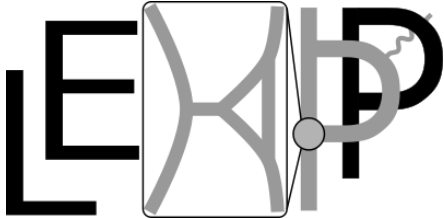
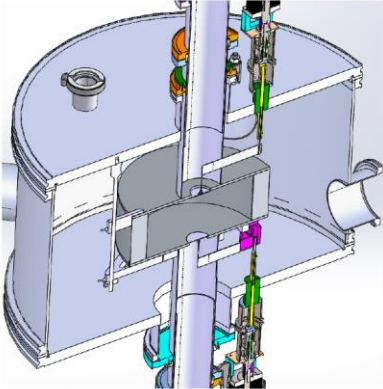


From the lowest energies to the highest: constraints on CP violation from permanent electric dipole moments



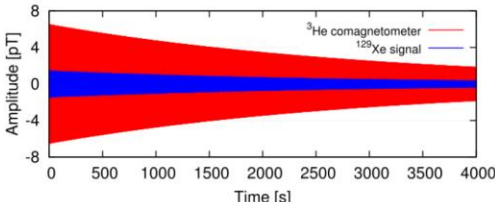
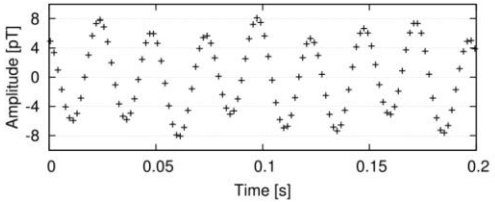
MPIK Gentner Colloquium, 30 November 2022
Skyler Degenkolb

Case 1: trapped neutron interferometry



Low-Energy Precision Physics
Physikalisches Institut
Universität Heidelberg

Case 2: nuclear spin gyroscopes



Typical(?) experiment/theory interaction



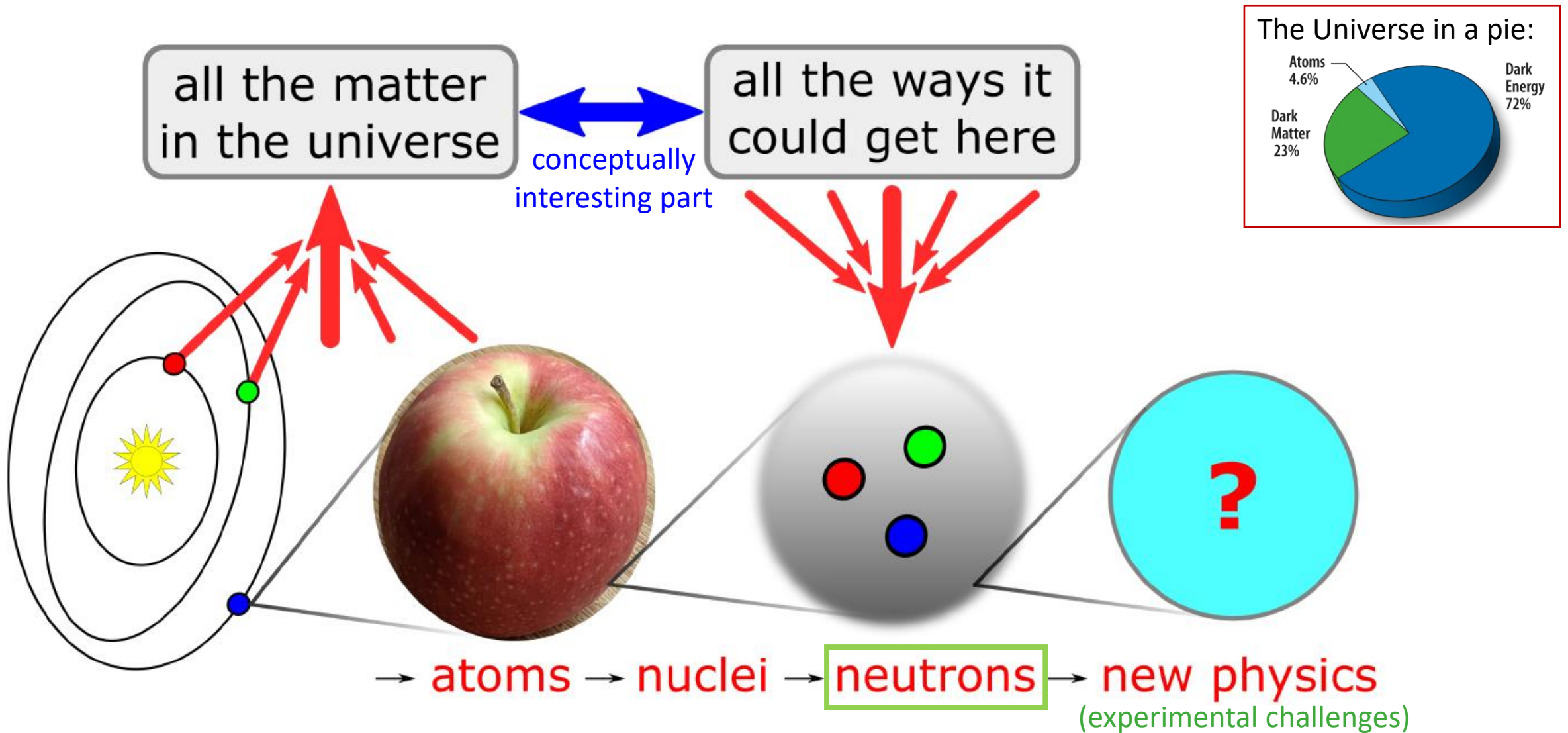
TC: "How many EDM experiments do we need?"

MJRM: "Only the one that discovers an EDM"

TC: "Oh – so now I know which one to work on...?"



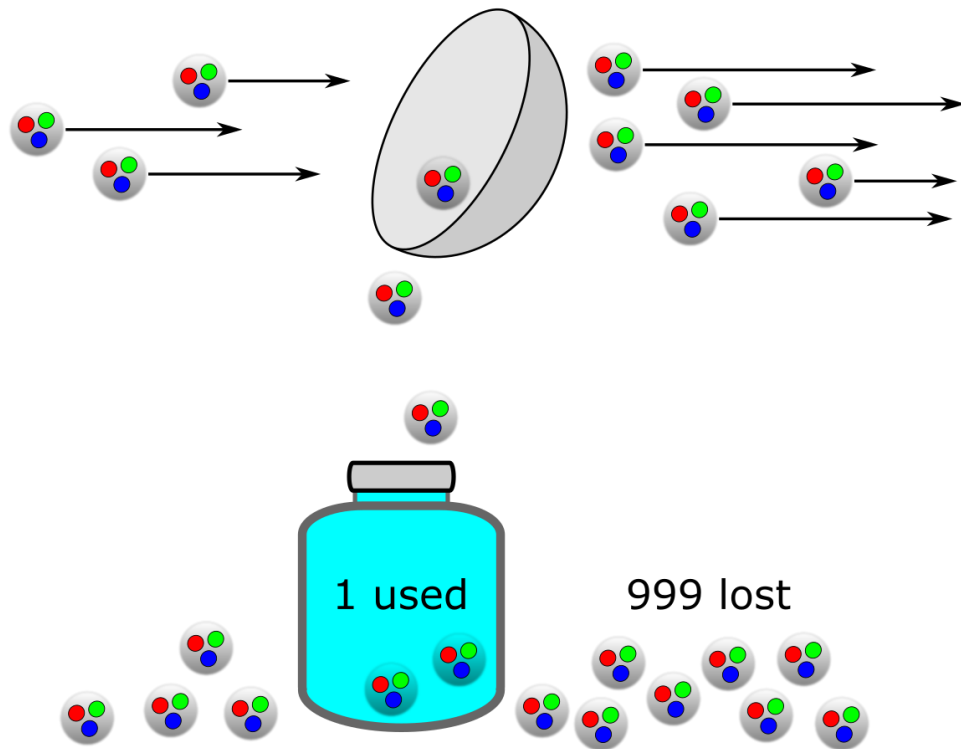
Our motivation, and a sense of scale



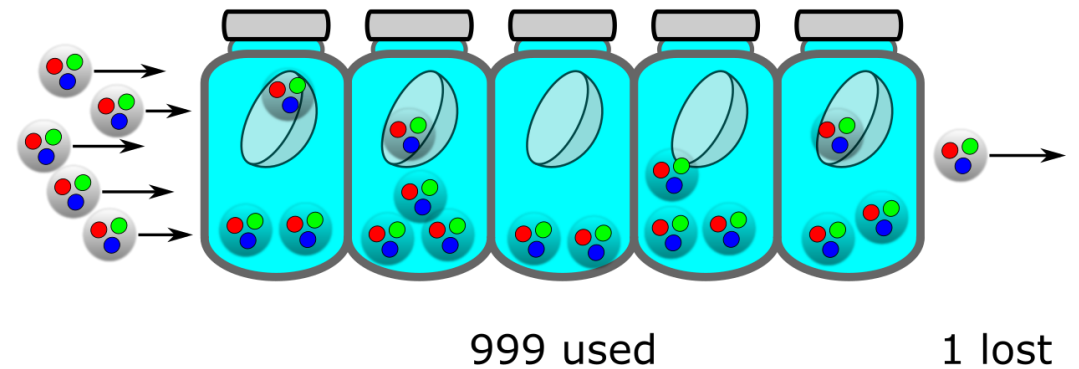
1

Challenge the first: statistics

State of the art: catch/pour
...with 0.1% success



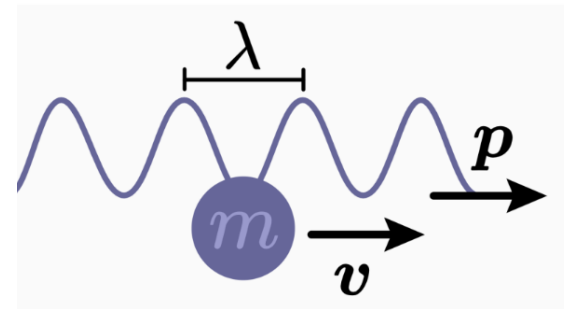
New approach: catch them
all, directly in many bottles



2 Challenge the second: observation time

“Never measure anything but frequency”
–Arthur Schawlow (1981 Physics Nobel Prize)

$$\delta\omega \sim \frac{1}{\delta t}$$

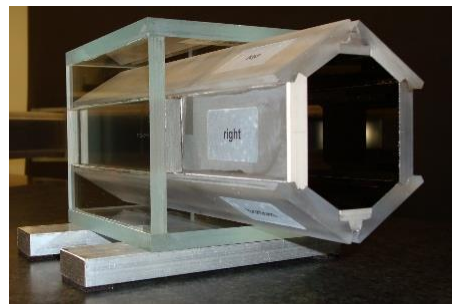


But... how to store or cool ensembles?

Wave optics, with massive particles!

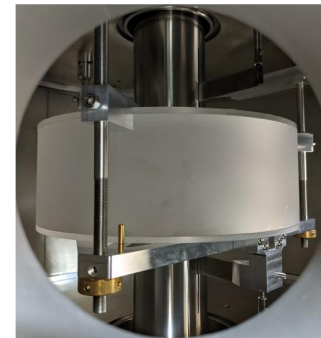
“Cold” beams: O(500 m/s)

particles fly through most experiments in milliseconds



 S-DH

S-DH



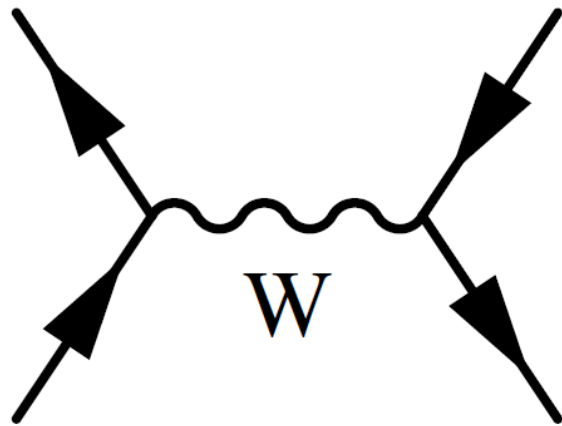
“Ultracold” traps: O(5 m/s)

particles stored for minutes ($>10^5$ ms)

3 Challenge: does this approach even make sense?

Well, suppose the scale of new physics is far above the SM...

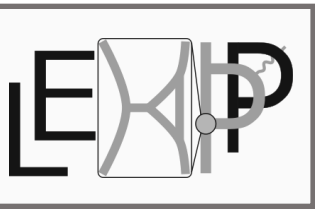
...or imagine we couldn't access the heavy gauge bosons we already know



$$\frac{-i}{k^2 - m_W^2} \left(g_{\mu\nu} - \frac{k_\mu k_\nu}{m_W^2} \right) \sim \frac{1}{m_W^2}$$

↑ “resonance”
 ↑ “high energy”
 ↑ $k \approx 0$

If the scale of new physics is \gg TeV, it looks the same whether we probe it at TeV or neV!



Orders of Magnitude

Some recent experimental EDM limits:

$$\begin{aligned} n: |d| &< 1.8 \times 10^{-26} \text{ e cm (90\% C.L.)} \\ {}^{129}\text{Xe}: |d| &< 1.4 \times 10^{-27} \text{ e cm (95\% C.L.)} \\ \text{ThO}: |d| &< 1.1 \times 10^{-29} \text{ e cm (90\% C.L.)} \end{aligned}$$

$$10^{-26} \text{ e cm} \times \frac{1 \text{ MV}}{m} \times \frac{1}{2\pi\hbar} = 24 \text{ nHz}$$

$10^{-22} \text{ eV sensitivity}$
(cf. 10^{13} eV for new physics scale)

Another example of jumping across several orders:

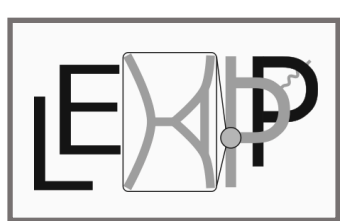
neutron capture $\rightarrow 10^{6-7} \text{ eV products}$

nEDM systematics $\rightarrow 10^{-7} \text{ eV shifts}$

Some handy conversion factors:

$$\begin{aligned} 1 \text{ neV} &= 1 \frac{\text{GeV}}{c^2} \times 1 \text{ cm} \times g \\ 1 \text{ e cm} &= 10^{13} \text{ e fm} \end{aligned}$$

... so how to get 10^{-14} in energy (or equivalent in momentum)?



“Permanent Electric Dipole Moment” = ?

Quantum eigenfrequencies:

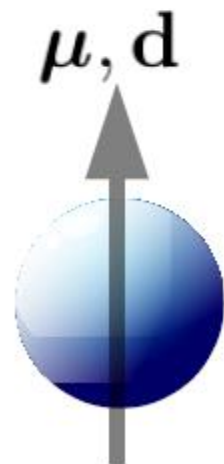
$$\hbar\omega_E \propto -d\mathbf{S} \cdot \mathbf{E}$$

$$\hbar\omega_B \propto -\mu\mathbf{S} \cdot \mathbf{B}$$

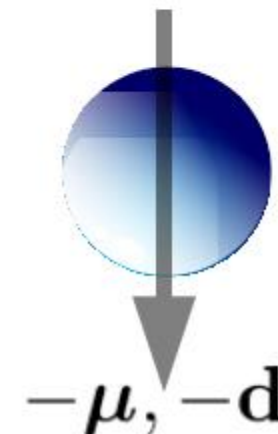
Classical moments:

$$\mathbf{d} = \int \mathbf{r}\rho(\mathbf{r})d\mathbf{r}$$

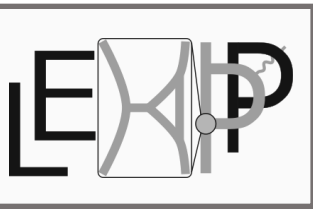
$$\boldsymbol{\mu} = \frac{1}{2} \int \mathbf{r} \times \mathbf{J}(\mathbf{r})d\mathbf{r}$$



$\pm\mathbf{S} \rightarrow \mp\mathbf{S}$



*Beware
of pictures
like this!*



“Permanent Electric Dipole Moment” = ?

Quantum eigenfrequencies:

$$\hbar\omega_E \propto -d\mathbf{S} \cdot \mathbf{E}$$

$$\hbar\omega_B \propto -\mu\mathbf{S} \cdot \mathbf{B}$$

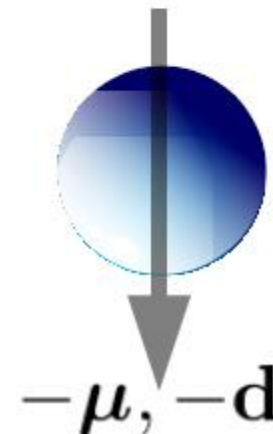
Classical moments:

$$\mathbf{d} = \int \mathbf{r}\rho(\mathbf{r})d\mathbf{r}$$

$$\boldsymbol{\mu} = \frac{1}{2} \int \mathbf{r} \times \mathbf{J}(\mathbf{r})d\mathbf{r}$$



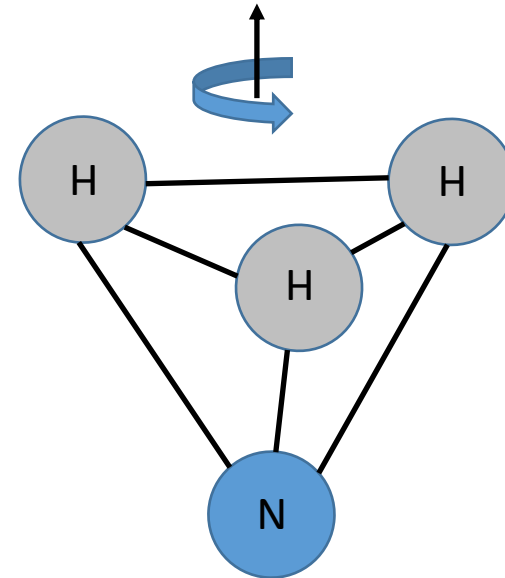
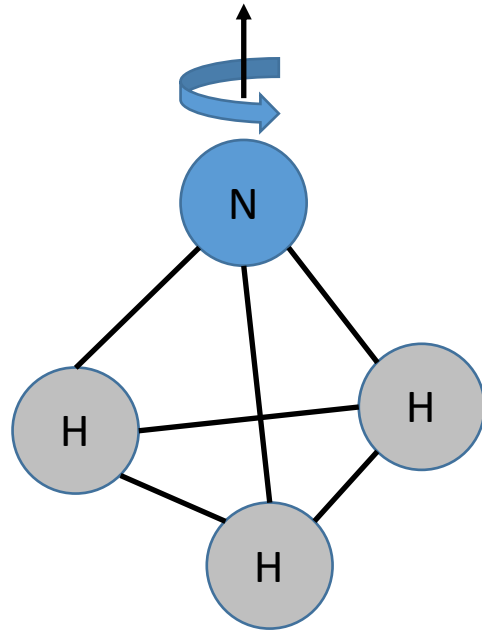
$\pm\mathbf{S} \rightarrow \mp\mathbf{S}$



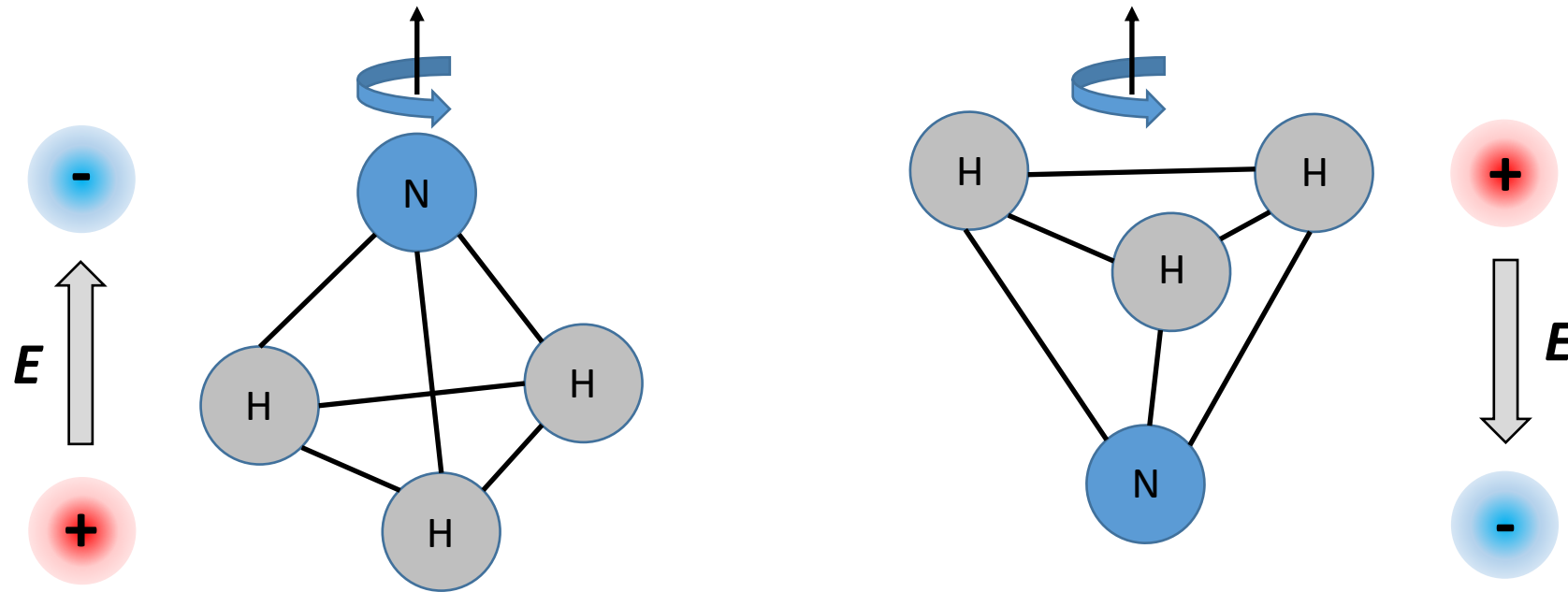
$T\mathbf{r}T^{-1} = \mathbf{r}$
$T\mathbf{p}T^{-1} = -\mathbf{p}$
$T\boldsymbol{\sigma}T^{-1} = -\boldsymbol{\sigma}$

Is it different from a molecular dipole?

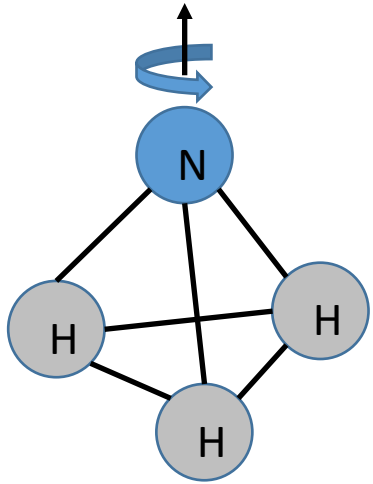
...or, "a warm-up for non-relativistic quantum methods"



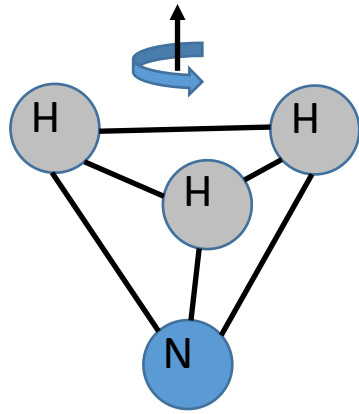
What happens in an electric field?



