

Probing Lepton Number Violation in Double Beta Decay and at the LHC



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in collaboration with

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MPIK Heidelberg, 21 January 2013

Overview

- Neutrinos
 - Oscillations
 - Absolute Mass
- Neutrinoless Double Beta Decay
 - Light Neutrino Exchange
 - New Physics Mechanisms
- Neutrino Mass Models
 - Effective Mass and Seesaw
 - Minimal Left-Right Symmetry
- LFV and LNV at the LHC
- Conclusion

Neutrino Oscillations

- **Neutrino interaction states different from mass eigenstates**

Neutrino flavour can change through propagation

$$\nu_i = \sum_{\alpha} U_{i\alpha} \nu_{\alpha}, \quad \nu_i(t) = e^{-i(E_i t - p_i x)} \nu_i$$

$$\Rightarrow P_{\alpha \rightarrow \beta} = \sin^2(2\theta) \sin^2\left(1.27 \frac{\Delta m^2}{\text{eV}^2} \frac{L/\text{km}}{E/\text{GeV}}\right)$$

- **Solar neutrino oscillations**

Large mixing

- **Atmospheric oscillations**

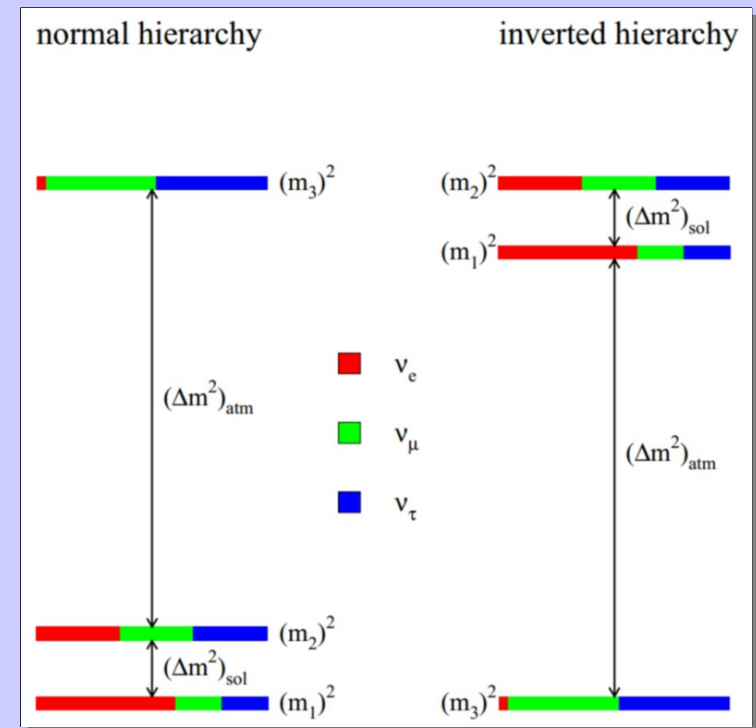
$\frac{5}{8}$ Maximal mixing

- **Reactor and accelerator neutrinos**

$$\sin^2(2\theta_{13}) = 0.092 \pm 0.021$$

- **Experimental unknowns and anomalies**

CP violation? Sign of Δm_{23} ? Sterile Neutrinos?



