

Looking for New Naturally Aligned Higgs Bosons at the LHC

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Based on (i) R. Battye, G. Brawn, A.P., JHEP1108 (2011) 020 (73p),
(ii) A.P., Phys. Lett. B706 (2012) 465; PRD93 (2016) 075012,
(iii) P.S.B. Dev, A.P., JHEP1412 (2014) 024,
(iv) P.S.B. Dev, A.P. et al, work in progress.

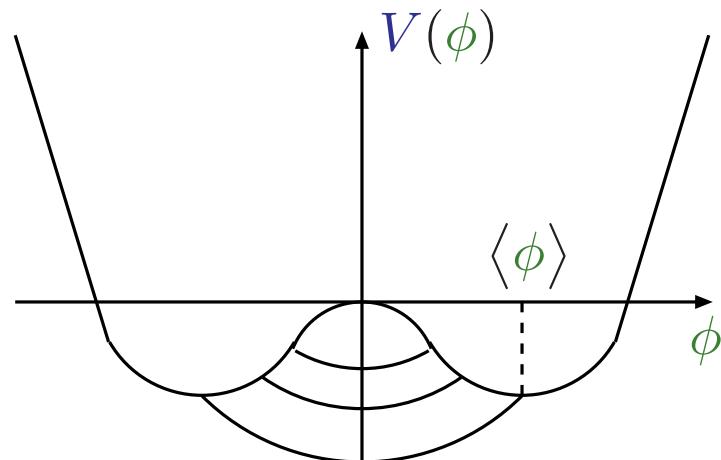
Outline:

- The Standard Theory of Electroweak Symmetry Breaking: SM
- The Two Higgs Doublet Model (2HDM) Potential
- SM Alignment in the 2HDM
- Symmetries of the 2HDM Potential
- Phenomenological Implications at the LHC
- Conclusions

- The Standard Theory of Electroweak Symmetry Breaking

Higgs Mechanism in the SM: $SU(3)_{\text{colour}} \otimes SU(2)_L \otimes U(1)_Y$

[P. W. Higgs '64; F. Englert, R. Brout '64.]



Higgs potential $V(\phi)$

$$V(\phi) = -m^2 \phi^\dagger \phi + \lambda (\phi^\dagger \phi)^2 .$$

Ground state:

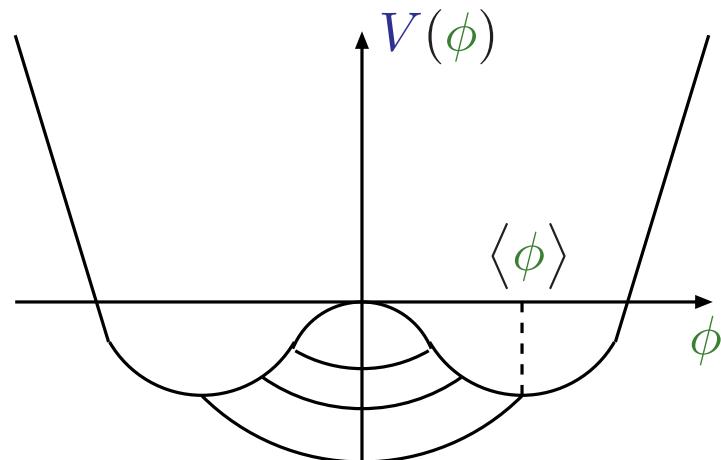
$$\langle \phi \rangle = \sqrt{\frac{m^2}{2\lambda}} \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

carries weak charge, but no electric charge and colour.

- The Standard Theory of Electroweak Symmetry Breaking

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Custodial Symmetry of the SM with $g' = Y_f = 0$ and $V(\phi)$:

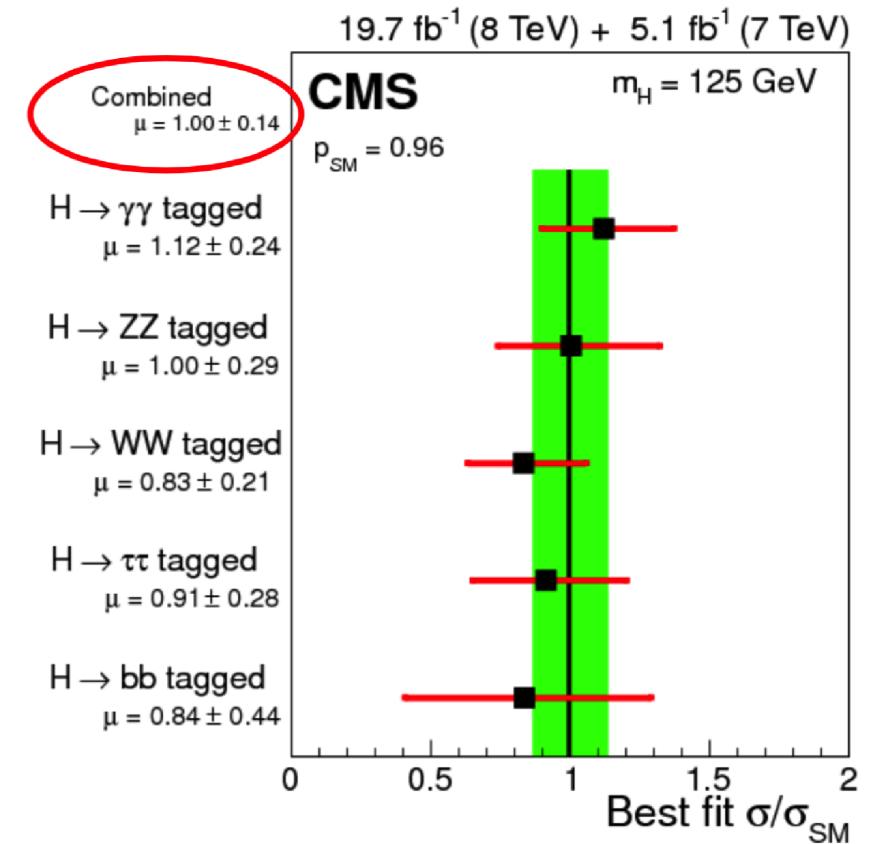
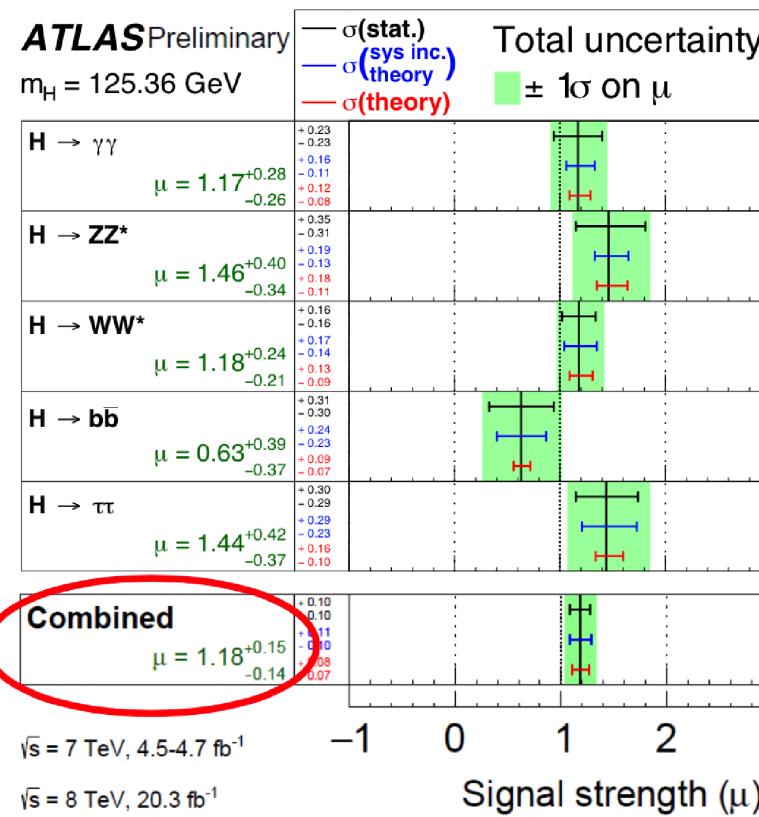
[P. Sikivie, L. Susskind, M. B. Voloshin, V. I. Zakharov '80.]

$$\Phi \equiv (\phi, i\sigma^2 \phi^*) \mapsto \Phi' \equiv U_L \Phi U_C ,$$

with $U_L \in SU(2)_L$ and $U_C \in SU(2)_C$, and $SU(2)_L \otimes SU(2)_C / \mathbb{Z}_2 \simeq SO(4)$.

