

Lepton Number Violation and the Baryon Asymmetry of the Universe

- HEINRICH PÄS -



tu
dortmund

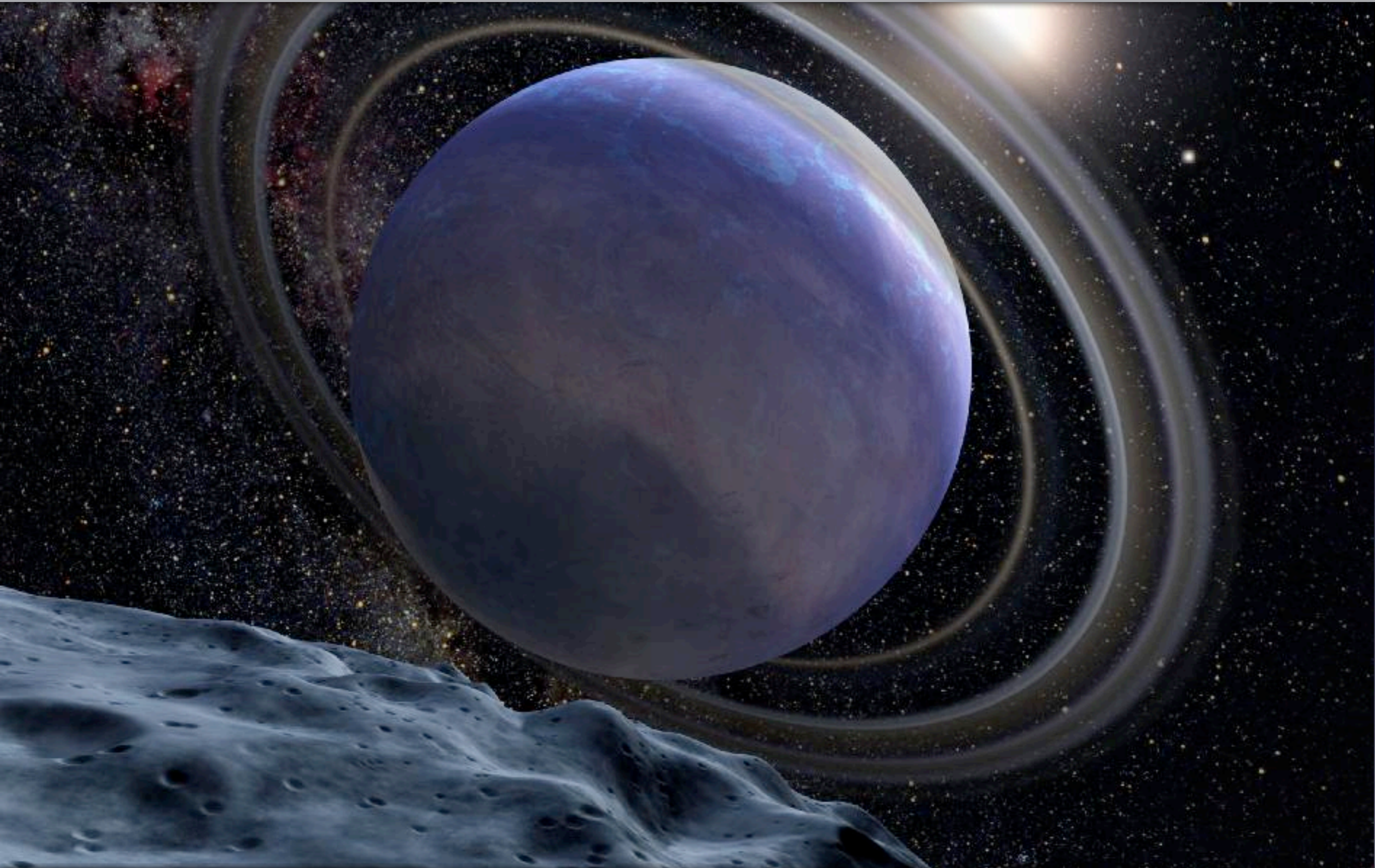
*MPI Seminar
Max-Planck-Institut für Kernphysik Heidelberg
December 14, 2015*

Outline

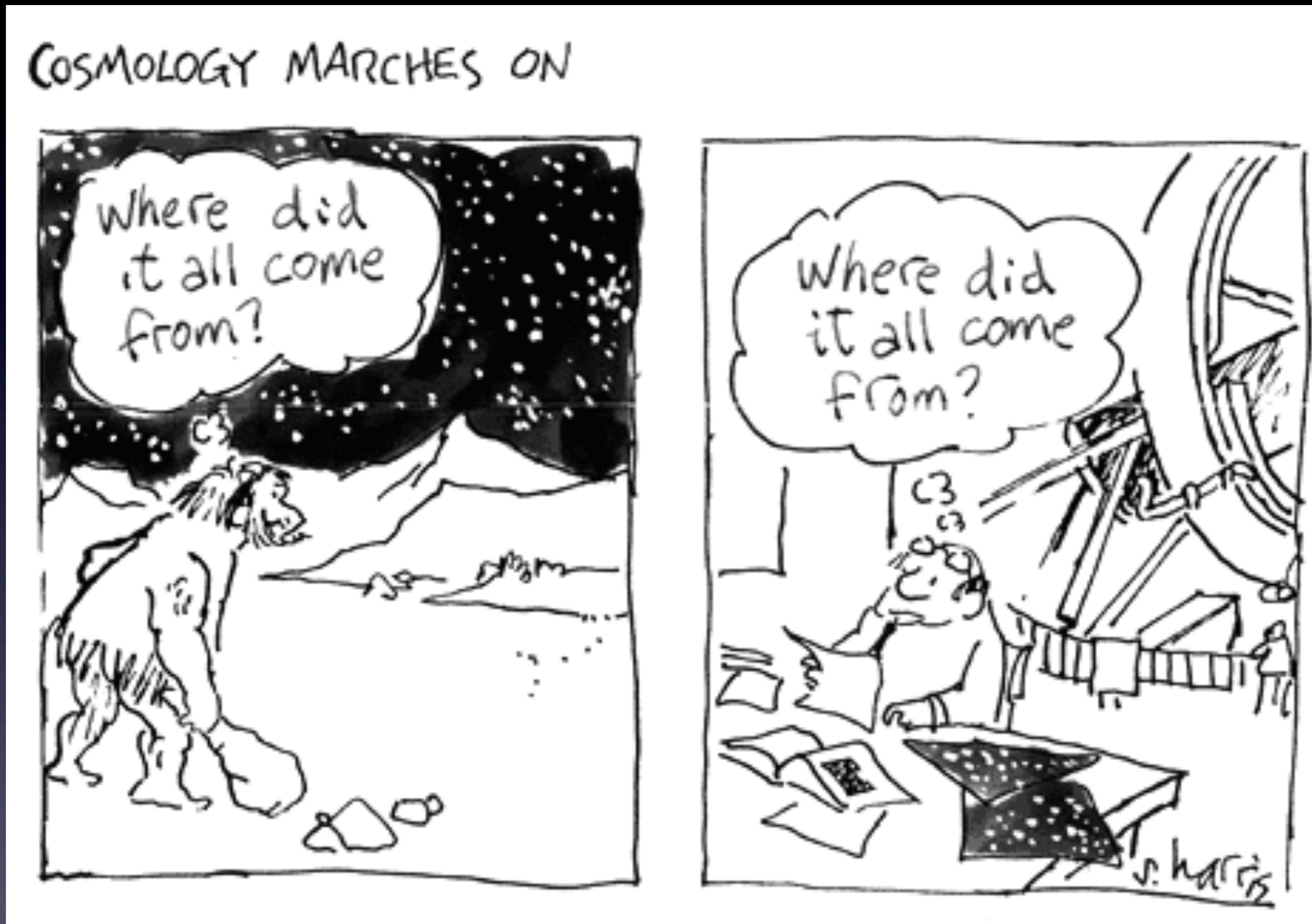
- ▶ The Baryon Asymmetry of the Universe
- ▶ Lepton Number Violation (LNV)
- ▶ Double Beta Decay
- ▶ LNV at the LHC
- ▶ Baryon Number Washout
- ▶ Conclusions

The Baryon Asymmetry of the Universe

Why is there something and not nothing?



The progress of cosmology in the last 30,000 years



Evidence for a Baryon Asymmetry

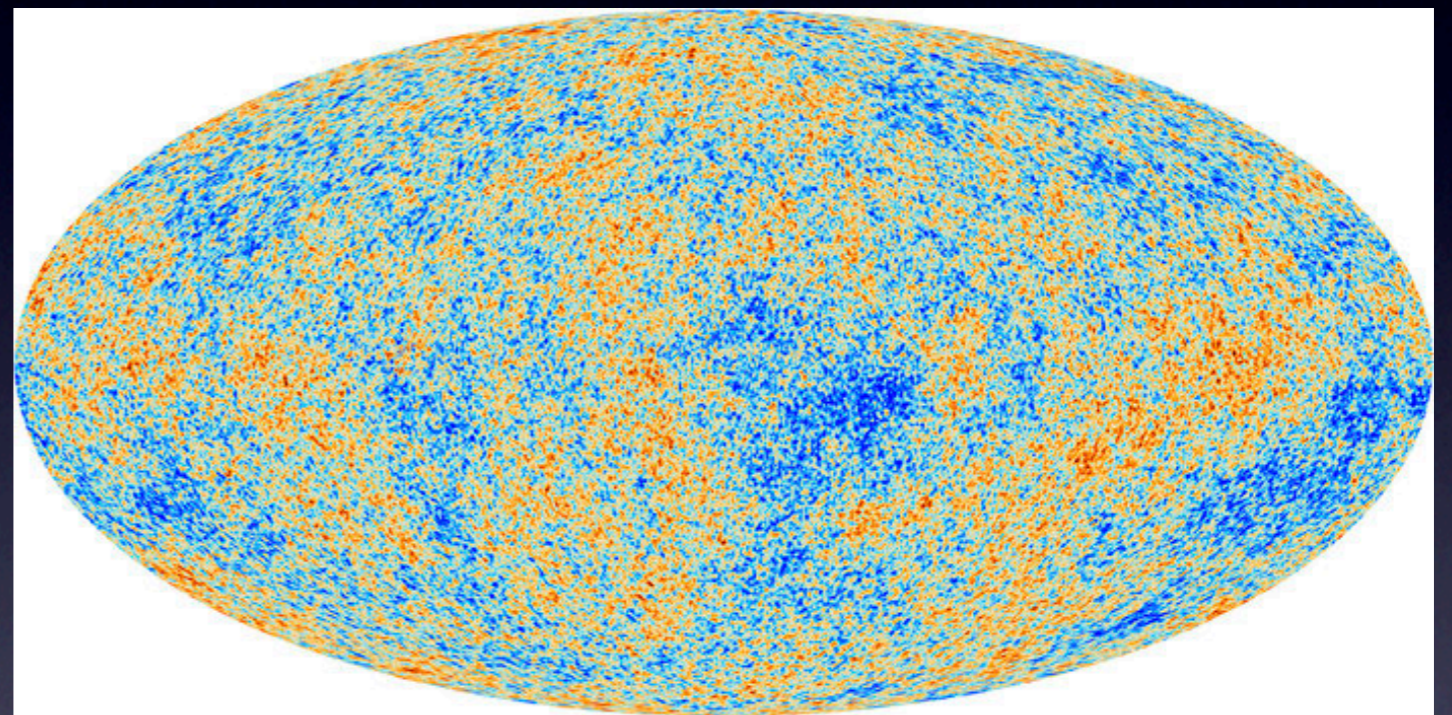
- ▶ Spectrum of anti-protons in cosmic radiation (BESS/ balloon in 35 km height) **consistent with production** from cosmic primaries



- ▶ **No anti-helium found** (AMS spectrometers on space shuttle discovery and International Space Station ISS)

Evidence for a Baryon Asymmetry

- ▶ **No traces of annihilation radiation** in the local galaxy cluster

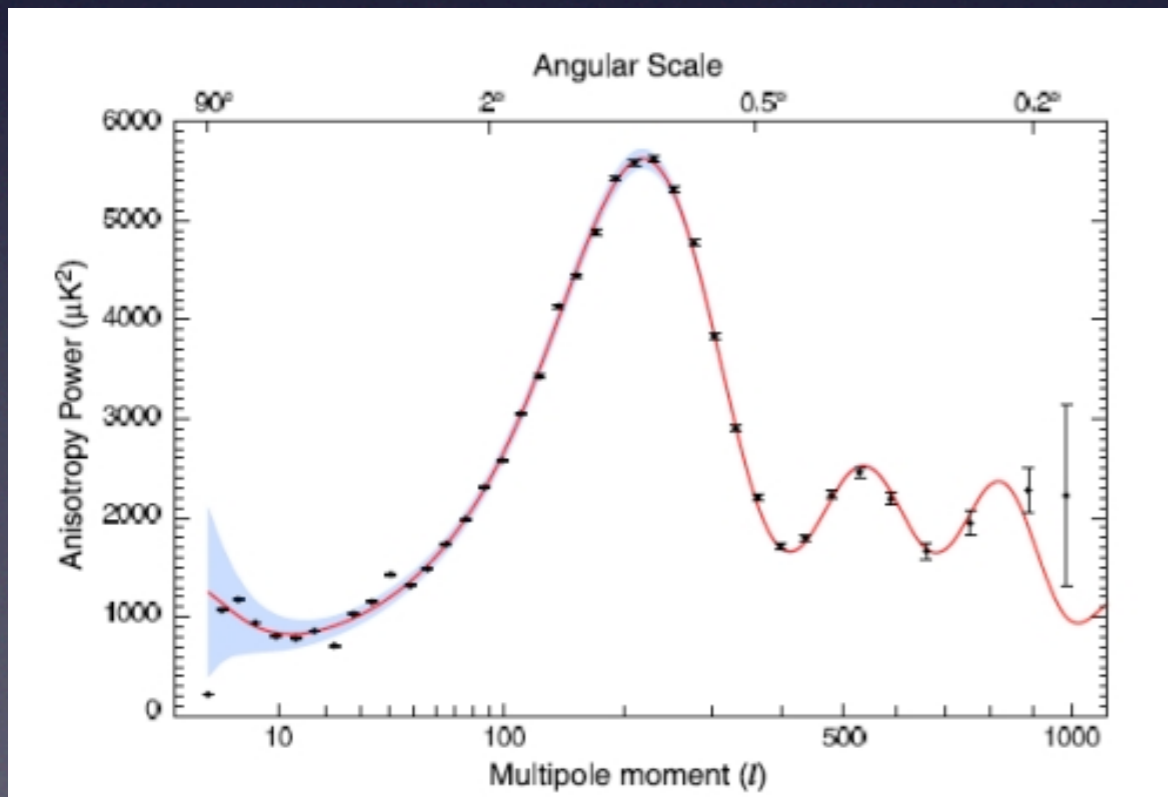
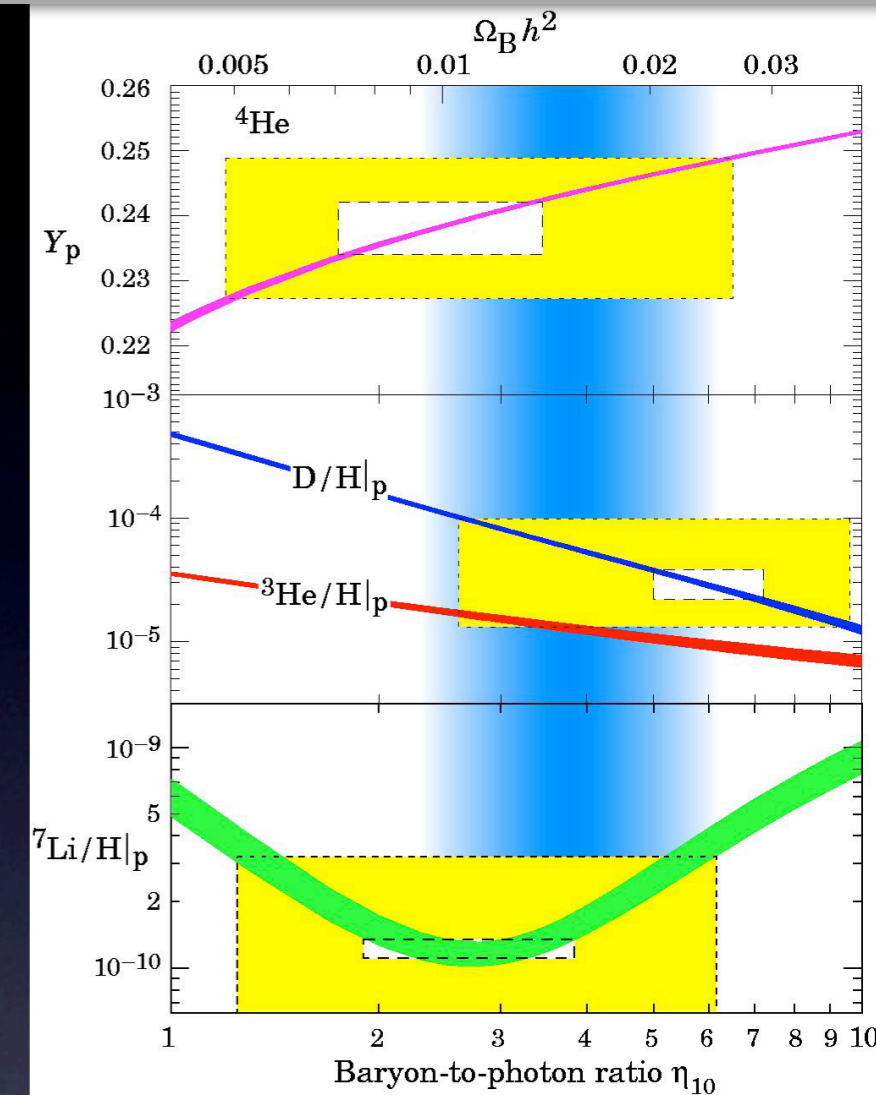