Find the eigenstates



$|1\rangle$

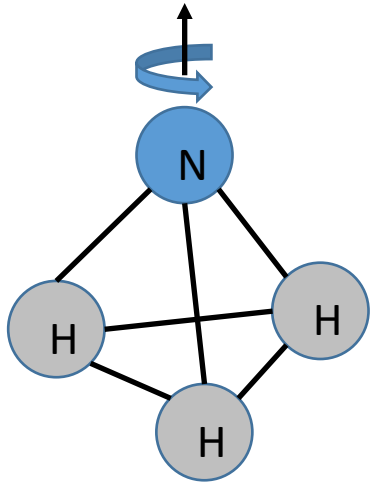


$|2\rangle$

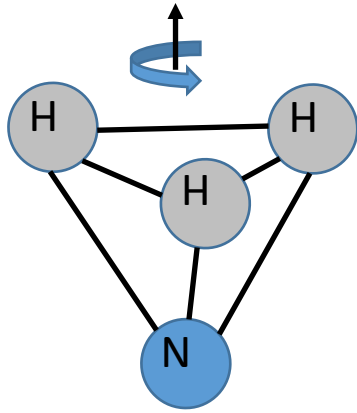
The *energy* eigenstates are:

$$\frac{1}{\sqrt{2}} (|1\rangle \pm |2\rangle)$$

Check the limiting cases



$|1\rangle$



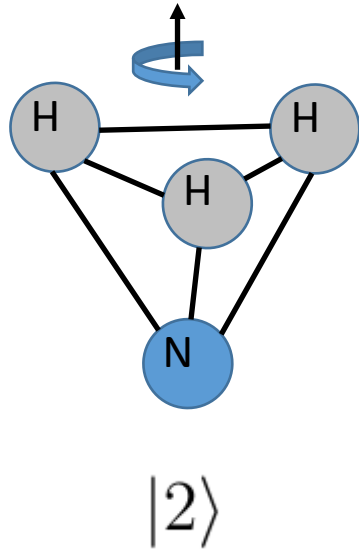
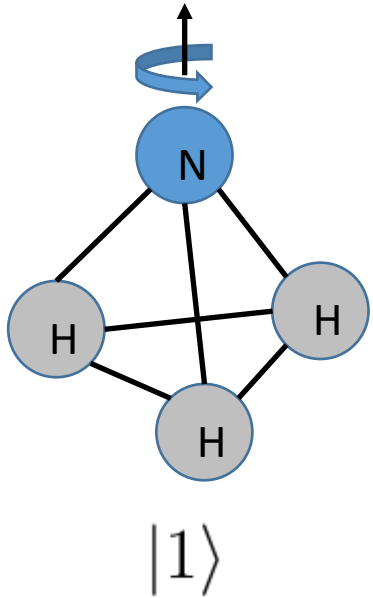
$|2\rangle$

$$E_{\pm} = E_0 + \sqrt{A^2 + d^2 E^2}$$

The *energy* eigenstates are:

$$\frac{1}{\sqrt{2}} (|1\rangle \pm |2\rangle)$$

No surprises, actually



$$E_{\pm} = E_0 + \sqrt{A^2 + d^2 E^2}$$

↙ ↘

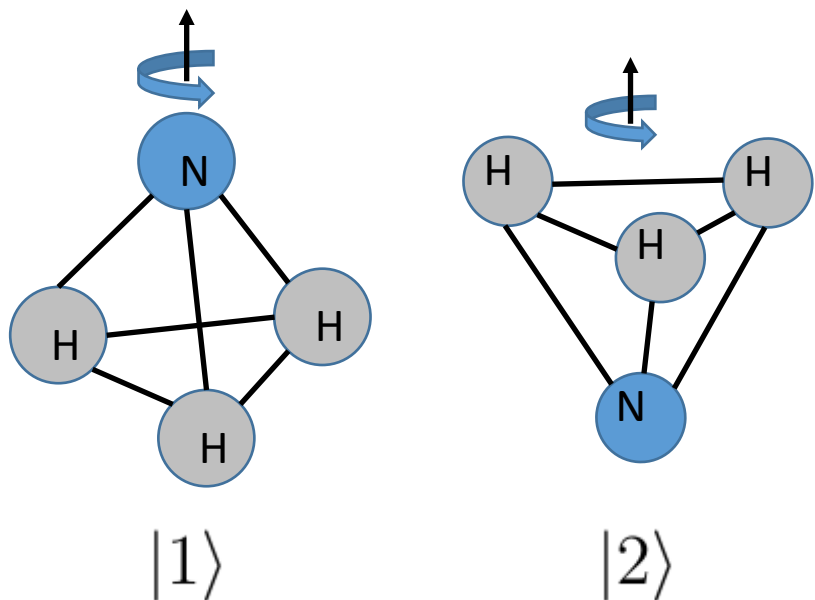
$$E_0 \pm dE \qquad E_0 \pm A \pm \frac{d^2 E^2}{2A}$$

$(dE \gg A)$ $(dE \ll A)$

The *energy* eigenstates are:

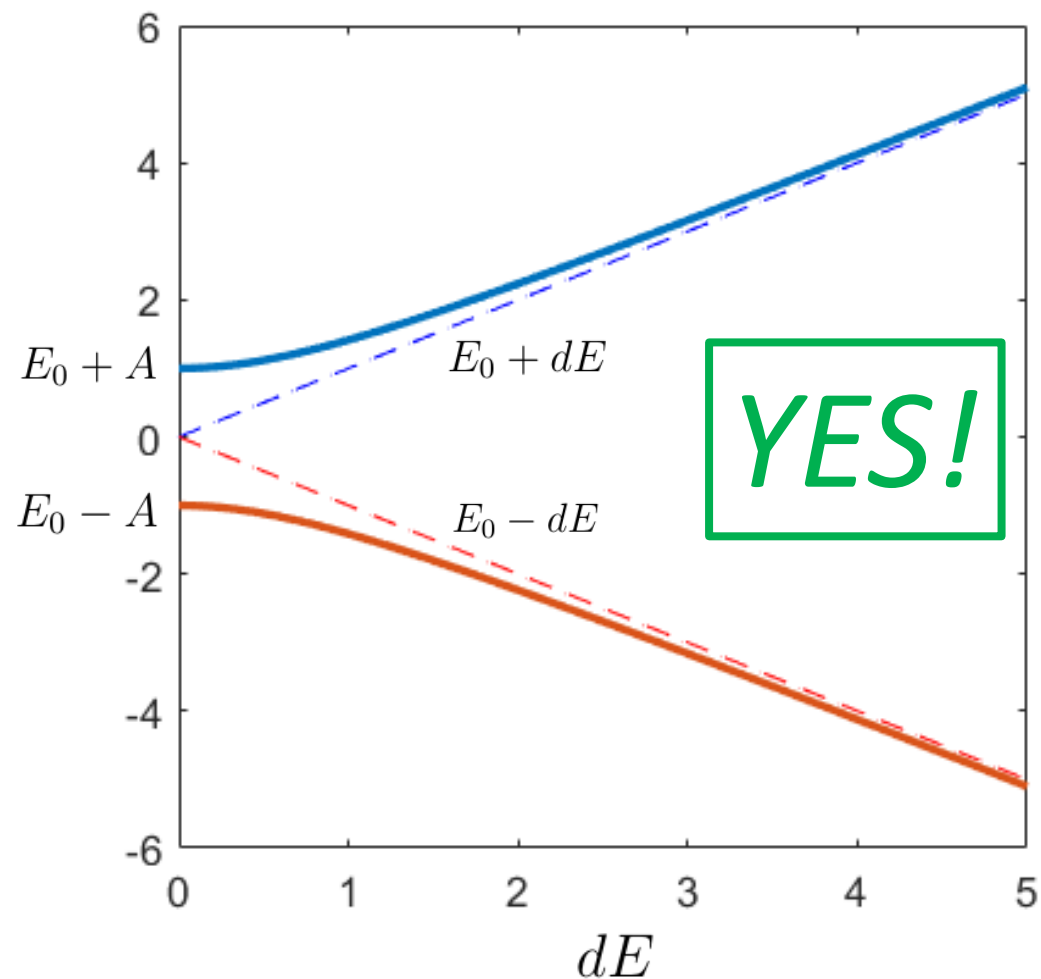
$$\frac{1}{\sqrt{2}} (|1\rangle \pm |2\rangle)$$

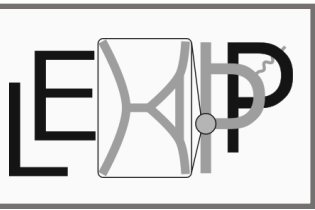
So... is it different from a molecular dipole?



The *energy* eigenstates are:

$$\frac{1}{\sqrt{2}} (|1\rangle \pm |2\rangle)$$





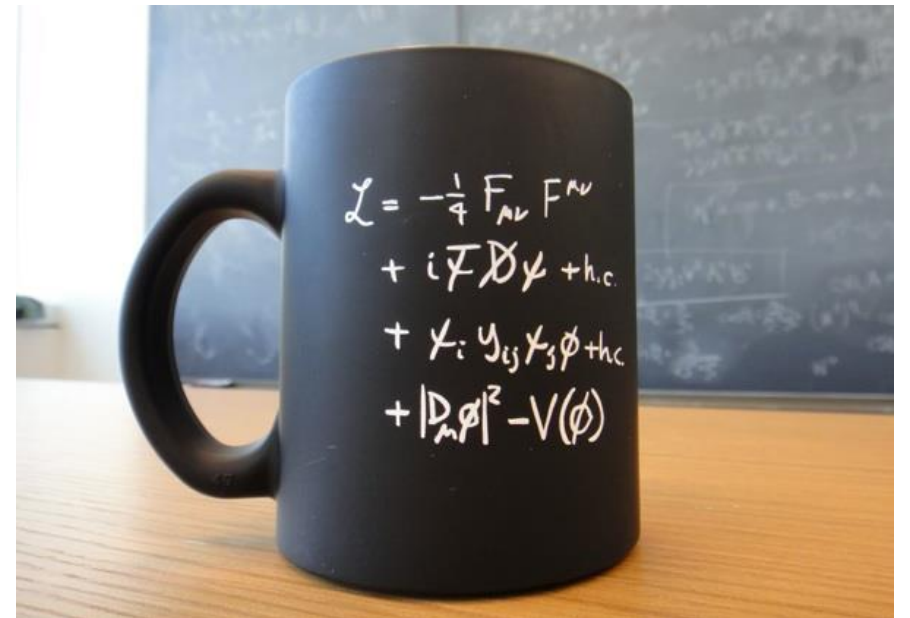
New Physics, in Familiar Terms

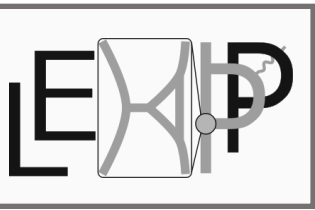
$$\mathcal{L}_{\text{fermion}} = -\frac{\mu}{2}\bar{\psi}\sigma^{\mu\nu}F_{\mu\nu}\psi - i\frac{d}{2}\bar{\psi}\sigma^{\mu\nu}\gamma^5 F_{\mu\nu}\psi$$

↓
MDM

↓
EDM

- Non-conservation of P and T already apparent in EDM term
- Consistency with zero vs. consistency with SM





A Taxonomy of Form Factors*

*which are not just for composite particles!

$$\mathcal{L}_{\text{fermion}} = -\frac{\mu}{2}\bar{\psi}\sigma^{\mu\nu}F_{\mu\nu}\psi - i\frac{d}{2}\bar{\psi}\sigma^{\mu\nu}\gamma^5F_{\mu\nu}\psi$$

↓
MDM

↓
EDM

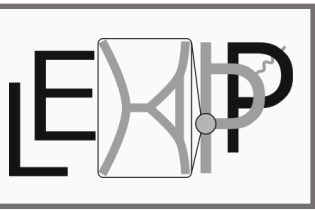
$$\begin{aligned}\langle p_f | j^\mu | p_i \rangle = & \bar{u}(p_f) \left[F_1(q^2)\gamma^\mu \right. \\ & + \frac{i\sigma^{\mu\nu}}{2m}q_\nu F_2(q^2) \\ & + i\epsilon^{\mu\nu\rho\sigma}\sigma_{\rho\sigma}q_\nu F_3(q^2) \\ & \left. + \frac{1}{2m}\left(q^\mu - \frac{q^2}{2m}\gamma^\mu\right)\gamma_5 F_4(q^2) \right] u(p_i)\end{aligned}$$

$$d = -\frac{F_3(0)}{2m}$$

$$Q = F_1(0)$$

$$\mu = \frac{F_1(0) + F_2(0)}{2m}$$

$$a = F_4(0)$$



A Taxonomy of Form Factors

$$\mathcal{L}_{\text{fermion}} = \underbrace{-\frac{\mu}{2} \bar{\psi} \sigma^{\mu\nu} F_{\mu\nu} \psi}_{\text{MDM}} \underbrace{- i \frac{d}{2} \bar{\psi} \sigma^{\mu\nu} \gamma^5 F_{\mu\nu} \psi}_{\text{EDM}}$$

$$\langle p_f | j^\mu | p_i \rangle = \bar{u}(p_f) \left[\begin{array}{l} F_1(q^2) \gamma^\mu \\ + \frac{i \sigma^{\mu\nu}}{2m} q_\nu F_2(q^2) \\ + i \epsilon^{\mu\nu\rho\sigma} \sigma_{\rho\sigma} q_\nu F_3(q^2) \\ + \frac{1}{2m} \left(q^\mu - \frac{q^2}{2m} \gamma^\mu \right) \gamma_5 F_4(q^2) \end{array} \right] u(p_i)$$

$$d = -\frac{F_3(0)}{2m}$$

$$\mu = \frac{F_1(0) + F_2(0)}{2m}$$

$$Q = F_1(0)$$

$$a = F_4(0)$$

Summary of Motivation

Problem: 1 extra baryon in every 10^9

- Asymmetry: $\eta = \frac{n_B - n_{\bar{B}}}{n_\gamma}$
- n_γ comes from CMB decoupling
- Actually normalize to entropy density s , since universe expands

Requirements: Sakharov's criteria

- Baryon-number (B) violation
- C and CP -violation
- Departure from thermal equilibrium

Solution: ???

- No complex antimatter nuclei
- No annihilation fronts
- No adequate symmetry-breaking

Prediction:

- New CP -violating physics
- Coupling to Standard Model baryons
- Polarization of bound states

Dimensional Analysis

Naïve estimate for generic new physics:

$$d_n \propto \frac{m_q}{\Lambda^2} \cdot e \cdot \phi_{\text{CPV}}$$

Current experiments: $10^{-26} e \text{ cm} \rightarrow \Lambda \sim 10 - 100 \text{ TeV}$

Analysis: neutron in Global context

Naïve estimate for generic new physics:

$$d_n \propto \frac{m_q}{\Lambda^2} \cdot e \cdot \phi_{\text{CPV}}$$

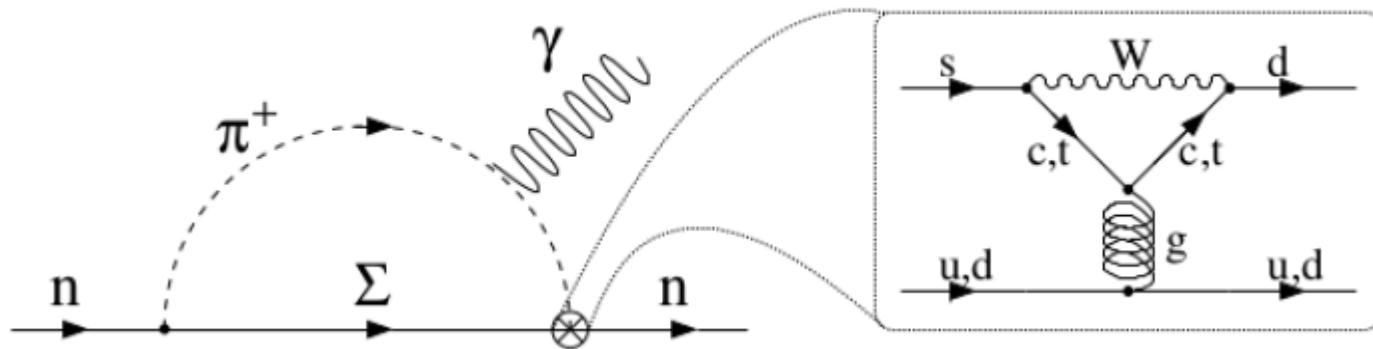
Current experiments: $10^{-26} e \text{ cm}$ $\rightarrow \Lambda \sim 10 - 100 \text{ TeV}$

Standard Model CKM: $10^{-32} e \text{ cm}$

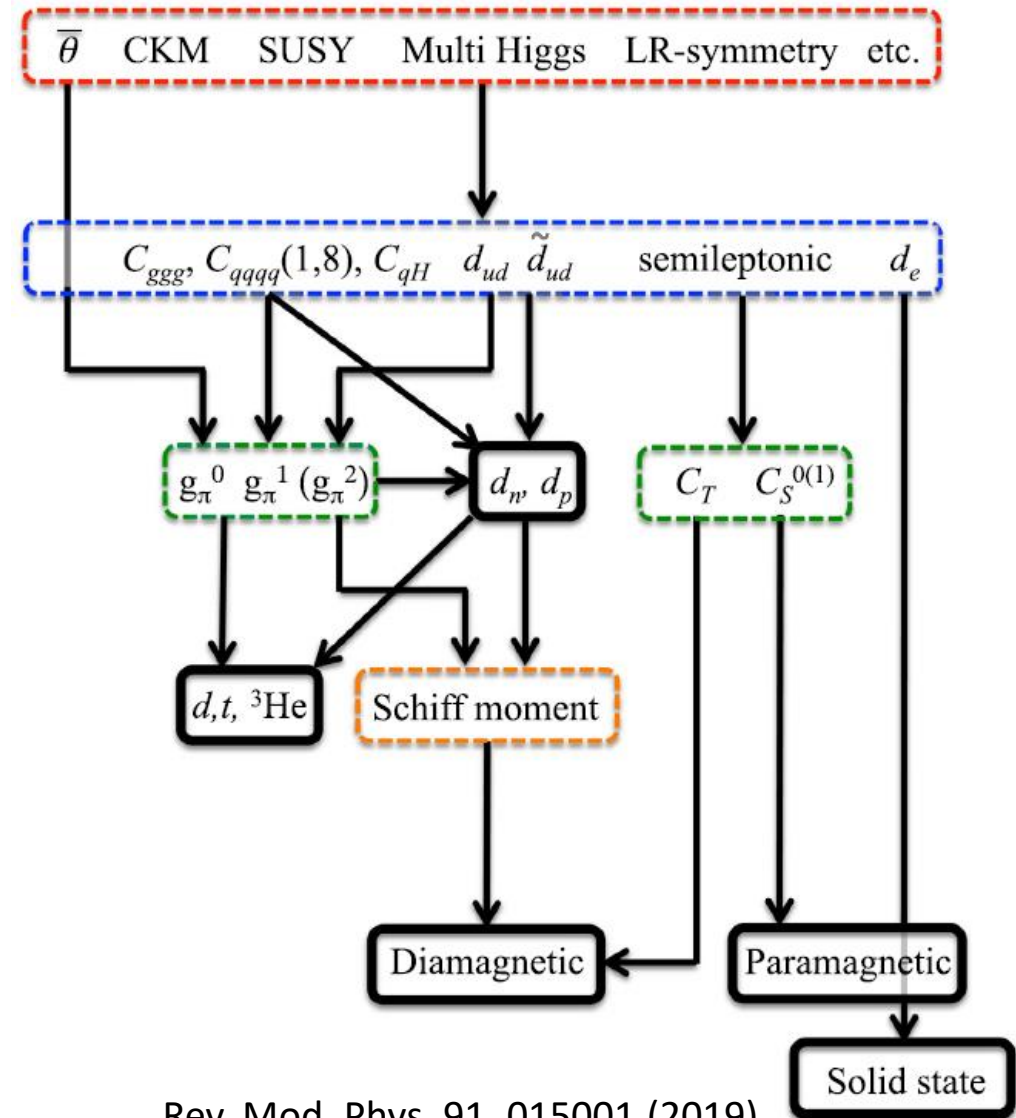
Standard Model QCD: ???

$$\rightarrow d_n \approx (10^{-16} e \text{ cm}) \bar{\theta}$$

Neutron EDM from CP-violating pion couplings:



Pospelov & Ritz, *Annals of Physics* 318 (2005): 119-169



Rev. Mod. Phys. 91, 015001 (2019)

Analysis: neutron in Global context

Define a matrix a_{ij} according to $d_i = \sum_j a_{ij} C_j$, e.g.,

$$d_n \approx \bar{d}_n^{\text{SR}} + 1.6 \times 10^{-14} \bar{g}_\pi^{(0)} - 8.6 \times 10^{-16} \bar{g}_\pi^{(1)} + 1.5 \times 10^{-18} C_T^{(0)}$$

for global analysis at the atomic/nuclear level.

Lattice calculations would* also give us some control at the hadronic level:

$$d_n^{\text{SR}} = g_T^{(n,u)} d_u + g_T^{(n,d)} d_d + g_T^{(n,s)} d_s$$

$$-(0.55 \pm 0.28) e \tilde{d}_u - (1.1 \pm 0.55) e \tilde{d}_d$$

$$+ \text{Weinberg} + 4\text{-fermion}$$

FLAG 2021, 5-10% for u, d

*QCD sum rules

*Naïve dim. analysis

“Global analysis” (hadronic/nuclear)

values: *Rev. Mod. Phys.* **91**, 015001 (2019)

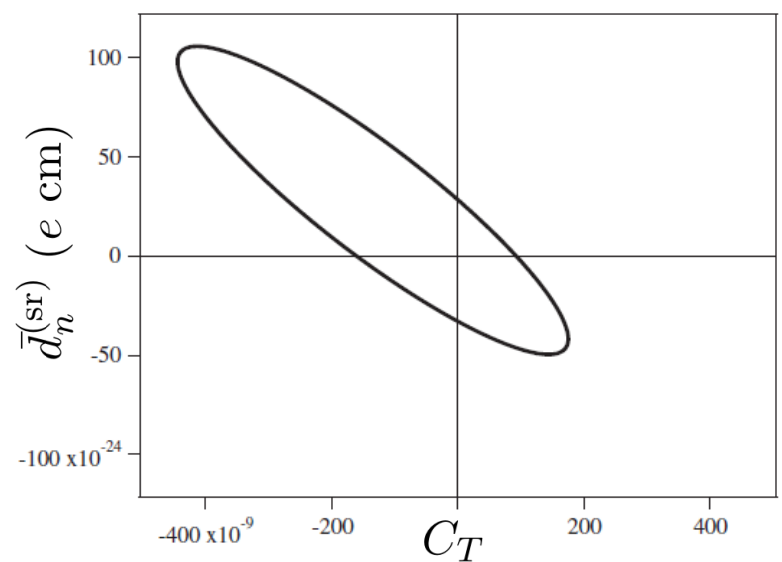
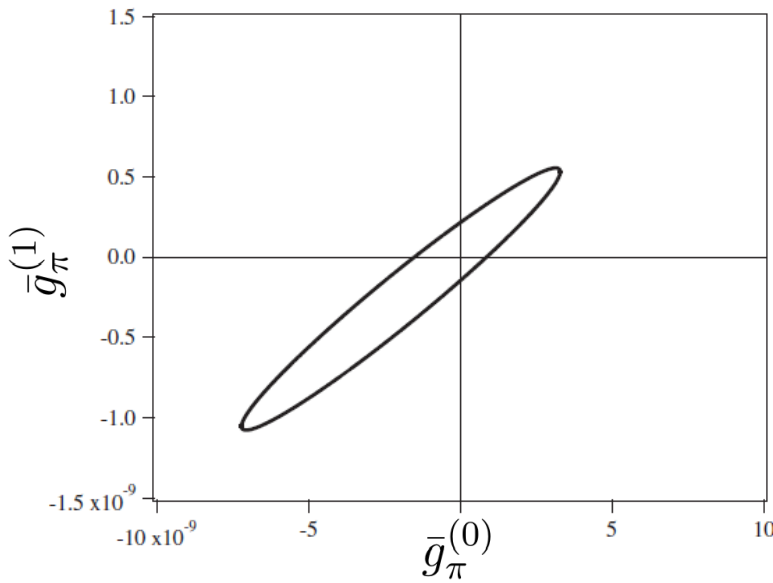
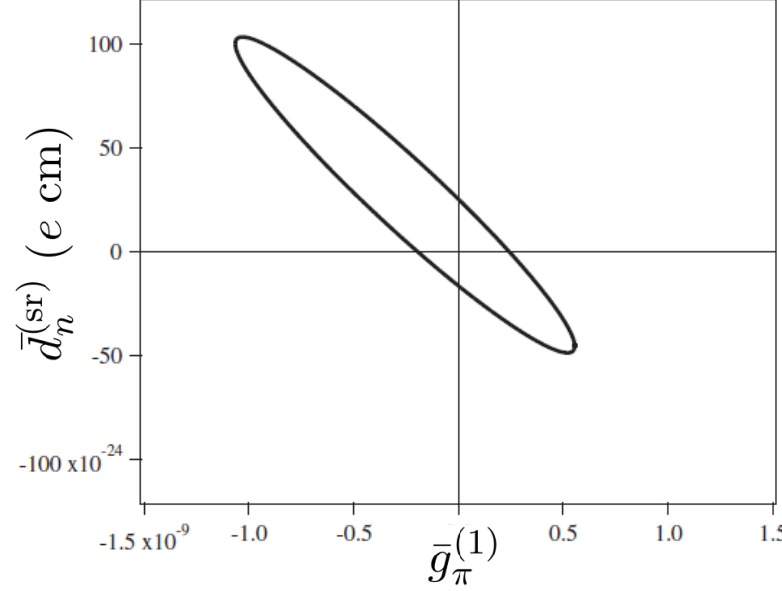
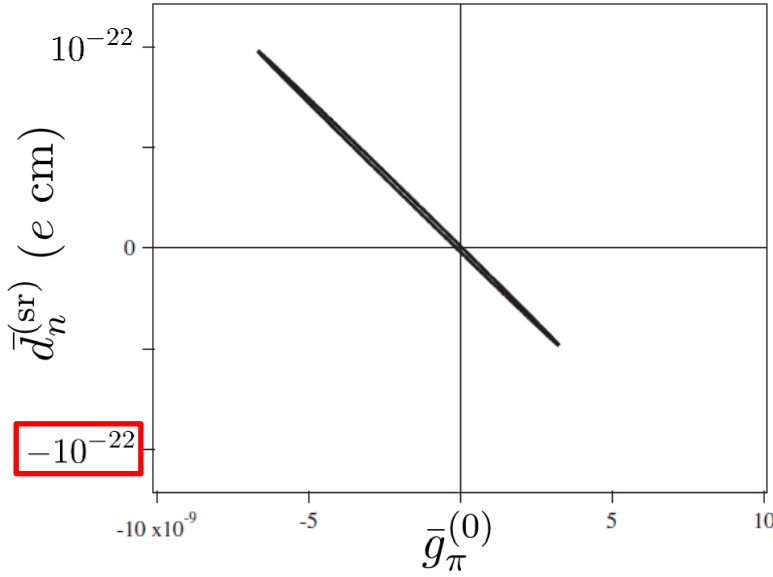
Define a matrix a_{ij} according to $d_i = \sum_j a_{ij} C_j$,

	d_n	d_{Xe}	d_{Hg}	d_{Ra}	
$\alpha_{ij} =$	1.0	1.6×10^{-14}	-8.6×10^{-16}	1.5×10^{-18}	$\bar{d}_n^{(sr)}$
	1.9×10^{-5}	-8.6×10^{-21}	-2.1×10^{-19}	-6.1×10^{-21}	$\bar{g}_\pi^{(0)}$
	-5.7×10^{-4}	-1.3×10^{-17}	1.9×10^{-17}	3.1×10^{-20}	$\bar{g}_\pi^{(1)}$
	1.2×10^{-2}	2.2×10^{-15}	-8.1×10^{-15}	-8.4×10^{-19}	$C_T^{(0)}$

