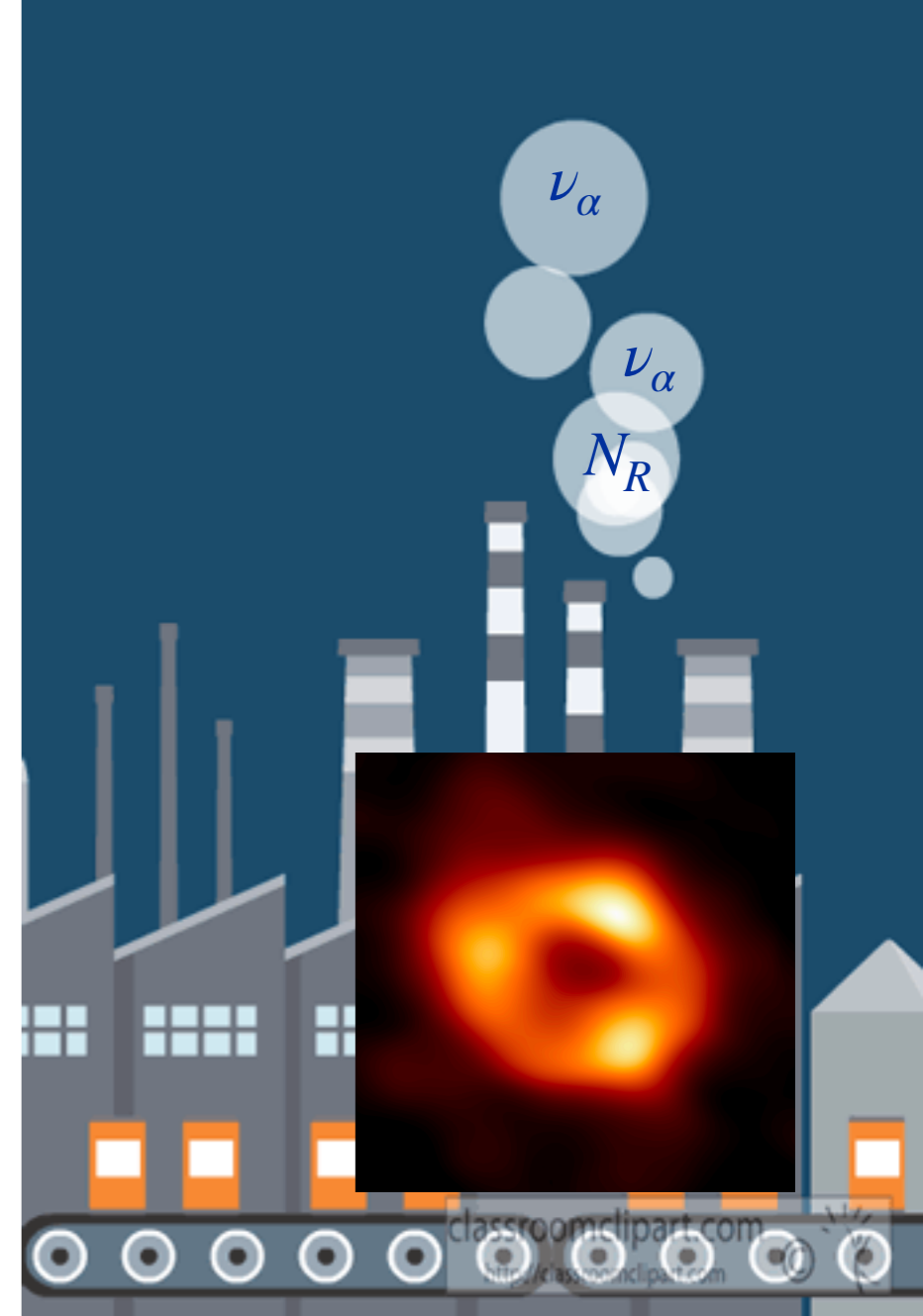


Primordial Black Holes as Particle Factories and Implications for the Early Universe

Yuber F. Perez-Gonzalez

Particle and Astroparticle Theory Seminar
Max-Planck-Institute für Kernphysik
June 26th, 2023



EHT Collaboration



What's the impact of having evaporating Primordial Black Holes in the Early Universe?

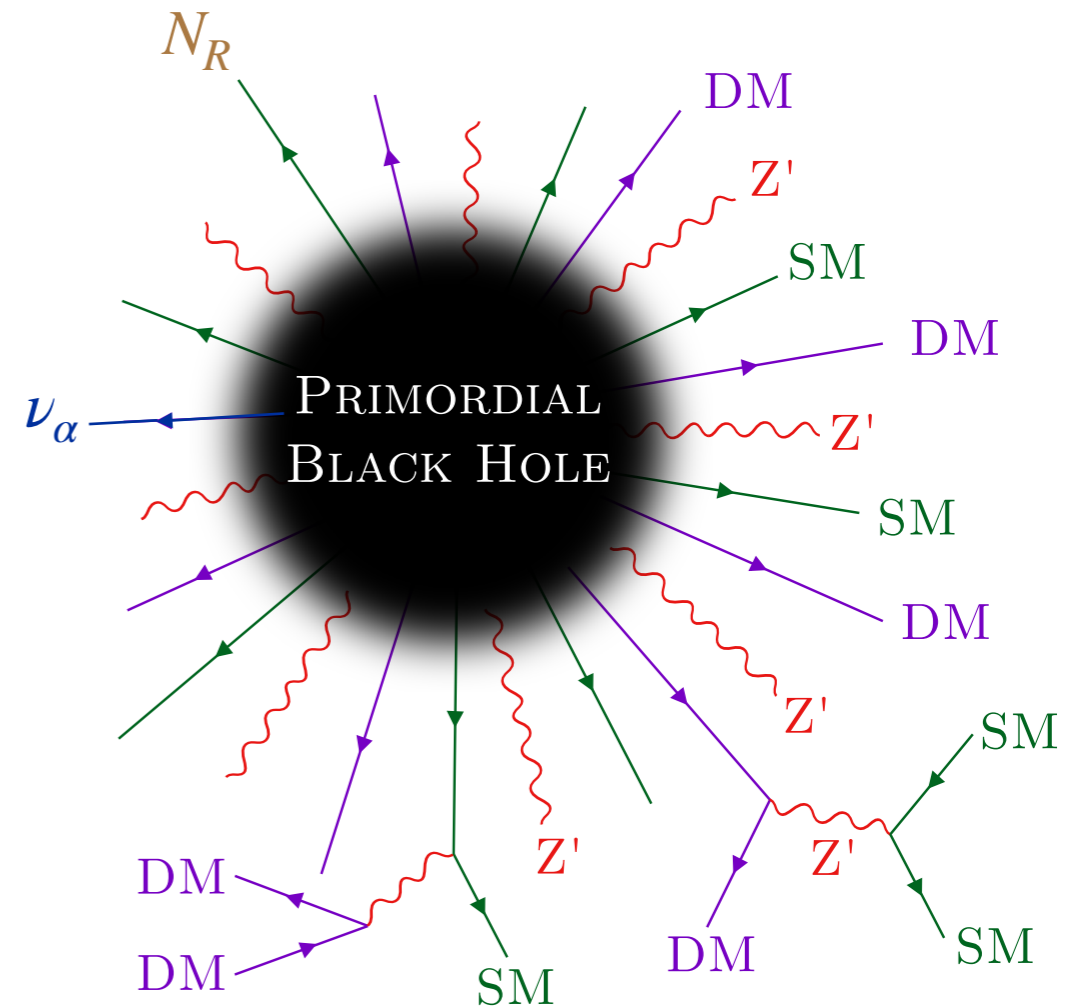
❖ *Introduction:*

- PBHs — Formation/Evaporation

❖ *First Act: Neutrinos and Leptogenesis*

❖ *Second Act: Dark Matter*

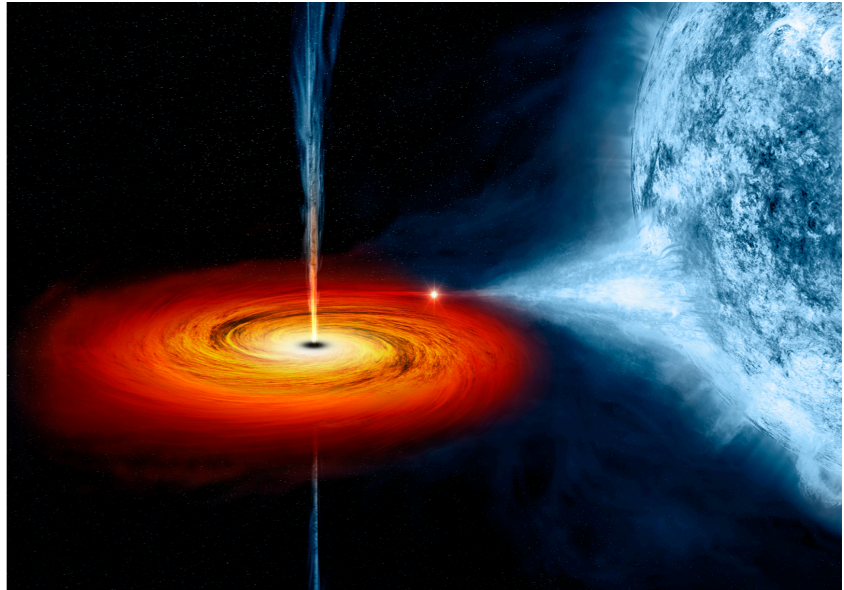
❖ *Third Act: High energy Neutrinos and evaporating PBHs*



Primordial Black Holes (PBH)

Astrophysical Black Holes

$$M \gtrsim 3M_{\odot}$$



Formation

Lighter Black Holes

$$r_S = 2GM$$

Large densities

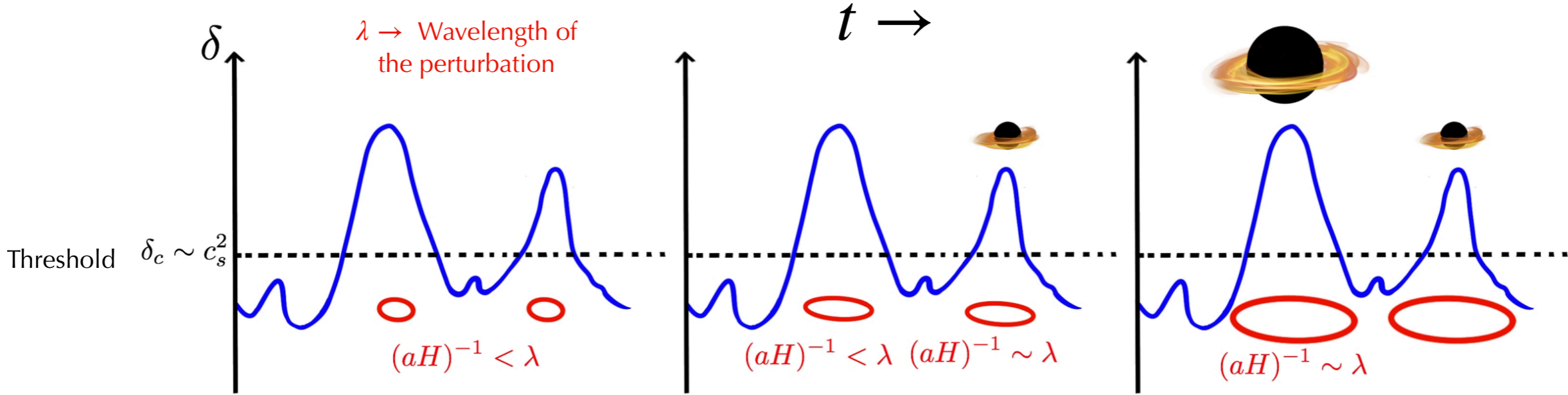
$$M_i \sim \frac{t}{G} \sim 10^{15} \left(\frac{t}{10^{-23} \text{ s}} \right) \text{ g}$$

Formation

- ❖ Bubble collisions
- ❖ Pressure reduction
- ❖ Collapse of density fluctuations

$$M_{\text{BH},i} \propto \underbrace{\frac{4\pi}{3} \rho \frac{1}{H^3}}_{\text{Mass contained in the horizon}}$$

$$\delta = \frac{\rho - \bar{\rho}}{\bar{\rho}}$$



Taken from Villanueva-Domingo, Mena, Palomares-Ruiz 2103.12087

Fraction of the total energy in PBH

$$\beta = \frac{\rho_{\text{PBH}}}{\rho_{\text{tot}}}$$

Assume a monochromatic mass distribution

All PBHs with the same mass

PBH form when the density perturbation enters the Hubble horizon

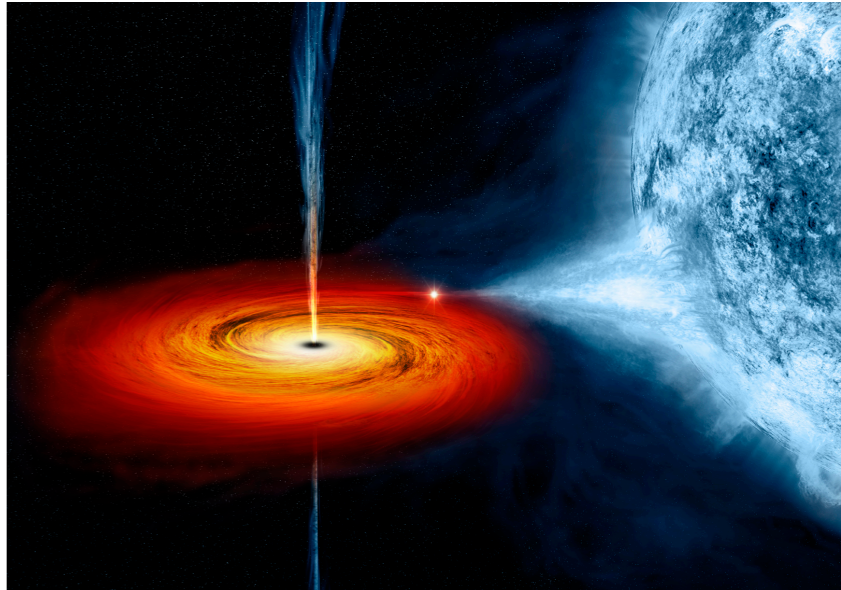
$$M_{\text{BH},i}, \beta$$

Carr et al. 2002.12778

Primordial Black Holes (PBH)

Astrophysical Black Holes

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Formation

Lighter Black Holes

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Large densities

$$M_i \sim \frac{t}{G} \sim 10^{15} \left(\frac{t}{10^{-23} \text{ s}} \right) \text{ g}$$

Assume a monochromatic mass distribution

$$M_{\text{BH},i}, \beta$$

Evaporation

Quantum effects are important

$$r_S \sim \lambda_C$$

$$r_S \sim 0.015 \text{ fm} \left(\frac{M}{10^{13} \text{ g}} \right)$$

Black Holes evaporate by thermal emission

Hawking, 1975

All degrees of freedom!

Evaporation — Schwarzschild BHs

Described by M_{BH}

Hawking
Instantaneous
Spectrum

$$\frac{d^2 N_i}{d\omega dt} = \frac{g_i}{2\pi^2} \frac{s_i \Gamma(M, \omega, \mu_i)}{\exp[\omega/T] - (-1)^{2s_i}}$$

Absorption
probability

$$s_i \Gamma(M, \omega, \mu_i) = \frac{1}{\pi} \sigma_{s_i}(M, \omega, \mu_i) \omega^2$$

BH Temperature

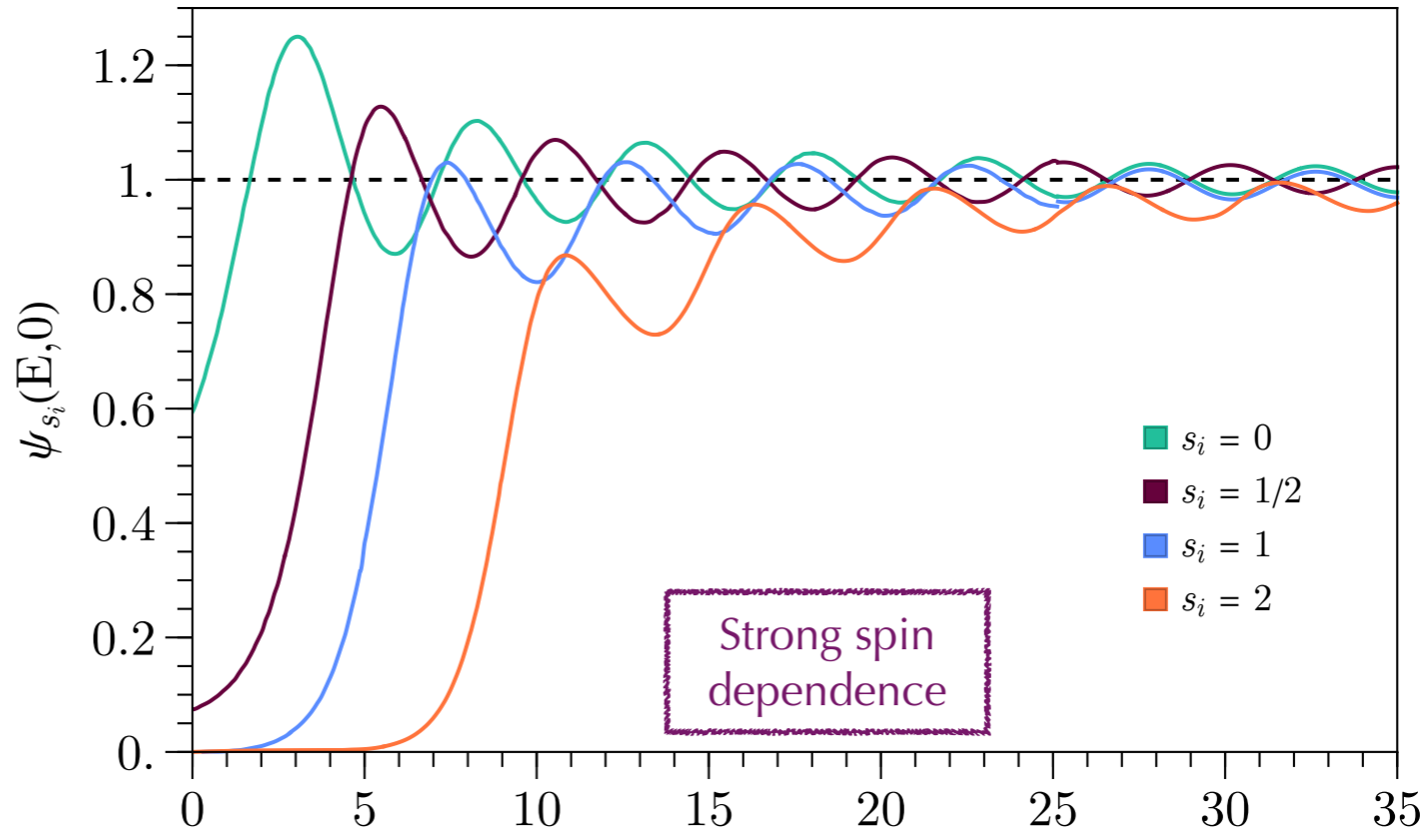
$$T = \frac{\hbar c^3}{8\pi G M k}$$

$$\sim 1 \text{ GeV} \left(\frac{10^{13} \text{ g}}{M} \right)$$



*Hic depositum est, quod mortale fuit Isaaci Newtoni

Reduced Absorption Cross Section



$$\psi_{s_i}(\omega) \equiv \frac{\sigma_{s_i}(\omega)}{27\pi G^2 M_{\text{BH}}^2}$$

ω/T

Only $\ell \geq s_i$ modes

Evaporation — Kerr BHs

Described by $M_{\text{BH}}, a_{\star} = JM_p^2/M_{\text{BH}}^2 \in [0,1)$

Axisymmetric case

$$\frac{d^2 N_i}{d\omega dt} = \frac{g_i}{2\pi} \sum_{l=s_i} \sum_{m=-l}^l \underbrace{\frac{s_i \Gamma_{lm}}{\exp(\varpi_m/T_{\text{BH}}) - (-1)^{2s_i}}}_{\frac{d^2 \mathcal{N}_{ilm}}{d\omega dt}}$$

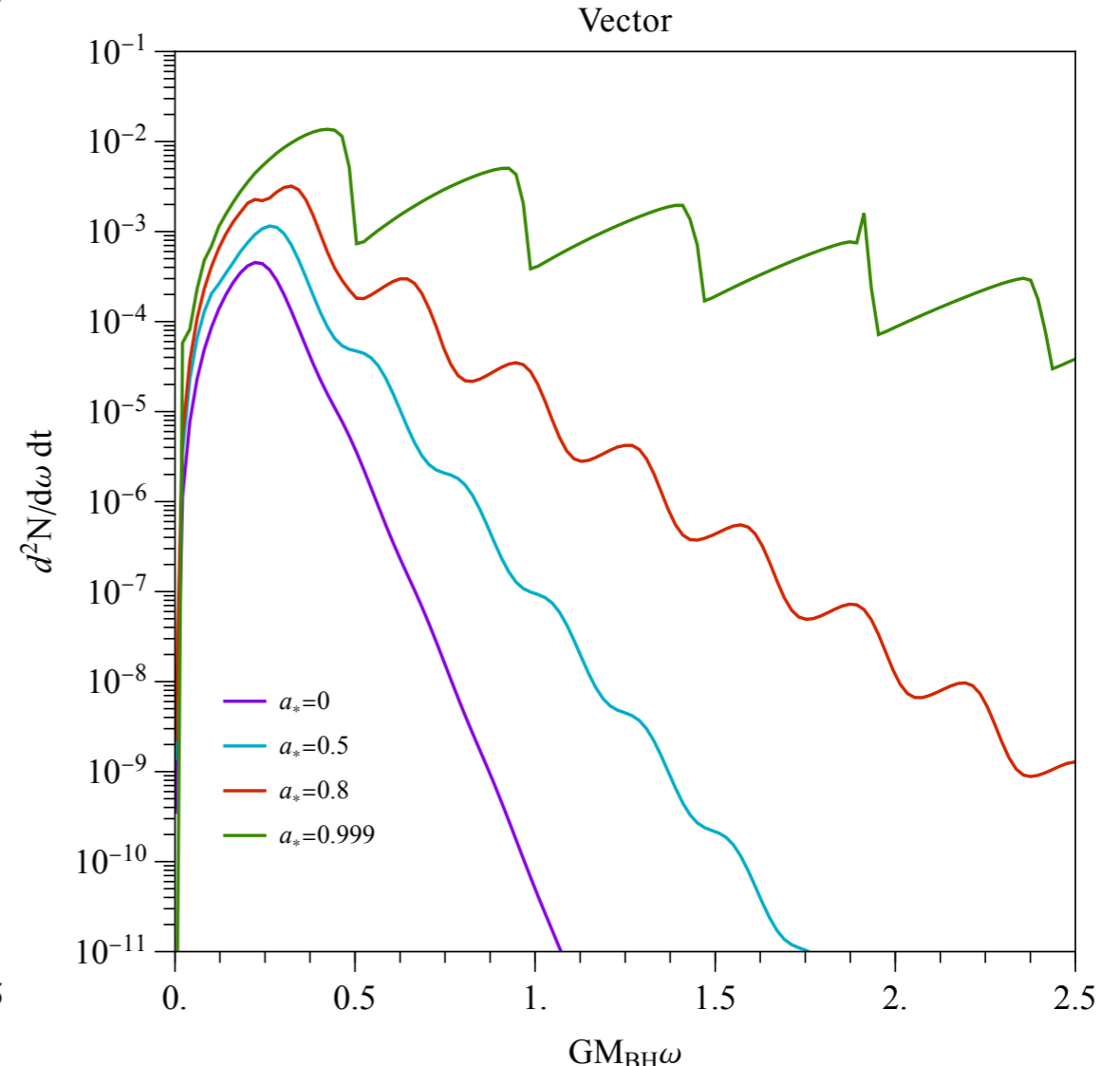
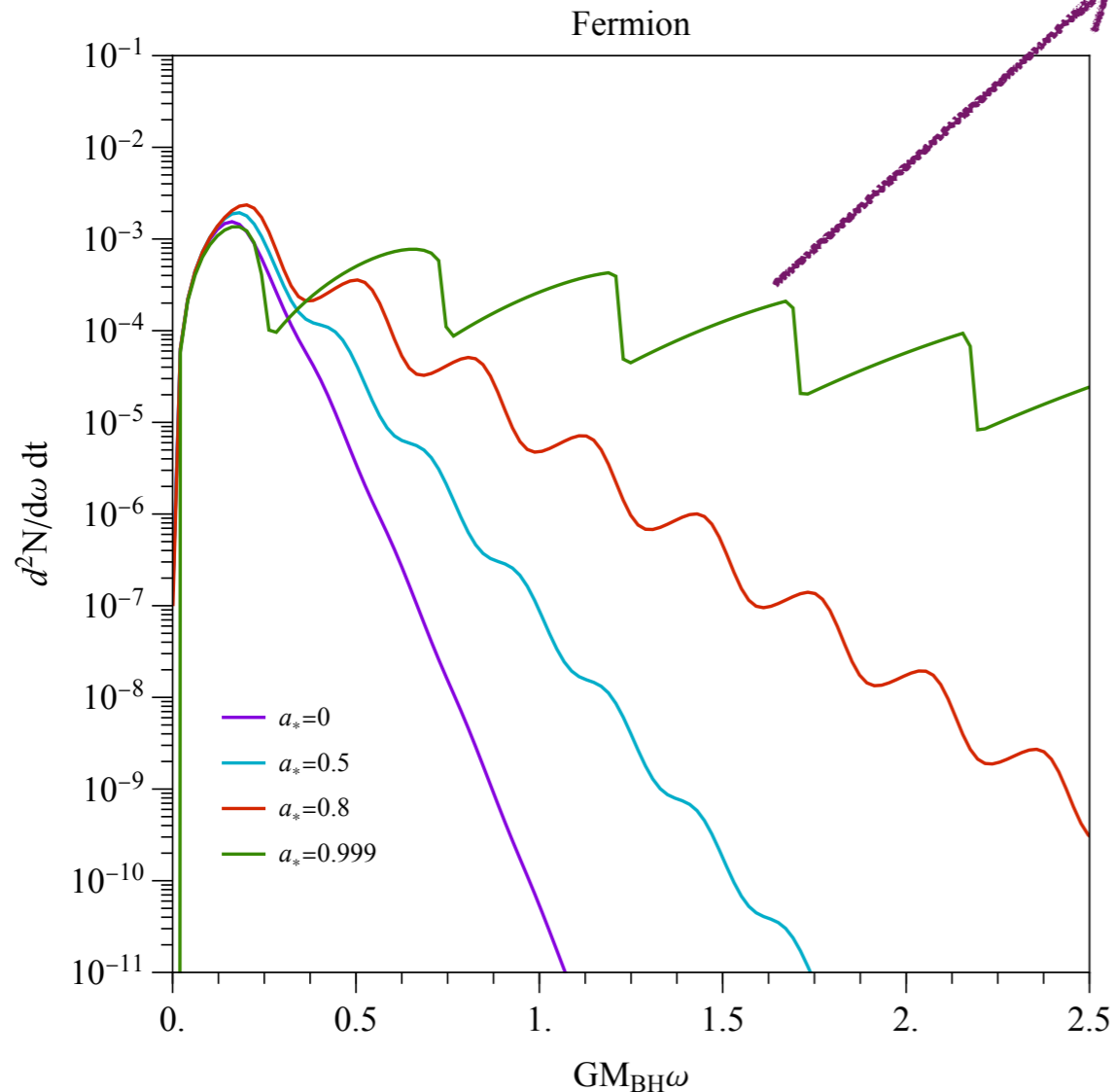
$$T = \frac{1}{4\pi GM} \frac{\sqrt{1-a_{\star}^2}}{1+\sqrt{1-a_{\star}^2}}$$