- ▶ **No distortion of the cosmic microwave background**

Size of the Baryon Asymmetry

- ▶ Big Bang Nucleosynthesis: synthesis of $p, n \rightarrow D, {}^3\text{He}, {}^4\text{He}, {}^7\text{Li}$ depends on baryon and radiation density

$$\eta_B = \frac{n_B - n_{\bar{B}}}{n_\gamma}$$

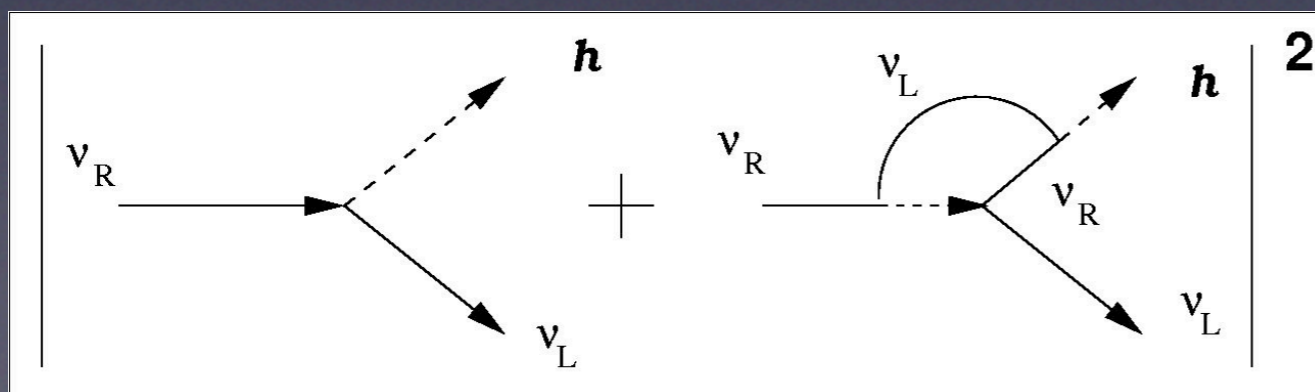


- ▶ CMB: comparing heights 1st peak (gravitation $\rightarrow\rightarrow$ gas) vs. 2nd peak (gravitation \leftrightarrow gas)

$$\eta_B^{\text{obs}} = (6.20 \pm 0.15) \times 10^{-10}$$

Baryogenesis

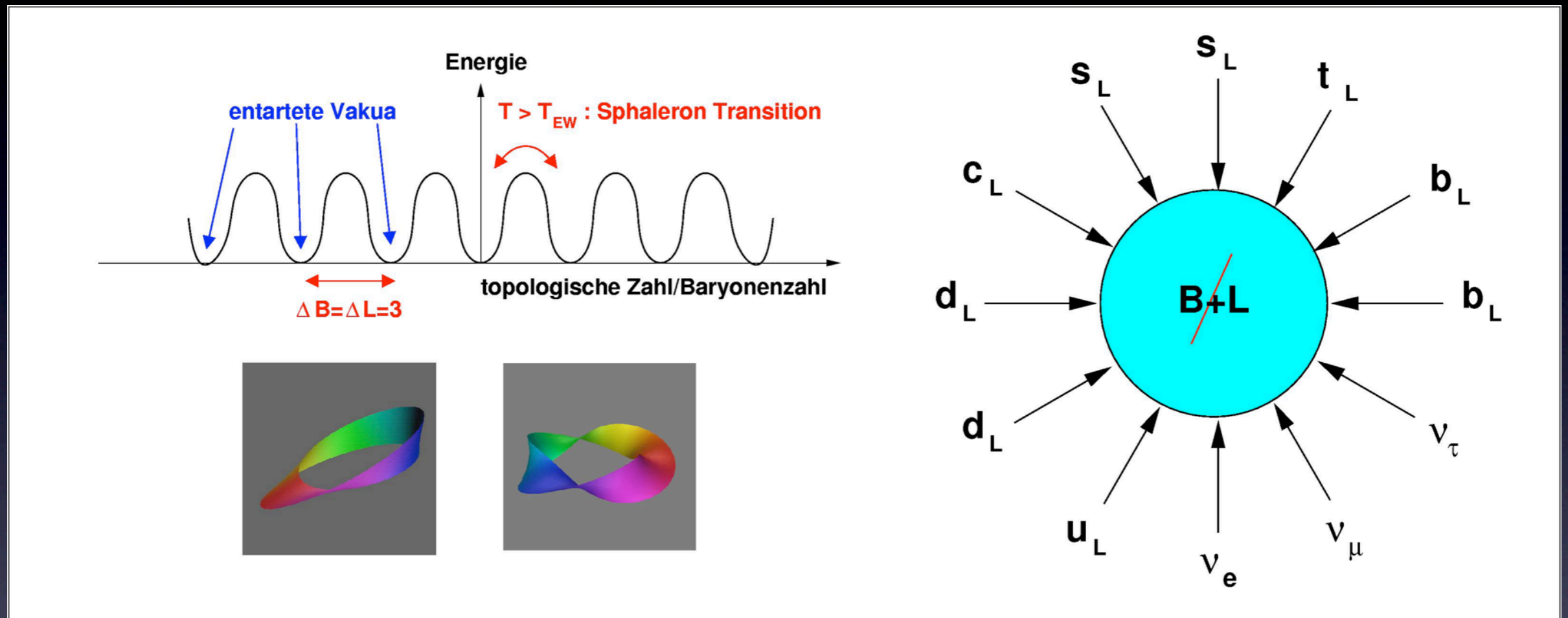
- ▶ Baryon Asymmetry as **initial condition**?
- ▶ **Cosmic inflation**: $a(t) \sim \exp(\Lambda/3)t$
 - Universe flat, homogenous and **empty**
 - necessity of Baryogenesis **after** reheating
- ▶ 3 **Sakharov conditions**:
 - Non-Equilibrium
 - C and CP violation
 - Baryon Number Violation
- ▶ Prominent example – **Leptogenesis**:
right-handed neutrino decay in early Universe



→ **Lepton
Asymmetry**
[Fukugita, Yanagida, 1986]

Sphalerons

Sphalerons: B-Violation within the Standard Model



Topologically different field configurations \Rightarrow
degenerate vacua with different baryon numbers [t'Hooft, 1976]

$T > T_{EW} \Rightarrow$ Transitions between vacua, $B+L$ Violation

[Kuzmin, Rubakov, Shaposhnikov, 1985]

Baryogenesis

- ▶ Electroweak baryogenesis via Sphalerons?
In the SM: EW phase transition not rapid enough (non-equilibrium), 2 competing vacua during phase transition necessary → depends on Higgs quartic coupling → only possible for Higgs masses < 70 GeV
- ▶ **Physics beyond SM needed!**
- ▶ Leptogenesis: B-L Asymmetry + B+L violating Sphalerons
⇒ B Asymmetry

Baryogenesis

Not simple:

“Das Nichts nichtet - the Nothing noths”

(Martin Heidegger)

BUT: This is **NO** talk on Baryogenesis !

So let's assume we somehow created a
Baryon Asymmetry and let's go on with
particle physics

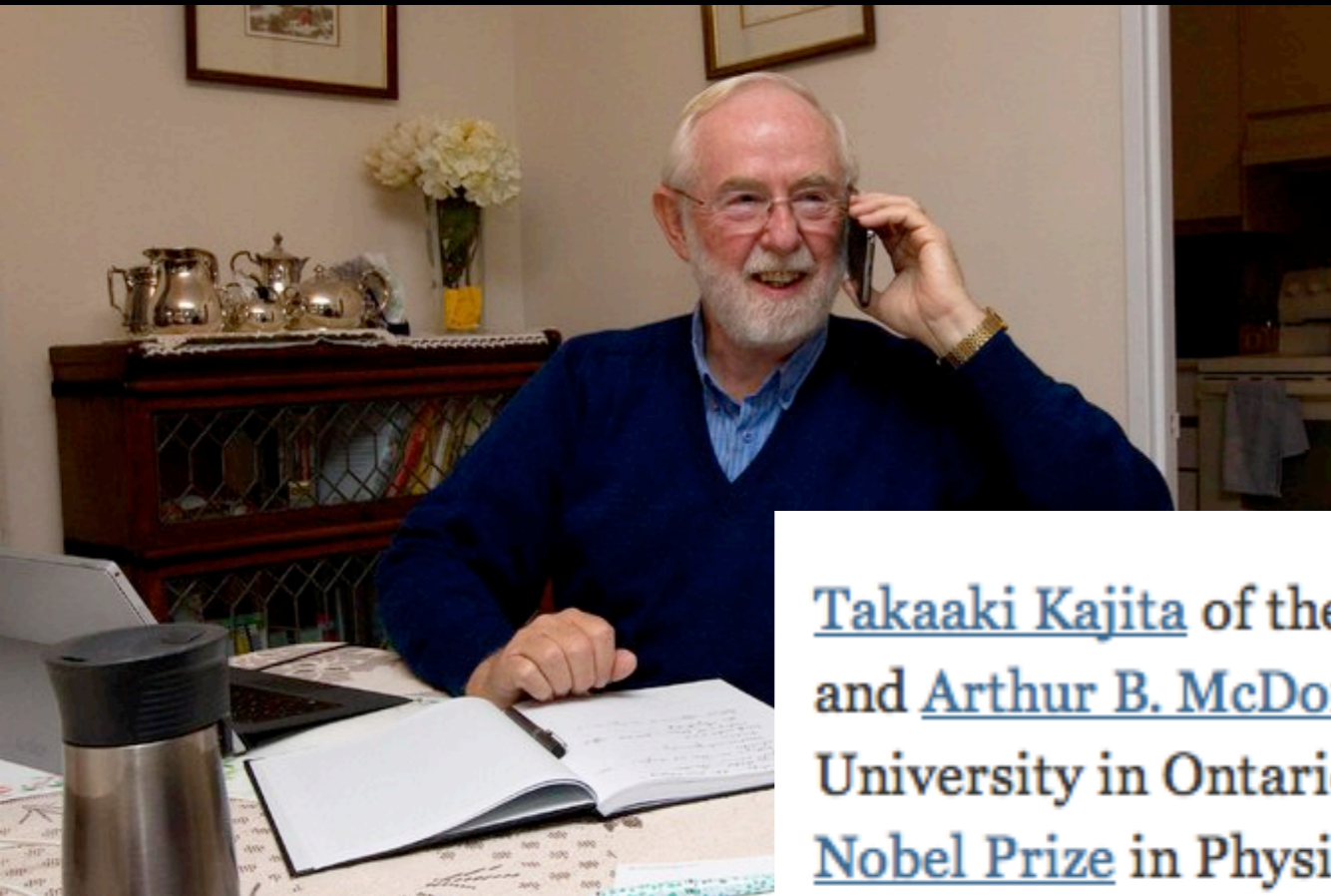
Lepton

Number

Violation

Why is Lepton Number Violation interesting ?

The New York Times



Takaaki Kajita of the University of Tokyo and Arthur B. McDonald of Queen's University in Ontario were awarded the Nobel Prize in Physics on Tuesday for discovering that the enigmatic subatomic particles known as neutrinos have mass.



“Uncharacteristically for a physics conference people gave the speaker a standing ovation. I stood up too. Having survived every experimental challenge since the late 1970s the Standard Model had finally fallen. The results showed that at the very least the theory is incomplete.”

*Hitoshi Murayama (UC Berkeley)
about the Neutrino-98-Konferenz*

Bullshit ?

Why is LNV interesting ?

Why is a non-zero ν mass physics beyond SM ?

EITHER - OR

$$m_D \overline{\nu}_L \nu_R$$

$$m_M \overline{\nu}_L^c \nu_L$$

$$m_M \overline{\nu}_R^c \nu_R$$

OR

~~$$m_M \overline{\nu}_R^c \nu_R$$~~

LNV

LNV

e.g.

