



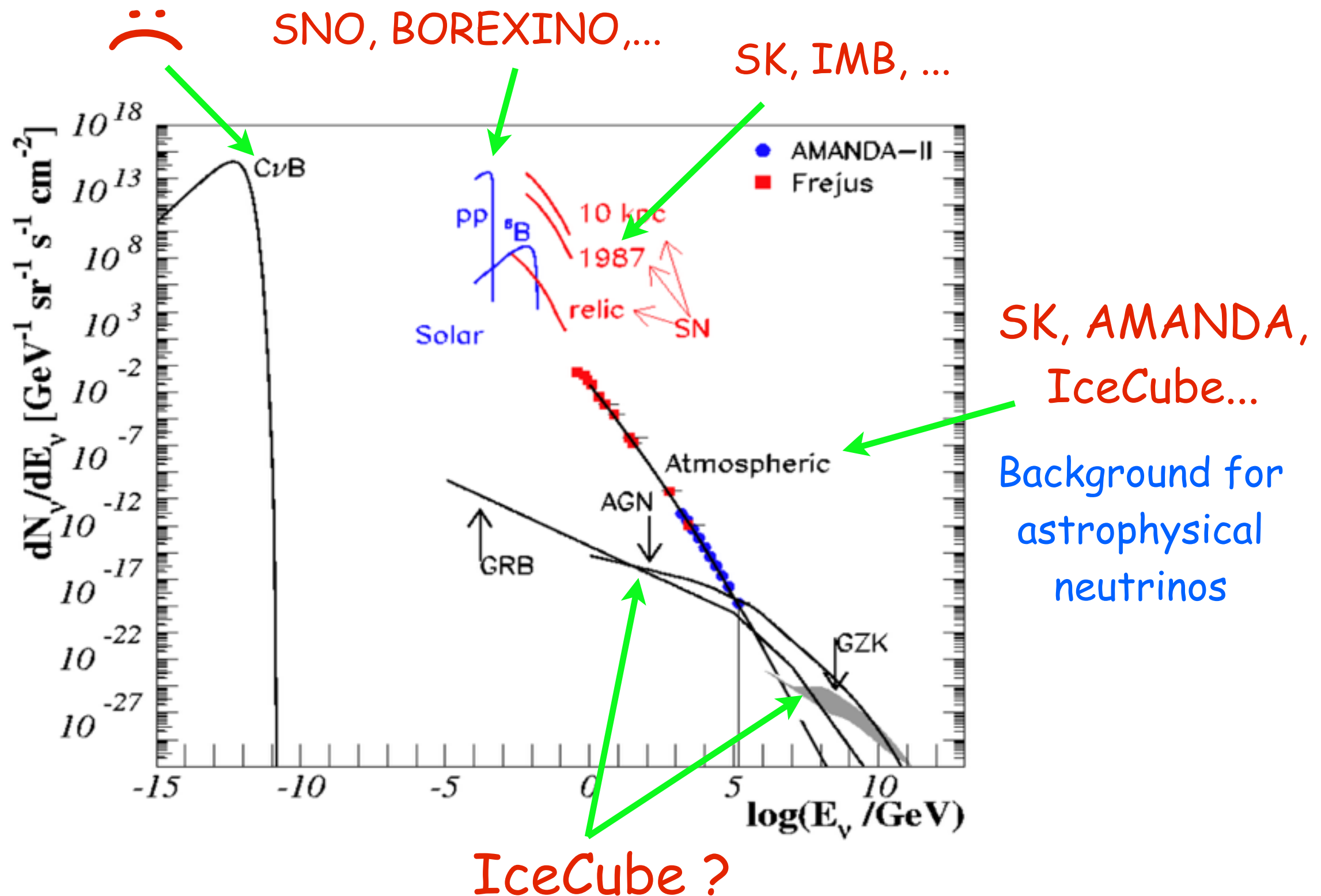
Interpreting the IceCube events by decaying dark matter hints and constraints

Arman Esmaili

Laboratori Nazionali del Gran Sasso (LNGS)
Theory Group

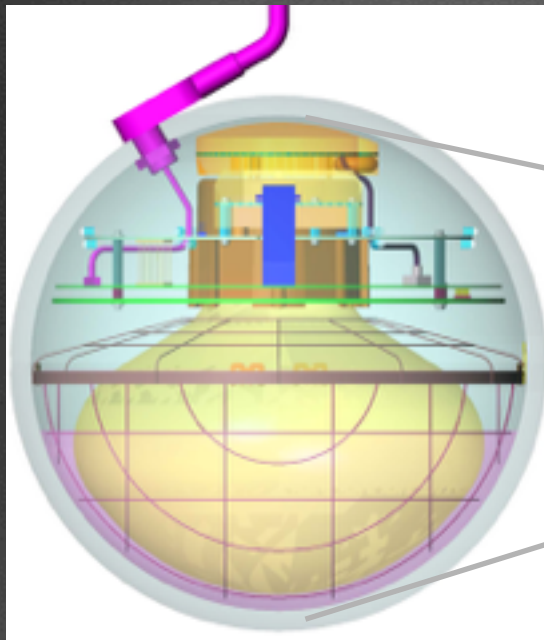
17/Sep/2015

Neutrino Sky

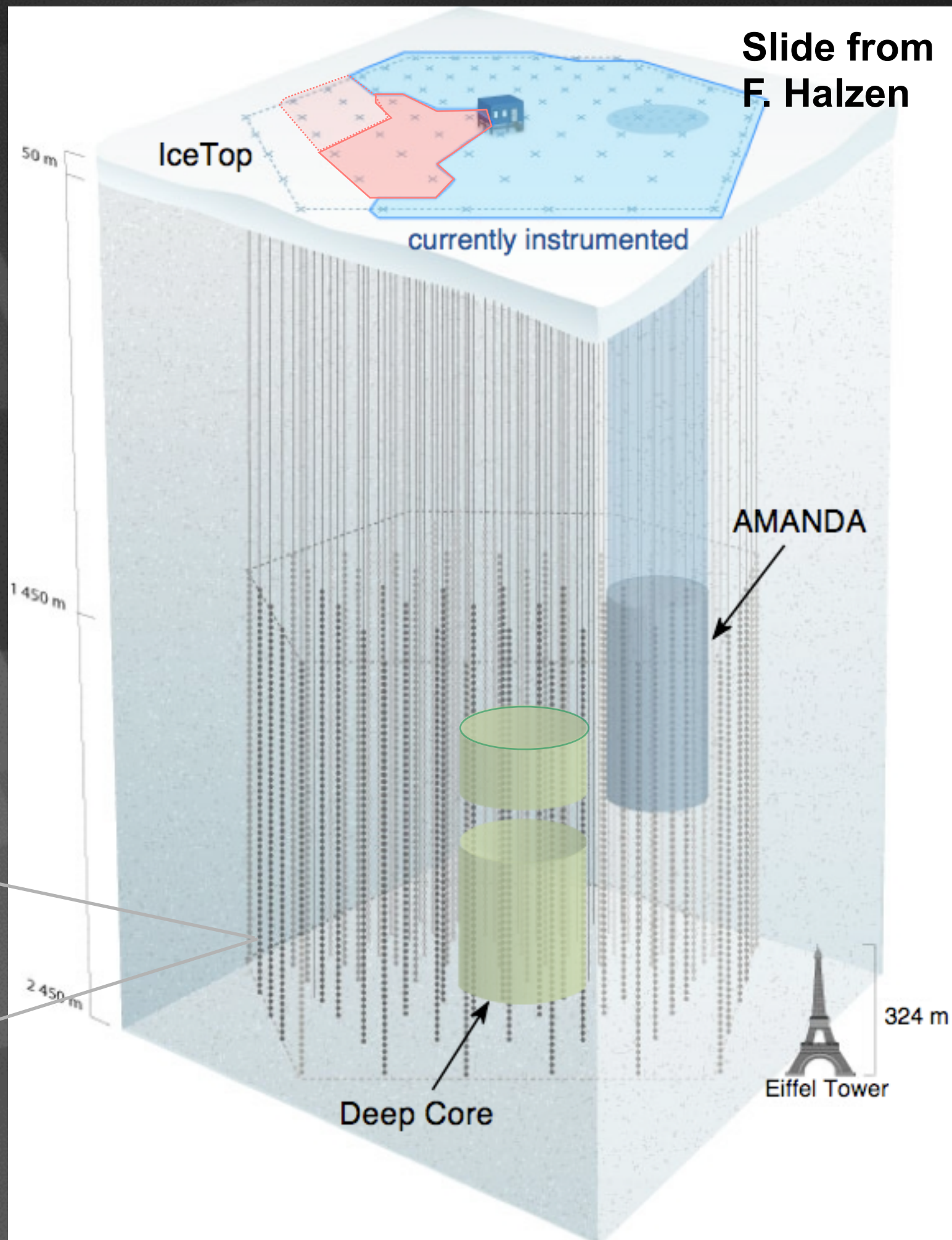


IceCube / Deep Core

- 5320 optical modules on 86 strings (+ IceTop)
- detects ~ 220 neutrinos and 1.7×10^8 muons per day
- threshold 10 GeV
- angular resolution < 1 degree

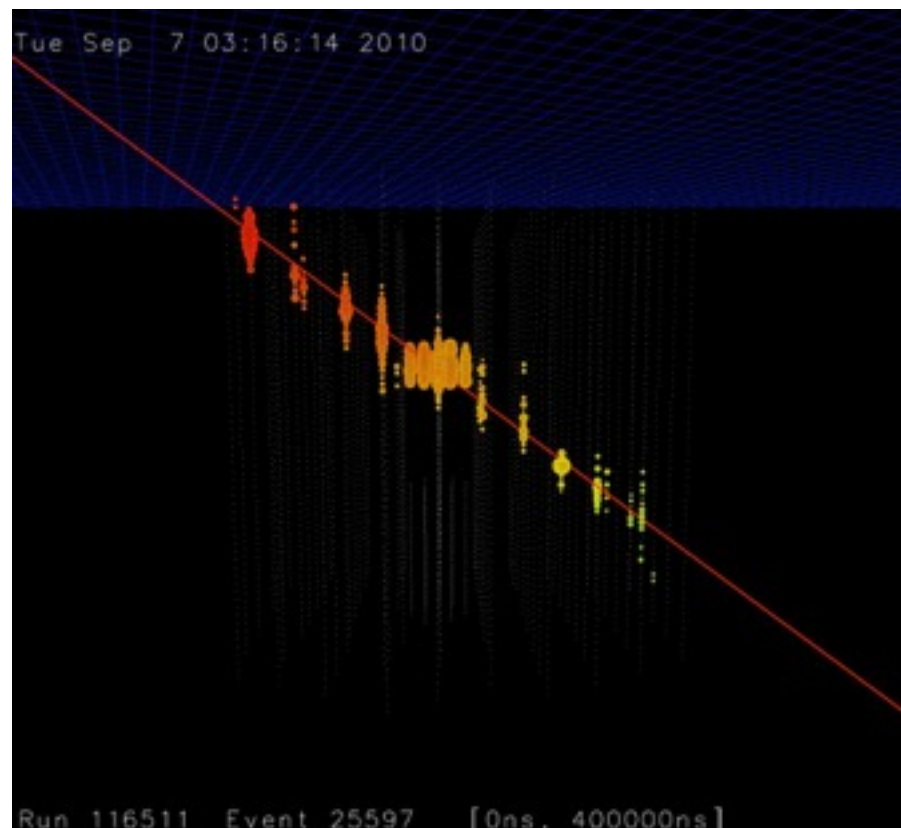
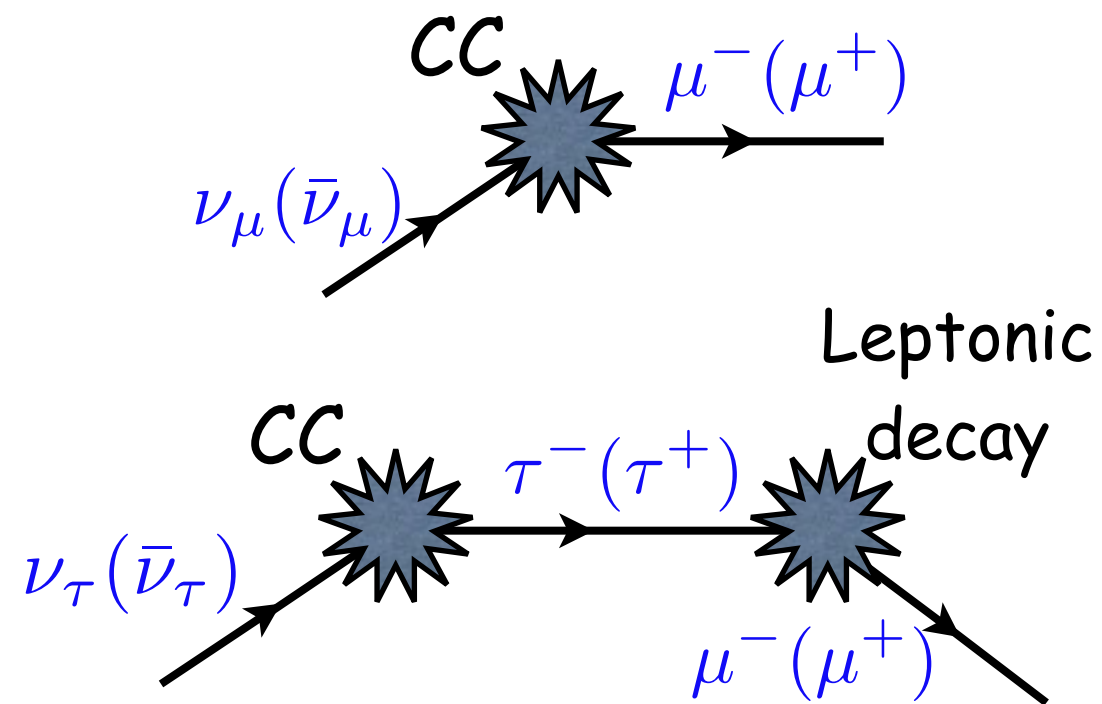


Digital Optical Module (DOM)

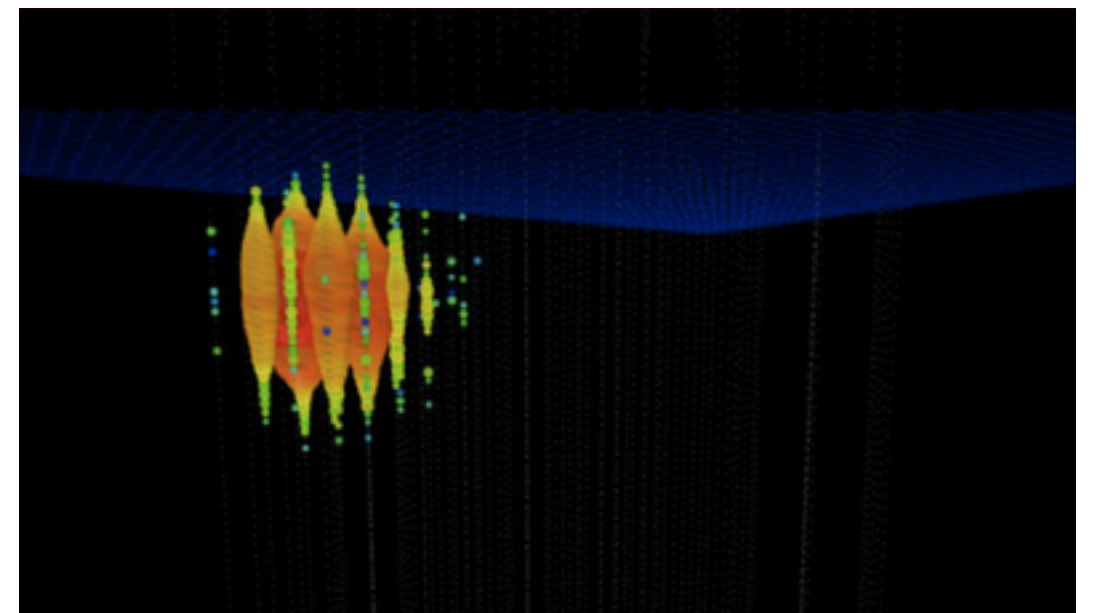
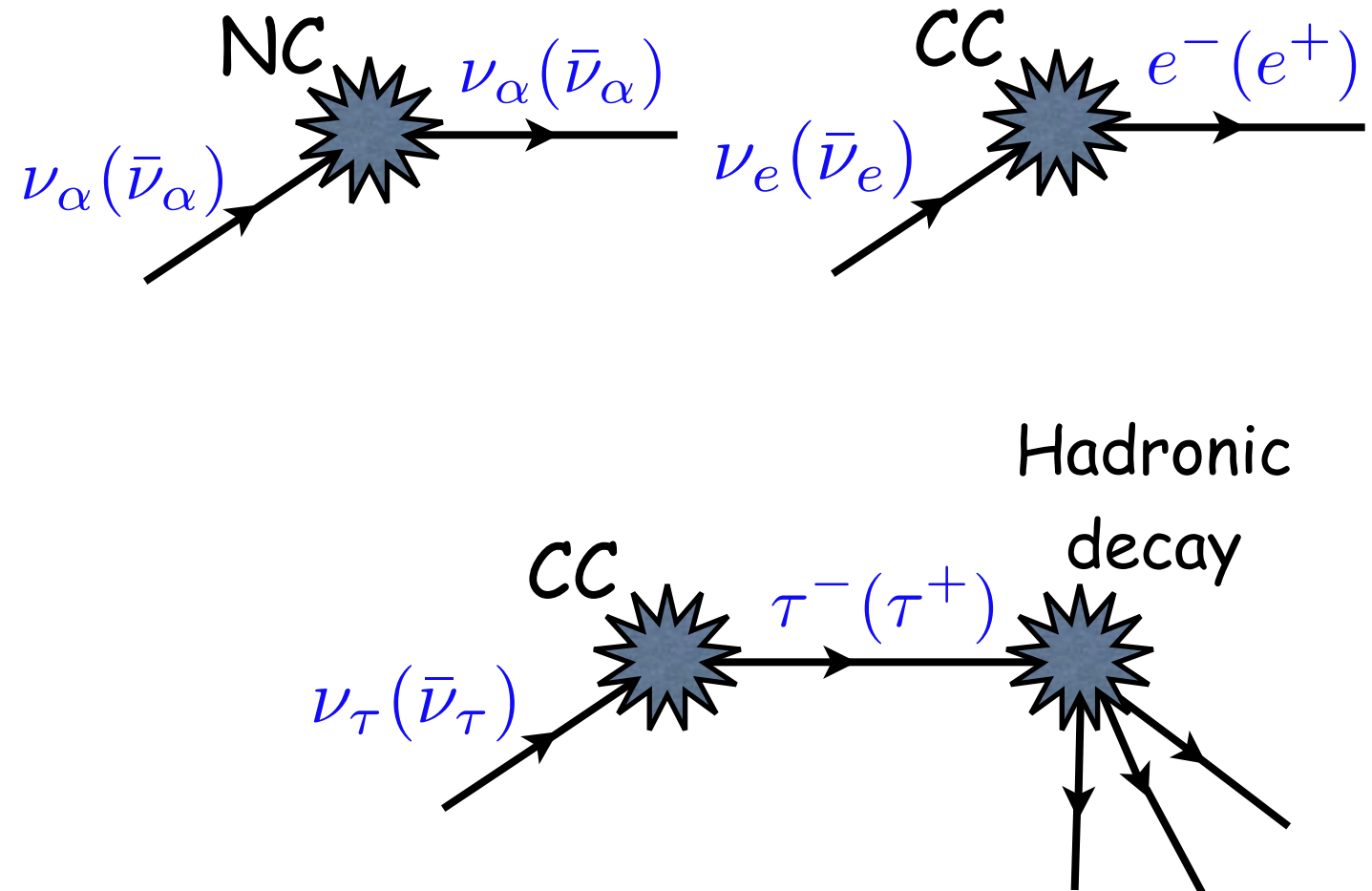


Flavoring at IceCube

muon-track events



cascade events



figures from
IceCube
website

Flavoring at IceCube

muon-track events

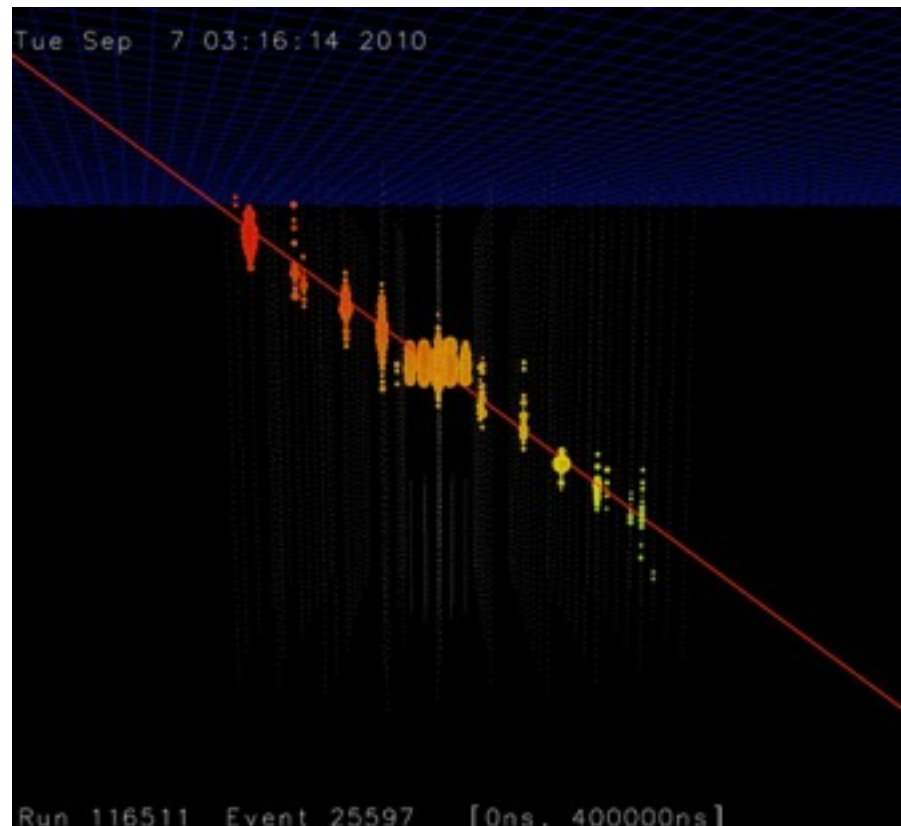
great angular
resolution ($< 1^\circ$)

moderate energy
resolution ($\sigma_E \sim E$)

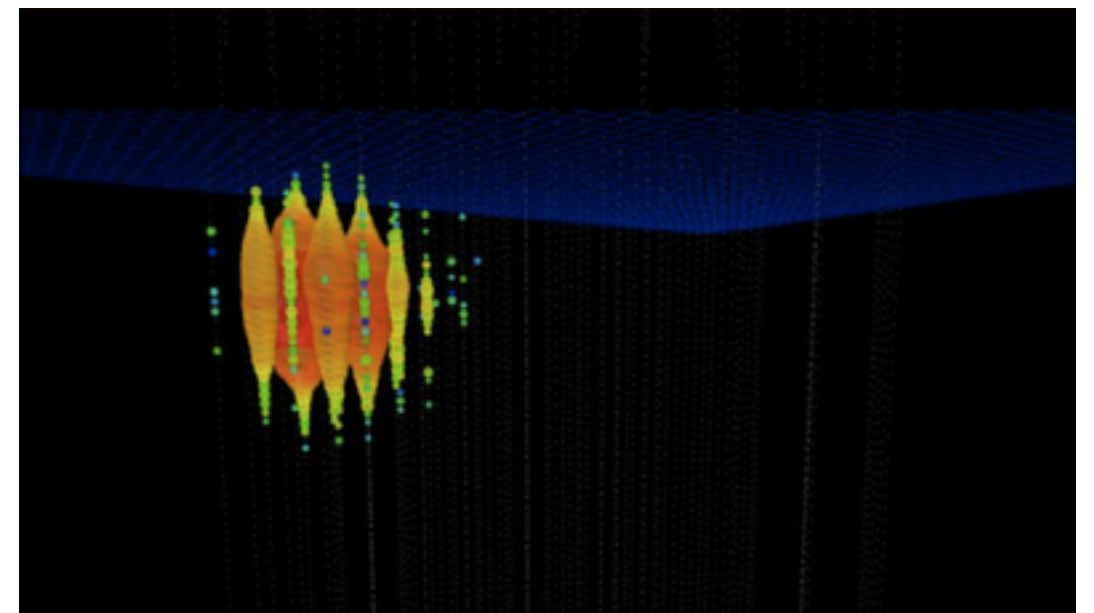
cascade events

poor angular resolution
($< 10^\circ - 20^\circ$)

great energy resolution
($\sigma_E \sim 0.15 \times E$)

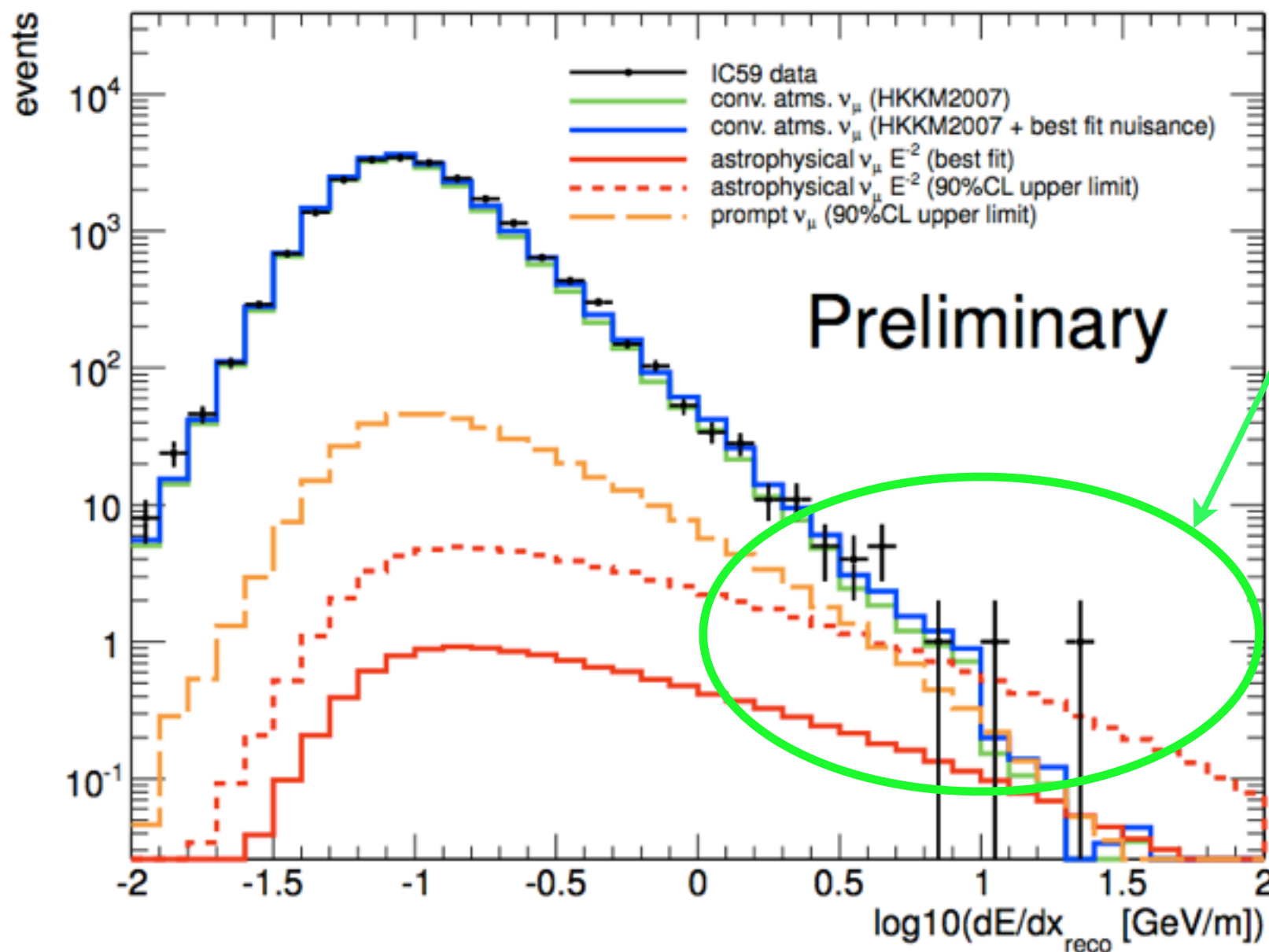


figures from
IceCube
website



Mission for IceCube began !

✓ muon-track events at IceCube-59, 348 days livetime.



excess in high energy tail
($\sim 300 \text{ TeV}$)

prompt atm neutrinos ?
or
astrophysical neutrinos?

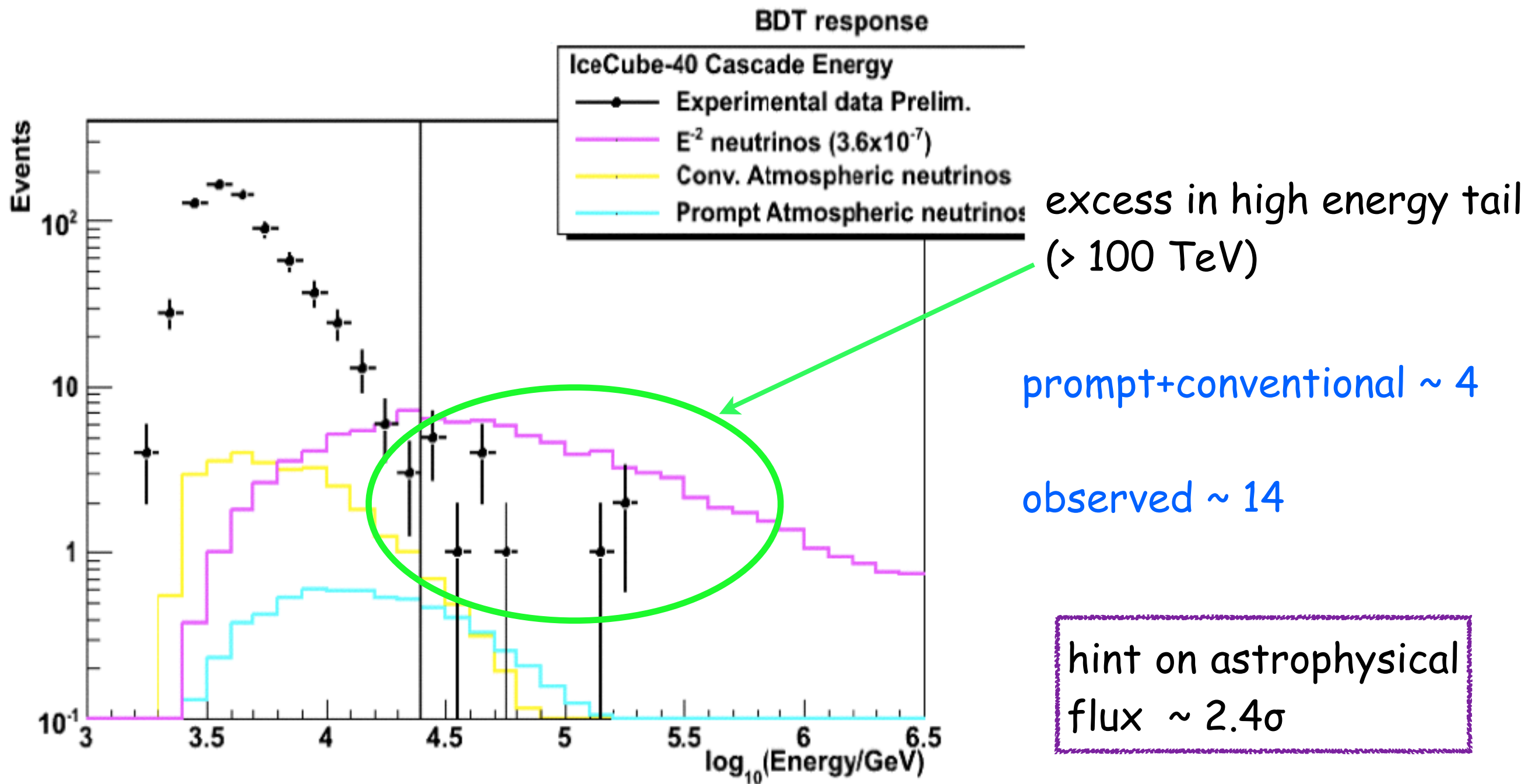
analysis shows preference
to astrophysical origin of
excess, significance $\sim 2.1\sigma$

A. Schukraft [IceCube Collaboration]

Nucl. Phys. Proc. Suppl. 266 (2013) [arXiv:1302.0127]

Mission for IceCube began !

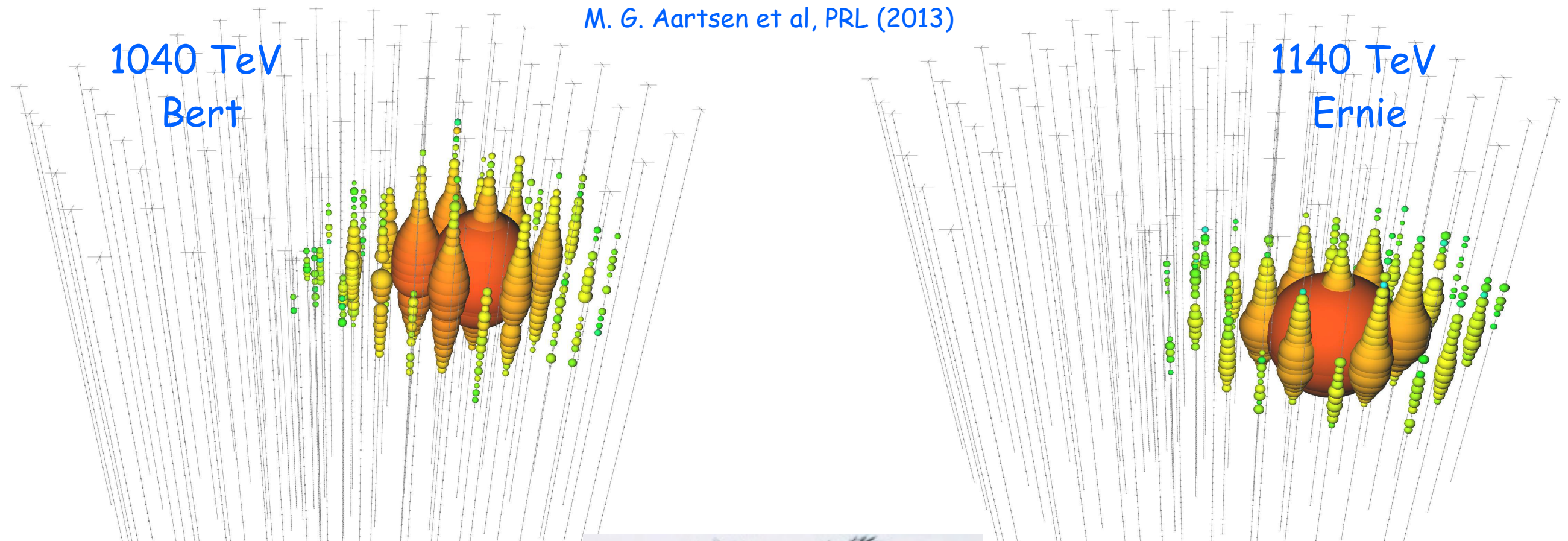
- ✓ cascade events at IceCube-40, 367 days livetime.



E. Middell, PhD thesis

Mission for IceCube began !

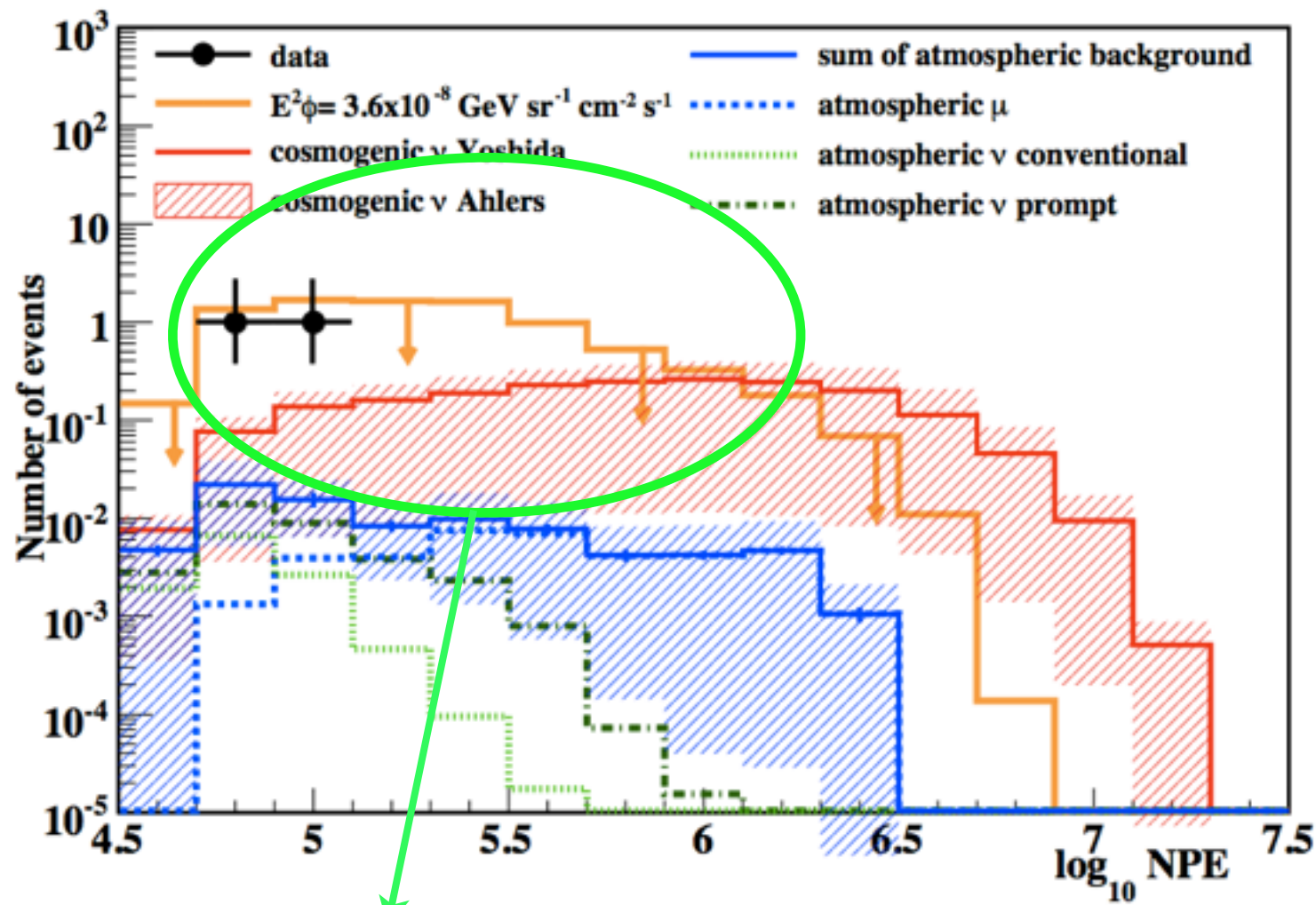
- ✓ The two PeV cascade events, 616 days livetime



Mission for IceCube began !

✓ The two PeV cascade events, 616 days livetime

M. G. Aartsen et al. [IceCube Collaboration],
Phys. Rev. Lett. 111 (2013), [arXiv:1304.5356]



it is like a cut-off at ~ PeV

expected bkg. (conventional+prompt)
 $\sim 0.08(-0.057)(+0.041)$ sys.

excess of events $\sim 2.8\sigma$

GZK ? too low energy, more events
should be seen in higher energies

astrophysical ? an E^{-2} spectrum
would give ~ 9 more events in
higher energies

flavor composition ?
NC of ν_α or CC of ν_e

demands more statistics

isotropy ?

HESE analysis

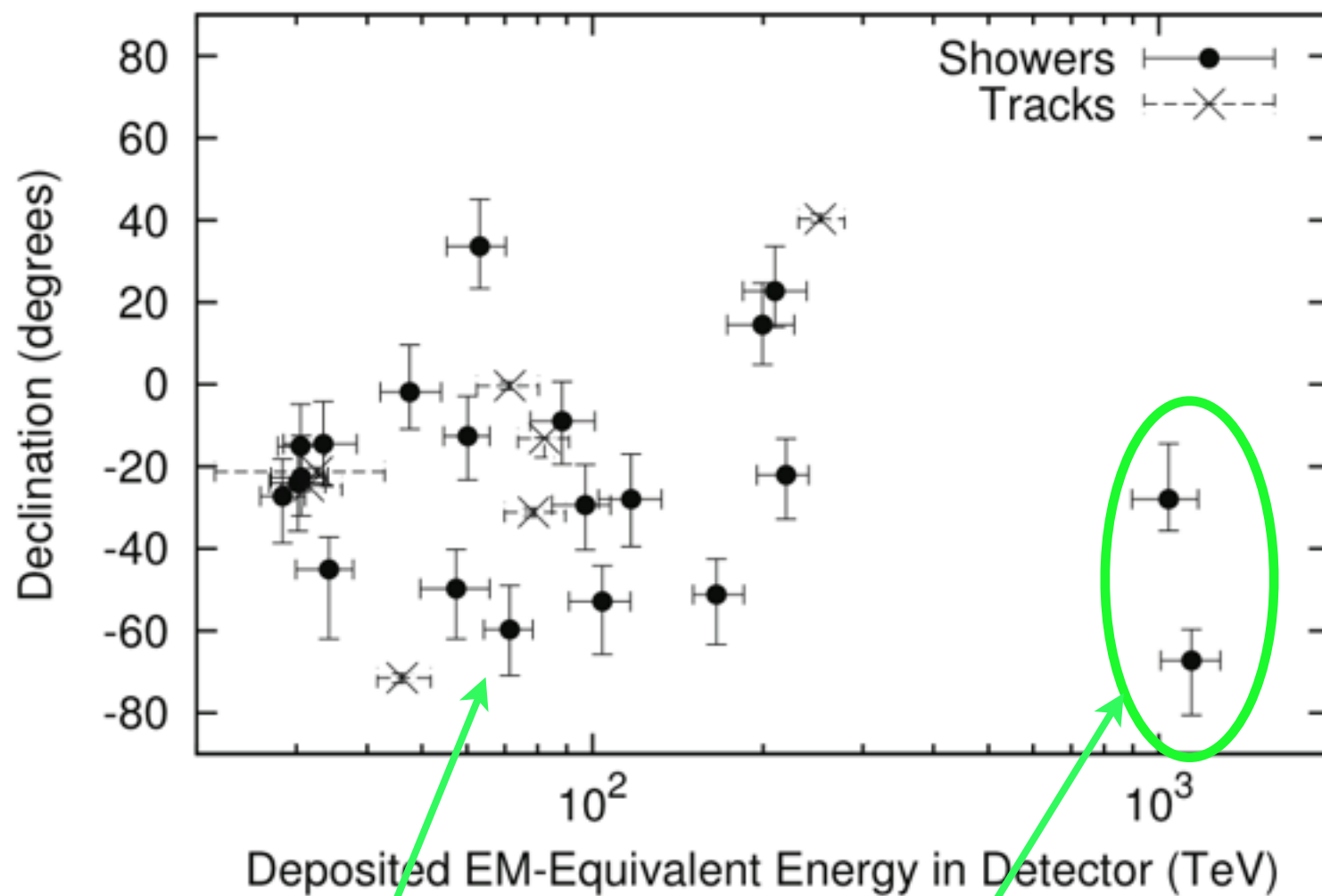
- ✓ select events interacting inside the detector only
- ✓ no light in the veto region
- ✓ veto for atmospheric muons and neutrinos (which are typically accompanied by muons)
- ✓ energy measurement: total absorption calorimetry



Mission for IceCube began !

✓ Looking for lower energy contained events, 662 days livetime

M. G. Aartsen et al. [IceCube Collaboration],
Science 342 (2013), [arXiv:1311.5238]



✓ 26 more events

✓ all the new events are lower in energy

✓ expected bkg. (conventional+prompt)
 $\sim 10.6(+4.5)(-3.5)$ sys.

excess of events $\sim 4.3\sigma$

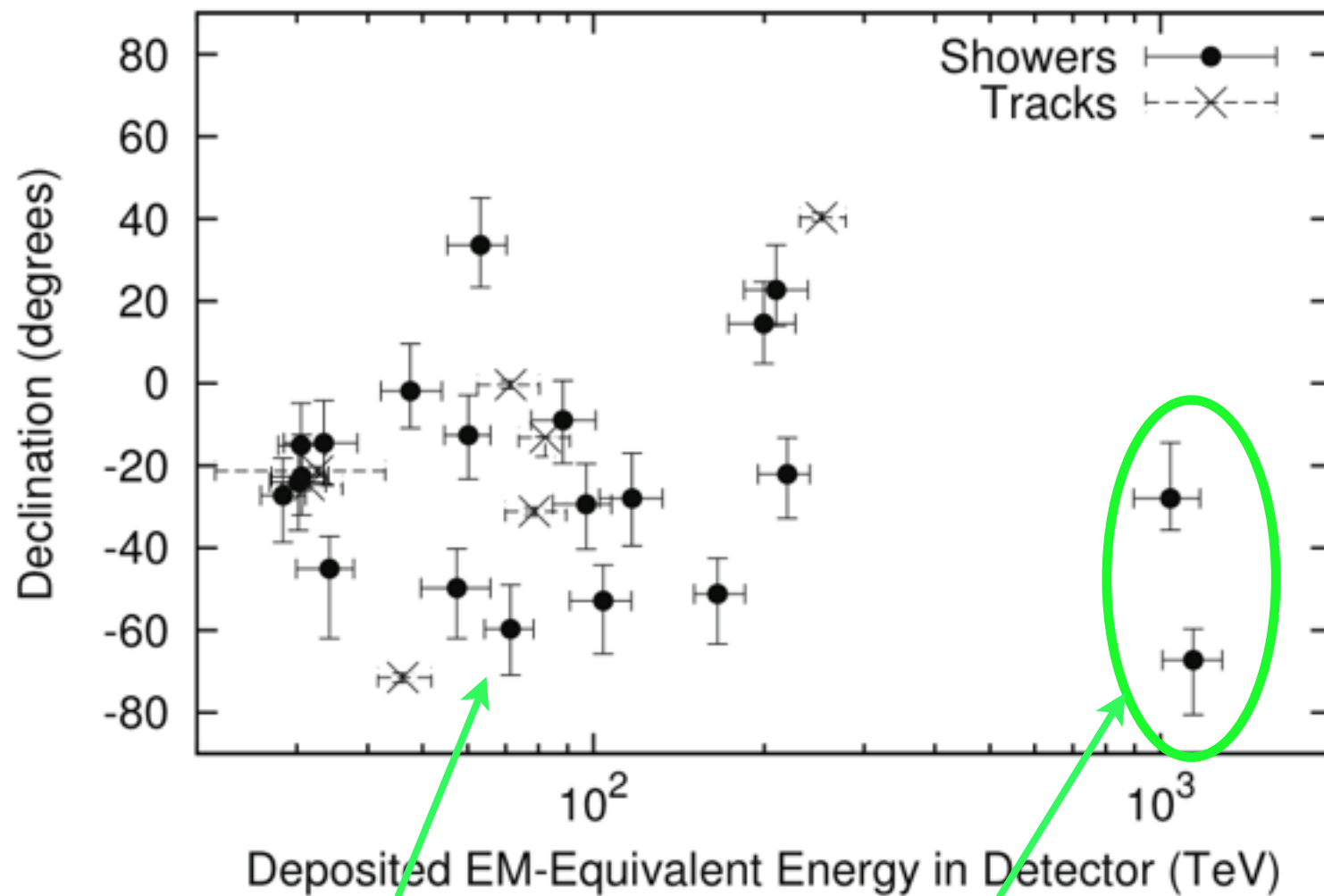
The whole family!

previous PeV
cascade events
(Bert and Ernie)

Mission for IceCube began !

