

A Leptophilic Model Explaining Dark Matter and Neutrino Masses at Low Energies

Michael A. Schmidt

Institute for Particle Physics Phenomenology
Durham

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based on:
Y. Farzan, S. Pascoli, MS [1005.5323] and [1006.xxxx]

Outline

1 Model

2 Lepton Sector

3 Dark Matter

- Dark Matter Annihilation
- Dark Matter Direct Detection
- DAMA/CoGeNT/CDMS-II?
- Indirect Detection

4 More Phenomenology

5 Comments on Alternatives

6 Conclusions

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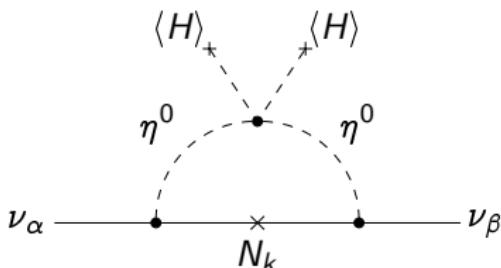
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Radiative Mass Generation

Ma Model_[Ma (2006)]



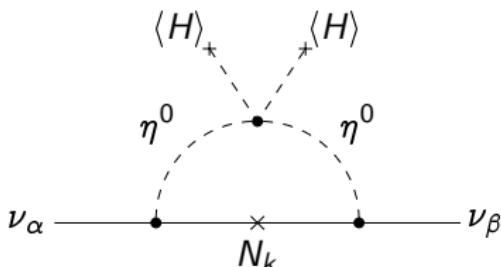
$$(m_\nu)_{\alpha\beta} = \sum_k \frac{Y_{\alpha k} Y_{\beta k} M_k}{16\pi^2} \left[\frac{m_R^2}{m_R^2 - M_k^2} \ln \frac{m_R^2}{M_k^2} - \frac{m_I^2}{m_I^2 - M_k^2} \ln \frac{m_I^2}{M_k^2} \right]$$

Conditions for Majorana Neutrino Mass Term

- Particle in loop have to couple to $SU(2)_L$ doublet

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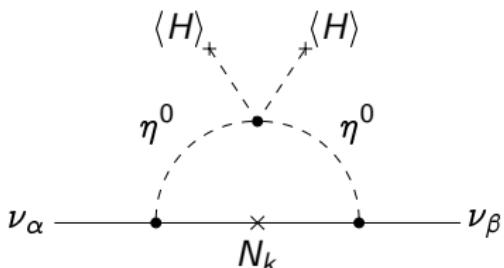
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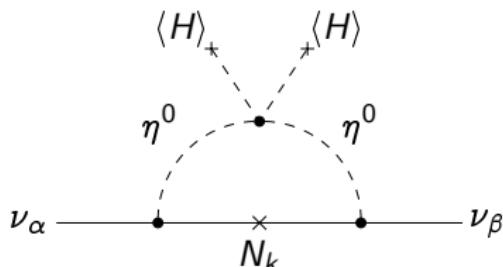
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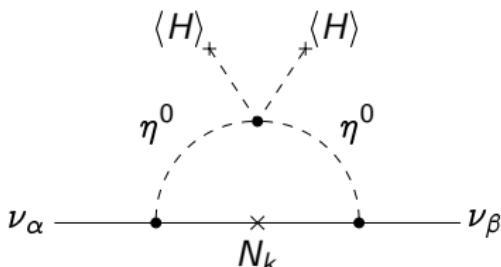
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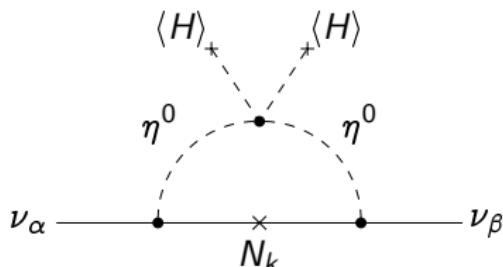
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 $\Rightarrow \eta^0$ stable particle \Rightarrow DM candidate

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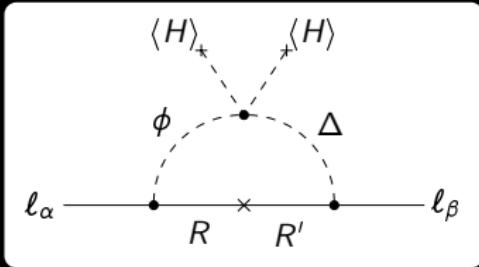
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- Discrete symmetry to avoid FCNCs and Dirac mass term
 $\Rightarrow \eta^0$ stable particle \Rightarrow DM candidate
- But no symmetry explanation for smallness of couplings

Model

Particle Content

	SU(2)	U(1)	U(1) $_X$	\mathbb{Z}_2
$\ell_L^{(i)}$	2	-1/2	0	+
R_R	2	-1/2	1	-
R'_R	2	1/2	-1	-
Δ	3	1	1	-
ϕ	1	0	-1	-



- Symmetry explanation for smallness of couplings $U(1) \rightarrow \mathbb{Z}_2$
- \Rightarrow here explicit, later spontaneous
- Symmetry protects smallness from large quantum corrections

Fermion Sector

$$- m_{RR} (R'^C)^\dagger \cdot R - g_\alpha \phi^\dagger R^\dagger \ell_{L\alpha} - \tilde{g}_\alpha \phi R^\dagger \ell_{L\alpha} - (\tilde{g}_\Delta)_\alpha R'^\dagger \cdot \Delta \cdot \ell_{L\alpha} + \text{h. c.}$$

▶ skip technical details

Particle Content and Symmetries

	SU(3) _c	SU(2) _L	U(1) _Y	U(1) _X [\mathbb{Z}_2]	U(1) _{L1}	U(1) _{L2}	U(1) _{L3}
$Q_L^{(i)}$	3	2	1/6				
$u_R^{(i)}$	3	1	2/3				
$d_R^{(i)}$	3	1	-1/3				
$\ell_L^{(i)}$	1	2	-1/2	0[+]	+1	-1	+1
$e_R^{(i)}$	1	1	-1	0[+]	-1	+1	-1
H	1	2	1/2				
R_R	1	2	-1/2	1[-]	+1	+1	+1
R'_R	1	2	1/2	-1[-]	-1	-1	-1
Δ	1	3	1	1[-]	0	0	-2
ϕ	1	1	0	-1[-]	0	0	0

