

Thoughts about the SUSY WIMP.

...and the state of SUSY generally...

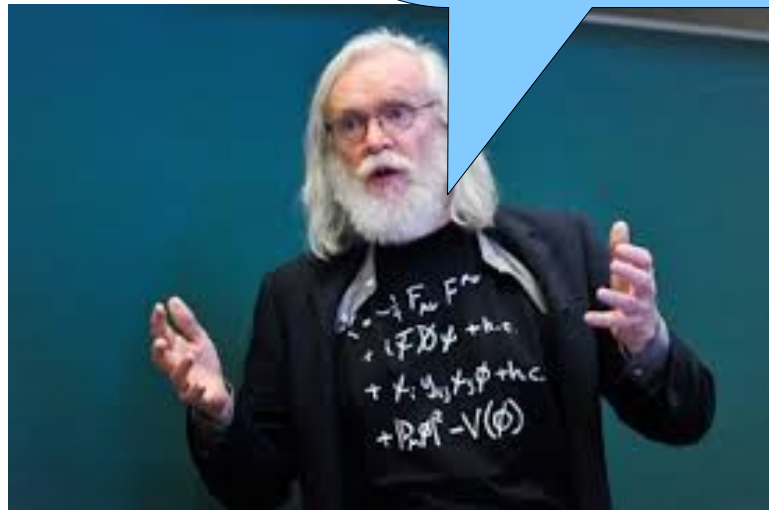
Kai Schmidt-Hoberg

“SUSY anywhere is better than SUSY nowhere!”

Largely based on

1603.09347
1701.03480

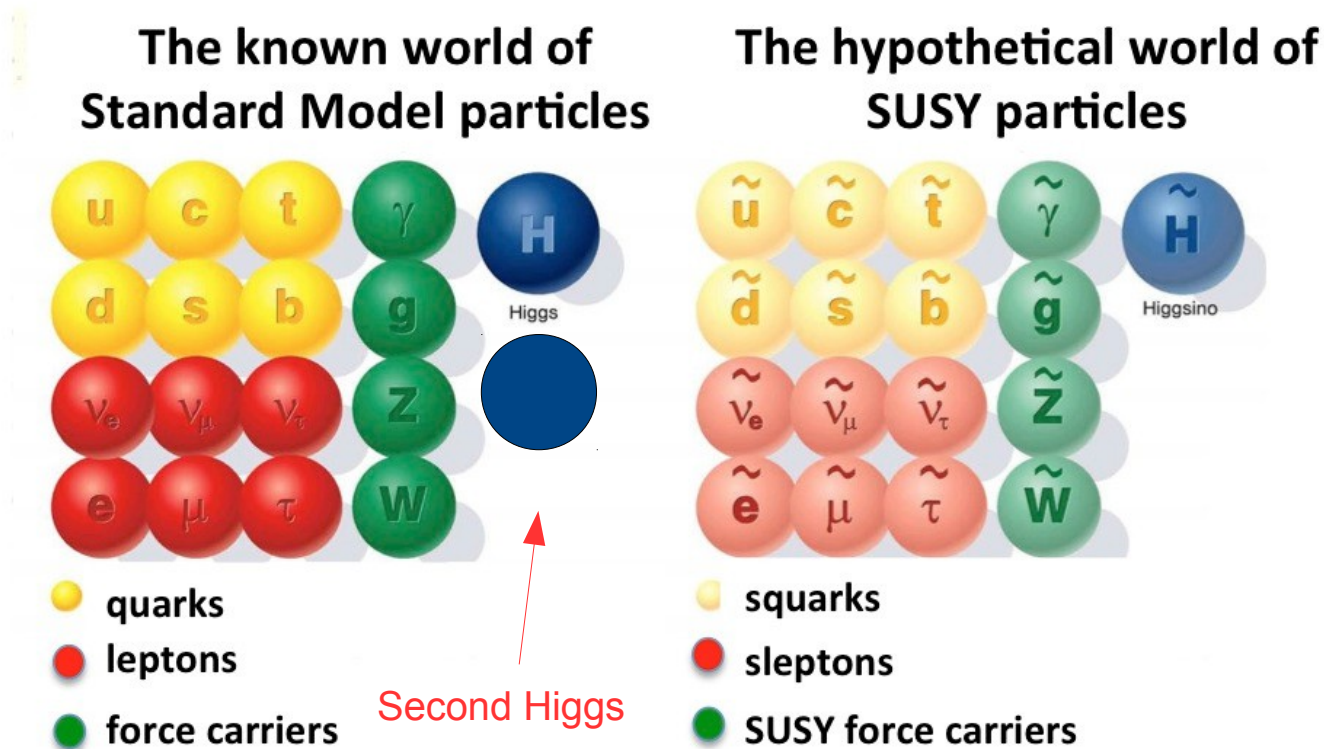
with G Ross and F Staub



MPIK theory seminar

SUSY → MSSM (this talk)

- > A SUSY partner for each SM particle with $\Delta s = 1/2$ with **the same mass**



- > SUSY broken in nature
- > Breaking mechanism unknown, parameterized by soft terms

- > Field content fixed: theory specified by superpotential and soft terms

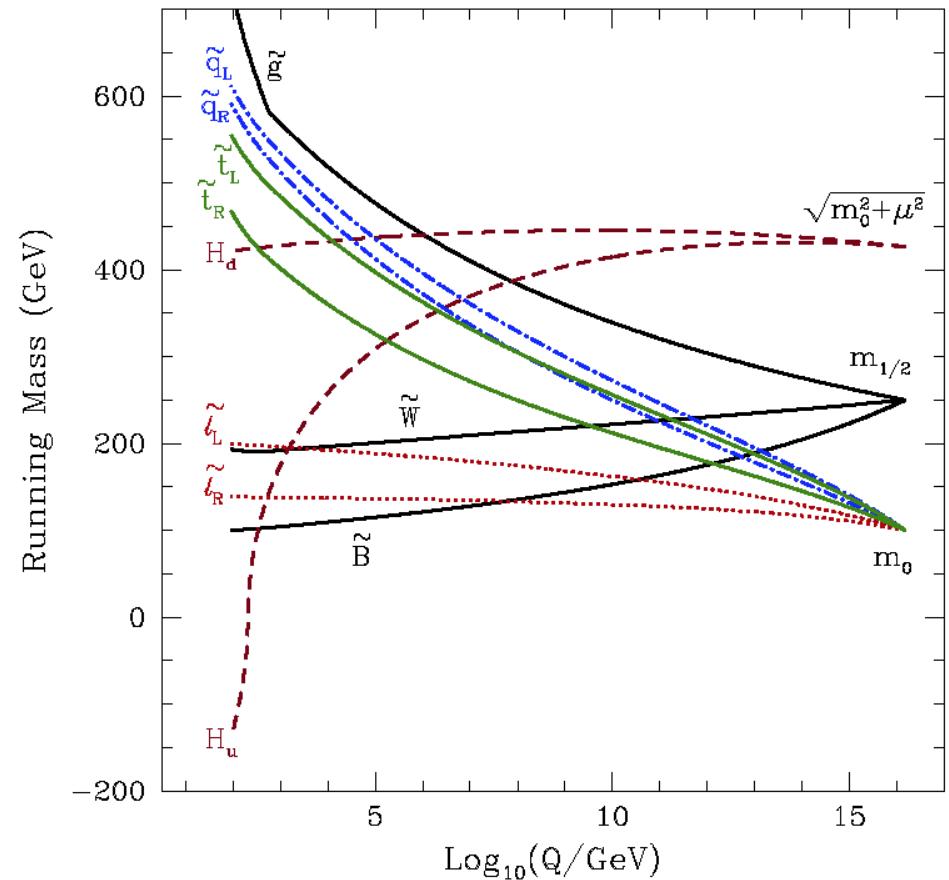
$$\begin{aligned}
 \mathcal{W} = & \mu H_u H_d + \kappa_i L_i H_u \\
 & + Y_e^{ij} H_d L_i E_j^c + Y_d^{ij} H_d Q_i D_j^c + Y_u^{ij} H_u Q_i U_j^c \quad \swarrow \text{superfields} \\
 & + \lambda_{ijk}^{(0)} L_i L_j E_k^c + \lambda_{ijk}^{(1)} L_i Q_j D_k^c + \lambda_{ijk}^{(2)} U_i^c D_j^c D_k^c \\
 & + \kappa_{ij}^{(0)} H_u L_i H_u L_j + \kappa_{ijk\ell}^{(1)} Q_i Q_j Q_k L_\ell + \kappa_{ijk\ell}^{(2)} U_i^c U_j^c D_k^c E_\ell^c
 \end{aligned}$$

$$\begin{aligned}
 L_{SB} = & -\frac{1}{2} \sum_a M_a \bar{\lambda}_a \lambda_a - \sum_i m_{\tilde{\Phi}_i}^2 |\tilde{\Phi}_i|^2 + \\
 & T_u H_u \tilde{Q} \tilde{u} + T_d H_d \tilde{Q} \tilde{d} + T_e H_d \tilde{L} \tilde{e} + B_\mu H_u H_d
 \end{aligned}$$

- > Many new parameters (>100) but likely not independent

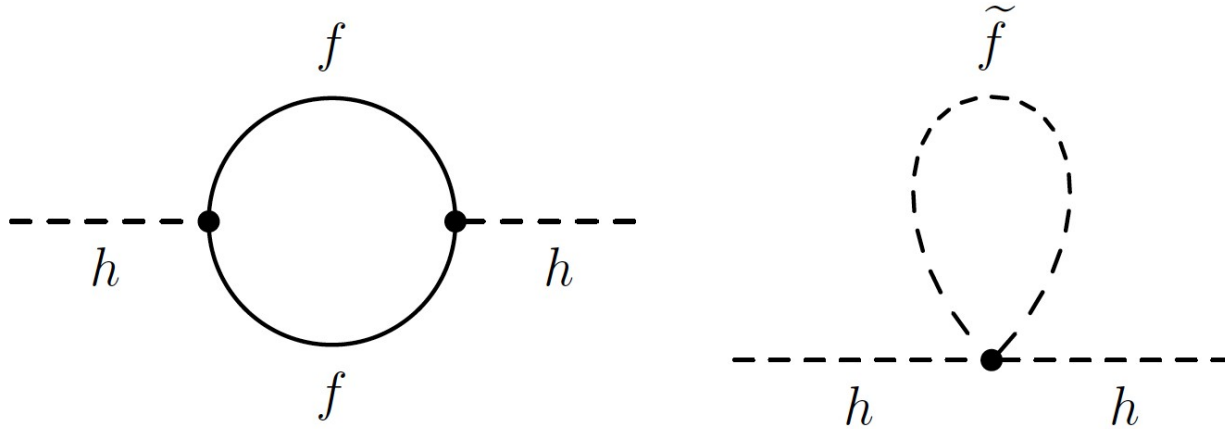
High scale SUSY

- > In many models SUSY breaking at high scale in hidden sector
- > Often some universality
- > Take into account running to predict SUSY spectrum at the electroweak scale.