Higgs Boson @ LHC: Signal Strength for Decay Modes

Signal strength: $\mu = \sigma_{\text{observed}}/\sigma_{\text{SM}}$



- Results consistent with SM

• The 2HDM Potential

[T. D. Lee '73]

$$\begin{aligned} V = & -\mu_1^2(\phi_1^\dagger \phi_1) - \mu_2^2(\phi_2^\dagger \phi_2) - m_{12}^2(\phi_1^\dagger \phi_2) - m_{12}^{*2}(\phi_2^\dagger \phi_1) \\ & + \lambda_1(\phi_1^\dagger \phi_1)^2 + \lambda_2(\phi_2^\dagger \phi_2)^2 + \lambda_3(\phi_1^\dagger \phi_1)(\phi_2^\dagger \phi_2) + \lambda_4(\phi_1^\dagger \phi_2)(\phi_2^\dagger \phi_1) \\ & + \frac{\lambda_5}{2}(\phi_1^\dagger \phi_2)^2 + \frac{\lambda_5^*}{2}(\phi_2^\dagger \phi_1)^2 + \lambda_6(\phi_1^\dagger \phi_1)(\phi_1^\dagger \phi_2) + \lambda_6^*(\phi_1^\dagger \phi_1)(\phi_2^\dagger \phi_1) \\ & + \lambda_7(\phi_2^\dagger \phi_2)(\phi_1^\dagger \phi_2) + \lambda_7^*(\phi_2^\dagger \phi_2)(\phi_2^\dagger \phi_1). \end{aligned}$$

V has 4 real mass parameters and 10 real quartic couplings.

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 & + \lambda_1(\phi_1^\dagger \phi_1)^2 + \lambda_2(\phi_2^\dagger \phi_2)^2 + \lambda_3(\phi_1^\dagger \phi_1)(\phi_2^\dagger \phi_2) + \lambda_4(\phi_1^\dagger \phi_2)(\phi_2^\dagger \phi_1) \\
 & + \frac{\lambda_5}{2}(\phi_1^\dagger \phi_2)^2 + \frac{\lambda_5^*}{2}(\phi_2^\dagger \phi_1)^2 + \lambda_6(\phi_1^\dagger \phi_1)(\phi_1^\dagger \phi_2) + \lambda_6^*(\phi_1^\dagger \phi_1)(\phi_2^\dagger \phi_1) \\
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Remark on the tree-level MSSM Higgs potential:

$$\begin{aligned}
 m_{12}^2 &= -B\mu, \quad \lambda_1 = \lambda_2 = -\frac{1}{8}(g_w^2 + g'^2), \quad \lambda_3 = -\frac{1}{4}(g_w^2 - g'^2), \\
 \lambda_4 &= \frac{1}{2}g_w^2, \quad \lambda_5 = \lambda_6 = \lambda_7 = 0.
 \end{aligned}$$

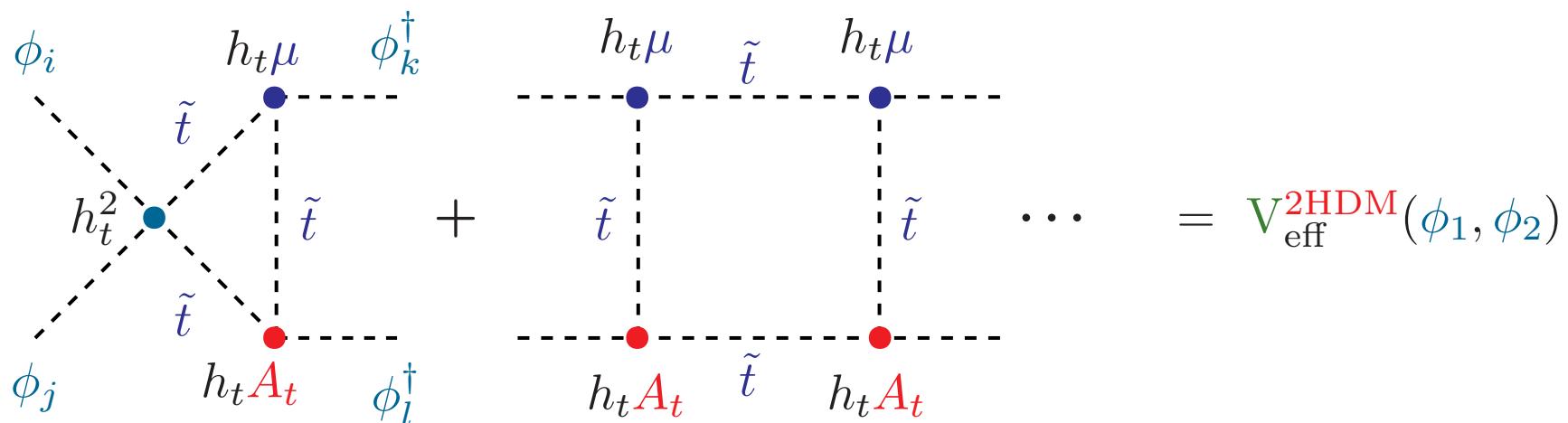
$\phi_2 \rightarrow e^{i \arg m_{12}^2} \phi_2 \implies$ Tree-level Higgs potential is invariant under CP

Loop effects are sizeable leading to a General 2HDM.

[J. Ellis, G. Ridolfi, F. Zwirner '91;
Y. Okada, M. Yamaguchi, T. Yanagida '91;
H.E. Haber, R. Hempfling '91.]

Radiative Higgs-sector CP Violation:

[A.P. '98; A.P., C.E.M. Wagner '99]



CP-violating terms are $\propto \text{Im} (m_{12}^{2*} \mu A_{t,b}) \neq 0$

Latest version:

CPsuperH2.3: J. S. Lee, M. Carena, J. Ellis, A.P., C. E. M. Wagner,
JHEP1602 (2016) 123.

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- Physical (CP-conserving) spectrum:

CP-even Higgs bosons H and h ; CP-odd scalar a ; charged scalars h^\pm .

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- Higgs coupling to gauge bosons $V = W, Z$:

$$g_{HVV} = \cos(\beta - \alpha) , \quad g_{hVV} = \sin(\beta - \alpha) ,$$

where $\tan \beta = \langle \phi_2 \rangle / \langle \phi_1 \rangle$ related to the diagonalization of the CP-odd mass matrix and the mixing angle α diagonalizes the CP-even mass matrix.

- **SM Alignment** in the 2HDM

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- **SM Alignment** $\beta \rightarrow \alpha$:

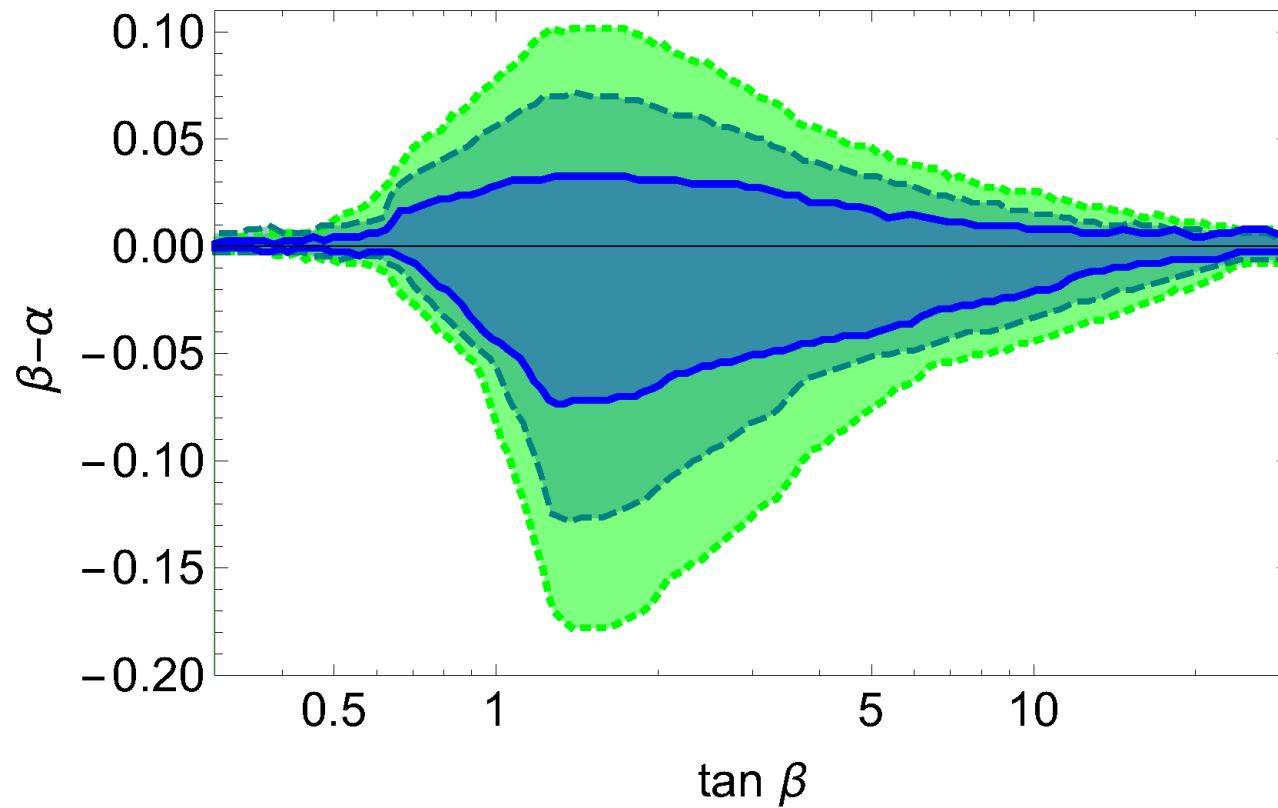
(i) **Decoupling**: $M_{h^\pm} \gg M_H$ [J. F. Gunion, H. E. Haber, Phys. Rev. D67 (2003) 075019.]

(ii) **Fine-tuning**: [P.S.B. Dev, A.P., JHEP1412 (2014) 024.]

$$\lambda_7 t_\beta^4 - (2\lambda_2 - \lambda_{345})t_\beta^3 + 3(\lambda_6 - \lambda_7)t_\beta^2 + (2\lambda_1 - \lambda_{345})t_\beta - \lambda_6 = 0$$

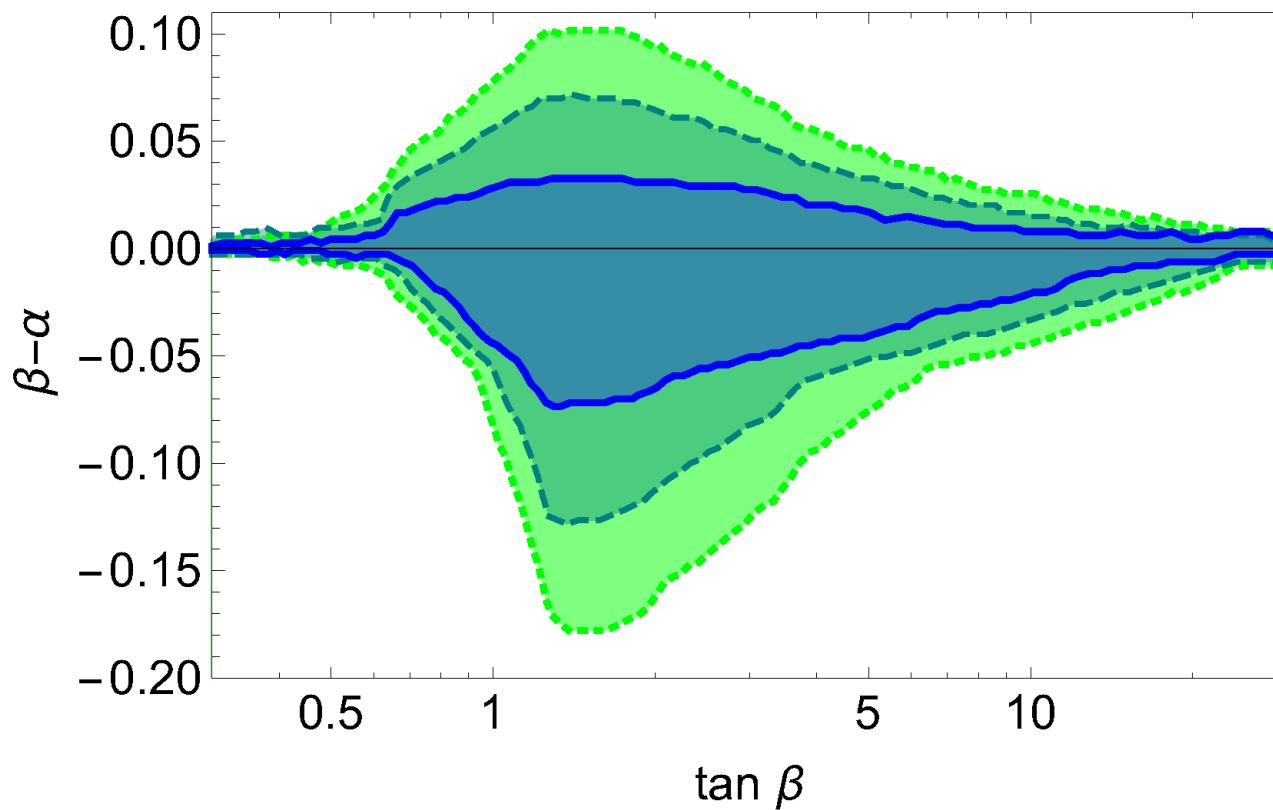
- Global Fit to the SM Alignment $\beta \rightarrow \alpha$:

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- **Natural Alignment (without decoupling or fine-tuning):**

[P.S.B. Dev, A.P., JHEP1412 (2014) 024.]

$$\lambda_1 = \lambda_2 = \lambda_{345}/2, \quad \lambda_6 = \lambda_7 = 0.$$

References (*an incomplete list on SM Alignment in the 2HDM*)

- **On the SM Higgs basis (also Decoupling of FCNC Effects):**
H. Georgi and D. V. Nanopoulos, Phys. Lett. B82 (1979) 95.
- **Alignment via Decoupling:**
J. F. Gunion, H. E. Haber, Phys. Rev. D67 (2003) 075019.
- I. F. Ginzburg, M. Krawczyk, Phys. Rev. D72 (2005) 115013.
- **Alignment via Fine-tuning:**
P. H. Chankowski, T. Farris, B. Grzadkowski, J. F. Gunion, J. Kalinowski, M. Krawczyk, Phys. Lett. B496 (2000) 195
- A. Delgado, G. Nardini, M. Quiros, JHEP **1307** (2013) 054.
- M. Carena, I. Low, N. R. Shah, C. E. M. Wagner, JHEP1404 (2014) 015.
- **Natural Alignment without Decoupling and without Fine-tuning:**
P.S.B. Dev, A.P., JHEP1412 (2014) 024.