Yukawa Couplings

Particle Content

	$U(1)_X$	\mathbb{Z}_2	$U(1)_{L1}$	$U(1)_{L2}$	$U(1)_{L3}$
$\ell_L^{(i)}$	0	+	+1	-1	+1
R_R	1	-	+1	+1	+1
R'_R	-1	-	-1	-1	-1
Δ	1	-	0	0	-2
ϕ	-1	-	0	0	0

- Dirac mass $m_{RR} \gtrsim 100 \text{ GeV}$

$$- m_{RR} (R'^C)^\dagger \cdot R - g_\alpha \phi^\dagger R^\dagger \ell_{L\alpha} - \tilde{g}_\alpha \phi R^\dagger \ell_{L\alpha} - (\tilde{g}_\Delta)_\alpha R'^\dagger \cdot \Delta \cdot \ell_{L\alpha} + \text{h. c.}$$

- $\cancel{\mu_2}$
- $\cancel{\mu_2}$, $U(1)_X$ breaking
- $\cancel{\mu_1}$, $U(1)_X$ breaking

Higgs Potential

Particle Content

	$U(1)_X$	\mathbb{Z}_2	$U(1)_{L1}$	$U(1)_{L2}$	$U(1)_{L3}$
Δ	1	-	0	0	-2
ϕ	-1	-	0	0	0

- DM coannihilation into SM Higgs boson h
- Direct mass terms

$$\mathcal{V} \supset m_\Delta^2 \text{Tr } \Delta^\dagger \Delta + m_\phi^2 \phi^\dagger \phi + \lambda_L \frac{v_H}{\sqrt{2}} h (\delta_1^2 + \delta_2^2)$$

$$+ \frac{2m_{\phi\Delta}^2}{v_H^2} H^T i\sigma_2 \Delta^\dagger H \phi^\dagger + \frac{2\tilde{m}_{\phi\Delta}^2}{v_H^2} H^T i\sigma_2 \Delta^\dagger H \phi + \tilde{m}_\phi^2 \phi^2 + \text{h. c.}$$

- $\cancel{U_3}$, $U(1)_\phi \times U(1)_\Delta \rightarrow U(1)_X$
- $\cancel{U_3}$, $U(1)_X \rightarrow \mathbb{Z}_2$, mixing ϕ and Δ
- $U(1)_X \rightarrow \mathbb{Z}_2$, mass splitting of $\text{Re}(\phi)$ and $\text{Im}(\phi)$

Neutral Scalar Masses

$$\phi = (\phi_1 + i \phi_2)/\sqrt{2}$$

$$\Delta^0 = (\Delta_1 + i \Delta_2)/\sqrt{2}$$

$$\begin{pmatrix} \delta_1 \\ \delta_2 \\ \delta_3 \\ \delta_4 \end{pmatrix} = \begin{pmatrix} \cos \alpha_1 & 0 & \sin \alpha_1 & 0 \\ 0 & \cos \alpha_2 & 0 & \sin \alpha_2 \\ -\sin \alpha_1 & 0 & \cos \alpha_1 & 0 \\ 0 & -\sin \alpha_2 & 0 & \cos \alpha_2 \end{pmatrix} \begin{pmatrix} \phi_1 \\ \phi_2 \\ \Delta_1 \\ \Delta_2 \end{pmatrix}$$

Neutral Scalar Masses

$$\phi = (\phi_1 + i \phi_2) / \sqrt{2}$$

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$$M_1^2 \simeq m_\phi^2 - \frac{m_{\phi\Delta}^4}{m_\Delta^2 - m_\phi^2} - \tilde{m}_\phi^2 - 2 \frac{m_{\phi\Delta}^2}{m_\Delta^2 - m_\phi^2} \tilde{m}_{\phi\Delta}^2$$

$$M_2^2 \simeq m_\phi^2 - \frac{m_{\phi\Delta}^4}{m_\Delta^2 - m_\phi^2} + \tilde{m}_\phi^2 + 2 \frac{m_{\phi\Delta}^2}{m_\Delta^2 - m_\phi^2} \tilde{m}_{\phi\Delta}^2$$

$$M_3^2 \simeq m_\Delta^2 + 2 \frac{m_{\phi\Delta}^2}{m_\Delta^2 - m_\phi^2} \tilde{m}_{\phi\Delta}^2$$

$$M_4^2 \simeq m_\Delta^2 - 2 \frac{m_{\phi\Delta}^2}{m_\Delta^2 - m_\phi^2} \tilde{m}_{\phi\Delta}^2$$

mixing angles: $|\tan 2\alpha_1| \simeq |\tan 2\alpha_2| \simeq 2m_{\phi\Delta}^2 / (m_\Delta^2 - m_\phi^2)$

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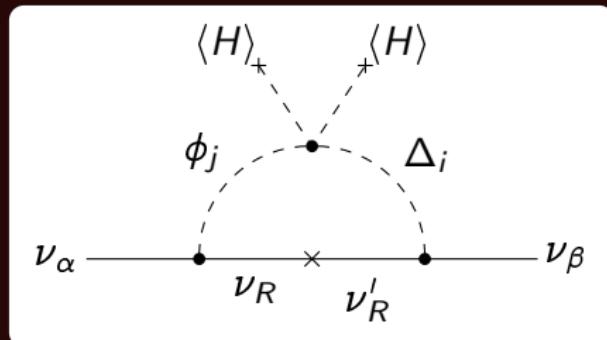
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Neutrino Masses

One Loop Diagram Generating Neutrino Masses



- neutral scalar mass eigenstates δ_i
- with scalar masses M_i
- α_1 mixing between $\delta_{1,3}$
- α_2 mixing between $\delta_{2,4}$

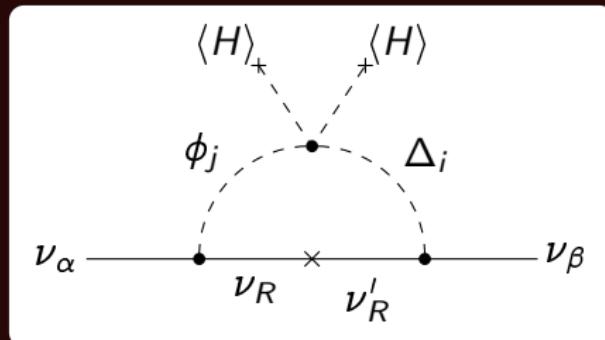
$$(m_\nu)_{\alpha\beta} = [g_\alpha(\tilde{g}_\Delta)_\beta + g_\beta(\tilde{g}_\Delta)_\alpha]\tilde{\eta} + [\tilde{g}_\alpha(\tilde{g}_\Delta)_\beta + \tilde{g}_\beta(\tilde{g}_\Delta)_\alpha]\eta$$

$$\tilde{\eta} = \frac{m_{RR}}{64\pi^2} \left(\frac{M_3^2}{m_{RR}^2 - M_3^2} \ln \frac{m_{RR}^2}{M_3^2} - \frac{M_1^2}{m_{RR}^2 - M_1^2} \ln \frac{m_{RR}^2}{M_1^2} \right) \sin 2\alpha_1 + [(1, 3) \rightarrow (2, 4)]$$

