

What is the 750 GeV $\gamma\gamma$ resonance?

... and what to do with it?

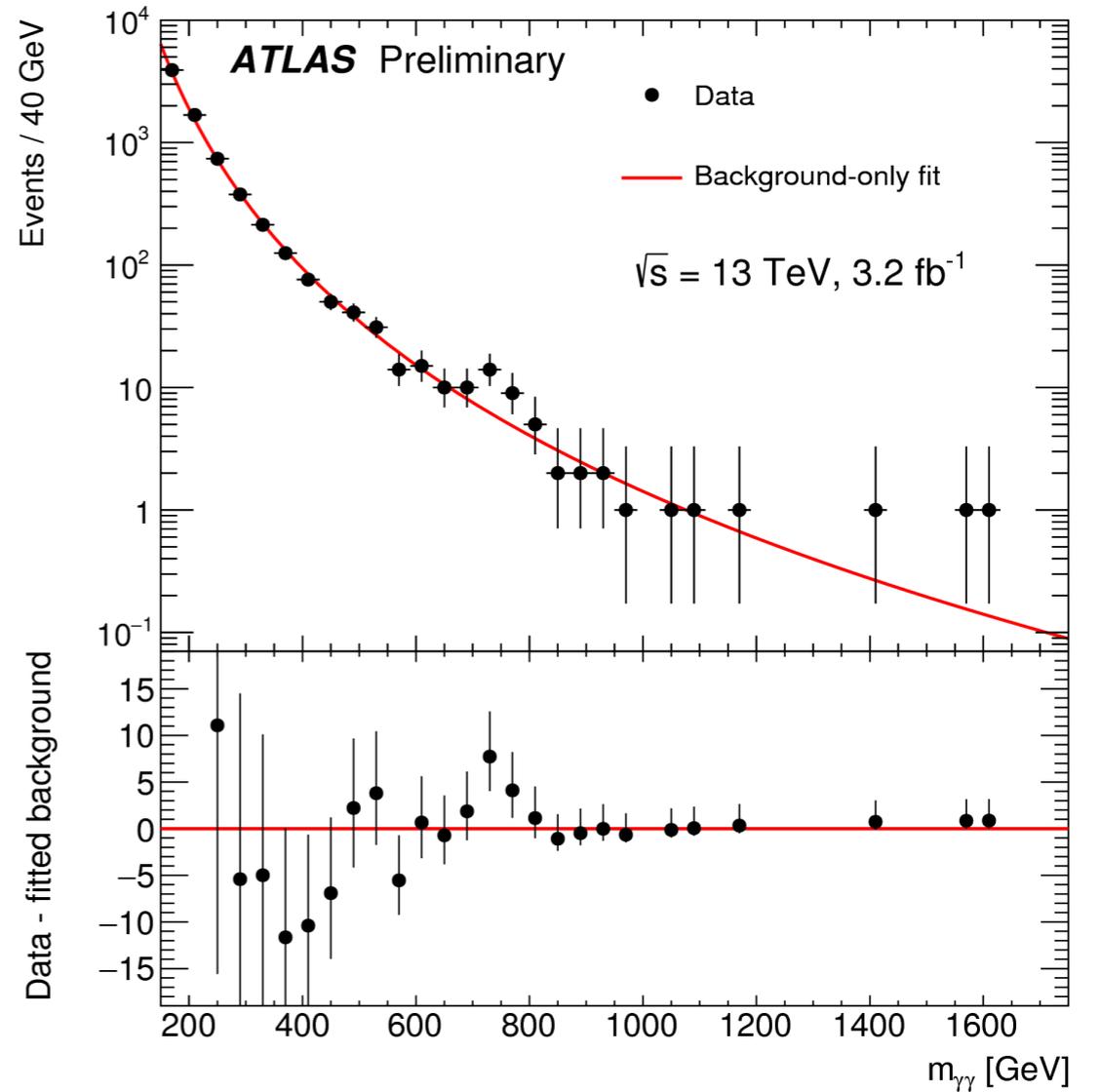
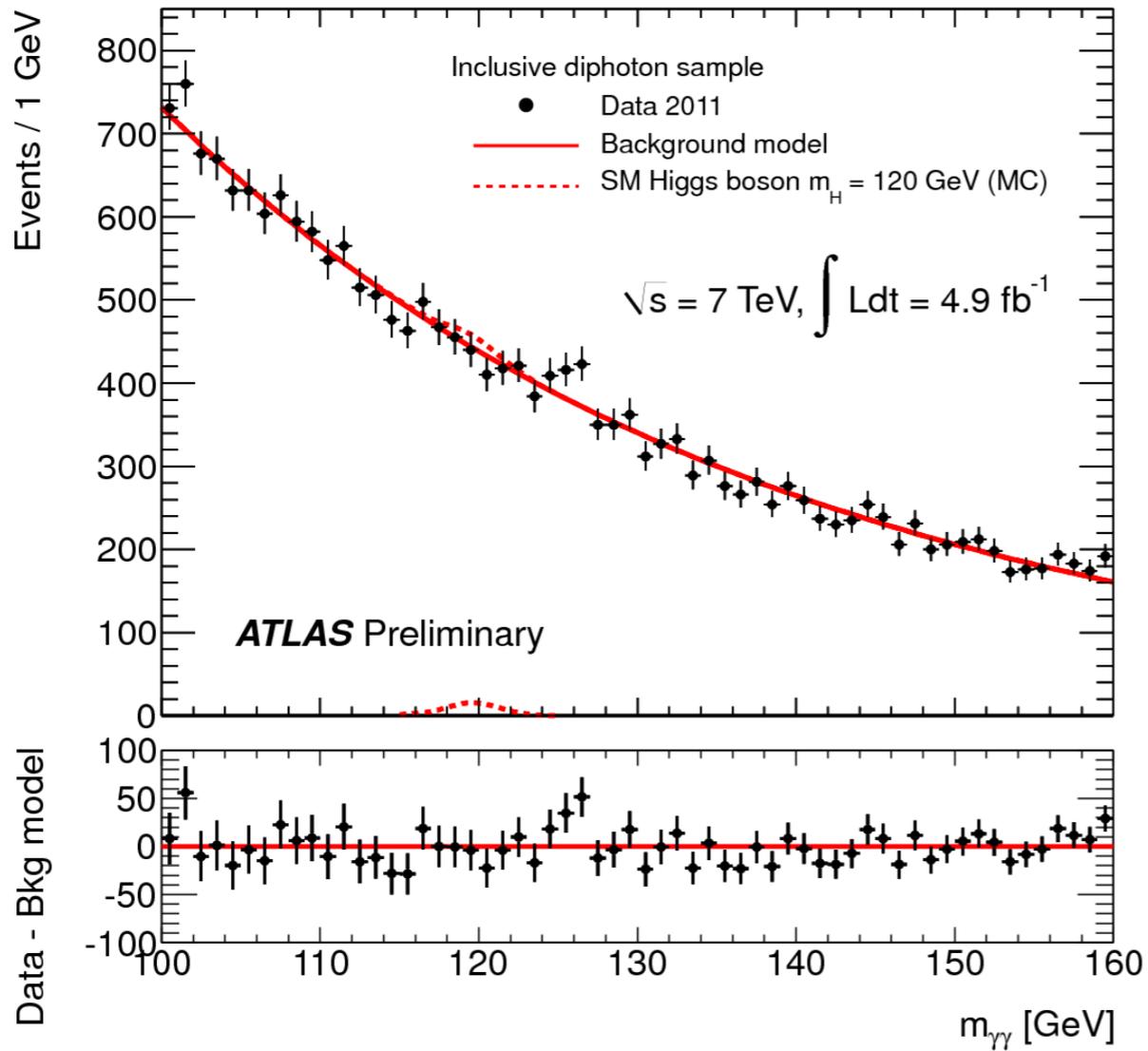
Roberto Franceschini (CERN)
May 30th 2016 - MPI Kernphysik (Heidelberg)

1512.04933, 1512.05330, 1604.06446

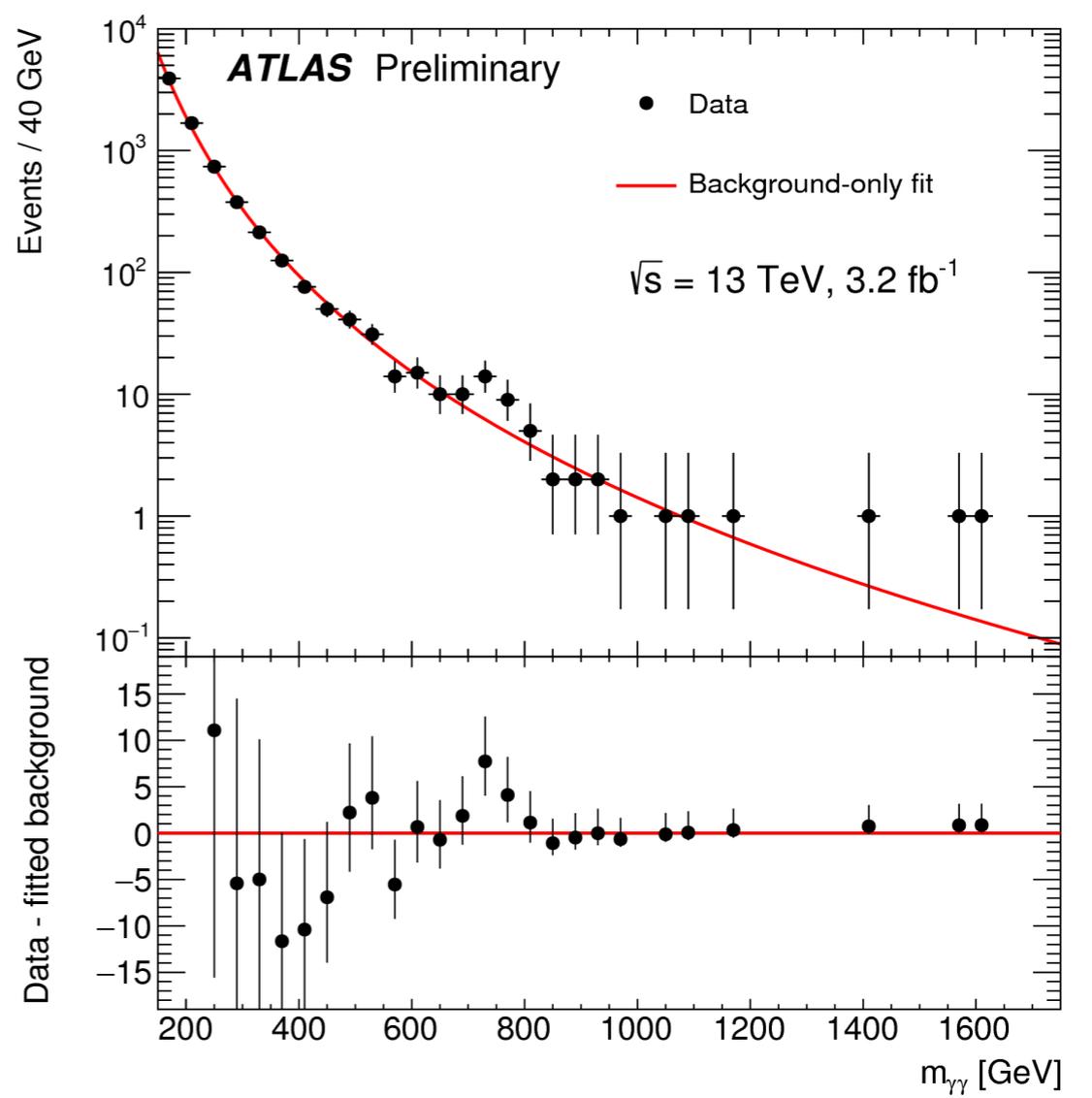
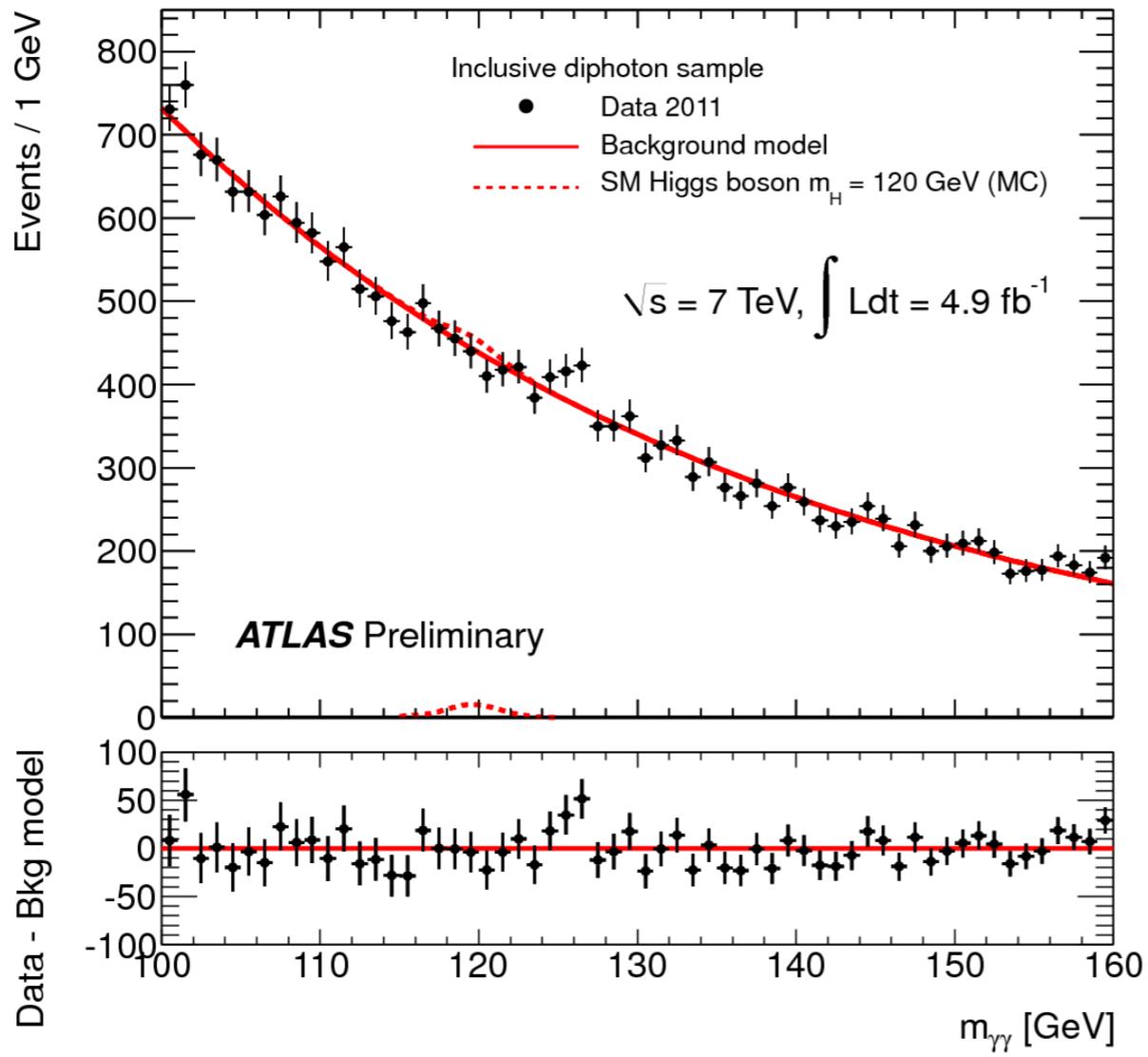
with

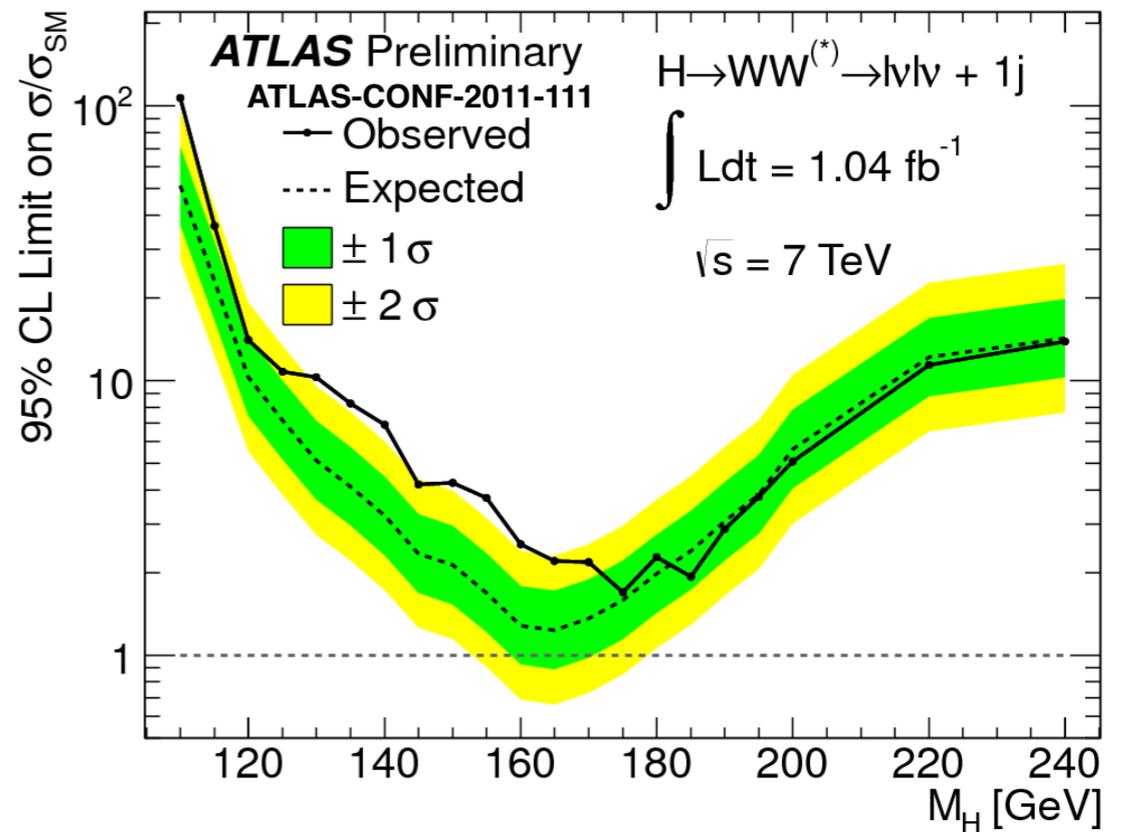
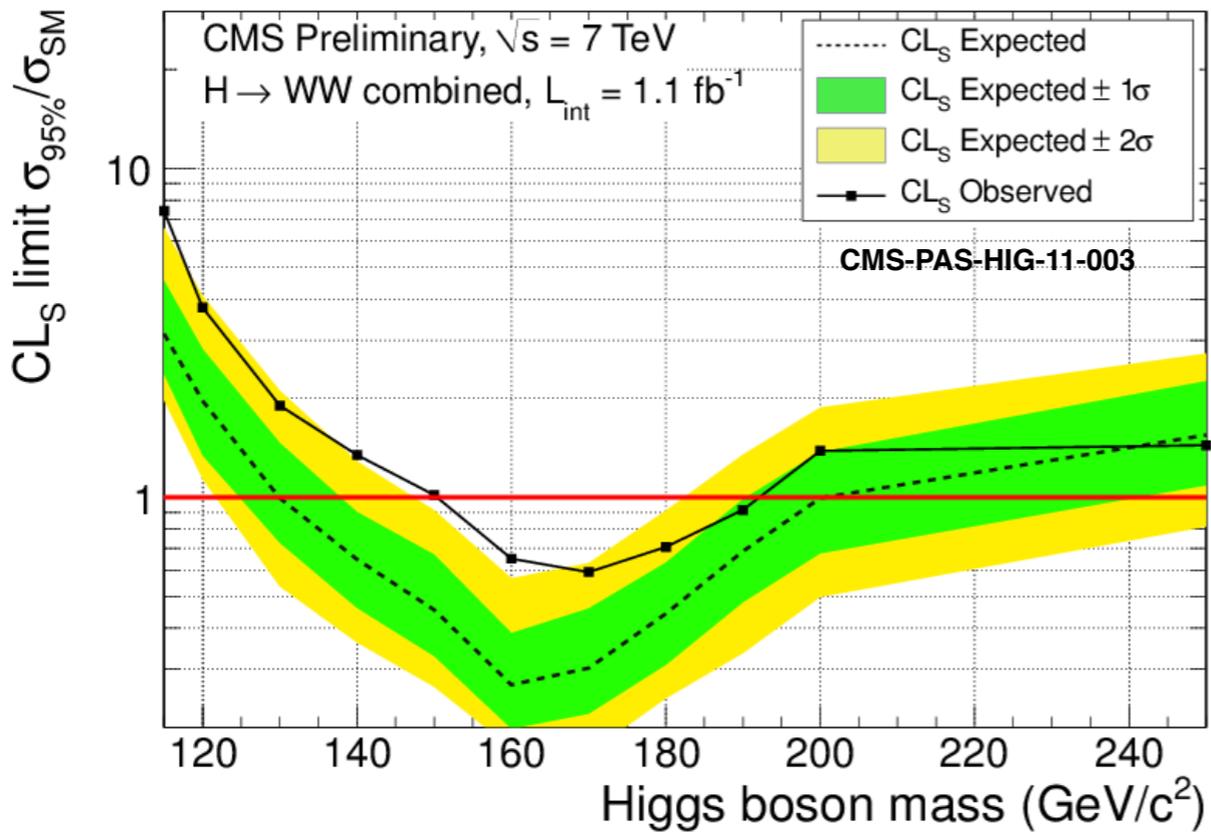
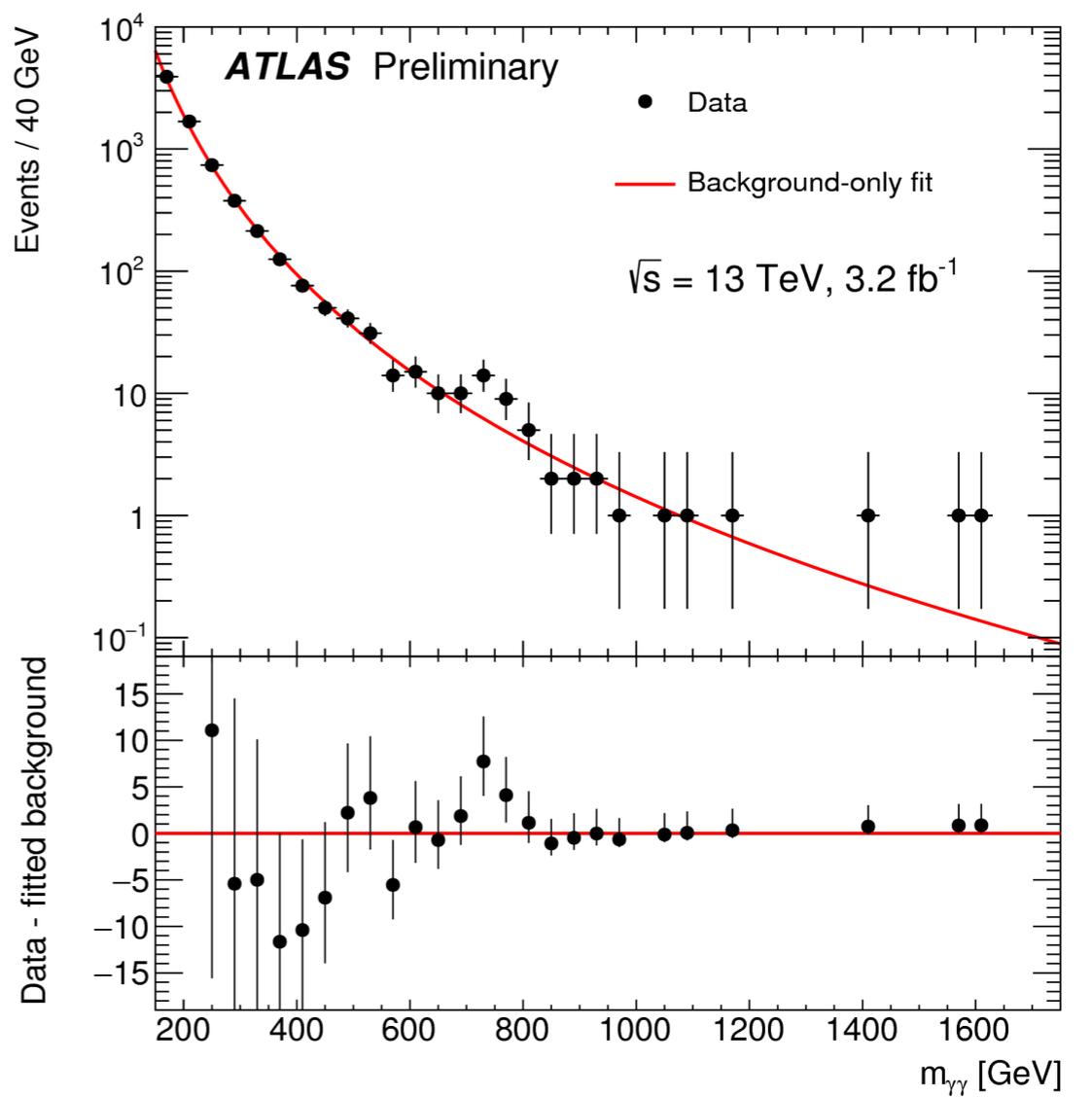
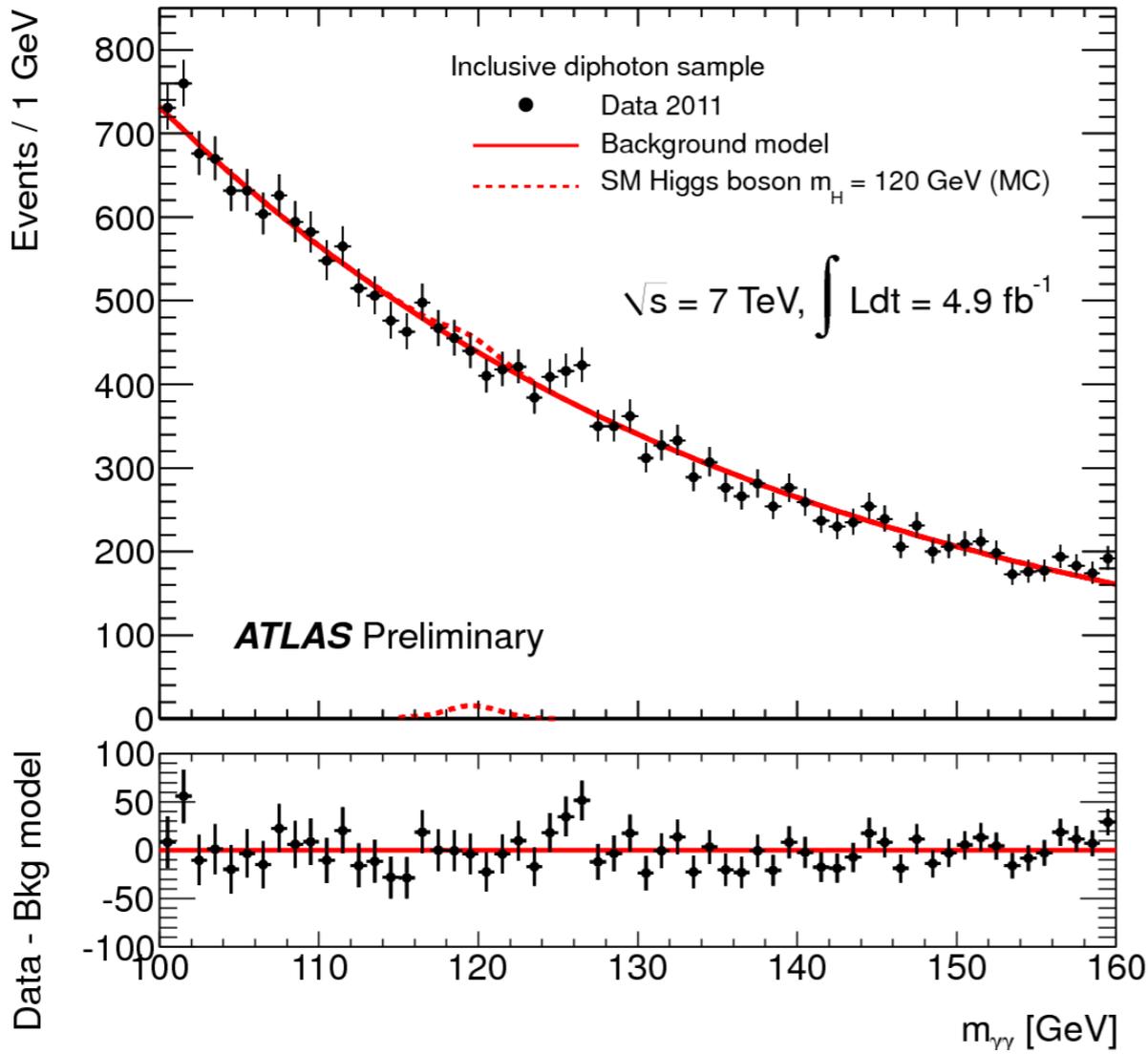
Bellazzini, Sala, Serra, Giudice, Kamenik, McCullough, Pomarol, Rattazzi, Redi, Riva, Strumia, and Torre

Déjà vu ?