Absolute Neutrino Mass

- Energy endpoint in Beta decay

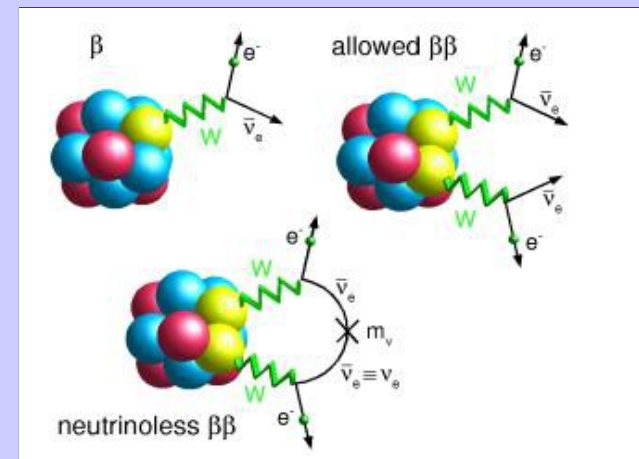
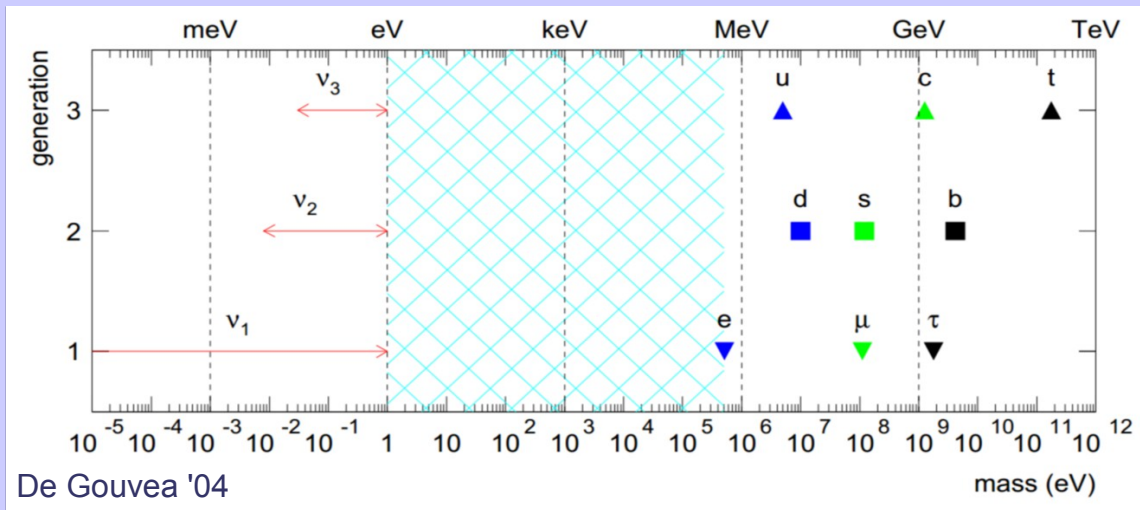
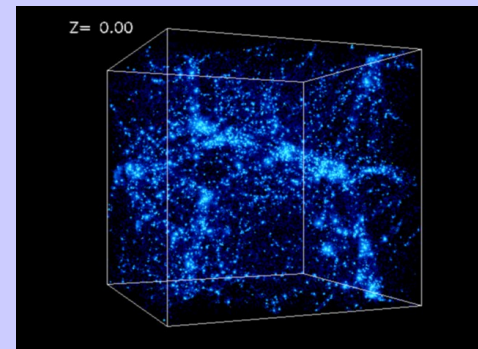
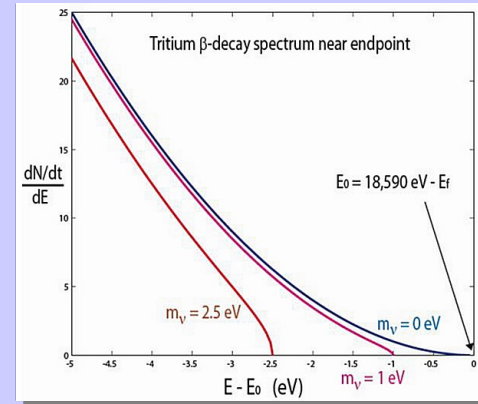
$$m_{\beta}^2 = \sum_i |U_{ei}|^2 m_i^2 < (2.2 \text{ eV})^2 \quad \text{KATRIN: } m_{\beta} \stackrel{5}{8} 0.2 \text{ eV}$$

- Impact on Large Scale Structure

$$\Sigma = \sum_i m_i < 0.4 - 1 \text{ eV}$$

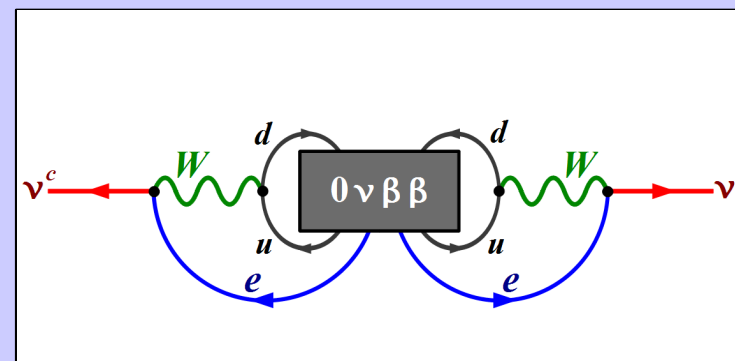
- Neutrinoless Double Beta Decay

$$m_{\beta\beta} = \left| \sum_i U_{ei}^2 m_{\nu_i} \right| < 0.2 - 2.0 \text{ eV} \quad \text{Future Experiments: } m_{\beta\beta} \stackrel{5}{8} 0.01 \text{ eV}$$



Neutrinoless Double Beta Decay

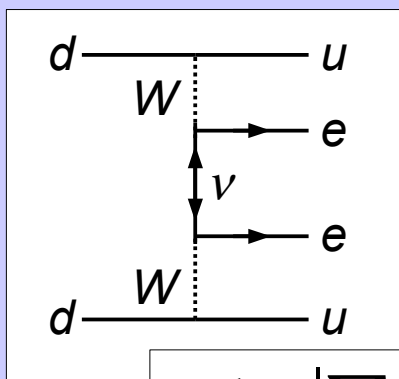
- **Process:** $(A, Z) \xrightarrow{23}_{11} (A, Z+2) + 2e^-$
- **Uncontroversial detection of $0\nu\beta\beta$ of utmost importance**
 - Prove lepton number to be broken
 - Prove neutrinos to be Majorana particles (Schechter, Valle '82)



$$\delta m_\nu \approx \frac{1}{(16\pi^2)^4} \frac{\text{MeV}^5}{M_W^4} \approx 10^{-23} \text{ eV}$$

