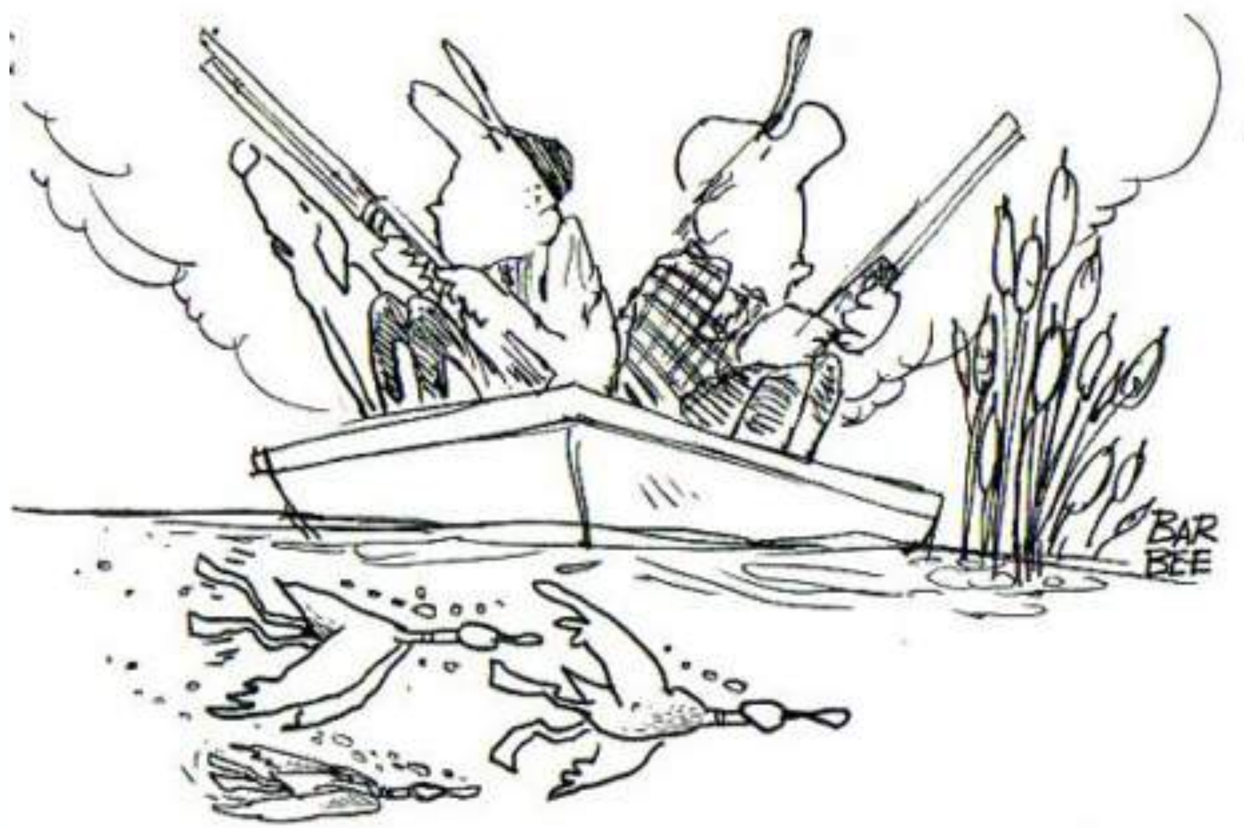
Hunting the Flavon

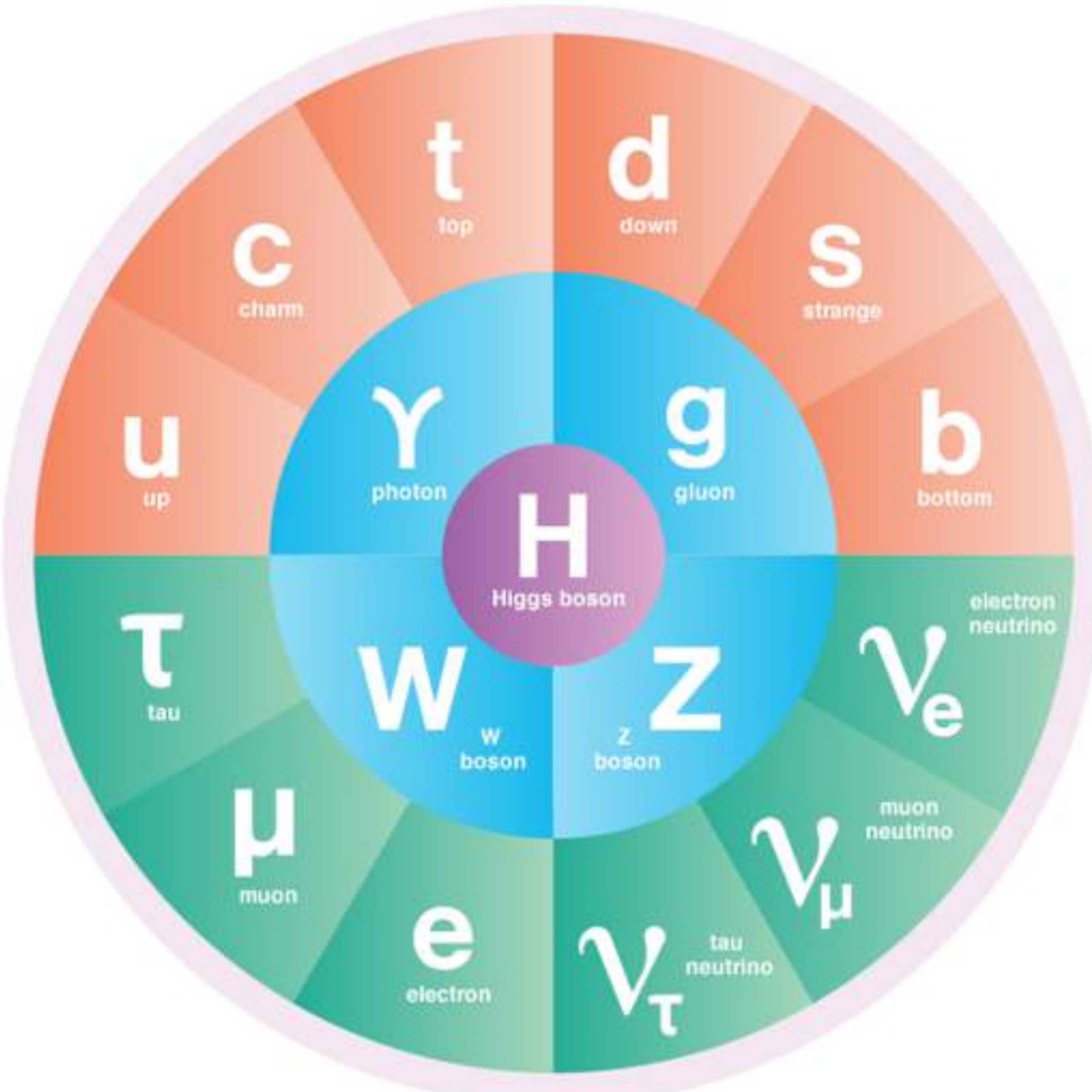
Martin Bauer



[MB, Gemmler, Carena, JHEP 1511, 016 (2015)]
[MB, Schell, Plehn, PRD 94, no. 5, 056003 (2016)]
[MB, Gemmler, Carena, 1512.03458]

Teilchen , Oct 24th 2016

The Standard Model of Particle Physics



Quarks

Spin 1/2
Charge 2/3 : Up type
Charge -1/3: Down type

Leptons

Spin 1/2
Charge -1: e, μ, τ
Charge 0: Neutrinos

Gauge Bosons

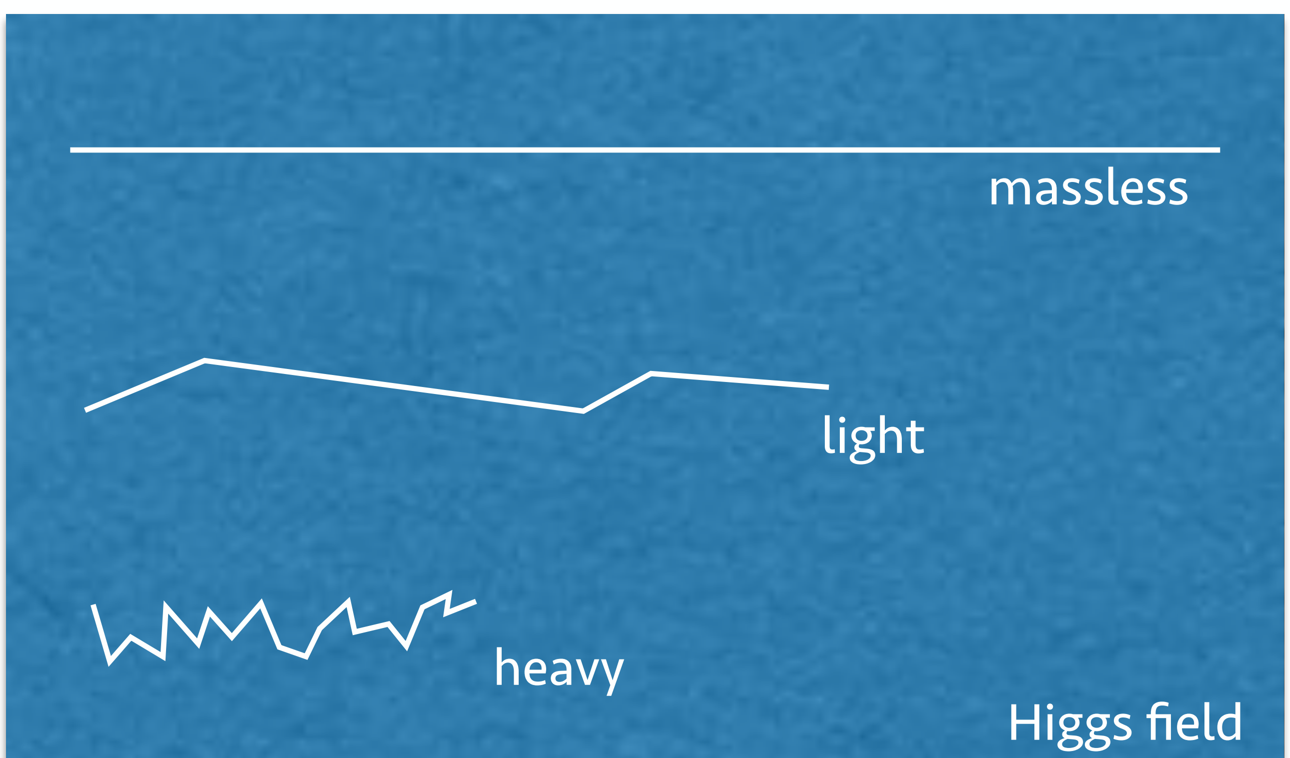
Spin 1
Charge 0 : g, γ, Z
Charge ±1: W

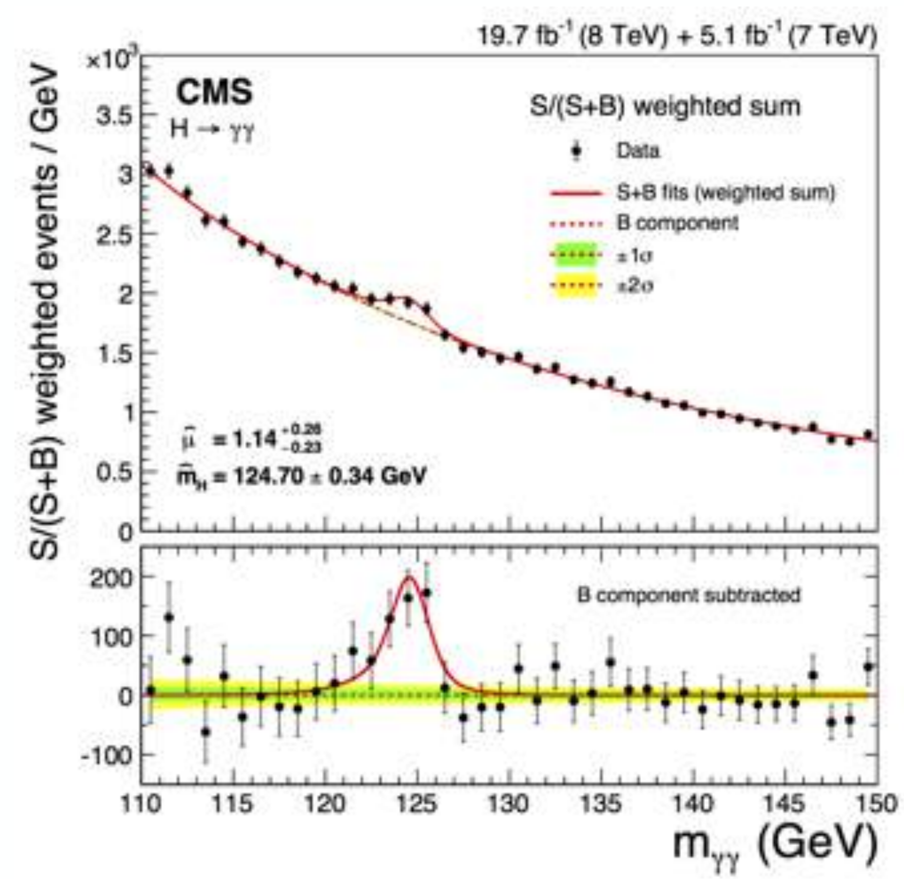
Higgs Boson

Spin 0
Charge 0

Why are elementary particles massive?

The Higgs mechanism explains why fundamental particles are massive.

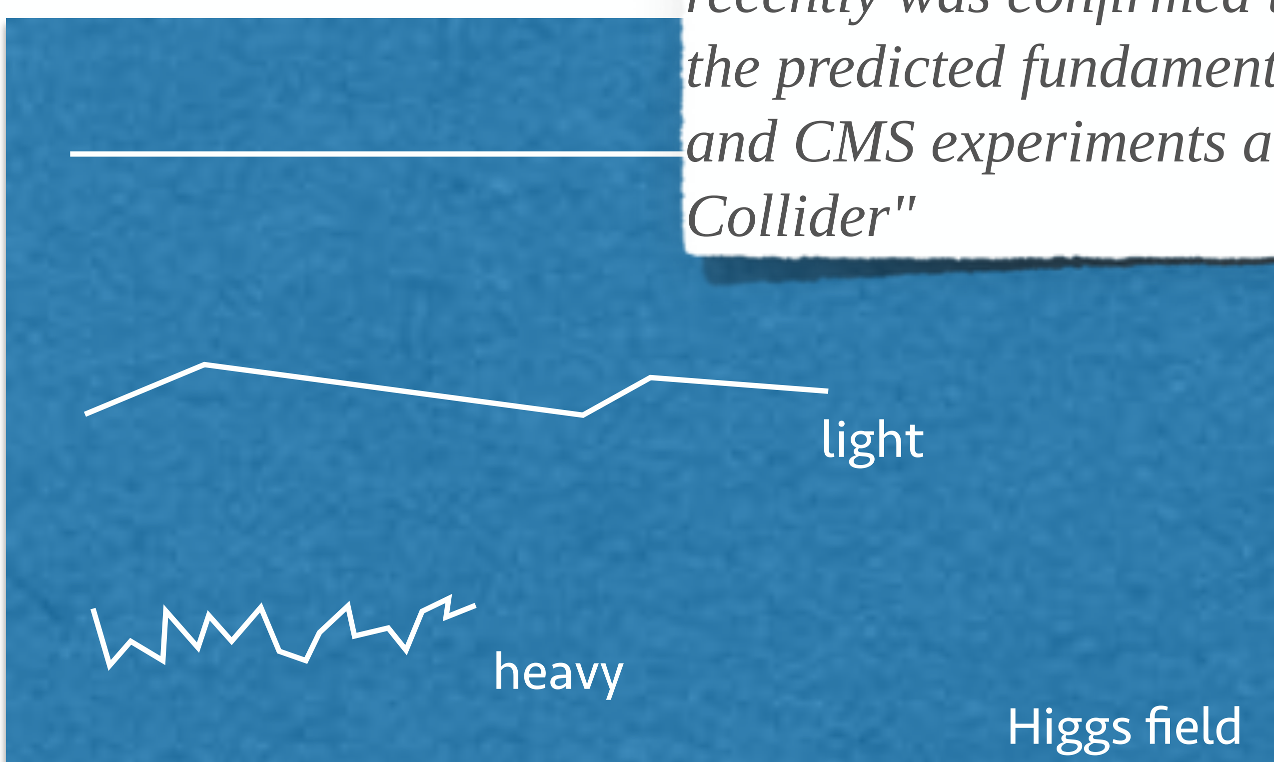


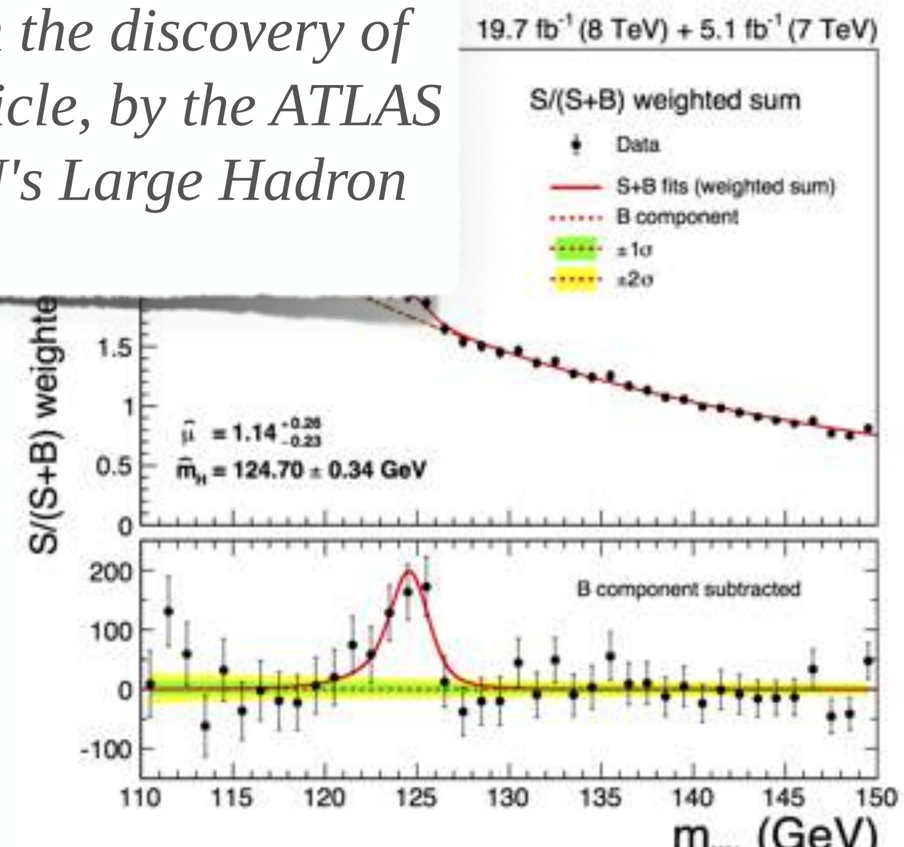


Why are elementary particles massive?

The Higgs mechanism explains why fundamental particles are massive.

The Nobel Prize in Physics 2013 was awarded jointly to François Englert and Peter W. Higgs *"for the theoretical discovery of* ***a mechanism that contributes to our understanding of the origin of mass of subatomic particles****, and which recently was confirmed through the discovery of the predicted fundamental particle, by the ATLAS and CMS experiments at CERN's Large Hadron Collider"*

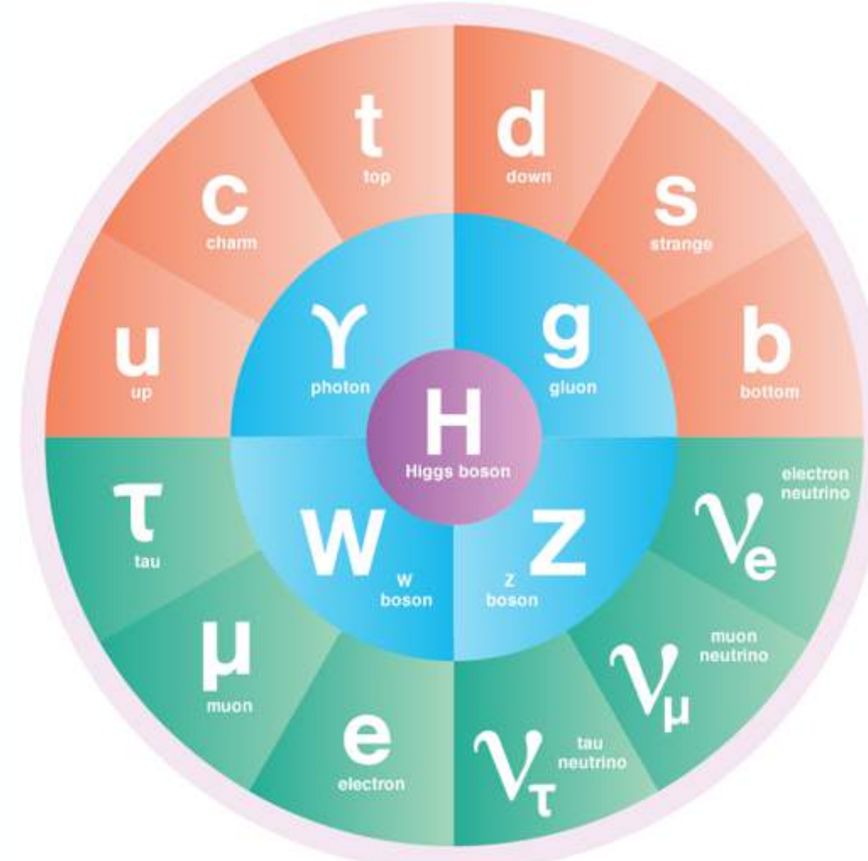


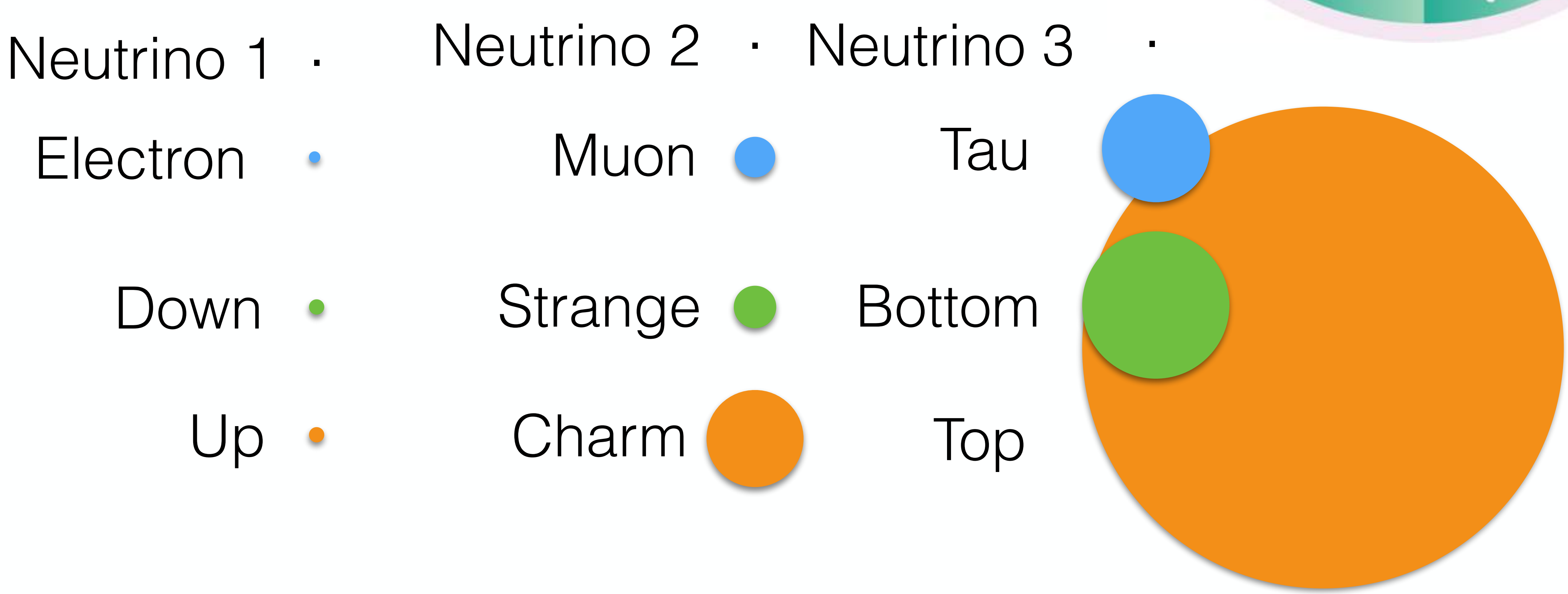


Why are fermion masses so different?

Fermion mass hierarchy:
At least 6 orders of magnitude

top mass = 170.000 x up mass





Take a lesson from history

Hierarchies → Fundamental Structure

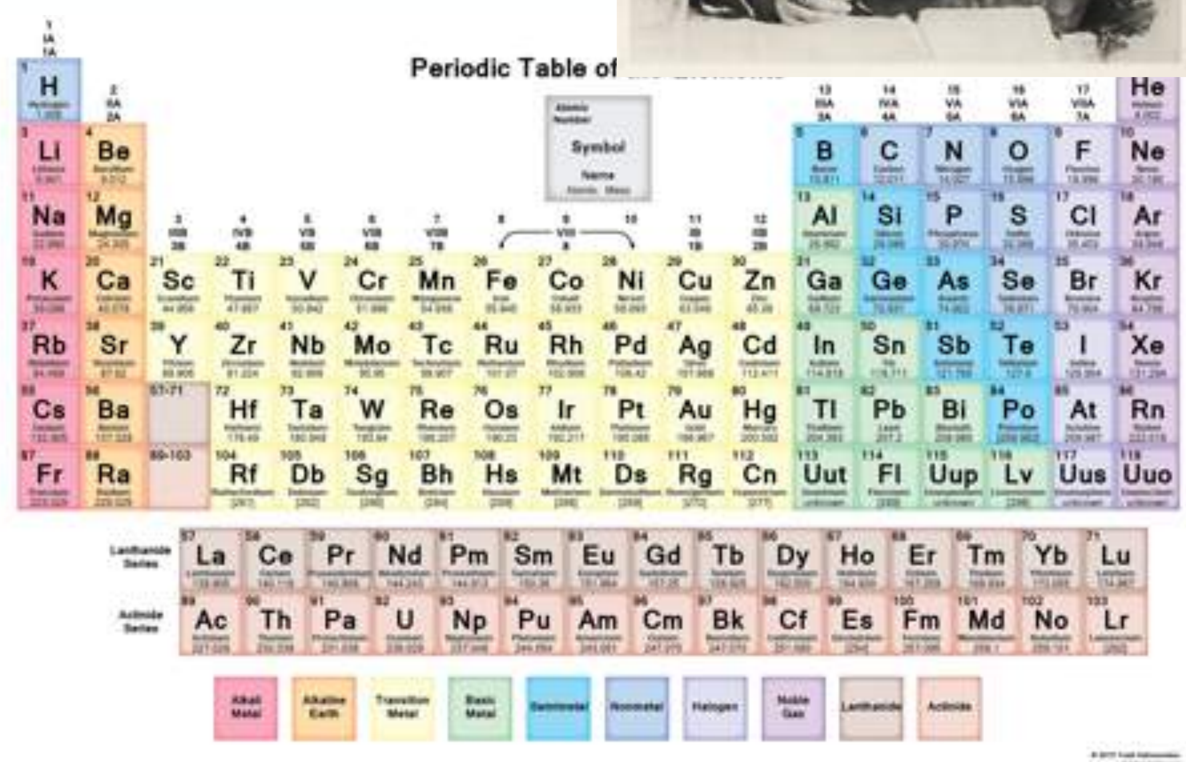
John Dalton, 1808

Take a lesson from history

Hierarchies → Fundamental Structure

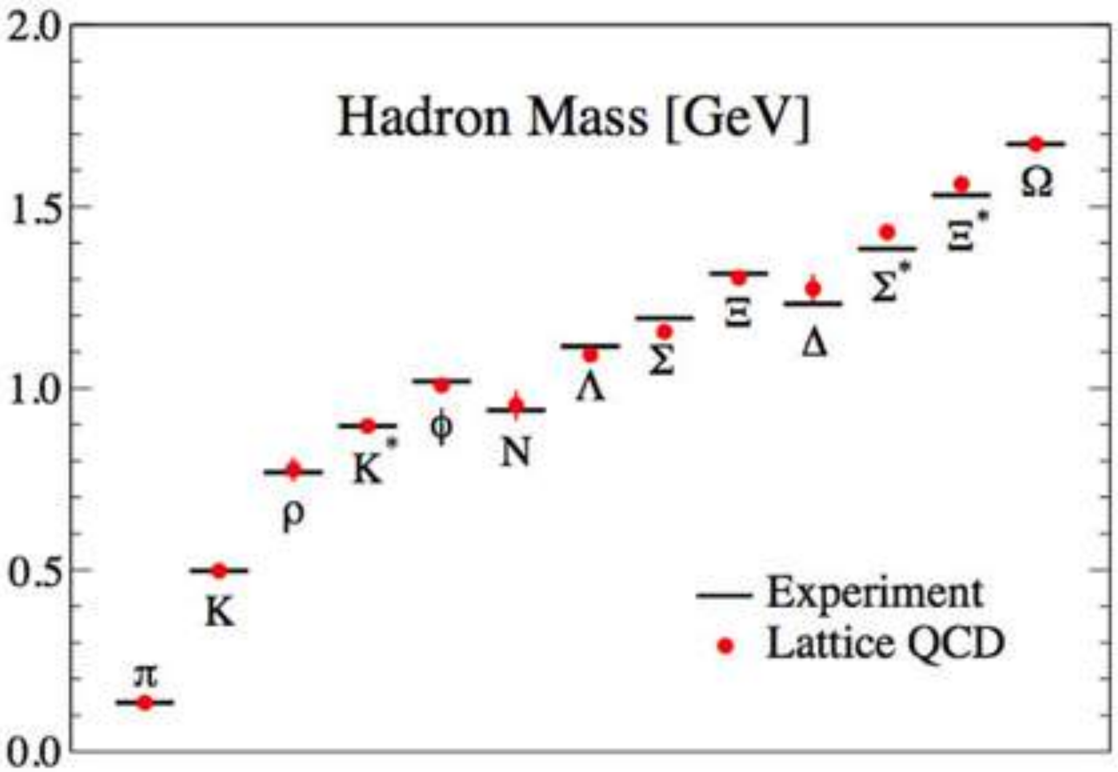
John Dalton, 1808

Dmitri Mendeleev, 1869



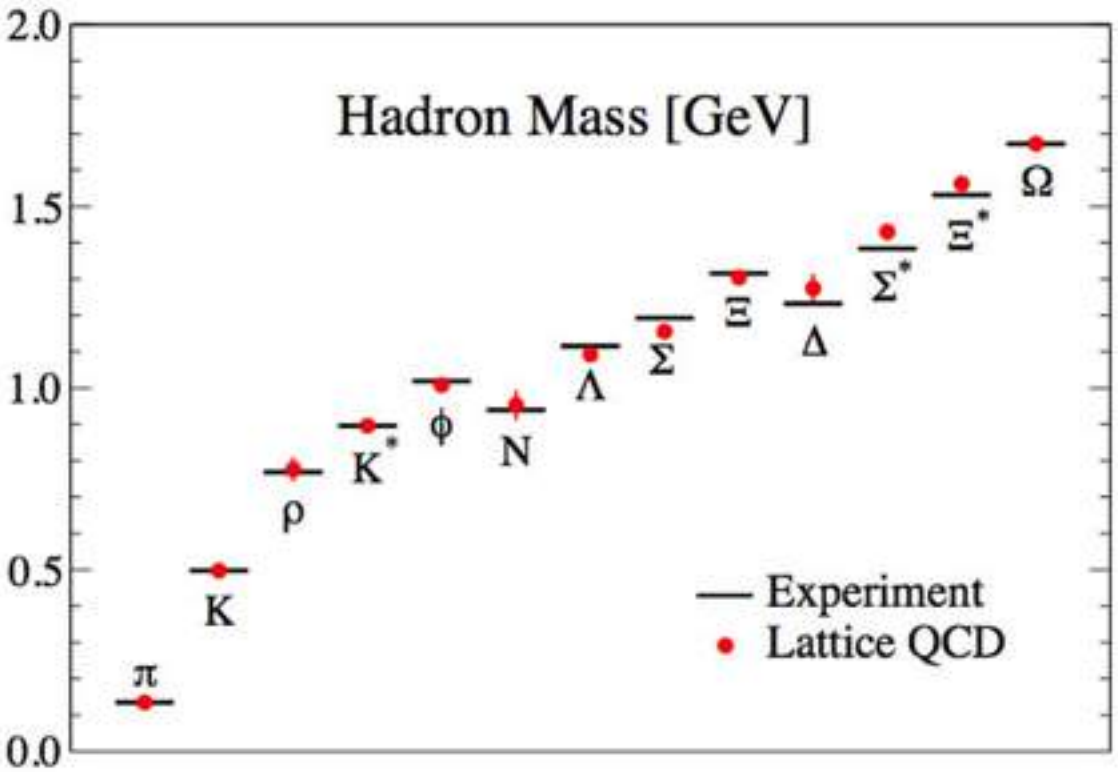
Take a lesson from history

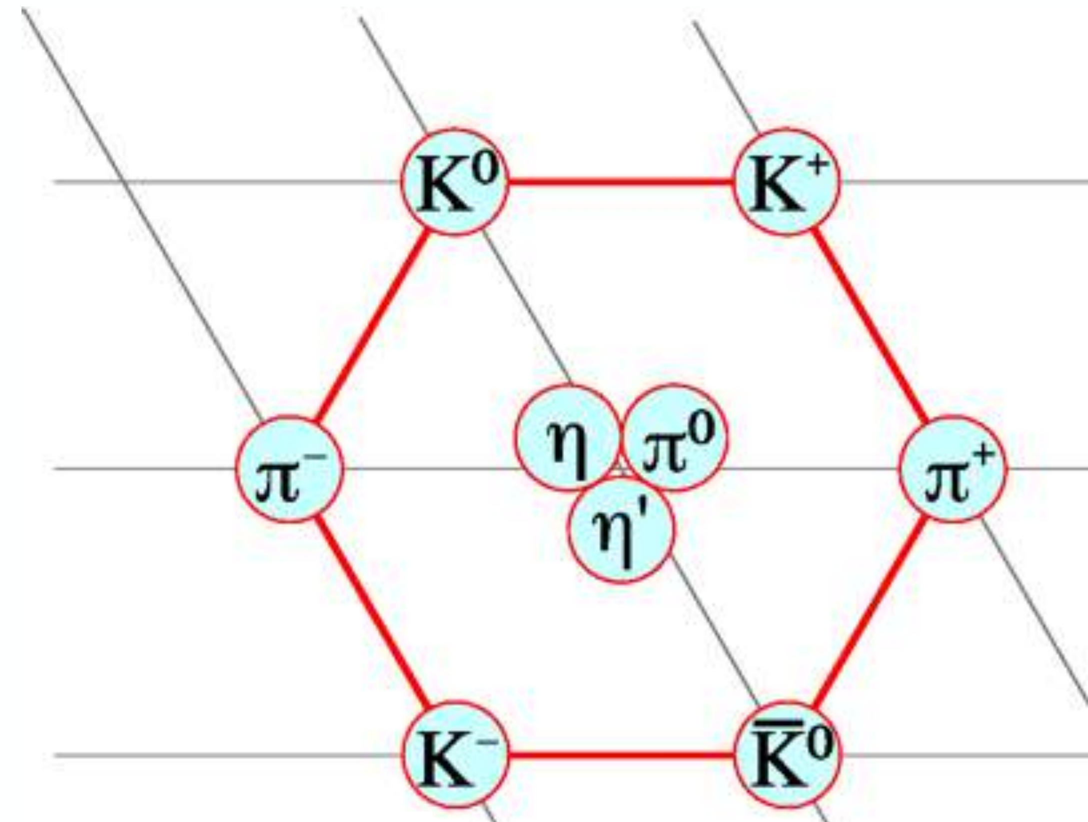
Hierarchies ⟶ Fundamental Structure



Take a lesson from history

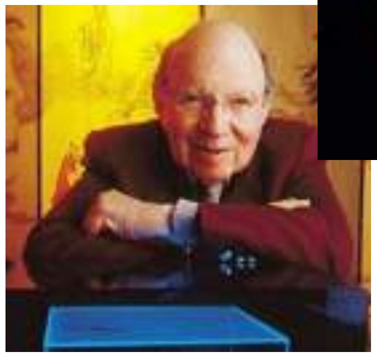
Hierarchies → Fundamental Structure





The eightfold way

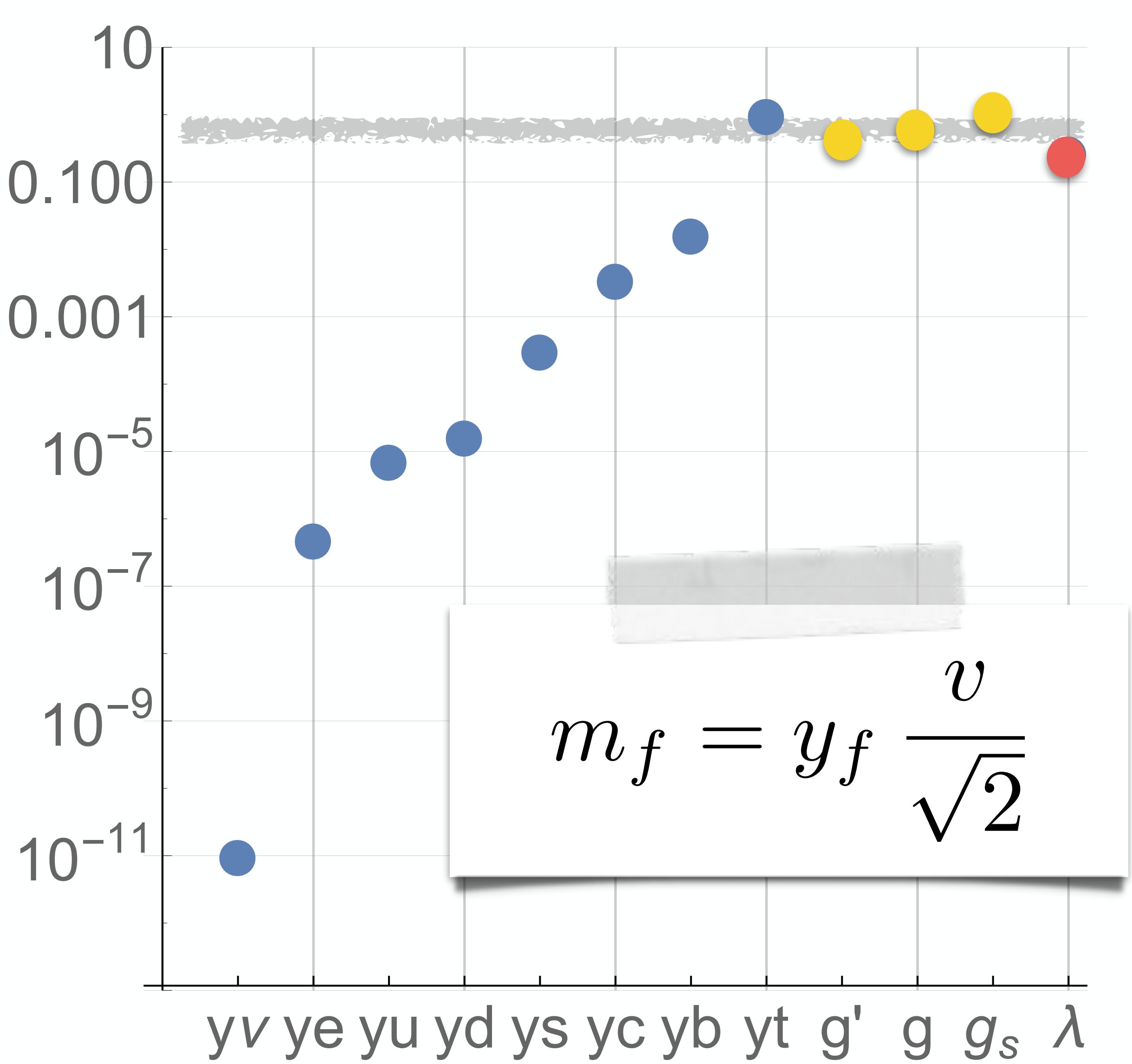
Fritzsch, Gell-Mann

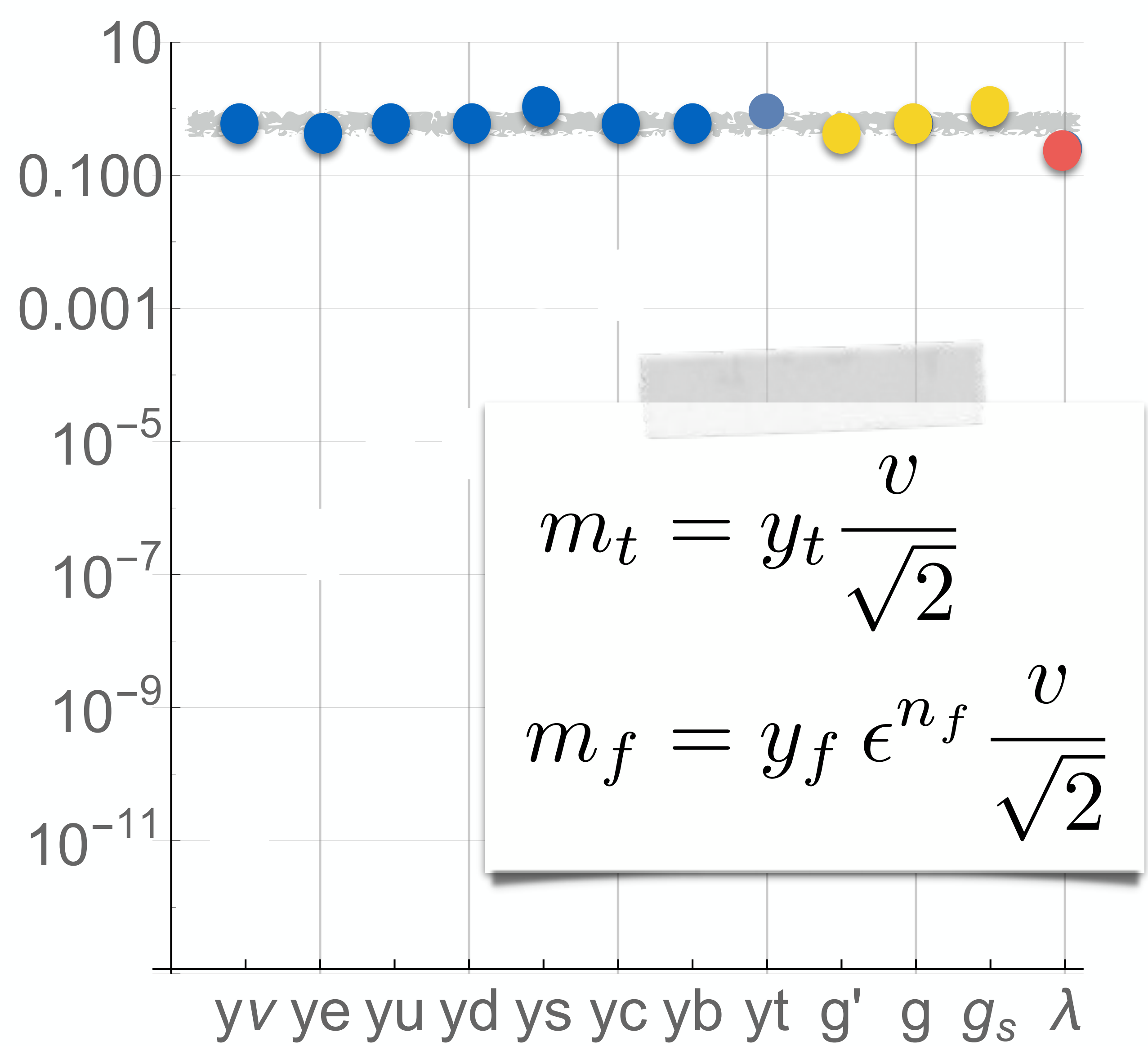


Ne’eman

Why are elementary particles so different?