Jamboree 2011-15





BSM means operating
in this moving field



BSM means operating
in this moving field



Join the dots

DM

Strong
CP

Flavor

ν

Weak/
Planck
Hierarchy

DE?

Join the dots

DM

Strong
CP

Flavor

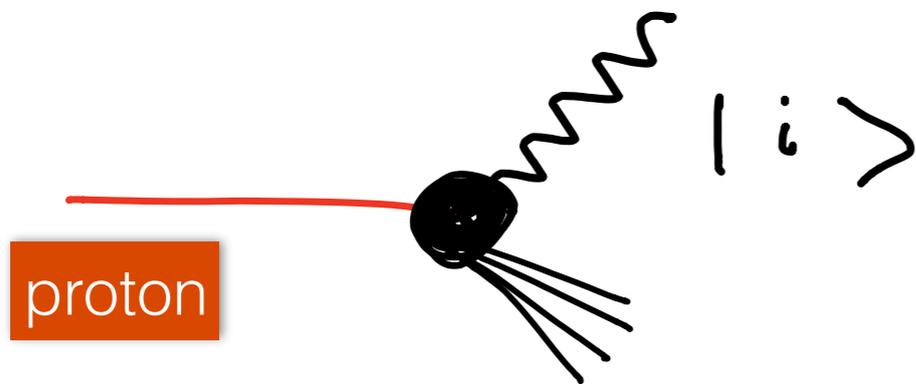
Overwhelming amount
of interpretations

ν

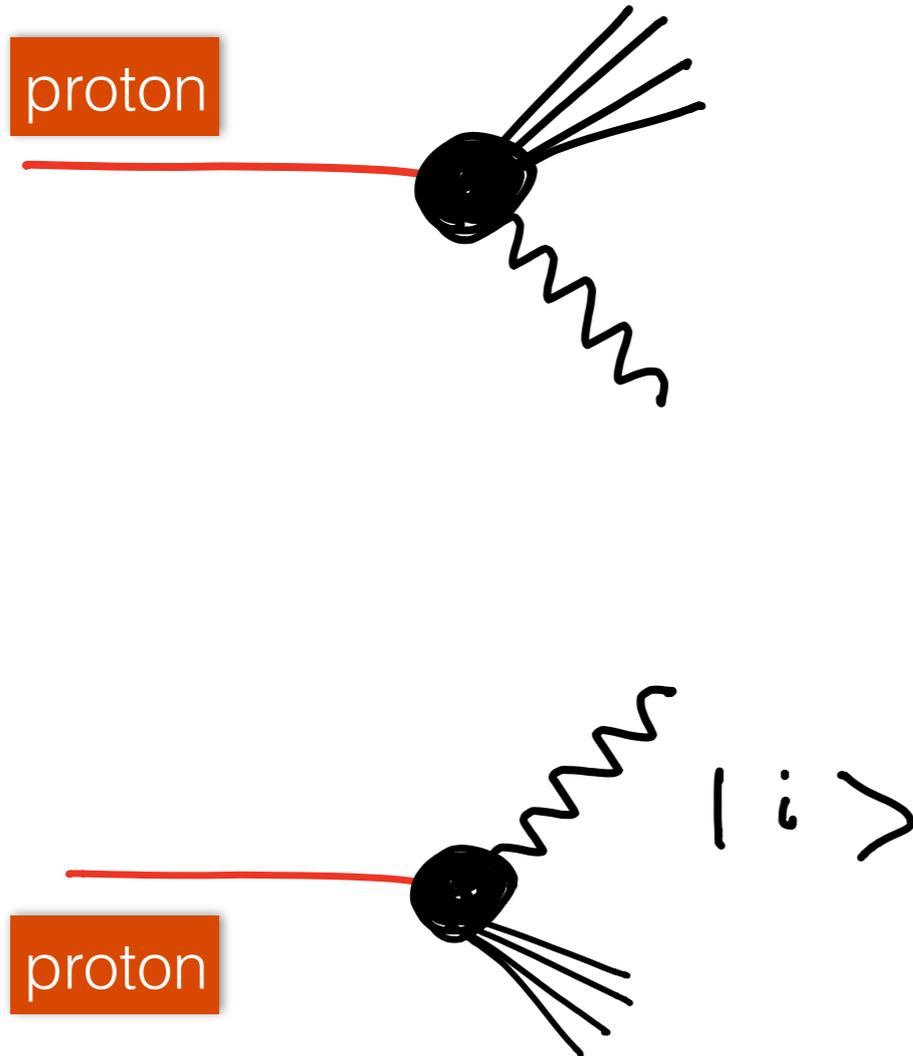
Weak/
Planck
Hierarchy

DE?

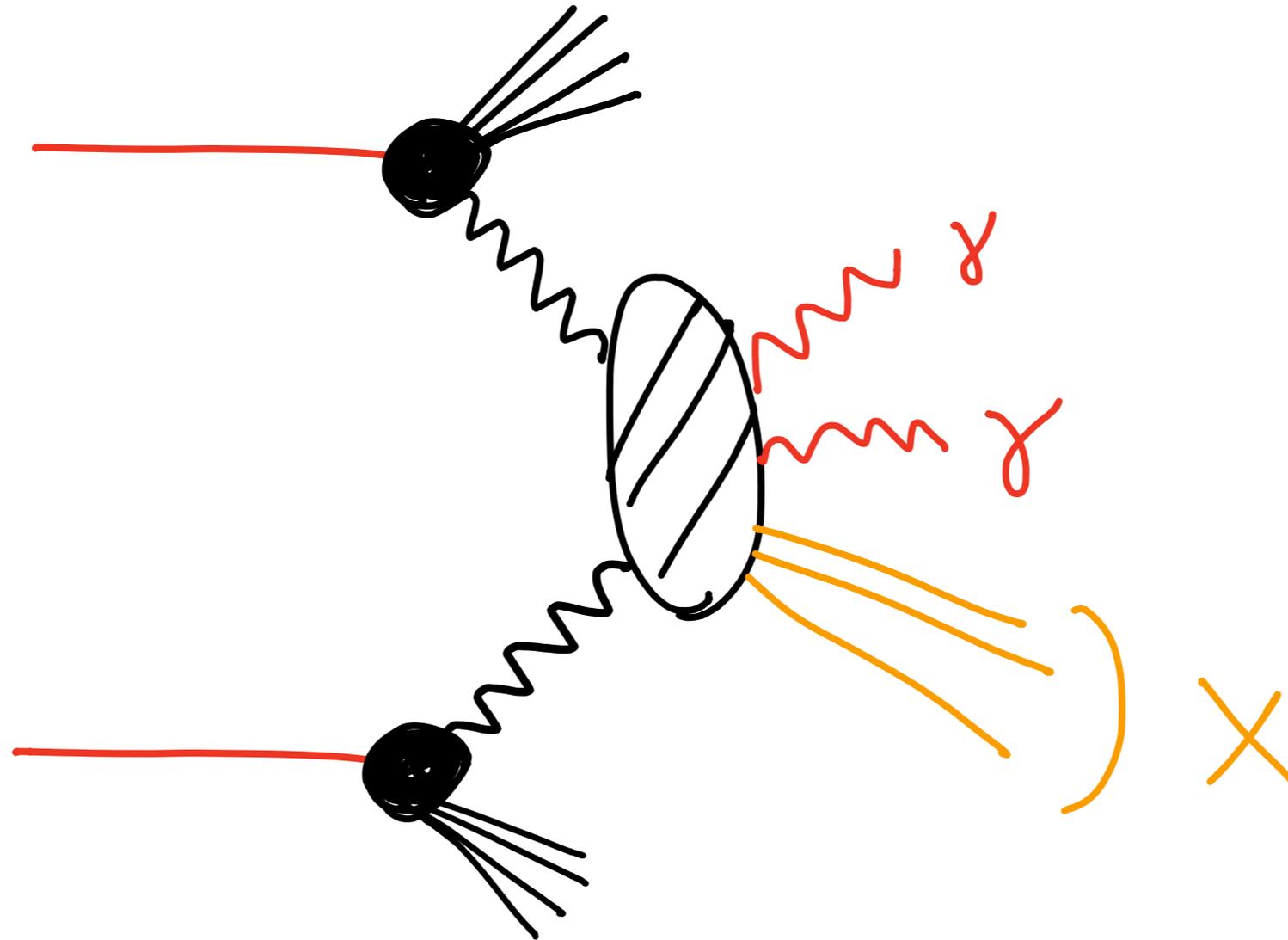
The basics



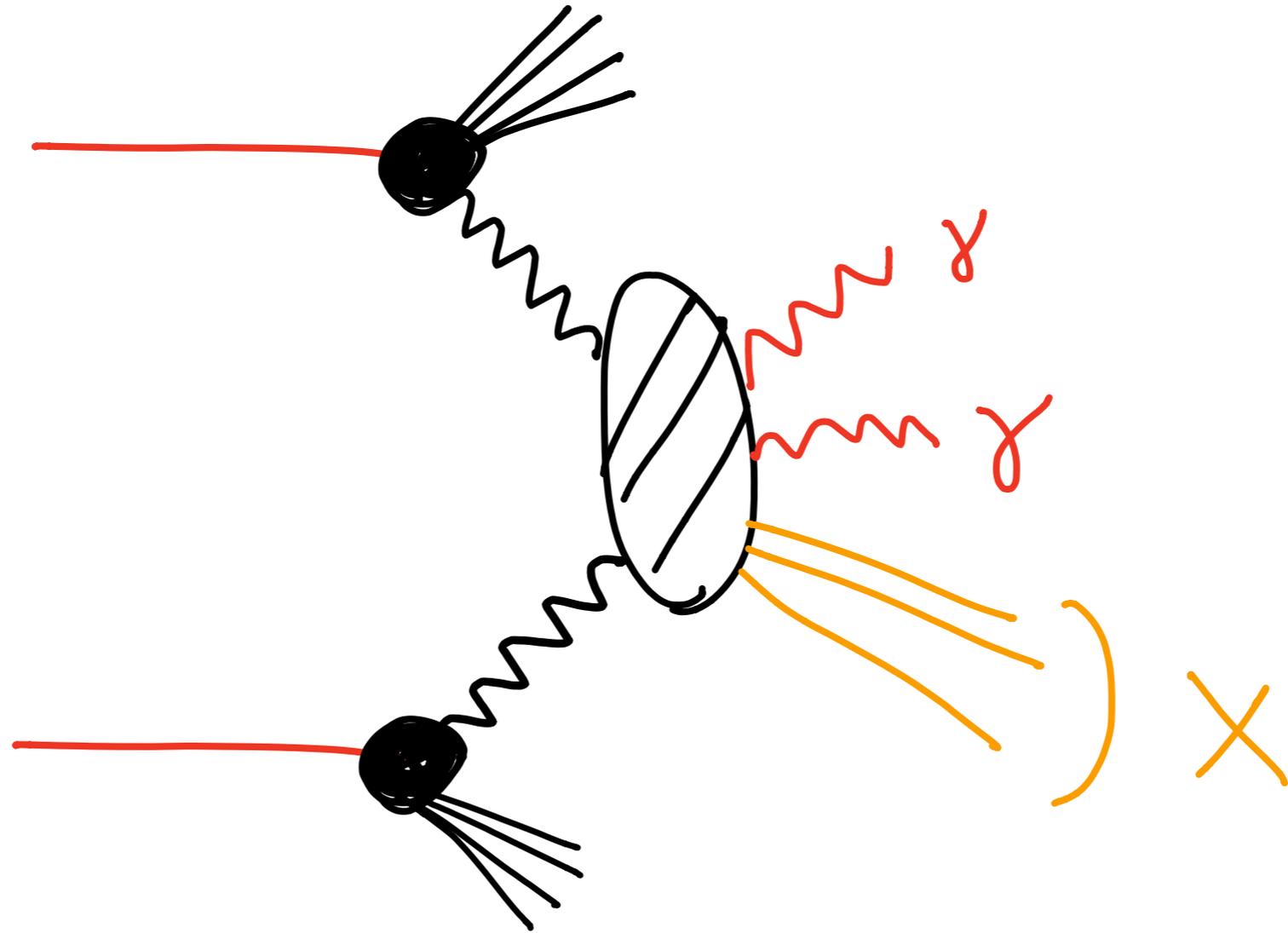
The basics



The basics

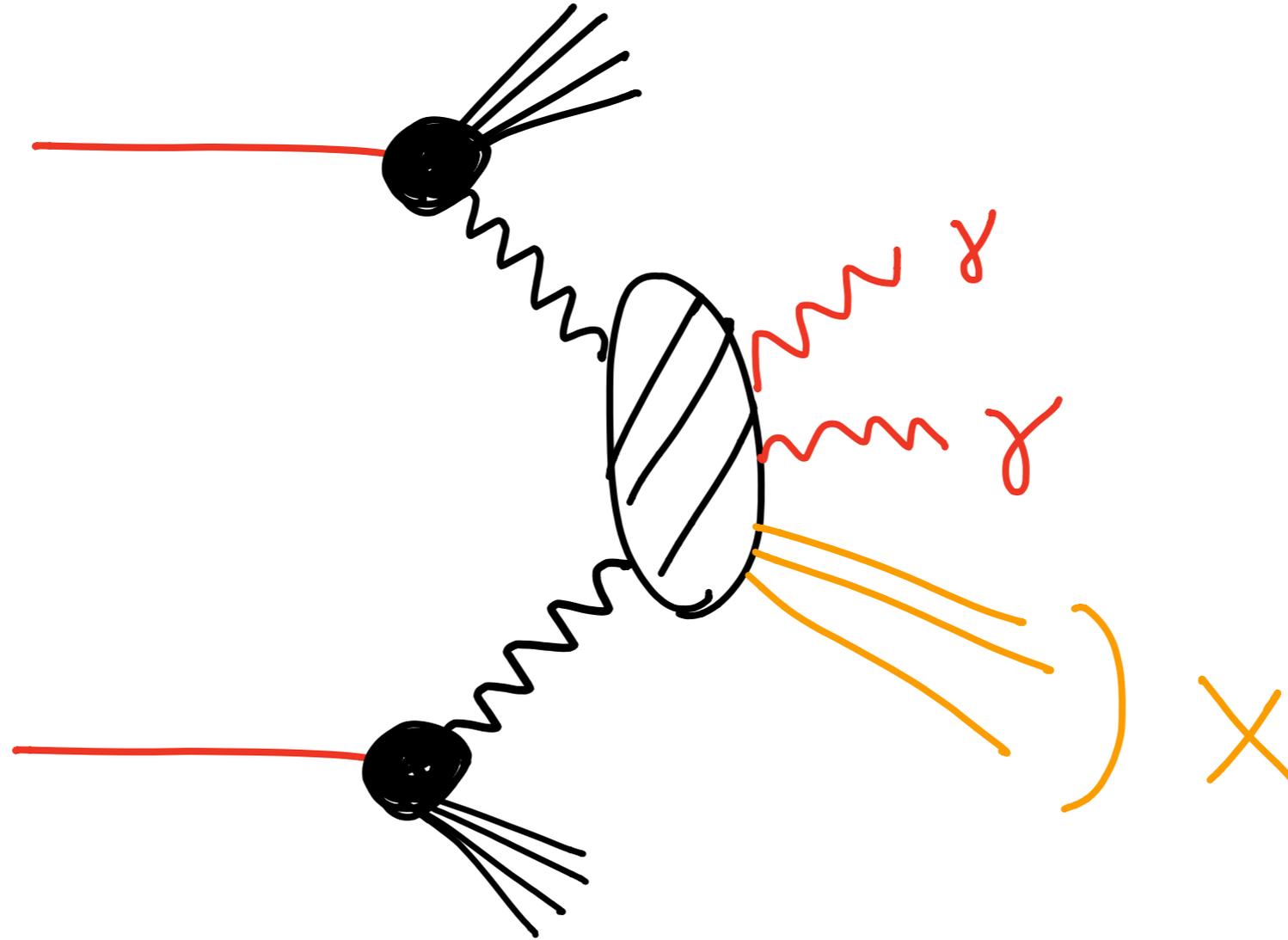


The basics



$|i\rangle = g, q, Q, \gamma, W, Z$

The basics



$|i\rangle = g, q, Q, \gamma, W, Z$

is there any X in $pp \rightarrow \gamma\gamma X$

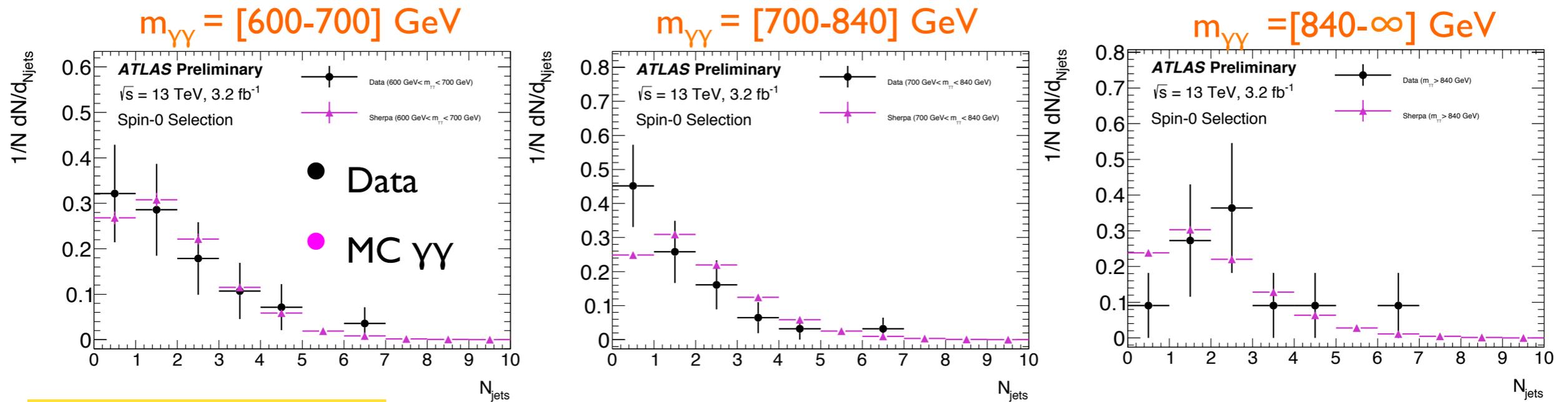
which proton constituents are initial states?

ATLAS-CONF-2016-018

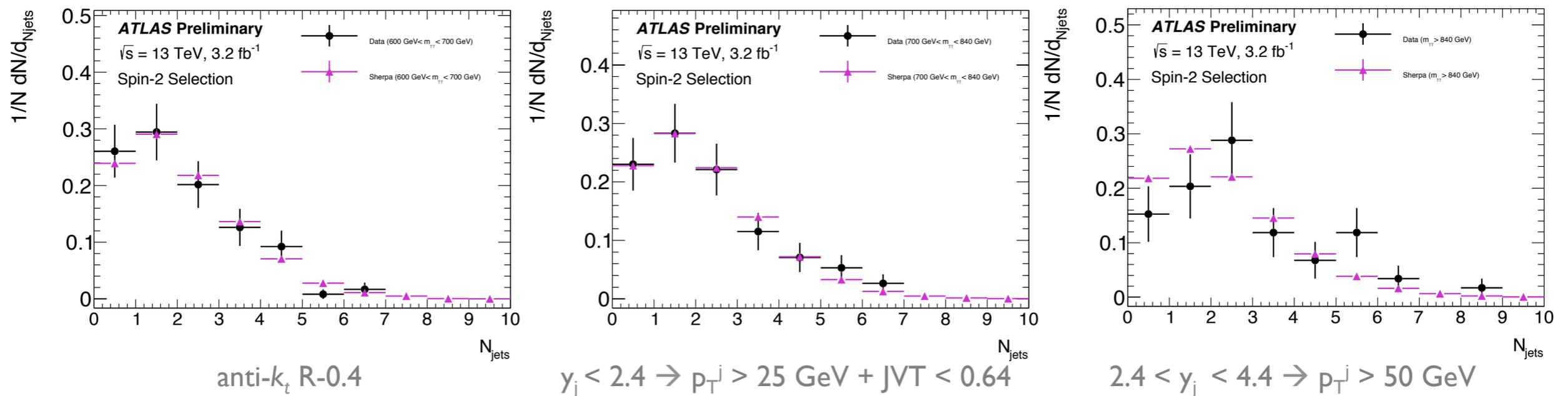
Properties of sideband and excess regions

N_{jets}

SPIN-0 ANALYSIS



SPIN-2 ANALYSIS

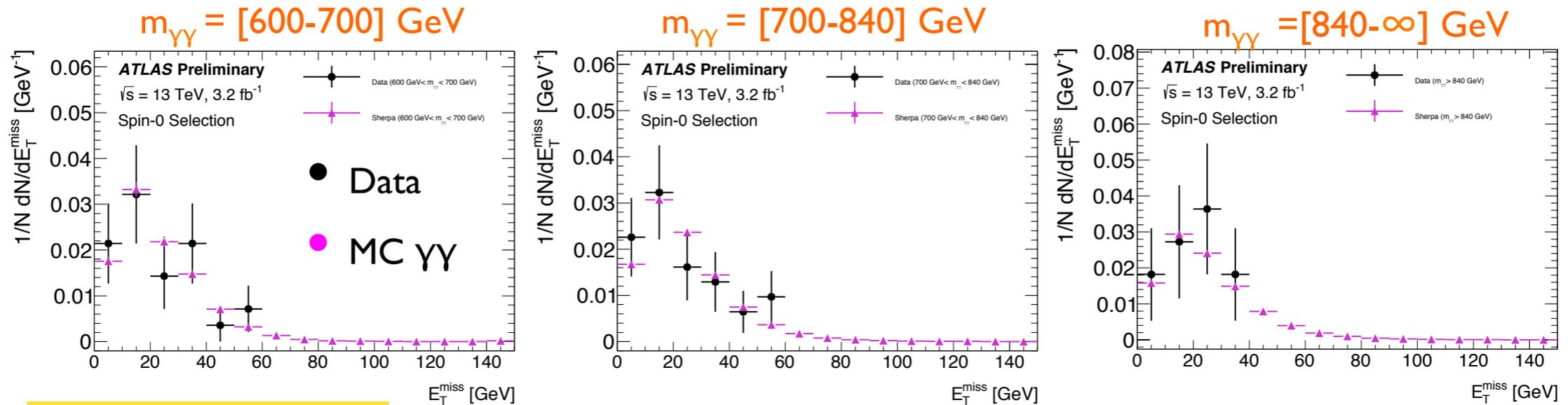


ATLAS-CONF-2016-018

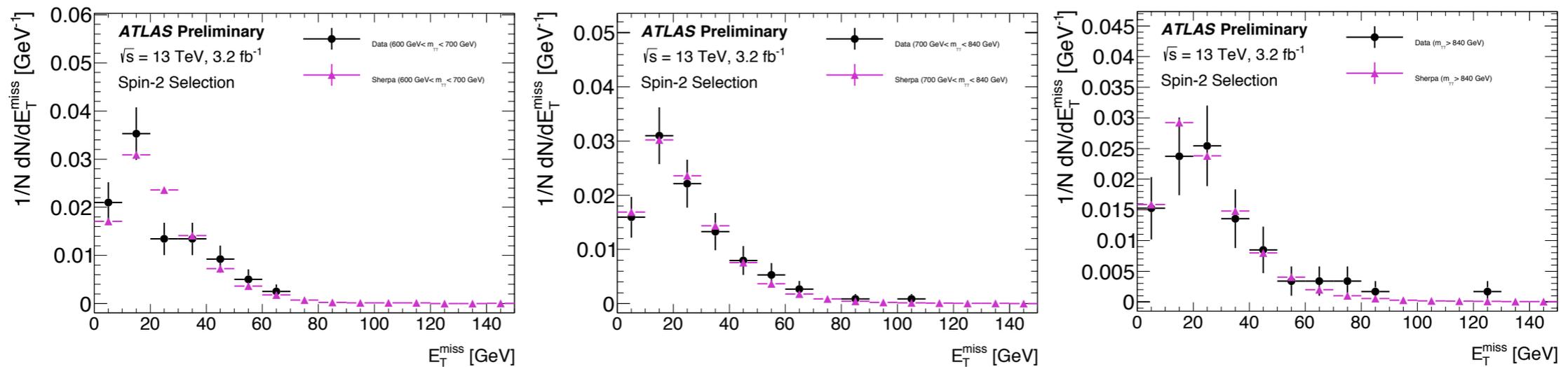
Properties of sideband and excess regions