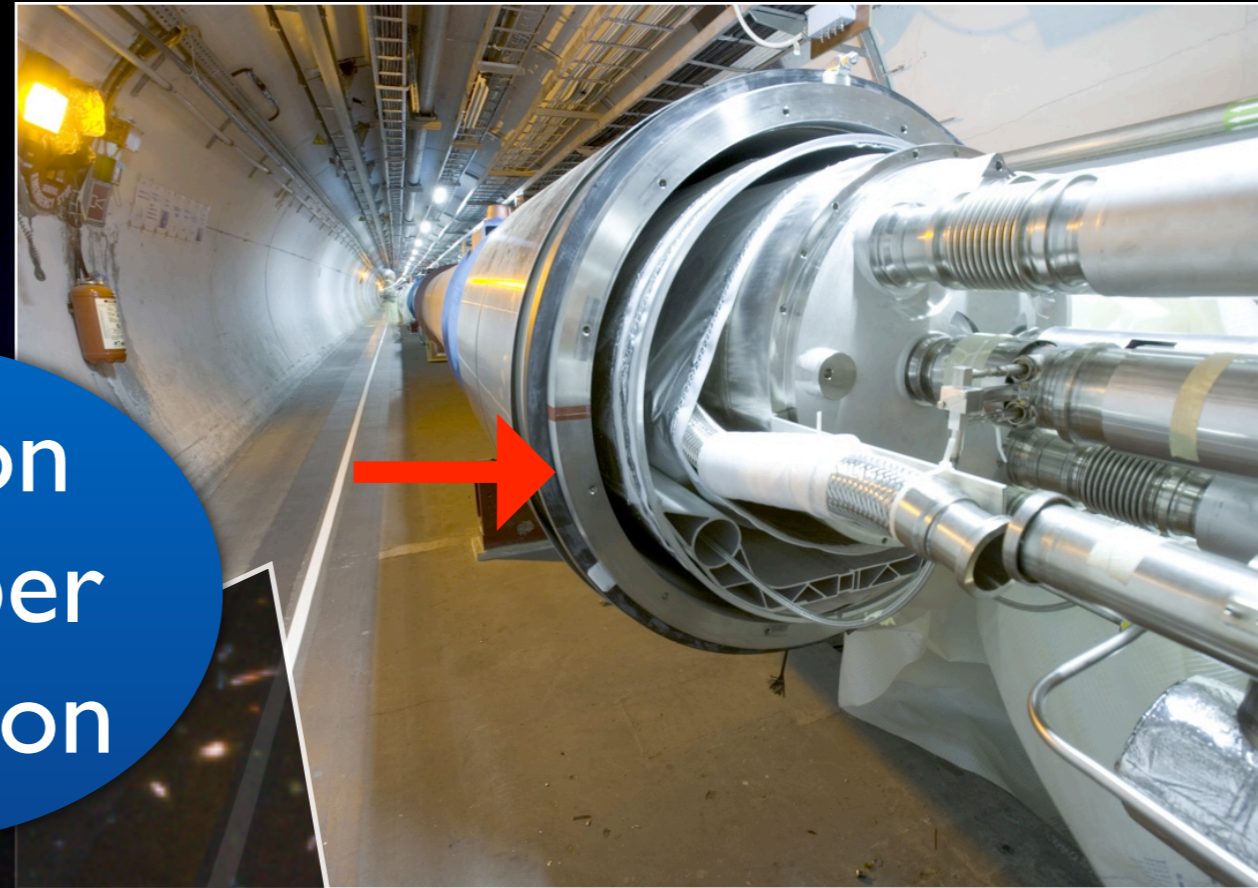
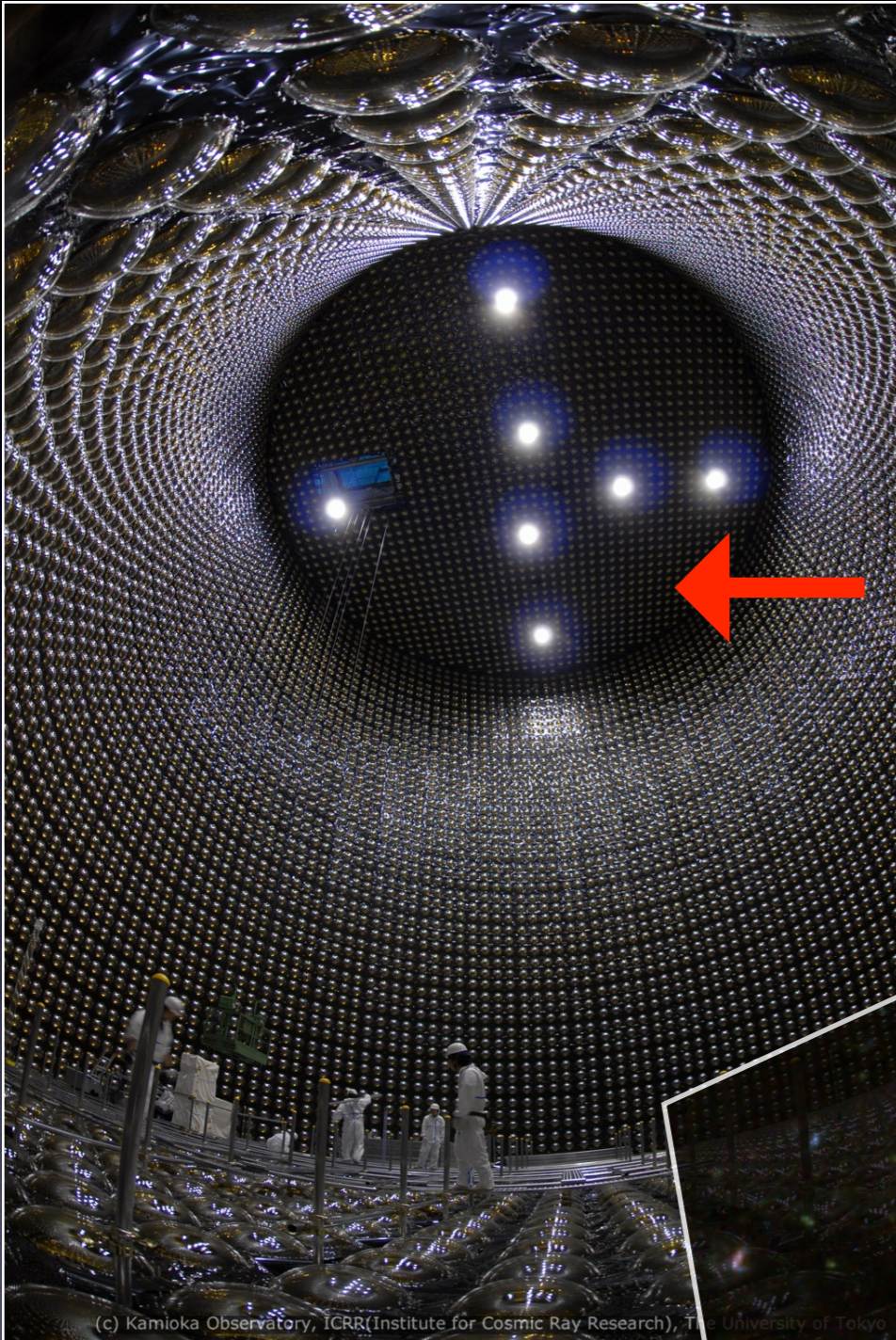
Seesaw-I

[Minkowski'77]

New Symmetry, e.g. Lepton Number

Lepton Number (Violation) is at the core of the link between ν mass & physics beyond SM!

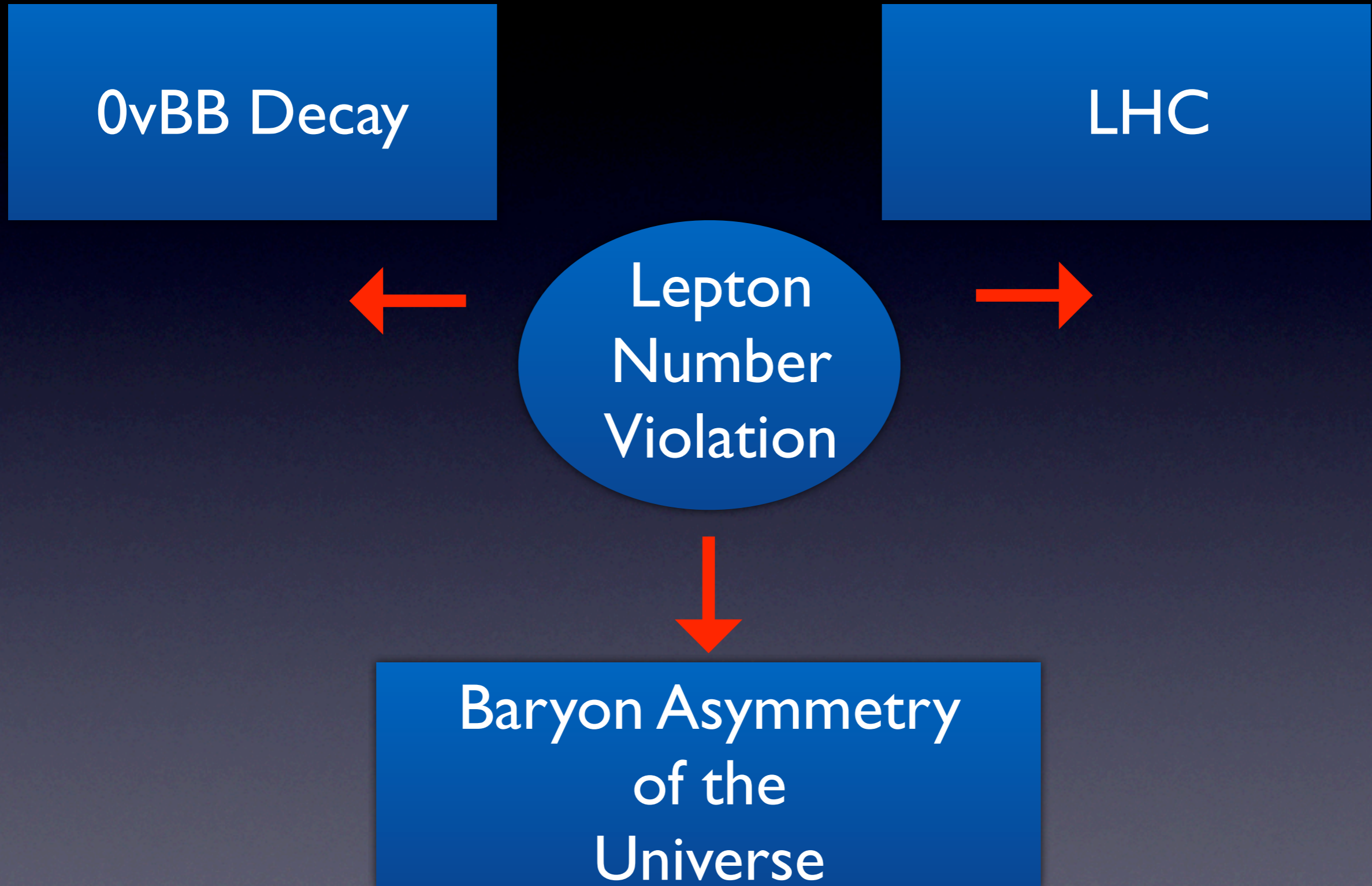
Key Message



Lepton
Number
Violation



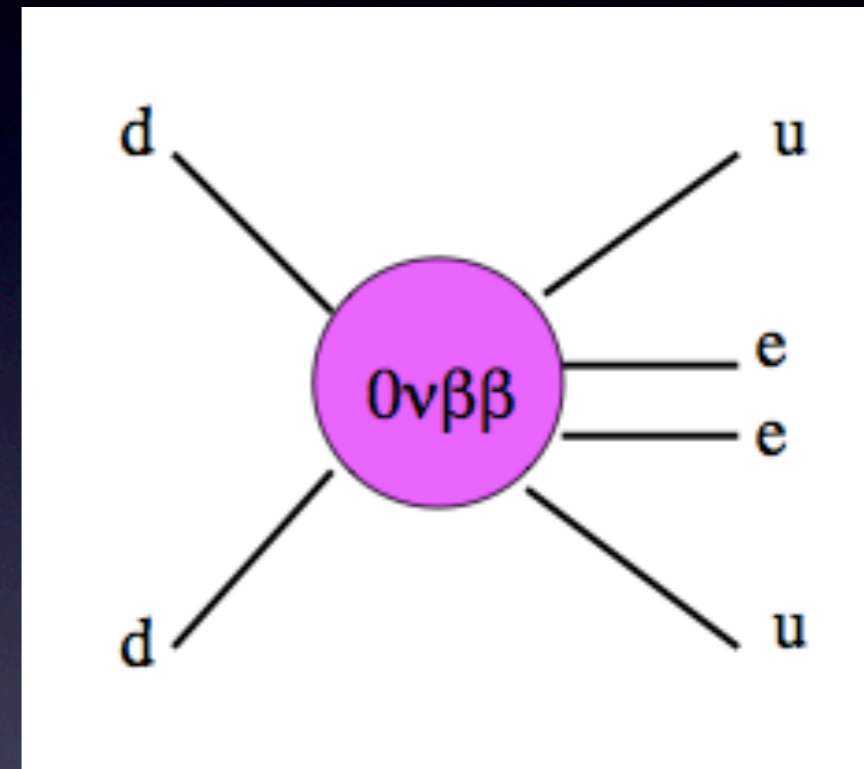
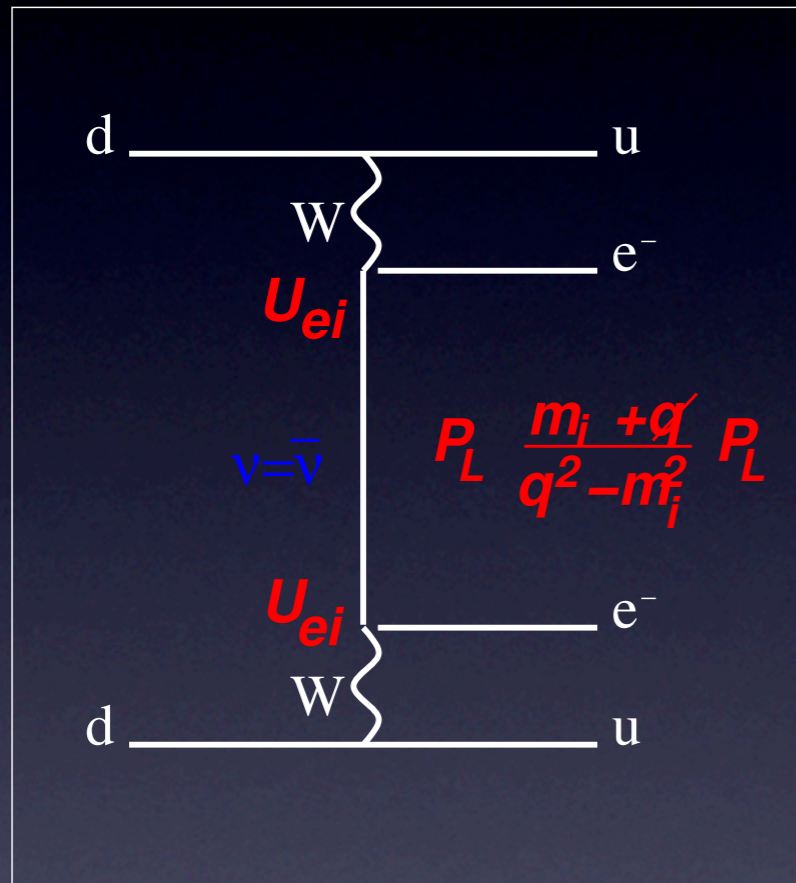
Key Message



