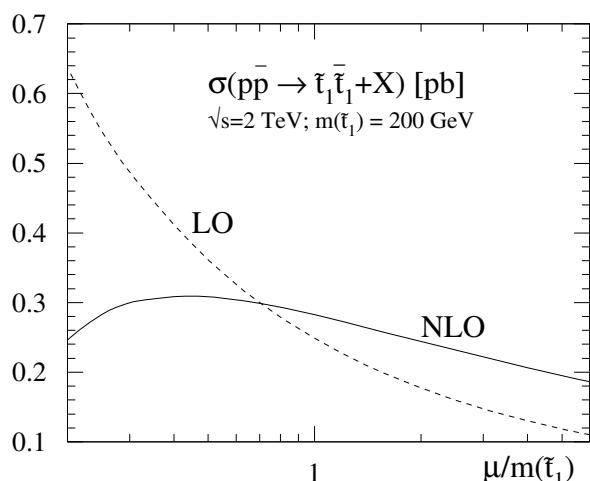
$$\hat{\sigma}_{\text{LO}}[gg] = \frac{\alpha_s^2 \pi}{s} \left[\beta \left(\frac{5}{48} + \frac{31m^2}{24s} \right) + \left(\frac{2m^2}{3s} + \frac{m^4}{6s^2} \right) \log \frac{1-\beta}{1+\beta} \right] \quad (\beta^2 = 1 - 4m^2/s)$$

⇒ no MSSM parameter dependence

- pp cross section: $\sigma = F_{\text{non-pert.}}(\mu) \otimes (C_0 \alpha_s^B(\mu) + C_1(\mu) \alpha_s^{B+1}(\mu) + \dots)_{\text{pert.}}$

Scale dependence through factorization of IR contributions and renormalization of UV contributions

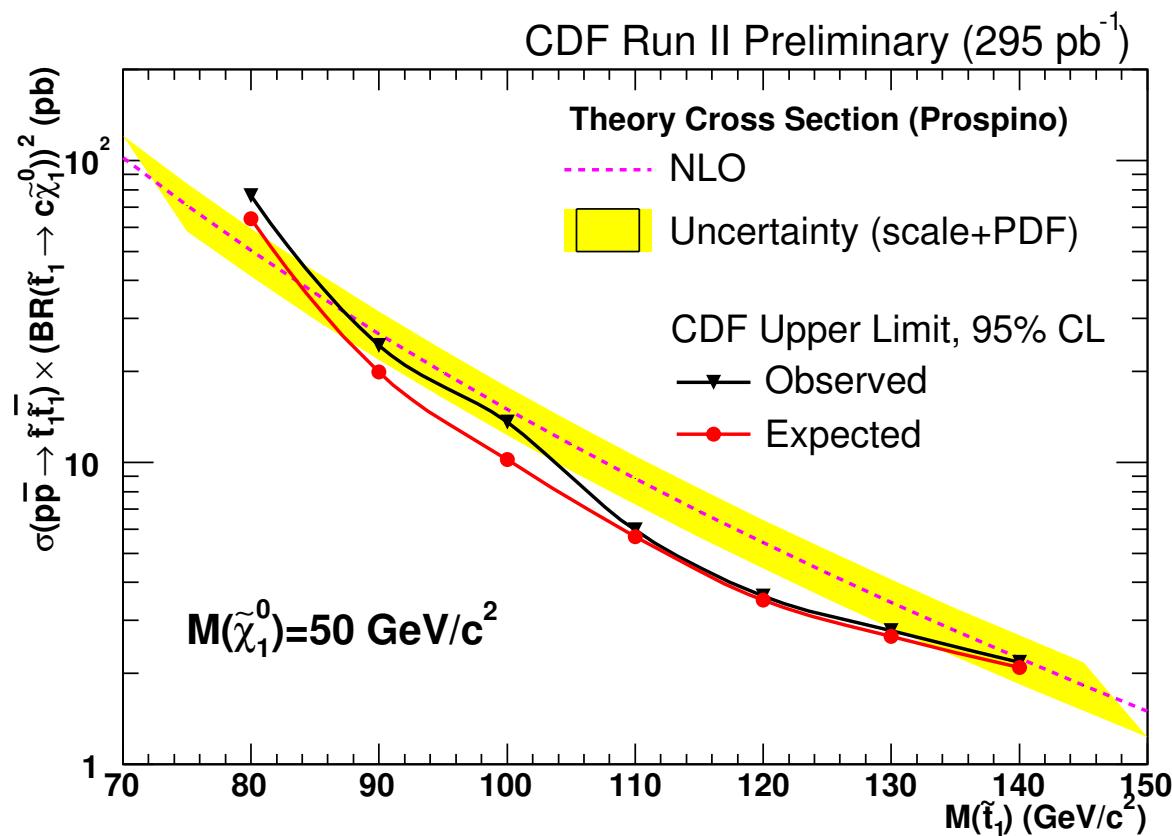
- NLO scale dependence (Beenakker, MK, Plehn, Spira, Zerwas)



→ theoretical uncertainty $\approx \pm 15\%$ at NLO

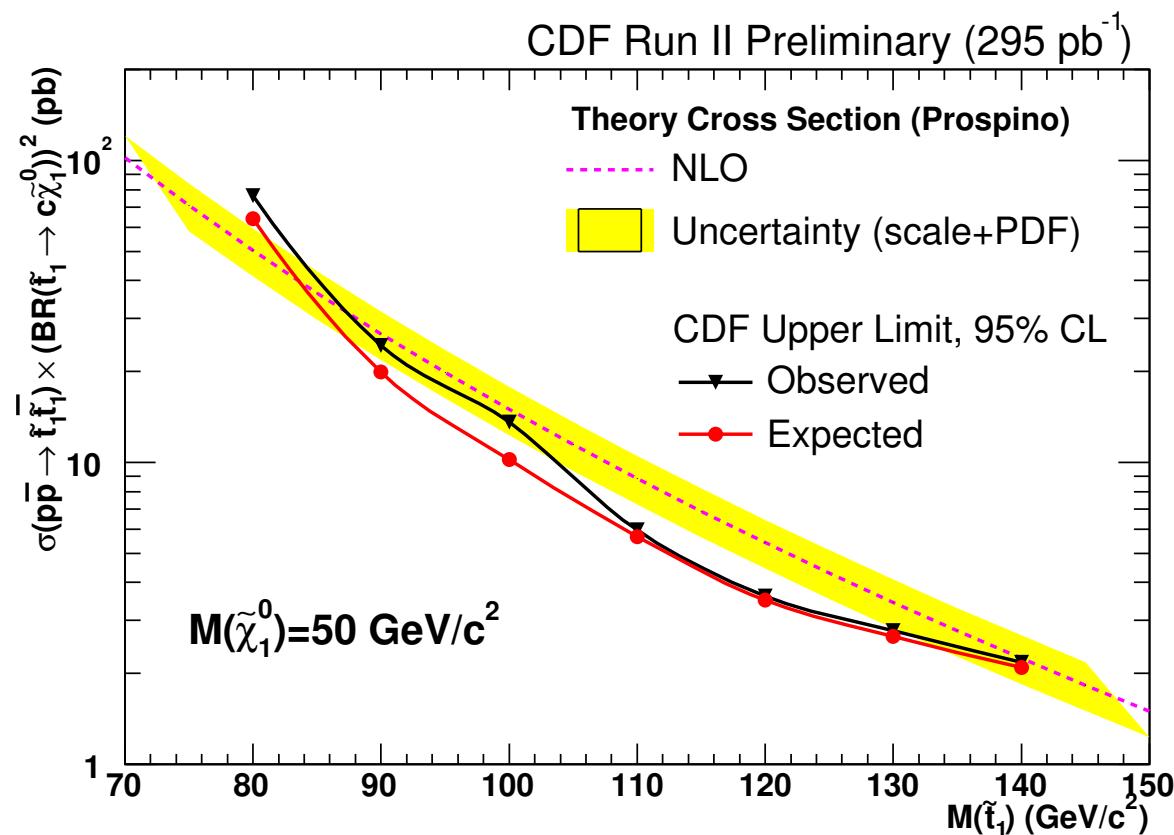
Top-squark searches

- Top-squark search in $p\bar{p} \rightarrow \tilde{t}_1 \tilde{t}_1 \rightarrow c\tilde{\chi}_1^0 \bar{c}\tilde{\chi}_1^0$ channel (CDF PRD 2007)



Top-squark searches

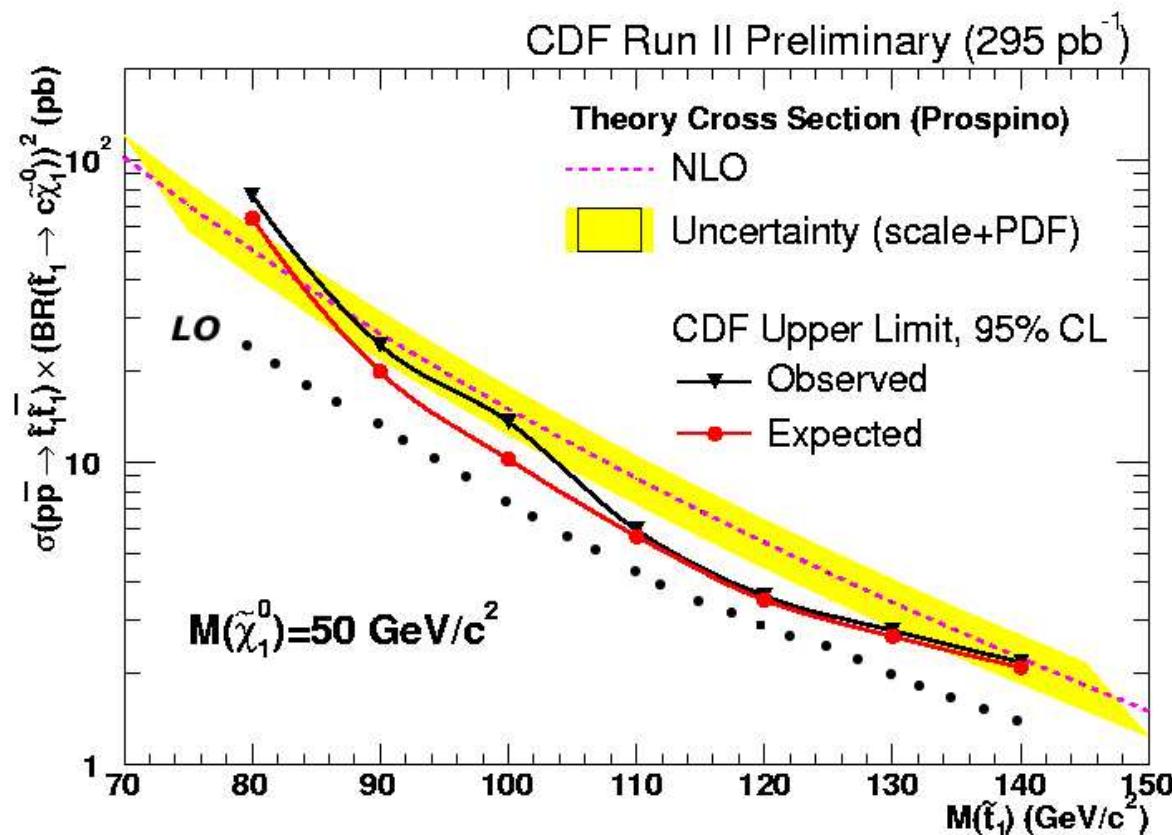
- Top-squark search in $p\bar{p} \rightarrow \tilde{t}_1 \tilde{t}_1 \rightarrow c\tilde{\chi}_1^0 \bar{c}\tilde{\chi}_1^0$ channel (CDF PRD 2007)



$$\Rightarrow 105 \text{ GeV} \leq M_{\tilde{t}_1} \leq 130 \text{ GeV}$$

Top-squark searches

- Top-squark search in $p\bar{p} \rightarrow \tilde{t}_1 \tilde{t}_1 \rightarrow c\tilde{\chi}_1^0 \bar{c}\tilde{\chi}_1^0$ channel (CDF PRD 2007)

