Vortices, Superfluid turbulence & Nonthermal Fixed Points in Bose Gases



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Thanks & credits to...



...my work group in Heidelberg:

Sebastian Bock Sebastian Erne Martin Gärttner Roman Hennig Markus Karl Steven Mathey Boris Nowak Nikolai Philipp Maximilian Schmidt Jan Schole Dénes Sexty Martin Trappe Jan Zill







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HELMHOLTZ



...my former students:

Cédric Bodet (\rightarrow NEC), Alexander Branschädel (\rightarrow KIT Karlsruhe), Stefan Keßler (\rightarrow U Erlangen), Matthias Kronenwett (\rightarrow R. Berger), **Christian Scheppach** (\rightarrow Cambridge, UK), Philipp Struck (\rightarrow Konstanz), Kristan Temme (\rightarrow Vienna)

Equilibration



Transient, metastable state e.g. Turbulence Non-thermal fixed point



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Classical Turbulence





Kinetic energy cascade

large scales (source) \rightarrow small scales (sink)



Classical Turbulence





Lewis F. Richardson (1881-1953)

Kinetic energy cascadeLewis R
(188)large scales (source) \rightarrow small scales (sink)

"Big whirls have little whirls that feed on their velocity, and little whirls have lesser whirls and so on to viscosity." (Richardson, 1920)



Classical Turbulence





Andrey N. Kolmogorov (1903-1987)

Kinetic energy cascade(19large scales (source) \rightarrow small scales (sink)

"Big whirls have little whirls that feed on their velocity, and little whirls have lesser whirls and so on to viscosity."

(Richardson, 1920)

Kolmogorov (1941)

 $E(k) \sim k^{-5/3}$

(dynamical critical phenomenon)



Wave turbulence

Wave Turbulence – e.g. on water



Theory prediction:

 $E_{\omega} \sim \omega^{-17/6}$

[Zakharov & Filonenko (67)]





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Imagine you had a balance equation for the radial flux





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Transport equation (Quantum Boltzmann eq.):

dilute Bose gas: $T_{kpqr} \equiv g = 4\pi a_0/m = const.$



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Radial transport equation (Quantum Boltzmann):

Stationary distribution $n(k,t) \equiv n(k)$ if $Q(k) \equiv Q$

This requires a particular scaling of $n(k) \sim k^{-\zeta}$



Wave turbulence

Stationary scaling n(k) within inertial region:





Wave turbulence in an ultracold Bose gas

Dilute ultracold Bose Gas

Gross-Pitaevskii Equation:

$$(g = 4\pi a_0/m)$$

$$i\frac{\partial\Psi(\boldsymbol{\rho},t)}{\partial t} = \left(-\frac{\nabla^2}{2} + g|\Psi(\boldsymbol{\rho},t)|^2\right)\Psi(\boldsymbol{\rho},t)$$

Momentum spectrum:

 $n(\mathbf{k}) = \langle \Psi^*(\mathbf{k})\Psi(\mathbf{k}) \rangle$



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J. Berges, A. Rothkopf, J. Schmidt, PRL **101** (08) 041603 C. Scheppach, J. Berges, TG PRA **81** (10) 033611



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Quantum Boltzmann breaks down for large *n*, once $|T_{kpqr}|n_k \gg O(1)$

$$\begin{split} \partial_t n(k) &= -\partial_k Q(k) \sim k^{d-1} J(k) \\ &= k^{d-1} d\Omega_k \int d^d p \, d^d q \, d^d r \, |T_{\mathbf{k}\mathbf{p}\mathbf{q}\mathbf{r}}|^2 \delta(\mathbf{k} + \mathbf{p} - \mathbf{q} - \mathbf{r}) \, \delta(\omega_{\mathbf{k}} + \omega_{\mathbf{p}} - \omega_{\mathbf{q}} - \omega_{\mathbf{r}}) \\ &\quad \text{coupling mom. conservation energy conservation} \\ &\times [(n_{\mathbf{k}} + 1)(n_{\mathbf{p}} + 1)n_{\mathbf{q}}n_{\mathbf{r}} - n_{\mathbf{k}}n_{\mathbf{p}}(n_{\mathbf{q}} + 1)(n_{\mathbf{r}} + 1)] \\ &\quad \text{in-scattering rate out-scattering rate} \end{split}$$