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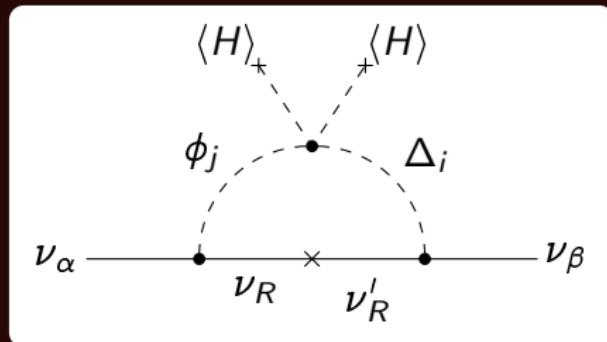
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$$\begin{aligned}\tilde{\eta} &\simeq \frac{m_{RR}}{16\pi^2} \left(\frac{\tilde{m}_\phi^2 m_{\phi\Delta}^2}{m_{RR}^2 m_\Delta^2} \left(\frac{m_{RR}^2}{m_{RR}^2 - m_\Delta^2} \ln \frac{m_{RR}^2}{m_\Delta^2} + 1 - \ln \frac{m_{RR}^2}{M_1^2} \right) - \frac{\tilde{m}_{\phi\Delta}^2}{m_{RR}^2 - m_\Delta^2} \ln \frac{m_{RR}^2}{m_\Delta^2} \right) \\ \eta &\simeq -\frac{m_{RR}}{16\pi^2} \frac{m_{\phi\Delta}^2}{m_{RR}^2 - m_\Delta^2} \ln \frac{m_{RR}^2}{m_\Delta^2}\end{aligned}$$

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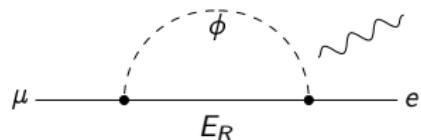
$$g\tilde{g}_\Delta \simeq 4.0 \times 10^{-6} \frac{m_\nu}{0.05 \text{ eV}} \frac{70 \text{ GeV}}{M_1} \frac{50 \text{ MeV}}{\delta} \frac{m_{RR}}{300 \text{ GeV}} \frac{0.1}{|\sin \alpha_1|} \left(\frac{m_{RR}^2}{m_{RR}^2 - m_\Delta^2} \dots \right)^{-1}$$

$$g\tilde{g}_\Delta \simeq 4.5 \times 10^{-6} \frac{m_\nu}{0.05 \text{ eV}} \frac{300 \text{ GeV}}{m_{RR}} \frac{1 \text{ GeV}^2}{\tilde{m}_{\phi\Delta}^2} \left(\frac{m_\Delta}{500 \text{ GeV}} \right)^2 \frac{m_{RR}^2 - m_\Delta^2}{m_\Delta^2} \left(\log \frac{m_{RR}^2}{m_\Delta^2} \right)^{-1}$$

$$\tilde{g}\tilde{g}_\Delta \simeq 1.8 \times 10^{-10} \frac{m_\nu}{0.05 \text{ eV}} \frac{300 \text{ GeV}}{m_{RR}} \frac{0.1}{\sin \alpha_1} \frac{m_{RR}^2 - m_\Delta^2}{m_\Delta^2} \left(\log \frac{m_{RR}^2}{m_\Delta^2} \right)^{-1}$$

Lepton Flavour Violation and $(g-2)_\mu$

Lepton Flavour Violation



$$\text{Br}(\mu \rightarrow e\gamma) = 2.5 \cdot 10^{-9} \left(\frac{300 \text{ GeV}}{m_{RR}} \right)^4 \left| \frac{g_\mu^*}{0.1} \frac{g_e}{0.1} \right|^2$$

$$\text{Br}(\tau \rightarrow \alpha\gamma) = 4.5 \cdot 10^{-10} \left(\frac{300 \text{ GeV}}{m_{RR}} \right)^4 \left| \frac{g_\tau^*}{0.1} \frac{g_\alpha}{0.1} \right|^2$$

Experimental Limits [PDG 2009 (90% CL)]

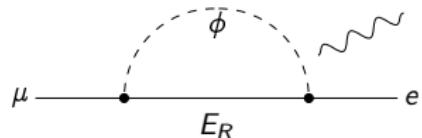
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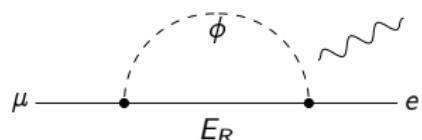
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Solutions

- $m_{RR}/g > 6 \text{ TeV}$
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(allowed by flavour structure)

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Anomalous Magnetic Moment of Muon

$$\delta(g-2)_\mu/2 \sim 10^{-11} \left(\frac{300 \text{ GeV}}{m_{RR}} \right)^4 |g_\mu|^2 \lesssim \text{exp. uncertainty}$$