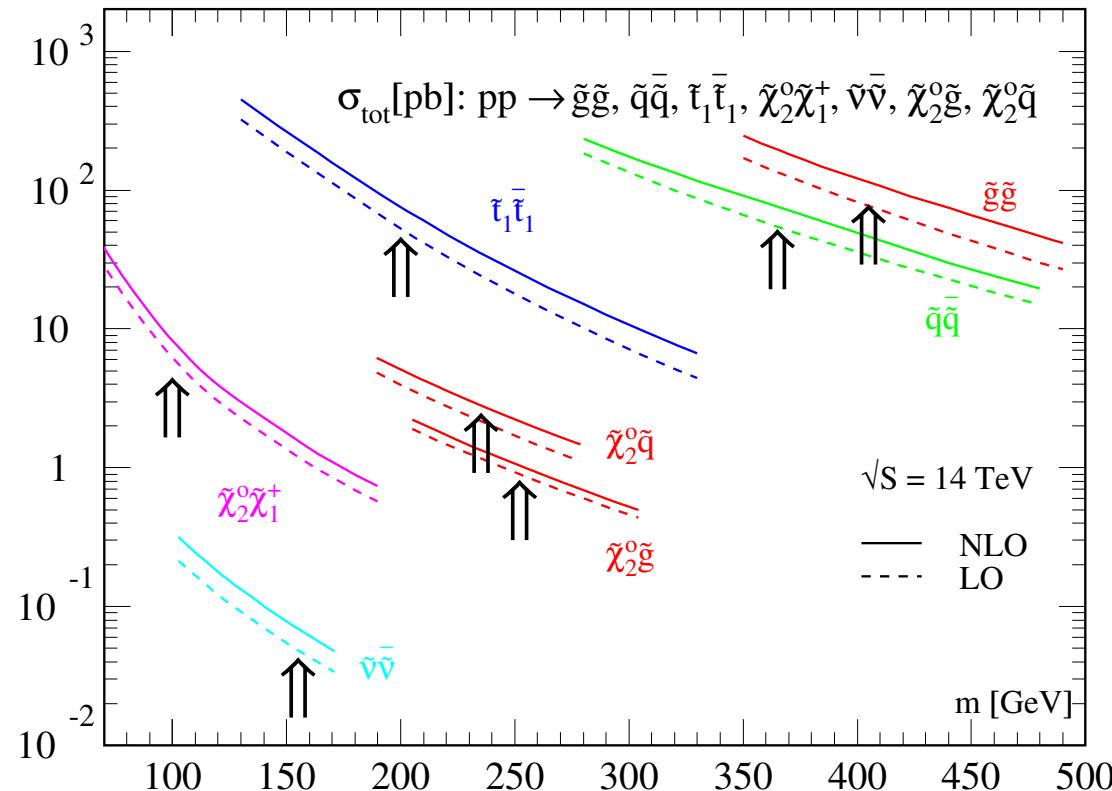


→ no exclusion based on LO cross sections

MSSM particle production at the LHC

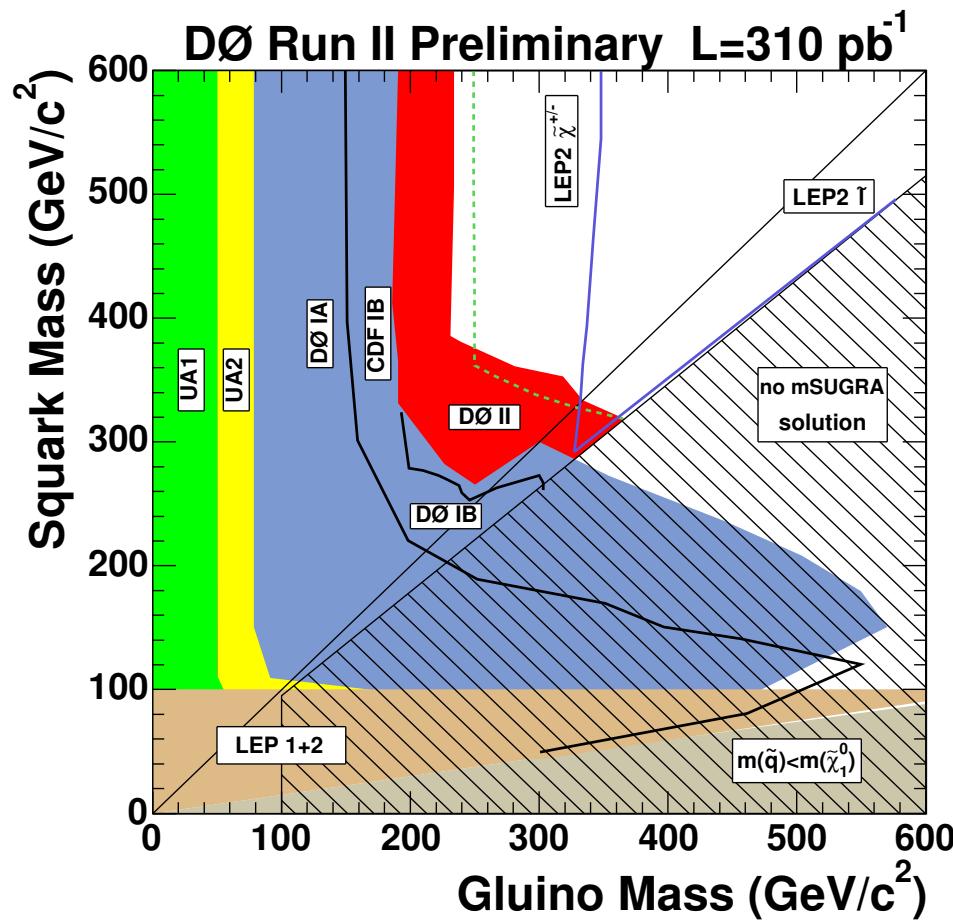
● NLO SUSY-QCD corrections for MSSM particle production at hadron colliders

→ public code PROSPINO (Beenakker, Höpker, MK, Plehn, Spira, Zerwas)



References: Beenakker, Höpker, Spira, Zerwas, 1995, 1997; Beenakker, MK, Plehn, Spira, Zerwas, 1998; Baer, Hall, Reno, 1998; Beenakker, Klasen, MK, Plehn, Spira, Zerwas, 1999; Beenakker, MK, Plehn, Spira, Zerwas, 2000; Berger, Klasen, Tait, 1999-2002; Beenakker, MK, Plehn, Spira, Zerwas, 2007

Current limits on sparticle masses



● mass limits (roughly)

$M_{\tilde{g}}$	$\gtrsim 250$ GeV
$M_{\tilde{q}} \approx M_{\text{gluino}}$	$\gtrsim 400$ GeV
$M_{\tilde{t}_1}$	$\gtrsim 100$ GeV
$M_{\tilde{\chi}_1^0}$	$\gtrsim 50$ GeV
$M_{\tilde{\chi}_1^\pm}$	$\gtrsim 100$ GeV
M_{sleptons}	$\gtrsim 100$ GeV

Precision calculations for BSM physics at the LHC

- no excess over SM expectations
 - exclude models / set limits on model parameters
- excess in inclusive jets + $E_{T,\text{miss}}$ signal (early LHC phase)
 - discriminate BSM models
- exploring BSM models (later LHC phase)
 - determine masses & spins

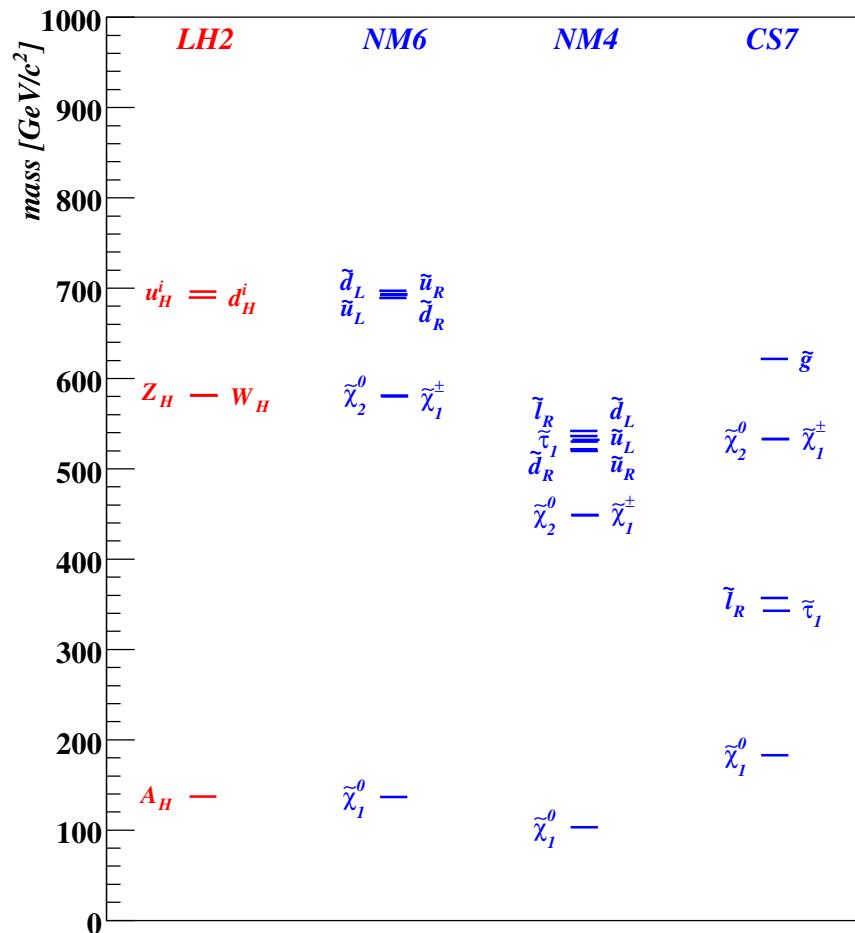
Precision calculations for BSM physics at the LHC