✓ Looking for lower energy contained events, 662 days livetime

M. G. Aartsen et al. [IceCube Collaboration],
Science 342 (2013), [arXiv:1311.5238]



✓ 26 more events

✓ all the new events are lower in energy

✓ expected bkg. (conventional+prompt)
 $\sim 10.6(+4.5)(-3.5)$ sys.

excess of events $\sim 4.3\sigma$

The whole family!

previous PeV
cascade events
(Bert and Ernie)

which one?

atmospheric ?

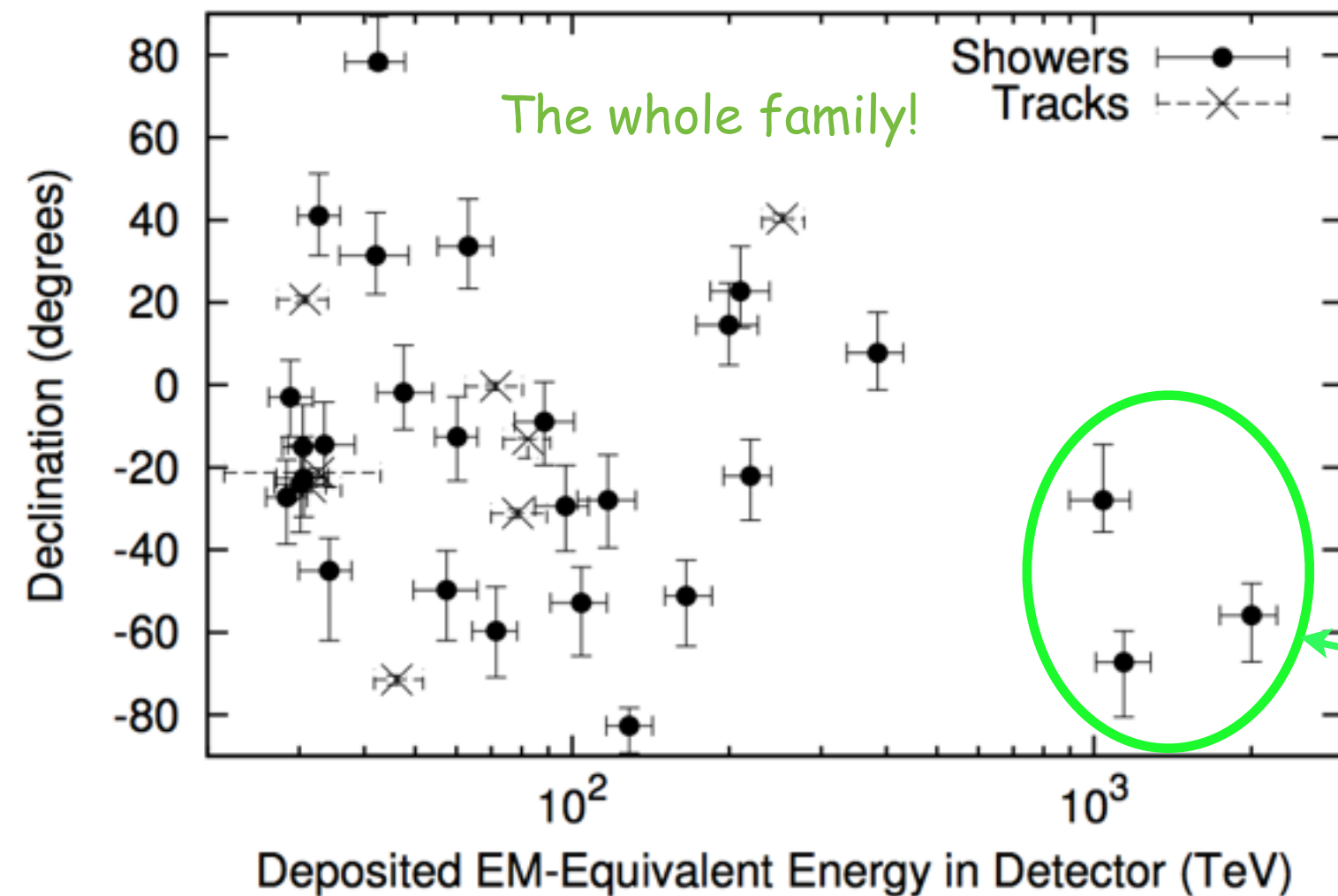
astrophysical ?

or something else ?

Mission for IceCube began !

✓ Looking for lower energy contained events, 988 days livetime

M. G. Aartsen et al. [IceCube Collaboration],
PRL 113 (2014), [arXiv:1405.5303]



✓ totally 37 events

✓ three events with energy $\sim$ PeV

✓ expected bkg. (conventional+prompt)
 $\sim 15.6(+10.1)(-5.8)$ sys.

excess of events $\sim 5.7\sigma$

which one?
atmospheric ?
astrophysical ?
or something else ?



Mission for IceCube began !

✓ Looking for lower energy contained events, 1347 days livetime

IPA 2015

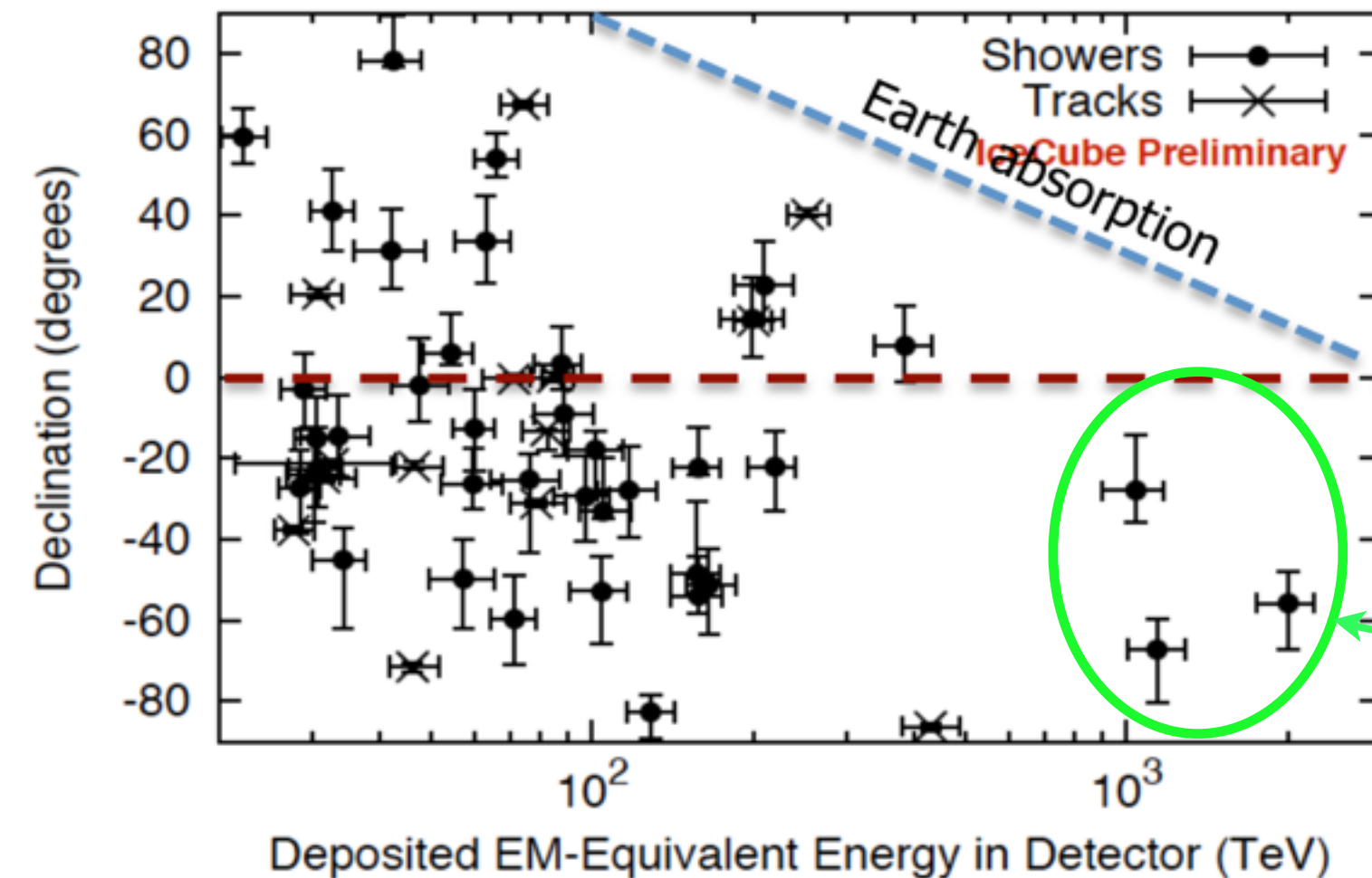
The whole family!

✓ totally 54 events

✓ still three events with energy $\sim$ PeV

4 years of data

excess of events $\sim 7\sigma$

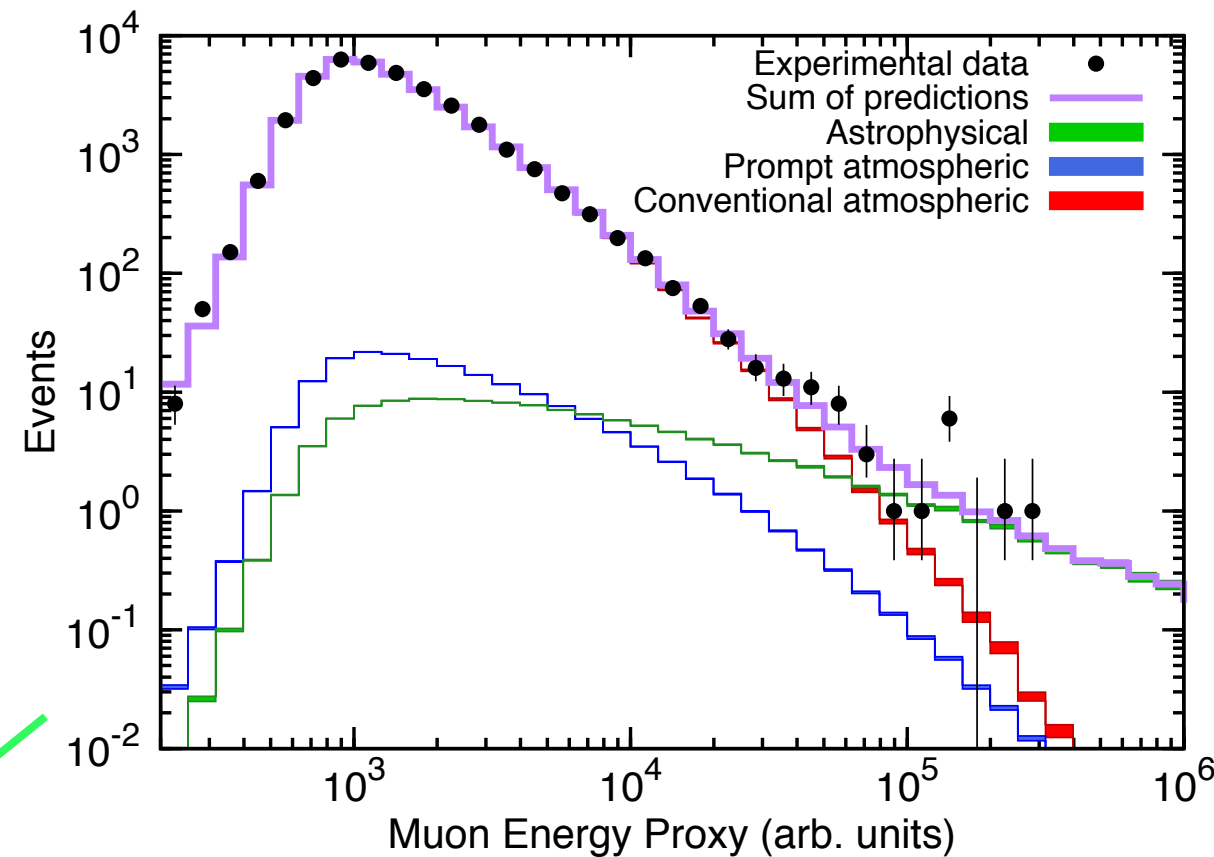


which one?
atmospheric ?
astrophysical ?
or something else ?



Mission for IceCube began !

✓ Looking for muon-track events, 660 days livetime



$$\Phi(E_\nu) = 9.9^{+3.9}_{-3.4} \times 10^{-19} \left(\frac{E_\nu}{100 \text{ TeV}} \right)^{-2} \text{ GeV}^{-1} \text{ cm}^{-2} \text{ sr}^{-1} \text{ s}^{-1}, \text{ fitting the index: } 2.2 \pm 0.2$$

lowering the energy
threshold of HESE
analysis to 1 TeV

$$\Phi(E_\nu) = 2.06^{+0.4}_{-0.3} \times 10^{-18} \left(\frac{E_\nu}{100 \text{ TeV}} \right)^{-2.46 \pm 0.12} \text{ GeV}^{-1} \text{ cm}^{-2} \text{ sr}^{-1} \text{ s}^{-1}$$

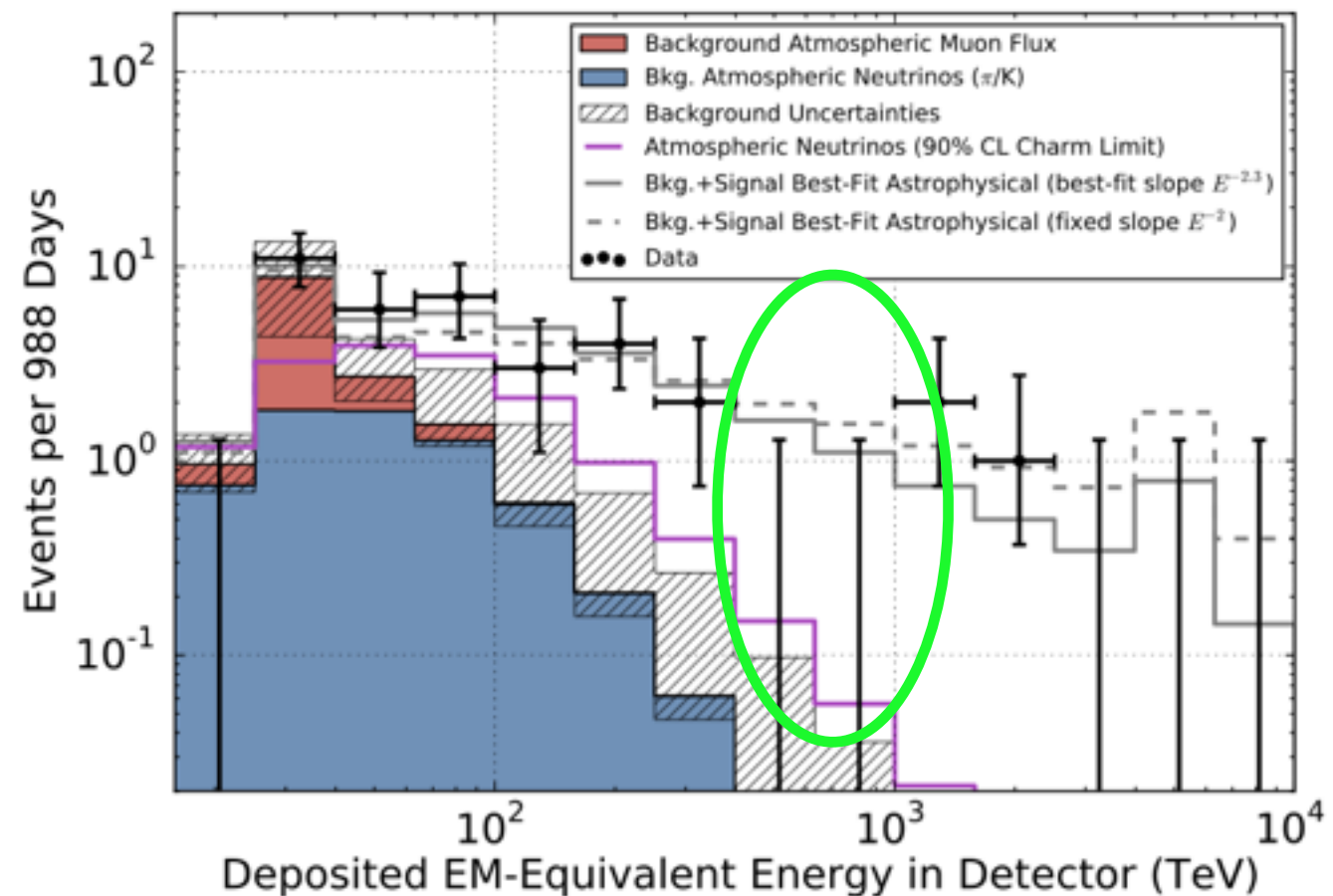
IceCube data

Problems with the astrophysical/atm interpretation of IceCube data

IceCube data

Problems with the astrophysical/atm interpretation of IceCube data

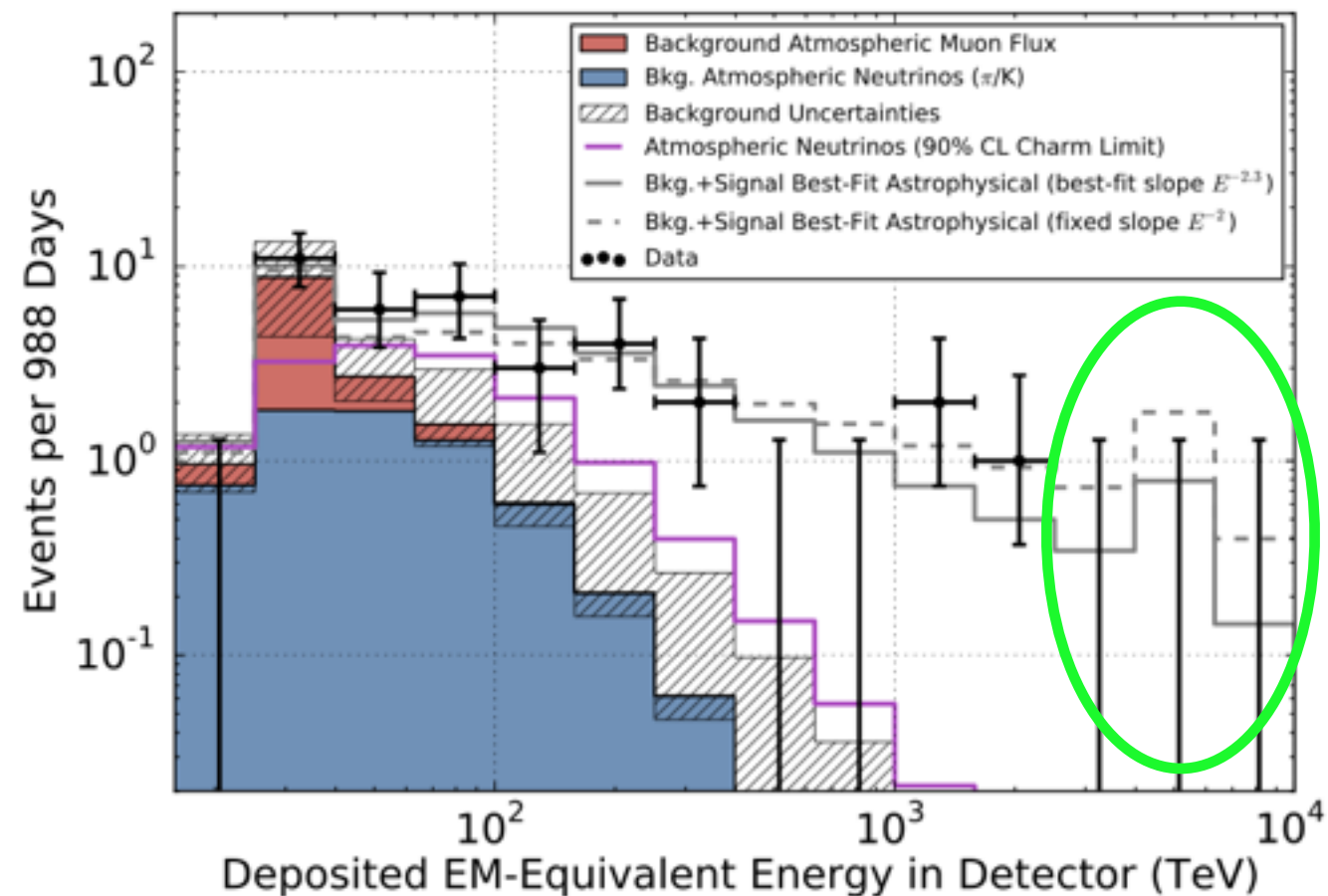
- ✓ deficit of events in the energy range $\sim (400 - 1000)$ TeV



IceCube data

Problems with the astrophysical/atm interpretation of IceCube data

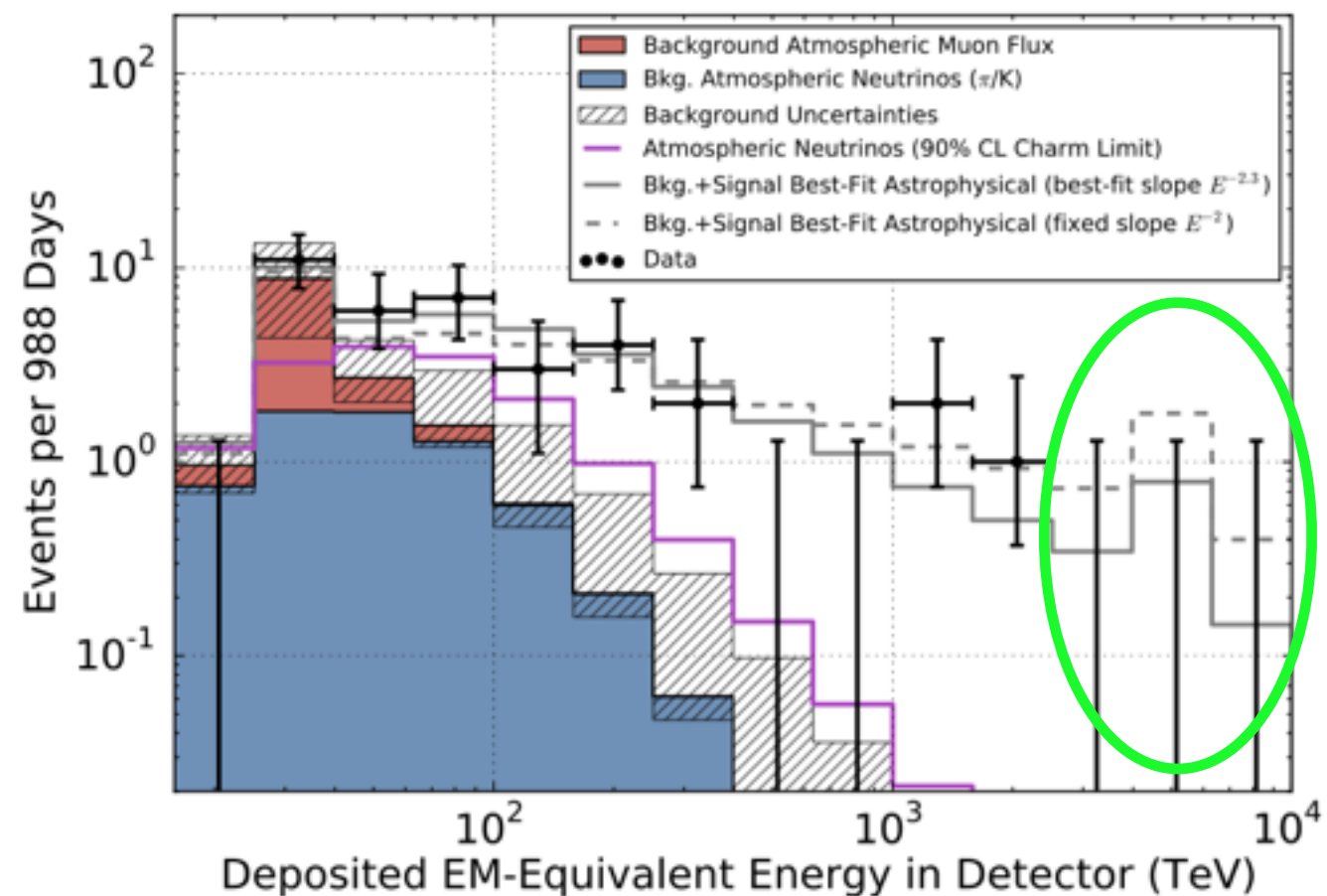
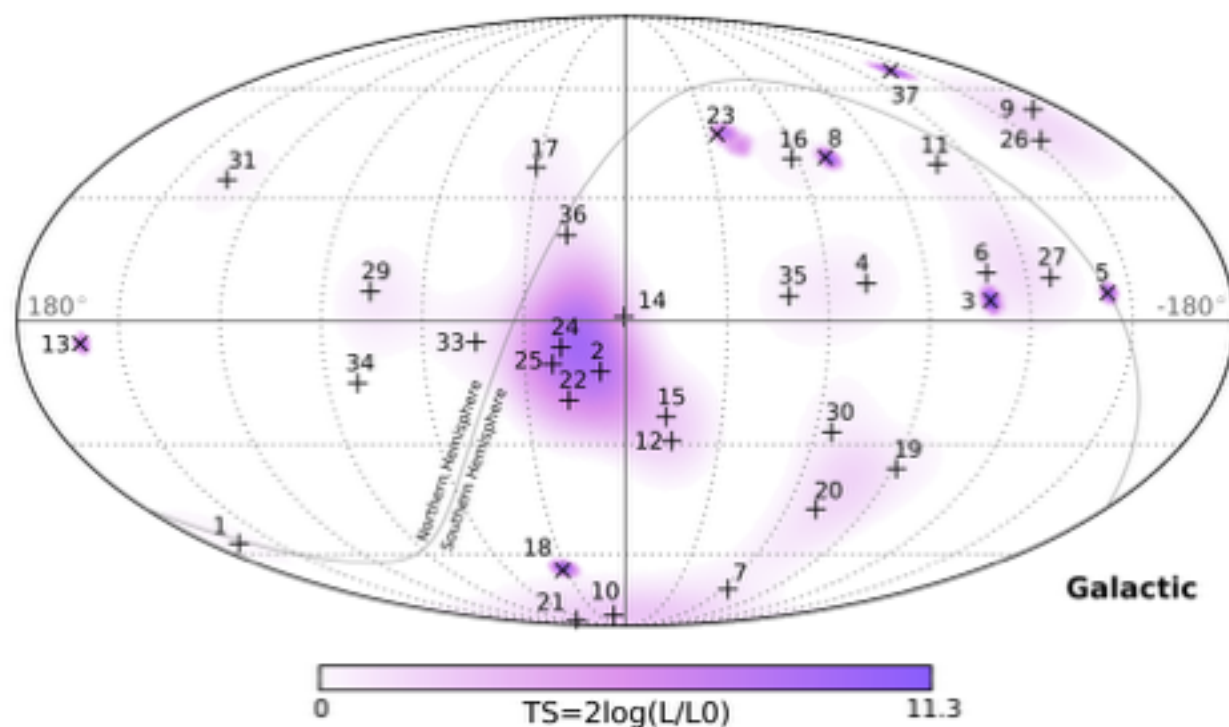
- ✓ deficit of events in the energy range $\sim (400 - 1000)$ TeV
- ✓ cut-off in events: no events observed with energy > 2 PeV



IceCube data

Problems with the astrophysical/atm interpretation of IceCube data

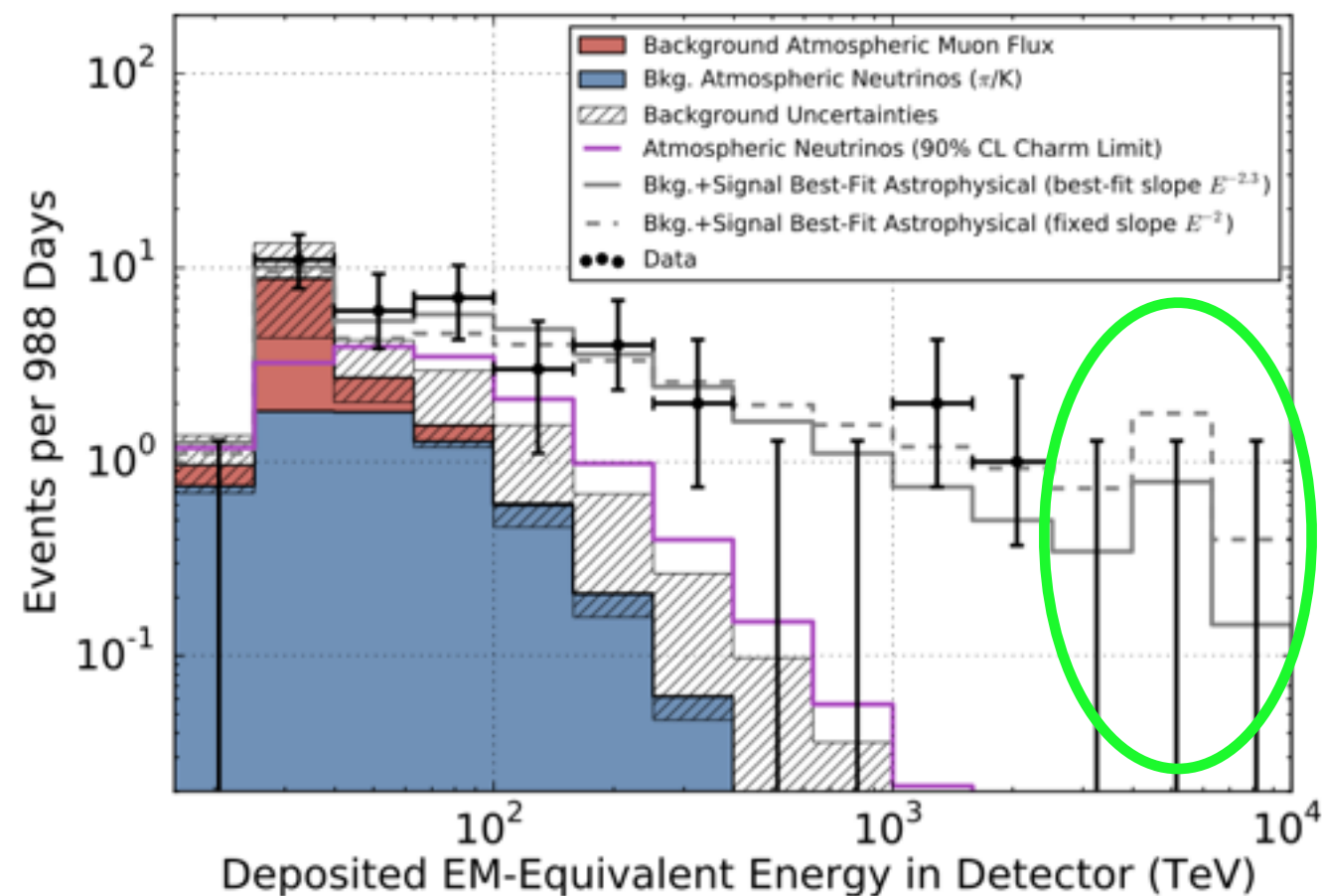
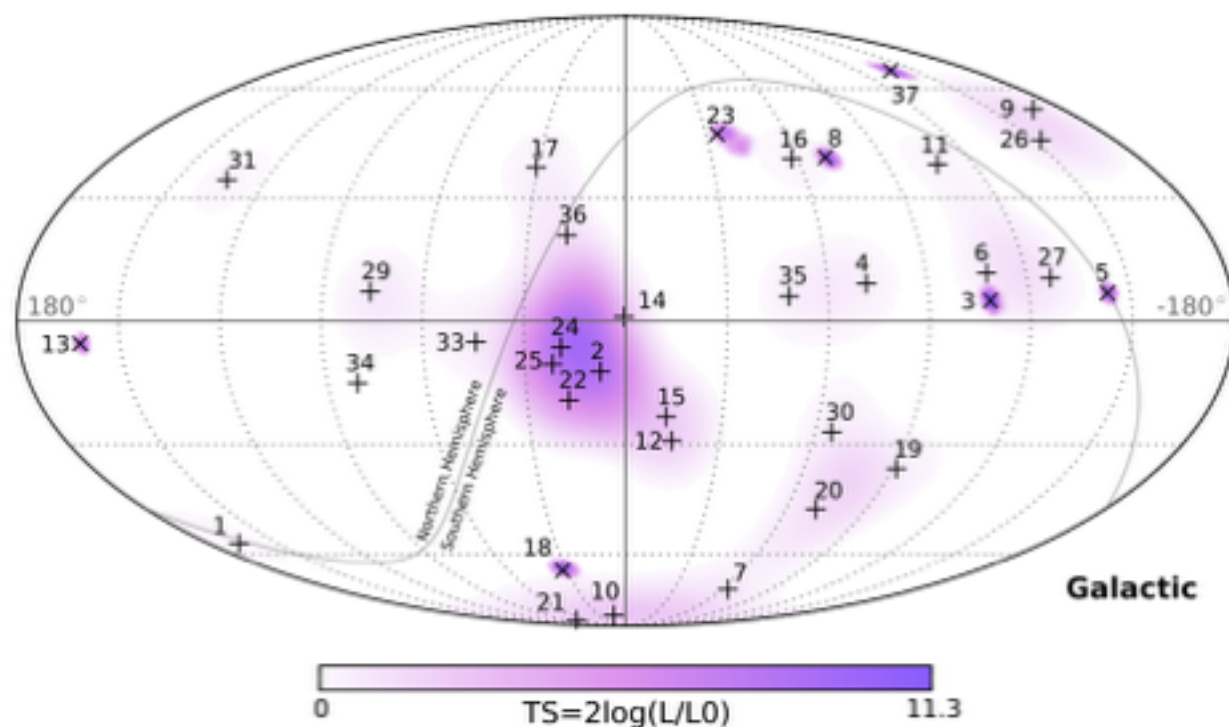
- ✓ deficit of events in the energy range $\sim (400 - 1000)$ TeV
- ✓ cut-off in events: no events observed with energy > 2 PeV
- ✓ angular distribution of events show mild anisotropies (enhanced toward GC)



IceCube data

Problems with the astrophysical/atm interpretation of IceCube data

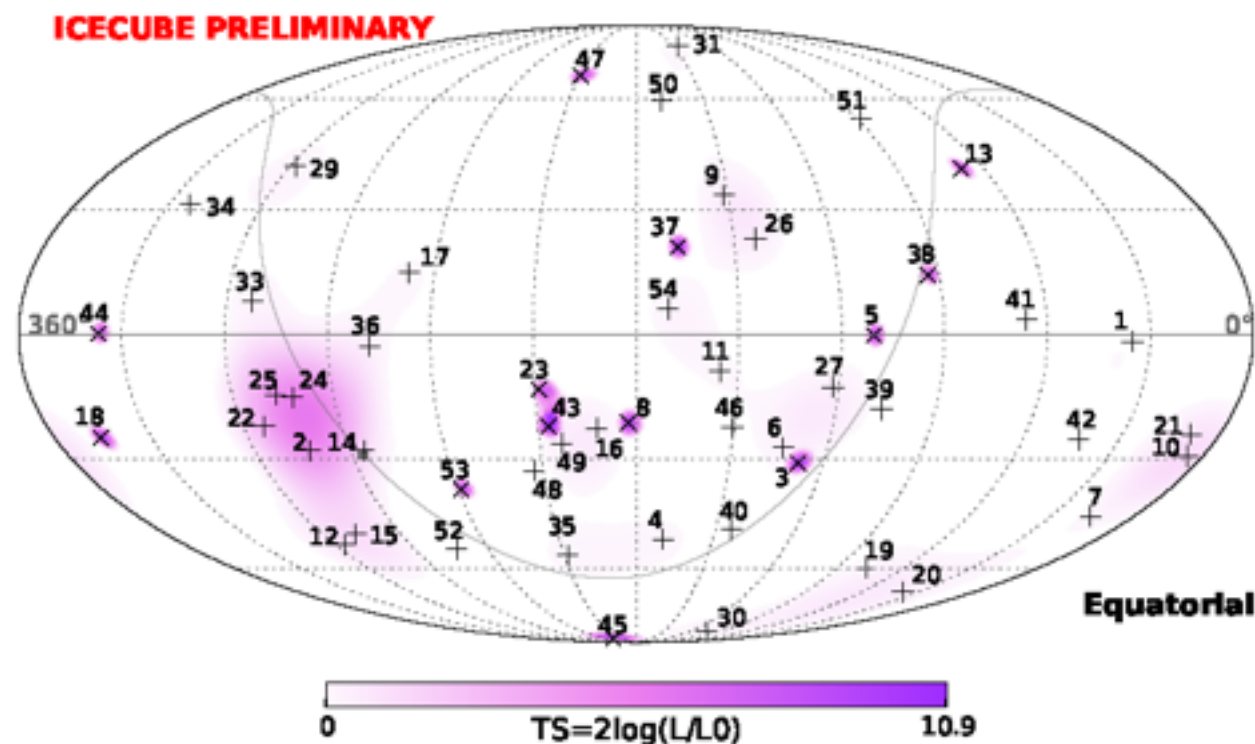
- ✓ deficit of events in the energy range $\sim (400 - 1000)$ TeV
- ✓ cut-off in events: no events observed with energy > 2 PeV
- ✓ angular distribution of events show mild anisotropies (enhanced toward GC)
- ⚠ none of the above-mentioned issues are significant



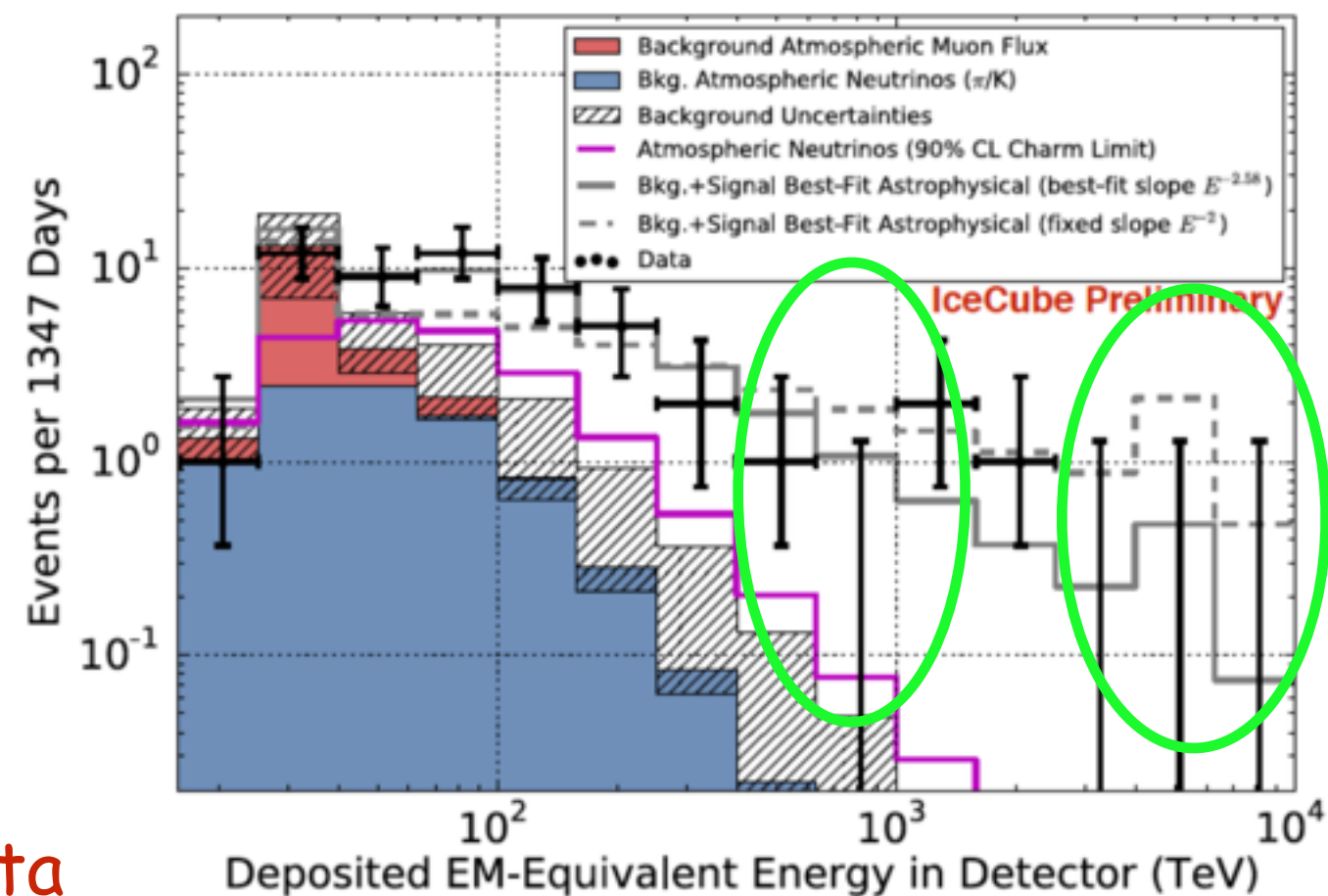
IceCube data

Problems with the astrophysical/atm interpretation of IceCube data

- ✓ deficit of events in the energy range $\sim (400 - 1000)$ TeV
- ✓ cut-off in events: no events observed with energy > 2 PeV
- ✓ angular distribution of events show mild anisotropies (enhanced toward GC)
- ⚠ none of the above-mentioned issues are significant



4 years of data



Interpretations of IceCube data

✓ "Conventional" interpretations of IceCube data

Cosmic ray sources

GRBs

Galaxy clusters

Star-forming galaxies

AGNs

Fermi bubbles

Galactic Center activities

-
-
-

M. D. Kistler, T. Stanev and H. Yuksel, arXiv:1301.1703 [astro-ph.HE]

K. Murase and K. Ioka, Phys. Rev. Lett. 111, no. 12, 121102 (2013) [arXiv:1306.2274 [astro-ph.HE]].

K. Murase, M. Ahlers and B. C. Lacki, Phys. Rev. D 88, no. 12, 121301 (2013) [arXiv:1306.3417 [astro-ph.HE]].

L. A. Anchordoqui, H. Goldberg, M. H. Lynch, A. V. Olinto, T. C. Paul and T. J. Weiler, arXiv:1306.5021 [astro-ph.HE].

R. Laha, J. F. Beacom, B. Dasgupta, S. Horiuchi and K. Murase, Phys. Rev. D 88, 043009 (2013) [arXiv:1306.2309 [astro-ph.HE]].

S. Razzaque, Phys. Rev. D 88, 081302 (2013) [arXiv:1309.2756 [astro-ph.HE]].

C. Y. Chen, P. S. Bhupal Dev and A. Soni, Phys. Rev. D 89, no. 3, 033012 (2014) [arXiv:1309.1764 [hep-ph]].

M. Ahlers and K. Murase, Phys. Rev. D 90, 023010 (2014) [arXiv:1309.4077 [astro-ph.HE]].

I. Tamborra, S. Ando and K. Murase, JCAP 1409, no. 09, 043 (2014) [arXiv:1404.1189 [astro-ph.HE]].

M. Kachelriess and S. Ostapchenko, Phys. Rev. D 90, 083002 (2014) [arXiv:1405.3797 [astro-ph.HE]].

M. Ahlers and F. Halzen, arXiv:1406.2160 [astro-ph.HE].

Y. Bai, A. J. Barger, V. Barger, R. Lu, A. D. Peterson and J. Salvado, Phys. Rev. D 90, 063012 (2014) [arXiv:1407.2243 [astro-ph.HE]].

A. Bhattacharya, R. Enberg, M. H. Reno and I. Sarcevic, arXiv:1407.2985 [astro-ph.HE].

C. Lunardini, S. Razzaque, K. T. Theodoseou and L. Yang, Phys. Rev. D 90, 023016 (2014) [arXiv:1311.7188 [astro-ph.HE]].

For a review

L. A. Anchordoqui, V. Barger, I. Cholis, H. Goldberg, D. Hooper, A. Kusenko, J. G. Learned and D. Marfatia et al., Journal of High Energy Astrophysics 1-2, 1 (2014) [arXiv:1312.6587 [astro-ph.HE]].

-
-
-

Interpretations of IceCube data

✓ "New Physics" interpretations of IceCube data

Lepto-quarks

Secret neutrino interactions

resonant absorption on
cosmic neutrino background

supermassive long-lived
particles

Dark matter decay

-
-
-

Y. Ema, R. Jinno and T. Moroi, Phys. Lett. B 733, 120 (2014) [arXiv:1312.3501 [hep-ph]].

K. Ioka and K. Murase, PTEP 2014, (2014) [arXiv:1404.2279 [astro-ph.HE]].

K. C. Y. Ng and J. F. Beacom, Phys. Rev. D 90, 065035 (2014) [arXiv:1404.2288 [astro-ph.HE]].

M. Ibe and K. Kaneta, Phys. Rev. D 90, 053011 (2014) [arXiv:1407.2848 [hep-ph]].

V. Barger and W. Y. Keung, Phys. Lett. B 727, 190 (2013) [arXiv:1305.6907 [hep-ph]].

B. Feldstein, A. Kusenko, S. Matsumoto and T. T. Yanagida, Phys. Rev. D 88, no. 1, 015004 (2013) [arXiv:1303.7320 [hep-ph]].