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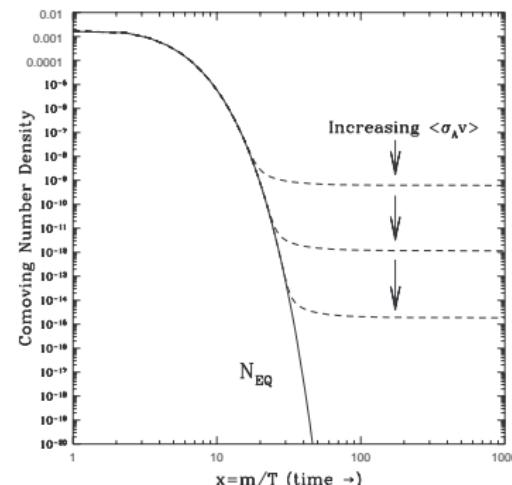
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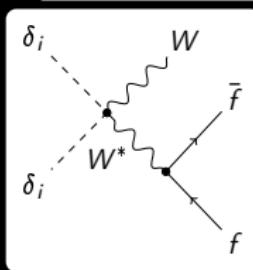
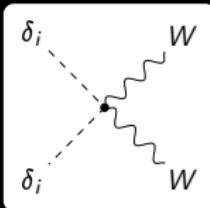
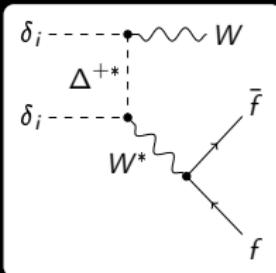
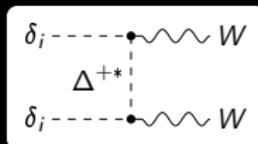
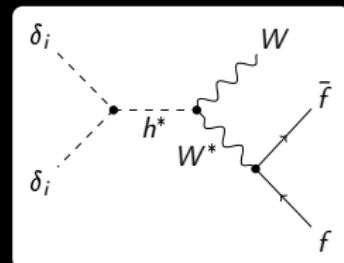
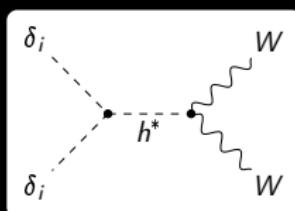
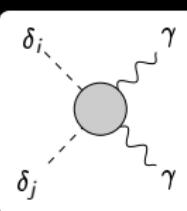
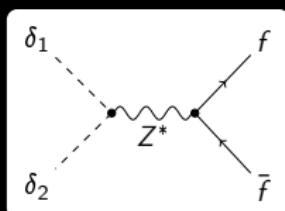
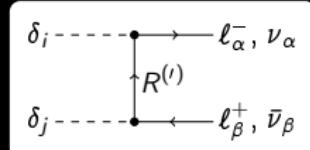
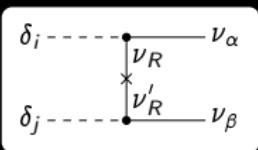
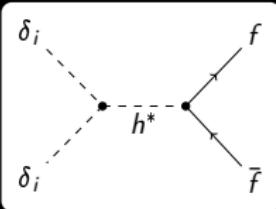
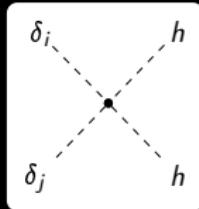
Dark Matter Freeze Out

- Assumption: thermal production after inflation
- Annihilation rate related to production rate
- Quasi-degenerate scalar masses
⇒ both species have to be considered
- $\sigma_{12} \ll \sigma_{11}, \sigma_{22} \Rightarrow \delta_1$ and δ_2 produced and later $\delta_2 \rightarrow \delta_1 \nu \bar{\nu}$
- $\Gamma(\delta_2 \xrightarrow{Z} \delta_1 \nu \bar{\nu}) \approx 14 \left(\frac{\delta}{50 \text{ MeV}} \right)^5 \left(\frac{\sin \alpha_1}{0.1} \right)^4 \text{ sec}^{-1}$
- $\sum_{i=1}^2 \langle \sigma(\delta_i \delta_i \rightarrow \dots) v \rangle = 3 \cdot 10^{-26} \frac{\text{cm}^3}{\text{sec}}$

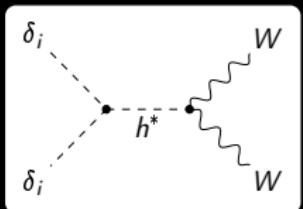
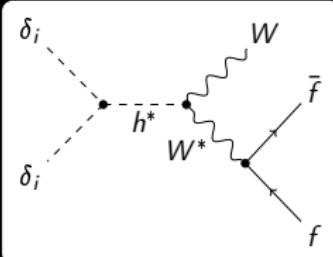
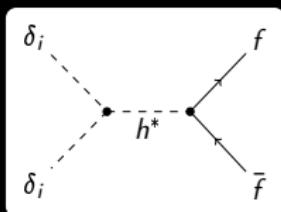
Thermal freezeout one species



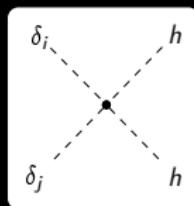
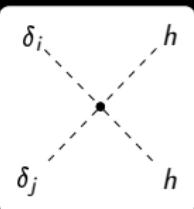
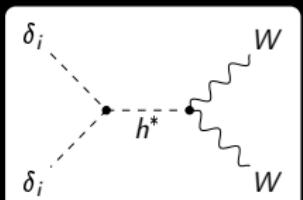
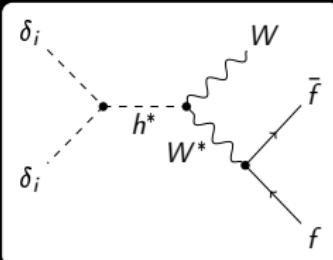
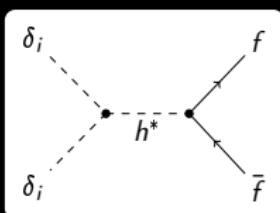
Dark Matter Annihilation Channels



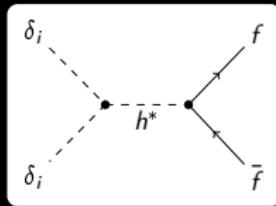
Dark Matter Annihilation



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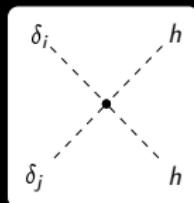
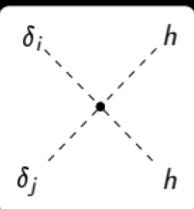
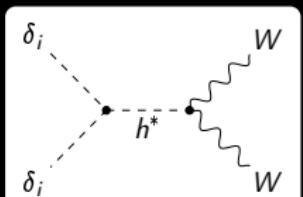
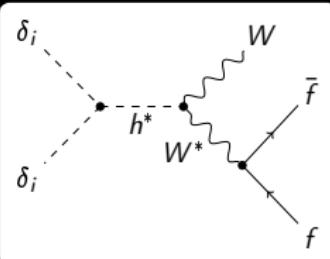
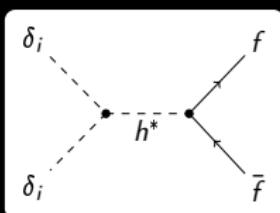


$$\Rightarrow \langle \sigma_{hh}^H v \rangle \simeq \frac{|\lambda_L|^2 (M_1^2 - m_h^2)^{1/2}}{16\pi M_1^3}$$