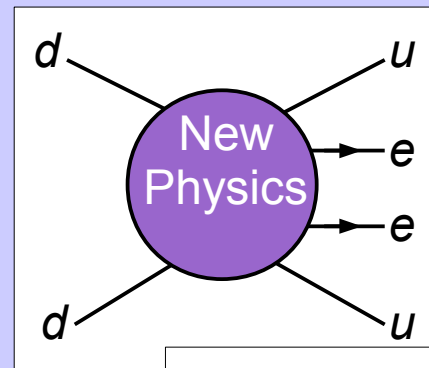
- **Which mechanism triggers the decay?**

Light Neutrino Exchange
(LH Current, Mass Mechanism)



$$T_{1/2}^{-1} \propto \left| \sum_i U_{ei}^2 m_{\nu_i} \right|$$

General Effective Operator

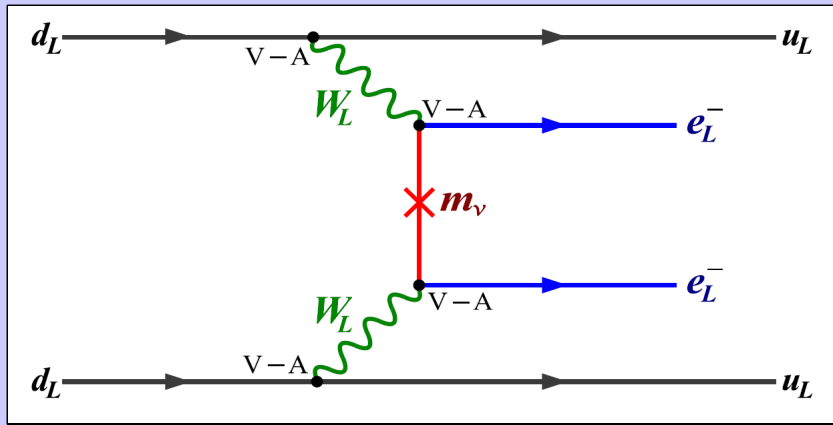


$$\bar{u}\bar{u}\bar{e}\bar{e}dd/M^5$$

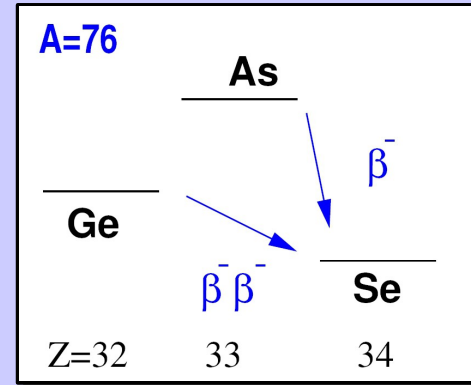
$$T_{1/2}^{-1} \approx 10^{25} \text{ y} \rightarrow M \approx 1 \text{ TeV}$$

Light Neutrino Exchange

- Standard Mass Mechanism



Heidelberg-Moscow
 $T_{1/2}({}^{76}\text{Ge}) \stackrel{5}{\approx} 1.9 \times 10^{25} \text{ y}$
 $m_\nu \stackrel{5}{\approx} (0.3 - 0.6) \text{ eV}$



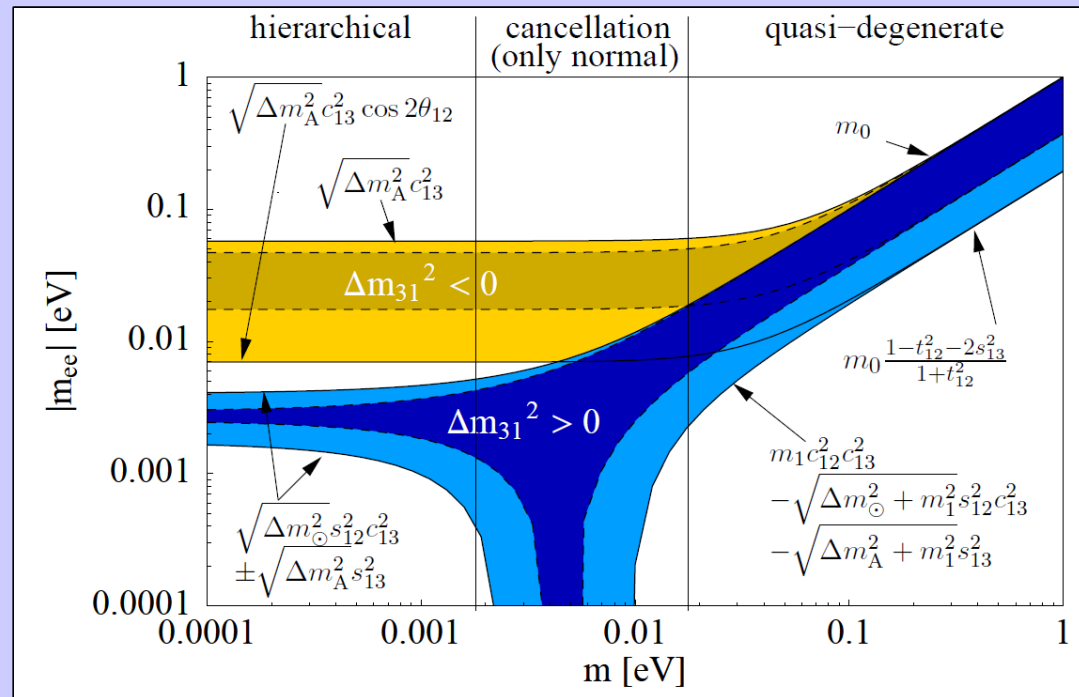
Lindner, Merle, Rodejohann (2005)