E_T^{miss}

SPIN-0 ANALYSIS



SPIN-2 ANALYSIS

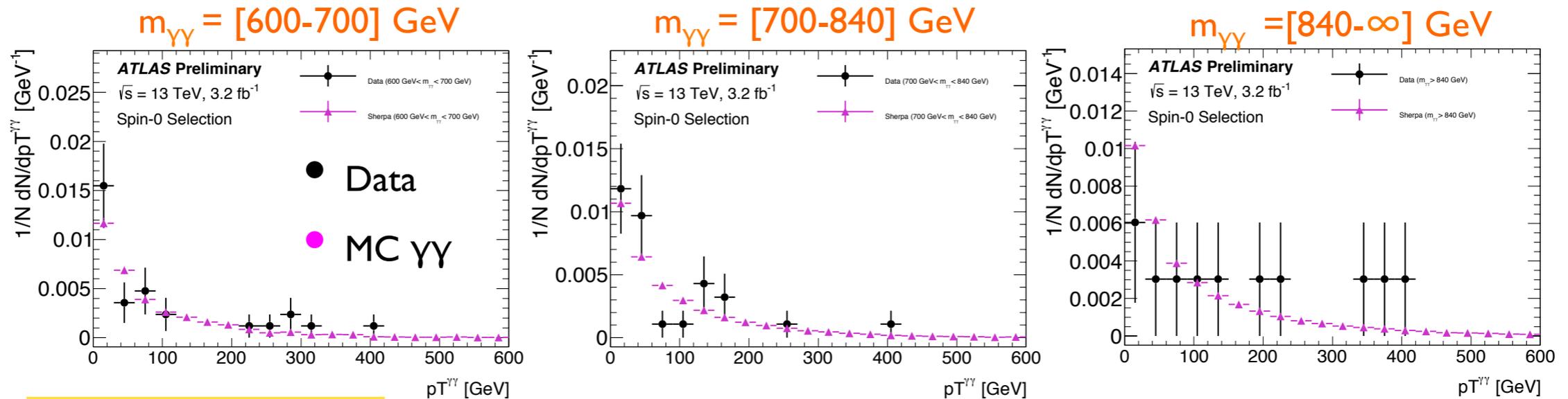


ATLAS-CONF-2016-018

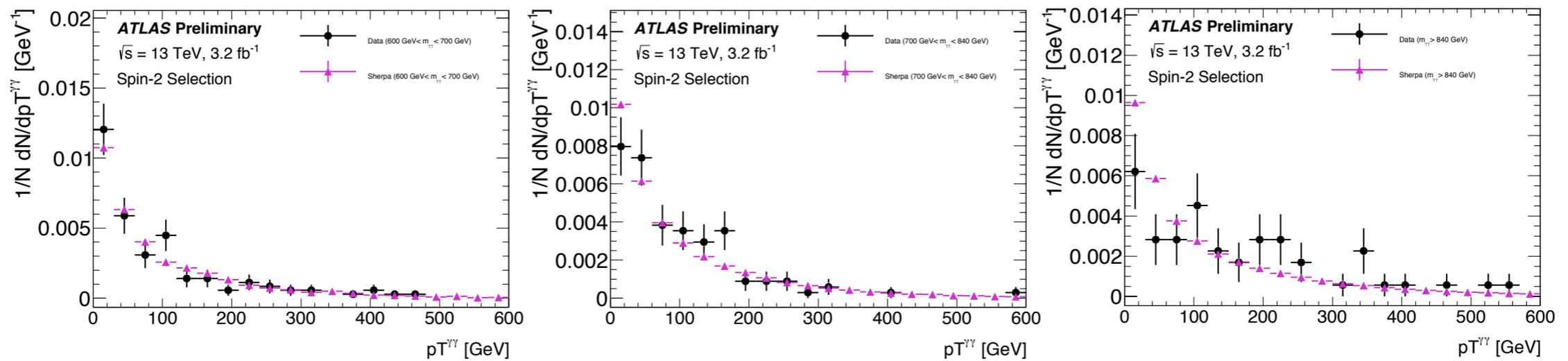
Properties of sideband and excess regions