References (*an incomplete list on symmetries in the 2HDM*)

- **Spontaneous CP Violation:** T. D. Lee, Phys. Rev. D **8** (1973) 1226.
- **Z_2 symmetry:** S. L. Glashow, S. Weinberg, Phys. Rev. D **15** (1977) 1958.
- **Inert Z_2 symmetry:** N. G. Deshpande, E. Ma, Phys. Rev. D **18** (1978) 2574.
- **PQ U(1) symmetry:** R. D. Peccei, H. R. Quinn, Phys. Rev. Lett. **38** (1977) 1440.
- **Custodial $SU(2)_L$ -preserving symmetry:**
P. Sikivie, L. Susskind, M. B. Voloshin, V. I. Zakharov, Nucl. Phys. B **173** (1980) 189.
- **Bilinear formalism:**
M. Maniatis, A. von Manteuffel, O. Nachtmann, F. Nagel, EPJC **48** (2006) 805;
C. C. Nishi, Phys. Rev. D **74** (2006) 036003 [Erratum-ibid. D **76** (2007) 119901].
- **$SU(2)_L \otimes U(1)_Y$ -preserving symmetries:**
I. P. Ivanov, Phys. Rev. D **75** (2007) 035001 [Erratum-ibid. D **76** (2007) 039902].
- **Hypercustodial $SU(2)_L$ -preserving symmetries:**
R. A. Battye, G. D. Brawn, A.P., JHEP1108 (2011) 020.
- **On completeness and uniqueness of classification:**
A.P., Phys. Lett. B **706** (2012) 465.

- Symmetries of the 2HDM Potential

[R. A. Battye, G. D. Brawn, A.P., JHEP1108 (2011) 020.]

Introduce the $SU(2)_L$ -covariant 8D complex field multiplet

$$\Phi = \begin{pmatrix} \phi_1 \\ \phi_2 \\ i\sigma^2\phi_1^* \\ i\sigma^2\phi_2^* \end{pmatrix}, \quad \text{with } U_L \in SU(2)_L : \Phi \mapsto \Phi' = U_L \Phi.$$

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Φ satisfies the **Majorana constraint**

$$\Phi = C \Phi^*,$$

where C is the **charge conjugation 8D matrix**

$$C = \sigma^2 \otimes \sigma^0 \otimes \sigma^2 = \begin{pmatrix} \mathbf{0}_4 & \mathbf{1}_4 \\ -\mathbf{1}_4 & \mathbf{0}_4 \end{pmatrix} \otimes (-i\sigma_2).$$

- The $SO(1,5)$ Bilinear Formalism

Introduce the *null 6-Vector*

$$R^A = \Phi^\dagger \Sigma^A \Phi = \begin{pmatrix} \phi_1^\dagger \phi_1 + \phi_2^\dagger \phi_2 \\ \phi_1^\dagger \phi_2 + \phi_2^\dagger \phi_1 \\ -i [\phi_1^\dagger \phi_2 - \phi_2^\dagger \phi_1] \\ \phi_1^\dagger \phi_1 - \phi_2^\dagger \phi_2 \\ \phi_1^\top i\sigma^2 \phi_2 - \phi_2^\dagger i\sigma^2 \phi_1^* \\ -i [\phi_1^\top i\sigma^2 \phi_2 + \phi_2^\dagger i\sigma^2 \phi_1^*] \end{pmatrix},$$

with $A = \mu, 4, 5$ and

$$\Sigma^\mu = \frac{1}{2} \begin{pmatrix} \sigma^\mu & \mathbf{0}_2 \\ \mathbf{0}_2 & (\sigma^\mu)^\top \end{pmatrix} \otimes \sigma^0,$$

$$\Sigma^4 = \frac{1}{2} \begin{pmatrix} \mathbf{0}_2 & i\sigma^2 \\ -i\sigma^2 & \mathbf{0}_2 \end{pmatrix} \otimes \sigma^0, \quad \Sigma^5 = \frac{1}{2} \begin{pmatrix} \mathbf{0}_2 & -\sigma^2 \\ -\sigma^2 & \mathbf{0}_2 \end{pmatrix} \otimes \sigma^0.$$

- The 2HDM Potential in the SO(1,5) Formalism

$$V = -\frac{1}{2} M_A R^A + \frac{1}{4} L_{AB} R^A R^B ,$$

with

$$M_A = (\mu_1^2 + \mu_2^2, \quad 2\text{Re}(m_{12}^2), \quad -2\text{Im}(m_{12}^2), \quad \mu_1^2 - \mu_2^2, \quad 0, \quad 0) ,$$

$$L_{AB} = \begin{pmatrix} \lambda_1 + \lambda_2 + \lambda_3 & \text{Re}(\lambda_6 + \lambda_7) & -\text{Im}(\lambda_6 + \lambda_7) & \lambda_1 - \lambda_2 & 0 & 0 \\ \text{Re}(\lambda_6 + \lambda_7) & \lambda_4 + \text{Re}(\lambda_5) & -\text{Im}(\lambda_5) & \text{Re}(\lambda_6 - \lambda_7) & 0 & 0 \\ -\text{Im}(\lambda_6 + \lambda_7) & -\text{Im}(\lambda_5) & \lambda_4 - \text{Re}(\lambda_5) & -\text{Im}(\lambda_6 - \lambda_7) & 0 & 0 \\ \lambda_1 - \lambda_2 & \text{Re}(\lambda_6 - \lambda_7) & -\text{Im}(\lambda_6 - \lambda_7) & \lambda_1 + \lambda_2 - \lambda_3 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \end{pmatrix} .$$

- The 2HDM Potential in the SO(1,5) Formalism

$$V = -\frac{1}{2} M_A R^A + \frac{1}{4} L_{AB} R^A R^B ,$$

with

$$M_A = (\mu_1^2 + \mu_2^2, \quad 2\text{Re}(m_{12}^2), \quad -2\text{Im}(m_{12}^2), \quad \mu_1^2 - \mu_2^2, \quad 0, \quad 0) ,$$

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- Unitary Field Transformations:

[A.P., Phys. Lett. B706 (2012) 465.]

$$\text{Sp}(4) : \quad \Phi \mapsto \Phi' = U \Phi , \quad \text{with} \quad U \in \text{U}(4) \quad \underline{\text{and}} \quad UCU^\top = C$$

$$\text{SO}(5) : \quad R^I \mapsto R'^I = O^I_J R^J , \quad \text{with} \quad O \in \text{SO}(5) \subset \text{SO}(1,5)$$

$$\implies \text{SO}(5) \sim \text{Sp}(4)/\mathbf{Z}_2$$

- **Symmetries of the $U(1)_Y$ -Invariant 2HDM Potential**

$SO(5)$ -diagonally reduced basis: $\text{Im } \lambda_5 = 0$ and $\lambda_6 = \lambda_7$.

The 2HDM potential exhibits a total of 13 accidental symmetries:

Symmetry	μ_1^2	μ_2^2	m_{12}^2	λ_1	λ_2	λ_3	λ_4	$\text{Re } \lambda_5$	$\lambda_6 = \lambda_7$
$(Z_2)^2 \times SO(2)$	–	–	0	–	–	–	–	–	0
$O(2) \times O(2)$	–	–	0	–	–	–	–	0	0
✓ $O(3) \times O(2)$	–	μ_1^2	0	–	λ_1	–	$2\lambda_1 - \lambda_3$	0	0
$Z_2 \times O(2)$	–	–	Real	–	–	–	–	–	Real
$(Z_2)^3 \times O(2)$	–	μ_1^2	0	–	λ_1	–	–	–	0
✓ $Z_2 \times [O(2)]^2$	–	μ_1^2	0	–	λ_1	–	–	$2\lambda_1 - \lambda_{34}$	0
✓ $SO(5)$	–	μ_1^2	0	–	λ_1	$2\lambda_1$	0	0	0
$Z_2 \times O(4)$	–	μ_1^2	0	–	λ_1	–	0	0	0
$SO(4)$	–	–	0	–	–	–	0	0	0
$O(2) \times O(3)$	–	μ_1^2	0	–	λ_1	$2\lambda_1$	–	0	0
$(Z_2)^2 \times SO(3)$	–	μ_1^2	0	–	λ_1	–	–	$\pm\lambda_4$	0
$Z_2 \times O(3)$	–	μ_1^2	Real	–	λ_1	–	–	λ_4	Real
$SO(3)$	–	–	Real	–	–	–	–	λ_4	Real

✓: Natural SM Alignment



[P.S.B. Dev, A.P., JHEP1412 (2014) 024.]

• Symmetry Breaking Scenarios and pseudo-Goldstone Bosons

[A.P., Phys. Lett. B706 (2012) 465.]