Why we like(d) SUSY

- > **Hierarchy problem:** stabilizes the weak against the Planck scale



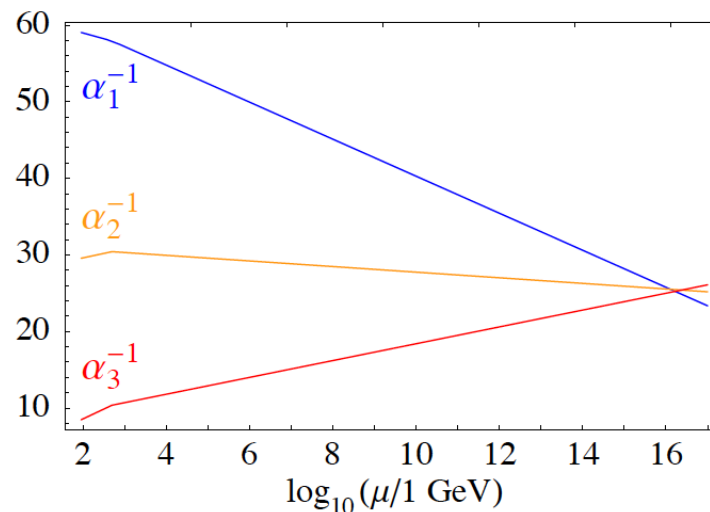
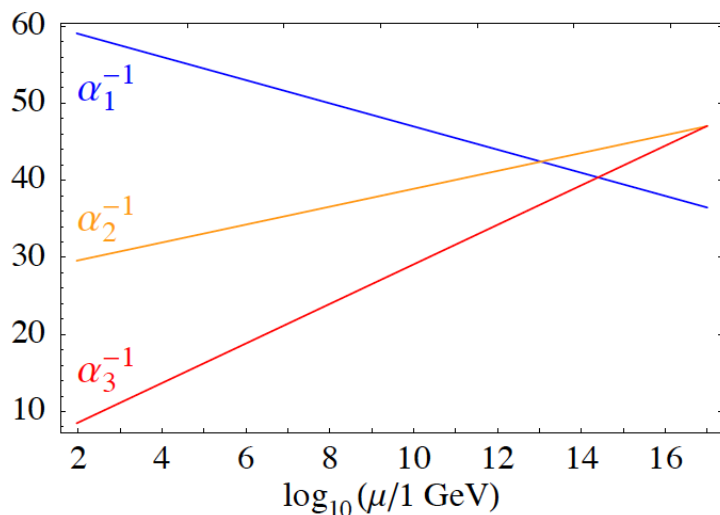
Why we like(d) SUSY

- > **Hierarchy problem:** stabilizes the weak against the Planck scale
- > **Dark matter:** If lightest SUSY particle stable \rightarrow dark matter candidate



Why we like(d) SUSY

- > **Hierarchy problem:** stabilizes the weak against the Planck scale
- > **Dark matter:** If lightest SUSY particle stable \rightarrow dark matter candidate
- > **Gauge coupling unification:**



Why we like(d) SUSY

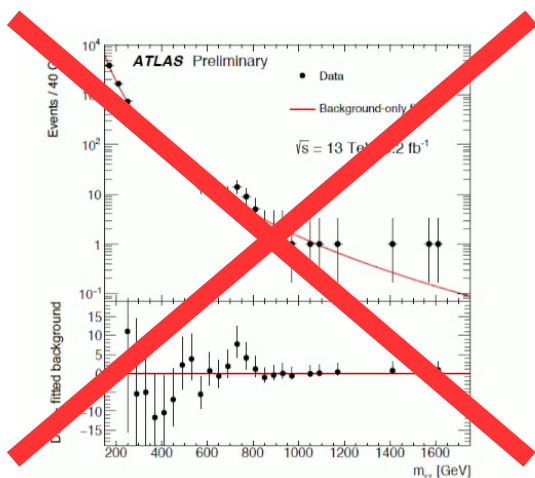
- > **Hierarchy problem:** stabilizes the weak against the Planck scale
- > **Dark matter:** If lightest SUSY particle stable \rightarrow dark matter candidate
- > **Gauge coupling unification:**
- > **A 125 GeV Higgs boson:** Additional hint for SUSY?

...a TC guy from Odense still owes me...



Why we like(d) SUSY

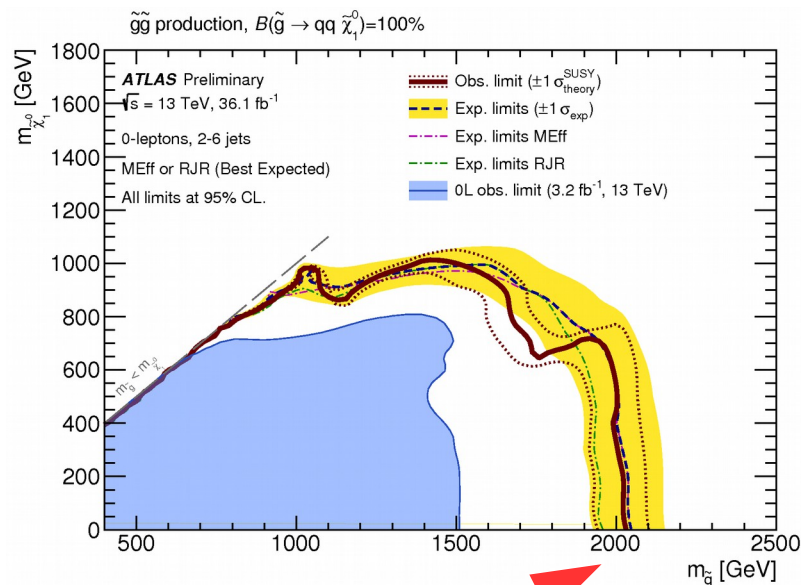
- > **Hierarchy problem:** stabilizes the weak against the Planck scale
- > **Dark matter:** If lightest SUSY particle stable \rightarrow dark matter candidate
- > **Gauge coupling unification:**
- > **A 125 GeV Higgs boson:** Additional hint for SUSY?
- > Also hard to get 750 GeV diphoton excess ;-)



Why we like(d) SUSY

- > Hierarchy problem: stabilizes the weak against the Planck scale
- > Dark matter: If lightest SUSY particle stable \rightarrow dark matter candidate
- > Gauge coupling unification:
- > A 125 GeV Higgs boson: Additional hint for SUSY?

> So why do people get worried?



$\sim 2 \text{ TeV}$

Why we like(d) SUSY

- > Hierarchy problem: stabilizes the weak against the Planck scale
- > Dark matter: If lightest SUSY particle stable \rightarrow dark matter candidate
- > Gauge coupling unification:
- > A 125 GeV Higgs boson: Additional hint for SUSY?
- > So why do people get worried?

NATURALNESS!



The LSP – a DM candidate.

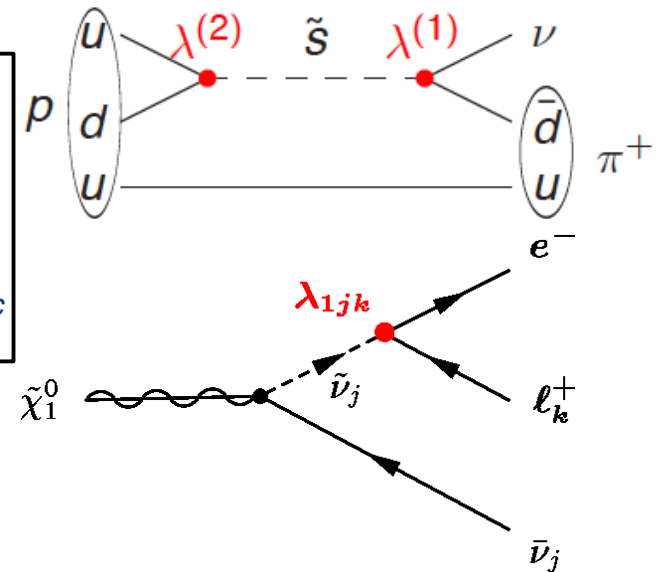
- Neutralinos are mixtures of bino, Wino and Higgsinos
- EM and colour neutral → potentially interesting dark matter candidates



The LSP – a DM candidate.

- Neutralinos are mixtures of bino, Wino and Higgsinos
- EM and colour neutral → potentially interesting dark matter candidates

$$\begin{aligned} \mathcal{W} = & \mu H_u H_d + \kappa_i L_i H_u \\ & + Y_e^{ij} H_d L_i E_j^c + Y_d^{ij} H_d Q_i D_j^c + Y_u^{ij} H_u Q_i U_j^c \\ & + \lambda_{ijk}^{(0)} L_i L_j E_k^c + \lambda_{ijk}^{(1)} L_i Q_j D_k^c + \lambda_{ijk}^{(2)} U_i^c D_j^c D_k^c \\ & + \kappa_{ij}^{(0)} H_u L_i H_u L_j + \kappa_{ijkl}^{(1)} Q_i Q_j Q_k L_\ell + \kappa_{ijkl}^{(2)} U_i^c U_j^c D_k^c E_\ell^c \end{aligned}$$

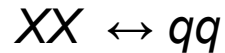


- With only SUSY and gauge invariance: extra terms leading to proton and dark matter decay → need additional symmetry.
- Standard assumption: R-parity conservation (good enough for dark matter – not good enough for the proton, need a better symmetry such as Z_4^R)

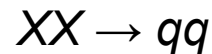
The relic abundance

- Neutralinos are produced in the early universe via thermal freeze out

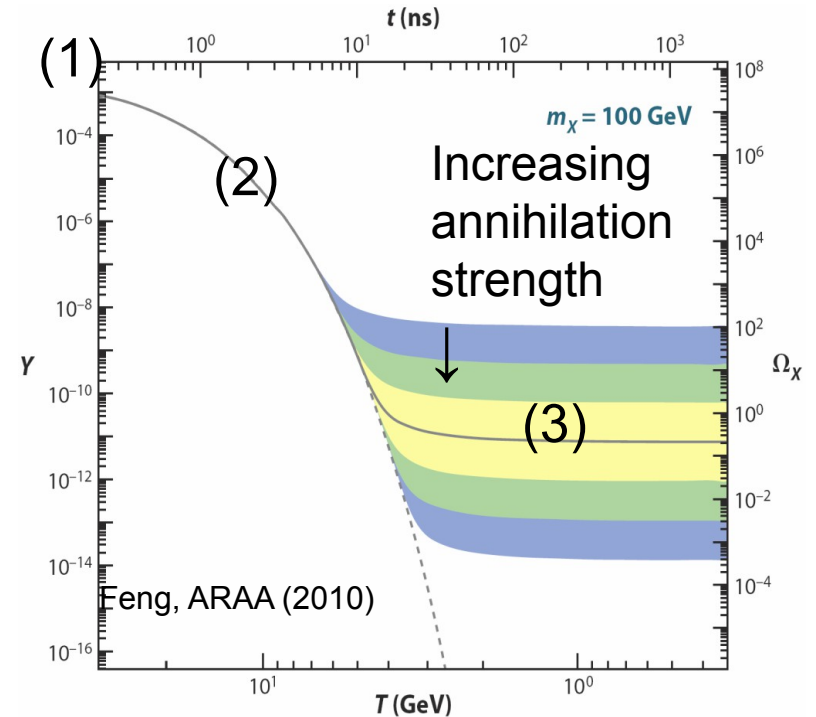
(1) Neutralino X is initially in thermal equilibrium:



(2) Universe cools:



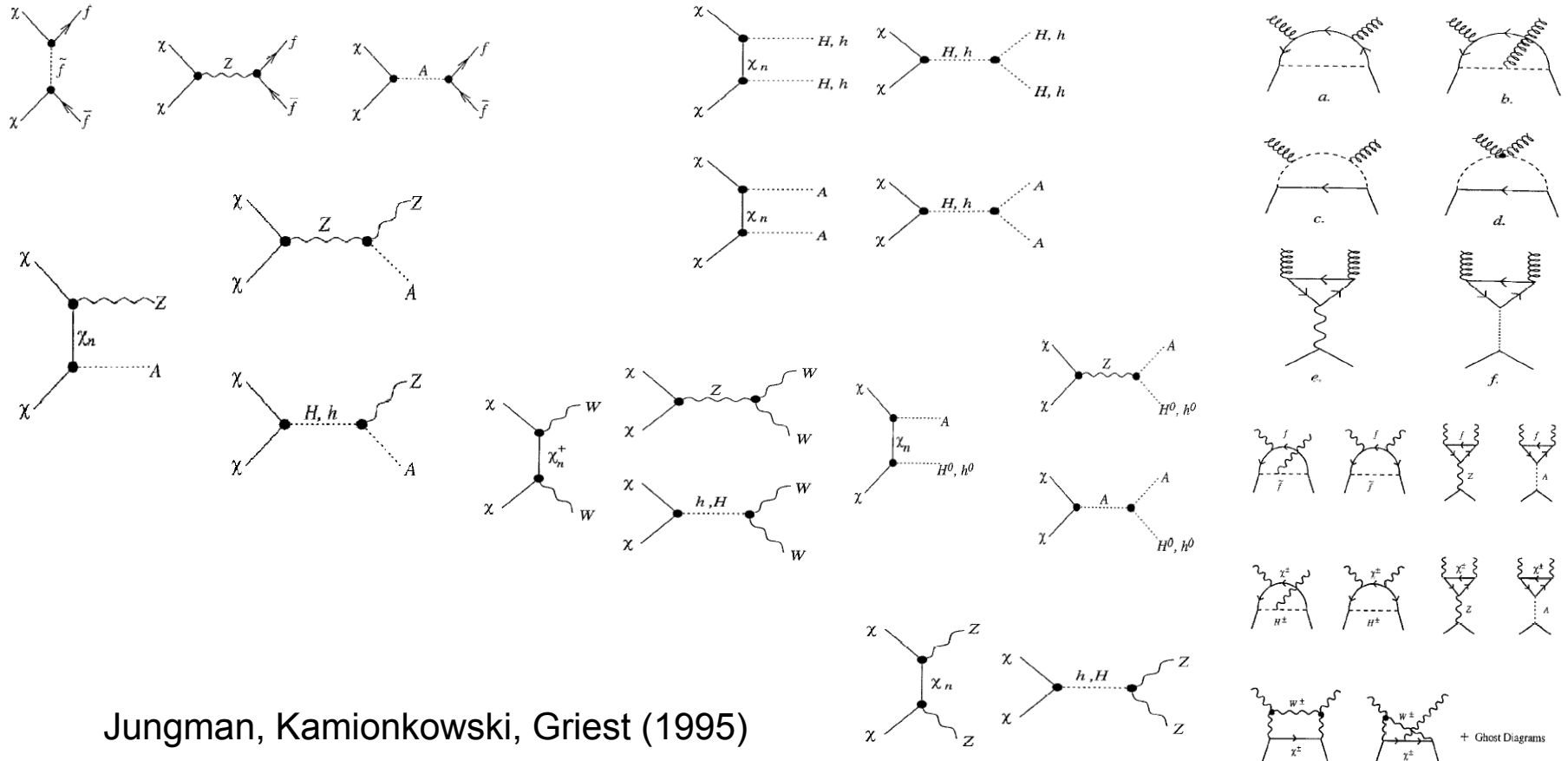
(3) Universe expands:



- Relic abundance depends on annihilation cross section

The relic abundance

- Neutralinos are produced in the early universe via thermal freeze out



Jungman, Kamionkowski, Griest (1995)

DM naturalness in the MSSM

➤ How naturally can the dark matter relic abundance be achieved?

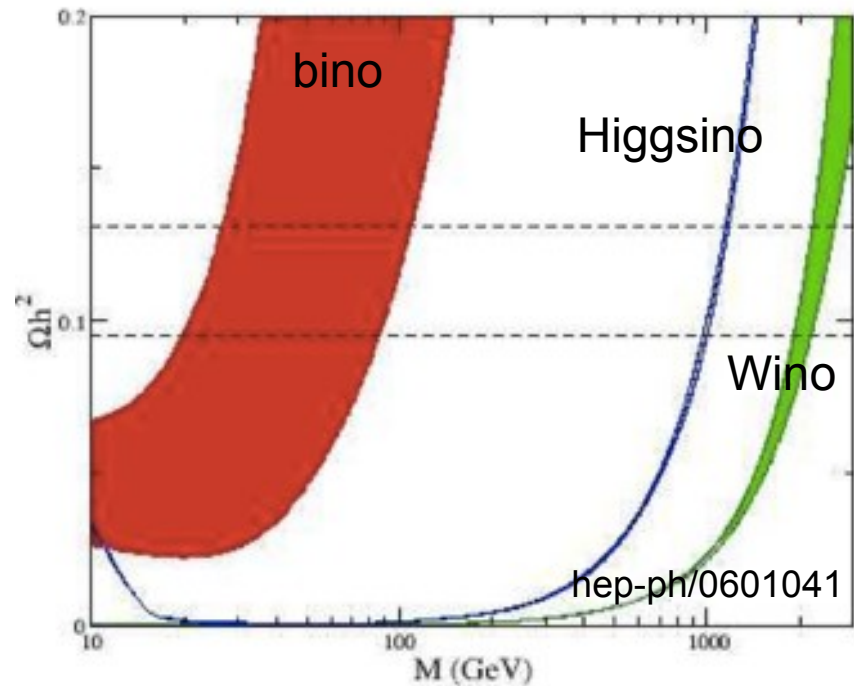
Often universal gaugino masses assumed at high scale, at low scale $M_3:M_2:M_1 \sim 6:2:1 \rightarrow$ bino LSP

Bino: Typically need to finely tune relic density via coannihilations or resonances :-)

Crucially depends on assumption of SUSY breaking terms!
Other patterns possible...

2-3 TeV Wino challenged by ID
Mariangela Lisanti et al [1307.4082](#)

1 TeV Higgsino looking good :-)



EW naturalness in the MSSM

- > How naturally can we achieve the correct Higgs vev?
- > Electroweak vev (or M_Z) determined by SUSY parameters (from minimization condition for scalar potential)

$$\frac{m_Z^2}{2} = \frac{m_{H_d}^2 - m_{H_u}^2 \tan^2 \beta}{\tan^2 \beta - 1} - \mu^2 \simeq -m_{H_u}^2 - \mu^2$$

- > Cancellation (tuning) needed for large SUSY masses
- > How to quantify this?

$$\Delta_p \equiv \frac{\partial \ln v^2}{\partial \ln p} = \frac{p}{v^2} \frac{\partial v^2}{\partial p} \quad \text{'sensitivity measure'}$$

- > Large Δ implies large tuning

Caveats of the sensitivity measure

- > What fundamental parameters should be included (and what are the fundamental parameters)?
- > Also depends on parameterization of fundamental parameters
- > At which scale?
- > It measures sensitivity rather than 'tuning' – can be different
- > The acceptable values Δ depends on taste – no absolute measure

→ **While for a given definition it can be calculated precisely, its physical interpretation is somewhat blurred.**

The usual story

- > What does this tell us about a natural SUSY spectrum?
- > μ is a superpotential parameter and hardly runs: $\mu_{\text{EW}} \sim \mu_{\text{GUT}}$

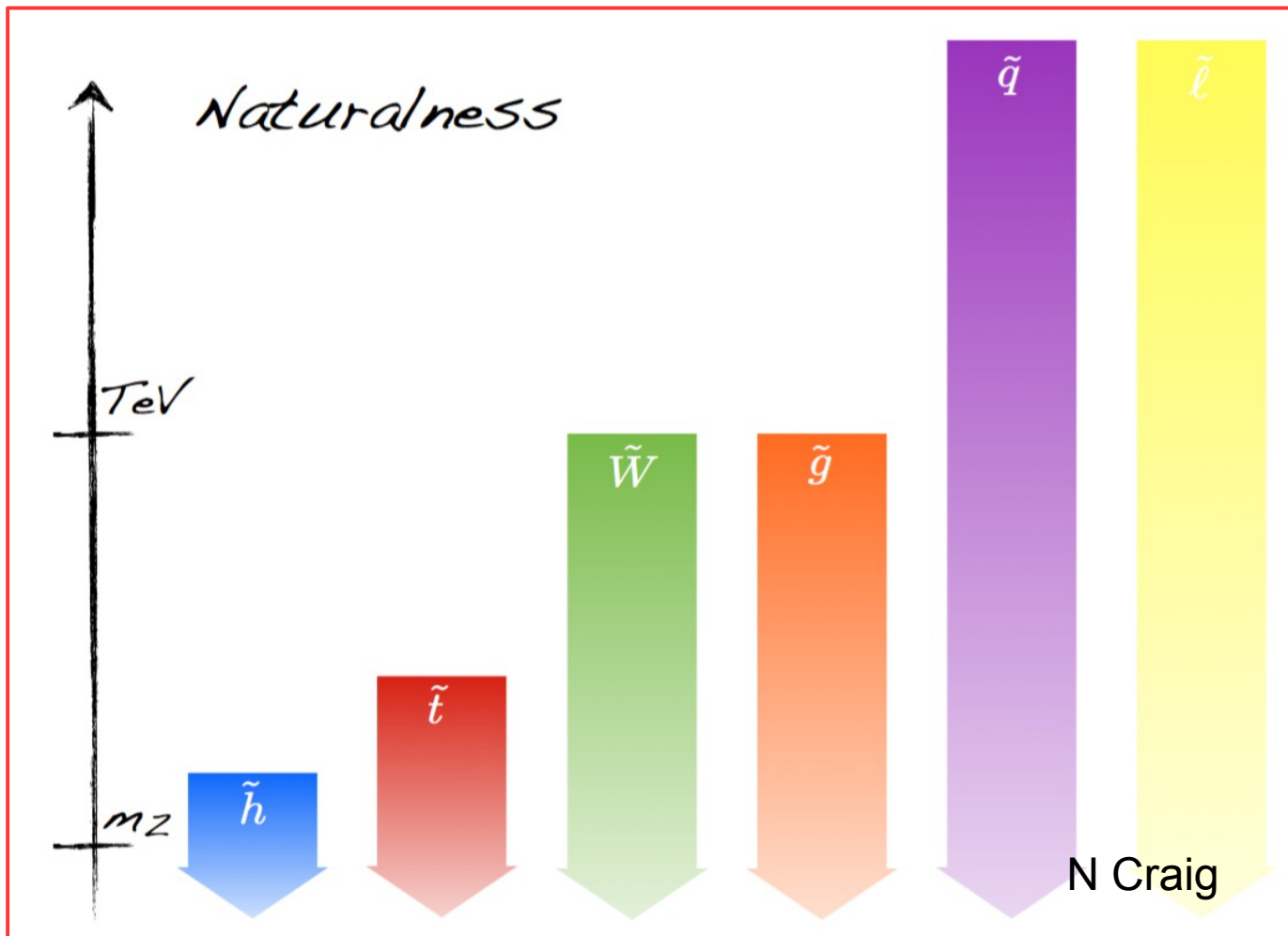
$$\text{Higgsino mass} \sim \mu \sim 1 \text{ TeV} - \Delta_\mu \sim \frac{2\mu^2}{M_Z^2} \sim 250$$