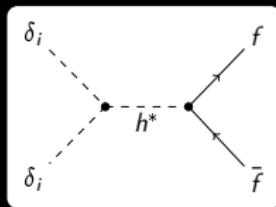


$$\Rightarrow \langle \sigma_{f\bar{f}}^H v \rangle \simeq N_c \frac{|\lambda_L|^2}{\pi} \frac{m_f^2}{(4M_1^2 - m_h^2)^2} \frac{(M_1^2 - m_f^2)^{3/2}}{M_1^3}$$

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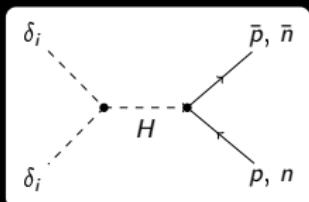
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In general for Higgs mediated annihilation:

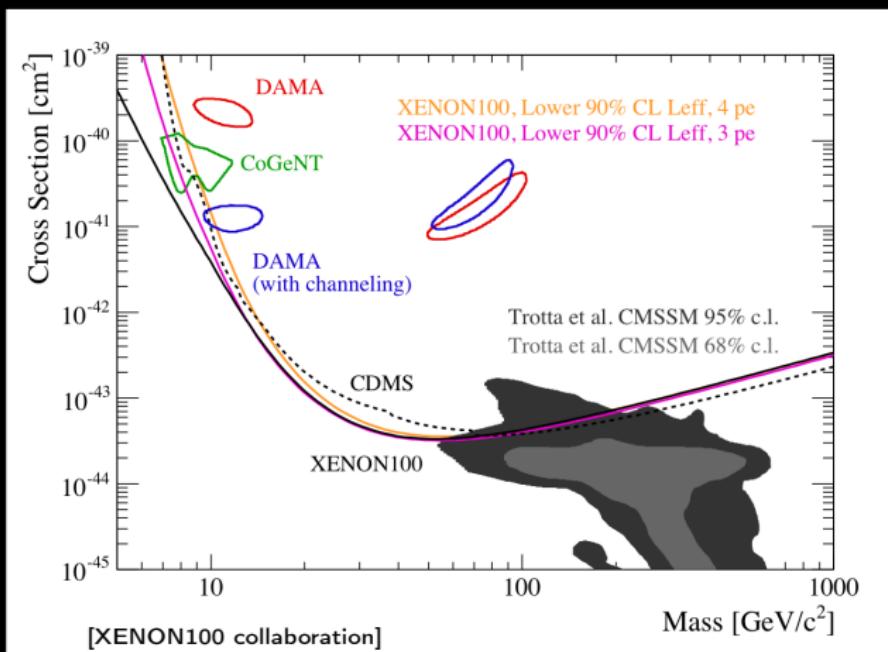
$$\langle \sigma(\delta_1 \delta_1 \rightarrow h^* \rightarrow \dots)_{Hv} \rangle = (2m_h \Gamma(h \rightarrow \dots))|_{m_h \rightarrow 2M_1} \frac{1}{4M_1^2} \frac{4|\lambda_L|^2 v_H^2}{(4M_1^2 - m_h^2)^2}$$

Using HDecay: $\Gamma|_{2 \times 70 \text{ GeV}} = 8.3 \text{ MeV} \Rightarrow \lambda_L \approx 0.07$

Direct DM Detection

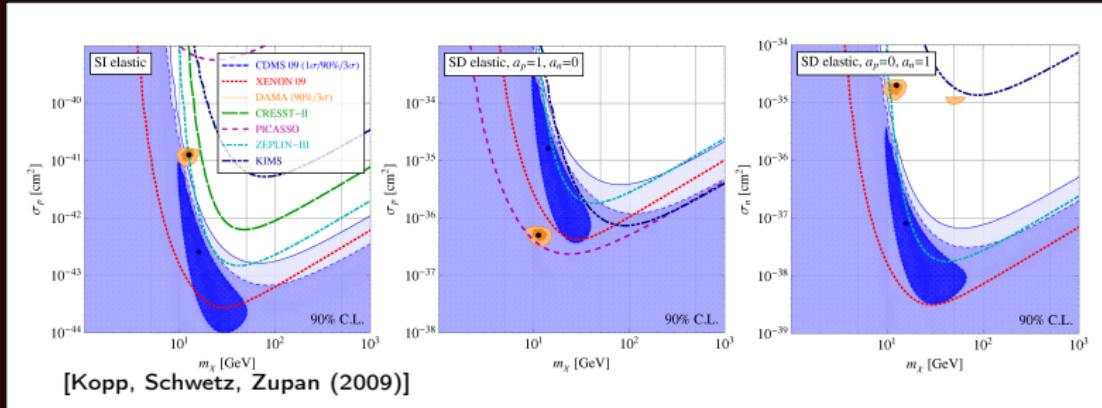


$$\sigma_n = \frac{|\lambda_L|^2}{\pi} \frac{\mu_{\delta_1 n}^2 m_p^2}{M_1^2 m_h^4} f^2$$
$$\approx 5.2 \times 10^{-44} \left(\frac{\lambda_L}{0.07} \right)^2 \left(\frac{70 \text{ GeV}}{M_1} \right)^2 \left(\frac{120 \text{ GeV}}{m_h} \right)^4 \left(\frac{f}{0.3} \right)^2 \text{ cm}^2$$



Elastic SI Scattering

CDMS-II



- $M_1 \sim 20\text{--}50$ GeV and $\sigma_n \sim 10^{-44}$ cm 2 – 10^{-43} cm 2
- for $M_1 = 50$ GeV, $m_h = 120$ GeV $\Rightarrow \sigma_n \simeq 5.4 \times 10^{-44} \left(\frac{f}{0.14} \right)^2$ cm 2

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CoGeNT/DAMA

- $7 \text{ GeV} \lesssim M_1 \lesssim 11 \text{ GeV}$ with $\sigma_n \sim 10^{-41} \text{ cm}^2\text{--}10^{-40} \text{ cm}^2$

$$\sigma_n \approx 1.3 \times 10^{-40} \left(\frac{f}{0.3}\right)^2 \left(\frac{8 \text{ GeV}}{M_1}\right)^2 \text{ cm}^2$$

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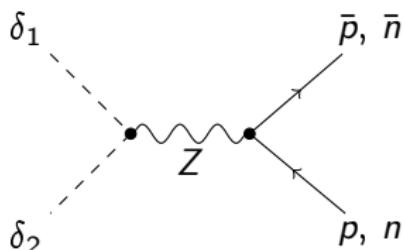
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However, a proper analysis is required to make definite statements, whether it is possible at all!