- excess in inclusive jets + $E_{T,\text{miss}}$ signal (early LHC phase)
→ discriminate BSM models

Precision calculations for BSM physics at the LHC

- excess in inclusive jets + $E_{T,\text{miss}}$ signal (early LHC phase)
→ discriminate BSM models

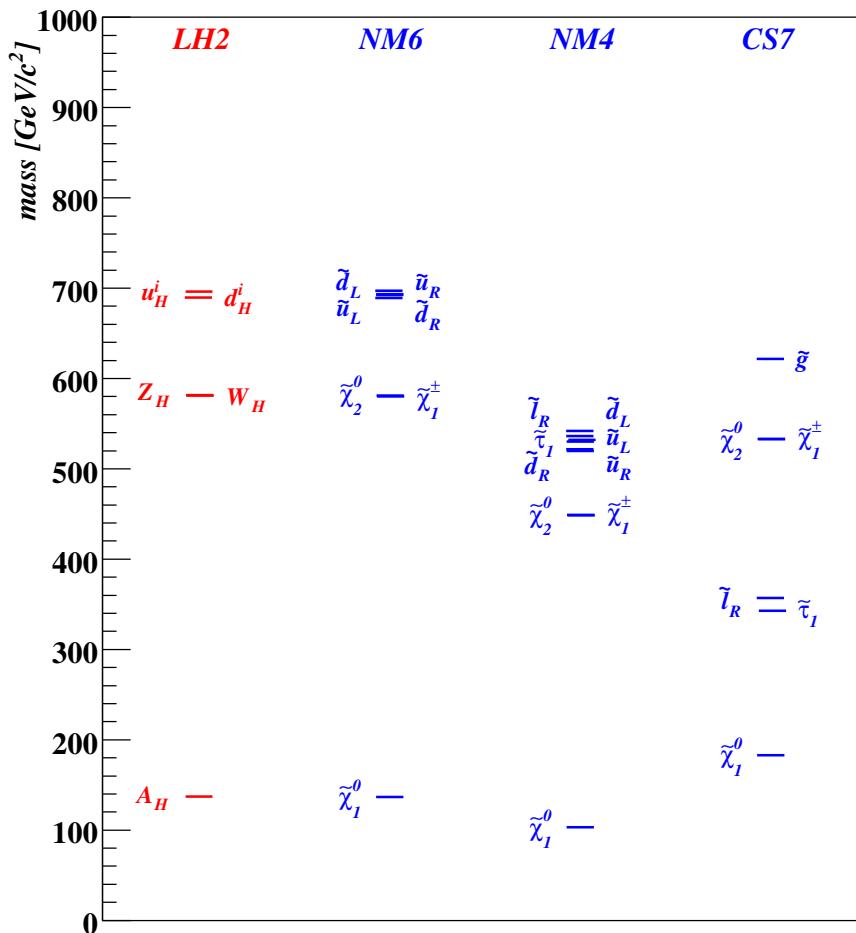
“look-alike models” (Hubisz, Lykken, Pierini, Spiropulu)



Precision calculations for BSM physics at the LHC

- excess in inclusive jets + $E_{T,\text{miss}}$ signal (early LHC phase)
→ discriminate BSM models

“look-alike models” (Hubisz, Lykken, Pierini, Spiropulu)



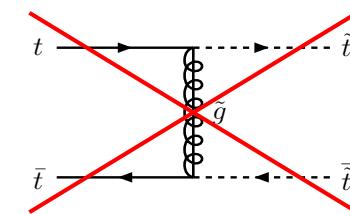
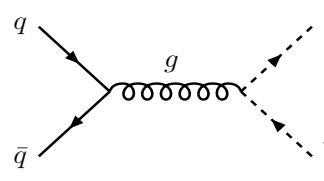
- Little Higgs model LH2 and SUSY models NM4 and CS7 give same number of events after cuts
- “twin models” LH2 and NM6 (SUSY) have different cross sections and event counts
→ distinguish models with same spectrum but different spins through event count

Precision calculations for BSM physics at the LHC

- excess in inclusive jets + $E_{T,\text{miss}}$ signal (early LHC phase)
→ discriminate BSM models

“ancient wisdom”: cross sections depend on spin, e.g.

$$q + \bar{q} \rightarrow \tilde{q} + \bar{\tilde{q}}$$



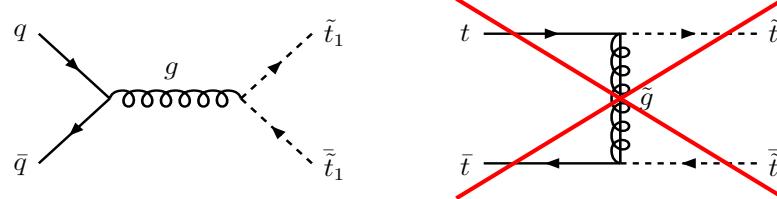
$$\hat{\sigma}_{\text{LO}}[q\bar{q} \rightarrow \tilde{q}\bar{\tilde{q}}] = \frac{\alpha_s^2 \pi}{s} \frac{4}{27} \beta^3 \quad (\beta^2 = 1 - 4m^2/s)$$

Precision calculations for BSM physics at the LHC

- excess in inclusive jets + $E_{T,\text{miss}}$ signal (early LHC phase)
→ discriminate BSM models

“ancient wisdom”: cross sections depend on spin, e.g.

$$q + \bar{q} \rightarrow \tilde{q} + \bar{\tilde{q}}$$



$$\hat{\sigma}_{\text{LO}}[q\bar{q} \rightarrow \tilde{q}\bar{\tilde{q}}] = \frac{\alpha_s^2 \pi}{s} \frac{4}{27} \beta^3 \quad (\beta^2 = 1 - 4m^2/s)$$

c.f. top production:

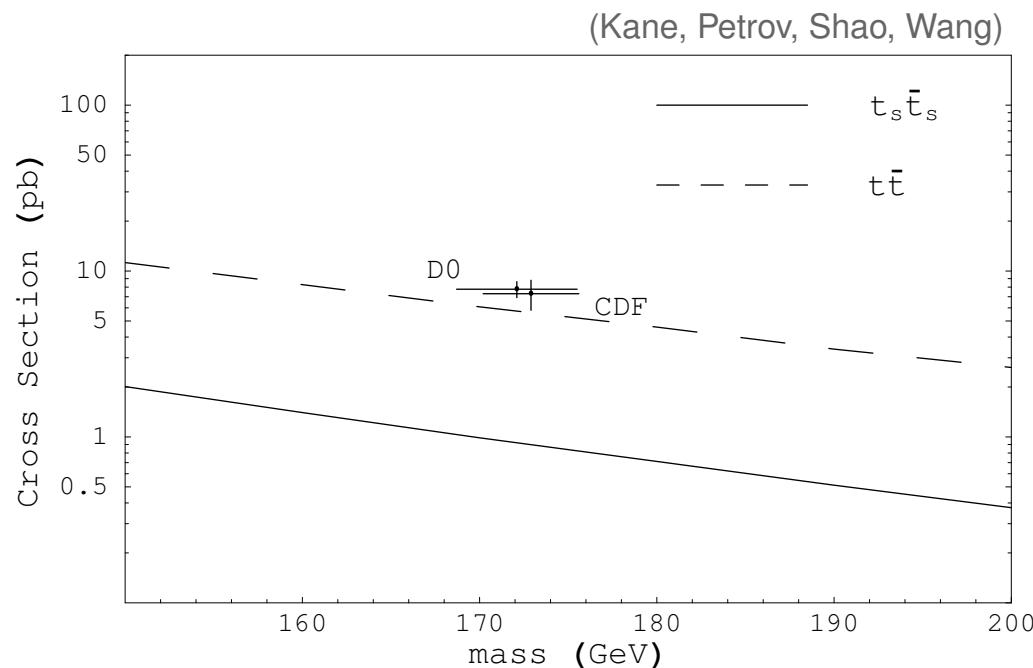
$$\hat{\sigma}_{\text{LO}}[q\bar{q} \rightarrow Q\bar{Q}] = \frac{\alpha_s^2 \pi}{s} \frac{8}{27} \frac{(s + 2m^2)}{s^2} \beta$$

→ $\sigma^{\text{top}}/\sigma^{\text{stop}} \sim 10$ at the Tevatron

Precision calculations for BSM physics at the LHC

- excess in inclusive jets + $E_{T,\text{miss}}$ signal (early LHC phase)
→ discriminate BSM models

“ancient wisdom”: cross sections depend on spin

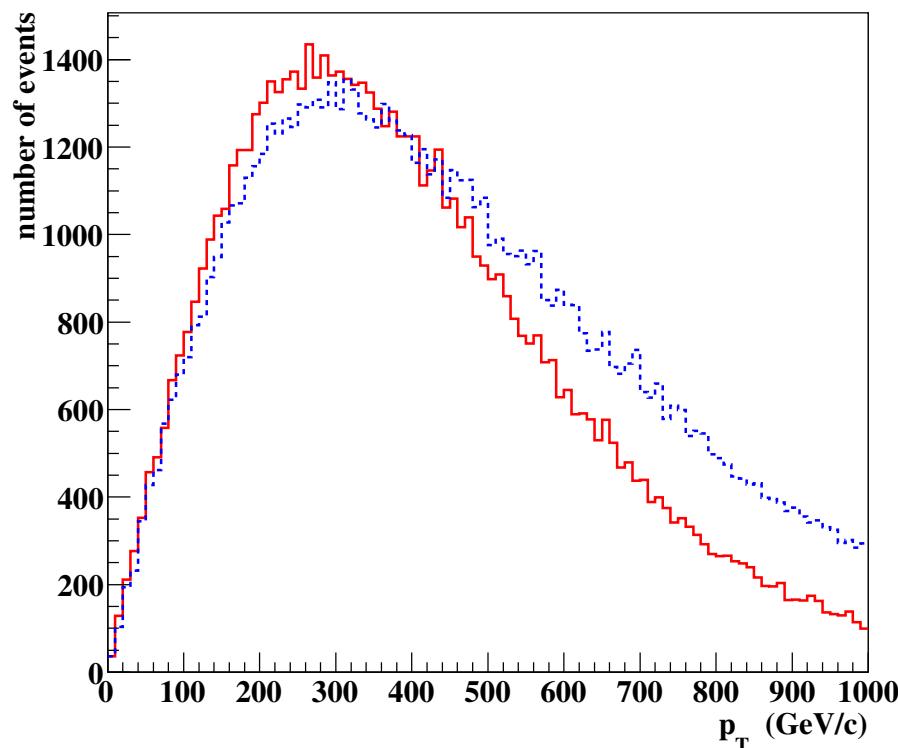


- no production of scalar quarks (assuming same mass and couplings as top...)