BH Temperature

$$\varpi_m = \omega - m\Omega$$

Explicit dependence on m
 $\Omega \rightarrow$ angular velocity

Enhanced emission of
 $l = m$ modes



Evaporation

$$\tilde{M}_{\text{BH}}^{\text{max}}(\xi) \equiv M_{\text{BH}}(\xi)/M_{\text{BH}}^{\text{in}}$$

$$a_{\star, \text{in}} \rightarrow 1$$

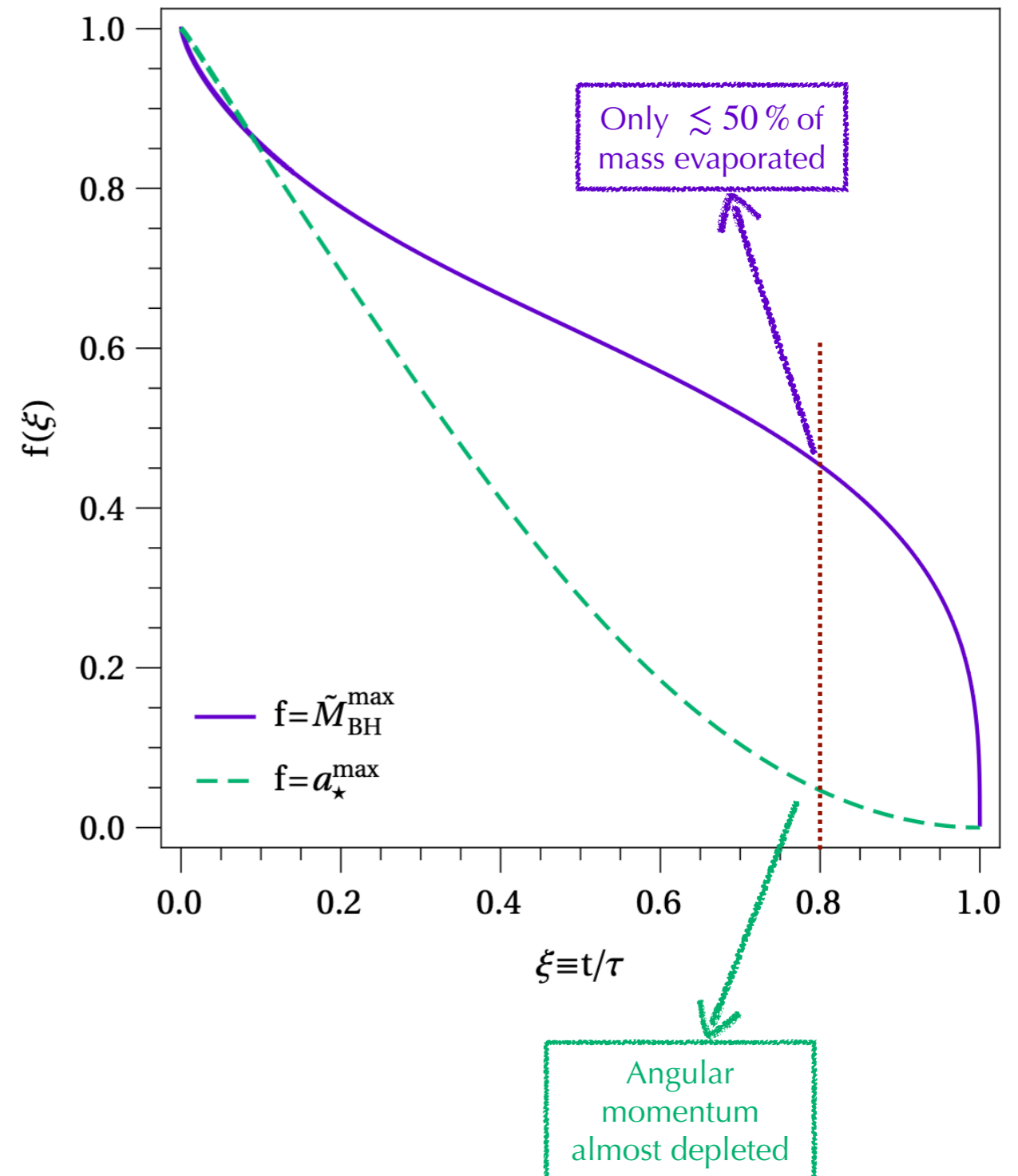
BHs lose mass over time

$$\frac{dM_{\text{BH}}}{dt} = - \underbrace{\varepsilon(M_{\text{BH}}, a_{\star})}_{\text{Evaporation function}} \frac{M_P^4}{M_{\text{BH}}^2}$$

$$\frac{da_{\star}}{dt} = - a_{\star} [\gamma(M_{\text{BH}}, a_{\star}) - 2\varepsilon(M_{\text{BH}}, a_{\star})] \frac{M_P^4}{M_{\text{BH}}^3}$$

Angular momentum depleted faster than mass

If there are some PBH still around they *might* have a small angular momentum



Evaporation

BHs lose mass over time

$$\frac{dM_{\text{BH}}}{dt} = - \varepsilon(M_{\text{BH}}, a_{\star}) \frac{M_{\text{P}}^4}{M_{\text{BH}}^2}$$

$$\frac{da_{\star}}{dt} = - a_{\star} [\gamma(M_{\text{BH}}, a_{\star}) - 2\varepsilon(M_{\text{BH}}, a_{\star})] \frac{M_{\text{P}}^4}{M_{\text{BH}}^3}$$

Depends on the set of *all* existing dofs

$$\varepsilon = \sum_{i=\text{all}} \frac{g_i}{2\pi^2} \int_0^{\infty} \sum_{l=s_i} \sum_{m=-l}^l \frac{d^2 \mathcal{N}_{ilm}}{d\omega dt} \omega d\omega$$

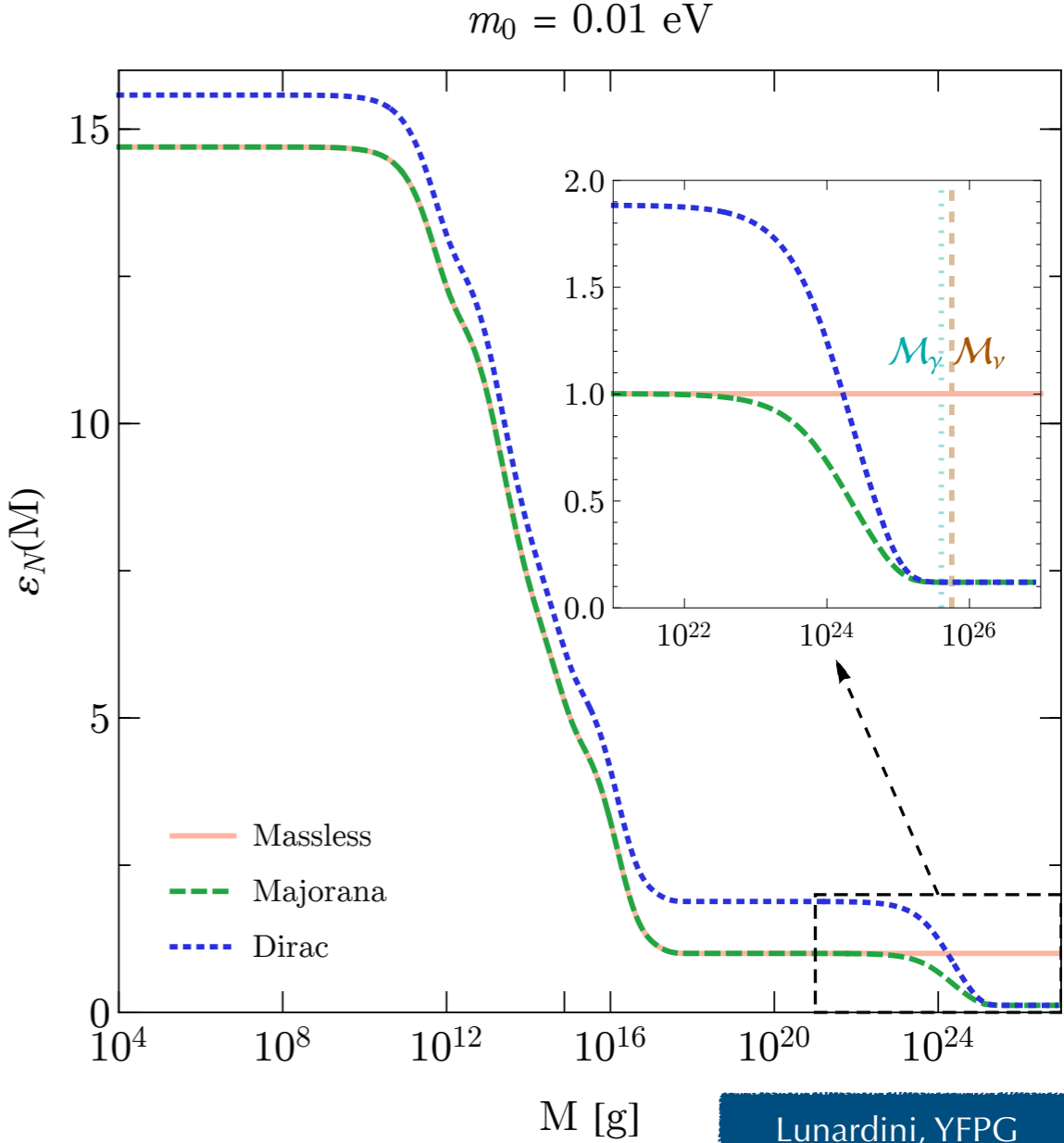
$$\gamma = \sum_{i=\text{all}} \frac{g_i}{2\pi^2} \int_0^{\infty} \sum_{l=s_i} \sum_{m=-l}^l \frac{d^2 \mathcal{N}_{ilm}}{d\omega dt} m d\omega$$

J. MacGibbon, 1991

BH lifetime for a Schwarzschild BH

$$\tau = \frac{1}{M_{\text{P}}^4} \int_{M_{\text{P}}}^{M_i} \frac{M^2 dM}{\varepsilon(M)}$$

Planck mass



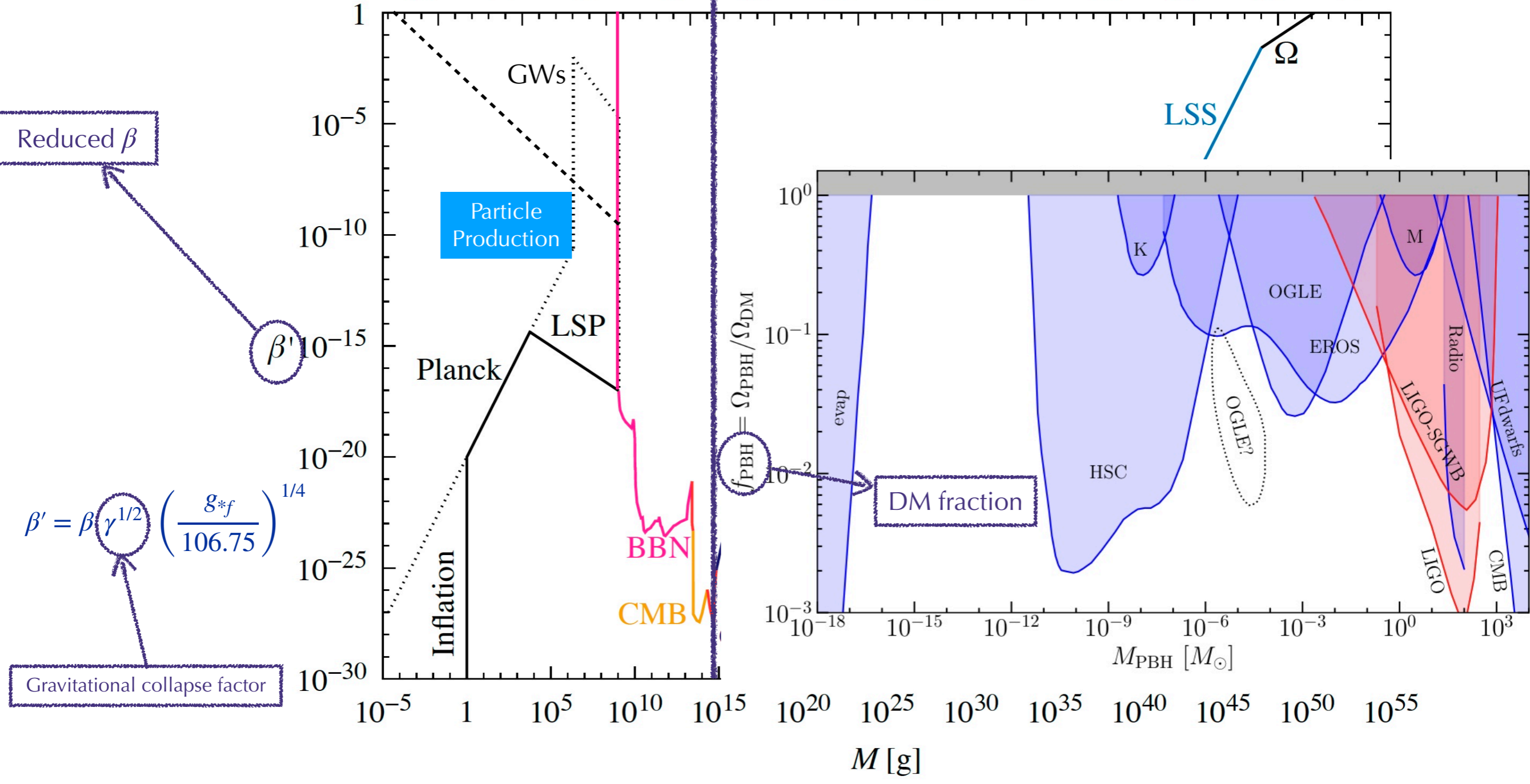
Lunardini, YFPG JCAP08(2020)014

BH with $M_i = M_{\odot}$:

$$\tau \sim 10^{67} \text{ years}$$

$$\sim 10^{57} * \text{age of the Universe}$$

Evaporated ← M/M_\odot → (Part of) Dark Matter?



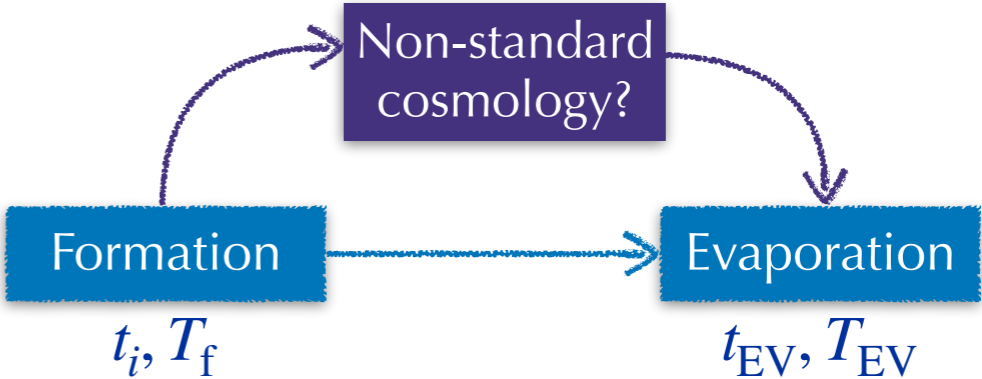
We focus on the region $M_{\text{BH},i} \leq 10^9 \text{ g}$