Y. Bai, R. Lu and J. Salvado, arXiv:1311.5864 [hep-ph].

A. Bhattacharya, M. H. Reno and I. Sarcevic, JHEP 1406, 110 (2014) [arXiv:1403.1862 [hep-ph]].

J. Zavala, Phys. Rev. D 89, 123516 (2014) [arXiv:1404.2932 [astro-ph.HE]].

A. Bhattacharya, R. Gandhi and A. Gupta, arXiv:1407.3280 [hep-ph].

C. Rott, K. Kohri and S. C. Park, arXiv:1408.4575 [hep-ph].

T. Higaki, R. Kitano and R. Sato, JHEP 1407, 044 (2014) [arXiv:1405.0013 [hep-ph]].

A. Esmaili and P. D. Serpico, JCAP 1311, 054 (2013), arXiv:1308.1105

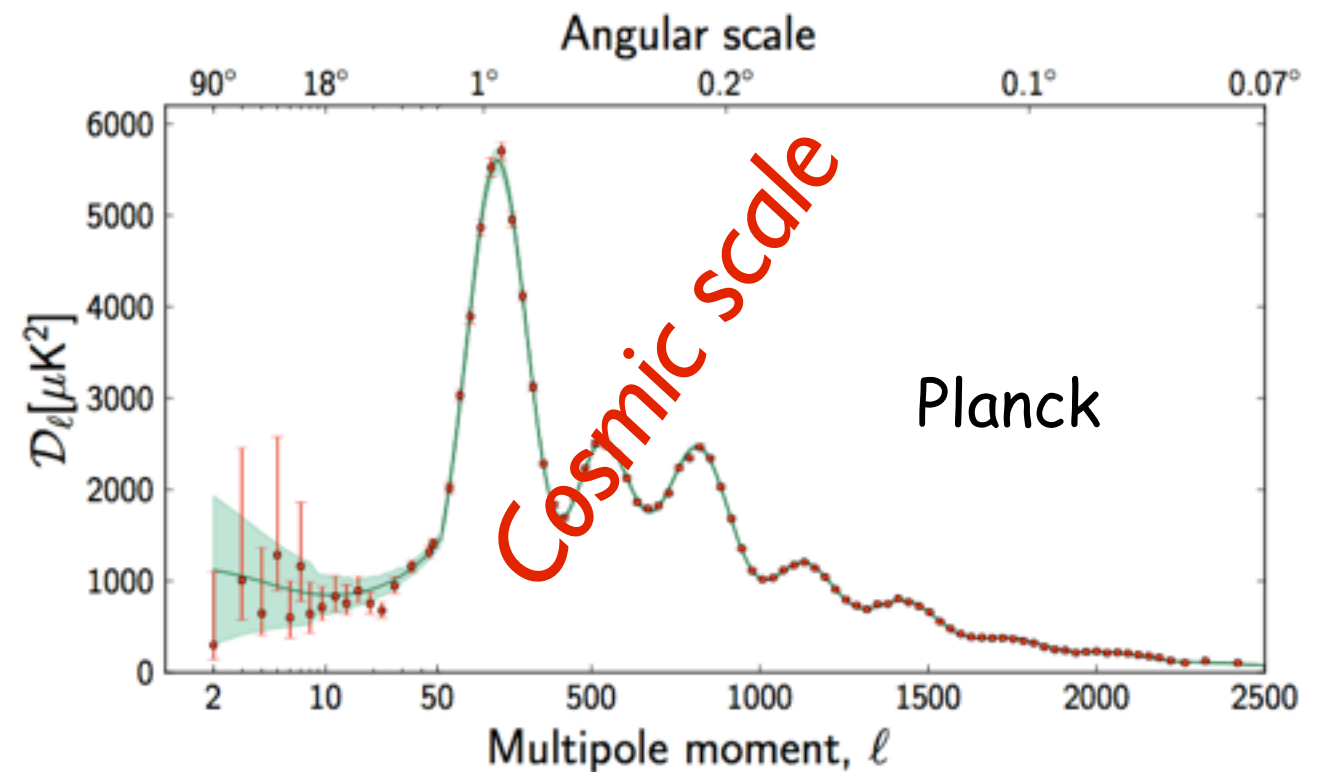
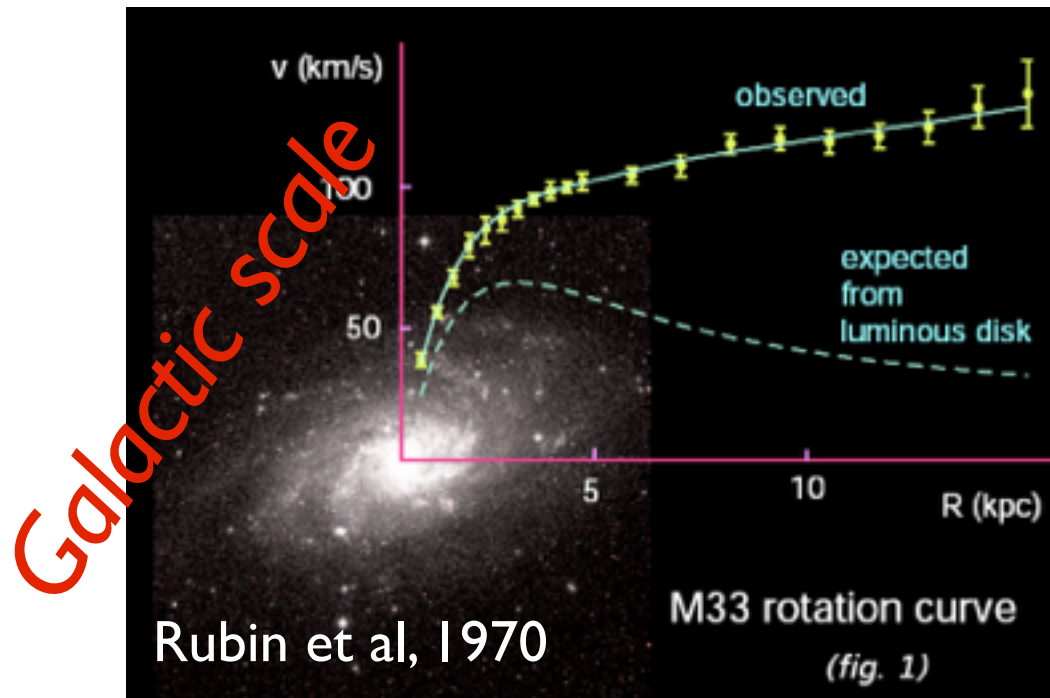
A. Esmaili, S. K. Kang and P. D. Serpico, JCAP 1412, 054 (2014), arXiv:1410.5979

A. Esmaili and P. D. Serpico, arXiv:1505.06486

-
-

A note on Dark Matter

DM exist!



A note on Dark Matter

DM exist!

What We Know?

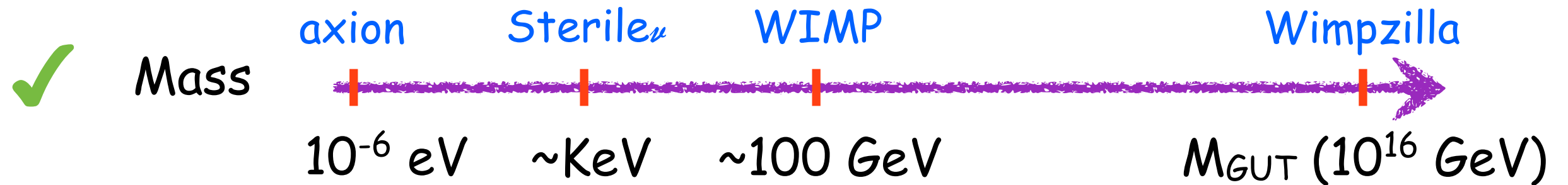
- ✓ Non Baryonic
- ✓ No electric and color charges
- ✓ Cold (or perhaps warm)
- ✓ Long lived (not necessarily stable)

All of these come from gravitational effects

A note on Dark Matter

DM exist!

What We Do Not Know?



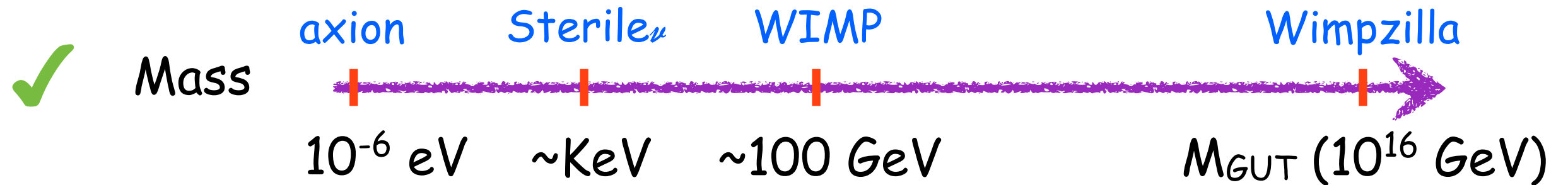
⚠ "WIMP" paradigm ?

Note that WIMP paradigm is a "particle physics" conjecture, needs to be validated at colliders

A note on Dark Matter

DM exist!

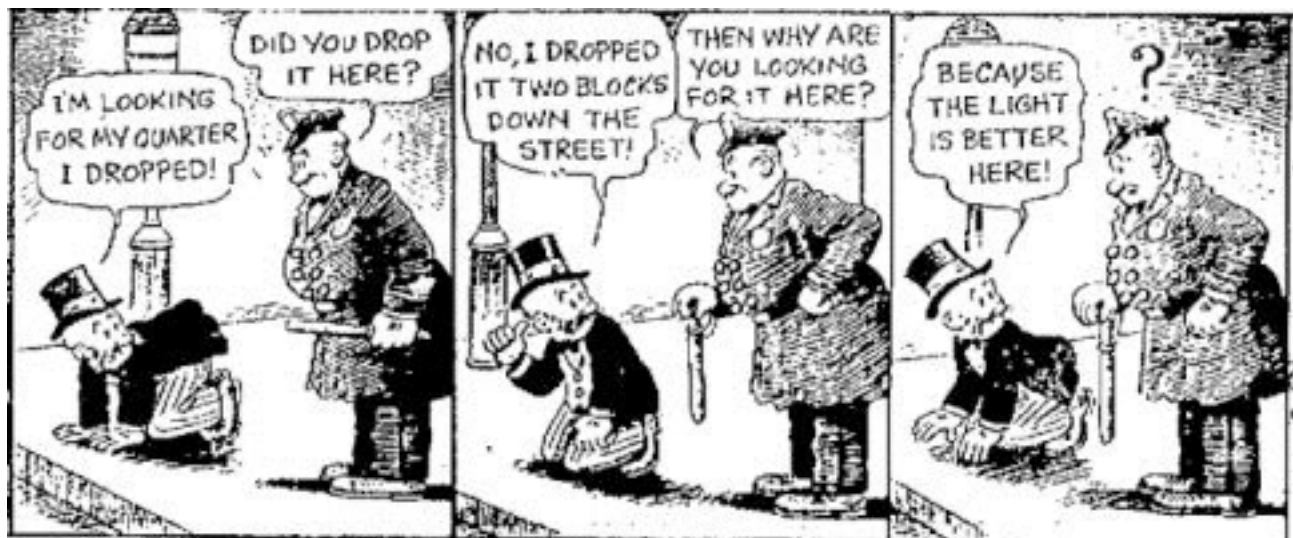
What We Do Not Know?



⚠ "WIMP" paradigm ?

Note that WIMP paradigm is a "particle physics" conjecture, needs to be validated at colliders

caution: streetlight effect



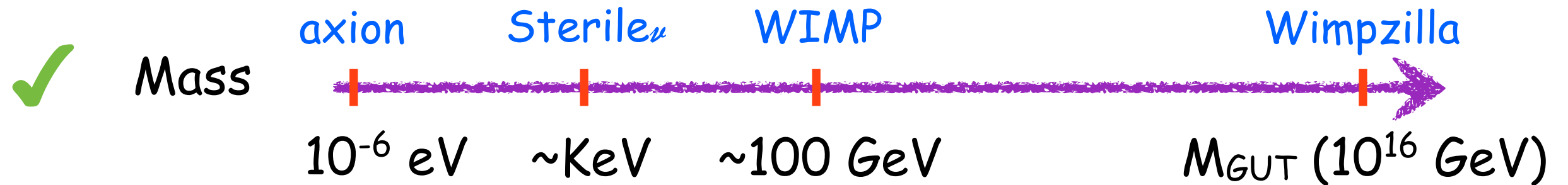
Mulla
Nasreddin



A note on Dark Matter

DM exist!

What We Do Not Know?



⚠ "WIMP" paradigm ?

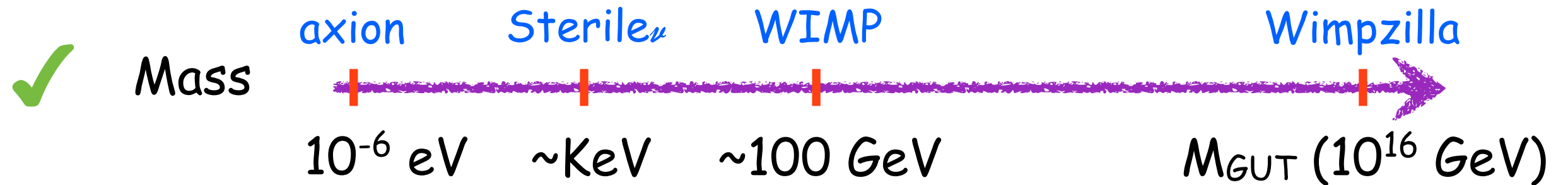
Note that WIMP paradigm is a "particle physics" conjecture, needs to be validated at colliders



A note on Dark Matter

DM exist!

What We Do Not Know?



⚠ "WIMP" paradigm ?

Note that WIMP paradigm is a "particle physics" conjecture, needs to be validated at colliders

✓ Lifetime: stable (∞) or
 $\tau_{\text{DM}} > 4.3 \times 10^{17} \text{ s}$ (age of Universe)

$\tau_{\text{DM}} > 2.2 \times 10^{19} \text{ s}$ (CMB) Y. Gong and X. Chen, PRD77 (2008), arXiv:0802.2296

✓ Possible decay and/or annihilation channels

✓ ...

Limits on lifetime from neutrino experiments

before recent IceCube data

$$16 - 2.5 \times 10^3 \text{ TeV}$$

AMANDA

$$340 - 2 \times 10^5 \text{ TeV}$$

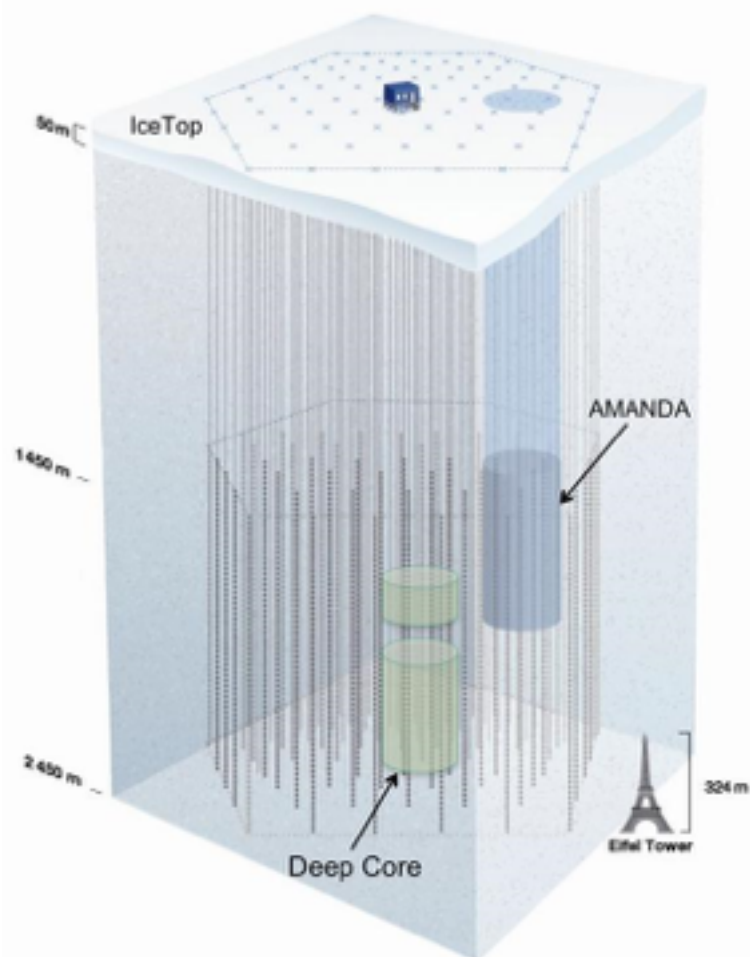
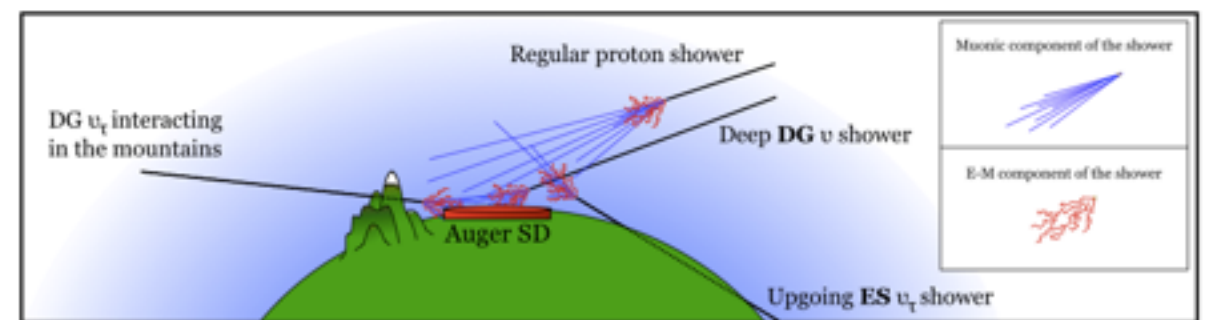
IceCube-22

$$3 \times 10^3 - 6.3 \times 10^6 \text{ TeV}$$

IceCube-40

$$10^5 - 10^8 \text{ TeV}$$

Auger



$$10^6 - 3.2 \times 10^{11} \text{ TeV}$$

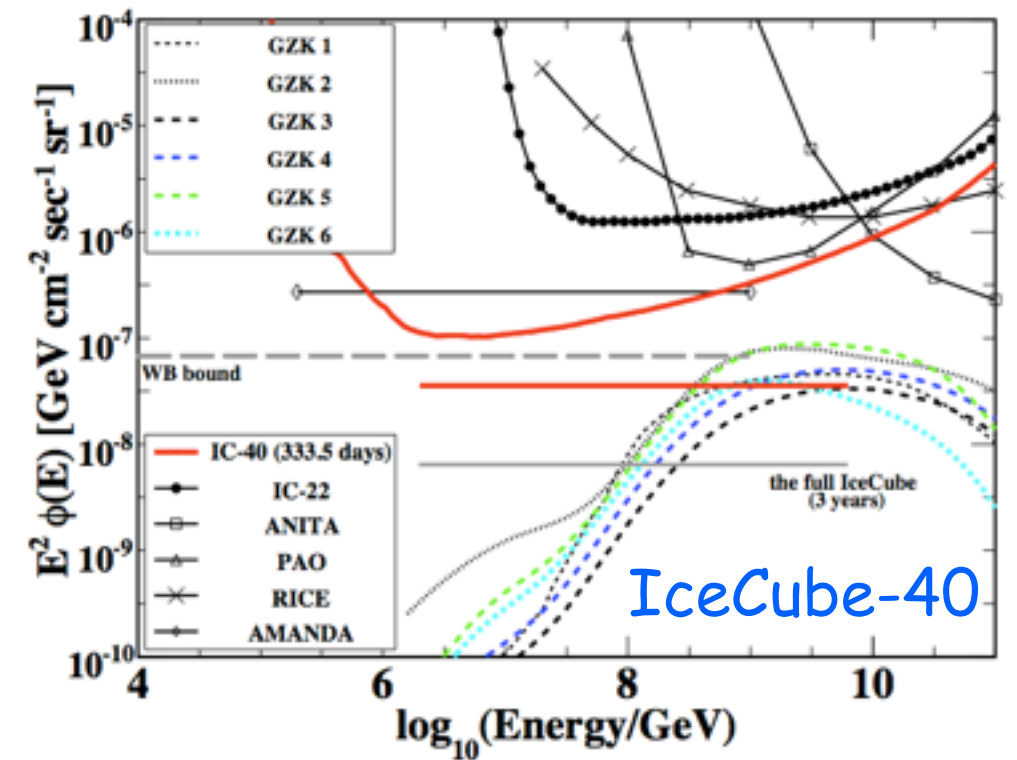
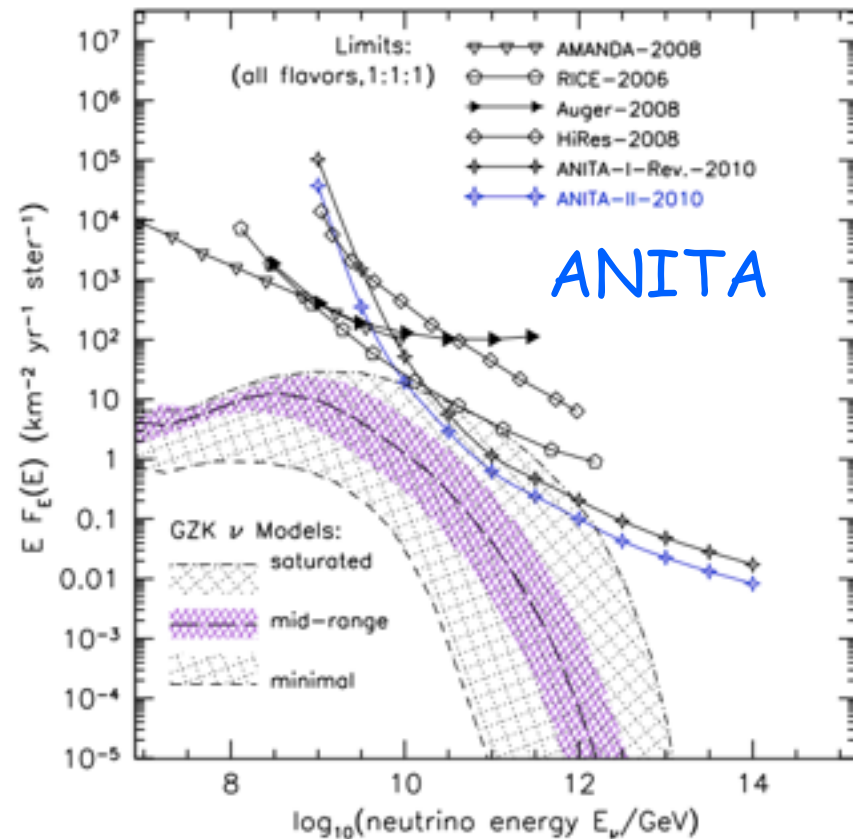
ANITA



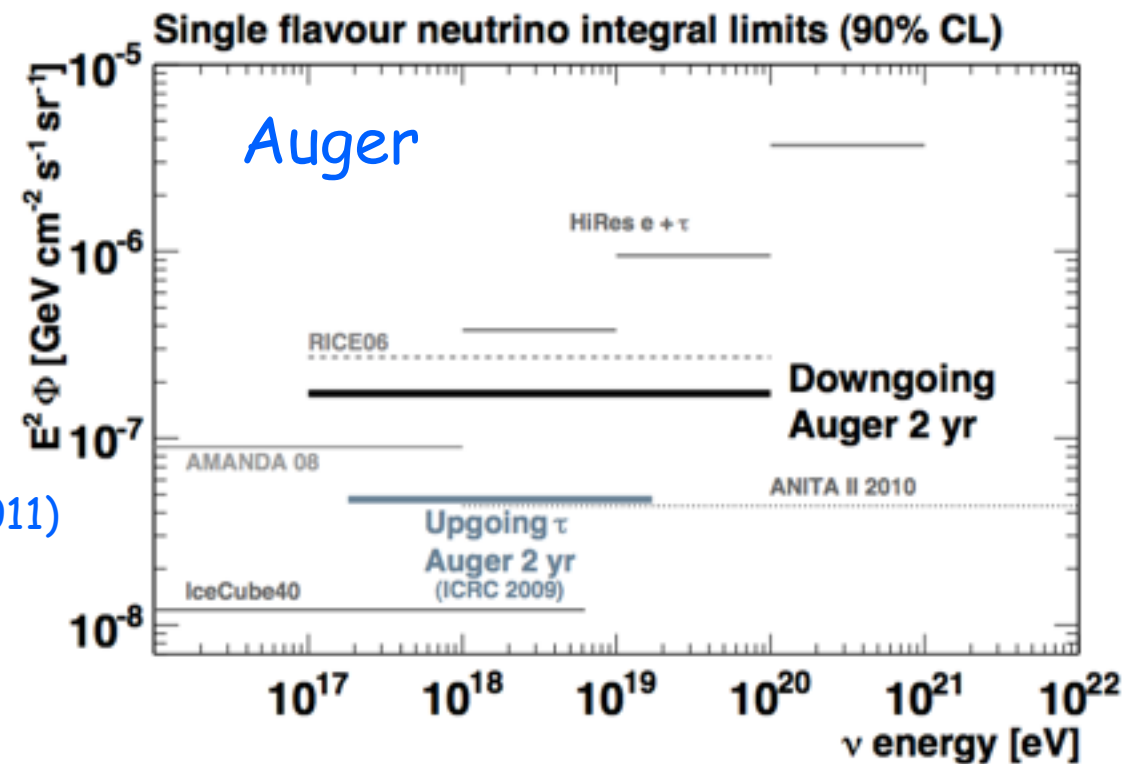
Experiments

R. Abbasi et al. [IceCube Collaboration], Phys.Rev.D83 (2011)
arXiv:1103.4250

P. Gorham et al. [ANITA Collaboration], Phys.Rev.D85 (2012)
arXiv:1011.5004, arXiv:1003.2961



P. Abreu et al. [Pierre Auger Collaboration], Phys.Rev.D84 (2011)
arXiv:1202.1493

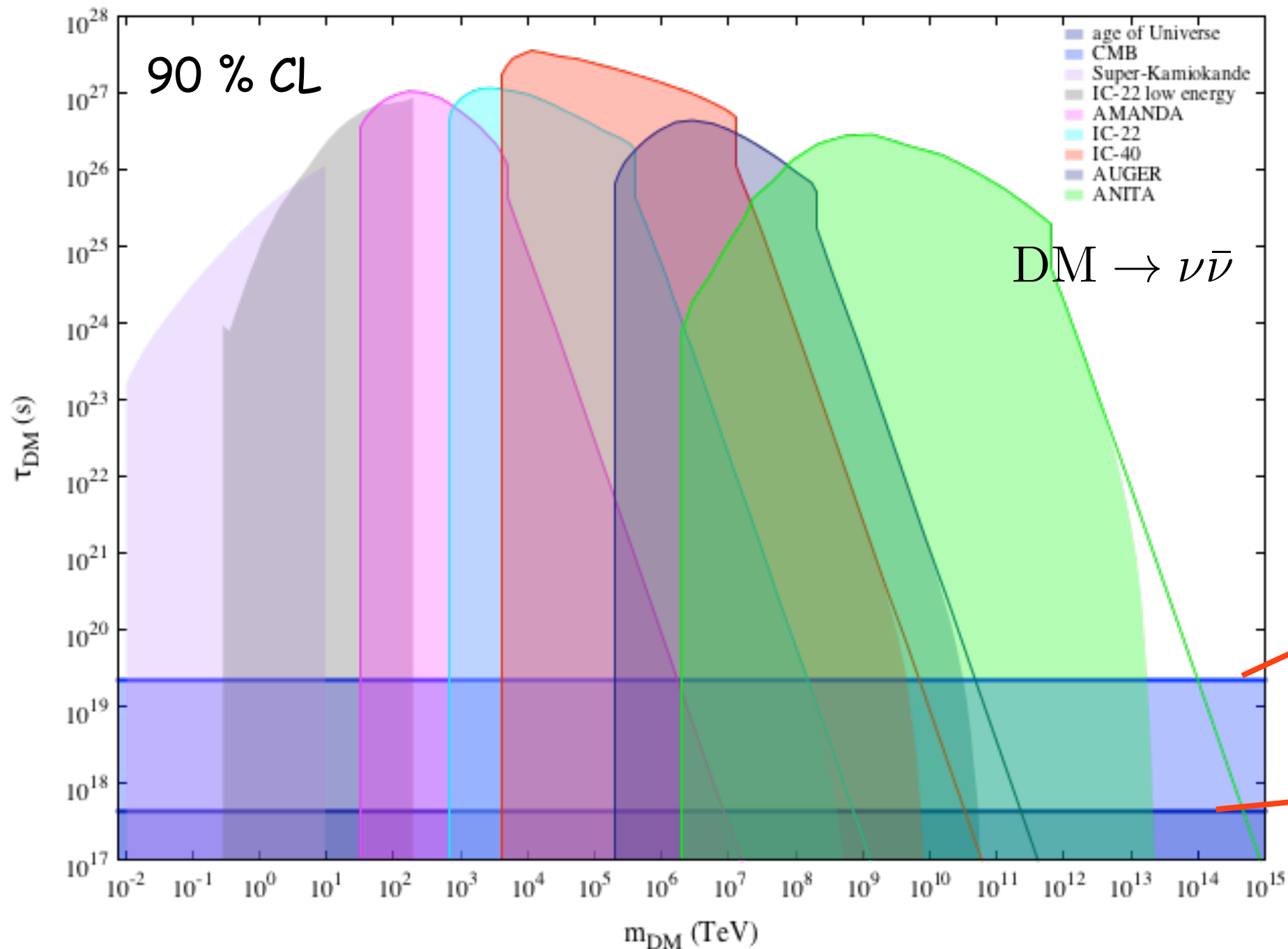


Limits on lifetime from neutrino experiments

before recent IceCube data

✓ Lifetime: stable (∞) or

A.E., Alejandro Ibarra and Orlando L. G. Peres
JCAP (2012) [arXiv: 1205.5281]



Using AMANDA, IceCube,
Auger and ANITA data

$\tau_{\text{DM}} > 2.2 \times 10^{19} \text{ s} \text{ (CMB)}$

$\tau_{\text{DM}} > 4.3 \times 10^{17} \text{ s}$
(age of Universe)

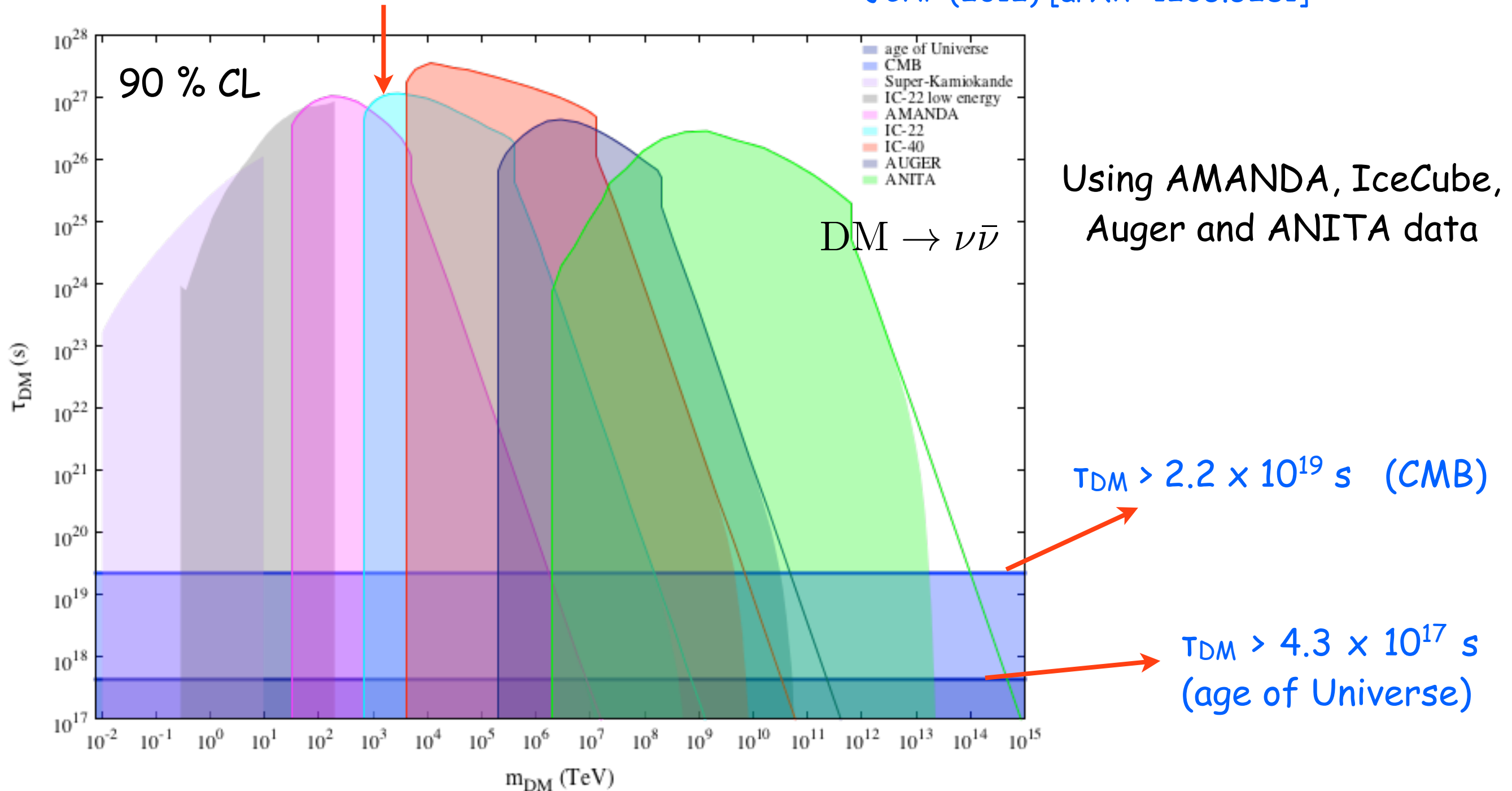
Limits on lifetime from neutrino experiments

before recent IceCube data



Lifetime: stable (∞) or
this talk

A.E., Alejandro Ibarra and Orlando L. G. Peres
JCAP (2012) [arXiv: 1205.5281]



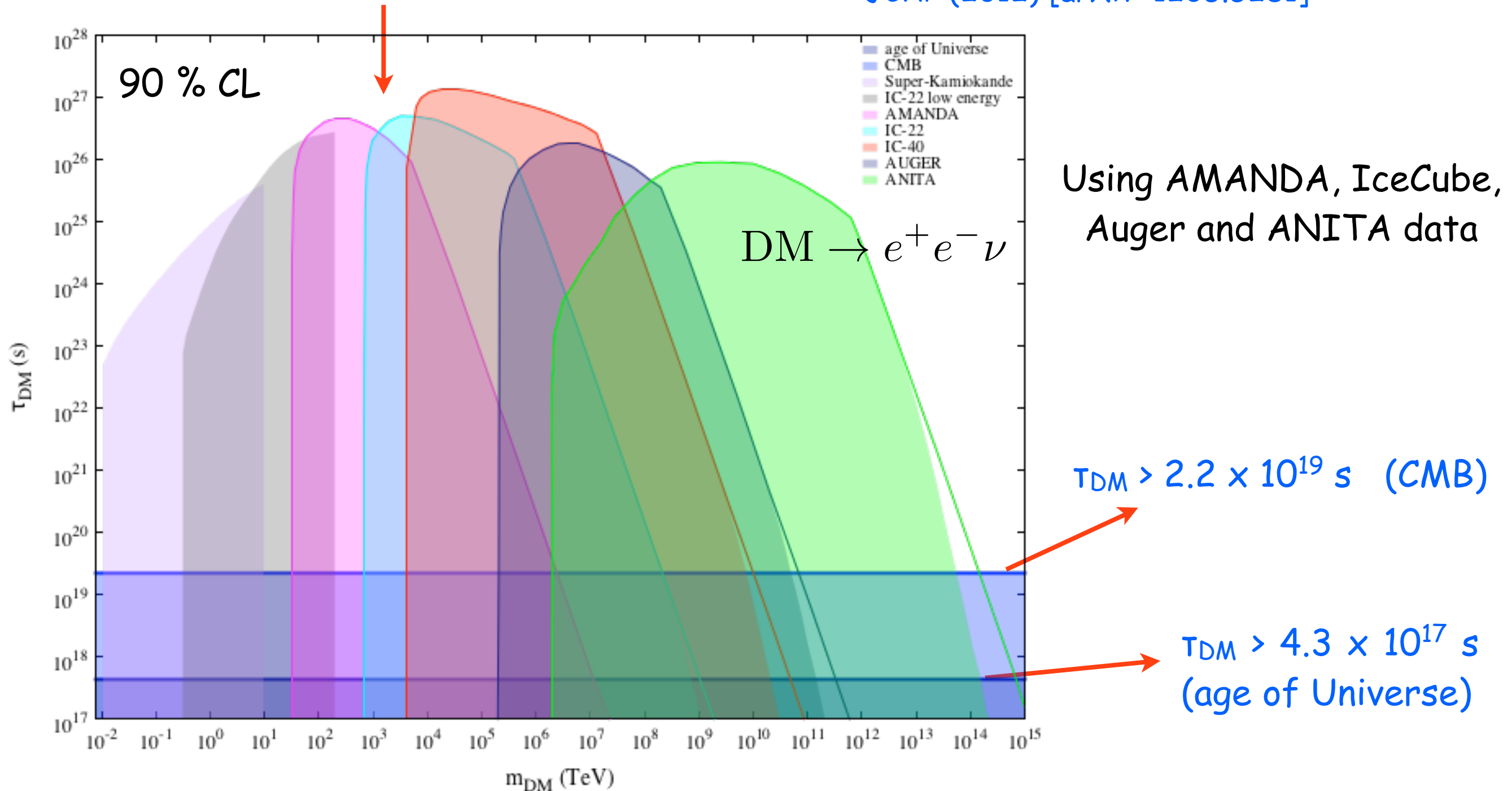
Limits on lifetime from neutrino experiments

before recent IceCube data



Lifetime: stable (∞) or
this talk

A.E., Alejandro Ibarra and Orlando L. G. Peres
JCAP (2012) [arXiv: 1205.5281]



Interpreting the IceCube events by decaying dark matter

Two main diagnostics:

- ✓ Energy distribution
- ✓ Angular distribution

Energy distribution of neutrinos from decaying DM



Galactic contribution:

$$\frac{dJ_h}{dE_\nu}(l, b) = \frac{1}{4\pi m_{\text{DM}} \tau_{\text{DM}}} \frac{dN_\nu}{dE_\nu} \int_0^\infty ds \, \rho_h[r(s, l, b)]$$

NFW

$$\rho_{\text{halo}}(r) \simeq \frac{\rho_h}{r/r_c (1 + r/r_c)^2}$$

$$r(s, l, b) = \sqrt{s^2 + R_\odot^2 - 2sR_\odot \cos b \cos l}$$

Energy distribution of neutrinos from decaying DM



Galactic contribution:

$$\frac{dJ_h}{dE_\nu}(l, b) = \frac{1}{4\pi m_{\text{DM}} \tau_{\text{DM}}} \frac{dN_\nu}{dE_\nu} \int_0^\infty ds \rho_h[r(s, l, b)]$$

NFW $\rho_{\text{halo}}(r) \simeq \frac{\rho_h}{r/r_c(1 + r/r_c)^2}$

$r(s, l, b) = \sqrt{s^2 + R_\odot^2 - 2sR_\odot \cos b \cos l}$



extragalactic contribution:

$$\frac{dJ_{\text{eg}}}{dE_\nu} = \frac{\Omega_{\text{DM}} \rho_c}{4\pi m_{\text{DM}} \tau_{\text{DM}}} \int_0^\infty dz \frac{1}{H(z)} \frac{dN_\nu}{dE_\nu} [(1+z)E_\nu]$$

Energy distribution of neutrinos from decaying DM

✓ Galactic contribution:

$$\frac{dJ_h}{dE_\nu}(l, b) = \frac{1}{4\pi m_{\text{DM}} \tau_{\text{DM}}} \frac{dN_\nu}{dE_\nu} \int_0^\infty ds \rho_h[r(s, l, b)]$$

NFW $\rho_{\text{halo}}(r) \simeq \frac{\rho_h}{r/r_c(1 + r/r_c)^2}$

$r(s, l, b) = \sqrt{s^2 + R_\odot^2 - 2sR_\odot \cos b \cos l}$

✓ extragalactic contribution:

$$\frac{dJ_{\text{eg}}}{dE_\nu} = \frac{\Omega_{\text{DM}} \rho_c}{4\pi m_{\text{DM}} \tau_{\text{DM}}} \int_0^\infty dz \frac{1}{H(z)} \frac{dN_\nu}{dE_\nu} [(1+z)E_\nu]$$

M. Cirelli et. al., JCAP (2011)

energy spectrum of neutrinos
at production point
(including the EW corrections)

quarks

$$\frac{dN_\nu}{dE_\nu} = (1 - b_H) \left. \frac{dN_\nu}{dE_\nu} \right|_S + b_H \left. \frac{dN_\nu}{dE_\nu} \right|_H$$

neutrinos,
charged leptons

Energy distribution of neutrinos from decaying DM

✓ Galactic contribution:

$$\frac{dJ_h}{dE_\nu}(l, b) = \frac{1}{4\pi m_{\text{DM}} \tau_{\text{DM}}} \frac{dN_\nu}{dE_\nu} \int_0^\infty ds \rho_h[r(s, l, b)]$$

NFW $\rho_{\text{halo}}(r) \simeq \frac{\rho_h}{r/r_c(1 + r/r_c)^2}$

$r(s, l, b) = \sqrt{s^2 + R_\odot^2 - 2sR_\odot \cos b \cos l}$

✓ extragalactic contribution:

$$\frac{dJ_{\text{eg}}}{dE_\nu} = \frac{\Omega_{\text{DM}} \rho_c}{4\pi m_{\text{DM}} \tau_{\text{DM}}} \int_0^\infty dz \frac{1}{H(z)} \frac{dN_\nu}{dE_\nu} [(1+z)E_\nu]$$

M. Cirelli et. al., JCAP (2011)

energy spectrum of neutrinos
at production point
(including the EW corrections)

quarks

neutrinos,
charged leptons

$$\frac{dN_\nu}{dE_\nu} = (1 - b_H) \left. \frac{dN_\nu}{dE_\nu} \right|_S + b_H \left. \frac{dN_\nu}{dE_\nu} \right|_H$$

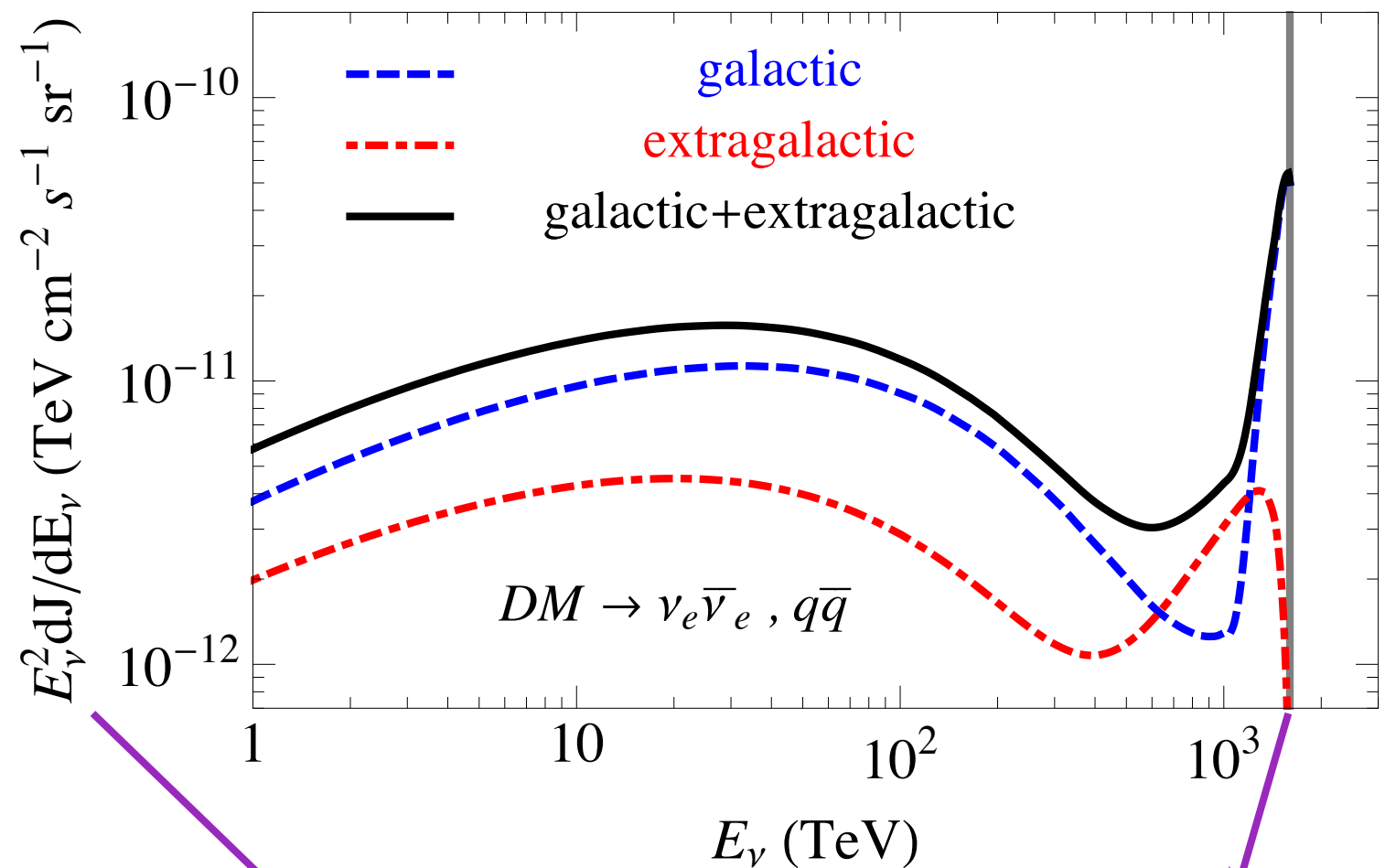
at the Earth $\begin{pmatrix} J_e \\ J_\mu \\ J_\tau \end{pmatrix} = \begin{pmatrix} P_{ee} & P_{e\mu} & P_{e\tau} \\ & P_{\mu\mu} & P_{\mu\tau} \\ & & P_{\tau\tau} \end{pmatrix} \begin{pmatrix} I_e \\ I_\mu \\ I_\tau \end{pmatrix}$ production point

decoherent oscillation

Flux of neutrinos from decaying DM

✓ an example:

A. E., Pasquale D. Serpico,
JCAP (2013) [arXiv:1308.1105]



$$(U_e + U_\mu + U_\tau)/3$$

$$m_{DM}/2 = 1.6 \text{ PeV}$$

$$b_H = 0.12 \text{ and } \tau_{DM} = 2 \times 10^{27} \text{ s}$$

Flux of neutrinos from decaying DM

✓ an example:

intriguing features:

a cut-off at $m_{\text{DM}}/2$

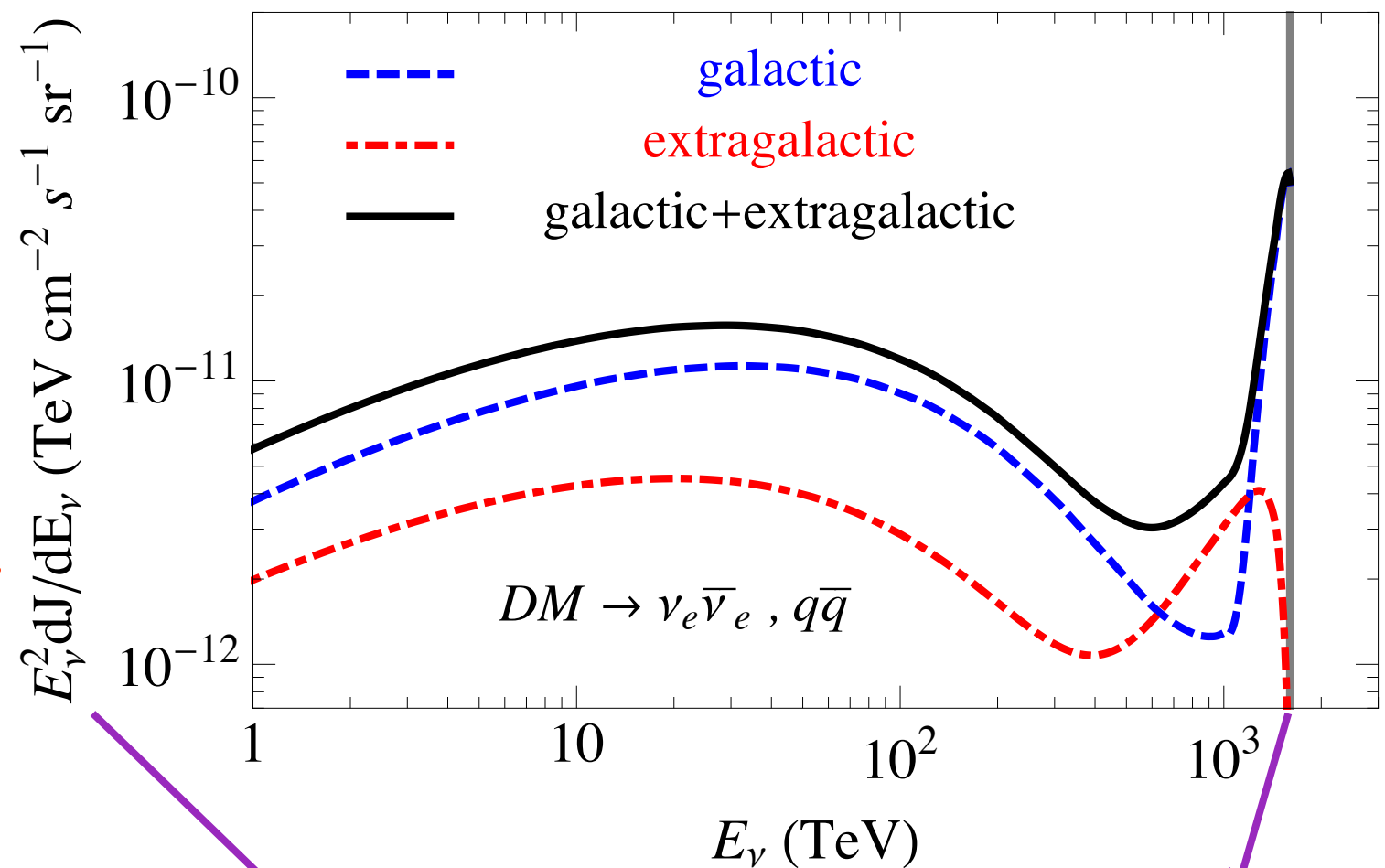
a peak in $\sim \text{PeV}$

a dip in $\sim (0.4-1) \text{ PeV}$

populated spectrum in $< 0.4 \text{ PeV}$

due to soft channel and EW cascades

A. E., Pasquale D. Serpico,
JCAP (2013) [arXiv:1308.1105]



$$(\nu_e + \nu_\mu + \nu_\tau)/3$$

$$m_{\text{DM}}/2 = 1.6 \text{ PeV}$$

$$b_H = 0.12 \text{ and } \tau_{\text{DM}} = 2 \times 10^{27} \text{ s}$$

Flux of neutrinos from decaying DM

✓ an example:

intriguing features:

a cut-off at $m_{\text{DM}}/2$

a peak in $\sim \text{PeV}$

a dip in $\sim (0.4-1) \text{ PeV}$

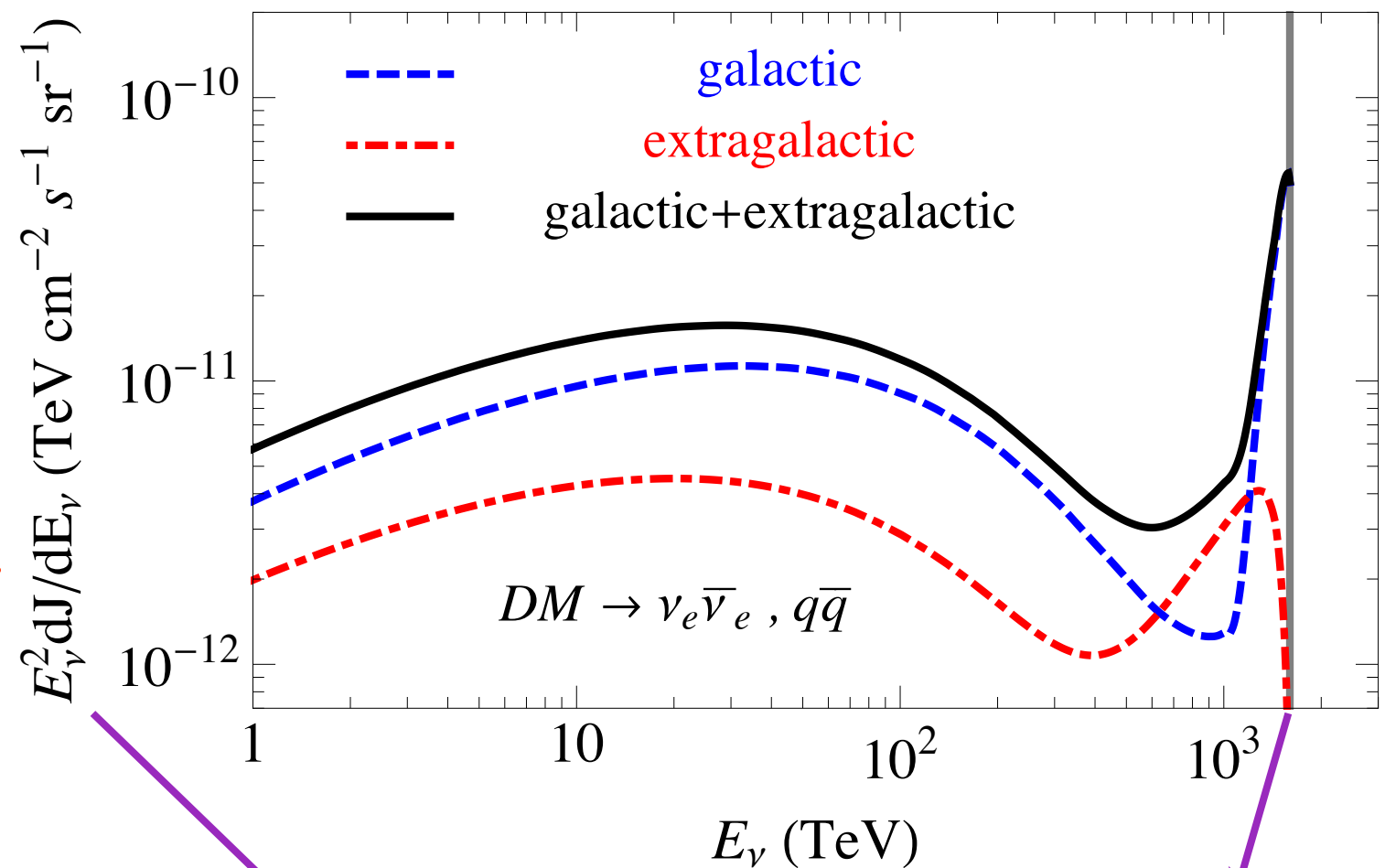
populated spectrum in $< 0.4 \text{ PeV}$

due to soft channel and EW cascades

b_H controls the peak height at $\sim \text{PeV}$

τ_{DM} controls the low energy population

A. E., Pasquale D. Serpico,
JCAP (2013) [arXiv:1308.1105]



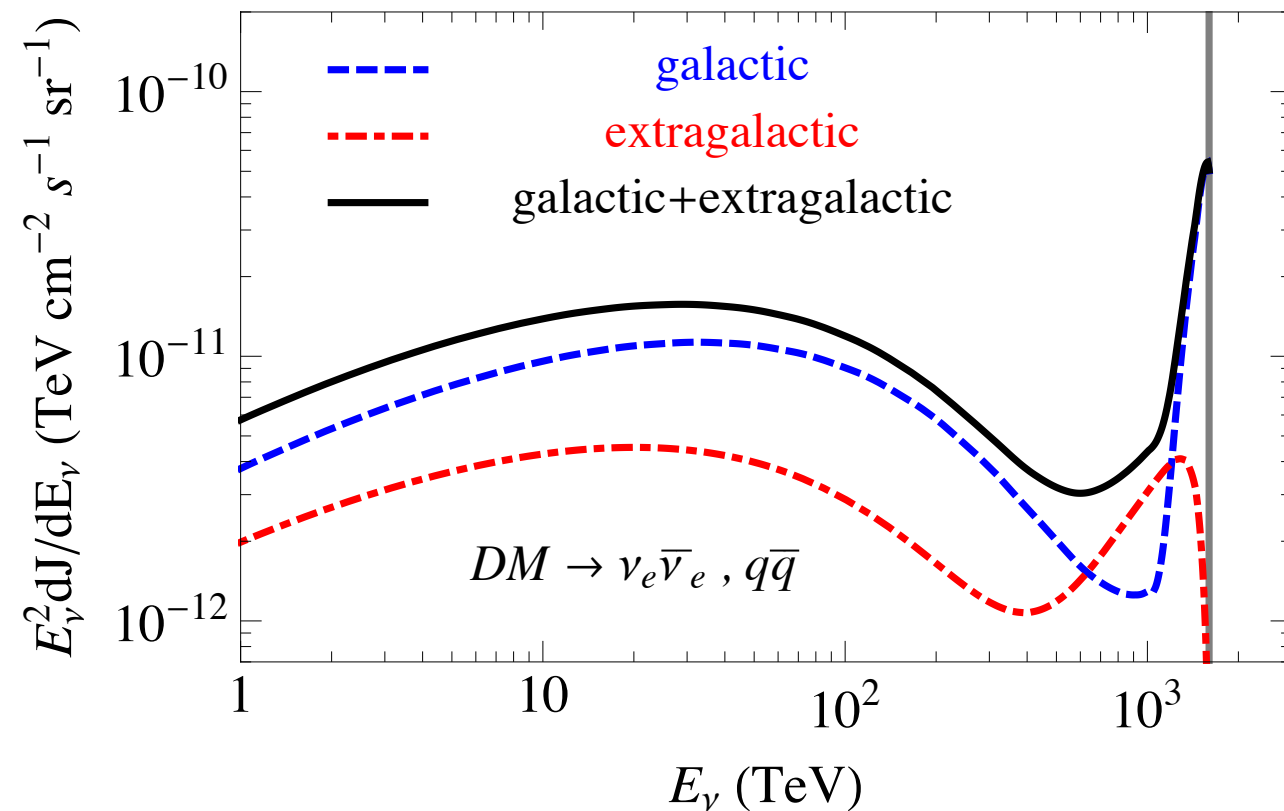
$$(U_e + U_\mu + U_\tau)/3$$

$$m_{\text{DM}}/2 = 1.6 \text{ PeV}$$

$$b_H = 0.12 \text{ and } \tau_{\text{DM}} = 2 \times 10^{27} \text{ s}$$

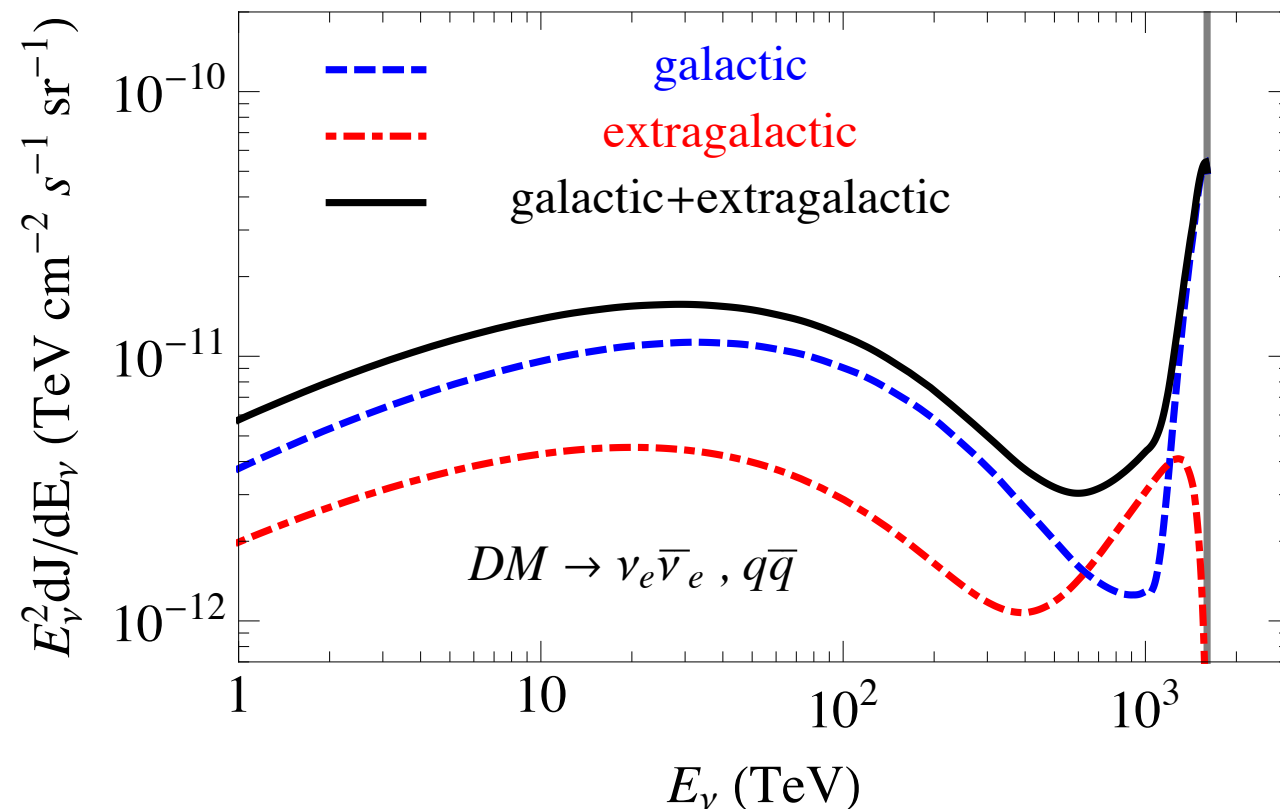
Flux of neutrinos from decaying DM

✓ fine-tuned decay channels ?

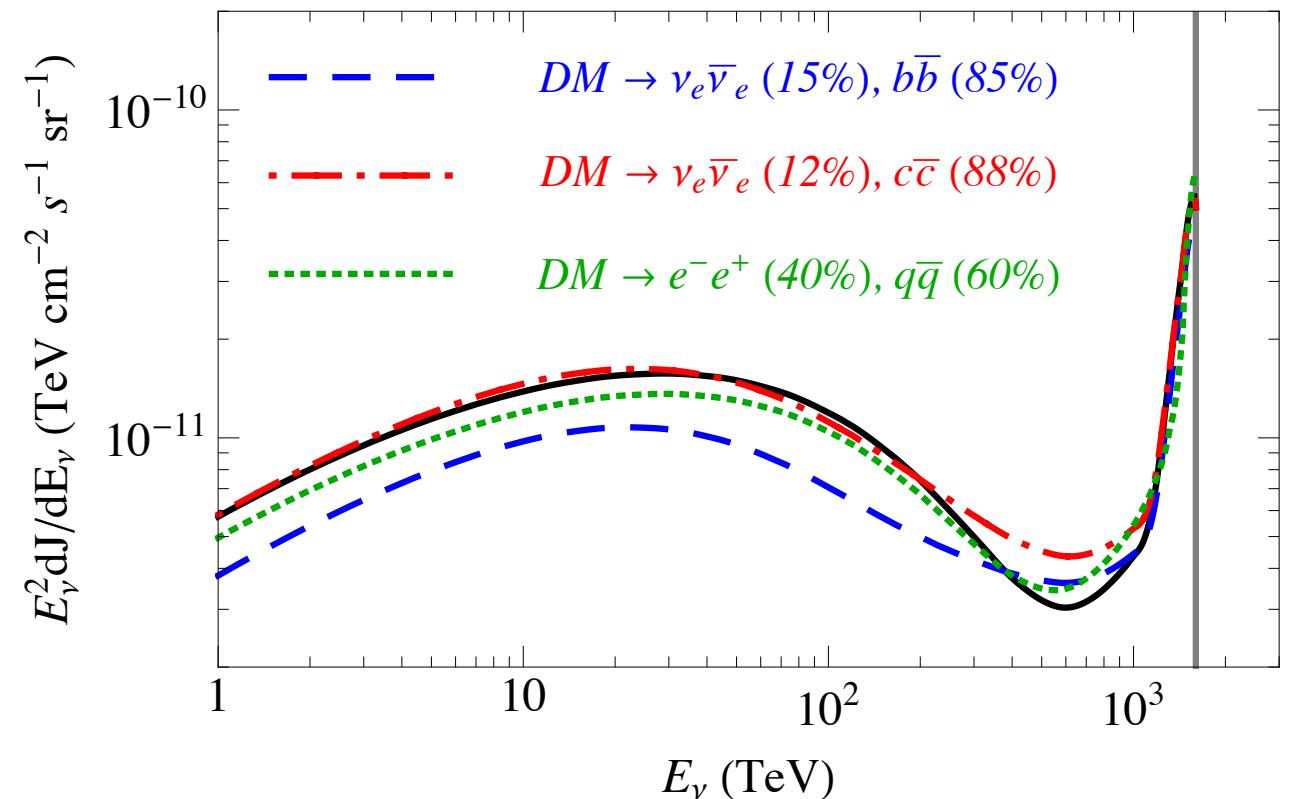


Flux of neutrinos from decaying DM

✓ fine-tuned decay channels ?



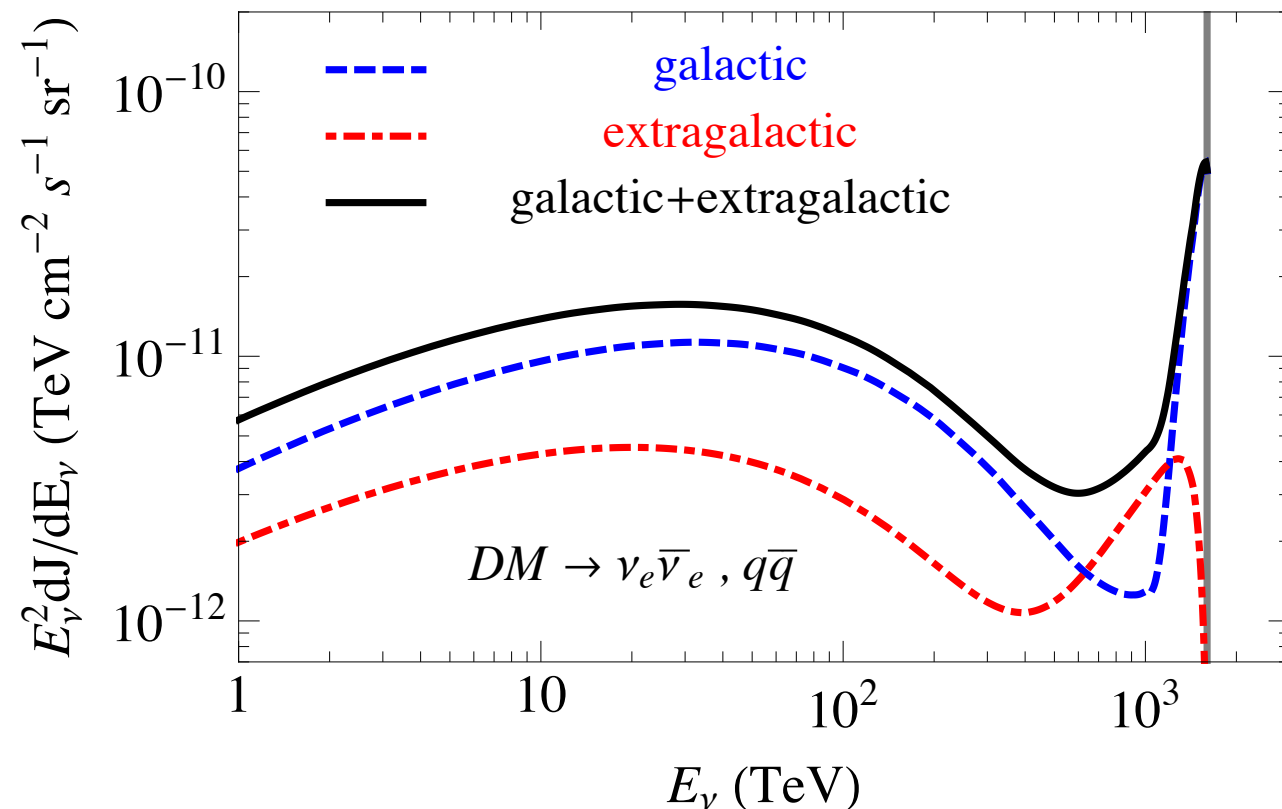
$$\tau_{DM} = (1-3) \times 10^{27} \text{ s}$$



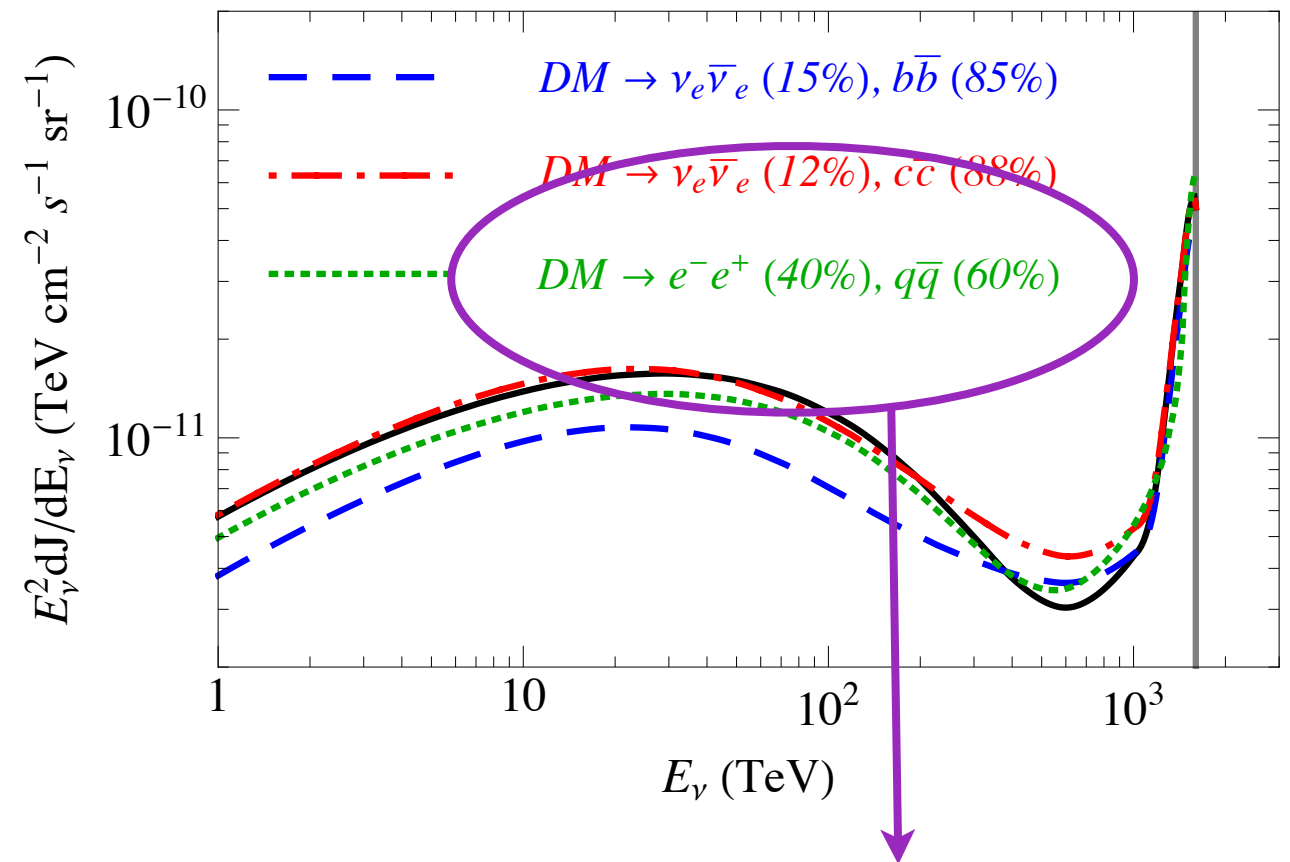
the intriguing features are generic

Flux of neutrinos from decaying DM

✓ fine-tuned decay channels ?



$$\tau_{DM} = (1-3) \times 10^{27} \text{ s}$$



the crucial role of EW cascades

the intriguing features are generic

Confronting with energy distribution of IceCube data

two years data set

✓ branching ratio b_H gives the two PeV events

✓ soft channel and lifetime τ_{DM} gives the upturn in low energy

the value of τ_{DM} is compatible with the bounds derived from neutrinos and gamma rays

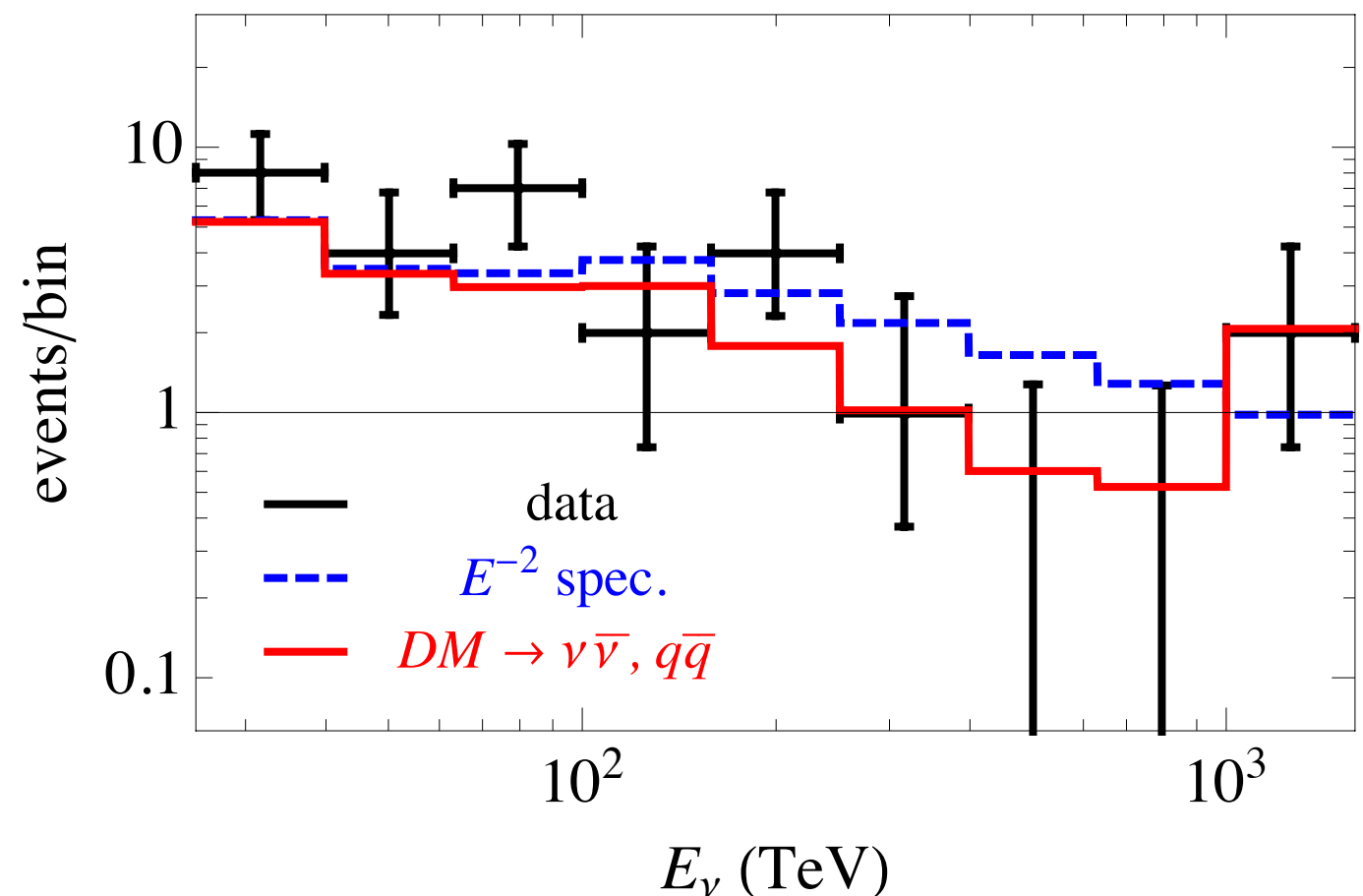
✓ natural explanation for the lack of events $> \text{PeV}$

the value of m_{DM} can be changed within the current uncertainty of the highest energy events

✓ the low energy bins contain large bkg. contribution

the important discriminators of DM vs astrophysical model are high energy bins, where clearly data shows preference to DM model

$$b_H = 0.12 \text{ and } \tau_{DM} = 2 \times 10^{27} \text{ s}$$



✓ different decay channels lead to qualitatively same result

Confronting with energy distribution of IceCube data

three years data set

SM sector  Dark sector

portal type:

$$\mathcal{L}_{\text{portal}} = \frac{\mathcal{O}_{\text{SM}} \mathcal{O}_{\text{DM}}}{\Lambda^{d-4}}$$

Confronting with energy distribution of IceCube data

three years data set

SM sector  Dark sector

portal type:

$$\mathcal{L}_{\text{portal}} = \frac{\mathcal{O}_{\text{SM}} \mathcal{O}_{\text{DM}}}{\Lambda^{d-4}}$$

“neutrino” portal:

$$\mathcal{O}_{\text{SM}} \rightarrow H L$$

A. Falkowski, J. Juknevich and J. Shelton
[arXiv:0908.1790](https://arxiv.org/abs/0908.1790)

Confronting with energy distribution of IceCube data

three years data set

SM sector $\longleftrightarrow$ Dark sector

portal type:

$$\mathcal{L}_{\text{portal}} = \frac{\mathcal{O}_{\text{SM}} \mathcal{O}_{\text{DM}}}{\Lambda^{d-4}}$$

“neutrino” portal:

$$\mathcal{O}_{\text{SM}} \rightarrow H L$$

A. Falkowski, J. Juknevich and J. Shelton
arXiv:0908.1790



$d = 4 :$

$$\mathcal{O}_{\text{DM}} \rightarrow N$$

heavy sterile neutrino, DM candidate

T. Higaki, R. Kitano and R. Sato, JHEP (2014)
arXiv:1405.0013

Confronting with energy distribution of IceCube data

three years data set

SM sector $\longleftrightarrow$ Dark sector

portal type:

$$\mathcal{L}_{\text{portal}} = \frac{\mathcal{O}_{\text{SM}} \mathcal{O}_{\text{DM}}}{\Lambda^{d-4}}$$

“neutrino” portal:

$$\mathcal{O}_{\text{SM}} \rightarrow H L$$

A. Falkowski, J. Juknevich and J. Shelton
arXiv:0908.1790



$d = 4 :$

$$\mathcal{O}_{\text{DM}} \rightarrow N$$

heavy sterile neutrino, DM candidate

T. Higaki, R. Kitano and R. Sato, JHEP (2014)
arXiv:1405.0013

UV completion:

$$SU(3)_c \times SU(2)_L \times U(1)_Y \times U(1)_{B-L}$$

“Higgs” field ϕ_{B-L} plays the role of inflaton

$$m_\phi \sim 10^{13} \text{ GeV}$$

$$T_R \sim 10^7 \text{ GeV}$$

Confronting with energy distribution of IceCube data

three years data set

Leptogenesis: $\phi \rightarrow N_2 N_2$ $M_2 \sim 10^{12} \text{ GeV} \xrightarrow{\text{green arrow}} \frac{n_B}{s} \sim 10^{-10}$

DM abundance: $\Omega_{N_1} \simeq 0.2 \left(\frac{M_1}{4 \text{ PeV}} \right)^3 \left(\frac{T_R}{3 \times 10^7 \text{ GeV}} \right)^{-1}$

DM lifetime: $\tau_{N_1} \simeq 8 \times 10^{28} \text{ s} \left(\frac{M_1}{1 \text{ PeV}} \right)^{-1} \left(\frac{10^{-29}}{|y_N|^2} \right)$

DM decay channels: $\text{Br}(\ell^\pm W^\mp) = 2\text{Br}(\nu_\ell Z) = 2\text{Br}(\nu_\ell h) = |U_{\ell 1}|^2$ **NH**

$\text{Br}(\ell^\pm W^\mp) = 2\text{Br}(\nu_\ell Z) = 2\text{Br}(\nu_\ell h) = |U_{\ell 3}|^2$ **IH**

Confronting with energy distribution of IceCube data

three years data set

IH

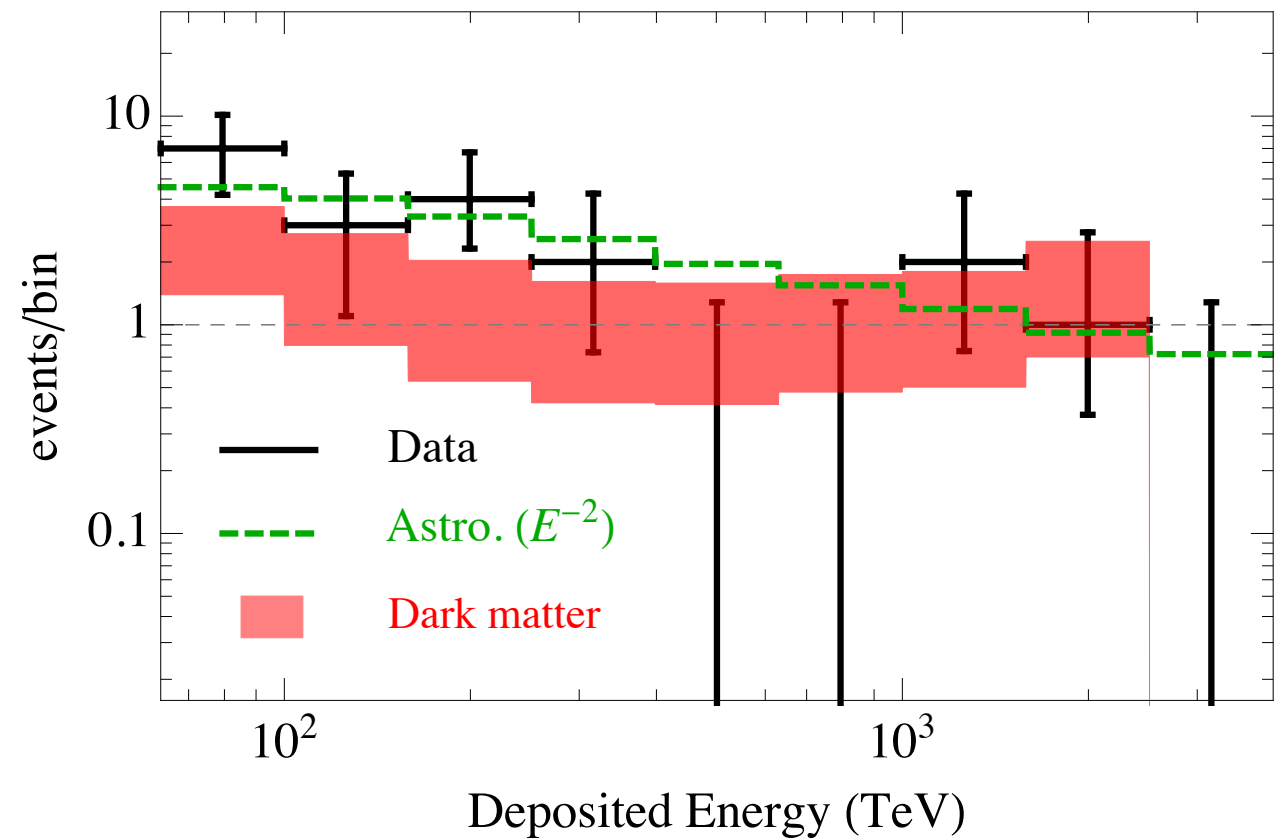
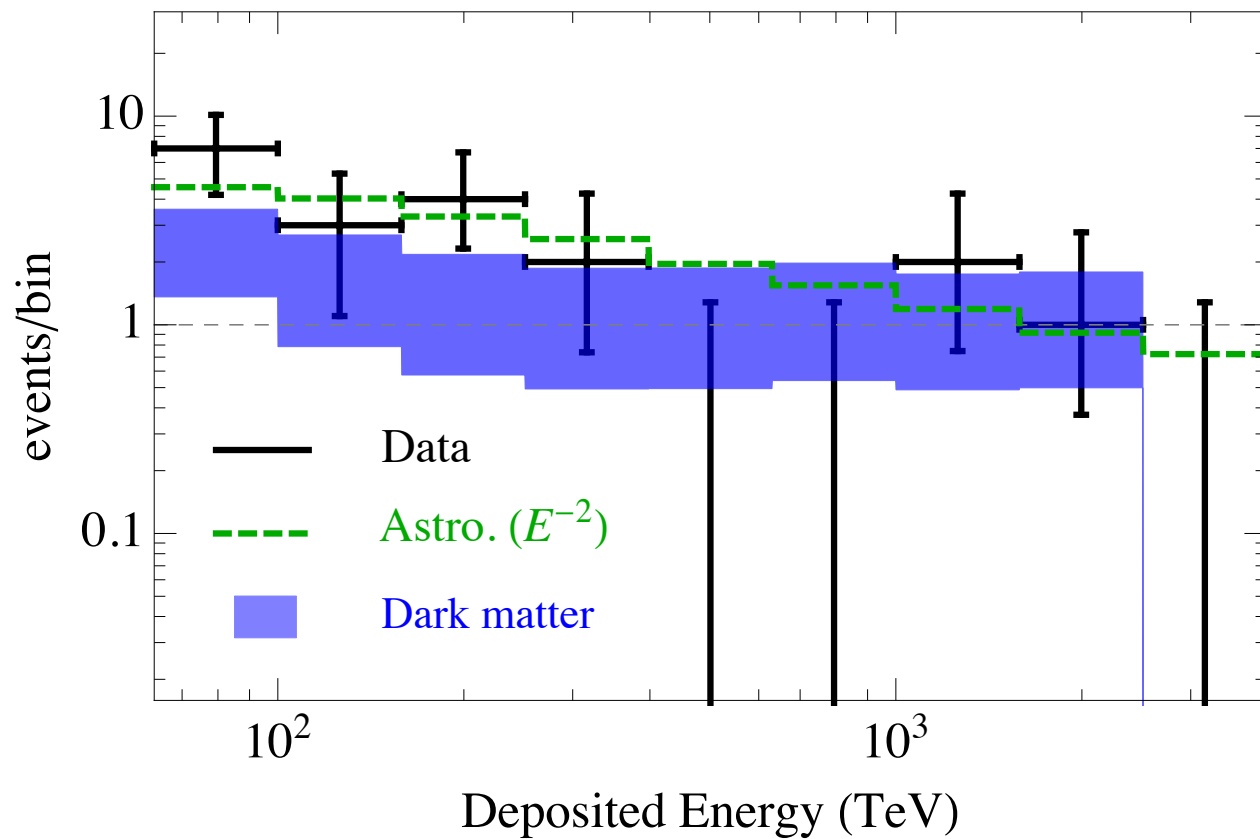
$$\tau_{\text{DM}} = 1.1 \times 10^{28} \text{ s}$$

shaded: $\pm 1\sigma$

NH

$$\tau_{\text{DM}} = 7.3 \times 10^{27} \text{ s}$$

shaded: $\pm 1\sigma$



$$m_{\text{DM}} = 4 \text{ PeV}$$

Confronting with energy distribution of IceCube data

three years data set

SM sector $\longleftrightarrow$ Dark sector

portal type:

$$\mathcal{L}_{\text{portal}} = \frac{\mathcal{O}_{\text{SM}} \mathcal{O}_{\text{DM}}}{\Lambda^{d-4}}$$

“neutrino” portal:

$$\mathcal{O}_{\text{SM}} \rightarrow H L$$

A. Falkowski, J. Juknevich and J. Shelton
arXiv:0908.1790 [hep-ph].

✓ $d = 4:$ $\mathcal{O}_{\text{DM}} \rightarrow N$

production mechanism:

$$m_\phi \gg m_N$$

inflaton decay

$$m_\phi \ll m_N$$

freeze-in

$$g\phi NN, \quad g \simeq 10^{-6}$$

Confronting with energy distribution of IceCube data

three years data set

SM sector $\longleftrightarrow$ Dark sector

portal type:

$$\mathcal{L}_{\text{portal}} = \frac{\mathcal{O}_{\text{SM}} \mathcal{O}_{\text{DM}}}{\Lambda^{d-4}}$$

“neutrino” portal:

$$\mathcal{O}_{\text{SM}} \rightarrow H L$$

A. Falkowski, J. Juknevich and J. Shelton
arXiv:0908.1790 [hep-ph].

✓ $d = 5$: $\mathcal{O}_{\text{DM}} \rightarrow \chi \phi$ singlet fermion and scalar
(Asymmetric DM)

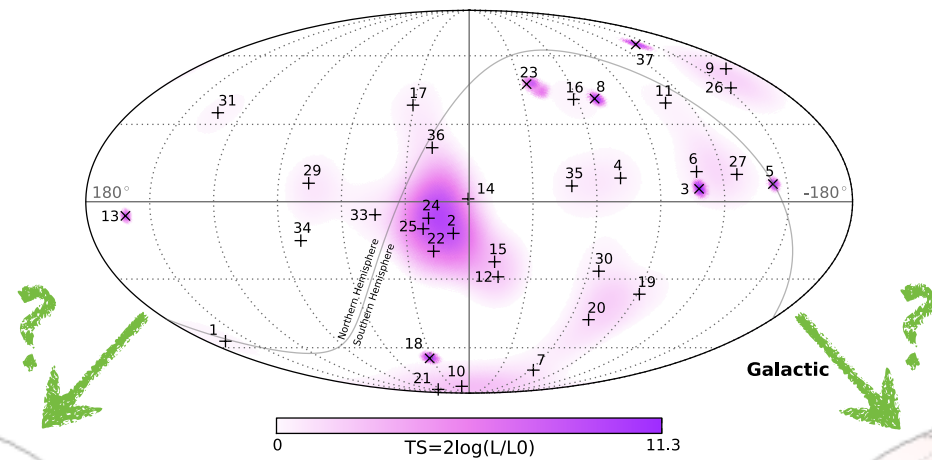
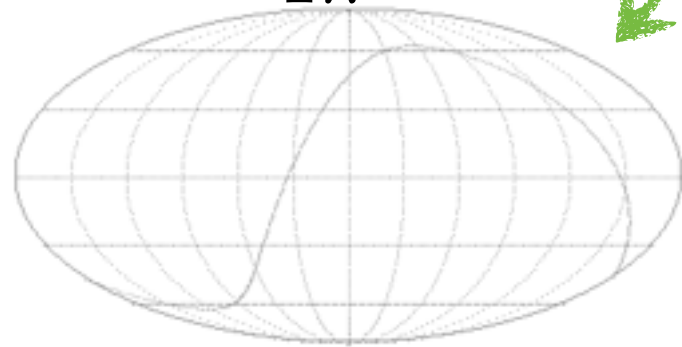
✓ $d = 6$: other portals

✓ For $d > 4$ there are more freedom in branching ratios. We have shown that for the most constrained model ($d=4$) a good fit to the data can be obtained. Obviously better fits can be achieved for $d > 4$.