- > “Natural SUSY requires light Higgsino”
- > What about the m_{H_u} part?
- > Loop effects introduce a large sensitivity to stop and gluino masses

$$\delta m_{H_u}^2 = -\frac{3y_t^2}{4\pi^2} m_{\tilde{t}}^2 \ln(\Lambda/m_{\tilde{t}})$$

$$\delta m_{\tilde{t}}^2 = \frac{2g_s^2}{3\pi^2} m_{\tilde{g}}^2 \ln(\Lambda/m_{\tilde{g}})$$

The 'natural SUSY' spectrum



EW naturalness in the MSSM

- Starting from the high scale, all soft terms contribute to m_{H_u} and m_Z

$$\begin{aligned} m_Z^2 \simeq & -2.18\mu^2 + 3.84M_3^2 + 0.32M_3M_2 + 0.047M_1M_3 - 0.42M_2^2 \\ & + 0.011M_2M_1 - 0.012M_1^2 - 0.65M_3A_t - 0.15M_2A_t \\ & - 0.025M_1A_t + 0.22A_t^2 + 0.004M_3A_b \\ & - 1.27m_{H_u}^2 - 0.053m_{H_d}^2 \\ & + 0.73m_{Q_3}^2 + 0.57m_{U_3}^2 + 0.049m_{D_3}^2 - 0.052m_{L_3}^2 + 0.053m_{E_3}^2 \\ & + 0.051m_{Q_2}^2 - 0.11m_{U_2}^2 + 0.051m_{D_2}^2 - 0.052m_{L_2}^2 + 0.053m_{E_2}^2 \\ & + 0.051m_{Q_1}^2 - 0.11m_{U_1}^2 + 0.051m_{D_1}^2 - 0.052m_{L_1}^2 + 0.053m_{E_1}^2, \end{aligned}$$

- We don't just want m_{H_u} to be small, but every contribution to it. Assuming no correlations among the terms, need rather light stops and gluinos

EW naturalness in the MSSM

- Starting from the high scale, all soft terms contribute to m_{H_u} and m_Z

$$m_Z^2 \simeq -2.18\mu^2 + 3.84M_3^2 + 0.32M_3M_2 + 0.047M_1M_3 - 0.42M_2^2 \\ + 0.011M_2M_1 - 0.012M_1^2 - 0.65M_3A_t - 0.15M_2A_t \\ - 0.025M_1A_t + 0.22A_t^2 + 0.004M_3A_b$$

$$\begin{aligned} & -1.27m_{H_u}^2 - 0.053m_{H_d}^2 \\ & + 0.73m_{Q_3}^2 + 0.57m_{U_3}^2 + 0.049m_{D_3}^2 - 0.052m_{L_3}^2 + 0.053m_{E_3}^2 \\ & + 0.051m_{Q_2}^2 - 0.11m_{U_2}^2 + 0.051m_{D_2}^2 - 0.052m_{L_2}^2 + 0.053m_{E_2}^2 \\ & + 0.051m_{Q_1}^2 - 0.11m_{U_1}^2 + 0.051m_{D_1}^2 - 0.052m_{L_1}^2 + 0.053m_{E_1}^2 \end{aligned}$$

$$m_i^2 = m_0^2$$

$$\sim 0.01 m_0^2$$

- We don't just want m_{H_u} to be small, but every contribution to it. Assuming no correlations among the terms, need rather light stops and gluinos
- But we know correlations should be present...
- Example: the scalar focus point.

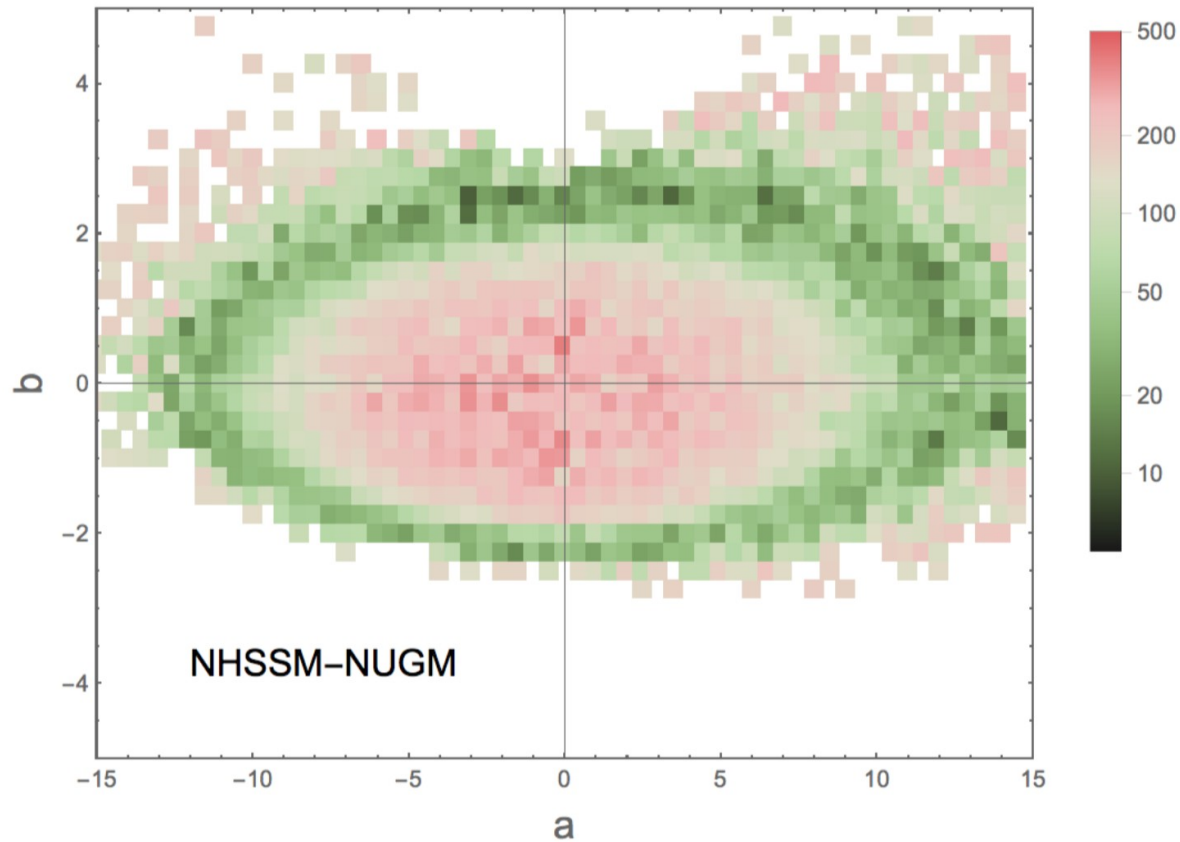
The gaugino focus point

- Assume fixed ratios of gaugino masses
- Possible also in GUTs

$$M_1 = a \cdot m_{1/2}$$

$$M_2 = b \cdot m_{1/2}$$

$$M_3 = m_{1/2}$$



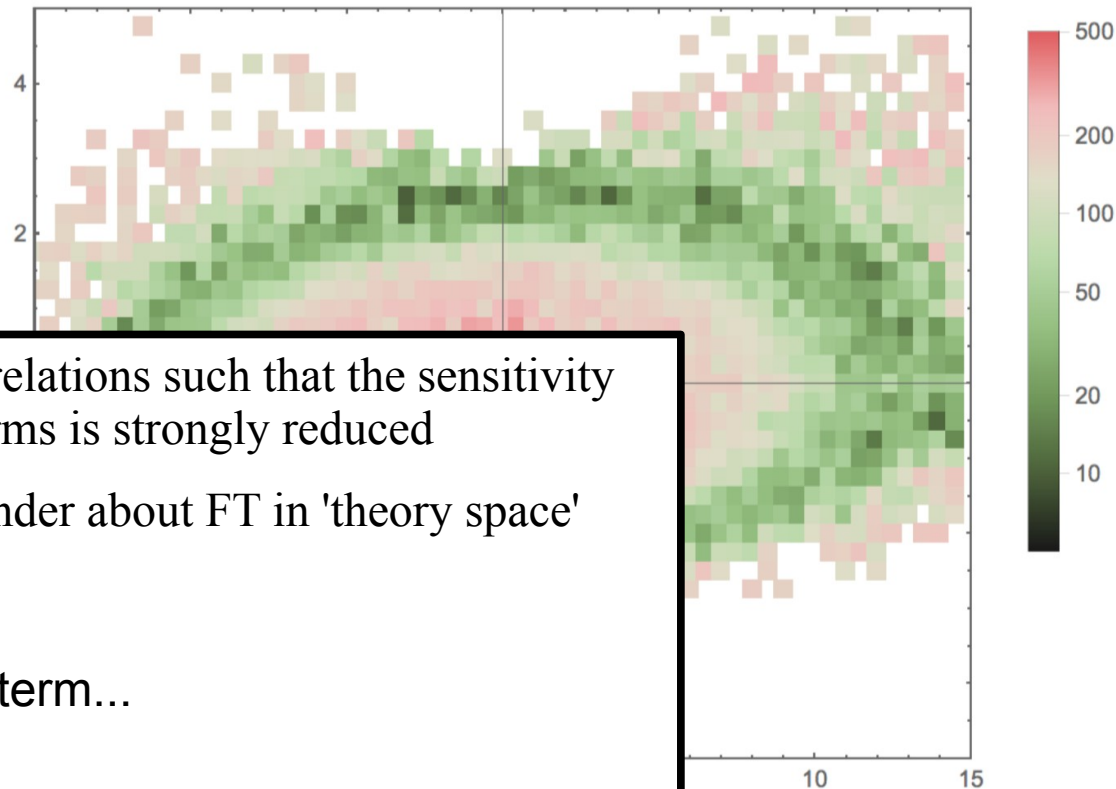
The gaugino focus point

- > Assume fixed ratios of gaugino masses
- > Possible also in GUTs

$$M_1 = a \cdot m_{1/2}$$

$$M_2 = b \cdot m_{1/2}$$

$$M_3 = m_{1/2}$$



- > Can find correlations such that the sensitivity to the soft terms is strongly reduced
- > You may wonder about FT in 'theory space'
- > μ not a soft term...

Comment on loop corrections

- > So far assumed tree-level relation for EWSB condition

$$\left. \frac{\partial V^{(L)}}{\partial v_u} \right|_{\tan\beta \rightarrow \infty} \equiv 0 = (m_{h_u}^2 + \mu^2 + \frac{1}{8}(g_1^2 + g_2^2)v^2)v + \Sigma_u$$

How to parametrise Σ_u ?

① $\Sigma_u \equiv v\Sigma_{uu}$

$$\rightarrow \frac{1}{2}M_Z^2 = -|\mu|^2 - m_{H_u}^2 + \Sigma_{uu}$$

no change in FT; only valid if Σ_{uu} is independent of v !

Comment on loop corrections

- > So far assumed tree-level relation for EWSB condition

$$\left. \frac{\partial V^{(L)}}{\partial v_u} \right|_{\tan\beta \rightarrow \infty} \equiv 0 = (m_{h_u}^2 + \mu^2 + \frac{1}{8}(g_1^2 + g_2^2)v^2)v + \Sigma_u$$

How to parametrise Σ_u ?