Carr et al. 2002.12778
Domènech et al. 2012.08151

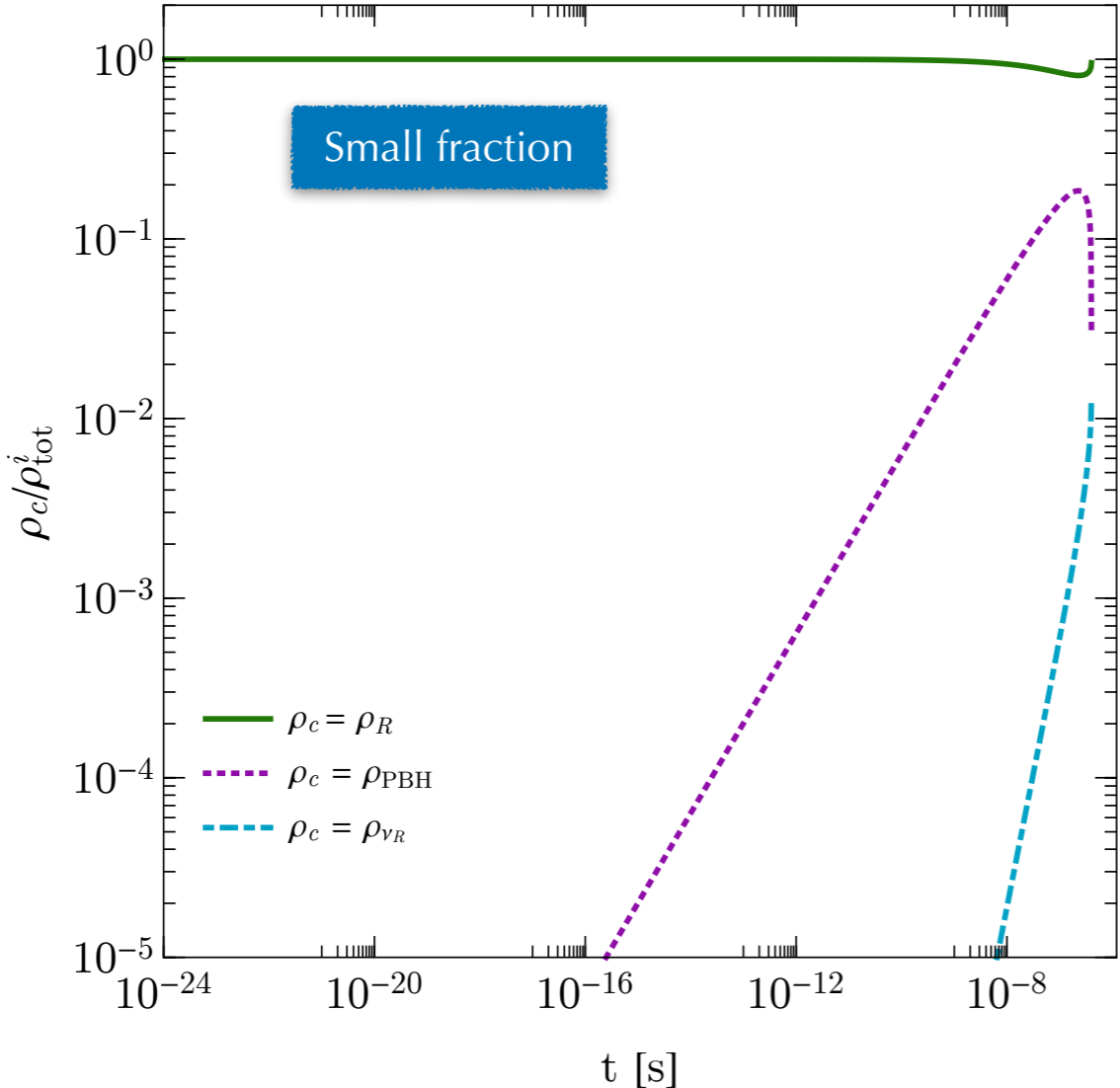
B. Kavanagh
10.5281/zenodo.3538999

PBHs: Rise and Fall

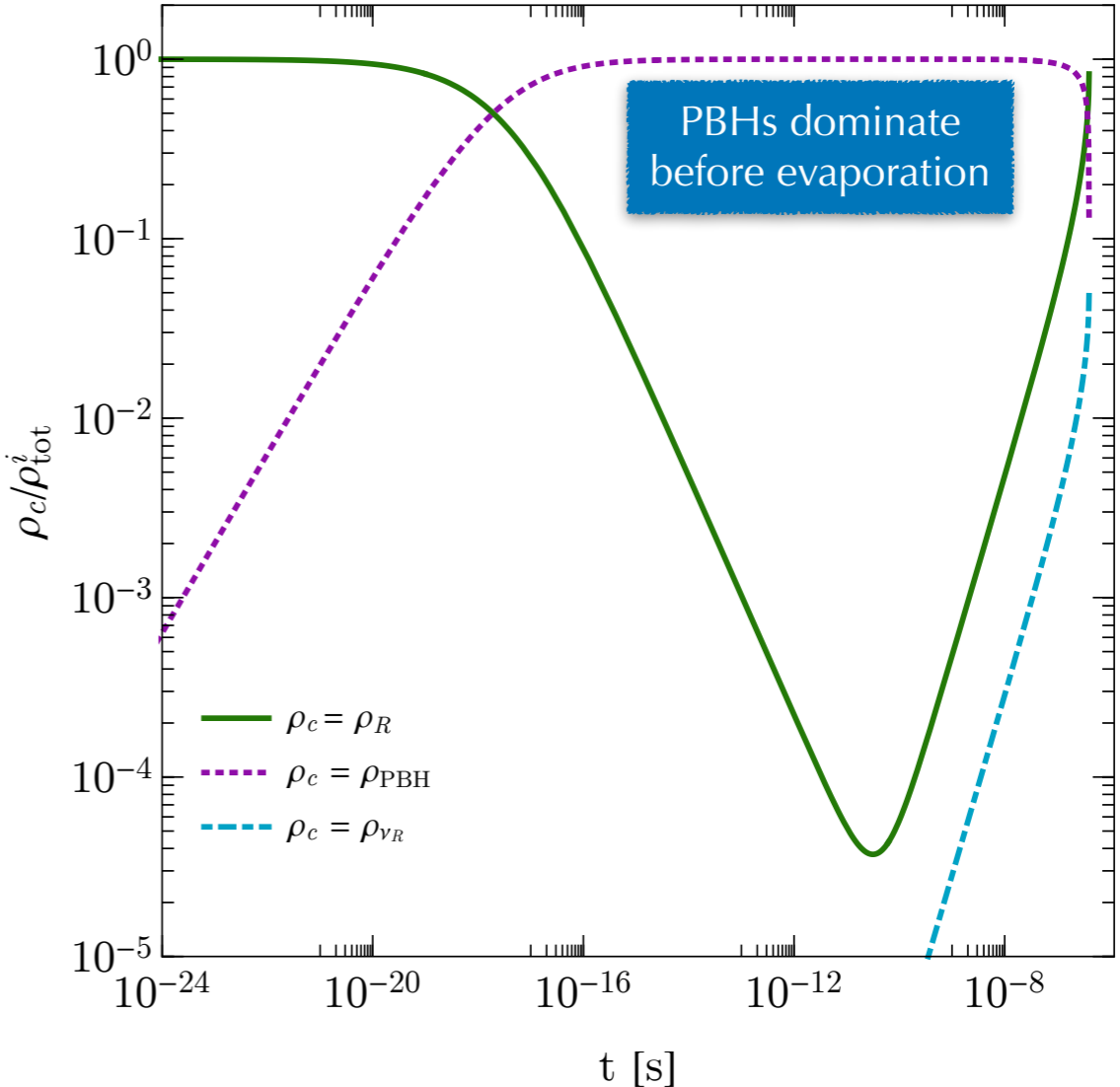
PBHs behave (almost) as matter



$$M_{\text{BH},i} = 10^7 \text{ g}, \beta = 10^{-13}$$



$$M_{\text{BH},i} = 10^7 \text{ g}, \beta = 10^{-7}$$



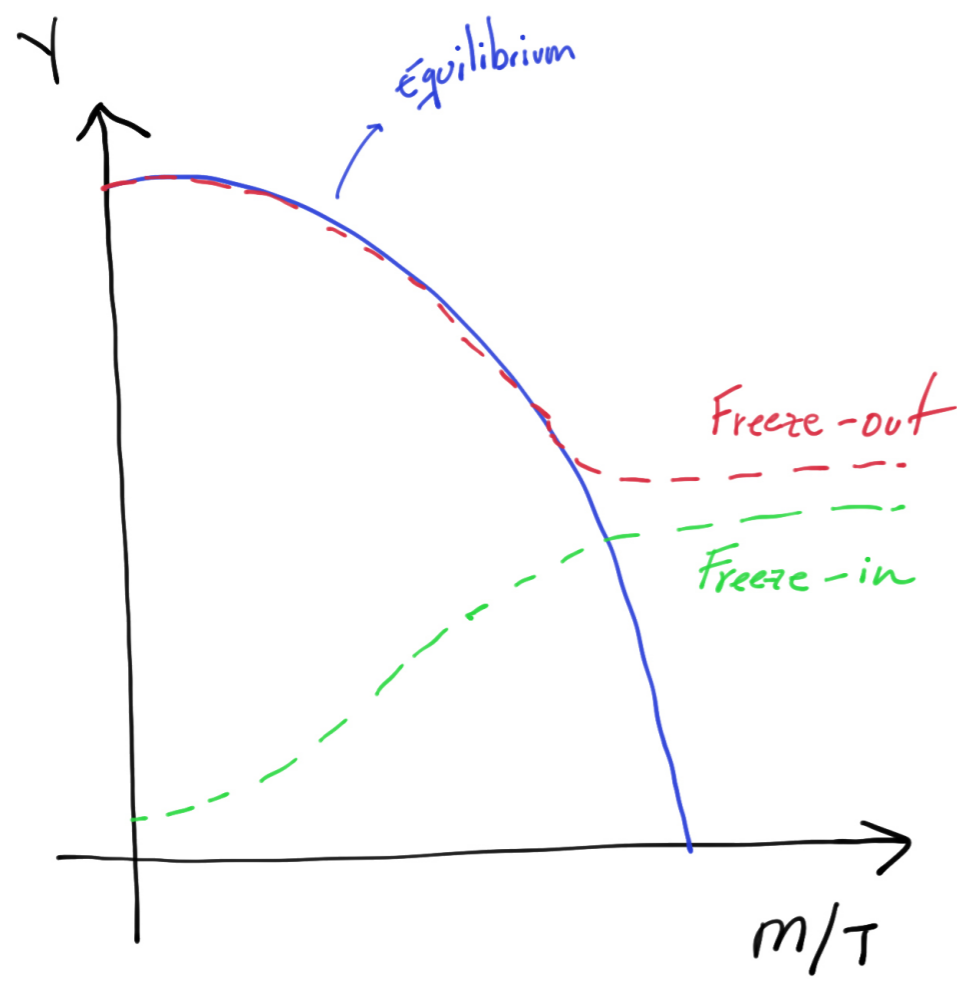
Initial fraction

$$\beta \gtrsim 2.5 \times 10^{-14} \left(\frac{g_*(T_f)}{106.75} \right)^{-\frac{1}{4}} \left(\frac{M_i}{10^8 \text{ g}} \right)^{-1} \left(\frac{\epsilon_D(M_i)}{15.35} \right)^{\frac{1}{2}}$$

Hooper et al,
1905.01301

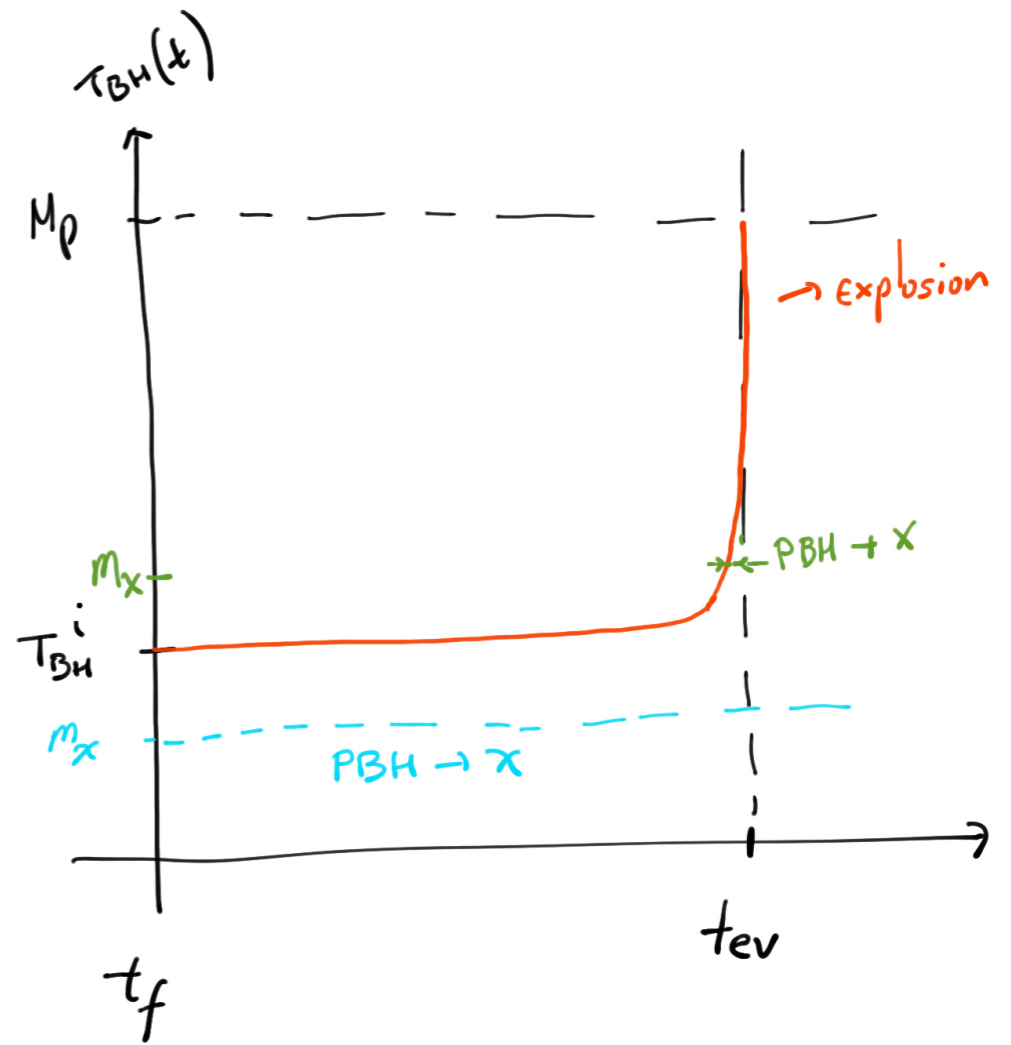
Why is interesting the particle production via evaporation?

(Thermal) Particle decoupling



Abundances depend on masses and interactions

Hawking Evaporation

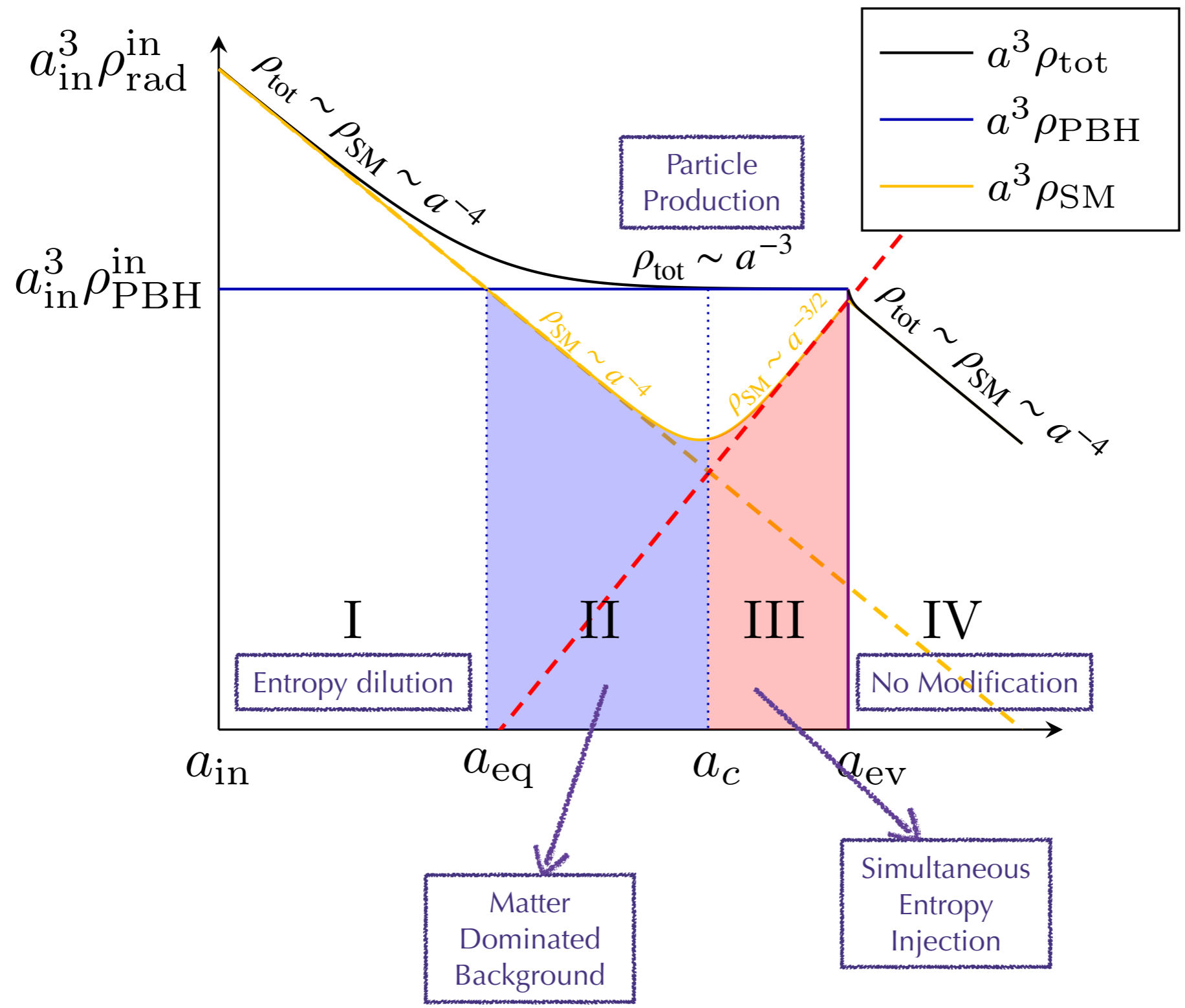


Abundances depend on masses and PBH parameters



- ❖ Modified Cosmology
- ❖ Particle production
- ❖ Entropy dilution

Why is interesting the particle production via evaporation?



Why is interesting the particle production via evaporation?

DM production
(Besides PBH-DM)

- ❖ Purely Gravitationally interacting DM
- ❖ Modify Freeze-In/Freeze-out mechanisms

Fujita et al, 1401.1909
 Morrison et al, 1812.10606
 Baldes et al, 2004.14773
 Masina, 2004.04740, 2103.13825
Cheek, Heurtier, YFPG, Turner
 2107.00013, 2107.00016,
 2212.03878



Baryon Asymmetry

- ❖ Modifying Baryogenesis scenarios
- ❖ **Leptogenesis scenarios**
- ❖ Producing a local asymmetry at PBH formation
- ❖ Connections with PBH-DM

Baumann, Steinhadt, Turok, 0703250
 Yamada and Iso, 1610.02586
 Fujita et al, 1401.1909
 Morrison et al, 1812.10606
 García-Bellido, Carr, Clesse, 1904.11482
 Hooper and Krnjaic, 2010.01134



Dark Radiation

- ❖ Production of hot gravitons
- ❖ Testable from future measurements on ΔN_{eff}

Hooper, Krnjaic, McDermott, 1905.01301
Lunardini, YFPG, 1910.07864
 Masina, 2004.04740, 2103.13825
Cheek, Heurtier, YFPG, Turner, 2207.09462

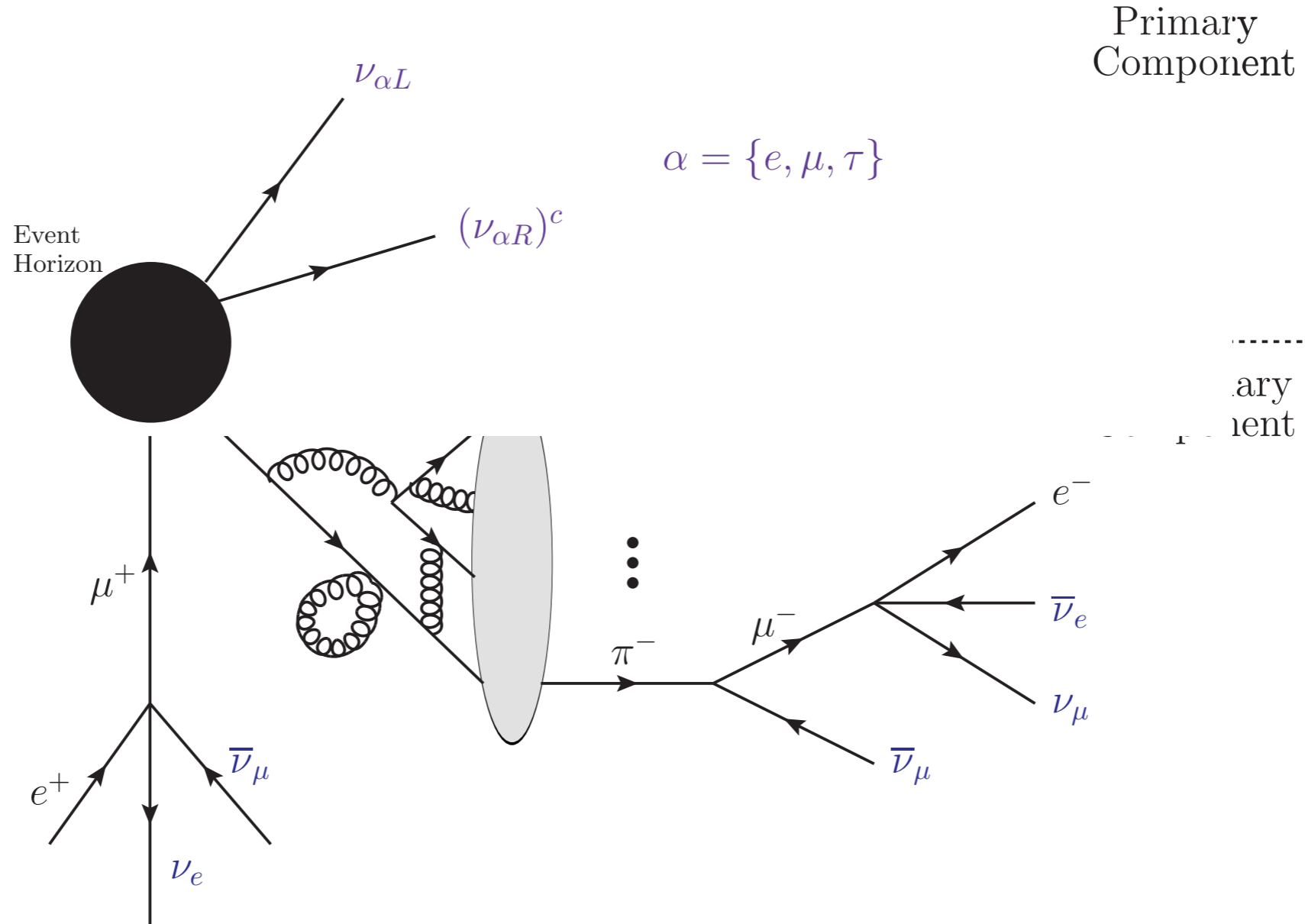
First Act

Neutrinos and Leptogenesis

Based on:
YFPG and Turner: 2010.03565
Bernal, Fong, YFPG, Turner 2203.08823

What about neutrino emission?

In the SM...



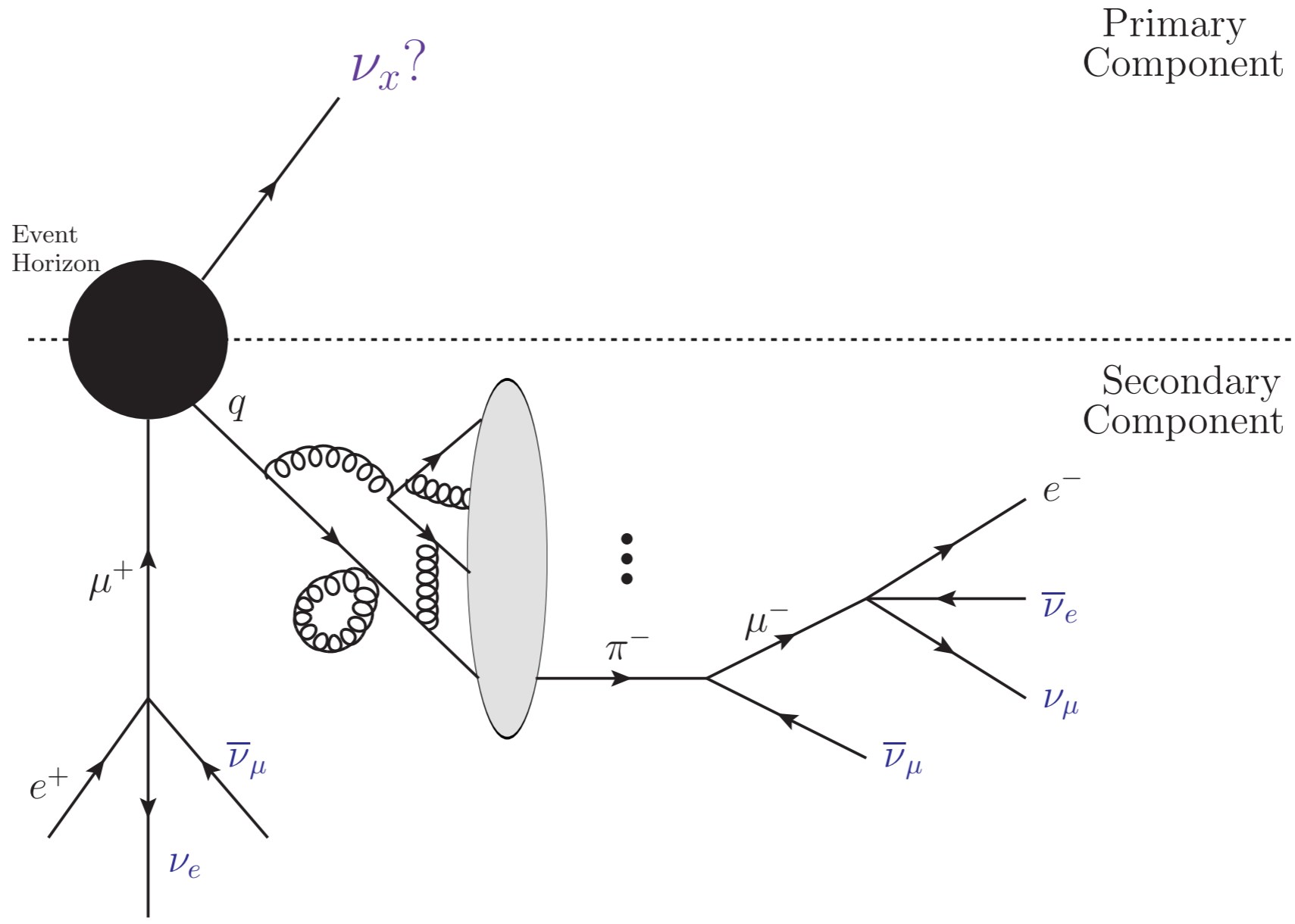
B. Carr, 1976

Constraints on the diffuse neutrino flux

J. MacGibbon, 1991
 F. Halzen, 9502268
 Bugaev, 0005295

What is the state of the emitted neutrino?

Neutrinos are massive



Primary Component

Secondary Component

Weak Interaction

What is the state of the emitted neutrino?

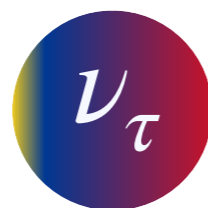
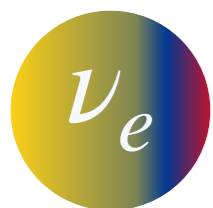
Weak interactions

$$n \rightarrow p^+ + e^- + \bar{\nu}_e$$

Interaction mediated by
a gauge boson

Associated with a
charged lepton

Flavor eigenstate



Hawking Effect

$$\langle 0_- | b_i^\dagger b_i | 0_- \rangle = \Gamma_{lm} \left[\exp(E_a/T_{\text{BH}} + 1) \right]^{-1}$$

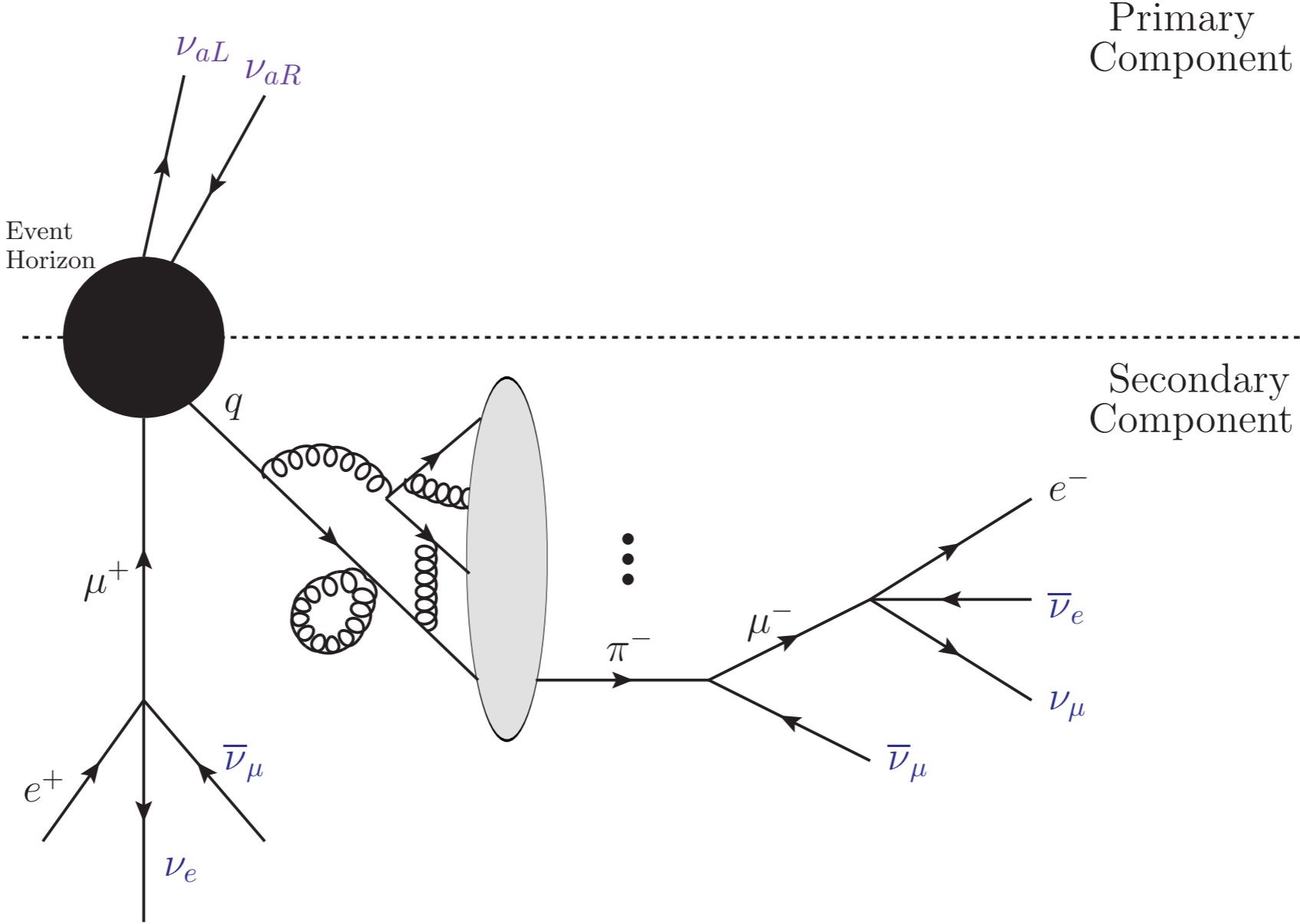
Particle definition in a curved
spacetime is observer dependent

Particles with well-
defined kinematical
properties

Mass eigenstate



Dirac vs Majorana



Majorana neutrinos

Phenomenological consequences?