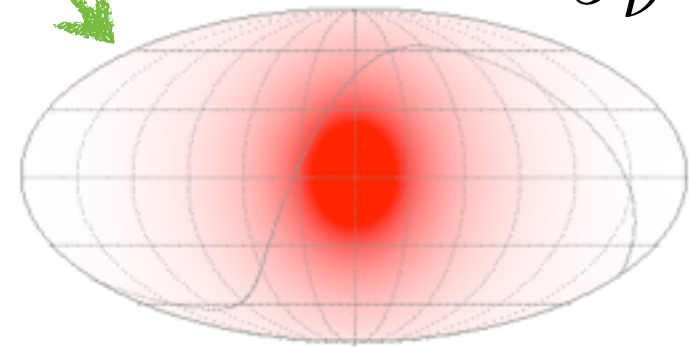
Angular distribution of neutrinos from decaying DM

✓ We would compare

$$p^{\text{iso}} = \frac{1}{4\pi}$$



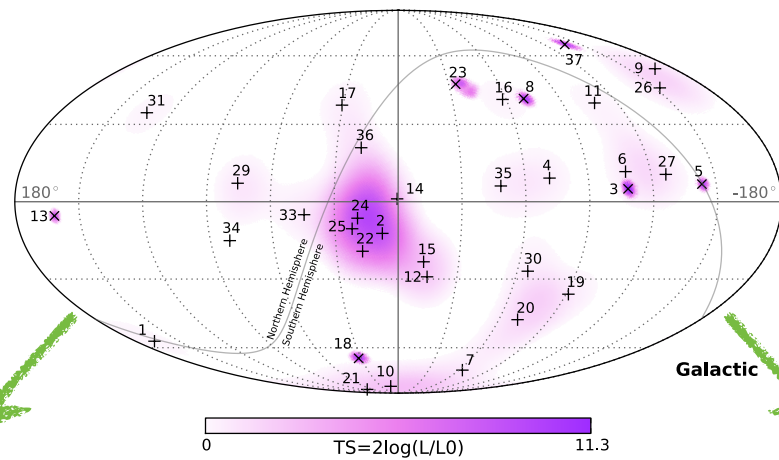
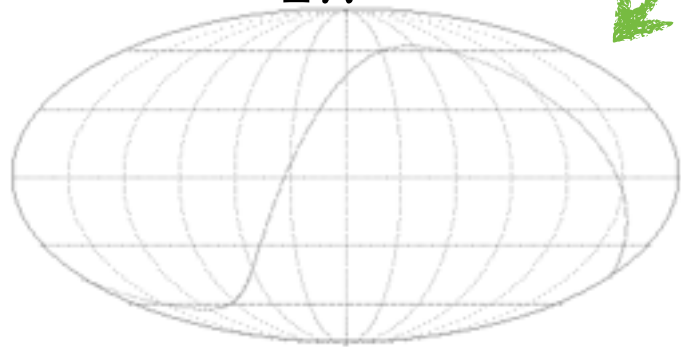
$$p^{\text{DM}} = \frac{1}{J_\nu} \frac{d^2 J_\nu}{db dl}$$



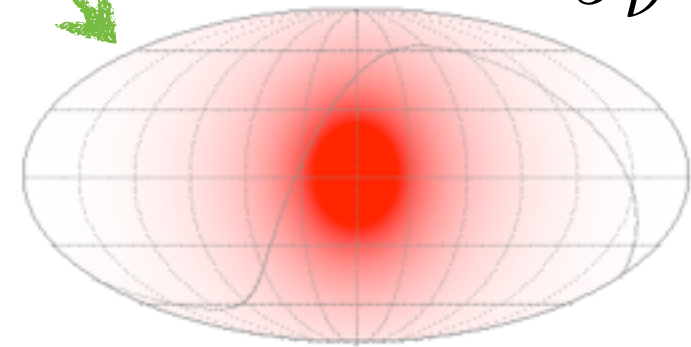
Angular distribution of neutrinos from decaying DM

✓ We would compare

$$p^{\text{iso}} = \frac{1}{4\pi}$$



$$p^{\text{DM}} = \frac{1}{J_\nu} \frac{d^2 J_\nu}{db dl}$$



PDF of data

$$p_i(b, l) = \frac{1}{2\pi\sigma_i^2} \exp \left[-\frac{|\vec{x} - \vec{x}_i|^2}{2\sigma_i^2} \right]$$

"flat sky"
approximation

PDF of
isotropic dis.

$$p^{\text{iso}} = \frac{1}{4\pi}$$

PDF of DM

$$p^{\text{DM}}(b, l) = \frac{1}{J_\nu} \frac{d^2 J_\nu}{db dl} = \frac{\int_0^\infty \rho[r(s, b, l)] ds + \Omega_{\text{DM}} \rho_c \beta}{4\pi(\eta + \Omega_{\text{DM}} \rho_c \beta)}$$

Angular distribution of neutrinos from decaying DM

✓ Likelihood analysis

Number of **signal** events

Test
Statistics

$$\text{TS}_{\text{like}} = 2 \sum_{i=1}^N (\ln f_i - \ln p_i^{\text{iso}}) = 2 \ln \left(\prod_{i=1}^N f_i \right) - 2N \ln \left(\frac{1}{4\pi} \right)$$

$$f_i = \int p_i(b, l) p^{\text{DM}}(b, l) \cos(b) \, db \, dl = \frac{1}{2\pi\sigma_i^2} \int e^{-\frac{|\vec{x}_i - \vec{x}|^2}{2\sigma_i^2}} p^{\text{DM}}(b, l) \cos(b) \, db \, dl$$

$N = 35$? too optimistic!

Angular distribution of neutrinos from decaying DM

✓ Likelihood analysis

Number of **signal** events

Test
Statistics

$$\text{TS}_{\text{like}} = 2 \sum_{i=1}^N (\ln f_i - \ln p_i^{\text{iso}}) = 2 \ln \left(\prod_{i=1}^N f_i \right) - 2N \ln \left(\frac{1}{4\pi} \right)$$

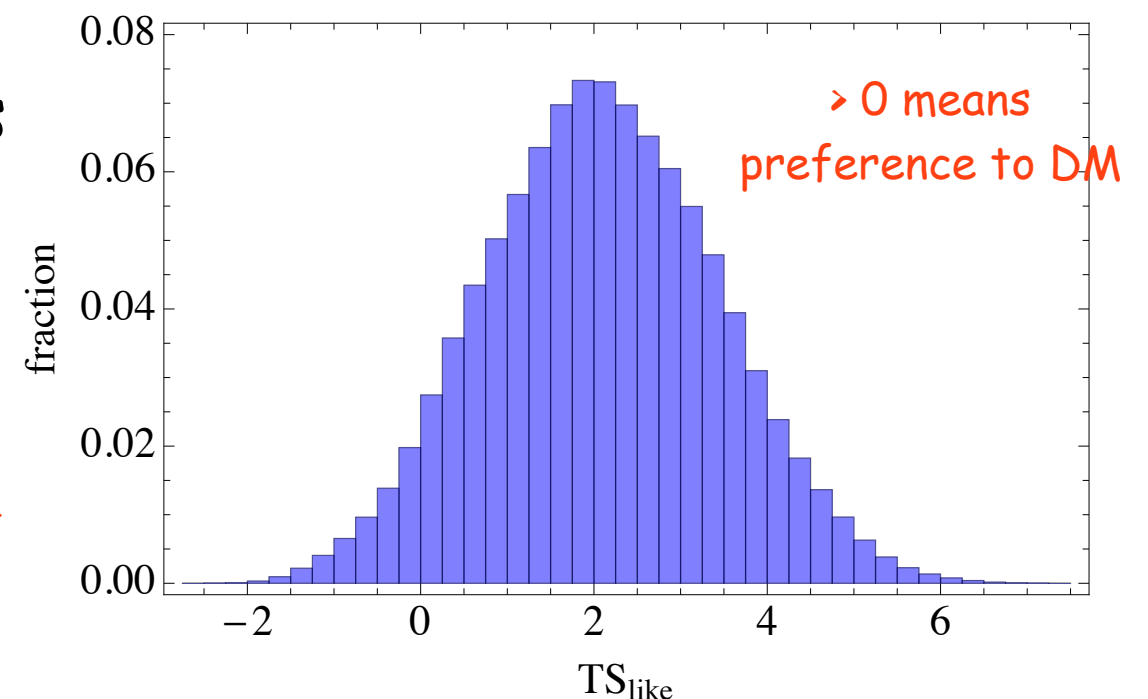
$$f_i = \int p_i(b, l) p^{\text{DM}}(b, l) \cos(b) \, db \, dl = \frac{1}{2\pi\sigma_i^2} \int e^{-\frac{|\vec{x}_i - \vec{x}|^2}{2\sigma_i^2}} p^{\text{DM}}(b, l) \cos(b) \, db \, dl$$

$N = 35$? too optimistic!

✓ let's assume $N_b = 15$ and all the events with $E > 150$ TeV are signal events

→ $\binom{26}{15}$ ways of selecting the bkg events among the low energy events

Distribution of TS_{like} for all these realizations
(mean value = 2.1)



Angular distribution of neutrinos from decaying DM

✓ Likelihood analysis

Number of **signal** events

Test
Statistics

$$\text{TS}_{\text{like}} = 2 \sum_{i=1}^N (\ln f_i - \ln p_i^{\text{iso}}) = 2 \ln \left(\prod_{i=1}^N f_i \right) - 2N \ln \left(\frac{1}{4\pi} \right)$$

$$f_i = \int p_i(b, l) p^{\text{DM}}(b, l) \cos(b) \, db \, dl = \frac{1}{2\pi\sigma_i^2} \int e^{-\frac{|\vec{x}_i - \vec{x}|^2}{2\sigma_i^2}} p^{\text{DM}}(b, l) \cos(b) \, db \, dl$$

$N = 35$? too optimistic!

✓ let's assume $N_b = 15$ and all the events with $E > 150$ TeV are signal events

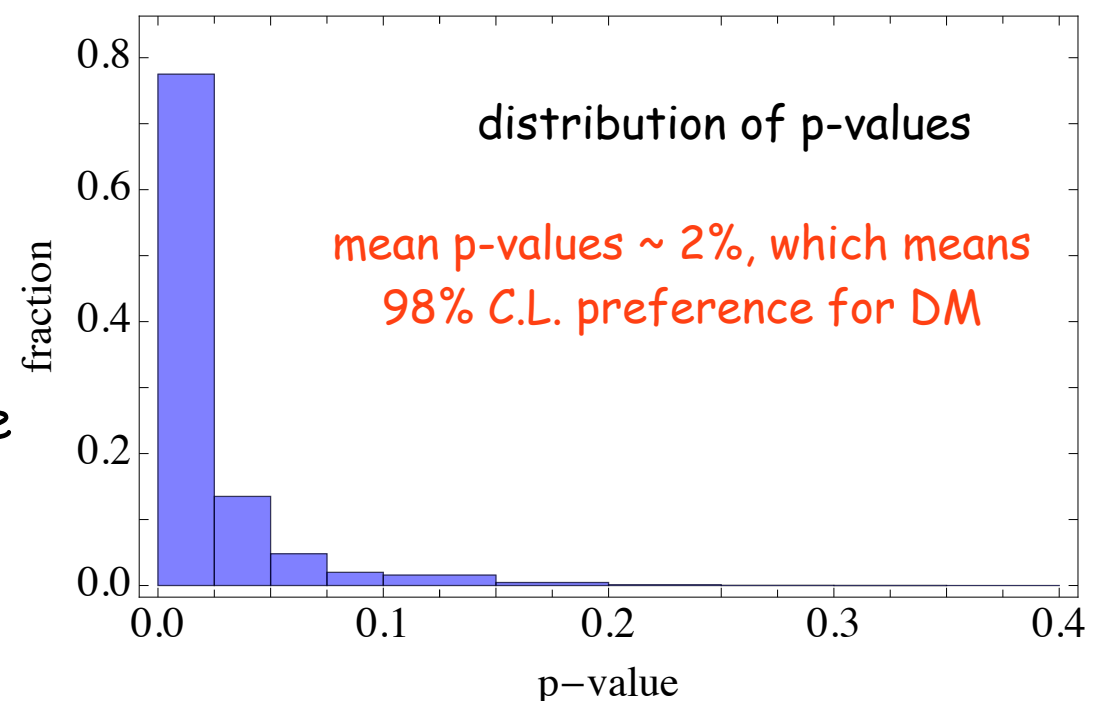
Quantifying the preference

generating a sample (10^5) of isotropically distributed set of 20 events



for each realization of bkg choosing, p-value is the fraction of generated events which have smaller TS_{like} than the one computed by observed data

p-value →



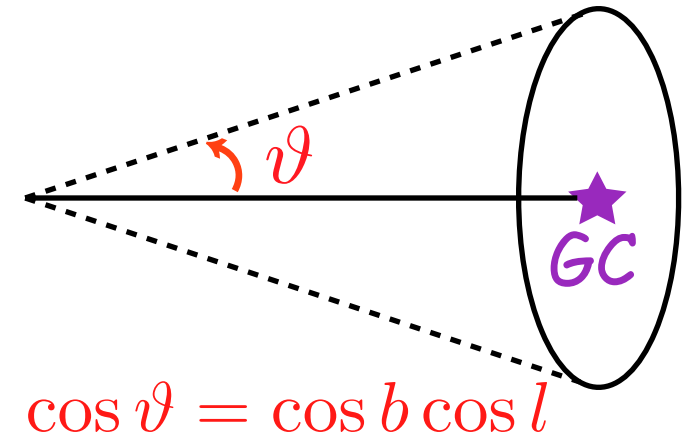
Angular distribution of neutrinos from decaying DM

✓ Kolmogorov-Smirnov test: a powerful non-parametric test

The 2-dim KS test have some ambiguities

$$p^{\text{iso}}(\vartheta) = \int_0^{2\pi} p^{\text{iso}}(\vartheta, \varphi) d\varphi = \int_0^{2\pi} \frac{1}{4\pi} d\varphi = \frac{1}{2}$$

$$p^{\text{DM}}(\vartheta) = \int_0^{2\pi} p^{\text{DM}}(\vartheta, \varphi) d\varphi = \frac{\int_0^\infty \rho[r(s, \vartheta)] ds + \Omega_{\text{DM}} \rho_c \beta}{2(\eta + \Omega_{\text{DM}} \rho_c \beta)}$$



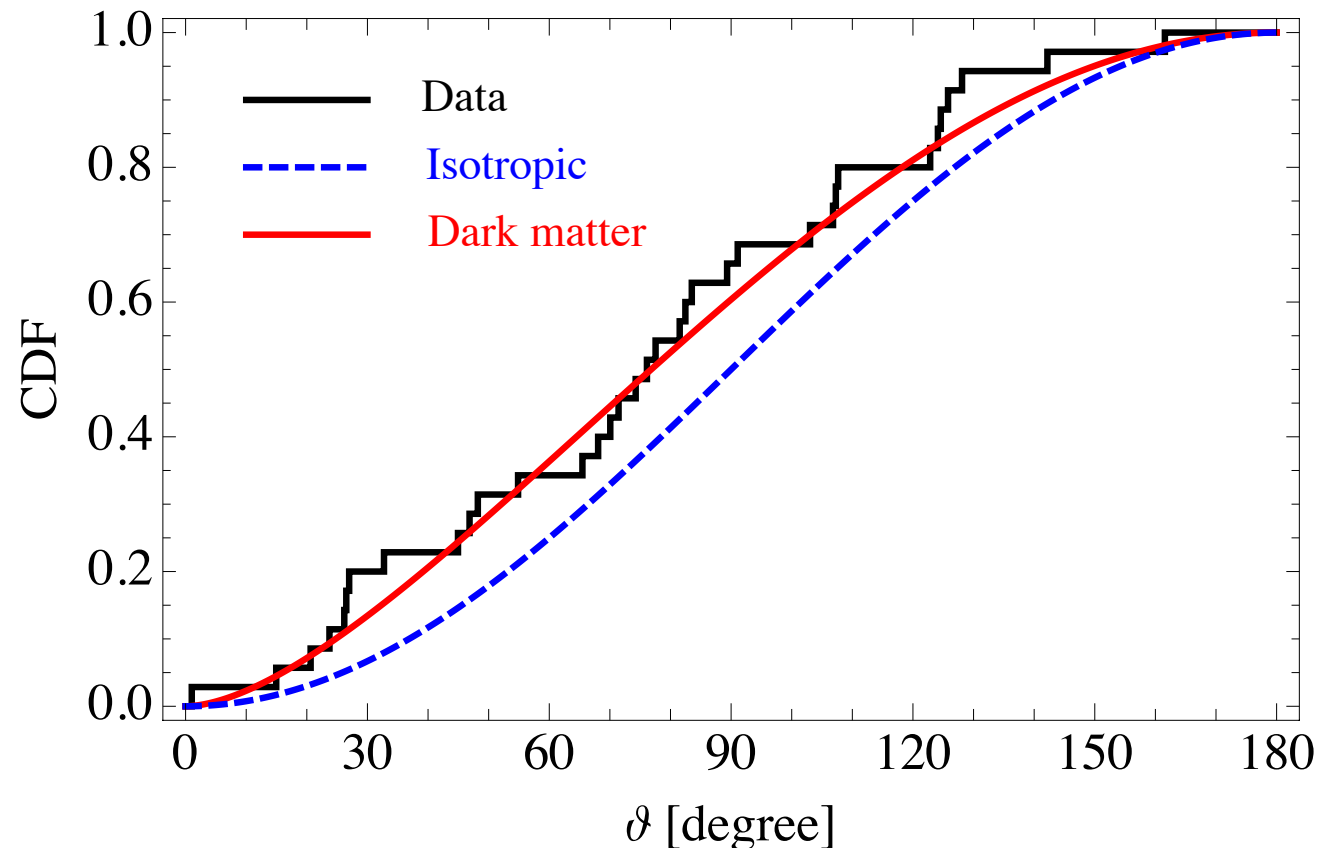
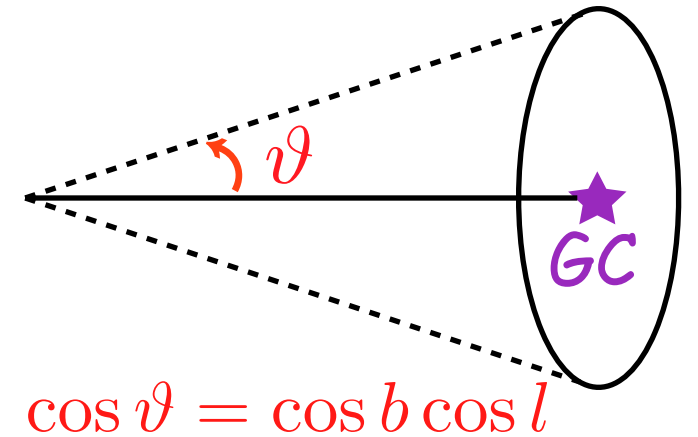
Angular distribution of neutrinos from decaying DM

✓ Kolmogorov-Smirnov test: a powerful non-parametric test

The 2-dim KS test have some ambiguities

$$p^{\text{iso}}(\vartheta) = \int_0^{2\pi} p^{\text{iso}}(\vartheta, \varphi) d\varphi = \int_0^{2\pi} \frac{1}{4\pi} d\varphi = \frac{1}{2}$$

$$p^{\text{DM}}(\vartheta) = \int_0^{2\pi} p^{\text{DM}}(\vartheta, \varphi) d\varphi = \frac{\int_0^\infty \rho[r(s, \vartheta)] ds + \Omega_{\text{DM}} \rho_c \beta}{2(\eta + \Omega_{\text{DM}} \rho_c \beta)}$$



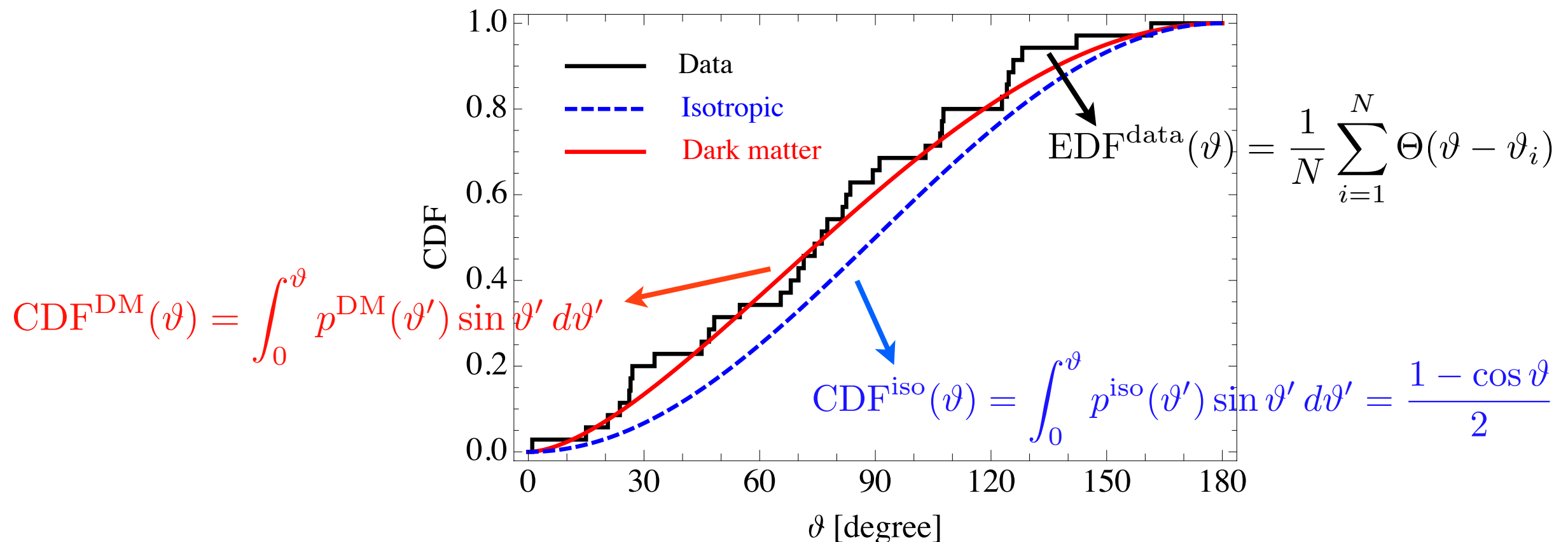
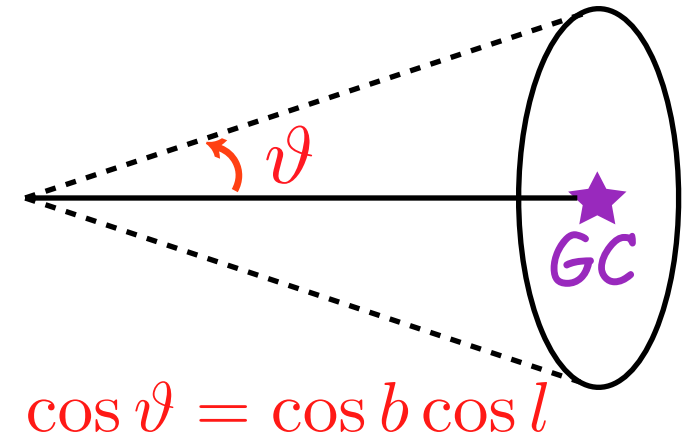
Angular distribution of neutrinos from decaying DM

✓ Kolmogorov-Smirnov test: a powerful non-parametric test

The 2-dim KS test have some ambiguities

$$p^{\text{iso}}(\vartheta) = \int_0^{2\pi} p^{\text{iso}}(\vartheta, \varphi) d\varphi = \int_0^{2\pi} \frac{1}{4\pi} d\varphi = \frac{1}{2}$$

$$p^{\text{DM}}(\vartheta) = \int_0^{2\pi} p^{\text{DM}}(\vartheta, \varphi) d\varphi = \frac{\int_0^\infty \rho[r(s, \vartheta)] ds + \Omega_{\text{DM}} \rho_c \beta}{2(\eta + \Omega_{\text{DM}} \rho_c \beta)}$$



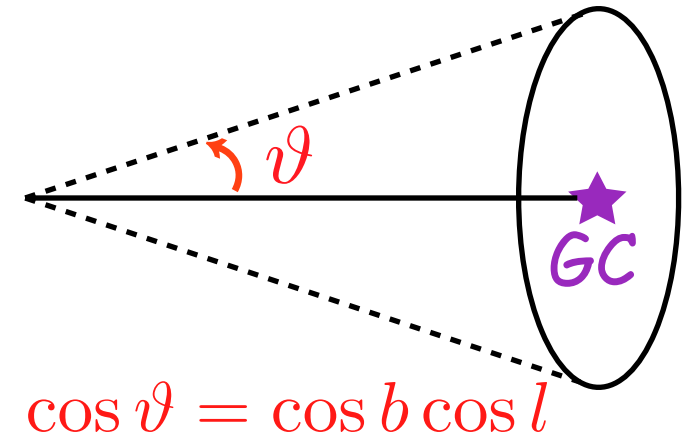
Angular distribution of neutrinos from decaying DM

✓ Kolmogorov-Smirnov test: a powerful non-parametric test

The 2-dim KS test have some ambiguities

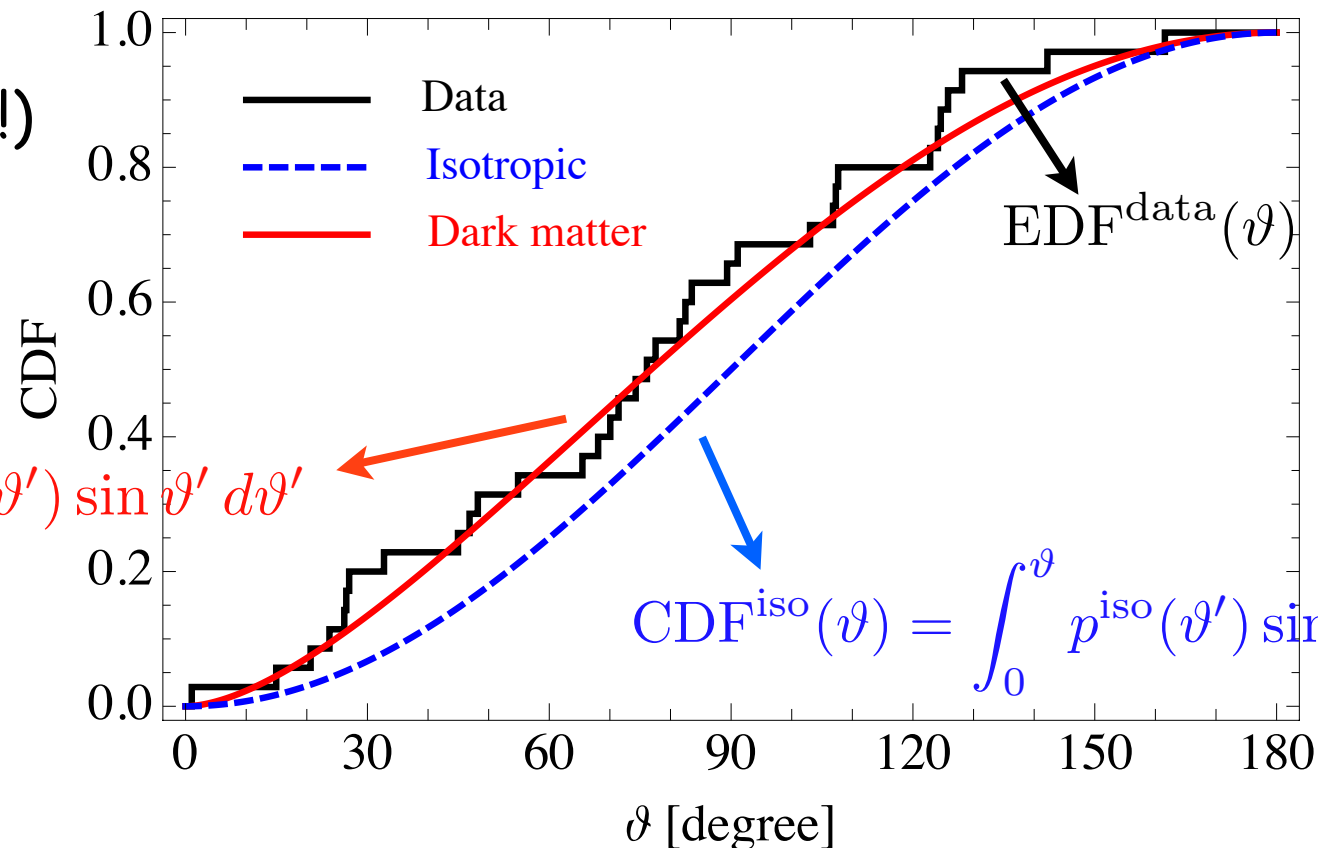
$$p^{\text{iso}}(\vartheta) = \int_0^{2\pi} p^{\text{iso}}(\vartheta, \varphi) d\varphi = \int_0^{2\pi} \frac{1}{4\pi} d\varphi = \frac{1}{2}$$

$$p^{\text{DM}}(\vartheta) = \int_0^{2\pi} p^{\text{DM}}(\vartheta, \varphi) d\varphi = \frac{\int_0^\infty \rho[r(s, \vartheta)] ds + \Omega_{\text{DM}} \rho_c \beta}{2(\eta + \Omega_{\text{DM}} \rho_c \beta)}$$



N = 35 (too optimistic!)

$$\text{CDF}^{\text{DM}}(\vartheta) = \int_0^\vartheta p^{\text{DM}}(\vartheta') \sin \vartheta' d\vartheta'$$

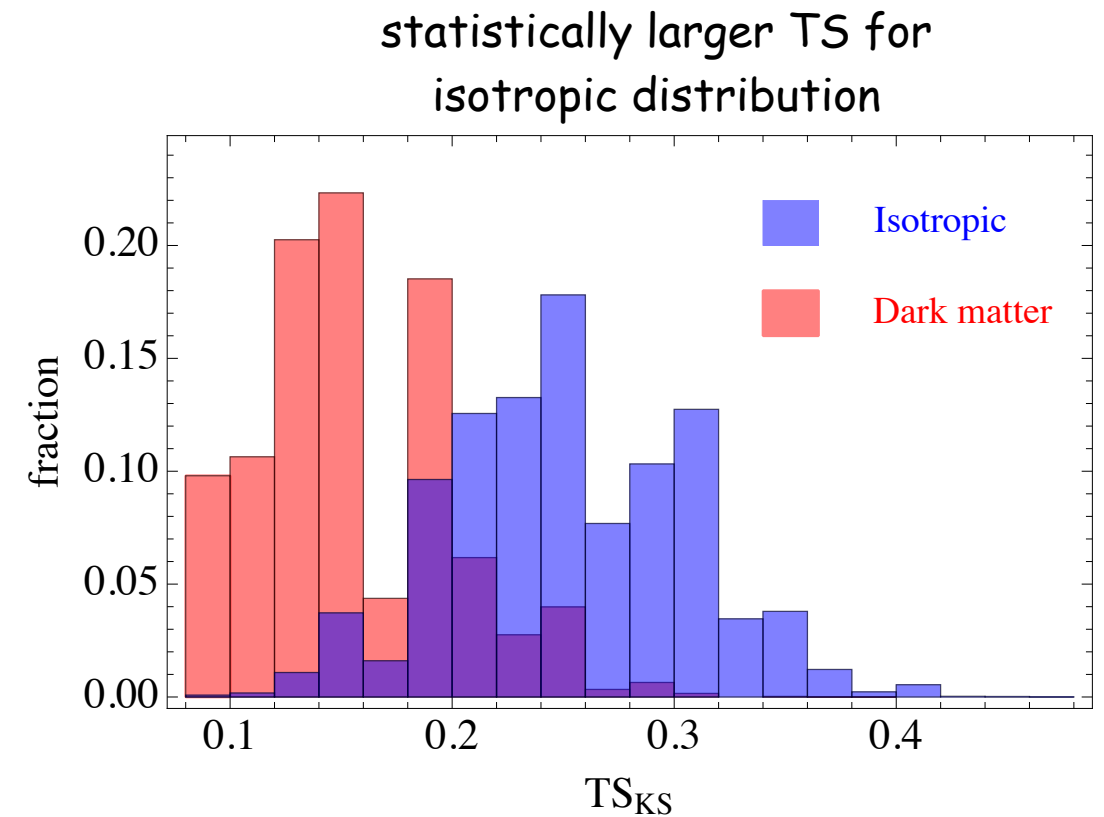


Angular distribution of neutrinos from decaying DM

✓ Kolmogorov-Smirnov test:

Test Statistics

$$TS_{KS} = \max_{1 \leq i \leq N} \left\{ CDF^{DM}(\vartheta_i) - \frac{i-1}{N}, \frac{i}{N} - CDF^{DM}(\vartheta_i) \right\}$$

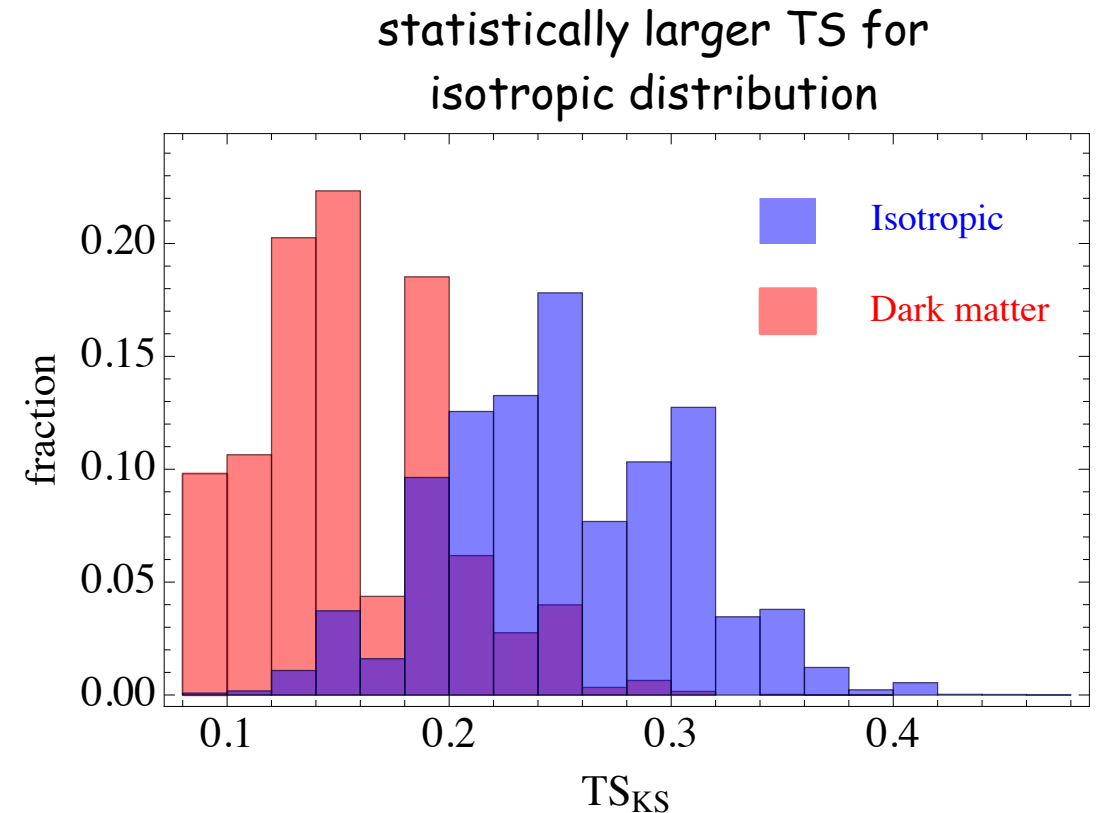


Angular distribution of neutrinos from decaying DM

✓ Kolmogorov-Smirnov test:

Test Statistics

$$TS_{KS} = \max_{1 \leq i \leq N} \left\{ CDF^{DM}(\vartheta_i) - \frac{i-1}{N}, \frac{i}{N} - CDF^{DM}(\vartheta_i) \right\}$$



again, generating a sample (10^5) of isotropically distributed set of 20 events

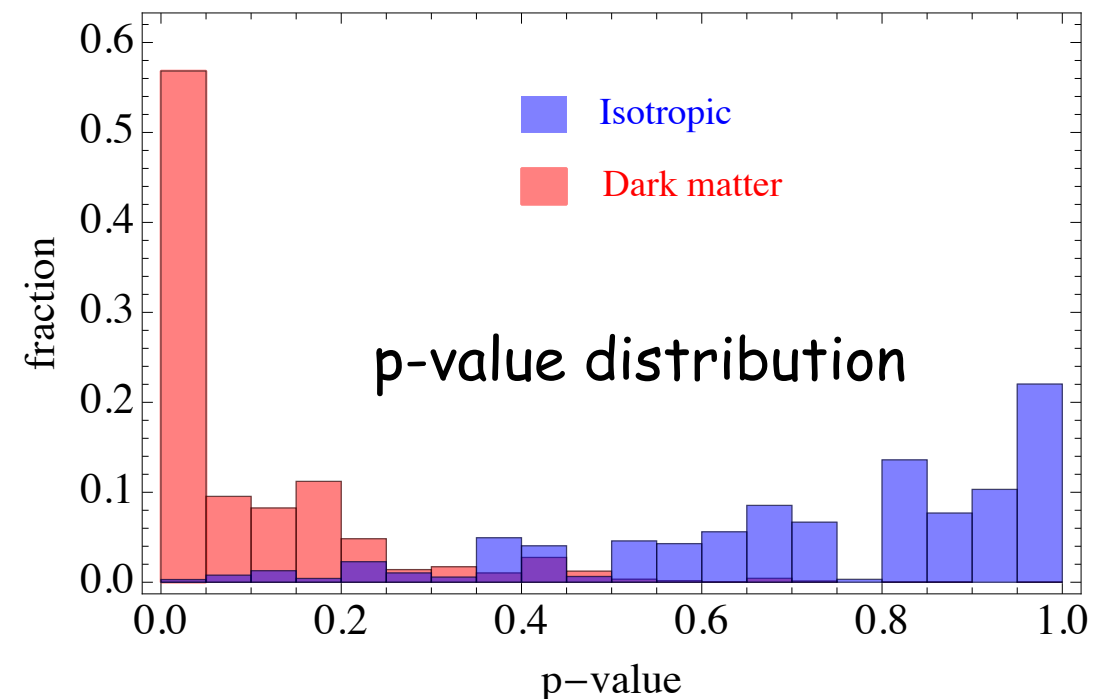


on the average, 10% of generated isotropic sample have smaller TS_{KS} than the values obtained for data vs DM dis.

for data vs isotropic dis. it is 73%



less than 2 σ preference for DM dis.



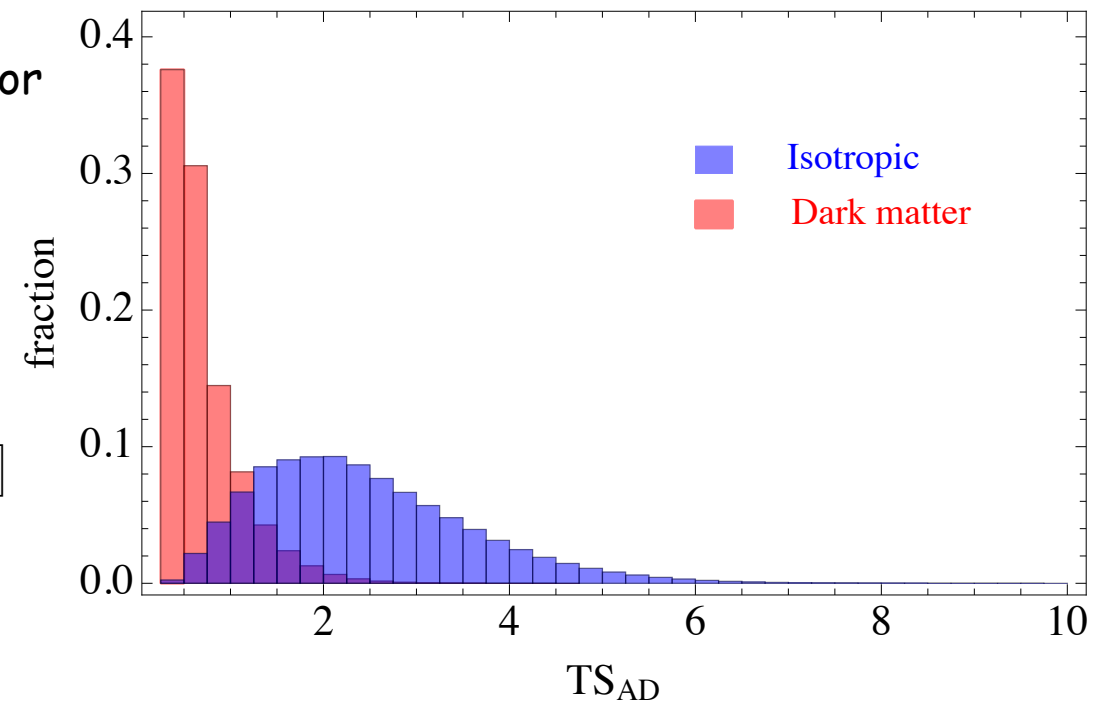
Angular distribution of neutrinos from decaying DM

- ✓ Anderson-Darling test: a powerful non-parametric test, especially sensitive to the end points

Test Statistics

$$TS_{AD} = -N - \frac{1}{N} \sum_{i=1}^N (2i - 1) [\ln(\text{CDF}^{\text{DM}}(\vartheta_i)) + \ln(1 - \text{CDF}^{\text{DM}}(\vartheta_{N+1-i}))]$$

statistically larger TS for
isotropic distribution



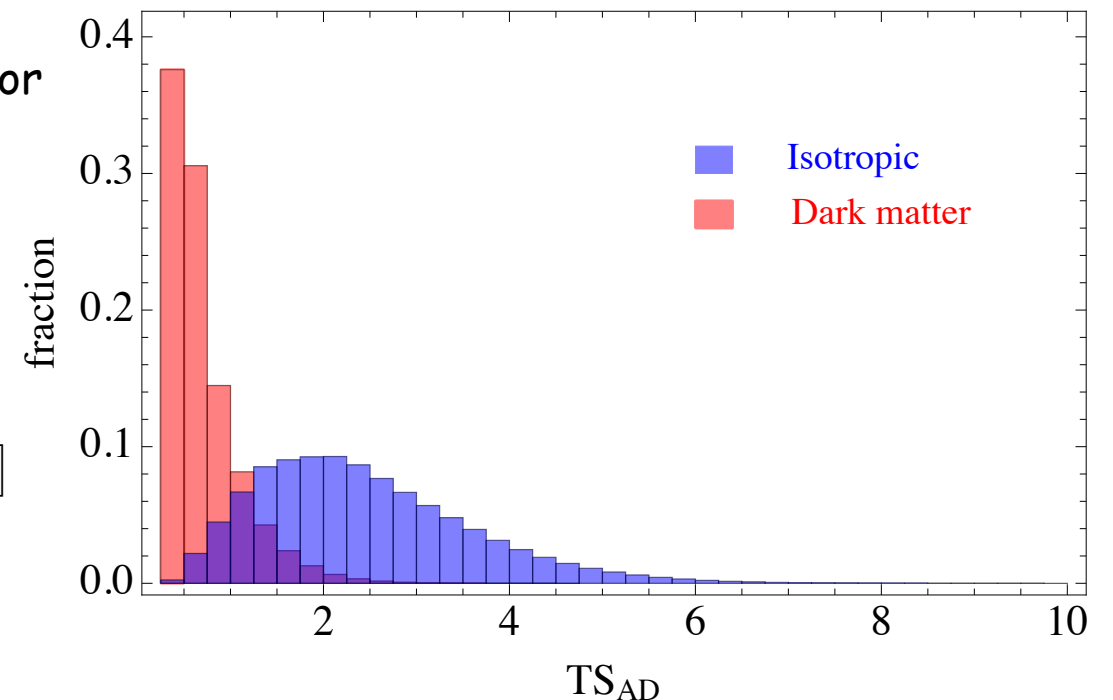
Angular distribution of neutrinos from decaying DM

✓ Anderson-Darling test: a powerful non-parametric test, especially sensitive to the end points

Test Statistics

$$TS_{AD} = -N - \frac{1}{N} \sum_{i=1}^N (2i - 1) [\ln(\text{CDF}^{\text{DM}}(\vartheta_i)) + \ln(1 - \text{CDF}^{\text{DM}}(\vartheta_{N+1-i}))]$$

statistically larger TS for
isotropic distribution



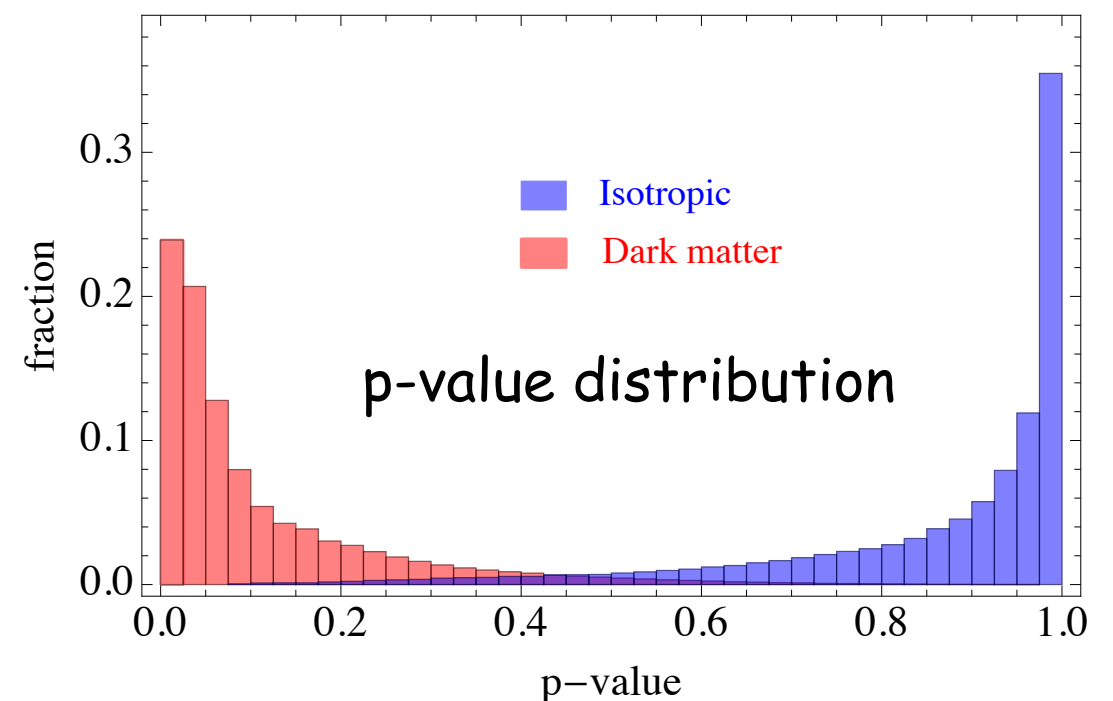
again, generating a sample (10^5) of isotropically
distributed set of 20 events



on the average, 11% of generated isotropic
sample have smaller TS_{KS} than the values
obtained for data vs DM dis.
for data vs isotropic dis. it is 86%



less than 2σ preference for DM dis.