- Decay Rate

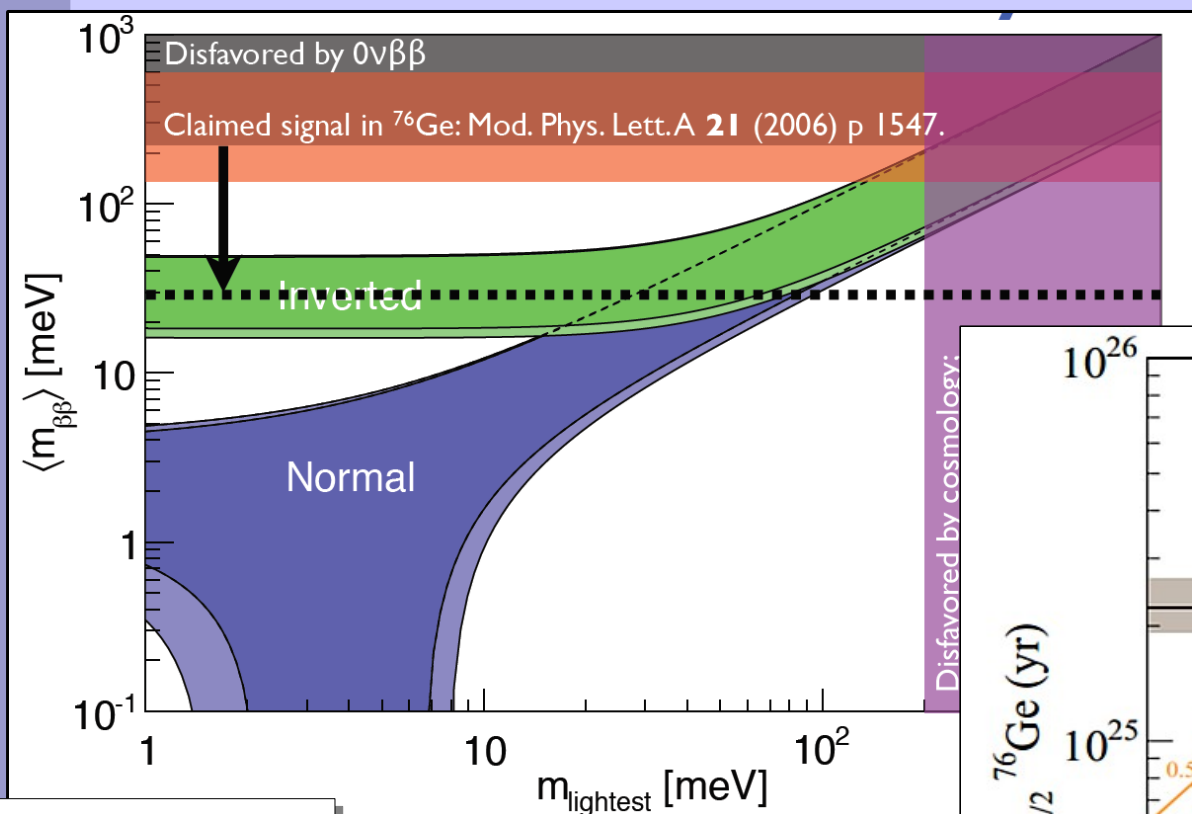
$$\Gamma = T_{1/2}^{-1} = \frac{m_{\beta\beta}^2}{m_e^2} G^{0\nu} |M^{0\nu}|^2$$