① $\Sigma_u \equiv v\Sigma_{uu}$

$$\rightarrow \frac{1}{2}M_Z^2 = -|\mu|^2 - m_{H_u}^2 + \Sigma_{uu}$$

no change in FT; only valid if Σ_{uu} is independent of v !

② $\Sigma_u \equiv \Sigma_1 v + \Sigma_2 v^2 + \Sigma_3 v^3$

$$\rightarrow \Delta_\mu = \frac{8\mu^2}{(g_1^2 + g_2^2 + 8\Sigma_3)v^2 + 4\Sigma_2 v}$$

Comment on loop corrections

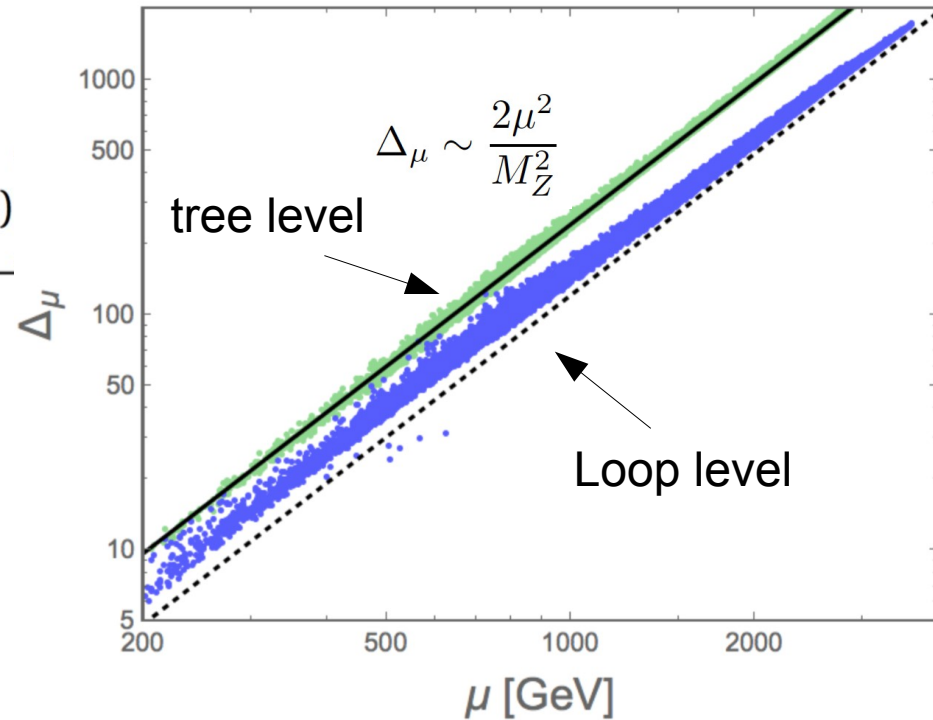
- > Result for leading stop correction

$$\Sigma_1 = - \frac{3m_t^2 Y_t^2 \left(-2 \log \left(\frac{2m_t^2 + v^2 Y_t^2}{Q^2} \right) + 2 + \log(4) \right)}{32\pi^2} \Delta_\mu$$

$$\Sigma_2 = 0$$

$$\Sigma_3 = \frac{3Y_t^4 \left(\log \left(\frac{2m_t^2 + v^2 Y_t^2}{Q^2} \right) - \log \left(\frac{v^2 Y_t^2}{Q^2} \right) \right)}{32\pi^2}$$

$$r^{FT} \equiv \frac{\Delta_\mu^{\text{Loop}}}{\Delta_\mu^{\text{Tree}}} \simeq \left(1 + \frac{3}{4\pi^2} \frac{Y_t^4}{g_1^2 + g_2^2} \log \left(\frac{m_t^2}{m_{\tilde{t}}^2} \right) \right)^{-1} \simeq \frac{1}{2}$$



- > Reduction of about ½ when including loop corrections

Comment on loop corrections

- > Result for leading stop correction

$$\Sigma_1 = - \frac{3m_t^2 Y_t^2 \left(-2 \log \left(\frac{2m_t^2 + v^2 Y_t^2}{Q^2} \right) + 2 + \log(4) \right)}{32\pi^2}$$

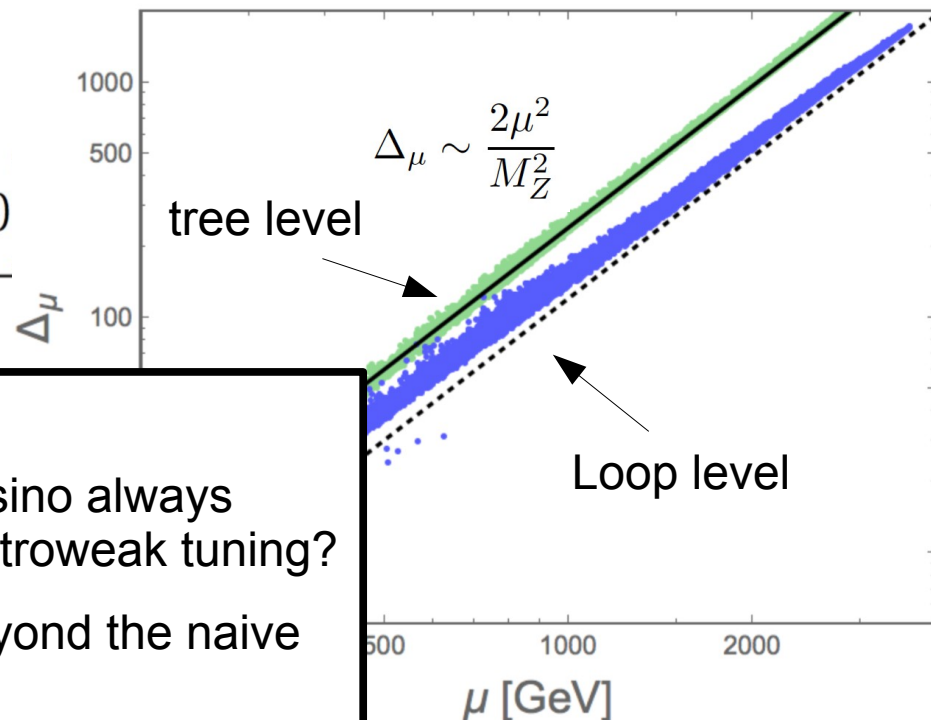
$$\Sigma_2 = 0$$

$$\Sigma_3 = \frac{3Y_t^4}{16\pi^2} \left(\log \left(\frac{2m_t^2 + v^2 Y_t^2}{Q^2} \right) + 2 + \log(4) \right)$$

$$r^{FT} \equiv \frac{\Delta_\mu^{\text{Loop}}}{\Delta_\mu^{\text{Tree}}} \simeq \left(1 + \frac{\Delta_\mu^{\text{Loop}}}{\Delta_\mu^{\text{Tree}}} \right)$$

- > Reduction of

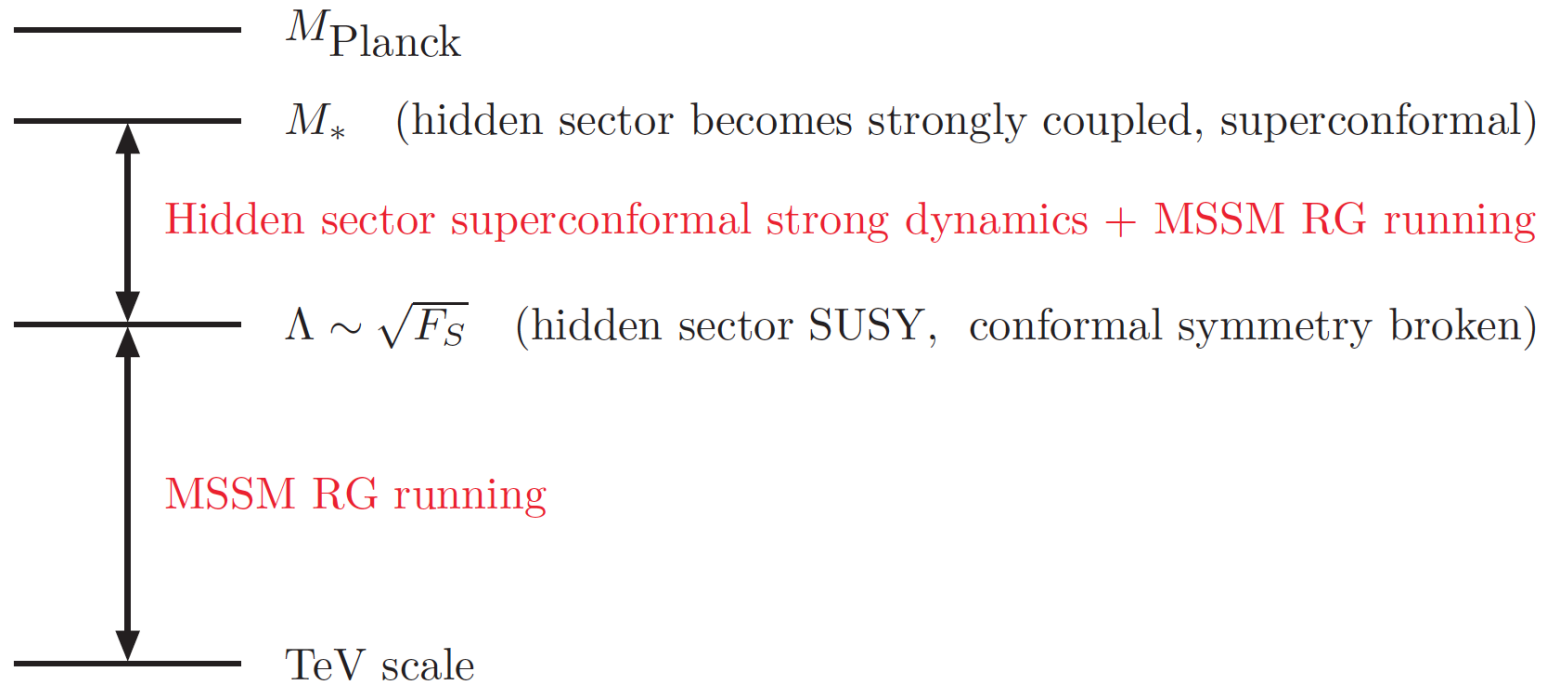
- > μ not a soft term...
- > So does a 1 TeV Higgsino always require significant electroweak tuning?
- > Two possible ways beyond the naive FT arguments
- > The μ -term could be correlated with the soft terms ($\rightarrow \mu$ problem)
- > The Higgsino mass could arise from something beyond the μ -term