...and invert it:

$$\begin{pmatrix} \bar{d}_n^{(sr)} \\ \bar{g}_\pi^{(0)} \\ \bar{g}_\pi^{(1)} \\ C_T^{(0)} \end{pmatrix} = \begin{pmatrix} 1.0 & 1.6 \times 10^{-14} & -8.6 \times 10^{-16} & 1.5 \times 10^{-18} \\ 1.9 \times 10^{-5} & -8.6 \times 10^{-21} & -2.1 \times 10^{-19} & -6.1 \times 10^{-21} \\ -5.7 \times 10^{-4} & -1.3 \times 10^{-17} & 1.9 \times 10^{-17} & 3.1 \times 10^{-20} \\ 1.2 \times 10^{-2} & 2.2 \times 10^{-15} & -8.1 \times 10^{-15} & -8.4 \times 10^{-19} \end{pmatrix} \begin{pmatrix} d_n \\ d_{Xe} \\ d_{Hg} \\ d_{Ra} \end{pmatrix}$$

Status from 2019: (hadronic/nuclear)

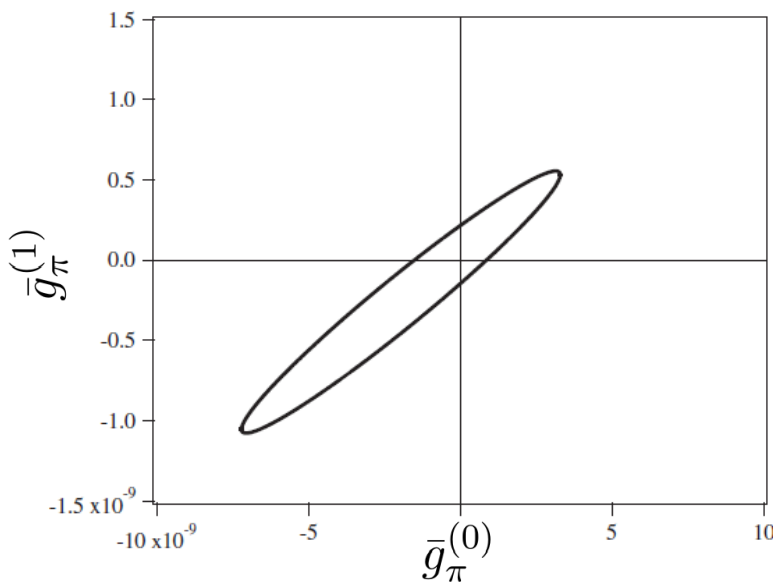
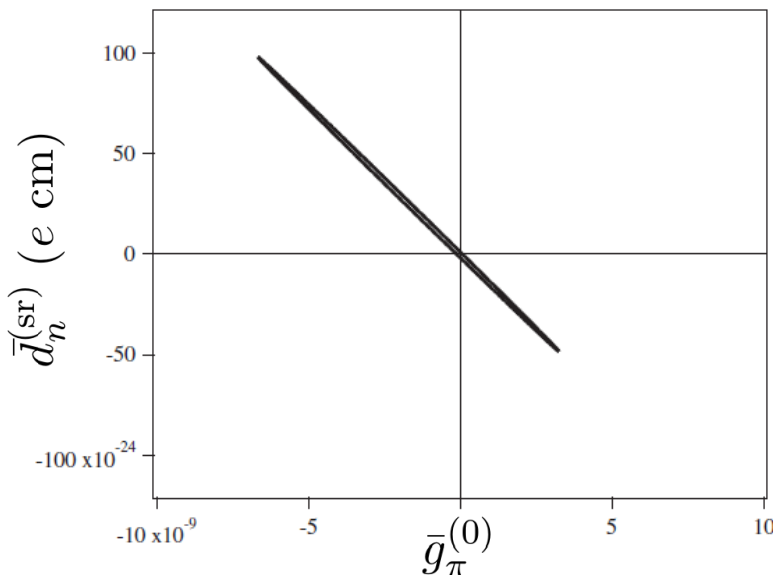


Global Analysis: T. Chupp, M. Ramsey-Musolf
Rev. Mod. Phys. **91**, 015001 (2019)
Phys. Rev. C **91**, 035502 (2015)

“Sole source” limits:

LE parameter	System	95% u.l.
d_e	ThO	$9.2 \times 10^{-29} e \text{ cm}$
C_S	ThO	8.6×10^{-9}
C_T	^{199}Hg	3.6×10^{-10}
$\bar{g}_\pi^{(0)}$	^{199}Hg	3.8×10^{-12}
$\bar{g}_\pi^{(1)}$	^{199}Hg	3.8×10^{-13}
$\bar{g}_\pi^{(2)}$	^{199}Hg	2.6×10^{-11}
\bar{d}_n^{sr}	Neutron	$3.3 \times 10^{-26} e \text{ cm}$
\bar{d}_p^{sr}	TiF	$8.7 \times 10^{-23} e \text{ cm}$
\bar{d}_p^{sr}	^{199}Hg	$2.0 \times 10^{-25} e \text{ cm}$
Other parameters		
d_d	$\approx 3/4 d_n$	$2.5 \times 10^{-26} e \text{ cm}$
$\bar{\theta}$	$\approx \bar{g}_\pi^{(0)} / (0.015)$	2.5×10^{-10}
$\tilde{d}_d - \tilde{d}_u$	$5 \times 10^{-15} \bar{g}_\pi^{(1)} e \text{ cm}$	$2 \times 10^{-27} e \text{ cm}$

Updates now in progress...



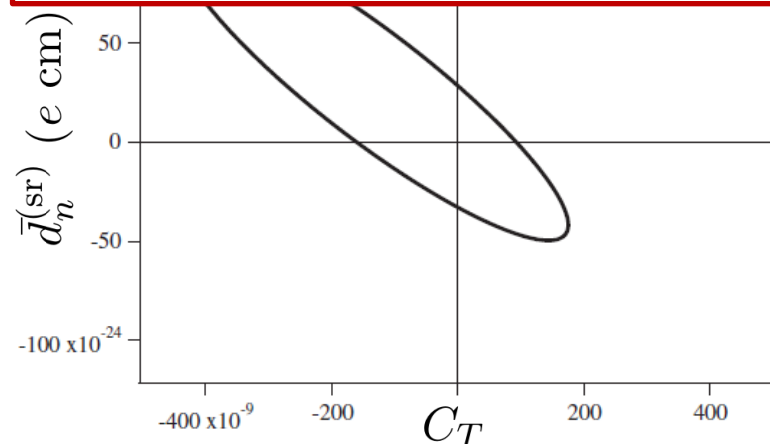
Since then:

- n : $|d| < 1.8 \times 10^{-26}$ e cm (90% C.L.) PSI: Phys. Rev. Lett. 124, 081803 (2020)
- ^{129}Xe : $|d| < 1.4 \times 10^{-27}$ e cm (95% C.L.) HeXe: Phys. Rev. Lett. **123**, 143003 (2019)
- ThO: $|d| < 1.1 \times 10^{-29}$ e cm (90% C.L.) ACME: Nature **562**, 355–360 (2018)

“Sole source” limits:

To match constraining power of ^{199}Hg ($|d| < 7.4 \times 10^{-30}$ e cm, 95% C.L.)

- ^{129}Xe needs factor 100 (i.e., approaching 10^{-30} e cm)
- ^{225}Ra needs factor 500 (i.e., approaching 10^{-26} e cm)
- ...or new species: $^{223/221}\text{Rn}$, ^{227}Ac even with worse sensitivity!



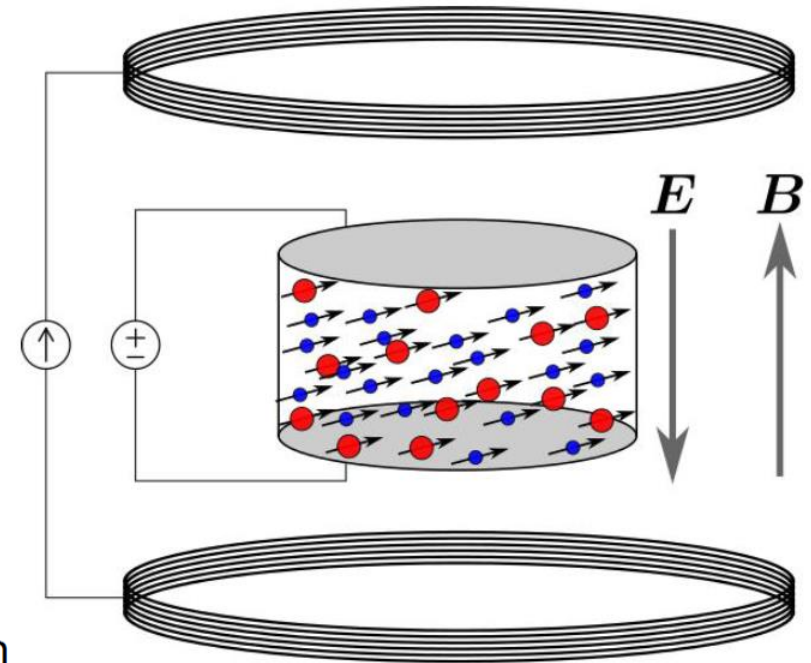
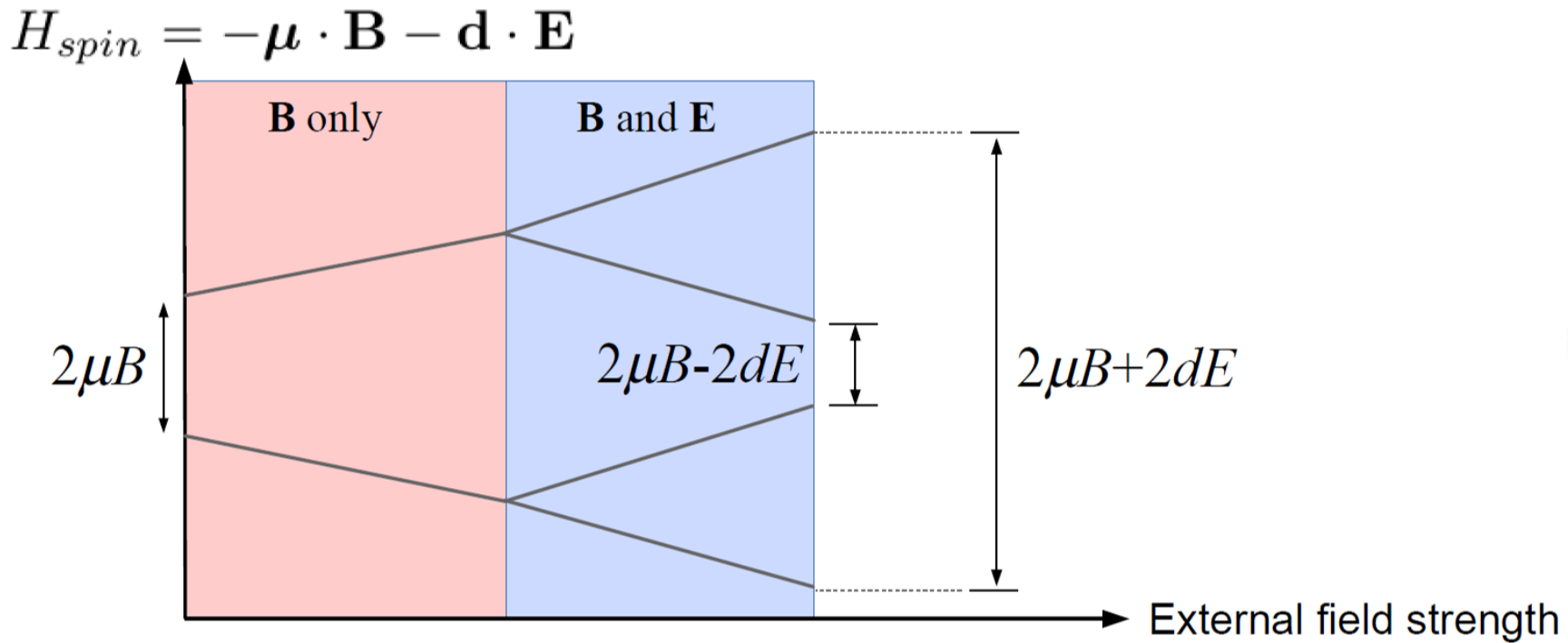
$\bar{g}_\pi^{(1)}$	^{199}Hg	3.8×10^{-13}
$\bar{g}_\pi^{(2)}$	^{199}Hg	2.6×10^{-11}
\bar{d}_n^{sr}	Neutron	3.3×10^{-26} e cm
\bar{d}_p^{sr}	TIF	8.7×10^{-23} e cm
\bar{d}_p^{sr}	^{199}Hg	2.0×10^{-25} e cm
Other parameters		
d_d	$\approx 3/4 d_n$	2.5×10^{-26} e cm
$\bar{\theta}$	$\approx \bar{g}_\pi^{(0)} / (0.015)$	2.5×10^{-10}
$\tilde{d}_d - \tilde{d}_u$	$5 \times 10^{-15} \bar{g}_\pi^{(1)}$ e cm	2×10^{-27} e cm

Many Parameters / Many Experiments

Sensitivity: System:	Paramagnetic	Diamagnetic	"Particle"
Trap	Tl, Cs, PbO, HfF ⁺ , Fr, BaF, ...	¹⁹⁹ Hg, ¹²⁹ Xe, ²²⁵ Ra, Rn, Pa, RaO, ...	n (ultra-cold)
Beam	YbF, ThO, WC	TlF	n (cold)
Storage ring	TaO ⁺	?	p, d, ³ He ⁺⁺ , μ, ...

Other: solid state (Gd₃Ga₅O₁₂, Eu_{0.5}Ba_{0.5}TiO₃), colliders (τ, Λ, ν, ...), crystal (n scattering on quartz), ...

How could you measure an EDM?

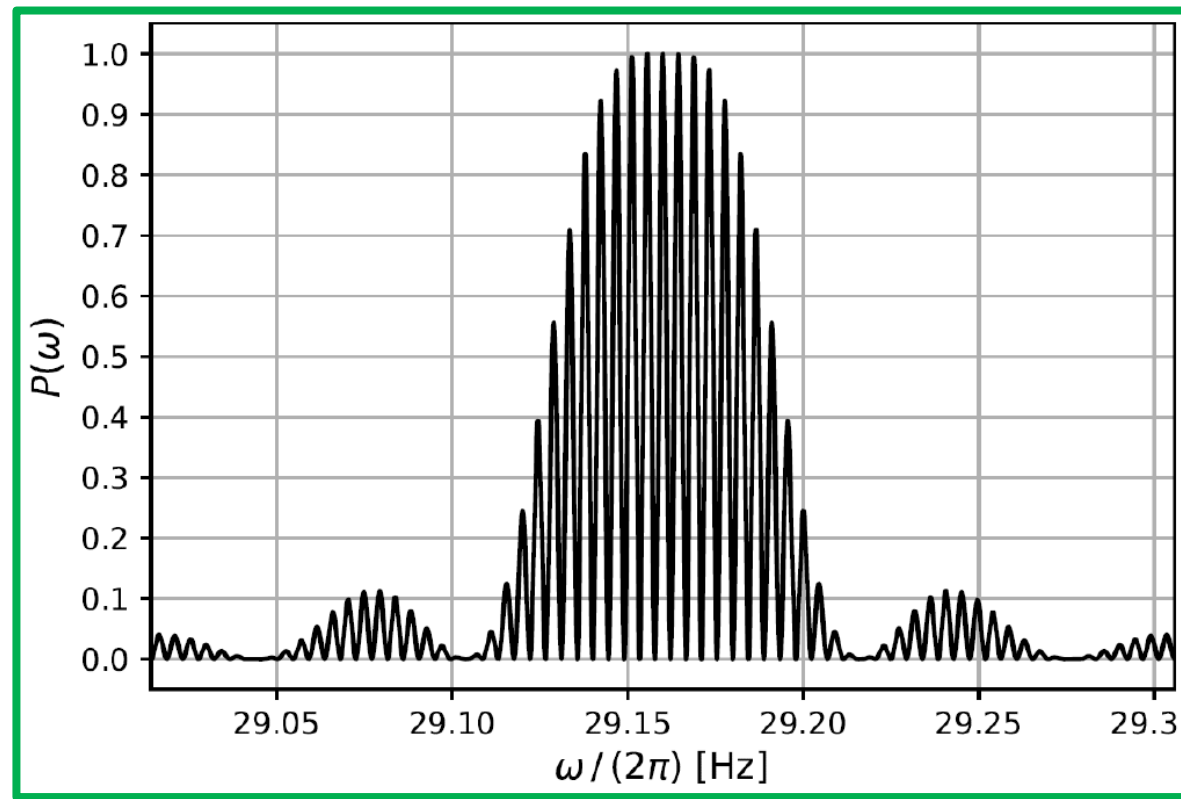
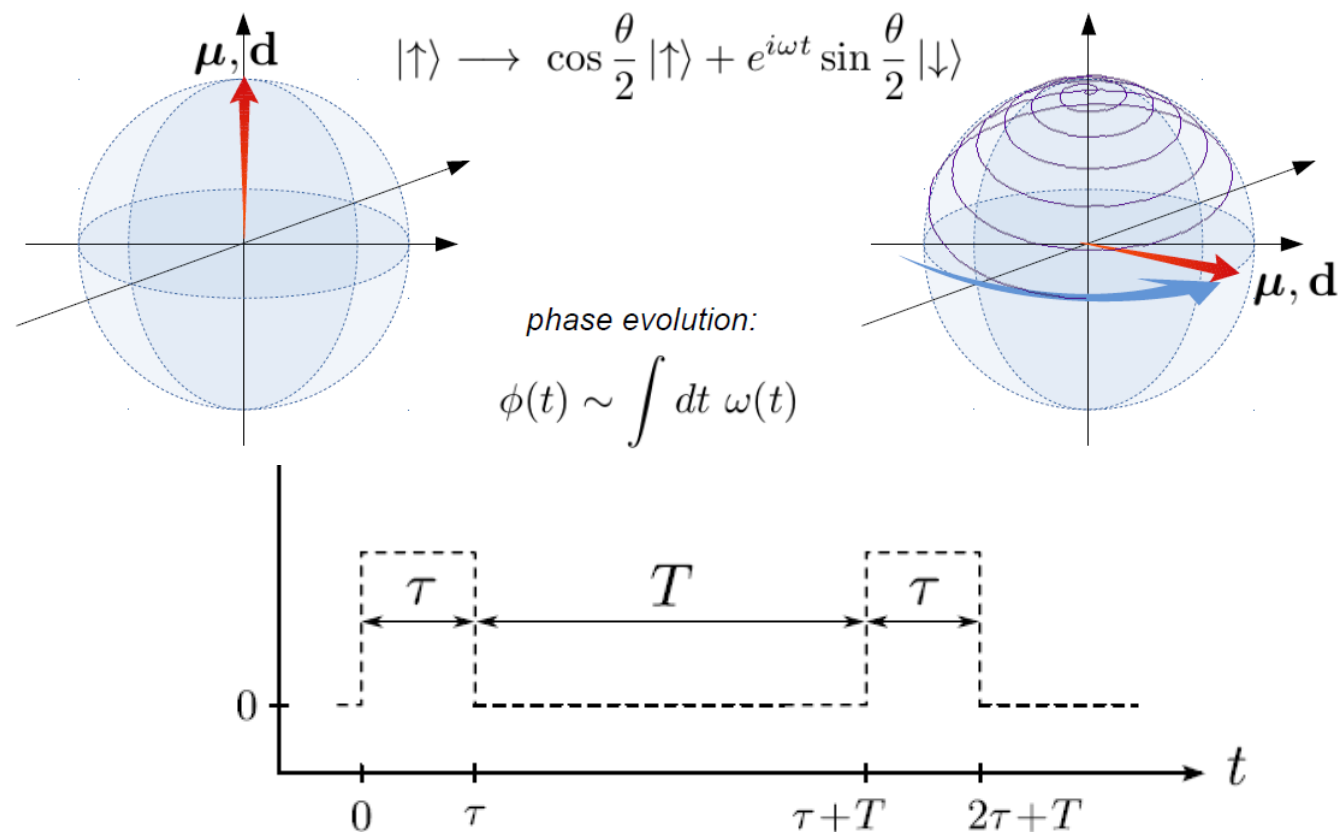


$$\hbar(\omega_+ - \omega_-) = 4dE$$

...up to drift, gradients, etc.