Dirac neutrinos

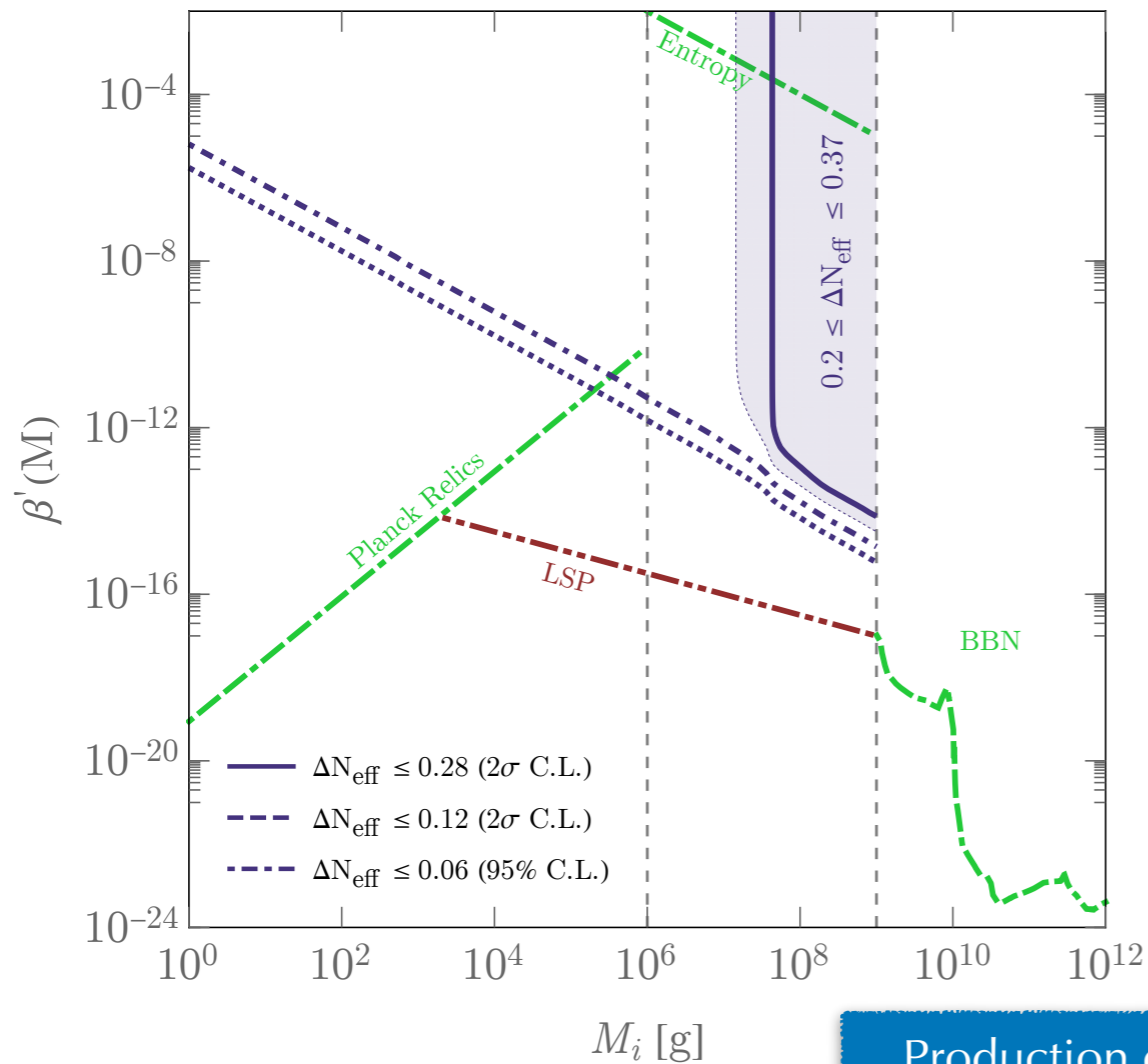
VS

Majorana neutrinos

$$\sigma_{\text{abs}}^{\nu}(+1/2) = \sigma_{\text{abs}}^{\nu}(-1/2)$$

Unruh, 1976

No helicity suppression



Production of light RH neutrinos!

Heavy RH neutrinos

PBH-induced Leptogenesis

Baumann, Steinhadt, Turok, 0703250
 Fujita et al, 1401.1909
 Morrison et al, 1812.10606
 Hooper and Krnjaic, 2010.01134

YFPG and Jessica Turner,
 PRD 104(2021) 103021
 Bernal, Fong, YFPG, Turner
 2203.08823

Cecilia Lunardini, YFPG
 JCAP08(2020)014

Thermal Leptogenesis in a nutshell

Baryon asymmetry

$$\eta_{\text{CMB}} = (6.23 \pm 0.17) \times 10^{-10}$$

$$\eta_{\text{BBN}} = (6.08 \pm 0.06) \times 10^{-10}$$

Type I seesaw:

$$-\mathcal{L} \supset \frac{1}{2} M_{N_i} \bar{N}_i^c N_i + \bar{\ell}_\alpha H^* Y_{\alpha i} N_i + \text{H.c.}$$

$$m_\nu \sim \frac{Y_\nu^2 v^2}{M_N} \sim \mathcal{O}(0.1 \text{ eV})$$

Sakharov Conditions

- ❖ Baryon and Lepton number violation
- ❖ CP violation
- ❖ Departure from thermal equilibrium

$$Y_\nu = \frac{1}{v} U_{\text{PMNS}} \sqrt{m_\nu} R^T M_N^T$$

Casas, Ibarra, 2001

Boltzmann Equations

$$\frac{dn_{N_i}}{dz} = D_i (n_{N_i} - n_{N_i}^{\text{eq}})$$

$$\frac{dn_{\text{B-L}}}{dz} = \sum_i \epsilon_i D_i (n_{N_i} - n_{N_i}^{\text{eq}}) - \mathcal{W}_i n_{\text{B-L}}$$

Washout

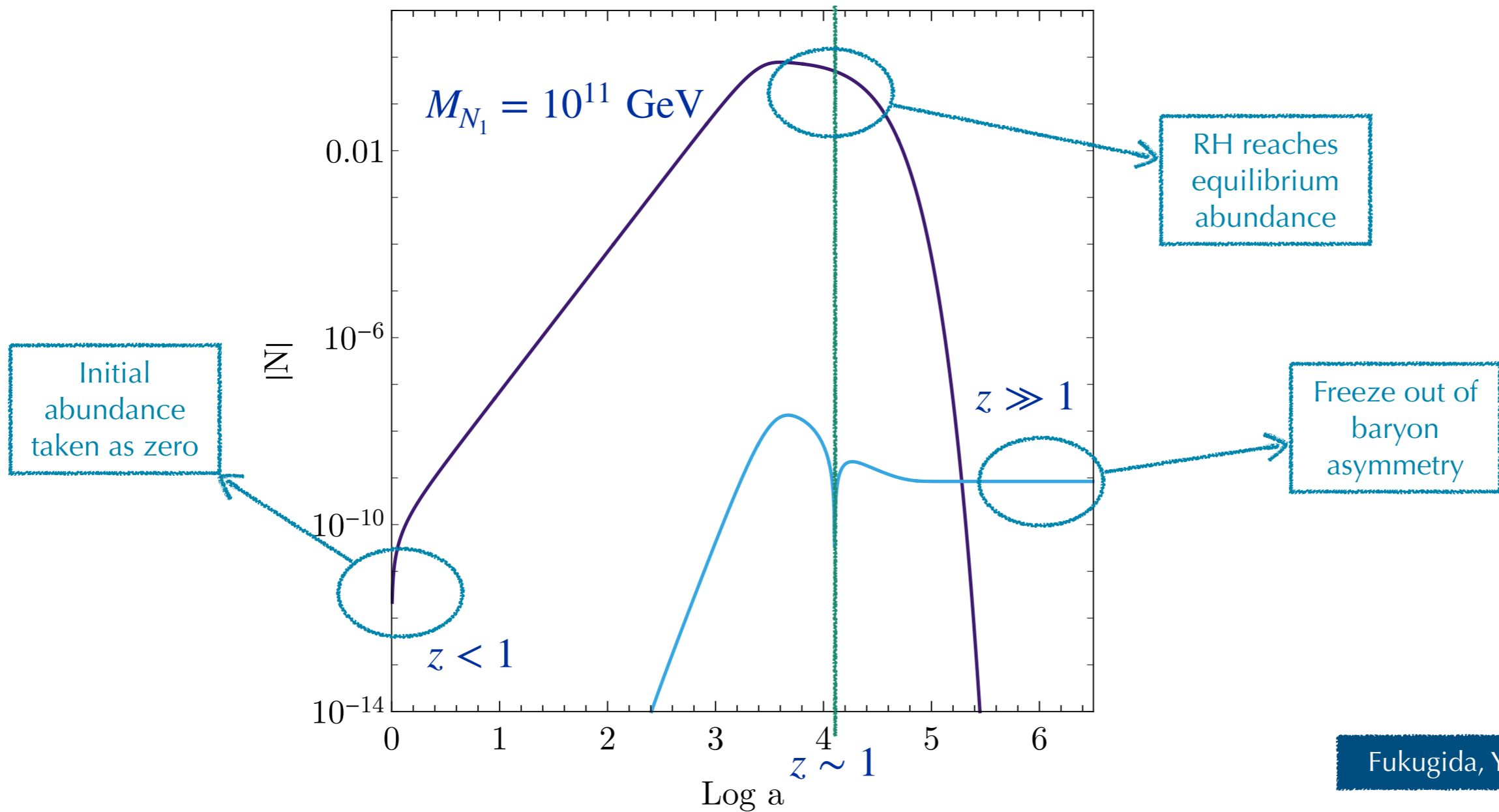
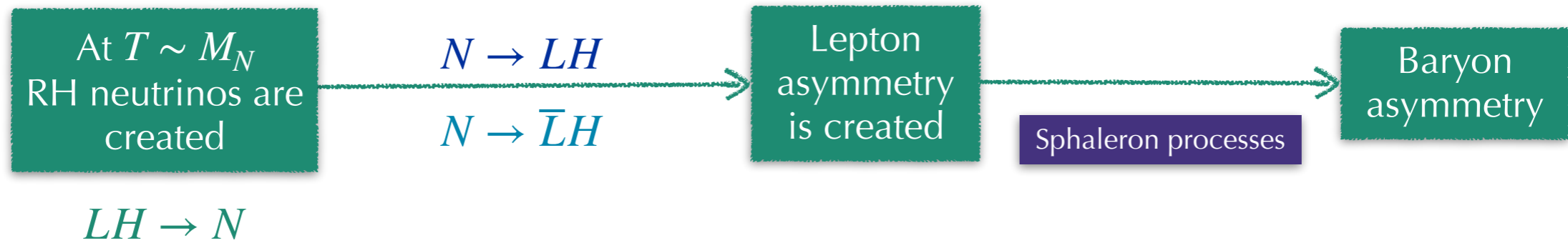
$$LH \leftrightarrow \bar{L}H$$

$$\Delta L = 2$$

Fukugida, Yanagida, '86

Thermal Leptogenesis in a nutshell

$$z \equiv \frac{M_{N_1}}{T}$$



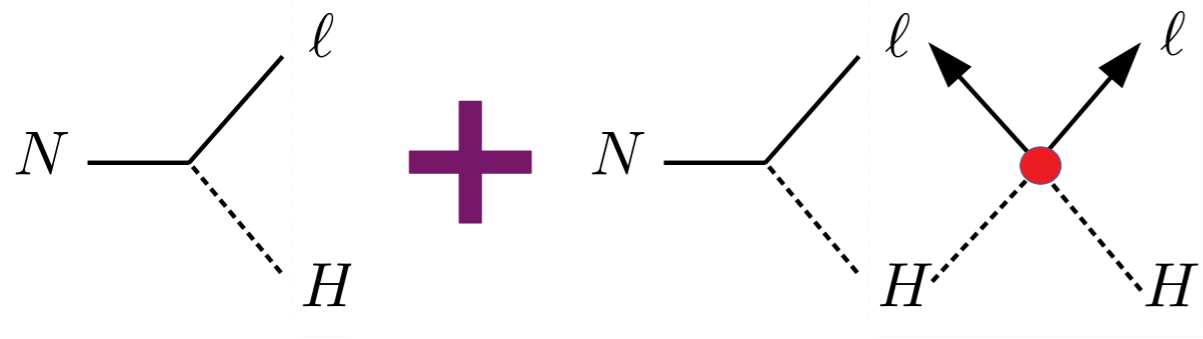
Fukugida, Yanagida, '86

High Scale Leptogenesis

$$M_N \gtrsim 10^{12} \text{ GeV}$$

$$M_{N_1} < M_{N_2} < M_{N_3}$$

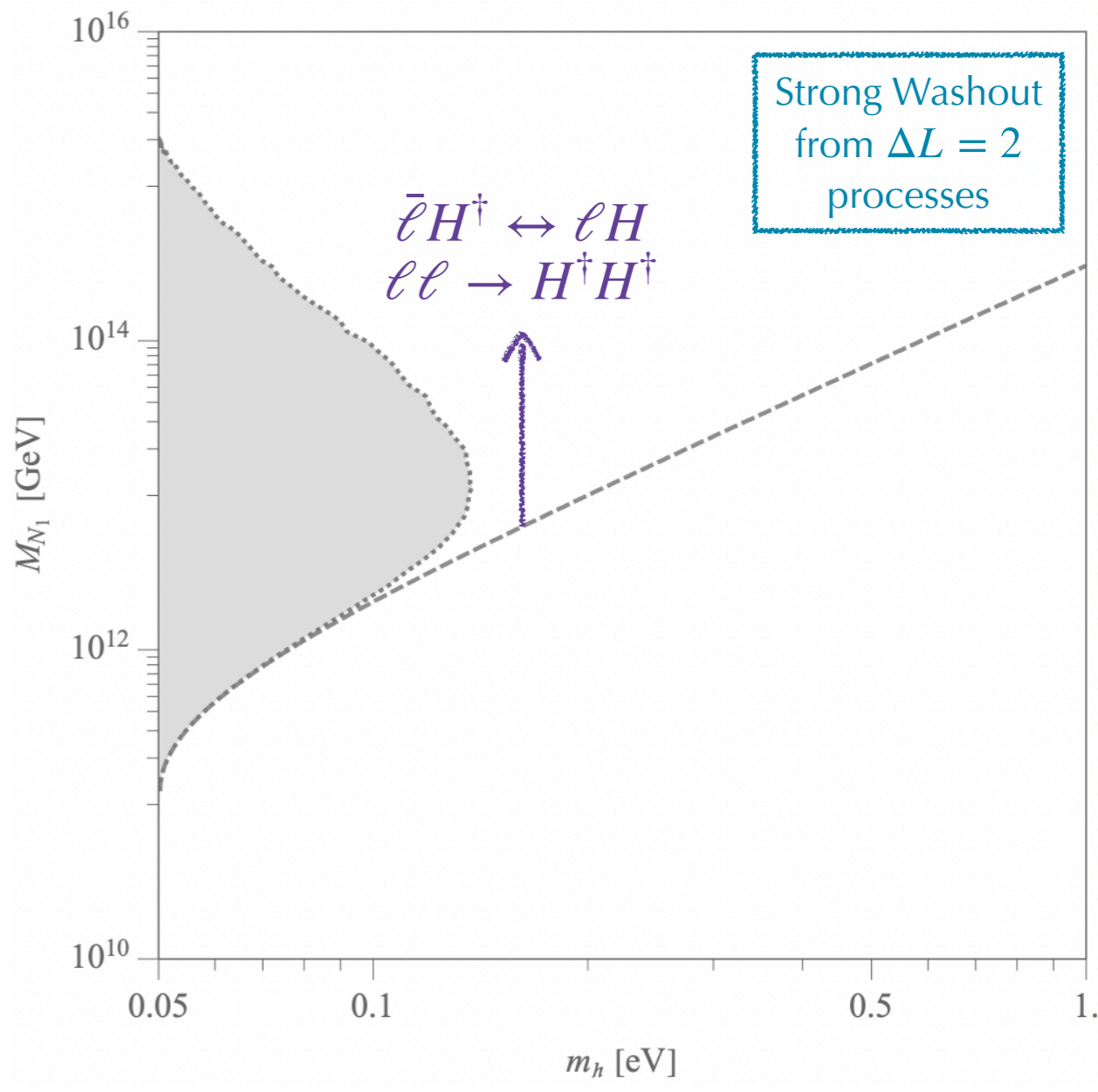
CP violation from interference



$$\frac{\epsilon}{M_N} \propto \frac{m_\nu}{v^2} \propto \sqrt{\sigma_{\Delta L=2}}$$

$\Delta L = 2$ interactions would be in equilibrium if

$$M_{N_1}^T \gtrsim 4 \left(\frac{0.1 \text{ eV}}{m_\nu} \right) 10^{12} \text{ GeV}$$



$m_h =$ Heaviest neutrino mass

Maximizing over Yukawa parameters

How to save HSL?

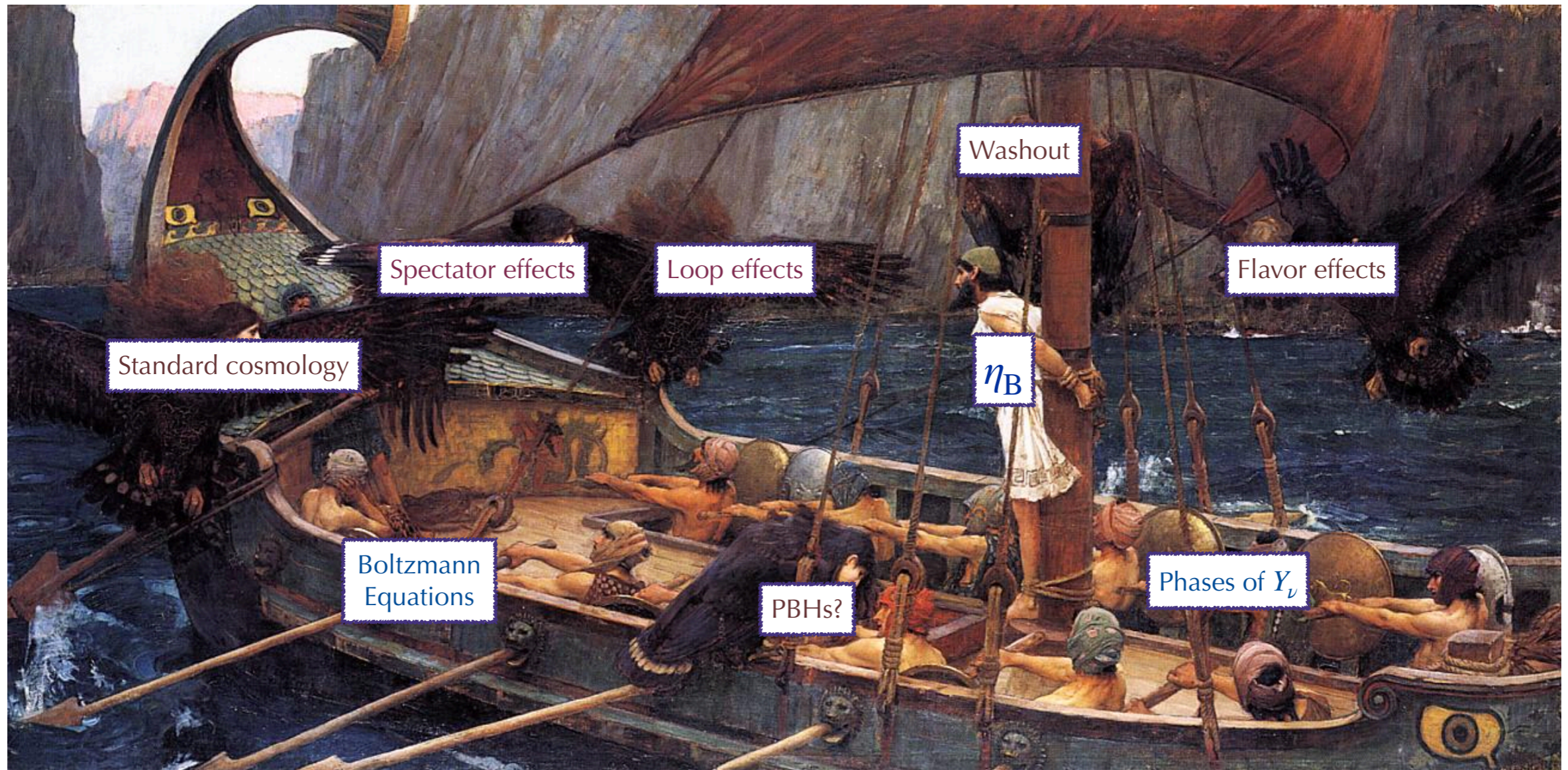
Produce RHNs after washout process have frozen out?

PBHs!

Giudice *et al.*, 2004
Buchmuller *et al.*, 2005

Intermezzo.

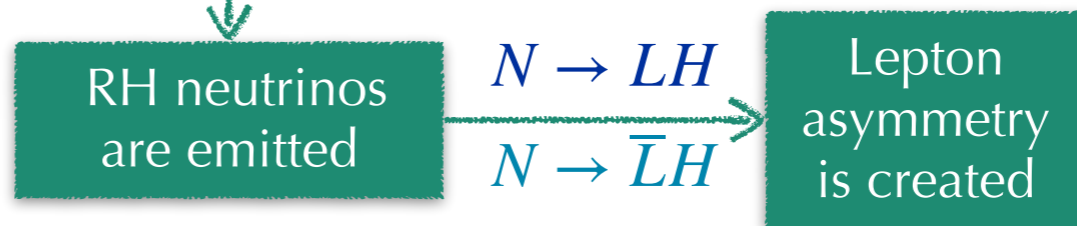
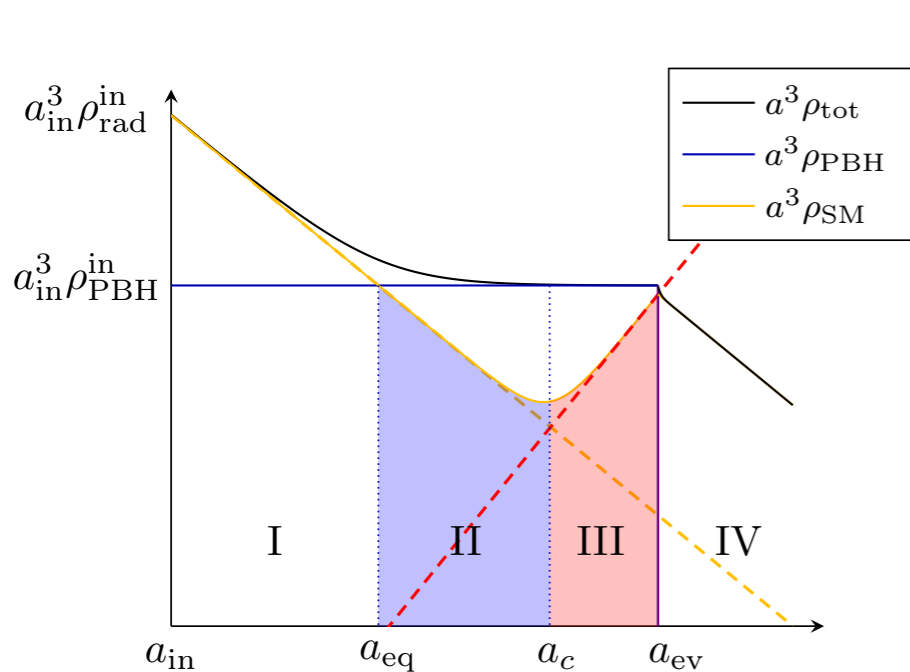
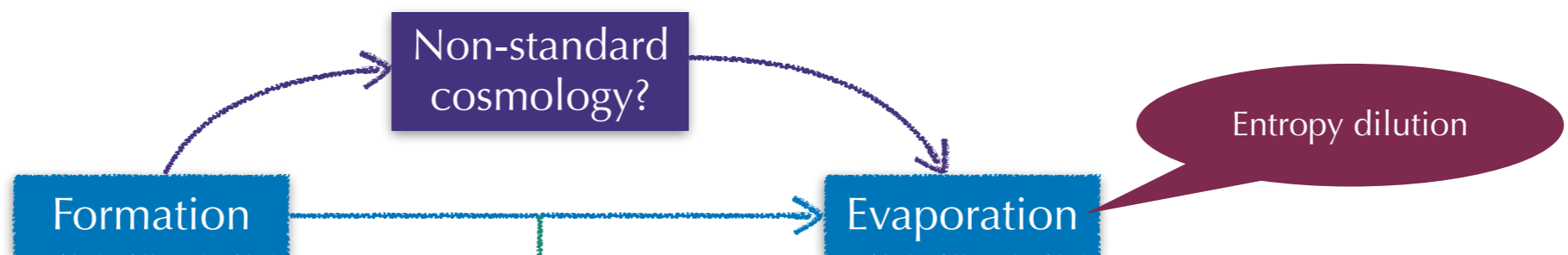
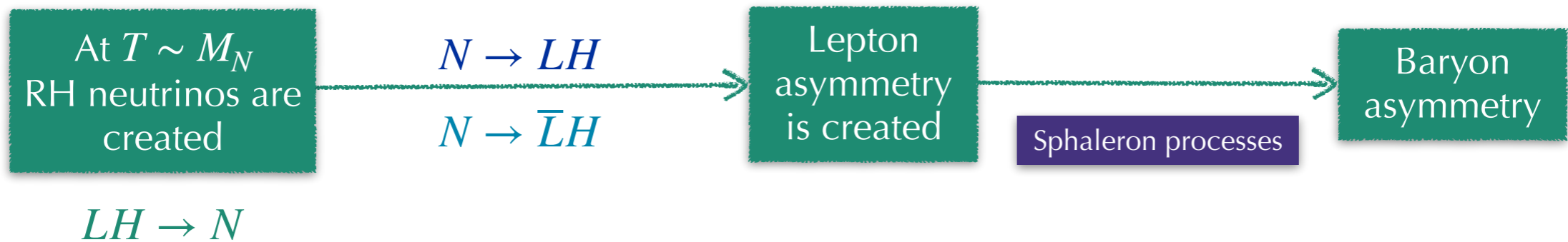
Universal LeptogeneSiS Equation Solver (ULYSSES)