A further correlation

- > The μ -term could be correlated with the soft terms

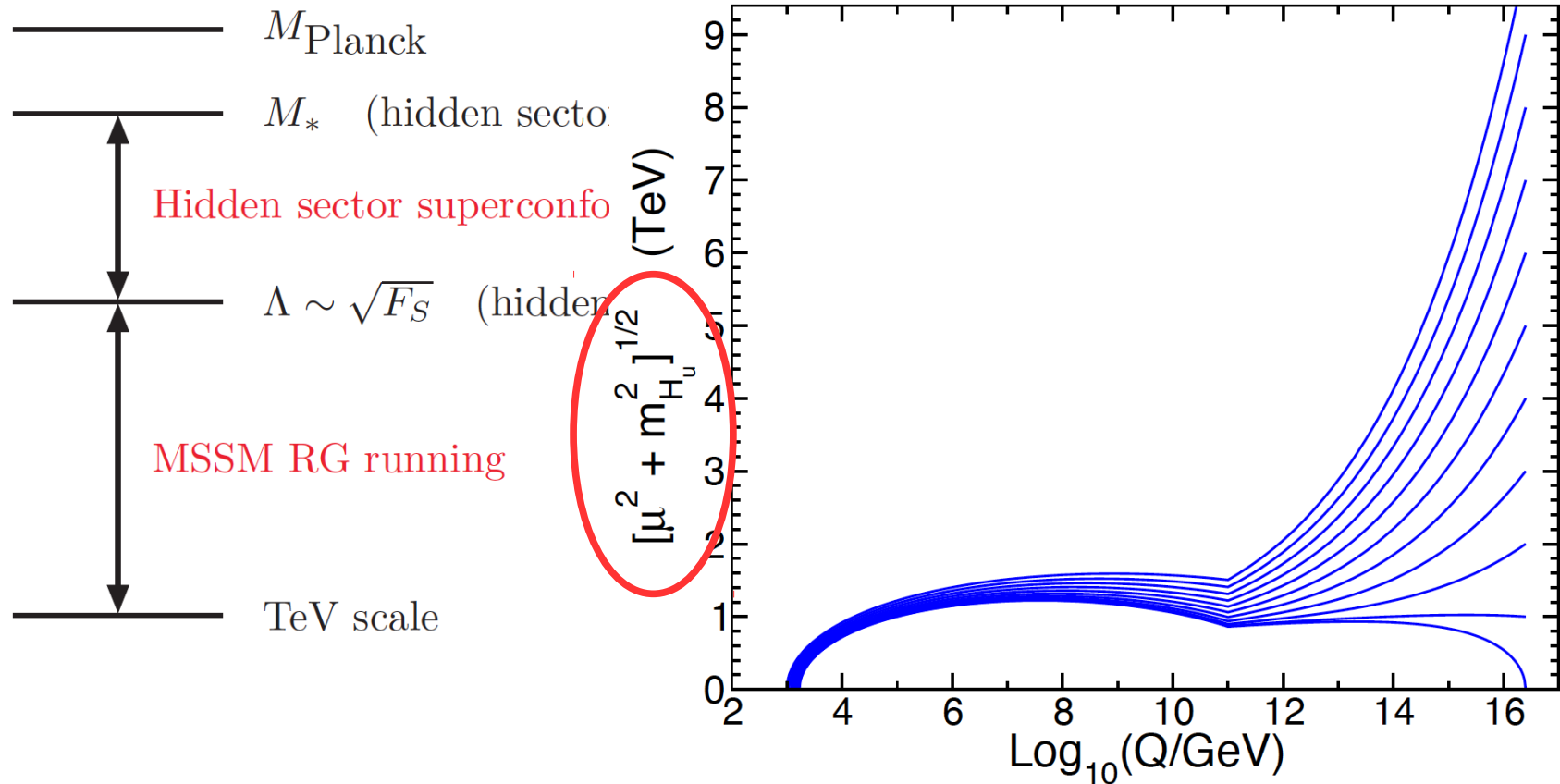
S Martin, 1712.05806



A further correlation

- > The μ -term could be correlated with the soft terms

S Martin, 1712.05806



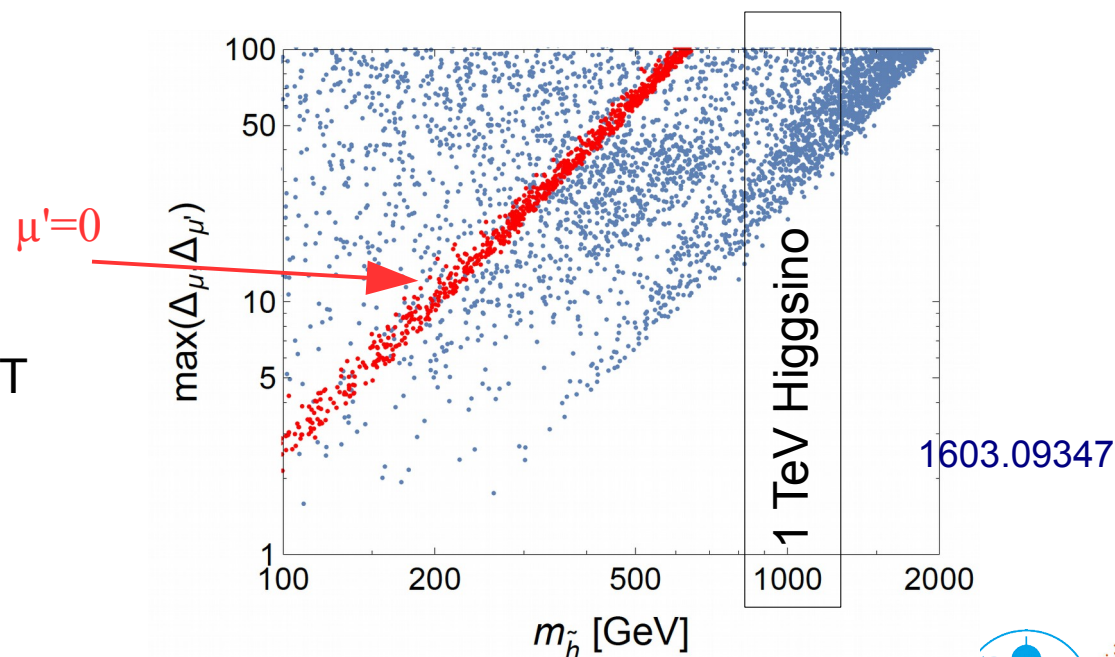
A new contribution to the Higgsino mass

- Non-standard SUSY breaking terms (in the classification of S Martin: 'maybe-soft')

$$\mathcal{L}_{NH} = \mu' \tilde{h}_d \tilde{h}_u + T'_{u,ij} h_d^* \tilde{u}_{R,i}^* \tilde{q}_j + T'_{d,ij} h_u^* \tilde{d}_{R,i}^* \tilde{q}_j + T'_{e,ij} h_u^* \tilde{e}_{R,i}^* \tilde{l}_j + \text{h.c.}$$

- μ' contributes to the Higgsino mass ($m_h \sim \mu + \mu'$) but does not enter the tadpole equations

- Can significantly reduce FT



Embedding this into a model

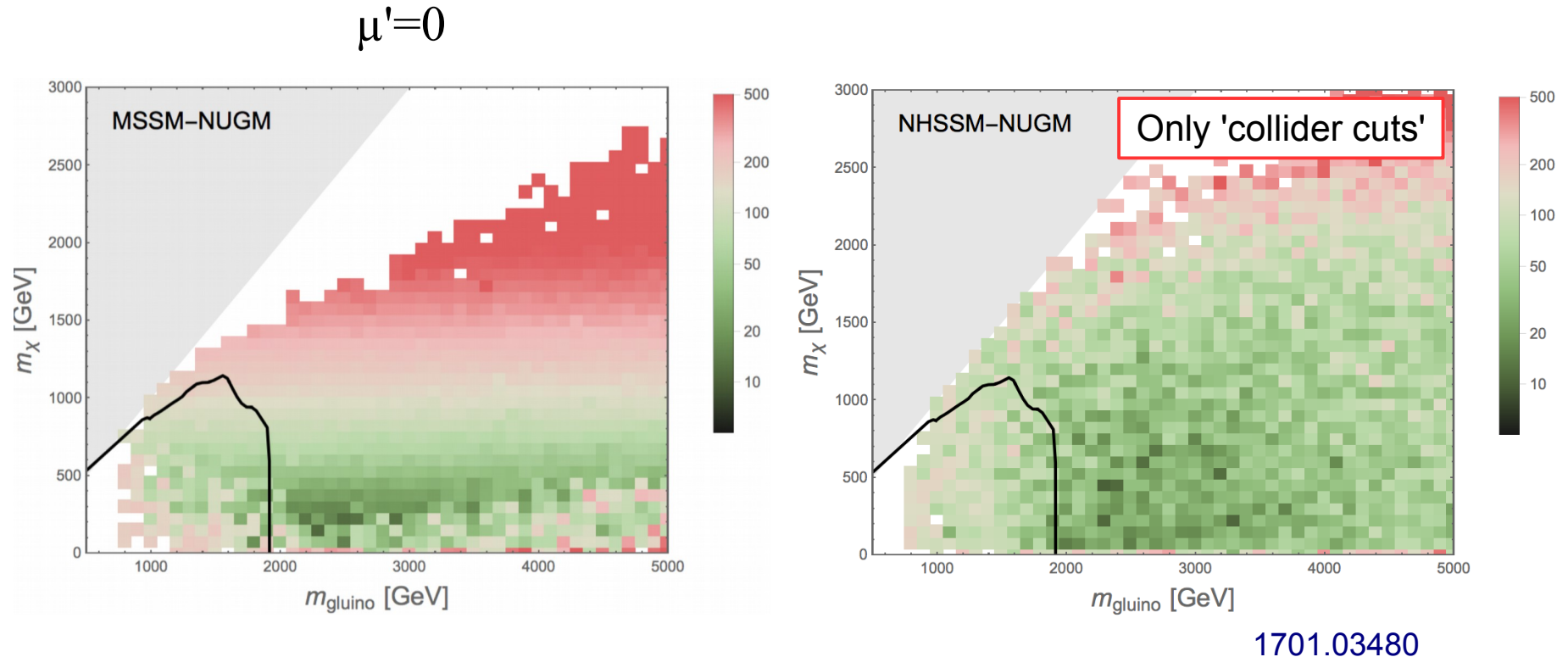
- Studied different MSSM variants with GUT boundary conditions

	$m_{h_u}^2$	$m_{h_d}^2$	M_1	M_2	M_3	μ'	A'_0
CMSSM	m_0^2	m_0^2	$m_{1/2}$	$m_{1/2}$	$m_{1/2}$	-	-
MSSM-NUHM	$m_{h_u}^2$	$m_{h_d}^2$	$m_{1/2}$	$m_{1/2}$	$m_{1/2}$	-	-
MSSM-NUGM	m_0^2	m_0^2	$a \cdot m_{1/2}$	$b \cdot m_{1/2}$	$m_{1/2}$	-	-
CNHSSM	m_0^2	m_0^2	$m_{1/2}$	$m_{1/2}$	$m_{1/2}$	μ'	A'_0
NHSSM-NUHM	$m_{h_u}^2$	$m_{h_d}^2$	$m_{1/2}$	$m_{1/2}$	$m_{1/2}$	μ'	A'_0
NHSSM-NUGM	m_0^2	m_0^2	$a \cdot m_{1/2}$	$b \cdot m_{1/2}$	$m_{1/2}$	μ'	A'_0