Time-Domain Interferometry

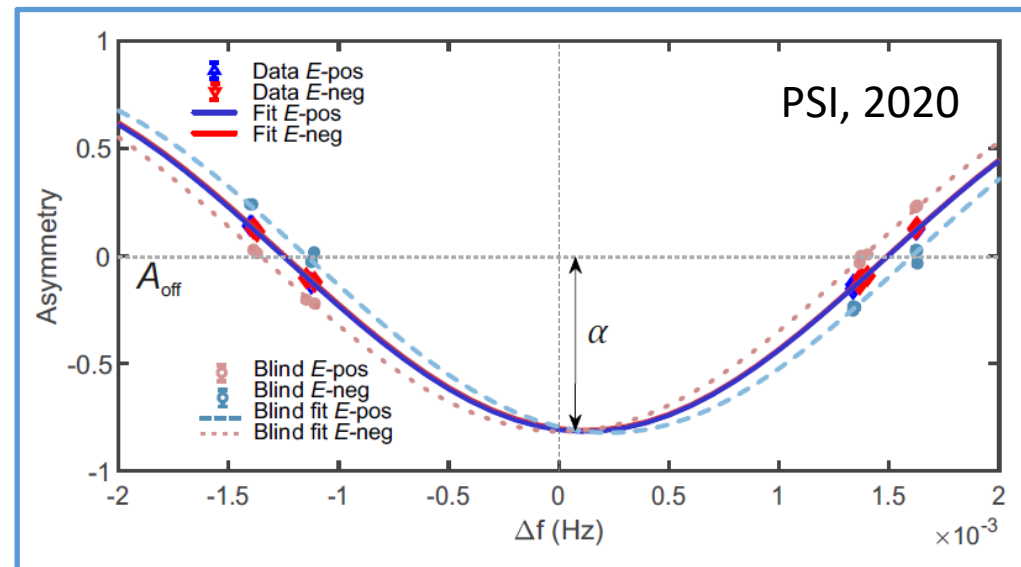
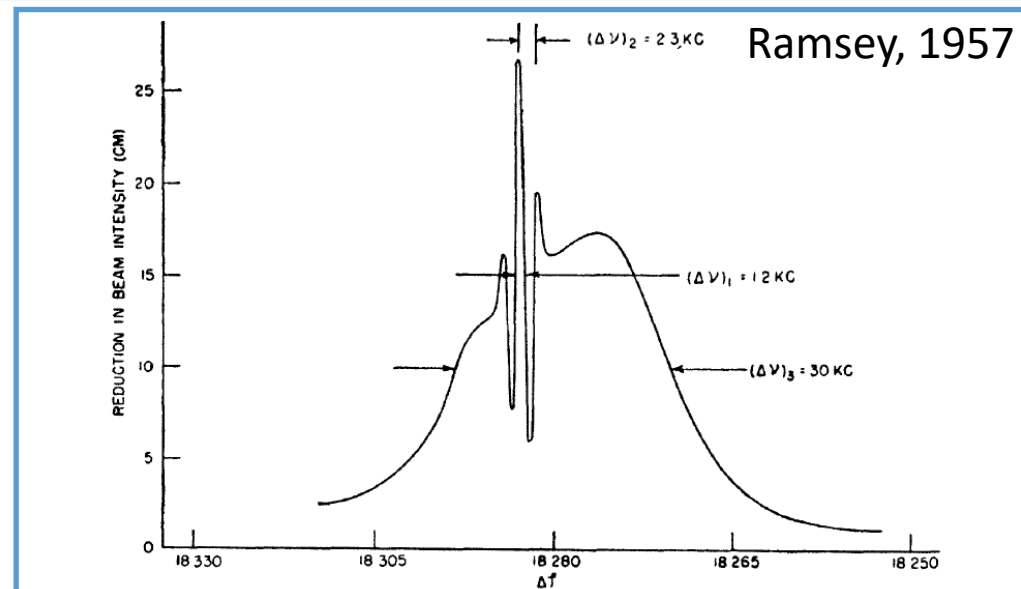
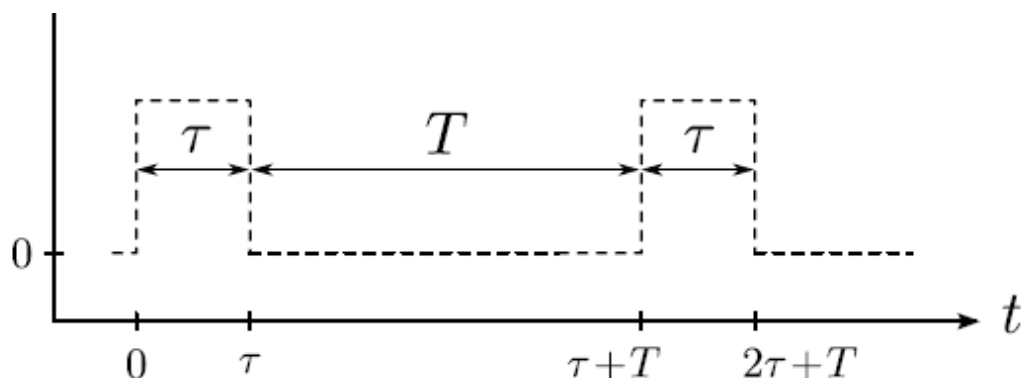
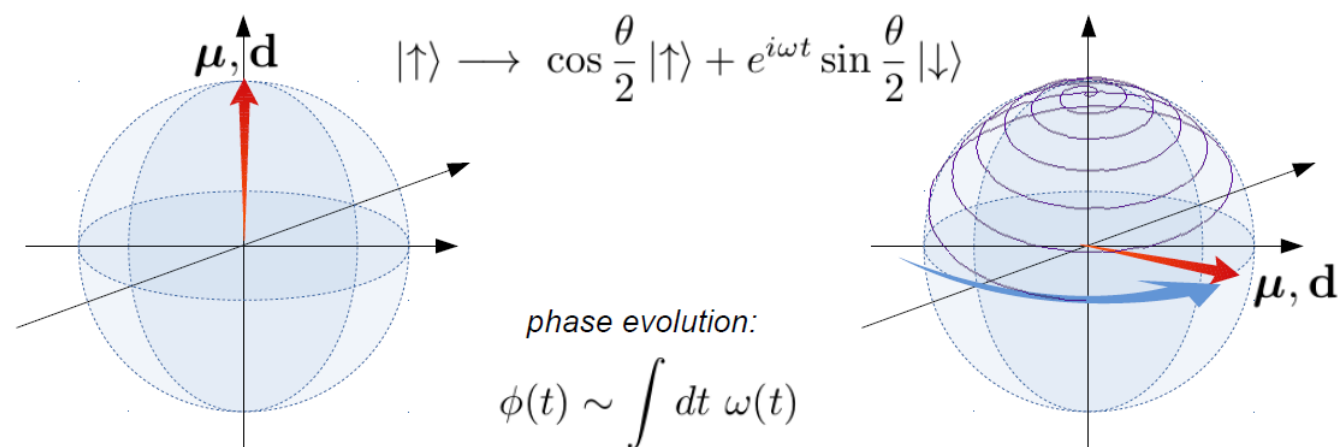
Ramsey's method to measure frequencies*:



*we'll come back to *frequency vs. phase*

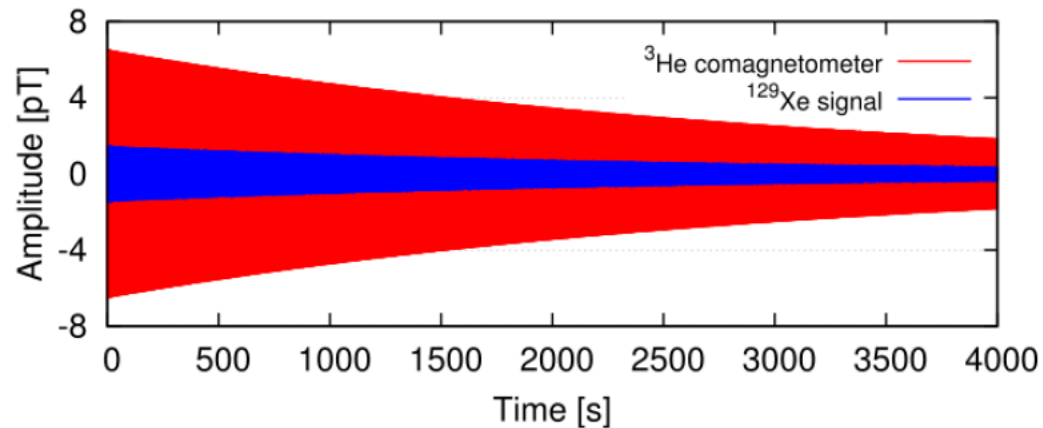
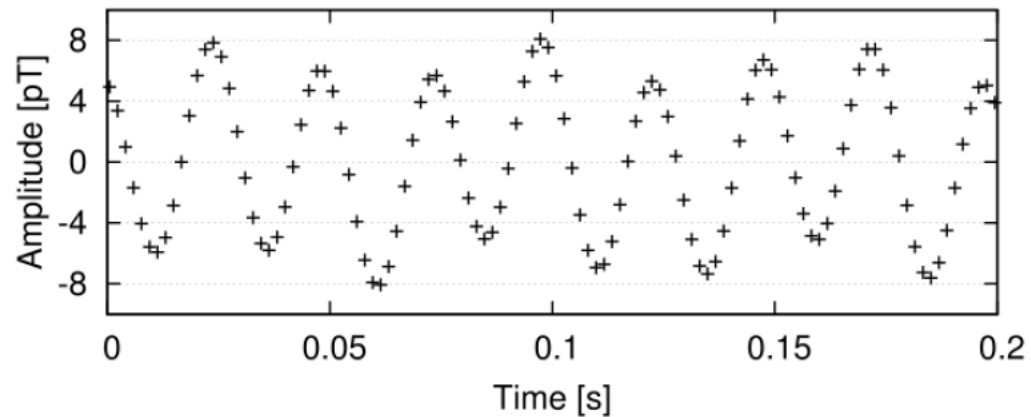
How could you measure an EDM?

Ramsey's method to measure frequencies*:



How could you measure an EDM?

What if we could measure continuously?



$$\delta d \sim \frac{h}{2ET_2} \frac{1}{S/N}$$

“phase noise” limit

↓

$$\frac{h}{2ET_2} \frac{1}{\sqrt{\phi_n T_2}}$$

↑

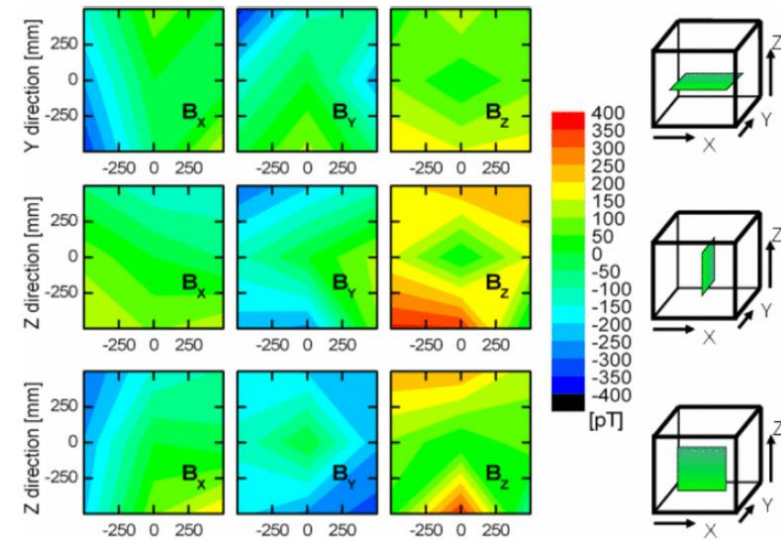
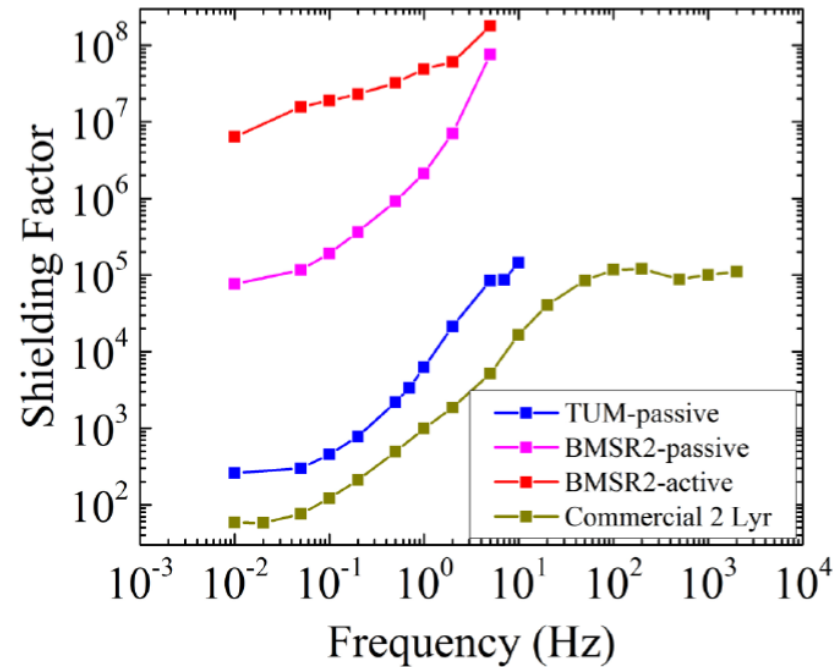
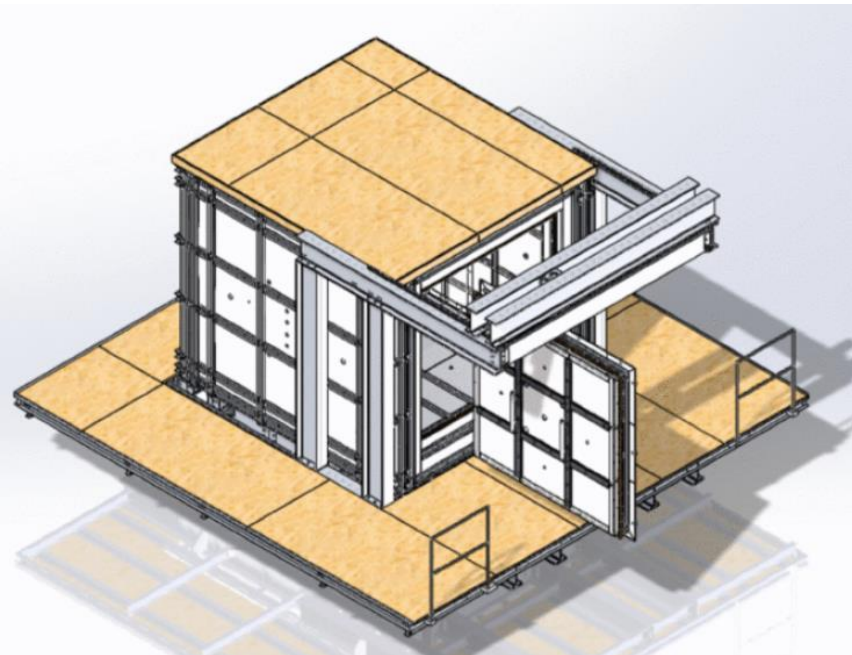
$$\frac{h}{2ET_2} \frac{1}{\sqrt{N_n}}$$

“count rate” limit

The HeXe Experiment: ^{129}Xe



Use the best magnetic shields available (at least to start with...)



Rev. Sci. Instrum. **85**, 075106 (2014)
J. Appl. Phys. **117**, 183903 (2015)

The HeXe Experiment



We do *not* expect a large Schiff enhancement in ^{129}Xe

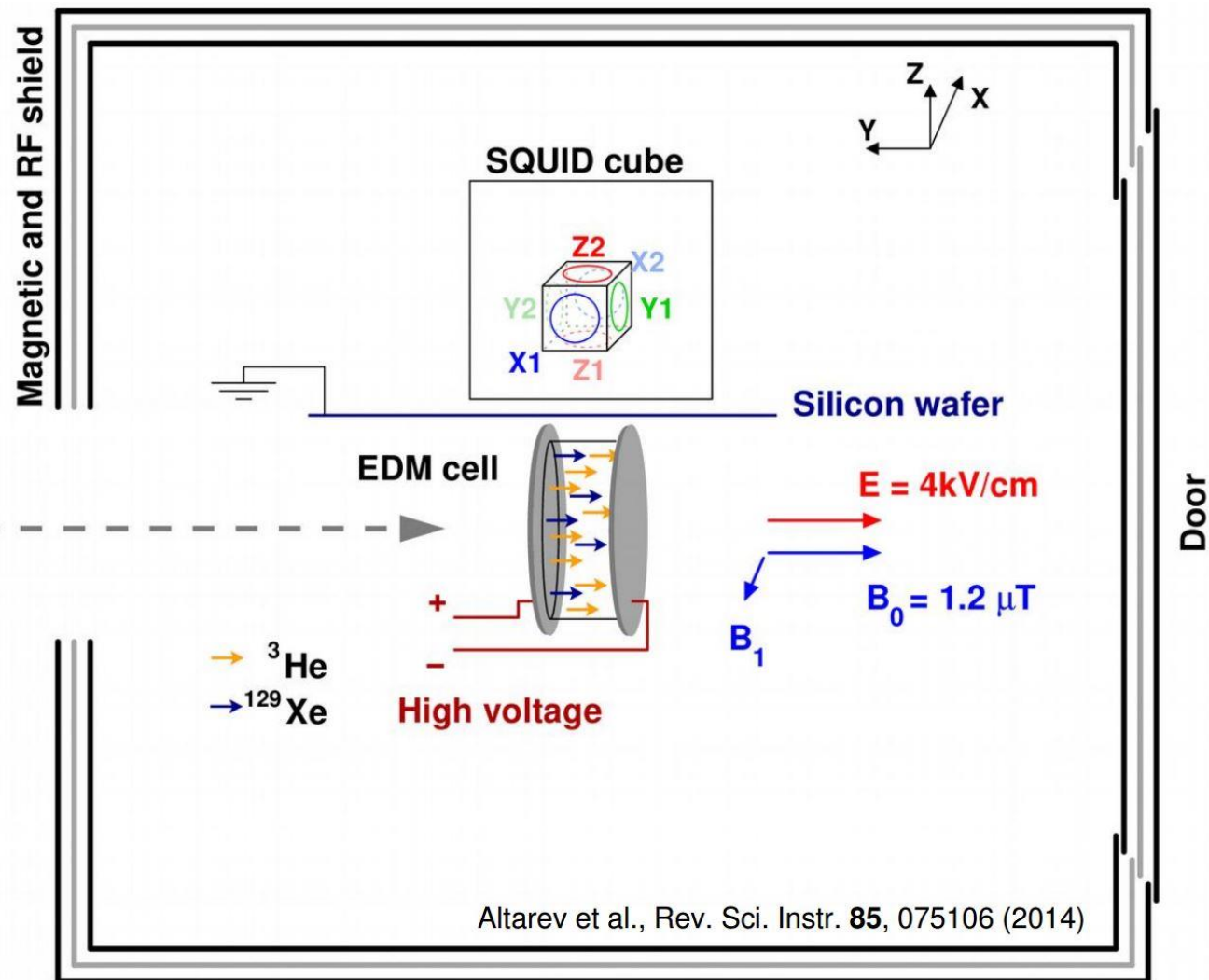
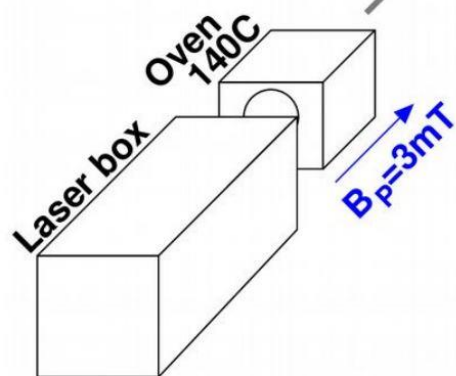
$$S = s_N d_N + \frac{m_N g_A}{F_\pi} [a_0 \bar{g}_\pi^{(0)} + a_1 \bar{g}_\pi^{(1)} + a_2 \bar{g}_\pi^{(2)}]$$

$$\rightarrow d_A(\text{dia}) = \kappa_S S - \underline{[k_{C_T}^{(0)} C_T^{(0)} + k_{C_T}^{(1)} C_T^{(1)}]}$$

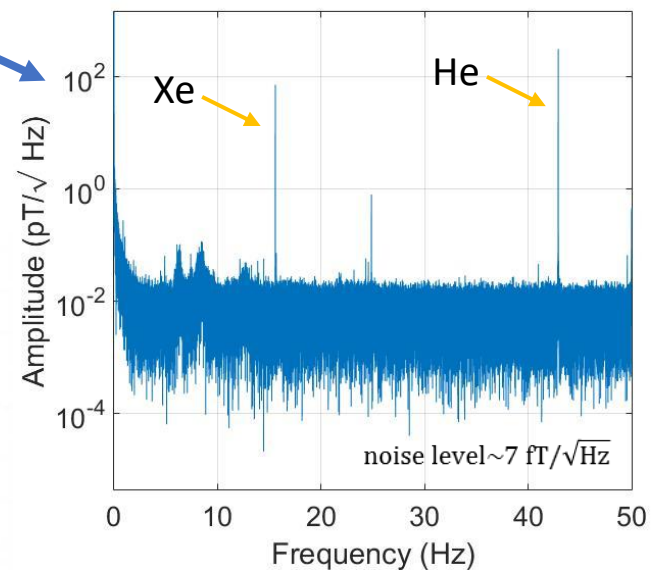
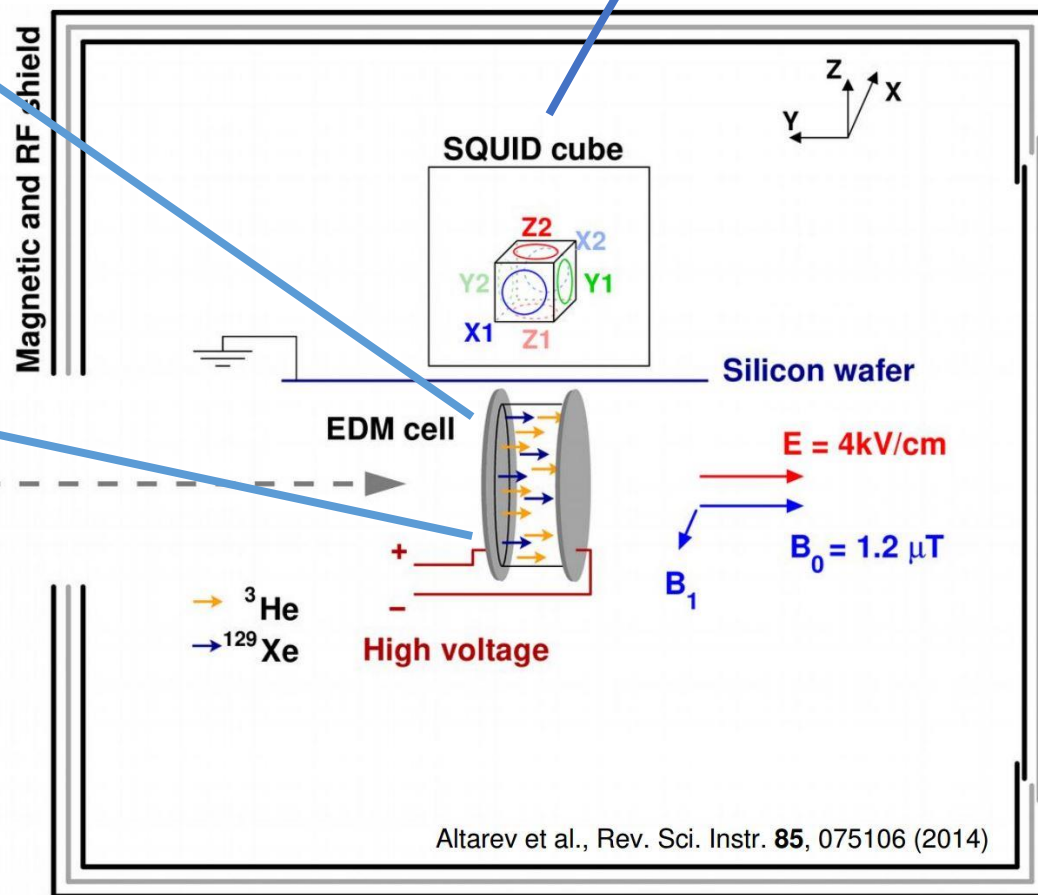
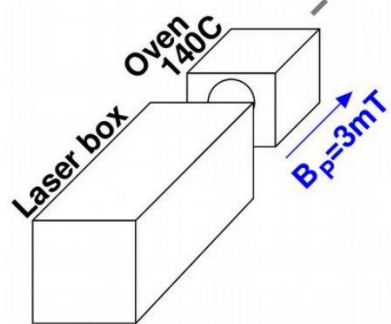
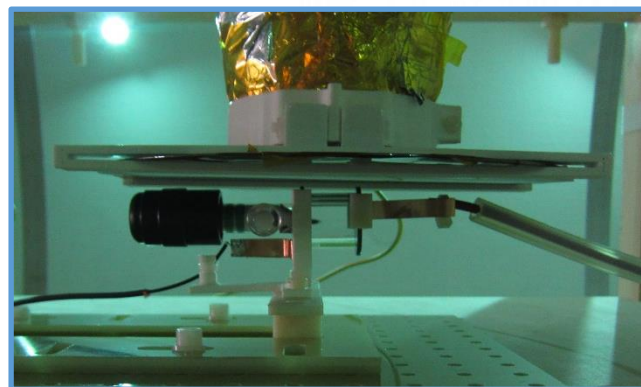
Octupole-deformed nuclei can have enhanced EDMs:

$$S \propto \frac{\eta \beta_1 \beta_3^2 A^{\frac{2}{3}} r_0^3}{E_+ - E_-}$$

...but new challenges for those.