No	Symmetry	Generators $T^a \leftrightarrow K^a$	Discrete Group Elements	Maximally Broken SO(5) Generators	Number of Pseudo- Goldstone Bosons
1	$Z_2 \times O(2)$	T^0	D_{CP1}	–	0
2	$(Z_2)^2 \times SO(2)$	T^0	D_{Z_2}	–	0
3	$(Z_2)^3 \times O(2)$	T^0	D_{CP2}	–	0
4	$O(2) \times O(2)$	T^3, T^0	–	T^3	1 (a)
✓ 5	$Z_2 \times [O(2)]^2$	T^2, T^0	D_{CP1}	T^2	1 (h)
✓ 6	$O(3) \times O(2)$	$T^{1,2,3}, T^0$	–	$T^{1,2}$	2 (h, a)
7	$SO(3)$	$T^{0,4,6}$	–	$T^{4,6}$	2 (h^\pm)
8	$Z_2 \times O(3)$	$T^{0,4,6}$	$D_{Z_2} \cdot D_{CP2}$	$T^{4,6}$	2 (h^\pm)
9	$(Z_2)^2 \times SO(3)$	$T^{0,5,7}$	$D_{CP1} \cdot D_{CP2}$	$T^{5,7}$	2 (h^\pm)
10	$O(2) \times O(3)$	$T^3, T^{0,8,9}$	–	T^3	1 (a)
11	$SO(4)$	$T^{0,3,4,5,6,7}$	–	$T^{3,5,7}$	3 (a, h^\pm)
12	$Z_2 \times O(4)$	$T^{0,3,4,5,6,7}$	$D_{Z_2} \cdot D_{CP2}$	$T^{3,5,7}$	3 (a, h^\pm)
✓ 13	$SO(5)$	$T^{0,1,2,\dots,9}$	–	$T^{1,2,8,9}$	4 (h, a, h^\pm)

✓ : Natural SM Alignment \mapsto

[P.S.B. Dev, A.P., JHEP1412 (2014) 024.]

- Phenomenological Implications at the LHC

- Maximally Symmetric Two Higgs Doublet Model

[P.S.B. Dev, A.P., JHEP1412 (2014) 024.]

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Symmetry Group: $G_\Phi = \text{SU}(2)_L \otimes \text{Sp}(4)/\mathbf{Z}_2 \simeq \text{SU}(2)_L \otimes \text{SO}(5)$.

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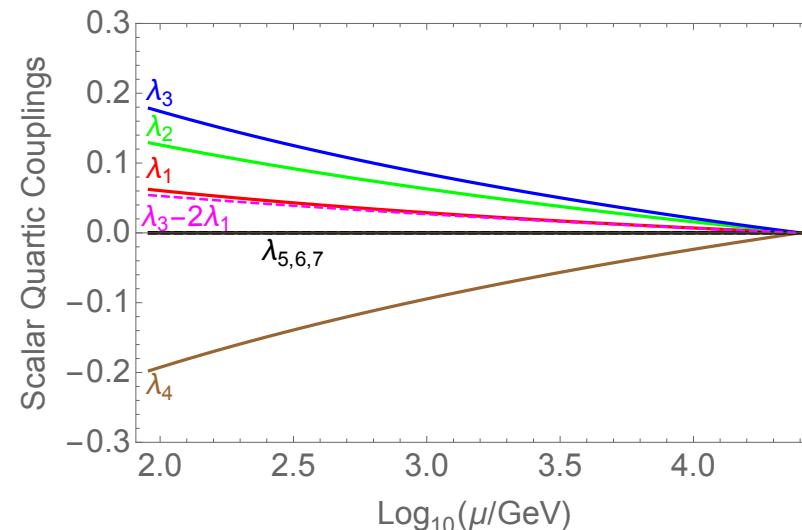
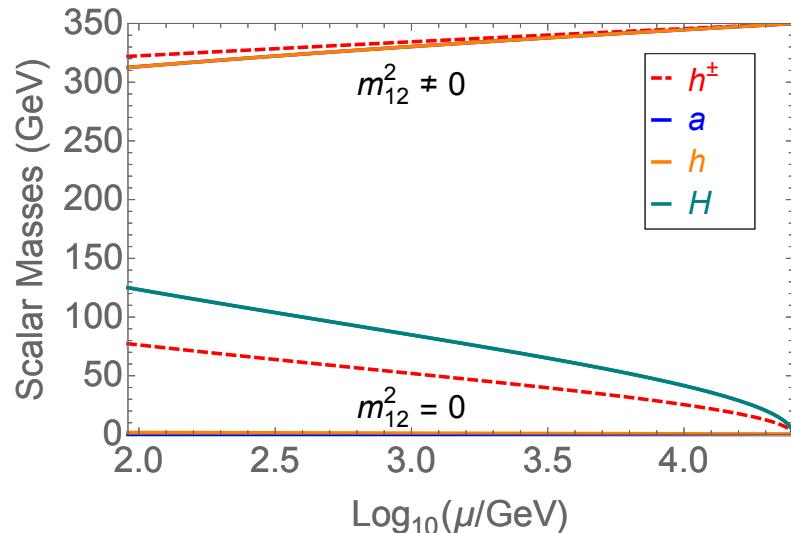
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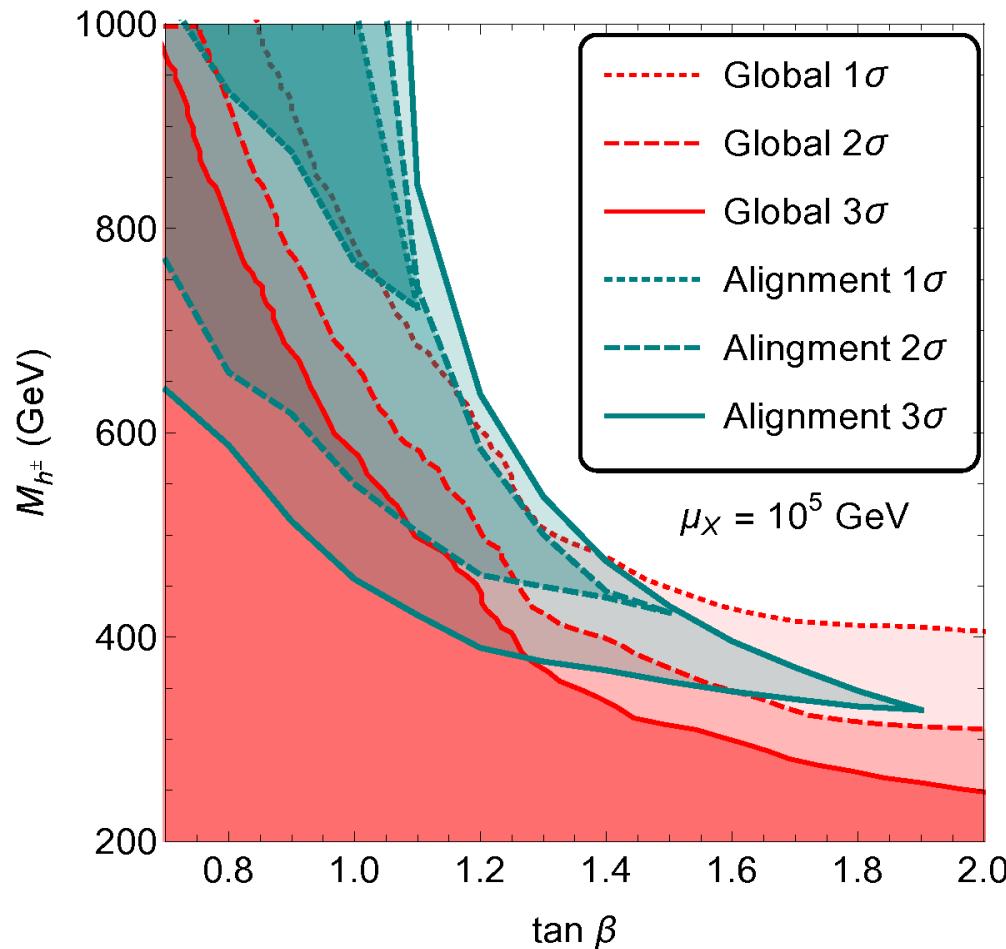
Breaking Effects: $-m_{12}^2 \phi_1^\dagger \phi_2$, U(1)_Y coupling g' , Yukawa couplings $h_{t,b}$.