Precision calculations for BSM physics at the LHC

- excess in inclusive jets + $E_{T,\text{miss}}$ signal (early LHC phase)
→ discriminate BSM models

how to discriminate “look-alike models”? (Hubisz, Lykken, Pierini, Spiropulu)

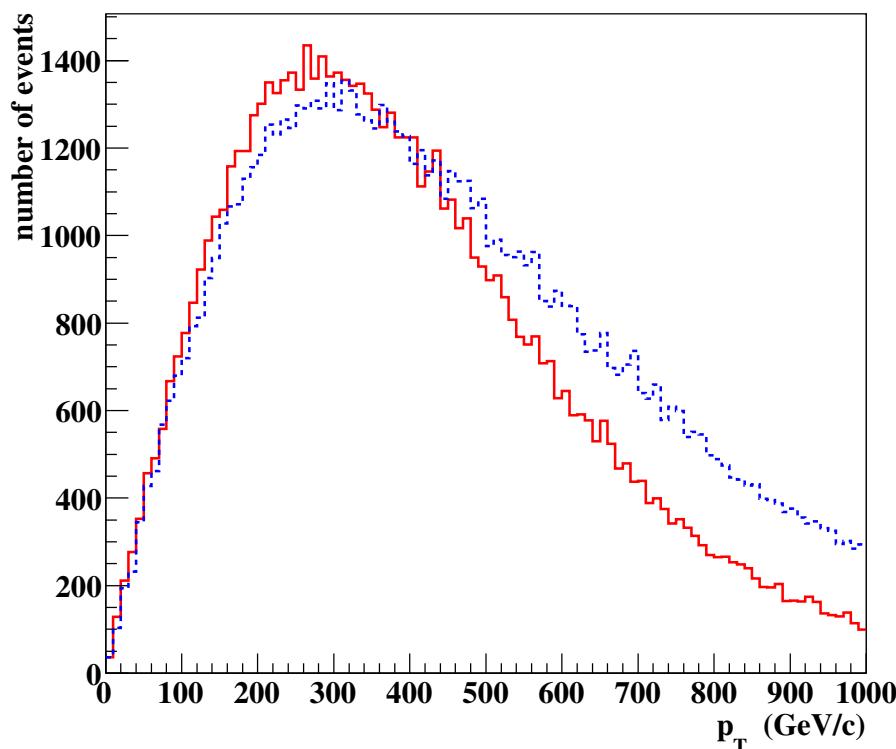


- consider ratios of event counts, e.g.
 $\sigma(p_T > 250, 500, \dots \text{ GeV})/\sigma$
→ systematic uncertainties cancel
→ normalization important

Precision calculations for BSM physics at the LHC

- excess in inclusive jets + $E_{T,\text{miss}}$ signal (early LHC phase)
→ discriminate BSM models

how to discriminate “look-alike models”? (Hubisz, Lykken, Pierini, Spiropulu)



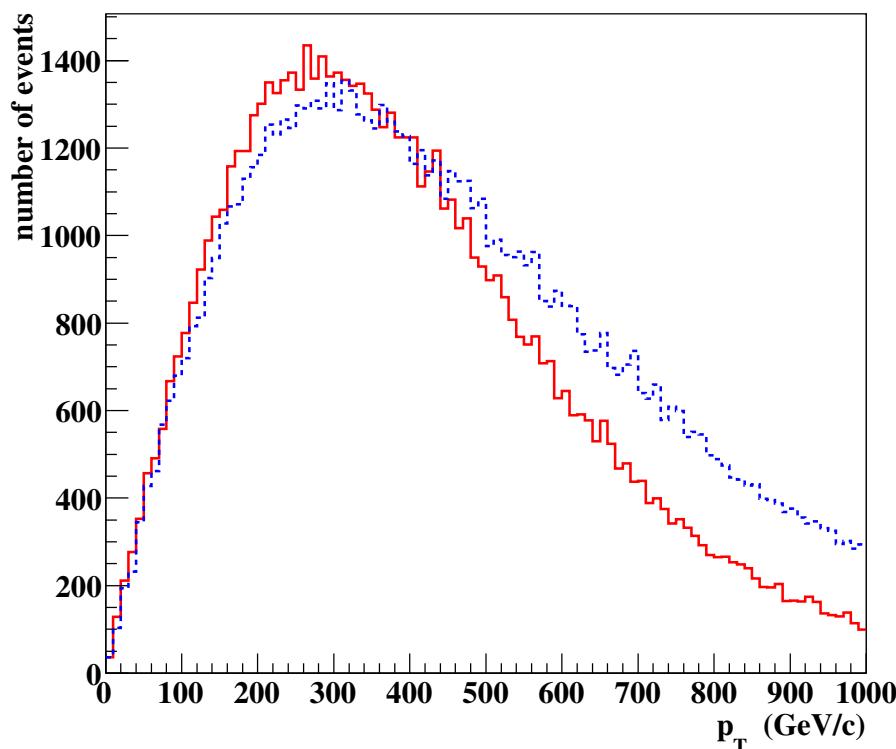
- consider ratios of event counts, e.g.
 $\sigma(p_T > 250, 500, \dots \text{ GeV})/\sigma$
→ systematic uncertainties cancel
→ normalization important
- NLO affects shape of distributions:

$$K \equiv \frac{\sigma^{\text{NLO}}(pp \rightarrow \tilde{t}\bar{t})}{\sigma^{\text{LO}}(pp \rightarrow \tilde{t}\bar{t})} = \begin{cases} 1.4 & (p_T > 100 \text{ GeV}) \\ 0.5 & (p_T > 1000 \text{ GeV}) \end{cases}$$

Precision calculations for BSM physics at the LHC

- excess in inclusive jets + $E_{T,\text{miss}}$ signal (early LHC phase)
→ discriminate BSM models

how to discriminate “look-alike models”? (Hubisz, Lykken, Pierini, Spiropulu)



- consider ratios of event counts, e.g.
 $\sigma(p_T > 250, 500, \dots \text{ GeV})/\sigma$
→ systematic uncertainties cancel
→ normalization important
- NLO affects shape of distributions:

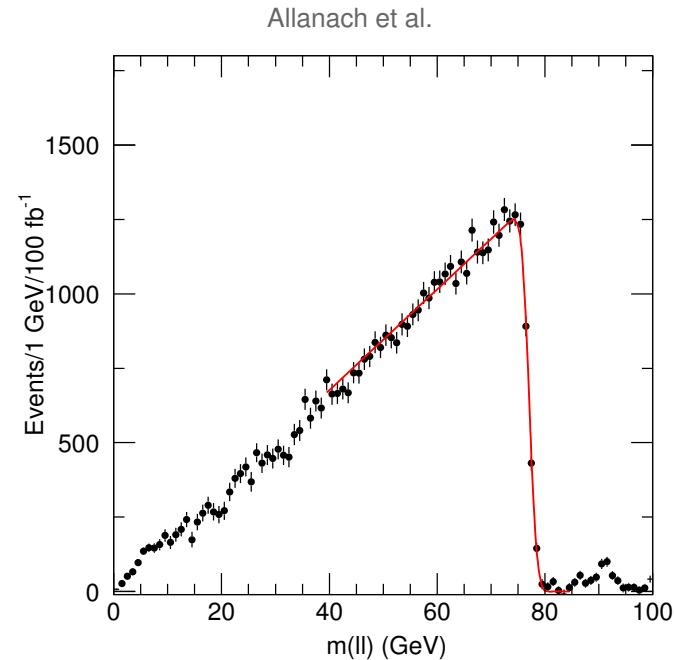
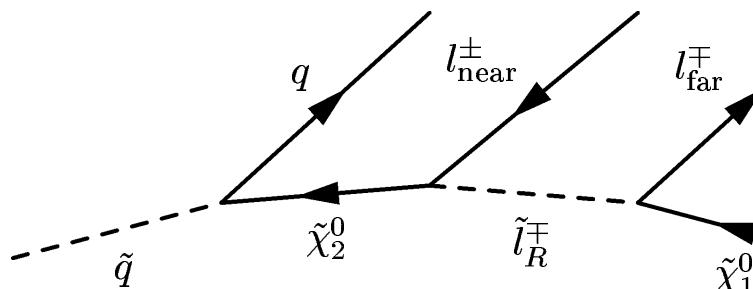
$$K \equiv \frac{\sigma^{\text{NLO}}(pp \rightarrow \tilde{t}\bar{t})}{\sigma^{\text{LO}}(pp \rightarrow \tilde{t}\bar{t})} = \begin{cases} 1.4 & (p_T > 100 \text{ GeV}) \\ 0.5 & (p_T > 1000 \text{ GeV}) \end{cases}$$

→ no proper tools to perform realistic NLO analysis of BSM decay chains

Precision calculations for BSM physics at the LHC

- exploring BSM models (later LHC phase)
→ determine masses & spins

Mass measurements from cascade decays, e.g.