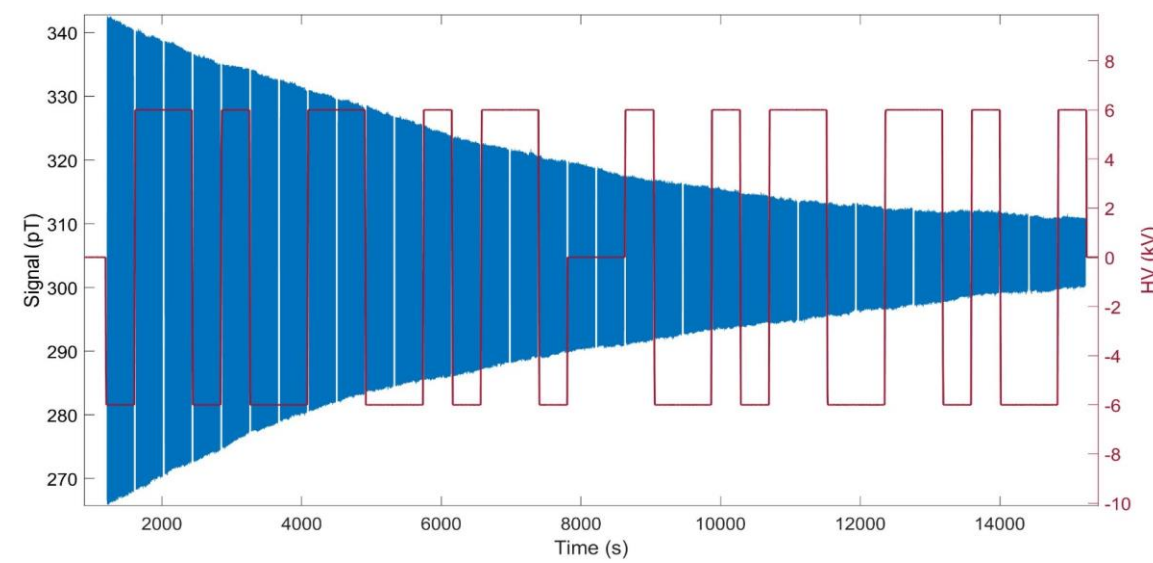
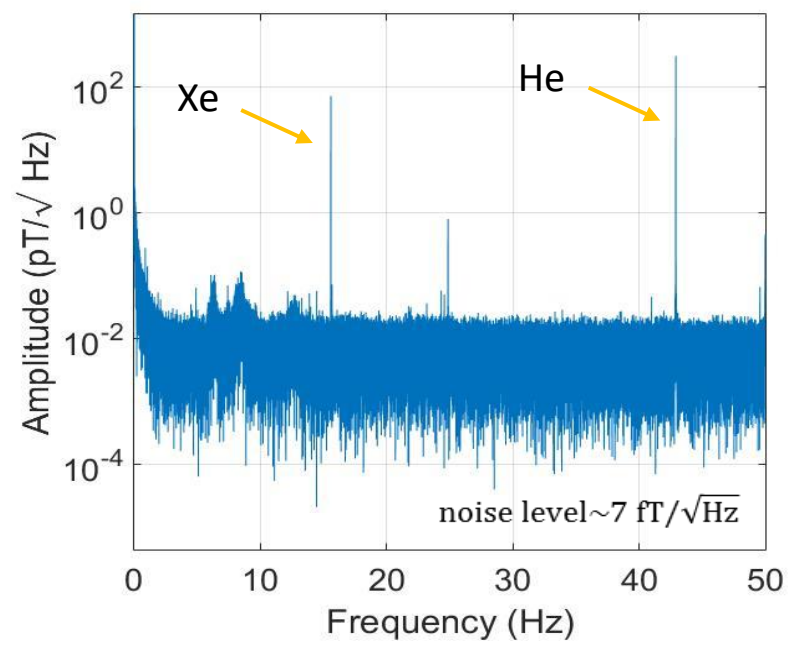
The HeXe Experiment



The HeXe Experiment



$E \sim 4 \text{ kV/cm}$
 $\tau \sim 4000 \text{ s}$
 $S \sim 20 \text{ pT}$
 $\epsilon \sim 8 \text{ fT}/\sqrt{\text{Hz}}$



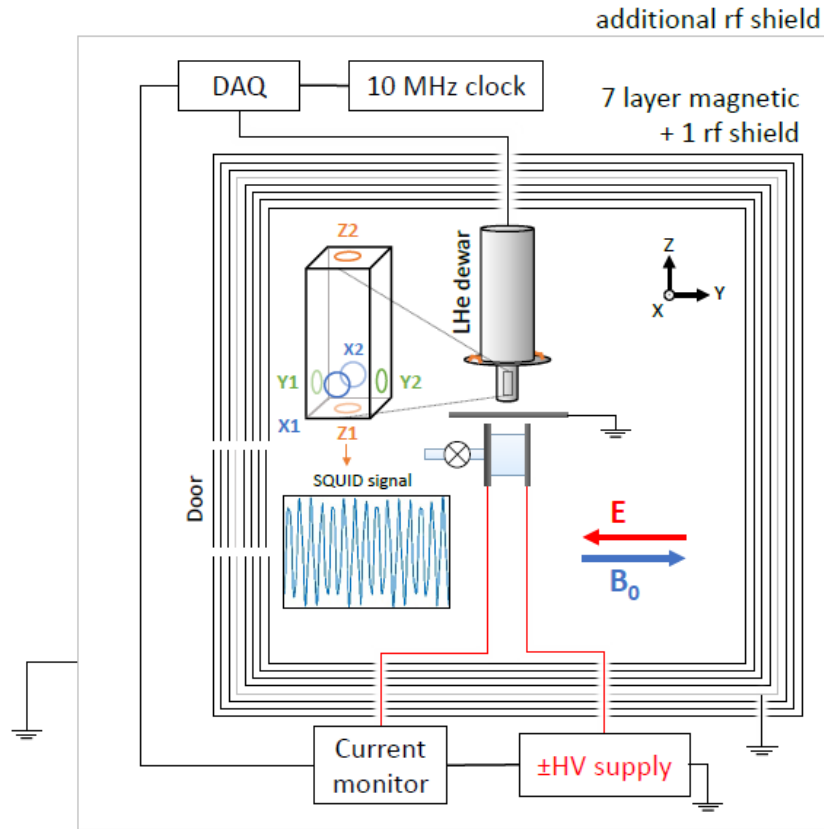
$$\delta\omega = \frac{1}{\tau(S/n)\sqrt{N}} = \frac{\epsilon\sqrt{f_{\text{BW}}}}{\tau S\sqrt{N}}$$

→ require nHz per run

$$\sigma_d = \frac{\hbar}{2E} \frac{\epsilon}{\tau^{3/2} S\sqrt{N}} = \frac{\hbar}{2E} \frac{\epsilon}{\tau S\sqrt{T}}$$

→ few $\times 10^{-27} \text{ e cm}/\sqrt{N}$

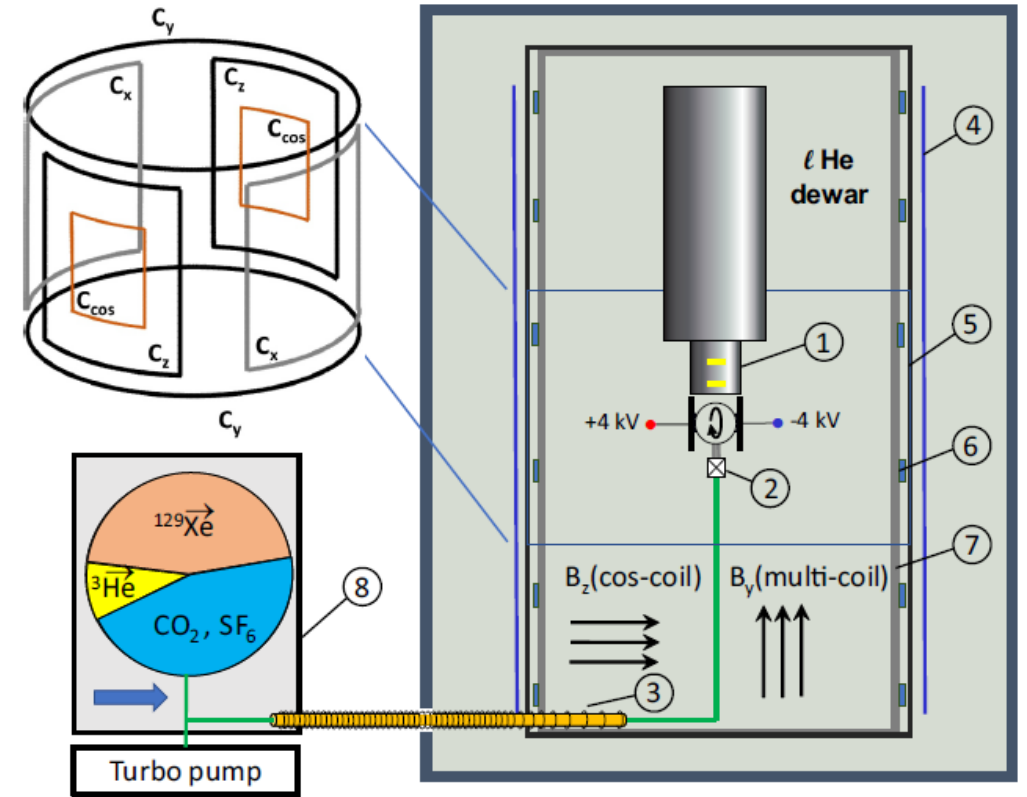
A rapidly-moving field!



Our result from HeXe:

$$d_A(^{129}\text{Xe}) = (1.4 \pm 6.6_{\text{stat}} \pm 2.0_{\text{syst}}) \times 10^{-28} \text{ e cm.}$$

Phys. Rev. Lett. **123**, 143003 (2019)

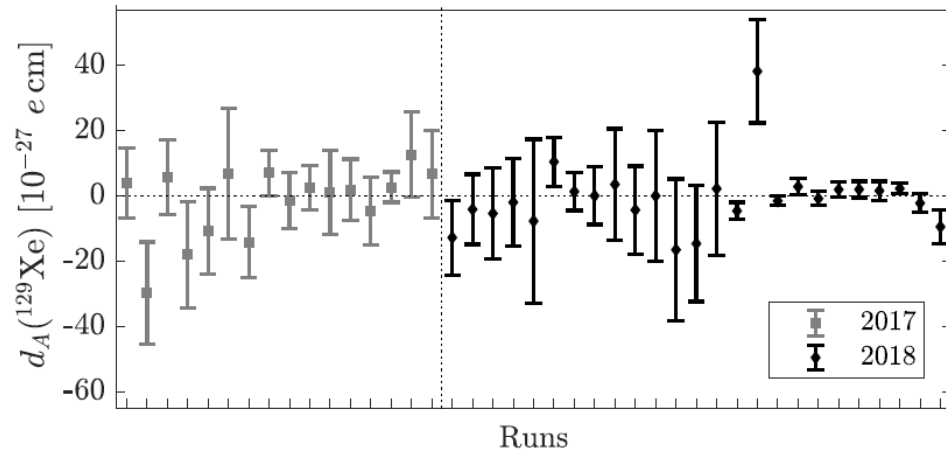


Near-simultaneous from MiXed:

$$d_{\text{Xe}} = (-4.7 \pm 6.4) \cdot 10^{-28} \text{ e cm}$$

Phys. Rev. A **100**, 022505 (2019)

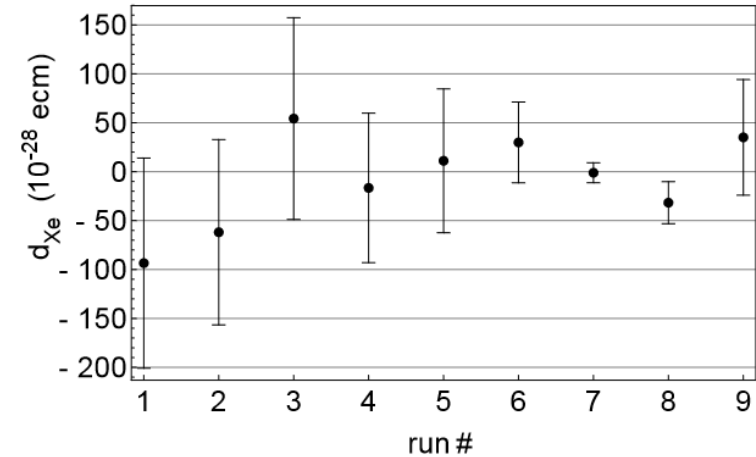
A rapidly-moving field!



Our result from HeXe:

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Phys. Rev. Lett. **123**, 143003 (2019)



Near-simultaneous from MiXed:

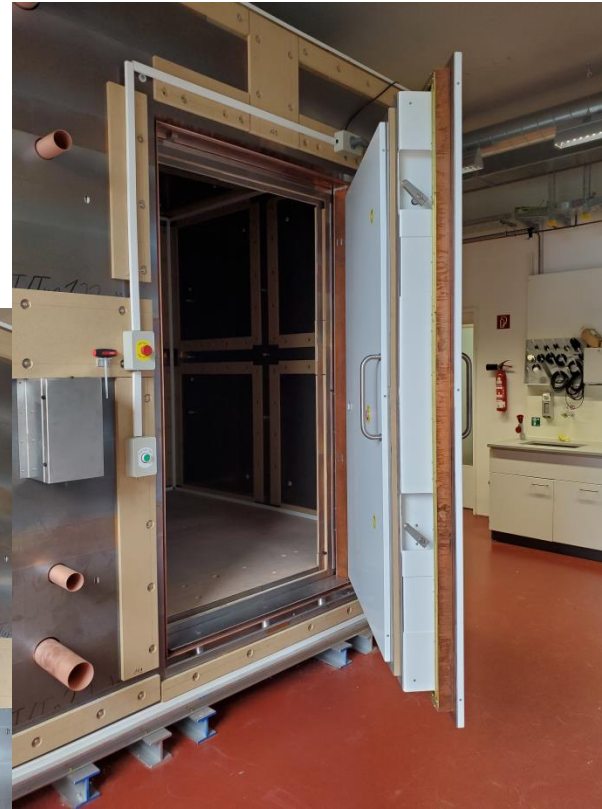
$$d_{\text{Xe}} = (-4.7 \pm 6.4) \cdot 10^{-28} \text{ e cm}$$

Phys. Rev. A **100**, 022505 (2019)

A United Future at the PI



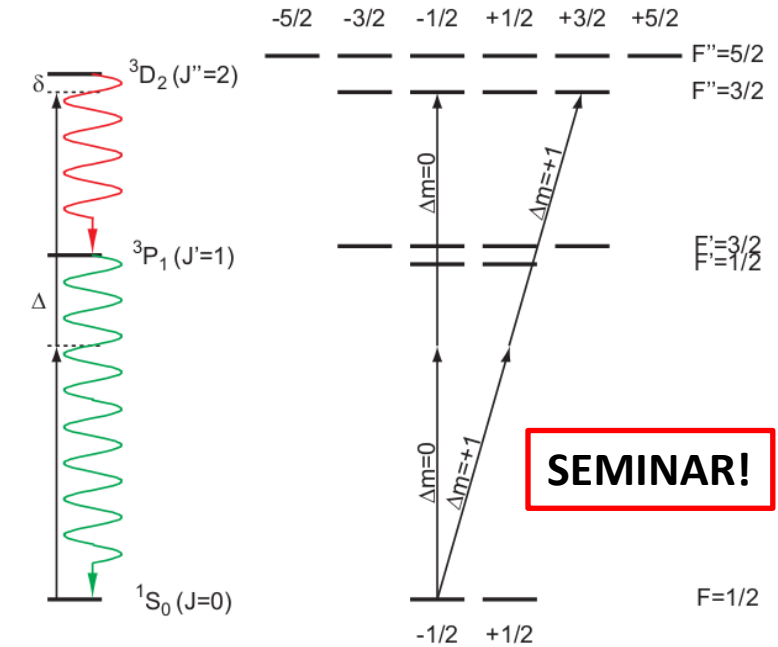
New magnetically shield room @HD!
Dedicated facility...



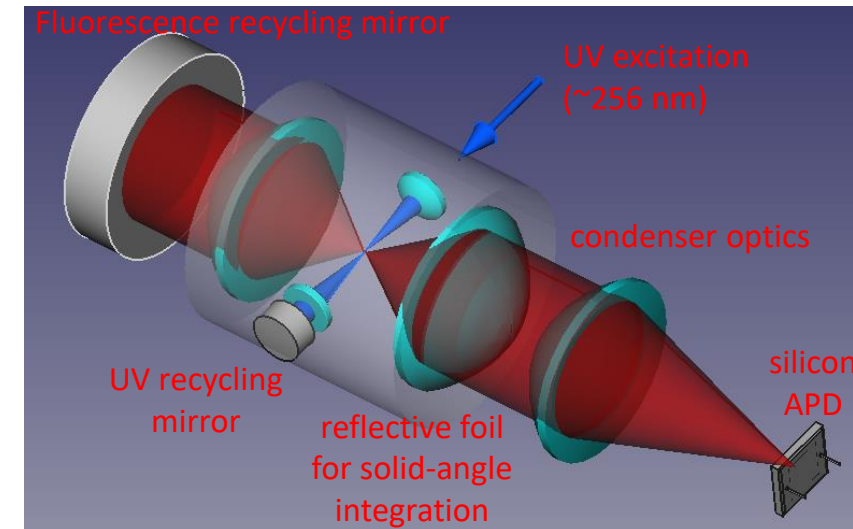
Next order-of-magnitude pursued by refining these now-known methods



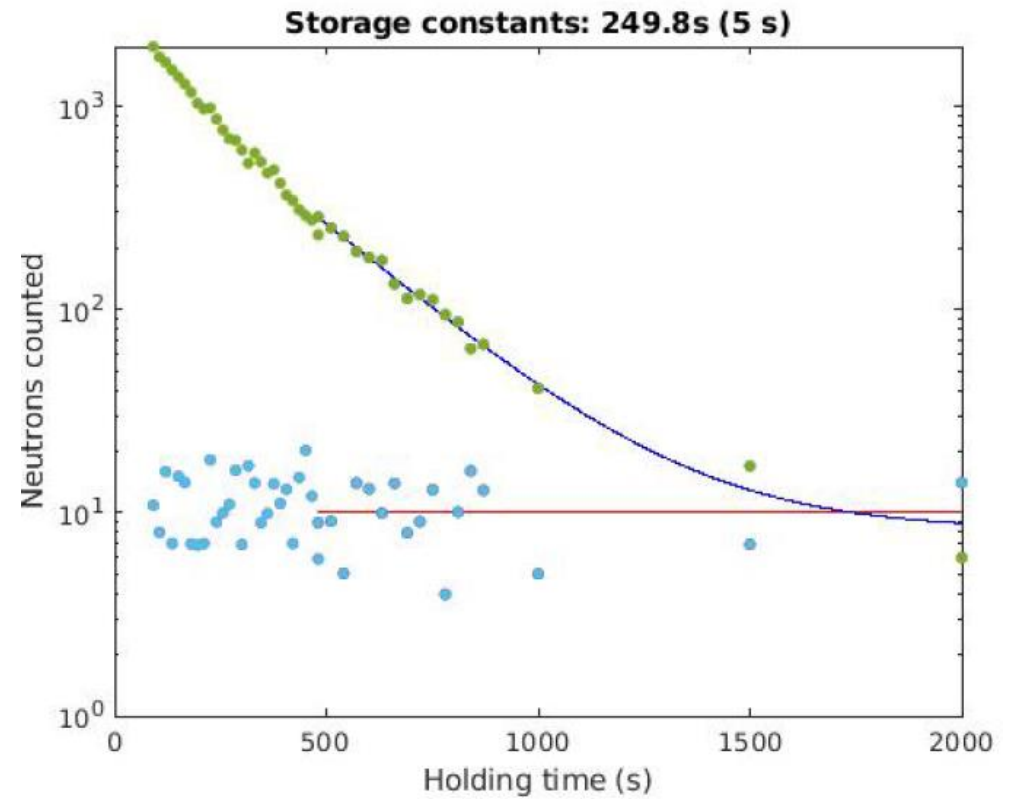
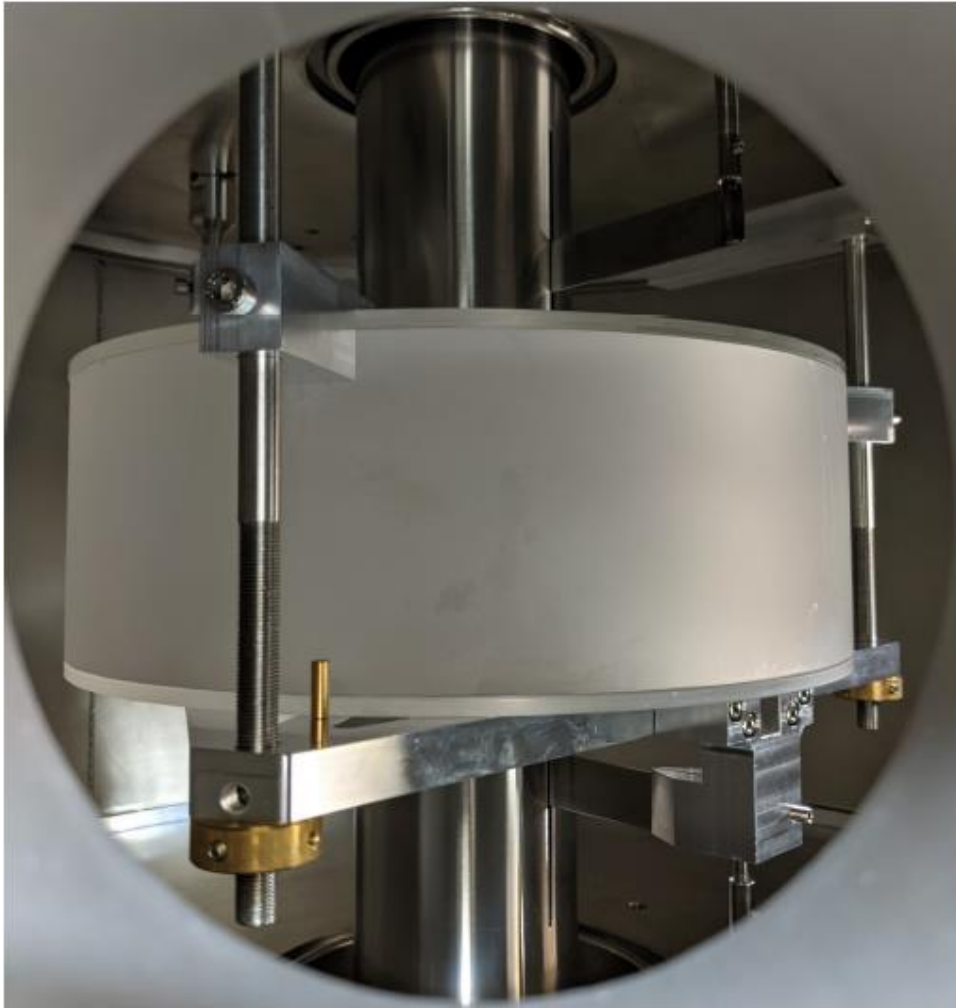
Laser spectroscopy may complement or eventually replace SQUIDS... new tools!



SEMINAR!

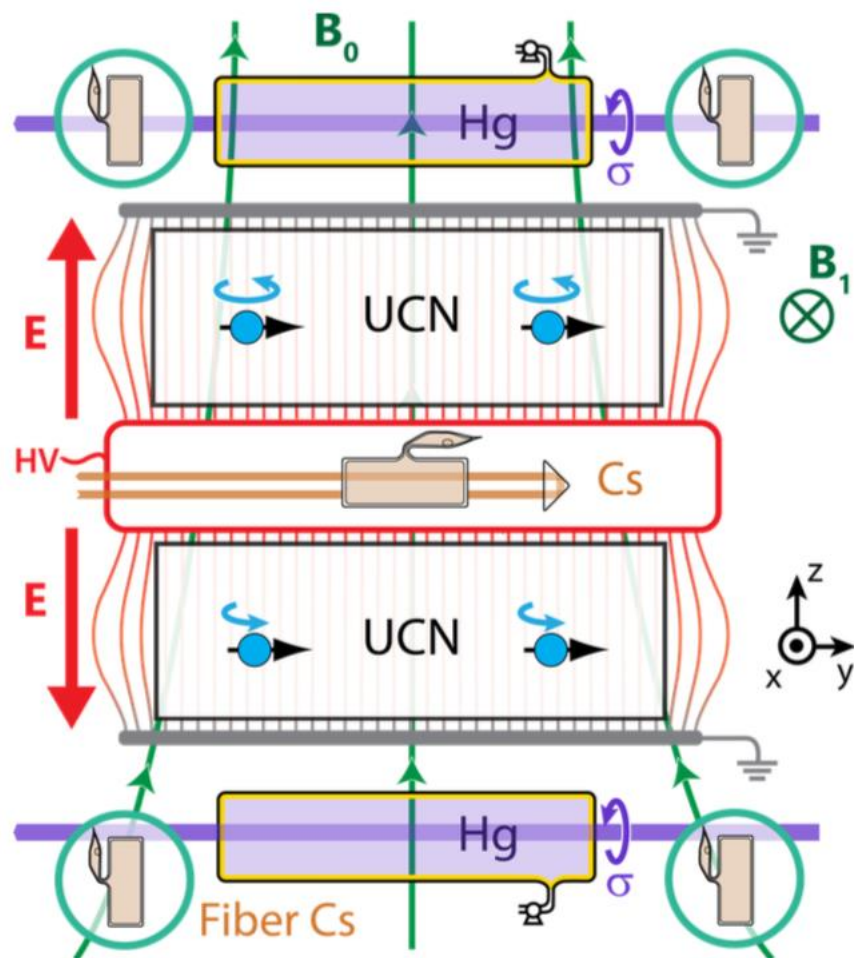


By contrast: neutrons disappear faster!