Angular distribution of neutrinos from decaying DM

✓ Some issues:

Angular resolution in KS
and AD tests?



even after shifting all the events
to higher ϑ values still the
preference for DM persist

Angular distribution of neutrinos from decaying DM

✓ Some issues:

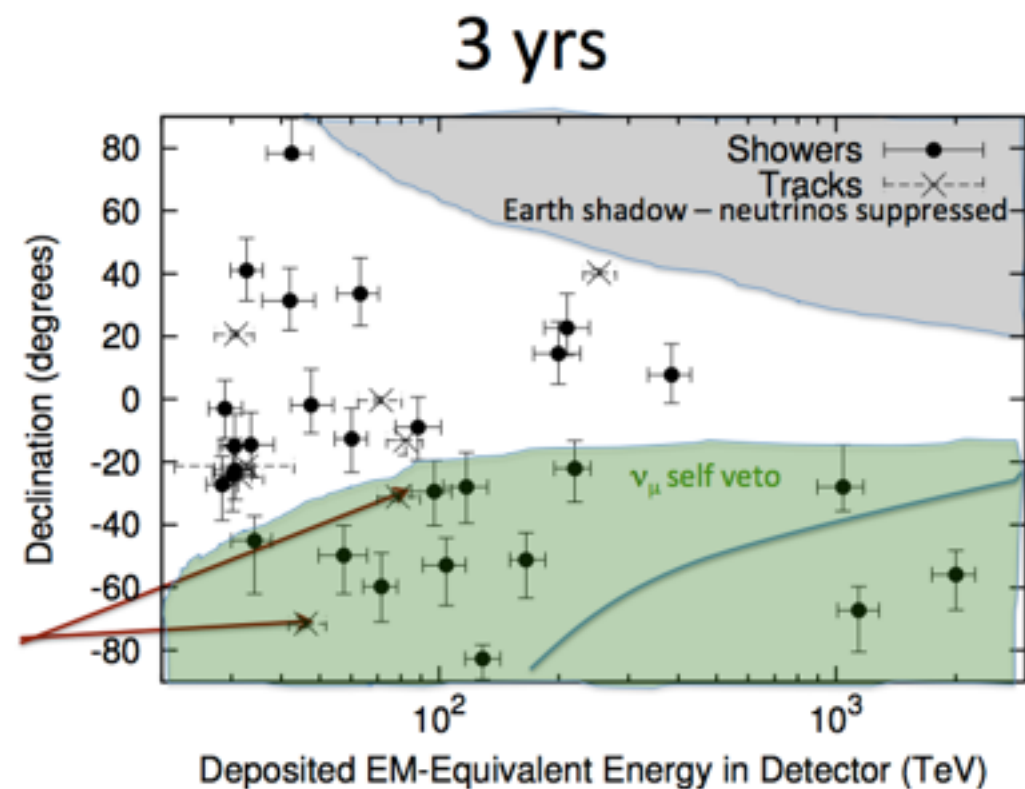
Angular resolution in KS
and AD tests?



even after shifting all the events
to higher ϑ values still the
preference for DM persist

Background rejection?

Figure from T. Gaisser



Angular distribution of neutrinos from decaying DM

✓ Some issues:

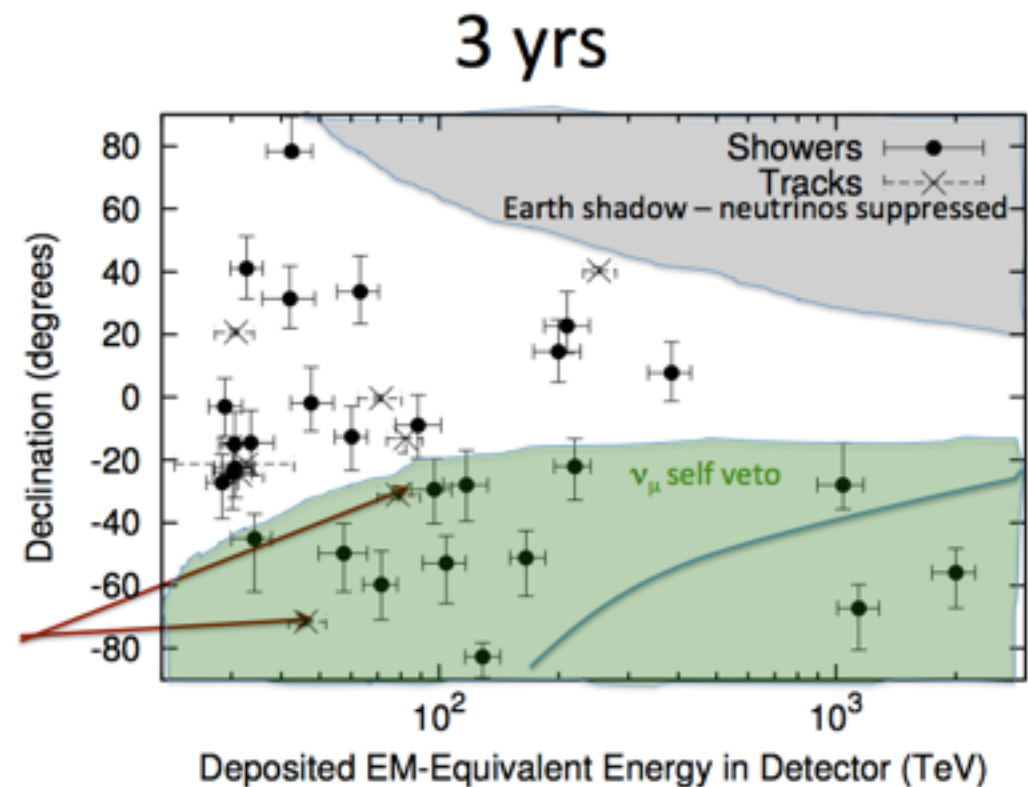
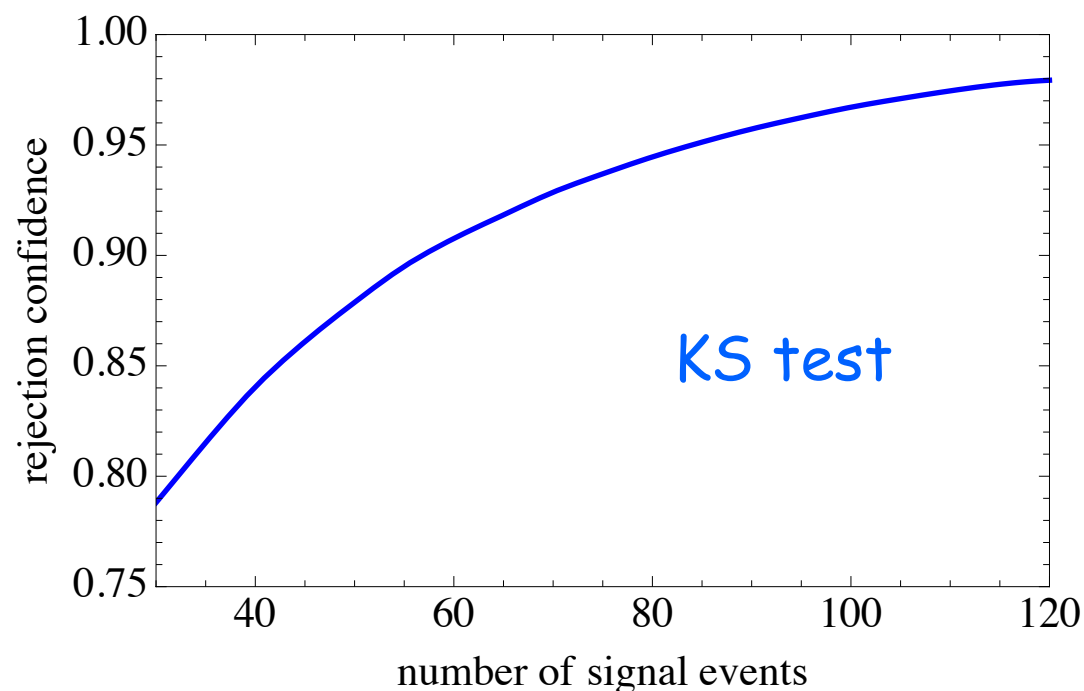
Angular resolution in KS
and AD tests?



even after shifting all the events
to higher ϑ values still the
preference for DM persist

Background rejection?

Figure from T. Gaisser



How many events are
needed for a 3 σ
discrimination?

Gamma ray bounds

Universe is opaque for
gamma-rays with $E > 1 \text{ TeV}$



cascades develop: gamma-ray
interaction with interstellar
radiation field and CMB



gamma-rays populate at
lower energies $< 10^{(2-3)} \text{ GeV}$

Gamma ray bounds

Universe is opaque for gamma-rays with $E > 1$ TeV $\rightarrow$ cascades develop: gamma-ray interaction with interstellar radiation field and CMB $\rightarrow$ gamma-rays populate at lower energies $< 10^{(2-3)}$ GeV

✓ Isotropic diffuse gamma-ray background by Fermi-LAT

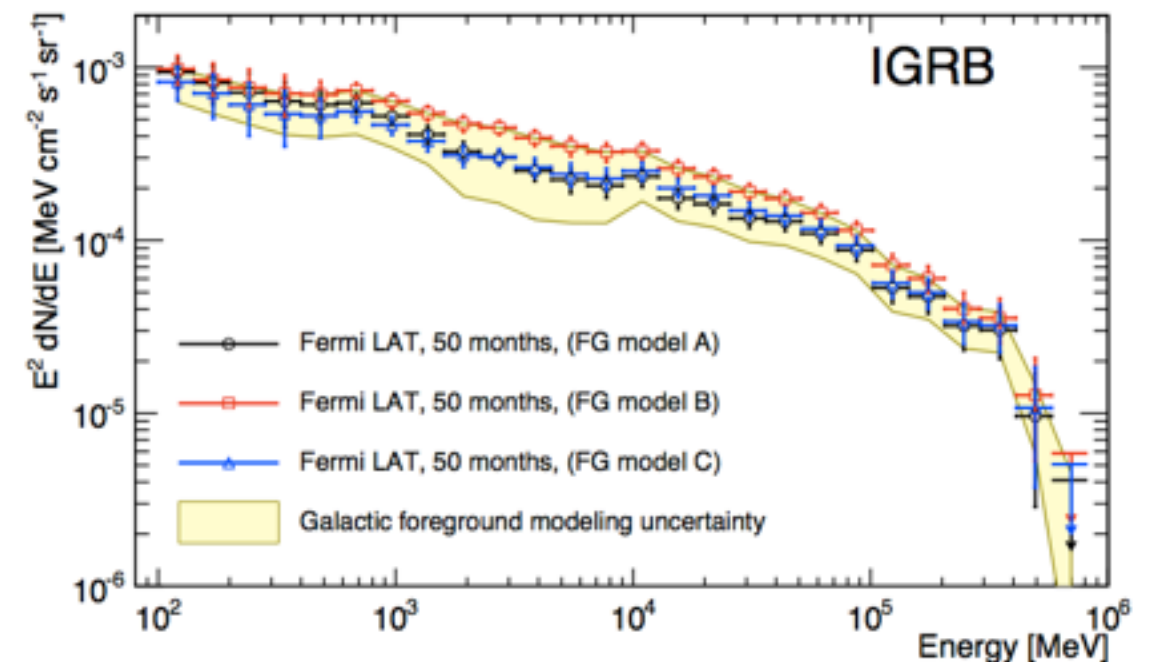
integrated energy density

$$\omega_\gamma = \frac{4\pi}{c} \int_{E_1}^{E_2} E_\gamma \frac{d\varphi_\gamma}{dE_\gamma} dE_\gamma \lesssim 4.4 \times 10^{-7} \text{ eV/cm}^3$$

$$E_1 \sim \mathcal{O}(1) \text{ GeV}$$

$$E_2 \sim \mathcal{O}(100) \text{ GeV}$$

M. Ackermann et al. [The Fermi LAT Collaboration],
arXiv:1410.3696 [astro-ph.HE].



Gamma ray bounds

Universe is opaque for gamma-rays with $E > 1$ TeV $\rightarrow$ cascades develop: gamma-ray interaction with interstellar radiation field and CMB $\rightarrow$ gamma-rays populate at lower energies $< 10^{(2-3)}$ GeV

✓ Isotropic diffuse gamma-ray background by Fermi-LAT

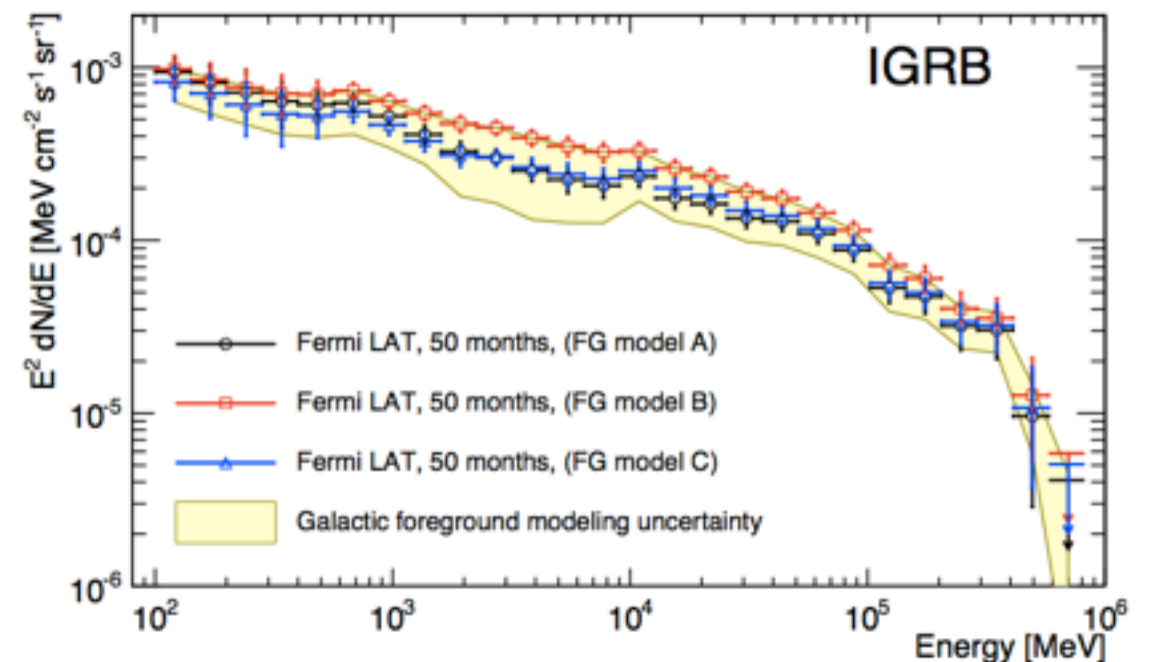
integrated energy density

$$\omega_\gamma = \frac{4\pi}{c} \int_{E_1}^{E_2} E_\gamma \frac{d\varphi_\gamma}{dE_\gamma} dE_\gamma \lesssim 4.4 \times 10^{-7} \text{ eV/cm}^3$$

$$E_1 \sim \mathcal{O}(1) \text{ GeV}$$

$$E_2 \sim \mathcal{O}(100) \text{ GeV}$$

M. Ackermann et al. [The Fermi LAT Collaboration], arXiv:1410.3696 [astro-ph.HE].



total electromagnetic
energy budget
(NH case)

$$\frac{4\pi}{c} \int \sum_{i=\text{gal,extragal}} \left[E_\gamma \frac{d\varphi_\gamma}{dE_\gamma} \right]^i + E_e \frac{d\varphi_{e^\pm}}{dE_e} \Big)^i dE \simeq 5.2 \times 10^{-8} \text{ eV/cm}^3$$



Gamma ray bounds

Universe is opaque for gamma-rays with $E > 1$ TeV



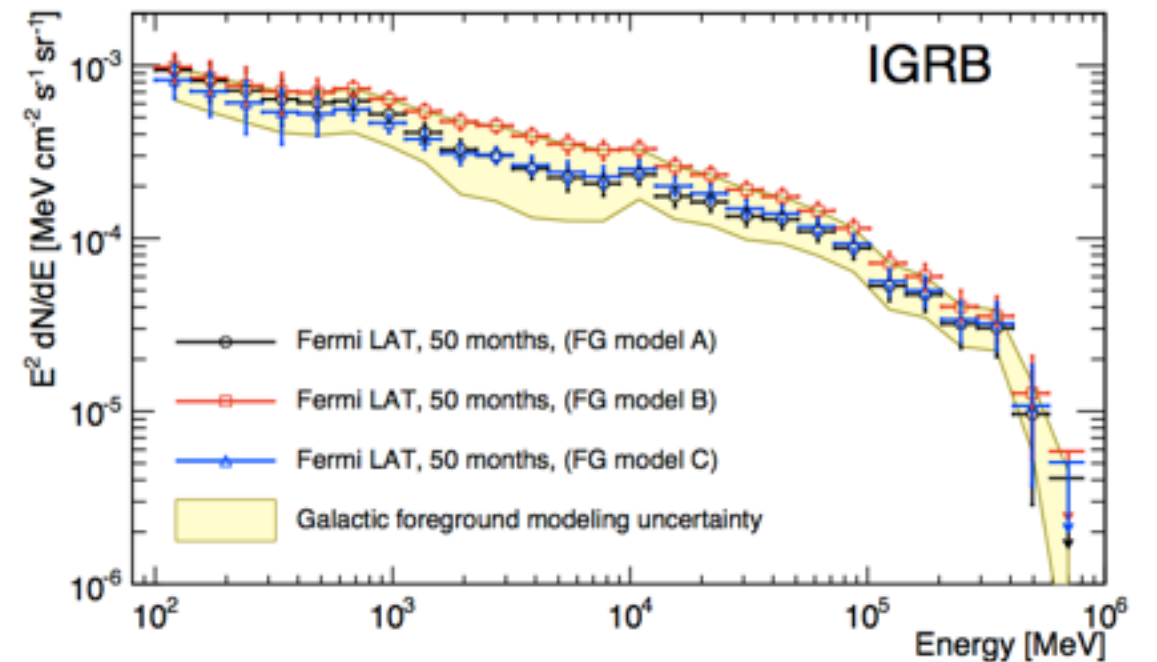
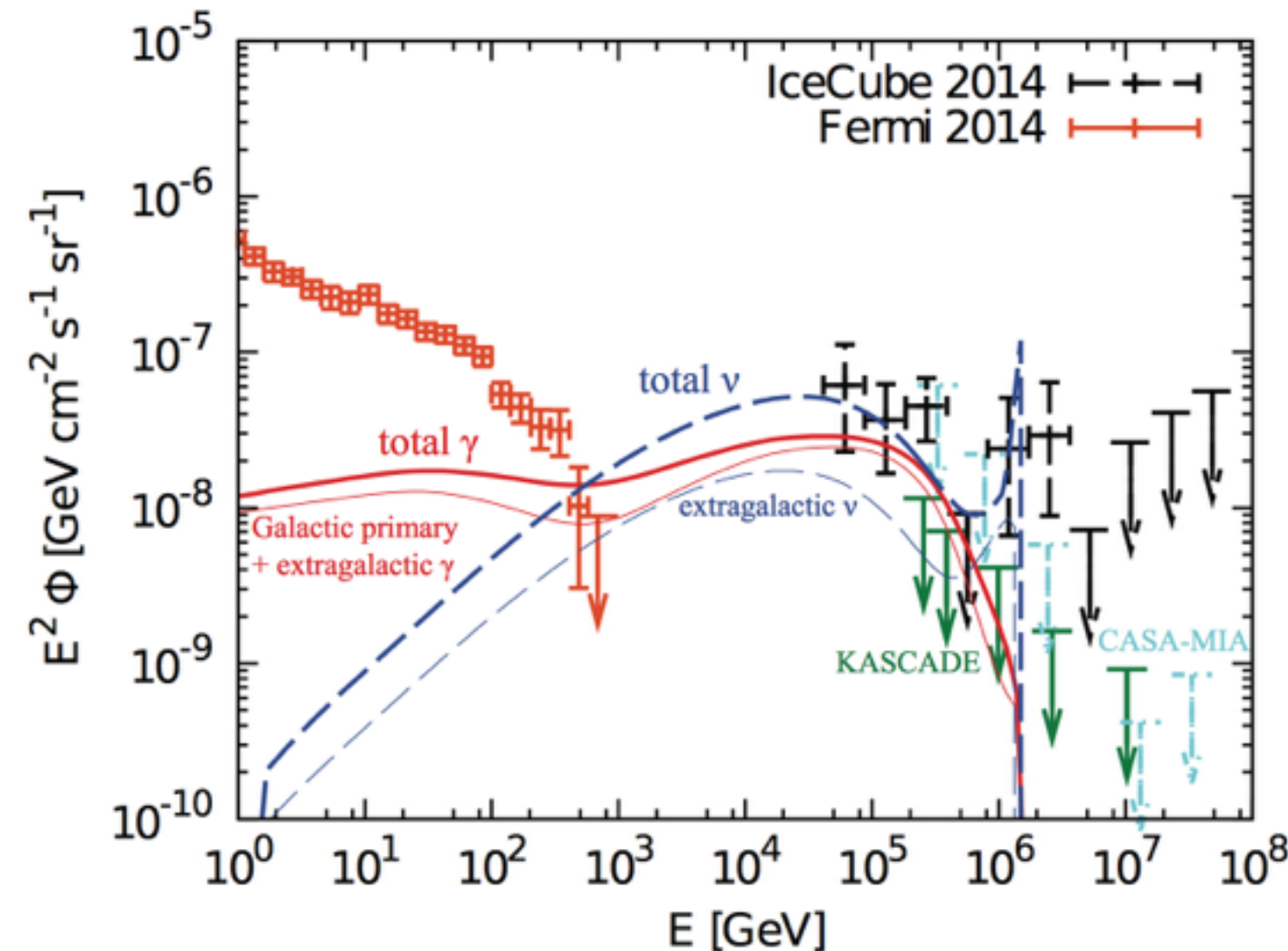
cascades develop: gamma-ray interaction with interstellar radiation field and CMB



gamma-rays populate at lower energies $< 10^{(2-3)}$ GeV

✓ Isotropic diffuse gamma-ray background by Fermi-LAT

M. Ackermann et al. [The Fermi LAT Collaboration], arXiv:1410.3696 [astro-ph.HE].



Murase, Laha, Ando, Ahlers, arXiv:1503.04663

Gamma ray bounds

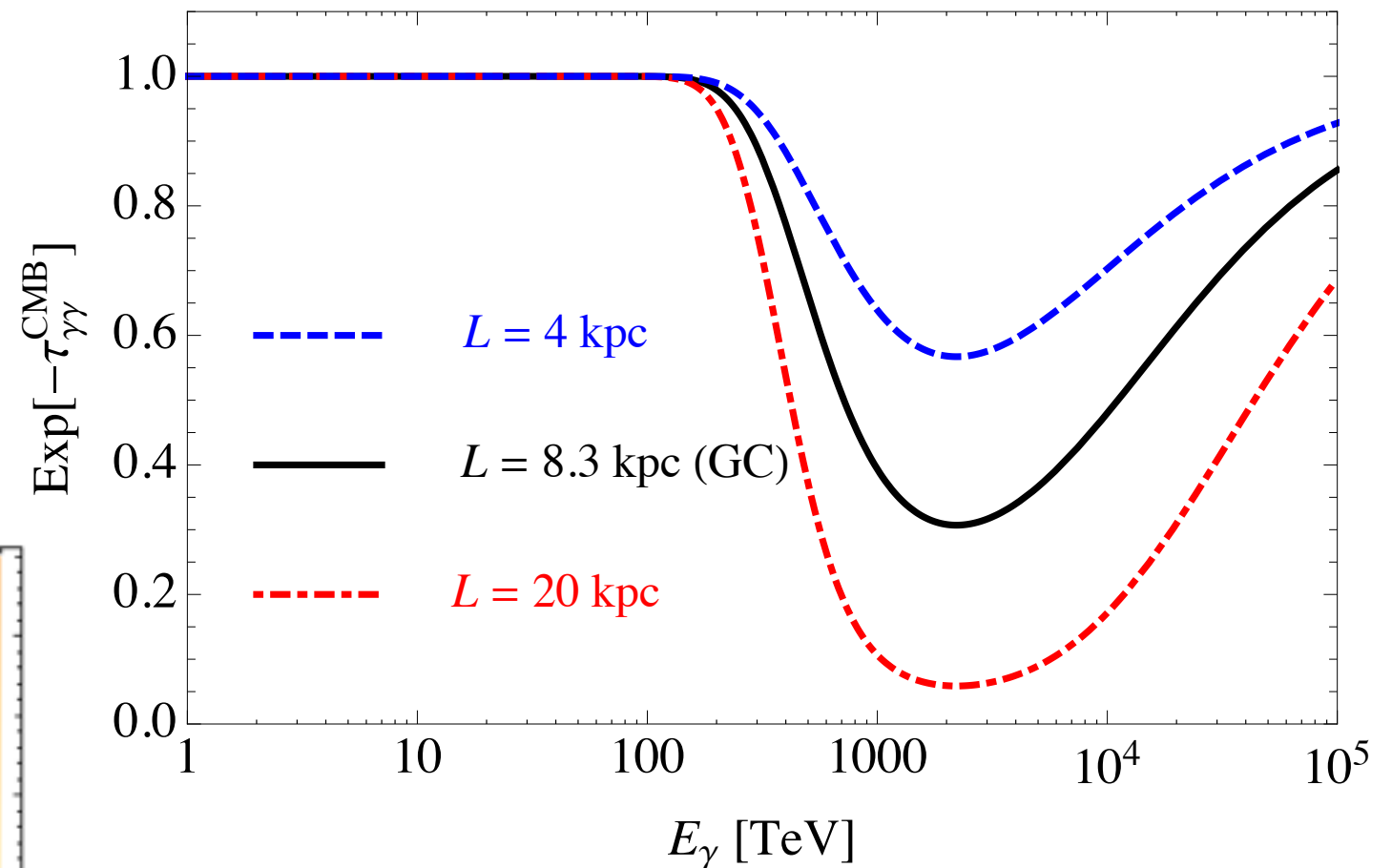
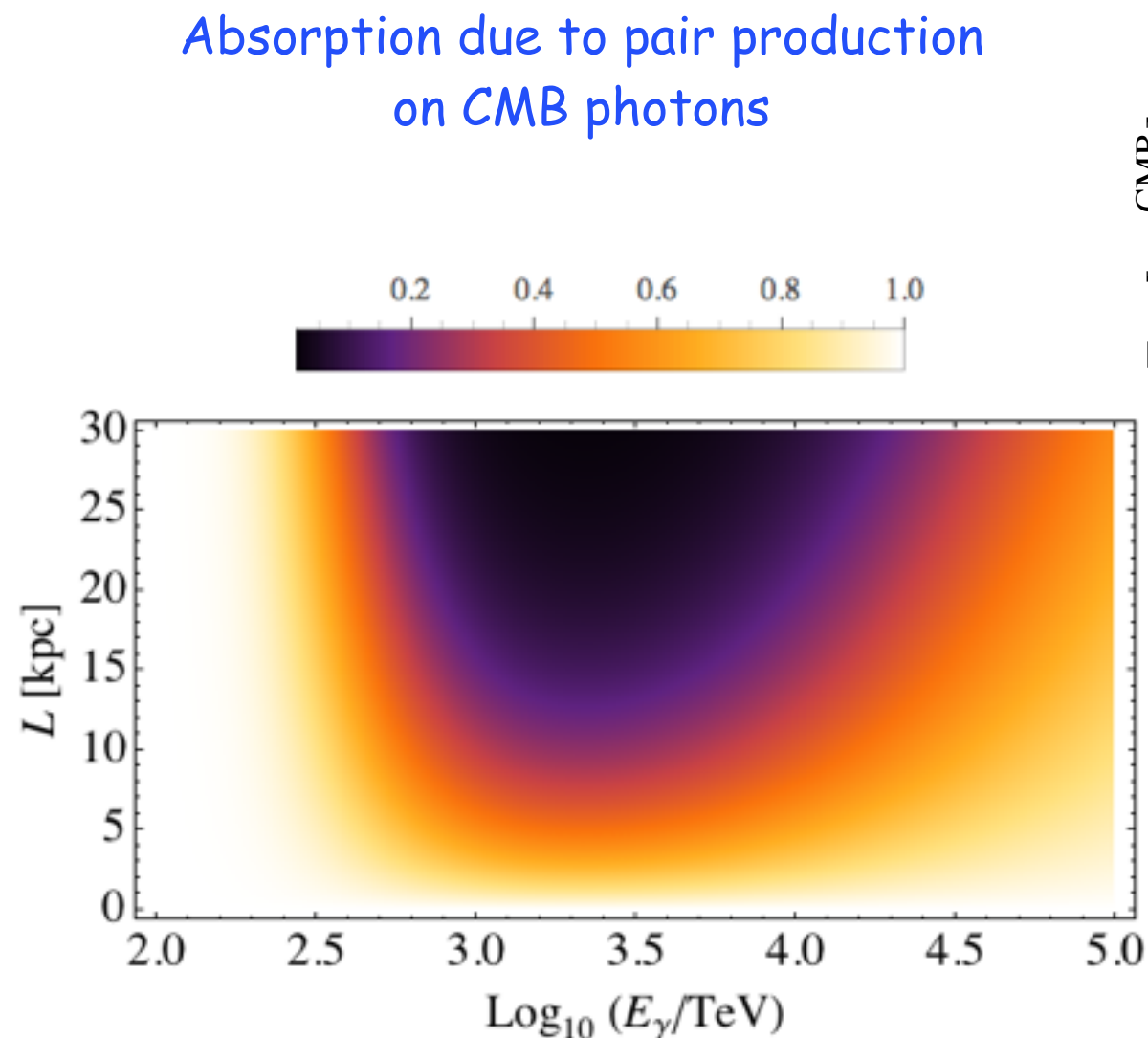
✓ Galactic component

at $\sim$ PeV, the absorption length of gamma-rays are comparable to Galactic distances



neither full absorption or cascade development, nor full transparency

A. E. and P. Serpico, arXiv:1505.06486



Absorption at $\sim$ PeV

Gamma ray bounds

✓ Galactic component

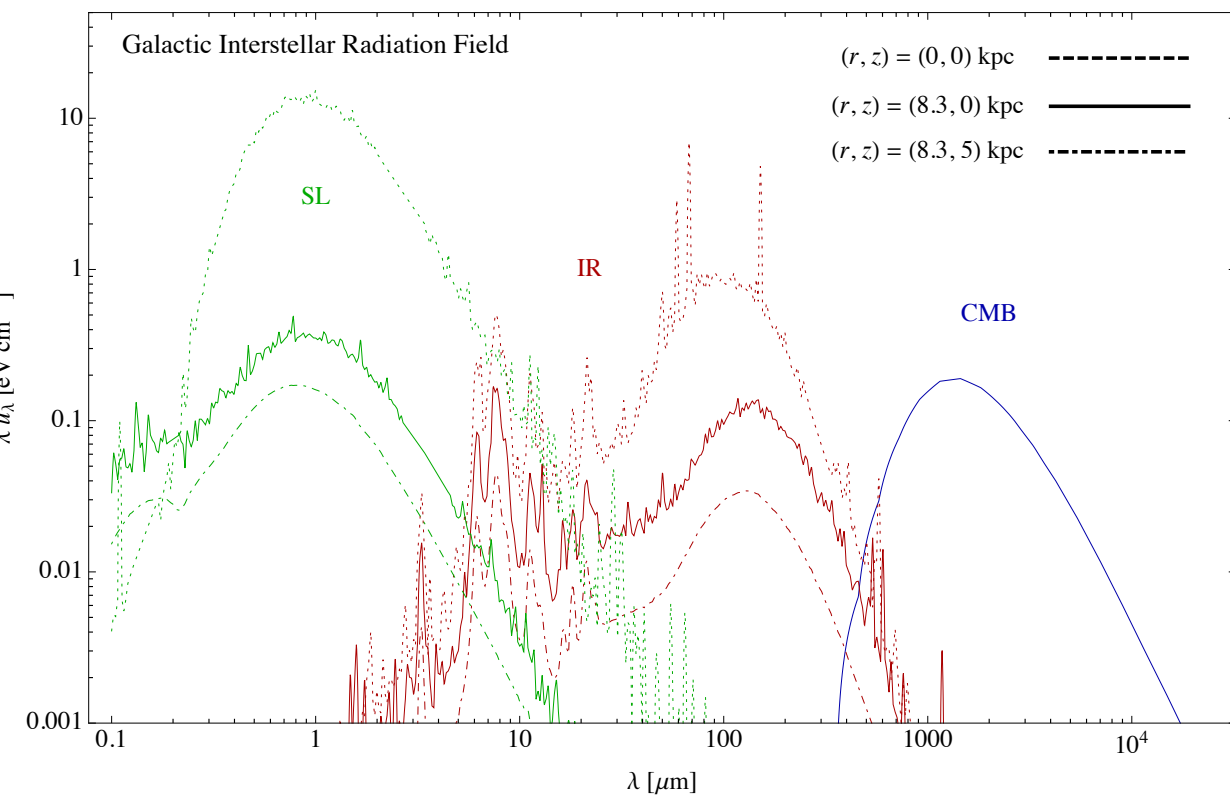
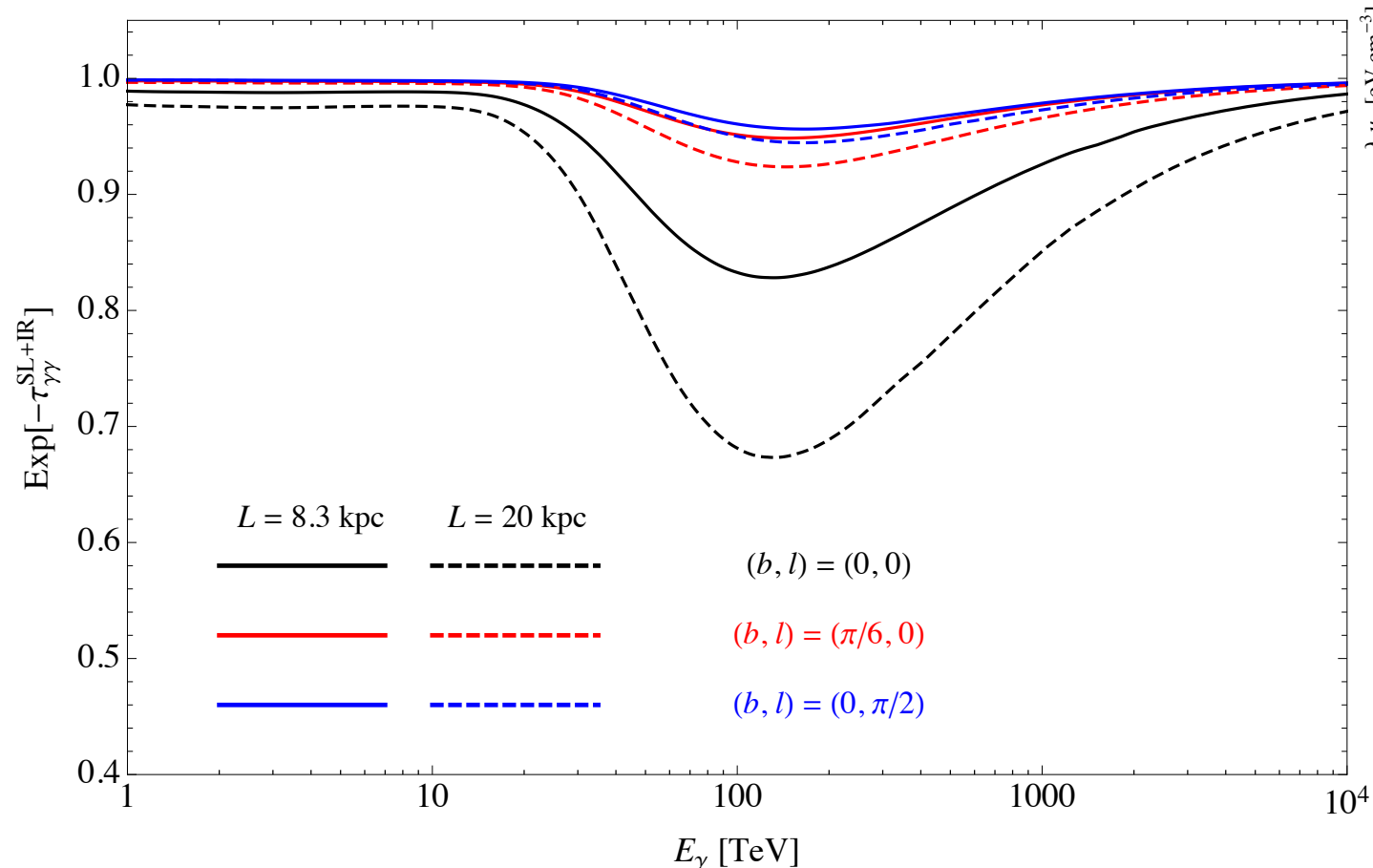
at $\sim \text{PeV}$, the absorption length of gamma-rays are comparable to Galactic distances



neither full absorption or cascade development, nor full transparency

A. E. and P. Serpico, arXiv:1505.06486

Absorption due to pair production
on SL+IR photons



Absorption at $\sim 100 \text{ TeV}$

Gamma ray bounds

✓ Galactic component

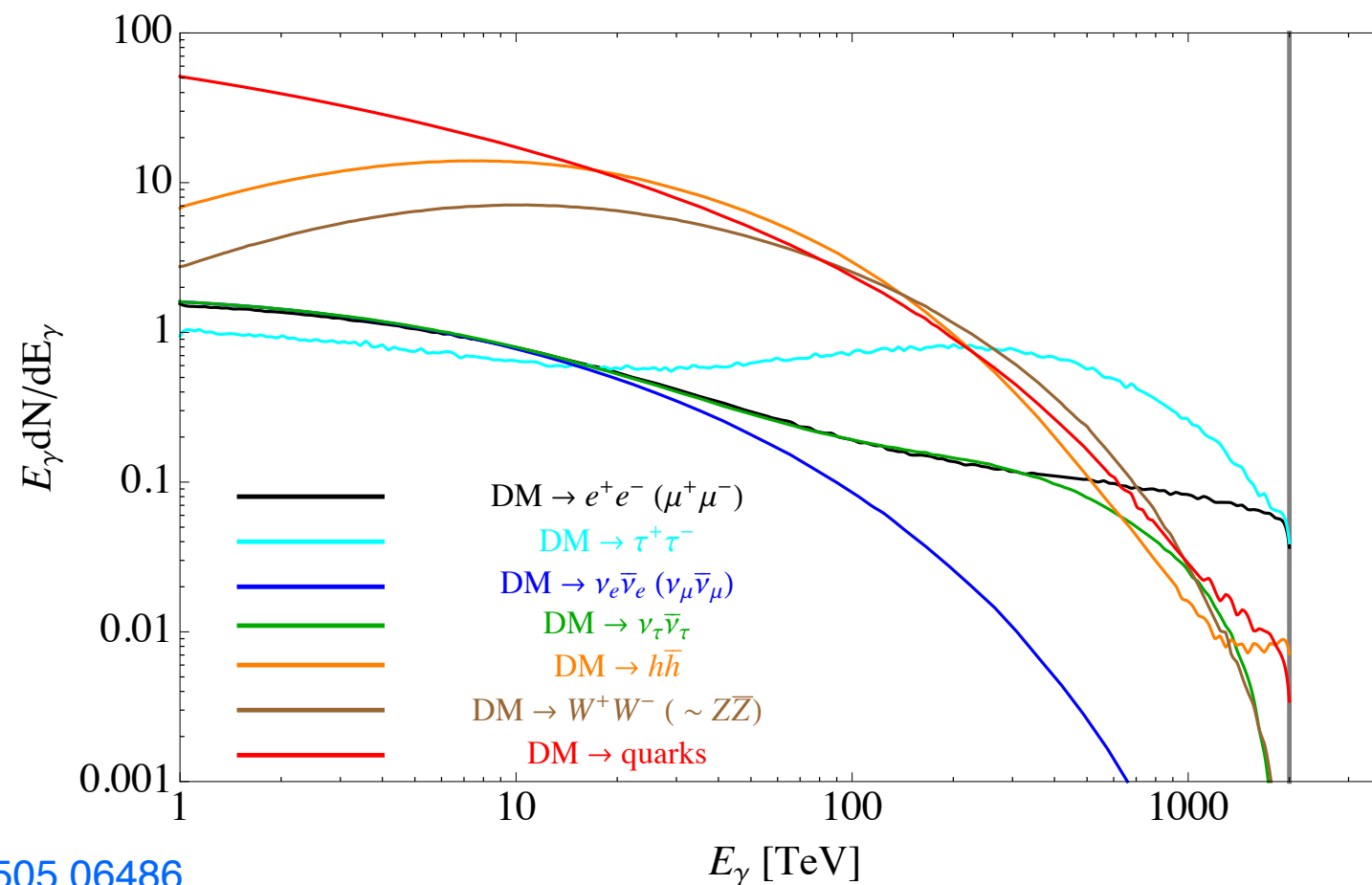
at $\sim \text{PeV}$, the absorption length of gamma-rays are comparable to Galactic distances



neither full absorption or cascade development, nor full transparency

Prompt component

$$\frac{d\Phi_\gamma}{dE_\gamma}(E_\gamma, b, l) = \frac{1}{4\pi m_{\text{DM}} \tau_{\text{DM}}} \frac{dN_\gamma}{dE_\gamma}(E_\gamma) \int_0^\infty \rho_h[\varrho(s, b, l)] e^{-\tau_{\gamma\gamma}(E_\gamma, s, b, l)} ds$$



calculated by
PYTHIA 8.2

Gamma ray bounds

✓ Galactic component

at $\sim \text{PeV}$, the absorption length of gamma-rays are comparable to Galactic distances

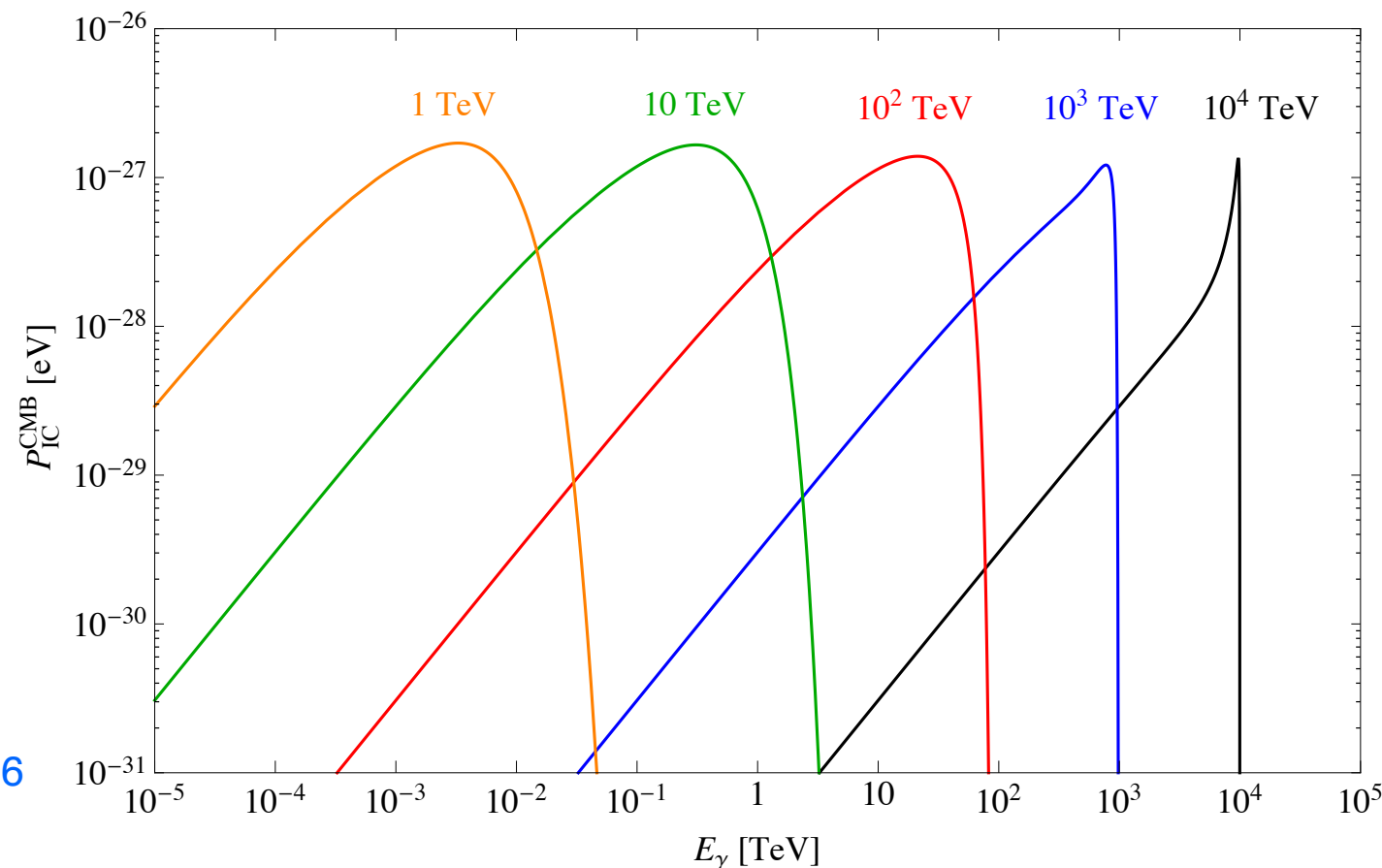


neither full absorption or cascade development, nor full transparency

inverse-Compton
component

$$\frac{d\Phi_{\text{IC}}}{dE_\gamma}(E_\gamma, b, l) = \frac{1}{4\pi E_\gamma} \int_0^\infty ds e^{-\tau_{\gamma\gamma}(E_\gamma, s, b, l)} \int_{m_e}^{m_{\text{DM}}/2} dE_e \frac{dn_e}{dE_e}(E_e, \varrho) P_{\text{IC}}(E_e, E_\gamma, \varrho)$$

$$P_{\text{IC}} = P_{\text{IC}}^{\text{CMB}} + P_{\text{IC}}^{\text{SL+IR}}$$



A. E. and P. Serpico, arXiv:1505.06486

Gamma ray bounds

✓ Galactic component

at $\sim$ PeV, the absorption length of gamma-rays are comparable to Galactic distances