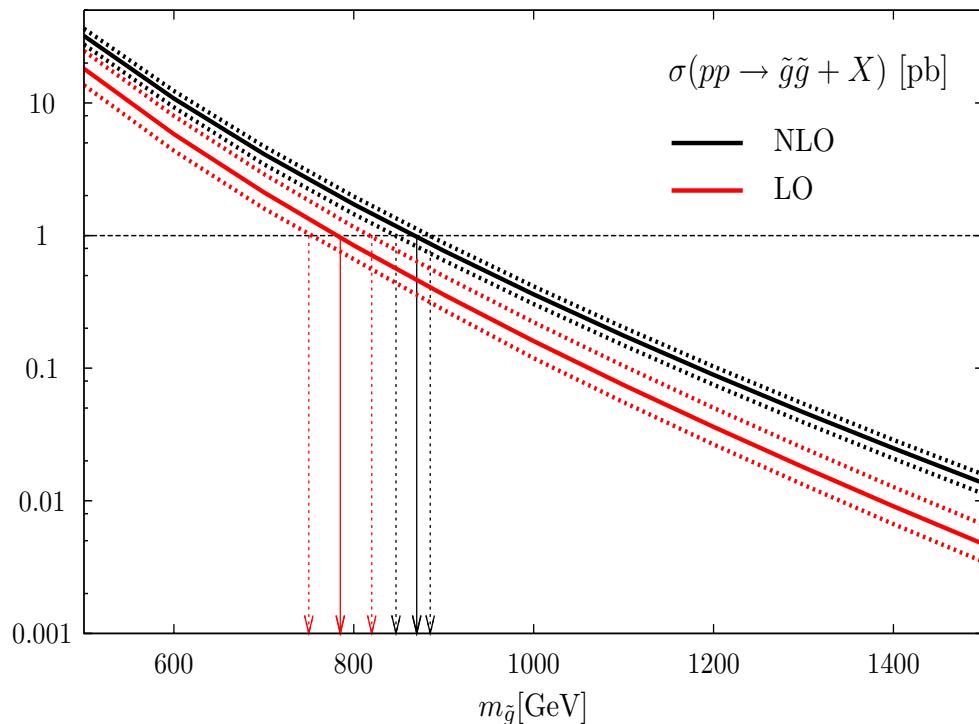


→ kinematic endpoint $(m_{ll}^{\max})^2 = (m_{\tilde{\chi}_2^0}^2 - m_{\tilde{l}_R^\pm}^2)(m_{\tilde{l}_R^\pm}^2 - m_{\tilde{\chi}_1^0}^2)/(m_{\tilde{l}_R^\pm}^2)$
sensitive to mass differences

Precision calculations for BSM physics at the LHC

- exploring BSM models (later LHC phase)
→ determine masses & spins

Mass measurements from total rate, e.g. gluino mass in SUSY models with heavy scalars (FP dark matter scenarios)

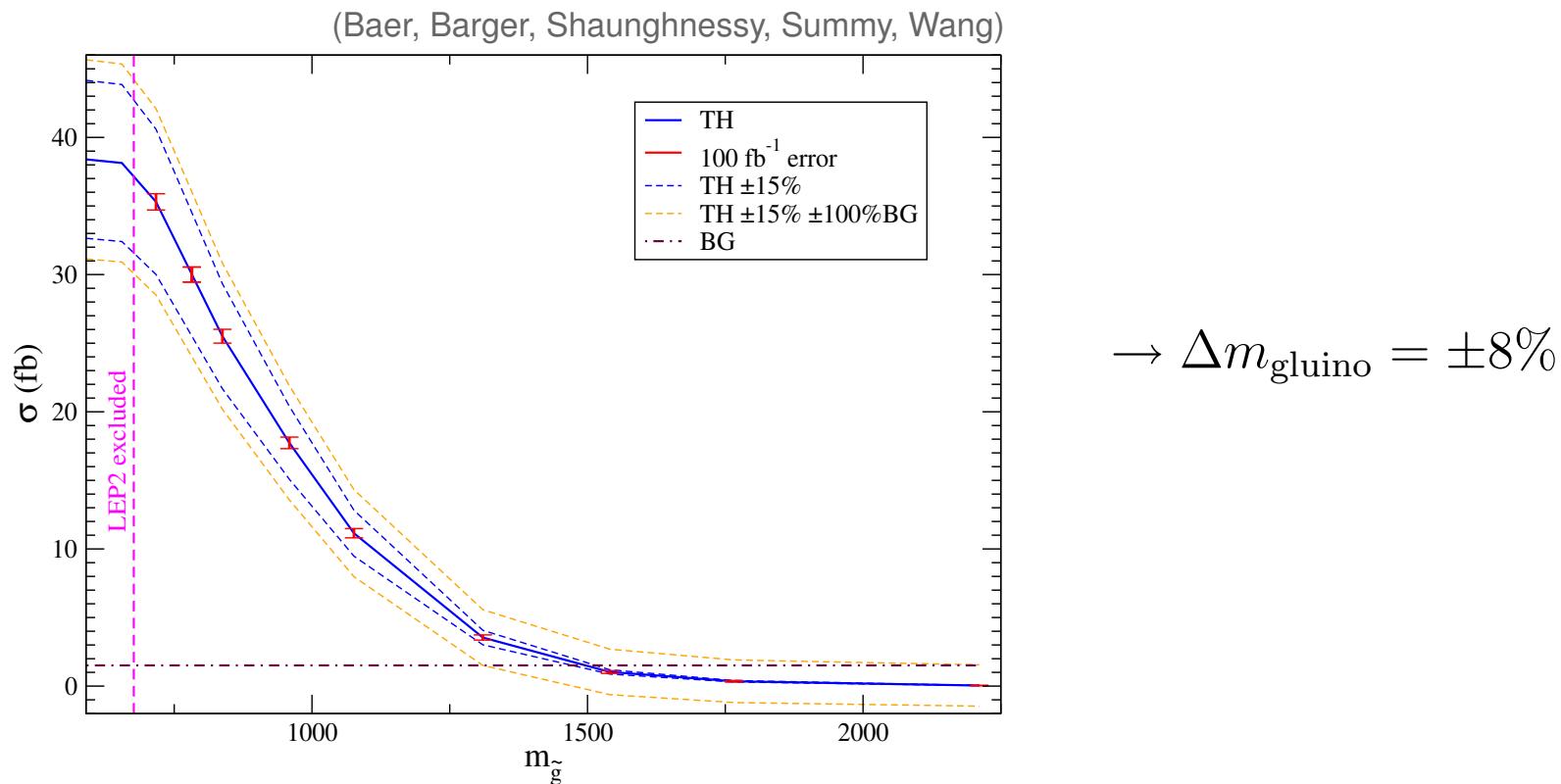


$$m_{\text{gluino}} = \begin{cases} 785 \pm 35 & (\text{LO}) \\ 870 \pm 15 & (\text{NLO}) \end{cases}$$

Precision calculations for BSM physics at the LHC

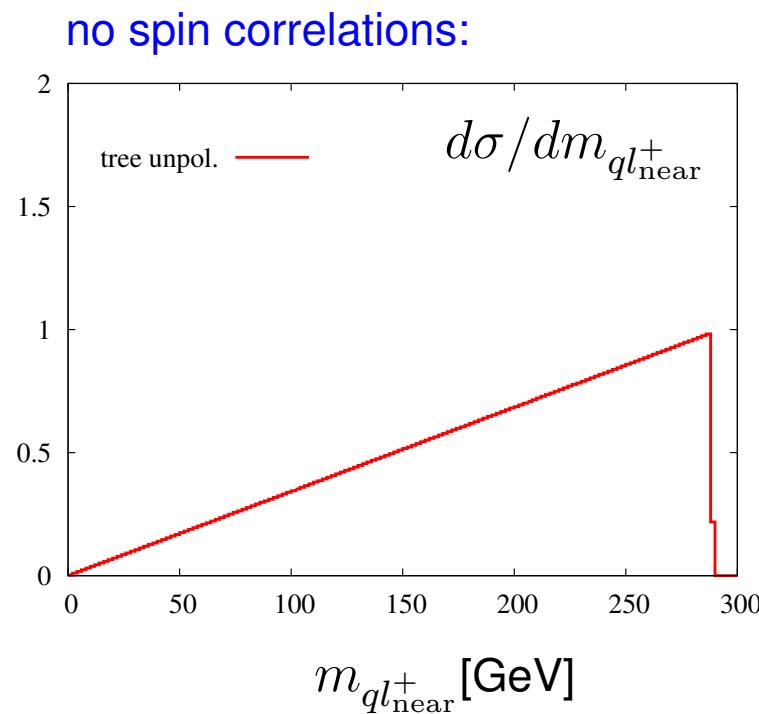
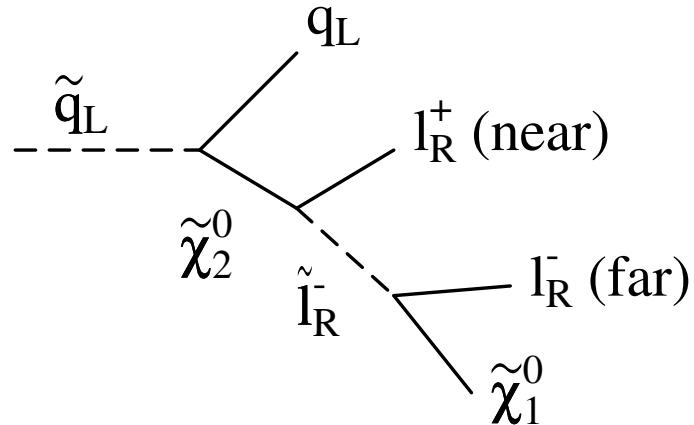
- exploring BSM models (later LHC phase)
→ determine masses & spins

Mass measurements from total rate, e.g. gluino mass in SUSY models with heavy scalars (FP dark matter scenarios)



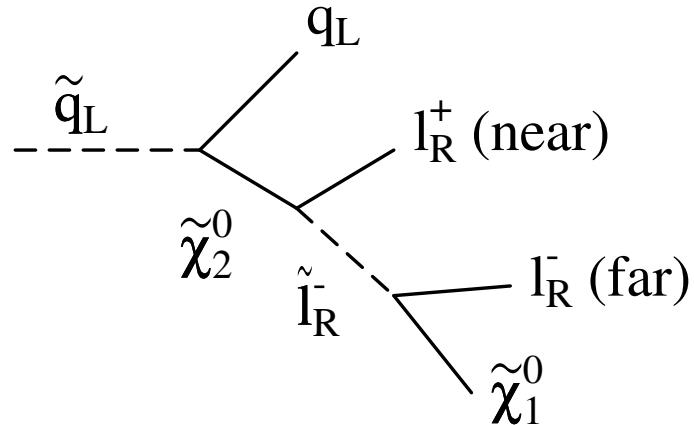
Precision calculations for BSM physics at the LHC