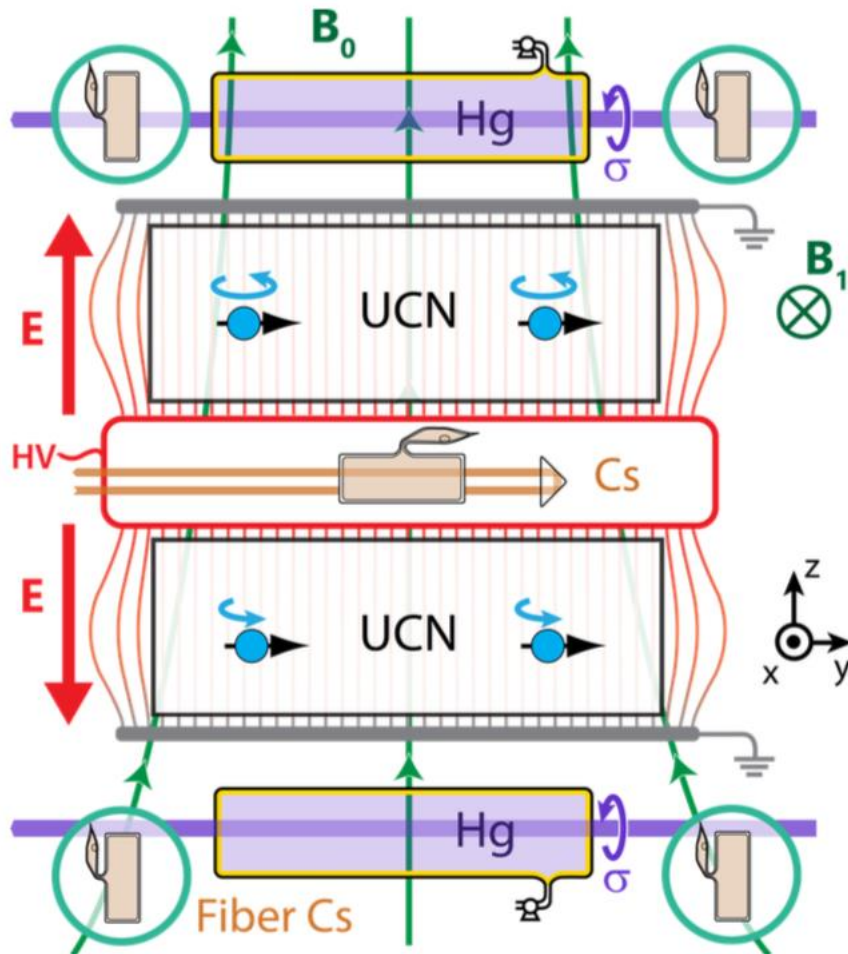
Much better at low temperature: Neulinger *et al.*, EPJA 58, 141 (2022)

The PanEDM Experiment



- Double chamber Ramsey interferometer at room temperature (but $T_{UCN} \sim 5\text{mK}$)
- ^{199}Hg magnetometers with few-fT resolution
- Cs magnetometers (also at high voltage)
- Magnetic shielding factor: 6×10^6 at 1 mHz
- Simultaneous spin detection for up/down
- SuperSUN UCN source at ILL in 2 phases:
 - Phase I: unpolarized UCN with 80 neV peak
 - Phase II: polarized UCN, magnetic storage
- Ongoing installation of parts, commissioning with UCN production in 2023-2024

Much lower statistics!



Statistical sensitivity:

$$\sigma(d_n) \gtrsim \frac{\hbar}{2\alpha|\mathbf{E}|T\sqrt{N}}$$

Frequency measurements:

$$|\delta\omega| = \frac{|dE|}{\hbar F}$$

SuperSUN

Phase I

Saturated source

density [cm^{-3}]

330

Diluted density [cm^{-3}]

63

Density in cells [cm^{-3}]

3.9

PanEDM Sensitivity [$1\sigma, e\text{ cm}$]

Per run

5.5×10^{-25}

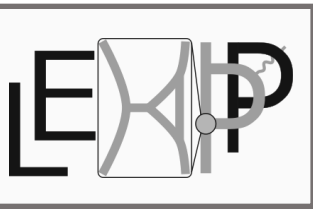
Per day

3.8×10^{-26}

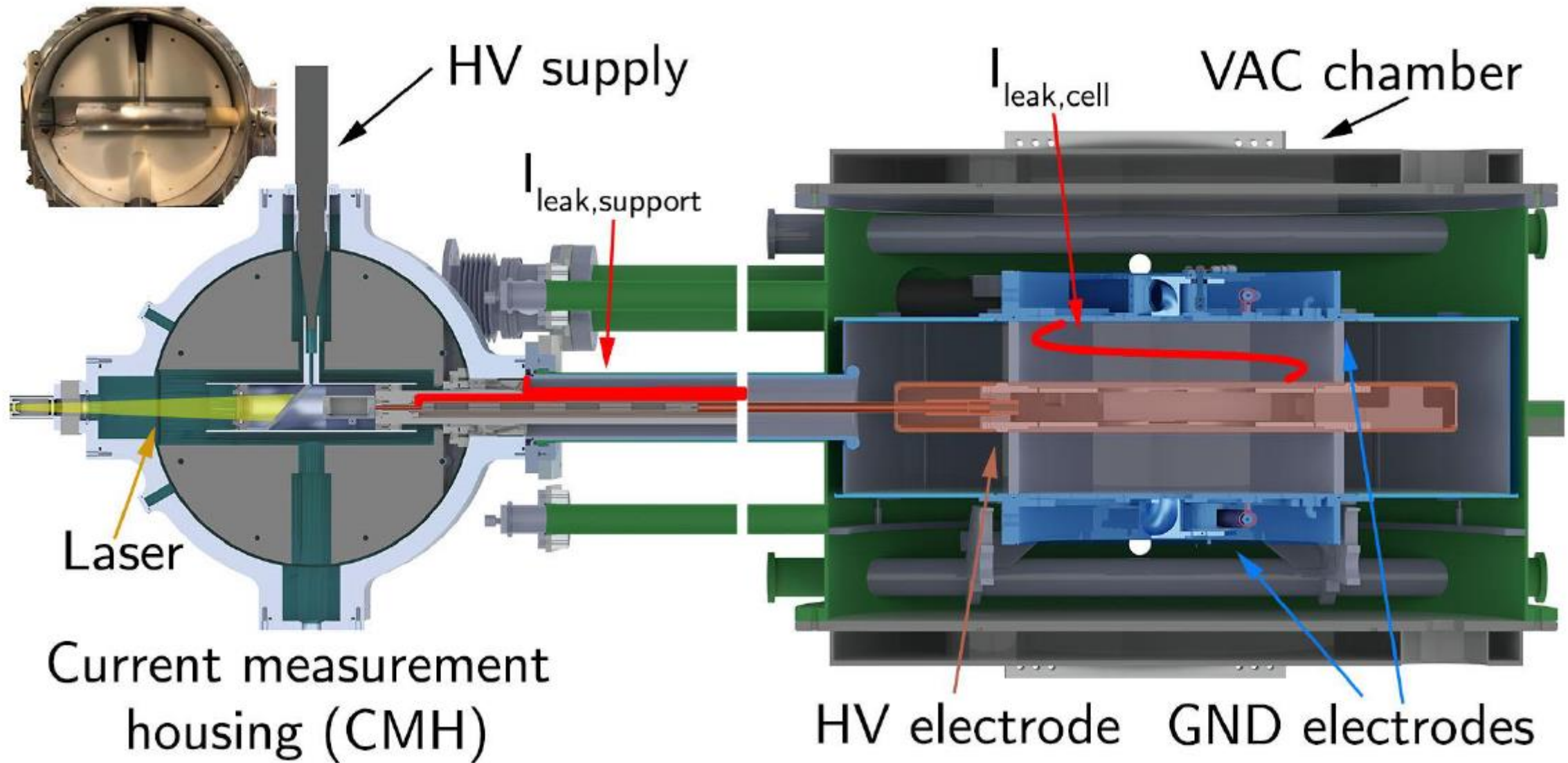
Per 100 days

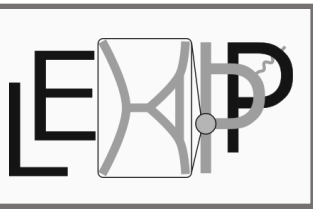
3.8×10^{-27}

$$\Delta E \Delta t \geq \hbar/2$$

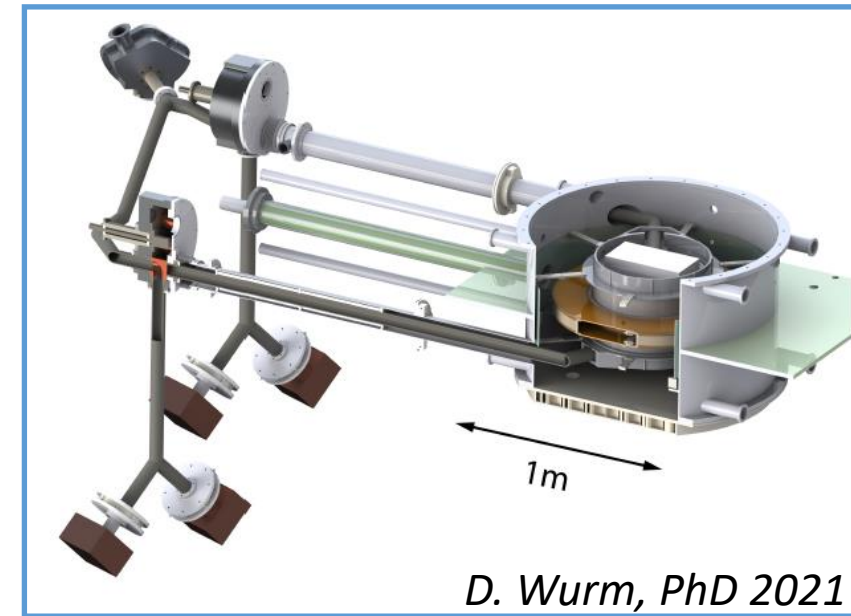
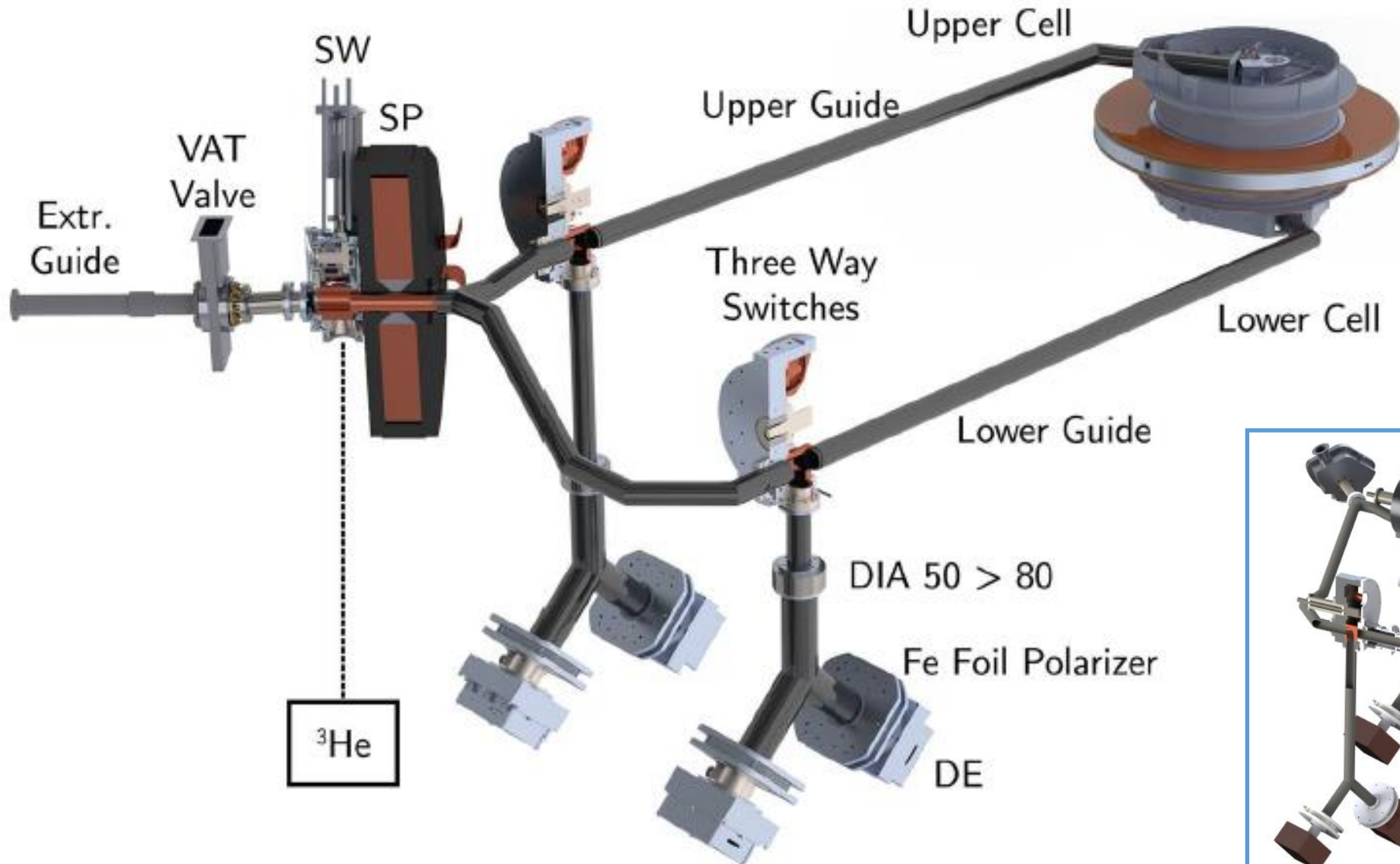


The PanEDM Experiment



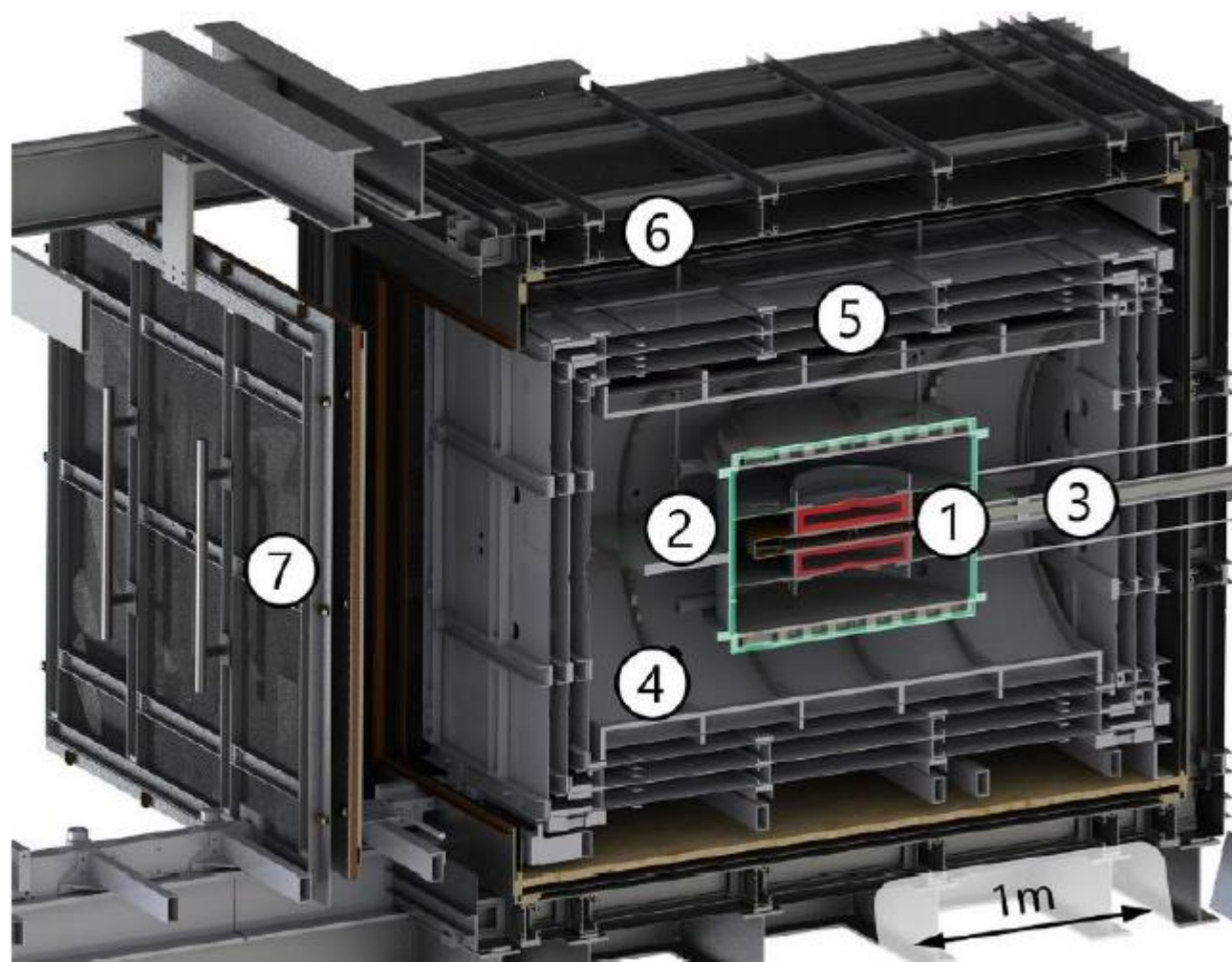


The PanEDM Experiment



D. Wurm, PhD 2021

The PanEDM Experiment



1: EDM cells

3: HV feed

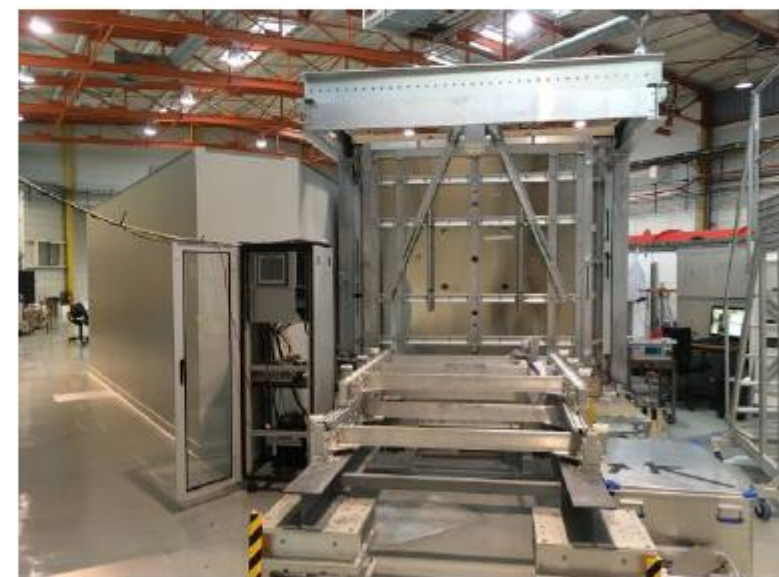
5: Inner shield

7: Outer shield door

2: Vac. Chamber

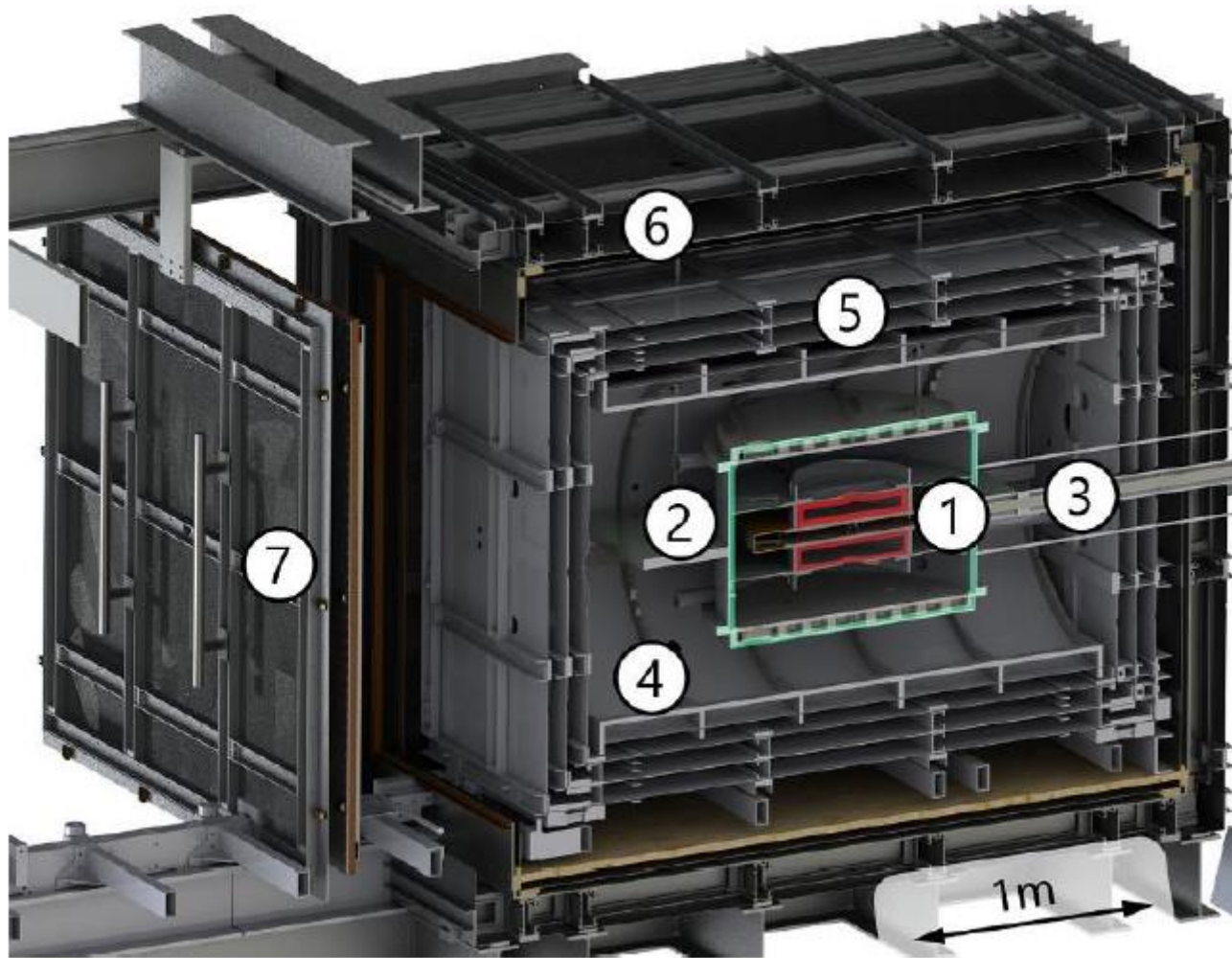
4: B_0 & B_1 coil

6: Outer shield

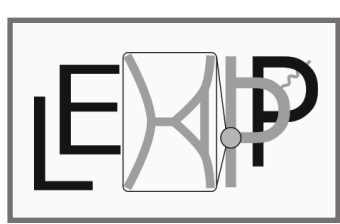


PanEDM @ ILL, 2021

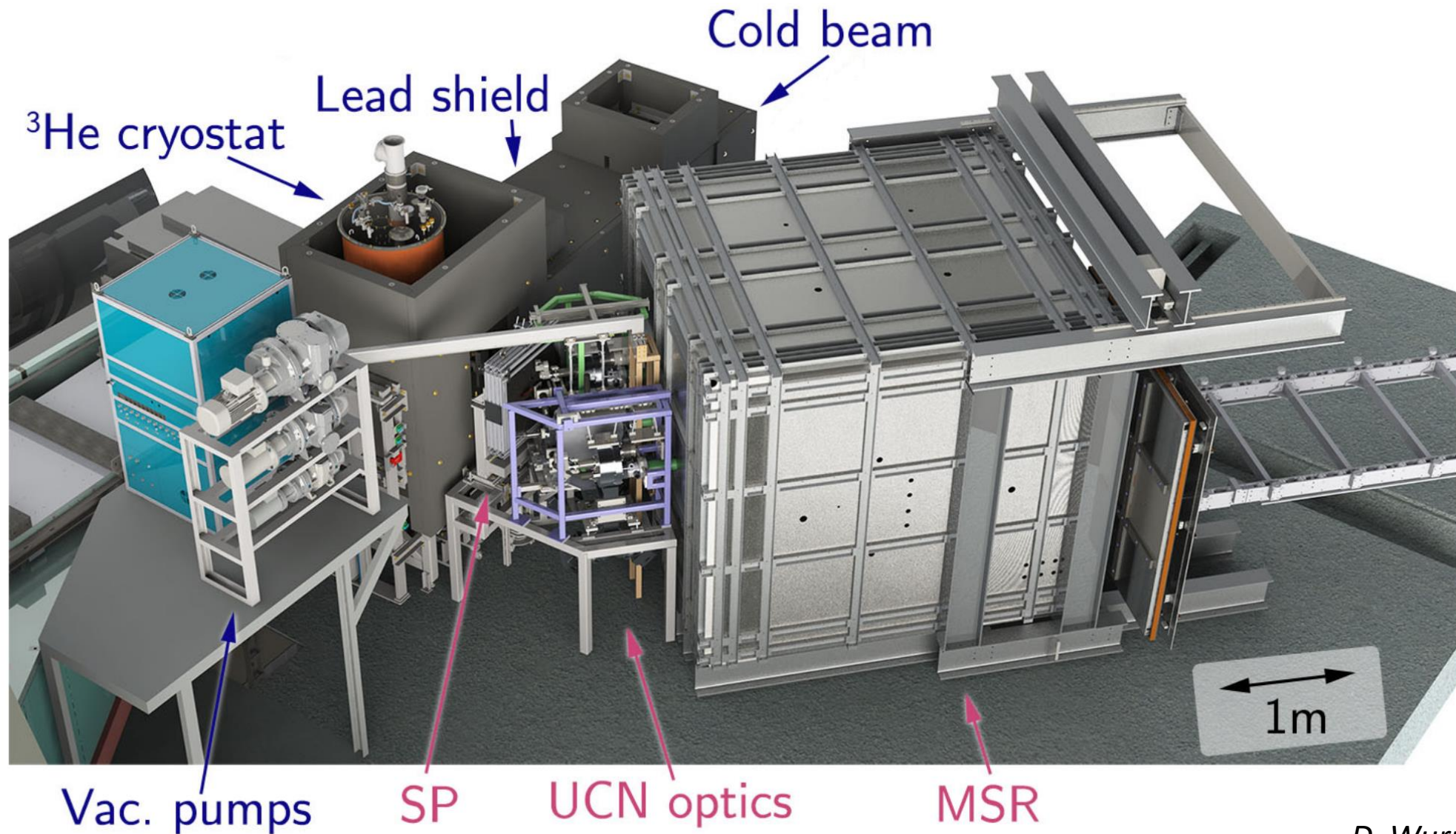
The PanEDM Experiment

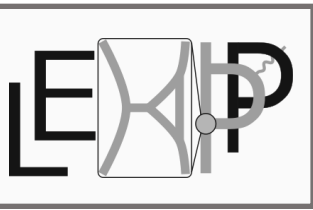


Rev. Sci. Instr. 85(7), 075106 (2014)
J. Appl. Phys. 117(18), 183903 (2015)