Double Beta Decay

Probing LNV

Most prominent: $0\nu\beta\beta$ decay



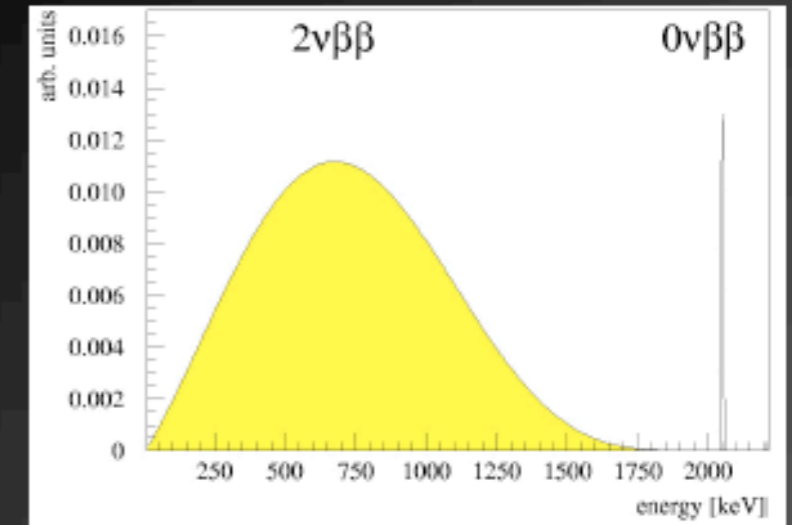
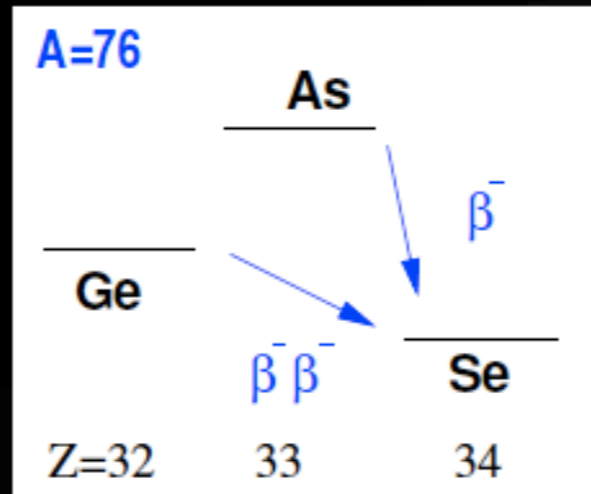
$\sim m_\nu$

“Mass Mechanism”

General Case

What is $0\nu\beta\beta$ decay?

$$2n \rightarrow 2p + 2e^-$$



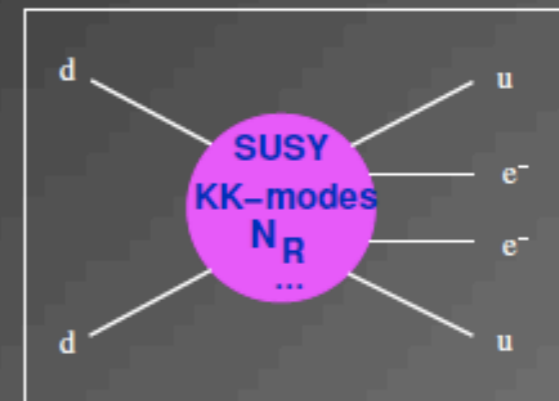
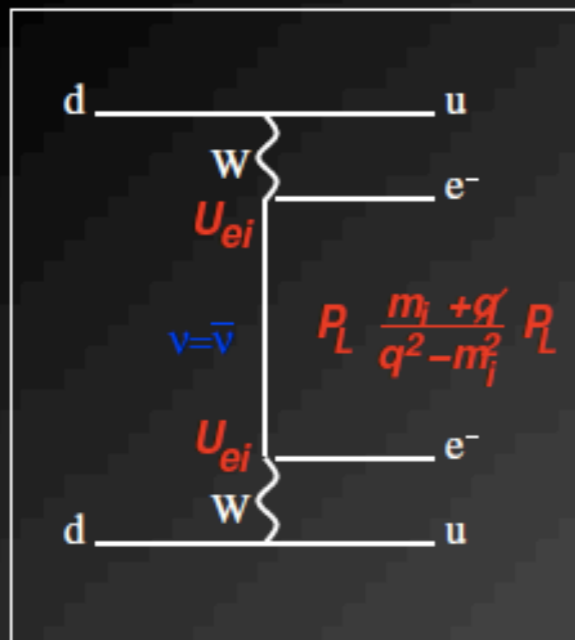
Mass mechanism:

$$[T_{1/2}^{0\nu}]^{-1} \propto \left| \sum_i U_{ei}^2 m_i \right|^2$$

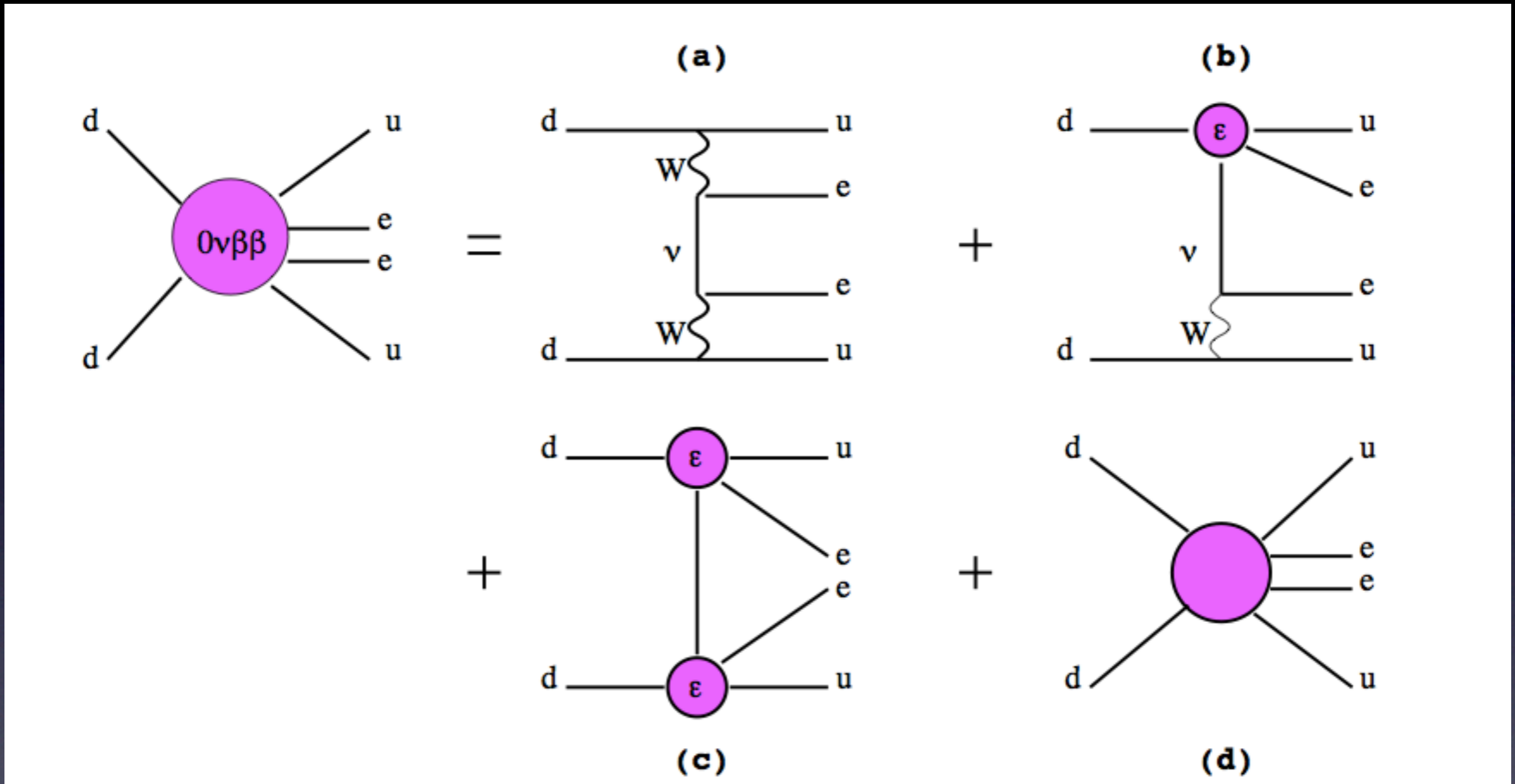
In general: Every operator

$$\bar{p} \bar{p} \bar{e} \bar{e} n n / M^5$$

will generate $0\nu\beta\beta$ decay

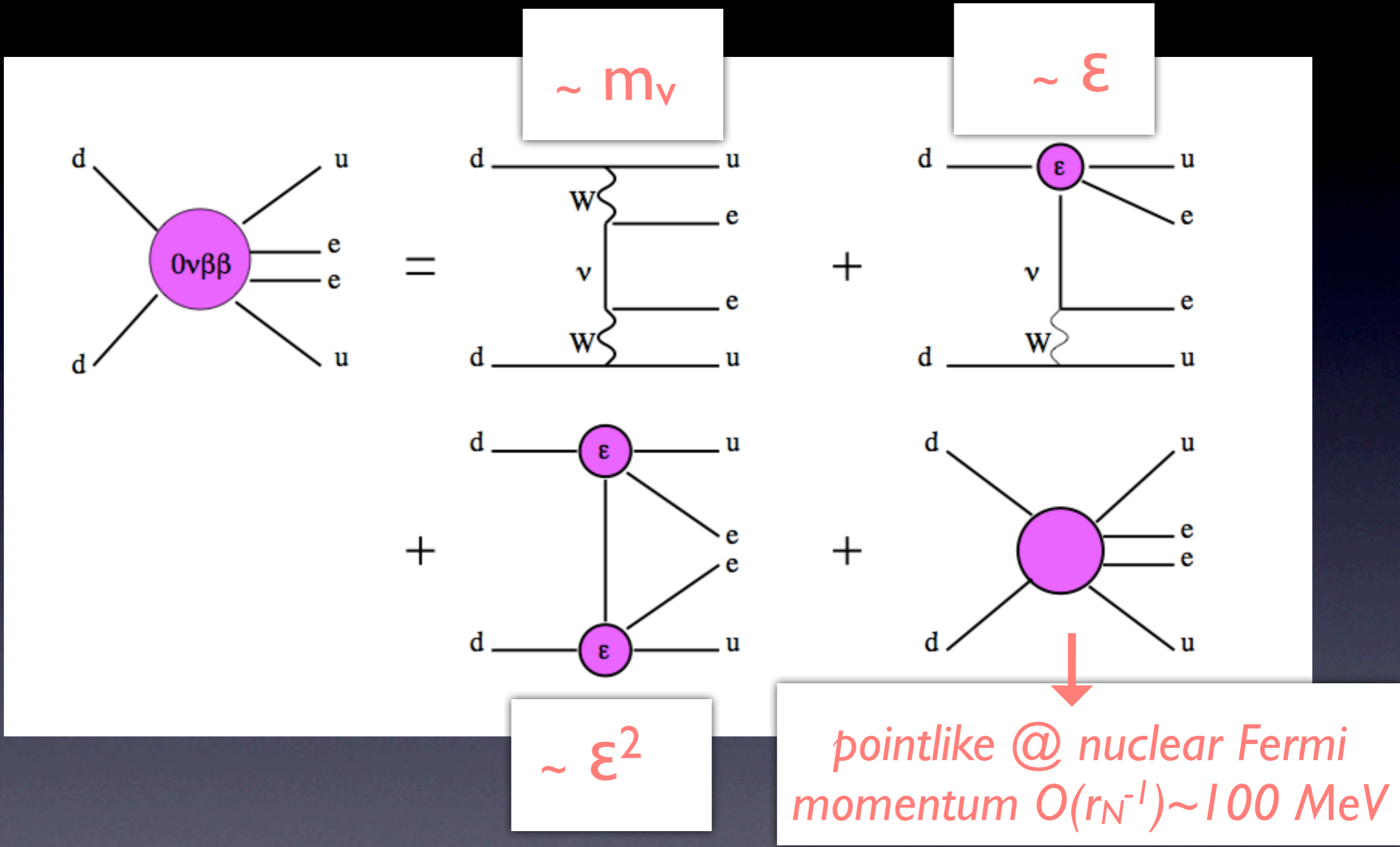


$0\nu\beta\beta$ Decay



[HP, Hirsch, Kovalenko, Klapdor-Kleingrothaus, PLB 1999 & 2001]

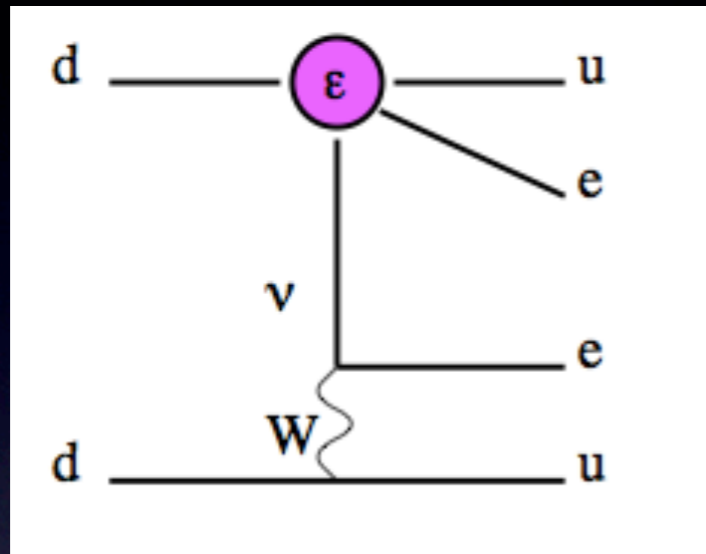
$0\nu\beta\beta$ Decay



[HP, Hirsch, Kovalenko, Klapdor-Kleingrothaus, PLB 1999 & 2001]

General parametrization for $0\nu\text{BB}$ Decay

Long-Range Part



ϵ_{V-A}^{V+A}	$4.4 \cdot 10^{-9}$
ϵ_{V+A}^{V+A}	$7.0 \cdot 10^{-7}$
ϵ_{S-P}^{S+P}	$1.1 \cdot 10^{-8}$
ϵ_{S+P}^{S+P}	$1.1 \cdot 10^{-8}$
ϵ_{TL}^{TR}	$6.4 \cdot 10^{-10}$
ϵ_{TR}^{TR}	$1.7 \cdot 10^{-9}$

$$\mathcal{L} = \frac{G_F}{\sqrt{2}} \left\{ j_{V-A}^\mu J_{V-A,\mu}^\dagger + \sum_{\alpha,\beta} \epsilon_\alpha^\beta j_\beta J_\alpha^\dagger \right\}$$

$$\mathcal{O}_{V-A} = \gamma^\mu (1 - \gamma_5)$$

$$\mathcal{O}_{V+A} = \gamma^\mu (1 + \gamma_5)$$

$$\mathcal{O}_{S-P} = (1 - \gamma_5)$$

$$\mathcal{O}_{S+P} = (1 + \gamma_5)$$

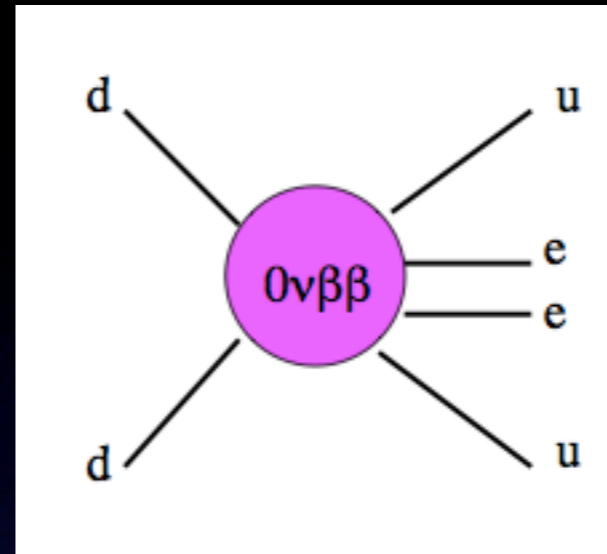
$$\mathcal{O}_{TL} = \frac{i}{2} [\gamma_\mu, \gamma_\nu] (1 - \gamma_5)$$

$$\mathcal{O}_{TR} = \frac{i}{2} [\gamma_\mu, \gamma_\nu] (1 + \gamma_5).$$

[HP, Hirsch, Kovalenko, Klapdor-Kleingrothaus, PLB 1999]