- exploring BSM models (later LHC phase)
→ determine masses & spins

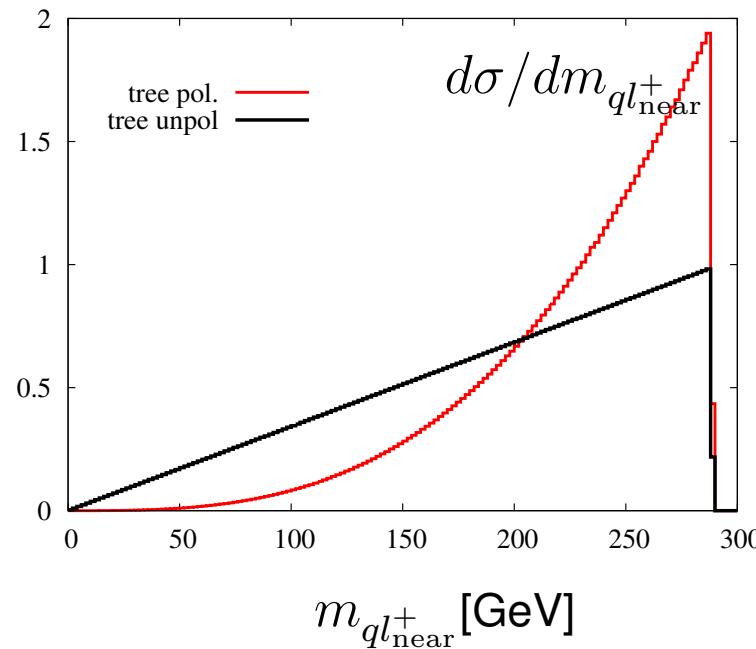


Precision calculations for BSM physics at the LHC

- exploring BSM models (later LHC phase)
→ determine masses & spins

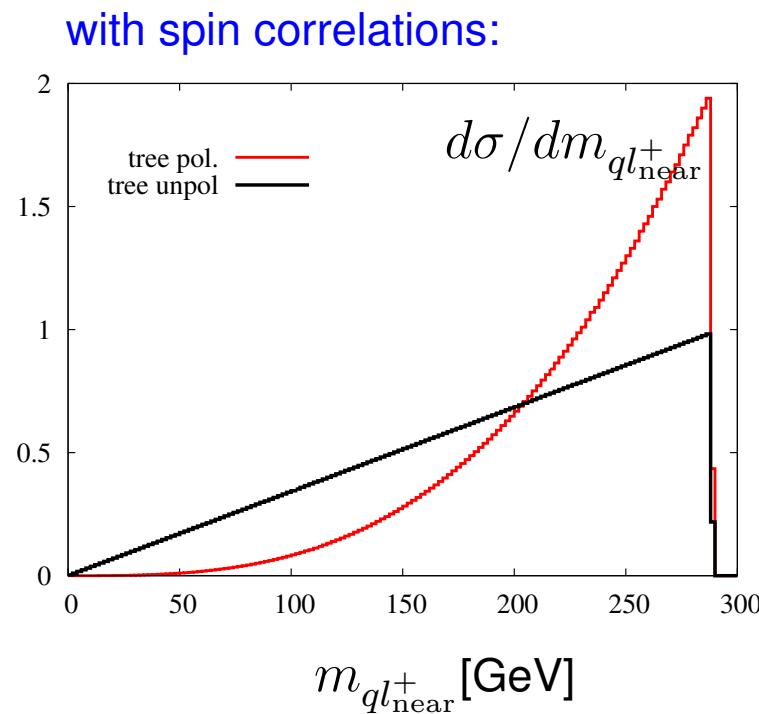
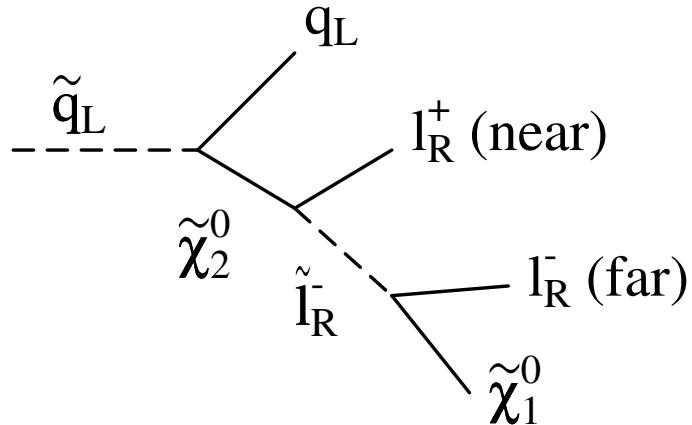


with spin correlations:



Precision calculations for BSM physics at the LHC

- exploring BSM models (later LHC phase)
→ determine masses & spins

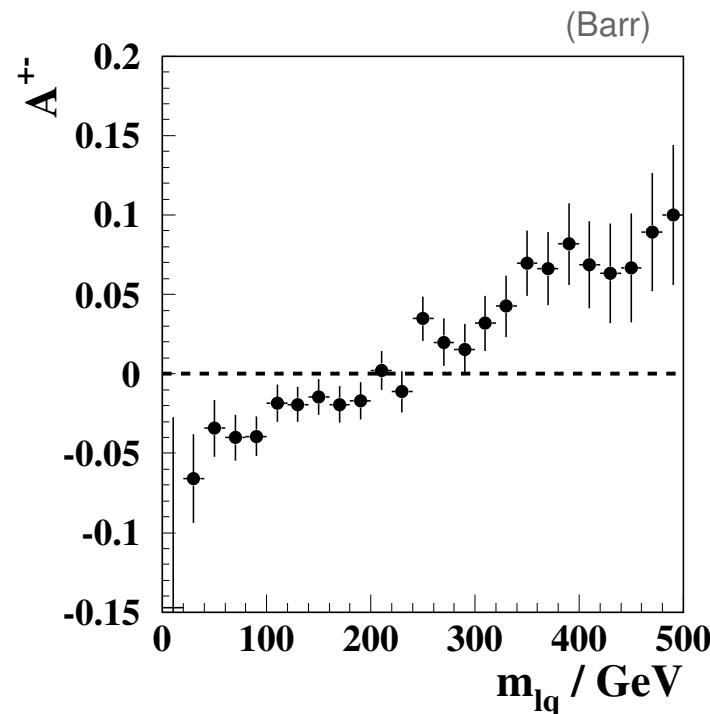
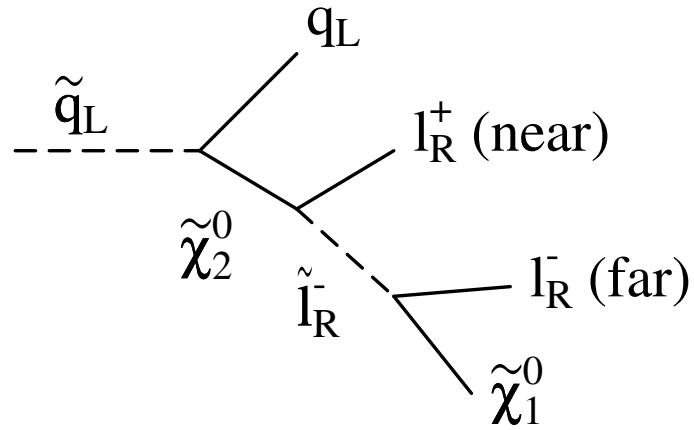


cannot determine $m_{ql_{\text{near}}^+}$ distribution experimentally

→ consider charge asymmetry $A = (d\sigma/dm_{ql^+} - d\sigma/dm_{ql^-})/(d\sigma/dm_{ql^+} + d\sigma/dm_{ql^-})$

Precision calculations for BSM physics at the LHC

- exploring BSM models (later LHC phase)
→ determine masses & spins

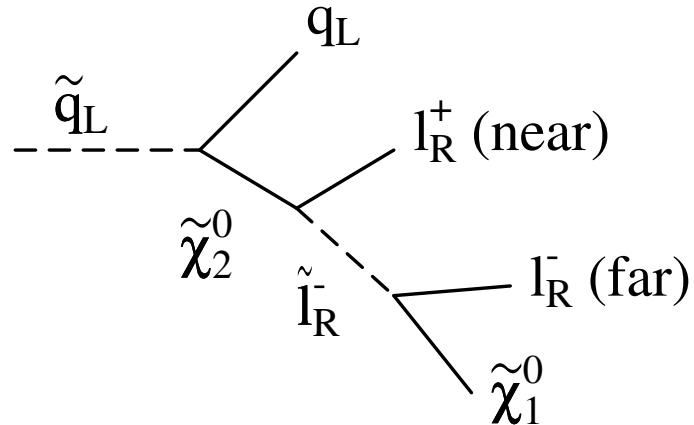


cannot determine $m_{ql_{\text{near}}^+}$ distribution experimentally

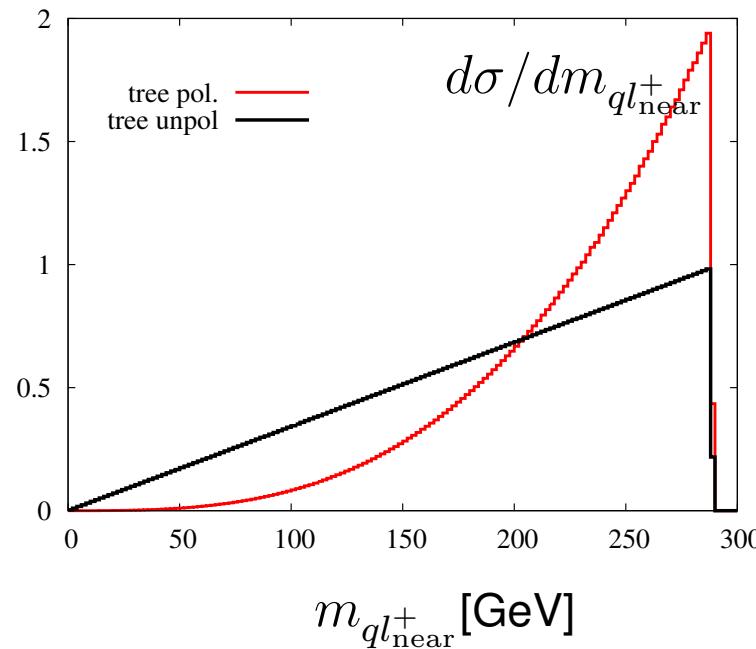
→ consider charge asymmetry $A = (d\sigma/dm_{ql^+} - d\sigma/dm_{ql^-})/(d\sigma/dm_{ql^+} + d\sigma/dm_{ql^-})$

Precision calculations for BSM physics at the LHC

- exploring BSM models (later LHC phase)
→ determine masses & spins

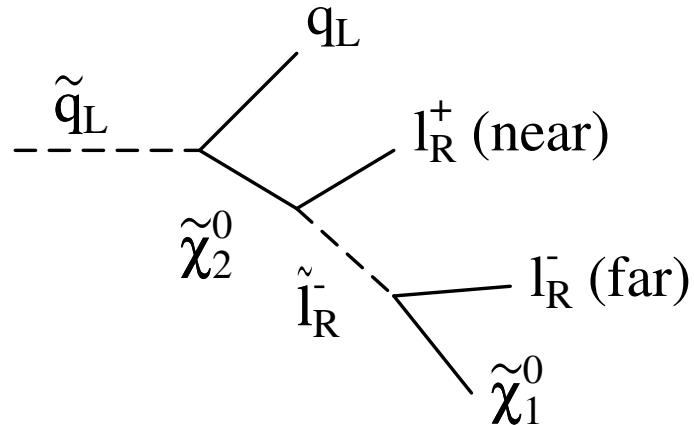


with spin correlations:

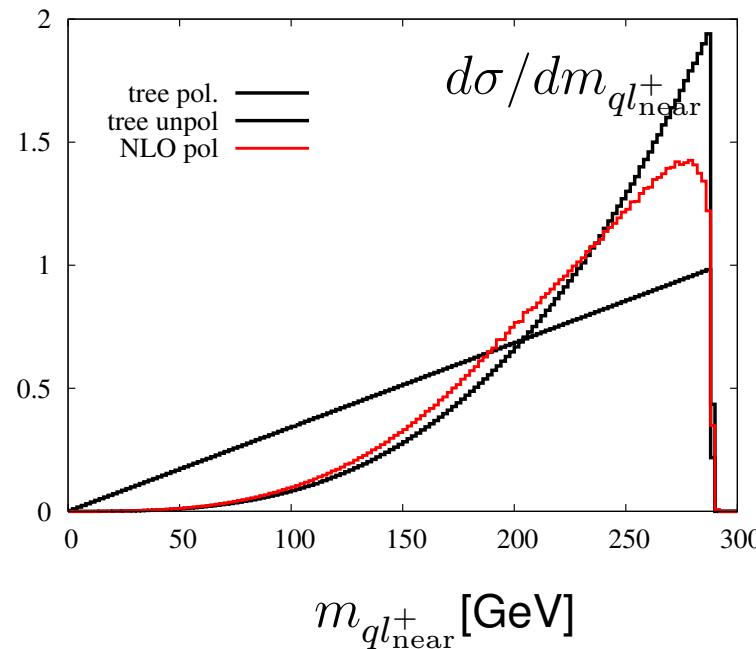


Precision calculations for BSM physics at the LHC

- exploring BSM models (later LHC phase)
→ determine masses & spins



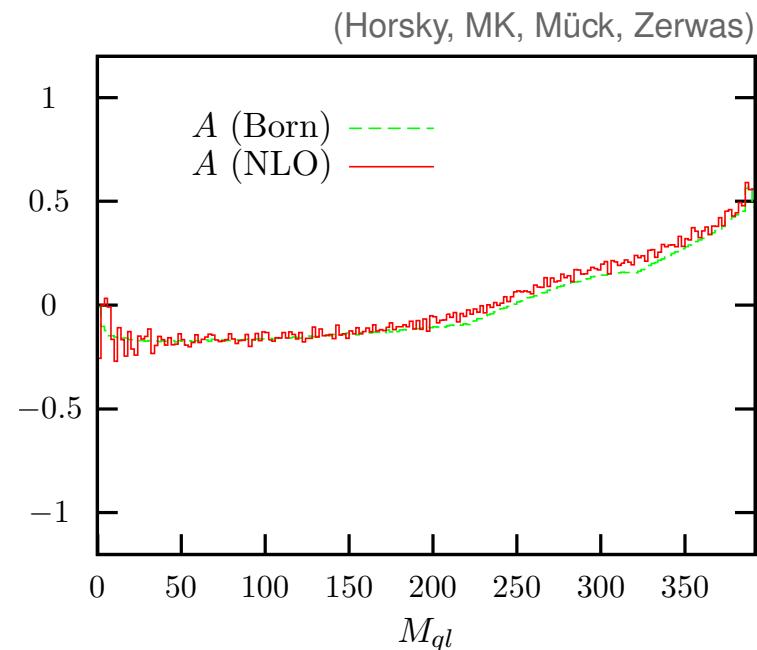
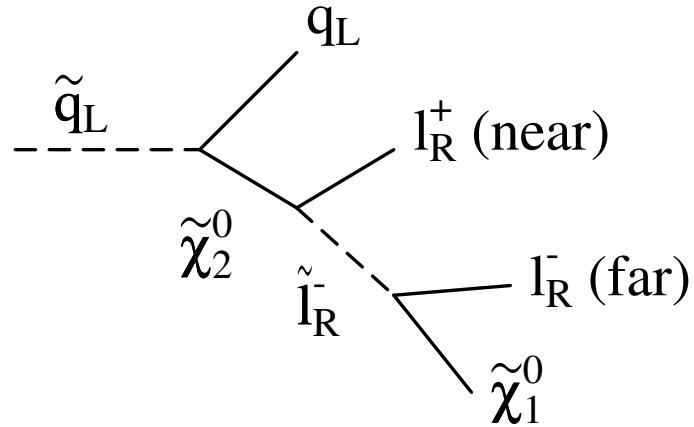
with spin correlations @ NLO:



→ radiative corrections affect shape of distributions (Horsky, MK, Mück, Zerwas)

Precision calculations for BSM physics at the LHC

- exploring BSM models (later LHC phase)
→ determine masses & spins



consider charge asymmetry $A = (d\sigma/dm_{ql+} - d\sigma/dm_{ql-})/(d\sigma/dm_{ql+} + d\sigma/dm_{ql-})$

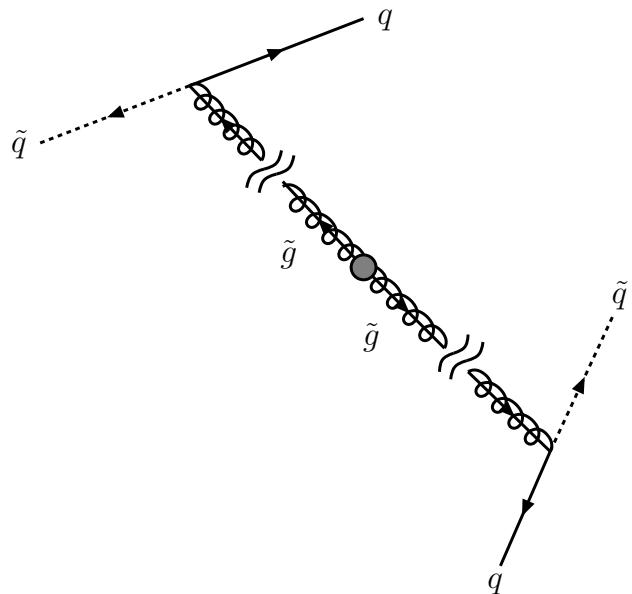
→ radiative corrections cancel in charge asymmetry

Precision calculations for BSM physics at the LHC

- exploring BSM models (later LHC phase)
→ determine masses & spins

consider invariant jet-mass distributions in decay chains like

$$pp \rightarrow \tilde{g} + \tilde{g} \rightarrow q\tilde{q}_R + q\tilde{q}_L \rightarrow qq\tilde{\chi}_1^0 + qq\tilde{\chi}_2^0$$



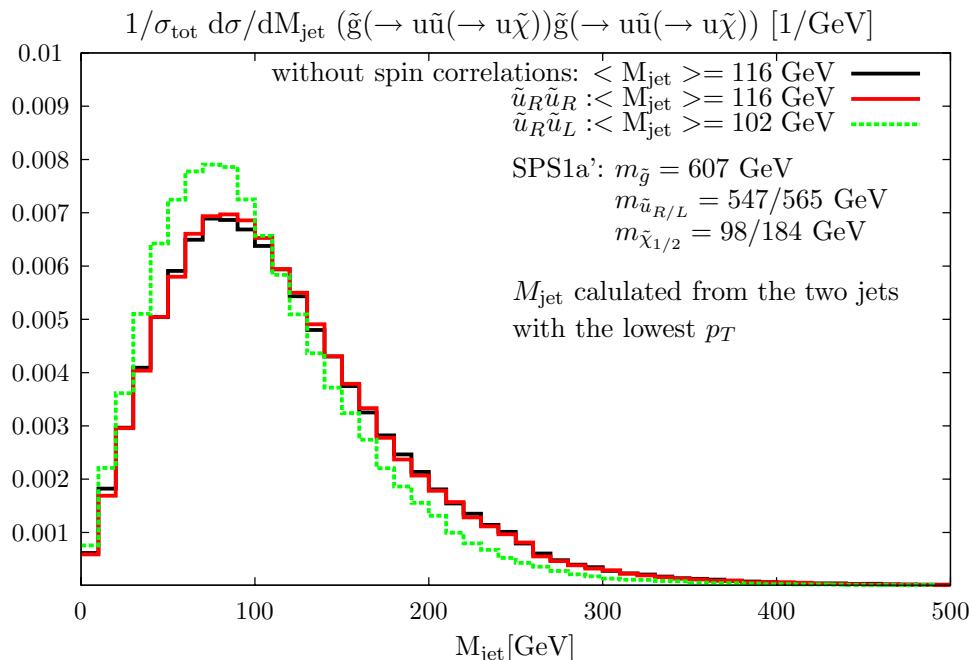
→ information on gluino spin in jet event shapes (MK, Popenda, Spira, Zerwas)

Precision calculations for BSM physics at the LHC

- exploring BSM models (later LHC phase)
→ determine masses & spins

consider invariant jet-mass distributions in decay chains like

$$pp \rightarrow \tilde{g} + \tilde{g} \rightarrow q\tilde{q}_R + q\tilde{q}_L \rightarrow qq\tilde{\chi}_1^0 + qq\tilde{\chi}_2^0$$



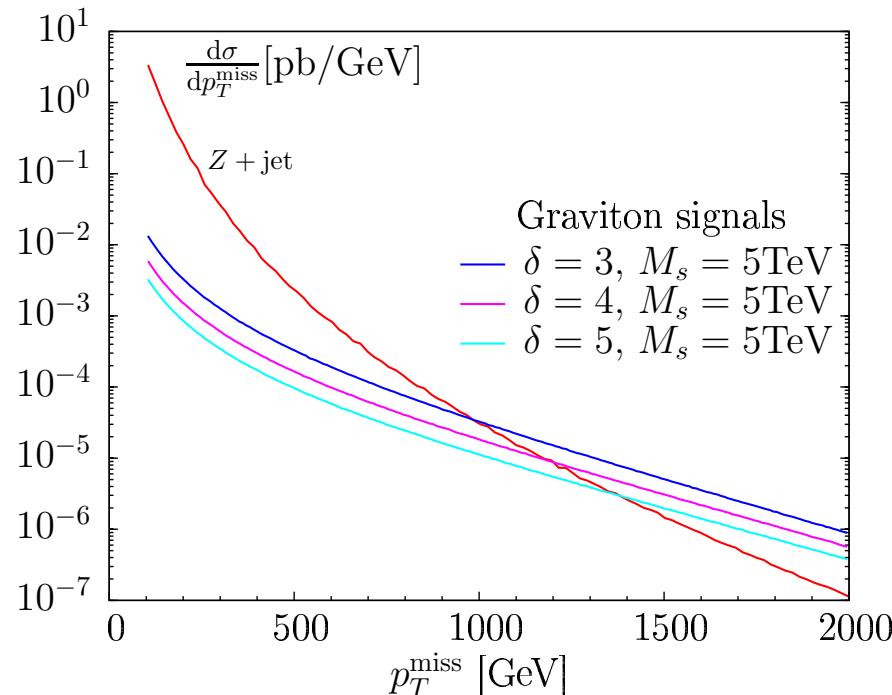
→ information on gluino spin in jet event shapes (MK, Popenda, Spira, Zerwas)