1701.03480

Results non-universal gaugino masses

- Region of small FT can be well beyond LHC reach



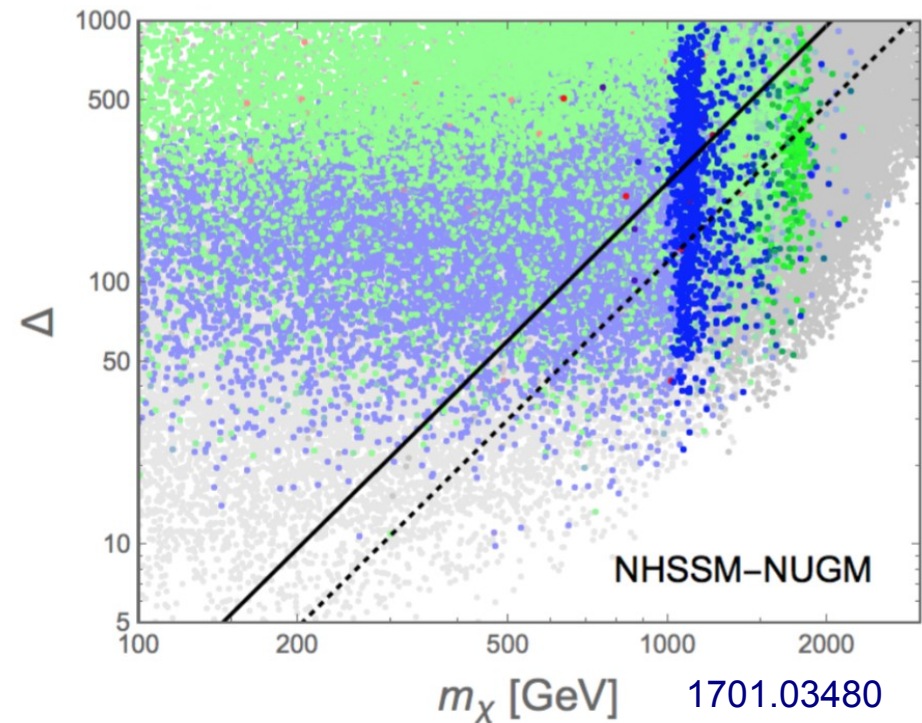
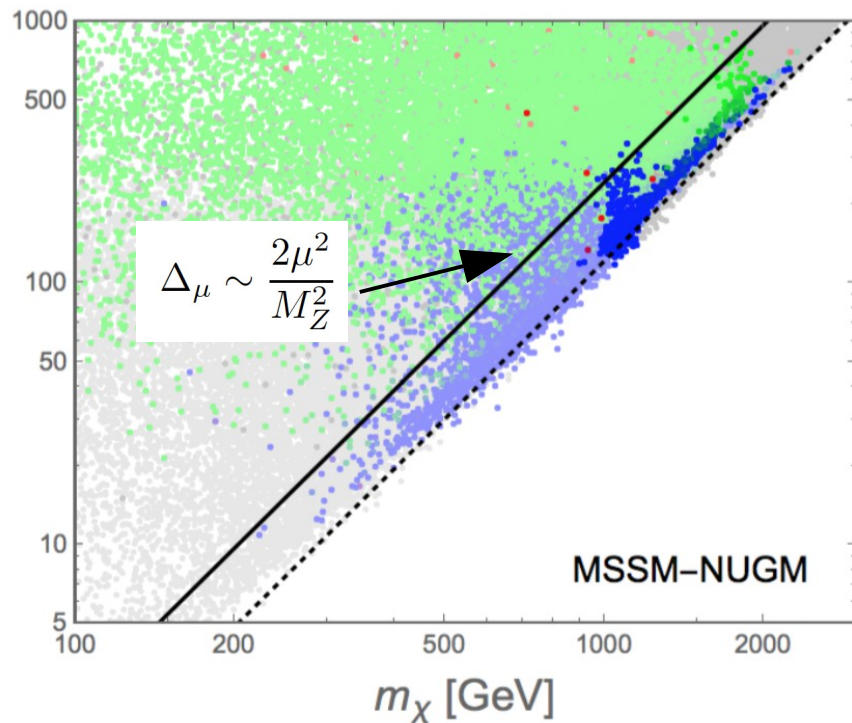
- Allowing for DM underabundance FT can be as small as 10.

Results non-universal gaugino masses

- > A 1 TeV Higgsino can be quite natural

$$\mu' = 0$$

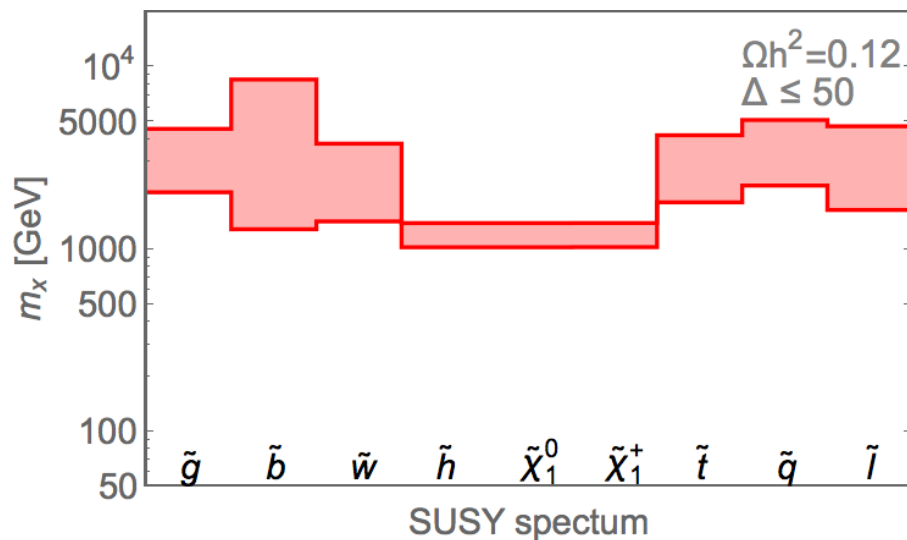
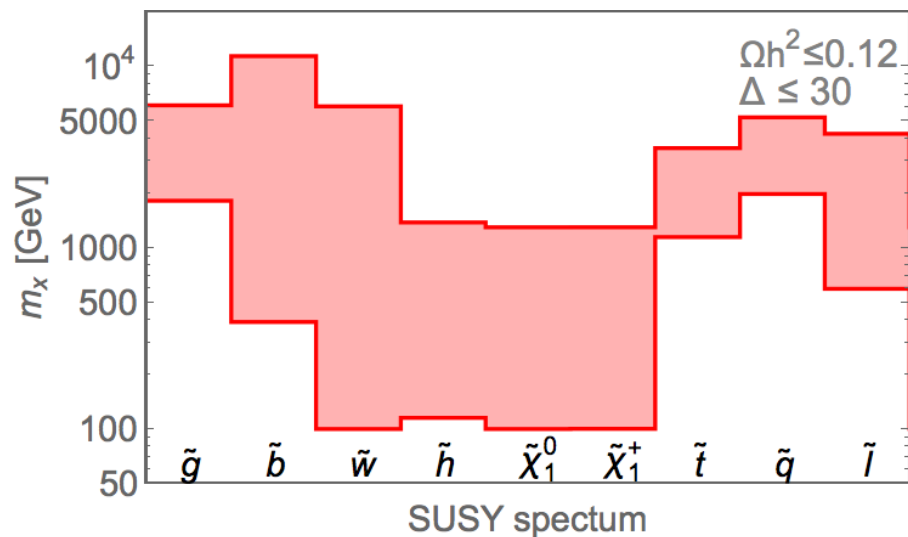
- bino
- Higgsino
- Wino



1701.03480

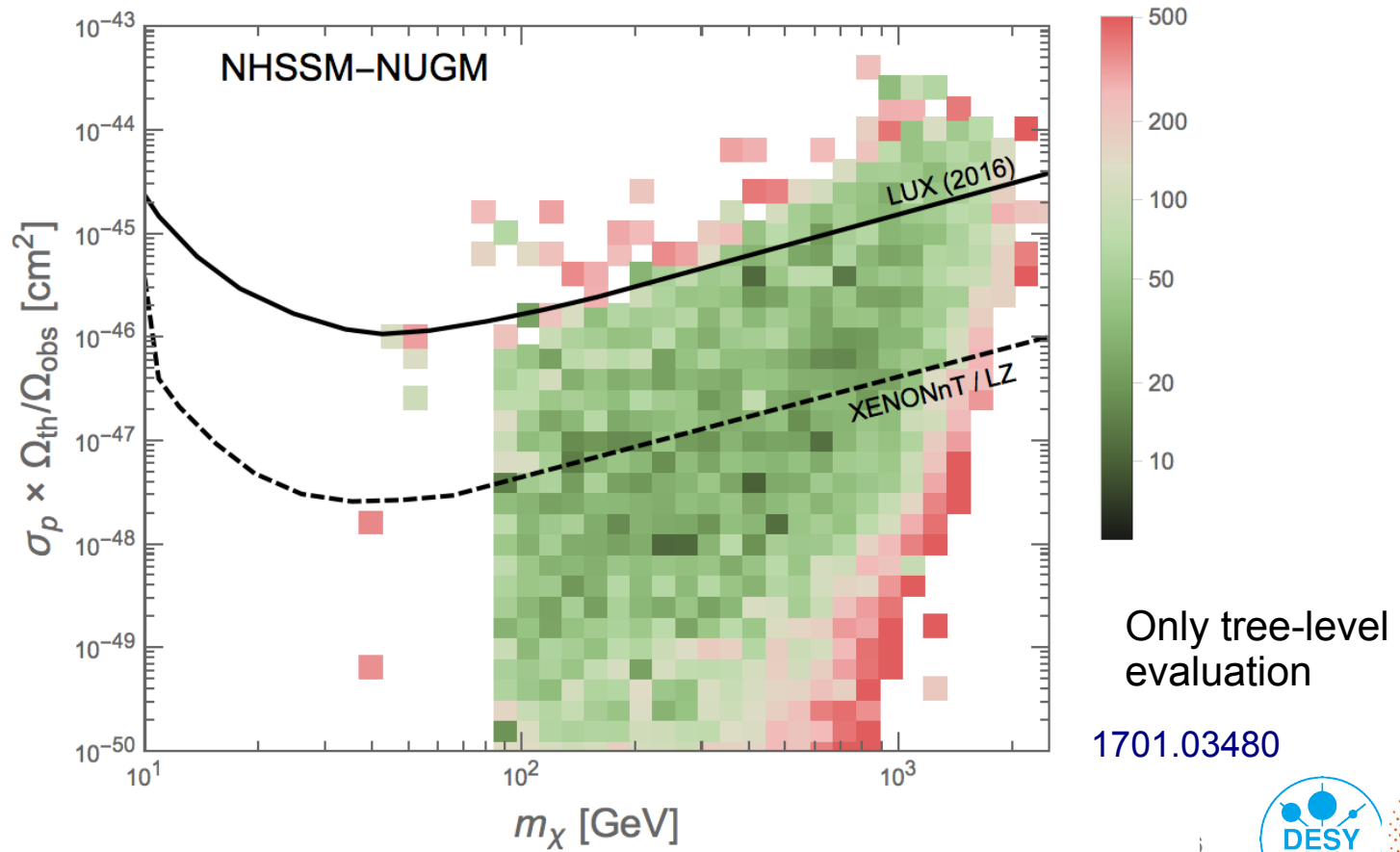
The SUSY spectrum

- What does the spectrum in the regions with low FT look like?