General parametrization for $0\nu\beta\beta$ Decay

Short-Range Part



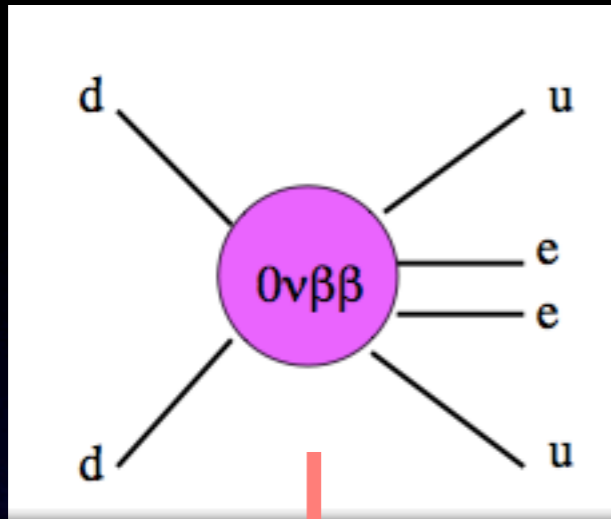
$$\mathcal{L} = \frac{G_F^2}{2} m_p^{-1} \{ \epsilon_1 J J j + \epsilon_2 J^{\mu\nu} J_{\mu\nu} j + \epsilon_3 J^\mu J_\mu j + \epsilon_4 J^\mu J_{\mu\nu} j^\nu + \epsilon_5 J^\mu J j_\mu + \epsilon_6 J^\mu J^\nu j_{\mu\nu} + \epsilon_7 J J^{\mu\nu} j_{\mu\nu} + \epsilon_8 J_{\mu\alpha} J^{\nu\alpha} j_\nu^\mu \},$$

$ \epsilon_1 $	$ \epsilon_2 $	$ \epsilon_3^{LLz} , \epsilon_3^{RRz} $	$ \epsilon_3^{LRz} , \epsilon_3^{RLz} $	$ \epsilon_4 $	$ \epsilon_5 $
$3 \cdot 10^{-7}$	$2 \cdot 10^{-9}$	$4 \cdot 10^{-8}$	$1 \cdot 10^{-8}$	$2 \cdot 10^{-8}$	$2 \cdot 10^{-7}$

[HP, Hirsch, Kovalenko, Klapdor-Kleingrothaus, PLB 2001]

LN at the LHC

$0\nu\beta\beta$ - LHC Complementarity



*pointlike @ nuclear Fermi
momentum $O(r_N^{-1}) \sim 100 \text{ MeV}$*



$d = 9$ operator



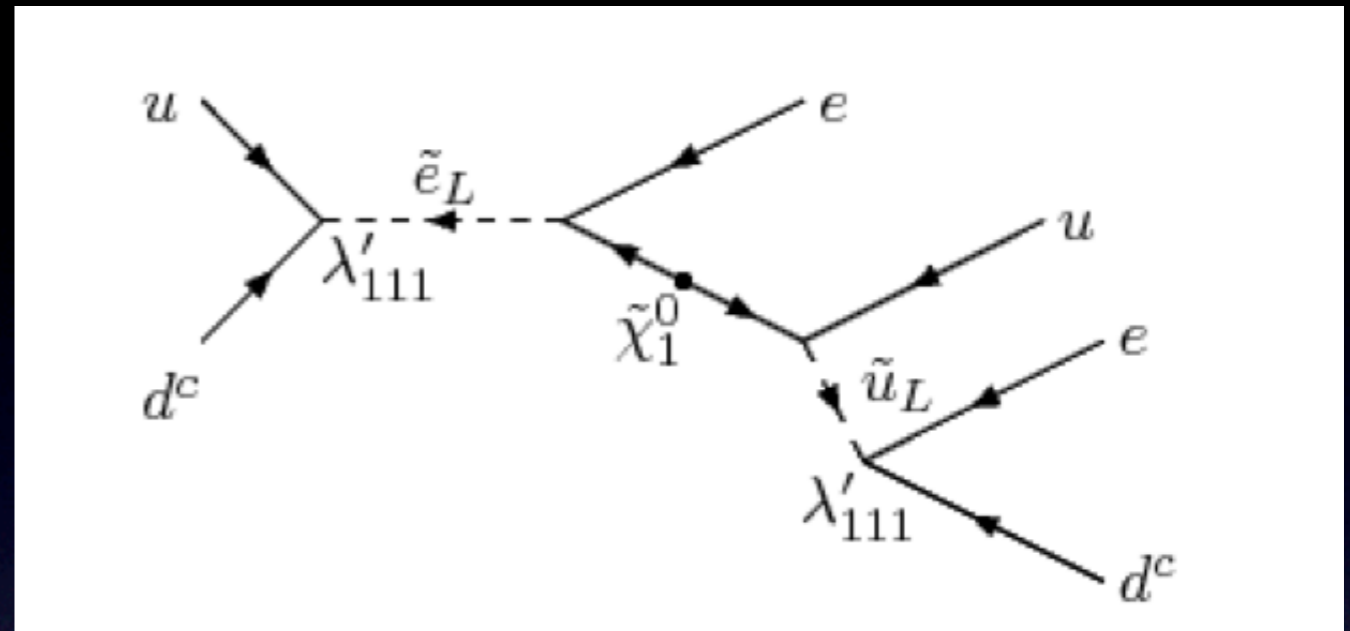
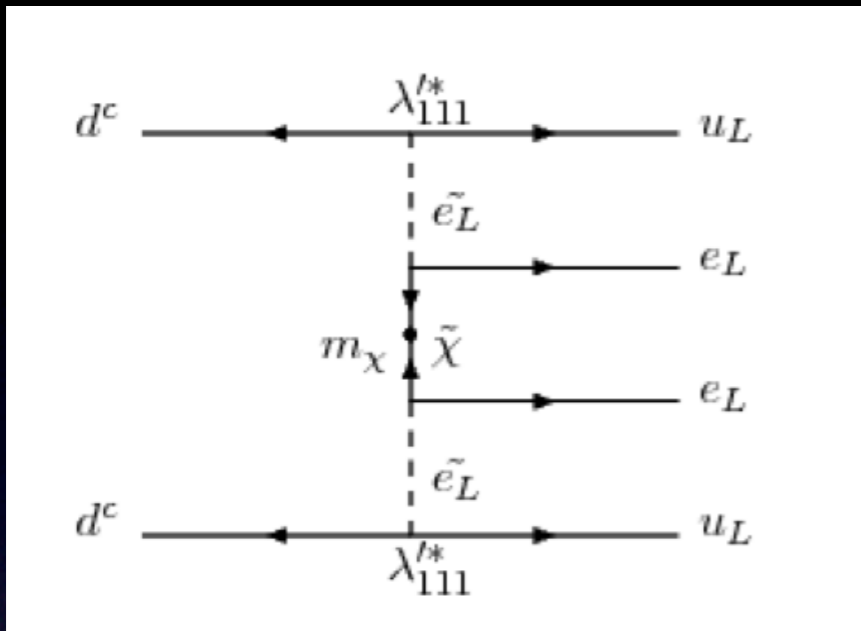
TeV scale particles



New Physics
@ the LHC!

$0\nu\beta\beta$ - LHC Complementarity!

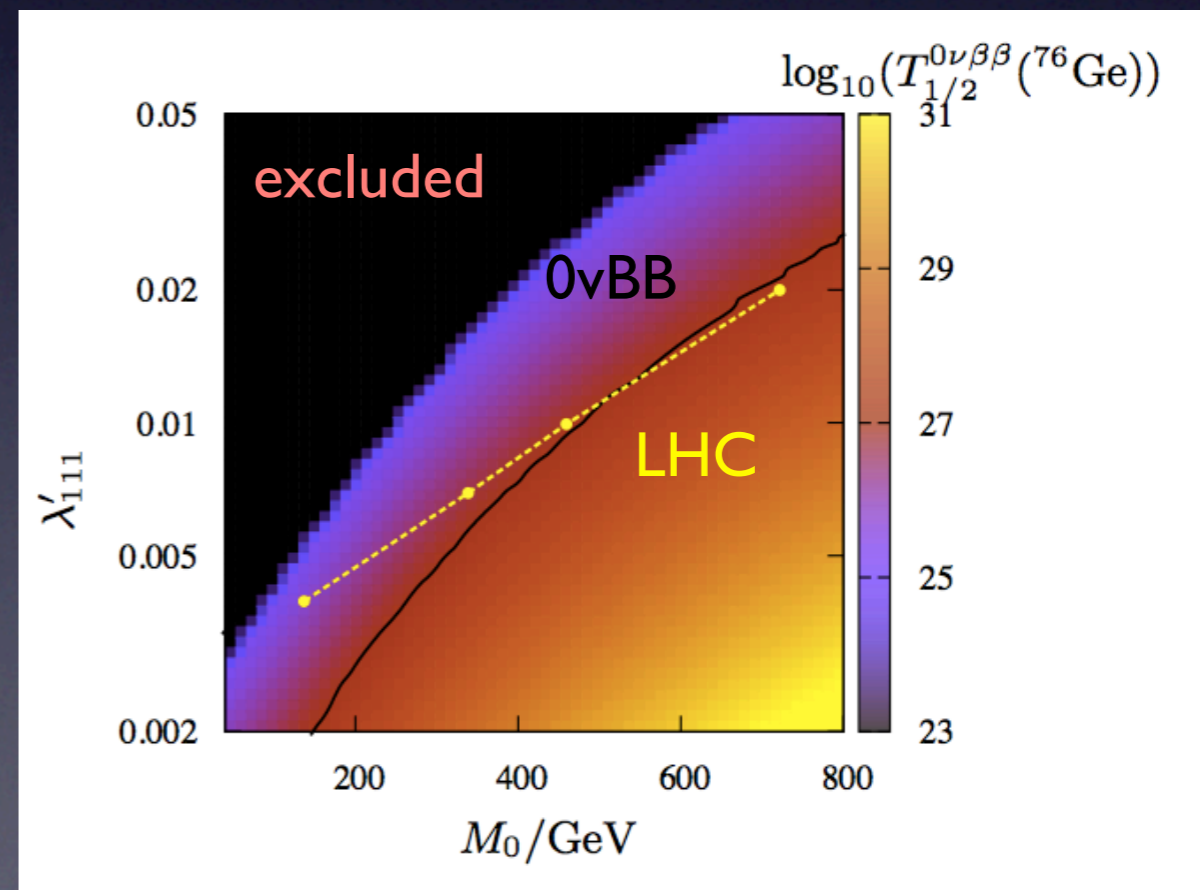
$0\nu\text{BB}$ @ LHC: Example RPV SUSY



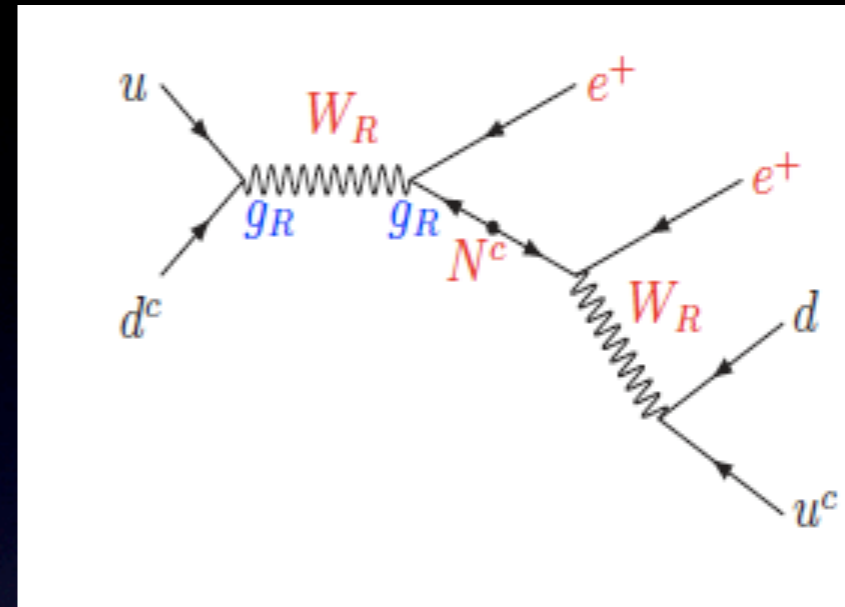
Like-Sign Di-Lepton Signal:
 10 fb^{-1} , $\sqrt{s} = 14 \text{ TeV}$
 SM + SUSY Background

$$M_{1/2} = 300 \text{ GeV} + 0.6M_0$$

[Allanach, Kom, HP, PRL 2009]



Also, e.g.
Left-Right
Symmetry:



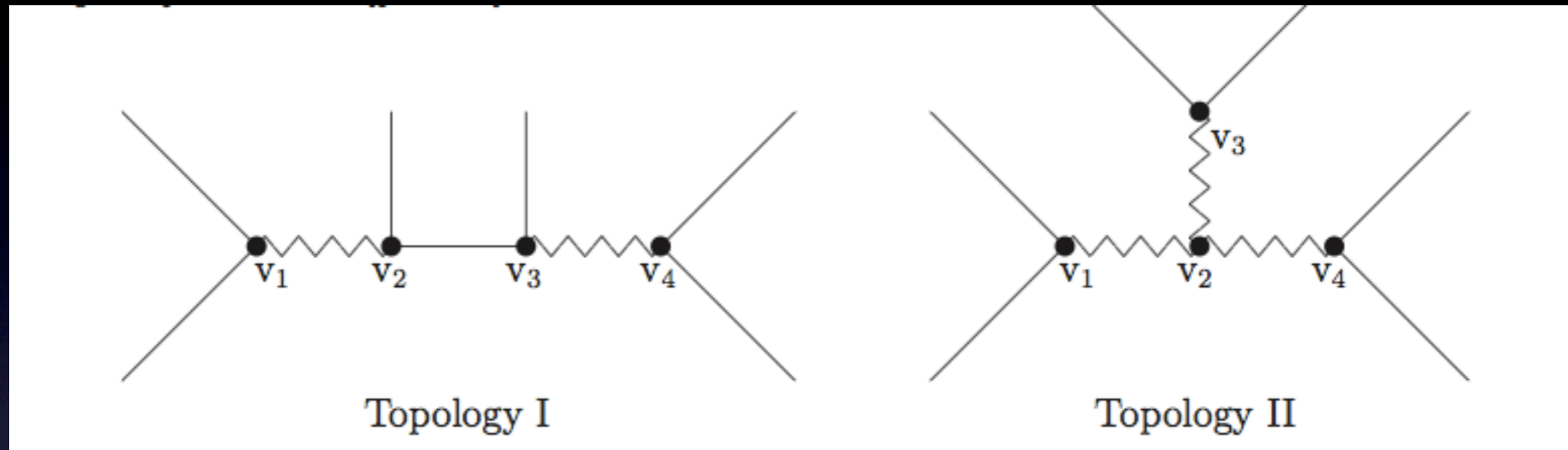
[Keung, Senjanovic 1983; Nemevsek, Nesti, Senjanovic, Tello 2001]

Model-independent approach also for LNV@LHC ?