Inelastic SI Scattering

Best fit in [Schwetz, Kopp, Zupan (2009)]: $M_1 = 10 \text{ GeV}$ with $\sigma_p \sim 1 \times 10^{-40} \text{ cm}^2$ which corresponds to $\sigma_n = 3.3 \times 10^{-40} \text{ cm}^2$



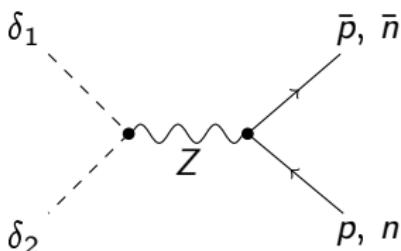
$$f_p/f_n = -(1 - 4 \sin^2 \theta_W) \approx -0.08$$

$$\sigma_n = \frac{8}{\pi} \sin^2 \alpha_1 \sin^2 \alpha_2 G_F^2 \mu_{\delta_1 n}^2$$

$$\simeq 6 \times 10^{-40} (\sin \alpha_1 \sin \alpha_2 / 0.07)^2 \text{ cm}^2$$

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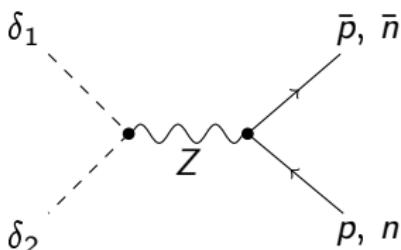
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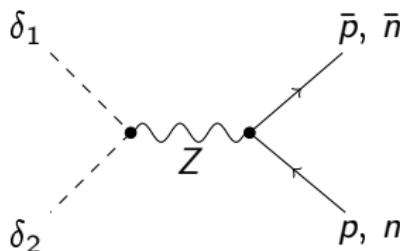
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- Increase of energy density for neutrinos

$$\frac{\Delta \rho_\nu}{\rho_\nu} \equiv \frac{\rho_\nu^f - \rho_\nu^i}{\rho_\nu} = \frac{\Omega_2 - \Omega_1}{\Omega_\nu} \approx \frac{\delta}{2M_1} \frac{\Omega_{\text{DM}}}{\Omega_\nu} \lesssim 1.2 \times 10^{-4},$$

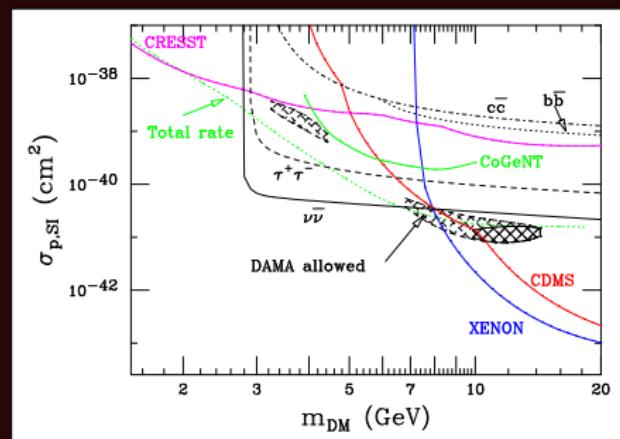
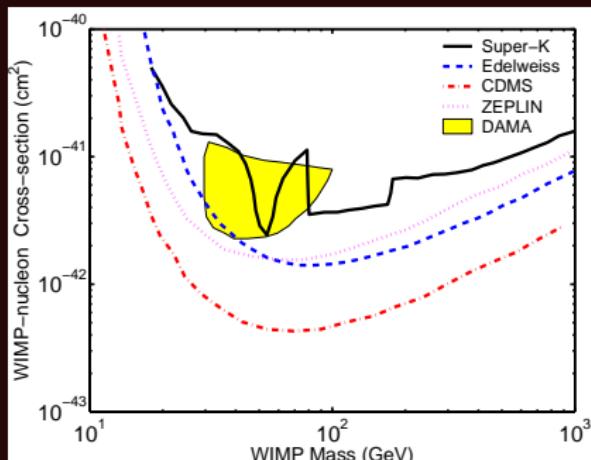
Neutrinos From Sun

- Number of WIMPs: $\dot{N} = C - AN^2 - EN$
- Capture rate
 $C(\rho_{DM}, \bar{\nu}, m_{DM}, \sigma) \simeq 1.3 \cdot 10^{25} \text{ sec}^{-1} \propto \rho_{DM} \sigma \bar{\nu}^{-1} m_{DM}^{-1}$
- Annihilation Rate $A = \langle \sigma v \rangle / V_{\text{eff}}$
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Bounds [Hooper, Petriello, Zurek, Kamionkowski (2008)]



Indirect Detection

- Possible DM signals in γ -rays (e.g. EGRET, Fermi-LAT), neutrinos (e.g. IceCube), positrons (e.g. PAMELA), anti-protons (e.g. PAMELA), anti-deuterons (e.g. AMS-02, GAPS)