A Granelli, K Moffat, YFPG,
H Schulz and J Turner,
arXiv: [2007.09150](https://arxiv.org/abs/2007.09150)
arXiv: [2301.05722](https://arxiv.org/abs/2301.05722)

- ❖ Leptogenesis via decays and resonant leptogenesis
- ❖ ARS Leptogenesis
- ❖ Easy parallelization
- ❖ Rapid evaluation
- ❖ Multidimensional scan of the parameter space

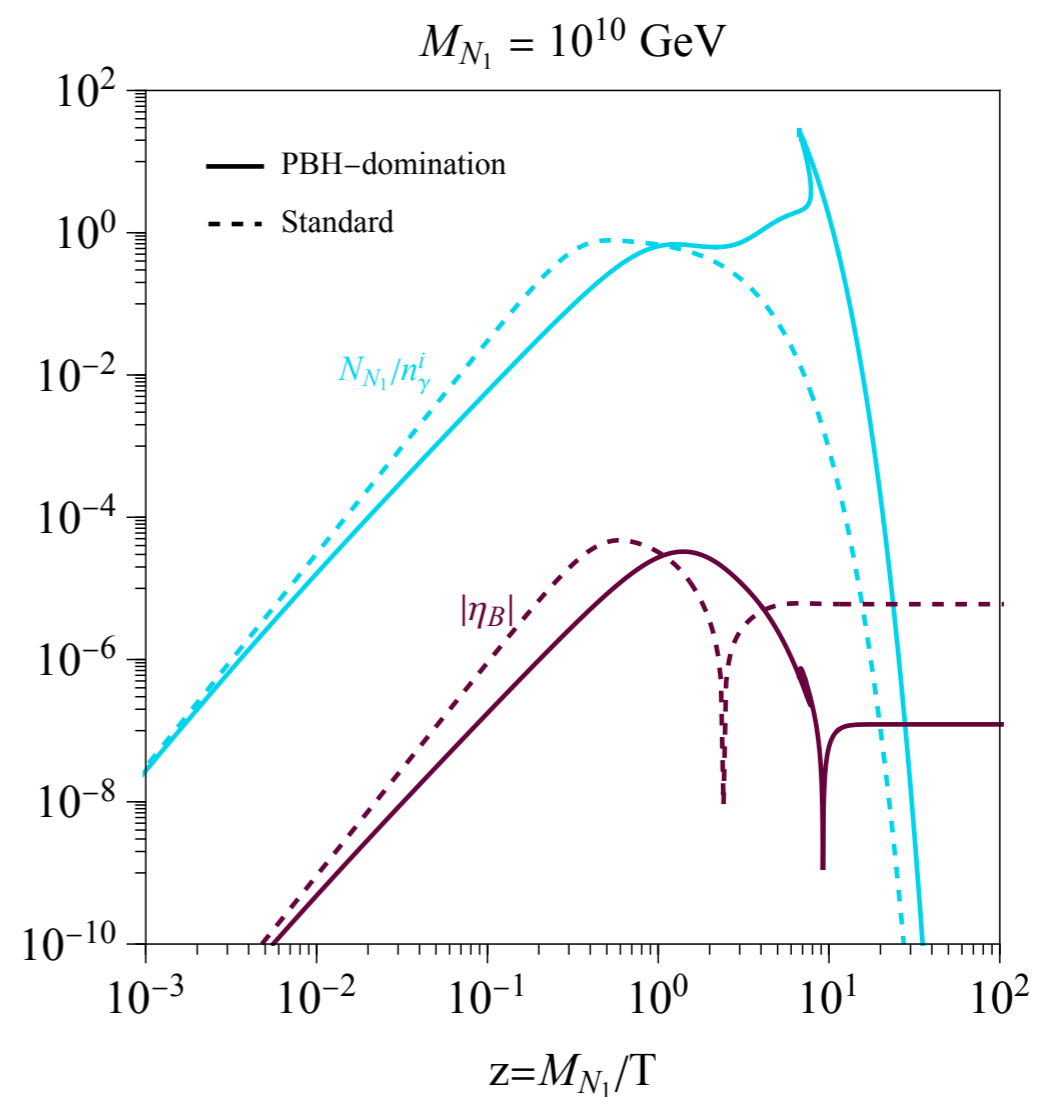
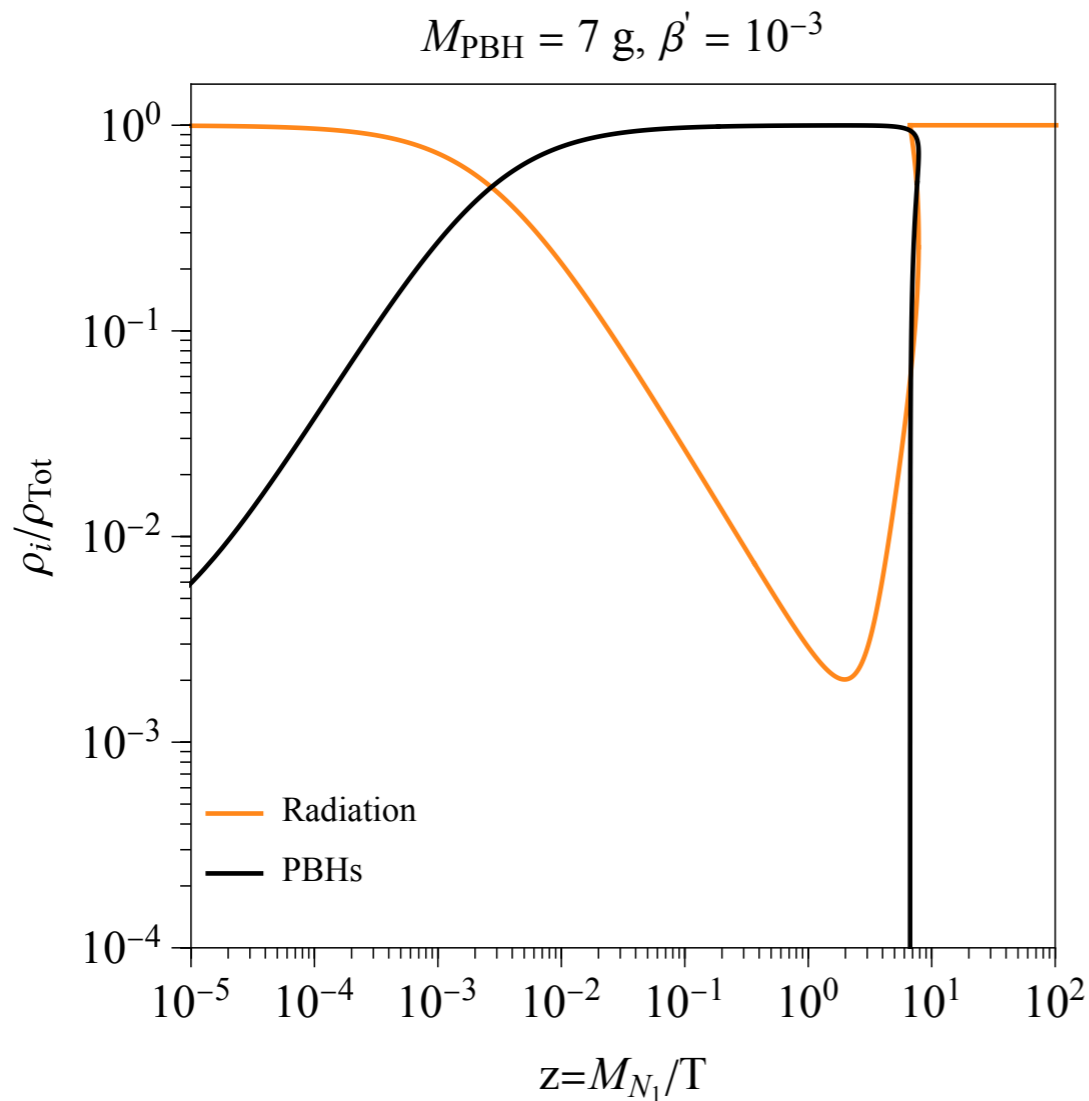
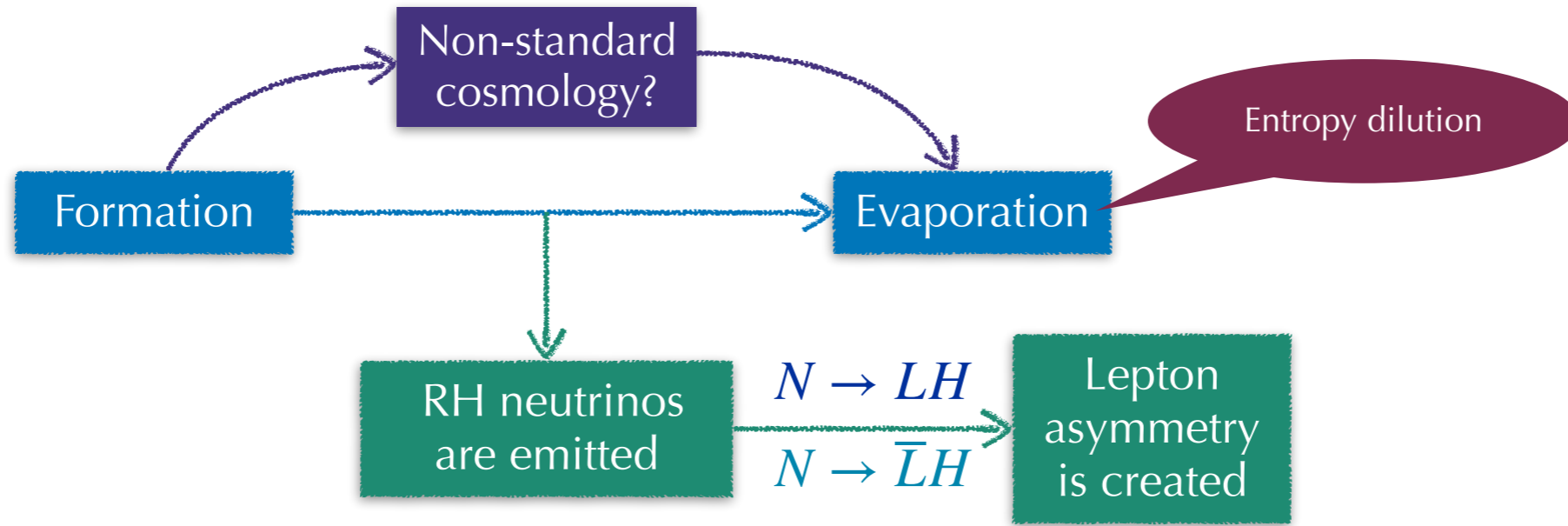
PBH + Leptogenesis



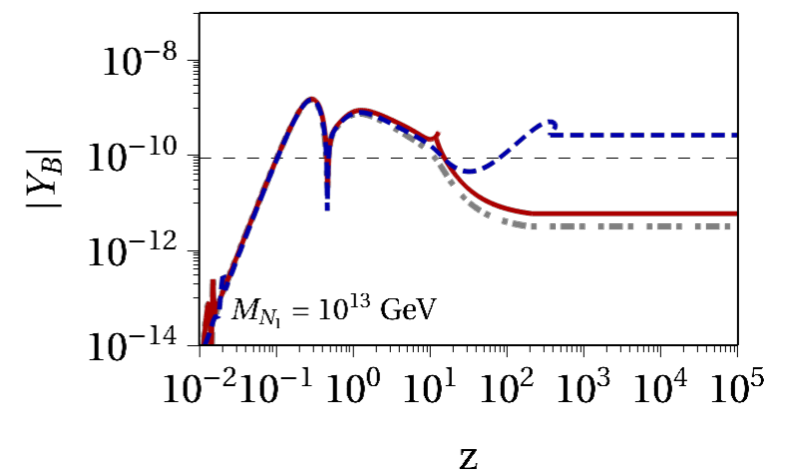
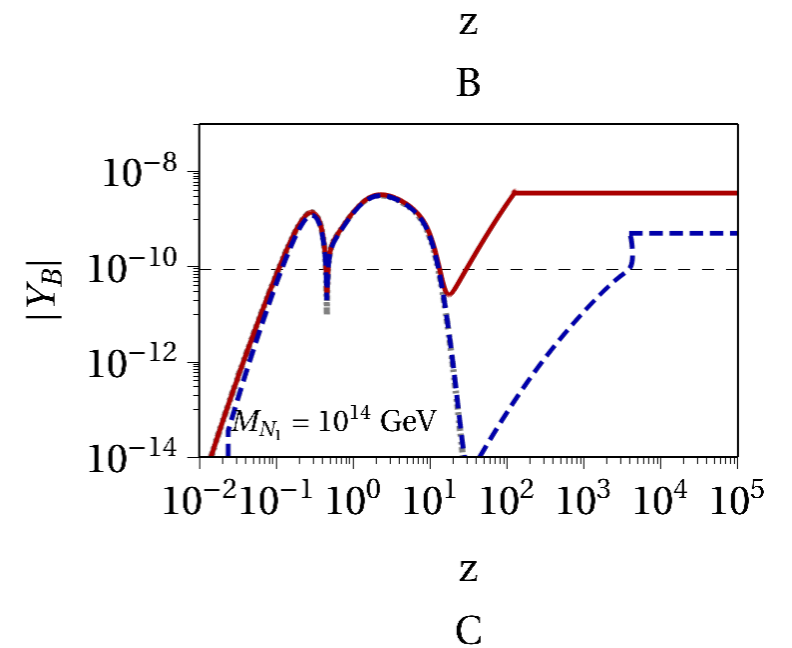
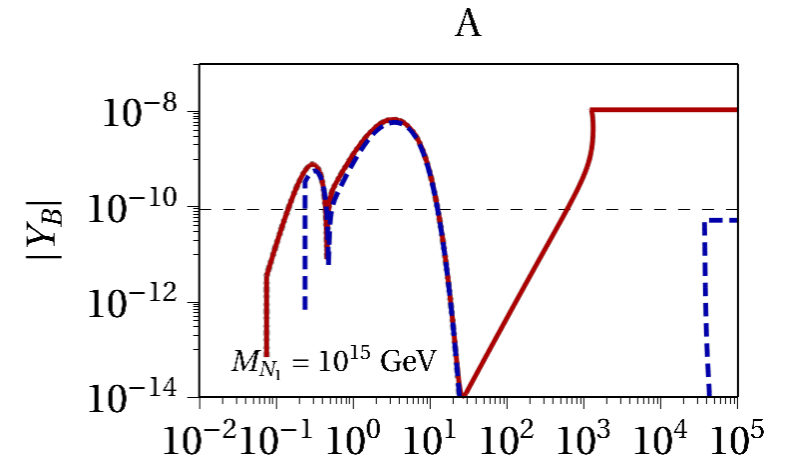
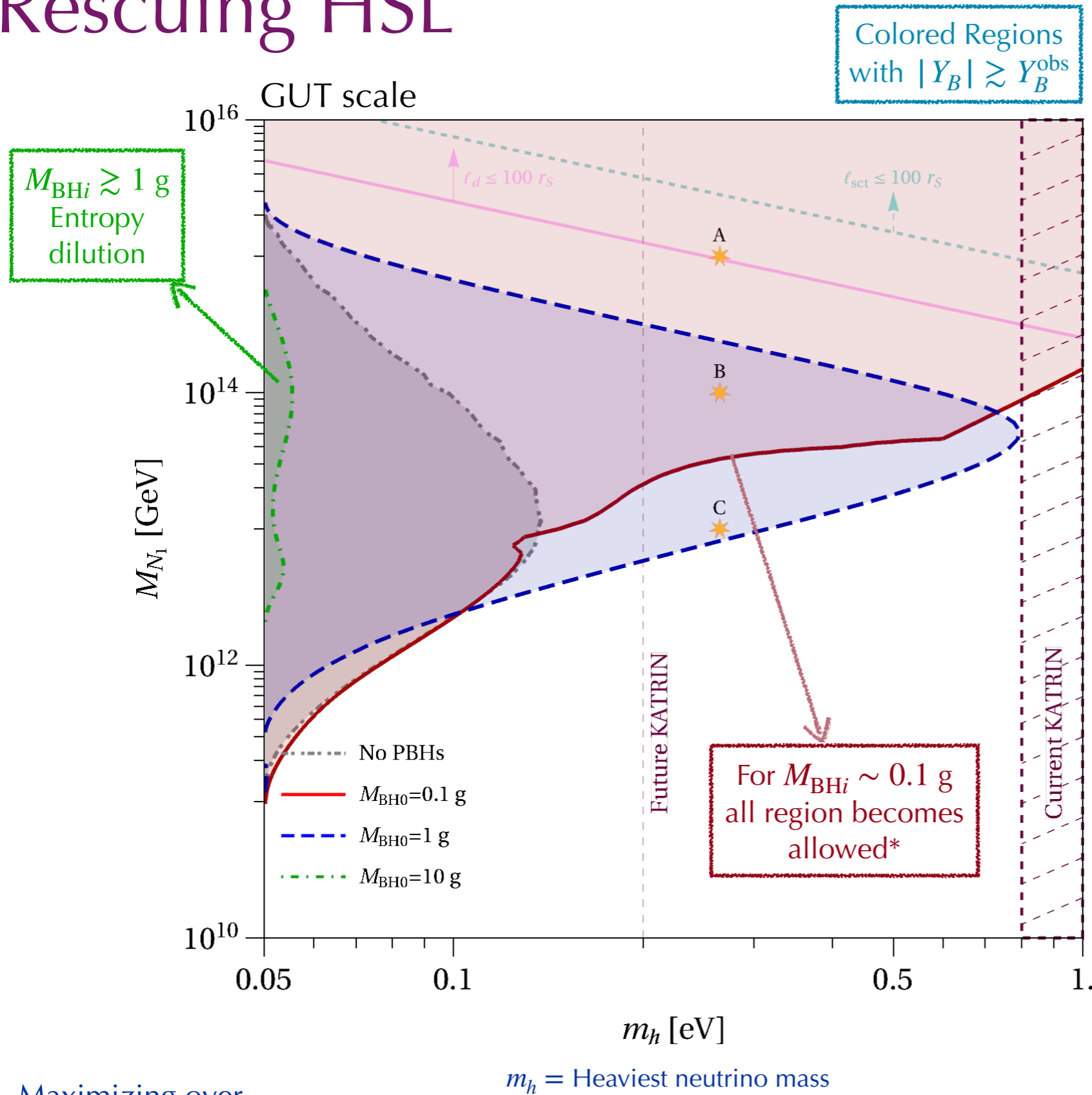
$$z_{BH} \equiv \frac{M_{N_1}}{T_{BH}}$$

PBH + Leptogenesis

$$z_{\text{BH}} \equiv \frac{M_{N_1}}{T_{\text{BH}}}$$



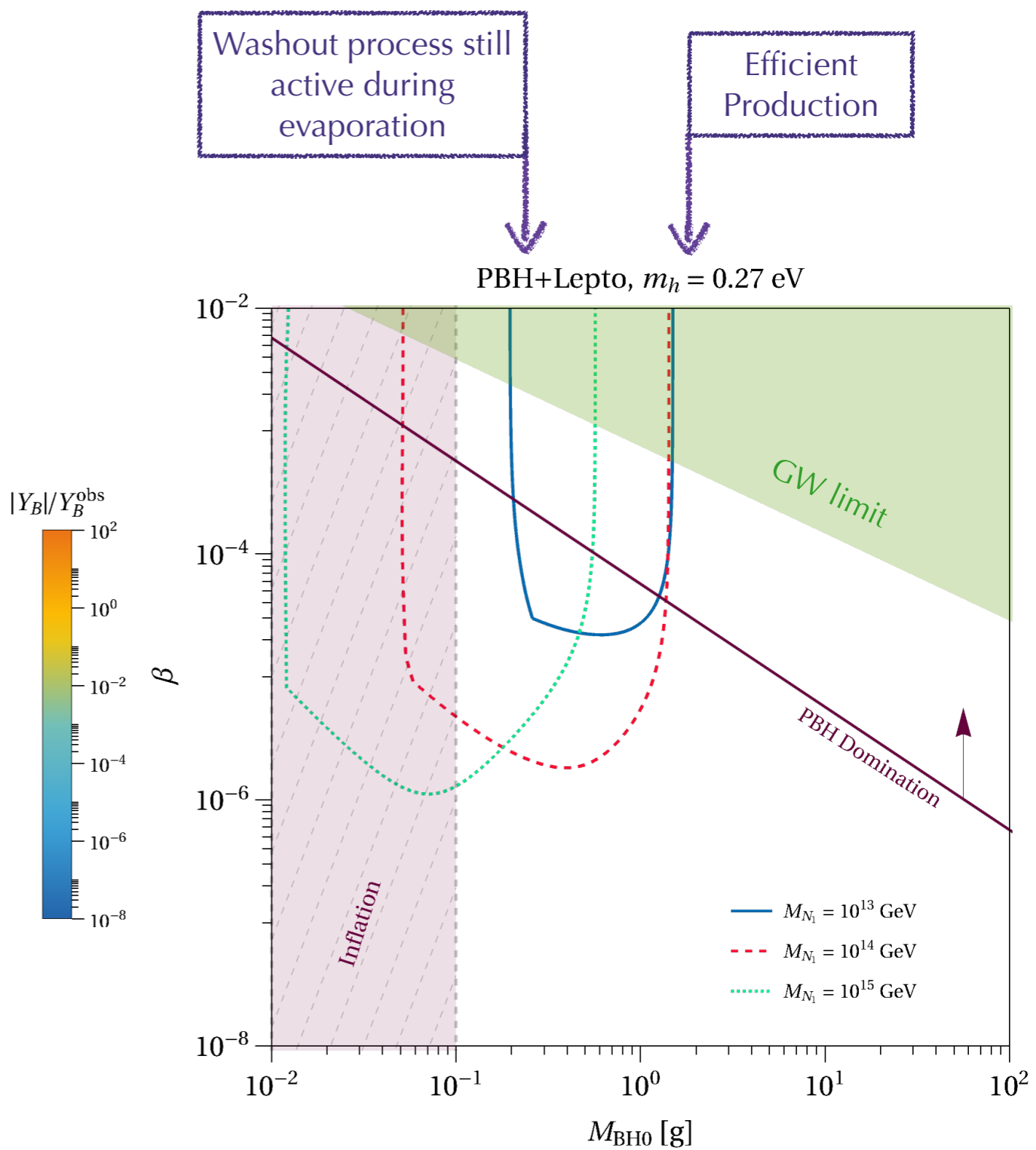
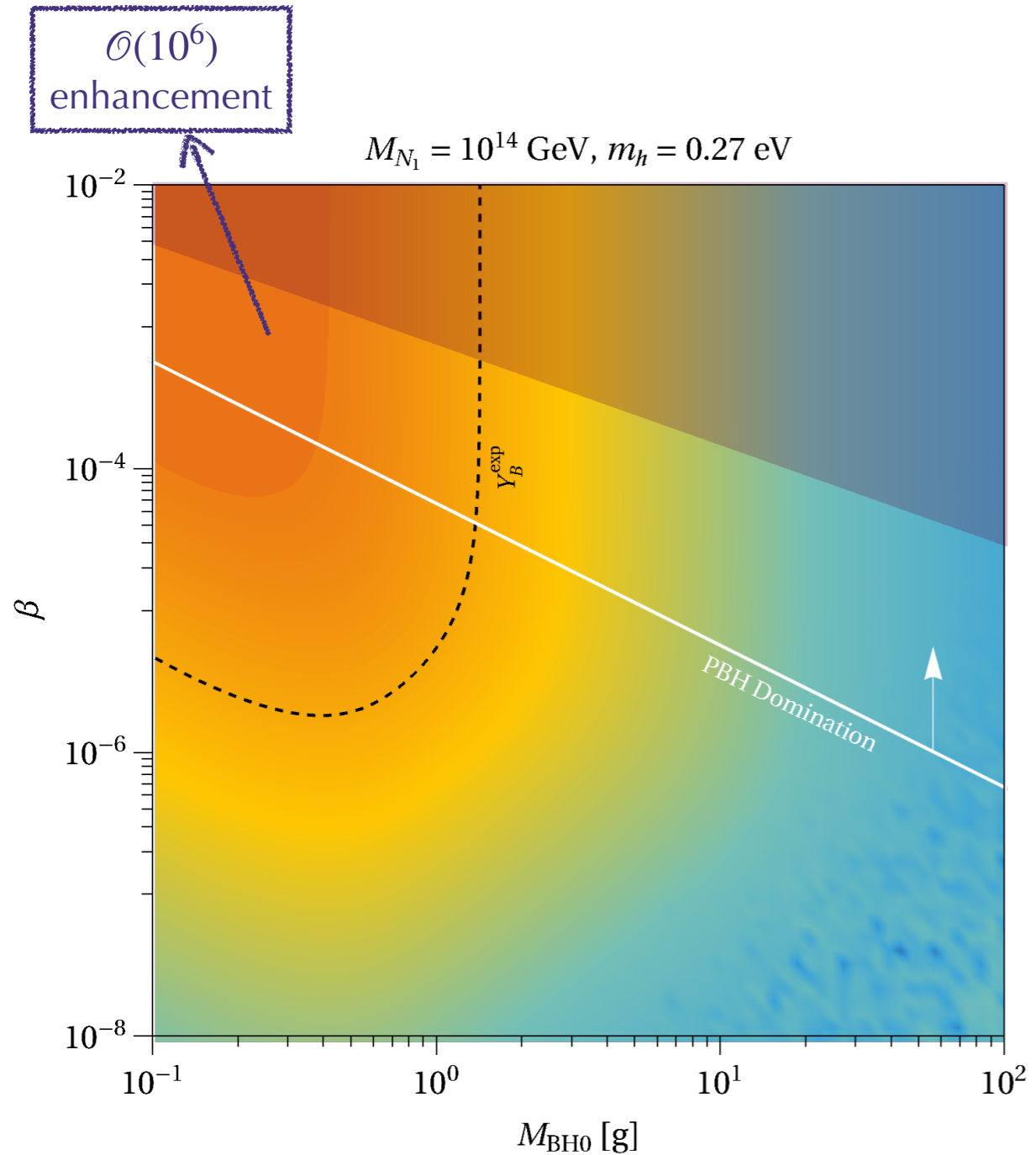
Rescuing HSL



Maximizing over Yukawa parameters

*Up to perturbativity

Rescuing HSL



Mass Distributions

Monochromatic approximation *too* approximated?