- Effective Mass

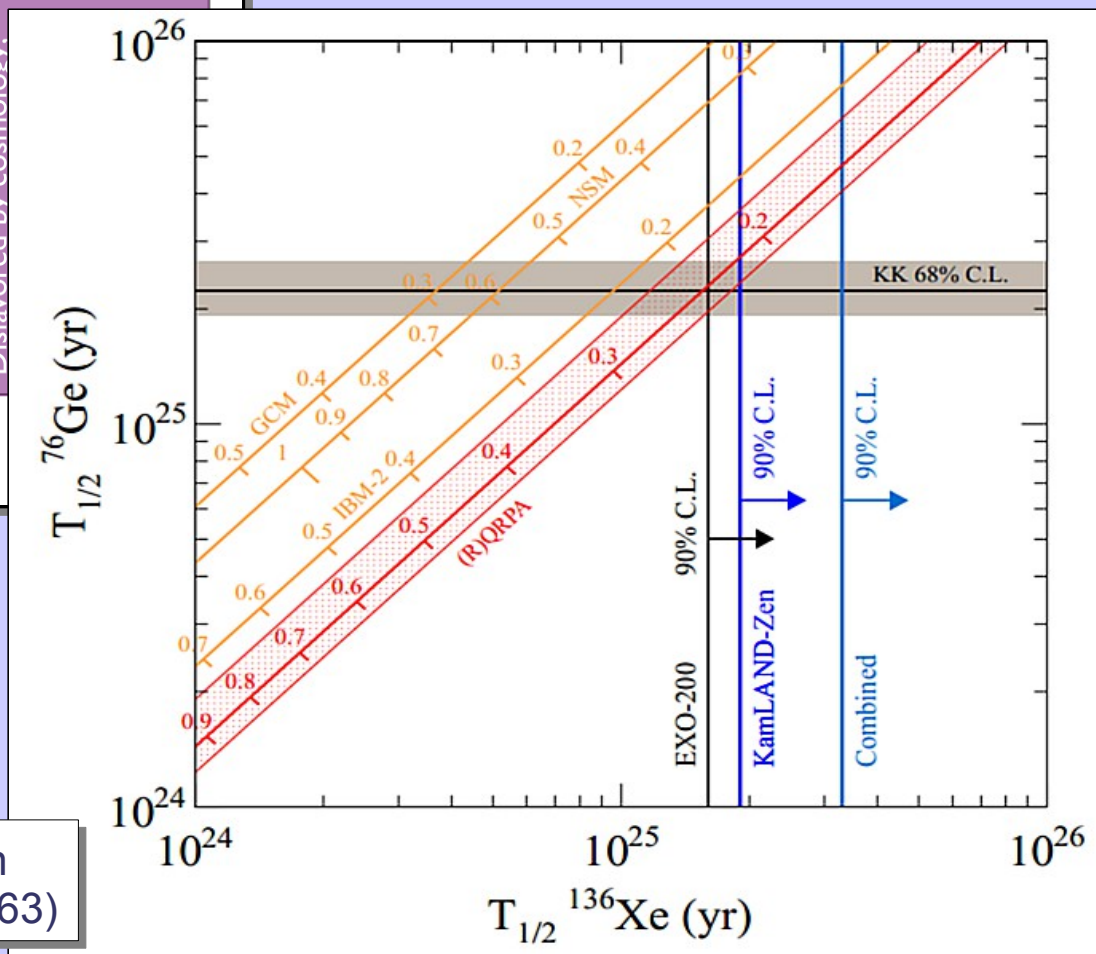
$$m_{\beta\beta} = \left| \sum_i U_{ei}^2 m_{\nu_i} \right| \equiv (m_\nu)_{ee}$$



Experimental Situation



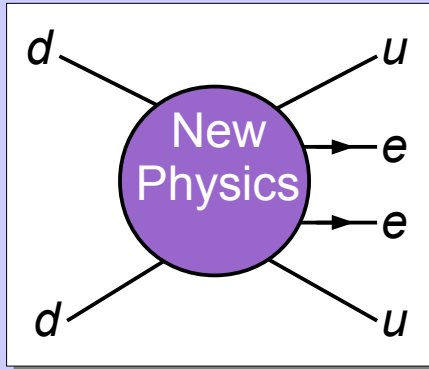
Detwiler (2012)