Cured by
effective many-body T-Matrix:
$$|T|^2 = g^2 \rightarrow |T_k^{MB}|^2 \sim \frac{g^2}{1 + (gkn_k)^2}$$



Dyn. QFT: Resummed Vertex

 $p = (p_0, \mathbf{p}):$

$$J(p) := \Sigma_{ab}^{\rho}(p) F_{ba}(p) - \Sigma_{ab}^{F}(p) \rho_{ba}(p) \stackrel{!}{=} 0$$

$$\Sigma_{ab}(x,y) = \frac{1}{a}$$

Vertex bubble resummation: (e.g. 2PI to NLO in 1/N)

$$\mathbf{M} \rightarrow \mathbf{M} = \mathbf{M} + \mathbf{M} +$$

[Dynamics: J. Berges, (02); G. Aarts et al., (02); Nonthermal fixed points: J. Berges, A. Rothkopf, J. Schmidt, PRL (08)]



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Bose gas in d spatial dimensions $n \sim k^{-\zeta}$



J. Berges, A. Rothkopf, J. Schmidt, PRL **101** (08) 041603; J. Berges, G. Hoffmeister, NPB 813, 383 (2009) C. Scheppach, J. Berges, TG PRA **81** (10) 033611



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Vortices in a superfluid ultracold Bose gas

Superfluid hydro of Bose-condensed Gas

The Gross-Pitaevskii Equation,

$$(g = 4\pi a_0/m)$$

$$i\frac{\partial\Psi(\boldsymbol{\rho},t)}{\partial t} = \left(-\frac{\nabla^2}{2} + g|\Psi(\boldsymbol{\rho},t)|^2\right)\Psi(\boldsymbol{\rho},t)$$

using defs.

 $\Psi(\boldsymbol{\rho}, t) = \sqrt{n(\boldsymbol{\rho}, t)} \exp[i\varphi(\boldsymbol{\rho}, t)]$ $Q = gn \qquad \mathbf{u}(\boldsymbol{\rho}, t) = \nabla\varphi(\boldsymbol{\rho}, t)$

can be written as

$$\frac{\partial}{\partial t}n + \nabla \cdot (n\mathbf{u}) = 0$$

$$\frac{\partial}{\partial t}\mathbf{u} + \mathbf{u} \cdot \nabla \mathbf{u} = -\nabla Q$$

Euler equation



Quantum Vortices



 $\Psi(\boldsymbol{\rho},t) = \sqrt{n(\boldsymbol{\rho},t)} \exp[i\varphi(\boldsymbol{\rho},t)]$ complex field

 $\mathbf{u}(\boldsymbol{\rho},t) = \nabla \varphi(\boldsymbol{\rho},t)$ velocity





Vortices in a Na condensate



J. R. Abo-Shaeer, C. Raman, J. M. Vogels, W. Ketterle 20 APRIL 2001 VOL 292 SCIENCE



Gross-Pitaevskii Simulations for an ultracold Bose gas

Movie 1: Phase evolution & Spectrum

 $\Psi(\boldsymbol{\rho},t) = \sqrt{n(\boldsymbol{\rho},t)} \exp[i\varphi(\boldsymbol{\rho},t)]$

 $n(k) = \langle \Psi^*(\mathbf{k}) \Psi(\mathbf{k}) \rangle \big|_{\text{angle average}}$





Movie by Jan Schole

Movie 2: Vortex "gas" & Spectrum

$$n(k) = \langle \Psi^*(\mathbf{k})\Psi(\mathbf{k}) \rangle \Big|_{\text{angle average}}$$