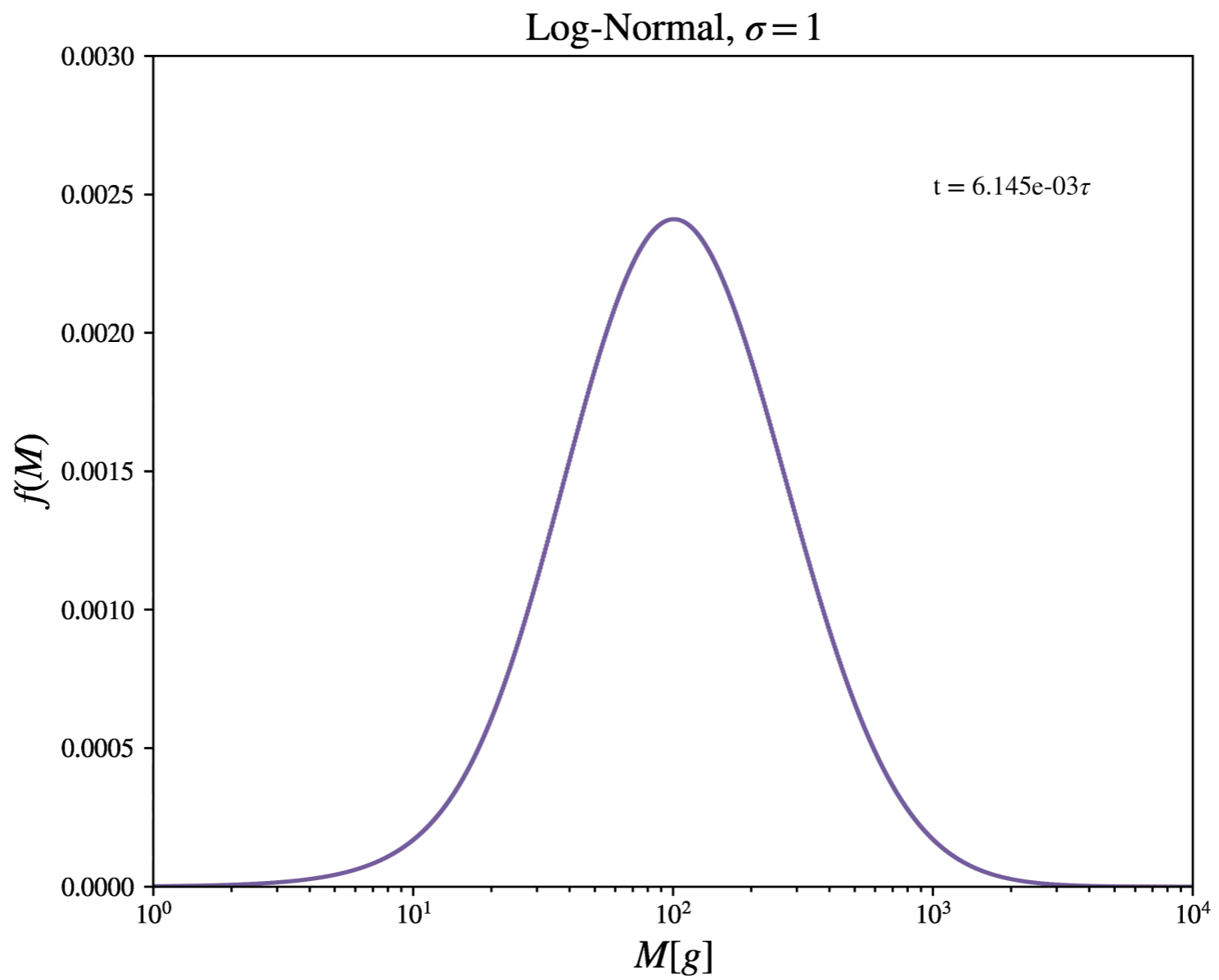
$$n_{\text{PBH}} = \int dM f(M) \quad f(M) = \frac{n_{\text{BH}}}{\sqrt{2\pi\sigma M}} \exp\left(-\frac{\log^2(M/M_c)}{2\sigma^2}\right)$$

Log-normal distribution

Dolgov, 93
Green, 2016
Kannike, 2017

Connection with different formation mechanisms?

Having PBHs with different masses could have a distinct impact on the previous results



Mass Distributions

Dolgov, 93
Green, 2016
Kannike, 2017

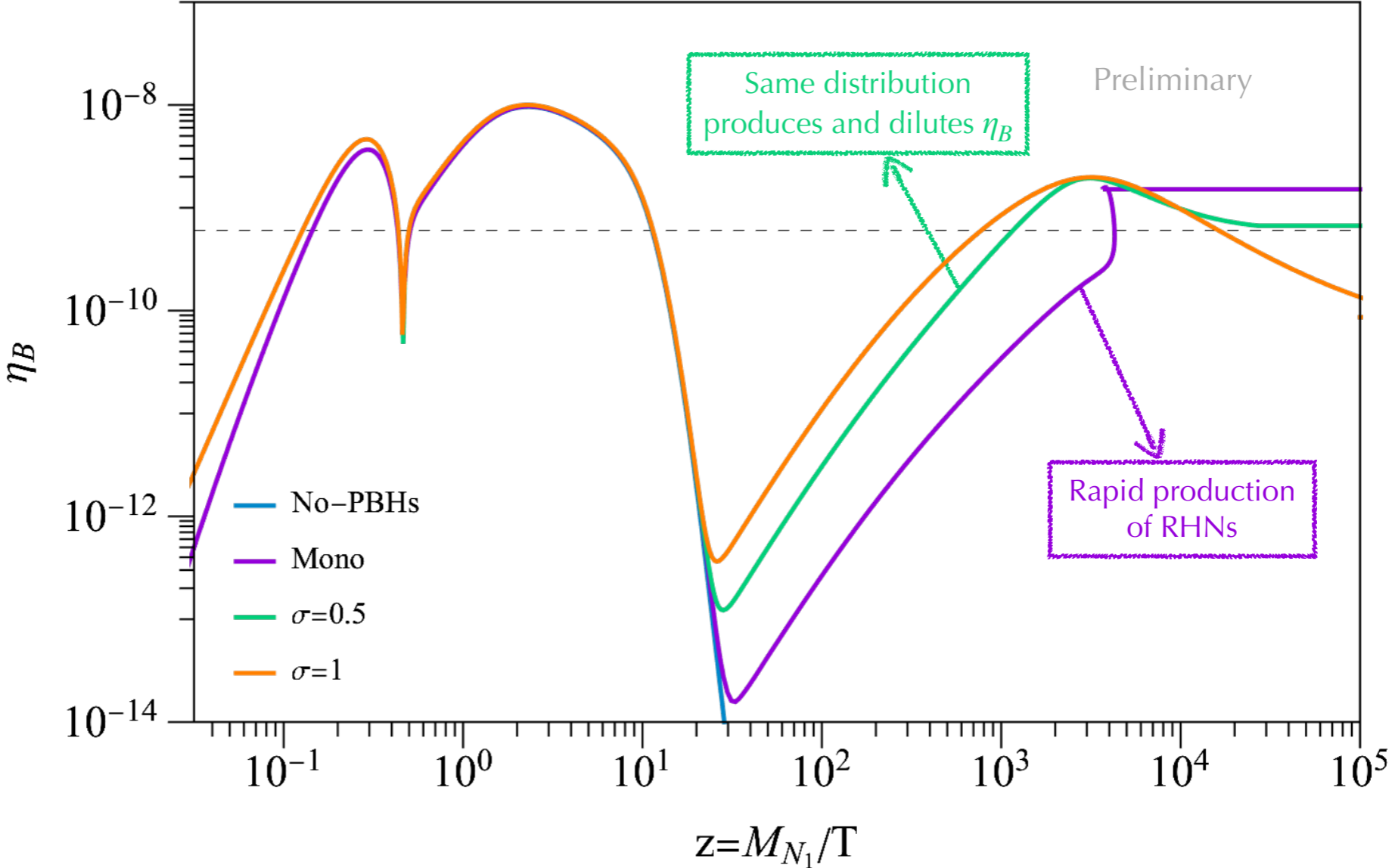
Monochromatic approximation *too* approximated?

$$n_{\text{PBH}} = \int dM f(M) \quad f(M) = \frac{n_{\text{BH}}}{\sqrt{2\pi\sigma M}} \exp\left(-\frac{\log^2(M/M_c)}{2\sigma^2}\right)$$

Log-normal distribution

$$M_c = 1 \text{ g}$$

Connection with different formation mechanisms?



Having PBHs with different masses could have a distinct impact on the previous results

Why is interesting the particle production via evaporation?

DM production
(Besides PBH-DM)

- ❖ Purely Gravitationally interacting DM
- ❖ Modify Freeze-In/Freeze-out mechanisms

Fujita et al, 1401.1909
 Morrison et al, 1812.10606
 Baldes et al, 2004.14773
 Masina, 2004.04740, 2103.13825
 Cheek, Heurtier, YFPG, Turner
 2107.00013, 2107.00016,
2212.03878
 .
 .
 .



Baryon Asymmetry

- ❖ Modifying Baryogenesis scenarios
- ❖ Leptogenesis scenarios
- ❖ Producing a local asymmetry at PBH formation
- ❖ Connections with PBH-DM

Baumann, Steinhadt, Turok, 0703250
 Yamada and Iso, 1610.02586
 Fujita et al, 1401.1909
 Morrison et al, 1812.10606
 García-Bellido, Carr, Clesse, 1904.11482
 Hooper and Krnjaic, 2010.01134
 .
 .
 .



Dark Radiation

- ❖ Production of hot gravitons
- ❖ Testable from future measurements on ΔN_{eff}

Hooper, Krnjaic, McDermott, 1905.01301
 Lunardini, YFPG, 1910.07864
 Masina, 2004.04740, 2103.13825
 Cheek, Heurtier, YFPG, Turner, 2207.09462
 .
 .

Second Act Dark Matter

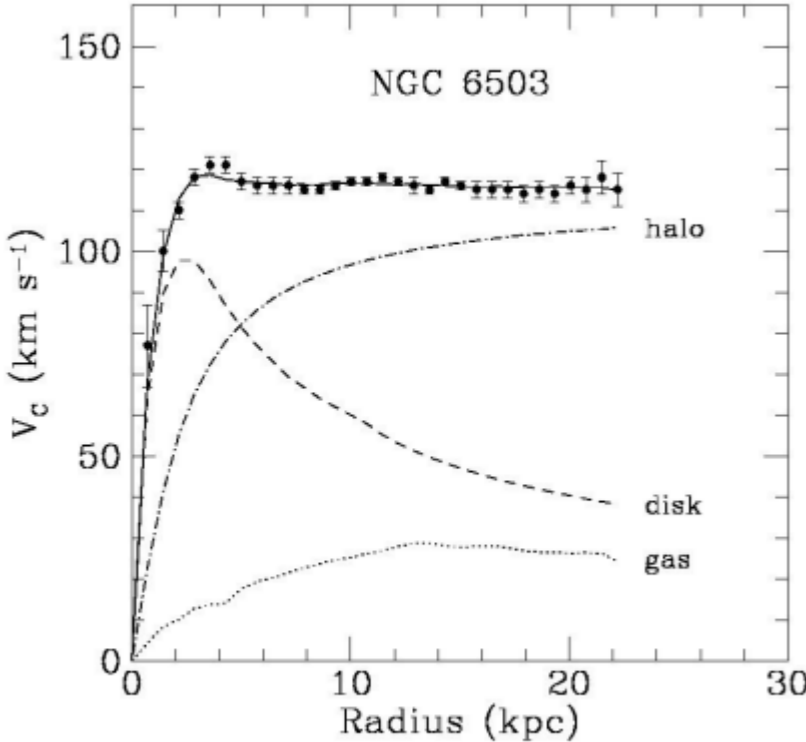
Based on:

Cheek, Heurtier, YFPG, Turner 2107.00013,
2107.00016, [2212.03878](#)

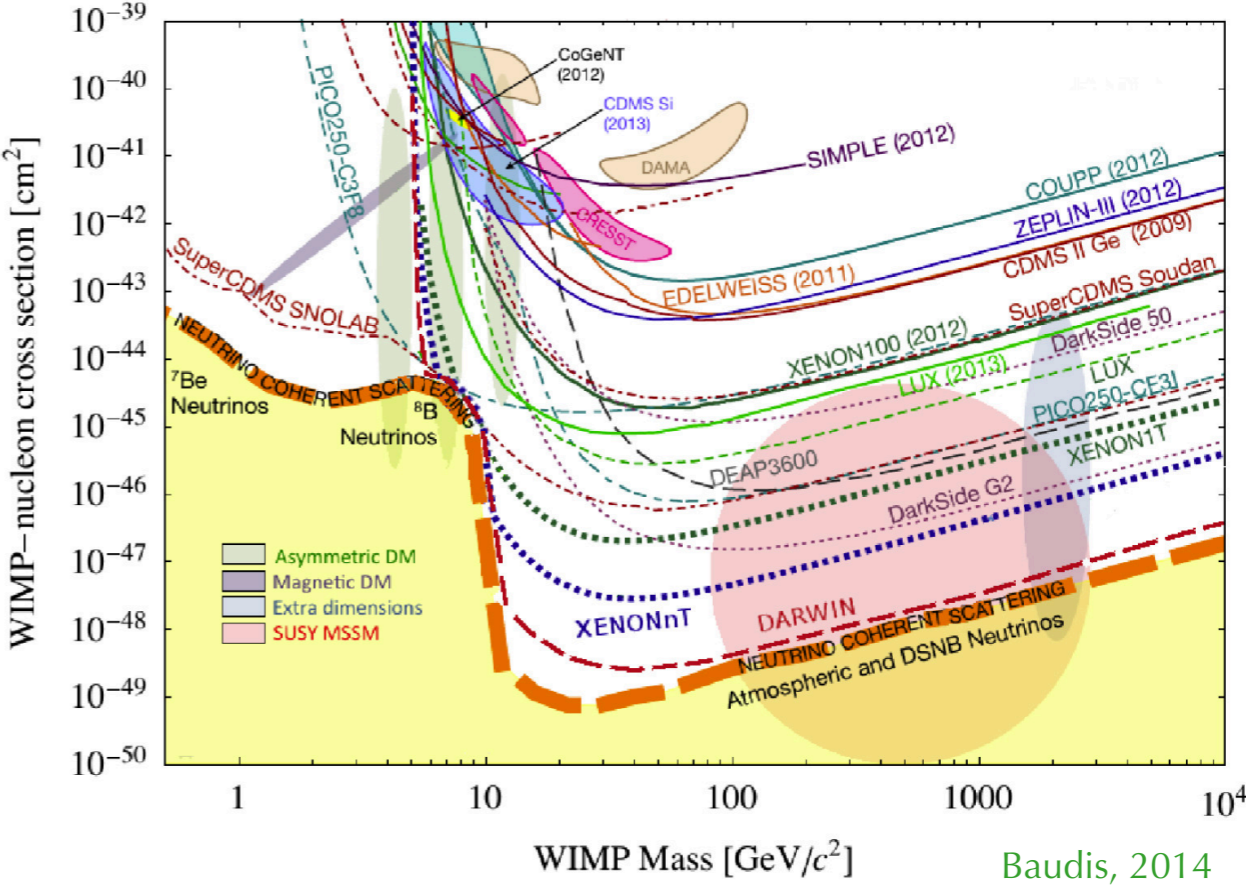
Cheek, Heurtier, YFPG, Turner, [2207.09462](#)

Only Gravitational Interacting DM

Plenty of evidence pointing to the existence of Dark Matter



We're searching for it in terrestrial experiments



Particle DM models **often** rely on BSM interactions to produce the relic abundance

Could the DM only interact gravitationally??

How to produce it in the Early Universe?

PBHs!

Gondolo *et al.*, 2020
 Bernal *et al.*, 2020
 Baldes *et al.*, 2020

Only Gravitational Interacting DM

Consider equations for DM number densities

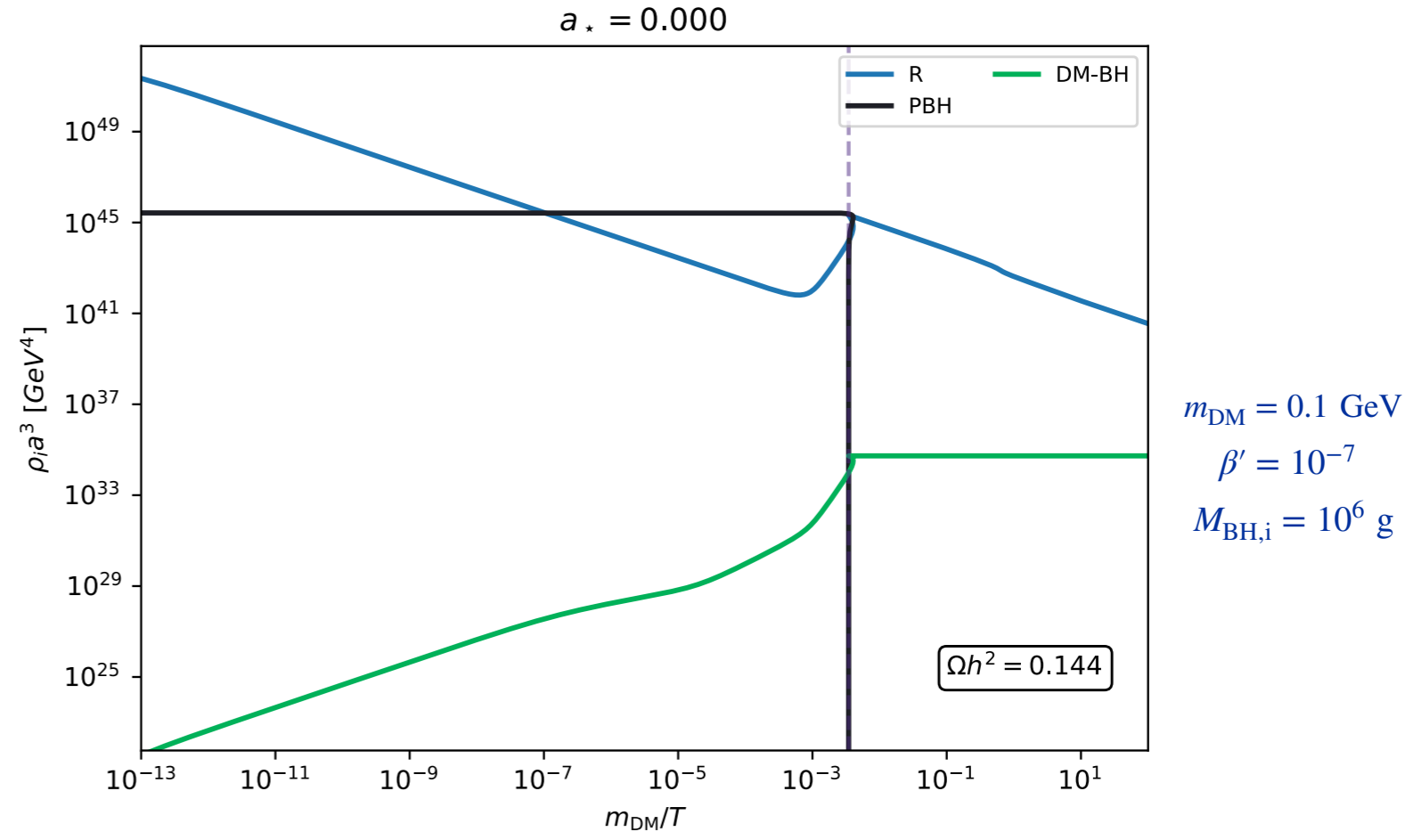
Boltzmann Equations

Evaporation

$$\dot{n}_{\text{DM}} + 3Hn_{\text{DM}} = g_i \int \frac{\partial f_i}{\partial t} \bigg|_{\text{PBH}} \frac{p^2 dp}{2\pi^2} = n_{\text{PBH}} \Gamma_{\text{BH} \rightarrow i}(M_{\text{BH}})$$

$$\dot{\rho}_{\text{SM}} + 4H\rho_{\text{SM}} = - \frac{d \log M}{dt} \bigg|_{\text{SM}} \rho_{\text{PBH}}$$

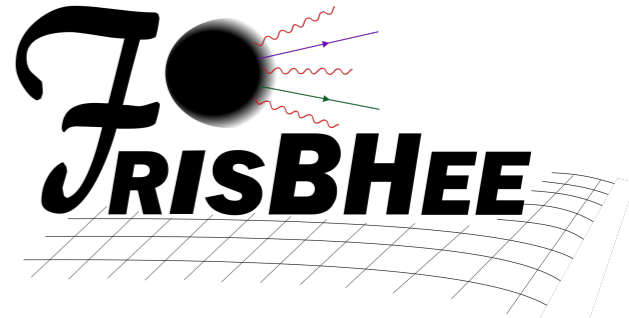
$$\dot{\rho}_{\text{PBH}} + 3H\rho_{\text{PBH}} = \frac{d \log M_{\text{BH}}}{dt} \rho_{\text{PBH}}$$



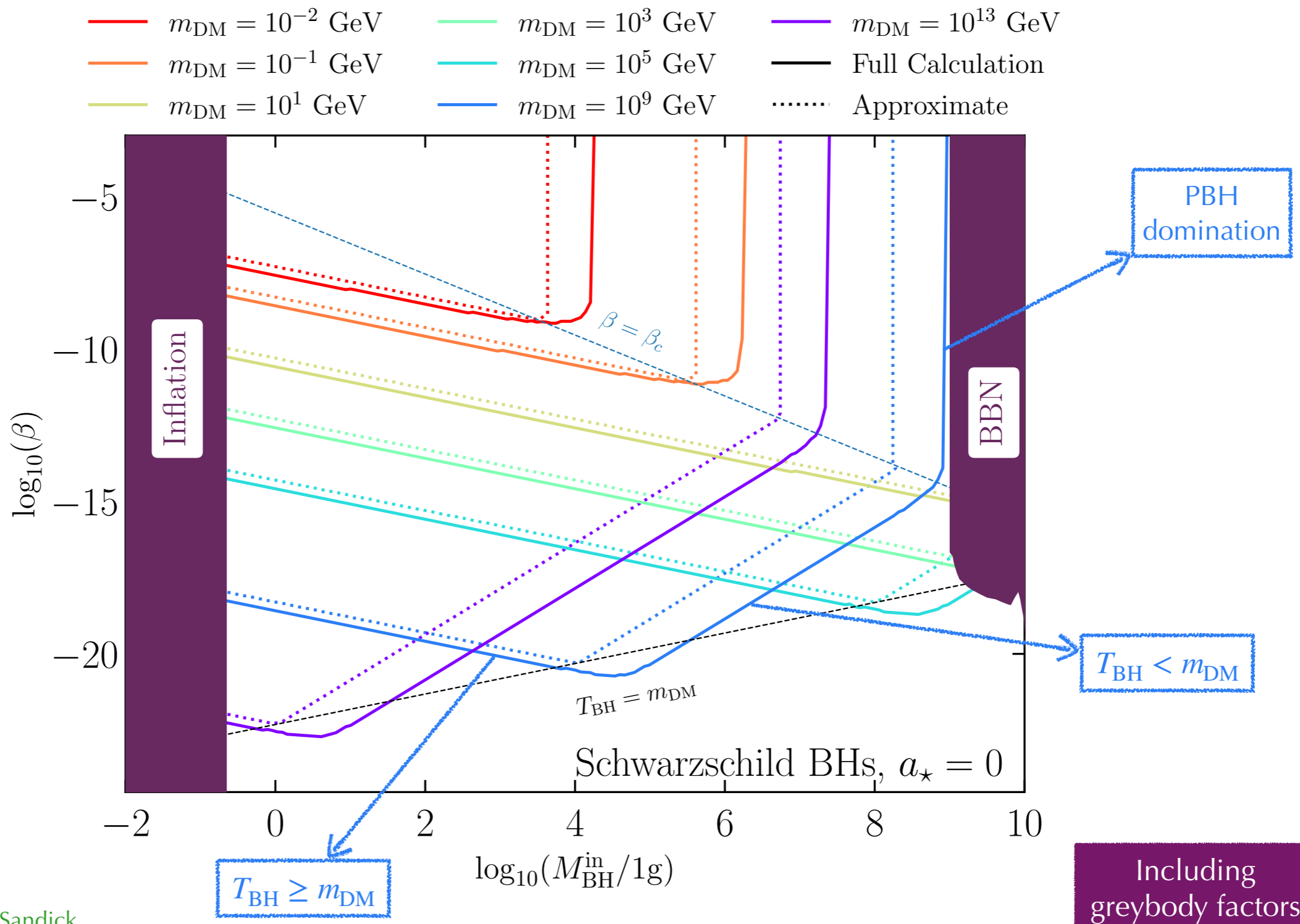
Particle emission rate

$$\Gamma_{\text{BH} \rightarrow i}(M_{\text{BH}}) = \int dp \frac{d^2 \mathcal{N}_i}{dp dt} \quad \frac{3H^2 M_p^2}{8\pi} = \rho_{\text{SM}} + \rho_{\text{DS}} + \rho_{\text{PBH}}$$