$p_T^{\gamma\gamma}$

SPIN-0 ANALYSIS



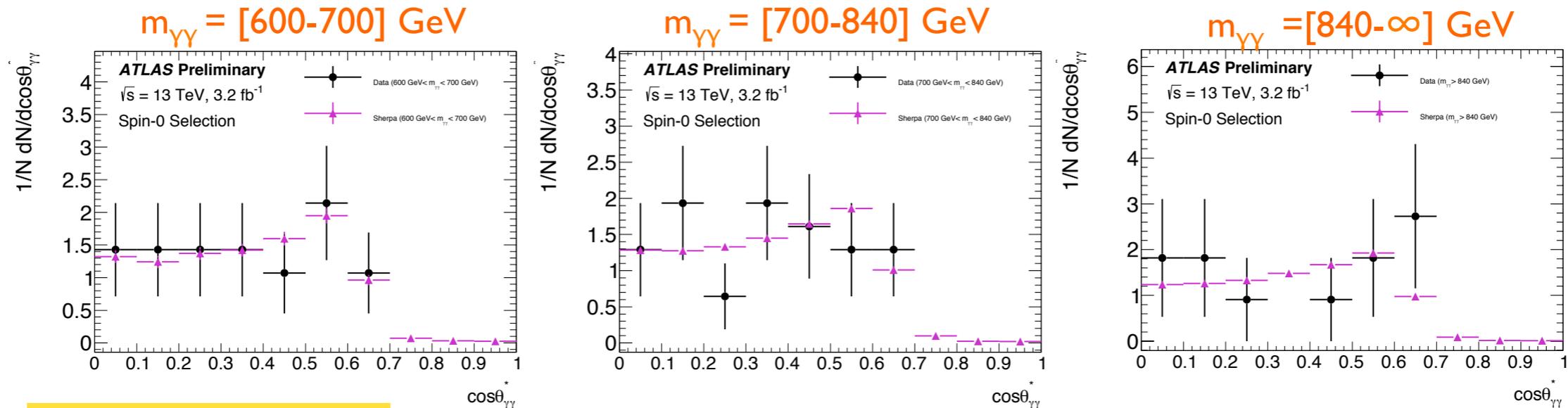
SPIN-2 ANALYSIS



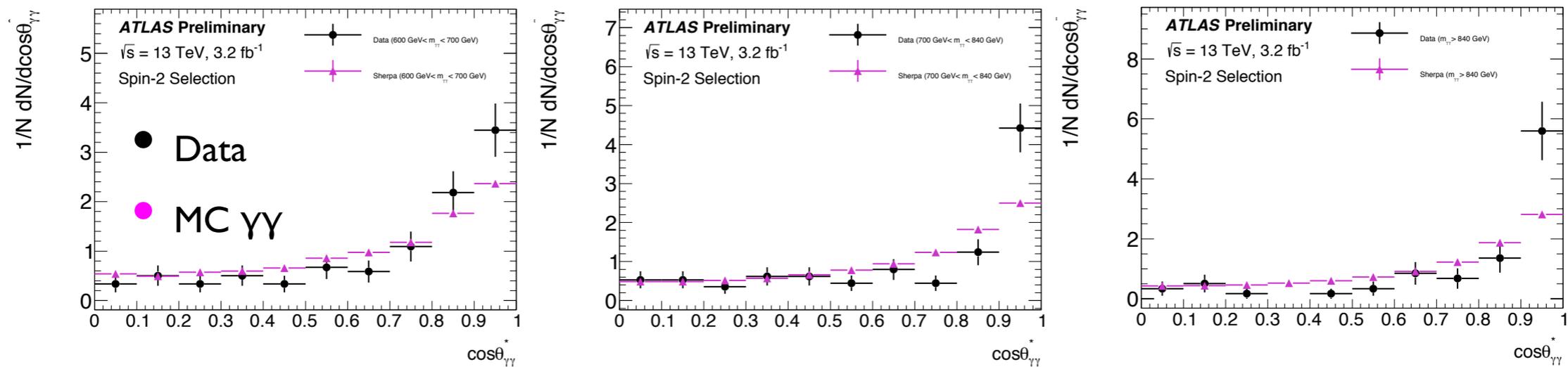
ATLAS-CONF-2016-018

Properties of sideband and excess regions $\cos\theta_{\gamma\gamma}^*$

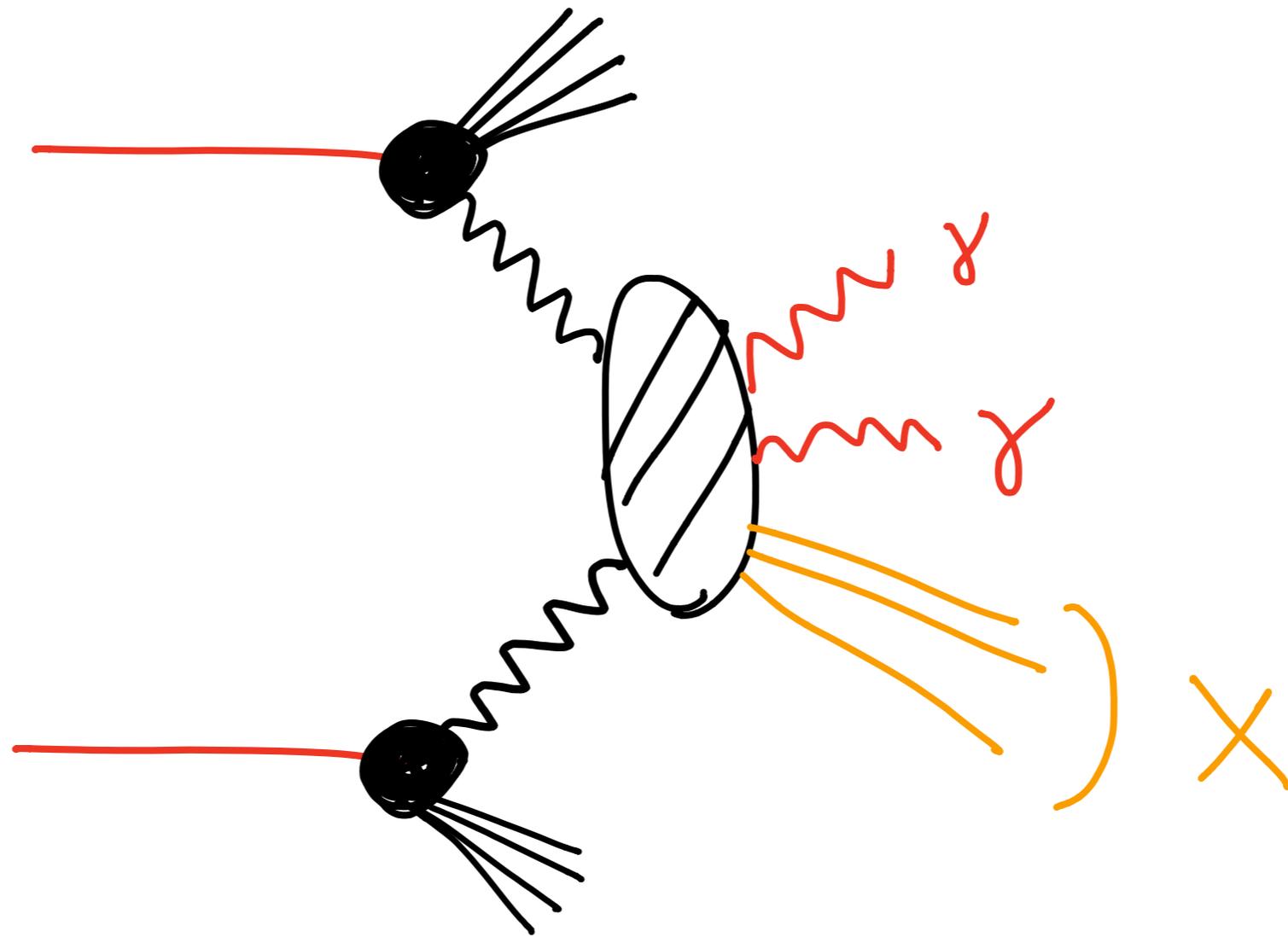
SPIN-0 ANALYSIS



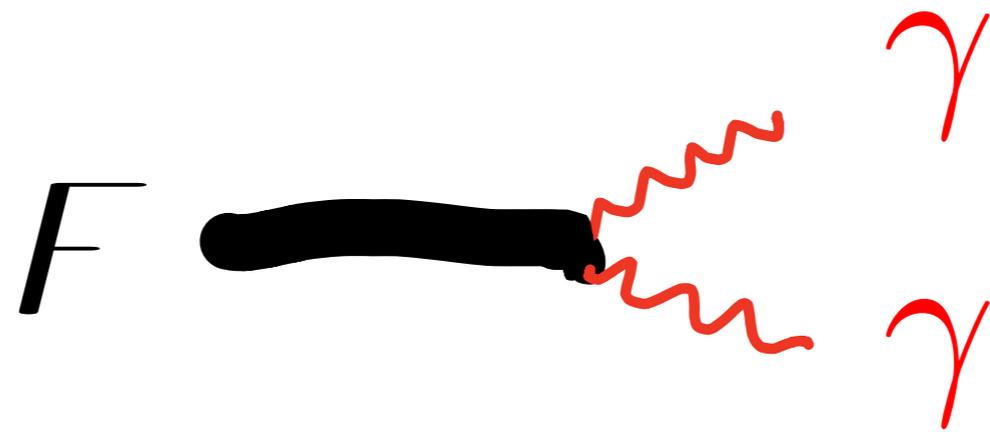
SPIN-2 ANALYSIS



simple resonance

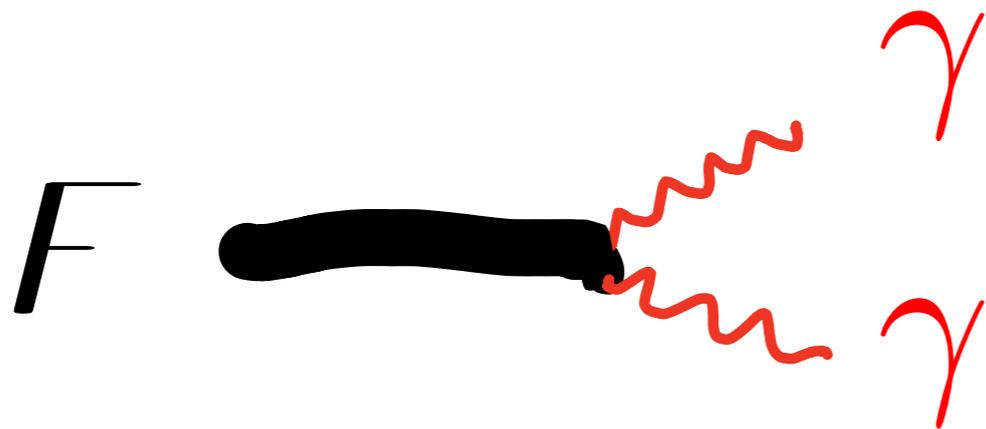


simple resonance



$$F \rightarrow \gamma\gamma$$

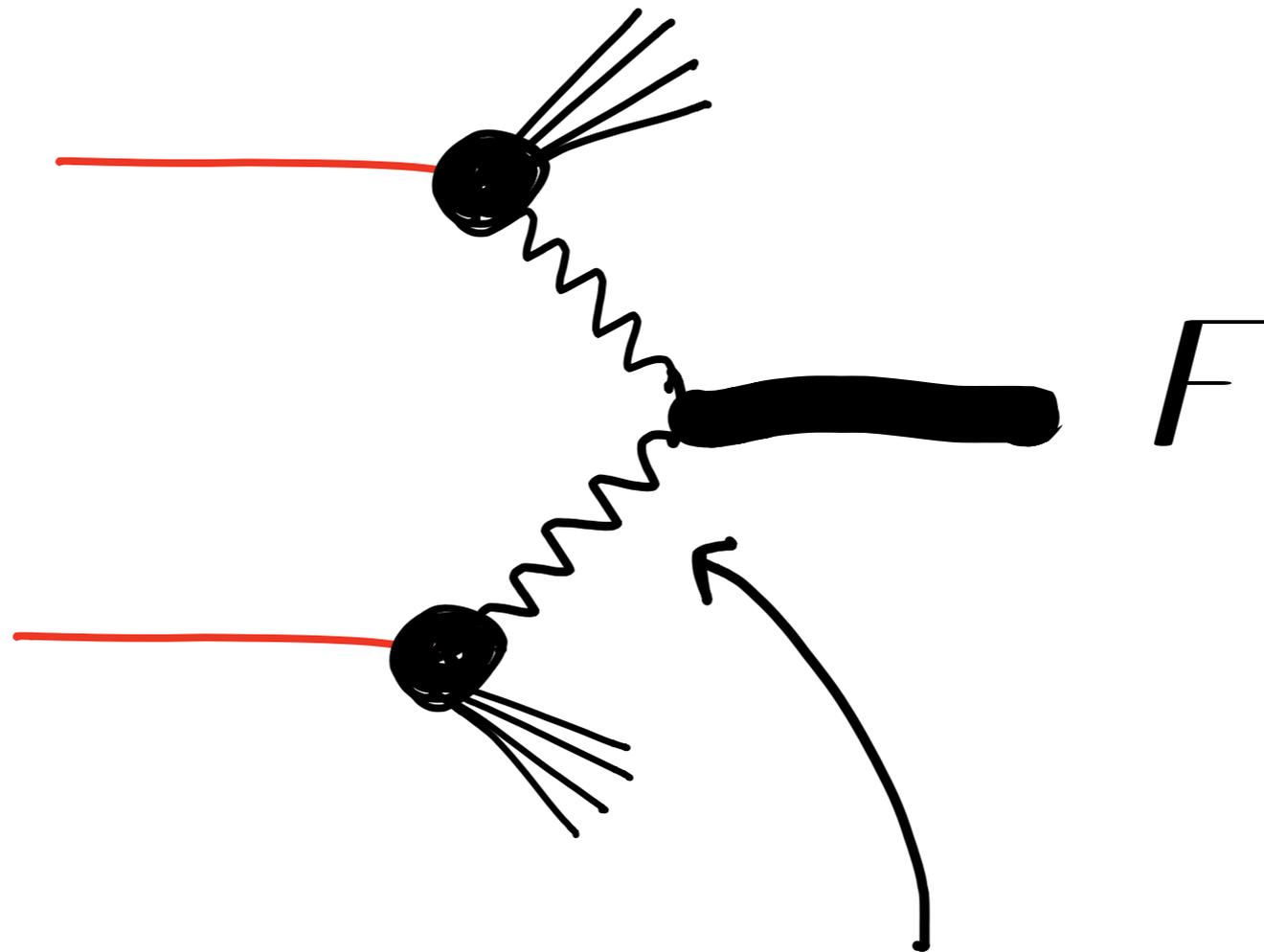
simple resonance



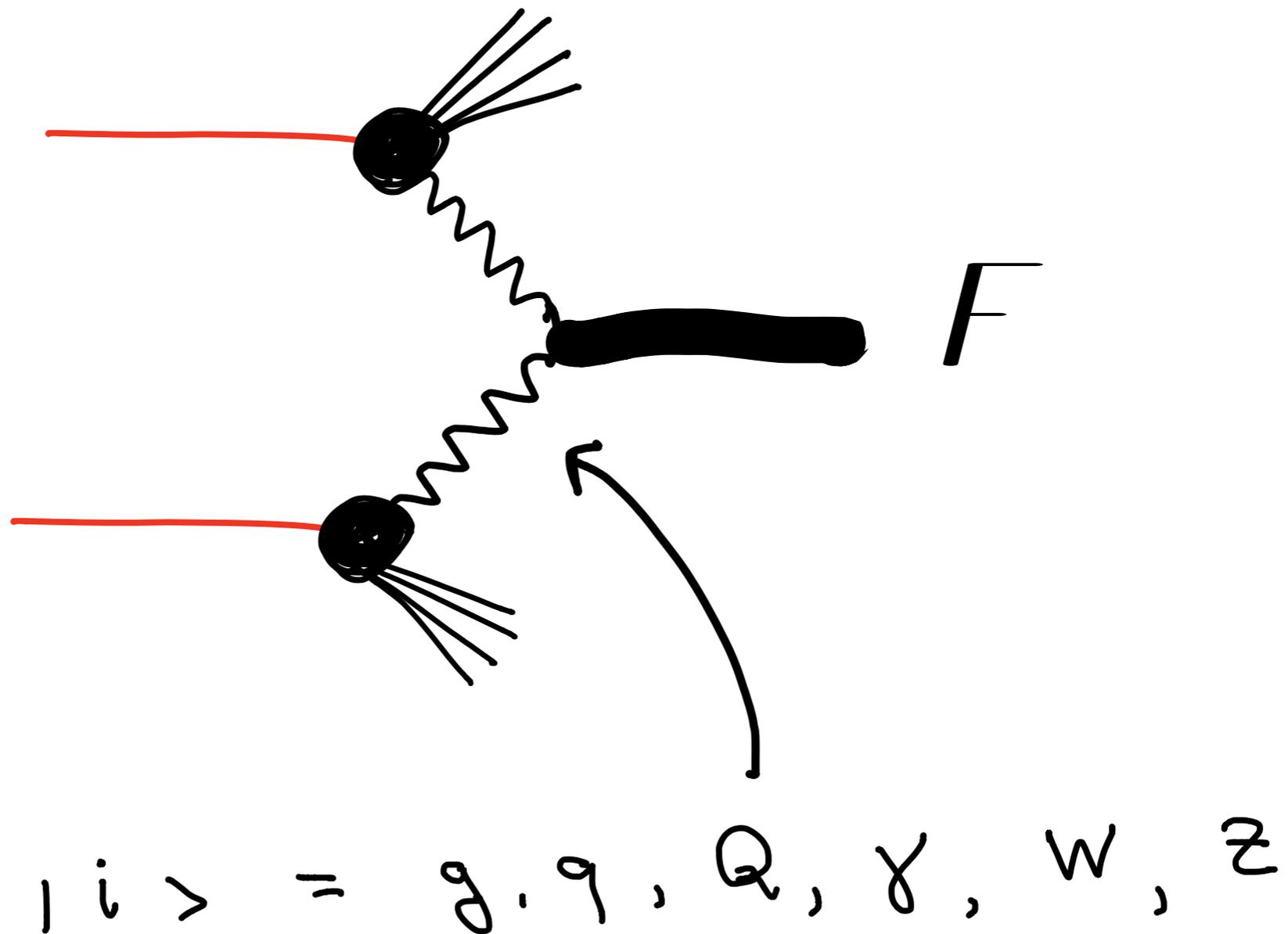
$$F \rightarrow \gamma\gamma$$

- spin
- CP (eigenstate?)
- flavor properties
- SU(2) charge

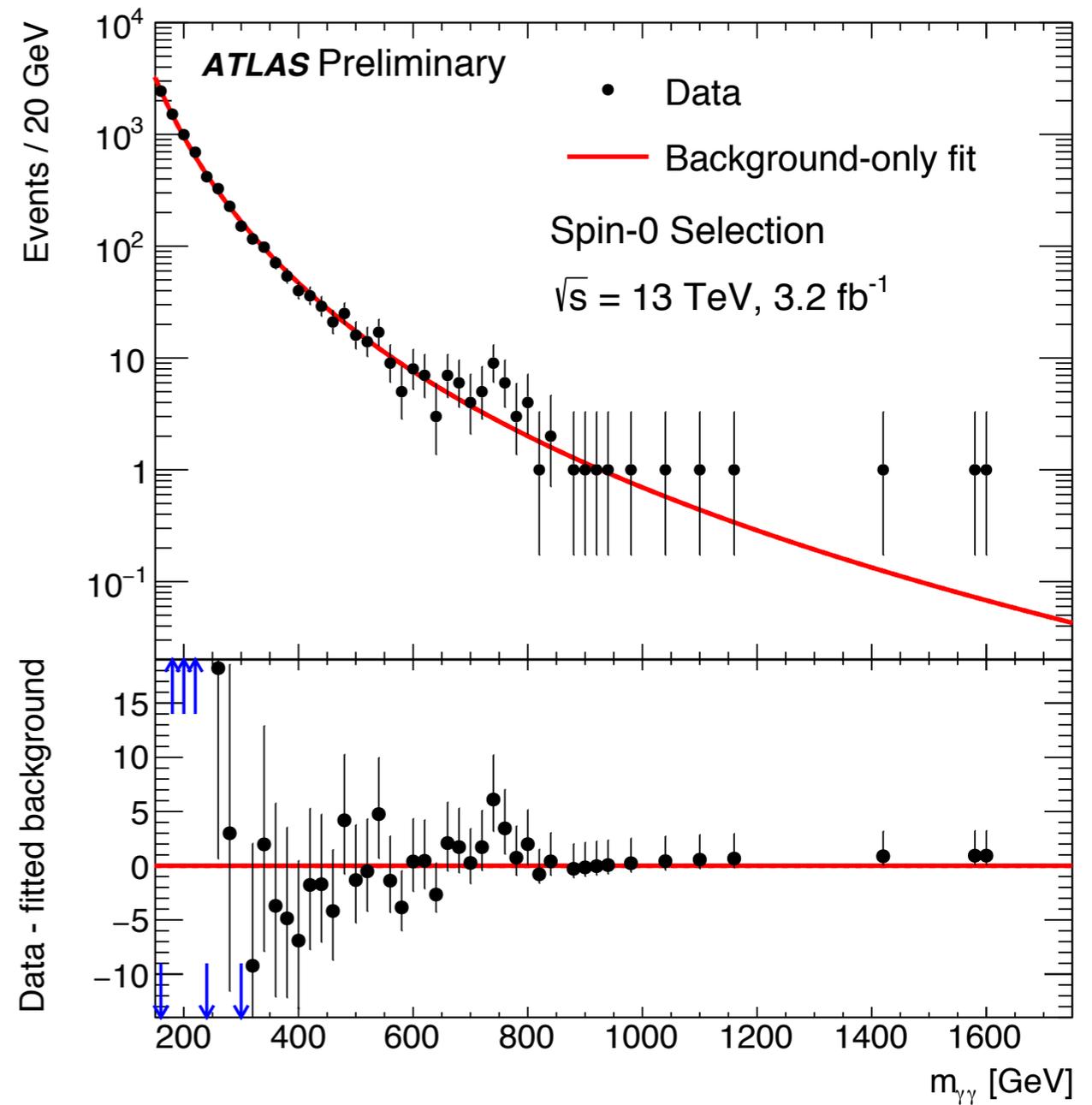
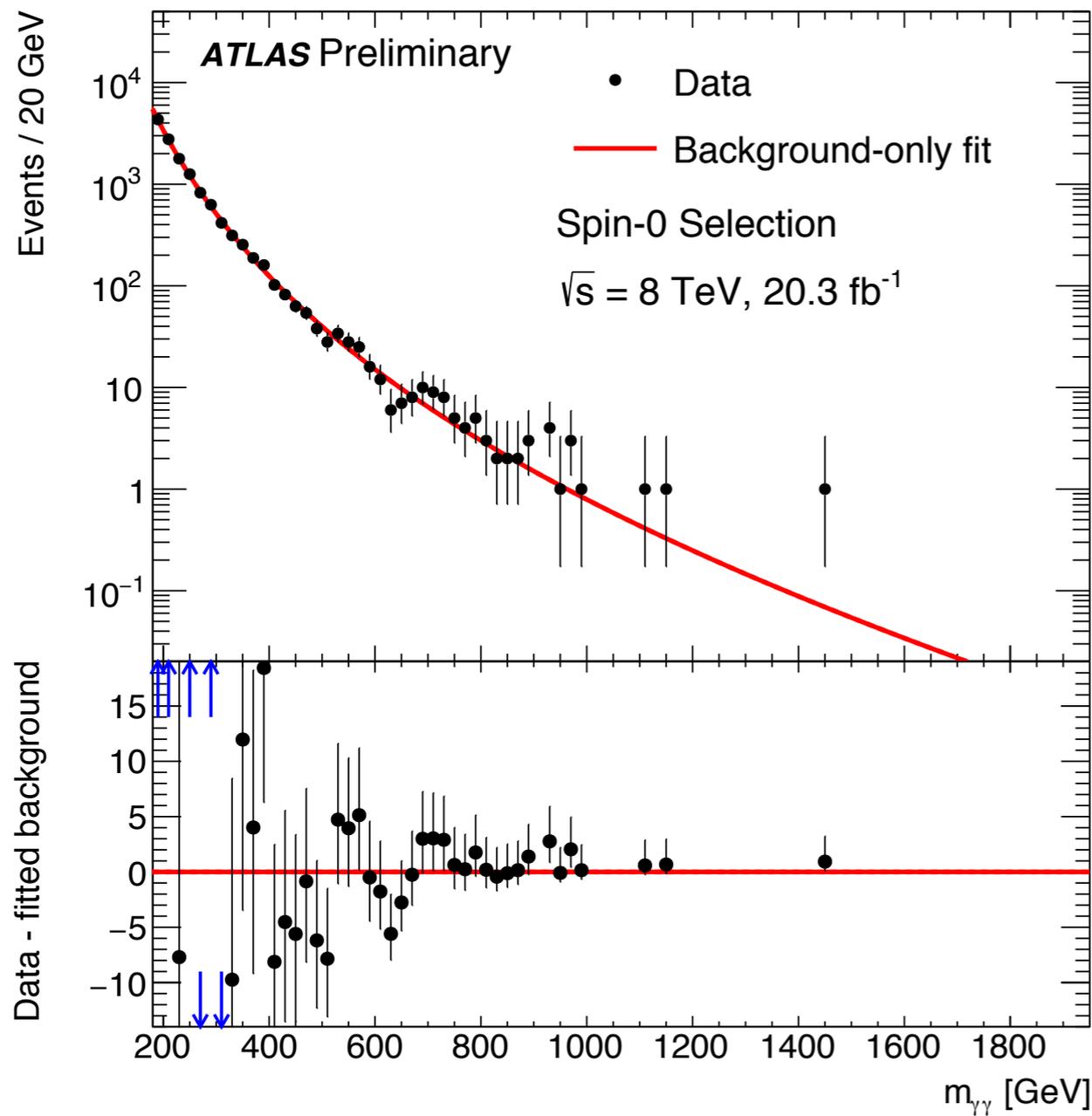
simple resonance



simple resonance

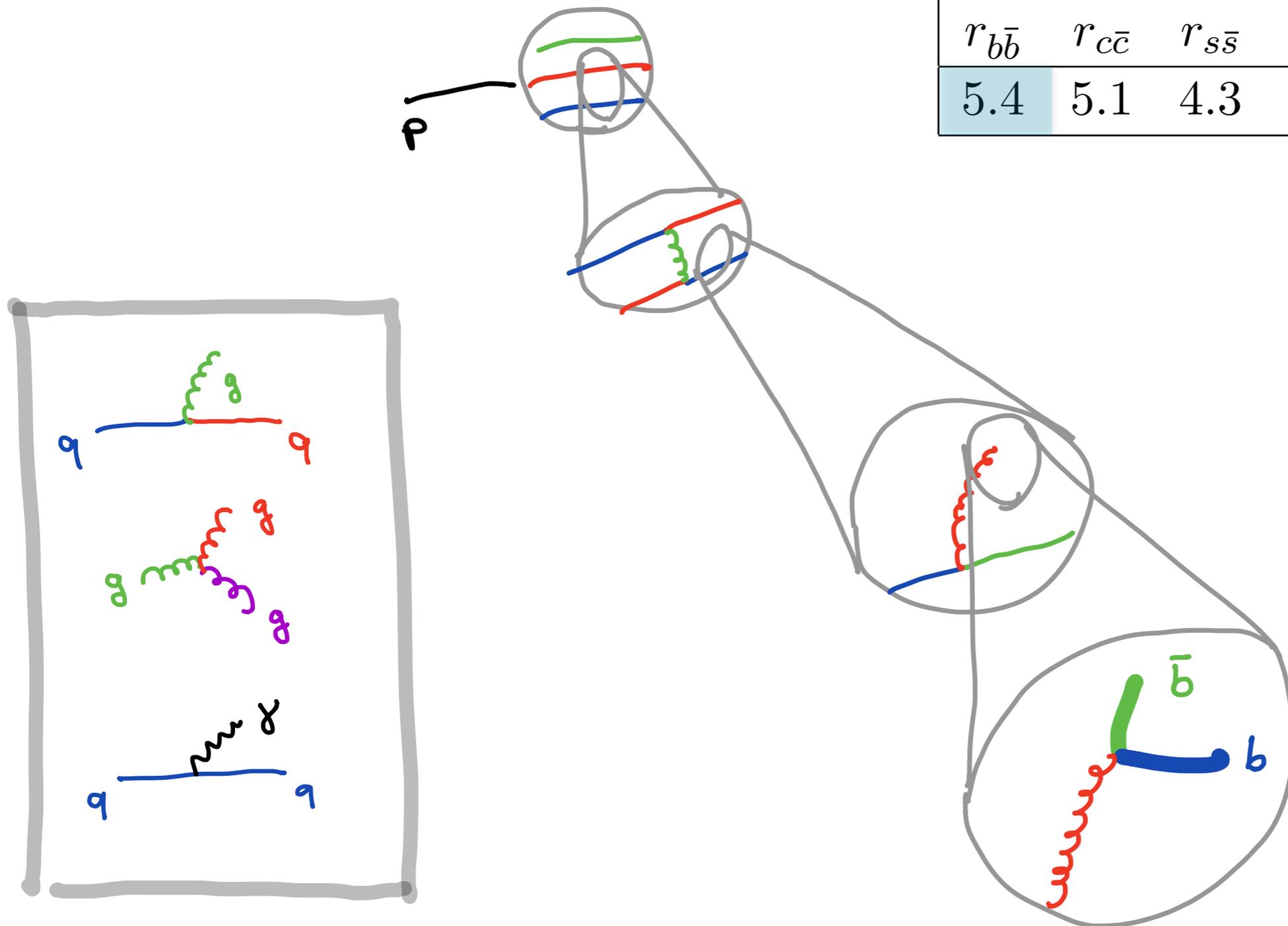


Comparison with 8 TeV



Comparison with 8 TeV

$r_{b\bar{b}}$	$r_{c\bar{c}}$	$r_{s\bar{s}}$	$r_{d\bar{d}}$	$r_{u\bar{u}}$	r_{gg}	$r_{\gamma\gamma}$
5.4	5.1	4.3	2.7	2.5	4.7	1.9



Spin

1602.02793

1603.04248

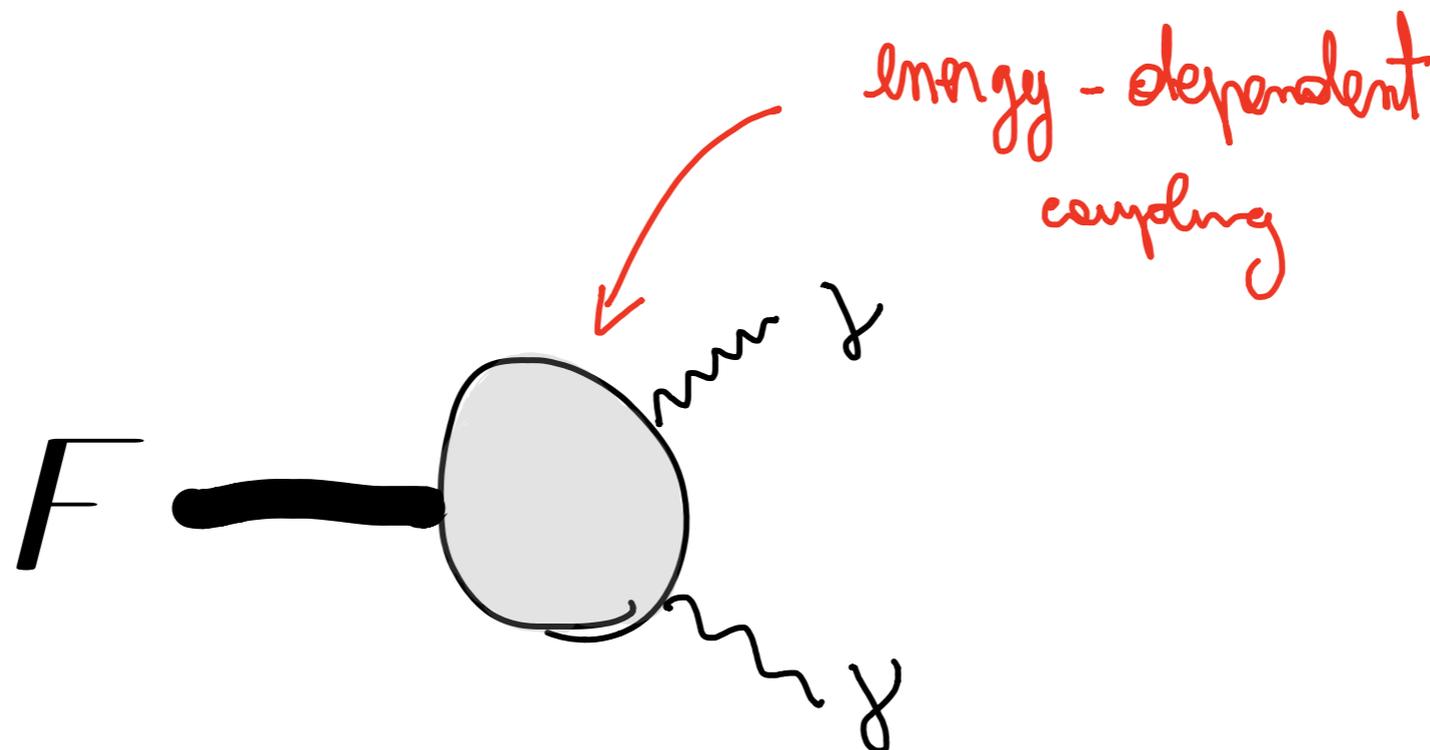
1604.06948

Spin-0: dim-5 $F \cdot F^{\mu\nu} F_{\mu\nu}$

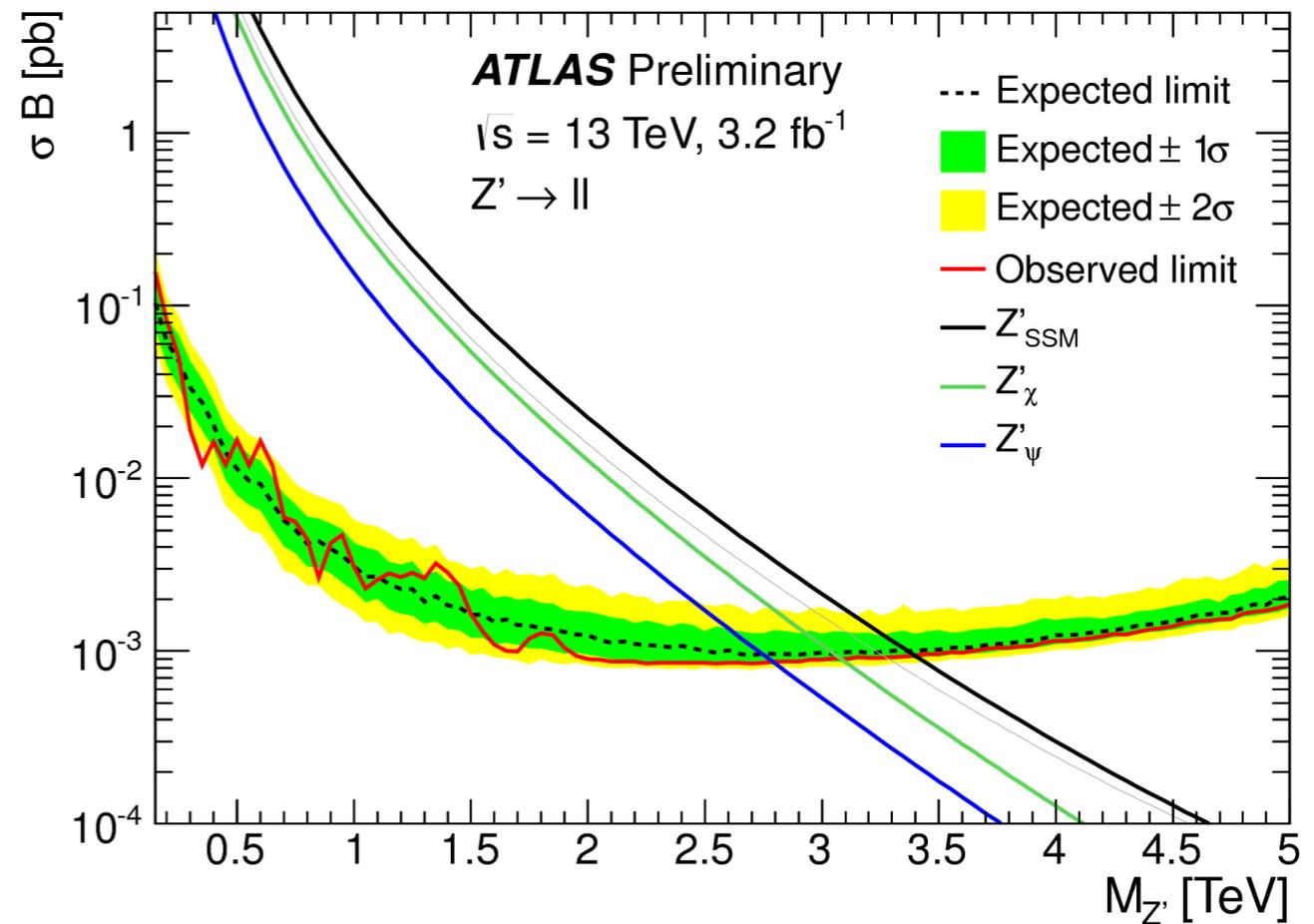
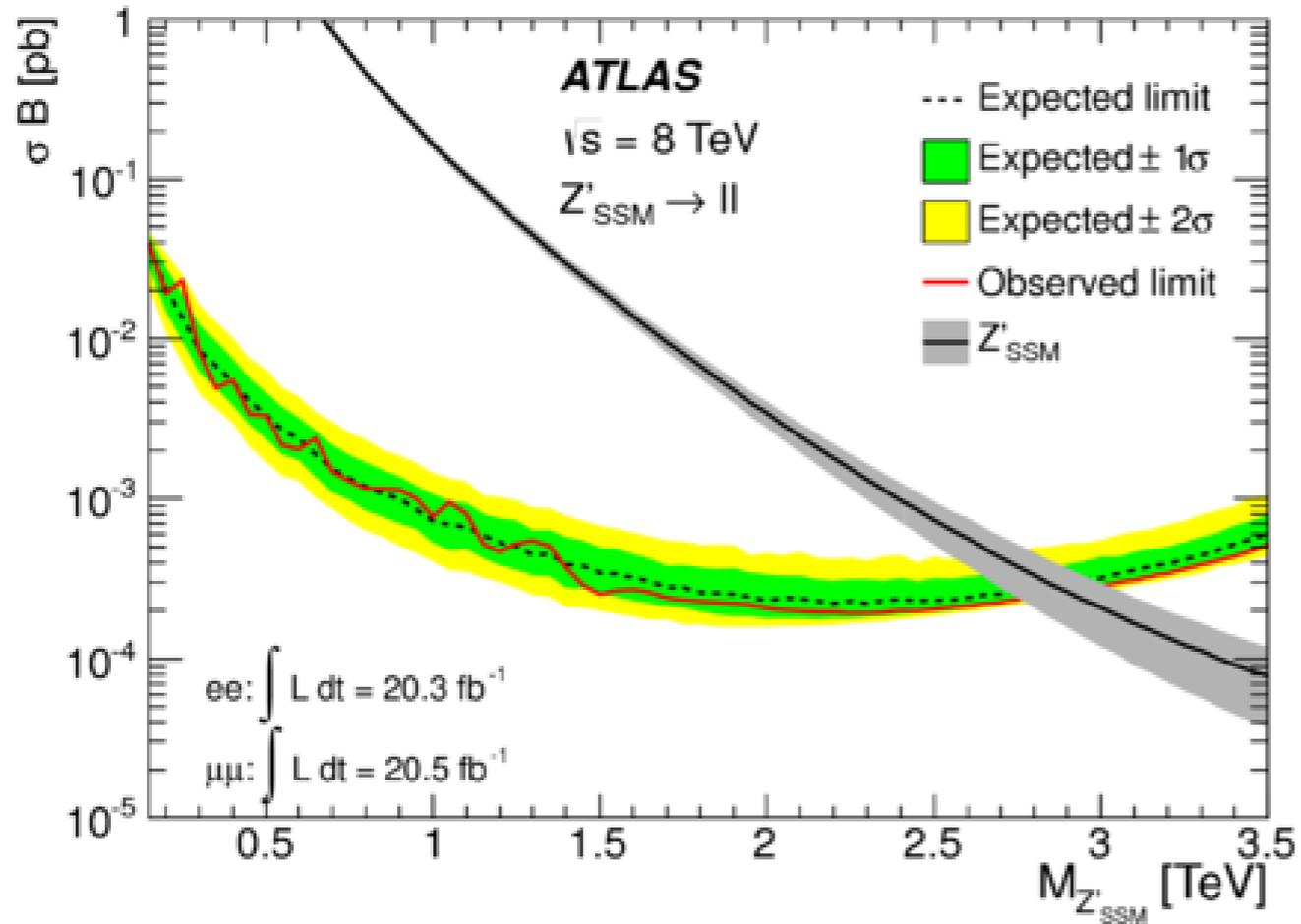
Spin-1: Landau-Yang

Spin-2: Tree-level coupling (e.g. $h_{\mu\nu} \cdot T^{\mu\nu}$)

...



Spin-2 and leptons

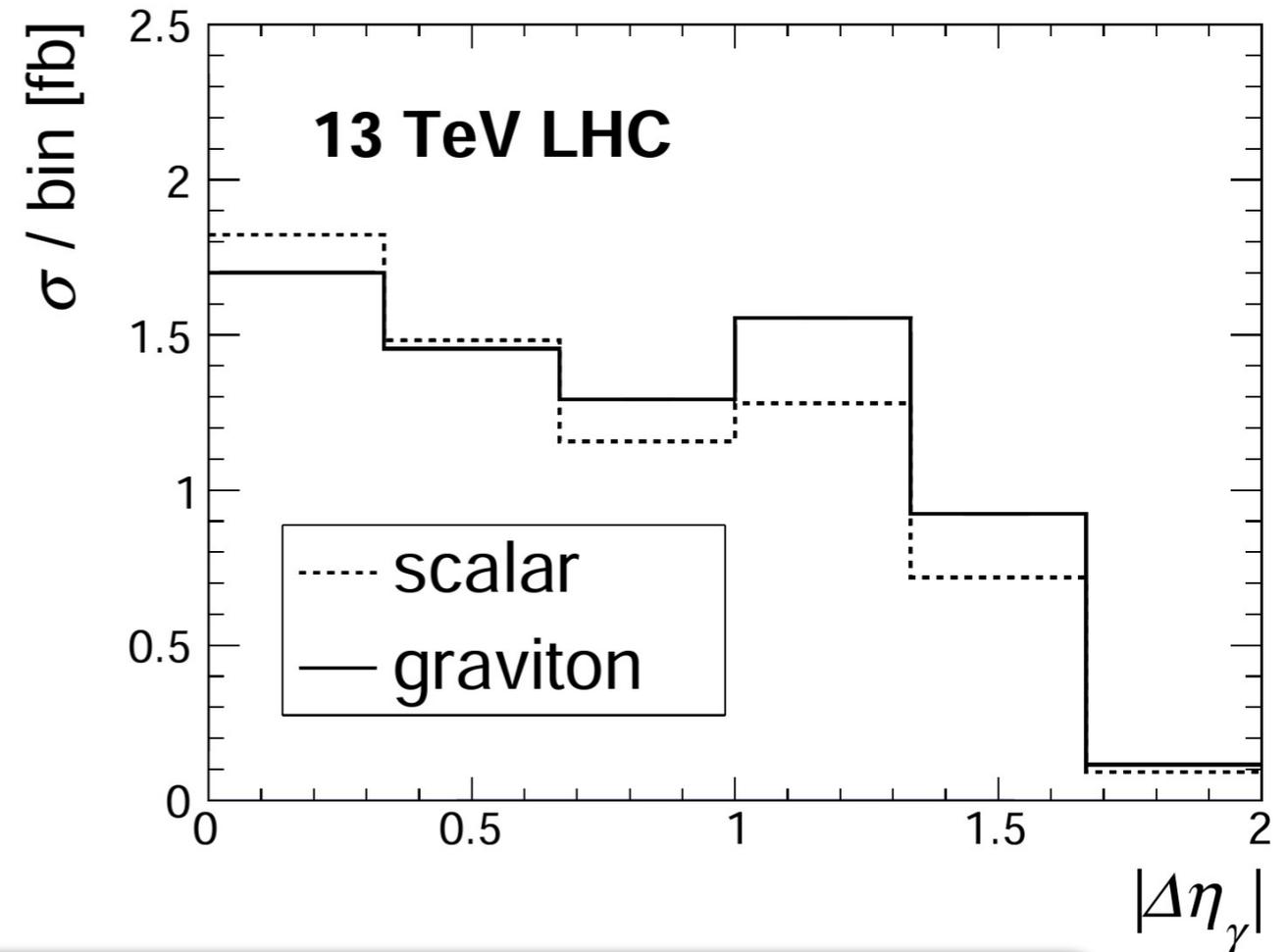
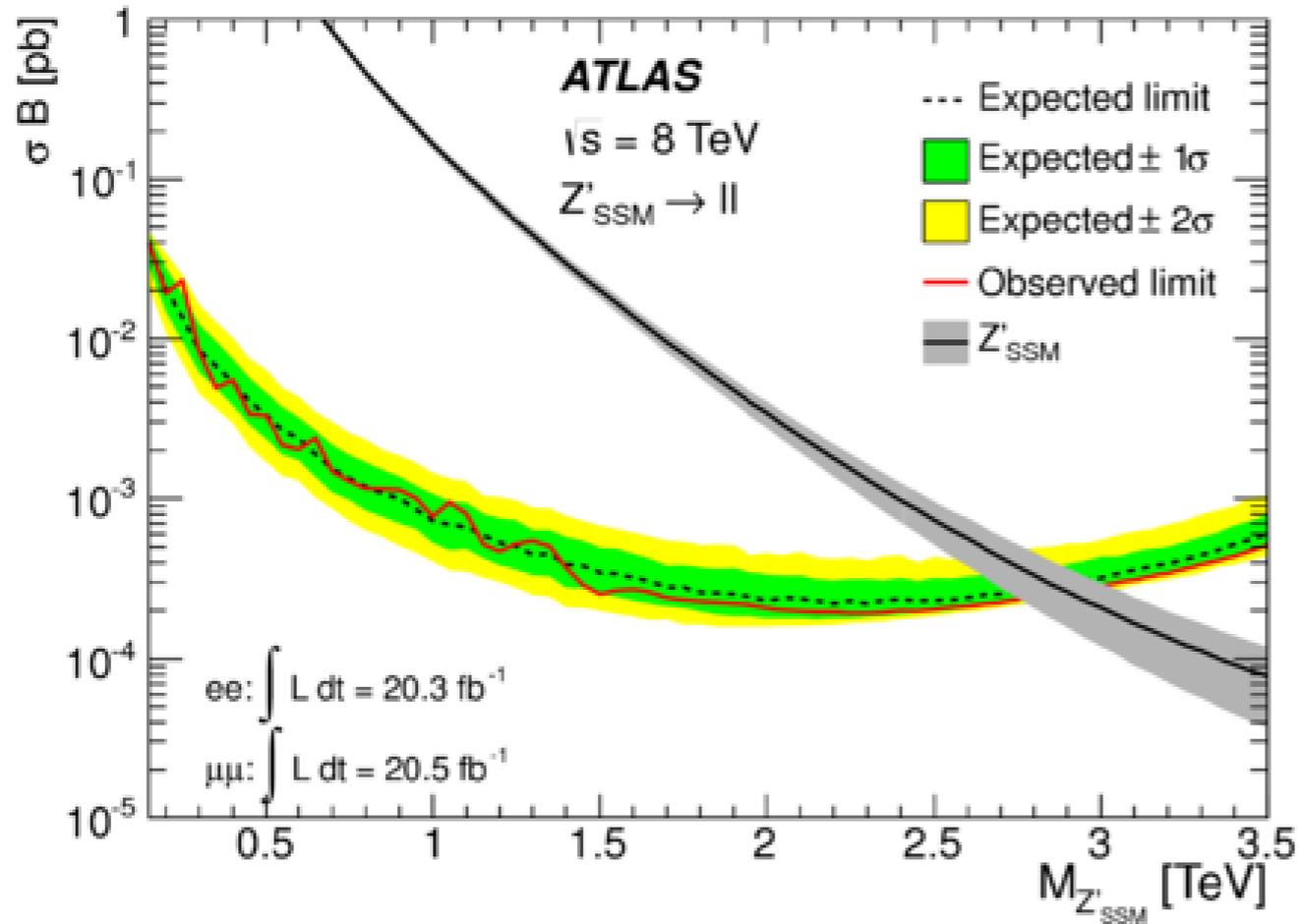