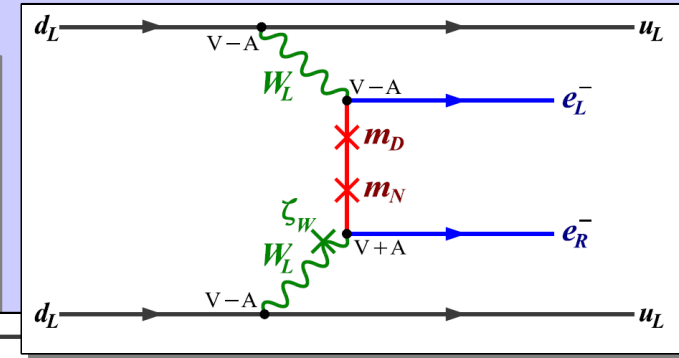
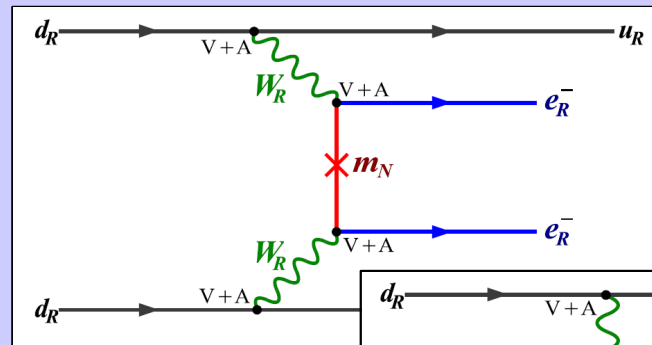
KamLAND-Zen
(arXiv:1211.3863)

New Physics Contributions to $0\nu\beta\beta$

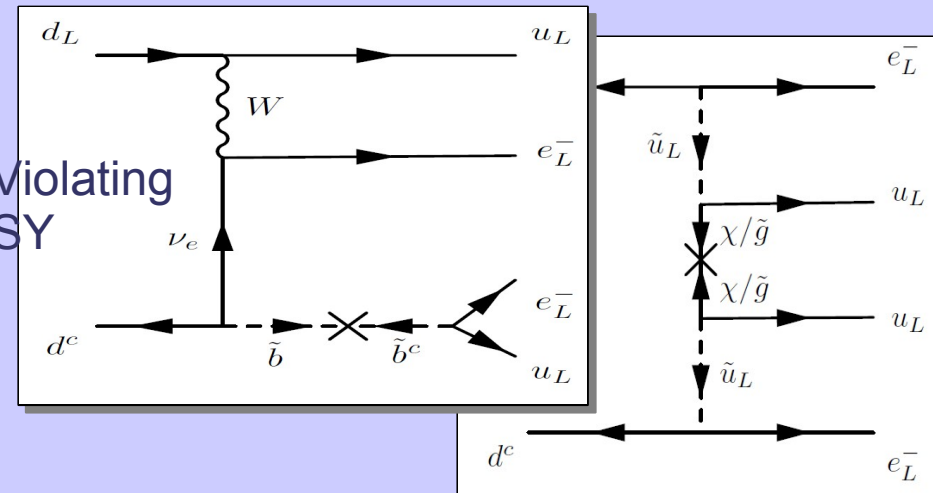
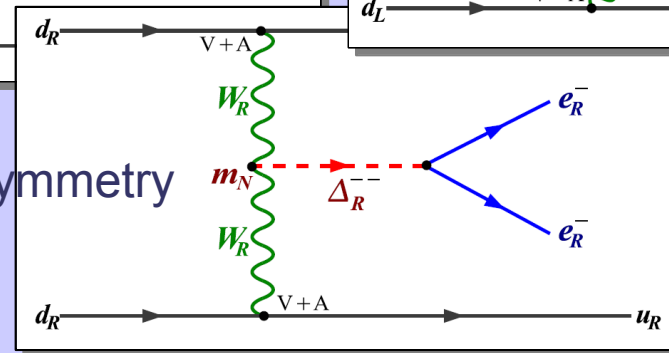
• Plethora of New Physics Scenarios



=



Left-Right Symmetry



$$\Gamma = T_{1/2}^{-1} = \epsilon_{\text{NP}}^2 G_{\text{NP}}^{0\nu} |M_{\text{NP}}^{0\nu}|^2$$

Extra Dimensions

Majorons

R-Parity Violating SUSY

Leptoquarks

...

Effective Mass and Seesaw Mechanism

- Effective operator for Majorana neutrino mass

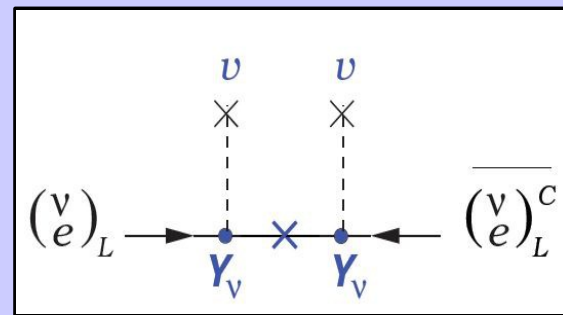
$$L = \frac{1}{2} \frac{h_{ij}}{\Lambda_{\text{LNV}}} (\bar{L}_i^c \cdot \tilde{H}) (\tilde{H}^T \cdot L_j) \rightarrow \frac{1}{2} (m_\nu)_{ij} \bar{\nu}_i^c \nu_j$$

Unique dim-5 Operator

- Seesaw Mechanism

Add right-handed neutrinos to the Standard Model particle content, $M \stackrel{\text{S}}{\sim} 10^{14}$ GeV

$$L = L_{\text{SM}} - \frac{1}{2} \bar{\nu}_R M \nu_R^c + \bar{\nu}_R Y_\nu L \cdot H_u$$