$$\mathcal{L} = \mu_h^2 H^\dagger H + \lambda(H^\dagger H)^2 + \mathcal{L}_{\text{gauge}} + \mathcal{L}_{\text{Yukawa}} + \frac{\mathcal{L}_5}{\Lambda} + \frac{\mathcal{L}_6}{\Lambda^2} + \cdots$$

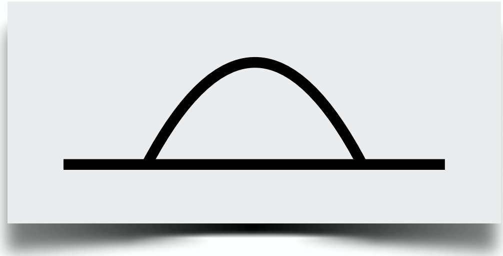




Theories of Flavor

1. Loop Induced

[Georgi, Glashow '72]

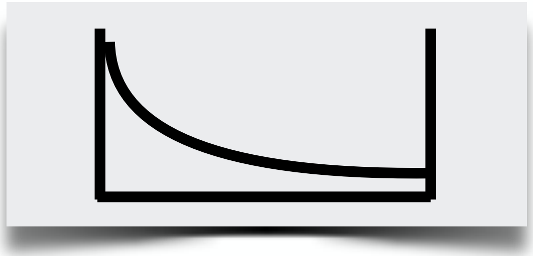


$$\epsilon \approx \frac{1}{16\pi^2} \ln\left(\frac{m_s^2}{\Lambda^2}\right)$$

2. Extra Dimensions

[Grossmann, Neubert 9912408]

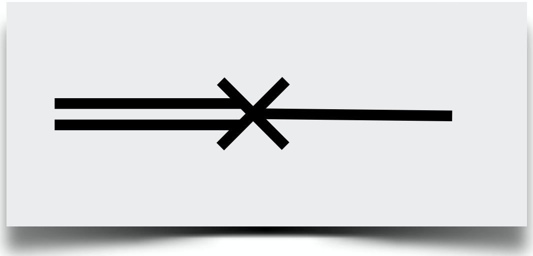
[Gherghetta, Pomarol 0003129]



$$\epsilon^n \approx e^{(c_L - c_R)\ln\frac{\Lambda_{\rm IR}}{\Lambda_{\rm UV}}}$$

3. Partial Compositeness

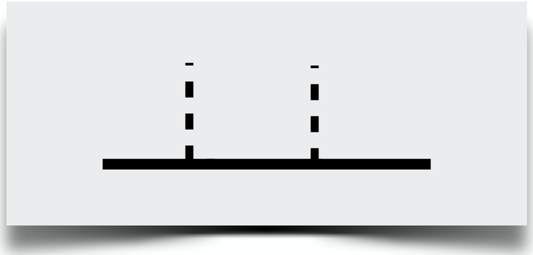
[Kaplan '91]



$$\epsilon^\gamma = \left(\frac{m}{M_B}\right)^\gamma$$

4. Froggatt Nielsen

[Froggatt, Nielsen '79]



$$\epsilon \approx \frac{f}{\Lambda}$$

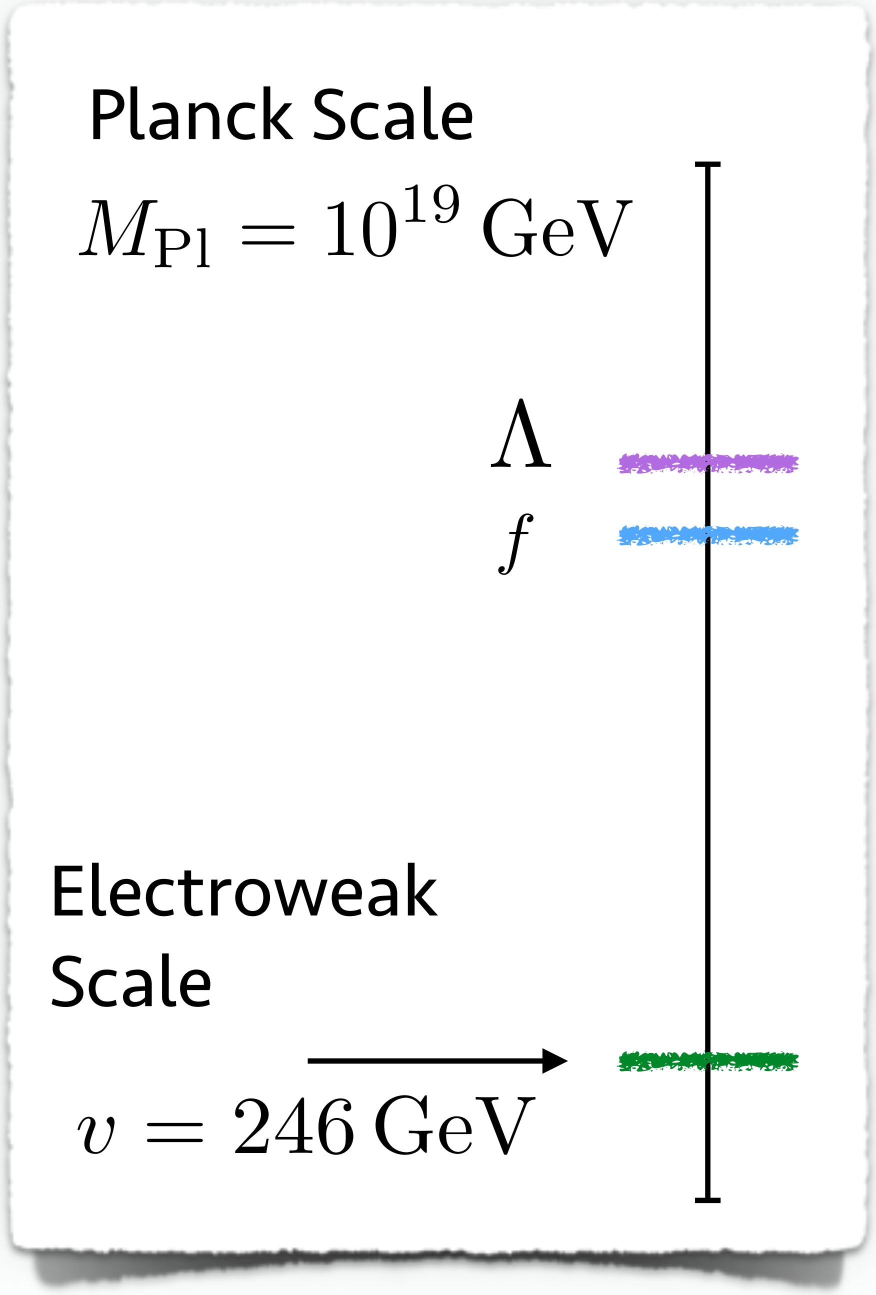
Can we discover this mechanism?

Illustration: $$y_t \bar{t} H t + y_f \left(\frac{S}{\Lambda}\right)^1 \bar{b} H b + \ldots$$

$$\langle S \rangle = f \qquad \Rightarrow \qquad y_b = \epsilon \, y_f$$

and $$\epsilon = \frac{f}{\Lambda} = 0.23$$

In general, the flavor scale can be arbitrarily high!



Why should the flavor scale be low?

- Why not?
- A link to Baryogenesis?
- A link to Dark Matter?
- It could be related to the electroweak scale

Talk by Iason Baldes 14:15
phil16

Why should the flavor scale be low?

- Why not?
- A link to Baryogenesis?
- A link to Dark Matter?
- It could be related to the electroweak scale

Flavon Potential

Scalar potential leads to a flavor breaking minimum

$$-\mathcal{L}_{\text{potential}} = -\mu_S^2\, S^\dagger S + \lambda_S\, (S^\dagger S)^2 + b\, (S^2 + S^{\dagger 2}) + \lambda_{HS}(S^\dagger S)(H^\dagger H) + V(H)\,.$$

Breaks the flavor symmetry

Two degrees of freedom

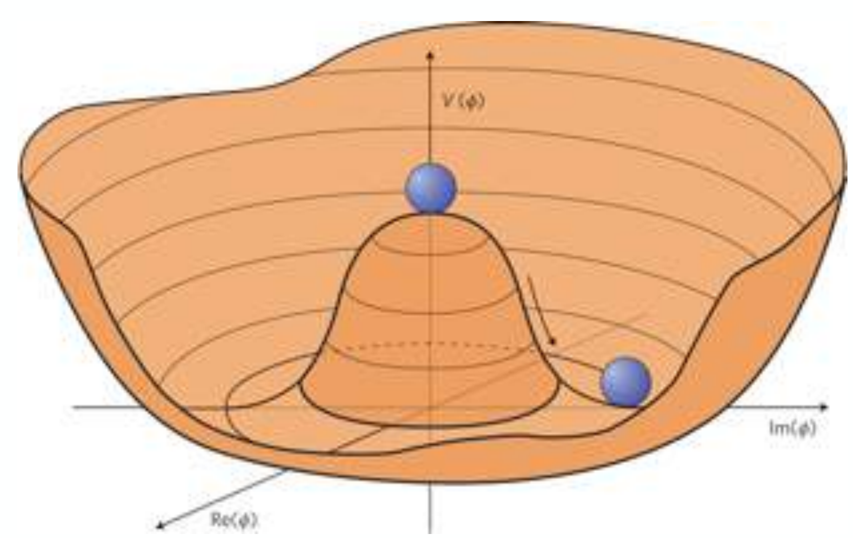
$$S(x) = \frac{f + s(x) + i\, a(x)}{\sqrt{2}}$$

With masses

$$m_s = \sqrt{\mu_s^2 - 2b}\, f$$

$$m_a = \sqrt{2b}$$

$$m_a < m_s \approx f < \Lambda$$



Yukawa Couplings

Yukawa couplings for quarks and leptons

$$\mathcal{L}_{\text{Yukawa}} = y^d_{ij} \left(\frac{S}{\Lambda}\right)^{n^d_{ij}} \overline{Q}_i \, H \, d_{R_j} + y^u_{ij} \left(\frac{S}{\Lambda}\right)^{n^u_{ij}} \overline{Q}_i \, \widetilde{H} \, u_{R_j}$$
$$+ \, y^\ell_{ij} \left(\frac{S}{\Lambda}\right)^{n^\ell_{ij}} \overline{L}_i \, H \, \ell_{R_j} + y^\nu_{ij} \left(\frac{S}{\Lambda}\right)^{n^\nu_{ij}} \overline{L}_i \, \widetilde{H} \, \nu_{R_j} + \text{h.c.}$$

Exponents are fixed by U(1) flavor charges.

$$n^d_{ij} = a_{Q_i} - a_{d_j} - a_H$$
$$n^u_{ij} = a_{Q_i} - a_{u_j} + a_H$$

Masses and Mixings

$$\epsilon = \frac{f}{\Lambda} \equiv \frac{\langle S \rangle}{\Lambda} \equiv (V_{\mathrm{CKM}})_{12} \approx 0.23$$

Quark and Lepton Masses

$$m_t \approx \frac{v}{\sqrt{2}} \qquad \frac{m_b}{m_t} \approx \epsilon^3 \qquad \frac{m_c}{m_t} \approx \epsilon^4 \qquad \frac{m_s}{m_t} \approx \epsilon^5 \qquad \frac{m_d}{m_t} \approx \epsilon^7 \qquad \frac{m_u}{m_t} \approx \epsilon^8$$

$$\frac{m_\tau}{m_t} \approx \epsilon^3 \qquad \frac{m_\mu}{m_t} \approx \epsilon^5 \qquad \frac{m_e}{m_t} \approx \epsilon^8 \qquad \frac{m_{\nu_1}}{m_t} \approx \epsilon^{24} \qquad \frac{m_{\nu_2}}{m_t} \approx \epsilon^{21} \qquad \frac{m_{\nu_3}}{m_t} \approx \epsilon^{20}$$

and Mixings

$$V_{\mathrm{CKM}} \approx \begin{pmatrix} 1 & \epsilon & \epsilon^3 \\ \epsilon & 1 & \epsilon^2 \\ \epsilon^3 & \epsilon^2 & 1 \end{pmatrix} \qquad U_{\mathrm{PMNS}} \approx \begin{pmatrix} 1 & \epsilon & \epsilon \\ \epsilon & 1 & 1 \\ \epsilon & 1 & 1 \end{pmatrix}$$

Masses and Mixings

$$\epsilon = \frac{f}{\Lambda} \equiv \frac{\langle S \rangle}{\Lambda} \equiv (V_{\rm CKM})_{12} \approx 0.23$$

Quark and Lepton Masses

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and Mixings

After fixing the ratio of scales there are 2 free parameters:

$$\boxed{m_a \quad \text{and} \quad f}$$

Flavon couplings

Flavon Couplings are dictated by this structure

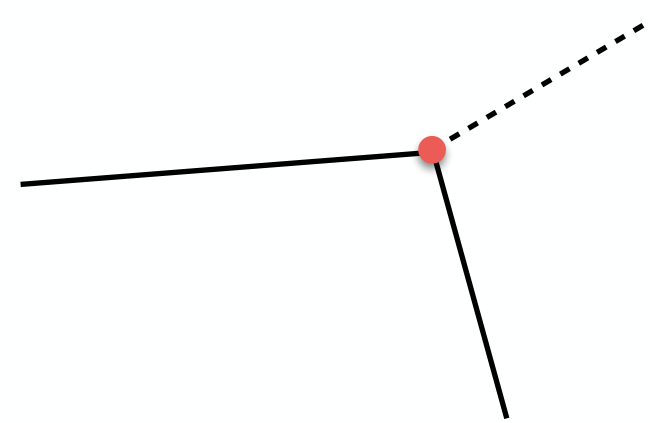
$$S(x) = \frac{f + s(x) + i\,a(x)}{\sqrt{2}}$$

$$g^u_{aij} = \frac{1}{f}\begin{pmatrix} 8m_u & \epsilon m_c & \epsilon^3 m_t \\ \epsilon^3 m_c & 4m_c & \epsilon^2 m_t \\ \epsilon^5 m_t & \epsilon^2 m_t & 0 \end{pmatrix}$$

$$g^d_{aij} = \frac{1}{f}\begin{pmatrix} 7m_d & \epsilon m_s & \epsilon^3 m_b \\ \epsilon m_s & 5m_s & \epsilon^2 m_b \\ \epsilon m_b & \epsilon^2 m_b & 3m_b \end{pmatrix}$$

$$g^\ell_{aij} = \frac{1}{f}\begin{pmatrix} 9m_e & \epsilon m_\mu & \epsilon m_\tau \\ \epsilon^3 m_\mu^3 & 5m_\mu & \epsilon^2 m_\tau \\ \epsilon^5 m_\tau & \epsilon^2 m_\tau & 3m_\tau \end{pmatrix}$$

small....potentially very small!



$$\sim \frac{m_i}{f}\epsilon$$

Flavon couplings

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$$S(x) = \frac{f + s(x) + i\, a(x)}{\sqrt{2}}$$

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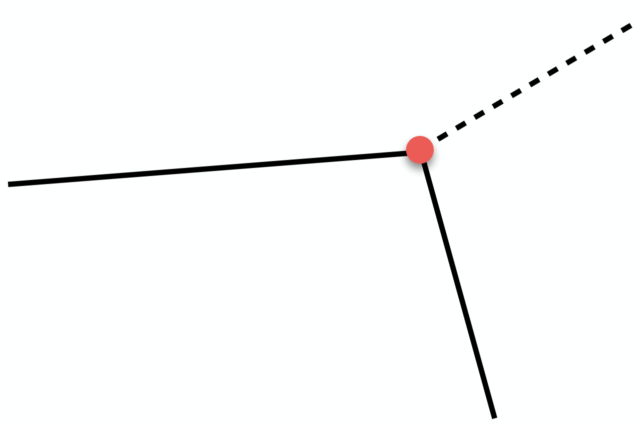
$$\boxed{g_{act} = \frac{1}{3\epsilon} g_{abb}}$$



$$\sim \frac{m_i}{f}\epsilon$$

Flavon couplings

Effects from flavon interactions lead to

- Quark Flavor Constraints
- Lepton Flavor Constraints
- Future Collider Constraints

Flavon couplings

Generate fundamental Yukawa couplings at the high scale and reproduce fermion masses and mixings in agreement with the SM.

$$m_{u_i} = (0.00138, 0.563, 150.1)\ \text{GeV}$$
$$m_{d_i} = (0.00342, 0.054, 2.29)\,\text{GeV}$$

$$|V_{\rm ckm}| = \begin{pmatrix} 0.974 & 0.226 & 0.0035 \\ 0.226 & 0.974 & 0.0388 \\ 0.011 & 0.037 & 0.999 \end{pmatrix}$$

Demand $|y_{ij}| \in [0.5, 1.5]$ with arbitrary phase.

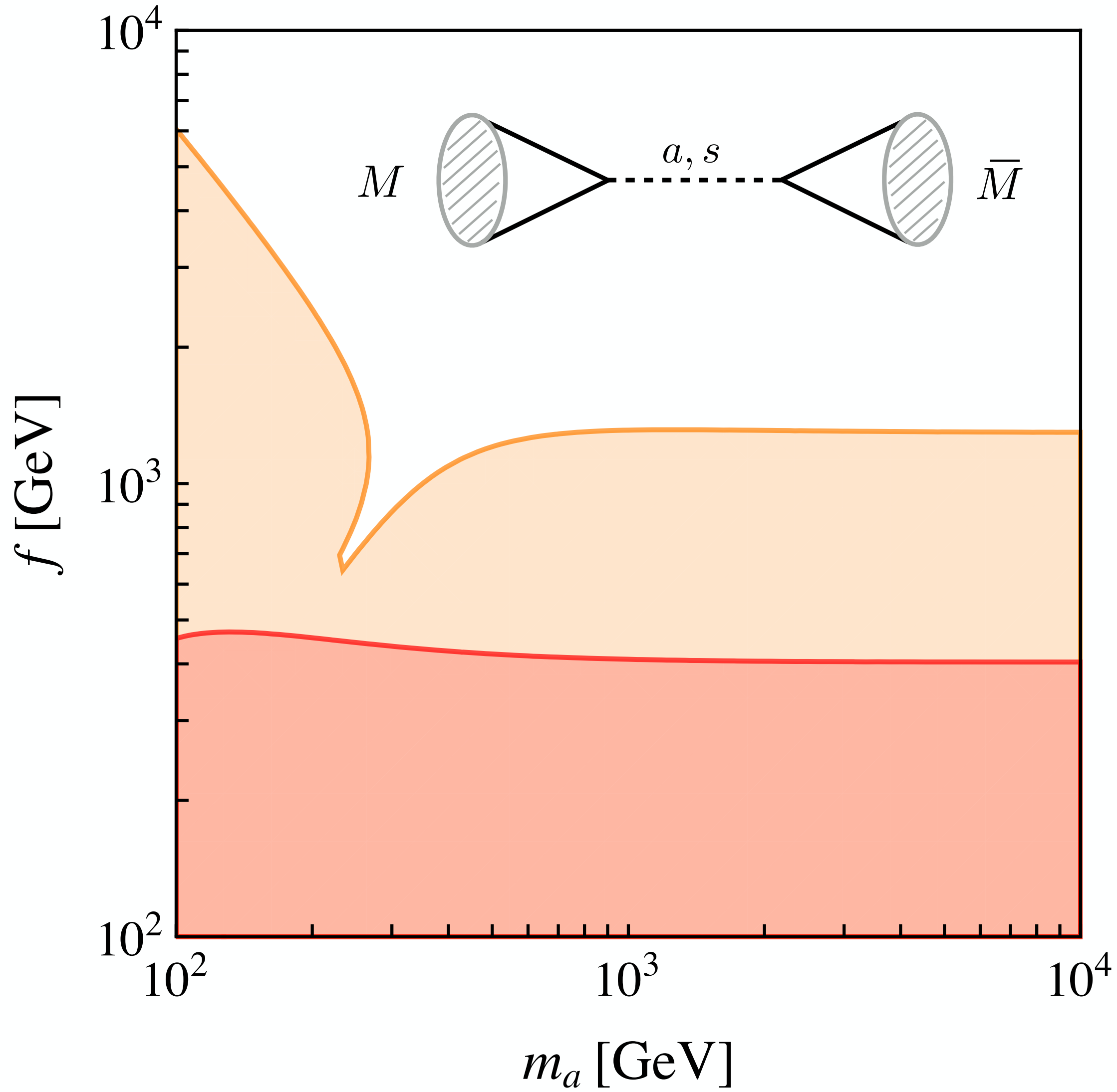
$$Y_u = \begin{pmatrix} 0.34 + 0.82i & -0.23 + 0.69i & 0.41 - 0.43i \\ -0.84 + 0.26i & -0.64 + 0.32i & 1.35 - 0.24i \\ 0.98 - 0.90i & -0.84 - 1.20i & 0.75 + 0.65i \end{pmatrix}$$

$$Y_d = \begin{pmatrix} 0.53 + 0.72i & 0.50 - 0.34i & 0.65 - 0.10i \\ 1.12 - 0.14i & 0.93 - 0.54i & -0.31 - 0.65i \\ -0.16 + 0.6i & -0.73 + 0.34i & 0.84 + 0.61i \end{pmatrix}$$

Quark Flavor constraints

$$\epsilon_K \,, \Delta m_K$$

$$\mathcal{H}_{\rm NP}^{\Delta F=2} = C_1^{ij}\,(\bar q_L^i\,\gamma_\mu\,q_L^j)^2 + \widetilde{C}_1^{ij}\,(\bar q_R^i\,\gamma_\mu\,q_R^j)^2 + C_2^{ij}\,(\bar q_R^i\,q_L^j)^2 + \widetilde{C}_2^{ij}\,(\bar q_L^i\,q_R^j)^2 + C_4^{ij}\,(\bar q_R^i\,q_L^j)\,(\bar q_L^i\,q_R^j) + C_5^{ij}\,(\bar q_L^i\,\gamma_\mu\,q_L^j)\,(\bar q_R^i\,\gamma^\mu q_R^j) + \text{h.c.}$$



$$C_2^{sd} = -(g_{ds}^*)^2\left(\frac{1}{m_s^2} - \frac{1}{m_a^2}\right),$$

$$\tilde{C}_2^{sd} = -g_{sd}^2\left(\frac{1}{m_s^2} - \frac{1}{m_a^2}\right),$$

$$C_4^{sd} = -\frac{g_{sd}g_{ds}}{2}\left(\frac{1}{m_s^2} + \frac{1}{m_a^2}\right).$$

Run down and match

$$C_{\epsilon_K} = \frac{\mathrm{Im}\langle K^0|\mathcal{H}^{\Delta F=2}|\bar K^0\rangle}{\mathrm{Im}\langle K^0|\mathcal{H}_{\rm SM}^{\Delta F=2}|\bar K^0\rangle}$$

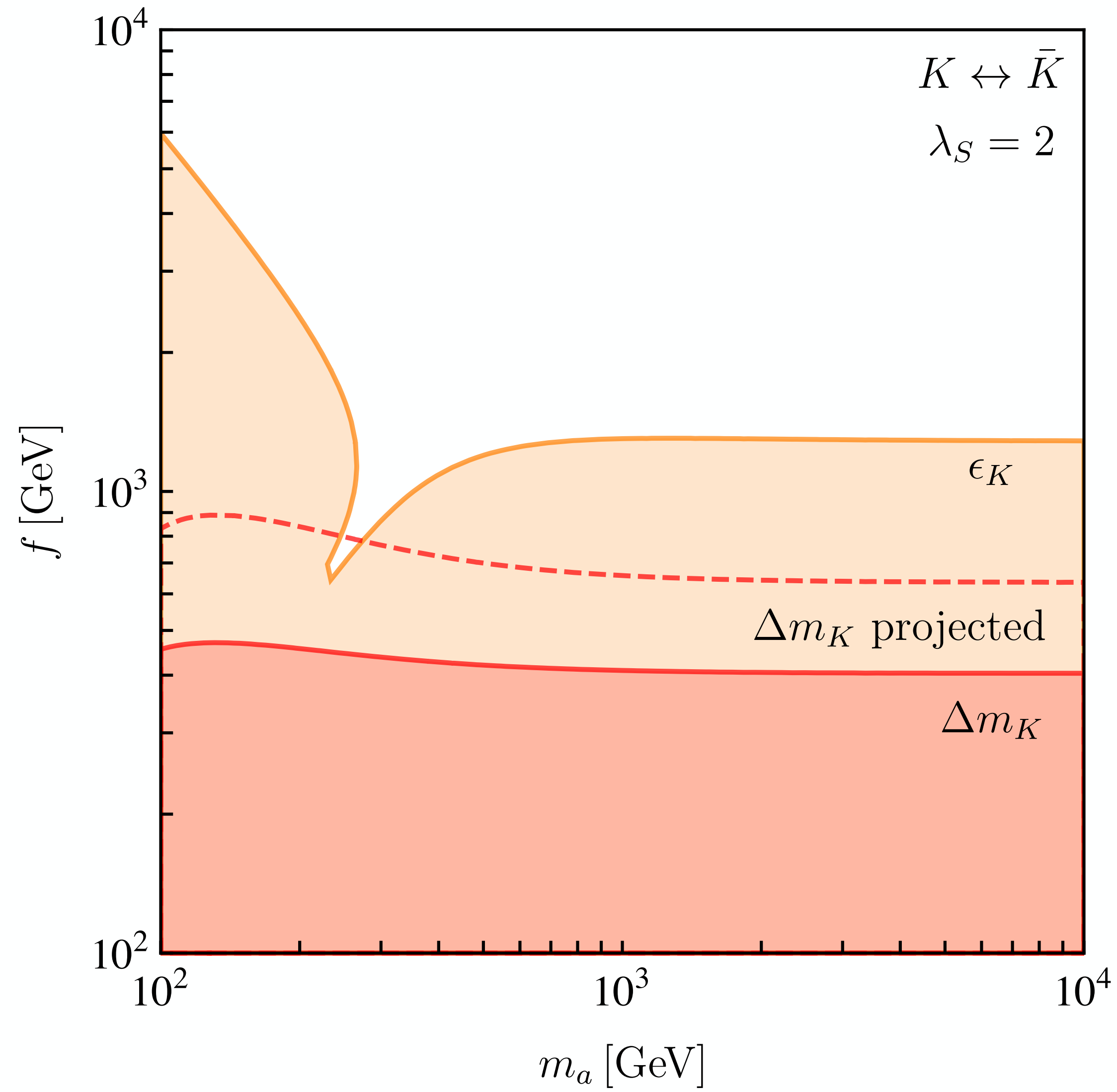
UTFIT

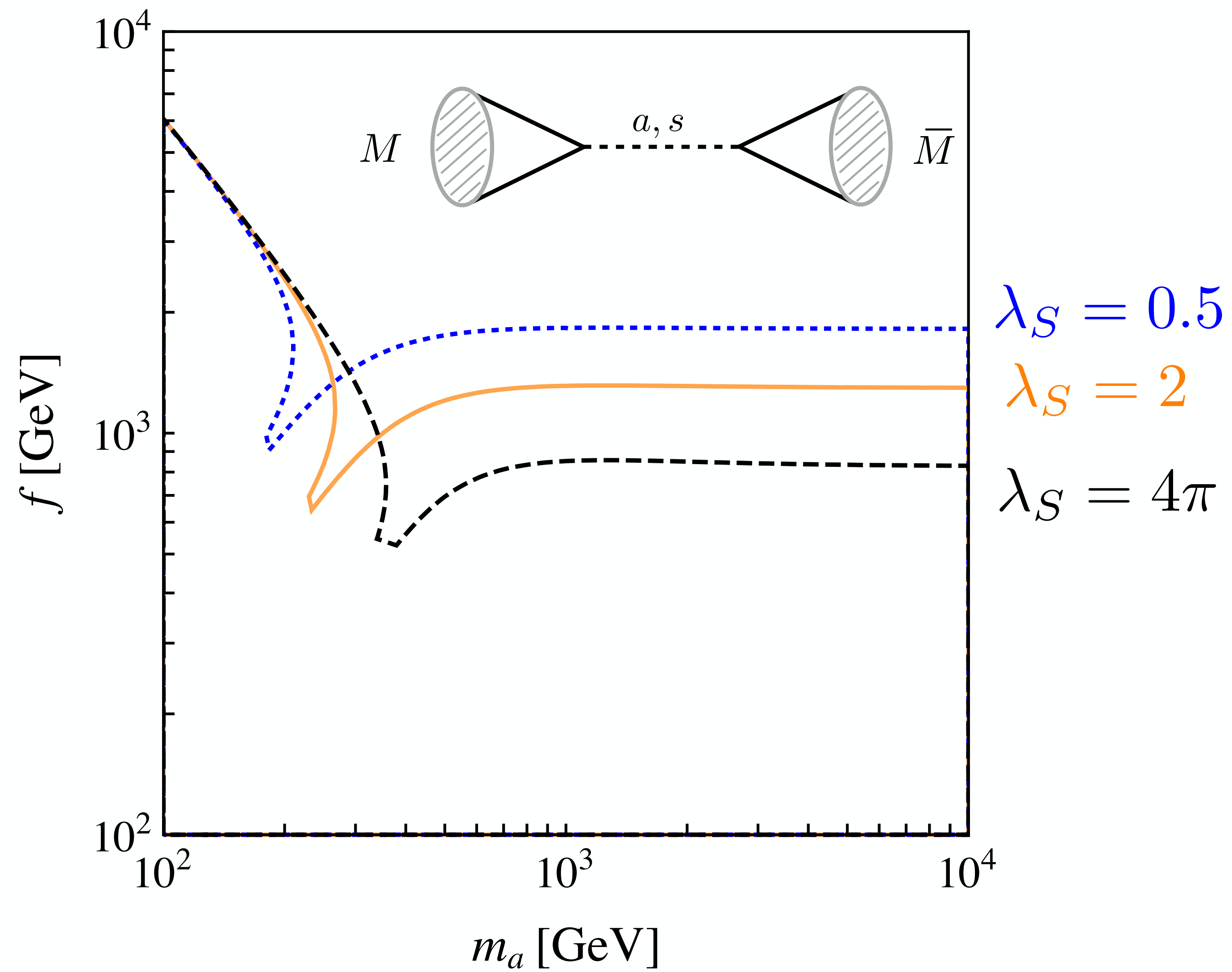
$$C_{\epsilon_K} = 1.05^{+0.36}_{-0.28} \quad @\,95\%\ CL\,, \qquad C_{\Delta_{m_K}} = 0.93^{+1.14}_{-0.42} \quad @\,95\%\ \text{CL}$$

Quark Flavor constraints

$\epsilon_K, \Delta m_K$

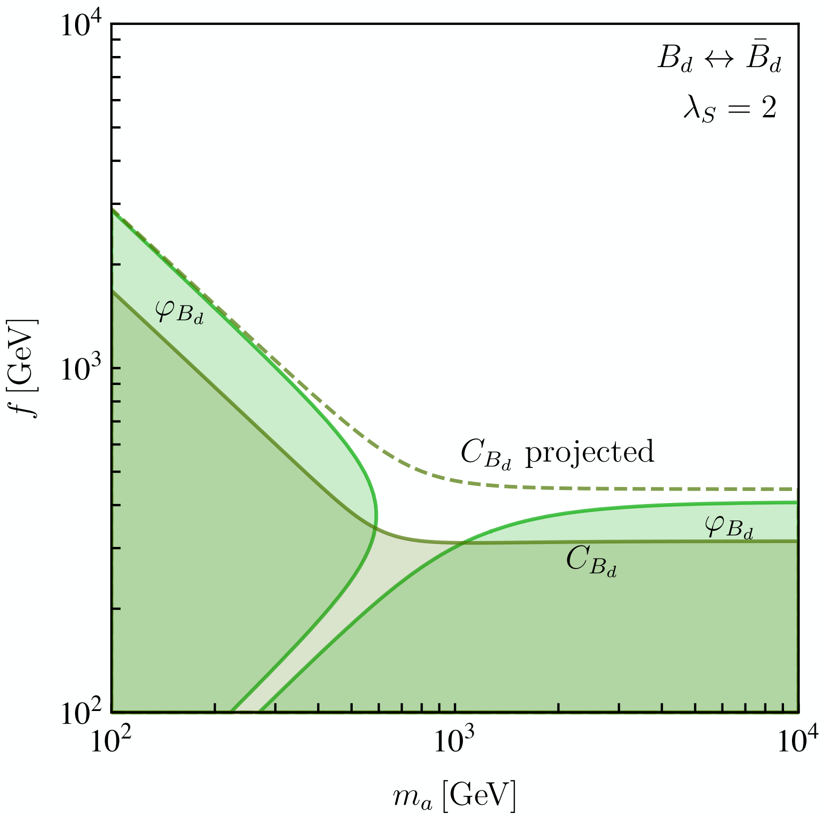
Varying the scalar quartic



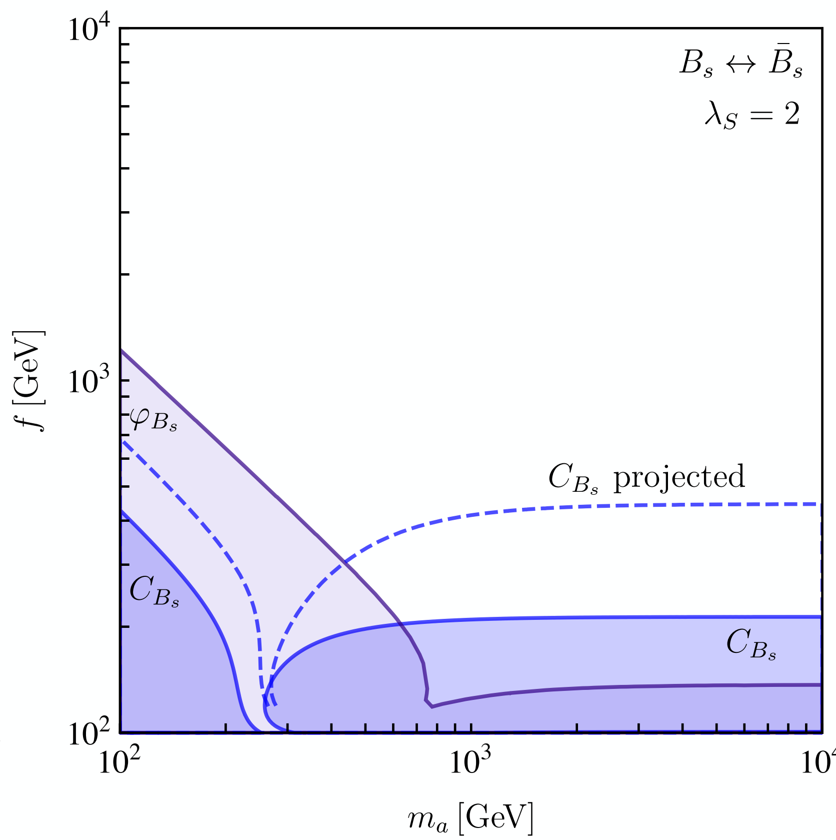


Quark Flavor constraints

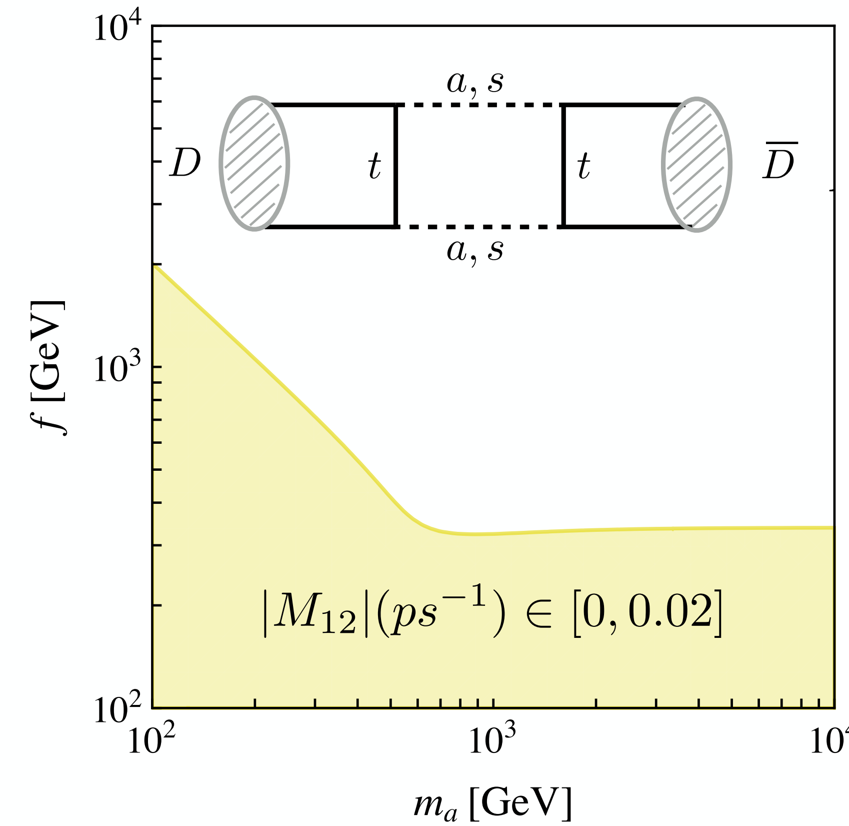
$B_d - \bar{B}_d$



$B_s - \bar{B}_s$



$D - \bar{D}$



$$C_{B_q} e^{2i\varphi_{B_q}} = \frac{\langle B_q | \mathcal{H}^{\Delta F=2} | \bar{B}_q \rangle}{\langle B_q | \mathcal{H}^{\Delta F=2}_{\text{SM}} | \bar{B}_q \rangle}$$

$$C_{B_d} = 1.07^{+0.36}_{-0.31} \quad @\,95\%\text{CL}\,, \qquad \varphi_{B_d} = -2.0^{+6.4}_{-6.0} \quad @\,95\%\text{CL}\,,$$

$$C_{B_s} = 1.052^{+0.178}_{-0.152} \quad @\,95\%\text{CL}\,, \qquad \varphi_{B_s} = 0.72^{+3.98}_{-2.28} \quad @\,95\%\text{CL}\,,$$

Using recent Lattice results

1602.03560 MILC & Fermilab Lattice

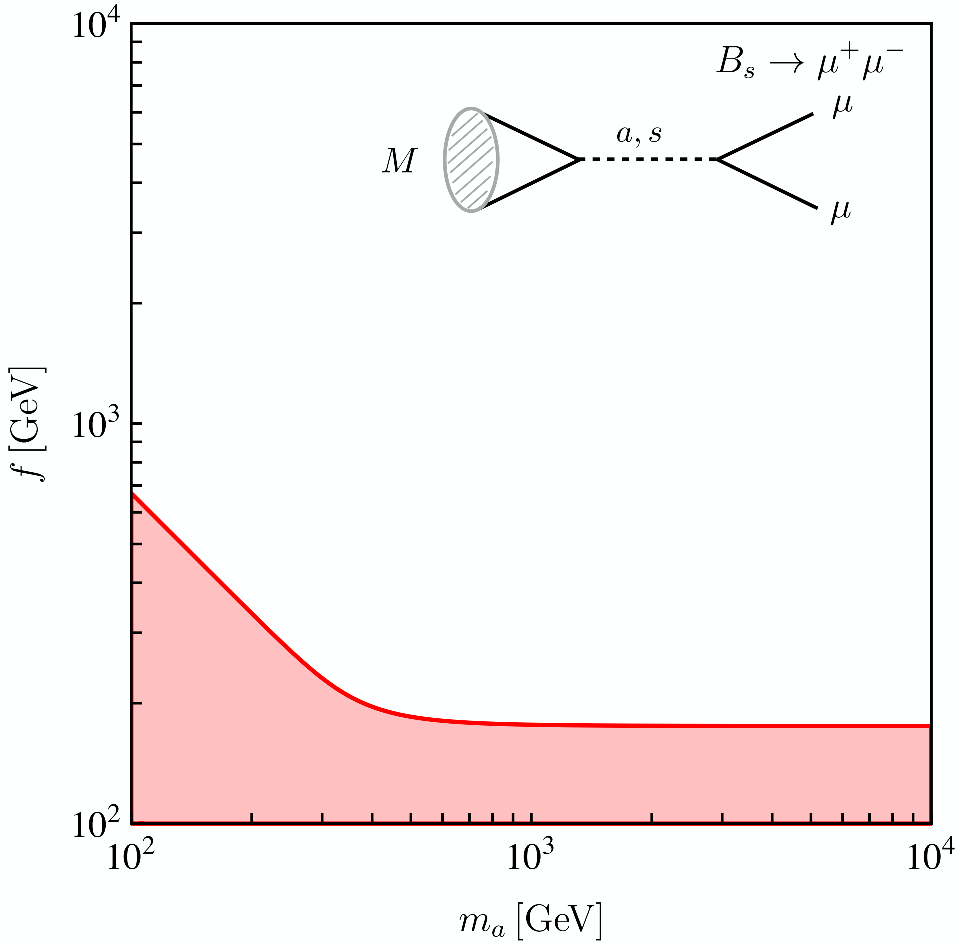
Quark Flavor constraints

$$\mathrm{BR}(M \to \ell\ell) = \frac{G_F^4 M_W^4}{8\pi^5} \beta\left(\frac{m_\ell}{M_M}\right) M_M f_M^2 m_\ell^2 \tau_M \left\{ \left| \frac{M_M^2 (C_P^{ij} - \tilde{C}_P^{ij})}{2m_\ell(m_i + m_j)} - C_A^{\mathrm{SM}} \right|^2 + \left| \frac{M_M^2 (C_S^{ij} - \tilde{C}_S^{ij})}{2m_\ell(m_i + m_j)} \right|^2 \beta^2 \left(\frac{m_\ell}{M_M}\right) \right\}$$

LHCb&CMS

$$\mathrm{BR}(B_s \to \mu^+\mu^-) = 2.8^{+0.7}_{-0.6} \times 10^{-9}$$

$$\mathrm{BR}(B_d \to \mu^+\mu^-) = 3.6 \pm 1.6 \times 10^{-10}$$



$$C_S^{ij} = \frac{\pi^2}{2G_F^2 M_W^2} \frac{2g_{\ell\ell} g_{ji}}{m_s^2}, \qquad \tilde{C}_S^{ij} = \frac{\pi^2}{2G_F^2 M_W^2} \frac{2g_{\ell\ell} g_{ij}}{m_s^2},$$

$$C_P^{ij} = \frac{\pi^2}{2G_F^2 M_W^2} \frac{2g_{\ell\ell} g_{ji}}{m_a^2}, \qquad \tilde{C}_P^{ij} = \frac{\pi^2}{2G_F^2 M_W^2} \frac{2g_{\ell\ell} g_{ij}}{m_a^2}.$$