Our code is publicly available in

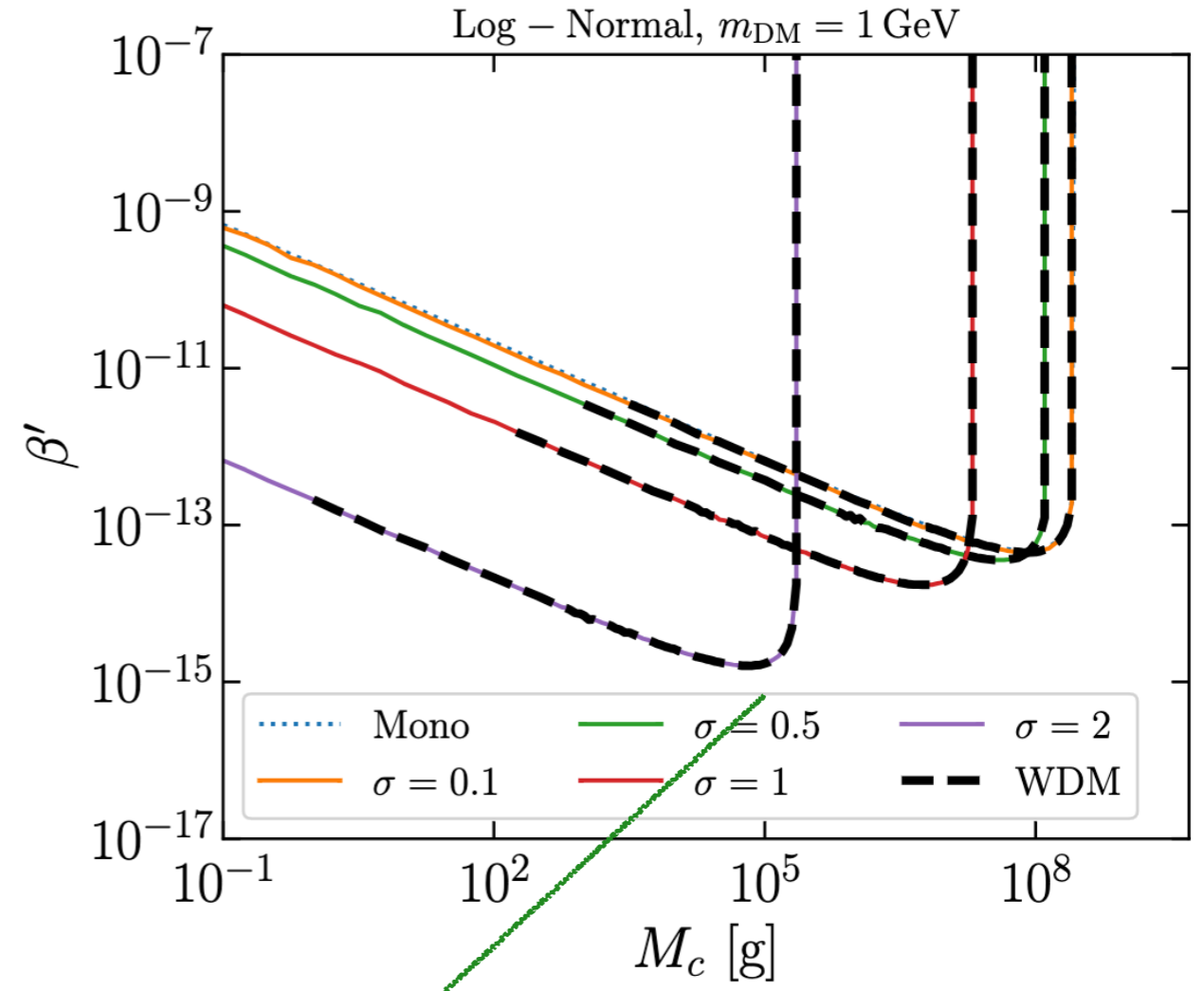
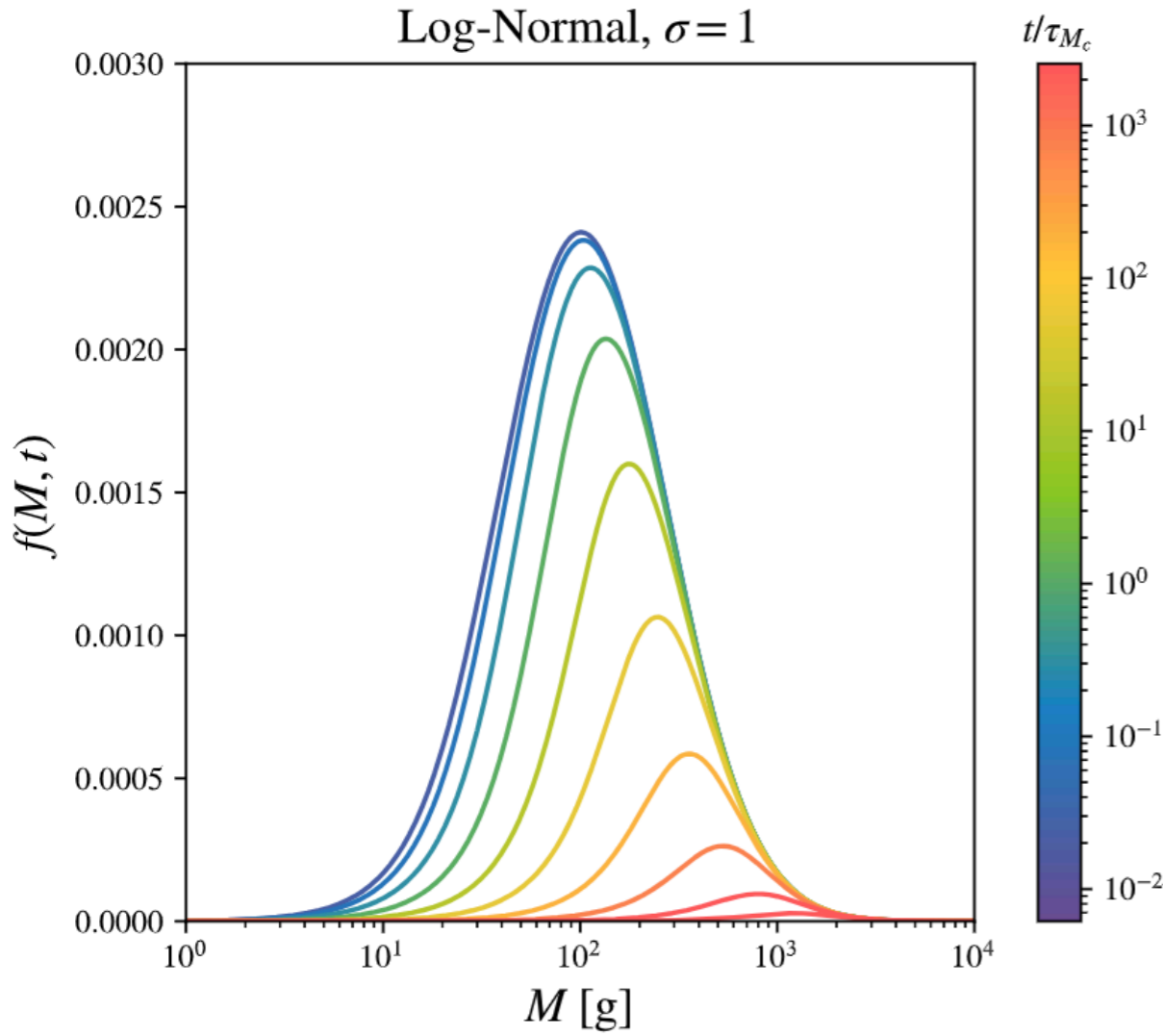


Only Gravitational Interacting DM



Dashed from Gondolo, Sandick, Shams Es Haghi 2009.02424

Mass distributions



$$N_{\text{DM}} \propto \begin{cases} \frac{M_{\text{BH}}^2}{M_p^2} & T_{\text{in}} > m_{\text{DM}} \\ \frac{M_p^2}{m_{\text{DM}}^2} & T_{\text{in}} < m_{\text{DM}} \end{cases}$$

Heavier PBHs produce more light DM

Testing these scenarios?

Gravitational Waves?

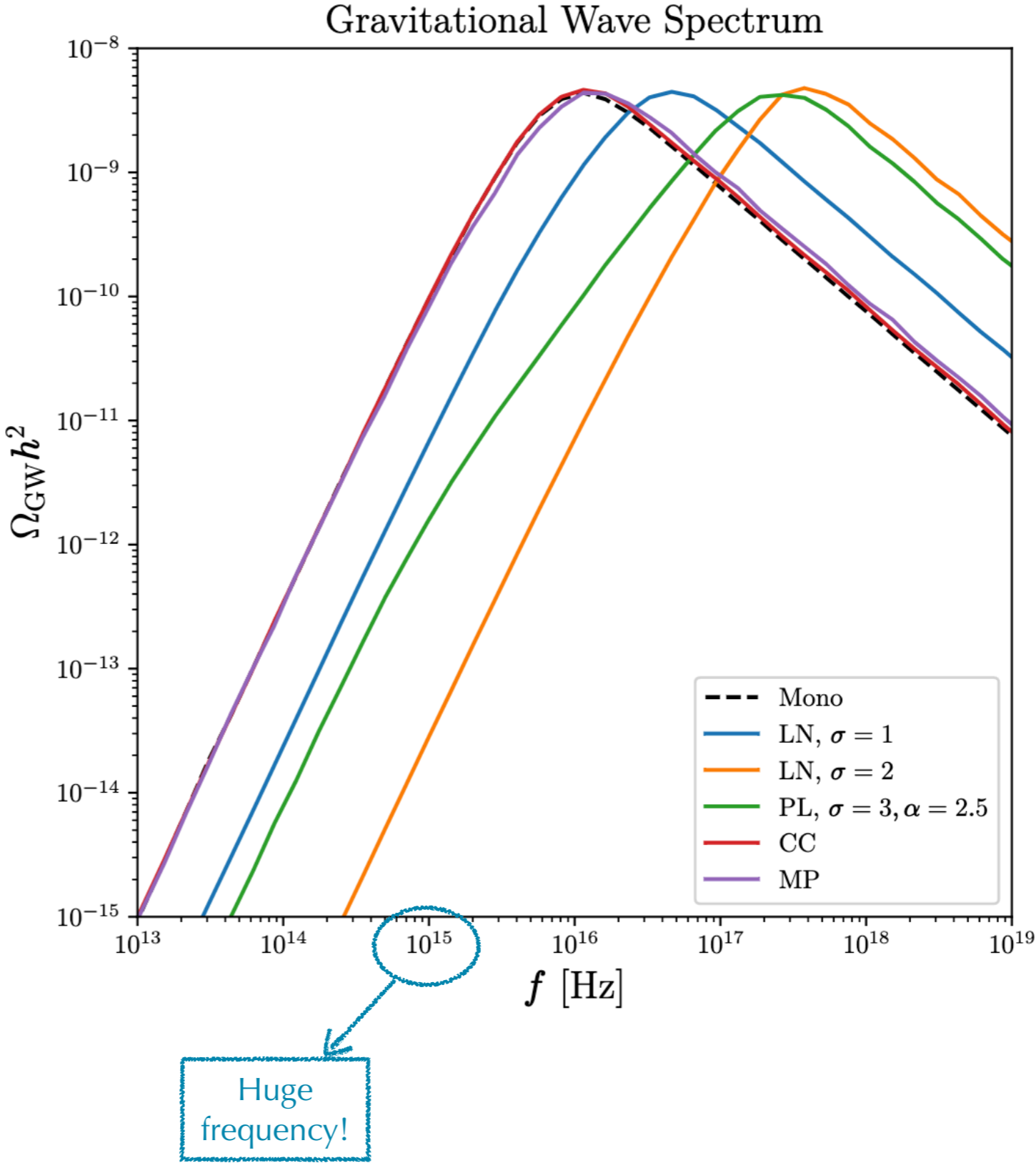
How to prove that the early Universe had a PBH-dominated era?

Example: GW produced by the evaporation

Arvanitaki, Geraci, 2013

Ito et al, 2020

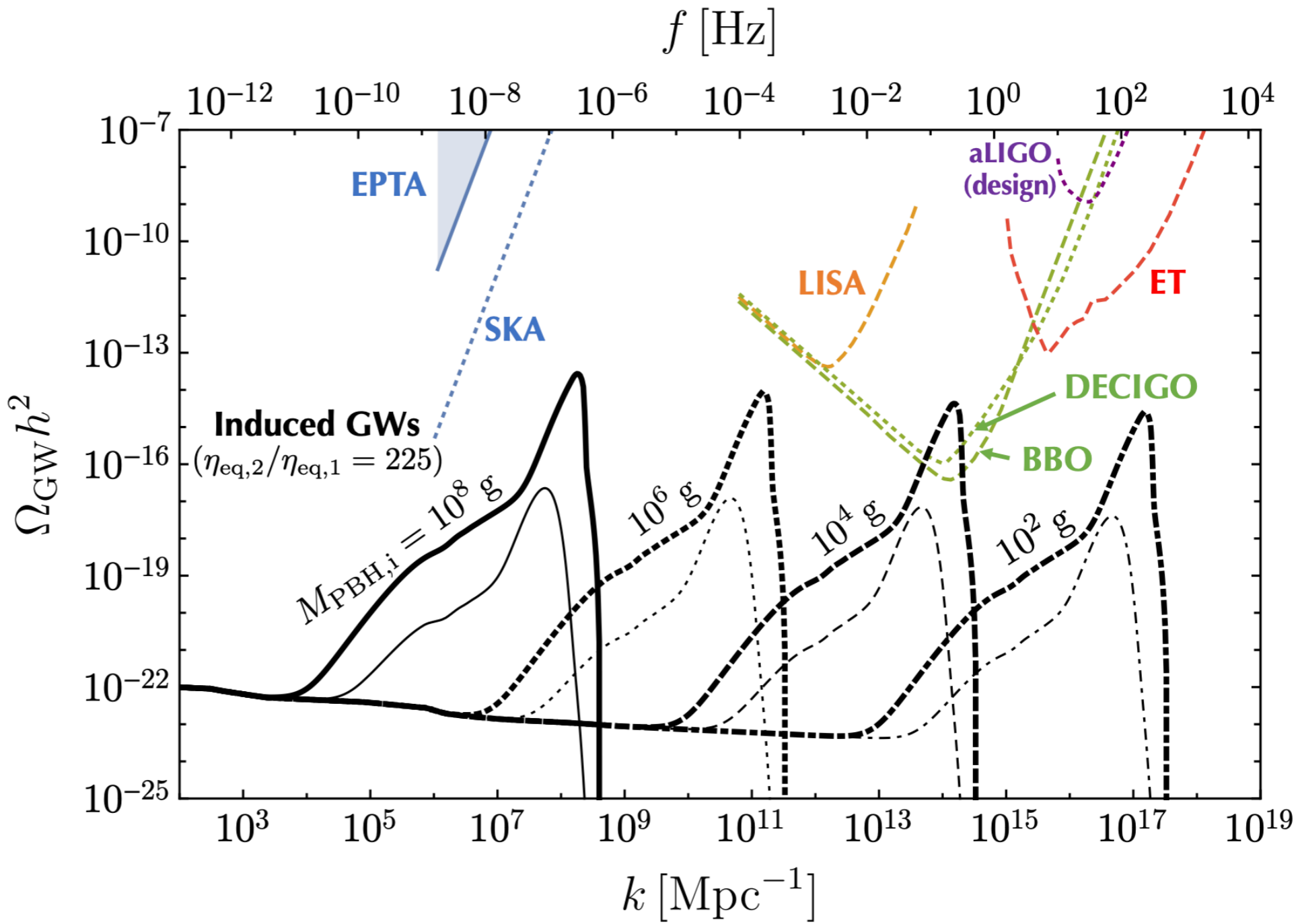
Chen et al, 2020



Gravitational Waves?

How to prove that the early Universe had a PBH-dominated era?

Poltergeist Mechanism



GWs are induced after the sudden transition from a matter to radiation dominated U

Inomata et al, 2003.10455

See also: 2205.06260

Third Act: High energy Neutrinos and Evaporating PBHs

Based on:
YFPG, 230X.XXXXX

Evaporating PBHs (EPBH)

- ❖ Perhaps some PBHs are evaporating today
- ❖ If this occurs close to Earth, we could see γ, ν 's, e^\pm
- ❖ Test BSM??
- ❖ How to measure PBH properties during evaporation?
- ❖ Kerr EPBHs \rightarrow Incorrect assumptions on Hawking spectrum

$$M_{\text{BH}}^{\text{in}} \sim 10^{15} \text{ g}$$

Baker, Thamm 2105.10506,
[2210.02805](#)

Capanema et al, [2110.05637](#)
Calzà, Rosa, [2210.06500](#)

Photons dominate the measurement

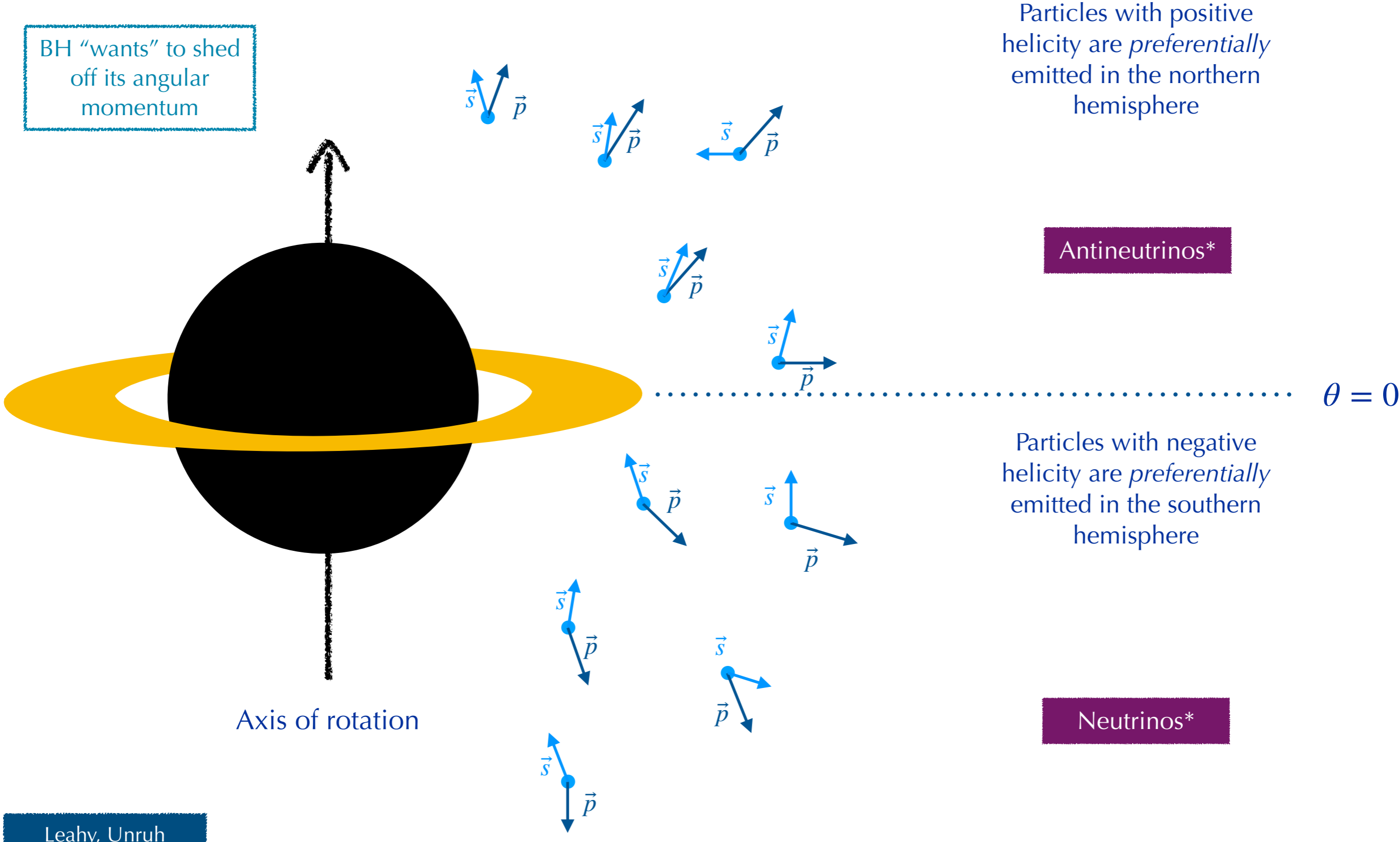
Anything to learn from neutrinos?



Parity Violation!!

How does this manifest in Hawking evaporation?