the absence of a signal in leptons points to a non-universally coupled spin-2

1603.08913
 1603.06980
 1603.08250
 1602.02793
 1603.05574

strong interactions at the TeV are typical in such frameworks

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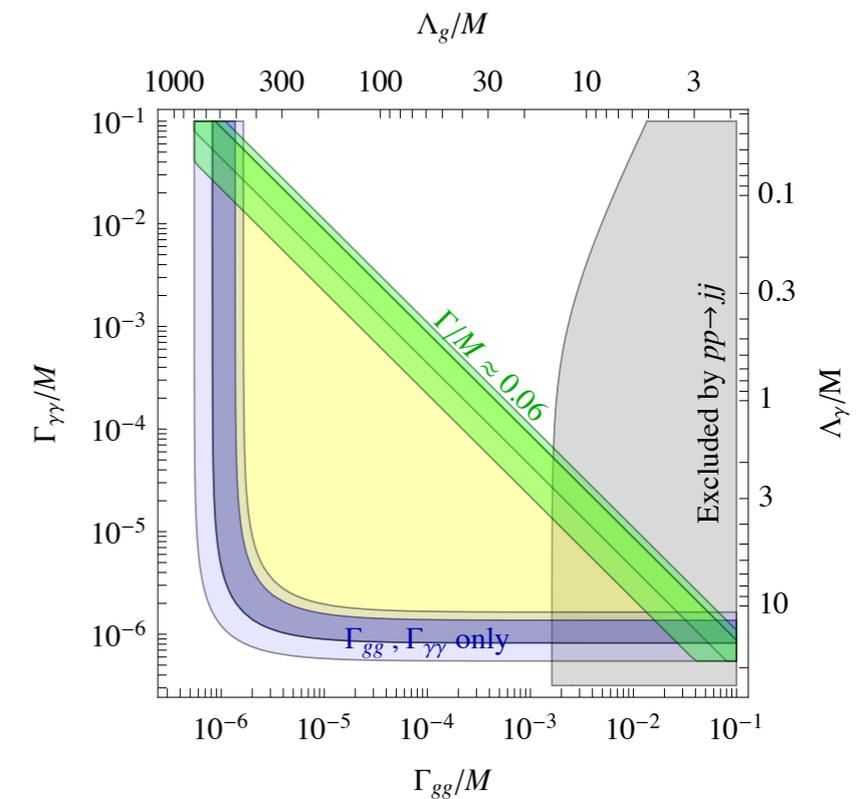
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Spin-0

focus on a weak singlet for simplicity

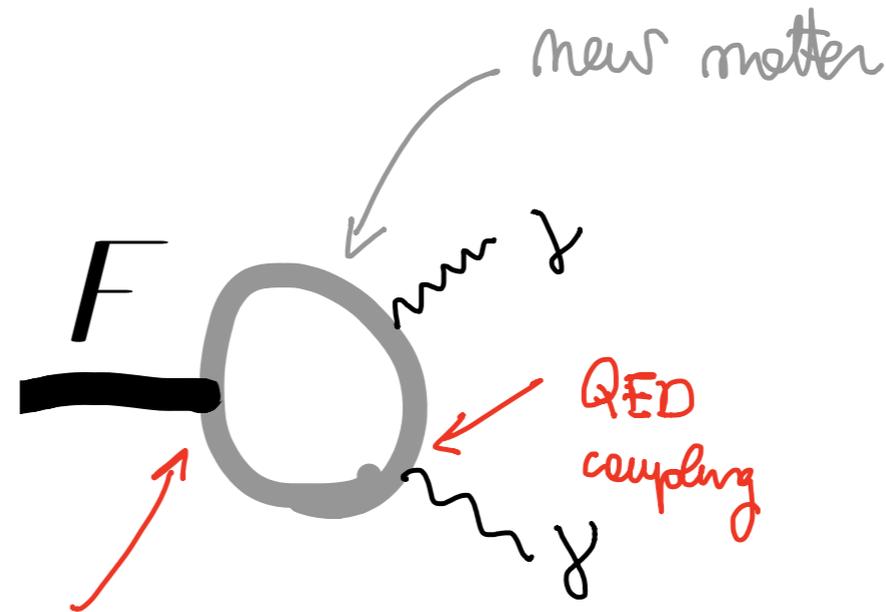
F

$$\sigma(pp \rightarrow S \rightarrow \gamma\gamma) = \frac{2J+1}{M\Gamma_S} \left[\sum_{\wp} C_{\wp\bar{\wp}} \Gamma(S \rightarrow \wp\bar{\wp}) \right] \Gamma(S \rightarrow \gamma\gamma)$$

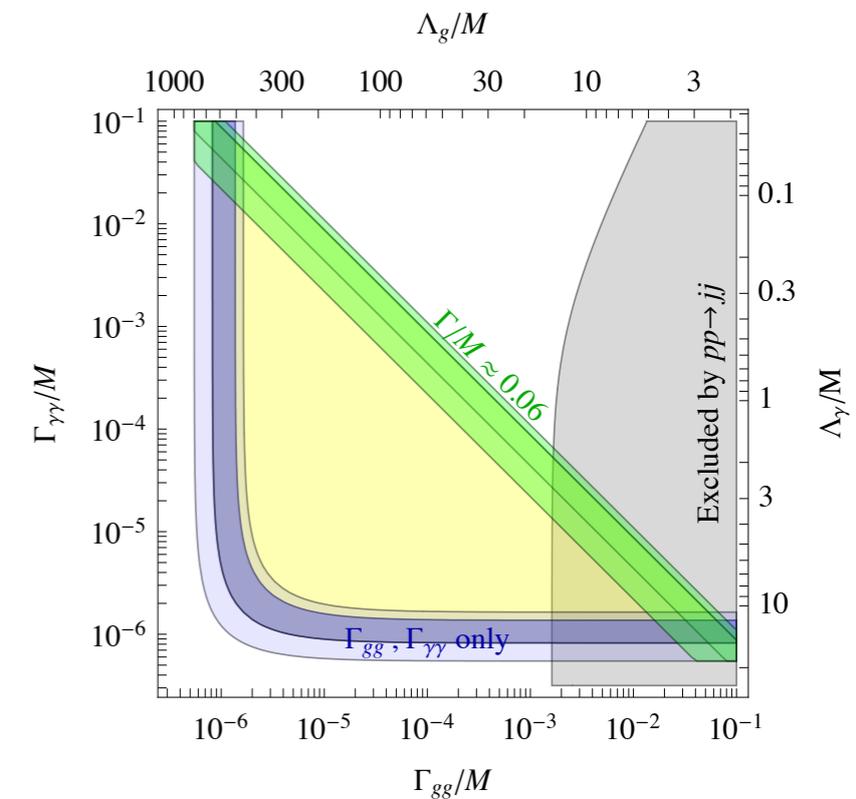


Spin-0

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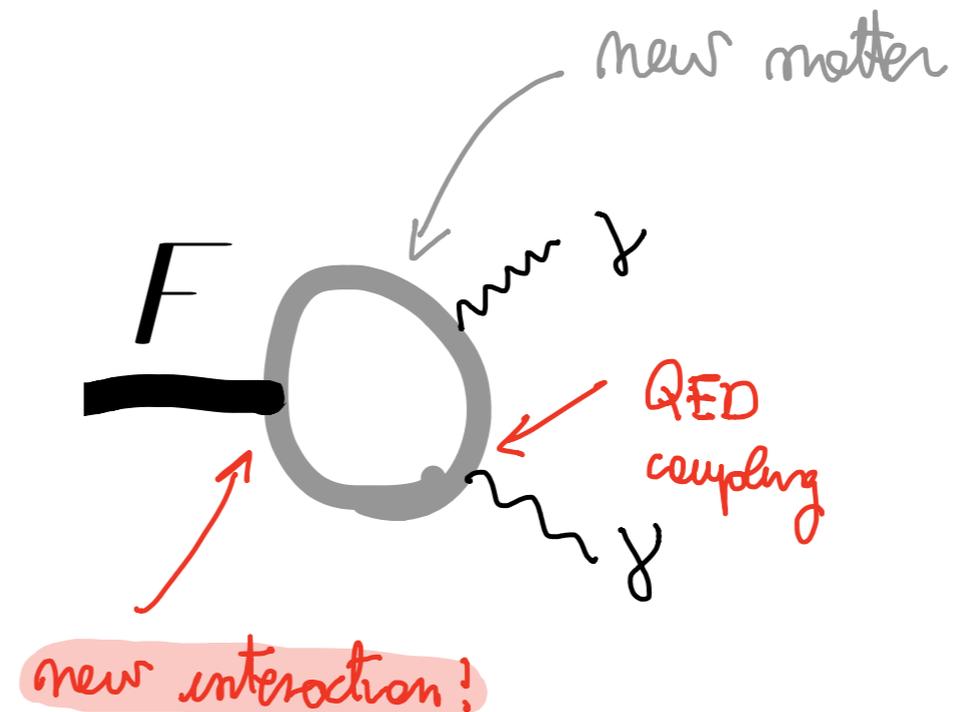


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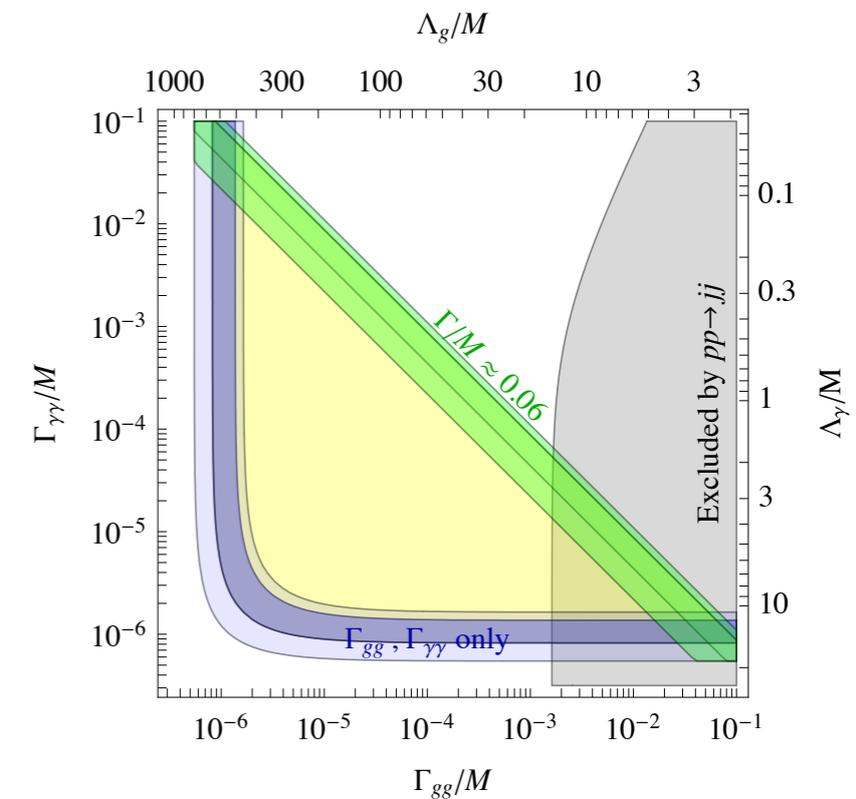


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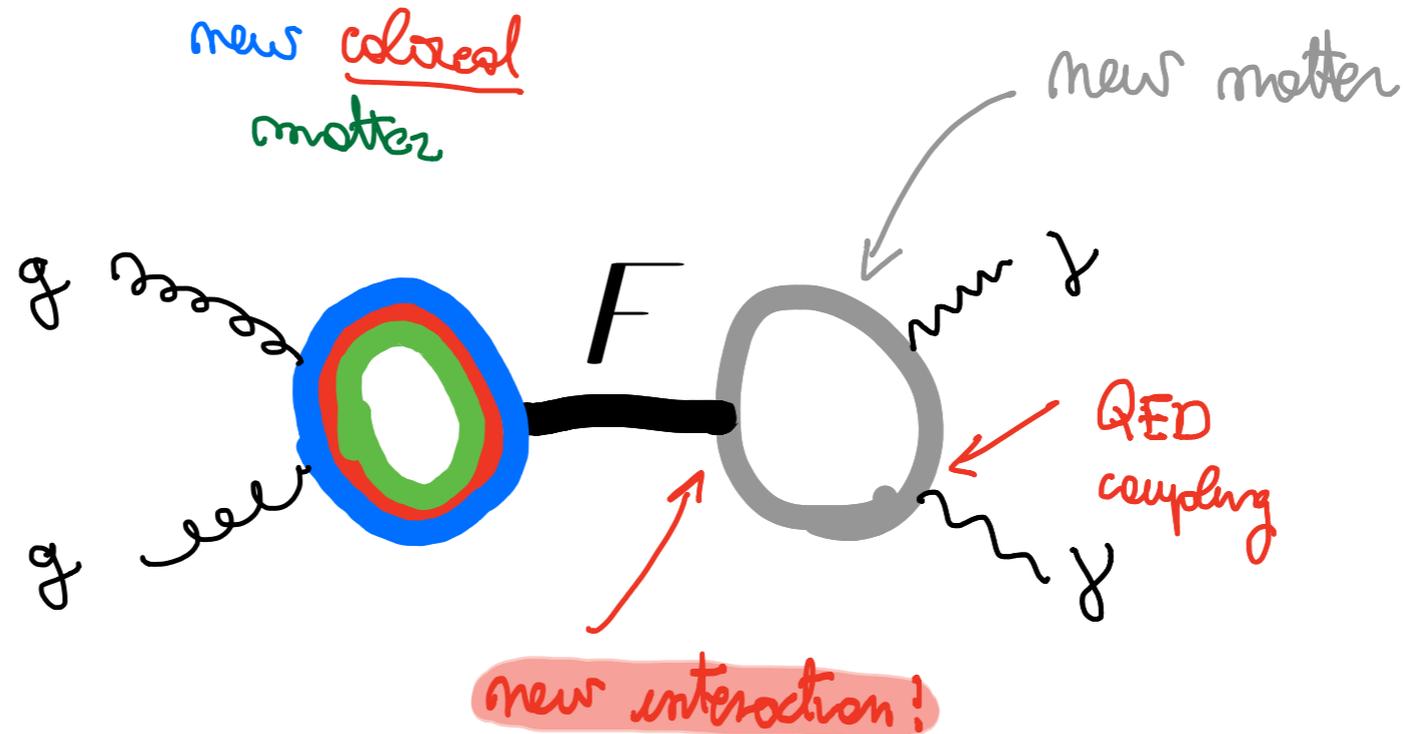


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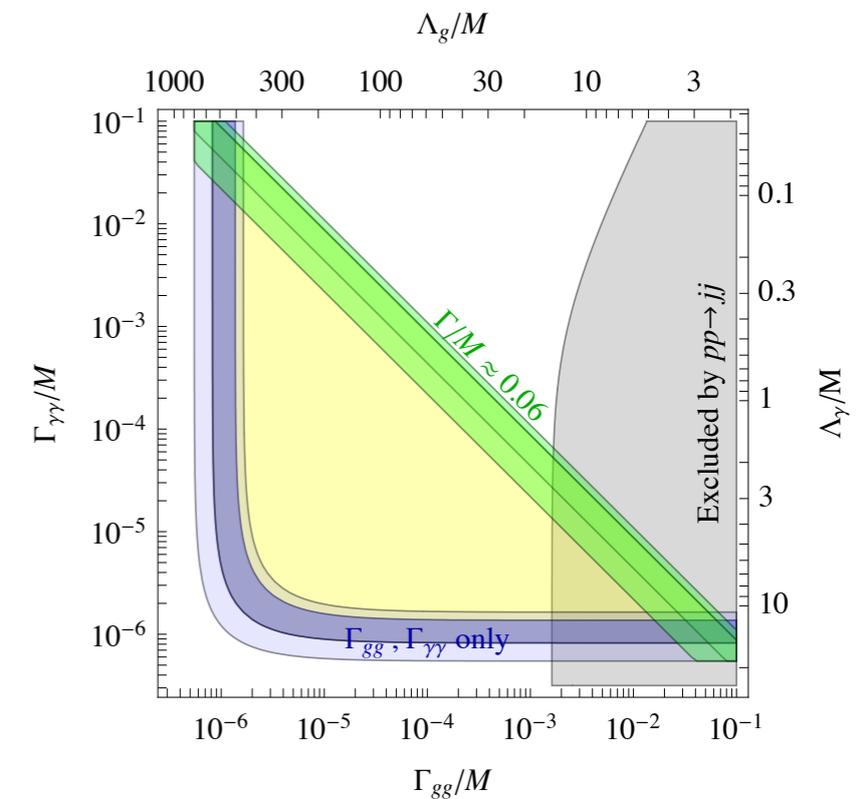


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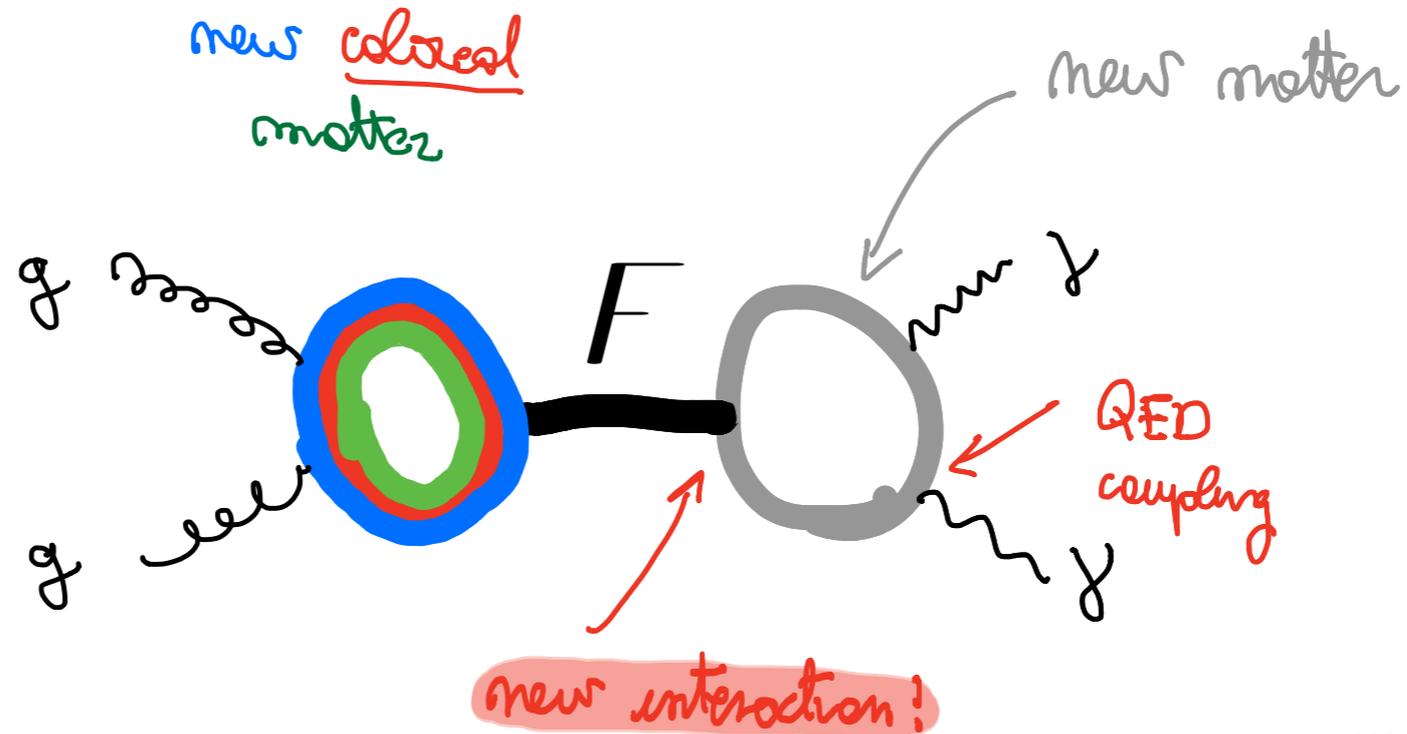


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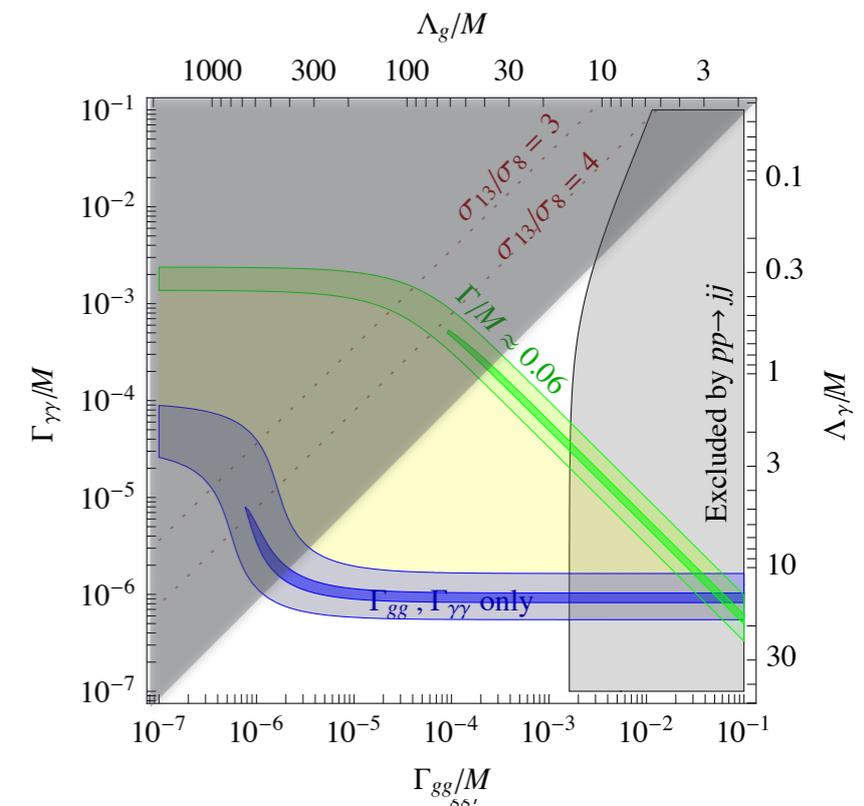