Precision calculations for BSM physics at the LHC

- exploring BSM models (later LHC phase)
→ determine masses & spins ... or the number of extra dimensions?

consider graviton production in ADD scenarios (Arkani-Hamed, Dimopoulos, Dvali)

$pp \rightarrow G + \text{jet} \rightarrow \text{monojet signature}$



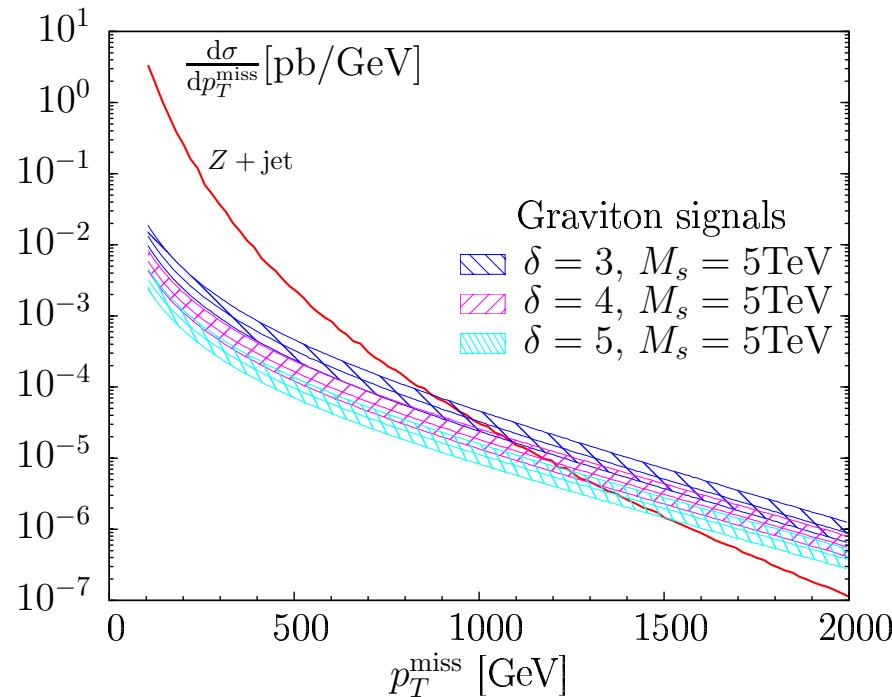
→ determine number of extra dimensions δ ?

Precision calculations for BSM physics at the LHC

- exploring BSM models (later LHC phase)
→ determine masses & spins ... or the number of extra dimensions?

consider graviton production in ADD scenarios (Arkani-Hamed, Dimopoulos, Dvali)

$pp \rightarrow G + \text{jet} \rightarrow \text{monojet signature}$



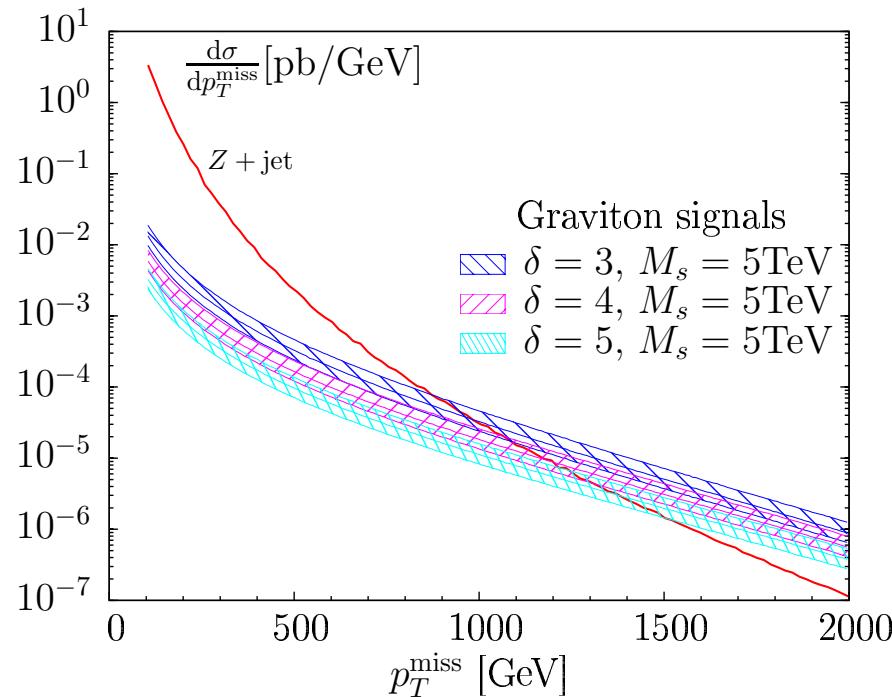
→ spoiled by QCD uncertainties...

Precision calculations for BSM physics at the LHC

- exploring BSM models (later LHC phase)
→ determine masses & spins ... or the number of extra dimensions?

consider graviton production in ADD scenarios (Arkani-Hamed, Dimopoulos, Dvali)

$pp \rightarrow G + \text{jet} \rightarrow \text{monojet signature}$



→ include NLO-QCD corrections (Karg, MK, Li, Zeppenfeld, in progress)

Progress in SUSY precision calculations at the LHC

- Sparticle production cross sections

- NLO QCD → [Prospino](#): Beenakker, Höpker, MK, Plehn, Spira, Zerwas
(cf. Baer, Hall, Reno; Berger, Klasen, Tait)
- threshold summation for $M_{\tilde{q}}, M_{\tilde{g}} \gtrsim 1$ TeV
(Bozzi, Fuks, Klasen; Kulesza, Motyka; Langenfeld, Moch;
Beenakker, Brensing, MK, Kulesza, Laenen, Niessen)
- electroweak corrections
(Hollik, Kollar, Mirabella, Trenkel; Bornhauser, Drees, Dreiner, Kim;
Beccaria, Macorini, Panizzi, Renard, Verzegnassi)
- beyond MSSM: NLO QCD for RPV SUSY
(Debajyoti Choudhury, Swapan Majhi, V. Ravindran; Yang, Li, Liu, Li; Dreiner, Grab, MK, Trenkel)

Progress in SUSY precision calculations at the LHC

● Sparticle production cross sections

- NLO QCD → [Prospino](#): Beenakker, Höpker, MK, Plehn, Spira, Zerwas
(cf. Baer, Hall, Reno; Berger, Klasen, Tait)
- threshold summation for $M_{\tilde{q}}, M_{\tilde{g}} \gtrsim 1$ TeV
(Bozzi, Fuks, Klasen; Kulesza, Motyka; Langenfeld, Moch;
Beenakker, Brensing, MK, Kulesza, Laenen, Niessen)
- electroweak corrections
(Hollik, Kollar, Mirabella, Trenkel; Bornhauser, Drees, Dreiner, Kim;
Beccaria, Macorini, Panizzi, Renard, Verzegnassi)
- beyond MSSM: NLO QCD for RPV SUSY
(Debajyoti Choudhury, Swapan Majhi, V. Ravindran; Yang, Li, Liu, Li; Dreiner, Grab, MK, Trenkel)

● Sparticle decays

- total rates → [Sdecay](#): (Mühlleitner, Djouadi, Mambrini) plus work by many authors.....
- only few results for shapes (Drees, Hollik, Xu; Horsky, MK, Mück, Zerwas)

Progress in SUSY precision calculations at the LHC

● Sparticle production cross sections

- NLO QCD → [Prospino](#): Beenakker, Höpker, MK, Plehn, Spira, Zerwas
(cf. Baer, Hall, Reno; Berger, Klasen, Tait)
- threshold summation for $M_{\tilde{q}}, M_{\tilde{g}} \gtrsim 1$ TeV
(Bozzi, Fuks, Klasen; Kulesza, Motyka; Langenfeld, Moch;
Beenakker, Brensing, MK, Kulesza, Laenen, Niessen)
- electroweak corrections
(Hollik, Kollar, Mirabella, Trenkel; Bornhauser, Drees, Dreiner, Kim;
Beccaria, Macorini, Panizzi, Renard, Verzegnassi)
- beyond MSSM: NLO QCD for RPV SUSY
(Debajyoti Choudhury, Swapan Majhi, V. Ravindran; Yang, Li, Liu, Li; Dreiner, Grab, MK, Trenkel)

● Sparticle decays

- total rates → [Sdecay](#): (Mühlleitner, Djouadi, Mambrini) plus work by many authors.....
- only few results for shapes (Drees, Hollik, Xu; Horsky, MK, Mück, Zerwas)

● Sparticle production \oplus decay: generic BSM cascades (SUSY, UEDs, LH,...)

- so far only tree level → challenge for LHC theory community...

Precision calculations at the LHC: conclusions & outlook ...

- Signal cross sections in the SM and MSSM are (or will rather soon be) known at NLO accuracy
 - theoretical uncertainty $\approx 15\%$ for inclusive cross sections
 - larger uncertainty for exclusive observables
 - can predict distributions and observables with cuts
 - in general no parton showers/hadronization

Precision calculations at the LHC: conclusions & outlook ...

- Signal cross sections in the SM and MSSM are (or will rather soon be) known at NLO accuracy
 - theoretical uncertainty $\approx 15\%$ for inclusive cross sections
 - larger uncertainty for exclusive observables
 - can predict distributions and observables with cuts
 - in general no parton showers/hadronization
- Progress in matching NLO calculations with parton showers and hadronization

Precision calculations at the LHC: conclusions & outlook ...

- Signal cross sections in the SM and MSSM are (or will rather soon be) known at NLO accuracy
 - theoretical uncertainty $\approx 15\%$ for inclusive cross sections
 - larger uncertainty for exclusive observables
 - can predict distributions and observables with cuts
 - in general no parton showers/hadronization
- Progress in matching NLO calculations with parton showers and hadronization
- Many backgrounds (multi-leg processes) are only known at LO
(Need breakthrough in techniques to do $pp \rightarrow> 4$ partons at NLO?)

Precision calculations at the LHC: conclusions & outlook ...

- Signal cross sections in the SM and MSSM are (or will rather soon be) known at NLO accuracy
 - theoretical uncertainty $\approx 15\%$ for inclusive cross sections
 - larger uncertainty for exclusive observables
 - can predict distributions and observables with cuts
 - in general no parton showers/hadronization
- Progress in matching NLO calculations with parton showers and hadronization
- Many backgrounds (multi-leg processes) are only known at LO
(Need breakthrough in techniques to do $pp \rightarrow> 4$ partons at NLO?)
- First LHC data will allow us to focus on the relevant processes...