Quasi-Degenerate Heavy Higgs Spectrum: $M_h^2 = M_a^2 + \lambda_5 v^2 \simeq M_{h^\pm}^2$

- **Constraints on** $(\tan \beta, M_{h^\pm})$

[P.S.B. Dev and A.P., JHEP1412 (2014) 024.]

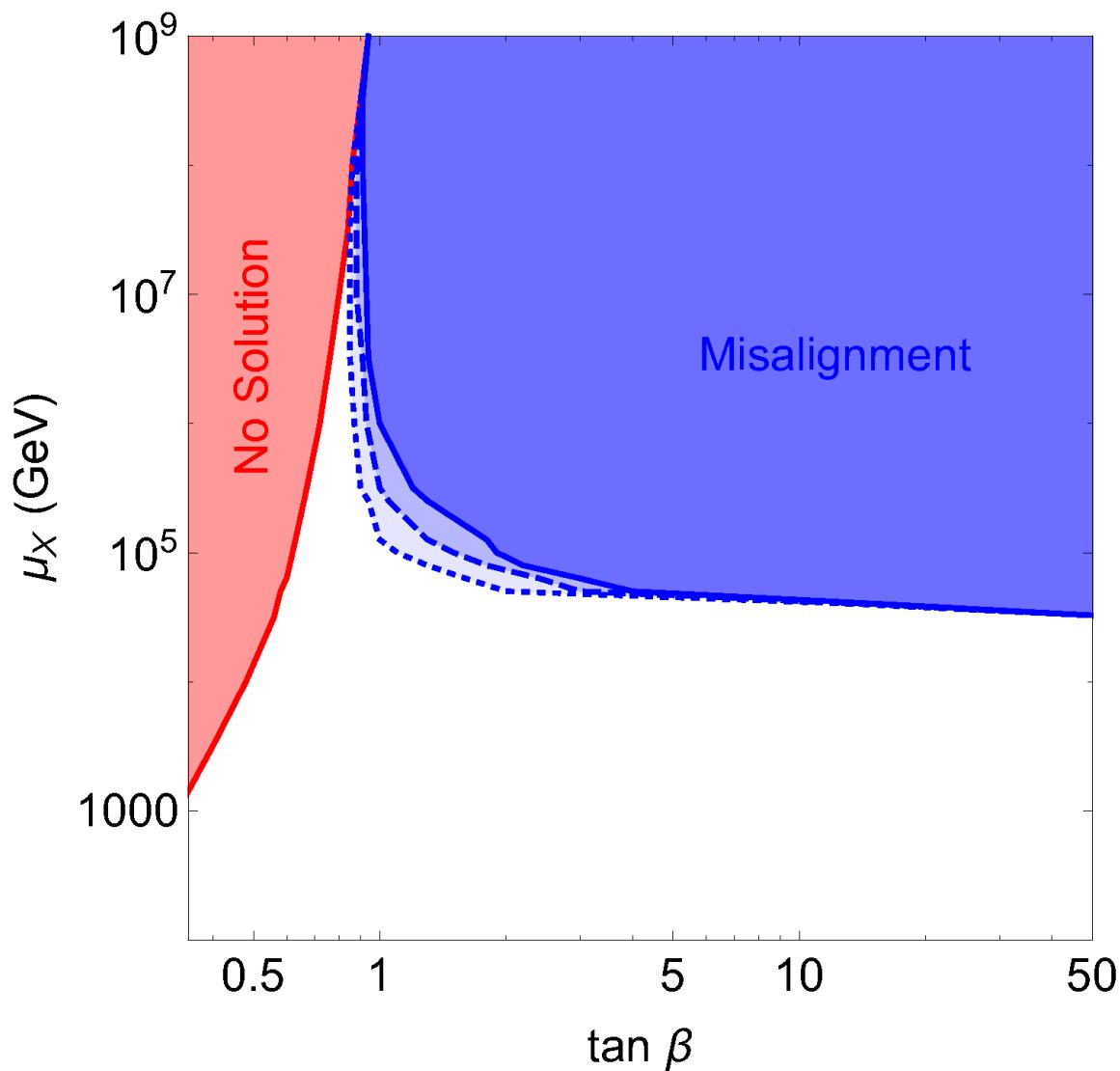


Constraints due to low-energy and LHC Higgs precision data.

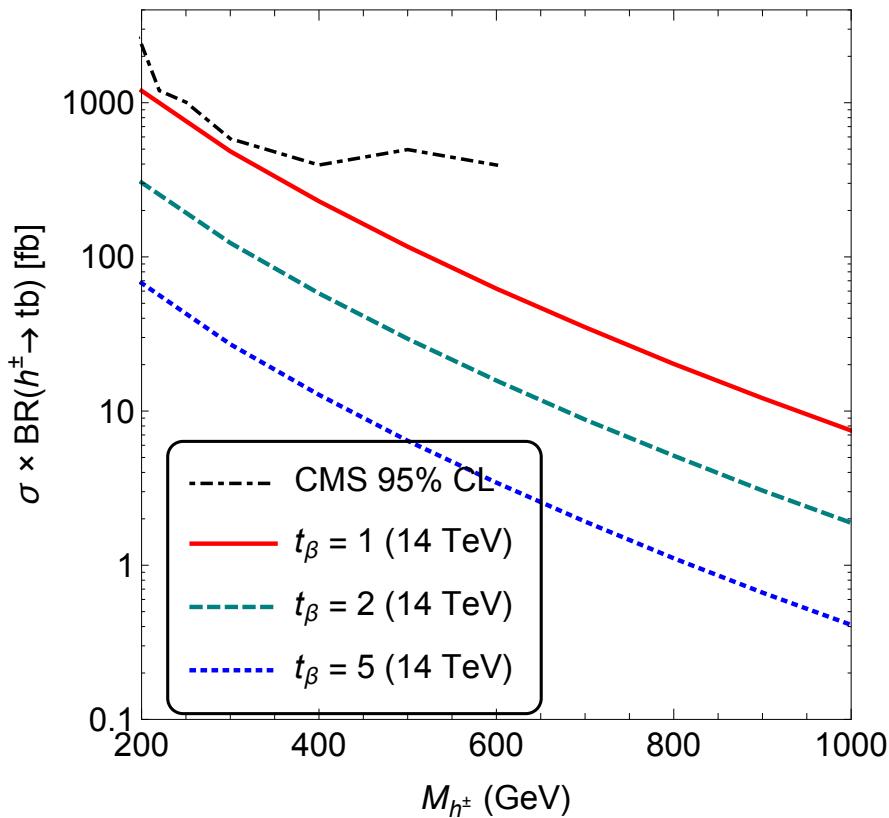
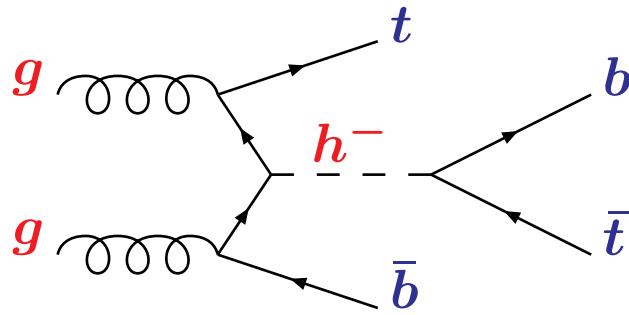
[See, e.g., J. Baglio, O. Eberhardt, U. Nierste, M. Wiebusch, PRD90 (2014) 015008;
K. Cheung, J.S. Lee, P.-Y. Tseng, JHEP1401 (2014) 085.]

- Misalignment Predictions

[P.S.B. Dev and A.P., JHEP1412 (2014) 024.]