- ▶ Necessary to go beyond effective field theory approach of [HP, Hirsch, Kovalenko, Klapdor-Kleingrothaus, PLB 1999]
- ▶ Open up vertices, systematic decomposition of the $d=9$ 0νBB operator [Bonnet, Hirsch, Ota, Winter, 2013]

Systematic decomposition of 0νBB Decay

2 Topologies:



— Fermion
∧ Scalar or Vector

→

VFV, SFS, SFV

SSS, VVV, VVS, SSV

⇒ Go through all gauge invariant possibilities giving

rise to $\bar{u} \bar{u} \bar{e} \bar{e} d d \frac{1}{M^5}$ [Bonnet, Hirsch, Ota, Winter, 2013]

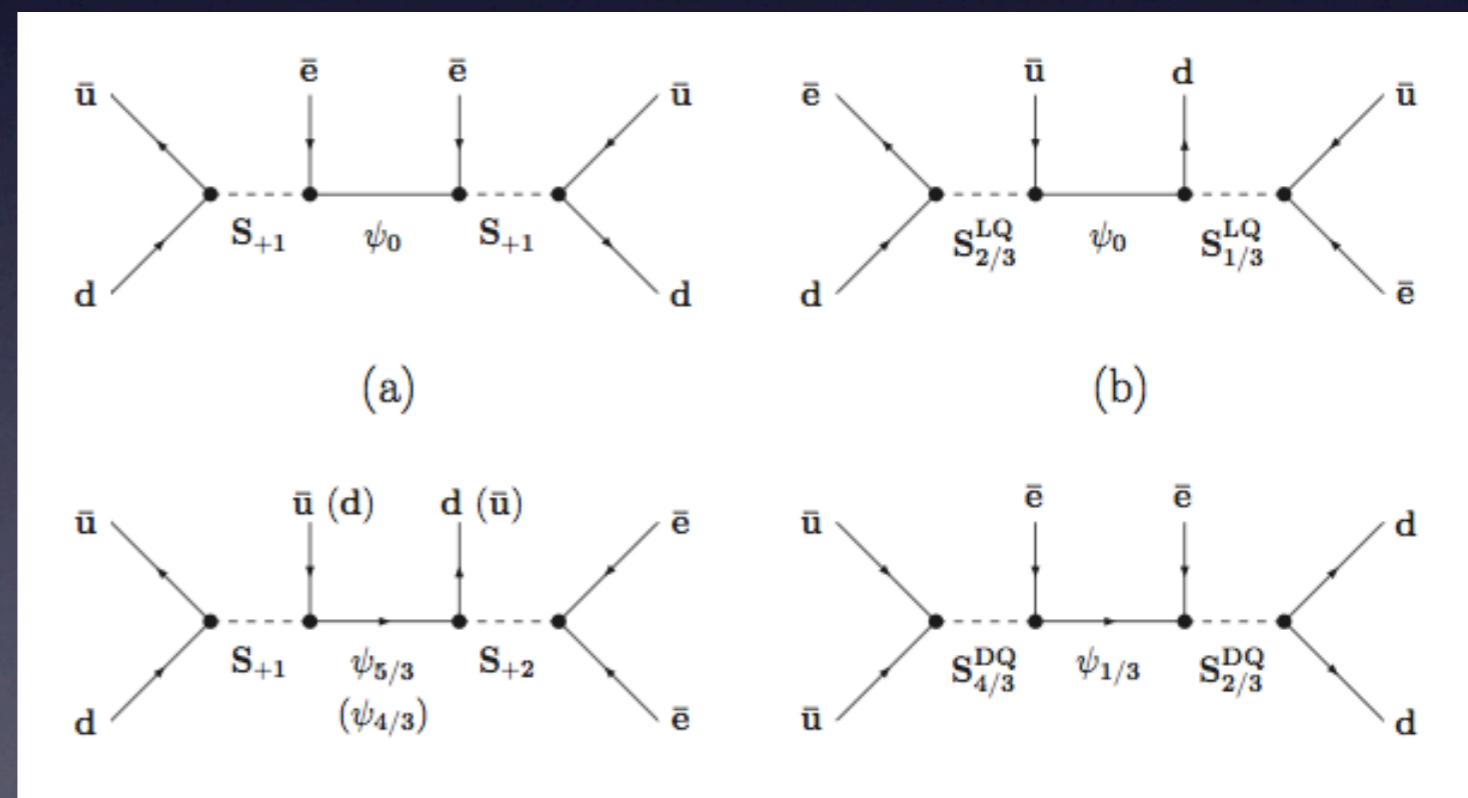
Model independent approach to $0\nu\text{BB}$ @LHC

Apply the decomposition of [Bonnet, Hirsch, Ota, Winter, 2013]
to $0\nu\text{BB}$ @ LHC:

- ▶ Results for Scalars and Vectors very similar
- ▶ Concentrate on Scalars & Topology 1:

S_{charge} ,
 S_{charge}^{LQ} ($L \neq 0, B \neq 0$),
 S_{charge}^{DQ} ($B = \pm 2/3$)

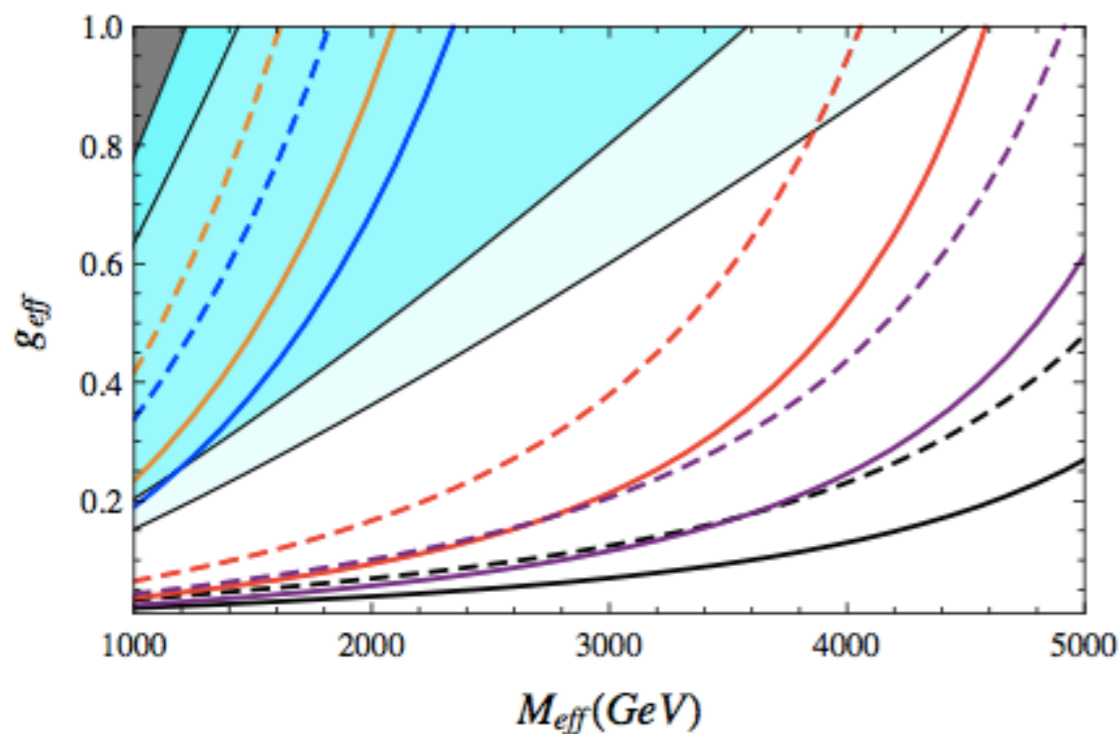
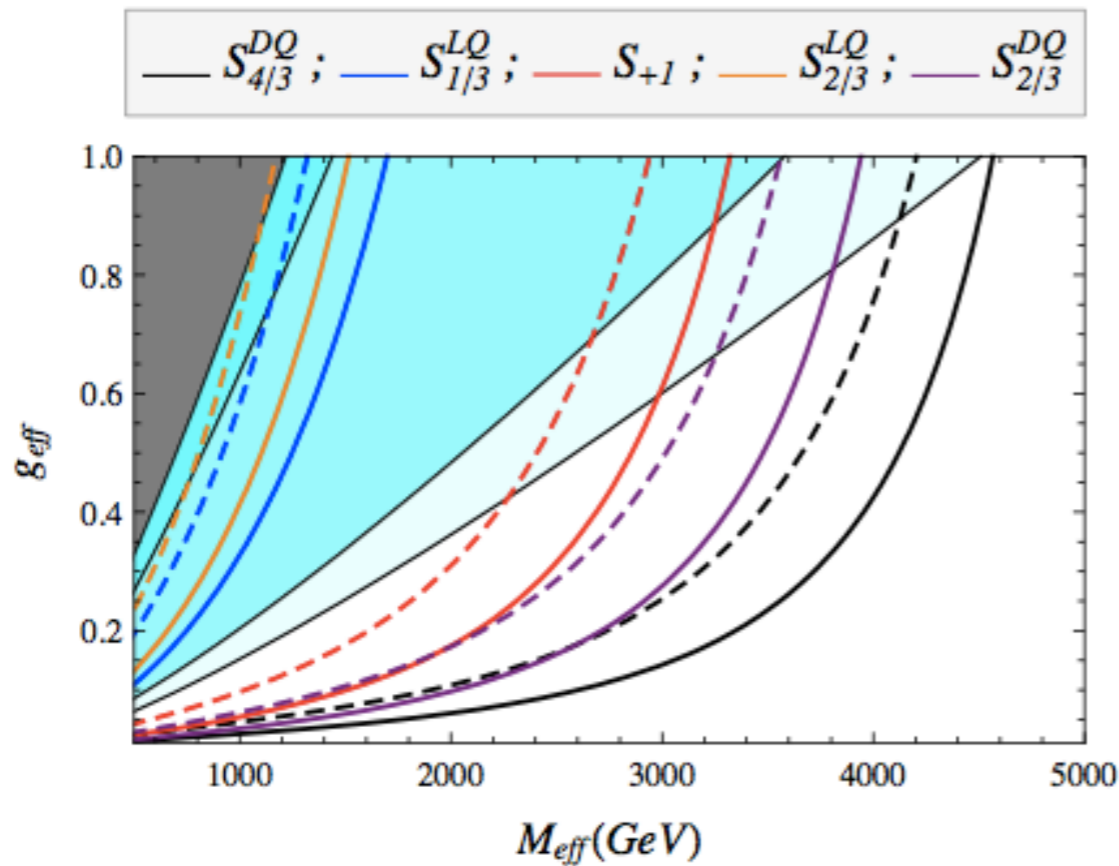
[Helo, Hirsch, Kovalenko, HP, 2013]



$0\nu\beta\beta$ versus LHC SSD at $\sqrt{s} = 14$ TeV

$m_\psi = 200$ GeV

$m_\psi = 1$ TeV



$0\nu\beta\beta$ - grey & blue areas:
present bound, $T_{1/2} > 10^{26}$ y
(smallest rate operator)

$T_{1/2} > 10^{26}$ y, $T_{1/2} > 10^{27}$ y
(largest rate operator)

$$\mathcal{A}^{0\nu\beta\beta} \propto \frac{g_1 g_2 g_3 g_4}{m_{S_i}^2 m_{\psi_q} m_{S_j}^2} \equiv \frac{g_{eff}^4}{M_{eff}^5}$$

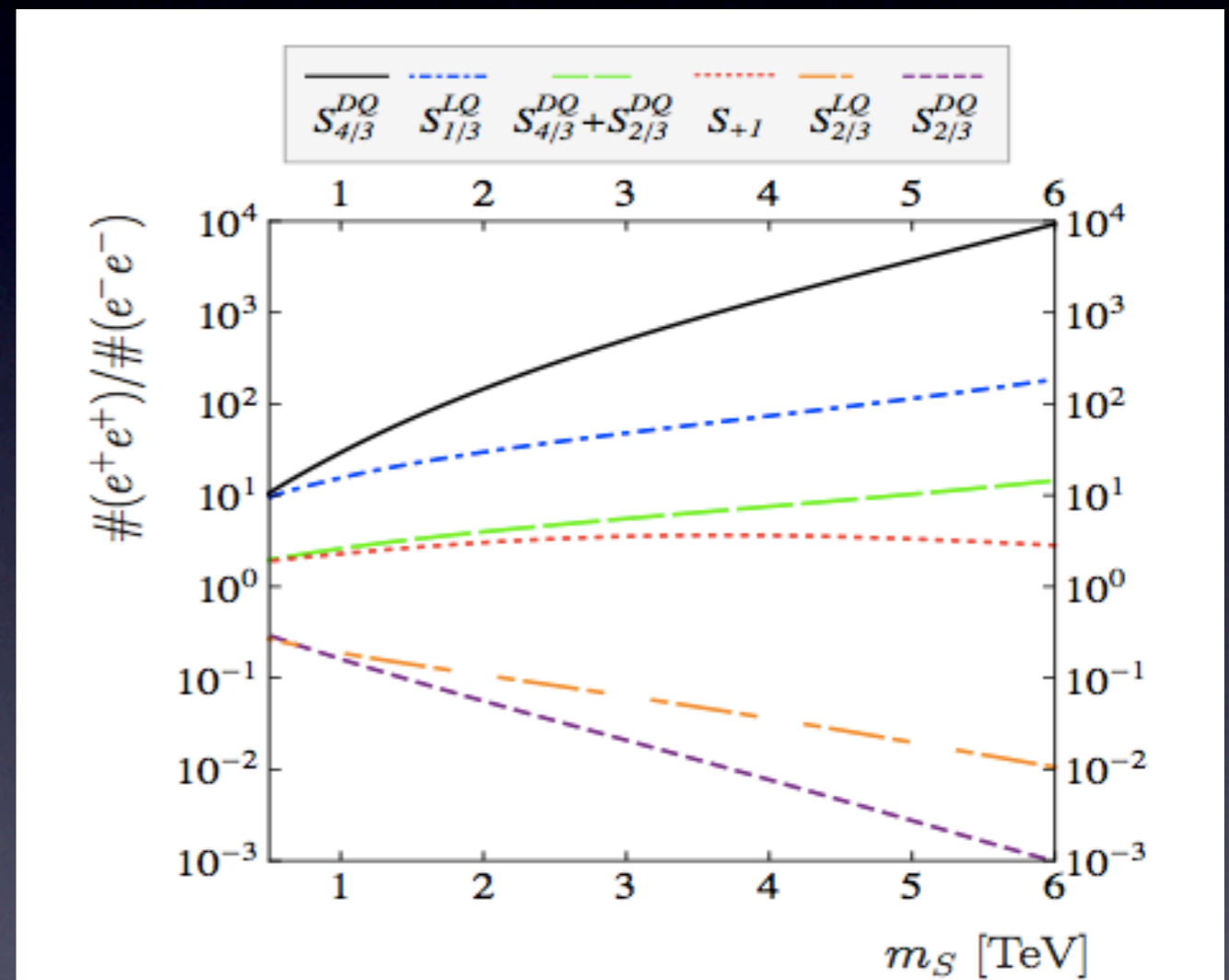
LHC: curved lines
dashed: $Br^{eff}(S \rightarrow eejj) = 10^{-2}$
solid: $Br^{eff}(S \rightarrow eejj) = 10^{-1}$

[Helo, Hirsch, Kovalenko, HP, 2013]

Discriminating $0\nu\beta\beta$ mechanisms: Charge Asymmetry

Moreover: LHC can **discriminate contributions**

- ▶ Invariant mass peaks (s-channel: $m^2_{eejj} \equiv m^2_S$)
- ▶ Charge asymmetry (due to different numbers of u, d, u^c, d^c in p and different diagrams)



[J.C. Helo, M. Hirsch, S. Kovalenko, HP, PRD88 (2013) 1 011901]

$0\nu\beta\beta$ vs LHC sensitivities

→ with the exception of Leptoquarks:

LHC more sensitive than $0\nu\beta\beta$ Decay !