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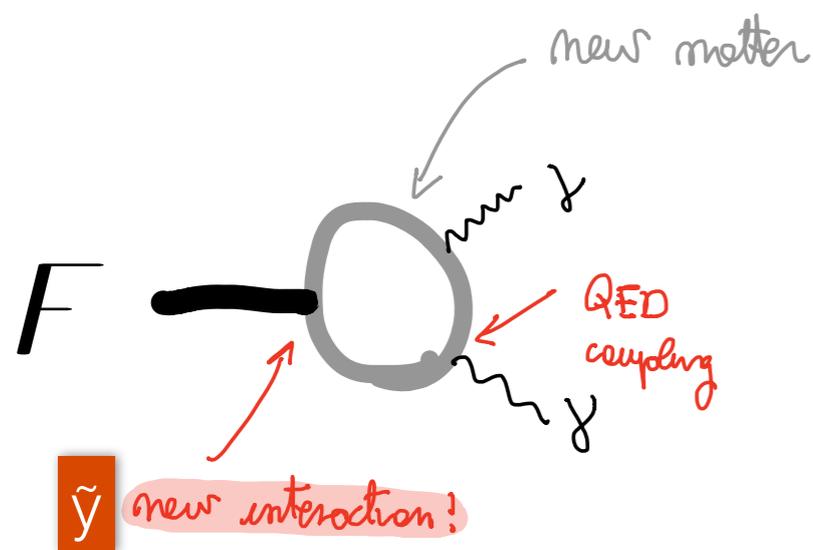


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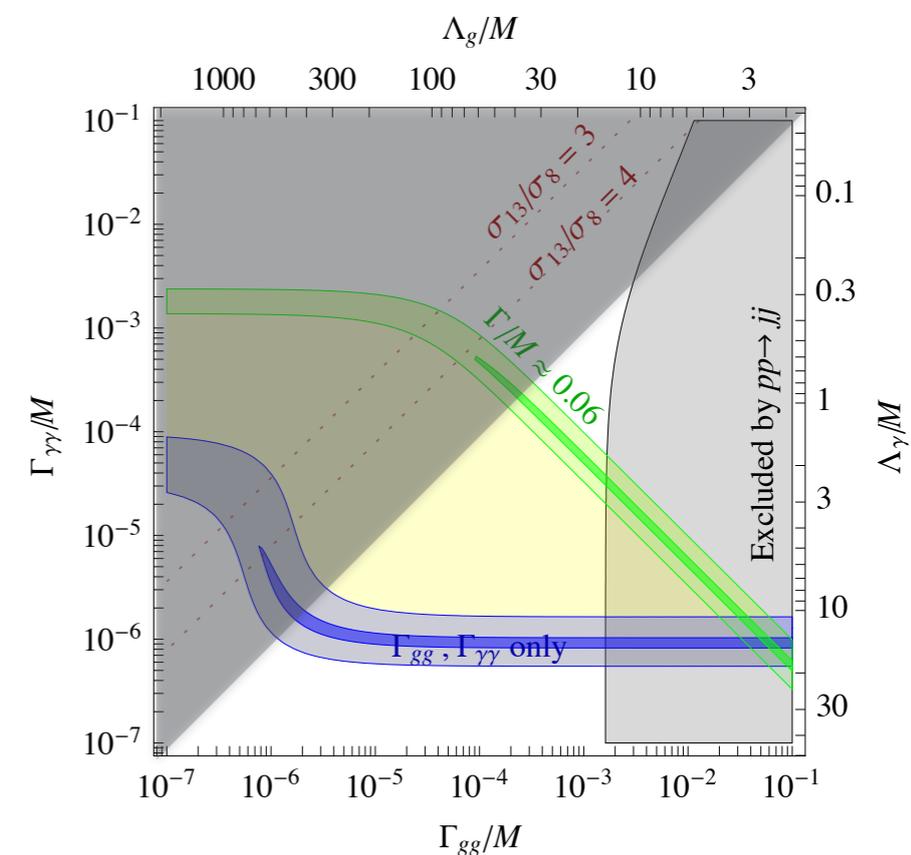
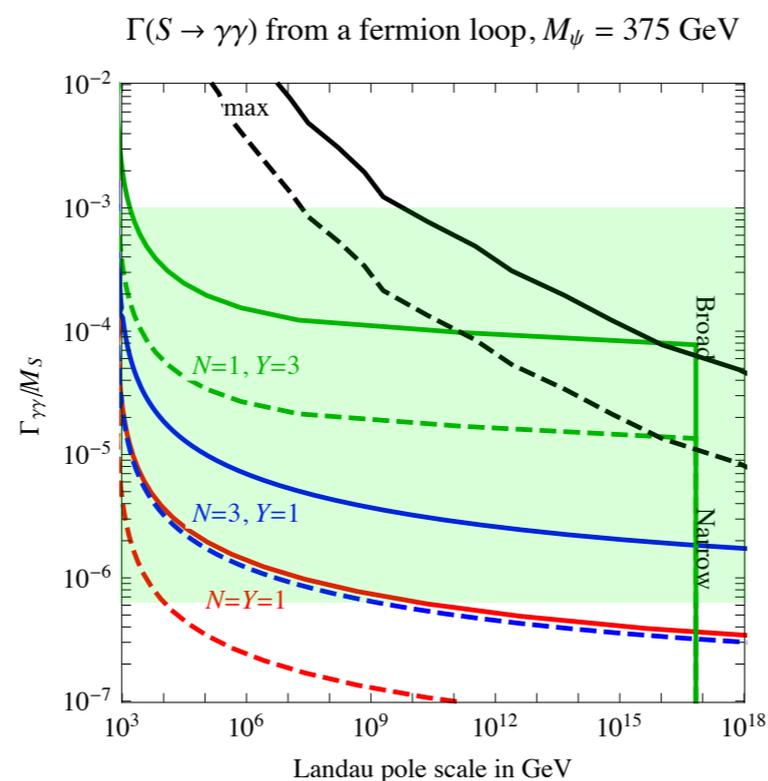


Spin-0

focus on a weak singlet for simplicity



$$\frac{\Gamma_{\gamma\gamma}}{M} \approx 0.6 \cdot 10^{-6} N^2 \tilde{y}^2 Y^4$$

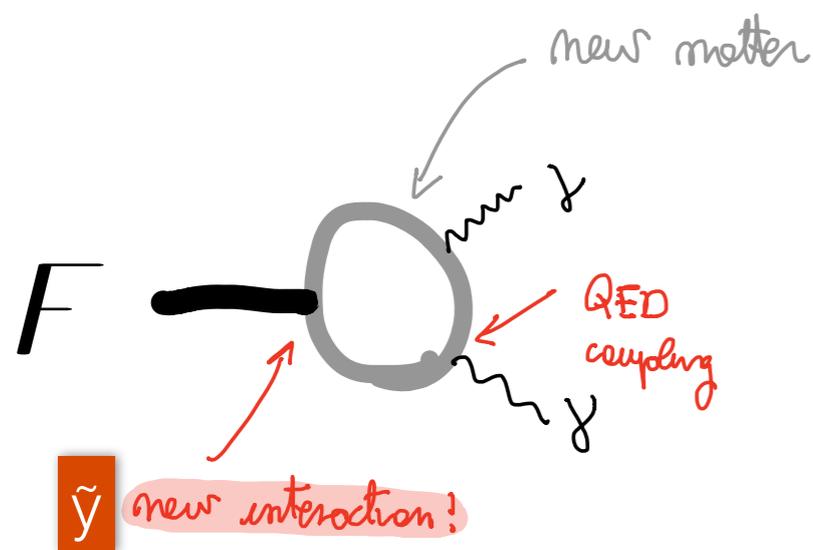


$$\frac{d\tilde{y}}{d \ln \mu} \propto \tilde{y}^3 (2Nd) - \text{gauge terms}$$

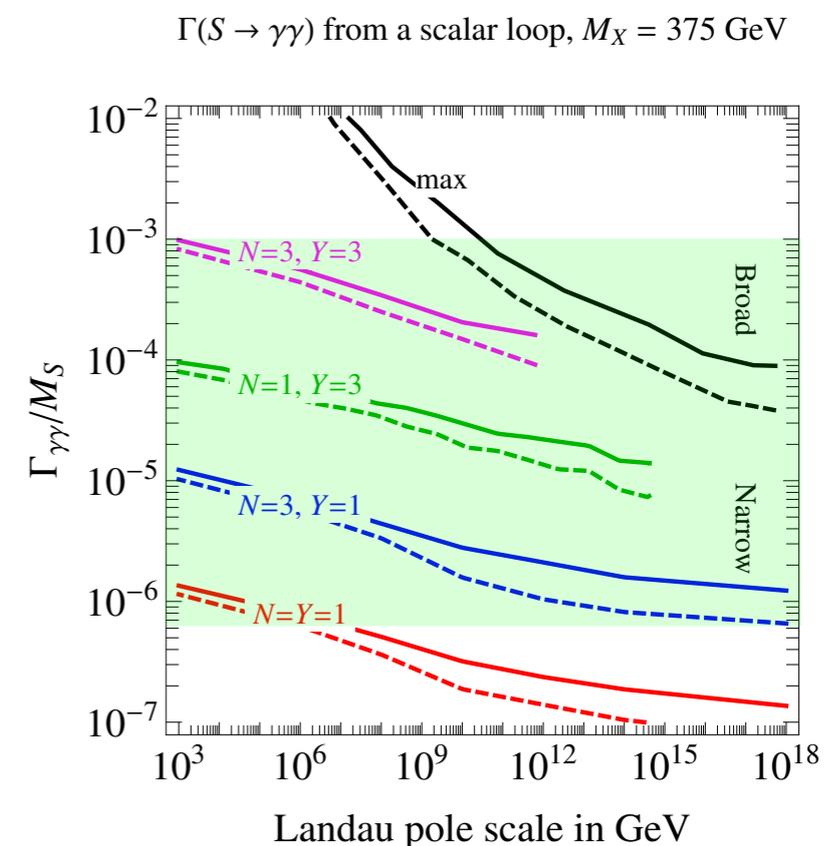
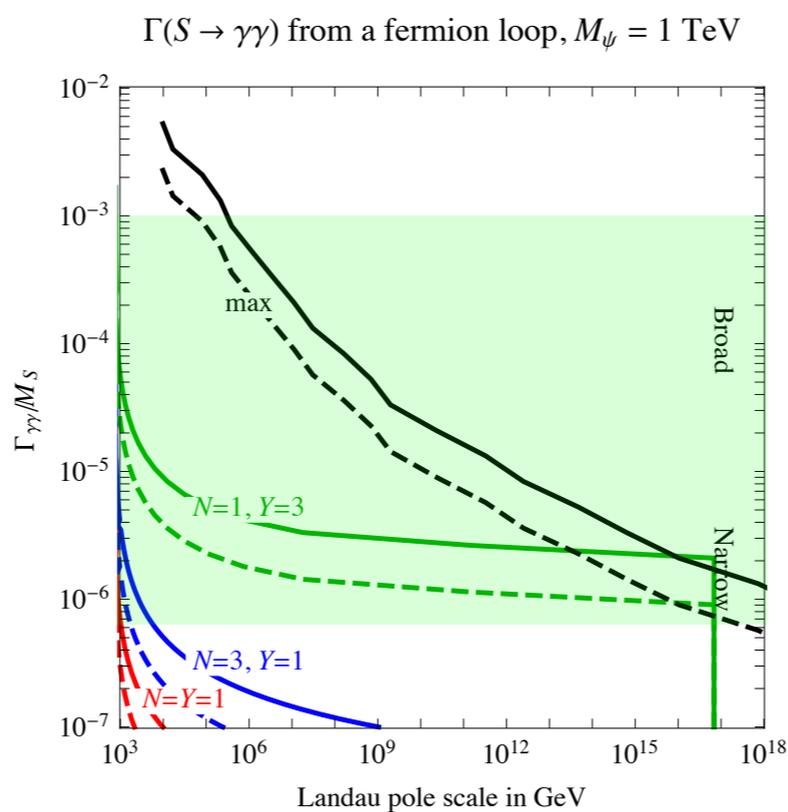
$$\frac{dg_i}{d \ln \mu} \propto g_i^3 \left(\frac{41}{6} + \underbrace{\frac{4}{3} Y^2 Nd}_{\text{new matter (fermions)}} \right)$$

Spin-0

focus on a weak singlet for simplicity



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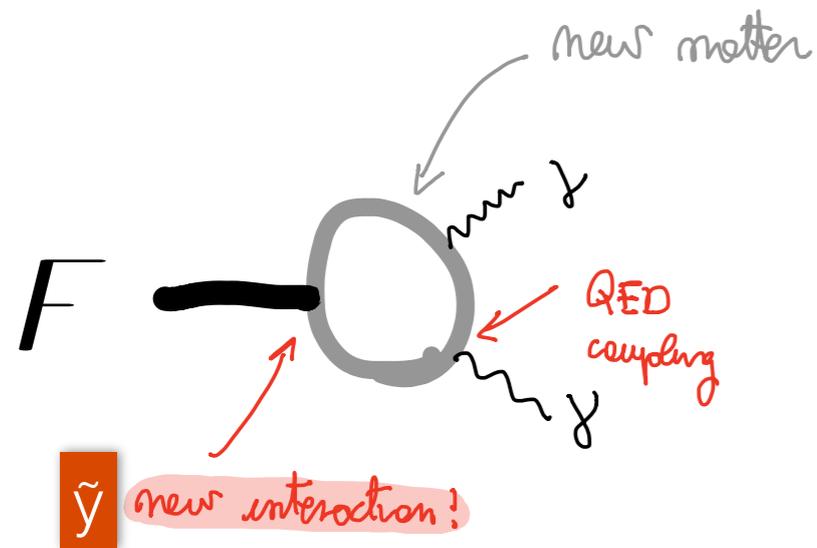


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$$\frac{dg_1}{d \ln \mu} \propto g_1^3 \left(\frac{41}{6} + \underbrace{\frac{4}{3} Y^2 N d}_{\text{new matter (fermions)}} \right)$$

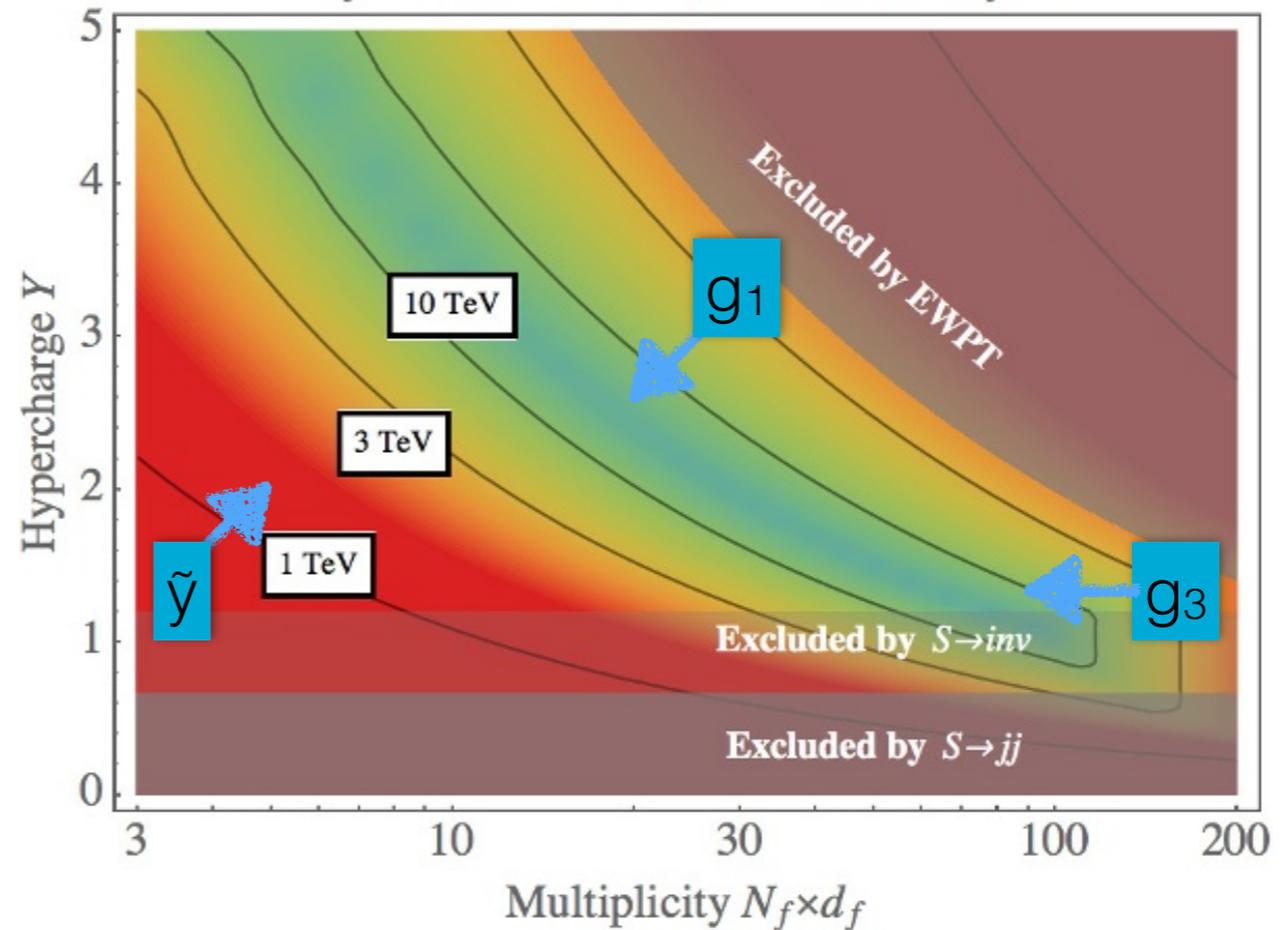
Spin-0

focus on a weak singlet for simplicity



$$\frac{\Gamma_{\gamma\gamma}}{M} \approx 0.6 \cdot 10^{-6} N^2 \tilde{y}^2 Y^4$$

$$M_f = 750 \text{ GeV}, \Gamma/M = 0.01, d_f = 3$$

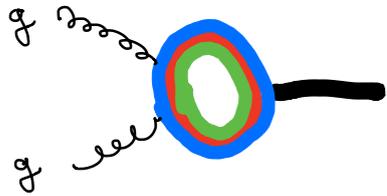


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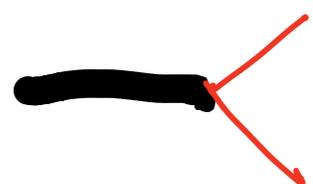
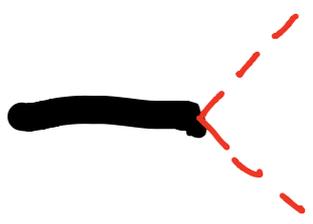
new strong
interactions

Other resonance channels



$\sigma(pp \rightarrow \gamma\gamma)$	$\sqrt{s} = 8 \text{ TeV}$		$\sqrt{s} = 13 \text{ TeV}$	
	narrow	broad	narrow	broad
CMS	$0.63 \pm 0.31 \text{ fb}$	$0.99 \pm 1.05 \text{ fb}$	$4.8 \pm 2.1 \text{ fb}$	$7.7 \pm 4.8 \text{ fb}$
ATLAS	$0.21 \pm 0.22 \text{ fb}$	$0.88 \pm 0.46 \text{ fb}$	$5.5 \pm 1.5 \text{ fb}$	$7.6 \pm 1.9 \text{ fb}$

final state f	σ at $\sqrt{s} = 8 \text{ TeV}$			σ at $\sqrt{s} = 13 \text{ TeV}$		
	observed	expected	ref.	observed	expected	ref.
$e^+e^-, \mu^+\mu^-$	$< 1.2 \text{ fb}$	$< 1.2 \text{ fb}$	[3]	$< 5 \text{ fb}$	$< 5 \text{ fb}$	[78]
$\tau^+\tau^-$	$< 12 \text{ fb}$	$< 15 \text{ fb}$	[3]	$< 60 \text{ fb}$	$< 67 \text{ fb}$	[79]
$Z\gamma$	$< 11 \text{ fb}$	$< 11 \text{ fb}$	[3]	$< 28 \text{ fb}$	$< 40 \text{ fb}$	[80]
ZZ	$< 12 \text{ fb}$	$< 20 \text{ fb}$	[3]	$< 200 \text{ fb}$	$< 220 \text{ fb}$	[81]
Zh	$< 19 \text{ fb}$	$< 28 \text{ fb}$	[3]	$< 116 \text{ fb}$	$< 116 \text{ fb}$	[82]
hh	$< 39 \text{ fb}$	$< 42 \text{ fb}$	[3]	$< 120 \text{ fb}$	$< 110 \text{ fb}$	[83]
W^+W^-	$< 40 \text{ fb}$	$< 70 \text{ fb}$	[3]	$< 300 \text{ fb}$	$< 300 \text{ fb}$	[84]
$t\bar{t}$	$< 450 \text{ fb}$	$< 600 \text{ fb}$	[3]			
invisible	$< 0.8 \text{ pb}$	-	[3]			
$b\bar{b}$	$\lesssim 1 \text{ pb}$	$\lesssim 1 \text{ pb}$	[3]			
jj	$\lesssim 2.5 \text{ pb}$	-	[3]			



	$\frac{\sigma_{13\text{TeV}}}{\sigma_{8\text{TeV}}}$	$10^2 \times$	r_{WW}^γ	r_{ZZ}^γ	$r_{Z\gamma}^\gamma$	r_{hh}^γ	$r_{t\bar{t}}^\gamma$	$r_{\tau\bar{\tau}}^\gamma$	$r_{\ell\bar{\ell}}^\gamma$	r_{gg}^γ
ATLAS	2.9 [5]*		3.0 [7]	13 [9]*	19 [10]*	4.1 [11]*	0.22 [13]	15 [15]	124 [17]*	0.14 [19]
CMS	4.0 [6]		0.5 [8]	4.6 [8]	—	2.8 [12]*	0.33 [14]	7.4 [16]	114 [18]*	0.083 [20]*

Strong interactions

- ✦ pseudo-Nambu-Goldstone bosons of:
 - A. internal symmetries
 - B. “susy”
 - C. conformal
- Quarkonium (and quirks)
- ...

Goldstone bosons

1512.05330

π^0 in QCD is an inspiring template

$$\pi^0 \rightarrow \gamma\gamma$$

unlike the QCD pion

1. the large absolute width $\Gamma(F \rightarrow \gamma\gamma)$ suggest lots of states ($\Gamma(F \rightarrow \gamma\gamma)/M \gtrsim 10^{-6}, 10^{-4}$ in most scenarios, vs $\Gamma(\pi^0 \rightarrow \gamma\gamma)/m_\pi \sim 10^{-7}$)
2. for a generic scalar expect a proportionally large contribution to the mass of F

shift symmetry protection to the mass of the GB!