- $gg \rightarrow t\bar{b}h^- \rightarrow t\bar{b}\bar{t}b$

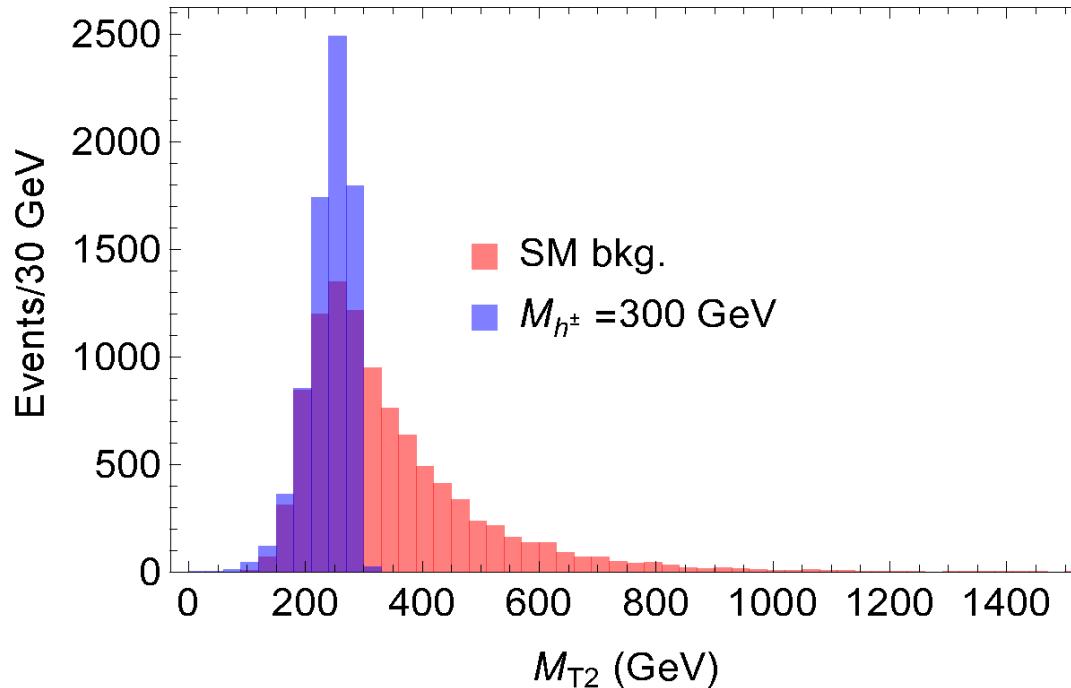


[CMS Collaboration, CMS-PAS-HIG-13-026;
P.S.B. Dev and A.P., JHEP1412 (2014) 024.]

$$\begin{aligned}
 p_T^\ell &> 20 \text{ GeV}, \\
 |\eta^\ell| &< 2.5, \\
 \Delta R^{\ell\ell} &> 0.4, \\
 M_{\ell\ell} &> 12 \text{ GeV}, \\
 |M_{\ell\ell} - M_Z| &> 10 \text{ GeV}, \\
 p_T^j &> 30 \text{ GeV}, \\
 |\eta^j| &< 2.4, \\
 \cancel{E}_T &> 40 \text{ GeV}.
 \end{aligned}$$

- M_{T2} Variable for reconstructing M_{h^\pm}

[C. G. Lester and D. J. Summers, Phys. Lett. B463 (1999) 99.]



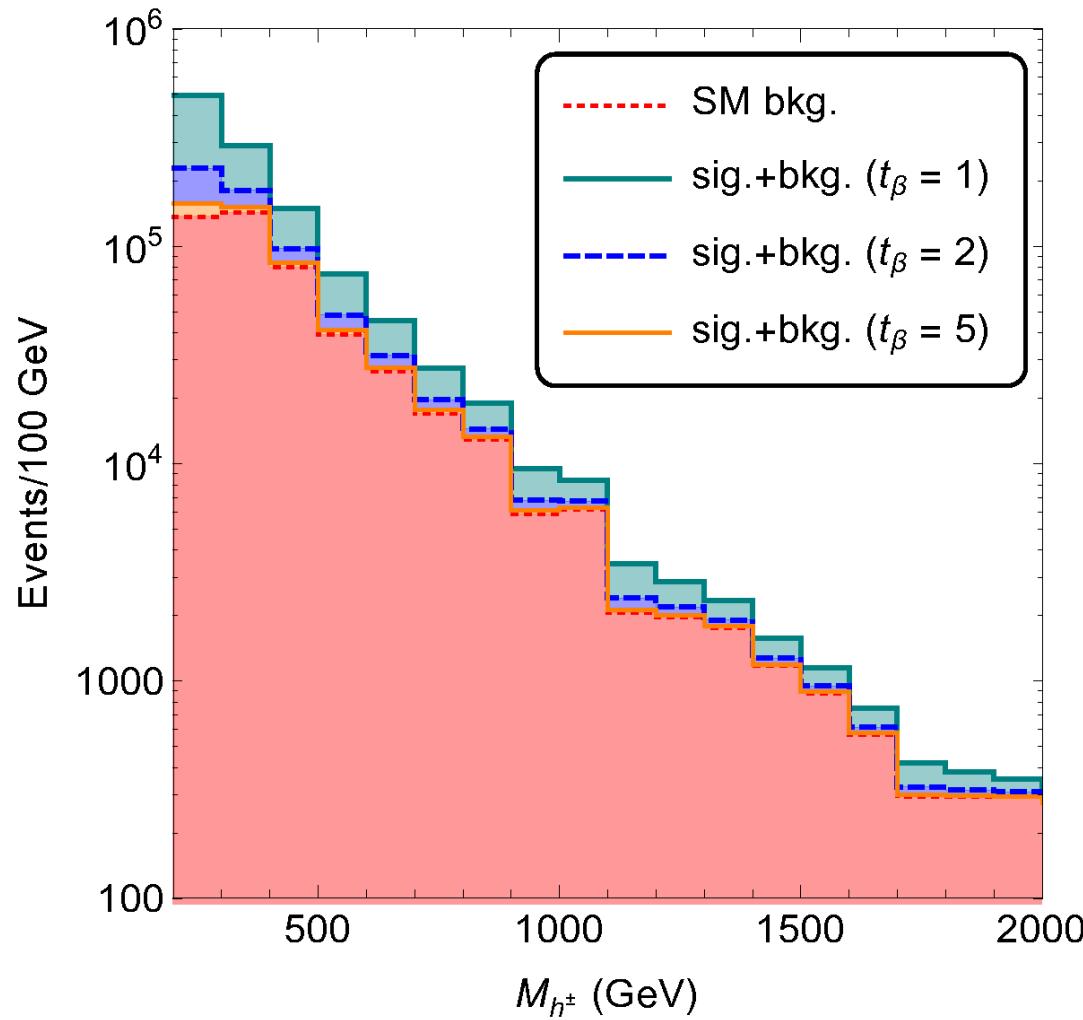
Decay chain: $gg \rightarrow h^\pm t b \rightarrow (\ell \nu_\ell bb) (\ell' \nu_{\ell'} b) b$

$$M_{T2} = \min_{\{p_{T_a} + p_{T_b} = p_T\}} \left[\max \{m_{T_a}, m_{T_b}\} \right] \lesssim M_{h^\pm}$$

$m_{T_{a,b}}$: usual transverse masses for $\{a\} = (\ell \nu_\ell bb)$ and $\{b\} = (\ell' \nu_{\ell'} bb)$.

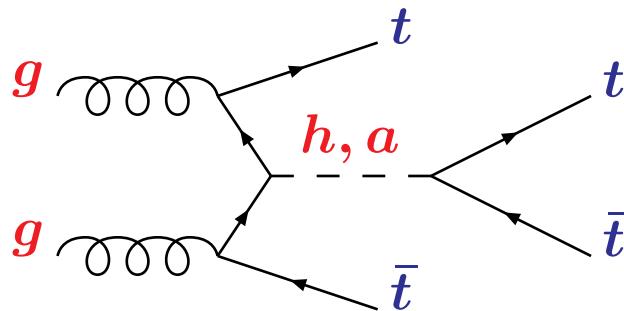
- Predicted number of events for h^\pm signal at $\sqrt{s} = 14$ TeV with 300 fb^{-1} integrated luminosity

[P.S.B. Dev and A.P., JHEP1412 (2014) 024.]