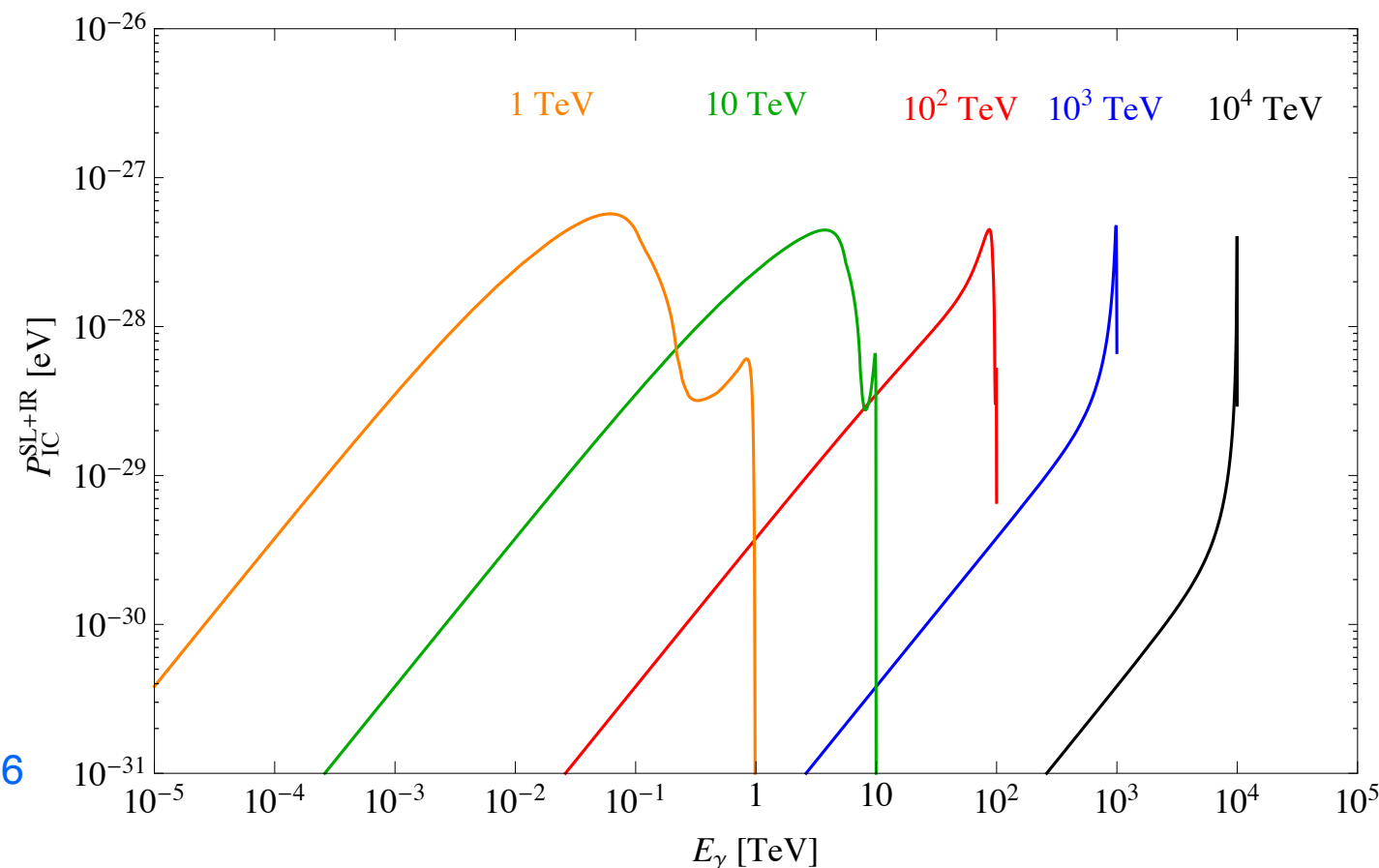


neither full absorption or cascade development, nor full transparency

inverse-Compton
component

$$\frac{d\Phi_{\text{IC}}}{dE_\gamma}(E_\gamma, b, l) = \frac{1}{4\pi E_\gamma} \int_0^\infty ds e^{-\tau_{\gamma\gamma}(E_\gamma, s, b, l)} \int_{m_e}^{m_{\text{DM}}/2} dE_e \frac{dn_e}{dE_e}(E_e, \varrho) P_{\text{IC}}(E_e, E_\gamma, \varrho)$$

$$P_{\text{IC}} = P_{\text{IC}}^{\text{CMB}} + P_{\text{IC}}^{\text{SL+IR}}$$



A. E. and P. Serpico, arXiv:1505.06486

Gamma ray bounds

✓ Galactic component

at $\sim \text{PeV}$, the absorption length of gamma-rays are comparable to Galactic distances

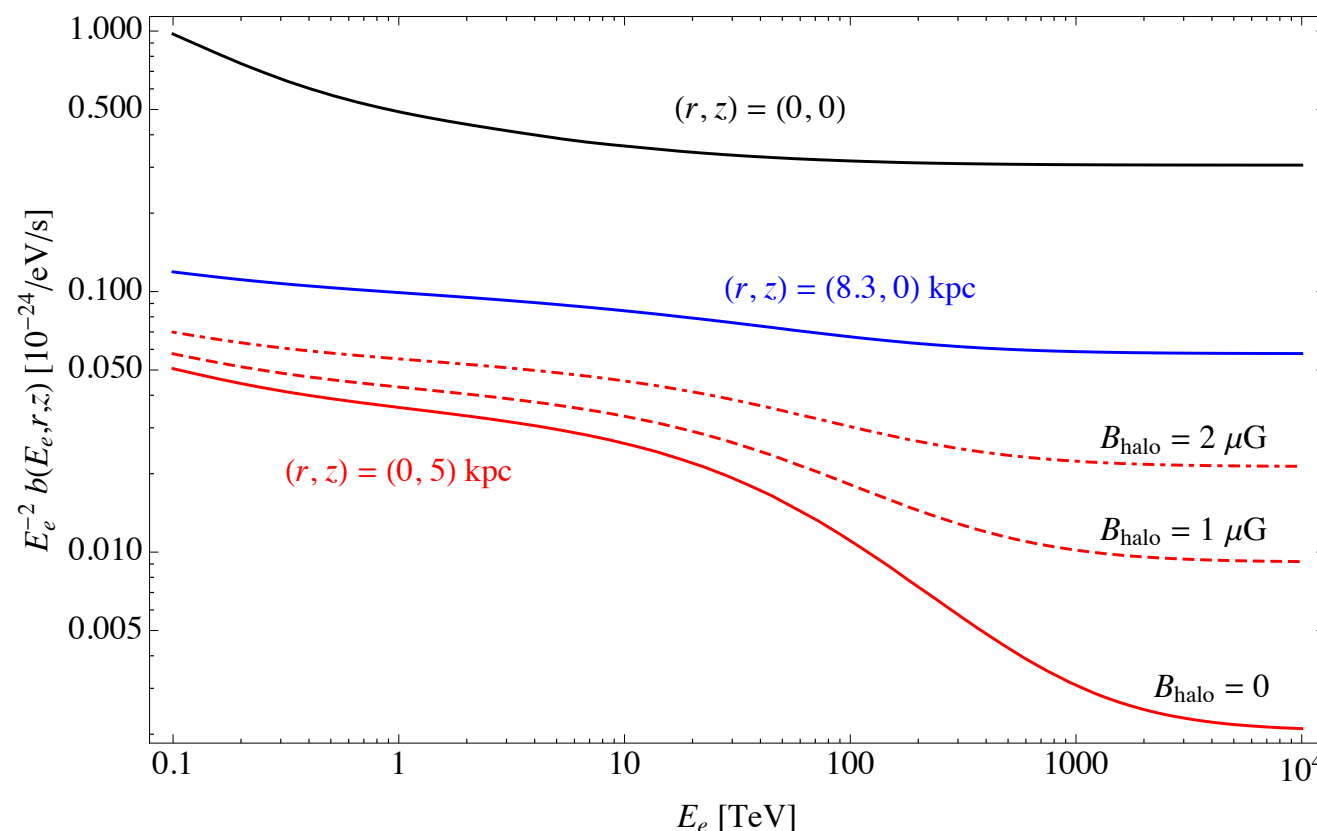


neither full absorption or cascade development, nor full transparency

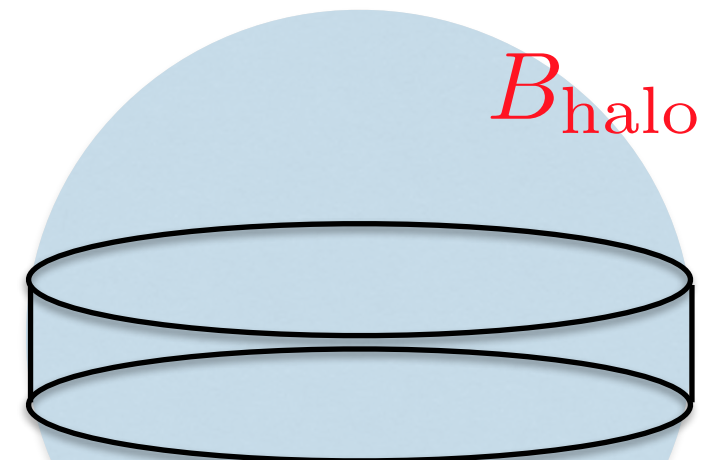
inverse-Compton component

$$\frac{d\Phi_{\text{IC}}}{dE_\gamma}(E_\gamma, b, l) = \frac{1}{4\pi E_\gamma} \int_0^\infty ds e^{-\tau_{\gamma\gamma}(E_\gamma, s, b, l)} \int_{m_e}^{m_{\text{DM}}/2} dE_e \frac{dn_e}{dE_e}(E_e, \varrho) P_{\text{IC}}(E_e, E_\gamma, \varrho)$$

$$\frac{dn_e}{dE_e}(E_e, \vec{x}) = \frac{1}{m_{\text{DM}} \tau_{\text{DM}}} \frac{\rho_h(\vec{x})}{b(E_e, \vec{x})} \int_{E_e}^{m_{\text{DM}}/2} \frac{dN_e}{dE'_e}(E'_e) I_{\text{diff}}(E_e, E'_e, \vec{x}) dE'_e$$



$$b(E_e, \vec{x}) \equiv -\frac{dE_e}{dt} = b_{\text{IC}}(E_e, \vec{x}) + b_{\text{syn}}(E_e, \vec{x})$$

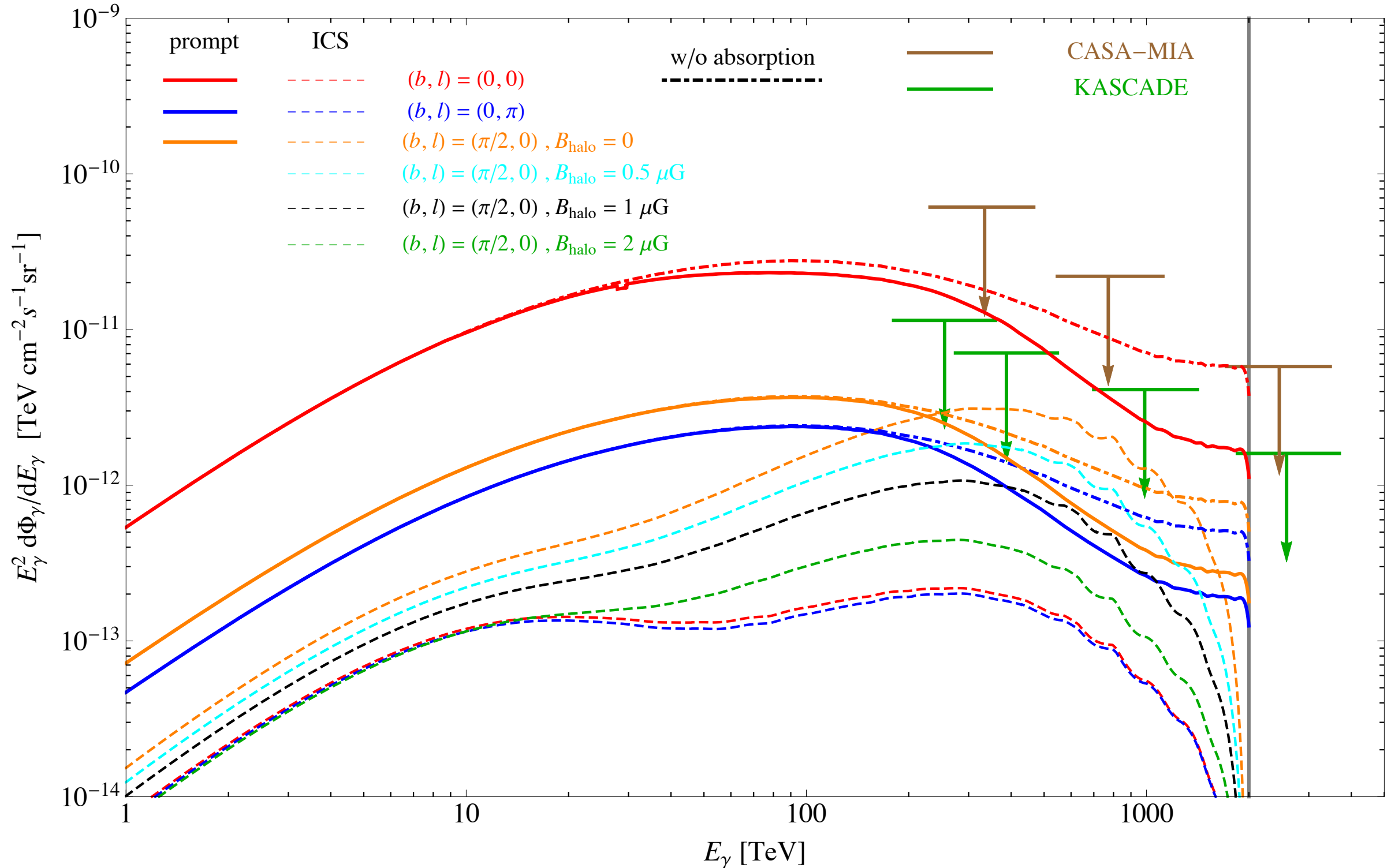


$$B_{\text{reg}}(\vec{x}) = B_0 \exp \left[-\frac{|r - R_\odot|}{r_B} - \frac{|z|}{z_B} \right]$$

Gamma ray bounds

✓ Galactic component

$$\tau_{\text{DM}} = 10^{28} \text{ s} \quad \text{and} \quad m_{\text{DM}} = 4 \text{ PeV}$$

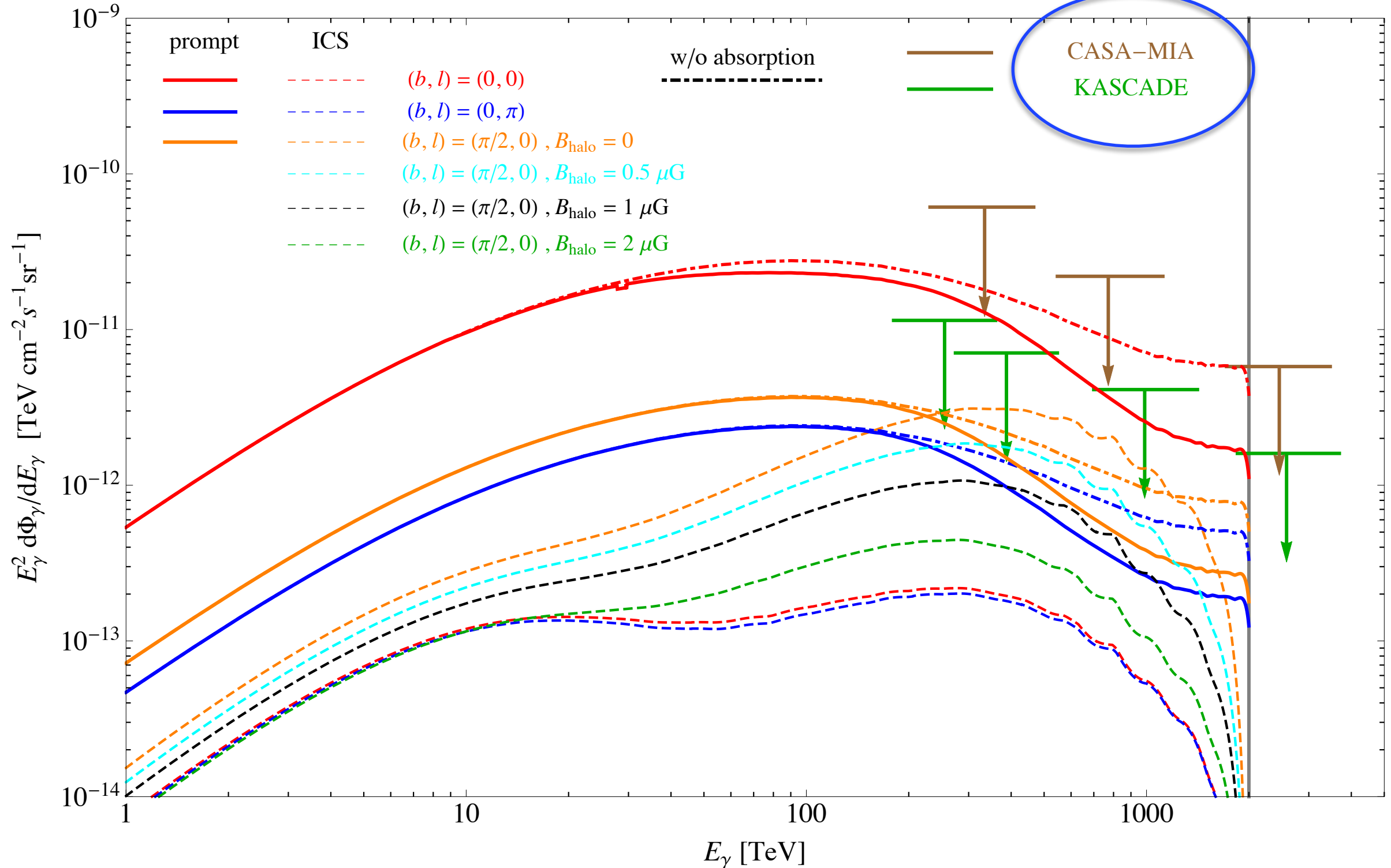


A. E. and P. Serpico, arXiv:1505.06486

Gamma ray bounds

✓ Galactic component

$$\tau_{\text{DM}} = 10^{28} \text{ s} \quad \text{and} \quad m_{\text{DM}} = 4 \text{ PeV}$$



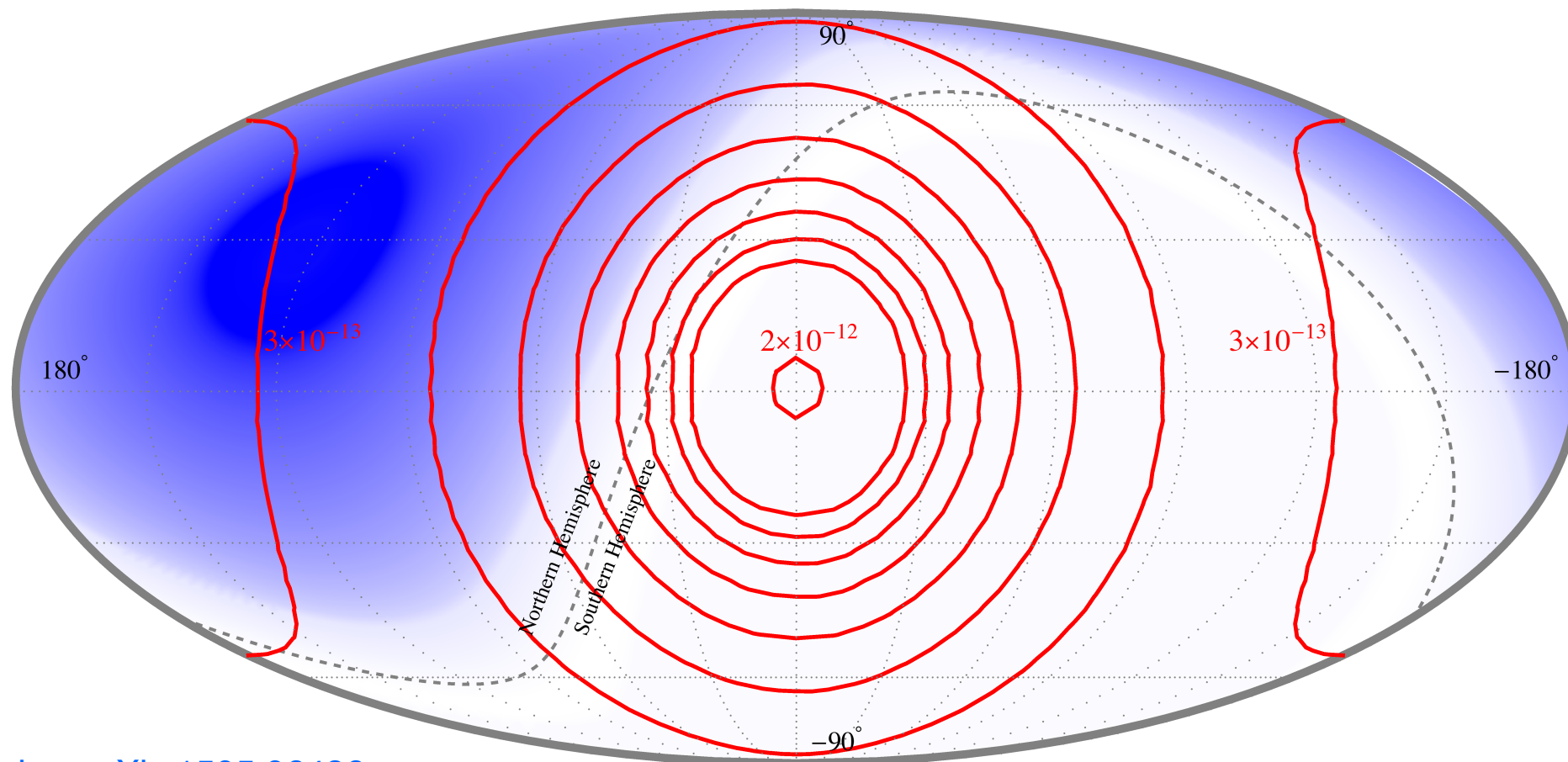
A. E. and P. Serpico, arXiv:1505.06486

Gamma ray bounds

✓ Galactic component



KASCADE

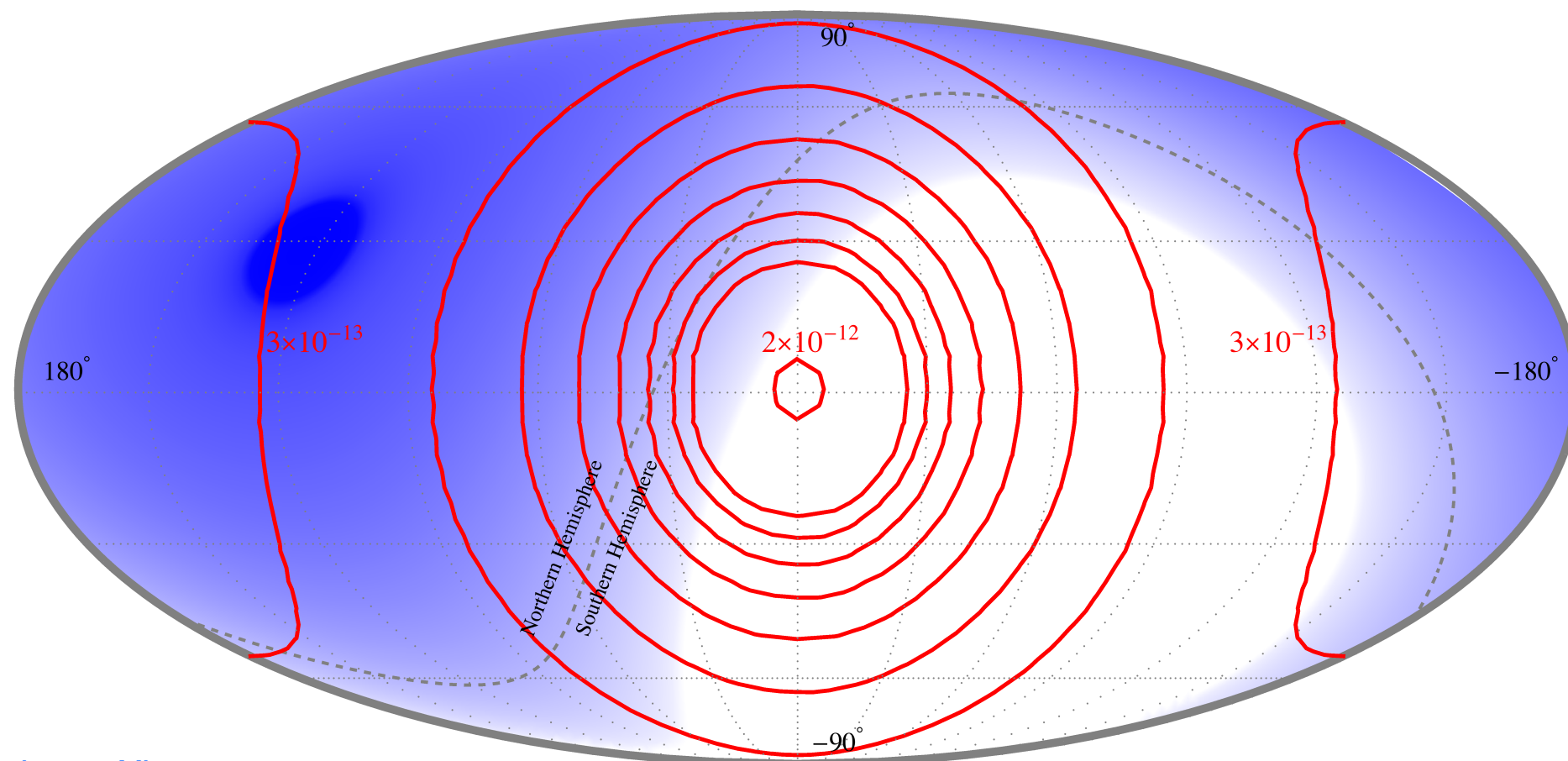


A. E. and P. Serpico, arXiv:1505.06486

Gamma ray bounds

✓ Galactic component

CASA-MIA



A. E. and P. Serpico, arXiv:1505.06486

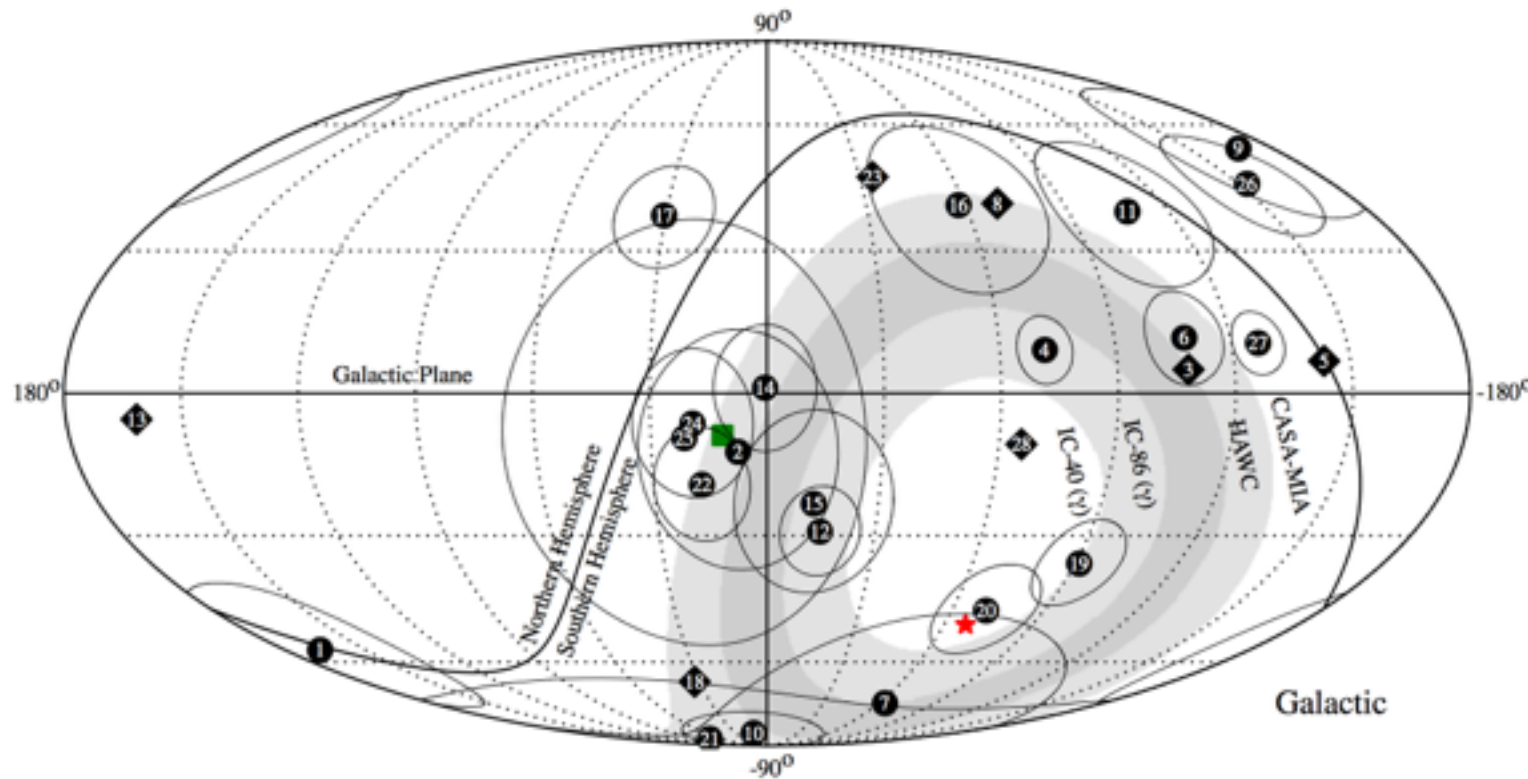
Gamma ray bounds

- ✓ Galactic component
- Future experiments

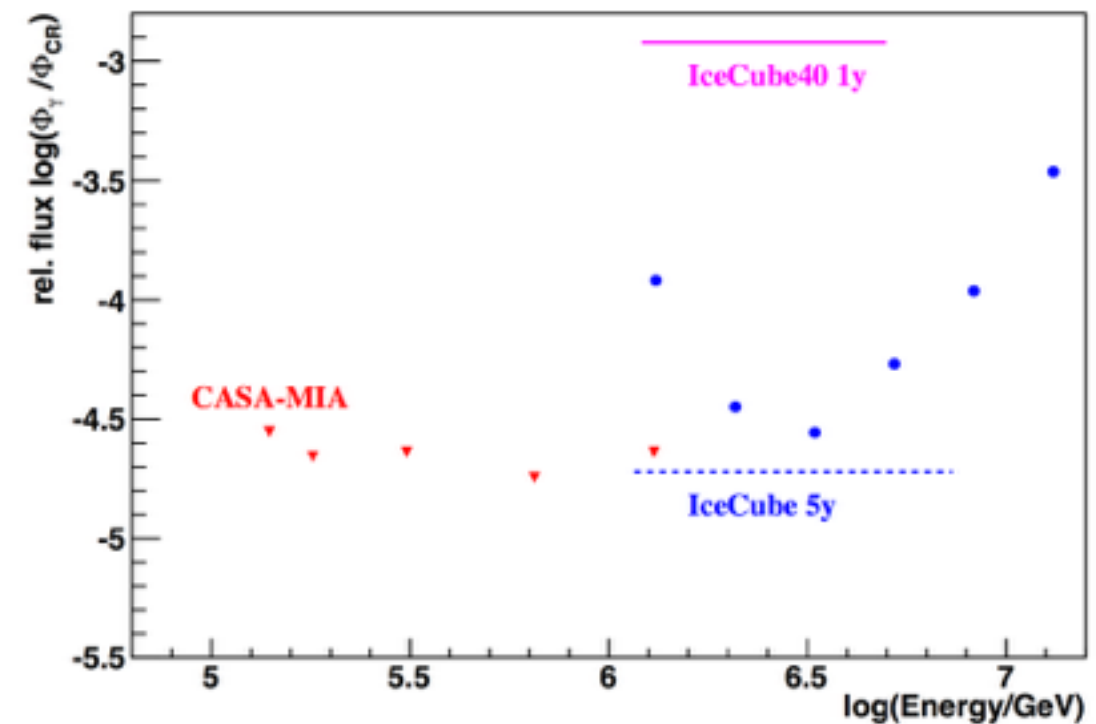
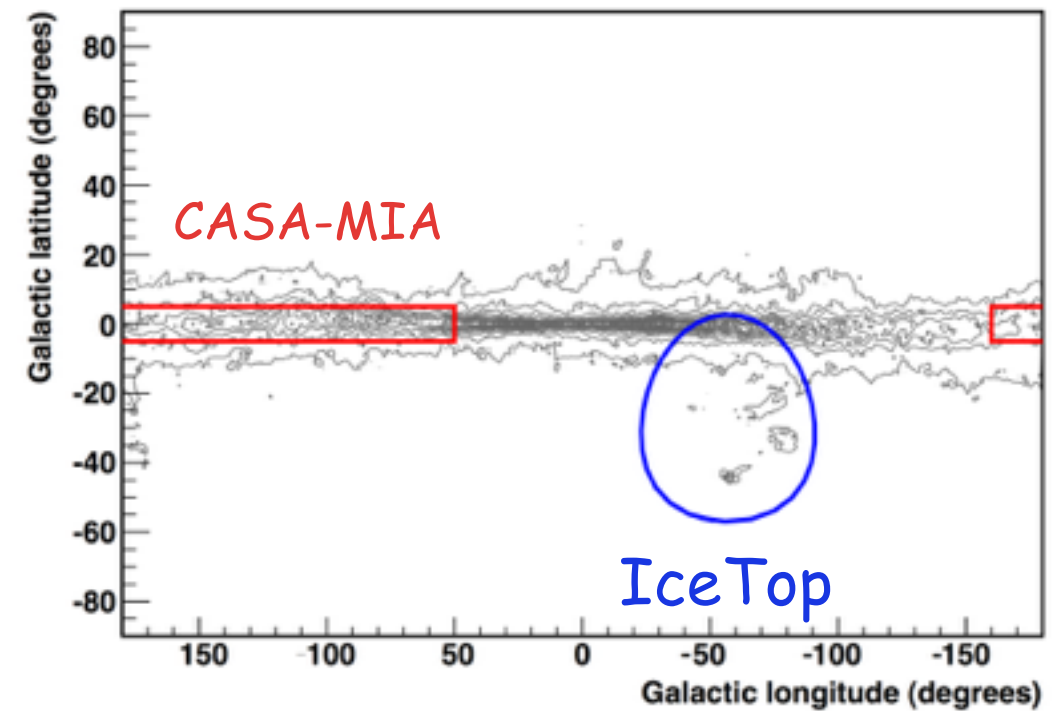
HAWC



M. Ahlers and K. Murase,
PRD (2014), arXiv:1309.4077



IceTop



Gamma ray bounds

✓ Galactic component

Anisotropy

$$a_\gamma = \frac{\left. \frac{d\Phi_\gamma}{dE_\gamma} \right|_{\text{GC}} - \left. \frac{d\Phi_\gamma}{dE_\gamma} \right|_{\text{anti-GC}}}{\frac{d\Phi_{\text{CR}}}{dE}}$$

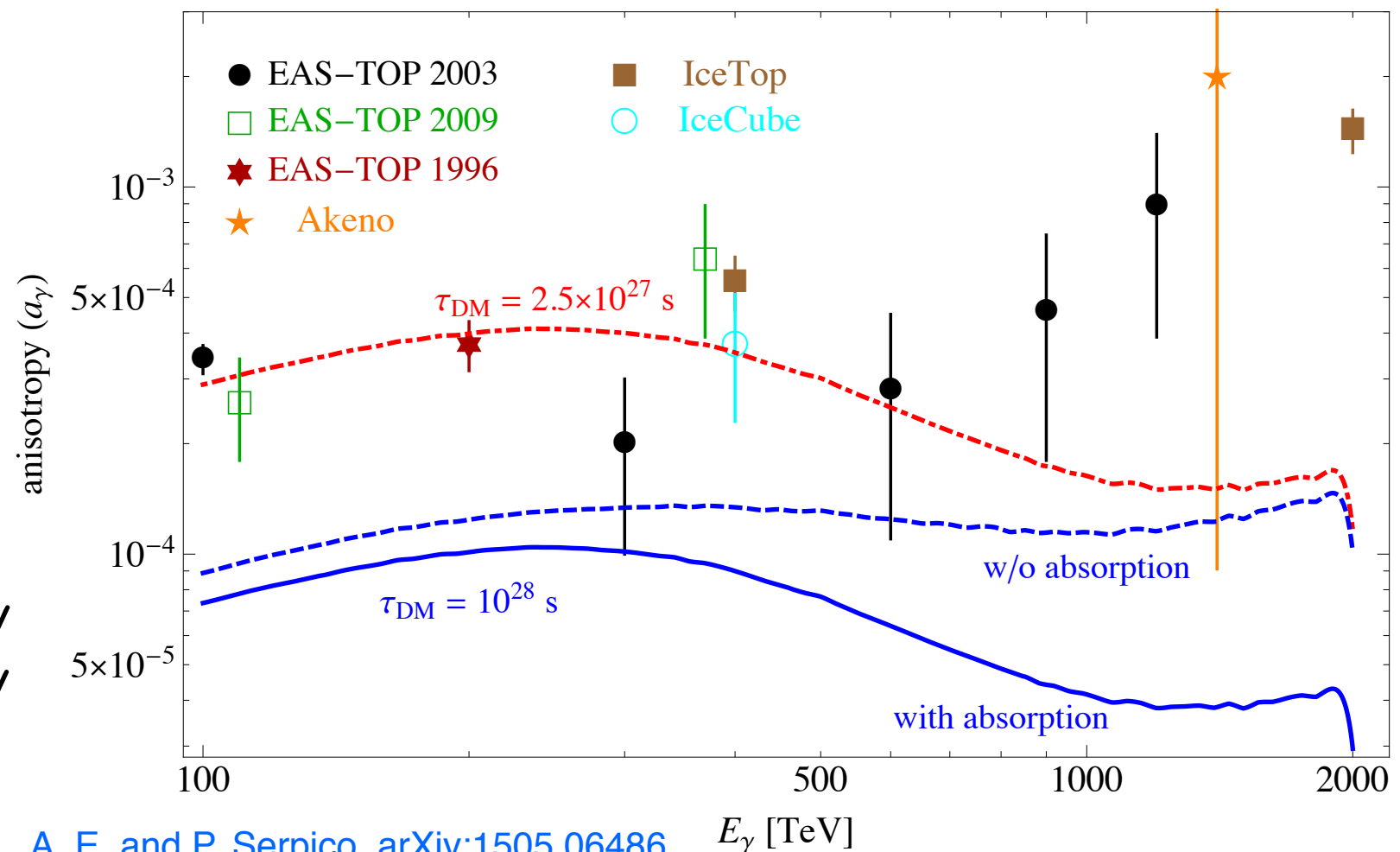
Total CR flux

✓ No need to γ /hadron discrimination

✓ Absorption suppress the anisotropy

✓ The bound 2.5×10^{27} s can be set

✓ Adding the phase info of anisotropy would improve the limits significantly



A. E. and P. Serpico, arXiv:1505.06486

Gamma ray bounds

✓ Galactic component

Anisotropy

$$a_\gamma = \frac{\left. \frac{d\Phi_\gamma}{dE_\gamma} \right|_{\text{GC}} - \left. \frac{d\Phi_\gamma}{dE_\gamma} \right|_{\text{anti-GC}}}{\frac{d\Phi_{\text{CR}}}{dE}}$$

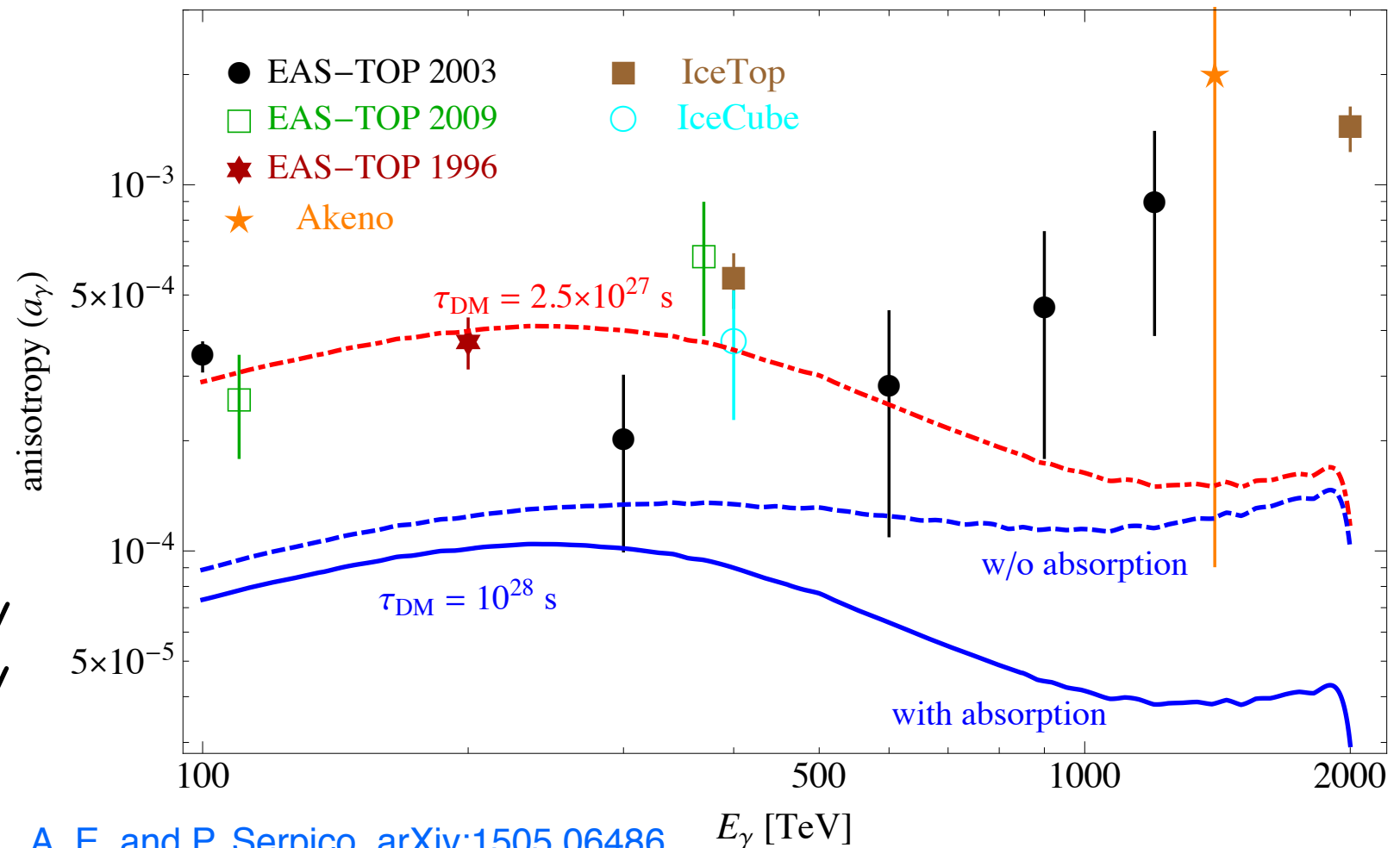
Total CR flux

✓ No need to γ /hadron discrimination

✓ Absorption suppress the anisotropy

✓ The bound 2.5×10^{27} s can be set

✓ Adding the phase info of anisotropy would improve the limits significantly



A. E. and P. Serpico, arXiv:1505.06486

What can we learn about DM if the IceCube events originate from conventional astrophysical flux?

constraining:

- ✓ DM lifetime
- ✓ annihilation cross section

Constraining DM properties

✓ DM lifetime

contribution of DM to the events in each bin should be smaller than N_{limit}

bin #	$\log_{10}(E_\nu/\text{TeV})$	$N_{\text{astro}}(E_\nu^{-2} \div E_\nu^{-2.3})$	N_{data}	$N_{\text{limit}} (E_\nu^{-2} \div E_\nu^{-2.3})$	N_{limit}
#1	1.4 – 1.6	9.46 ÷ 10	11	7.8 ÷ 7.46	16.6
#2	1.6 – 1.8	4.31 ÷ 5.3	6	6.53 ÷ 5.87	10.5
#3	1.8 – 2.0	4.55 ÷ 5.68	7	7.41 ÷ 6.58	11.8
#4	2.0 – 2.2	3.97 ÷ 4.82	3	3.98 ÷ 3.73	6.68
#5	2.2 – 2.4	3.32 ÷ 3.56	4	5.15 ÷ 5.01	8.00
#6	2.4 – 2.6	2.59 ÷ 2.42	2	3.65 ÷ 3.71	5.32
#7	2.6 – 2.8	1.96 ÷ 1.62	0	2.3 ÷ 2.3	2.3
#8	2.8 – 3.0	1.55 ÷ 1.1	0	2.3 ÷ 2.3	2.3
#9	3.0 – 3.2	1.2 ÷ 0.74	2	4.31 ÷ 4.64	5.32
#10	3.2 – 3.4	0.92 ÷ 0.5	1	3.3 ÷ 3.51	3.89
#11	3.4 – 3.6	0.73 ÷ 0.35	0	2.3 ÷ 2.3	2.3
#12	3.6 – 3.8	1.72 ÷ 0.76	0	2.3 ÷ 2.3	2.3

Poisson statistics:

at $q\%$ C.L.

$$\frac{q}{100} = \frac{\int_0^{N_{\text{limit}}^i} L(N_{\text{data}}^i, N) dN}{\int_0^\infty L(N_{\text{data}}^i, N) dN}$$

$$L(N_{\text{data}}^i, N) = \frac{(N + N_{\text{astro}}^i)^{N_{\text{data}}^i}}{N_{\text{data}}^i!} e^{-(N + N_{\text{astro}}^i)}$$