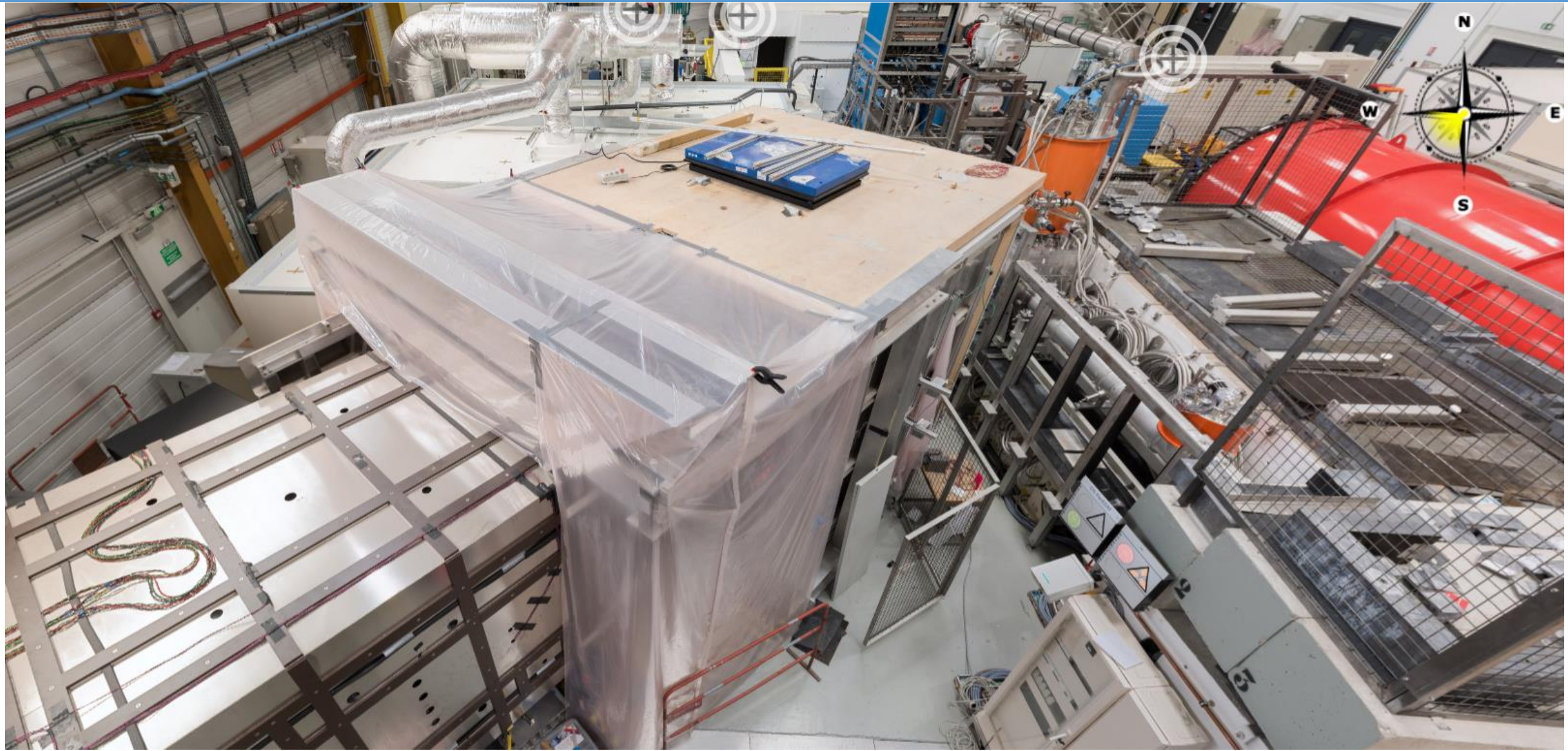


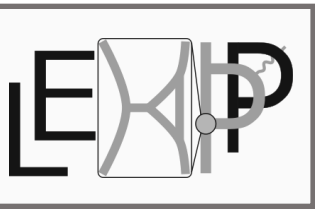
The SuperSUN-PanEDM Installation



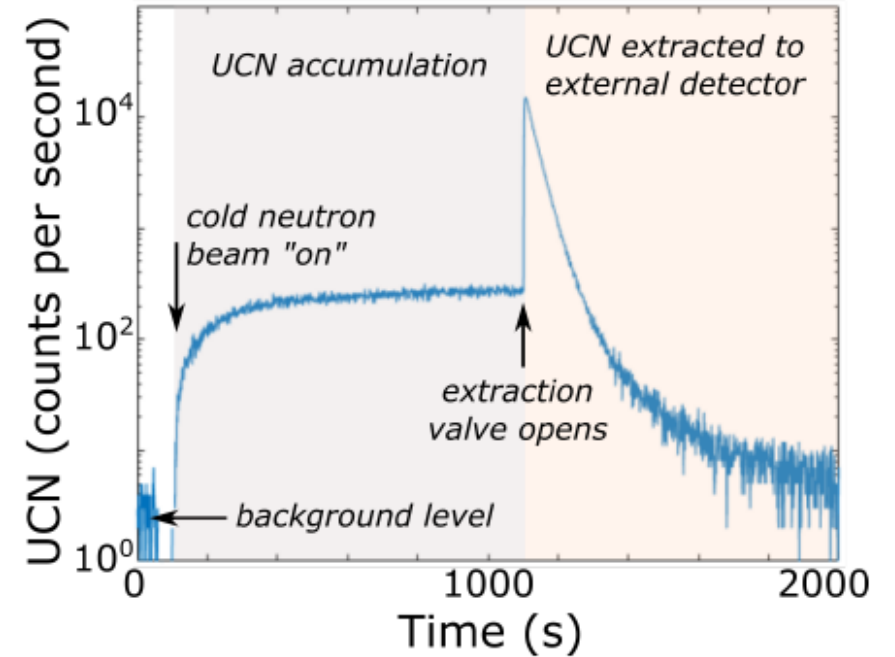
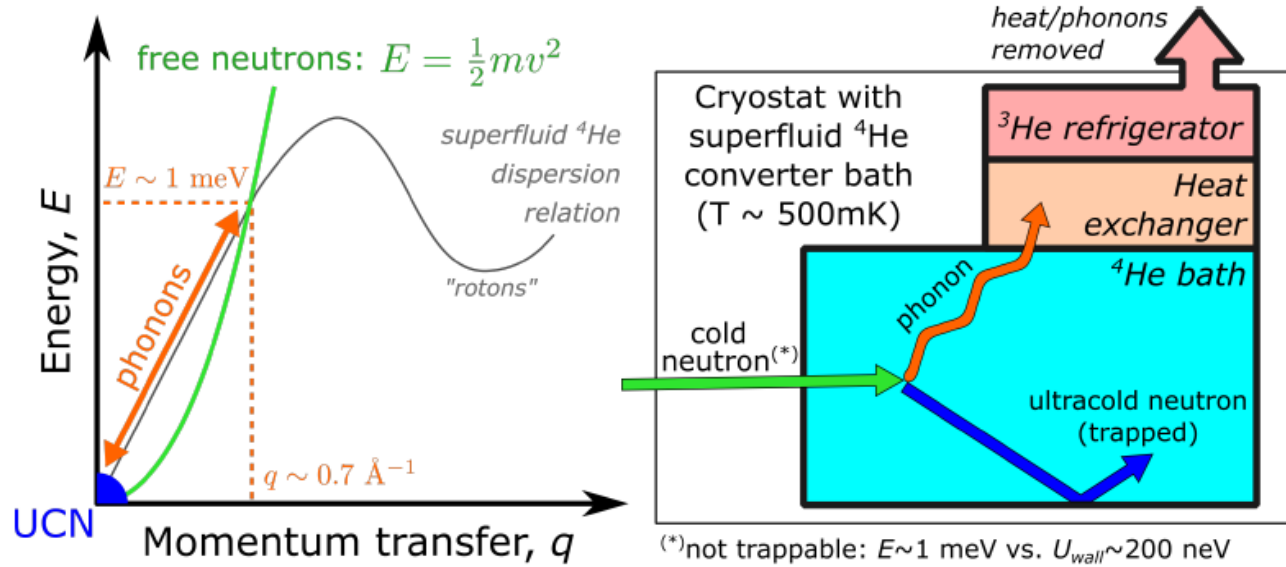


Reality always looks messier!





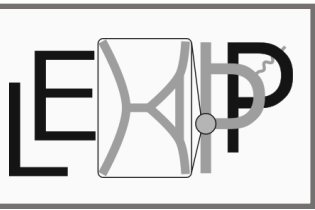
UCN and Production in He-II



Velocity	"Temperature"	Energy
$10^0 - 10^1 \text{ m/s}$	Ultracold	5 neV – 500 neV
$10^1 - 10^2 \text{ m/s}$	Very cold	$0.5 \mu\text{eV} - 50 \mu\text{eV}$
$10^2 - 10^3 \text{ m/s}$	Cold	$50 \mu\text{eV} - 5 \text{ meV}$
$2.2 \times 10^3 \text{ m/s}$	Thermal	25 meV
$2 \times 10^3 - 2 \times 10^4 \text{ m/s}$	Hot	20 meV – 2 eV

$$R \sim \left(\frac{5 \times 10^{-8}}{\text{cm}^3 \text{ s}} \right) \times \left. \frac{d\Phi}{d\lambda} \right|_{8.9\text{\AA}} \times \left(\frac{V_{\text{trap}}}{233 \text{ neV}} \right)^{\frac{3}{2}} \quad \text{production}$$

$$\frac{1}{\tau} = \frac{1}{\tau_{\beta}} + \frac{1}{\tau_{\text{up}}} + \frac{1}{\tau_{\text{capture}}} + \frac{1}{\tau_{\text{wall}}} + \dots \quad \text{loss}$$

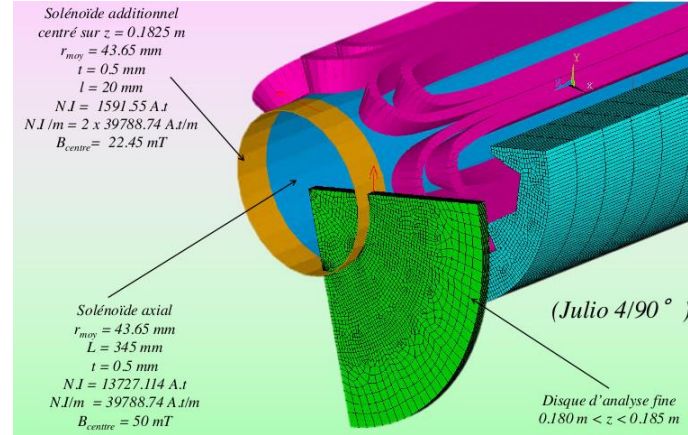


SuperSUN Neutron Source: Cutaway



^3He pumping

1K pot



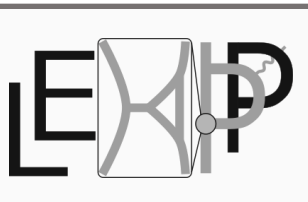
SC Octupole ~2.1T



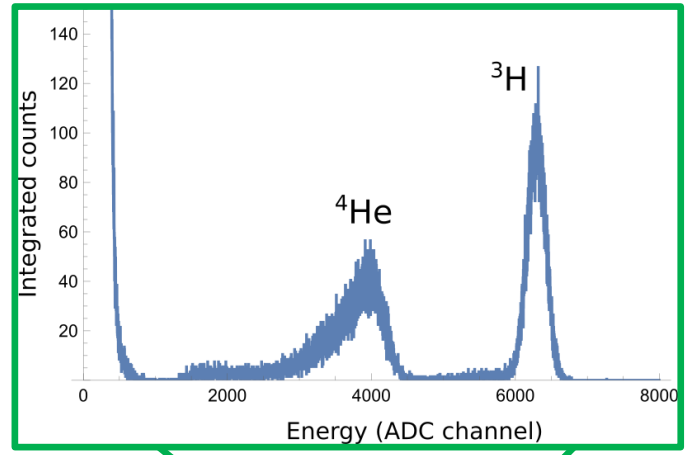
cryogenic CN guide

Isotopically pure ^4He

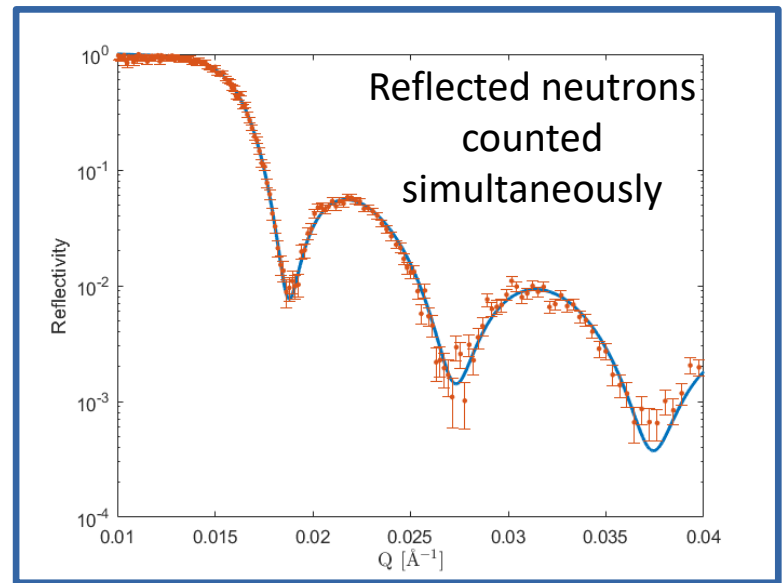
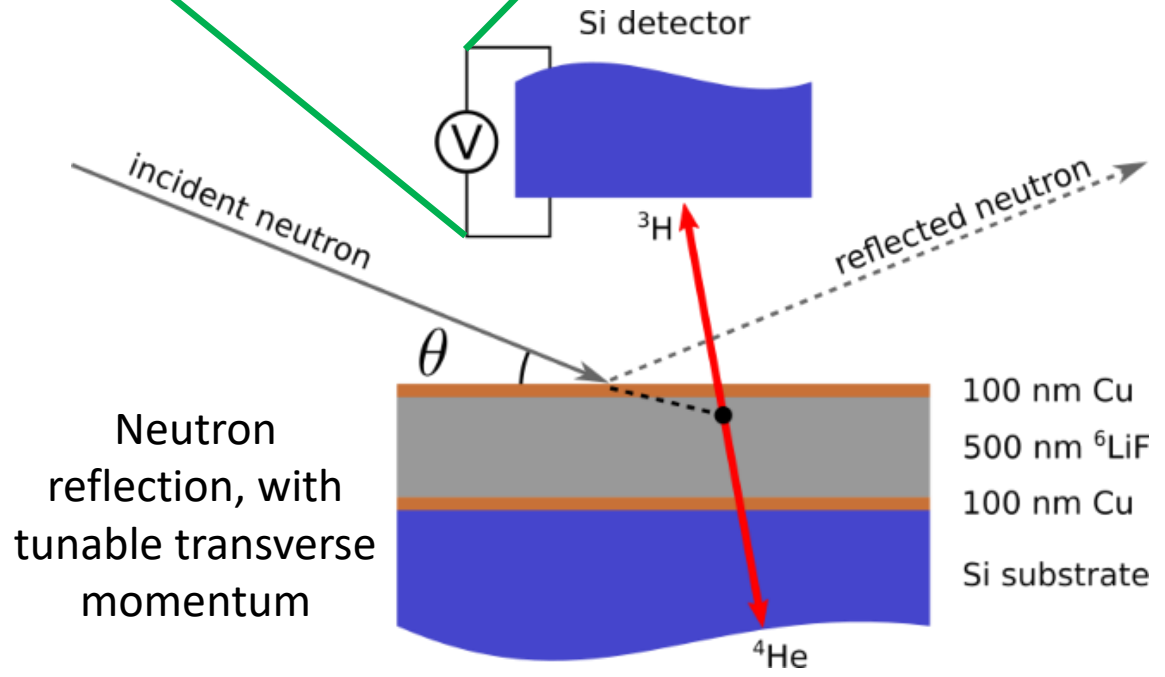
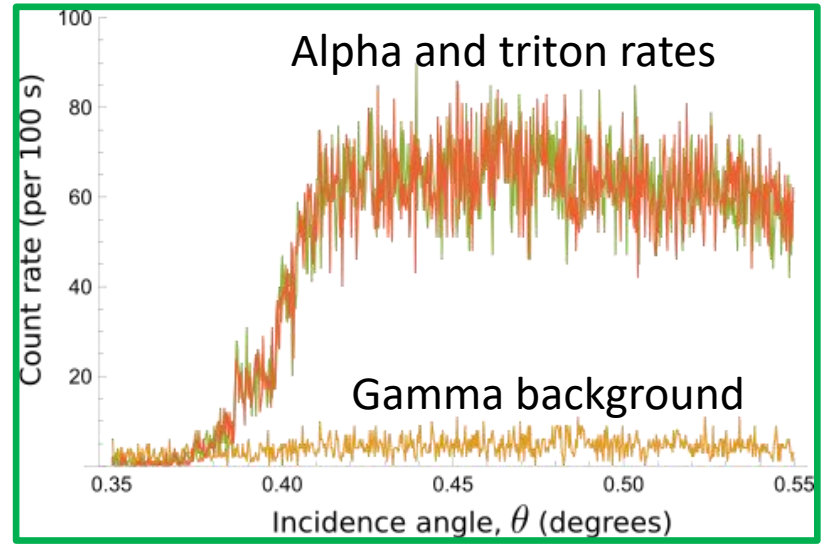
UCN out



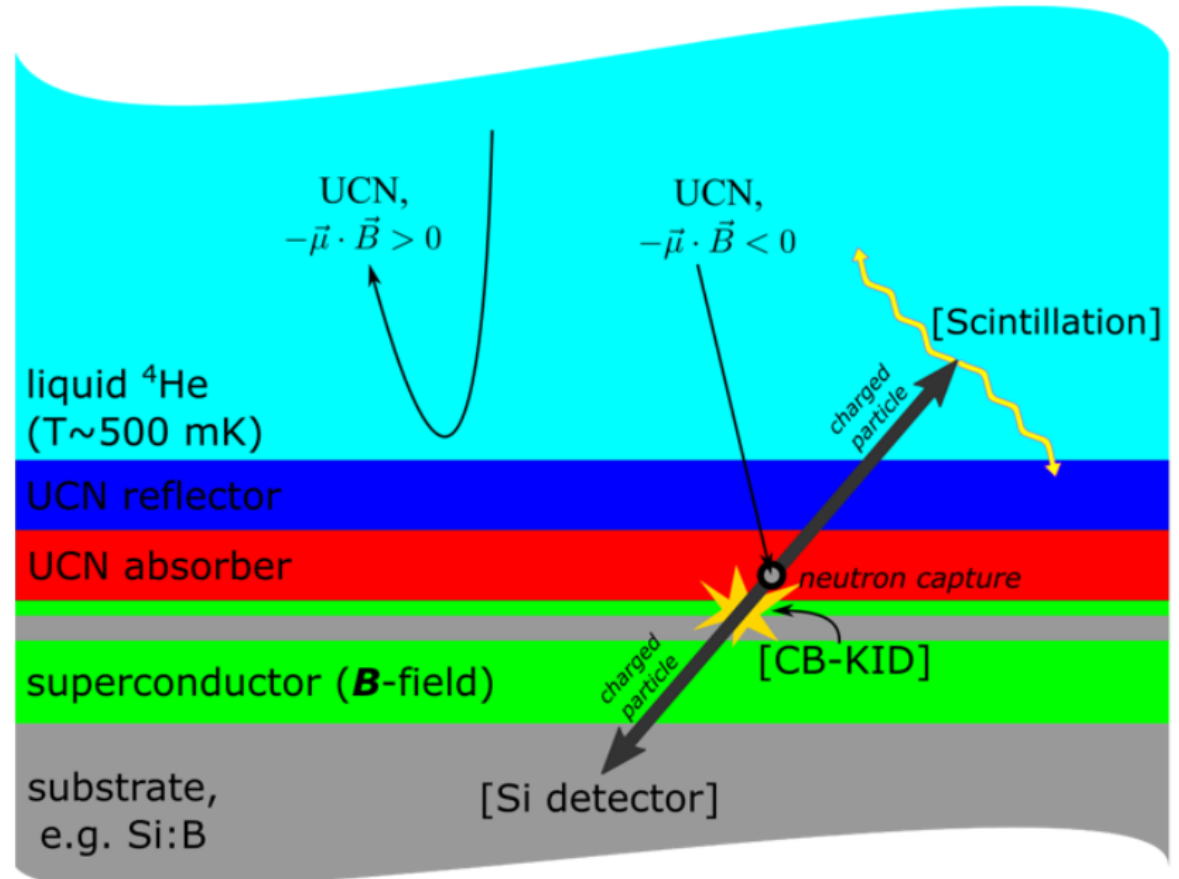
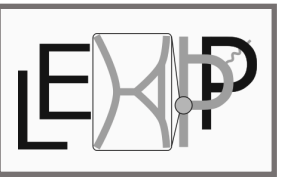
Proof-of-principle: quantum sensing



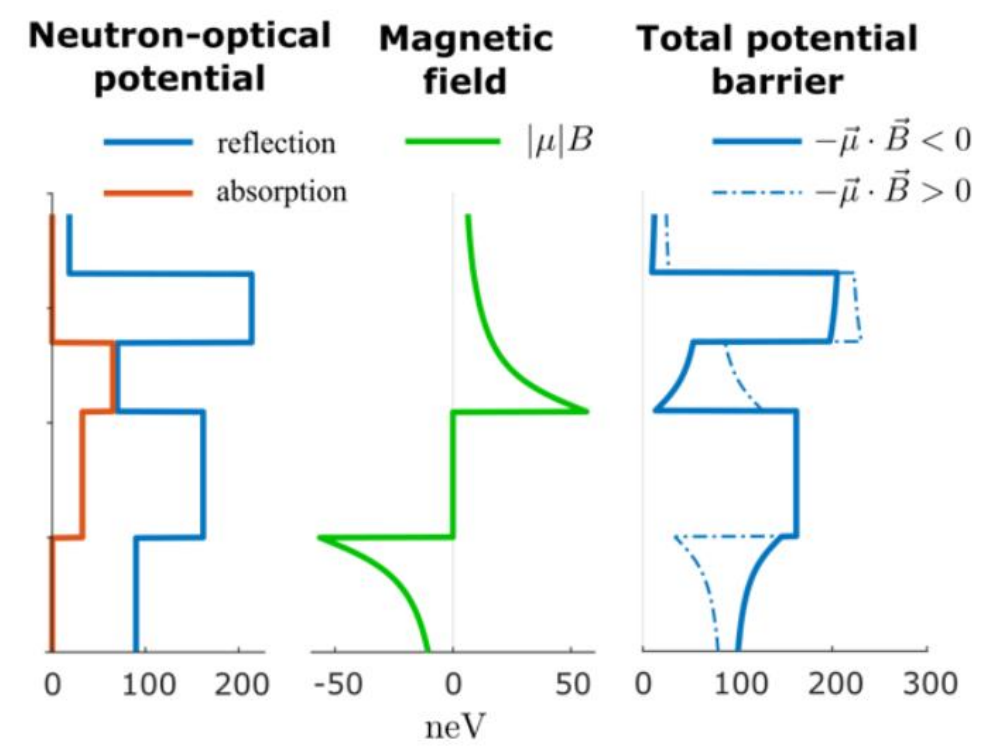
← Capture products counted with Si-detector →