- Light neutrino mass matrix at low energies

$$m_\nu = m_D^T M^{-1} m_D \text{ for } m_D \ll M_R \quad m_\nu \approx 0.1 \text{ eV} \left(\frac{m_D}{100 \text{ GeV}} \right)^2 \left(\frac{M}{10^{14} \text{ GeV}} \right)^{-1}$$

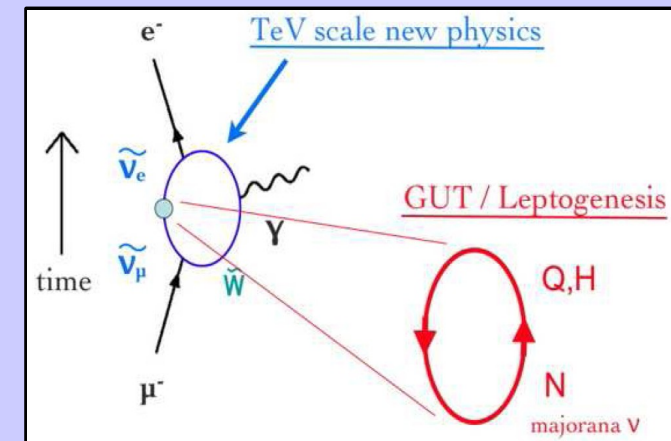
Problems of Seesaw Mechanism

- **Introduces high energy scale**
- **Right-handed neutrinos are singlets**
Couple only via small mixture with active neutrinos
- **Mechanism not testable with low energy observables**



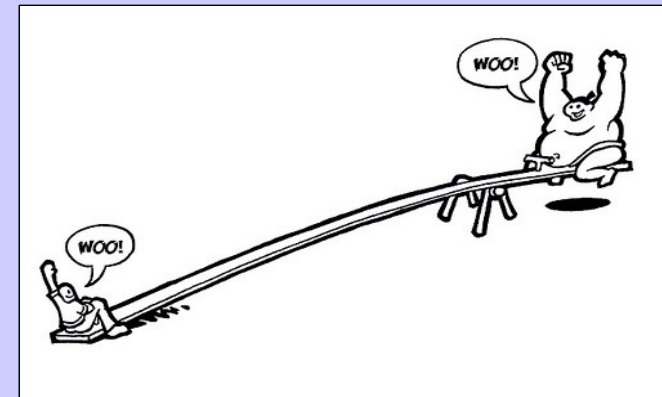
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- **Possible Solutions**
 - **SUSY Seesaw**
Testable LFV effects from sleptons



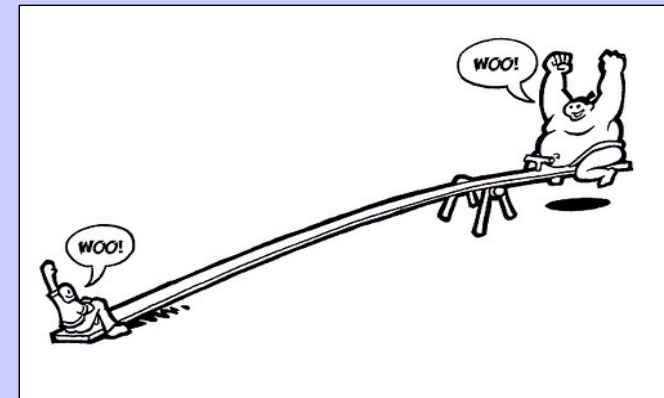
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 - **“Bent” Seesaw mechanisms**
LNV at low scale allows low mass of right-handed neutrinos



Problems of Seesaw Mechanism

- **Introduces high energy scale**
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 - **SUSY Seesaw**
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 - **“Bent” Seesaw mechanisms**
LNV at low scale allows low mass of right-handed neutrinos
 - **Left-Right symmetric models**
Right-handed neutrinos couple with gauge strength to charged leptons



Minimal Left-Right Symmetrical Model

- Based on

$$SU(3) \times SU(2)_L \times SU(2)_R \times U(1)_{B-L}$$

Pati & Salam '74
Mohapatra & Senjanovic '75

- Higgs Sector:**

Bidoublet (EW Breaking) + Left-handed Triplet + Right-handed Triplet (Breaking Lepton Number + Parity + $SU(2)_R$)

- Generate $N_i + W_R + Z_R$ masses

$$M_{N_i} \approx M_{W_R} \approx M_{Z_R} \approx \langle \Delta_R \rangle \approx 0.5 - 5 \text{ TeV}$$