$0\nu\beta\beta$ signal → LHC signal

- OR -

No LHC signal → no $0\nu\beta\beta$ signal

$0\nu\beta\beta$ is
Long Range
(e.g. mass-
mechanism)

→ how to find out?

A major problem

Uncontroversial detection of $0\nu\beta\beta$ decay: uttermost importance!

- prove lepton number to be broken in Nature
- prove neutrinos to be Majorana particles *Schechter and Valle, 1982*

However: it will immediately generate another puzzle:

which mechanism that triggers the decay?

Without identification of the underlying mechanism:

- experimental evidence for $0\nu\beta\beta$ decay will only provide ambiguous information about the concrete physics underlying the decay!
- No information about m_ν can be obtained from a measurement of the neutrinoless double beta decay half life!

Big question:

what is it?



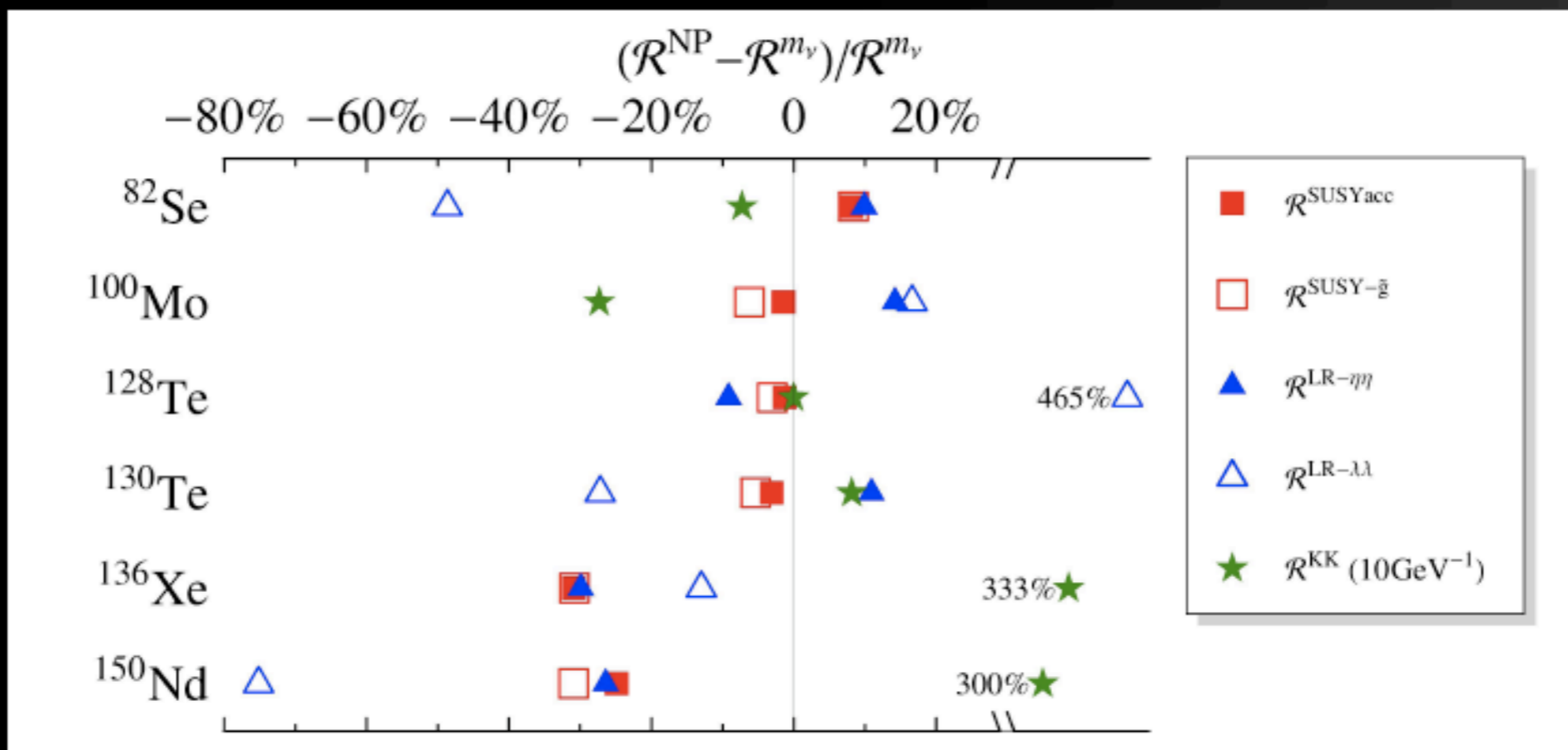
Half life ratios

- Concentrate on: different mechanisms result in different NMEs
- Problem: smaller NME for e.g. the mass mechanism as compared to any alternative new physics mechanism can be compensated by a larger value for the neutrino mass
- However: If one mechanism dominates $\rightarrow \langle m_\nu \rangle$ or ϵ_{NP} drops out in the ratio of experimentally determined half lives for two different emitter isotopes

$$\frac{T_{1/2}(^A X)}{T_{1/2}(^{76}\text{Ge})} = \frac{|\mathcal{M}(^{76}\text{Ge})|^2 G(^{76}\text{Ge})}{|\mathcal{M}(^A X)|^2 G(^A X)}$$

- \Rightarrow Half life ratios depend on the mechanism of double beta decay, but not on the new physics parameter!
- Compare with theoretical prediction for different mechanisms!
- Error in NME ratio can be reduced compared to theoretical error in one matrix element (cancellations of systematic effects)!

Half life ratios: Results



F. Deppisch, H. Päs, Phys. Rev. Lett. 98 (2007) 232501

Matrix elements calculated in the QRPA approach of

A. Staudt, K. Muto and H. V. Klapdor-Kleingrothaus, Europhys. Lett. **13**, 31 (1990); M. Hirsch, K. Muto, T. Oda and H. V. Klapdor-Kleingrothaus, Z. Phys. A **347**, 151 (1994)

or taken from literature using the same code

Half life ratios: results

- R_P SUSY contributions:

similar and rather small deviations

Most effectively discriminated by comparing ^{82}Se and ^{136}Xe (60% variation)

- Left-right symmetric models:

strong deviations for $\lambda\lambda$ combination, comparing ^{128}Te and ^{150}Nd :

$$T_{1/2}^{LR}/T_{1/2}^{m_\nu}[^{128}\text{Te}] \gtrsim 20 \times T_{1/2}^{LR}/T_{1/2}^{m_\nu}[^{150}\text{Nd}]$$

small deviations for $\eta\eta$ combination comparison of ^{100}Mo and ^{136}Xe yields a variation of 70 %

- Extra-dimensional neutrino models with a large brane shift parameter:

large deviations for ^{136}Xe and ^{150}Nd :

$$T_{1/2}^{KK}/T_{1/2}^{m_\nu}[^{150}\text{Nd}] \gtrsim 5 \times T_{1/2}^{KK}/T_{1/2}^{m_\nu}[^{100}\text{Mo}]$$

Caution: strong deformation of ^{150}Nd is ignored in most QRPA calculations

→ Simkovic, Pacearescu, Faessler, 2004

Nuclear matrix element uncertainties

- **Theoretical errors** of NME calculation **dominate** experimental errors \Rightarrow difficult to determine the confidence level with which either mechanism can be excluded to generate the observed double beta evidence!
- Assuming e.g. a **statistical distribution** of matrix element values \Leftrightarrow **relative variation of 60%** in $\mathcal{R}^{NP}(^A X)$ w.r.t. $\mathcal{R}^{m_\nu}(^A X)$ is **significant only** if NMEs would be known with an accuracy of **15%**! \rightarrow **unrealistic!**
- **Estimates of uncertainties vary**: factor 3-5 (spread of published values) to only **30%** (uncertainties inherent in QRPA) **Rodin, Faessler, Simkovic, Vogel, 2006**

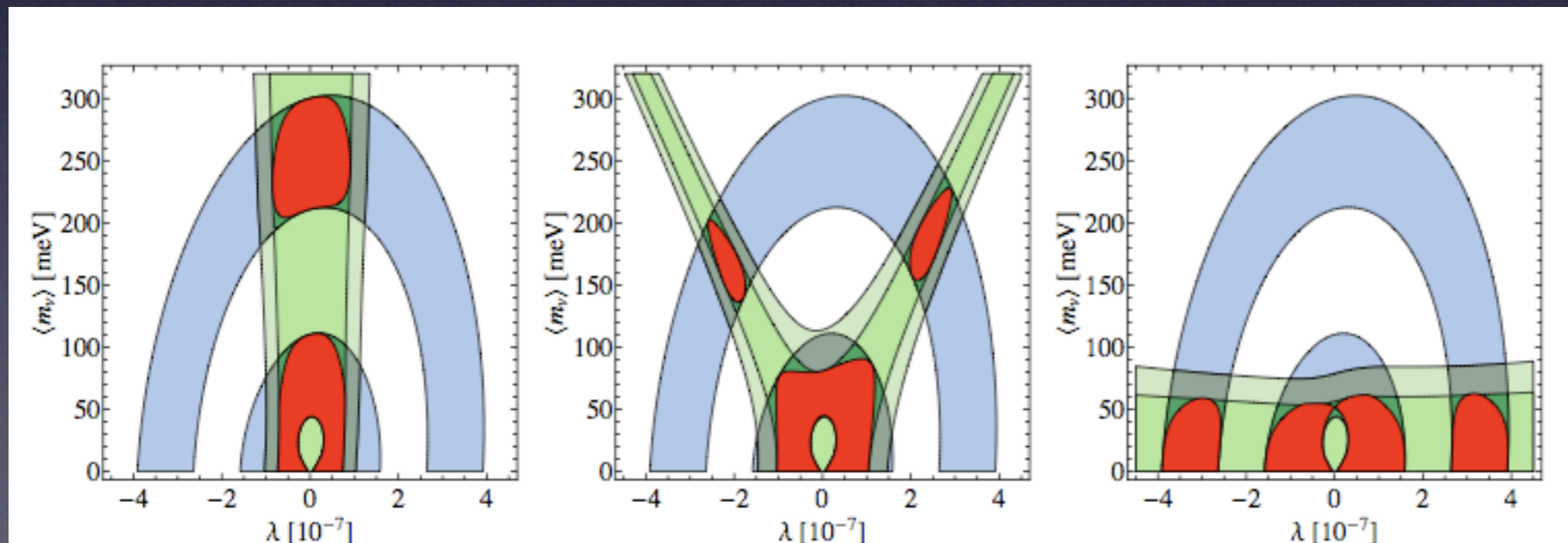
However:

- **significance will increase** if a **whole set of measurements** in different isotopes resembles the expected pattern
- **systematical effects** (like a too small g_{pp} in the pn-QRPA approach, a different g_A , higher-order terms, different model-space) **will cancel out**
- \rightarrow **check results with alternative codes!**
- \rightarrow **include pion exchange** which may be dominating in some of the models discussed!

$0\nu\beta\beta$: Pinning down the mechanism

Possibilities to disentangle at least some of the possible mechanisms:

- analysis of **angular correlations** between the emitted electrons
Doi, Kotani, Nishiura and Takasugi, 1983; Ali, Borisov and Zhuridov, 2006 & 2007
→ **few experiments** sensitive to electron tracks
- comparative study of $0\nu\beta\beta$ and $0\nu\beta^+$ with **electron capture (EC) decay**
Hirsch, Muto, Oda, Klapdor-Kleingrothaus, 1994
→ **small rates** and **experimental challenge** to observe the produced X-rays or Auger electrons
- study of **double beta decay to excited 0^+ states**
Simkovic, Nowak, Kaminski, Raduta, Faessler, 2001
→ **few experiments** sensitive to transitions to excited states.