"Quantum Sensing" for Neutrons

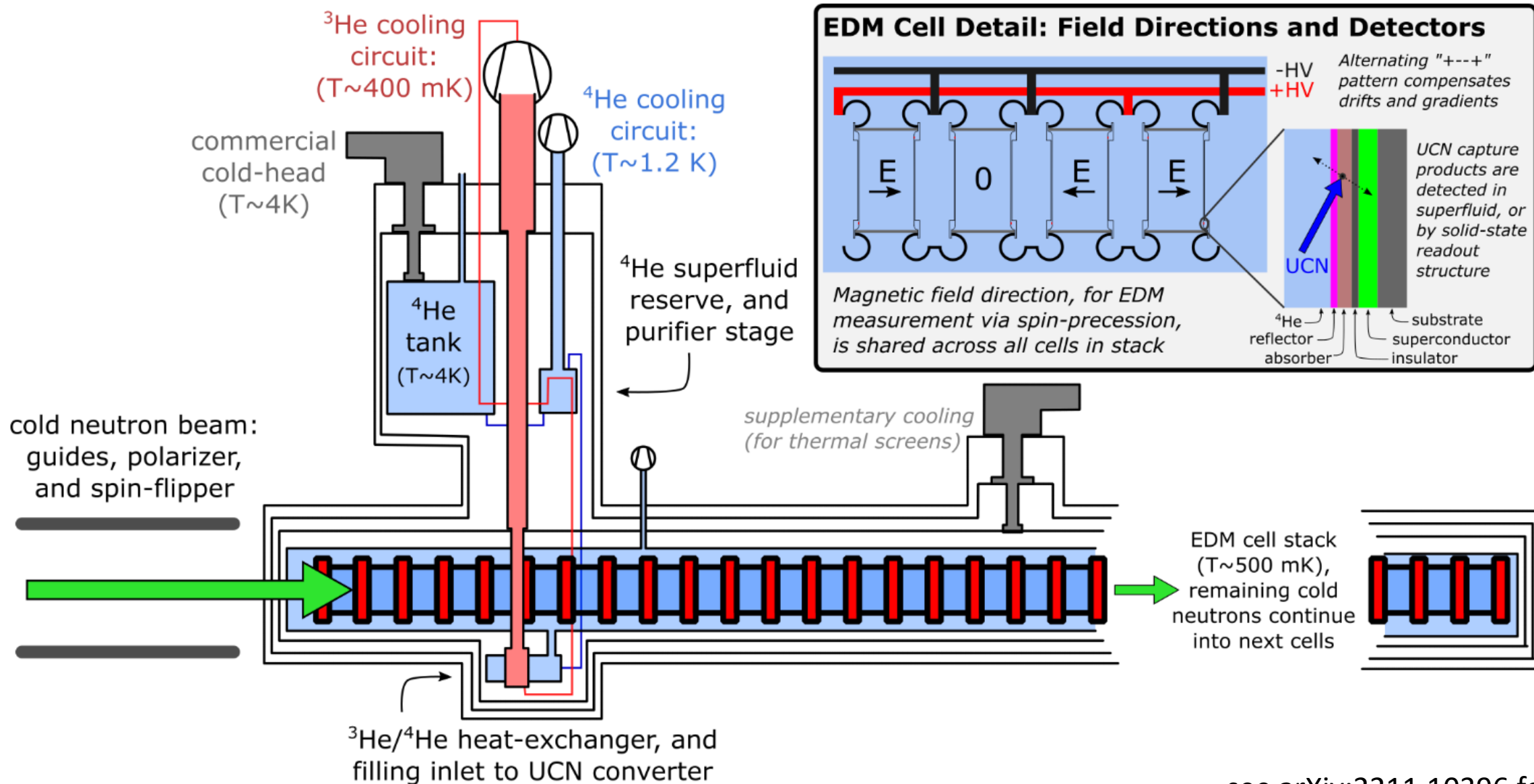
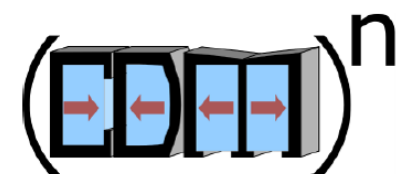


Spin-dependent shifts of the wall potential
(magnetic field increased for visibility)

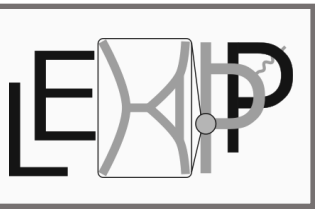




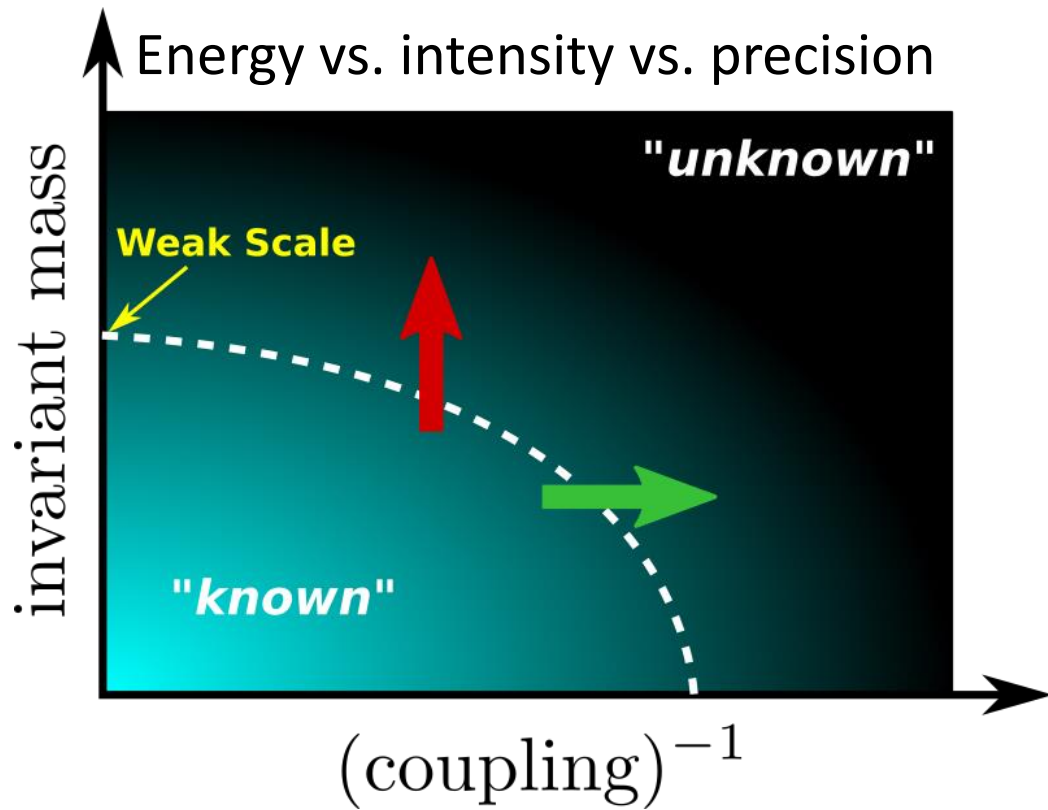
The next generation... scaling up!



...see arXiv:2211.10396 for details



Thematic Recap



1

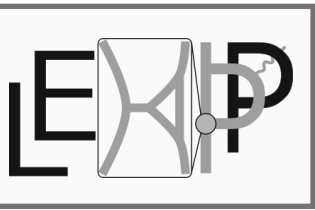
Statistics are the first key!

2

Observation time is the second

3

... and yes, it finally makes sense to follow the green arrow!



Questions?

EXPERIMENT

OUR NEW ~~TELESCOPE~~ WILL
ANSWER TWO KEY QUESTIONS:

- 1) WHY IS THERE ALL THIS MATTER?
- 2) CAN WE DO ANYTHING ABOUT IT?



what-if.xkcd.com

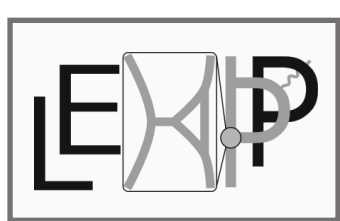
Special thanks to:

V. Cirigliano, J. de Vries, U. Schmidt

PI mechanical workshop
PI technical design office

Institut Laue-Langevin, NPP & SANE
BNC GINA team and user support

PanEDM collaboration
HeXe collaboration



Seeking students and Post-Docs!

WE WANT TO HIRE YOU TO
WRITE ON OUR COMPUTERS.

WE CAN OFFER YOU A
BUNCH OF PAYCHECKS!

THERE ARE
GHOSTS HERE.



xkcd.com

EDMs in the SM do not vanish

- CP violation from three sources (ignoring neutrinos):

$$\mathcal{L}_{\text{CPV}} = \mathcal{L}_{\text{CKM}} + \mathcal{L}_{\bar{\theta}} + \mathcal{L}_{\text{BSM}}$$

- CKM CP-violation (Standard Model):

$$\mathcal{L}_{\text{CKM}} = -\frac{ig_2}{\sqrt{2}} \sum_{p,q} V^{pq} \bar{U}_L^p W^+ D_L^q + \text{H.c.}$$

- Strong CP-violation (Standard Model):

$$\mathcal{L}_{\bar{\theta}} = -\frac{\alpha_S}{16\pi^2} \bar{\theta} \text{Tr}(G^{\mu\nu} \tilde{G}_{\mu\nu})$$

details:

Rev. Mod. Phys. **91**, 015001 (2019)

Phys. Rev. C **91**, 035502 (2015)

Prog. Part. Nucl. Phys. **71**, 21 (2013)

EDMs in the SM do not vanish

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$$\mathcal{L}_{\text{CKM}} = -\frac{ig_2}{\sqrt{2}} \sum_{p,q} V^{pq} \bar{U}_L^p W^+ D_L^q + \text{H.c.}$$

- Strong CP-violation (Standard Model)*:

$$\mathcal{L}_{\bar{\theta}} = -\frac{\alpha_S}{16\pi^2} \bar{\theta} \text{Tr}(G^{\mu\nu} \tilde{G}_{\mu\nu})$$

details:

Rev. Mod. Phys. **91**, 015001 (2019)

Phys. Rev. C **91**, 035502 (2015)

Prog. Part. Nucl. Phys. **71**, 21 (2013)

*recently called into question: arXiv:2205.15093, 2001.07152, 1912.03941, 2106.11369

Effective Field Theory for EDMs

General Effective Lagrangian:

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \frac{C^{(5)}}{\Lambda} O^{(5)} + \sum_i \frac{C_i^{(6)}}{\Lambda^2} O_i^{(6)} + \dots$$

Global Analysis: T. Chupp, M. Ramsey-Musolf
Rev. Mod. Phys. **91**, 015001 (2019)
Phys. Rev. C **91**, 035502 (2015)

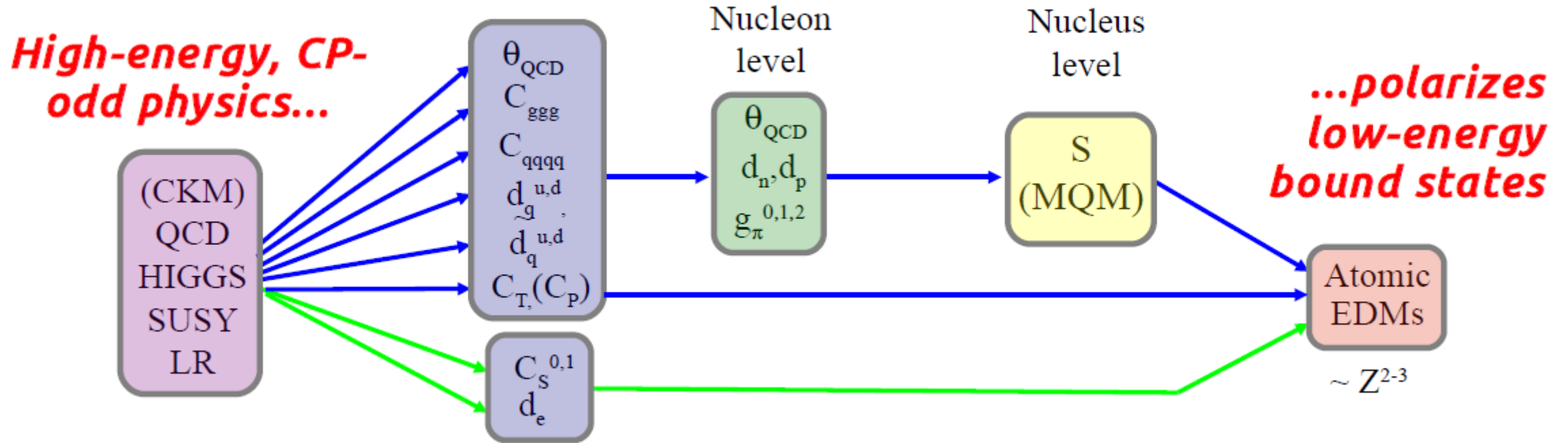
Dimension-Six terms for the neutron:

$$\begin{aligned} \mathcal{L}_{\text{eff}}^{(6)} = & -\frac{i}{2} \sum_{l,q} d_q \bar{q} \sigma_{\mu\nu} \gamma^5 F^{\mu\nu} q \\ & -\frac{i}{2} \sum_q \tilde{d}_q g_s \bar{q} \sigma_{\mu\nu} \gamma^5 G^{\mu\nu} q \\ & + d_W \frac{g_s}{6} G \tilde{G} G + \sum_i C_i^{(4f)} O_i^{(4f)} \end{aligned}$$

Prog. Part. Nucl. Phys. **71**, 21 (2013)

Wilson coefficient	Operator (dimension)	Number
$\bar{\theta}$	Theta term (4)	1
δ_e	Electron EDM (6)	1
$\text{Im } C_{\ell e q}^{(1,3)}, \text{Im } C_{\ell e q d}$	Semi-leptonic (6)	3
δ_q	Quark EDM (6)	2
$\tilde{\delta}_q$	Quark chromo EDM (6)	2
$C_{\tilde{G}}$	Three-gluon (6)	1
$\text{Im } C_{quqd}^{(1,8)}$	Four-quark (6)	2
$\text{Im } C_{\varphi ud}$	Induced four-quark (6)	1
Total		13

Interpreting EDM bounds



neutron: $\bar{d}_n^{sr}, \bar{g}_{\pi}^{(0)}, (\bar{g}_{\pi}^{(1)})$

diamagnetic: $\bar{g}_{\pi}^{(0,1)}, C_T^{0,1}$

paramagnetic: $d_e, C_S^{(0)}$

$$\mathcal{L}_{\pi NN} = \bar{N} \left[\bar{g}_{\pi NN}^{(0)} \vec{\tau} \cdot \vec{\pi} + \bar{g}_{\pi NN}^{(1)} \pi^0 + \bar{g}_{\pi NN}^{(2)} (3\tau_3 \pi^0 - \vec{\tau} \cdot \vec{\pi}) \right] N$$

$$\mathcal{L}_T = \frac{8G_F}{\sqrt{2}} \bar{e} \sigma^{\mu\nu} e \nu_{\nu} \bar{N} \left[C_T^{(0)} + C_T^{(1)} \tau_3 \right] S_{\mu} N$$

$$\mathcal{L}_S = -\frac{G_F}{\sqrt{2}} \bar{e} i \gamma_5 e \bar{N} \left[C_S^{(0)} + C_S^{(1)} \tau_3 \right] N$$

It's not so simple after all...

- **Schiff's theorem:** the field due to an EDM induces a displacement of the bound charges, which exactly cancels it*

$$H_0 = \sum \frac{p^2}{2m} + U(\mathbf{r})$$

Hamiltonian of the charge-system (no EDM)

*Schiff: *Phys. Rev.* **132**, 2194 (1963)
J. Engel: elegant formulation used here

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*Add constituent EDMs
As a perturbation...*

$$\mathbf{d}_{\text{tot}} = \sum_i \mathbf{d}_i \quad (\text{sum over constituents})$$

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(sum over constituents)

$$\begin{aligned} H &= H_0 - \sum \mathbf{d} \cdot \mathbf{E} \\ &= H_0 + \sum \mathbf{d} \cdot \frac{\nabla U(\mathbf{r})}{q} \\ &= H_0 + \sum \frac{i}{q} [\mathbf{d} \cdot \mathbf{p}, H_0] \end{aligned}$$

Now see what effect this has...

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Eigenstates receive an energy shift due to the perturbation:

$$\begin{aligned} |0\rangle \rightarrow |\tilde{0}\rangle &= |0\rangle + \sum_n \frac{|n\rangle \langle n| \sum \frac{i}{q} [\mathbf{d} \cdot \mathbf{p}, H_0] |0\rangle}{E_0 - E_n} \\ &= \left(1 + \sum \frac{i}{q} \mathbf{d} \cdot \mathbf{p} \right) |0\rangle \end{aligned}$$

It's not so simple after all...

- What is the total, observable, dipole moment after this shift?

$$\begin{aligned}\tilde{\mathbf{d}} &= \sum \mathbf{d} + \langle \tilde{0} | \sum q\mathbf{r} | \tilde{0} \rangle \\ &= \sum \mathbf{d} + \langle \tilde{0} | \left(1 - \sum \frac{i}{q} \mathbf{d} \cdot \mathbf{p} \right) \sum q\mathbf{r} \left(1 + \sum \frac{i}{q} \mathbf{d} \cdot \mathbf{p} \right) | \tilde{0} \rangle \\ &= \sum \mathbf{d} + i \langle 0 | \left[\sum q\mathbf{r}, \sum \frac{1}{q} \mathbf{d} \cdot \mathbf{p} \right] | 0 \rangle \\ &= \sum \mathbf{d} - \sum \mathbf{d} \\ &= 0\end{aligned}$$

But some details can save us!

- Schiff's theorem assumes:

- pointlike particles → *incorrect for nuclei*

$$\mathbf{S} = \frac{1}{10} \langle r^2 \mathbf{d} \rangle - \frac{1}{6Z} \langle r^2 \rangle \langle \mathbf{d} \rangle$$

...see Prog. Part. Nucl. Phys. **71**, 21 (2013)

- non-relativistic treatment → *incorrect for atomic electrons*

$$U_{\text{lab}} = -\mathbf{d}_{\text{lab}} \cdot \mathbf{E} = -\mathbf{d}_{\text{rest}} \cdot \mathbf{E} + \frac{\gamma}{1 + \gamma} (\boldsymbol{\beta} \cdot \mathbf{d})(\boldsymbol{\beta} \cdot \mathbf{E})$$

...see American Journal of Physics **75**, 532 (2007)

A Key Systematic Effect

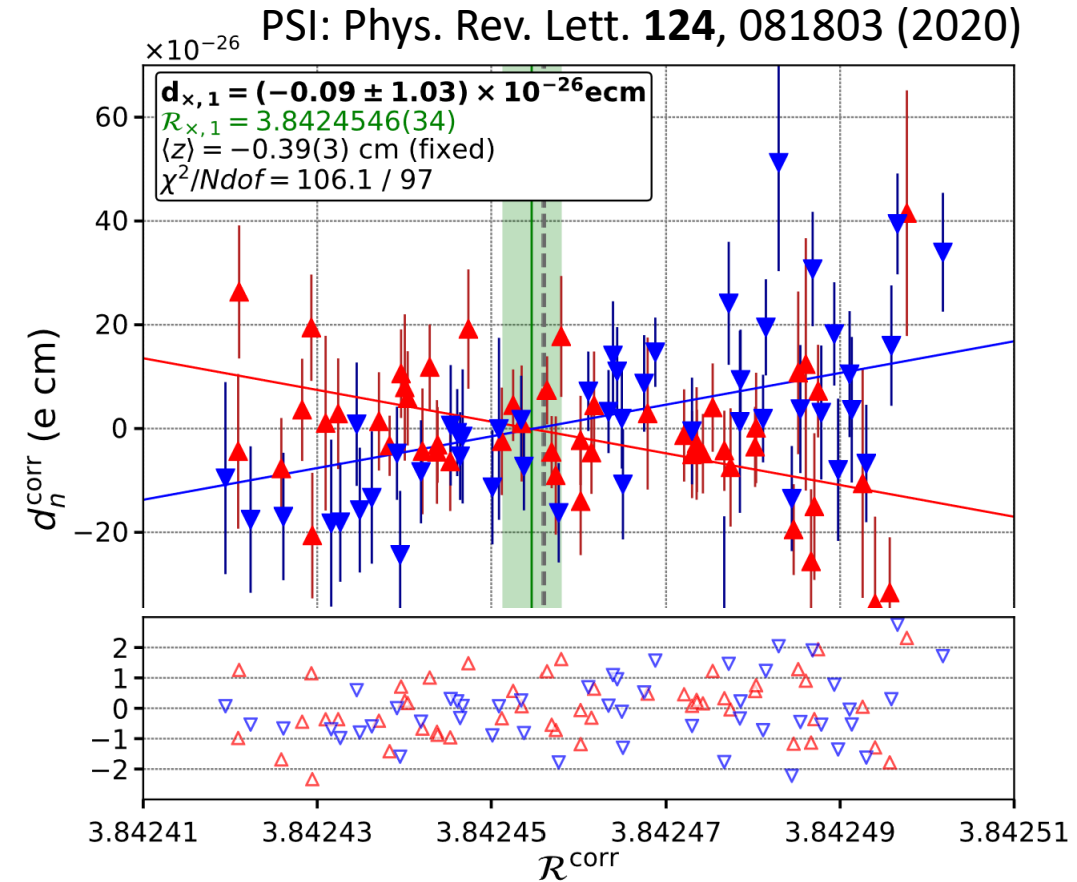
Motional magnetic field: $B = \frac{1}{c^2} \mathbf{E} \times \mathbf{v}$

Magnetic field gradients: $B_{\perp} = -\frac{R}{2} \frac{\partial B_z}{\partial z}$

Leading false EDM effect at 2nd order:

$$d_{\text{false}}^{(\text{slow})} = -\frac{\hbar}{4c^2} \left(\frac{v_{\perp}}{B_z} \right)^2 \frac{\partial B_z}{\partial z} \quad \text{Adiabatic (UCN)}$$

$$d_{\text{false}}^{(\text{fast})} = -\frac{\hbar}{8c^2} \gamma^2 R^2 \frac{\partial B_z}{\partial z} \quad \text{Diabatic (atomic magnetometer)}$$



$$R^{\text{corr}} = \left| \frac{\gamma_{\text{UCN}}}{\gamma_{\text{Hg}}} \right| \frac{1 + \delta_T + \delta_{\text{Earth}} + \delta^{\text{EDM}} + \delta_{\text{false}}^{\text{EDM}} + \dots}{1 + \delta_T + \delta_{\text{Earth}}}$$

Baryon asymmetry

Consider



$$\eta_p = \frac{n_p - n_{\bar{p}}}{n_\gamma}$$

Threshold energy/temperature:

$$T \gg \frac{2 \text{ GeV}}{k_B} \sim 2 \times 10^{13} \text{ K}$$

After freeze-out:

$$n_\gamma = \int_0^\infty \frac{d\omega}{c^3 \pi^2} \frac{\omega^2}{e^{\frac{\hbar\omega}{k_B T}} - 1} = 16\pi \left(\frac{k_B T}{hc} \right)^3 \zeta(3)$$

Statistics of the photon gas

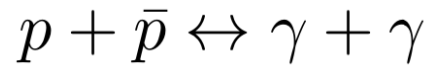
$$N_\gamma = \int d\omega \frac{g(\omega)}{e^{\frac{\hbar\omega}{k_B T}} - 1}$$

$$g(\omega) = \frac{V \omega^2}{c^3 \pi^2}$$

$$\langle E \rangle = \int_0^\infty \hbar\omega \frac{g(\omega)}{e^{\frac{\hbar\omega}{k_B T}} - 1}$$

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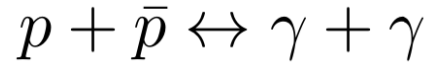
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$$\frac{1}{T} = \left(\frac{\partial S}{\partial E} \right)_V$$

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Photons:	411 / cm³	(observed)
Protons:	10 / m³	(predicted)
	0.2 / m³	(observed)