Goldstone bosons

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2. for a generic scalar expect a proportionally large contribution to the mass of F

shift symmetry protection to the mass of the GB!

coupling to all gauge bosons (can) naturally arise at the same order

couplings to fermions protected by chiral symmetry

$$\begin{aligned} \mathcal{L}_I^{\Phi=\eta} = & -i\frac{\eta}{f} (C_t m_t \bar{t} \gamma^5 t + C_b m_b \bar{b} \gamma^5 b + C_\tau m_\tau \bar{\tau} \gamma^5 \tau) \\ & -\frac{\eta}{f} \left(C_{gg} \frac{\alpha_3}{8\pi} G_{\mu\nu}^a \tilde{G}^{a\mu\nu} + C_{\gamma\gamma} \frac{\alpha_e}{8\pi} F_{\mu\nu} \tilde{F}_{\mu\nu} \right) \\ & -\frac{\eta}{f} \left(C_{WW} \frac{\alpha_2}{4\pi} W_{\mu\nu}^+ \tilde{W}^{-\mu\nu} + C_{ZZ} \frac{\alpha_2 \cos^2 \theta_W}{8\pi} Z_{\mu\nu} \tilde{Z}^{\mu\nu} + C_{Z\gamma} \frac{\alpha_e}{4\pi \tan \theta_W} Z_{\mu\nu} \tilde{F}^{\mu\nu} \right) \end{aligned}$$

CP odd

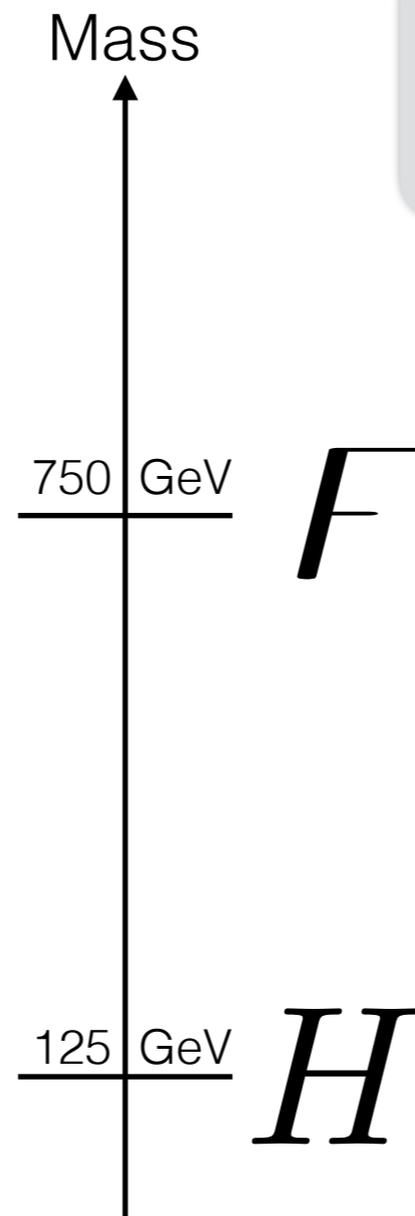
Higgs + F as Goldstone (η)

global symmetry breaking $G \rightarrow H$ (e.g. $SU(3)^2 \rightarrow SU(2)$)

1512.05330

F is (less) light compared to the scale of symmetry breaking

Higgs is light compared to the scale of symmetry breaking



$$\mathcal{L}_I^\eta = -\frac{\eta}{f} \left(iC_t \frac{\sqrt{2}m_t}{v} \bar{q}_L \tilde{H} t_R + h.c. + \dots \right) - \frac{\eta}{f} \left(C_G \frac{\alpha_3}{8\pi} G_{\mu\nu}^a \tilde{G}^{a\mu\nu} + C_W \frac{\alpha_2}{8\pi} W_{\mu\nu}^i \tilde{W}^{i\mu\nu} + C_B \frac{\alpha_1}{8\pi} B_{\mu\nu} \tilde{B}^{\mu\nu} \right)$$

$$\Gamma_{\gamma\gamma} = \left(\frac{\alpha_e}{8\pi f} \right)^2 \frac{m_\eta^3}{\pi} |C_{\gamma\gamma}^{\text{eff}}|^2, \\ C_{\gamma\gamma}^{\text{eff}} = C_{\gamma\gamma} + \frac{1}{2} C_t N_c^{(t)} Q_t^2 A_{1/2}(x_t),$$

$$\Gamma_{gg} = 8 \left(\frac{\alpha_3}{8\pi f} \right)^2 \frac{m_\eta^3}{\pi} |C_{gg}^{\text{eff}}|^2, \\ C_{gg}^{\text{eff}} = C_{gg} + \frac{1}{4} C_t A_{1/2}(x_t),$$

$$\Gamma_{t\bar{t}} \simeq 3 \frac{C_t^2}{8\pi} \frac{m_t^2 m_\eta}{f^2}$$

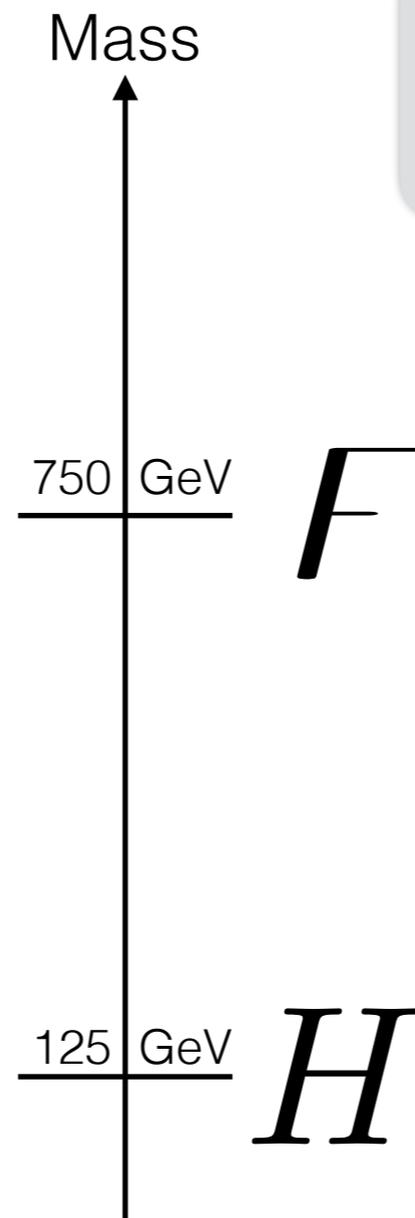
Higgs + F as Goldstone (η)

global symmetry breaking $G \rightarrow H$ (e.g. $SU(3)^2 \rightarrow SU(2)$)

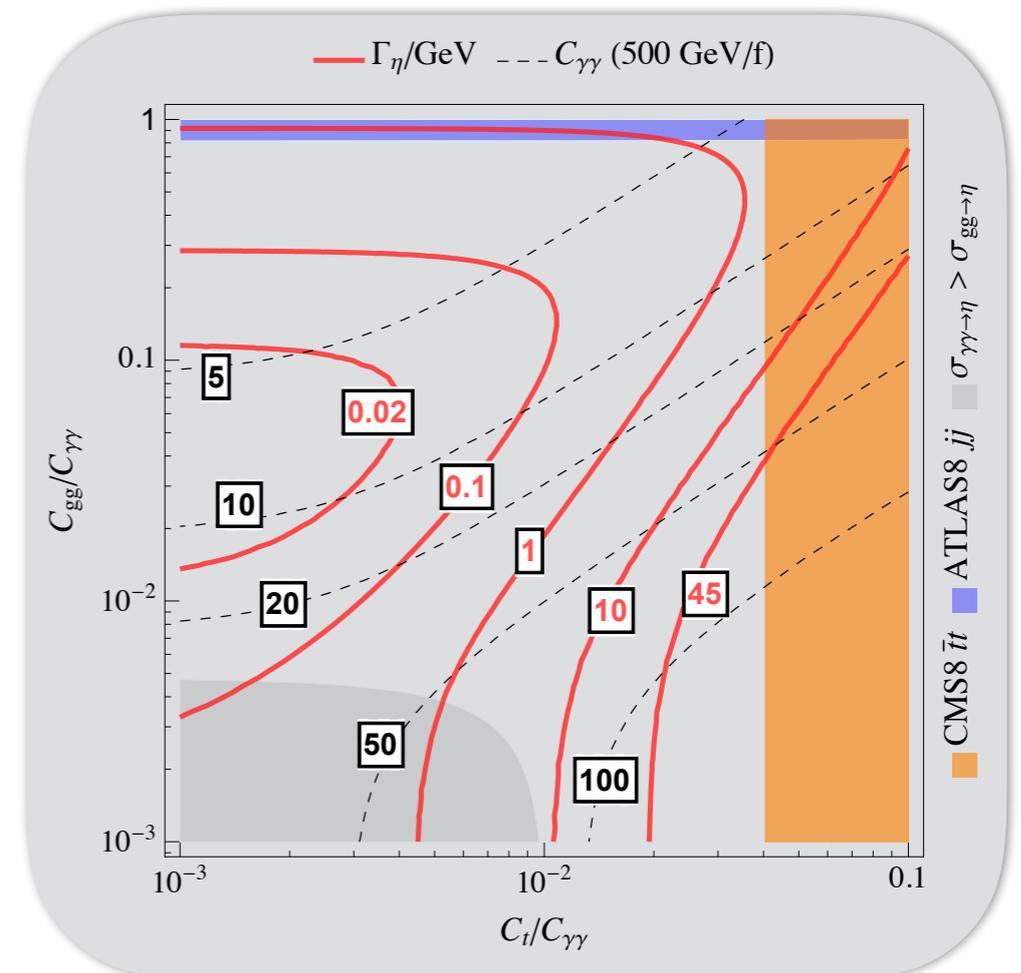
1512.05330

F is (less) light compared to the scale of symmetry breaking

Higgs is light compared to the scale of symmetry breaking



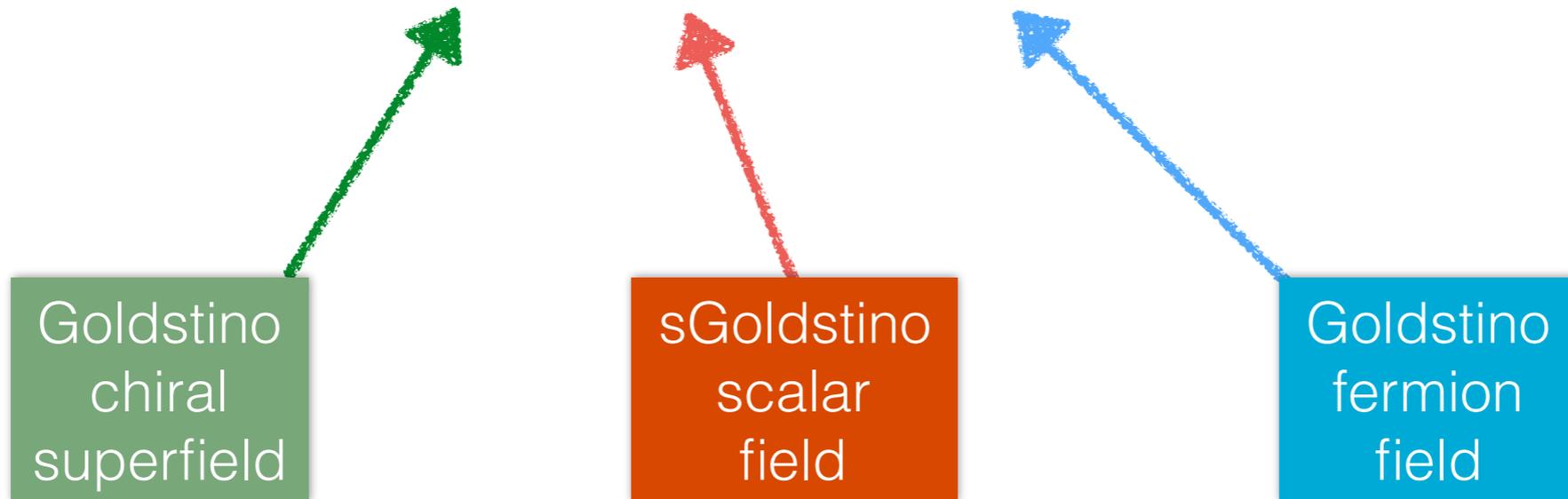
$$\mathcal{L}_I^\eta = -\frac{\eta}{f} \left(iC_t \frac{\sqrt{2}m_t}{v} \bar{q}_L \tilde{H} t_R + h.c. + \dots \right) - \frac{\eta}{f} \left(C_G \frac{\alpha_3}{8\pi} G_{\mu\nu}^a \tilde{G}^{a\mu\nu} + C_W \frac{\alpha_2}{8\pi} W_{\mu\nu}^i \tilde{W}^{i\mu\nu} + C_B \frac{\alpha_1}{8\pi} B_{\mu\nu} \tilde{B}^{\mu\nu} \right)$$



F as a sGoldstino (Φ)

1512.05330

$$X = \Phi + \sqrt{2}\theta G + \theta^2 F_X$$



$$-\frac{1}{2\mathcal{F}} \int d^2\theta X (m_1 W^\alpha W_\alpha + m_2 W^{\alpha a_2} W_{\alpha}^{a_2} + m_3 W^{\alpha a_3} W_{\alpha}^{a_3}) + h.c.$$

coupling of the sGoldstino are proportional to the SUSY breaking masses

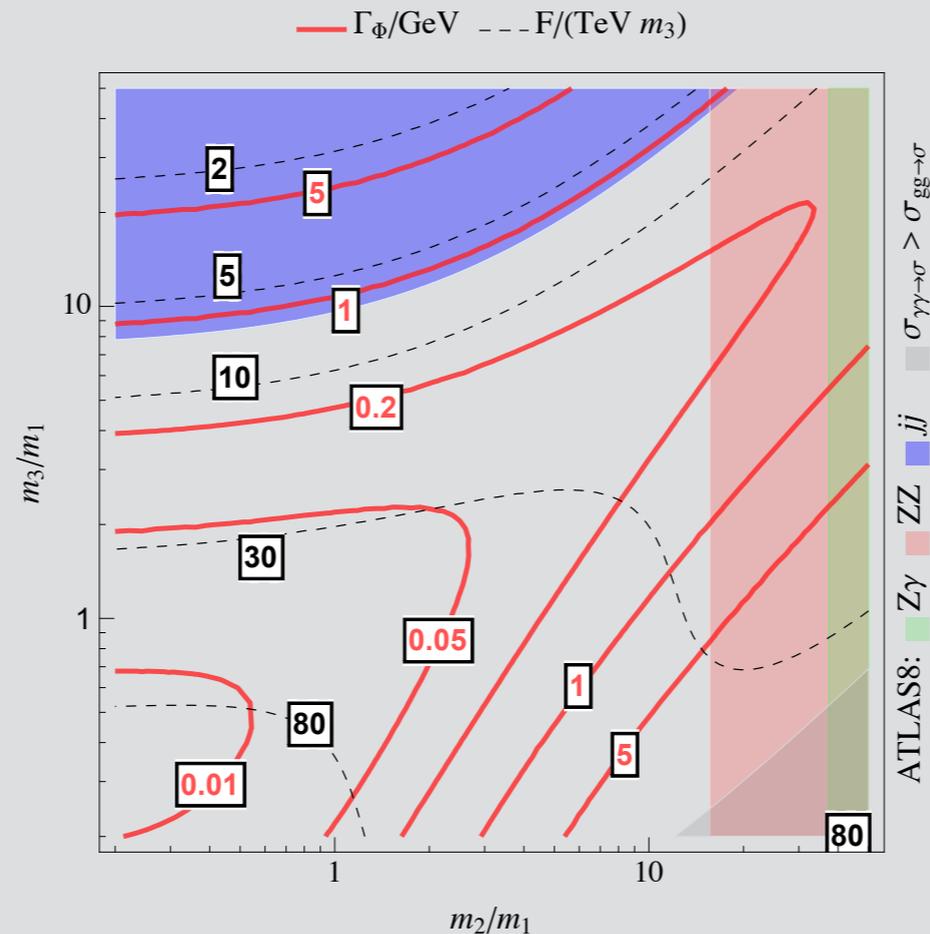
$$\Gamma_{gg} = \left(\frac{m_3}{2\mathcal{F}}\right)^2 \frac{m_\sigma^3}{\pi}, \quad \Gamma_{\gamma\gamma} = \frac{1}{2} \left(\frac{m_{\sigma\gamma\gamma}}{4\mathcal{F}}\right)^2 \frac{m_\sigma^3}{\pi},$$

negligible decay to 2 invisible Goldstino

$$\Gamma_{ZZ} \simeq \frac{1}{2} \left(\frac{m_{\sigma ZZ}}{4\mathcal{F}}\right)^2 \frac{m_\sigma^3}{\pi}, \quad \Gamma_{WW} \simeq \left(\frac{m_2}{4\mathcal{F}}\right)^2 \frac{m_\sigma^3}{\pi}, \quad \Gamma_{Z\gamma} \simeq \left(\frac{m_{\sigma Z\gamma}}{4\mathcal{F}}\right)^2 \frac{m_\sigma^3}{\pi}.$$

F as a sGoldstino (Φ)

1512.05330



coupling of the sGoldstino are proportional to the SUSY breaking masses

$$\Gamma_{gg} = \left(\frac{m_3}{2\mathcal{F}}\right)^2 \frac{m_\sigma^3}{\pi}, \quad \Gamma_{\gamma\gamma} = \frac{1}{2} \left(\frac{m_{\sigma\gamma\gamma}}{4\mathcal{F}}\right)^2 \frac{m_\sigma^3}{\pi},$$

$$\Gamma_{ZZ} \simeq \frac{1}{2} \left(\frac{m_{\sigma ZZ}}{4\mathcal{F}}\right)^2 \frac{m_\sigma^3}{\pi}, \quad \Gamma_{WW} \simeq \left(\frac{m_2}{4\mathcal{F}}\right)^2 \frac{m_\sigma^3}{\pi}, \quad \Gamma_{Z\gamma} \simeq \left(\frac{m_{\sigma Z\gamma}}{4\mathcal{F}}\right)^2 \frac{m_\sigma^3}{\pi}.$$

negligible decay to 2 invisible Goldstino

F as a Dilaton (σ)

spontaneous symmetry breaking of conformal symmetry

1512.05330

$\sigma T^{\mu}_{\mu} \sim$ masses + β -functions

$$\frac{\sigma}{f} T^{\mu}_{\mu}^{\text{CFT}} = \frac{\sigma}{f} \left(\frac{\alpha_3}{8\pi} \kappa_3 G_{\mu\nu}^2 + \kappa_e \frac{\alpha_e}{8\pi} F_{\mu\nu}^2 - y_t (1 + \epsilon_t) \bar{q}_L \tilde{H} t_R + h.c. + 2(1 + \epsilon_H) |D_{\mu} H|^2 + \dots \right)$$

tree-level coupling to Z and W (massive vectors)

loop-level coupling to gluon and photon (beta functions)

$$\Gamma_{ZZ} \simeq \Gamma_{WW}/2 \simeq \Gamma_{hh} \simeq \frac{m_{\sigma}^3}{32\pi f^2} \quad r_{ZZ}^{\gamma} \approx 0.05 \left(\frac{\kappa_e}{240} \right)^2$$

β functions of $O(16\pi^2) \rightarrow$ large dilaton mass correction

CP of F

1604.06446

critical importance

- greatest discriminator of many scenarios
- truly “fitfy-fitfy” (unlike for the Higgs boson CP)

shortcuts: $F \rightarrow HH$ or ZH

several challenges

important differences w.r.t the Higgs CP measurement

1. no $F \rightarrow 4$ fermion final state
2. $M_F \gg m_h \Rightarrow$ small boost

only $g^+ g^+ \rightarrow g^+ g^+ F$ scattering is sensitive to CP for other helicities $A \sim M_F$

Thrust and $\Delta\phi(jj)$

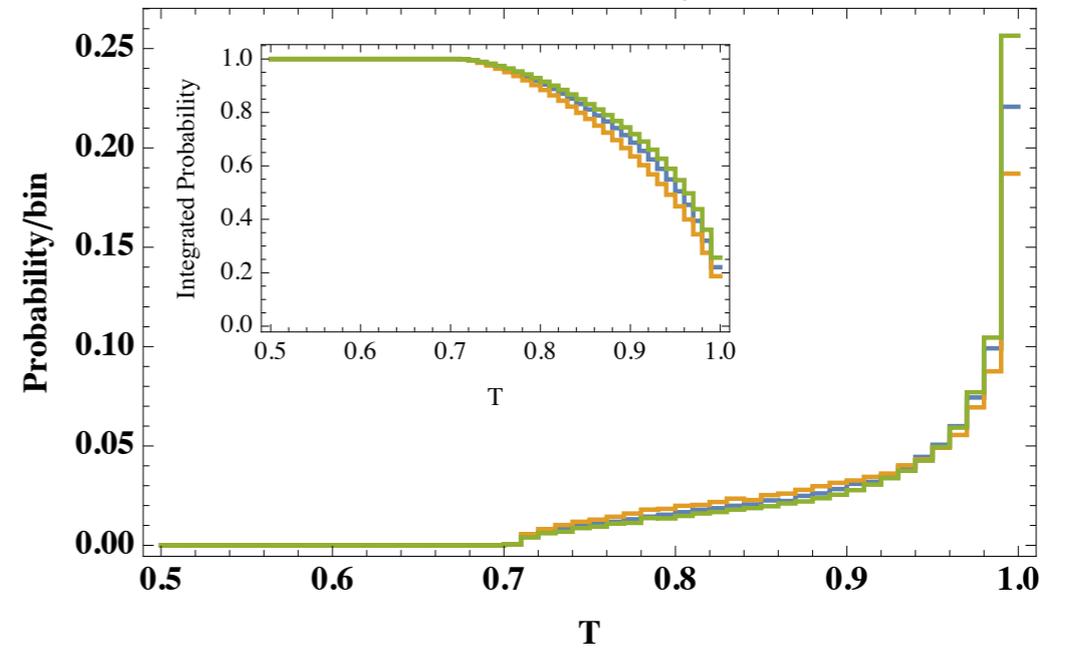
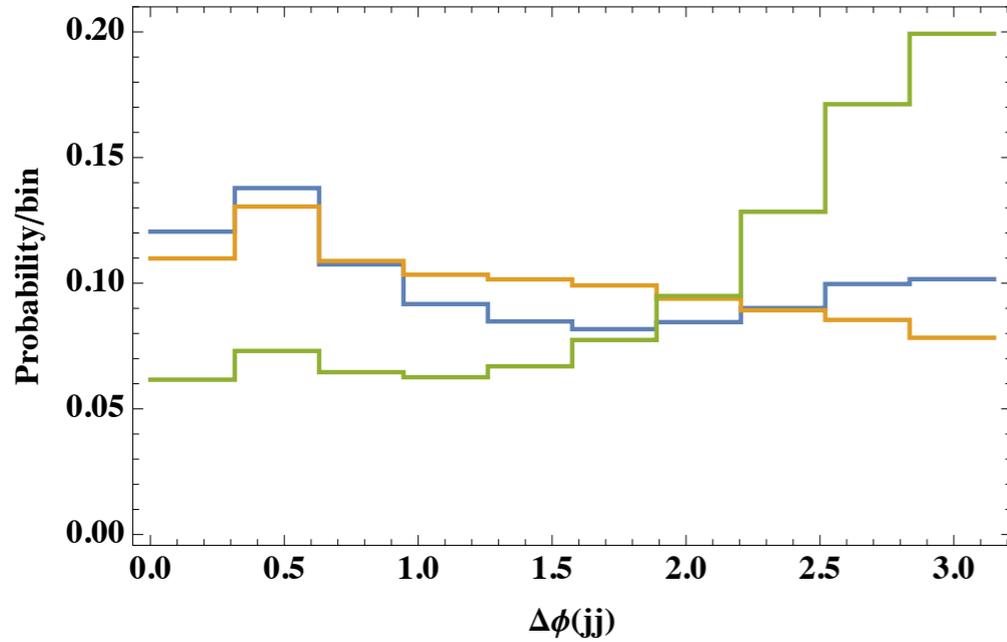
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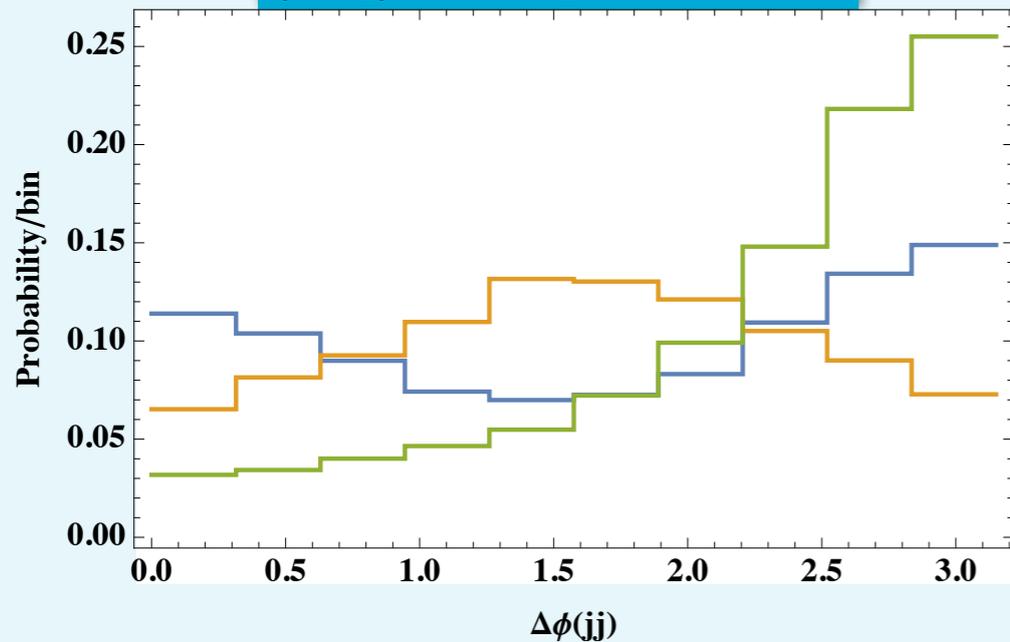
$pp \rightarrow F jj$

$\Delta\phi_{jj}$

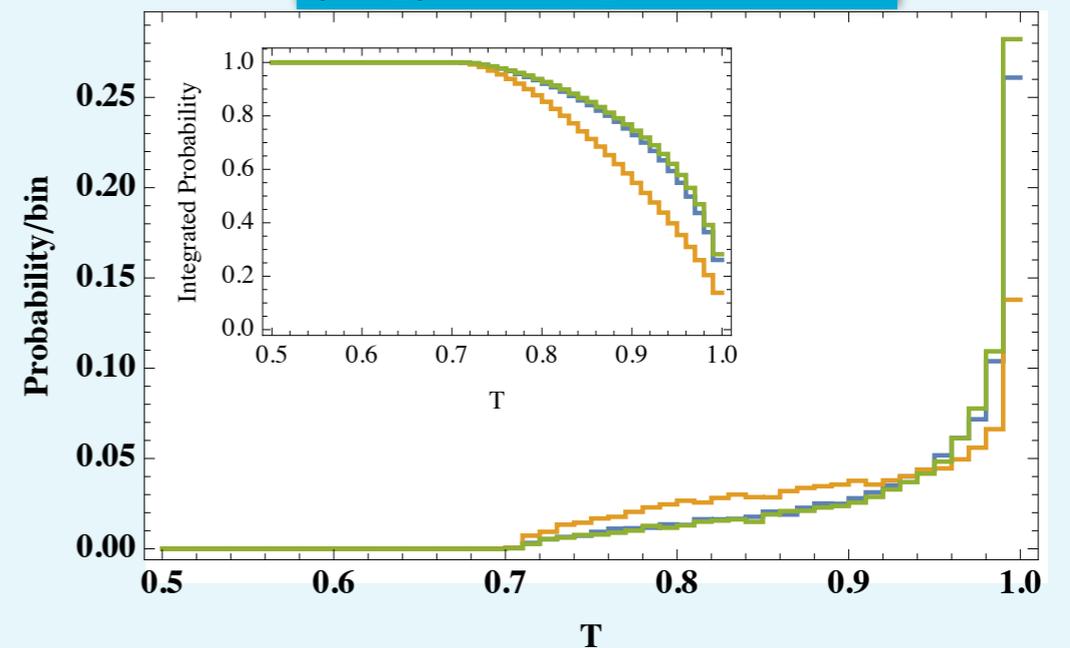
$$T = \max_n \frac{\sum_{i \in \text{jets}} |n \cdot p_i|}{\sum_{i \in \text{jets}} |p_i|}$$



$|\Delta\eta_{jj}| > 2.5 \quad m_{jj} > 500 \text{ GeV}$



$|\Delta\eta_{jj}| > 2.5 \quad m_{jj} > 500 \text{ GeV}$



Pair production of F

if strong interactions are behind F, multiple production is expected with no big suppression

$$U(\pi) = e^{i\pi/f} = 1 + \frac{\pi}{f} + \left(\frac{\pi}{f}\right)^2 + \dots$$