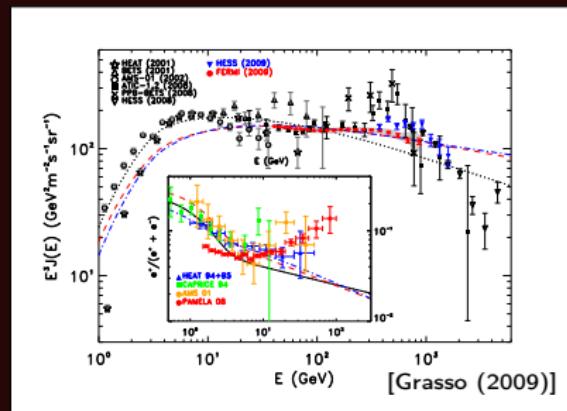
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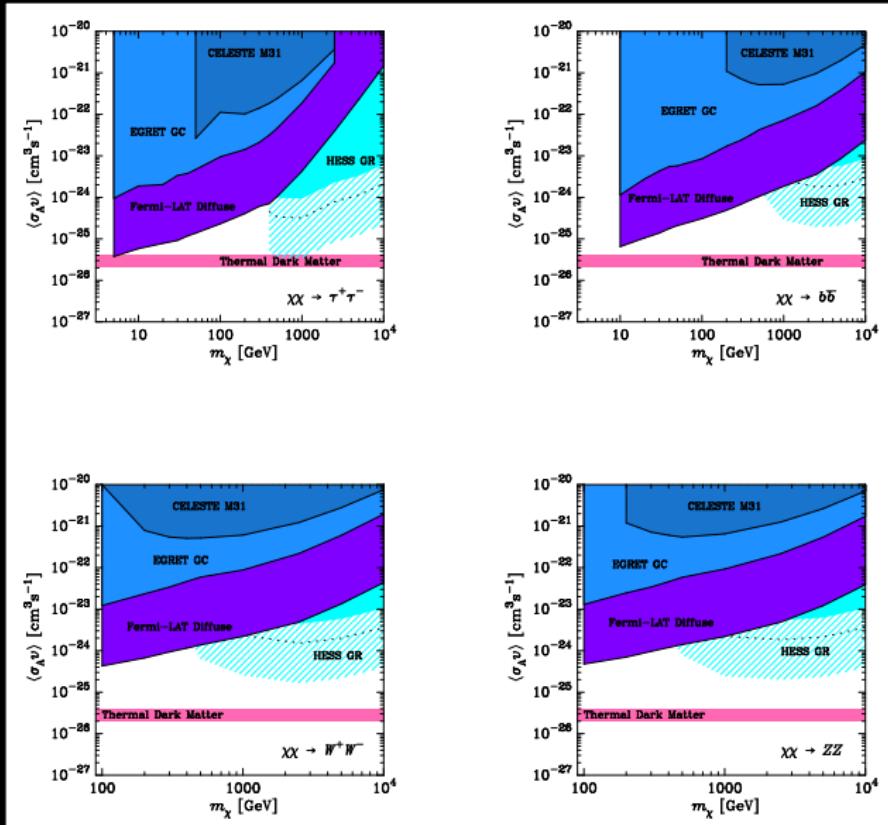
Signals from PAMELA, Fermi-LAT, . . .



- Positron excess in PAMELA
- Bump in charged lepton flux in ATIC/BESS, FERMI/HESS
- \Rightarrow can be explained by acceleration of secondaries
[Blasi, Serpico (2009); Ahlers, Mertsch, Sarkar (2009)] as well as pulsars
- Study of nuclei in cosmic rays can refute or confirm explanations

Constraints from Isotropic Diffuse γ -ray Background

[Abazajian, Agrawal, Chacko, Kilic (2010)]



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2 Lepton Sector

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4 More Phenomenology

5 Comments on Alternatives

6 Conclusions

Constraints

Electroweak Precision Tests: Higgs Triplet

$$\hat{S} = \frac{g_{\text{SU}(2)}^2}{24\pi^2} \xi \quad \hat{T} = \frac{25g_{\text{SU}(2)}^2}{576\pi^2} \frac{m_\Delta^2}{m_W^2} \xi^2$$

with $\xi := (m_{\Delta^{++}}^2 - m_\Delta^2)/m_\Delta^2$ and $m_{\Delta^{++}}^2 = m_\Delta^2 + 2m_{\Delta^+}^2$

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Invisible Z-Decay Width

- DM particle δ_1 couples to Z-boson via mixing of Δ^0
- If $M_1 + M_2 < m_Z$, the corresponding Z-decay width is

$$\Gamma(Z \rightarrow \delta_1 \delta_2) = \frac{G_F \sin^2 \alpha_1 \sin^2 \alpha_2}{6\sqrt{2}\pi} m_Z^3$$

- Bound on mixing angle in scalar sector: $\sin \alpha_1 \sin \alpha_2 < 0.07$,
i.e. $m_{\phi\Delta}^2 \ll m_\Delta^2$ (protected by $U(1)_\phi \times U(1)_\Delta$)

Collider Physics

Higgs Search

- Higgs might dominantly decay invisibly if $2M_1 < m_h$

$$H \rightarrow \delta_1 \delta_1, \quad H \rightarrow \delta_2 \delta_2 \rightarrow (\delta_1 \nu \bar{\nu})(\delta_1 \nu \bar{\nu})$$

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New Particles

- New particles accessible at LHC
- ... decay into the SM particles and DM \Rightarrow missing energy
- Mass relation of triplet $2m_{\Delta^+}^2 = m_{\Delta^{++}}^2 + m_{\Delta^+}^2$
- Expect small mass splitting $m_{\Delta^{++}}^2 - m_{\Delta^+}^2$
- Determination of g_α : $\text{Br}(E_R^- \rightarrow \ell_\alpha^- \delta_{1,2}) \propto |g_\alpha|^2$
- Determination of \tilde{g}_Δ : decay modes of Δ^+ and Δ^{++} , especially $\Gamma(\Delta^{++} \rightarrow \ell_\alpha^+ \ell_\beta^+ \delta_{1,2}) \propto |(\tilde{g}_\Delta)_\alpha g_\beta + (\tilde{g}_\Delta)_\beta g_\alpha|^2$

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Alternative Scenario [Boehm, Farzan, Hambye, Palomares-Ruiz, Pascoli (2006)]

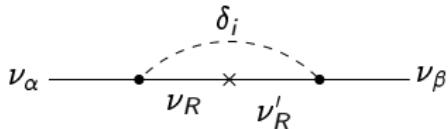
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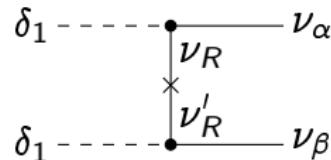
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Neutrino Mass generation



DM annihilation

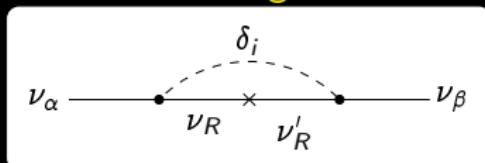


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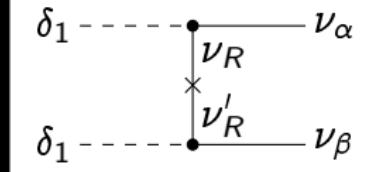
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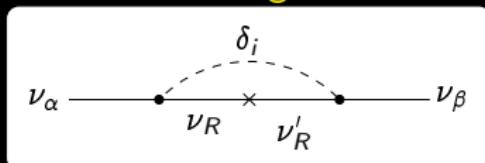


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 \Rightarrow light dark matter $M_1 \sim \mathcal{O}(\text{MeV})$