Quark Flavor constraints

$$\mathrm{BR}(M \to \ell\ell) = \frac{G_F^4 M_W^4}{8\pi^5} \beta\left(\frac{m_\ell}{M_M}\right) M_M f_M^2 m_\ell^2 \tau_M \left\{ \left| \frac{M_M^2 (C_P^{ij} - \tilde{C}_P^{ij})}{2m_\ell(m_i + m_j)} - C_A^{\mathrm{SM}} \right|^2 + \left| \frac{M_M^2 (C_S^{ij} - \tilde{C}_S^{ij})}{2m_\ell(m_i + m_j)} \right|^2 \beta^2 \left(\frac{m_\ell}{M_M}\right) \right\}$$

LHCb&CMS

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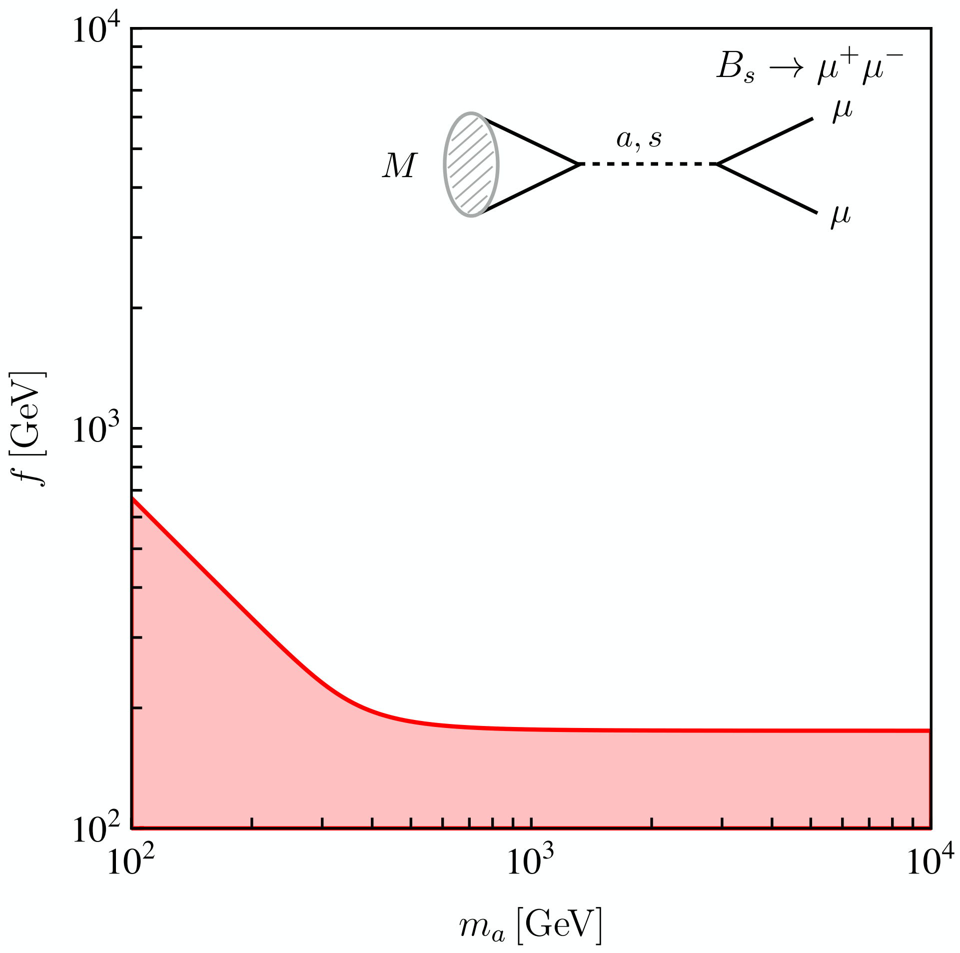
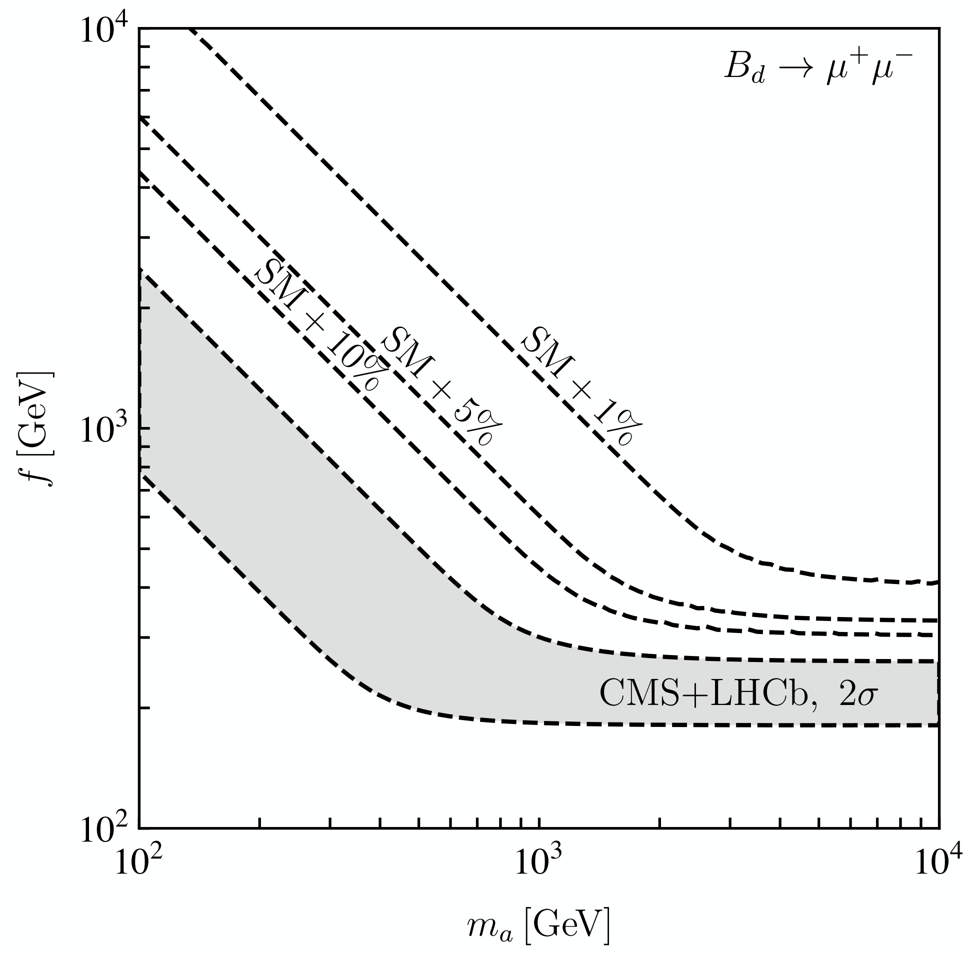


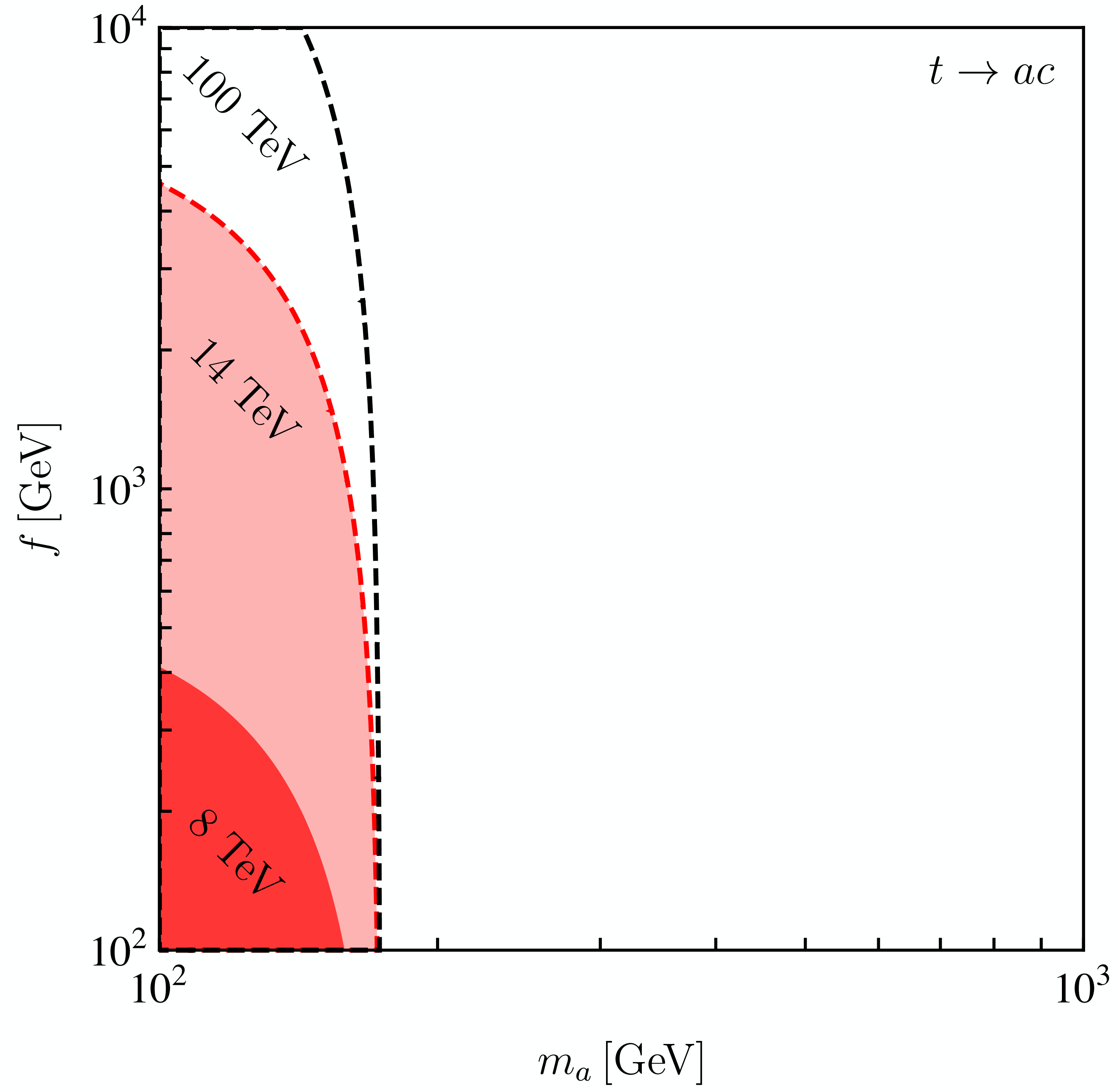


Top decays

LHC 8 $\mathrm{BR}(t \to ac) < 5.6 \times 10^{-3}$

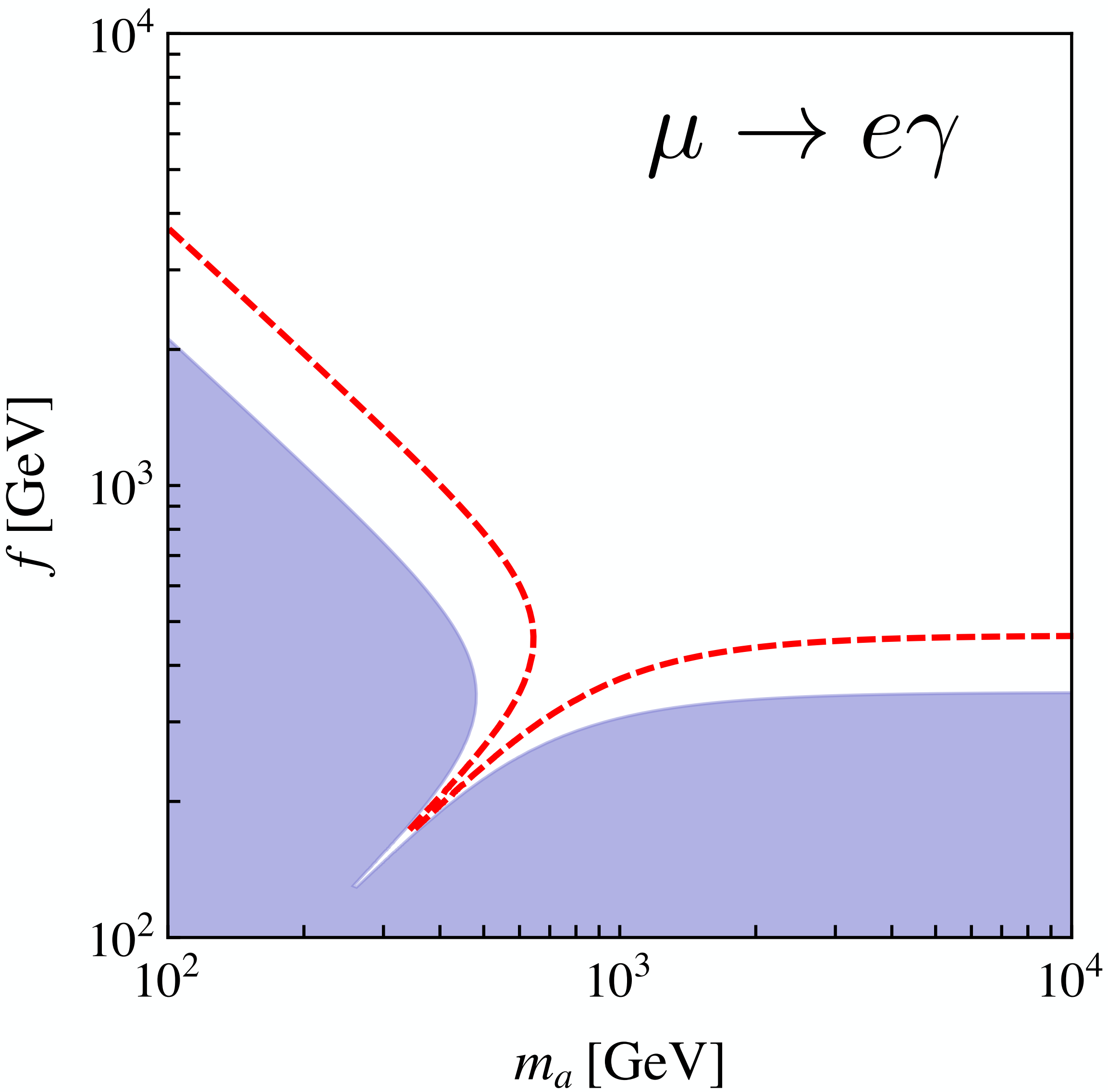
LHC 14 $\mathrm{BR}(t \to ac) < 4.5 \times 10^{-5}$

100TeV $\mathrm{BR}(t \to ac) < 1.5 \times 10^{-6}$



Lepton Flavor constraints

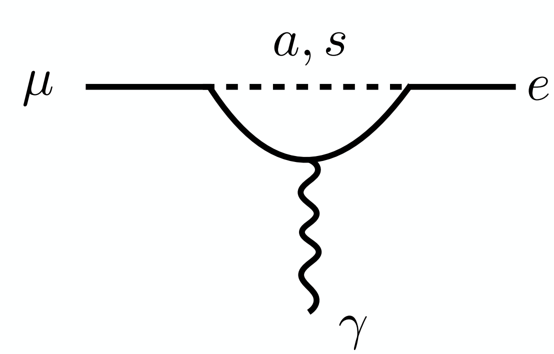
$$\mathcal{L}_{\rm eff} = m_{\ell'}\, C_T^L\, \bar{\ell}\sigma^{\rho\lambda} P_L\, \ell'\, F_{\rho\lambda} + m_{\ell'}\, C_T^R\, \bar{\ell}\sigma^{\rho\lambda} P_R\, \ell'\, F_{\rho\lambda}$$

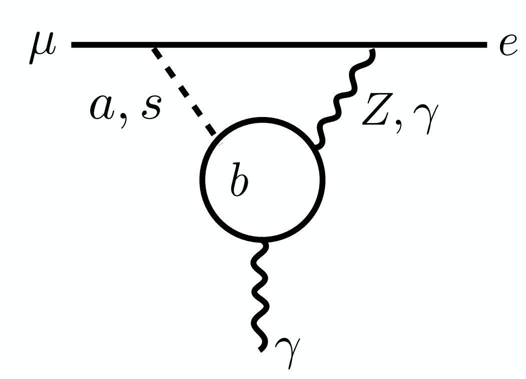


$$C_T^L = \frac{g}{32\pi^2} \sum_{k=e,\mu,\tau} \left[\frac{1}{6} \left(g^*_{\ell k} g_{\ell' k} + \frac{m_\ell}{m_k} g^*_{k\ell} g_{k\ell'} \right) \left(\frac{1}{m_s^2} - \frac{1}{m_a^2} \right) - g_{\ell k} g_{k\ell'} \frac{m_k}{m_{\ell'}} \left\{ \frac{1}{m_s^2} \left(\frac{3}{2} + \log \frac{m_{\ell'}^2}{m_s^2} \right) - \frac{1}{m_a^2} \left(\frac{3}{2} + \log \frac{m_{\ell'}^2}{m_a^2} \right) \right\} \right]$$

MEG I $\quad \mathrm{BR}(\mu \to e\gamma) < 5.7 \times 10^{-13}$

MEG II $\quad \mathrm{BR}(\mu \to e\gamma) < 6 \times 10^{-14}$





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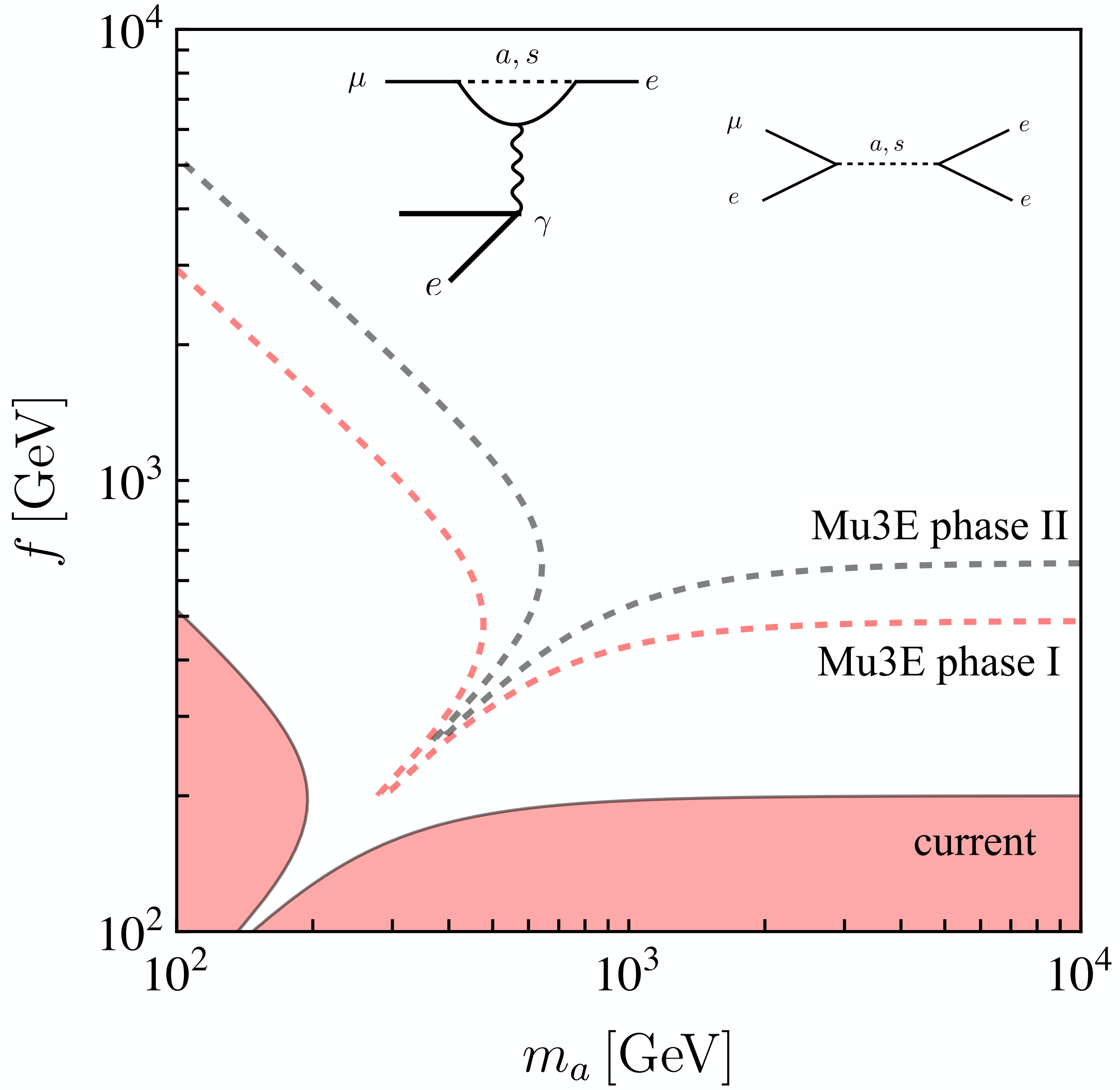
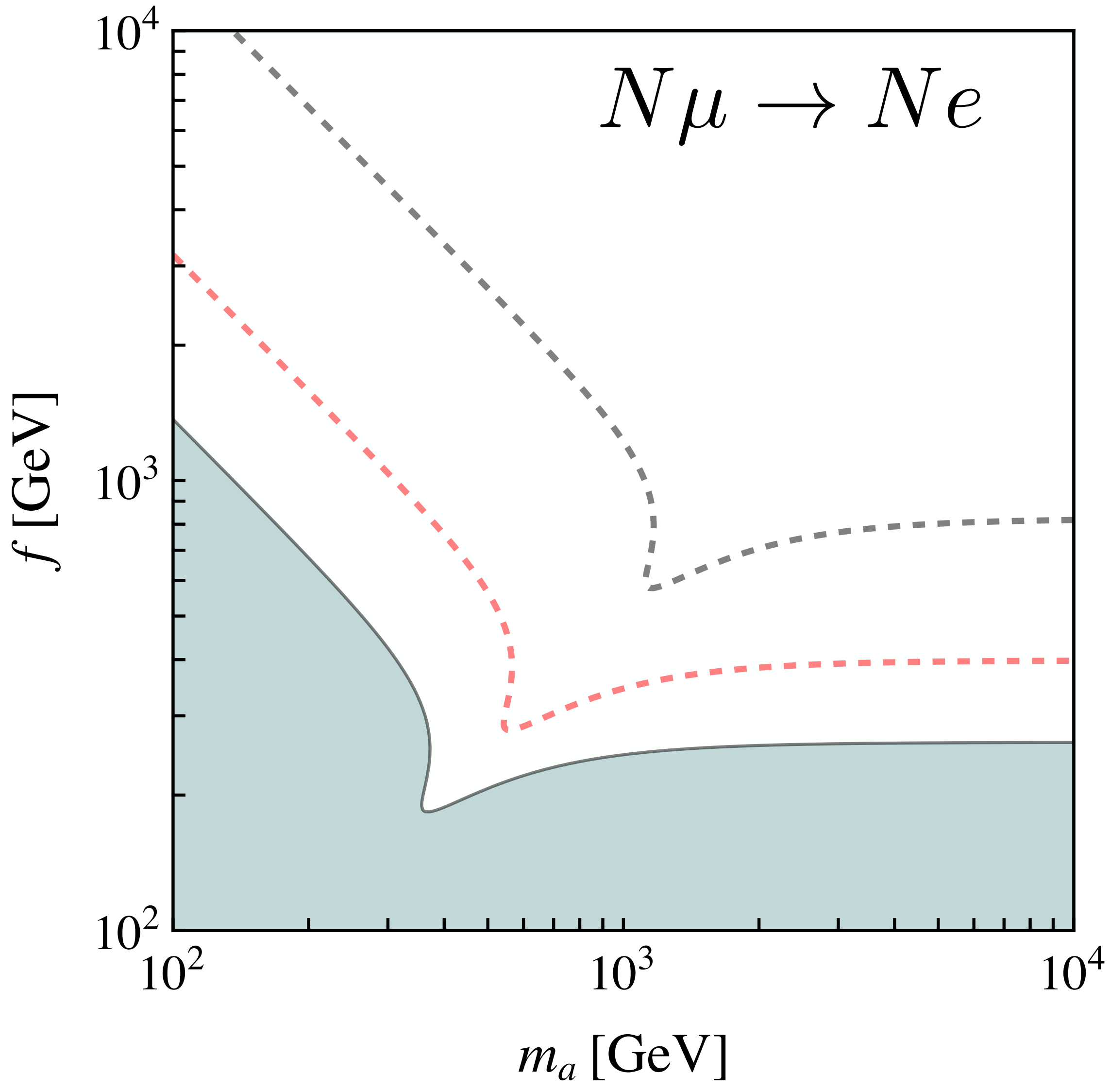


$$\text{BR}(\tau \to 3\mu) < 2.1 \cdot 10^{-8}\,,$$
$$\text{BR}(\tau \to 3e) < 2.7 \cdot 10^{-8}\,,$$
$$\text{BR}(\mu \to 3e) < 1.0 \cdot 10^{-12}\,.$$

Mu3E will improve this by 3-4 orders of magnitude!

Lepton Flavor constraints

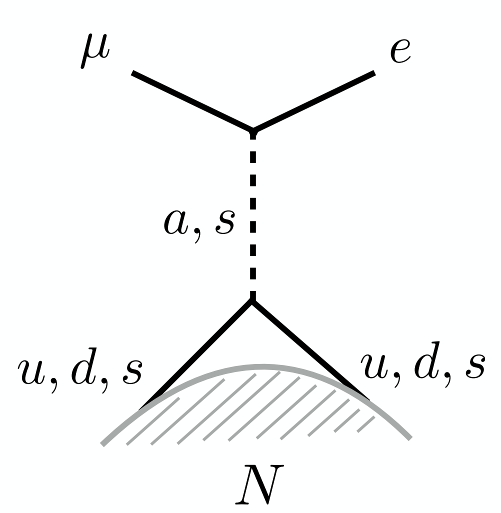
$$\mathcal{L}_{\rm eff} = C_{qq}^{VL}\, \bar{e}\gamma^\nu P_L \mu\, \bar{q}\gamma_\nu q + m_\mu m_q\, C_{qq}^{SL} \bar{e} P_R \mu\, \bar{q}q + m_\mu \alpha_s C_{gg}^{L}\, \bar{e} P_R \mu\, G_{\rho\nu} G^{\rho\nu} \; + R \leftrightarrow L,$$

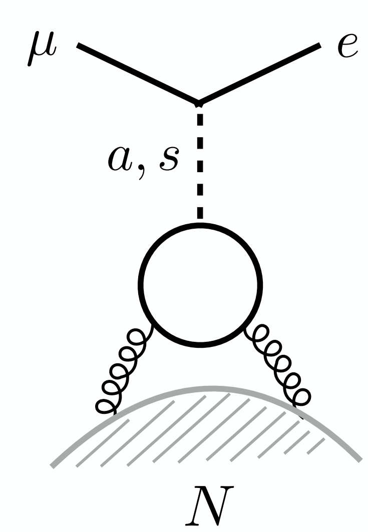


Sindrum II $\quad \mathrm{BR}(\mu \to e)^{\rm Au} < 7 \times 10^{-13}$

DeeMe $\quad \mathrm{BR}(\mu \to e)^{\rm Si} < 2 \times 10^{-14}$

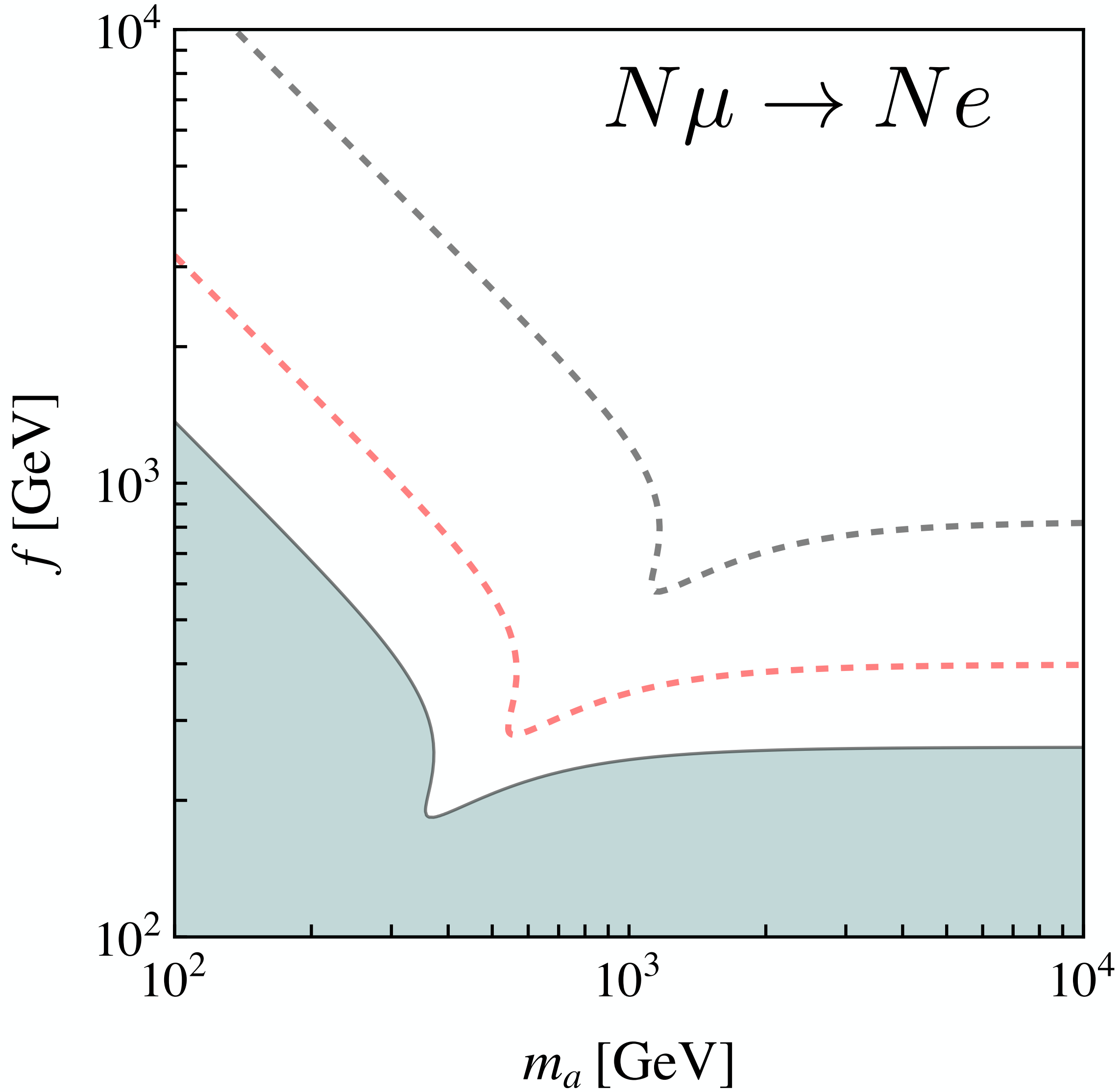
COMET $\quad \mathrm{BR}(\mu \to e)^{\rm Al} < 6 \times 10^{-17}$

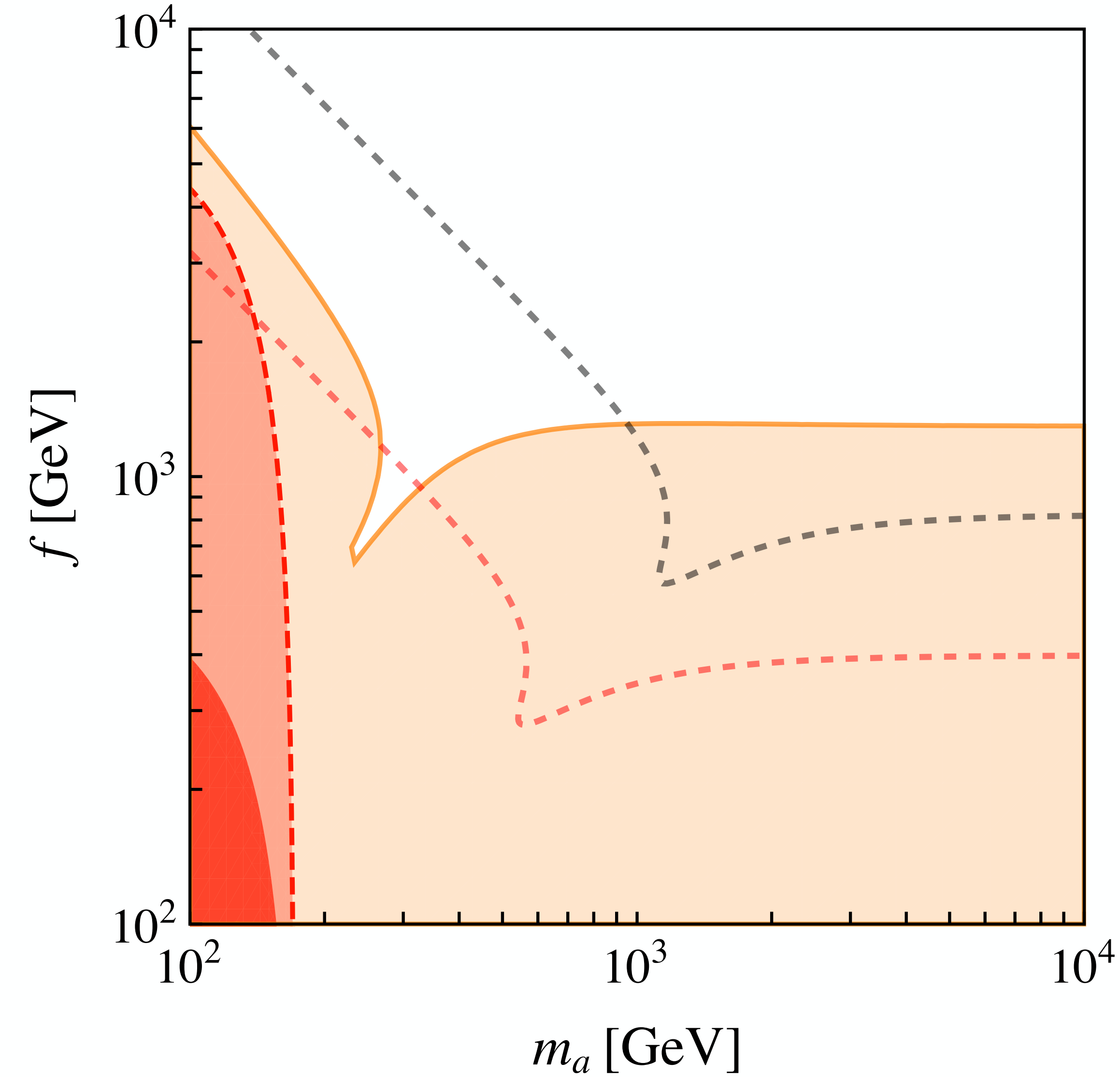




Lepton Flavor constraints

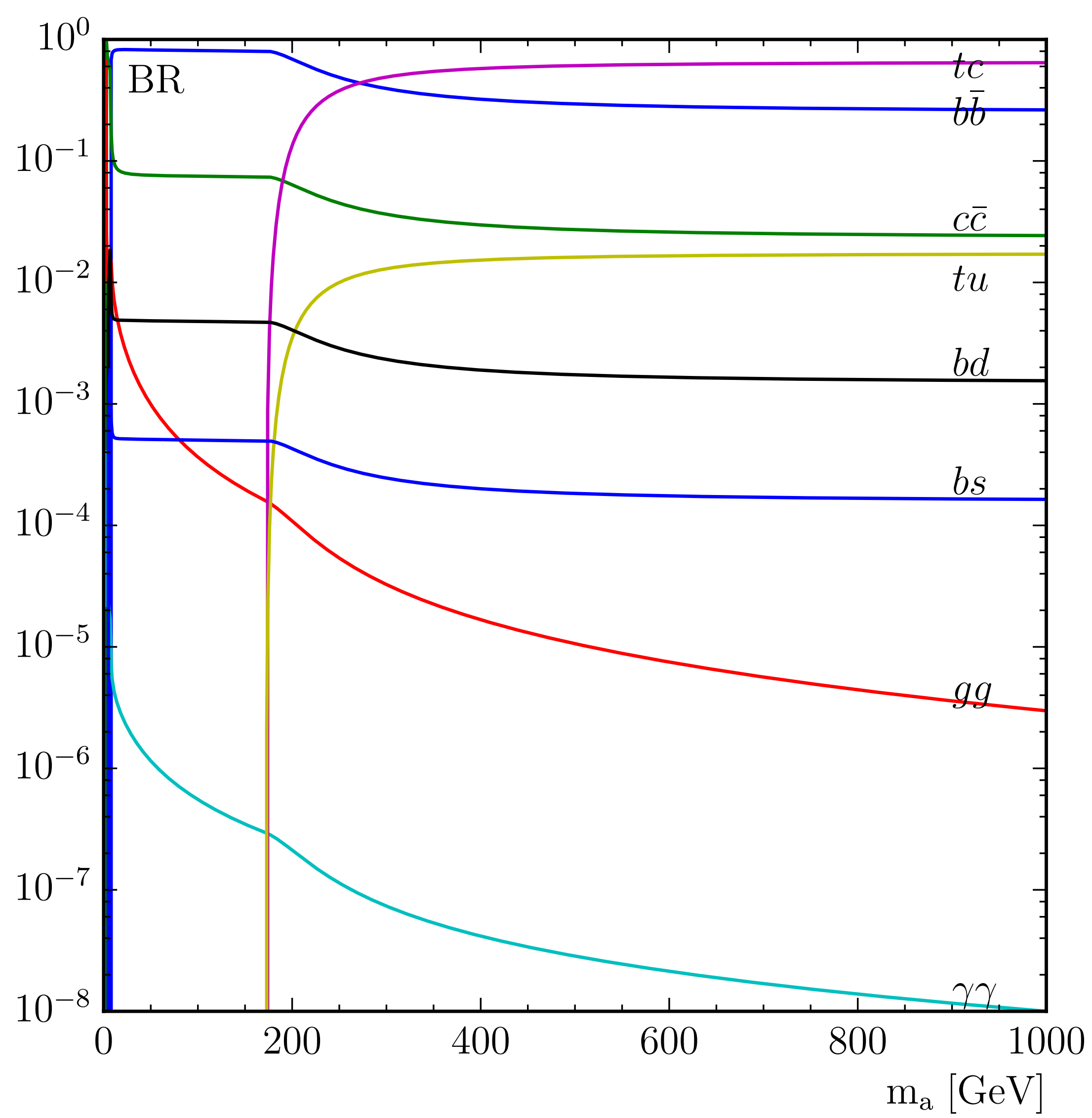
Golden Age of Lepton flavor violation:

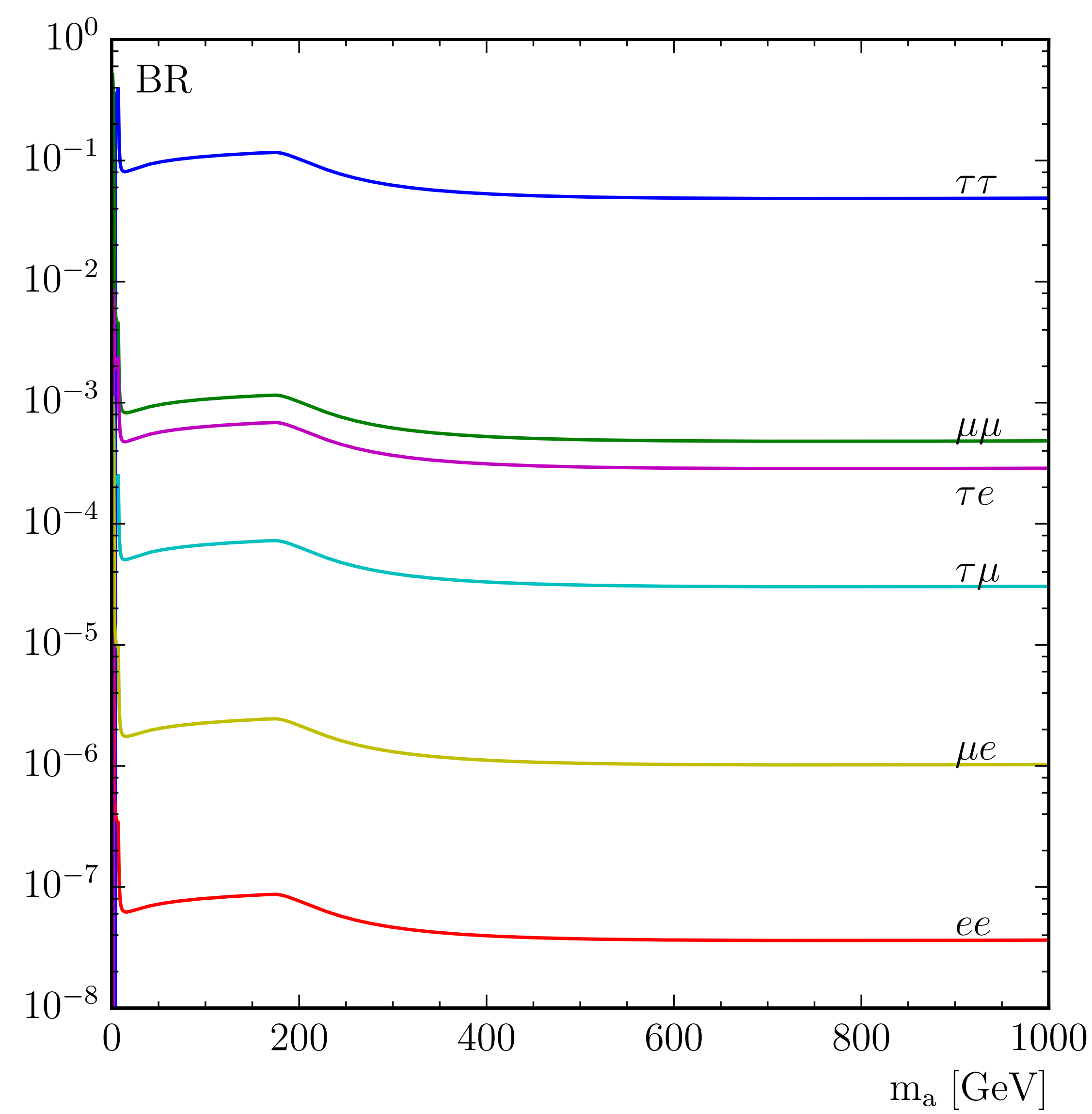




Future Collider Searches

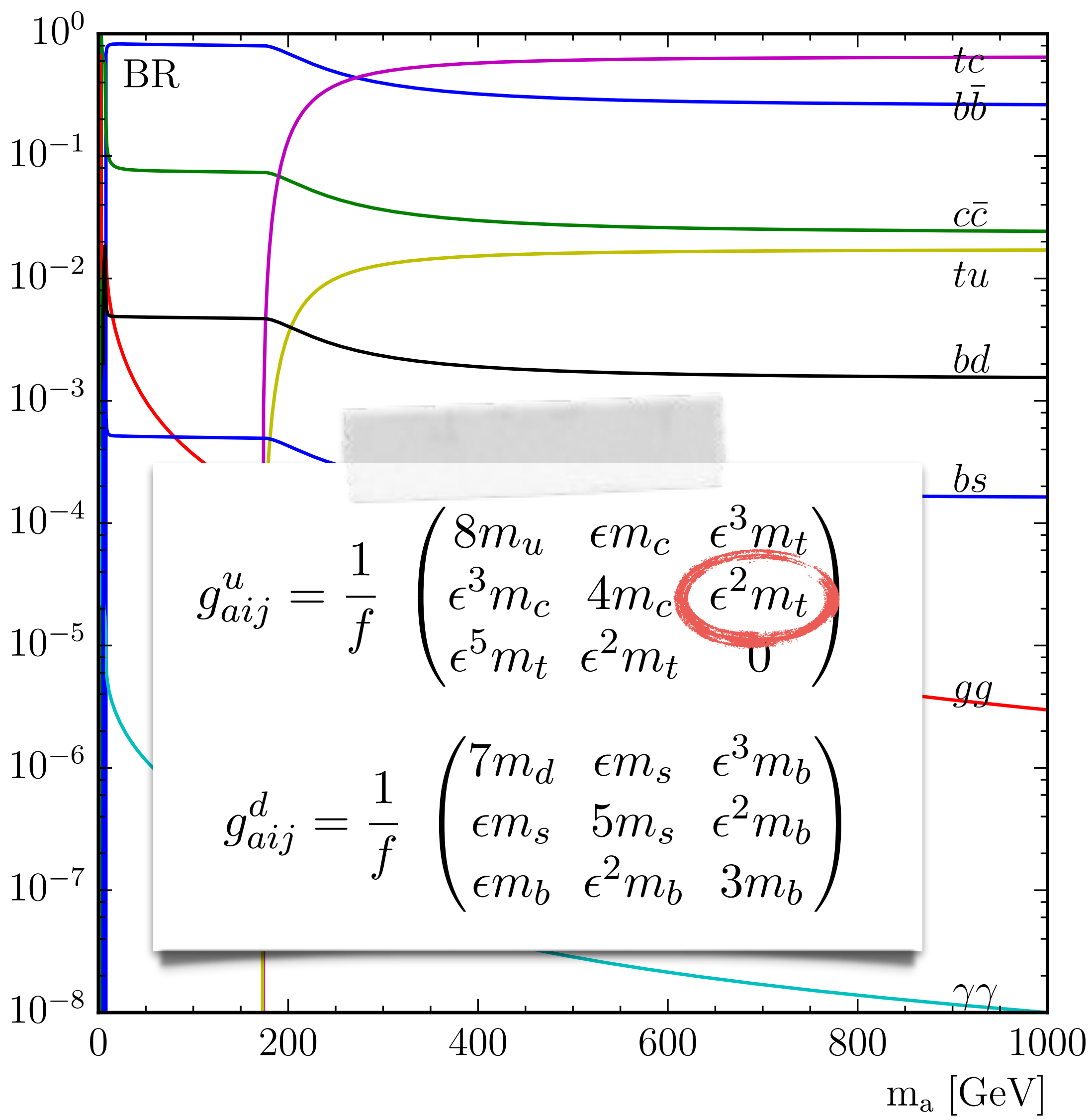
Branching Ratios





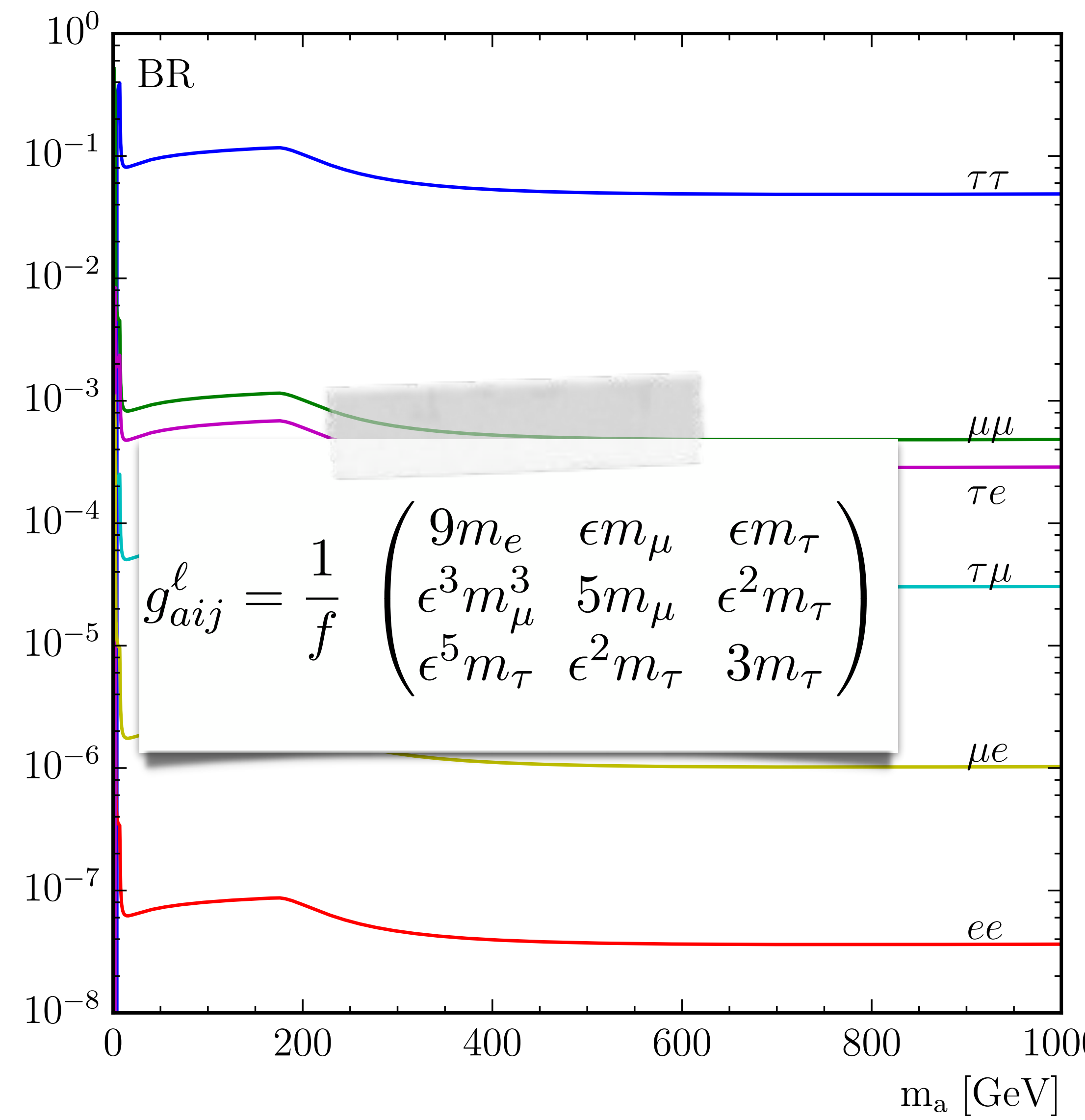
Future Collider Searches

Branching Ratios



$$g^{u}_{aij} = \frac{1}{f}\begin{pmatrix} 8m_u & \epsilon m_c & \epsilon^3 m_t \\ \epsilon^3 m_c & 4m_c & \epsilon^2 m_t \\ \epsilon^5 m_t & \epsilon^2 m_t & 0 \end{pmatrix}$$

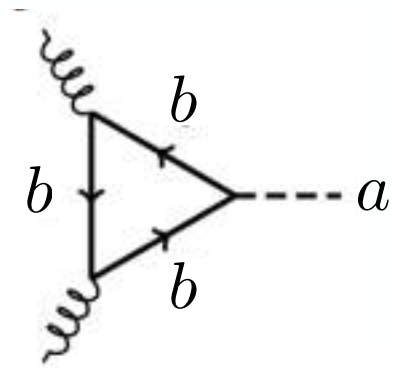
$$g^{d}_{aij} = \frac{1}{f}\begin{pmatrix} 7m_d & \epsilon m_s & \epsilon^3 m_b \\ \epsilon m_s & 5m_s & \epsilon^2 m_b \\ \epsilon m_b & \epsilon^2 m_b & 3m_b \end{pmatrix}$$

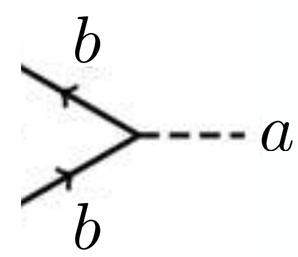


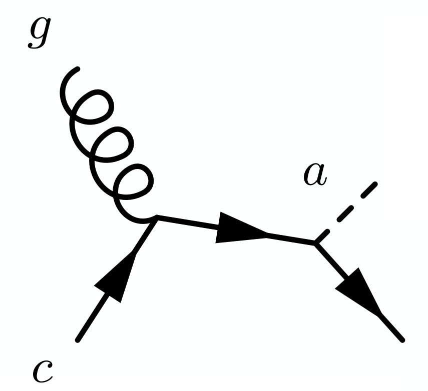
$$g^{\ell}_{aij} = \frac{1}{f}\begin{pmatrix} 9m_e & \epsilon m_\mu & \epsilon m_\tau \\ \epsilon^3 m_\mu^3 & 5m_\mu & \epsilon^2 m_\tau \\ \epsilon^5 m_\tau & \epsilon^2 m_\tau & 3m_\tau \end{pmatrix}$$

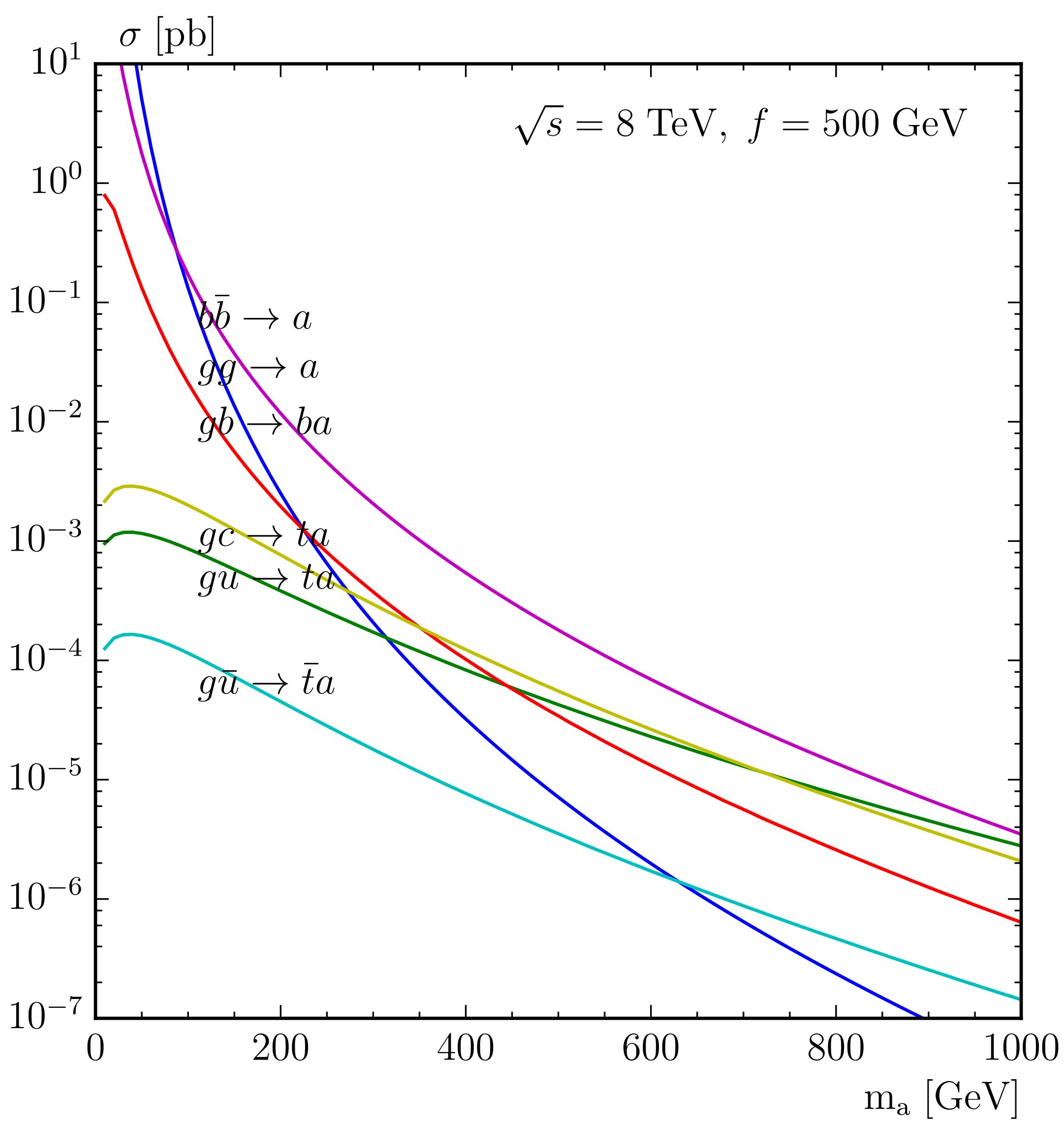
(Future) Collider Searches

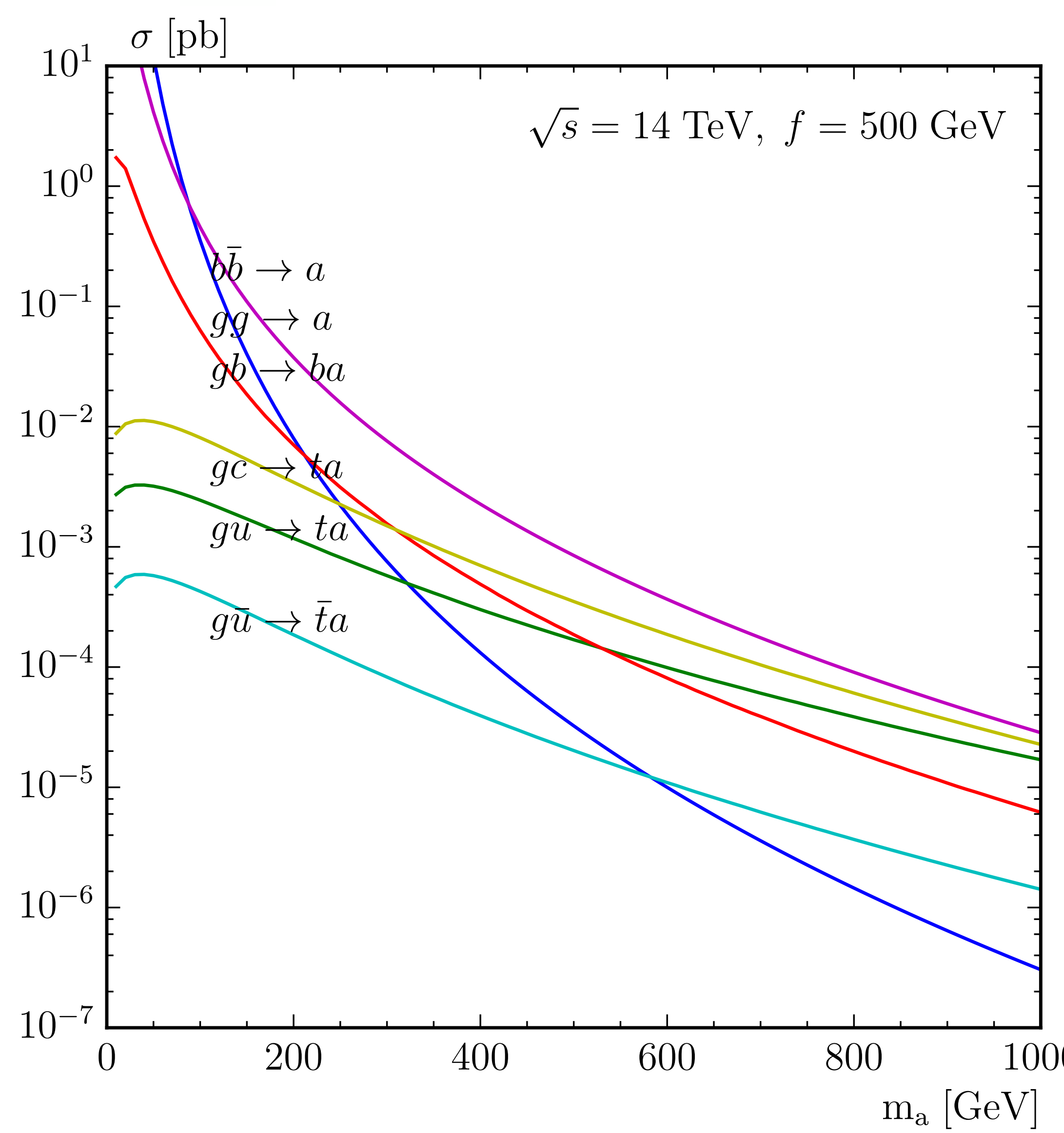
Production Cross Sections at LHC





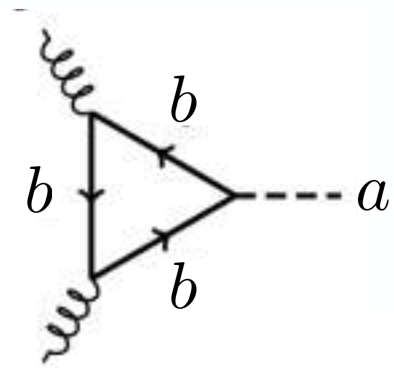


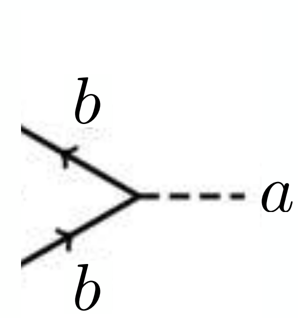


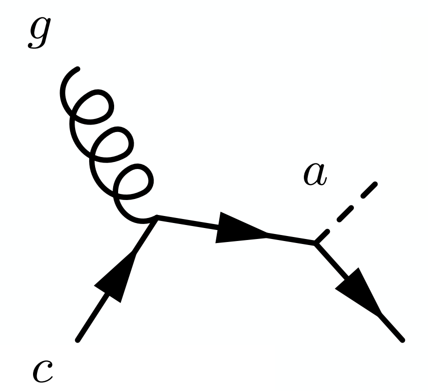


Future Collider Searches

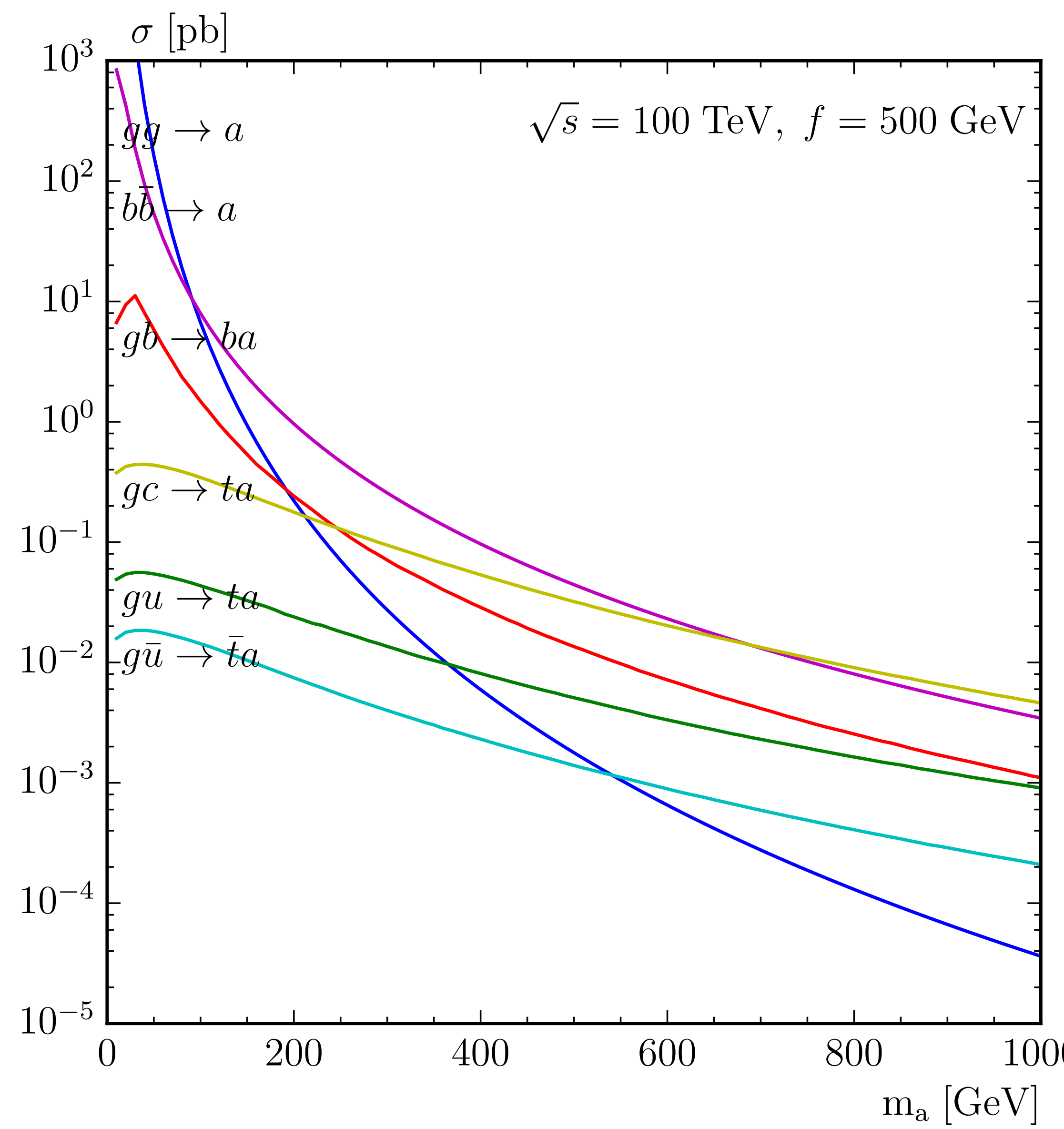
Production Cross Sections at 100 TeV



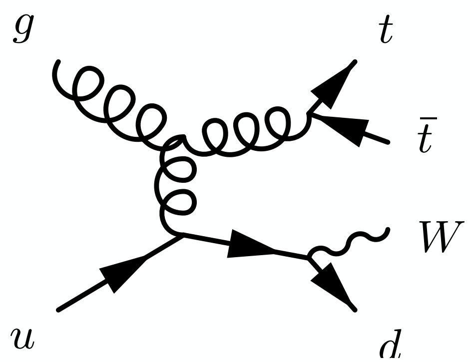


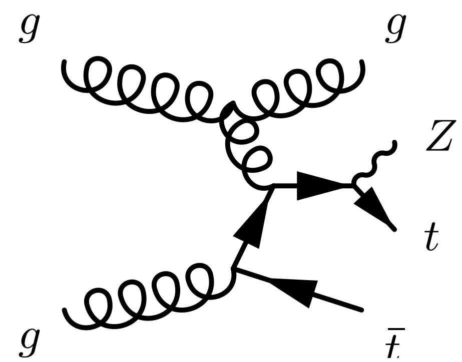


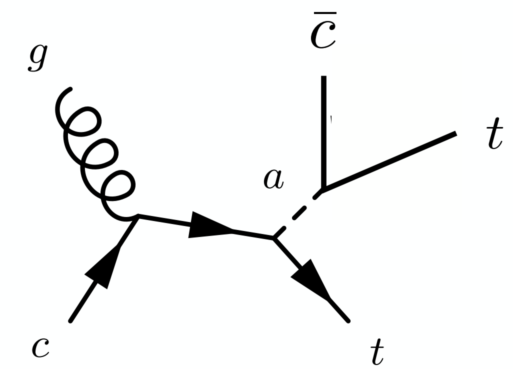
The huge background and the BRs make decays into taus and bs hopeless!



Future Collider Searches



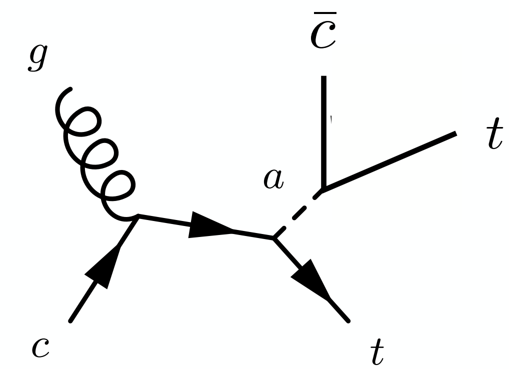




- 2 same-sign leptons(+) (2 hardest ones) with $p_T > 10$ GeV, $|\eta| < 2.5$ $R_{\text{iso}} = 0.2$
- if there is a 3rd lepton of different sign, veto events with $|m_{\ell_i^{(\text{ss})}\ell^{(\text{ds})}} - m_Z| < 15$ GeV
- require for hardest jet $p_T > 100$ GeV
- b-tagging: partonlevel b within $R < 0.3$, assumed efficiency 50 %
- require for the remaining jets $N_b \geq 2$
- $\not{p}_T > 50$ GeV
- minimize $R_{\ell_1 b_i} + R_{\ell_2 b_j}$ to define $(\ell b)_1$ and $(\ell b)_2$
- minimize $\Delta y((\ell b)_i, j)$ to define $(\ell b j)$ and (ℓb)
- calculate m_{T2}

Future Collider Searches

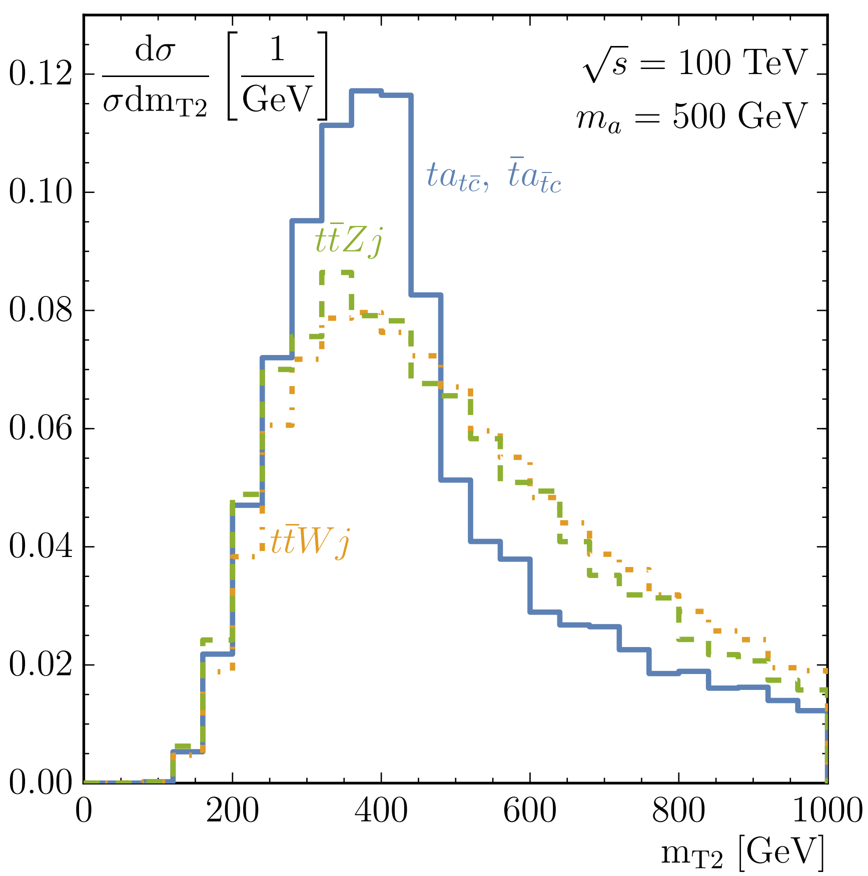
- b-$\bar{b}$ distinction

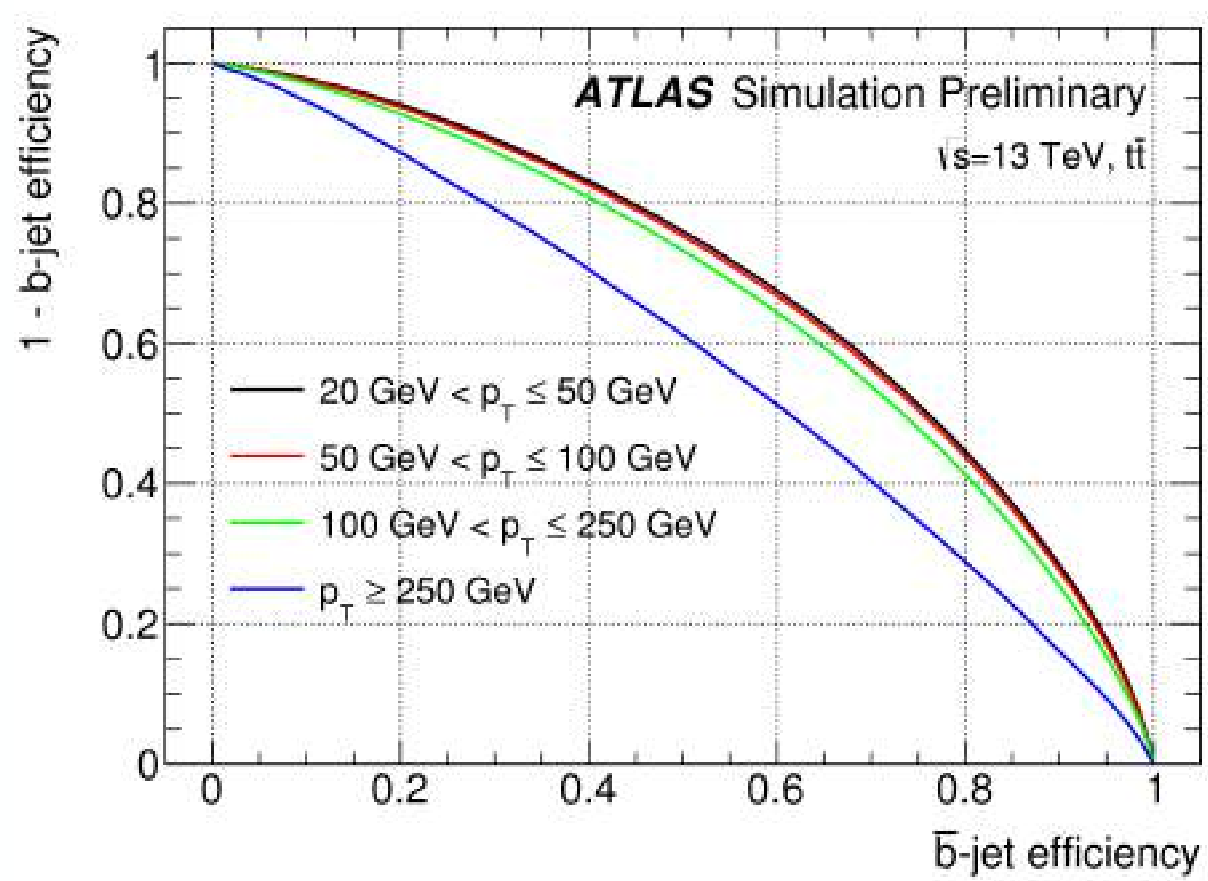


conservative $\epsilon_b = 0.2\,, \epsilon_{\bar{b}} = 0.06$

optimistic $\epsilon_b = 0.2\,, \epsilon_{\bar{b}} = 0.01$

ATL-PHYS-PUB-2015-040



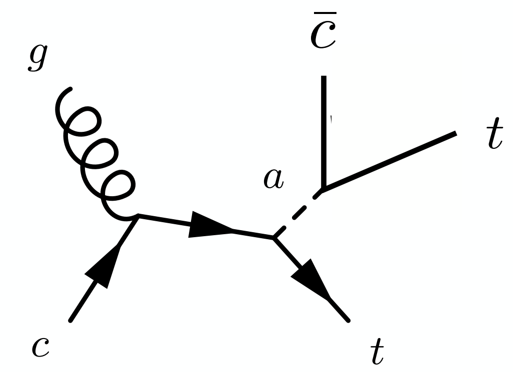


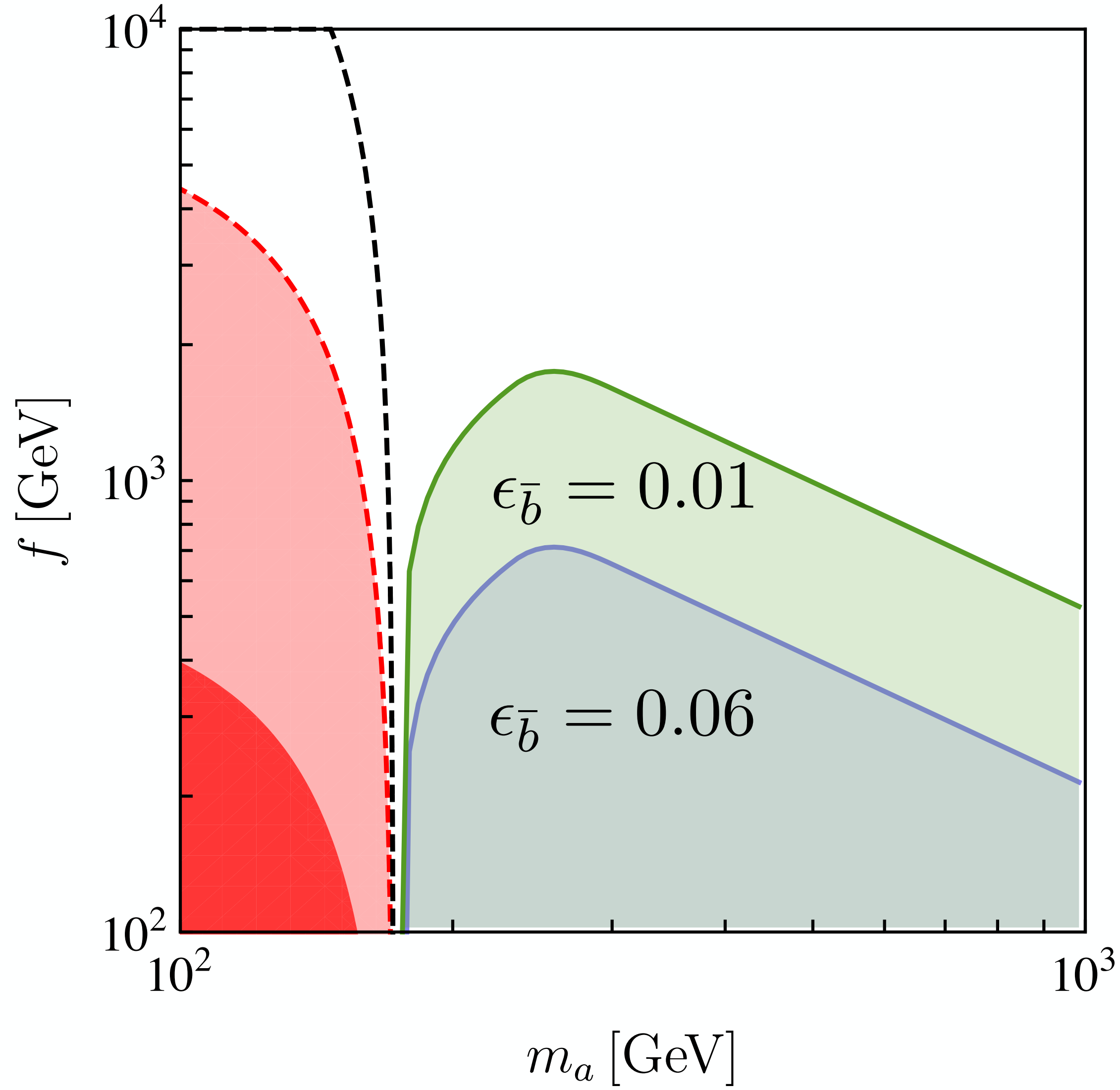
Future Collider Searches

- b-$\bar{b}$ distinction

conservative $\quad \epsilon_b = 0.2\,,\, \epsilon_{\bar{b}} = 0.06$

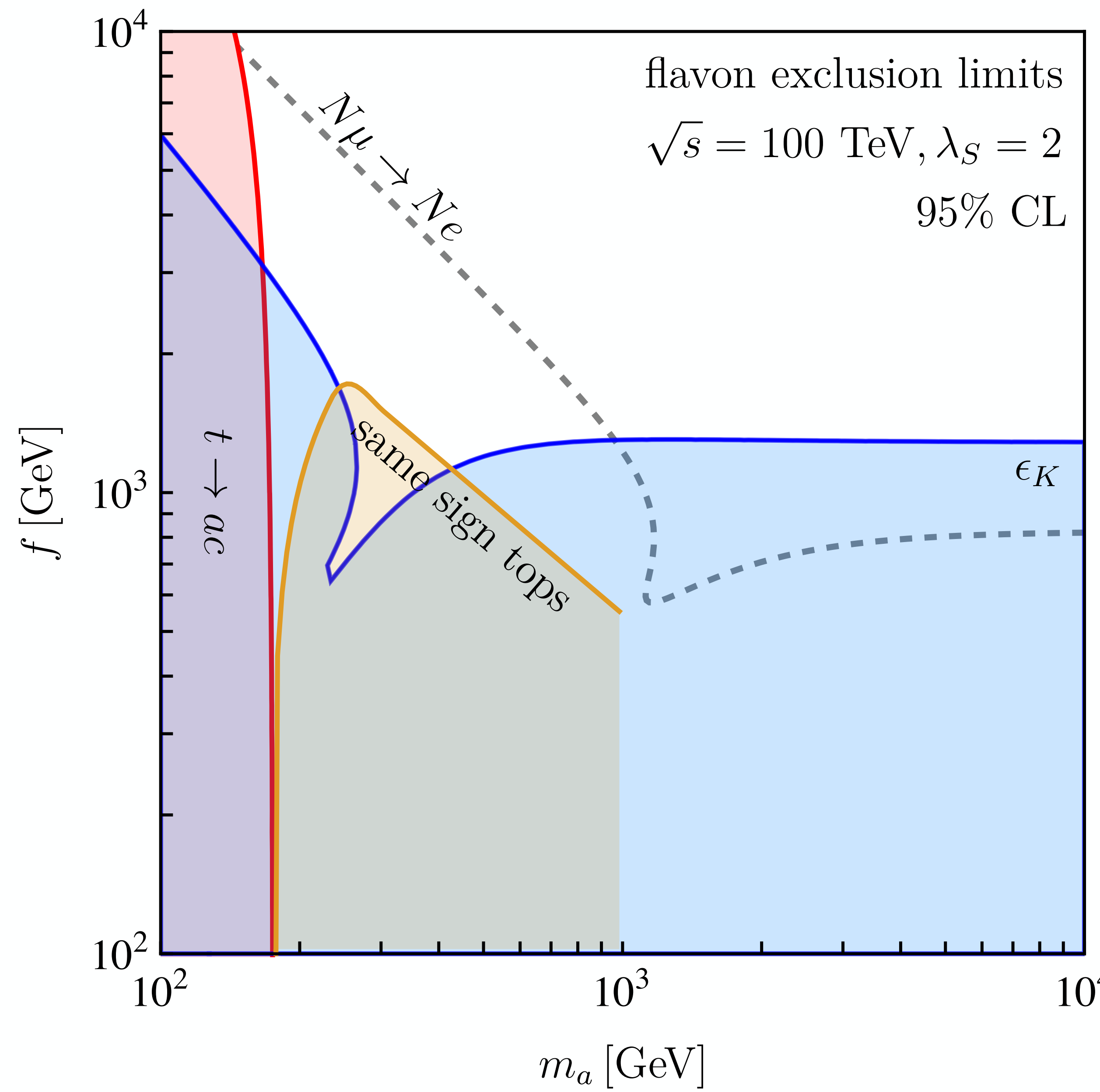
optimistic $\quad \epsilon_b = 0.2\,,\, \epsilon_{\bar{b}} = 0.01$





Future Collider Searches

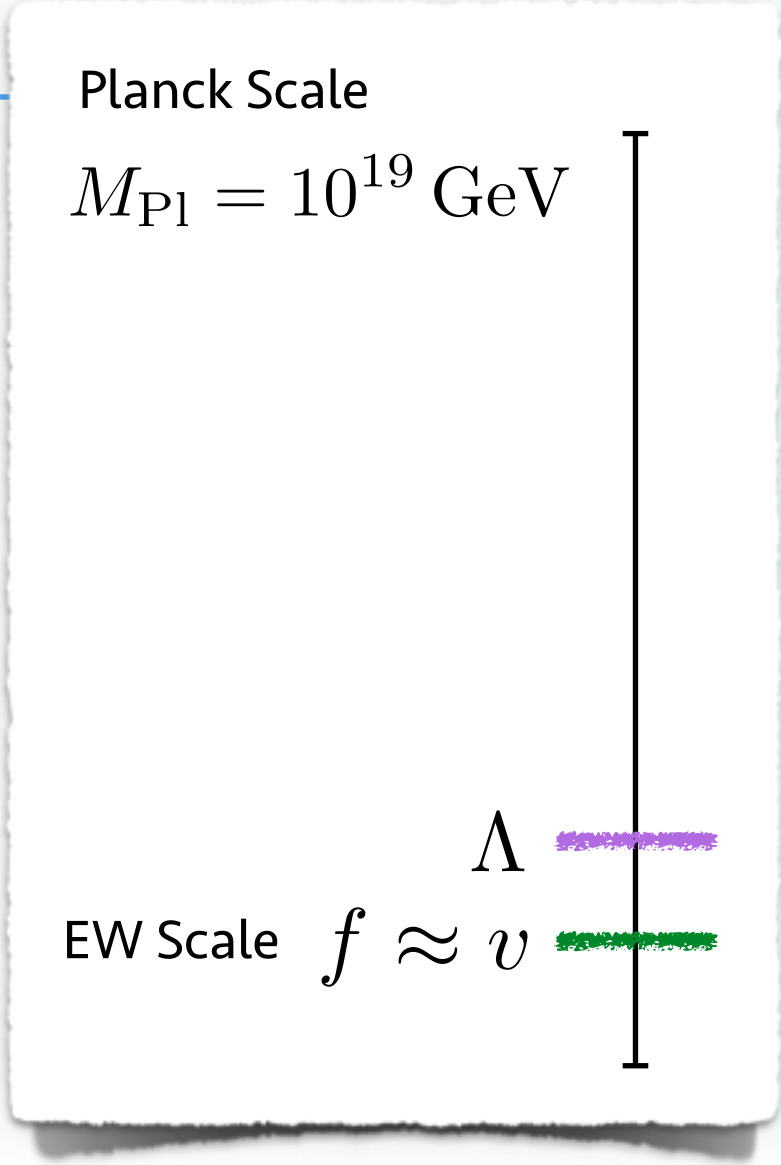
- Next generation lepton flavor experiments will cut deep into the parameter space
- A 100 TeV collider is our first semi-realistic shot at discovering a flavon



Why should the flavor scale be low?

- Why not?
- A link to Baryogenesis?
- A link to Dark Matter?
- It could be related to the electroweak scale

Why should the flavor scale be low?