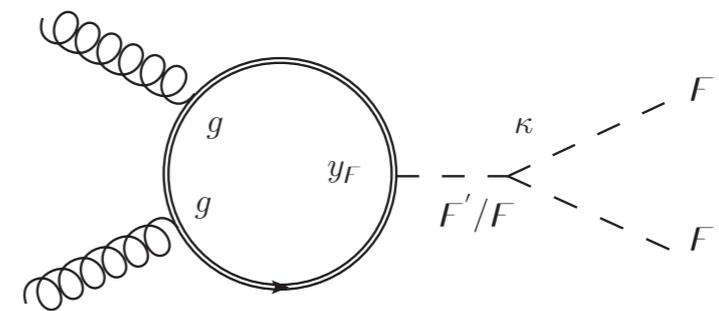
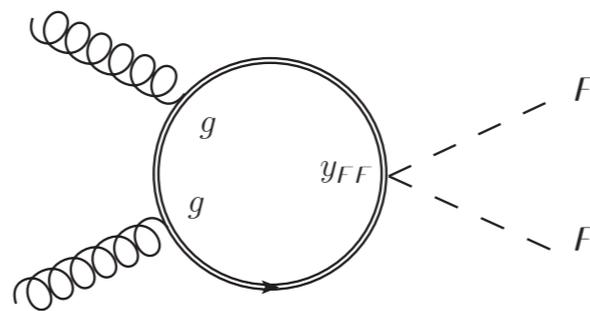
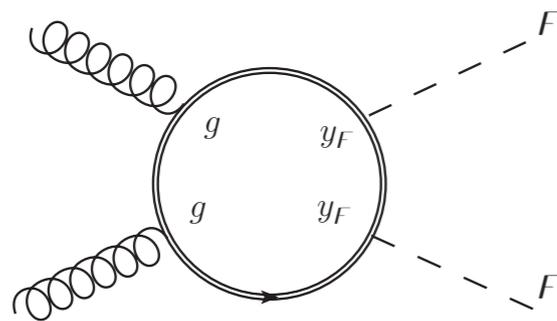
$$\mathcal{L} \sim \frac{c_5}{\Lambda} \cdot F \mathcal{L}_5 + \frac{c_6}{\Lambda^2} \cdot F^2 \mathcal{L}_6$$

$$\sigma_{FF}/\sigma_F \sim (10^{-2} - 10^{-4}) \cdot (c_6/c_5)^2$$

internal symmetries more easily suppress single production than double production

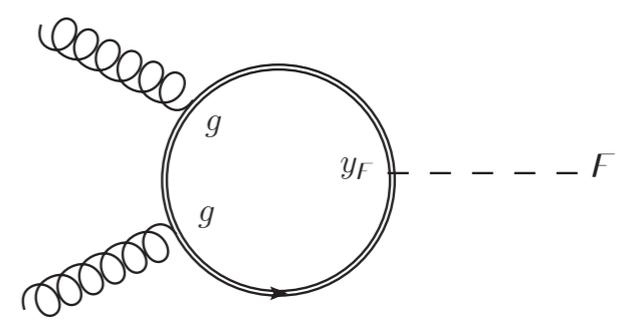
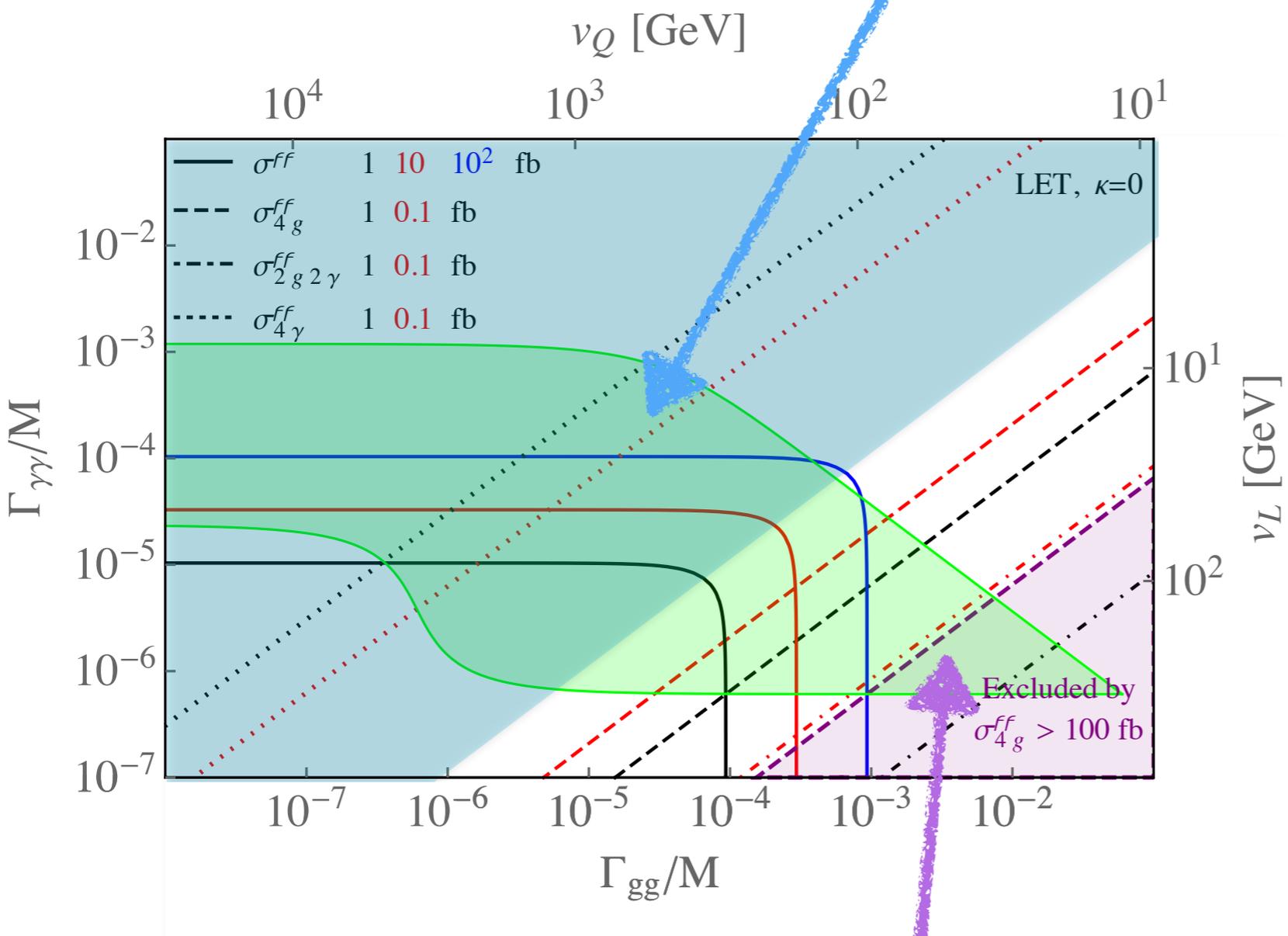
$$|F|^2 G_{\mu\nu} G^{\mu\nu}$$

F couples to	$\sigma_{FF}/\sigma_F = \sigma_{\gamma\gamma F}/2\sigma_{\gamma\gamma}$
$b\bar{b}$	$0.015\% (\text{TeV}/\Lambda)^2 (c_b^{(6)}/c_b)^2$
$d\bar{d}$	$0.050\% (\text{TeV}/\Lambda)^2 (c_d^{(6)}/c_d)^2$
GG	$0.13\% (\text{TeV}/\Lambda)^2 (c_{gg}^{(6)}/c_{gg})^2$
$\gamma\gamma$	$1.9\% (\text{TeV}/\Lambda)^2 (c_{\gamma\gamma}^{(6)}/c_{\gamma\gamma})^2$



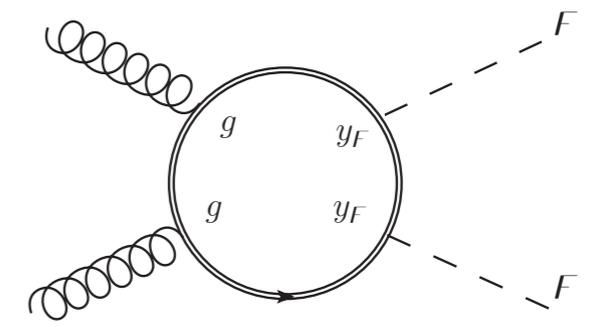
$$\sigma_{FF}/\sigma_F \sim y_F/M \sim 1/f$$

Large 4γ , probes $\gamma\gamma$ initial state



$$\Gamma_{\gamma\gamma} = \frac{\alpha^2 q_L^4 N_L^2 M_F^3}{144\pi^3 v_L^2}$$

$$\Gamma_{gg} = \frac{\alpha_3^2 N_Q^2 I_Q^2 M_F^3}{18\pi^3 v_Q^2}$$



2γ2j or ttγγ ???

Large 4j and 4t final states

$pp \rightarrow F \rightarrow$	4j	4γ	4t
LHC8	<100 fb	<26 fb	<70 fb

$v_Q \equiv M_Q/y_Q$ and $v_L \equiv M_L/y_L$

Conclusions

- $gg \rightarrow F \rightarrow \gamma \gamma$ is the most compelling
- **heavy quark** initial state is also possible (b, c, s)
- light, coupling to gauge bosons \rightarrow **Goldstone boson**
 - ☑ GB of internal symmetry linked to **strong EWSB**
 - ☑ SUSY partner of GB of **SUSY**
 - ☐ GB of conformal invariance
- **CP of F** is one of the most pressing questions
- **F pair production** is worth searching, great reward
- Outlook:
 - * more decay channels $F \rightarrow jj, tt, Z\gamma, ZZ, WW, HH, \text{invisible}, 3\&4\text{-body}$
 - * direct search of states in the loop
 - * other companions from the new theory of TeV physics

Conclusions

Citation: Particle Data Group, 2016 update

$$F (750_{000})$$

$$I(J^P) = ?(0^?)$$

J needs confirmation

OMITTED FROM SUMMARY TABLE

Needs confirmation.

F MASS

VALUE (GeV)	EVTS	DOCUMENT ID	TECN	COMMENT
750 ± 30	OUR AVERAGE	ATLAS, CMS		$pp \rightarrow F$

• • • We do not use the following data for average, fits, limits, etc. • • •

F WIDTH

VALUE (GeV)	CL%	DOCUMENT ID	TECN	COMMENT
<100	95	ATLAS, CMS		$pp \rightarrow F$

• • • We do not use the following data for average, fits, limits, etc. • • •

F DECAY MODES

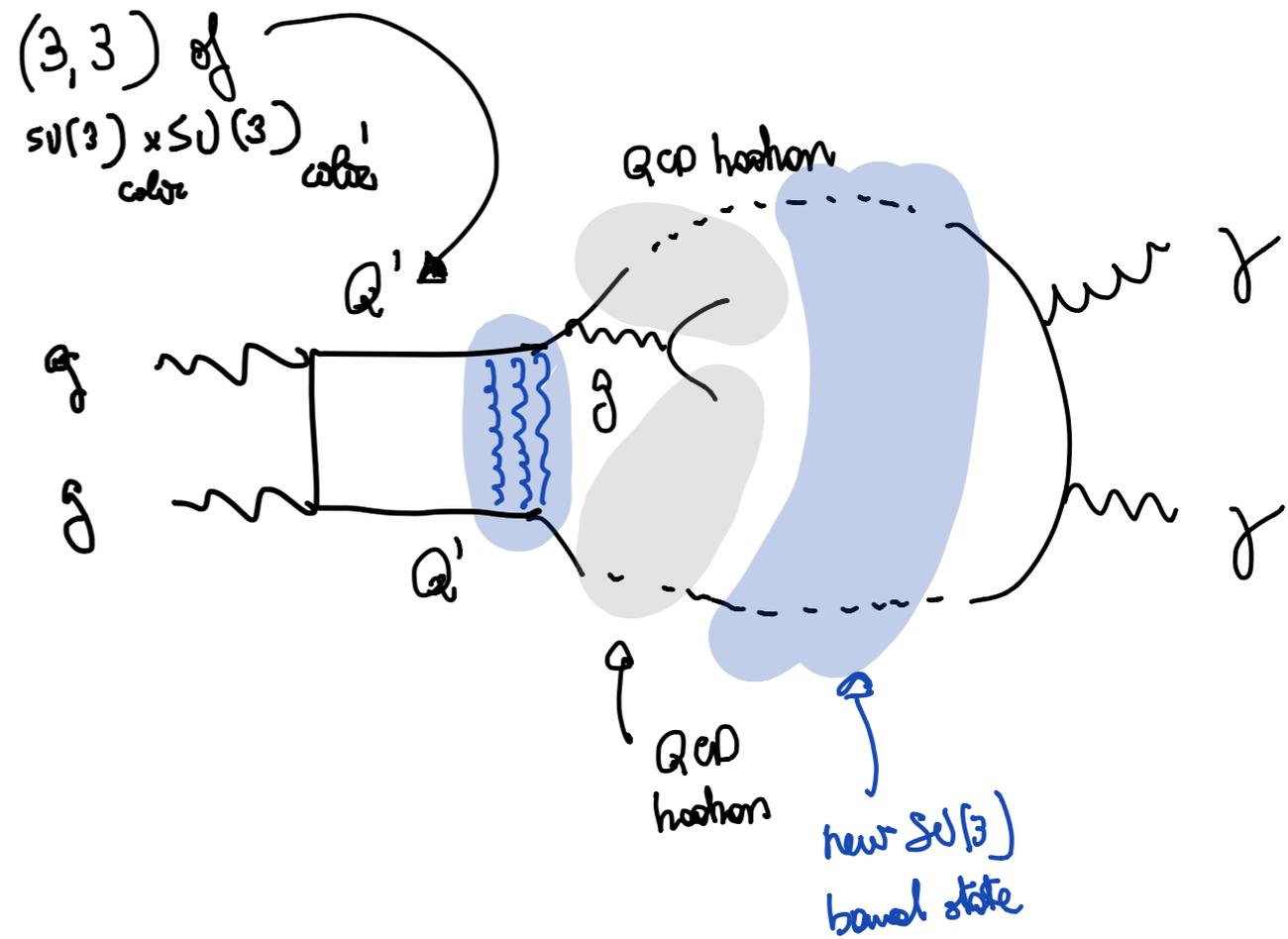
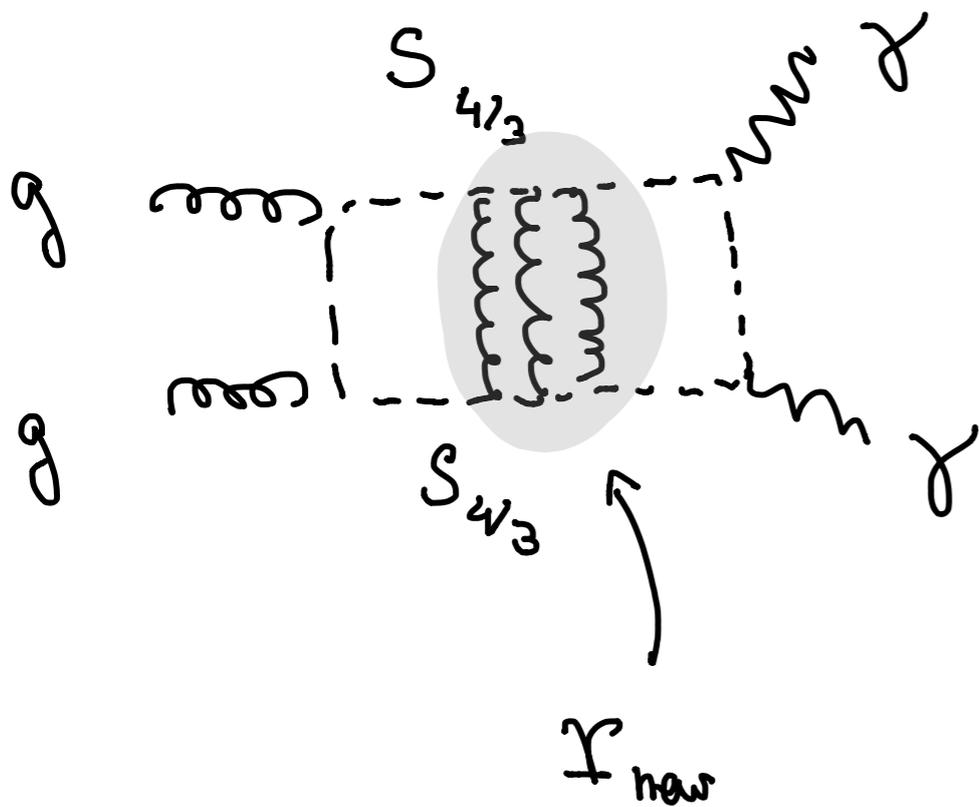
Mode	Fraction (Γ_i/Γ)
Γ_1 $\gamma\gamma$	seen(?)
Γ_2 $\gamma Z, ZZ, WW, jj$	expected
Γ_3 $t\bar{t}, b\bar{b}$, invisible	possible
Γ_4 3-body, 4-body	predicted

Thank you!

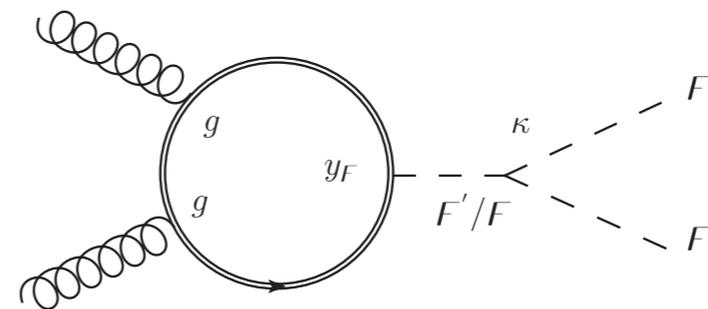
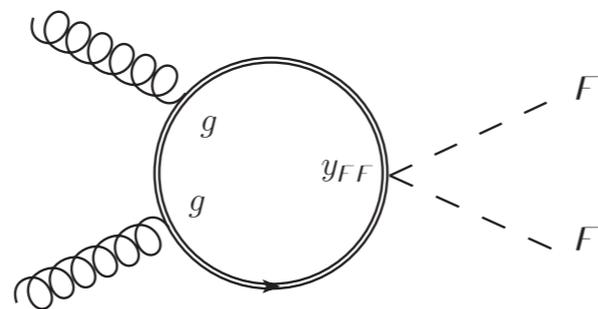
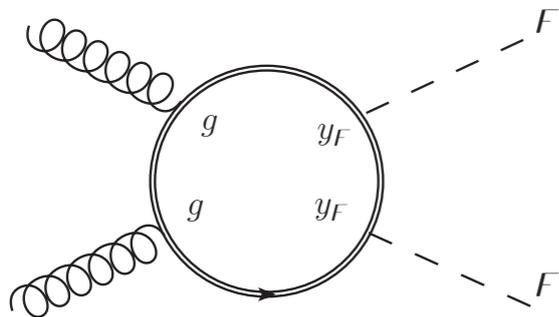
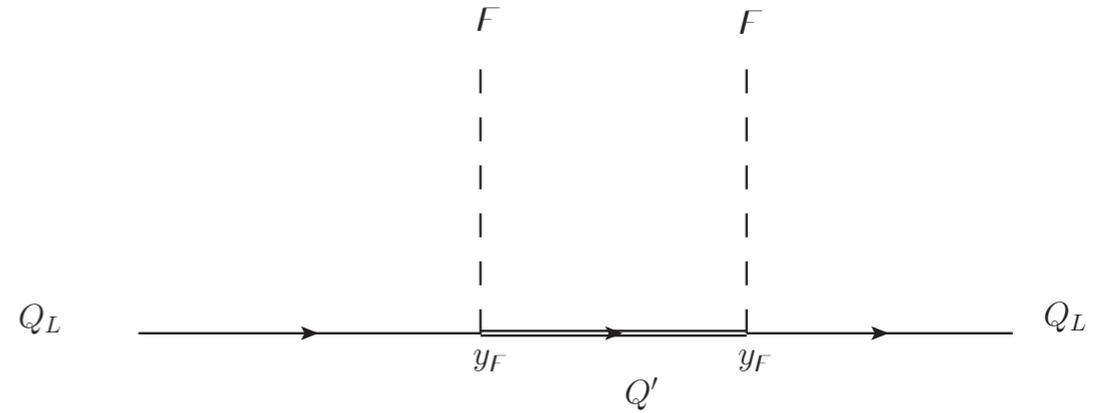
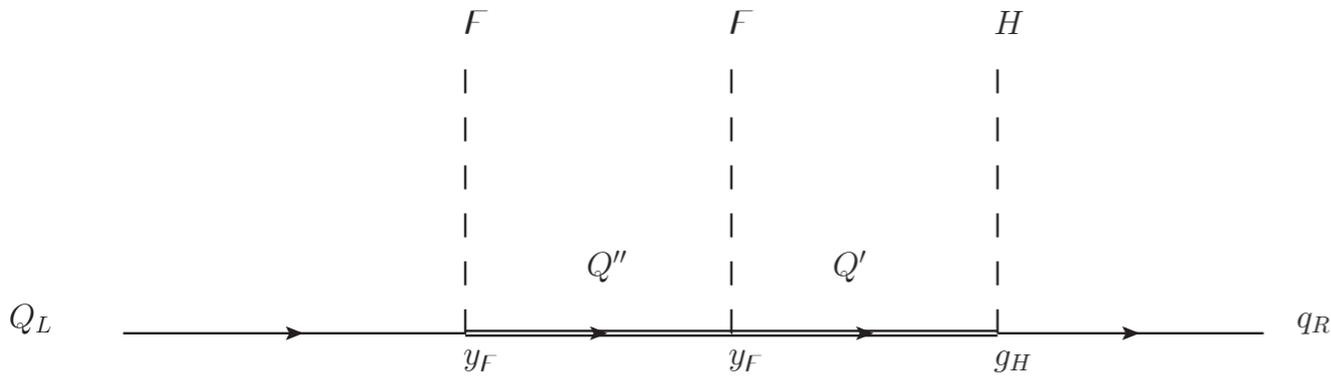
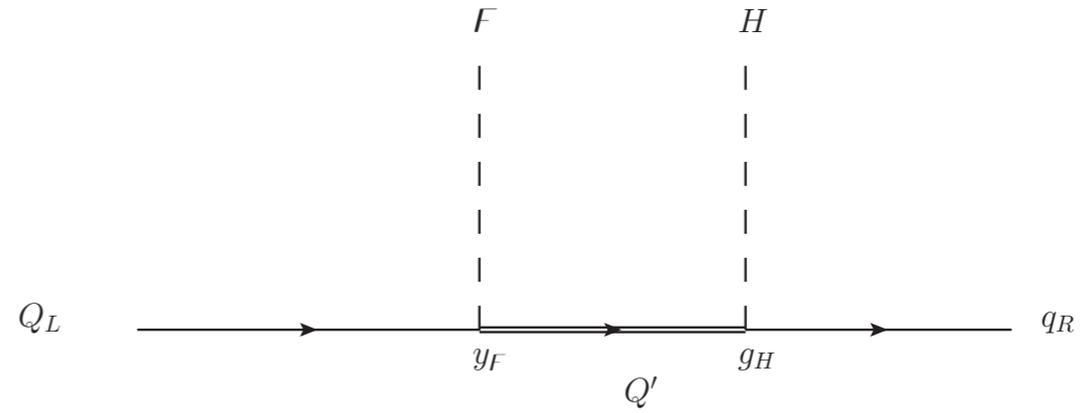
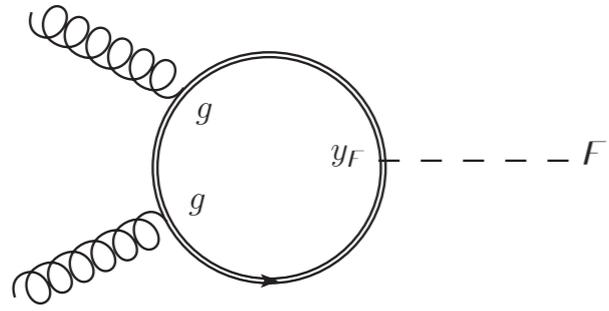
Quarkonium and Quirks

1602.08819

1604.06180



Pair production mechanism



QCD Goldstone bosons and Quarkonium

	BR	Mass
π^0	Photon Photon 0.98798	134.9766 MeV/c ²
K _L	Photon Photon 0.000547	497.648 MeV/c ²
K _S	Photon Photon 2.71×10^{-6}	497.648 MeV/c ²
η	Photon Photon 0.3938	547.51 MeV/c ²
$\eta'(958)$	Photon Photon 0.0212	957.78 MeV/c ²
f ₀ (980)	Photon Photon —	980. MeV/c ²
a ₀ ⁰ (980)	Photon Photon —	984.7 MeV/c ²
f ₀ (1370)	Photon Photon —	1350. MeV/c ²
D ₀	Photon Photon $0. \times 10^{-5}$	1864.5 MeV/c ²
D ₀ -bar	Photon Photon $0. \times 10^{-5}$	1864.5 MeV/c ²
f ₂ (1950)	Photon Photon —	1944. MeV/c ²
f ₂ (2300)	Photon Photon —	2297. MeV/c ²
$\eta_c(1S)$	Photon Photon 0.000240	2980.4 MeV/c ²
$\chi_{c0}(1P)$	Photon Photon 0.000235	3414.76 MeV/c ²
$\chi_{c2}(1P)$	Photon Photon 0.000243	3556.2 MeV/c ²
$\eta_c(2S)$	Photon Photon $0. \times 10^{-4}$	3638. MeV/c ²
$\psi(2S)$	Photon Photon $0. \times 10^{-4}$	3686.093 MeV/c ²