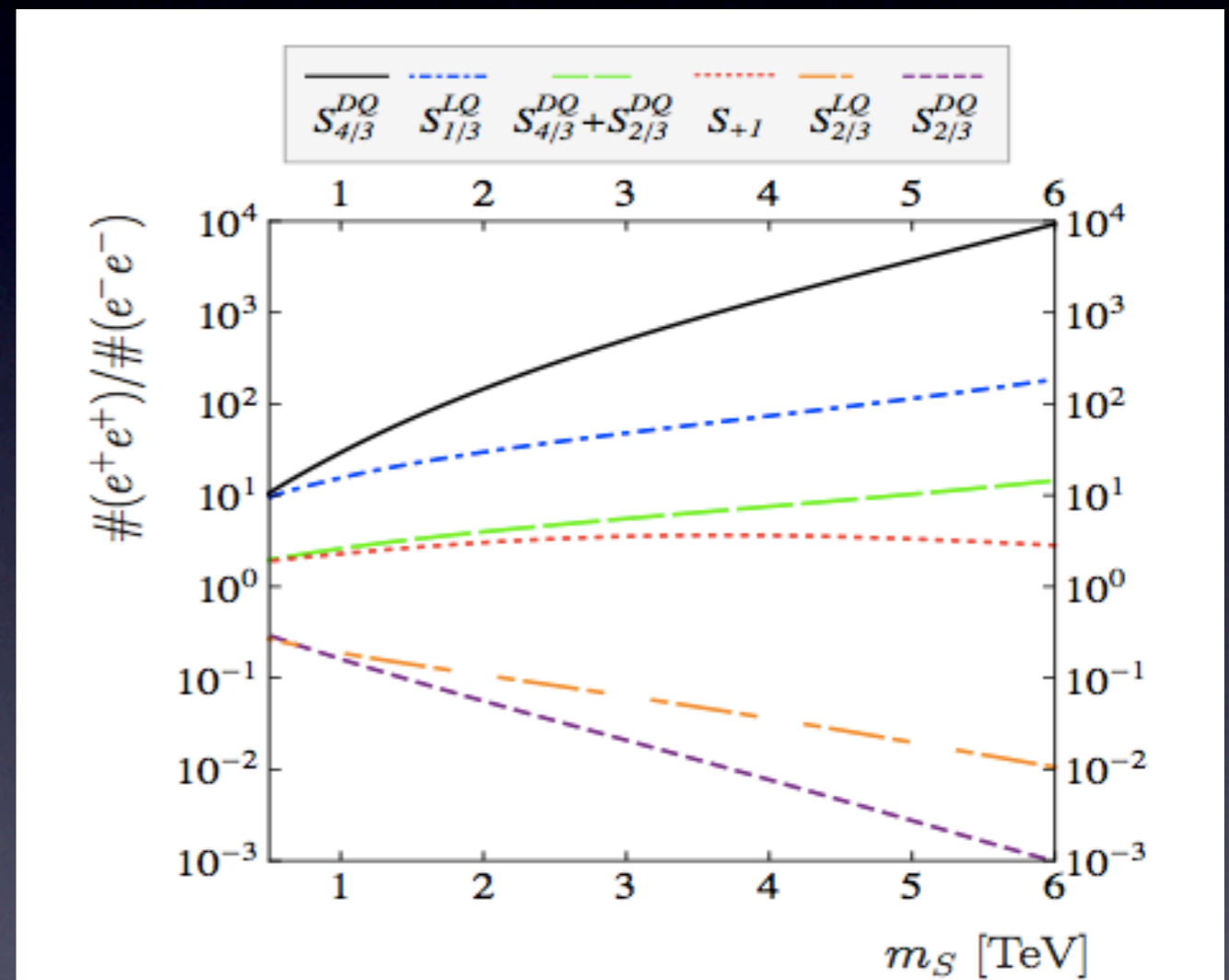


[F. Deppisch
+Super-NEMO,
2010]

Discriminating $0\nu\beta\beta$ mechanisms: Charge Asymmetry

Moreover: LHC can **discriminate contributions**

- ▶ Invariant mass peaks (s-channel: $m^2_{eejj} \equiv m^2_S$)
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[J.C. Helo, M. Hirsch, S. Kovalenko, HP, PRD88 (2013) 1 011901]

What's the connection with Baryogenesis ?

Baryon Number Washout

“Falsifying Leptogenesis at the LHC”

[F. Deppisch, J. Harz, M. Hirsch, PRL 112 (2014) 221601]

“Falsifying Leptogenesis at the LHC”

- ▶ LNV @ LHC
- ▶ Lower bound on washout of Lepton Number Asymmetry
- ▶ No out-of-equilibrium condition in early universe!

But **EVEN WORSE**: consider **Sphalerons**

Leptogenesis:

~~B-L~~

ν_R decay

+

~~B+L~~

Sphalerons



B Asymmetry

In Reverse:

~~B-L~~

e.g. LHC

+

~~B+L~~

Sphalerons



B washout

Large LNV @ LHC (or elsewhere) will washout ANY pre-existing Baryon Asymmetry, irrespective of the Baryogenesis mechanism (Leptogenesis, etc...)

Original Paper: [Fukugita, Yanagida, 1990]

“Sphaleron induced Baryon Number Non-conservation and a constraint on Majorana neutrino masses”

Also:

[Gelmini, Yanagida, 1992] keV-bound on ν mass

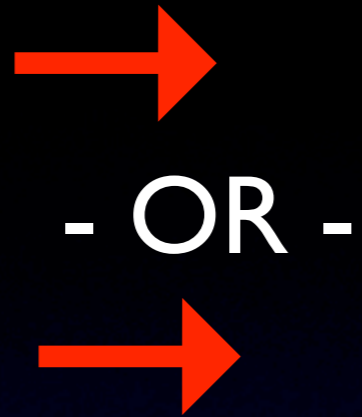
[Klapdor-Kleingrothaus, Kolb, Kuzmin, 1990] Bound on sneutrinos

[Hollenberg, HP, Schalla, 2011] Bound on 4th generation neutrinos

[...many others]

Conclusions I

$0\nu\beta\beta$



Long range
mechanism e.g. m_ν

Short range
mechanism



LNV
@ LHC

[Deppisch, Harz, Hirsch, Huang, HP, 2015]



Low-Scale
Baryogenesis



also detectable @ LHC?
“2 for one”

Conclusions II

High scale
baryogenesis



~~LN
@ LHC~~



If
 $0\nu\text{BB}$



Mass mechanism

m_ν



Very probably high scale origin of m_ν

[Deppisch, Harz, Hirsch, Huang, HP, 2015]

But:

$0\nu\beta\beta$ probes LNV only for the
electron flavor!

How to close this
Flavor Loophole?

Outlook

Closing the flavor
loophole:

combine $0\nu\text{BB}$
and LFV

Washout processes are in Thermal equilibrium for:

$$\frac{\Gamma_W}{H} \equiv \frac{c_D}{n_\gamma H} \frac{T^{2D-4}}{\Lambda_D^{2D-8}} = c'_D \frac{\Lambda_{\text{Pl}}}{\Lambda_D} \left(\frac{T}{\Lambda_D} \right)^{2D-9} \gtrsim 1, \quad (8)$$

with $c'_D = \pi^2 c_D / (3.3 \sqrt{g_*}) \approx 0.3 c_D$. This is the case in the temperature interval

$$\Lambda_D \left(\frac{\Lambda_D}{c'_D \Lambda_{\text{Pl}}} \right)^{\frac{1}{2D-9}} \equiv \lambda_D \lesssim T \lesssim \Lambda_D. \quad (9)$$

[Deppisch, Harz, Hirsch, Huang, HP, 2015]

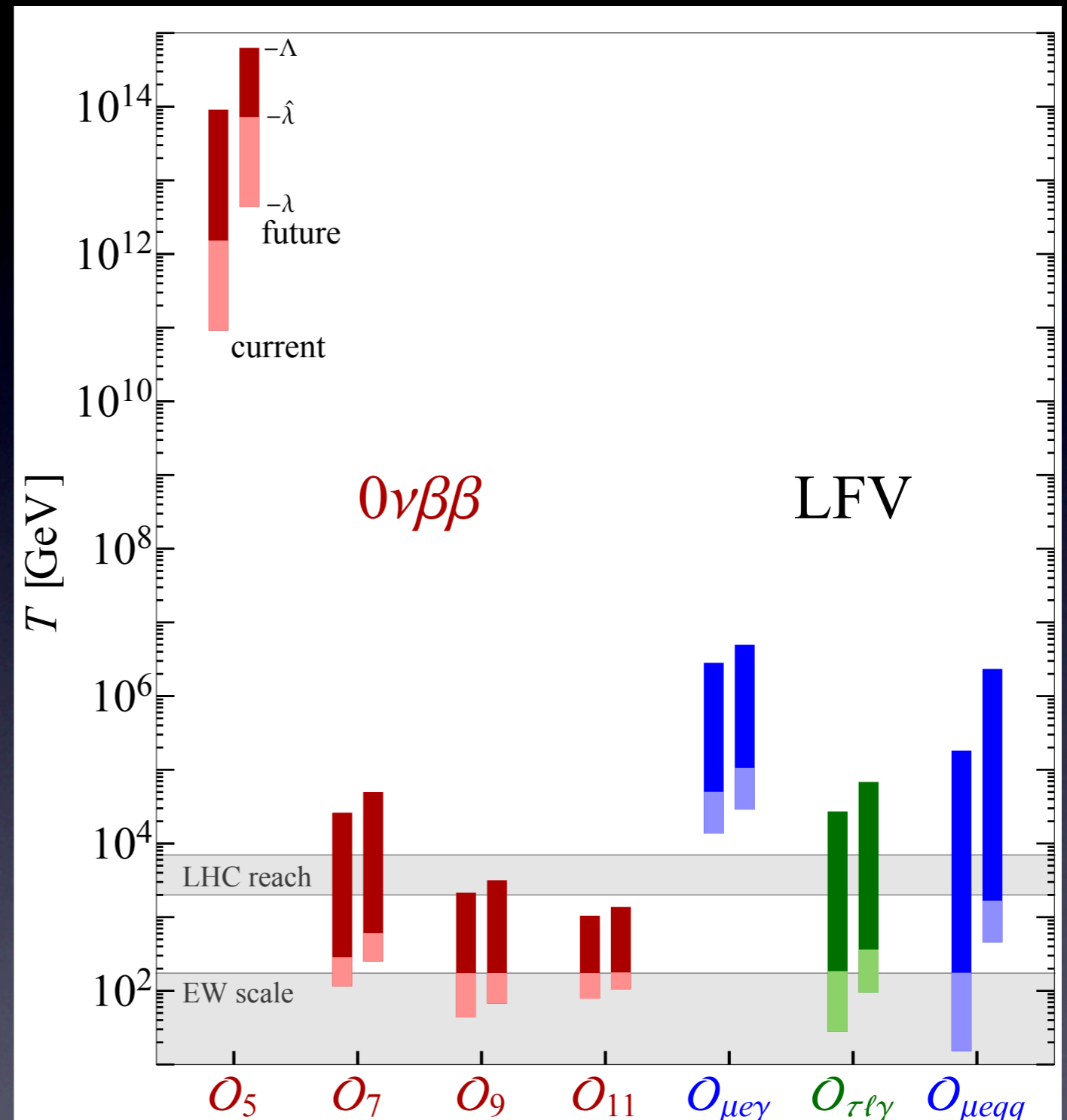
Outlook: closing the flavor loophole

combine $0\nu\text{BB}$ and LFV

Washout intervals:
 lower bound: equilibrium
 upper bound: EFT breakdown

$$T_{1/2} = 2.1 \times 10^{25} \text{ y}$$

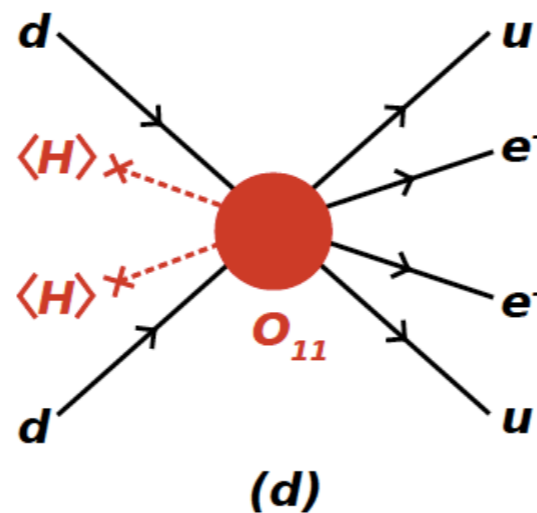
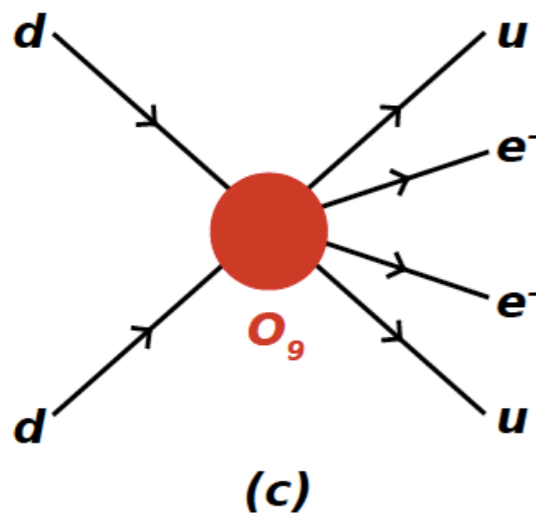
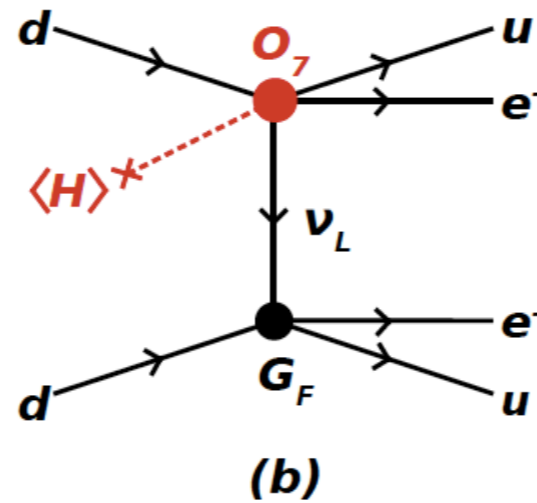
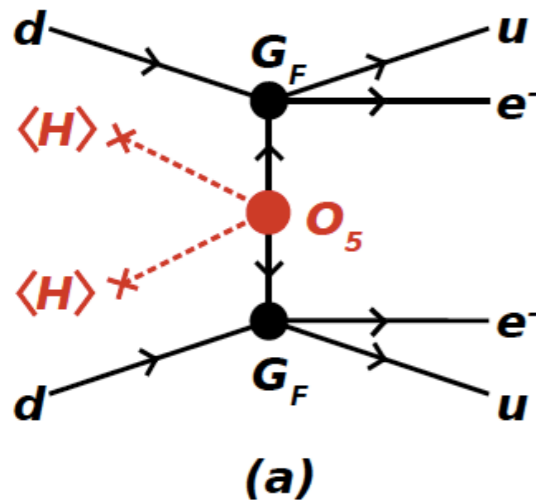
An asymmetry generated at scales above λ will be washed out if $0\nu\text{BB}$ is observed



[Deppisch, Harz, Hirsch, Huang, HP, 2015]

Outlook

Washout processes are in Thermal equilibrium for:



\mathcal{O}_D	λ_D^0 [GeV]	Λ_D^0 [GeV]
\mathcal{O}_5	9.2×10^{10}	9.1×10^{13}
\mathcal{O}_7	1.2×10^2	2.6×10^4
\mathcal{O}_9	4.3×10^1	2.1×10^3
\mathcal{O}_{11}	7.8×10^1	1.0×10^3

\mathcal{O}_i	λ_i^0 [GeV]	Λ_i^0 [GeV]
$\mathcal{O}_{\mu e \gamma}$	1.4×10^4	2.8×10^6
$\mathcal{O}_{\tau l \gamma}$	2.8×10^1	2.7×10^4
$\mathcal{O}_{\mu e q q}$	1.5×10^1	1.8×10^5

$$T_{1/2} = 2.1 \times 10^{25} \text{ y} \cdot (\Lambda_D / \Lambda_D^0)^{2d-8}$$

[Deppisch, Harz, Hirsch, Huang, HP, 2015]

Summary

$0\nu\text{BB}$



- EITHER -

$\text{LNV@LHC} \rightarrow$ Low scale Baryogenesis, “2 for 1”

- OR -

very probably high-scale origin of m_ν
(like vanilla type-I seesaw + leptogenesis)

[Deppisch, Harz, Hirsch, Huang, HP, 2015]

Loopholes exist and should be checked!

(Flavor restriction, conserved charges, etc...)

[Deppisch, Harz, Hirsch, 2014; Antaramian, Hall, Rasin, 1994...]