- Why not?
- A link to Baryogenesis?
- A link to Dark Matter?
- It could be related to the electroweak scale

Flavor from the Electroweak Scale?

$$y_b \left(\frac{\boxed{S}}{\Lambda}\right)^{n_b} \bar{Q}_L H b_R \;\rightarrow\; y_b \left(\frac{\boxed{H^\dagger H}}{\Lambda^2}\right)^{n_b} \bar{Q}_L H b_R$$

with $$\epsilon = \frac{v^2}{2\Lambda^2} = \frac{m_b}{m_t} \quad \Rightarrow \quad \Lambda \approx (5-6)\,v$$

Two drawbacks:

[Babu 033002]

[Giudice, Lebedev 0804.1753]

- The flavon is a flavor singlet
- The coupling to b quarks is

$$g_{hbb} \propto 3\frac{m_b}{v} \qquad \Gamma(h \to b\bar{b}) \approx 9 \times \Gamma(h \to b\bar{b})_{\rm SM}$$

Flavor from the Electroweak Scale

$$y_b \left(\frac{S}{\Lambda}\right)^{n_b} \bar{Q}_L H b_R \to y_b \left(\frac{H_u H_d}{\Lambda^2}\right)^{n_b} \bar{Q}_L H_d b_R$$

with $\epsilon = \dfrac{v_u v_d}{\Lambda^2} = \dfrac{m_b}{m_t} \quad \Rightarrow \quad \Lambda \approx (5-6)\, v \sqrt{\dfrac{\tan\beta}{1+\tan^2\beta}}$

$$\tan\beta = \mathcal{O}(1), \quad \Lambda \approx 1\mathrm{TeV}$$

Flavor from the Electroweak Scale

$$\mathcal{L}_{\text{Yuk}} = y^u_{ij} \left(\frac{H_u H_d}{\Lambda^2}\right)^{a_i - a_{u_j} - a_{H_u}} \bar{Q}_i H_u u_{Rj} + y^d_{ij} \left(\frac{H_u H_d}{\Lambda^2}\right)^{a_i - a_{d_j} - a_{H_d}} \bar{Q}_i H_d d_{Rj} + h.c.$$

11 Flavour charges, 8 + 2 conditions

$$m_t \approx \frac{v_u}{\sqrt{2}}\,, \quad \frac{m_b}{m_t} \approx \frac{m_c}{m_t} \approx \varepsilon^1\,, \quad \frac{m_s}{m_t} \approx \varepsilon^2\,, \quad \frac{m_d}{m_t} \approx \frac{m_u}{m_t} \approx \varepsilon^3\,.$$

$$V_{\text{CKM}} \approx \begin{pmatrix} 1 & 1 & \epsilon \\ 1 & 1 & \epsilon \\ \epsilon & \epsilon & 1 \end{pmatrix}$$

$$\begin{aligned} m_t &= 172\,\text{GeV} \\ m_b \approx m_c &\approx 2.9\,\text{GeV} \\ m_s &= 50\,\text{MeV} \\ m_u = m_d &\approx 1\,\text{MeV} \end{aligned}$$

A rescaling freedom remains

$$\begin{aligned} a_{H_u} &= 1\,, \\ a_{H_d} &= 0\,, \end{aligned} \qquad \begin{aligned} a_1 &= 2\,, & a_u &= -2\,, & a_d &= -1\,, \\ a_2 &= 2\,, & a_c &= 0\,, & a_s &= 0\,, \\ a_3 &= 1\,, & a_t &= 0\,, & a_b &= 0\,. \end{aligned}$$

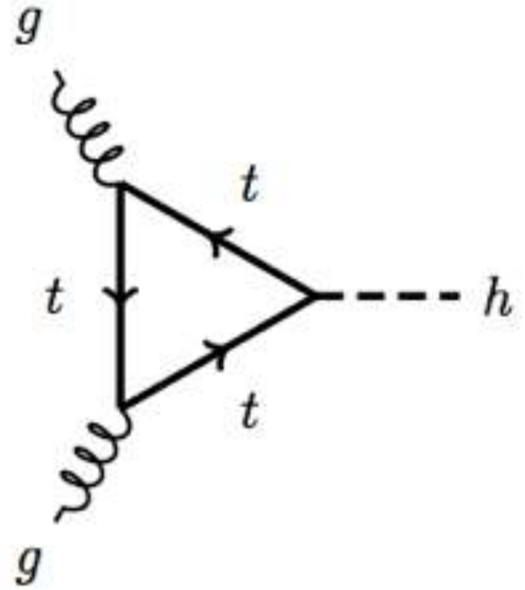
Higgs Couplings

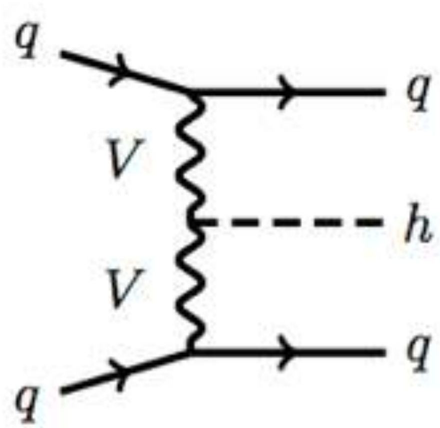
- Couplings are rescaled $g_{hVV} = \kappa_V \, g_{hVV}^{\mathrm{SM}}$ $g_{hff} = \kappa_f \, g_{hff}^{\mathrm{SM}}$
- To $W^{\pm}, Z$ fixed by gauge symmetry:

$$\kappa_V = \sin(\beta - \alpha)$$

- To the top: $\kappa_t = \kappa_t^{\mathrm{II\,HDM}} = \dfrac{\cos\alpha}{\sin\beta}$

Higgs Production like in a 2HDM of type II

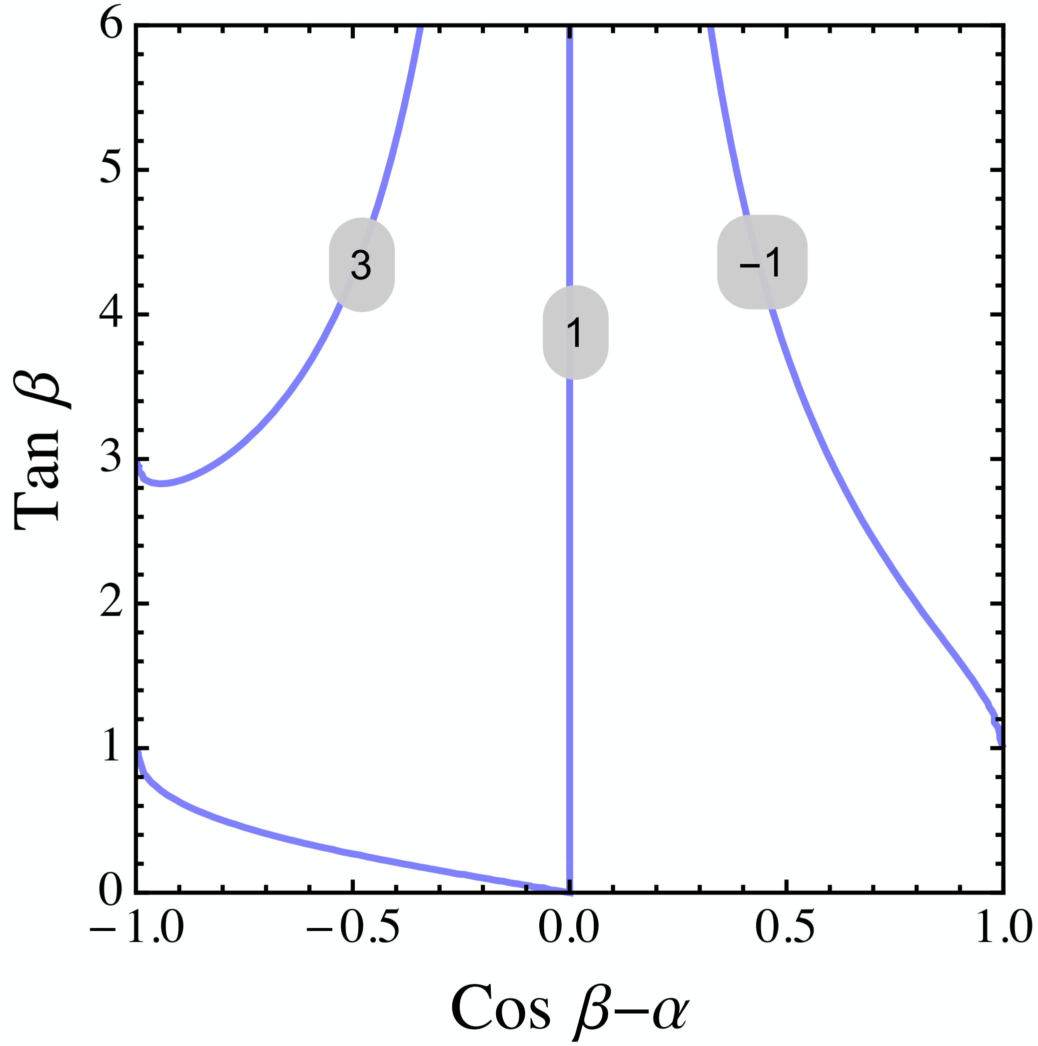




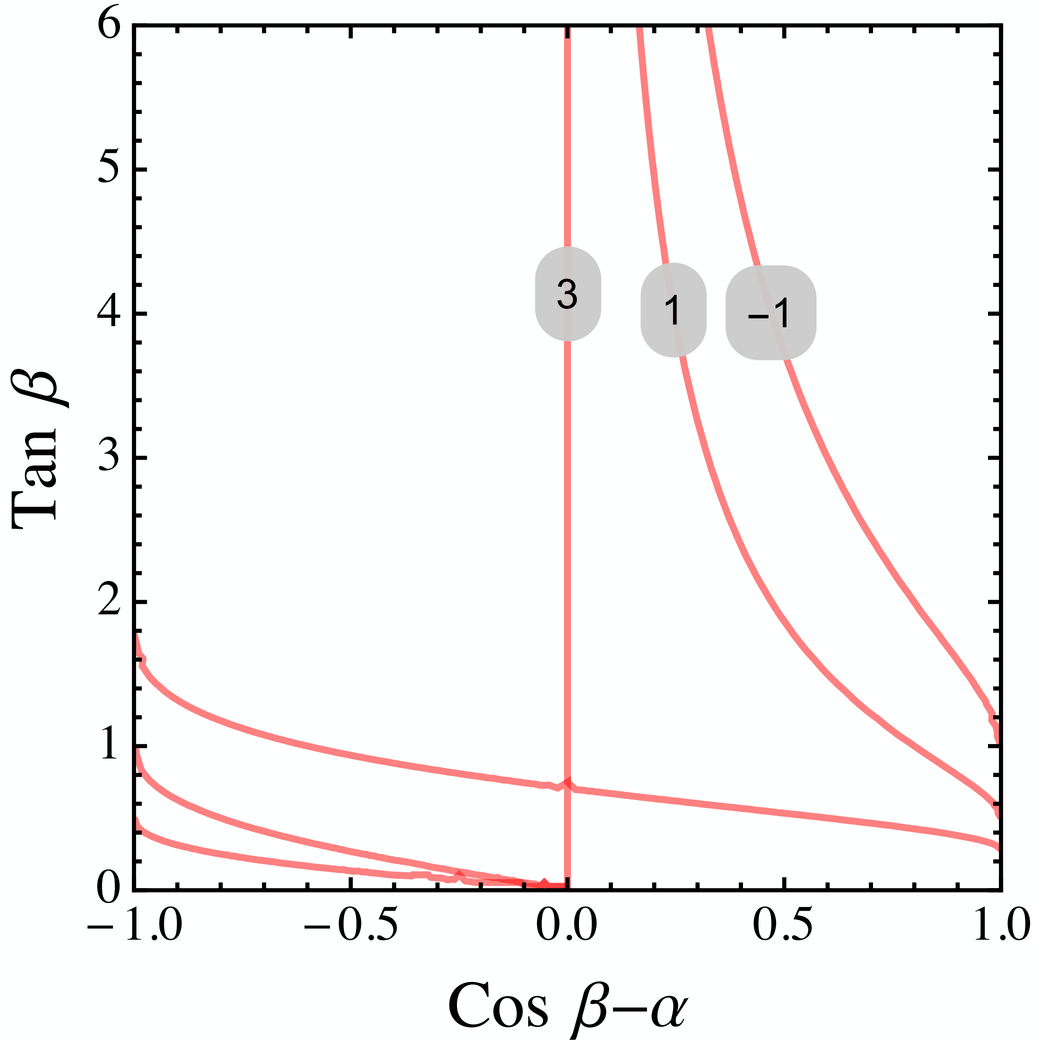
Higgs Couplings

- To the bottom:

$$\kappa_b^{\mathrm{II\,HDM}} = -\frac{\sin\alpha}{\cos\beta}$$



$$\kappa_b = -2\frac{\sin\alpha}{\cos\beta} + \frac{\cos\alpha}{\sin\beta}$$



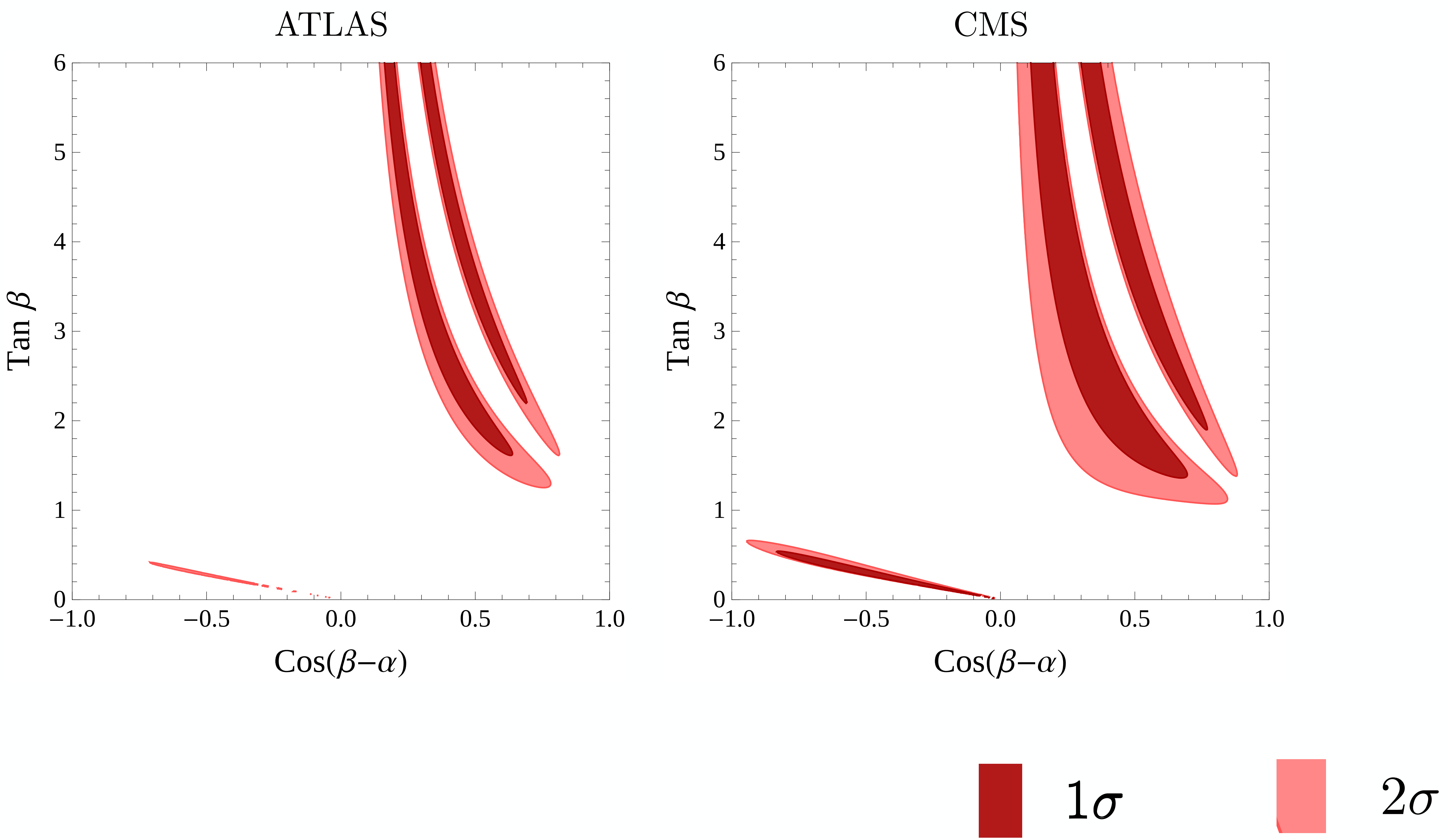
Global Higgs Fit

We performed a global Higgs fit to 8 different channels at ATLAS & CMS

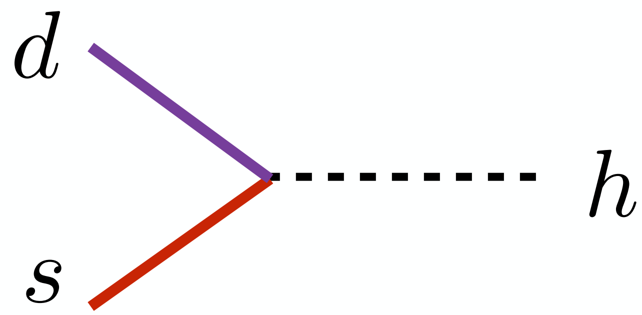
$$\mu_X = \frac{\sigma_{\text{prod}}}{\sigma^{\text{SM}}_{\text{prod}}} \frac{\Gamma_{h\to X}}{\Gamma^{\text{SM}}_{h\to X}} \frac{\Gamma^{\text{SM}}_{h,\,\text{tot}}}{\Gamma_h}$$

Decay Mode	Production Channels $\sigma_{gg\to h}$, $\sigma_{t\bar{t}\to h}$	Production Channels σ_{VBF}, σ_{VH}	Experiment
$h \to WW^*$	$\mu_W = 1.02^{+0.29}_{-0.26}$ [17] $\mu_W \simeq 0.75 \pm 0.35$ [18]	$\mu_W = 1.27^{+0.53}_{-0.45}$ [17] $\mu_W \simeq 0.7 \pm 0.85$ [18]	ATLAS CMS
$h \to ZZ^*$	$\mu_Z = 1.7^{+0.5}_{-0.4}$ [19] $\mu_Z = 0.8^{+0.46}_{-0.36}$ [20]	$\mu_Z = 0.3^{+1.6}_{-0.9}$ [19] $\mu_Z = 1.7^{+2.2}_{-2.1}$ [20]	ATLAS CMS
$h \to \gamma\gamma$	$\mu_\gamma = 1.32 \pm 0.38$ [21] $\mu_\gamma = 1.13^{+0.37}_{-0.31}$ [22]	$\mu_\gamma = 0.8 \pm 0.7$ [21] $\mu_\gamma = 1.16^{+0.63}_{-0.58}$ [22]	ATLAS CMS
$h \to \bar{b}b$	$\mu_b = 1.5 \pm 1.1$ [23] $\mu_b = 0.67^{+1.35}_{-1.33}$ [25]	$\mu_b = 0.52 \pm 0.32 \pm 0.24$ [24] $\mu_b = 1.0 \pm 0.5$ [26]	ATLAS CMS
$h \to \tau\tau$	$\mu_\tau = 2.0 \pm 0.8^{+1.2}_{-0.8} \pm 0.3$ [27] $\mu_\tau \simeq 0.5^{+0.8}_{-0.7}$ [28]	$\mu_\tau = 1.24^{+0.49}_{-0.45}{}^{+0.31}_{-0.29} \pm 0.08$ [27] $\mu_\tau \simeq 1.1^{+0.7}_{-0.5}$ [28]	ATLAS CMS

Global Higgs Fit



Flavor from the Electroweak Scale



$$\sim g_{hsd}$$

Universal function

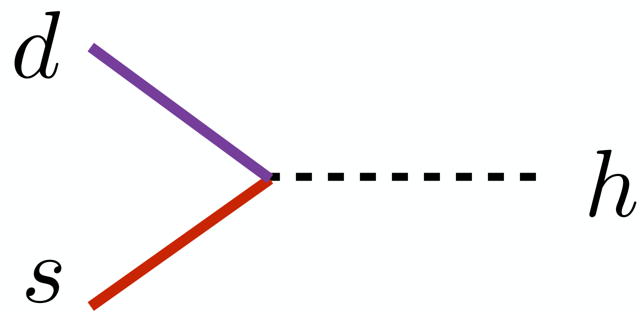
$$f^h(\alpha,\beta) = \frac{c_\alpha}{s_\beta} - \frac{s_\alpha}{c_\beta}$$

$$g_{hd_id_i} = \left(\frac{c_\alpha}{s_\alpha} + n_{d_i} f^h(\alpha,\beta)\right)\frac{m_{d_i}}{v}$$

$$g_{hd_id_j} = \left(\mathcal{Q}_{ij}\frac{m_{d_j}}{v} - \frac{m_{d_i}}{v}\mathcal{D}_{ij}\right) f^h(\alpha,\beta)$$

$$\mathcal{Q}^u \sim \mathcal{Q}^d \sim \begin{pmatrix} 2 & \varepsilon^2 & \varepsilon \\ \varepsilon^2 & 2 & \varepsilon \\ \varepsilon & \varepsilon & 1 \end{pmatrix}, \qquad \mathcal{U} \sim \begin{pmatrix} -2 & \varepsilon^2 & \varepsilon^2 \\ \varepsilon^2 & \varepsilon^2 & \varepsilon^4 \\ \varepsilon^2 & \varepsilon^4 & \varepsilon^4 \end{pmatrix}, \qquad \mathcal{D} \sim \begin{pmatrix} -1 & \varepsilon & \varepsilon \\ \varepsilon & \varepsilon^2 & \varepsilon^2 \\ \varepsilon & \varepsilon^2 & \varepsilon^2 \end{pmatrix}$$

Flavor from the Electroweak Scale



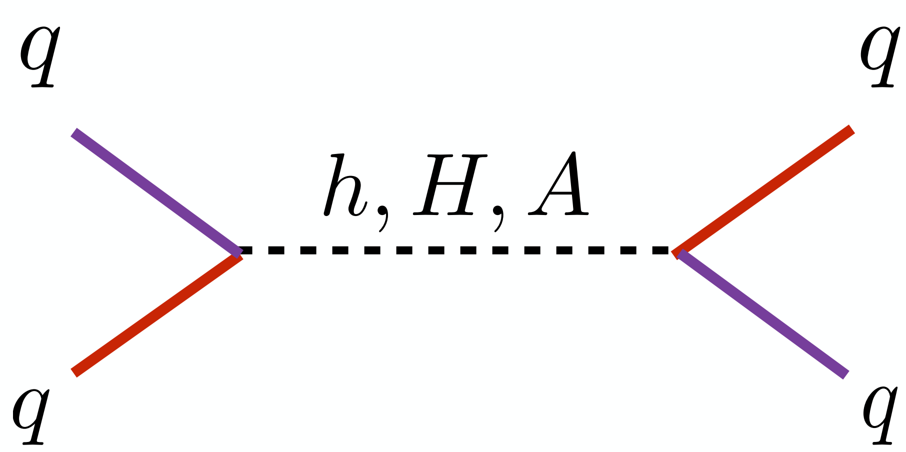
$$\sim \quad \frac{m_s}{v}\,\epsilon\, f^h(\alpha,\beta) \approx 3\times 10^{-6} \times f^h(\alpha,\beta)$$

$$g_{hd_id_i} = \left(\frac{c_\alpha}{s_\alpha} + n_{d_i} f^h(\alpha,\beta)\right)\frac{m_{d_i}}{v}$$

$$g_{hd_id_j} = \left(\mathcal{Q}_{ij}\frac{m_{d_j}}{v} - \frac{m_{d_i}}{v}\mathcal{D}_{ij}\right) f^h(\alpha,\beta)$$

$$\mathcal{Q}^u \sim \mathcal{Q}^d \sim \begin{pmatrix} 2 & \varepsilon^2 & \varepsilon \\ \varepsilon^2 & 2 & \varepsilon \\ \varepsilon & \varepsilon & 1 \end{pmatrix}, \qquad \mathcal{U} \sim \begin{pmatrix} -2 & \varepsilon^2 & \varepsilon^2 \\ \varepsilon^2 & \varepsilon^2 & \varepsilon^4 \\ \varepsilon^2 & \varepsilon^4 & \varepsilon^4 \end{pmatrix}, \qquad \mathcal{D} \sim \begin{pmatrix} -1 & \varepsilon & \varepsilon \\ \varepsilon & \varepsilon^2 & \varepsilon^2 \\ \varepsilon & \varepsilon^2 & \varepsilon^2 \end{pmatrix}$$

Flavor from the Electroweak Scale

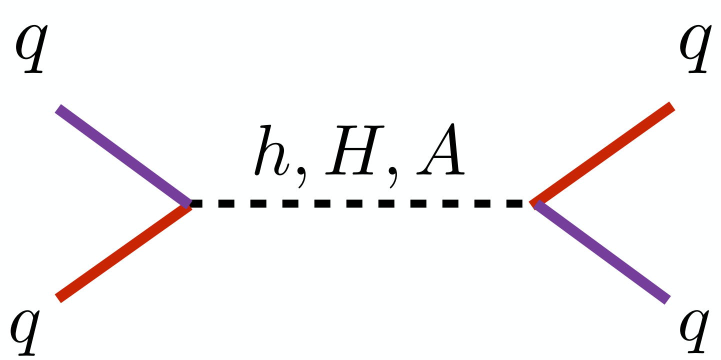


$$g_{hd_id_j} = g^h(\alpha,\beta)\left(\frac{m_d}{v}\right)_{ij} + f^h(\alpha,\beta)\left[\mathcal{Q}^d_{ij}\left(\frac{m_d}{v}\right)_{jj} - \left(\frac{m_d}{v}\right)_{ii}\mathcal{D}_{ij}\right]$$

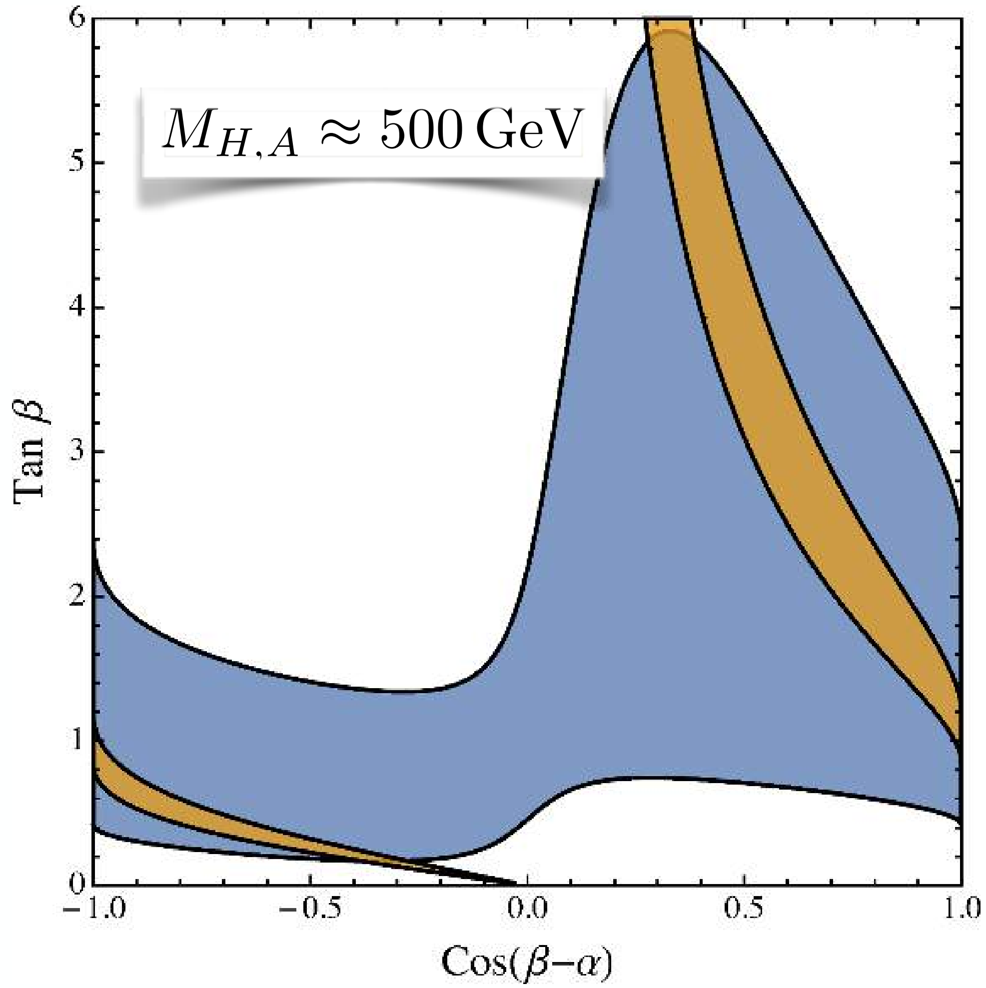
$$g_{Hd_id_j} = G^H(\alpha,\beta)\left(\frac{m_d}{v}\right)_{ij} + F^H(\alpha,\beta)\left[\mathcal{Q}^d_{ij}\left(\frac{m_d}{v}\right)_{jj} - \left(\frac{m_d}{v}\right)_{ii}\mathcal{D}^d_{ij}\right]$$

$$g_{Ad_id_j} = G^A(\alpha,\beta)\left(\frac{m_d}{v}\right)_{ij} + F^A(\alpha,\beta)\left[\mathcal{Q}^d_{ij}\left(\frac{m_d}{v}\right)_{jj} - \left(\frac{m_d}{v}\right)_{ii}\mathcal{D}^d_{ij}\right]$$

Constraints from Meson Mixing

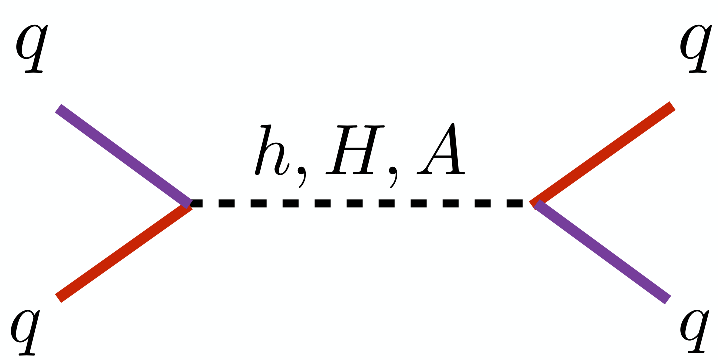


$$\approx \varepsilon^X \frac{m_q^2}{v^2} \left\{ \frac{f^h(\alpha,\beta)^2}{m_h^2} + \frac{F^H(\alpha,\beta)^2}{M_H^2} \pm \frac{F^A(\alpha,\beta)^2}{M_A^2} \right\}$$

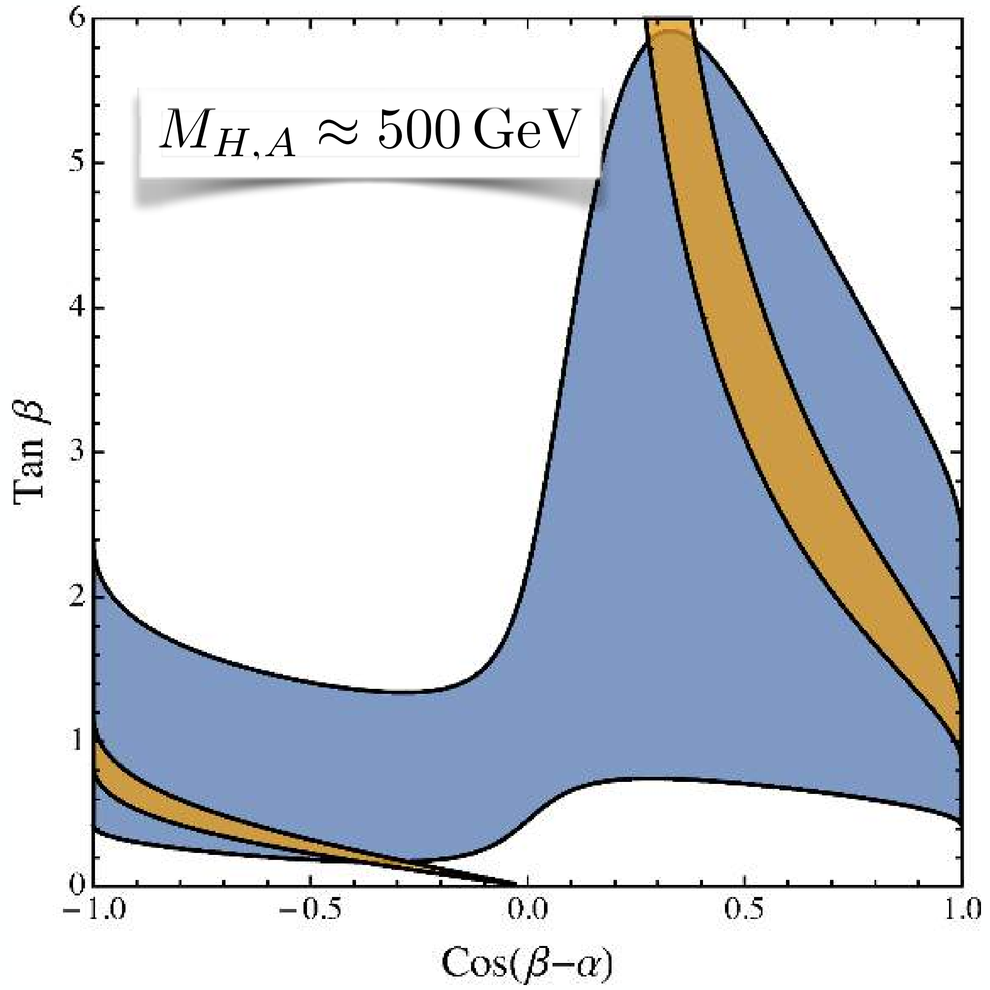


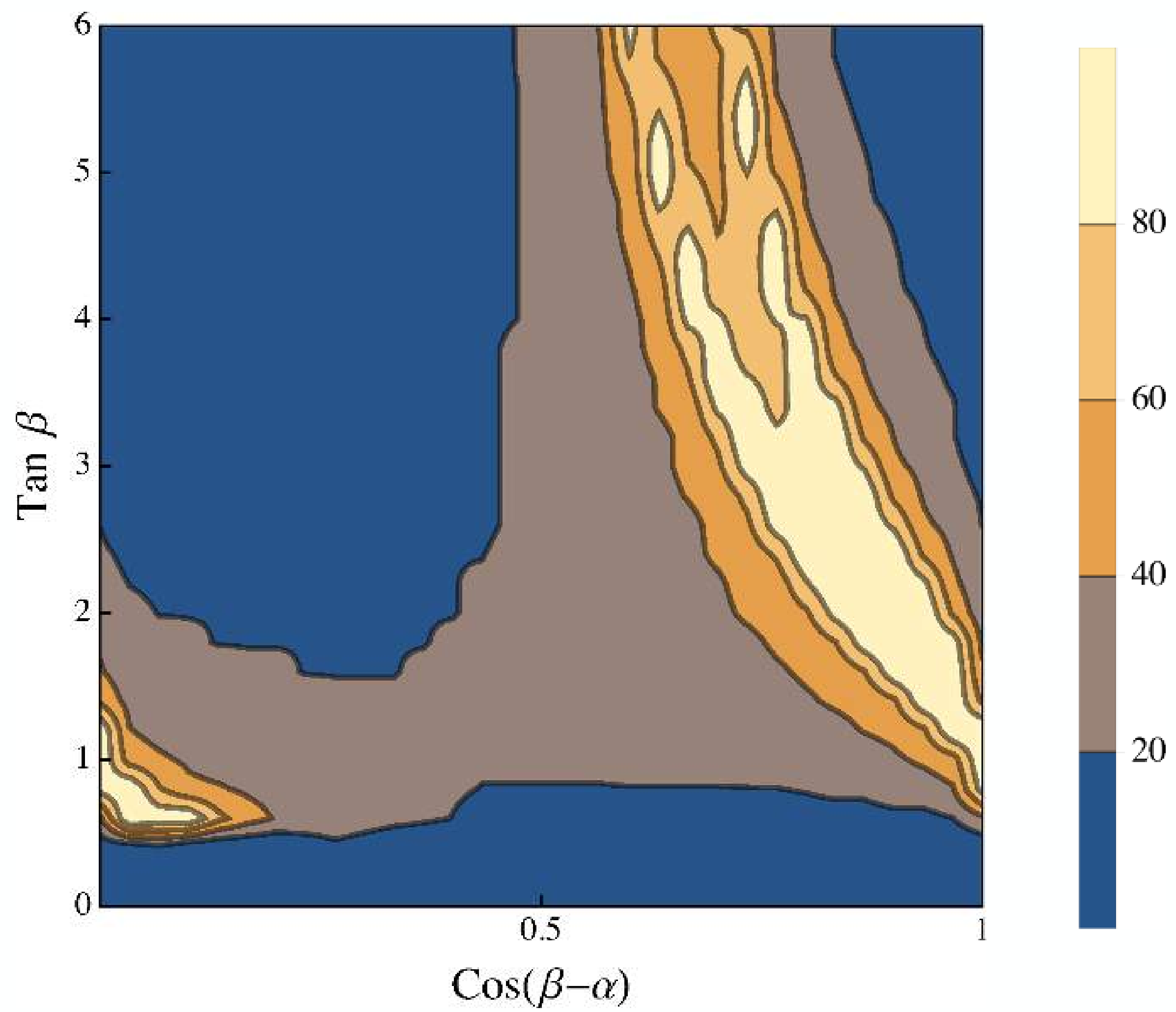
$$\begin{aligned}
\mathcal{H}_{\text{eff}}^{\Delta S=2} = {} & C_1^{sd}\,(\bar{s}_L\,\gamma_\mu\,d_L)^2 + \tilde{C}_1^{sd}\,(\bar{s}_R\,\gamma_\mu\,d_R)^2 \\
& + C_2^{sd}\,(\bar{s}_R\,d_L)^2 + \tilde{C}_2^{sd}\,(\bar{s}_L\,d_R)^2 \\
& + C_4^{sd}\,(\bar{s}_R\,d_L)\,(\bar{s}_L\,d_R) \\
& + C_5^{sd}\,(\bar{s}_L\,\gamma_\mu\,d_L)\,(\bar{s}_R\,\gamma^\mu d_R) \; + h.c.
\end{aligned}$$

$K^0 - \bar{K}^0$ Mixing

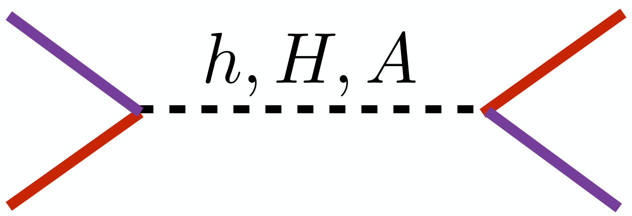


$$\approx \varepsilon^X \frac{m_q^2}{v^2} \left\{ \frac{f^h(\alpha,\beta)^2}{m_h^2} + \frac{F^H(\alpha,\beta)^2}{M_H^2} \pm \frac{F^A(\alpha,\beta)^2}{M_A^2} \right\}$$





$K^0 - \bar{K}^0$ Mixing



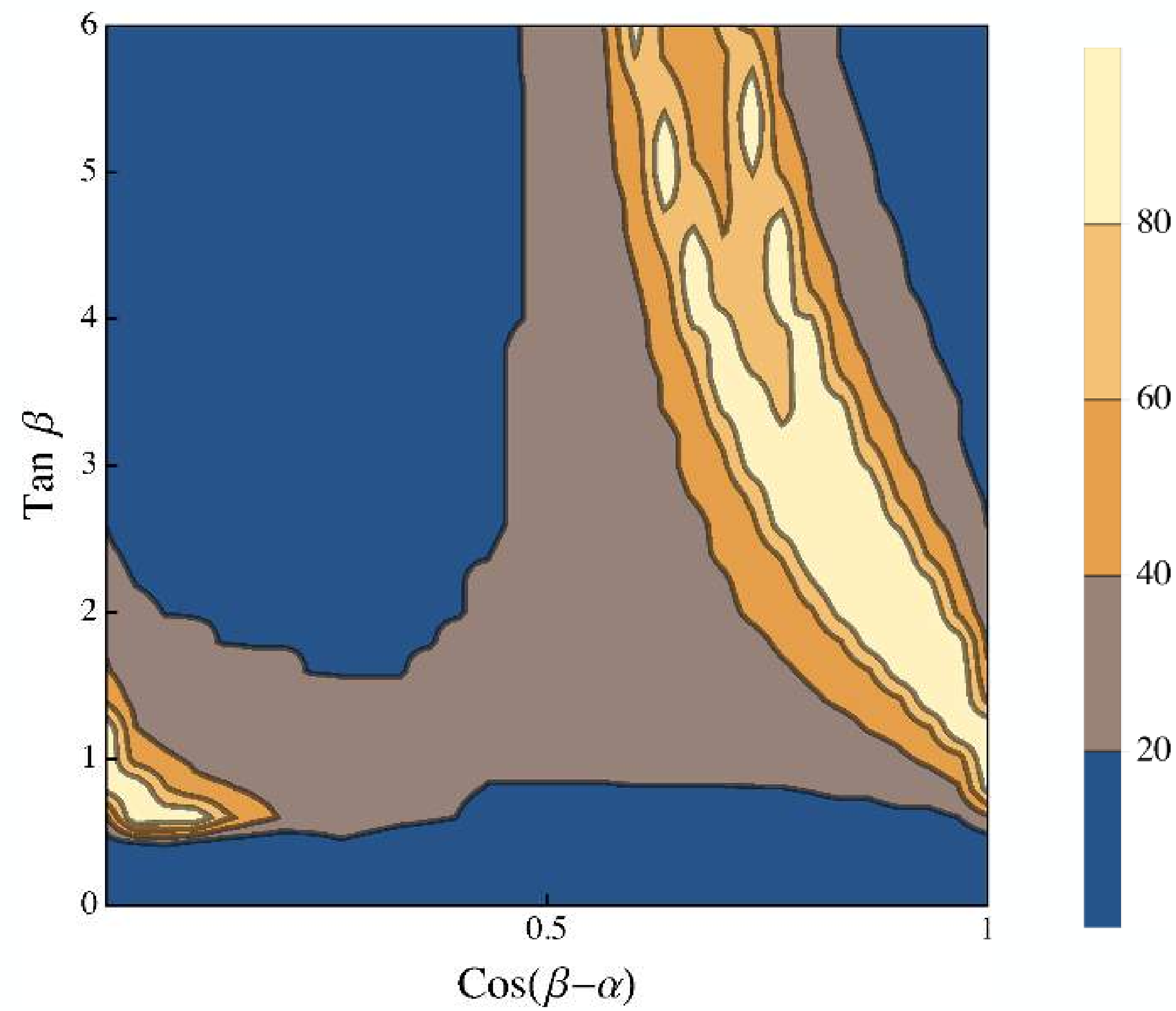
$$\approx \frac{c_i}{v^2}\left\{\frac{f^h(\alpha,\beta)^2}{m_h^2}+\frac{F^H(\alpha,\beta)^2}{M_h^2}\pm\frac{F^A(\alpha,\beta)^2}{M_A^2}\right\}$$

Flavor dependent Wilson Coefficient

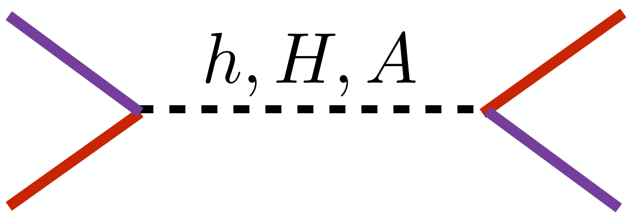
$$\mathcal{Q}^d \sim \begin{pmatrix} 2 & \boxed{\varepsilon^2} & \varepsilon \\ \varepsilon^2 & 2 & \varepsilon \\ \varepsilon & \varepsilon & 1 \end{pmatrix} \qquad \mathcal{D} \sim \begin{pmatrix} -1 & \boxed{\varepsilon} & \varepsilon \\ \varepsilon & \varepsilon^2 & \varepsilon^2 \\ \varepsilon & \varepsilon^2 & \varepsilon^2 \end{pmatrix}$$

$$c_2^{sd} = \varepsilon^4 m_s^2$$
$$\tilde{c}_2^{sd} = \varepsilon^2 m_s^2$$
$$c_4^{sd} = \varepsilon^3 m_s^2$$

$M_{H,A} \approx 500\,\mathrm{GeV}$



$B_s^0 - \bar{B}_s^0$ Mixing



$$\approx \frac{c_i}{v^2}\left\{\frac{f^h(\alpha,\beta)^2}{m_h^2} + \frac{F^H(\alpha,\beta)^2}{M_h^2} \pm \frac{F^A(\alpha,\beta)^2}{M_A^2}\right\}$$