Neutrino Emission Asymmetry



Leahy, Unruh
PRD 19(1979)3509

*in the ultrarelativistic limit

$\theta \rightarrow$ polar angle

Neutrino Emission Asymmetry

BH "wants" to shed off its angular momentum

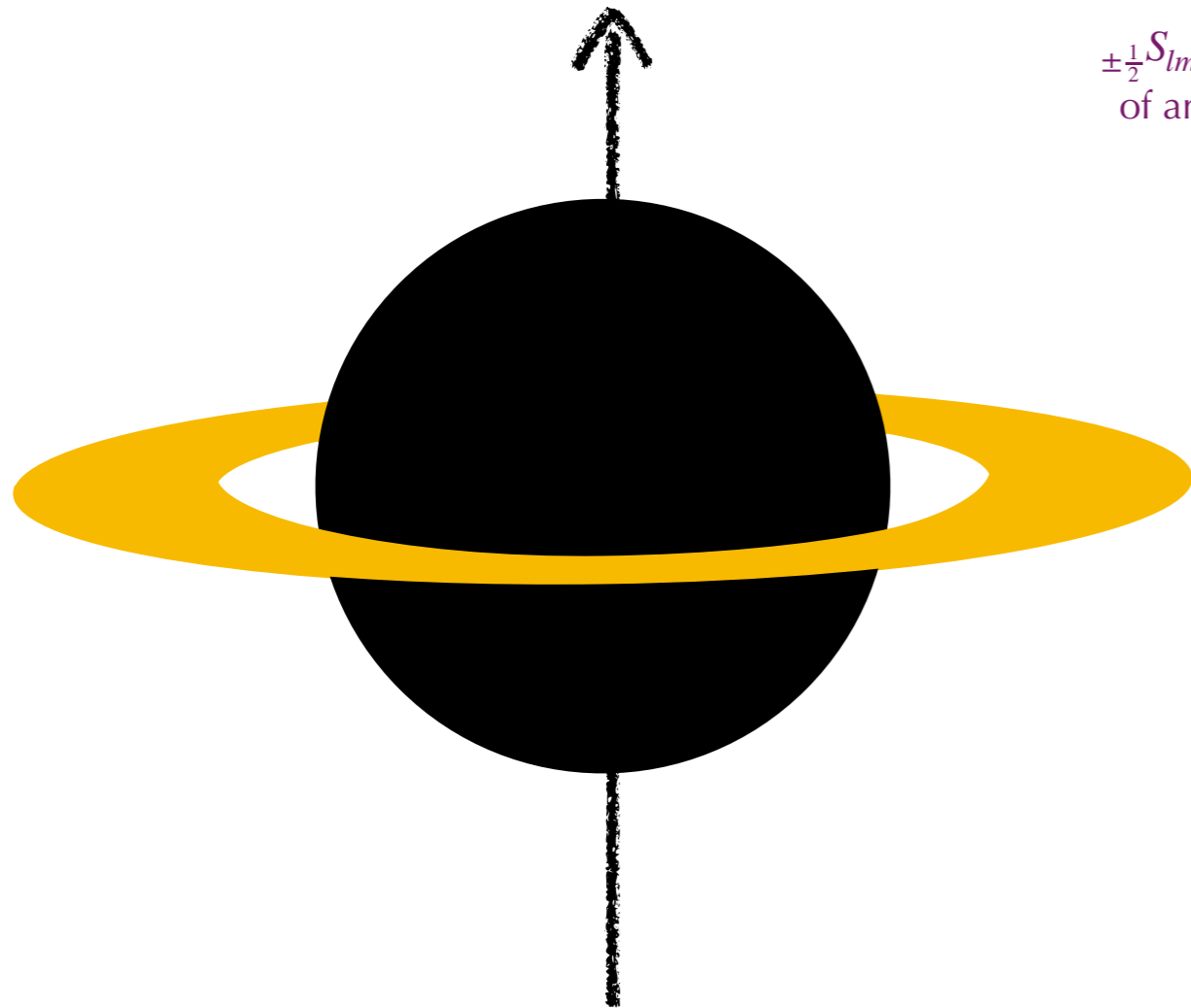
$$\mathcal{A} \equiv N_\nu - N_{\bar{\nu}}$$

$$\frac{d^3 \mathcal{A}}{d\omega dt d\Omega} = \frac{1}{4\pi} \sum_{l=1/2} \sum_{m=-l}^l \frac{s \Gamma_{lm}}{\exp(\varpi/T) + 1} \left\{ \left| S_{lm}(\theta) \right|^2 - \left| S_{lm}(\theta) \right|^2 \right\}$$

$\pm \frac{1}{2} S_{lm}(\theta) \rightarrow$ solutions of angular equation

"Neutrinos"

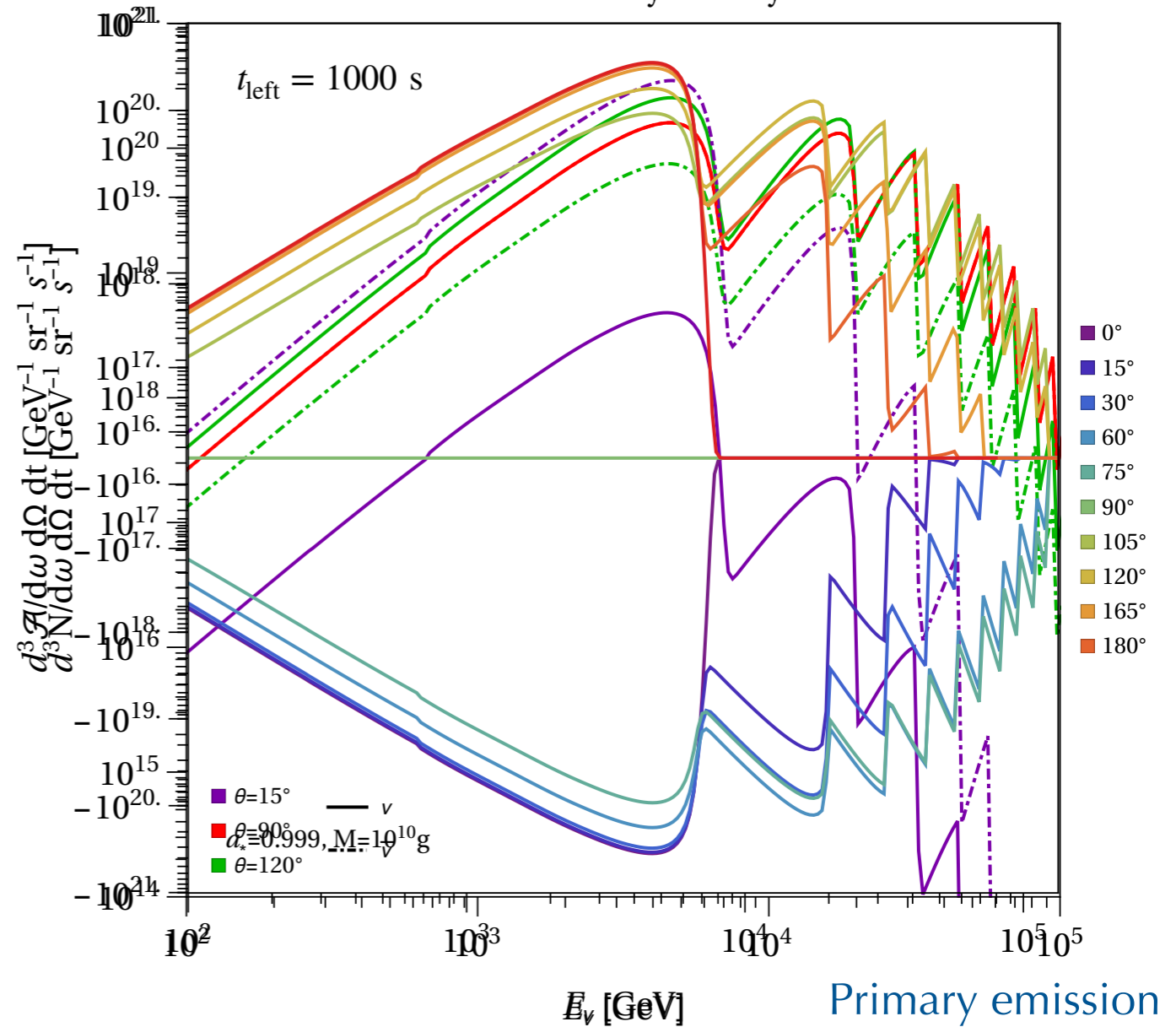
"Antineutrinos"



Axis of rotation

Could neutrinos tell us the spin of a PBH?

Neutrino Emission Asymmetry Fluxes



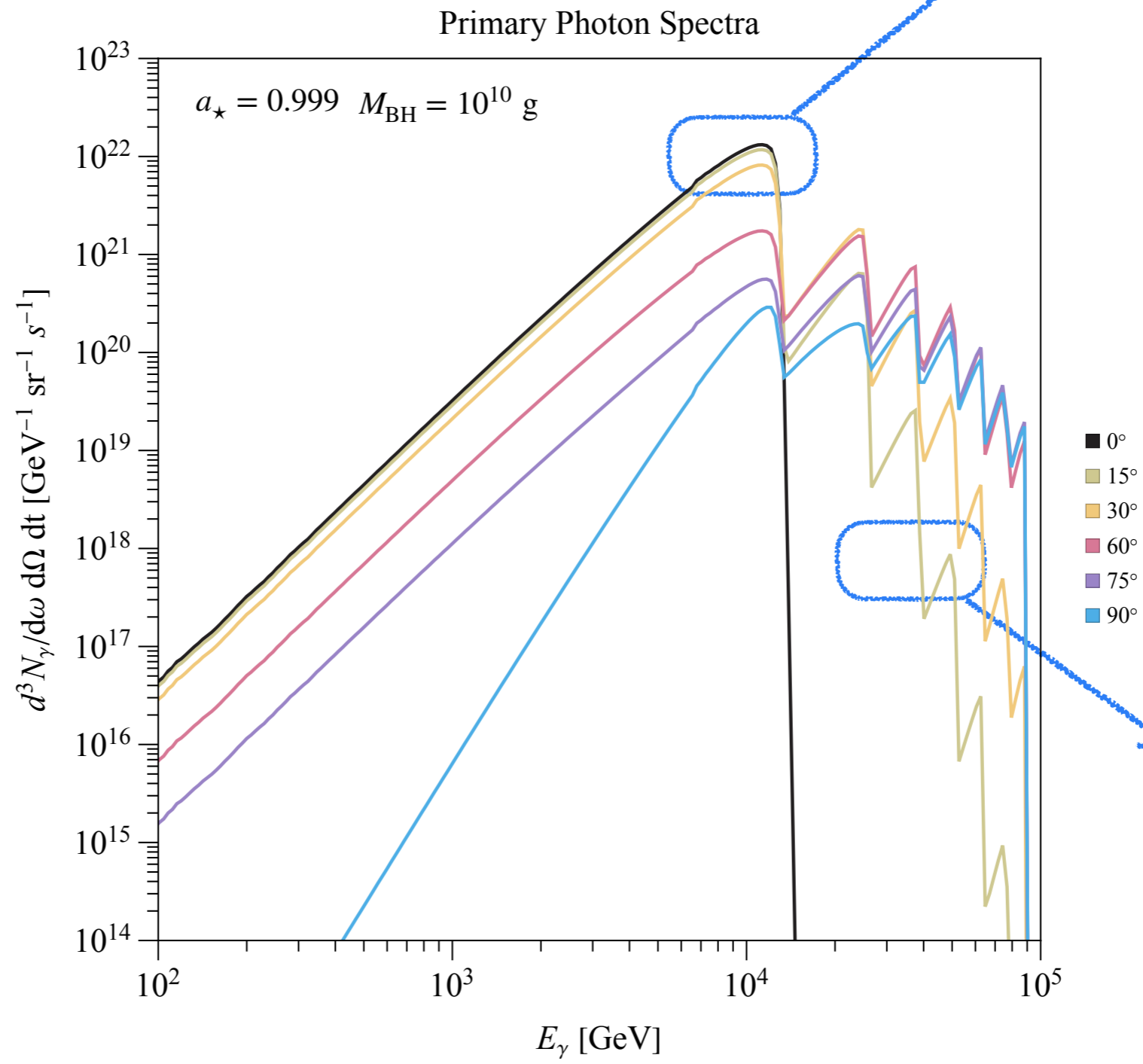
Photons?

BH "wants" to shed off its angular momentum

Emission of higher spin particles is enhanced

In the poles only $l = 1$ contribute

Also dependent on the polar angle

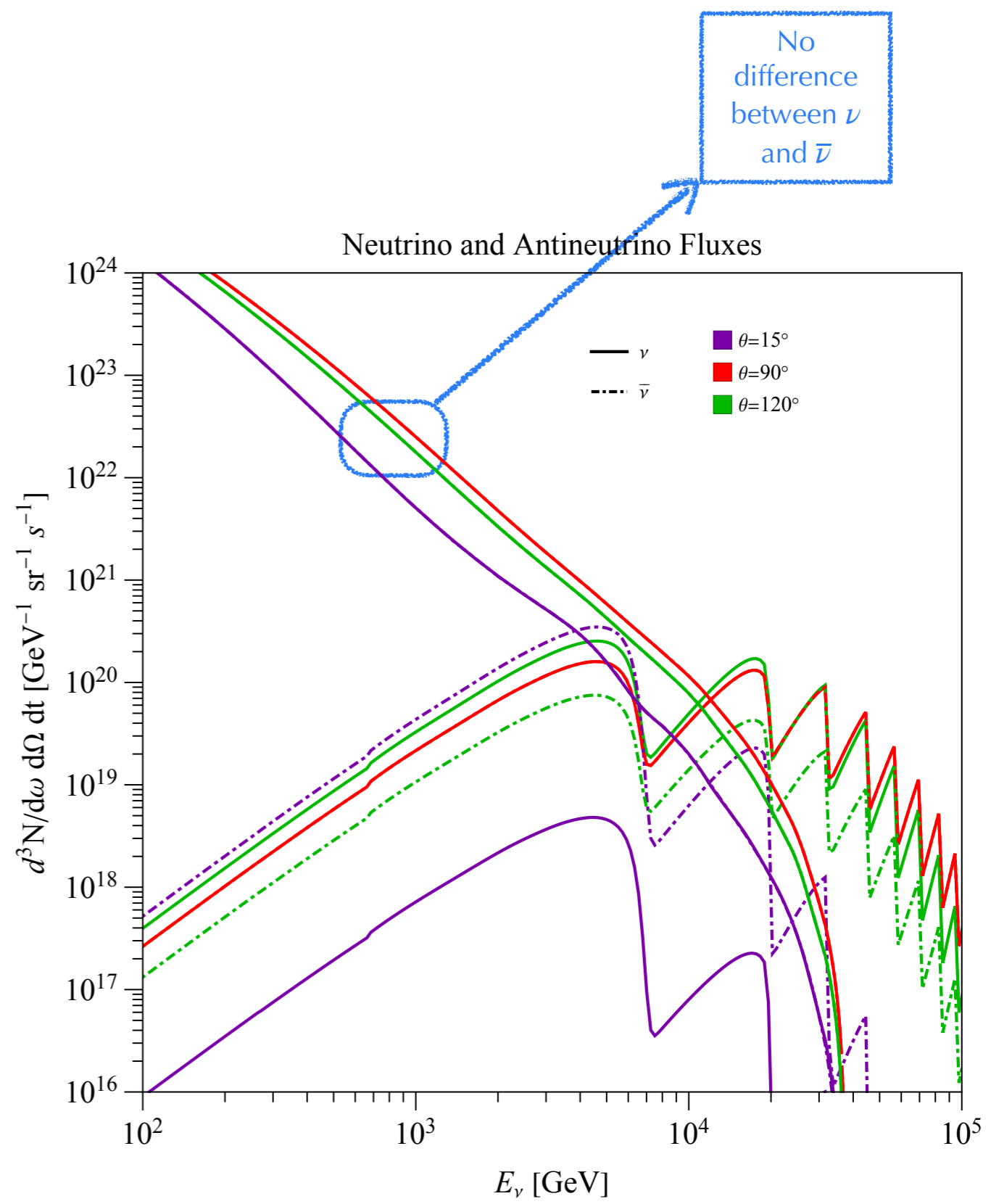
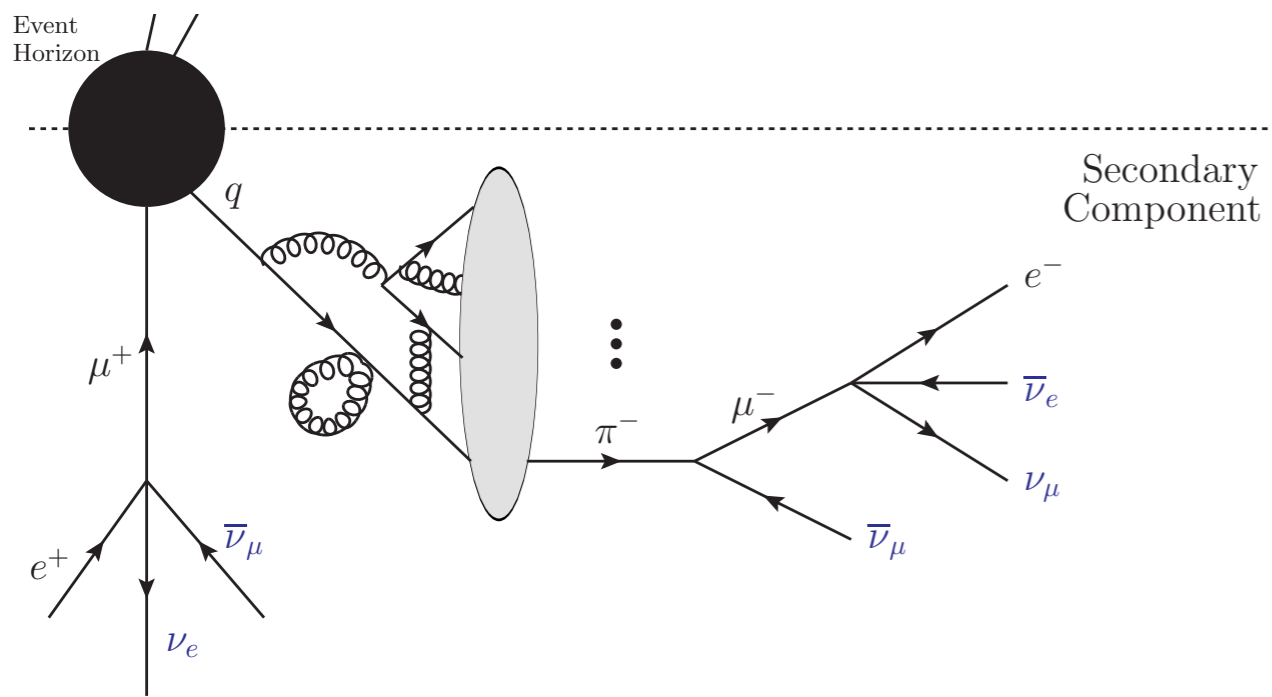


Modes start to contribute for larger θ

Secondaries?

$$\frac{d^3 N_{\nu(\gamma)}^{\text{sec}}}{d\omega dt d\Omega} = \int_0^\infty d\omega' \int d\Omega' \sum_i \frac{d^3 N_i}{d\omega' dt d\Omega'} \frac{d^2 n_{i \rightarrow \nu(\gamma)}}{d\omega d\Omega}$$

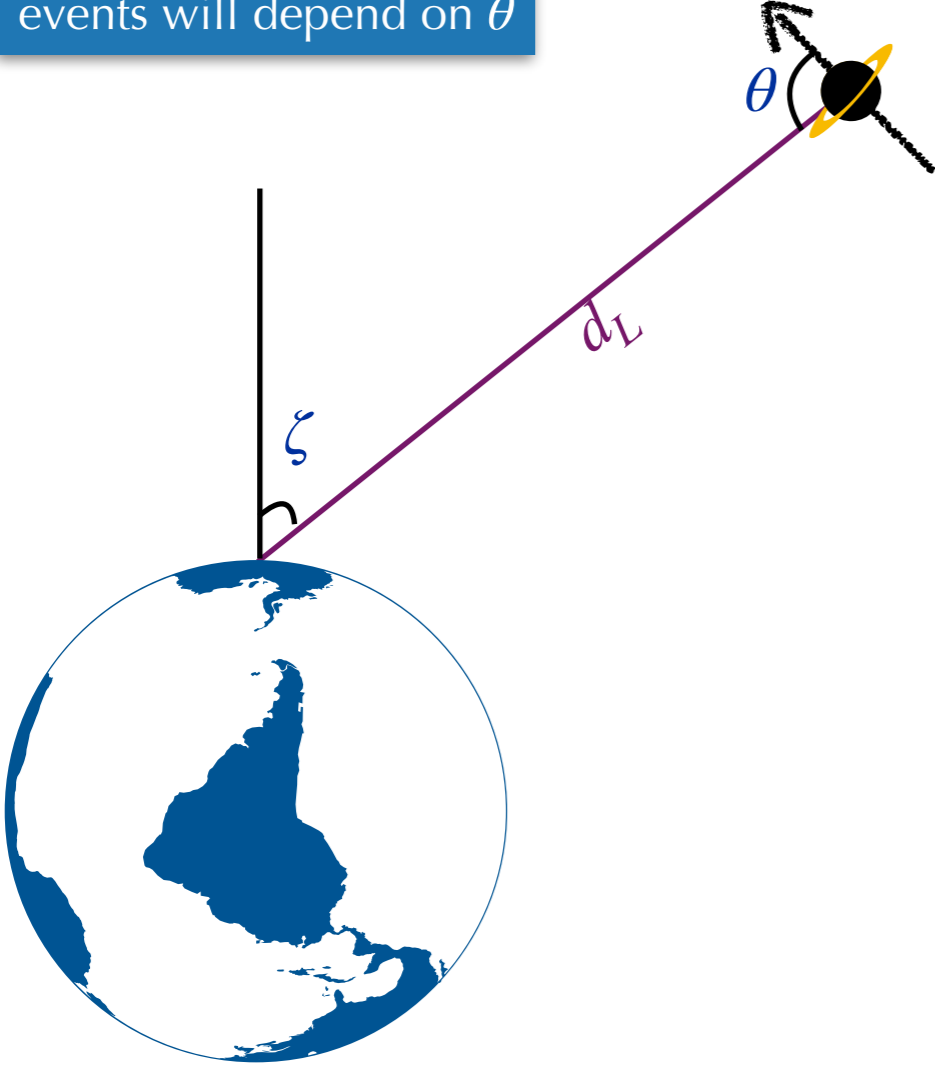
$\frac{d^2 n_{i \rightarrow \nu(\gamma)}}{d\omega d\Omega} \rightarrow$ energy and angular distribution of neutrinos/photons \rightarrow Pythia



Determining the angular momentum

Previous works ignored the dependence on θ

Neutrino - antineutrino events will depend on θ

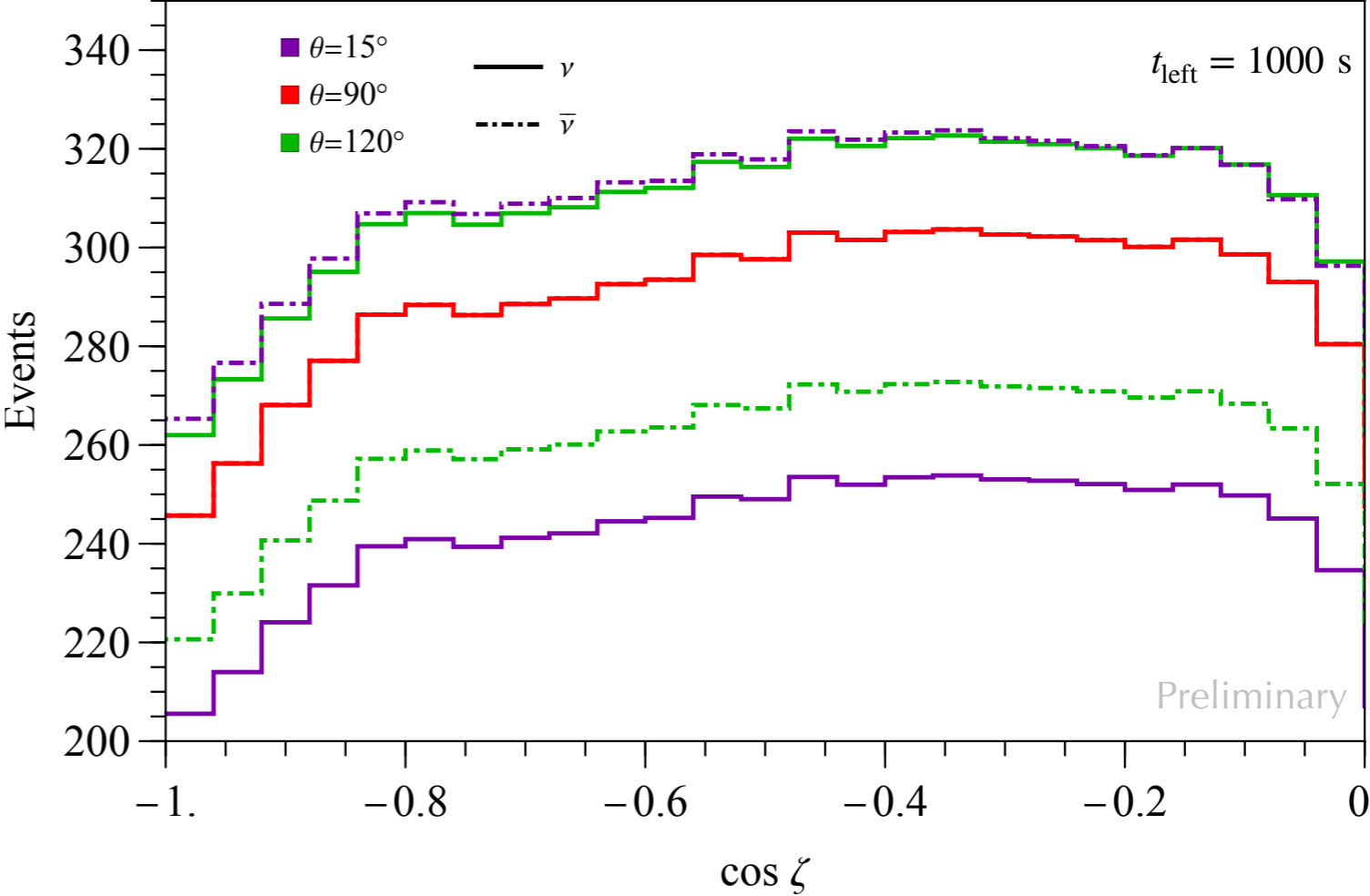


$$N_X(\theta) = \frac{1}{d_L^2} \int_{\omega_{\min}}^{\omega_{\max}} \int_0^{\tau} dt \frac{d^3 N_X}{d\omega dt d\Omega} A_{\text{eff}}(\omega, \zeta) d\omega$$

IceCube/HAWK effective area

$X = \nu_\mu, \bar{\nu}_\mu, \gamma$

$a_* = 0.999, d_L = 2 \times 10^{-4} \text{ pc}$



$d_L = 2 \times 10^{-4} \text{ pc} \approx$



Pale blue dot

Conclusions

- ❖ PBH evaporation offers a unique mechanism to produce particles in the Early Universe
- ❖ The effects are threefold:
 - Universal particle emission
 - Modifying the Cosmological Background
 - Entropy dilution
- ❖ We explored the effects on leptogenesis assuming the existence of a PBH population
- ❖ Future directions:
 - Relating to “more realistic” PBH formation mechanisms (connected to PBH-DM?)
 - Low scale leptogenesis? Sphalerons around PBHs after EWPT?
 - Kerr PBH \longrightarrow Additional interesting properties!
 - Anything else to learn by measuring neutrinos — antineutrinos in IC for and EPBH?

Thank you!