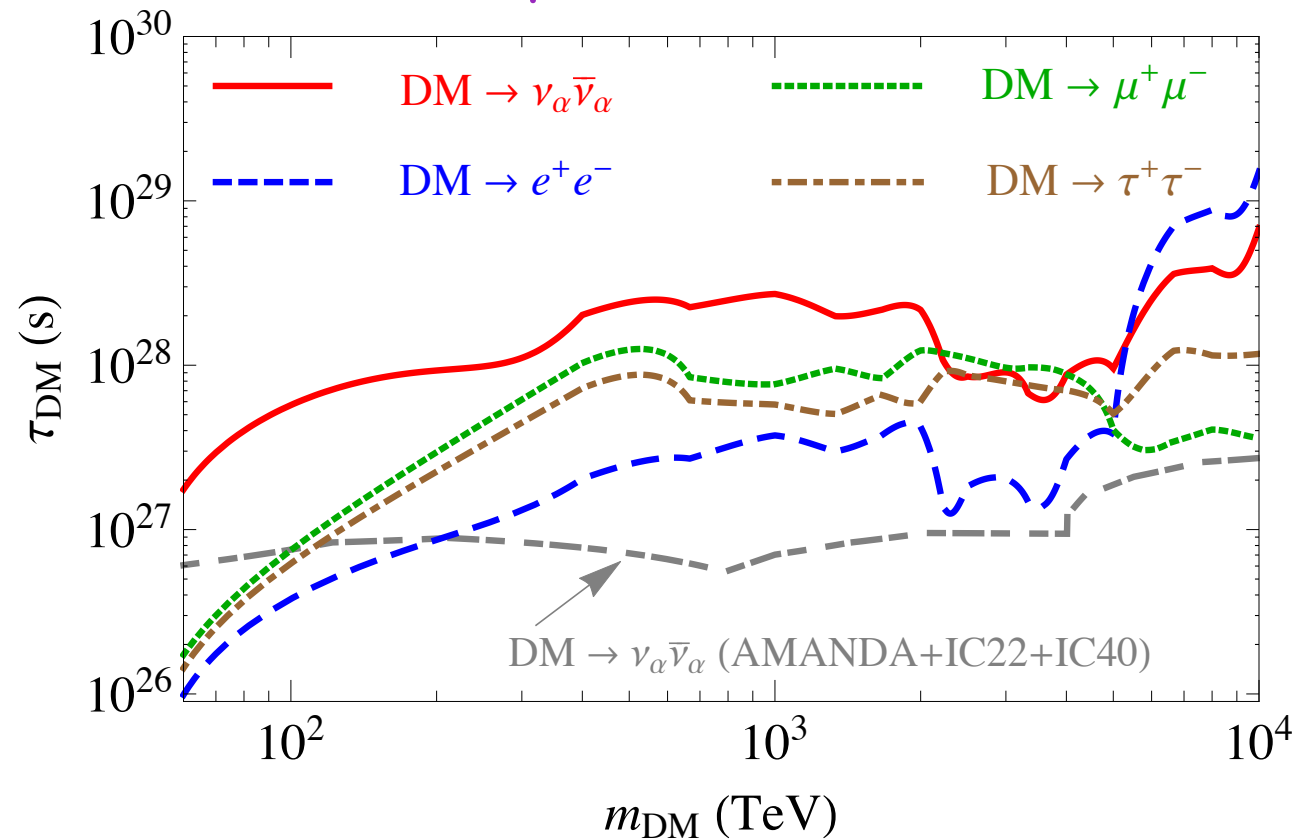
or

$$L(N_{\text{data}}^i, N) = \frac{(N)^{N_{\text{data}}^i}}{N_{\text{data}}^i!} e^{-N}$$

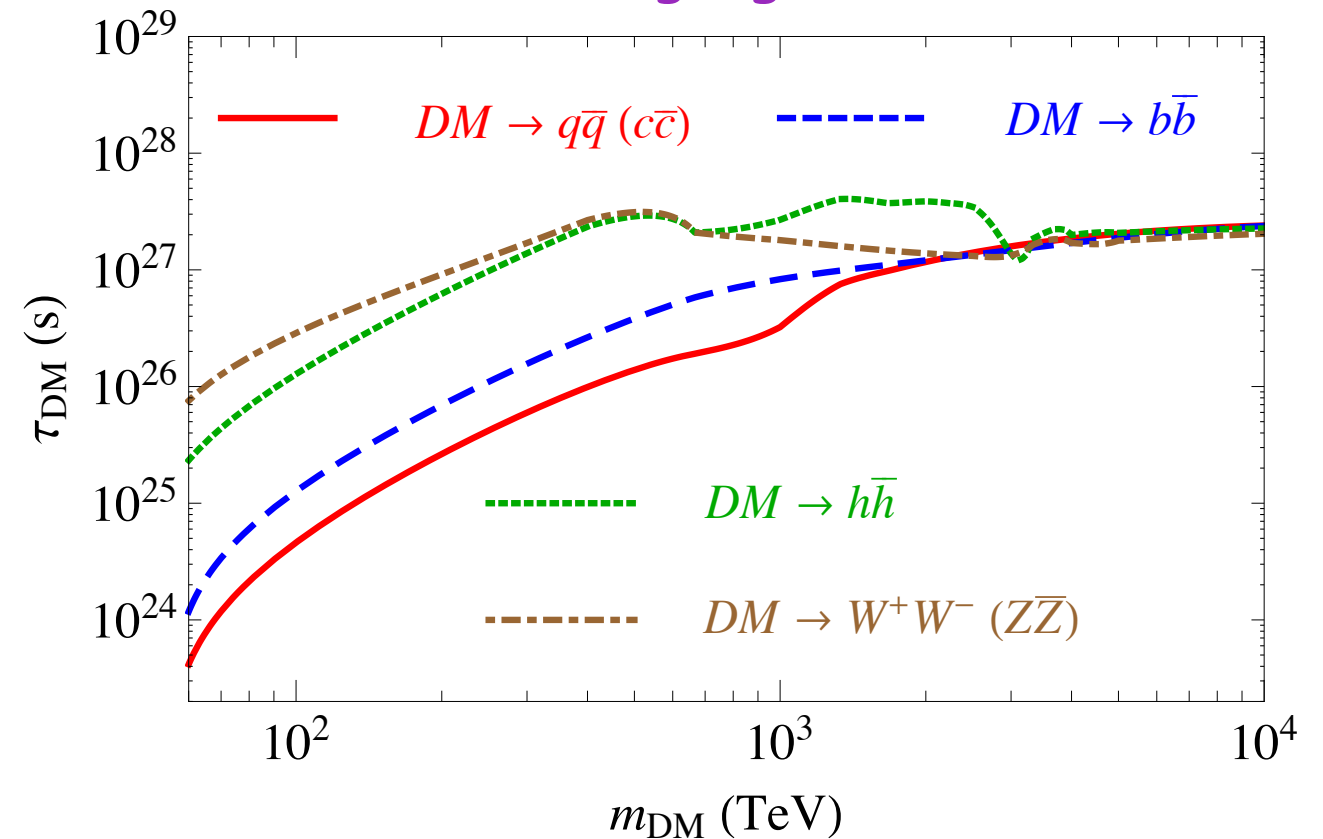
Constraining DM properties

✓ limits on DM lifetime (90% C.L.)

leptonic channels



hadronic/gauge channels

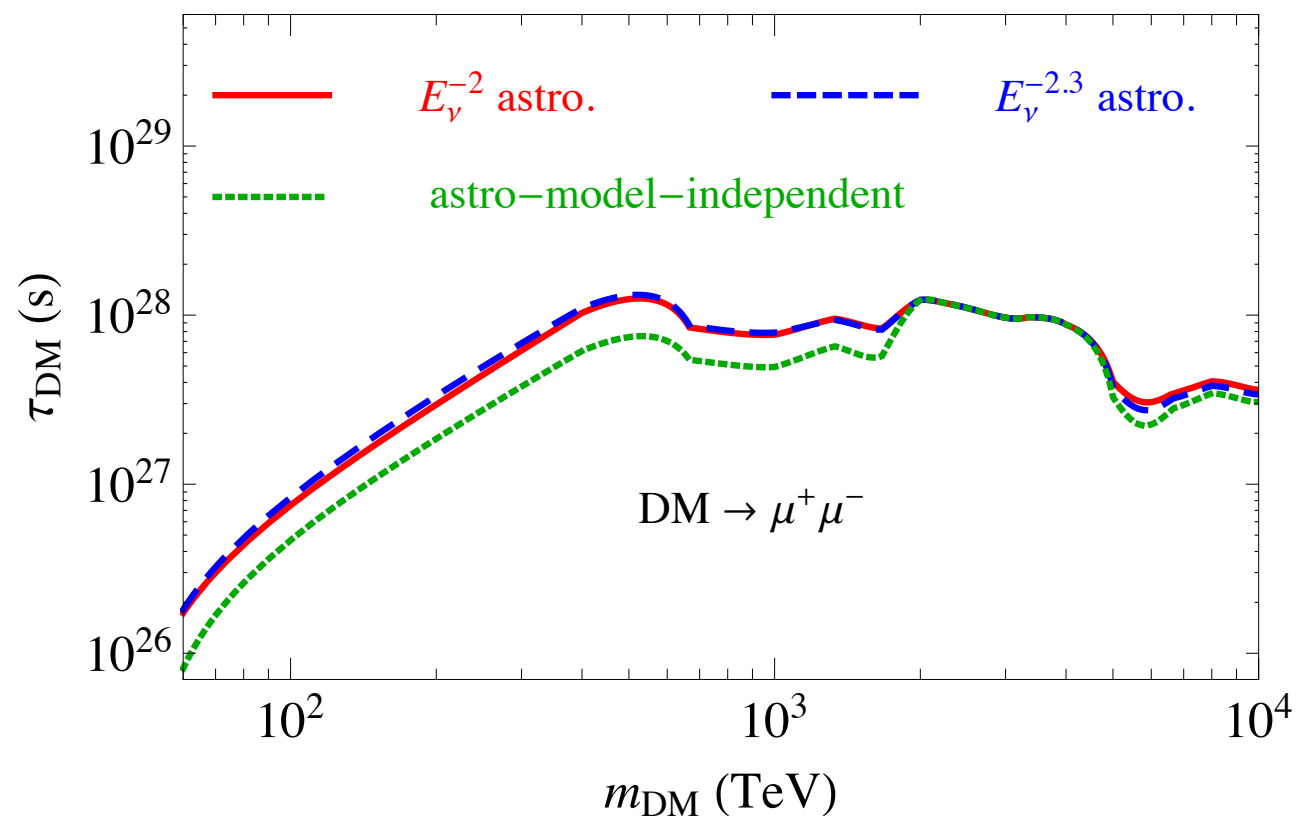
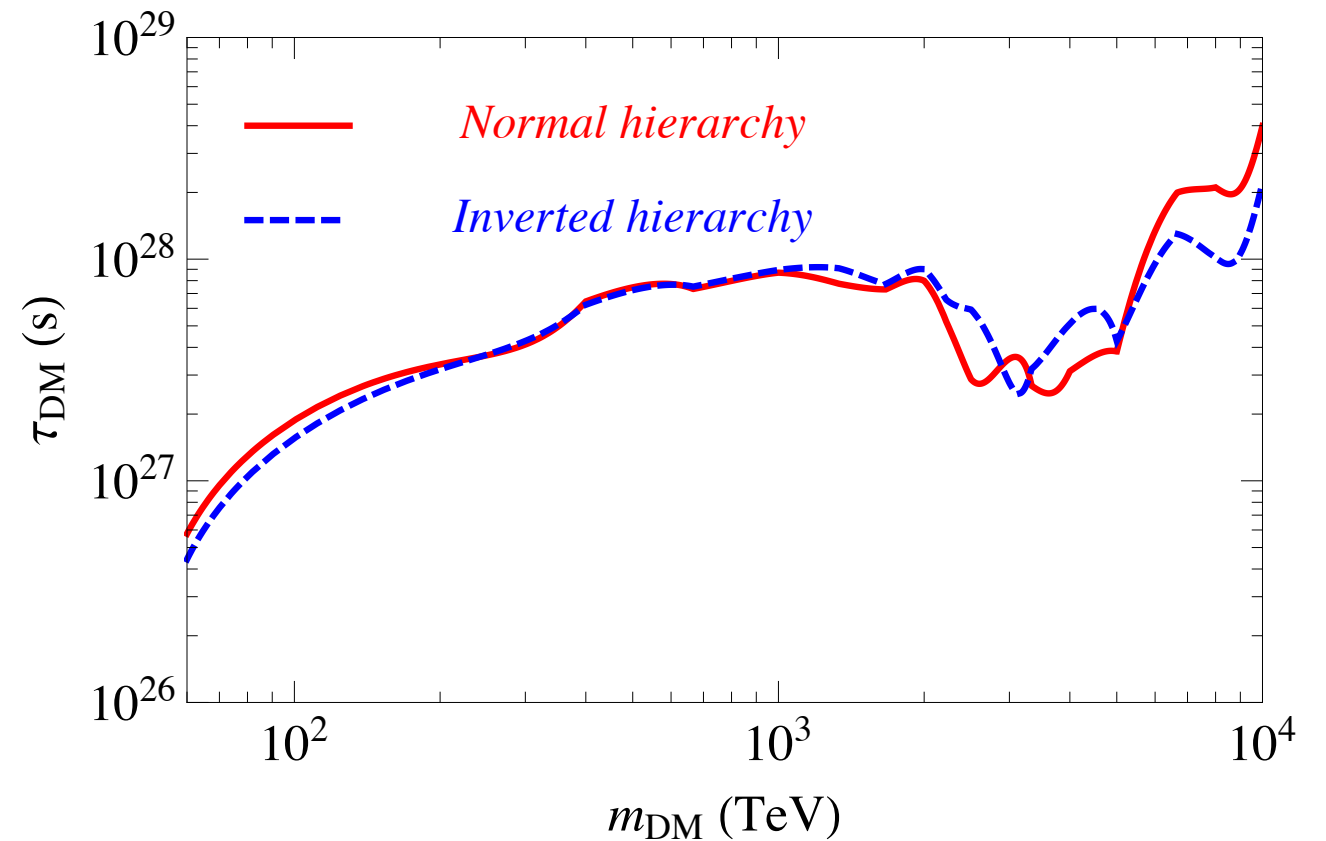


- ✓ at least one order of magnitude stronger lower limit on the DM lifetime, in the relevant DM mass range
- ✓ for a specific model, different channels should be scaled according to the corresponding branching ratios

Constraining DM properties

✓ limits on DM lifetime (90% C.L.)

NH and IH cases →



← dependence on the astro. model?

Constraining DM properties

✓ Annihilation cross section

The lower part (< 100 TeV) of the observed spectrum can be used to probe $\langle\sigma v\rangle$

Constraining DM properties

✓ Annihilation cross section

The lower part (< 100 TeV) of the observed spectrum can be used to probe $\langle\sigma v\rangle$

The isotropic components of neutrino flux from DM annihilation:

The residual isotropic flux from the Galactic halo (anti-GC direction)

$$\frac{dJ_{\text{iso}}^{\text{ann}}}{dE_\nu} = \frac{\langle\sigma v\rangle}{2} \frac{1}{4\pi m_{\text{DM}}^2} \frac{dN}{dE_\nu} (\text{l.o.s.})_{\text{anti-GC}} \quad \text{where} \quad (\text{l.o.s.})_{\text{anti-GC}} = \int_0^\infty \rho^2[r(s, b=0, l=\pi)] ds$$

Constraining DM properties

✓ Annihilation cross section

The lower part (< 100 TeV) of the observed spectrum can be used to probe $\langle\sigma v\rangle$

The isotropic components of neutrino flux from DM annihilation:

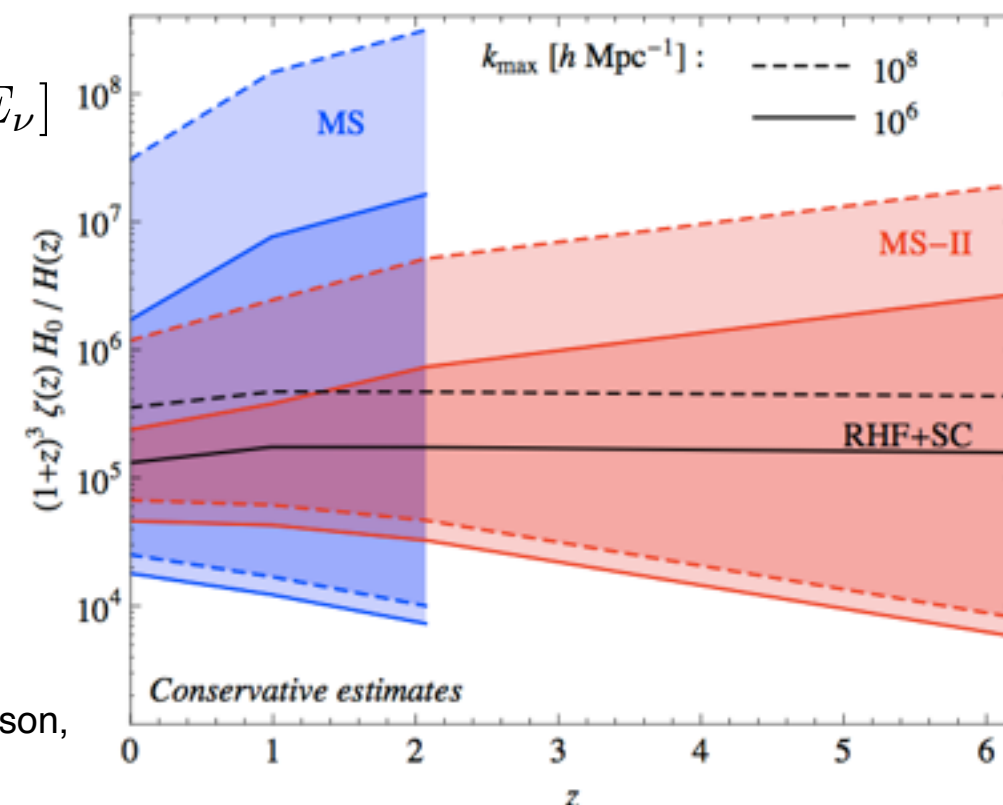
The residual isotropic flux from the Galactic halo (anti-GC direction)

$$\frac{dJ_{\text{iso}}^{\text{ann}}}{dE_\nu} = \frac{\langle\sigma v\rangle}{2} \frac{1}{4\pi m_{\text{DM}}^2} \frac{dN}{dE_\nu} (\text{l.o.s.})_{\text{anti-GC}} \quad \text{where } (\text{l.o.s.})_{\text{anti-GC}} = \int_0^\infty \rho^2[r(s, b=0, l=\pi)] ds$$

The cosmic flux from all redshift

$$\frac{dJ_{\text{cos}}^{\text{ann}}}{dE_\nu} = \frac{\langle\sigma v\rangle}{2} \frac{\Omega_{\text{DM}}^2 \rho_c^2}{4\pi m_{\text{DM}}^2} \frac{c}{H_0} \int_0^\infty \frac{(1+z)^3 \zeta(z) dz}{\sqrt{\Omega_m(1+z)^3 + \Omega_\Lambda}} \frac{dN}{dE_\nu} [(1+z) E_\nu]$$

$\zeta(z)$ flux multiplier (DM clustering)



E. Sefusatti, G. Zaharijas, P. D. Serpico, D. Theurel and M. Gustafsson,
Mon. Not. Roy. Astron. Soc. (2014) [arXiv:1401.2117].

Constraining DM properties

✓ upper limits on annihilation cross section $\langle\sigma v\rangle$ (90% C.L.)

minimum $\div$ maximum value used for $\zeta(z)$ unit of $\langle\sigma v\rangle$ is $10^{-22} \text{ cm}^3 \text{s}^{-1}$

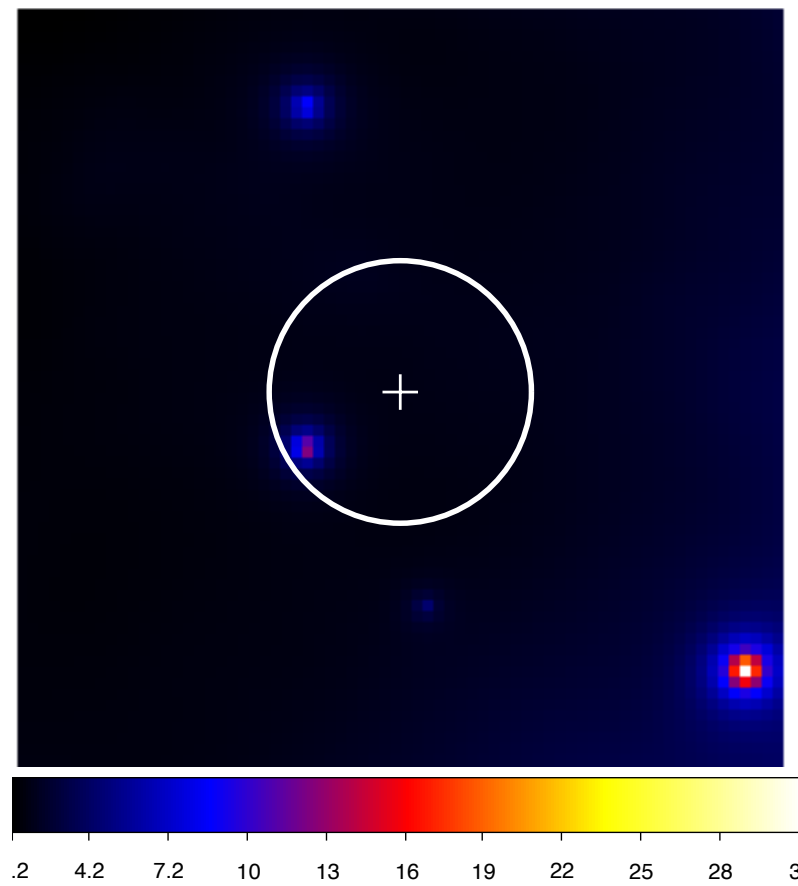
DM + DM $\rightarrow$ m_{DM}	100 TeV	50 TeV	30 TeV
$\nu_\alpha \bar{\nu}_\alpha$	1.39 $\div$ 0.22	1.21 $\div$ 0.36	2.44 $\div$ 0.88
$q\bar{q}$	489 $\div$ 84.5	1427 $\div$ 299	9934 $\div$ 4603
$b\bar{b}$	185 $\div$ 30.4	517 $\div$ 106	3514 $\div$ 1621
$c\bar{c}$	592 $\div$ 100	1708 $\div$ 348	11218 $\div$ 5215
e^+e^-	14.7 $\div$ 2.38	17.8 $\div$ 5.06	41.3 $\div$ 14.2
$\mu^+\mu^-$	4.47 $\div$ 0.65	9.06 $\div$ 1.6	23.7 $\div$ 9.23
$\tau^+\tau^-$	5.84 $\div$ 0.93	10.9 $\div$ 2.3	28.5 $\div$ 10.8
$h\bar{h}$	21.2 $\div$ 3.36	53.4 $\div$ 9.49	177 $\div$ 76.5
$Z\bar{Z}$	11.9 $\div$ 2.05	18.1 $\div$ 4.09	40.7 $\div$ 16.3
W^+W^-	14.4 $\div$ 2.4	23.7 $\div$ 4.96	54.5 $\div$ 22.3

✓ for some final states (neutrinos, charged leptons) the limit is a bit stronger than the unitary bound

IceCube / Fermi-LAT cross correlation

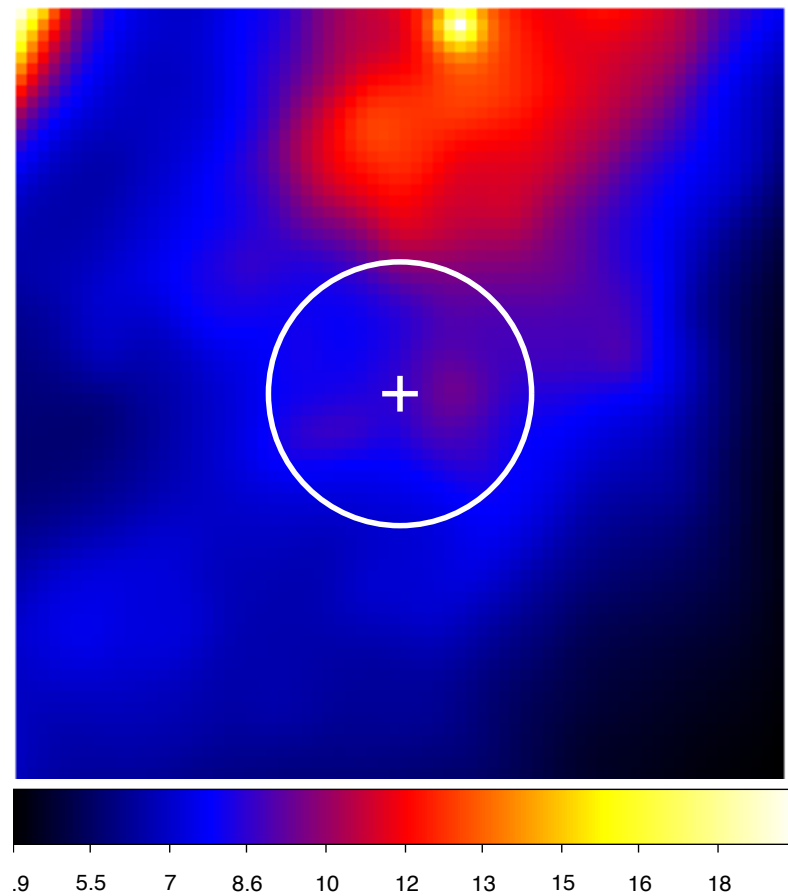
- ✓ in the 4-yr data set, 14 muon-tracks observed
- ✓ let's look at the Fermi-LAT skymap at the position of these events

3FGL J0725.8-0054



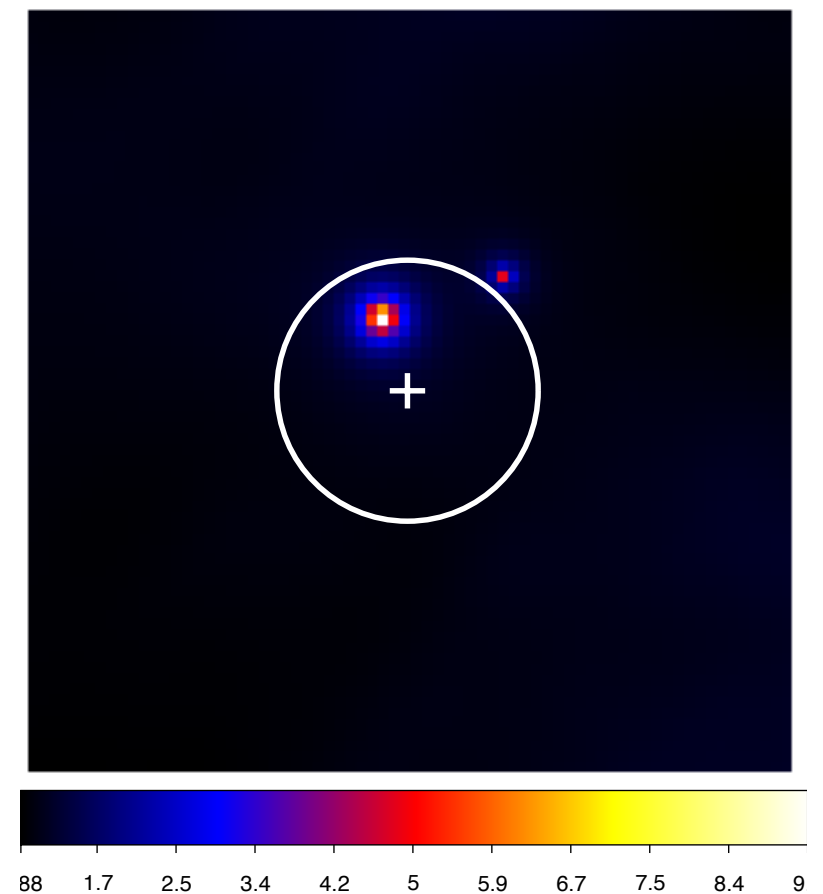
5

IC 443, SNR



38

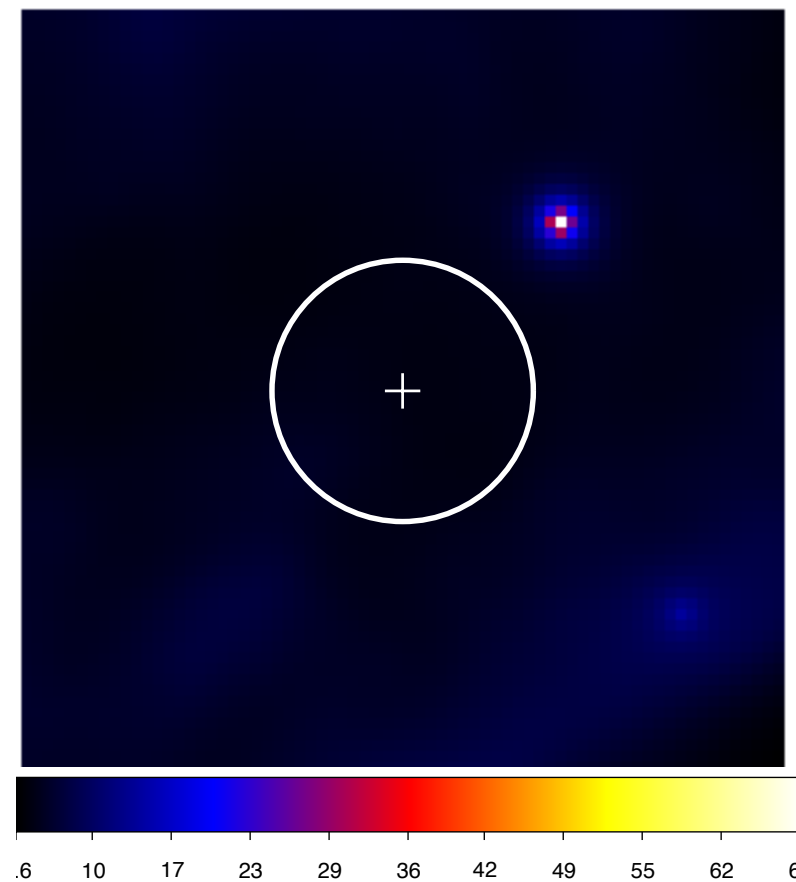
3FGL J2227.8+0040



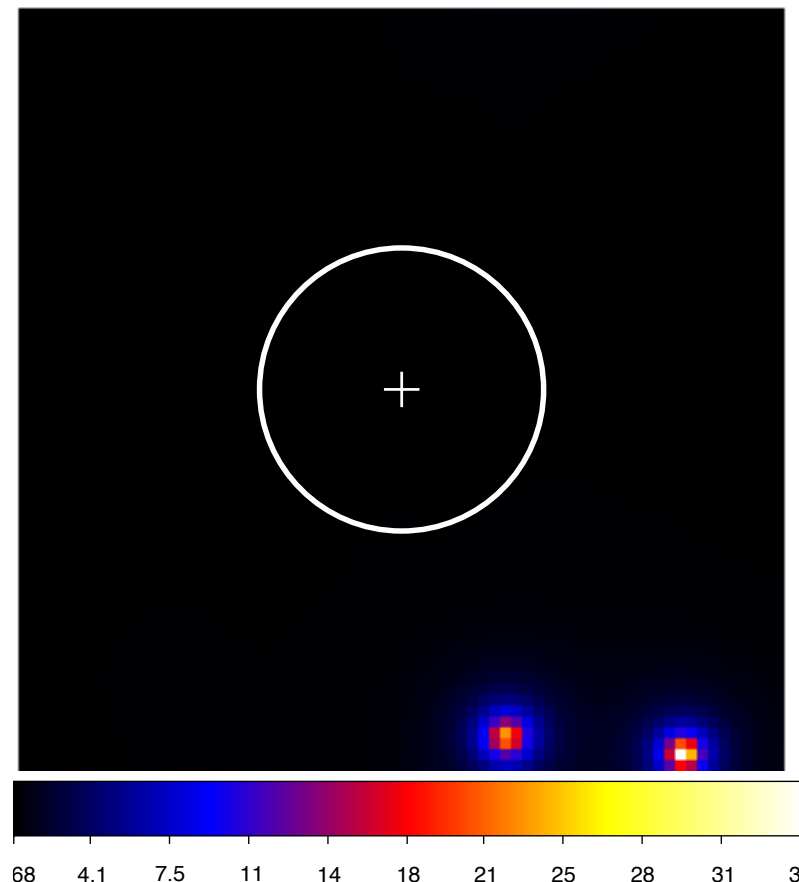
44

work in progress!!!

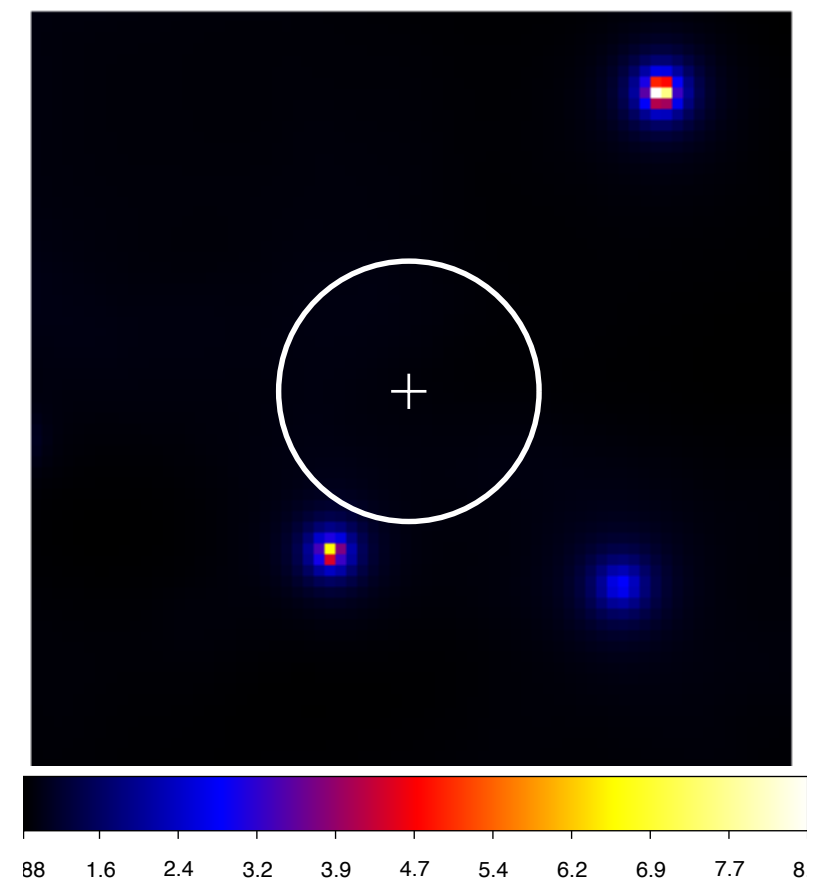
IceCube / Fermi-LAT cross correlation



13



18



47

work in progress!!!

conclusions

✓ The excess of events observed by IceCube in the energy range ~ 30 TeV - 2 PeV is an evidence for astrophysical flux or other "New Physics" induced fluxes

✓ Several features of the observed events motivate us for a DM interpretation: cut-off at ~ 2 PeV, a mild dip in the (400 - 1000) TeV and anisotropy.

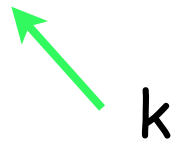
✓ We argued that a PeV-scale decaying DM, with generic decay channels, can naturally explain these features. The required lifetime is allowed by the current limits. Both the energy and angular distributions mildly prefer DM interpretation.

✓ With more statistics in the next few years, the DM interpretation of IceCube events can be tested. The gamma-ray flux expected in this scenario can be detected by the next generation of EAS detectors. Also, anisotropy measurements in the CR flux would be constraining.

Thank you !

Decaying Dark Matter

$$\Phi = k E_\nu^{-2}$$



	k	$E^{\min}-E^{\max}$ (TeV)	N_{bg}	N_{sig}	N_{limit}
AMANDA	7.4×10^{-8}	$16 - 2.5 \times 10^3$	6	7	5.4
IceCube-22	1.6×10^{-7}	$340 - 2 \times 10^5$	0.6	3	6.1
IceCube-40	3.6×10^{-8}	$2 \times 10^3 - 6.3 \times 10^6$	0.1	0	2.3
Auger	1.7×10^{-7}	$10^5 - 10^8$	0	0	2.3
ANITA	1.3×10^{-7}	$10^6 - 3.2 \times 10^{11}$	0.97	1	3.3

$$N_{\text{exp}} = T \Delta \Omega \sum_{\alpha} \left[\int_{E_{\nu}^{\min}}^{E_{\nu}^{\max}} \Phi_{\nu_{\alpha} + \bar{\nu}_{\alpha}} A_{\text{eff}}^{\alpha}(E_{\nu}) dE_{\nu} \right]$$

$$q/100 = \frac{\int_0^{N_{\text{limit}}} L(N_{\text{sig}}|N) dN}{\int_0^{\infty} L(N_{\text{sig}}|N) dN} \quad \text{where} \quad L(N_{\text{sig}}|N) = \frac{(N + N_{\text{bg}})^{N_{\text{sig}}}}{N_{\text{sig}}!} e^{-(N + N_{\text{bg}})}$$

Decaying Dark Matter

✓ extragalactic contribution:

$$H(z) = H_0 \sqrt{\Omega_\Lambda + \Omega_m(1+z)^3}$$

$$\frac{dJ_{\text{eg}}}{dE_\nu} = \frac{\Omega_{\text{DM}} \rho_c}{4\pi m_{\text{DM}} \tau_{\text{DM}}} \int_0^\infty dz \frac{1}{H(z)} \frac{dN_\nu}{dE_\nu} [(1+z)E_\nu] e^{-s_\nu(E_\nu, z)}$$

Opacity

$$s_\nu(E_\nu, z) = \begin{cases} 7.4 \times 10^{-17} (1+z)^{7/2} (E_\nu / \text{TeV}), & \text{for } 1 \ll z < z_{\text{eq}} \\ 1.7 \times 10^{-14} (1+z)^3 (E_\nu / \text{TeV}), & \text{for } z \gg z_{\text{eq}} \end{cases}$$

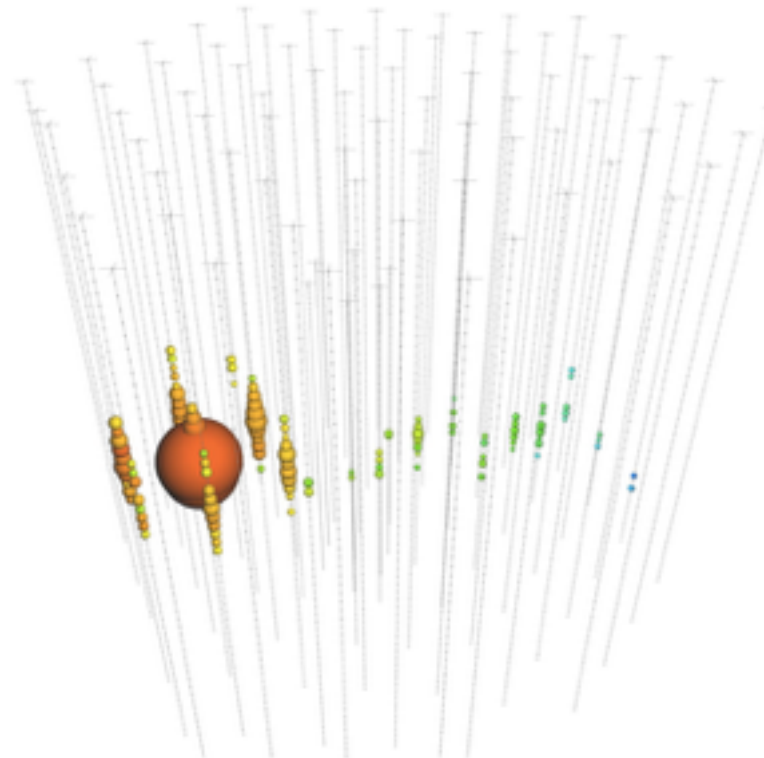
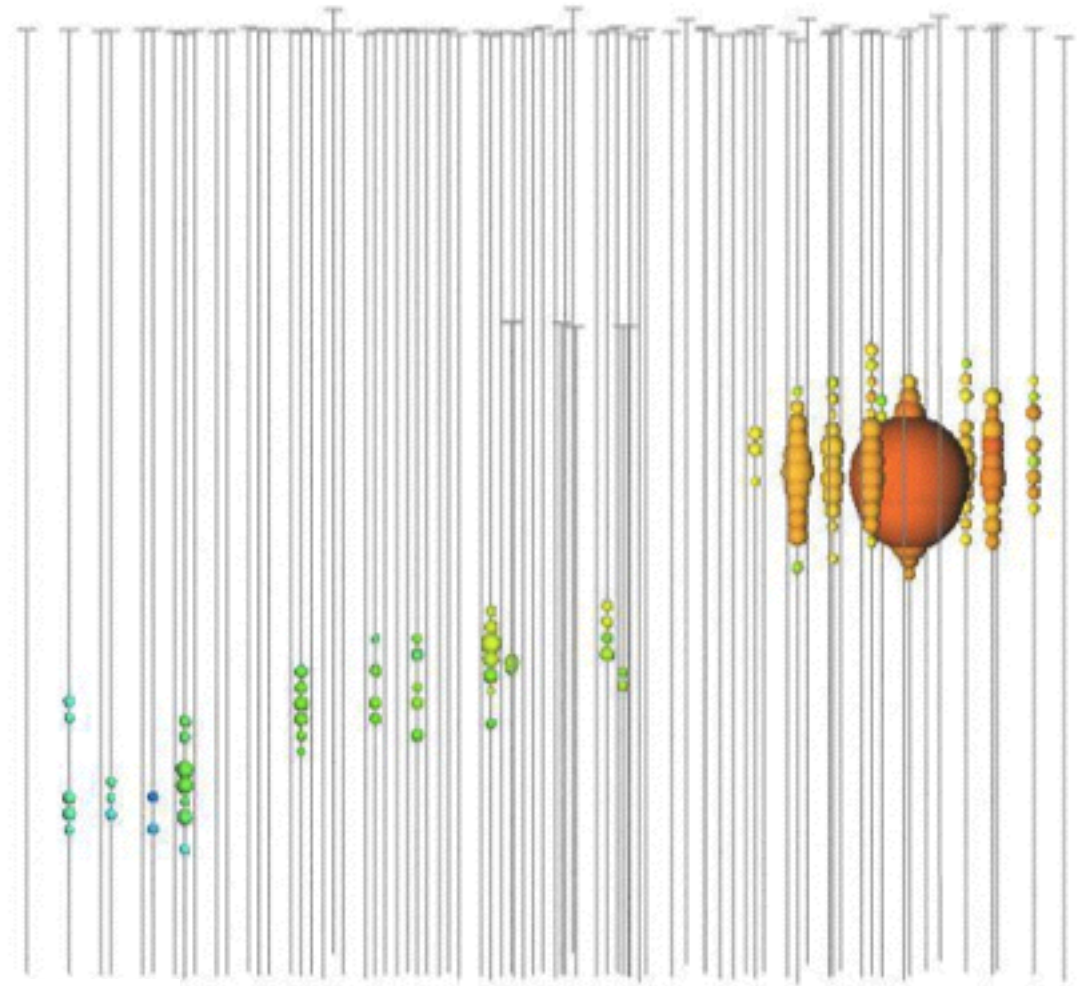
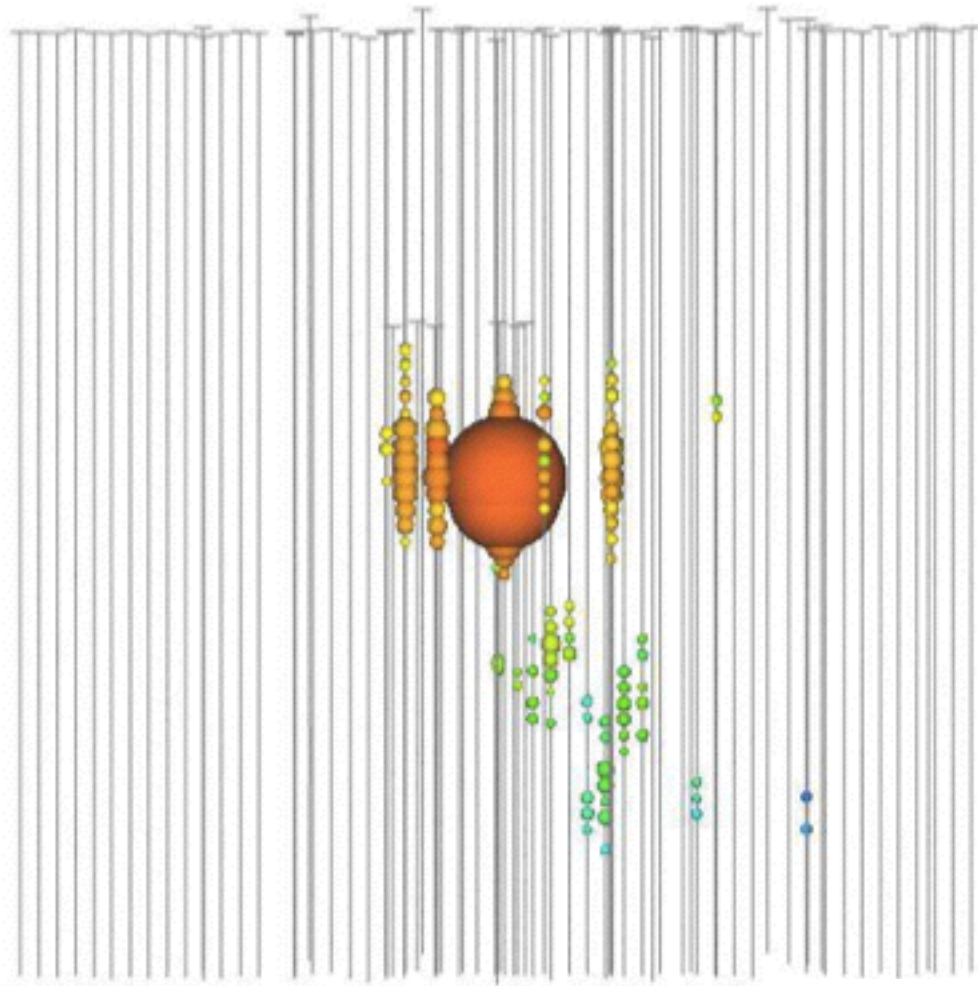
at the Earth

$$\begin{pmatrix} J_e \\ J_\mu \\ J_\tau \end{pmatrix} = \begin{pmatrix} P_{ee} & P_{e\mu} & P_{e\tau} \\ & P_{\mu\mu} & P_{\mu\tau} \\ & & P_{\tau\tau} \end{pmatrix} \begin{pmatrix} I_e \\ I_\mu \\ I_\tau \end{pmatrix}$$

production point

decoherent
oscillation

Event 8 (Track)



Gamma ray bounds

✓ Galactic component

at $\sim \text{PeV}$, the absorption length of gamma-rays
are comparable to Galactic distances



neither full absorption or cascade
development, nor full transparency

Gamma ray bounds

✓ Galactic component

at $\sim$ PeV, the absorption length of gamma-rays
are comparable to Galactic distances



neither full absorption or cascade
development, nor full transparency

energy budget consideration :

the integrated flux of γ -ray above 330 TeV must be below $1.0 \times 10^{-13} \text{ cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1}$

CASA-MIA bounds

the integrated flux of γ -ray above 775 TeV must be below $2.6 \times 10^{-14} \text{ cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1}$

Gamma ray bounds

✓ Galactic component

at $\sim$ PeV, the absorption length of gamma-rays are comparable to Galactic distances



neither full absorption or cascade development, nor full transparency

energy budget consideration :

the integrated flux of γ -ray above 330 TeV must be below $1.0 \times 10^{-13} \text{ cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1}$

CASA-MIA bounds

the integrated flux of γ -ray above 775 TeV must be below $2.6 \times 10^{-14} \text{ cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1}$

For DM decay
(NH case)
unattenuated flux



the integrated flux of γ -ray above 330 TeV is $1.2 \times 10^{-14} \text{ cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1}$

the integrated flux of γ -ray above 775 TeV is $8.8 \times 10^{-15} \text{ cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1}$