Flavor dependent Wilson Coefficient

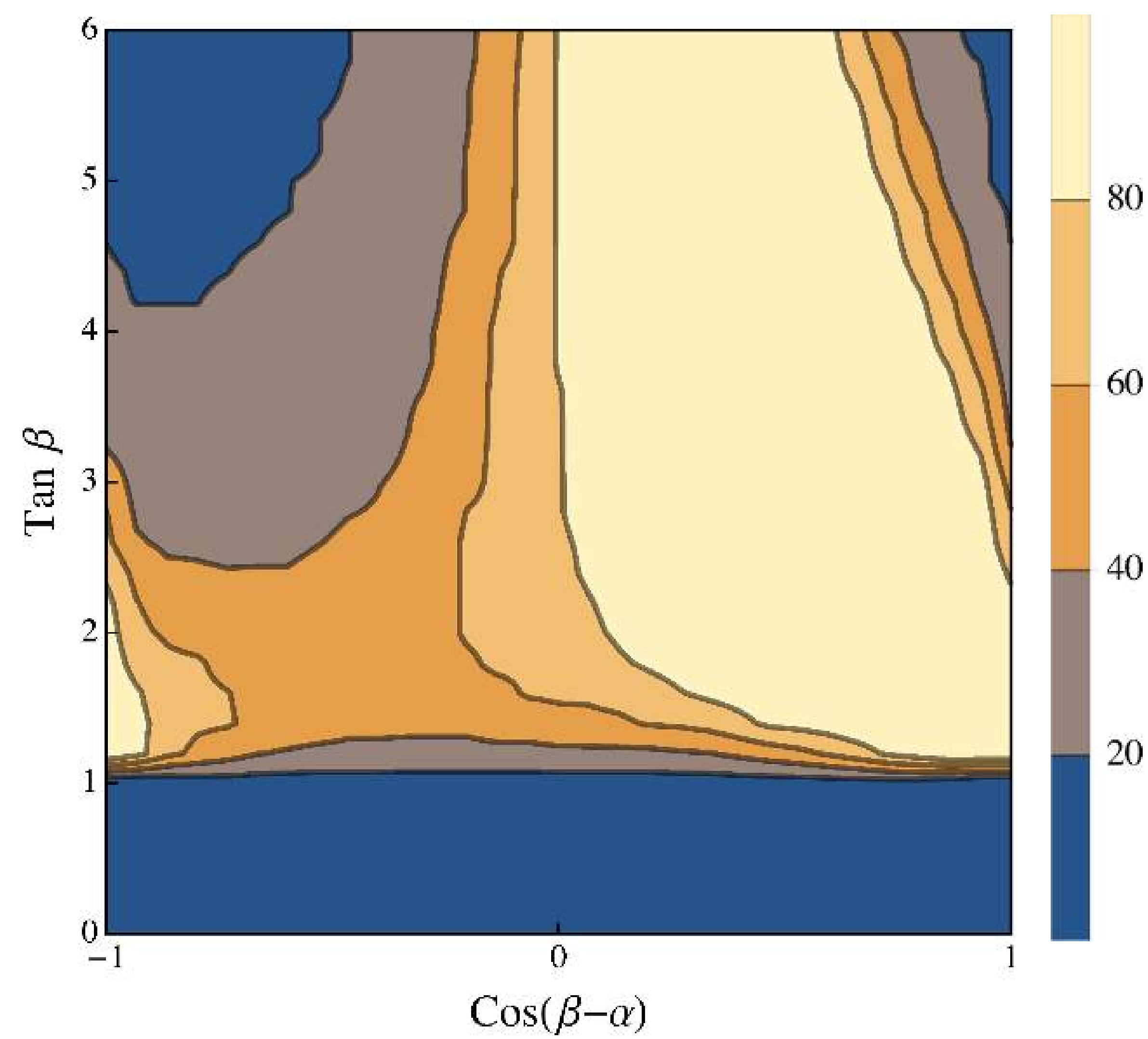
$$\mathcal{Q}^d \sim \begin{pmatrix} 2 & \varepsilon^2 & \varepsilon \\ \varepsilon^2 & 2 & \varepsilon \\ \varepsilon & \varepsilon & 1 \end{pmatrix} \qquad \mathcal{D} \sim \begin{pmatrix} -1 & \varepsilon & \varepsilon \\ \varepsilon & \varepsilon^2 & \varepsilon^2 \\ \varepsilon & \varepsilon^2 & \varepsilon^2 \end{pmatrix}$$

$M_{H,A} \approx 500\,\mathrm{GeV}$

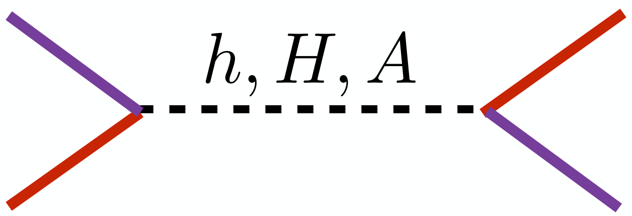
$$c_2^{bs} = \varepsilon^2 m_b^2$$

$$\tilde{c}_2^{bs} = \varepsilon^4 m_b^2$$

$$c_4^{bs} = \varepsilon^3 m_b^2$$



$B_d^0 - \bar{B}_d^0$ Mixing



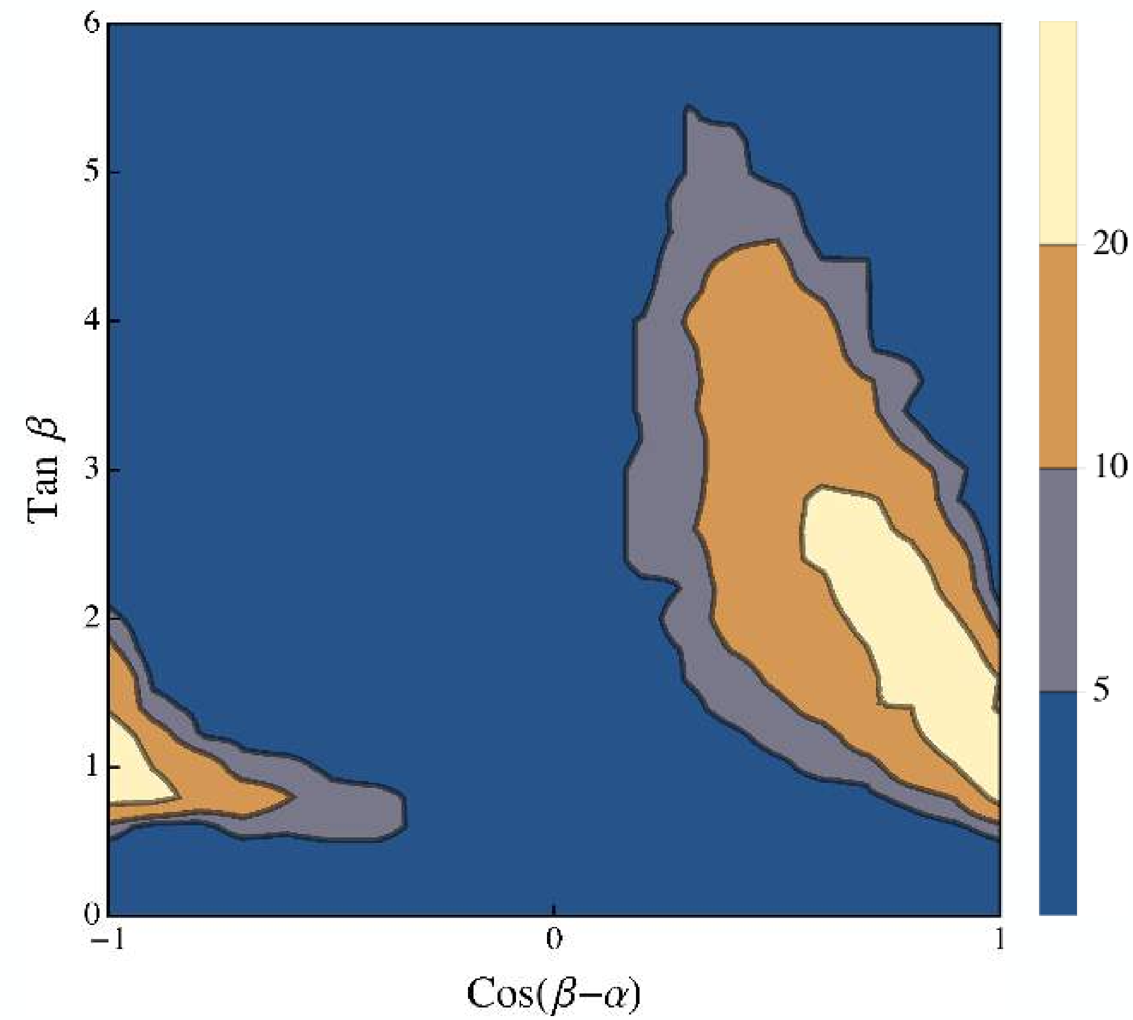
$$\approx \frac{c_i}{v^2}\left\{ \frac{f^h(\alpha,\beta)^2}{m_h^2} + \frac{F^H(\alpha,\beta)^2}{M_h^2} \pm \frac{F^A(\alpha,\beta)^2}{M_A^2} \right\}$$

Flavor dependent Wilson Coefficient

$$\mathcal{Q}^d \sim \begin{pmatrix} 2 & \varepsilon^2 & \varepsilon \\ \varepsilon^2 & 2 & \varepsilon \\ \boxed{\varepsilon} & \varepsilon & 1 \end{pmatrix} \qquad \mathcal{D} \sim \begin{pmatrix} -1 & \varepsilon & \varepsilon \\ \varepsilon & \varepsilon^2 & \varepsilon^2 \\ \boxed{\varepsilon} & \varepsilon^2 & \varepsilon^2 \end{pmatrix}$$

$M_{H,A} \approx 500\,\mathrm{GeV}$

$$c_2^{bd} = \varepsilon^2 m_b^2$$
$$\tilde{c}_2^{bd} = \varepsilon^2 m_b^2$$
$$c_4^{bd} = \varepsilon^2 m_b^2$$



Flavor Bounds

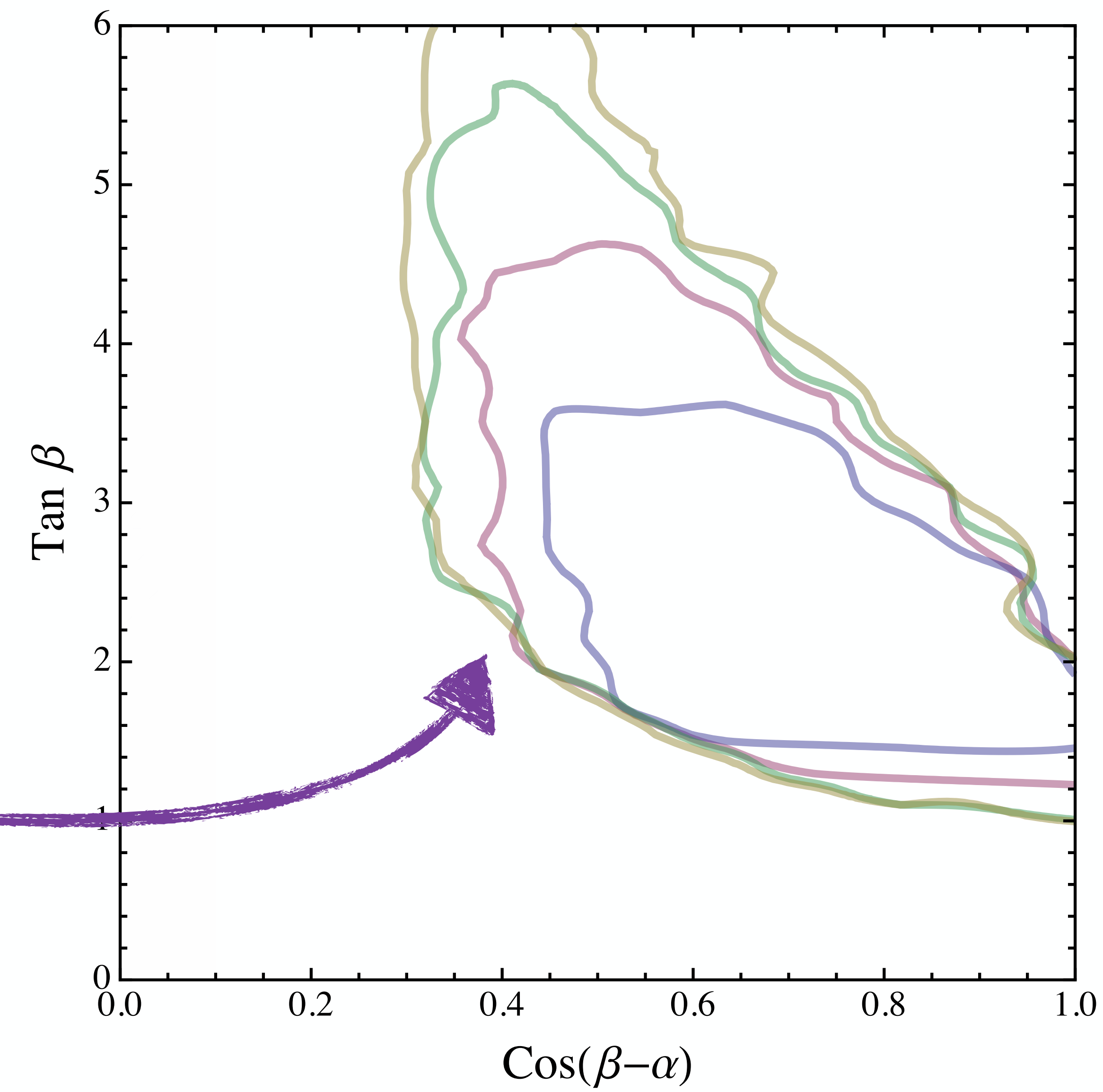
$$M \equiv M_H = M_A = M_{H^+}$$

$M = 700\,\mathrm{GeV}$

$M = 600\,\mathrm{GeV}$

$M = 500\,\mathrm{GeV}$

$M = 400\,\mathrm{GeV}$



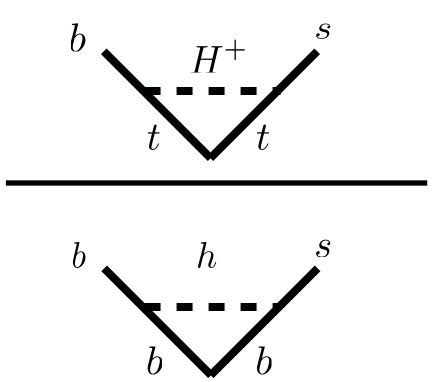
10% Contours

Flavor Bounds

- Contributions to rare leptonic decays depend on the lepton flavor sector
- Bounds from Loop induced processes like $b \to s\gamma$ put constraints on the mass of the charged scalar

$$M_{H^\pm} \gtrsim 358\,(480)\,\mathrm{GeV} \quad @\,99\%(95\%)\ \mathrm{CL}$$ [Misiak et al. 1503.01789]

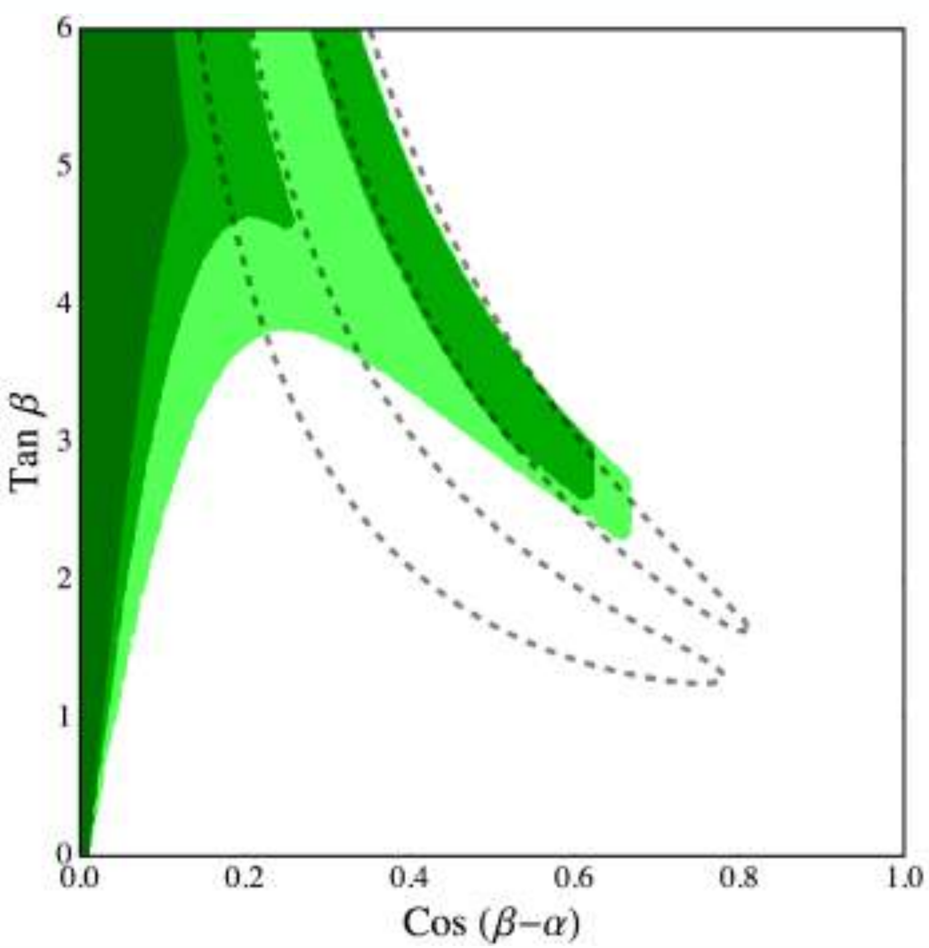
- Neutral Scalar contributions are typically much smaller

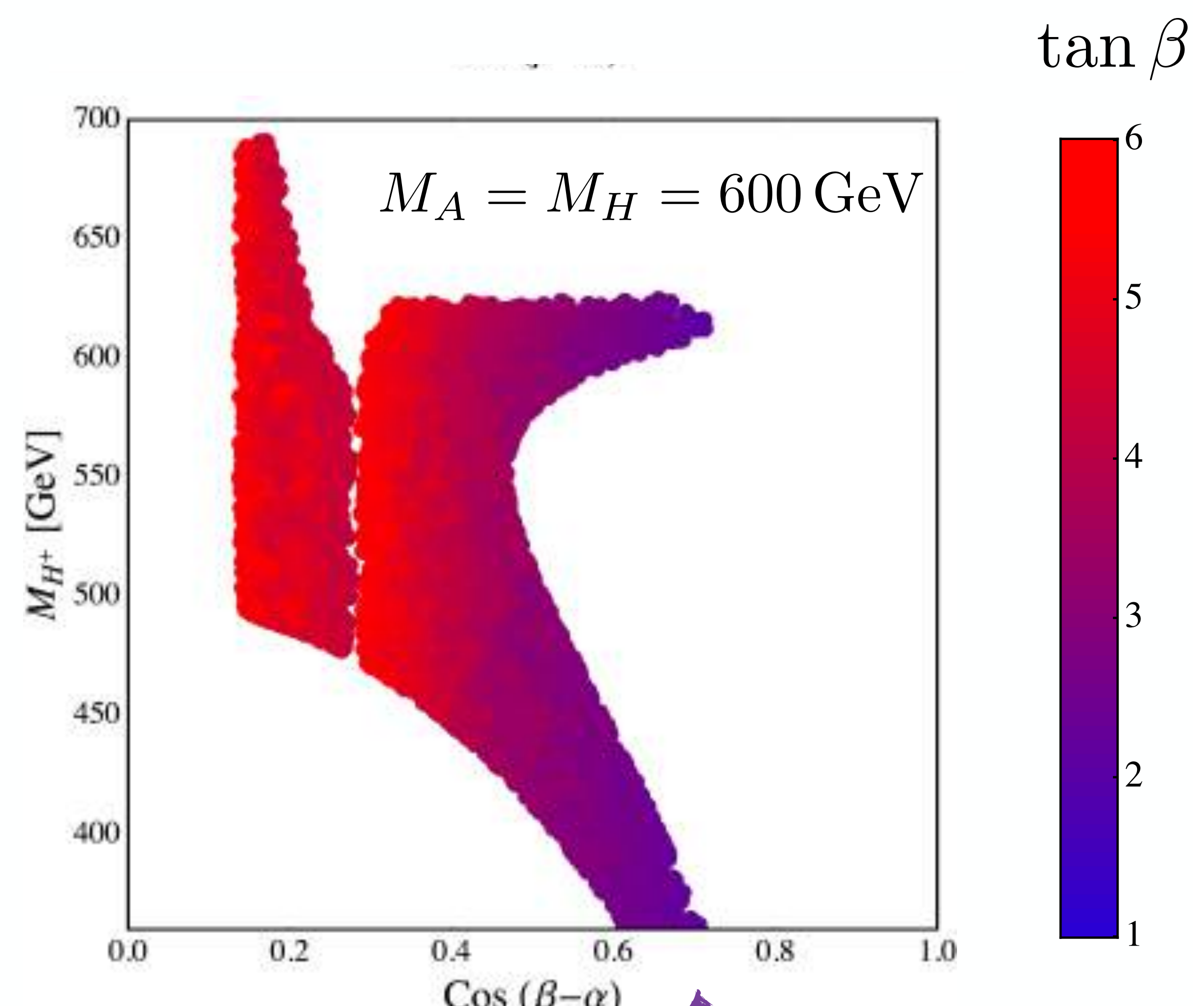


$$\frac{m_t\, V_{tb}\, V_{ts}^*}{m_b\, f(\alpha,\beta)\,\varepsilon} \approx \mathcal{O}(10^2 - 10^3)$$

Decoupling and EWPM

- The global Higgs fit (and Flavor bounds) demand sizable $\cos(\beta-\alpha)$
- Flavor bounds demand heavy extra scalars

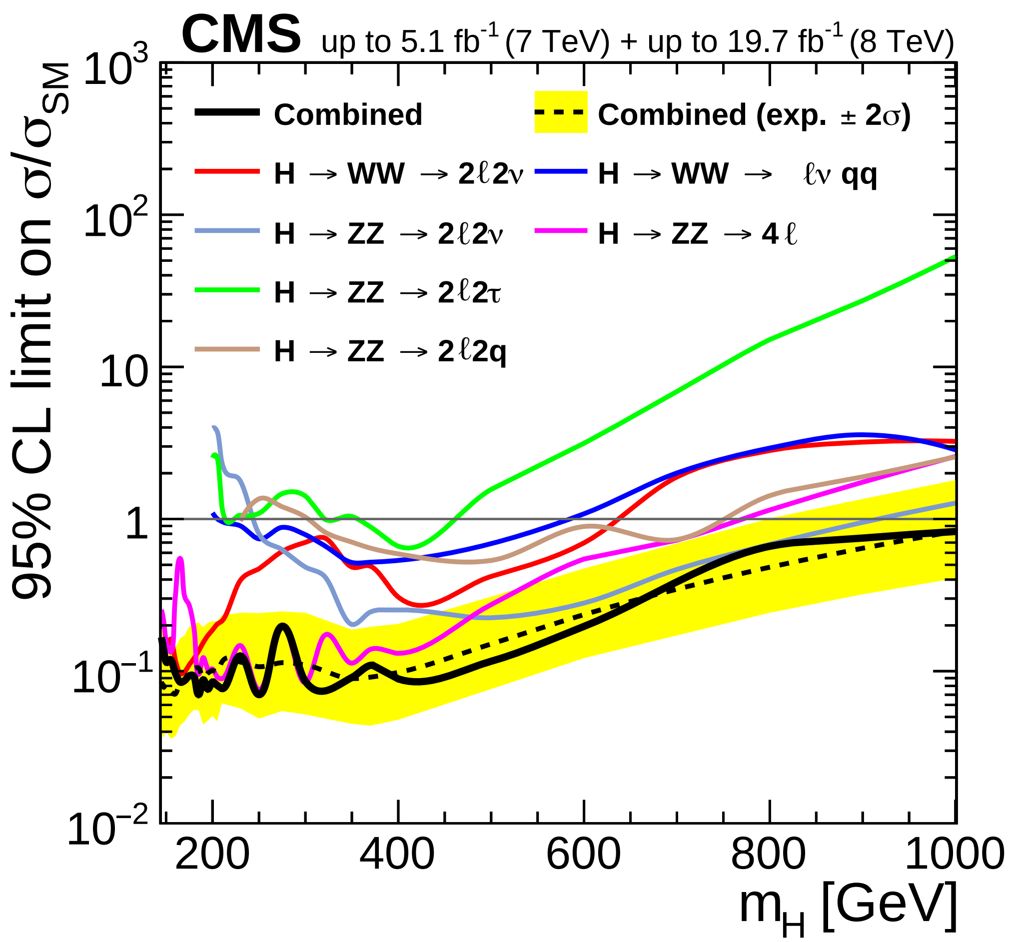


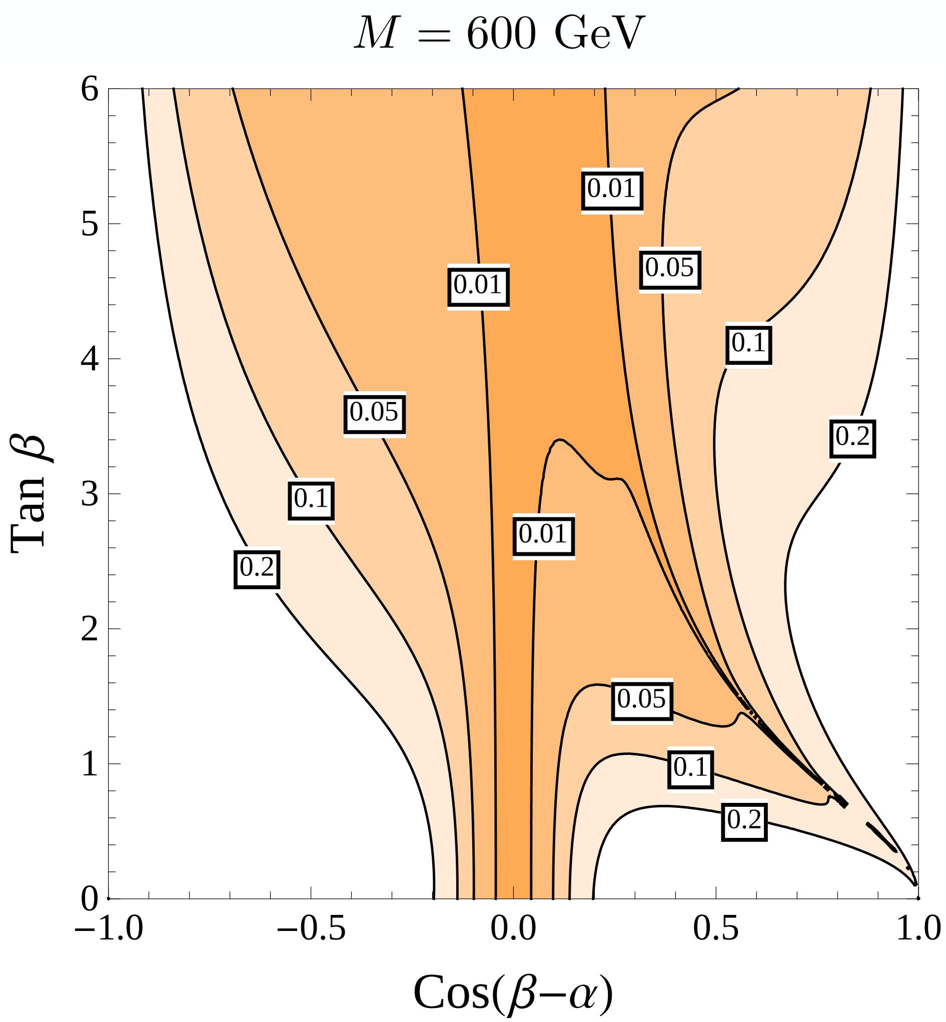


$M = 700\,\text{GeV}$ $M = 600\,\text{GeV}$ $M = 500\,\text{GeV}$

Includes the Higgs fit

Collider Searches for Heavy Higgses: $H \to ZZ$



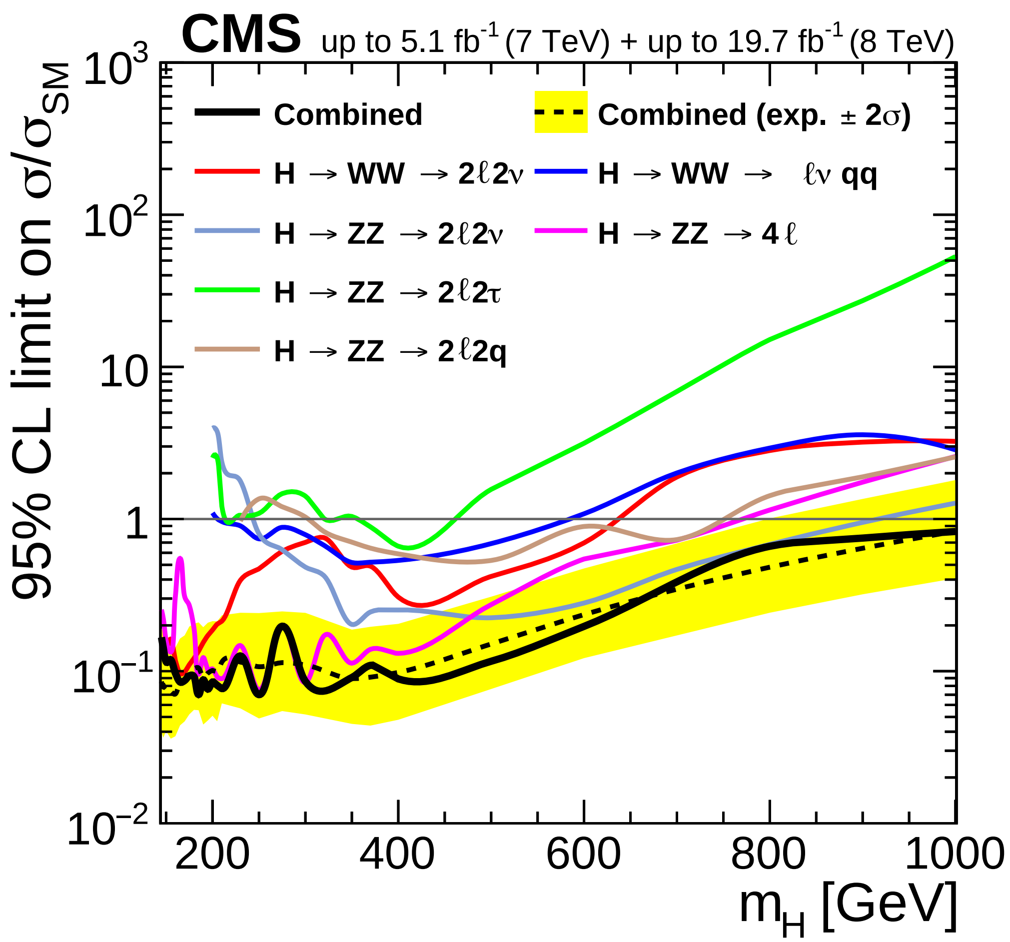


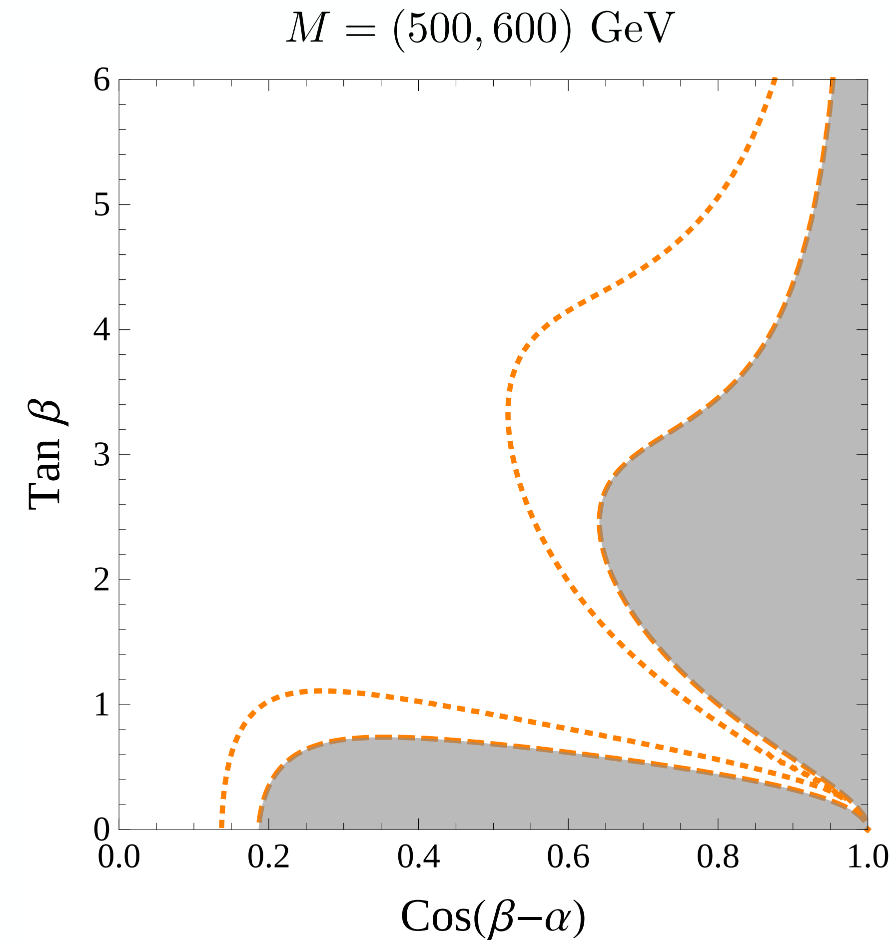
$$\frac{\sigma(gg\to H)\times \mathrm{Br}(H\to VV)}{(\sigma(gg\to H)\times \mathrm{Br}(H\to VV))_{\rm SM}} = (\kappa_t^H)^2\left(1+\xi_b^H\frac{\kappa_b^H}{\kappa_t^H}\right)^2(\kappa_V^H)^2\,\frac{\Gamma_H^{\rm SM}}{\Gamma_H}\,,$$

$$\frac{\sigma(pp\to qqH)\times \mathrm{Br}(H\to VV)}{(\sigma(pp\to qqH)\times \mathrm{Br}(H\to VV))_{\rm SM}} = (\kappa_V^H)^4\,\frac{\Gamma_H^{\rm SM}}{\Gamma_H}\,,$$

[CMS 1504.00936]

Collider Searches for Heavy Higgses: $H \to ZZ$



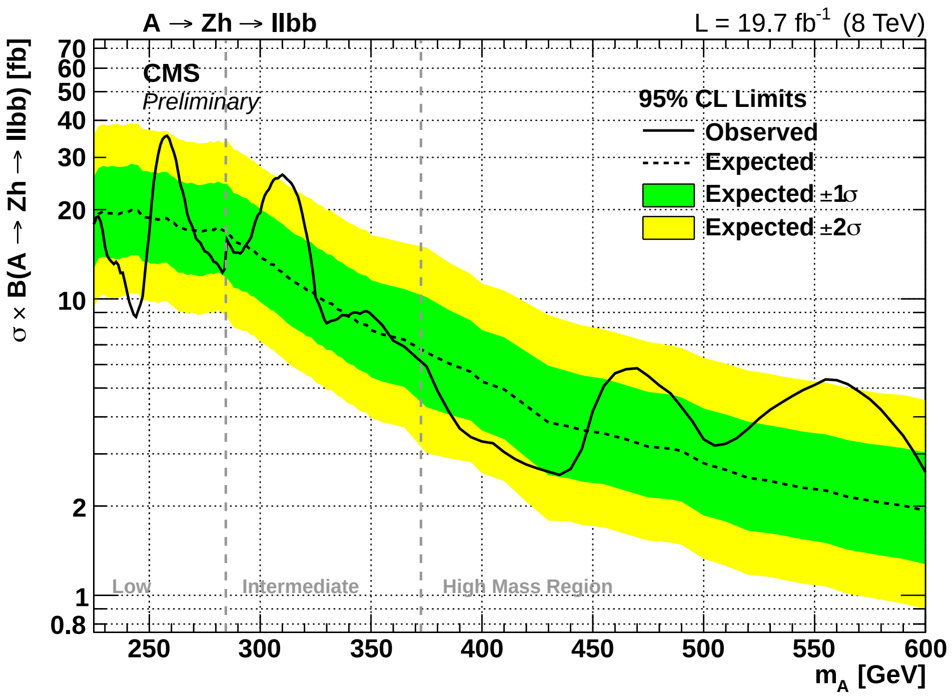


$$\frac{\sigma(gg \to H) \times \mathrm{Br}(H \to VV)}{(\sigma(gg \to H) \times \mathrm{Br}(H \to VV))_{\mathrm{SM}}} = (\kappa_t^H)^2 \left(1 + \xi_b^H \frac{\kappa_b^H}{\kappa_t^H}\right)^2 (\kappa_V^H)^2 \frac{\Gamma_H^{\mathrm{SM}}}{\Gamma_H},$$

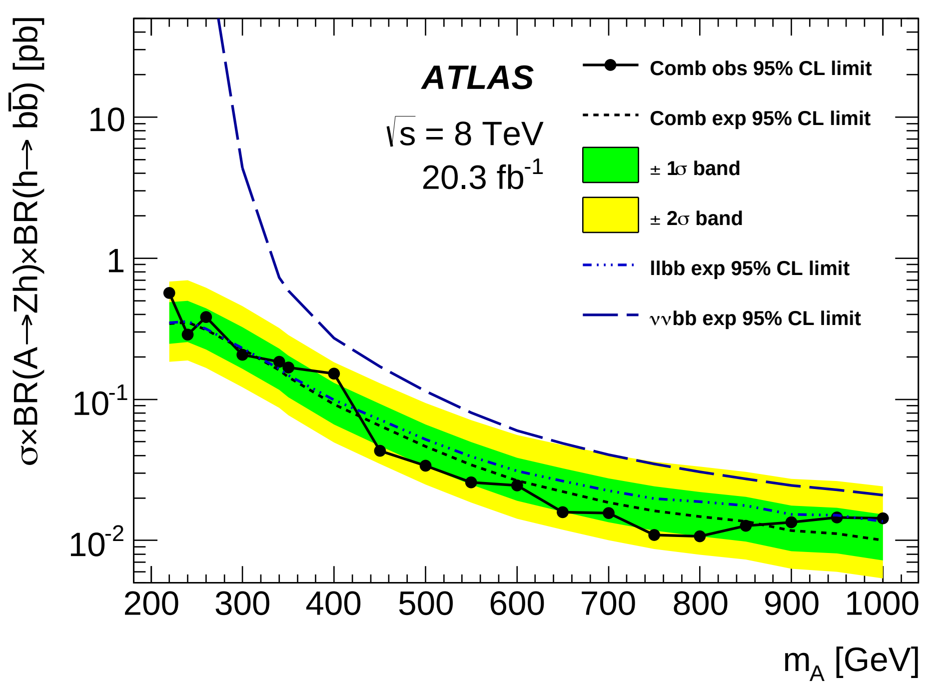
$$\frac{\sigma(pp \to qqH) \times \mathrm{Br}(H \to VV)}{(\sigma(pp \to qqH) \times \mathrm{Br}(H \to VV))_{\mathrm{SM}}} = (\kappa_V^H)^4 \frac{\Gamma_H^{\mathrm{SM}}}{\Gamma_H},$$

[CMS 1504.00936]

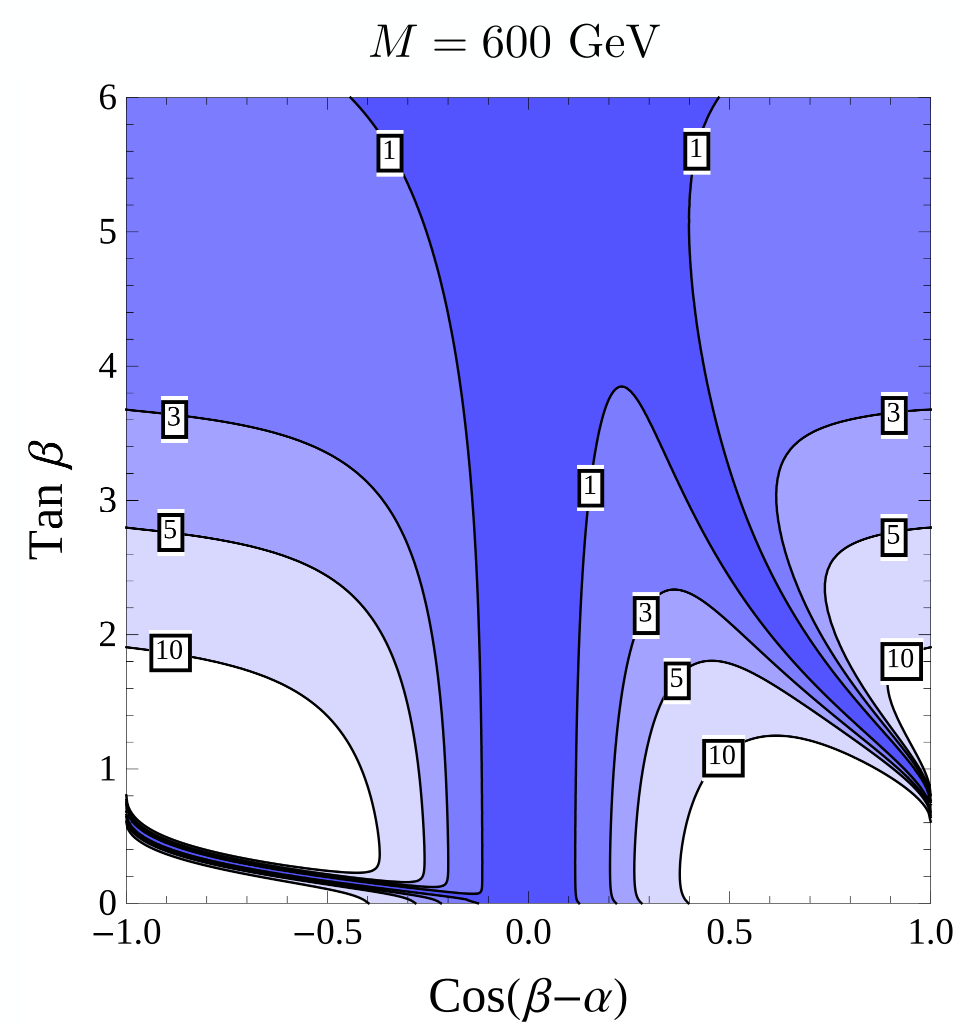
Collider Searches for Heavy Higgses: $A \to Zh$