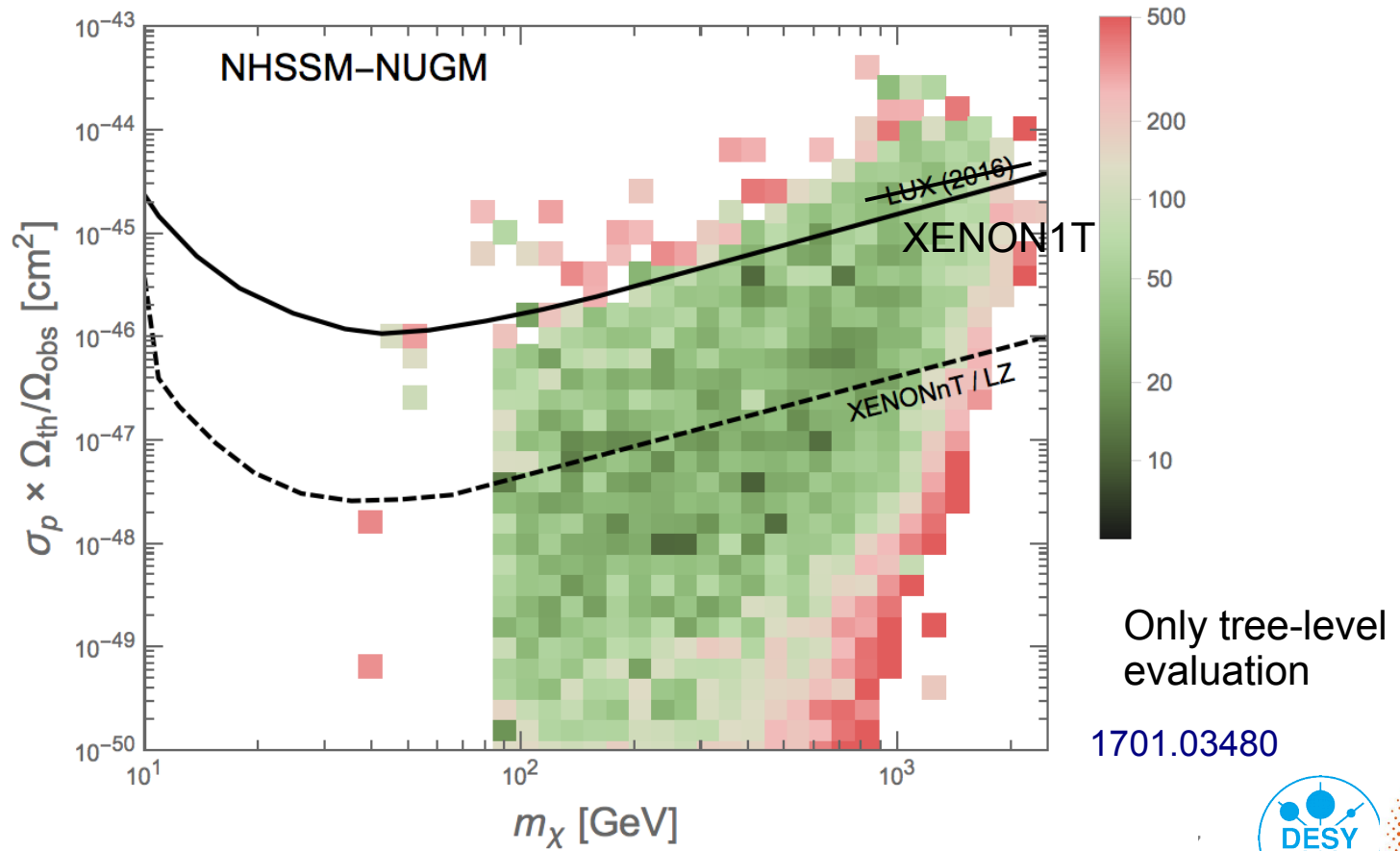
Prospects for direct detection

- Prospects for direct detection
- No lower bound on relic abundance (and rescaled) – other DM component



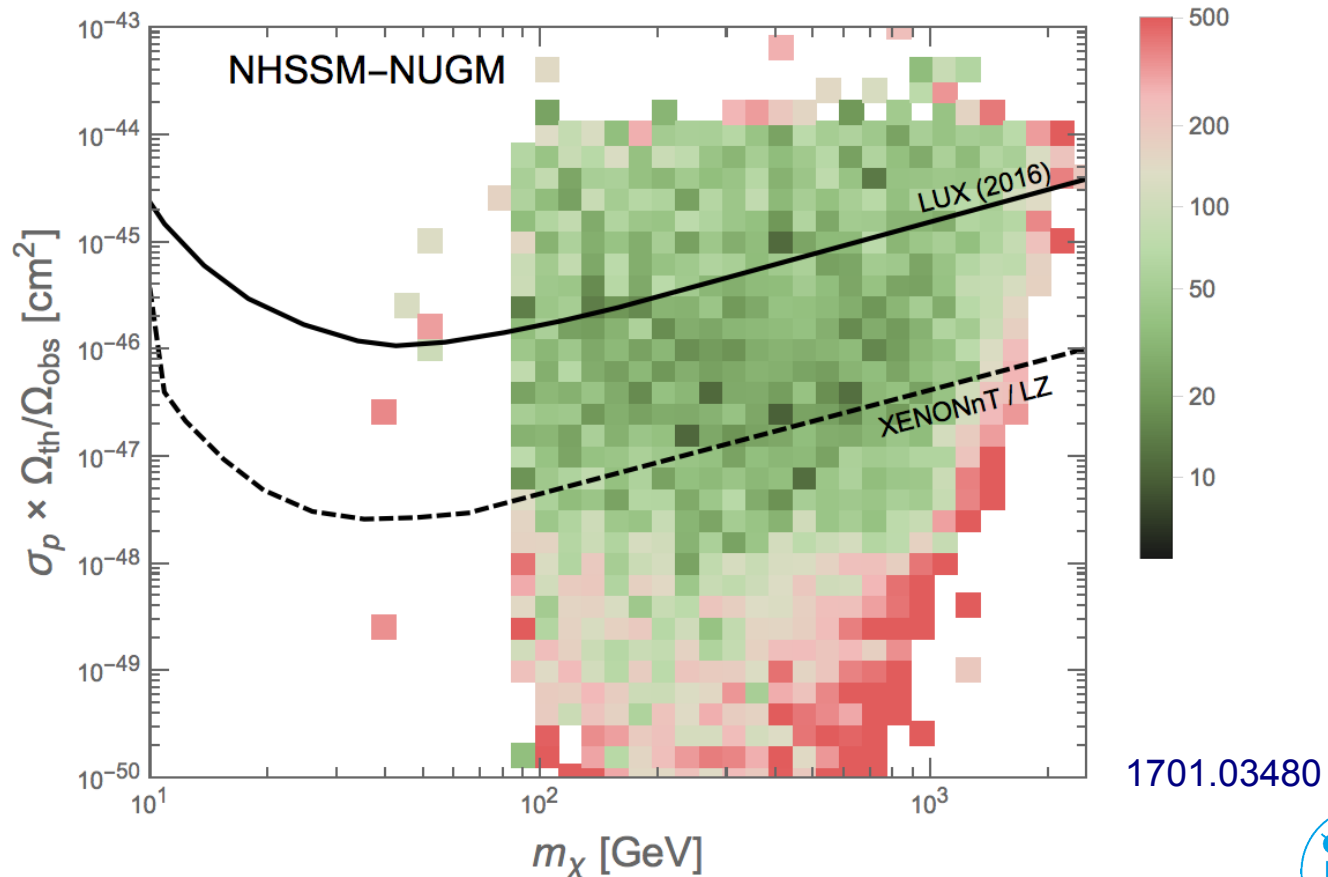
Prospects for direct detection

- Prospects for direct detection
- No lower bound on relic abundance (and rescaled) – other DM component



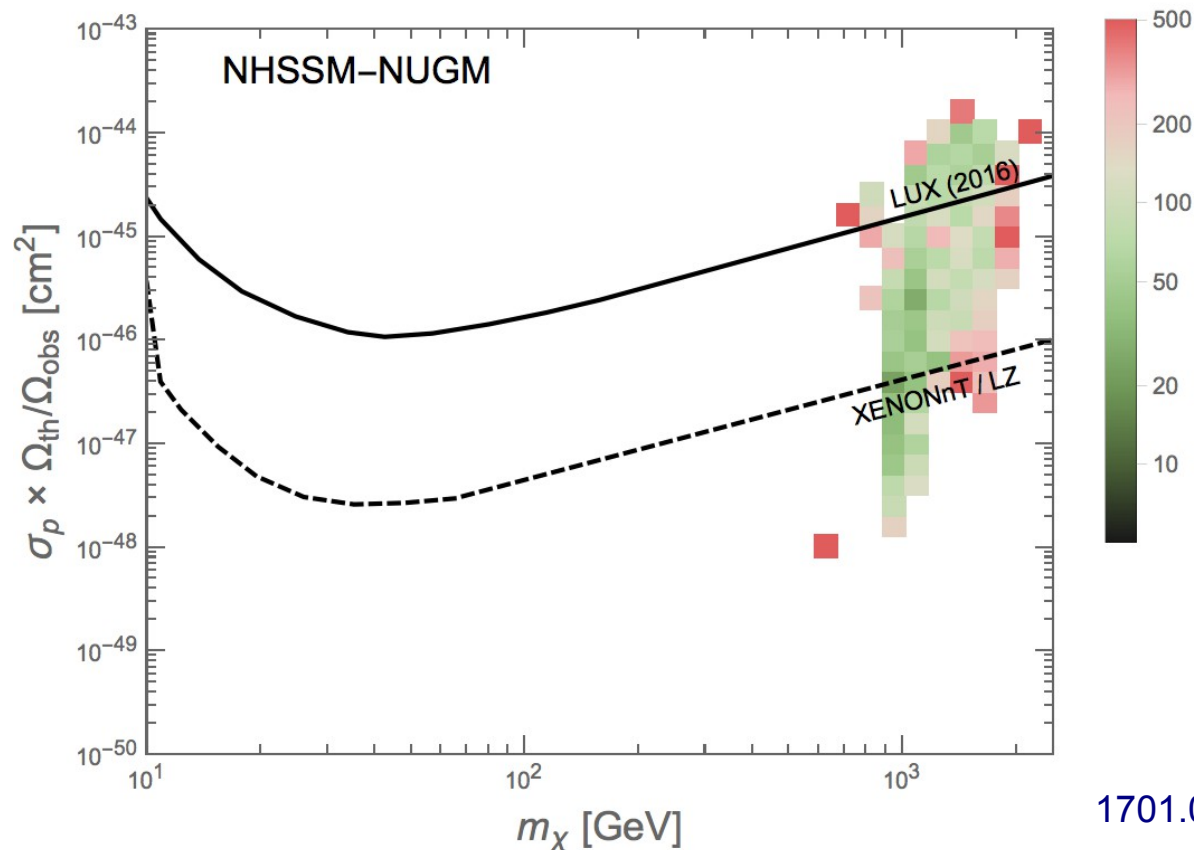
Prospects for direct detection

- Prospects for direct detection
- No lower bound on relic abundance (not rescaled) – non-thermal production (gravitino decay)



Prospects for direct detection

- Prospects for direct detection
- Correct (thermal) relic abundance



1701.03480

Summary

- > What looks unnatural from an IR perspective **might** still look natural from the UV
- > Extra Higgsino mass contribution μ' could help
- > To do: build a UV model



- > SUSY could well be beyond the LHC reach
- > Good chances at direct detection experiments to find it

Summary

- > What looks unnatural from an IR perspective **might** still look natural from the UV
- > Extra Higgsino mass contribution μ' could help
- > To do: build a UV model

Thank you!



- > SUSY could well be beyond the LHC reach
- > Good chances at direct detection experiments to find it