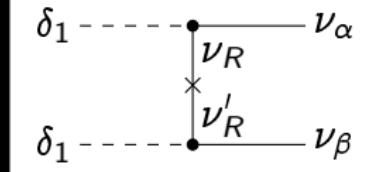
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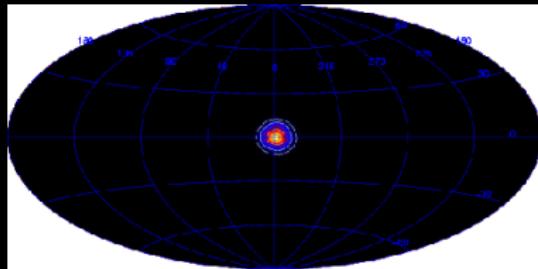
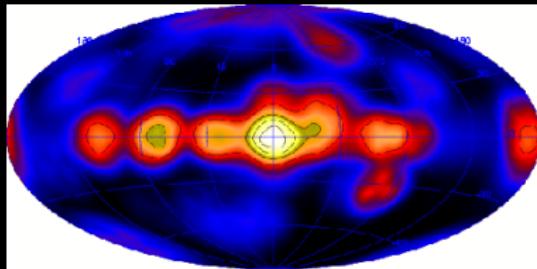
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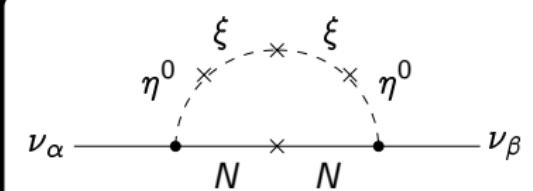
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- MeV scale dark matter might explain the 511 keV γ line measured by SPI/INTEGRAL [Weidenspointner (2007)]



A Gauged Model

Particle Content

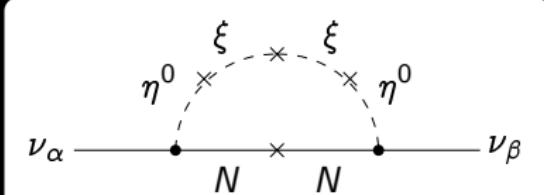
	SU(2)	U(1)	U(1) _X	\mathbb{Z}_2
$\ell_L^{(i)}$	2	-1	0	+
N_R	1	0	1	-
N_L	1	0	1	-
η	2	-1	-1	-
ξ	1	0	-1	-
ϕ	1	0	-2	+



A Gauged Model

Particle Content

	SU(2)	U(1)	U(1) _X	\mathbb{Z}_2
$\ell_L^{(i)}$	2	-1	0	+
N_R	1	0	1	-
N_L	1	0	1	-
η	2	-1	-1	-
ξ	1	0	-1	-
ϕ	1	0	-2	+



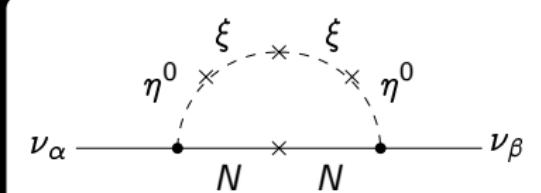
- Neutrino masses from pseudo Dirac neutrinos N

$$m_{\alpha\beta} \simeq \frac{\kappa_{\xi\phi} \lambda_{H\eta\xi\phi}^2}{16\pi^2} \frac{\langle\phi\rangle^2 \kappa_{H\eta\xi}^2}{m_\eta^4} \frac{\langle H\rangle^2}{m_N^2} \sum_{i,j} (Y_N)_{i\alpha} (Y_N)_{j\beta} f_{ij}(m_N, Y_{LL}, Y_{RR}, m_\xi)$$

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- LFV similar: $(m_N, m_{\eta^-})/Y_N \gtrsim 6 \text{ TeV}$ (unless special flavour structure)

More Phenomenology

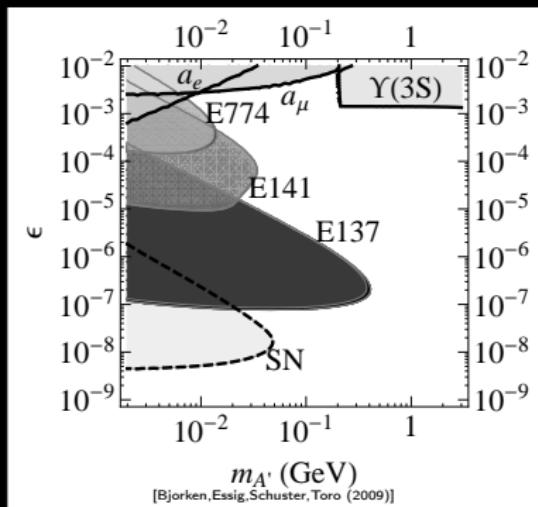
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Gauge kinetic Lagrangian

$$\mathcal{L}_{U(1)} = \frac{1}{4} F_Y^2 + \frac{1}{4} F_X^2 + \frac{\epsilon}{2} F_X F_Y$$

η in loop induces

$$\epsilon \simeq \frac{g_X g_Y}{16\pi^2} \ln \frac{\Lambda^2}{\mu_\eta^2}$$



More Phenomenology

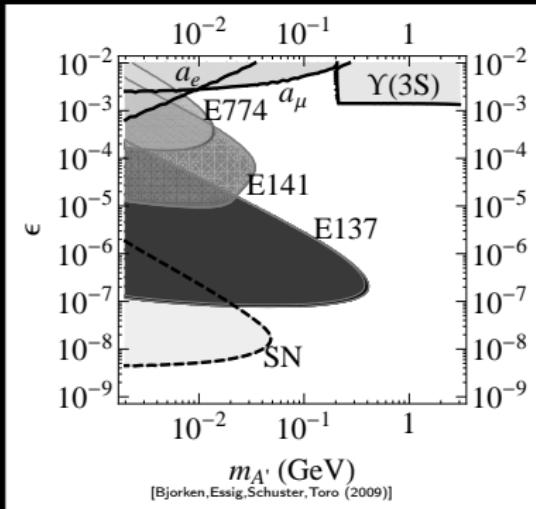
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Collider

- SM Higgs mixes with $U(1)_X$ -Higgs \Rightarrow Higgs mass bounds weakened
- Invisible Higgs decay like in all DM models in which Higgs exchange dominates

Outline

1 Model

2 Lepton Sector

3 Dark Matter

- Dark Matter Annihilation
- Dark Matter Direct Detection
- DAMA/CoGeNT/CDMS-II?
- Indirect Detection

4 More Phenomenology

5 Comments on Alternatives

6 Conclusions

Conclusions and Outlook

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- Neutrino mass can be generated radiatively at the TeV scale **and** linked to dark matter

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Outlook

- Study of a gauged model
- Leptogenesis in models of radiative neutrino mass generation

Thank you very much for your attention.