CMS PAS HIG-14-011



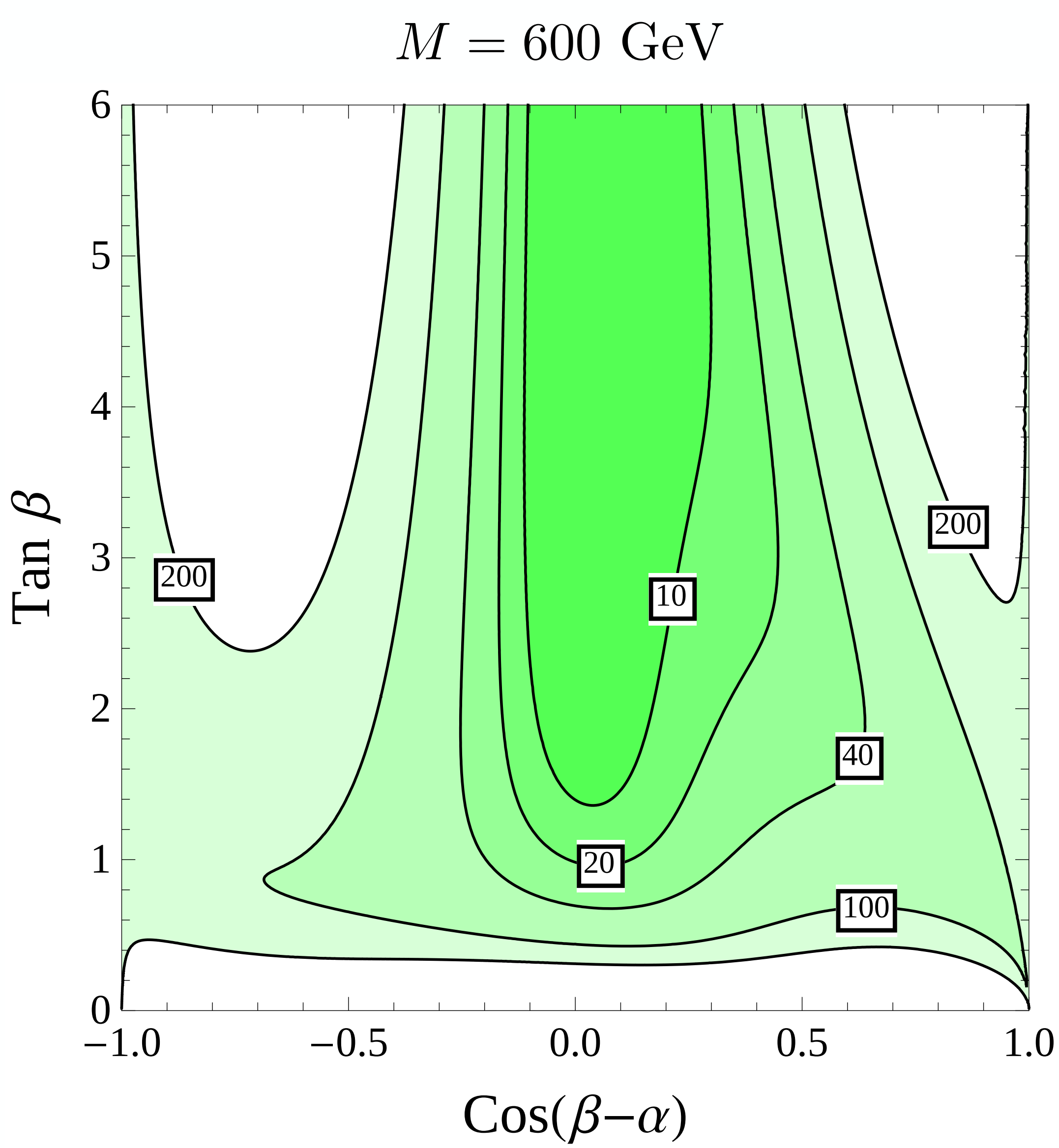
[ATLAS 1502.04478]



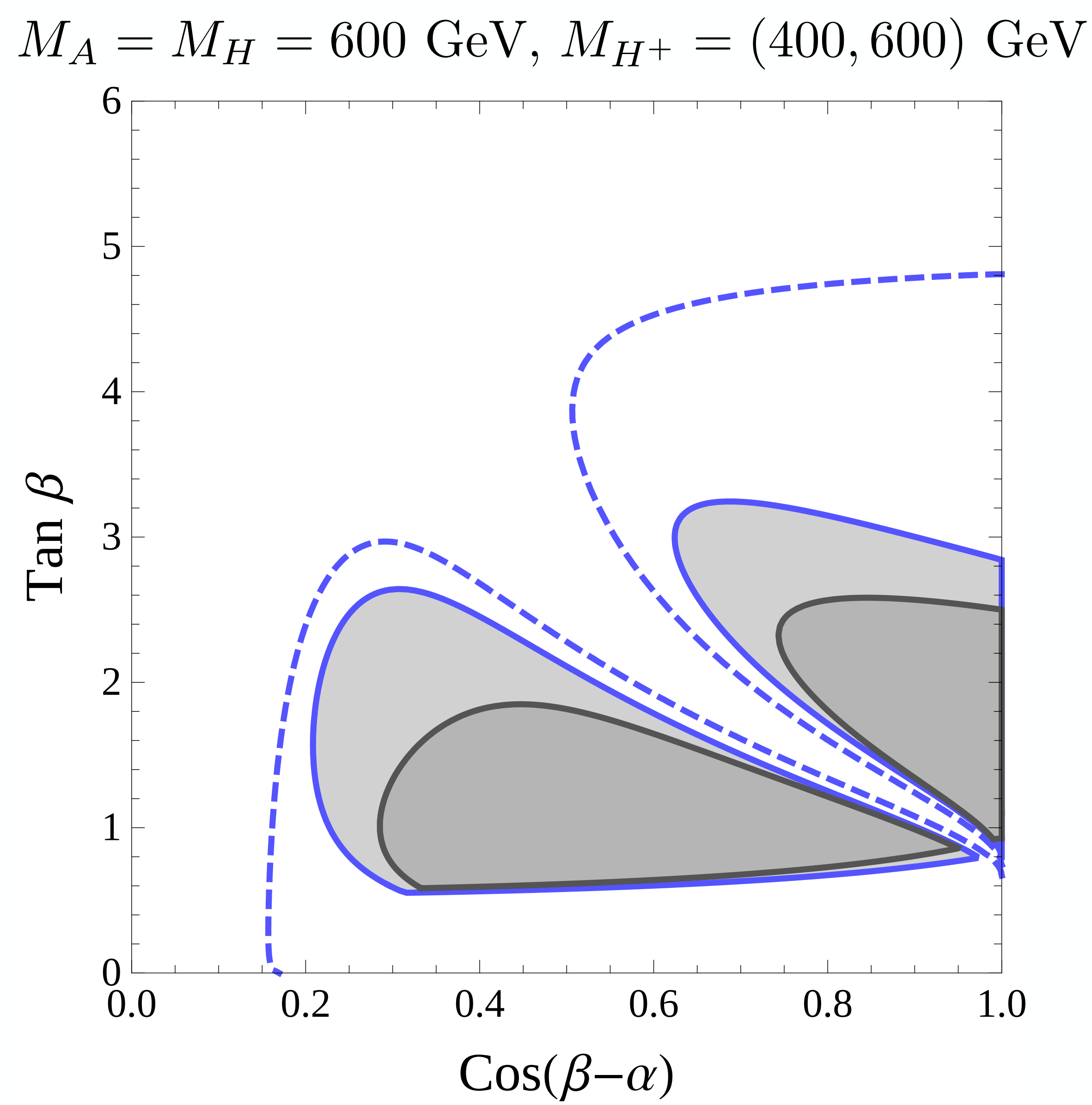
$$\sigma(gg \to A) \times \mathrm{Br}(A \to hZ \to \ell^+\ell^- b\bar{b})$$

Finite Width Effects

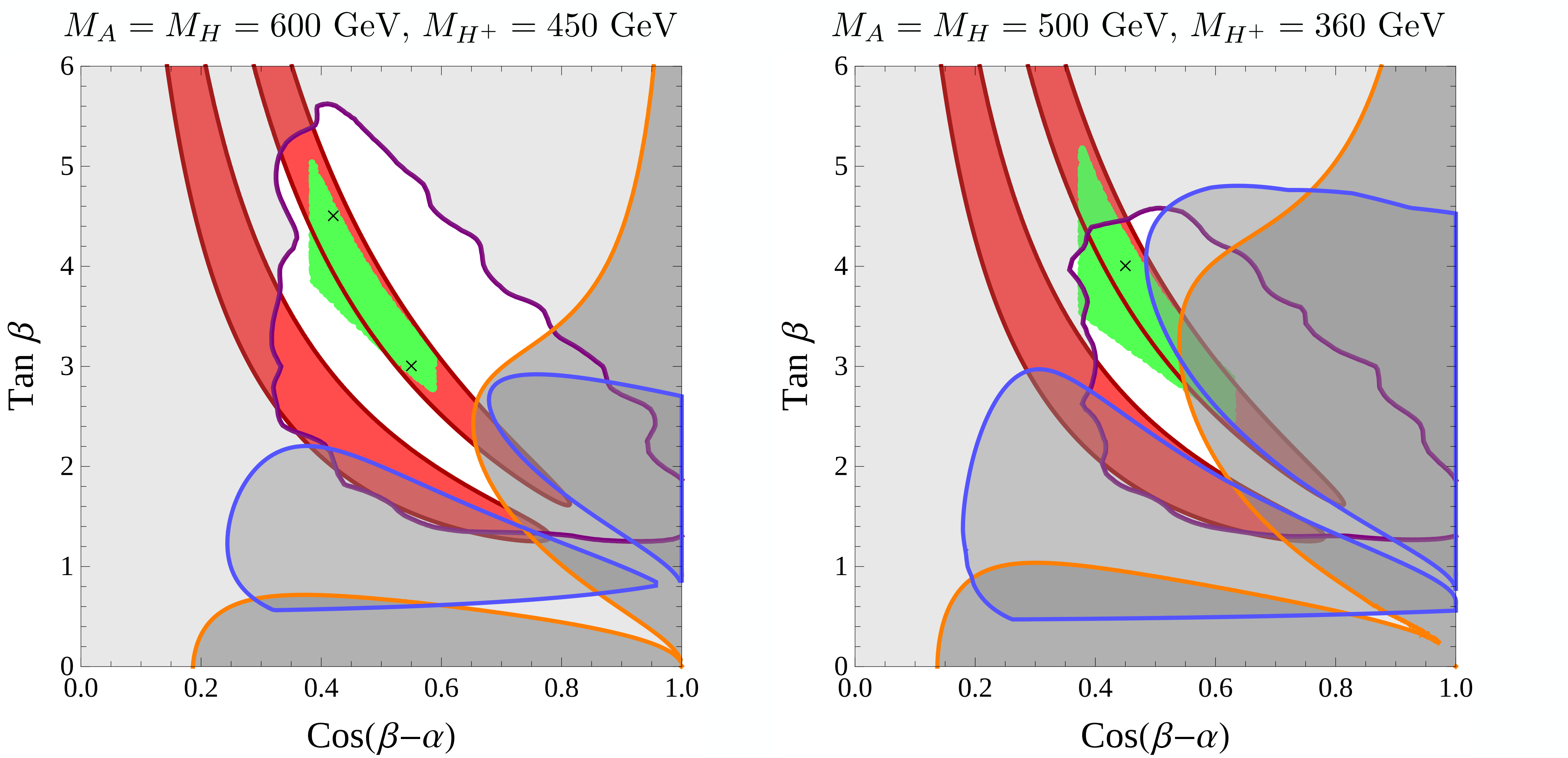
$\Gamma_A[\mathrm{GeV}]$



$\sigma(gg \to A) \times \mathrm{Br}(A \to hZ \to \ell^+\ell^- b\bar{b})$



Final Plots



Flavour constraint

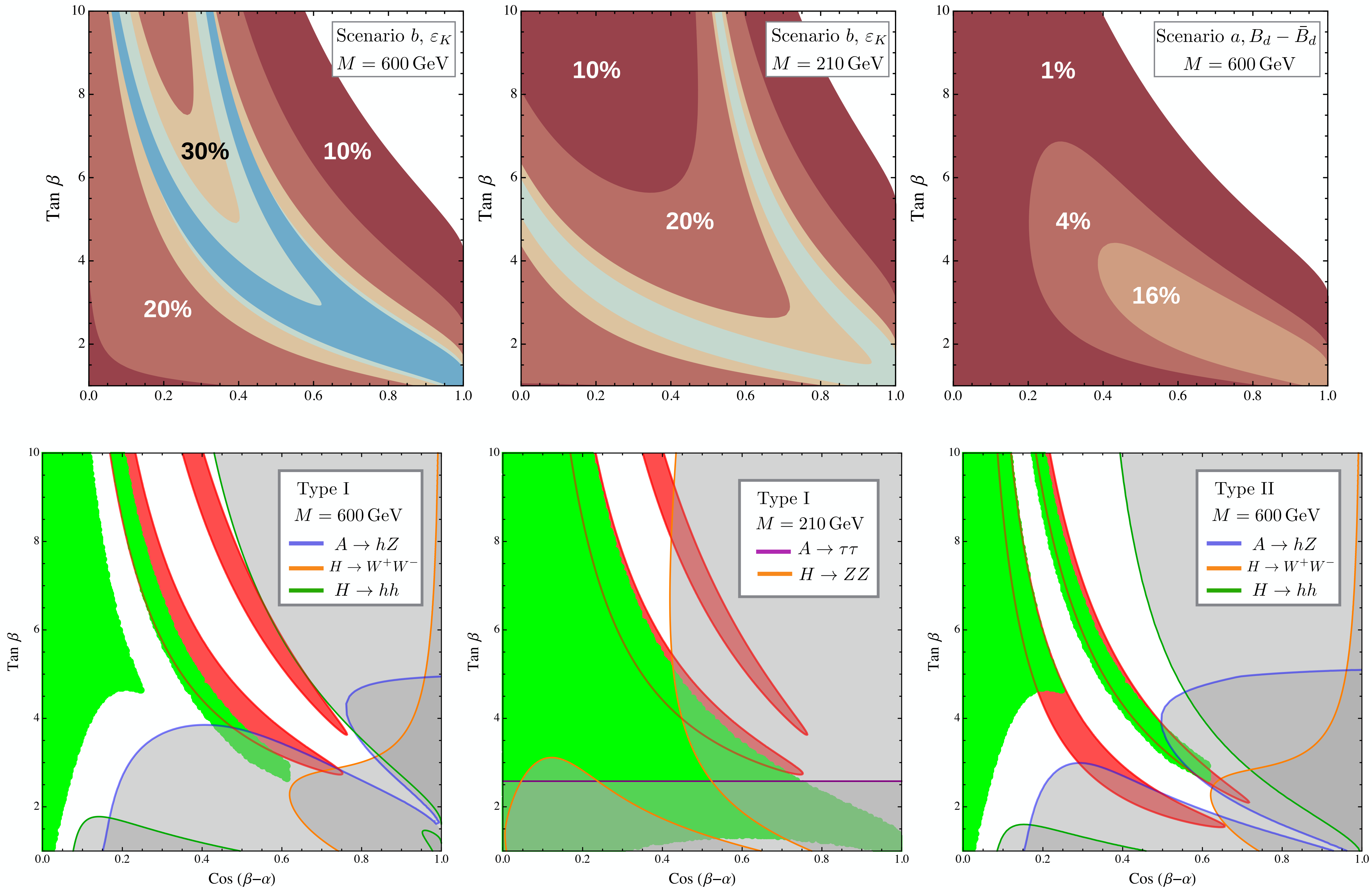
Pseudoscalar into

Heavy scalar into gauge bosons

Global Higgs fit

EWPB

Generalizations: Type I and Type II



Conclusions

- Electroweak scale flavor symmetries will be discovered or excluded by the LHC!

- A generic flavon is very hard to discover. A 100 TeV collider would be the first machine in history with a realistic shot.

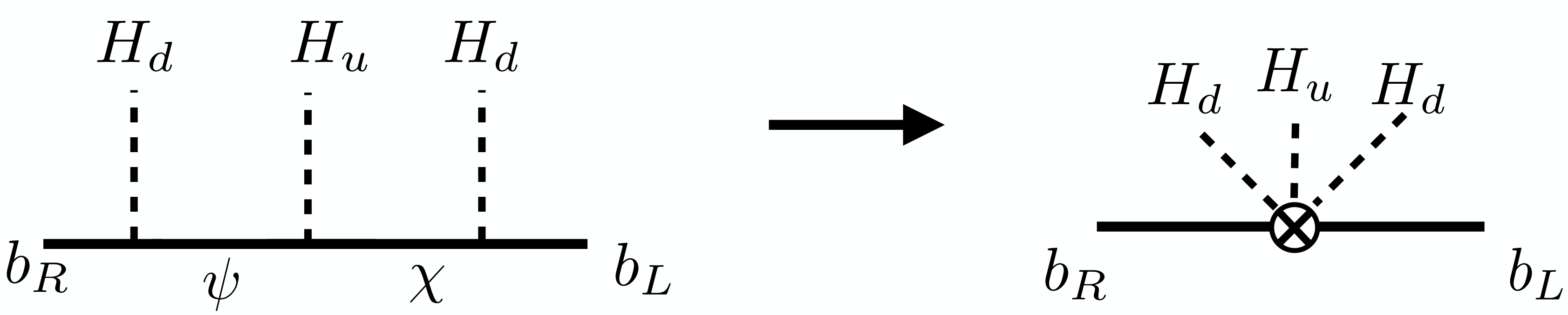
- The upcoming golden age of Lepton flavor will test the flavor structure in the lepton sector and improve on the bounds in the quark sector

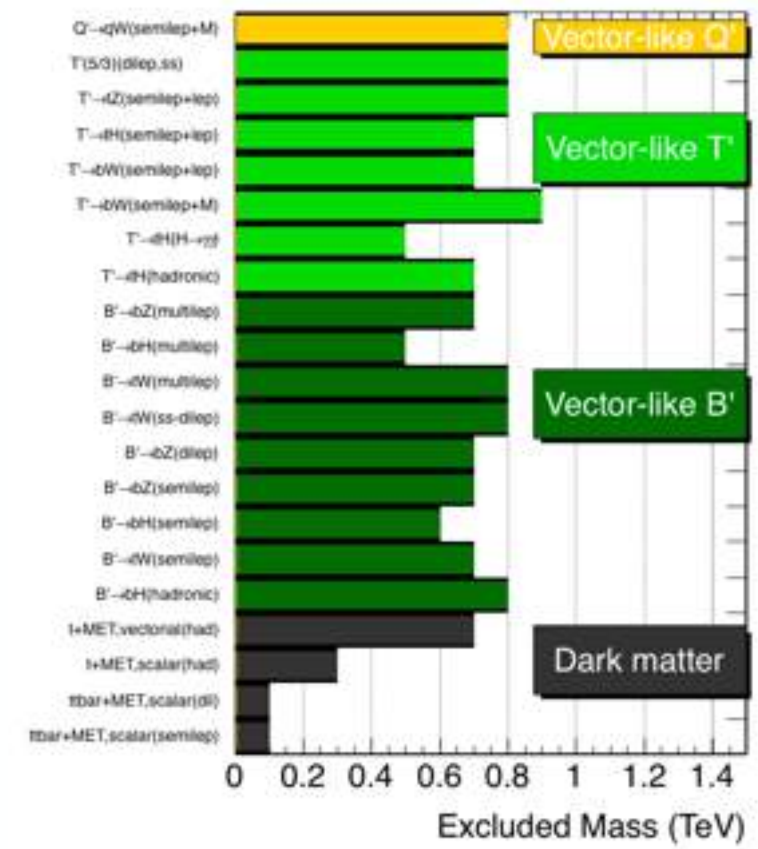
Thank you!

Backup

Collider Searches for Vector Fermions

Heavy Vector Quarks have to be at the TeV scale!





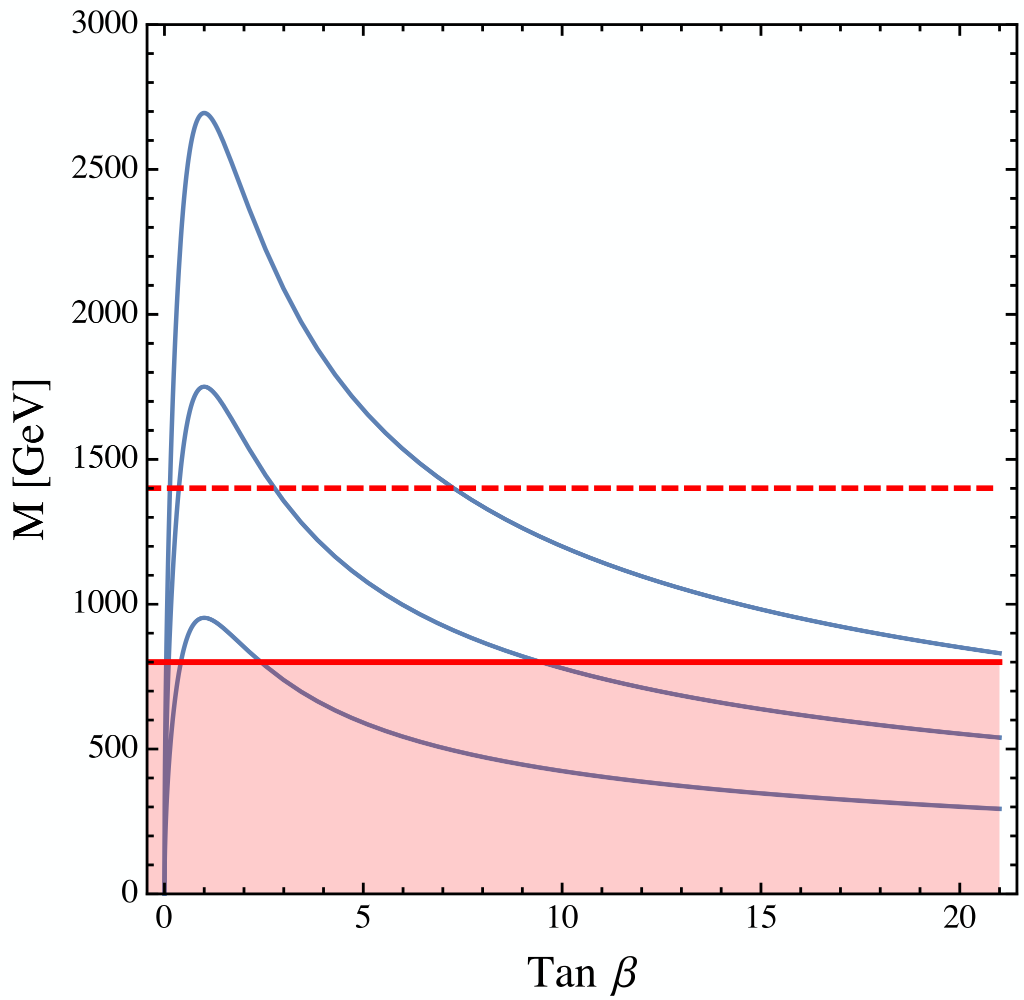
CMS summary 2014]

Heavy quarks	Vector-like quark $TT \rightarrow Ht + X$	1 e, μ	≥ 2 b, ≥ 4 j	Yes	14.3	T mass 790 GeV	T in (T,B) doublet	ATLAS-CONF-2013-018
	Vector-like quark $TT \rightarrow Wb + X$	1 e, μ	≥ 1 b, ≥ 3 j	Yes	14.3	T mass 670 GeV	isospin singlet	ATLAS-CONF-2013-060
	Vector-like quark $TT \rightarrow Zt + X$	2/$\geq$3 e, μ	$\geq$2/$\geq$1 b	–	20.3	T mass 735 GeV	T in (T,B) doublet	ATLAS-CONF-2014-036
	Vector-like quark $BB \rightarrow Zb + X$	2/$\geq$3 e, μ	$\geq$2/$\geq$1 b	–	20.3	B mass 755 GeV	B in (B,Y) doublet	ATLAS-CONF-2014-036
	Vector-like quark $BB \rightarrow Wt + X$	2 e, μ (SS)	≥ 1 b, ≥ 1 j	Yes	14.3	B mass 720 GeV	B in (T,B) doublet	ATLAS-CONF-2013-051

[ATLAS summary July 2014]

A UV Completion: Vectorlike Fermions

Heavy Vectorlike Quarks have to be at the TeV scale!



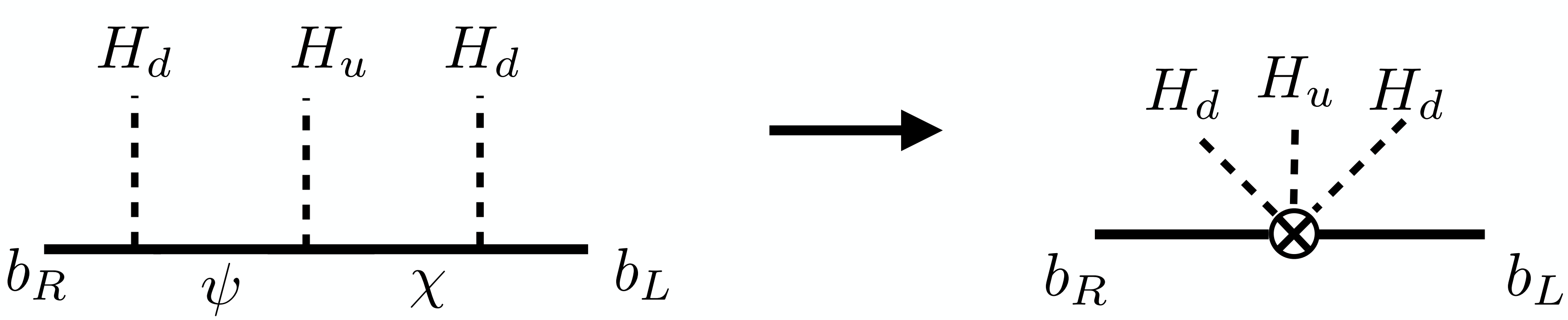
LHC14 projection

[Matsedonskyi, Panico, Wulzer 1409.0100]

LHC8

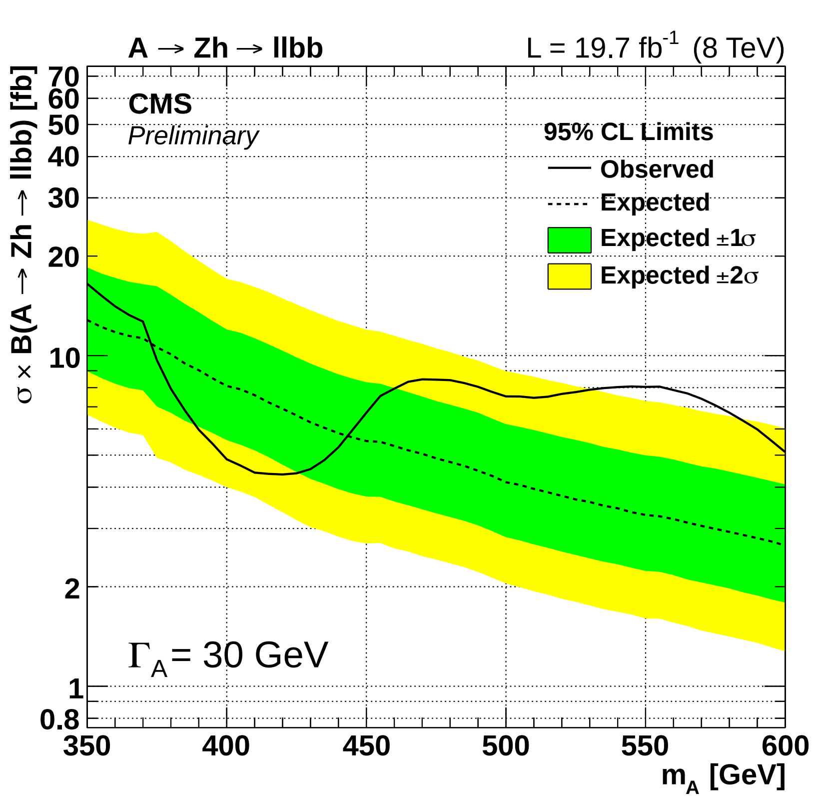
[CMS summary 2014]

[ATLAS summary July 2014]

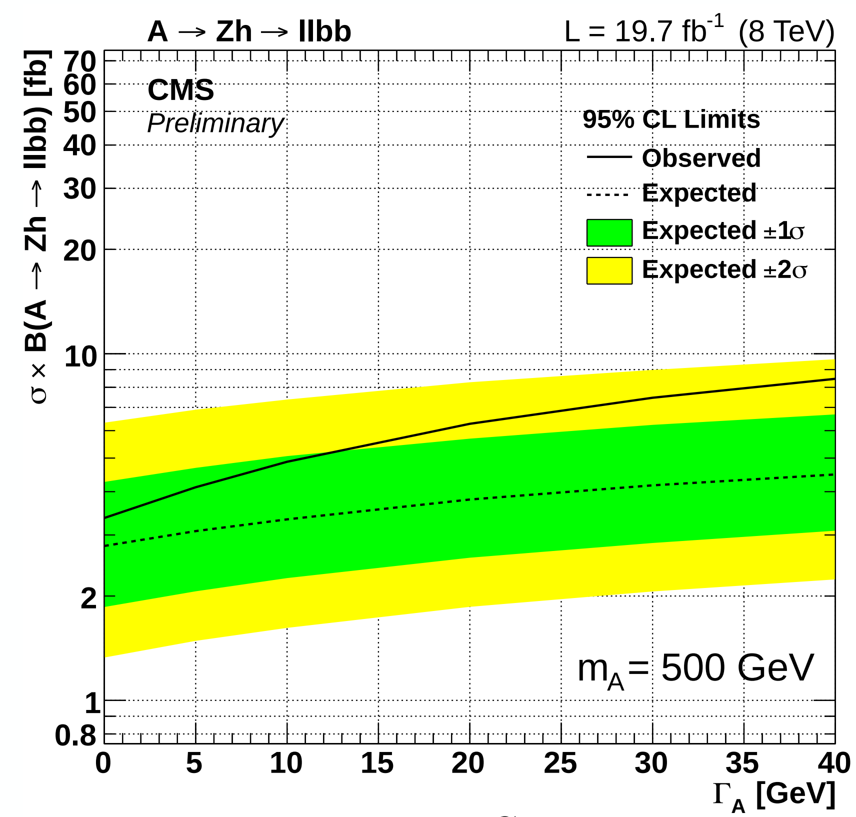


$$Y_b^{\rm eff} \equiv \varepsilon\, y^d = \frac{y_1\, y_2\, y_3}{M_\eta\, M_\psi}\, \frac{v_u v_d}{2}$$

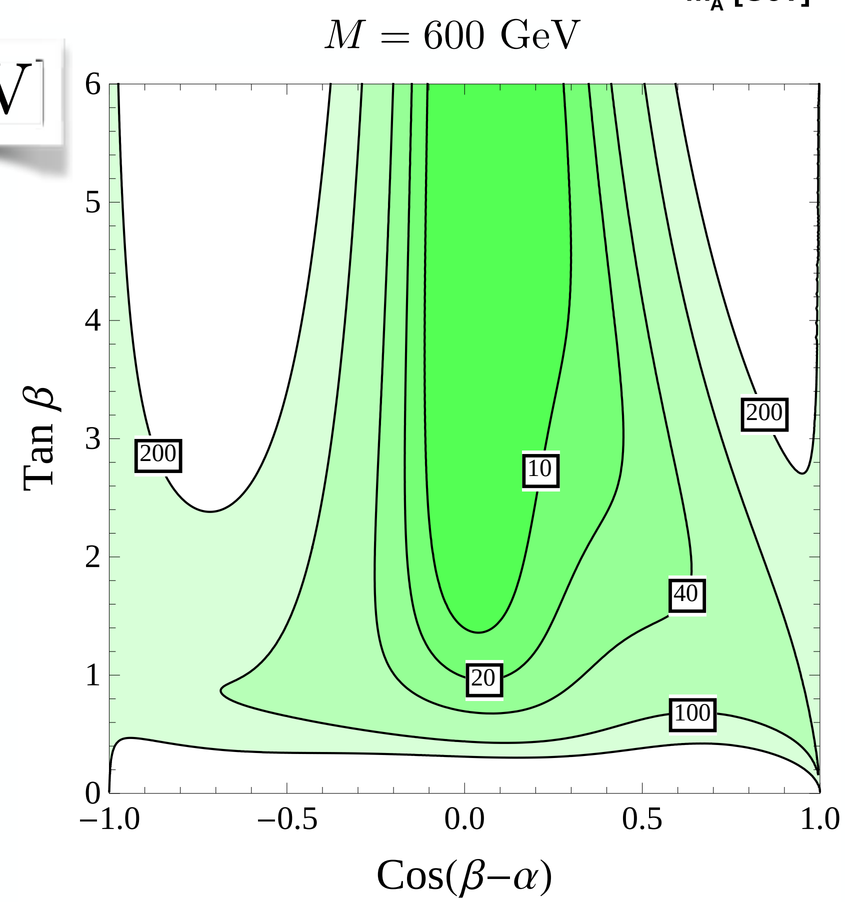
Finite Width Effects



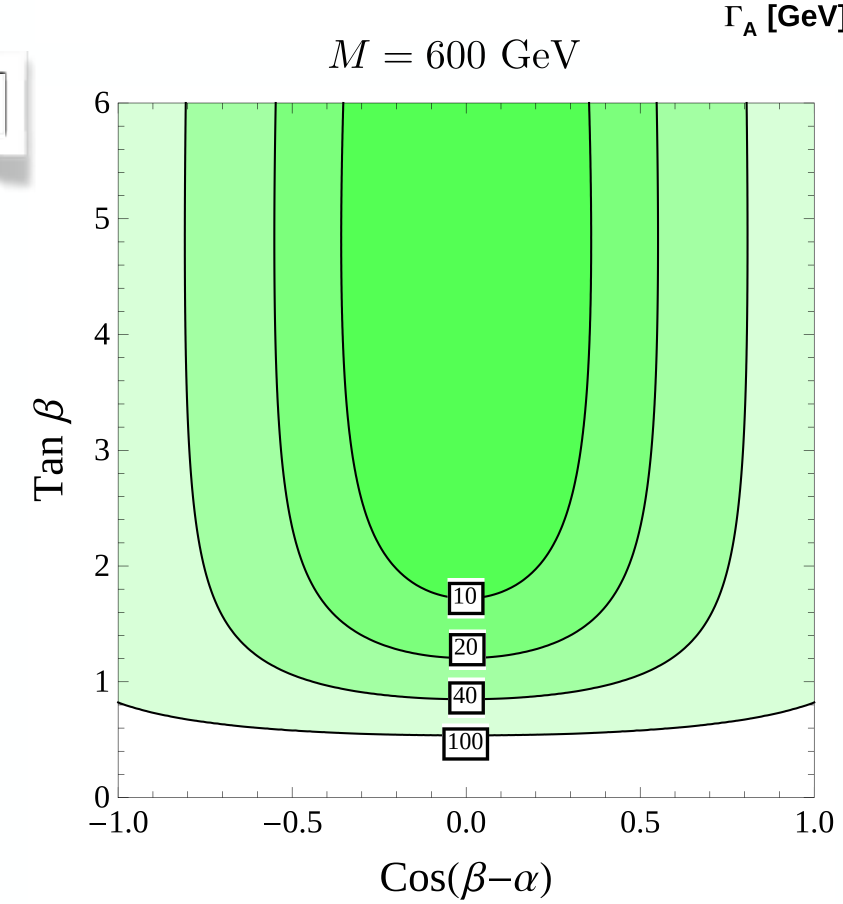
CMS PAS HIG-14-011



Γ_A[GeV]



Γ_H[GeV]



Why should the flavor scale be low?

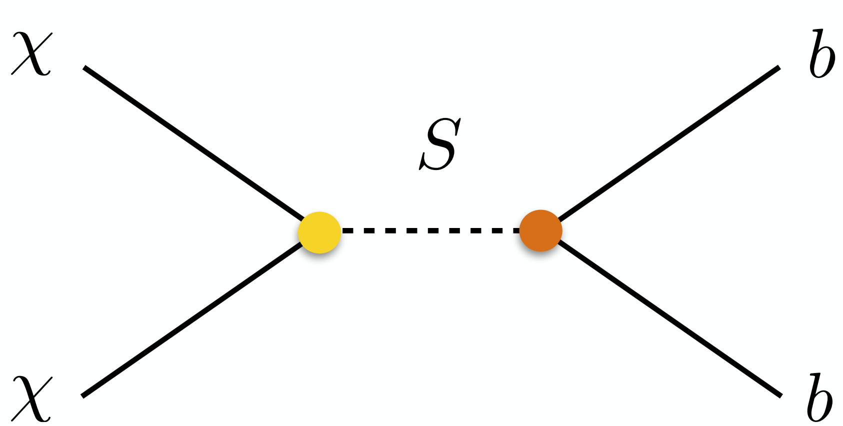
- Why not?
- A link to Baryogenesis?
- A link to Dark Matter?
- It could be related to the electroweak scale

A link to Dark Matter?

Dark matter gets ist mass from the flavon:

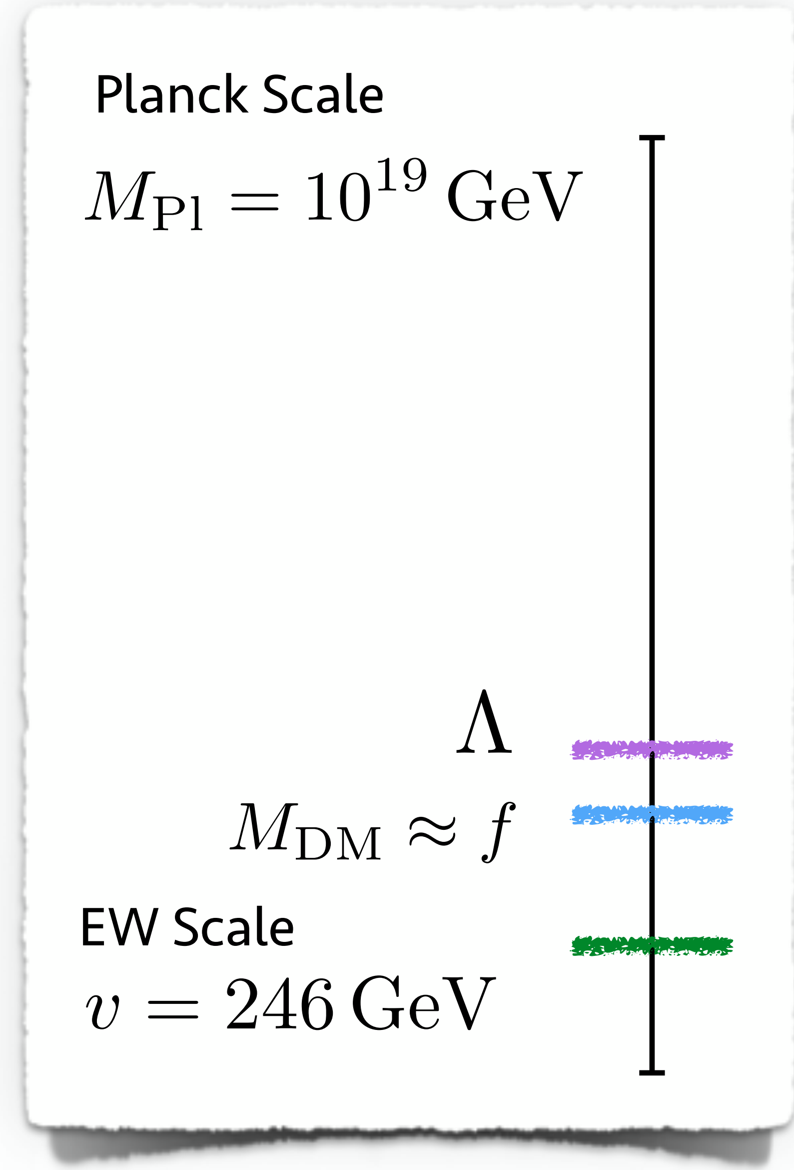
$$y_t \bar{t} H t + y_f \frac{S}{\Lambda} \bar{b} H b + y_\chi S \bar{\chi} \chi$$

$$\Rightarrow M_{\mathrm{DM}} = y_\chi f$$



Relic density

$$\Omega_\chi \propto \frac{M_{\mathrm{DM}}^2}{y_\chi^2 y_b^2} = \frac{f^2}{y_b^2}$$



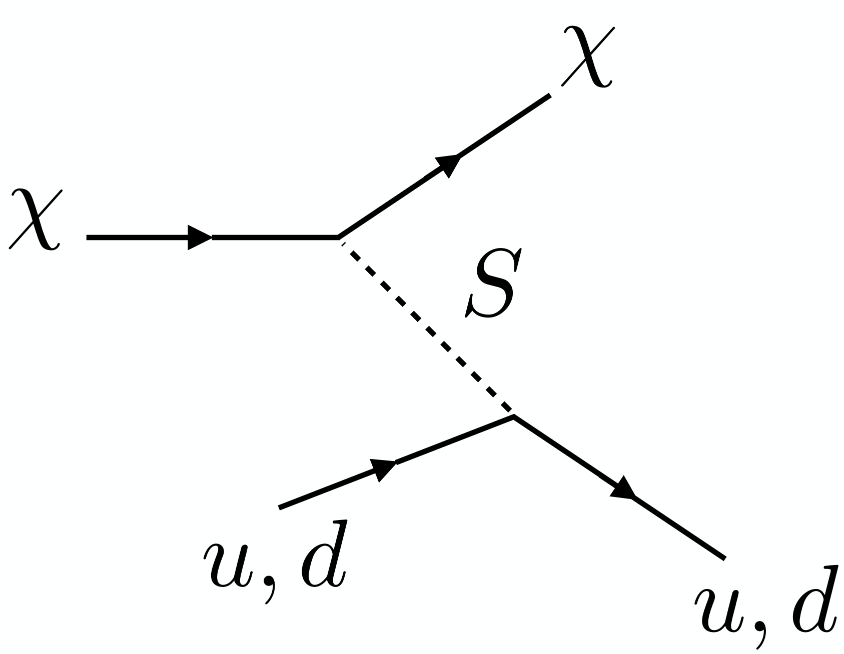
$$\Rightarrow 1\,\mathrm{TeV} \gtrsim f < 10\,\mathrm{TeV}$$

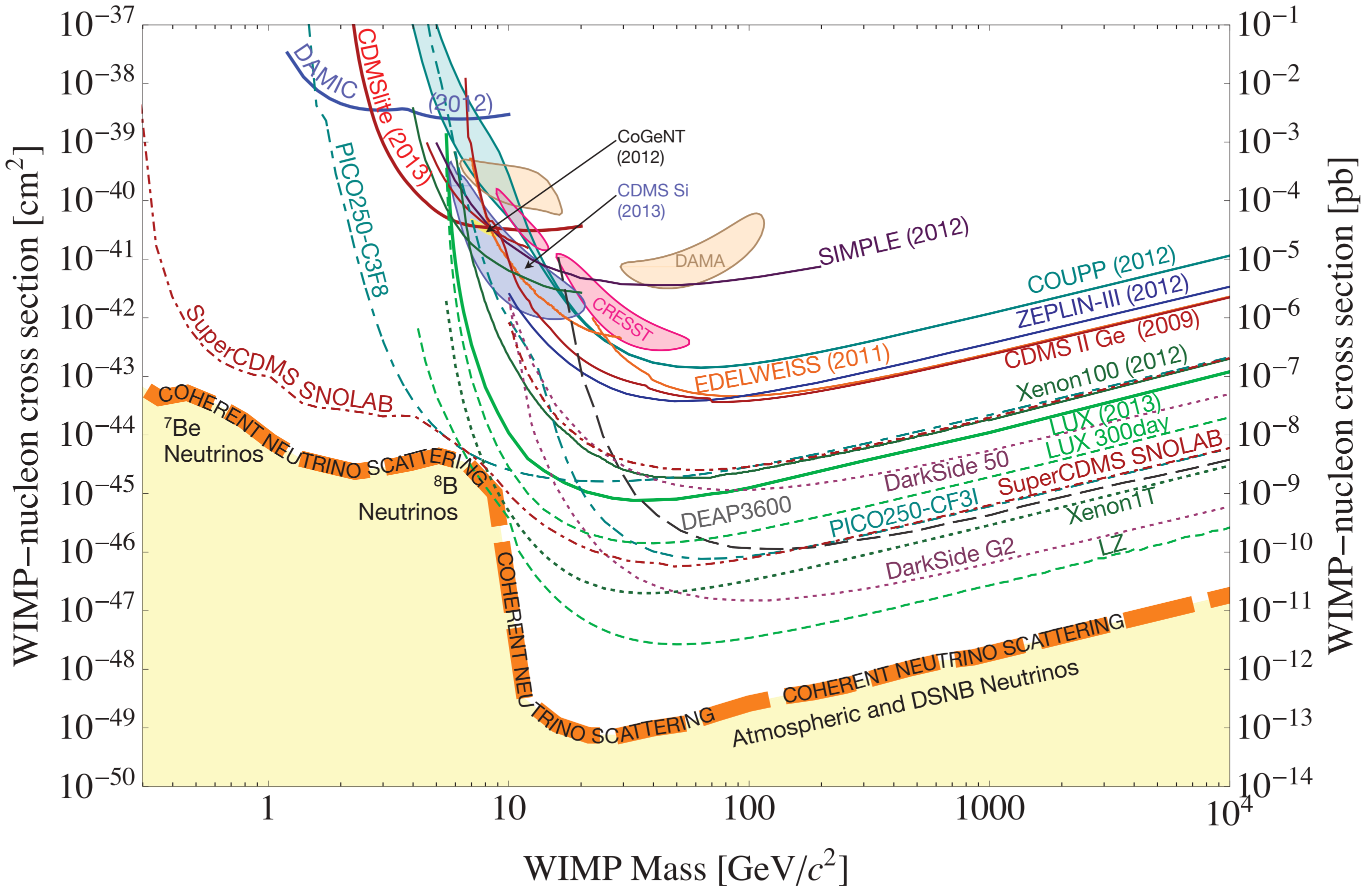
A link to Dark Matter?