B. Nowak, D. Sexty, TG, PRB 84: 020506(R), 2011

Spectrum in 2+1 D





Cascades in 2+1 D





Interpretation: Random vortex distributions

Point vortex model



B. Nowak, J. Schole, D. Sexty, TG, arXiv:1111.61XX [cond-mat.quant-gas]



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Point vortex model in 2+1 D







Simulations in 2+1 D





Vortex position correlations





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Dyn. QFT: Resummed Vertex

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Movie 3: Vortex Lines in 3+1 D

$$n(k) = \langle \Psi^*(\mathbf{k})\Psi(\mathbf{k}) \rangle \Big|_{\text{angle average}}$$



3+1 D simulations





Line vortex model in 3+1 D





Simulations in 3+1 D





Decomposition of flow





Acoustic turbulence





Non-thermal fixed point



thermal equilibrium



[Fig. courtesy: J. Berges '08]

Vortex tangles in Bose Einstein Condensates





Relativistic scalar field

Strong Turbulence

Simulations of the non-linear Klein-Gordon equation, O(2) symmetry

$$(\partial_t^2 - \partial_x^2)\varphi(x, t) + \lambda\varphi^3(x, t) = 0$$

Initial condition: Highly occupied zero mode, Unoccupied modes with k>0

(video)

See also: http://www.thphys.uni-heidelberg.de/~sexty/videos

TG, B. Nowak, D. Sexty, arXiv:1108.0541 [hep-ph]



Strong Turbulence = Charge Separation

Modulus of complex field $|\phi|$ vs. mean charge distribution



TG, B. Nowak, D. Sexty, arXiv:1108.0541 [hep-ph] cf. also Tkachev, Kofman, Starobinsky, Linde (1998)



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Strong Turbulence = Charge Separation

Charge density distribution

VS.

power spectrum





TG, B. Nowak, D. Sexty, arXiv:1108.0541 [hep-ph]



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Supplementary slides

Cascades in 2+1 D: Fluxes



Time evolution of vortex density





Decomposition of Energy

$$E_{tot} = \int \left(\frac{1}{2} |\nabla \sqrt{n}e^{-i\varphi}|^2 + \frac{1}{2}gn^2\right) d\rho$$

= $E_{kin} + E_q + E_{int}$
$$u(\rho, t) = \nabla \varphi(\rho, t)$$

$$E_{kin} = \frac{1}{2} \int |\sqrt{n}u|^2 d\rho = E_{kin}^i + E_{kin}^c$$

$$\nabla \times (\sqrt{n}u)^c = 0$$

$$\nabla \cdot (\sqrt{n}u)^i = 0$$

$$E_q = \frac{1}{2} \int (\nabla \sqrt{n})^2 d\rho$$



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Simulations in 2+1 D

 $E(\mathbf{k}) = \boldsymbol{\omega}(\mathbf{k})\mathbf{k}^{d-1}\mathbf{n}(\mathbf{k})$





Simulations in 2+1 D

 $E(\mathbf{k}) = \boldsymbol{\omega}(\mathbf{k})\mathbf{k}^{d-1}\mathbf{n}(\mathbf{k})$



Lewis Fry Richardson, FRS (1881-1953)

Big whirls have little whirls that feed on their velocity, and little whirls have lesser whirls and so on to viscosity.

(L.F. Richardson, The supply of energy from and to Atmospheric Eddies, 1920)

Great fleas have little fleas upon their backs to bite 'em, And little fleas have lesser fleas, and so ad infinitum. And the great fleas themselves, in turn, have greater fleas to go on; While these again have greater still, and greater still, and so on.

(Augustus de Morgan, A Budget of Paradoxes, 1872, p. 370)

So, naturalists observe, a flea Has smaller fleas that on him prey; And these have smaller still to bite 'em; And so proceed ad infinitum.

(Jonathan Swift: Poetry, a Rhapsody, 1733)