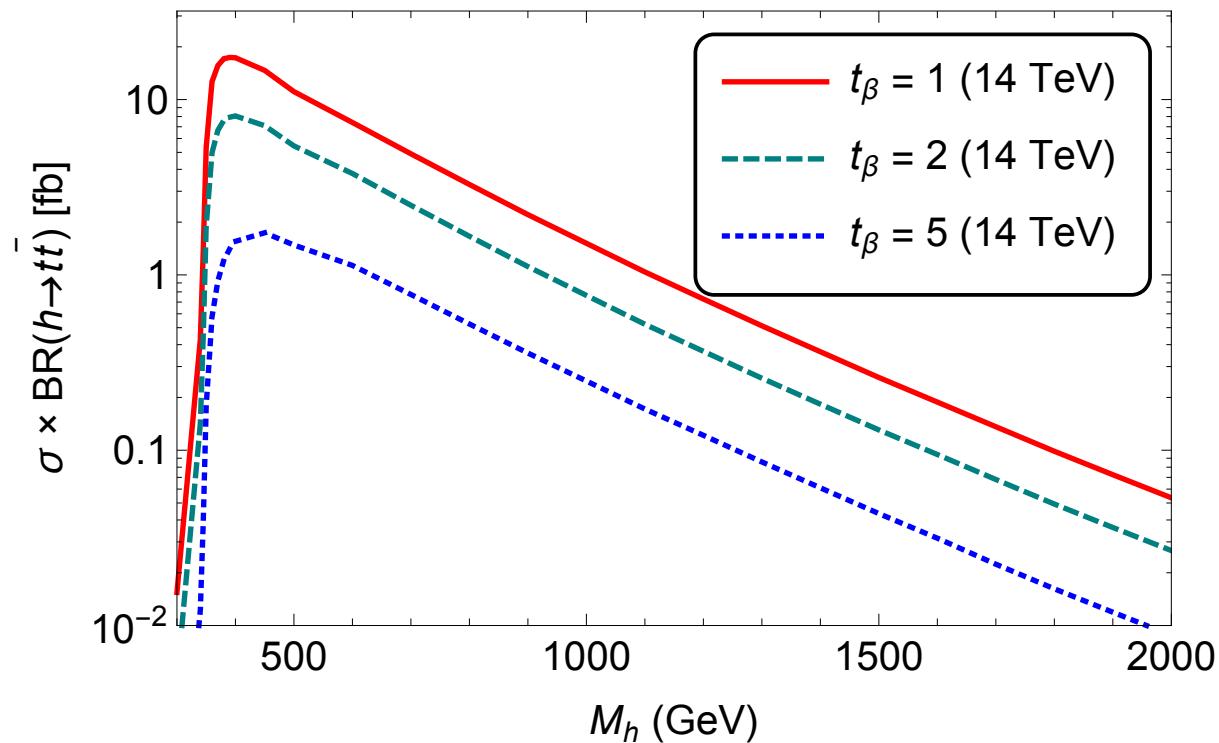


using MadGraph5.aMC@NLO.

- $gg \rightarrow t\bar{t}(h, a) \rightarrow t\bar{t}t\bar{t}$

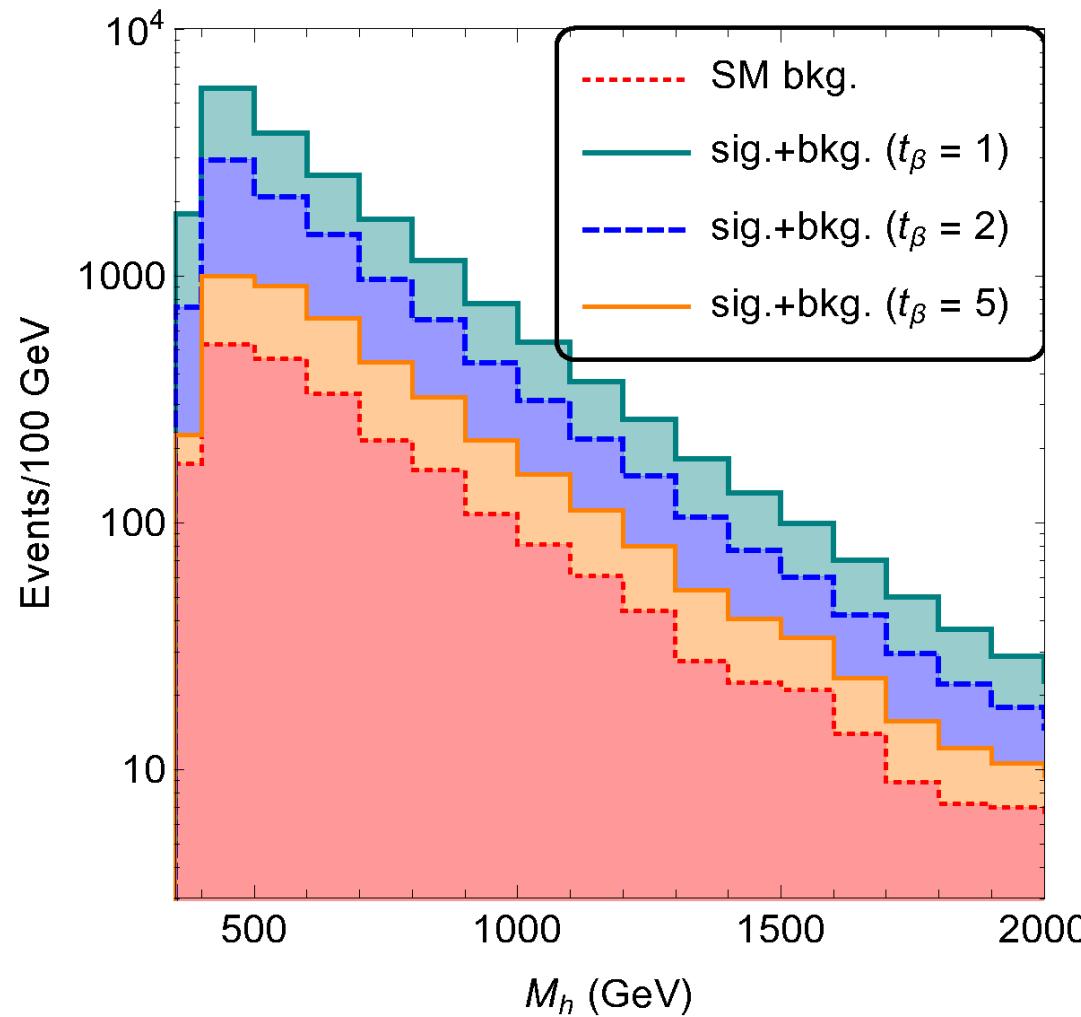


[P.S.B. Dev and A.P., JHEP1412 (2014) 024.]



- Predicted number of events for h, a signal at $\sqrt{s} = 14$ TeV with 300 fb^{-1} integrated luminosity

[P.S.B. Dev and A.P., JHEP1412 (2014) 024.]



using MadGraph5.aMC@NLO.

• Conclusions

- The 2HDM potential may exhibit up to 13 distinct discrete and global symmetries as derived in the SO(1,5) bilinear formalism.

- Maximal Symmetry Group for n HDM Potentials:

$$G_\Phi = \mathrm{SU}(2)_L \otimes \mathrm{Sp}(2n)/\mathbf{Z}_2.$$

- 3 syms. lead to Natural Alignment in 2HDM ($\lambda_1 = \lambda_2 = \lambda_{345}/2$):

(i) SO(5); (ii) O(3) \times O(2); (iii) Z₂ \times O(2) \times O(2).

- The Aligned Heavy Higgs Sector is Gaugophobic.

- Two-loop RG Effects give rise to misalignment predictions leading to lower limits on the charged Higgs mass M_{h^\pm} .

- Probing New Aligned Higgs Doublets via the production channels:

(i) $gg \rightarrow t\bar{t}(h, a) \rightarrow t\bar{t}t\bar{t}$; (ii) $gg \rightarrow t\bar{b}h^- \rightarrow t\bar{b}\bar{t}b$.