Dark matter gets ist mass from the flavon:

$$y_t \bar{t} H t + y_f \frac{S}{\Lambda} \bar{b} H b + y_\chi S \bar{\chi} \chi$$

$$\Rightarrow M_{\mathrm{DM}} = y_\chi f$$





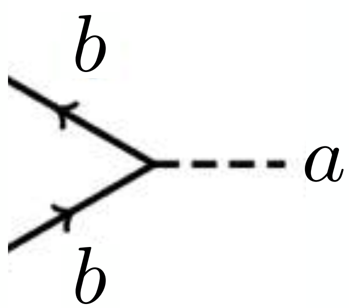
..work with Martin Klaasen

[MB, Klaasen, 160……]

Future Collider Searches

Resonance Searches

$$(\sigma \times \mathrm{BR})^{100}_{\mathrm{limit}} = \sqrt{\frac{\mathcal{L}_8}{\mathcal{L}_{100}}}\sqrt{\frac{\sigma^{100}_{\mathrm{BG}}}{\sigma^{8}_{\mathrm{BG}}}}(\sigma \times \mathrm{BR})^{8}_{\mathrm{limit}}$$

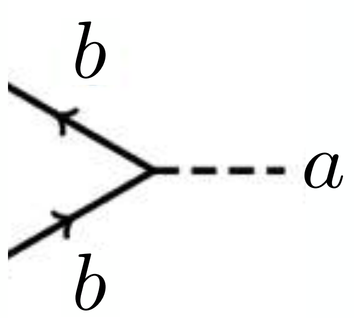


$m_X = 1000$ (500) GeV	ATLAS 8 TeV	CMS 8 TeV	scaled to 100 TeV	flavon
di-jet	0.34	0.18 (M=1300 GeV)	$3. \cdot 10^{-2}$	$9.1 \cdot 10^{-4}$
$\tau\tau$	$5 \cdot 10^{-3}$ (0.03)	$9 \cdot 10^{-3}$ (0.04)	$5. \cdot 10^{-4}$ $(3. \cdot 10^{-3})$	$1.7 \cdot 10^{-4}$ $(2.2 \cdot 10^{-3})$
di-muon	$7 \cdot 10^{-4}$ $(3 \cdot 10^{-3})$	$9 \cdot 10^{-4}$ $(3 \cdot 10^{-3})$	$6 \cdot 10^{-5}$ $(3. \cdot 10^{-4})$	$1.7 \cdot 10^{-6}$ $(2.2 \cdot 10^{-5})$
$\gamma\gamma$	$5 \cdot 10^{-3}$ $(8 \cdot 10^{-3})$	$(2 \cdot 10^{-3})$	$4 \cdot 10^{-4}$ $(7. \cdot 10^{-4})$	$3.5 \cdot 10^{-11}$ $(1.3 \cdot 10^{-9})$

Future Collider Searches

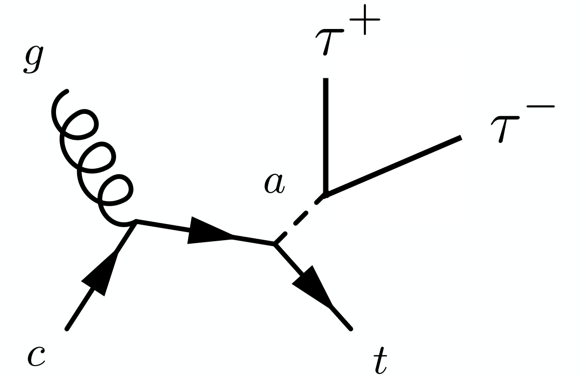
Resonance Searches

$$(\sigma \times \mathrm{BR})^{100}_{\mathrm{limit}} = \sqrt{\frac{\mathcal{L}_8}{\mathcal{L}_{100}}}\sqrt{\frac{\sigma^{100}_{\mathrm{BG}}}{\sigma^{8}_{\mathrm{BG}}}}(\sigma \times \mathrm{BR})^{8}_{\mathrm{limit}}$$



$m_X = 1000$ (500) GeV	ATLAS 8 TeV	CMS 8 TeV	scaled to 100 TeV	flavon
di-jet	0.34	0.18 (M=1300 GeV)	$3. \cdot 10^{-2}$	$9.1 \cdot 10^{-4}$
$\tau\tau$	$5 \cdot 10^{-3}$ (0.03)	$9 \cdot 10^{-3}$ (0.04)	$5. \cdot 10^{-4}$ $(3. \cdot 10^{-3})$	$1.7 \cdot 10^{-4}$ $(2.2 \cdot 10^{-3})$
di-muon	$7 \cdot 10^{-4}$ $(3 \cdot 10^{-3})$	$9 \cdot 10^{-4}$ $(3 \cdot 10^{-3})$	$6 \cdot 10^{-5}$ $(3. \cdot 10^{-4})$	$1.7 \cdot 10^{-6}$ $(2.2 \cdot 10^{-5})$
$\gamma\gamma$	$5 \cdot 10^{-3}$ $(8 \cdot 10^{-3})$	$(2 \cdot 10^{-3})$	$4 \cdot 10^{-4}$ $(7. \cdot 10^{-4})$	$3.5 \cdot 10^{-11}$ $(1.3 \cdot 10^{-9})$

Future Collider Searches



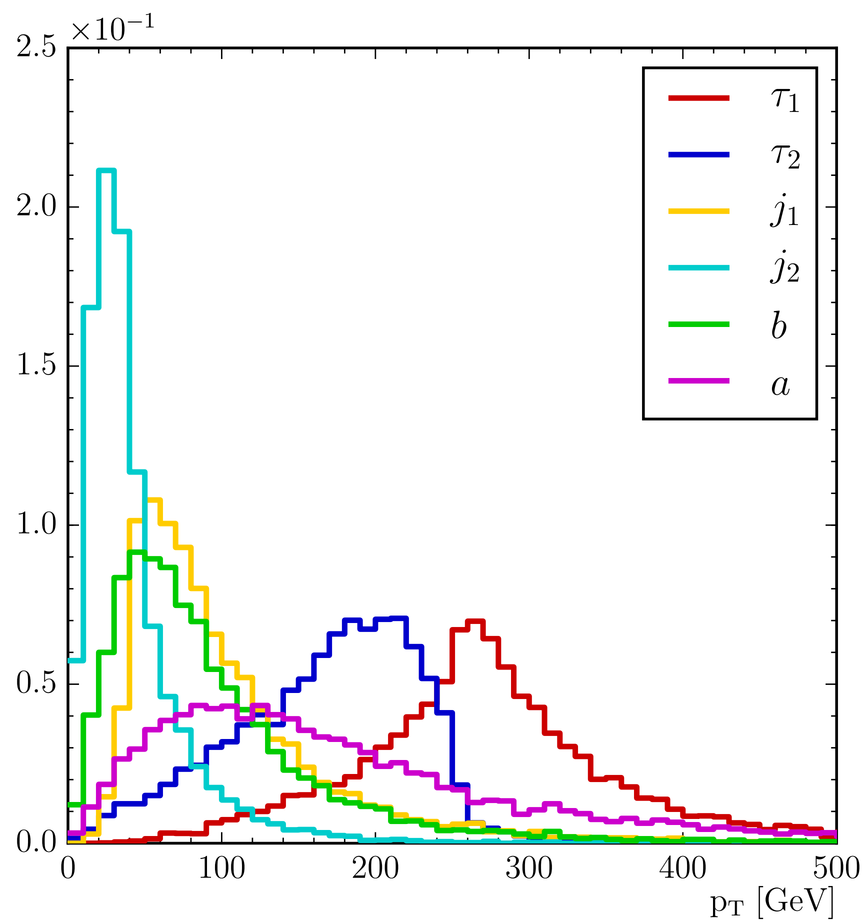
Associated Production

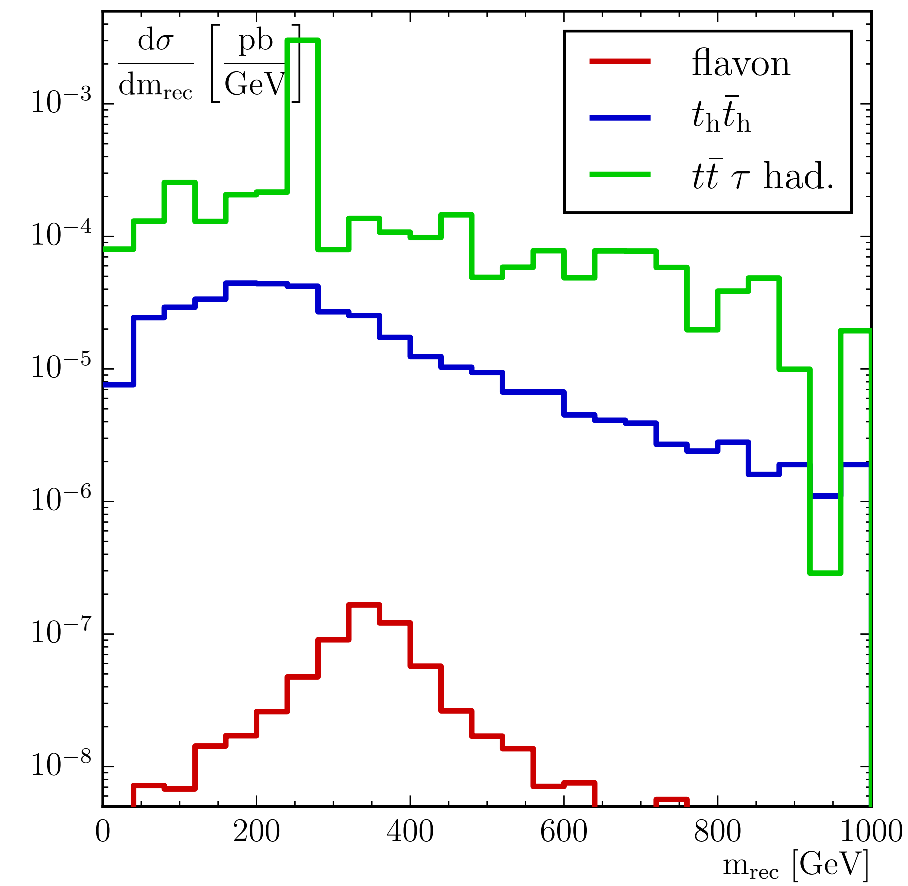
$$n_j \geq 5, \qquad n_\ell = 0, \qquad p_T^{j_1} > 150 \text{ GeV} \qquad m(j_3 j_4 j_5) \in [140, 190] \text{ GeV}$$

Tau tagging $\epsilon_\tau = 0.15$

Tau Misidentification $\epsilon_j = 10^{-3}$

hopeless…





Benchmark 1 : $M_A = M_H = 600$ GeV, $M_{H^+} = 450\,\text{GeV}$,
1a $\cos(\beta-\alpha) = 0.55$, $\tan\beta = 3$,
1b $\cos(\beta-\alpha) = 0.42$, $\tan\beta = 4.5$,

Light Higgs Couplings:

1a $\kappa_t = 1.02$, $\kappa_V = 0.84$, $\kappa_b = \kappa_\tau = -0.61$, $\kappa_c = 1.22$, $\kappa_s = -0.41$
1b $\kappa_t = 1.00$, $\kappa_V = 0.91$, $\kappa_b = \kappa_\tau = -0.96$, $\kappa_c = 1.02$, $\kappa_s = -0.95$

Higgs Signal Strength:

1a	μ_V	μ_γ	μ_b	μ_c
$\sigma_{gg\to h}$	1.38	1.21	0.74	2.95
$\sigma_{t\bar{t}\to h}$	1.33	1.17	0.71	2.84
$\sigma_{VBF}, \sigma_{VH}$	0.89	0.78	0.48	1.91

1b	μ_V	μ_γ	μ_b	μ_c
$\sigma_{gg\to h}$	0.96	0.91	1.09	1.22
$\sigma_{t\bar{t}\to h}$	0.90	0.85	1.02	1.14
$\sigma_{VBF}, \sigma_{VH}$	0.74	0.70	0.84	0.94

Heavy Scalar Production Cross Sections for 1a (1b):

8 TeV: $\sigma(gg \to A) = 78(36)$ fb, $\sigma(gg \to H) = 32(21)$ fb,
$\sigma(pp \to H^- t(b)) = 9(4)$ fb,
14 TeV: $\sigma(gg \to A) = 361(157)$ fb, $\sigma(gg \to H) = 166(97)$ fb,
$\sigma(pp \to H^- t(b)) = 63(25)$ fb,

Heavy Scalar Decay Modes:

A	Γ_i/Γ_A	
	1a	**1b**
Zh	70.2%	62%
W^-H^+	14.4%	21.8%
$b\bar{b}$	1.6%	5.2%
$t\bar{t}$	12.9%	8.7%
$\tau^+\tau^-$	0.2%	0.7%
$t\bar{c}$	0.4%	1.1%

H	Γ_i/Γ_H	
	1a	**1b**
WW	52.9%	43%
ZZ	25.6%	20.9%
hh	9.2%	16.9%
W^-H^+	6.8%	11.2%
$t\bar{t}$	3.9%	3.5%

H^+	Γ_i/Γ_{H^+}	
	1a	**1b**
hW	78.7%	81.5%
$t\bar{b}$	21.2%	18.2%
$\tau\nu$	0.048%	0.33%

Total Width for 1a (1b):

$\Gamma_h = 2.22\,(3.71)$ MeV, $\Gamma_A = 24.6\,(16.3)$ GeV, $\Gamma_H = 36.4\,(26.1)$ GeV,
$\Gamma_{H^+} = 10.2\,(5.8)$ GeV.

Benchmark 2 : $M_A = M_H = 500$ GeV, $M_{H^+} = 360\,\text{GeV}$,
$\cos(\beta-\alpha) = 0.45\,, \quad \tan\beta = 4,$

Light Higgs Couplings:

1b $\kappa_t = 1.01\,, \quad \kappa_V = 0.9\,, \quad \kappa_b = \kappa_\tau = -0.81\,, \quad \kappa_c = 1.1\,, \quad \kappa_s = -0.71$

Higgs Signal Strength:

2	μ_V	μ_γ	μ_b	μ_c
$\sigma_{gg\to h}$	1.15	1.07	0.94	1.76
$\sigma_{t\bar{t}\to h}$	1.09	1.02	0.90	1.67
$\sigma_{VBF}, \sigma_{VH}$	0.86	0.80	0.71	1.32

Heavy Scalar Production Cross Sections:

8 TeV: $\sigma(gg \to A) = 130$ fb, $\quad \sigma(gg \to H) = 53$ fb, $\quad \sigma(pp \to H^- t(b)) = 12$ fb,
14 TeV: $\sigma(gg \to A) = 546$ fb, $\quad \sigma(gg \to H) = 224$ fb, $\quad \sigma(pp \to H^- t(b)) = 66$ fb,

Heavy Scalar Decay Modes:

A	Γ_i/Γ_A
Zh	56.6%
W^-H^+	23.3%
$b\bar{b}$	5.3%
$t\bar{t}$	12.4%
$\tau^+\tau^-$	0.66%
$t\bar{c}$	1.1%

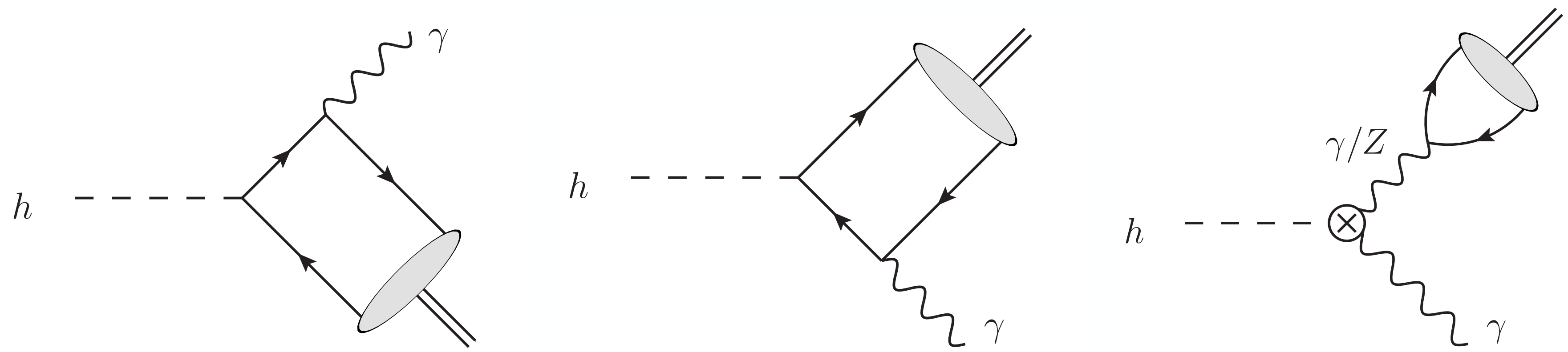
H	Γ_i/Γ_H
WW	45.4%
ZZ	21.8%
hh	11.5%
W^-H^+	12.6%
$t\bar{t}$	3.65%

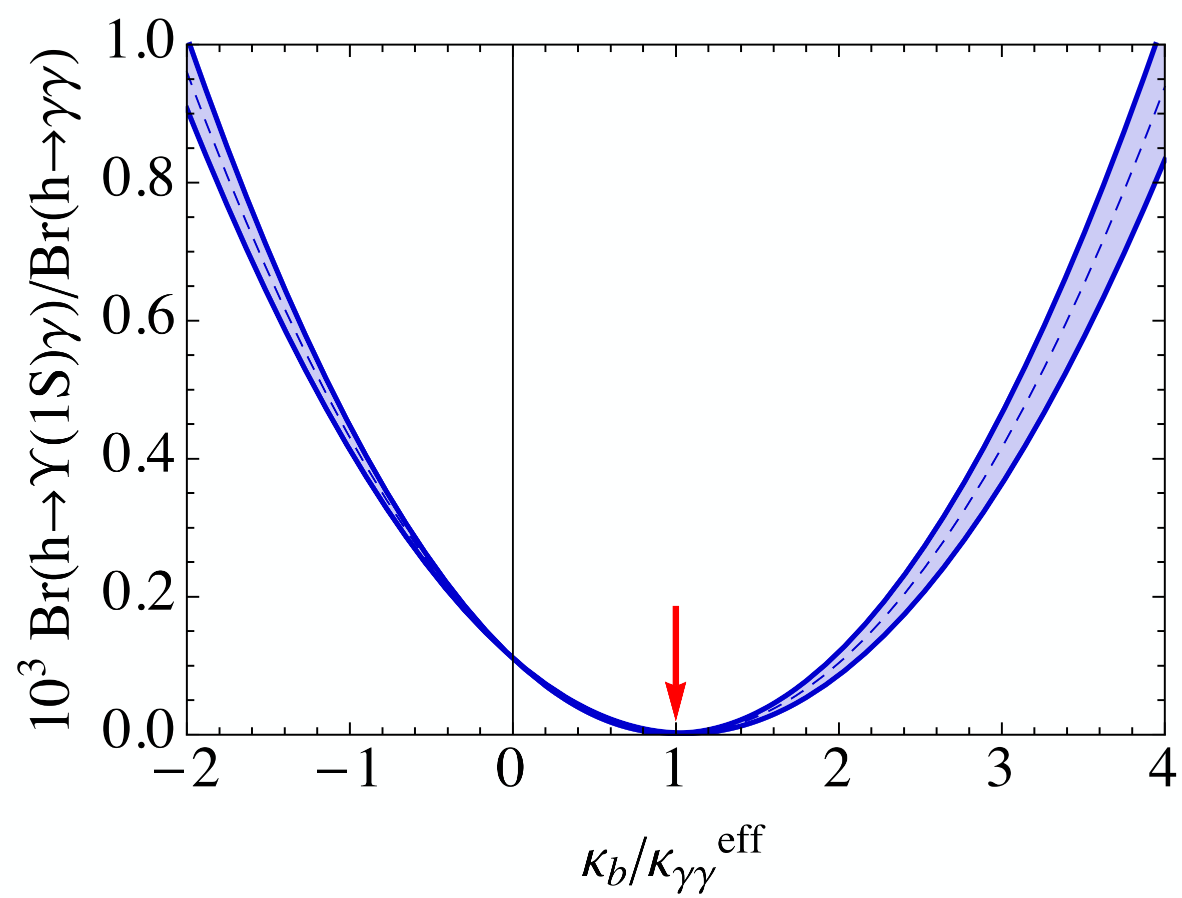
H^+	Γ_i/Γ_{H^+}
hW	71.8%
$t\bar{b}$	27.8%
$\tau\nu$	0.4%

Total Width:

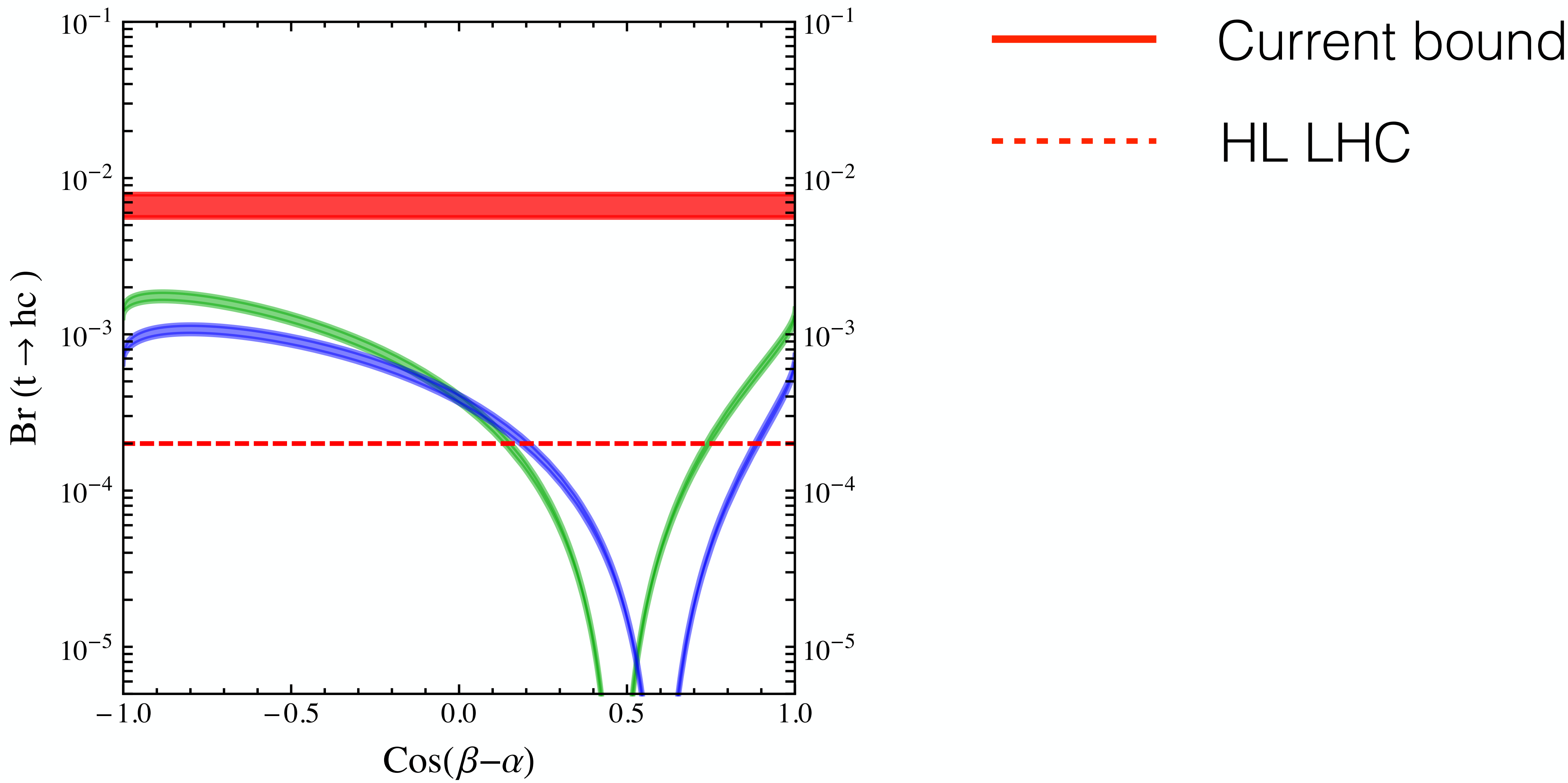
$\Gamma_h = 3$ MeV, $\quad \Gamma_A = 10.7$ GeV, $\quad \Gamma_H = 15.7$ GeV, $\quad \Gamma_{H^+} = 3$ GeV.

Hadronic Radiative Higgs Decays



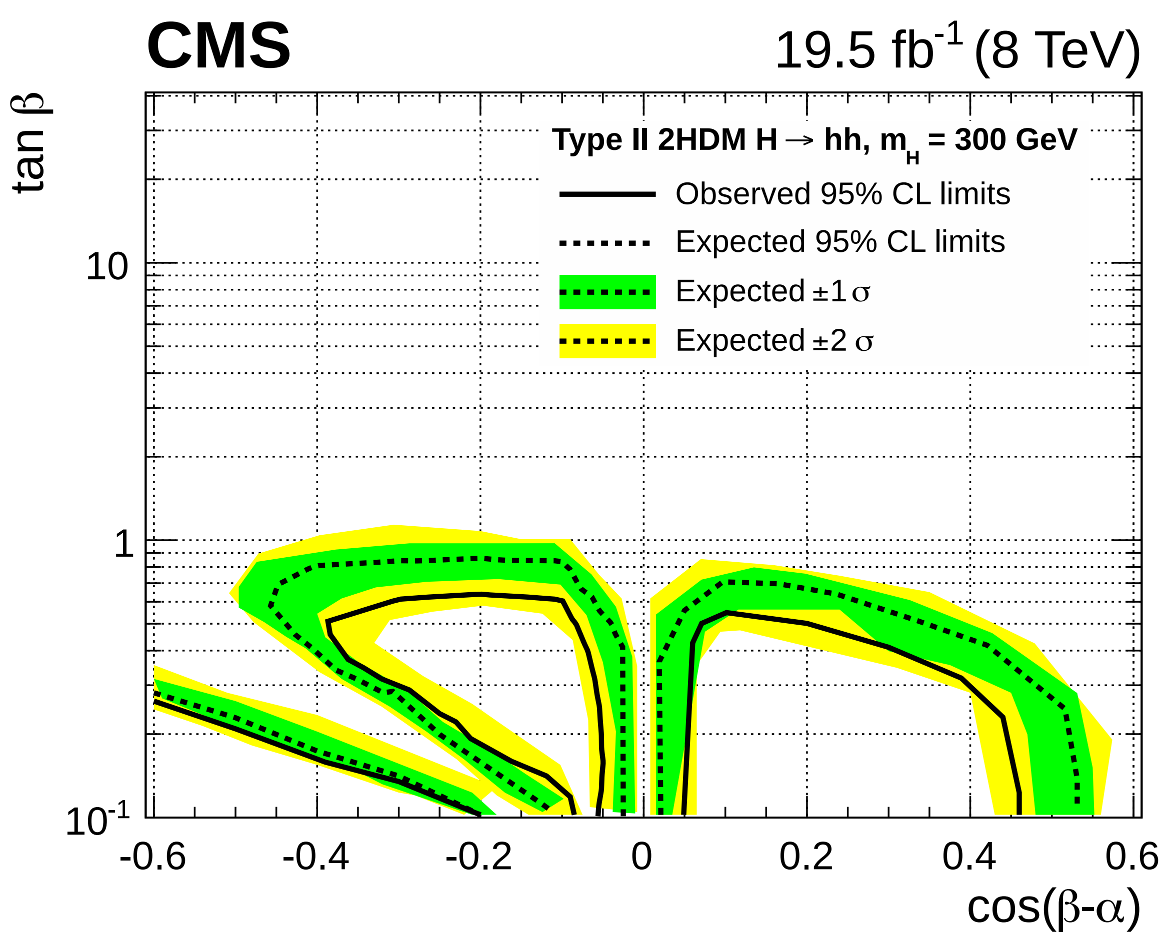


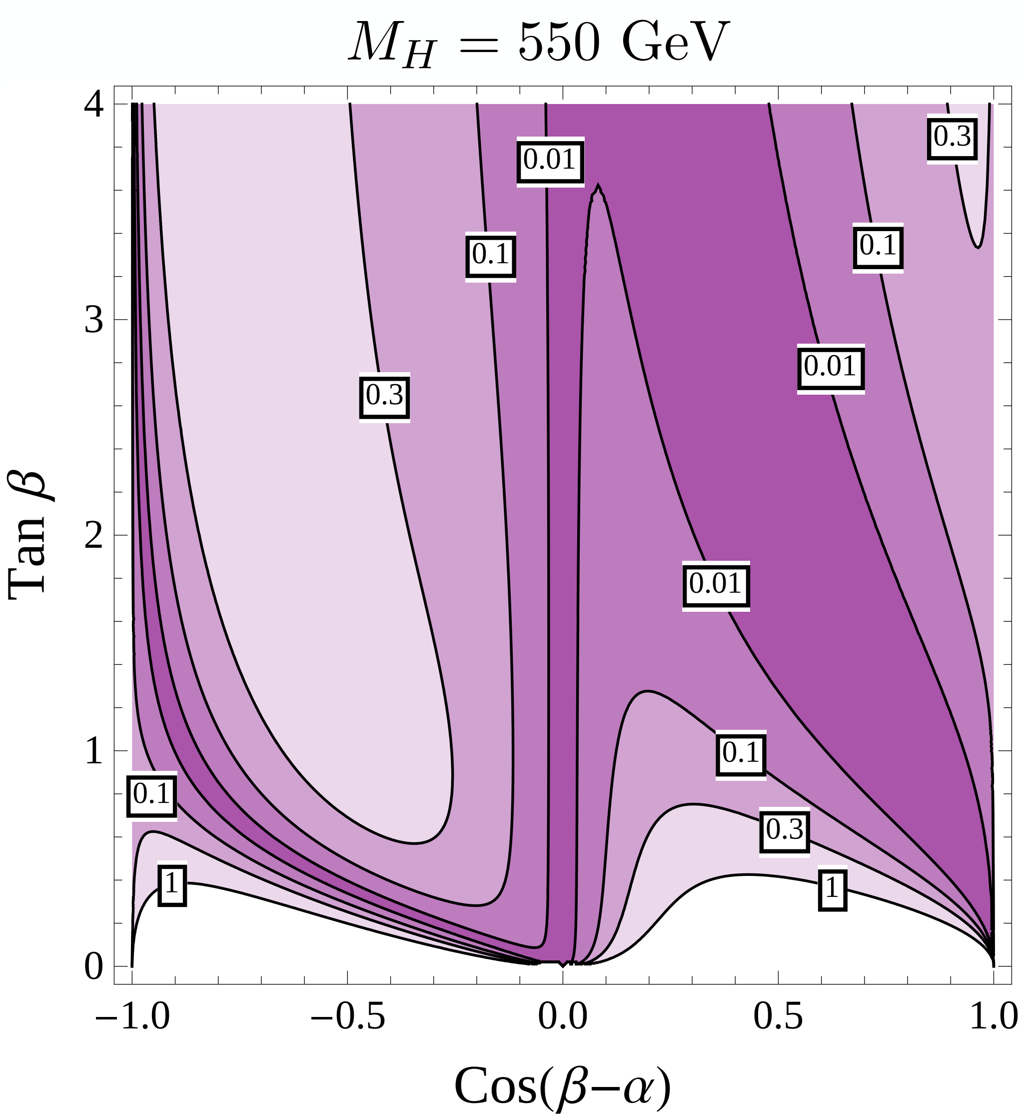
Flavor off-diagonal Top Decays



[ATLAS 1403.6293]   [CMS 1410.2751]

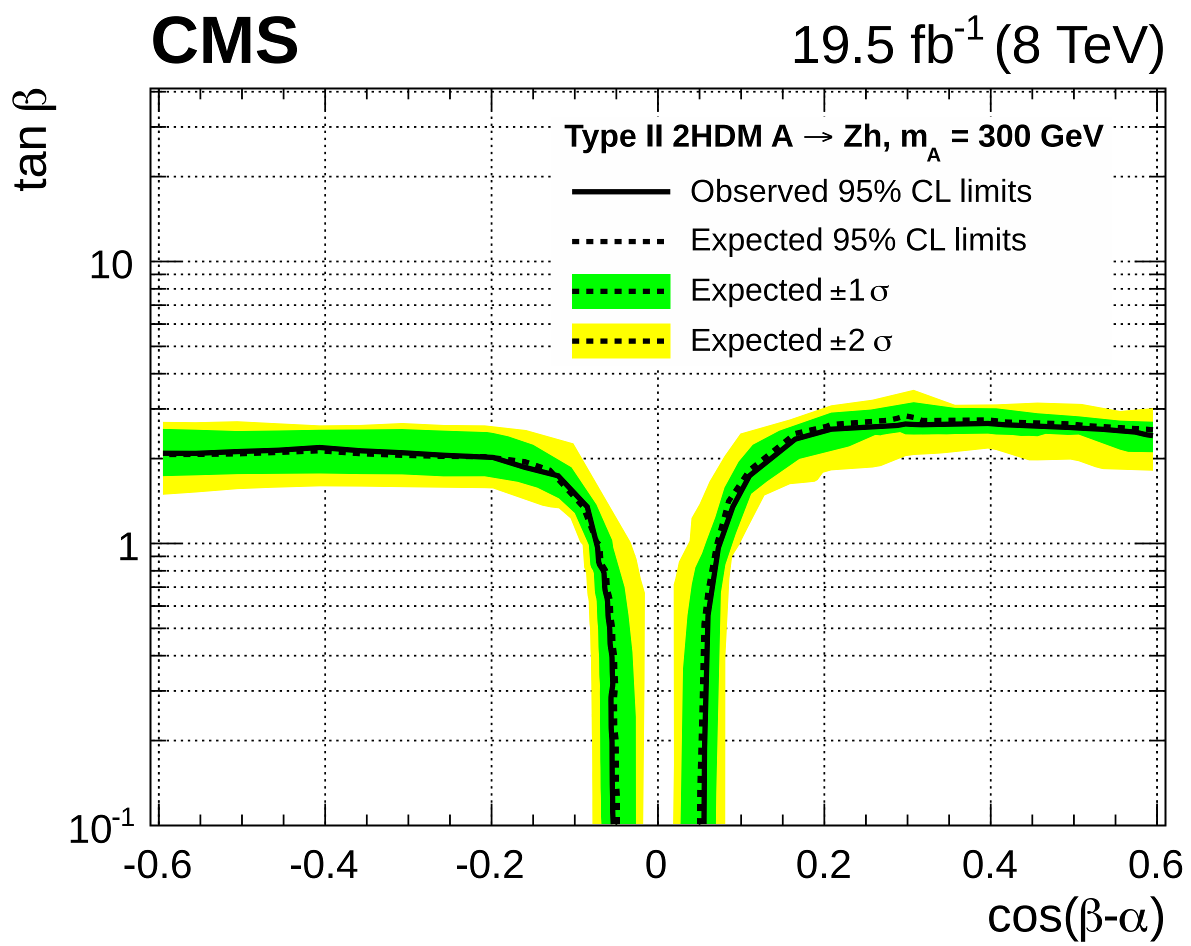
Heavy Higgs Searches

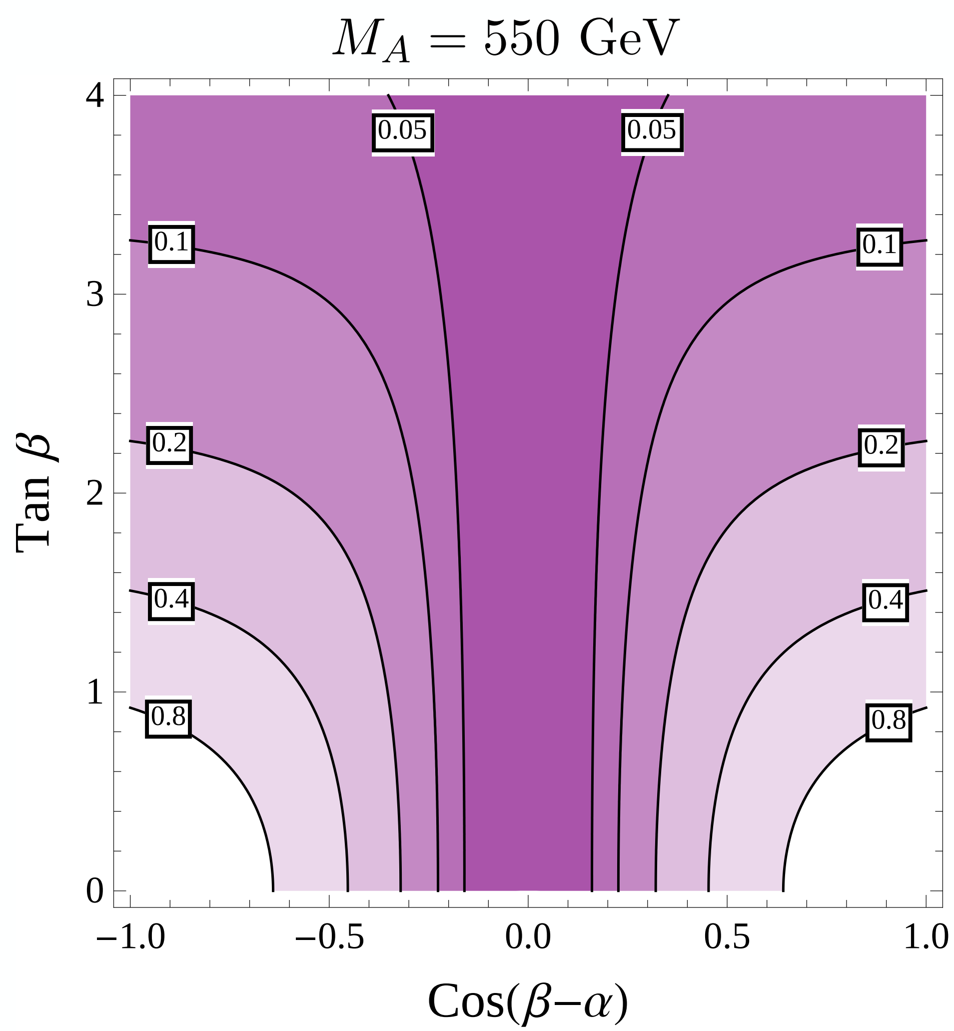




[CMS 1410.2751]

Heavy Higgs Searches





[CMS 1410.2751]

Flavor Structure

$$f^h(\alpha,\beta) = 2\sin(\beta-\alpha) + \cos(\beta-\alpha)\left(\frac{1}{\tan\beta} - \tan\beta\right)$$

$$F^H(\alpha,\beta) = 2\cos(\beta-\alpha) + \sin(\beta-\alpha)\left(\tan\beta - \frac{1}{\tan\beta}\right)$$

$$F^A(\alpha,\beta) = \left(\tan\beta + \frac{1}{\tan\beta}\right)$$

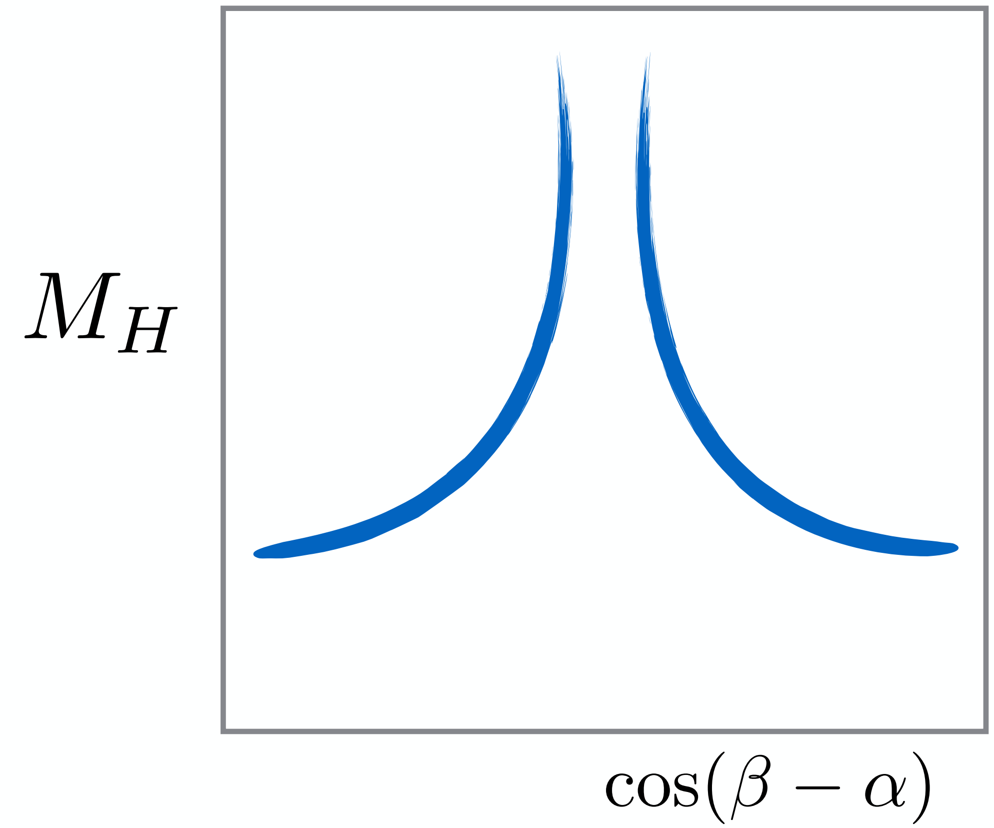
Collider Searches for Heavy Higgses

$$H_u = \frac{1}{\sqrt{2}} \begin{pmatrix} v_u + h_u + i a_u \\ \sqrt{2} H_u^- \end{pmatrix}$$

$$H_d = \frac{1}{\sqrt{2}} \begin{pmatrix} \sqrt{2} H_d^+ \\ v_d + h_d + i a_d \end{pmatrix}$$

$$8 \,\text{real fields} \quad - \quad W^\pm, Z^0$$

$$= h, H, A, H^\pm$$



The heavy Higgses cannot decouple!