- General Seesaw II Mechanism

$$M_\nu = \begin{pmatrix} M_L & M_D \\ M_D^T & M_R \end{pmatrix},$$

- Charged current weak interactions

Neglect any
Left-Right mixing

$$J_{W_1}^{\mu-} = \frac{g_L}{\sqrt{2}} (\bar{\nu}_i U_{li}^{LL} + \bar{N}_i^c U_{li}^{LR}) \gamma^\mu \ell_L + \frac{g_R}{\sqrt{2}} \sin \zeta_W (\bar{\nu}_i U_{li}^{RL} + \bar{N}_i U_{li}^{RR}) \gamma^\mu \ell_R,$$

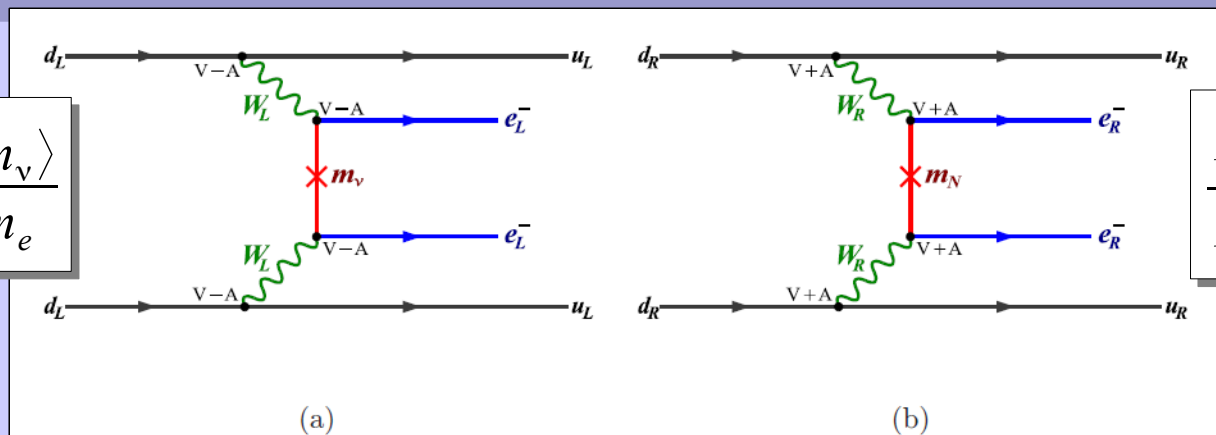
$$J_{W_2}^{\mu-} = -\frac{g_L}{\sqrt{2}} \sin \zeta_W (\bar{\nu}_i U_{li}^{LL} + \bar{N}_i U_{li}^{LR}) \gamma^\mu \ell_L + \frac{g_R}{\sqrt{2}} (\bar{N}_i U_{li}^{RR} + \bar{\nu}_i^c U_{li}^{RL}) \gamma^\mu \ell_R,$$

$$J_{W_L}^{\mu-} \approx \frac{g_L}{\sqrt{2}} U_{li} \bar{\nu}_i \gamma^\mu \ell_L,$$

$$J_{W_R}^{\mu-} \approx \frac{g_R}{\sqrt{2}} V_{li} \bar{N}_i \gamma^\mu \ell_R,$$

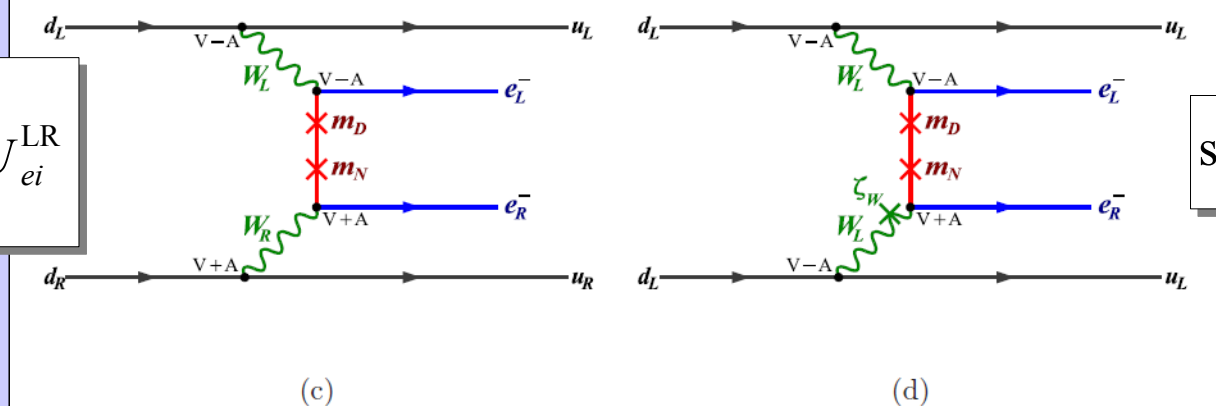
Neutrinoless Double Beta Decay in the LRSM

$$\sum_i (U_{ei}^{LL})^2 \frac{m_{\nu_i}}{m_e} = \frac{\langle m_\nu \rangle}{m_e}$$



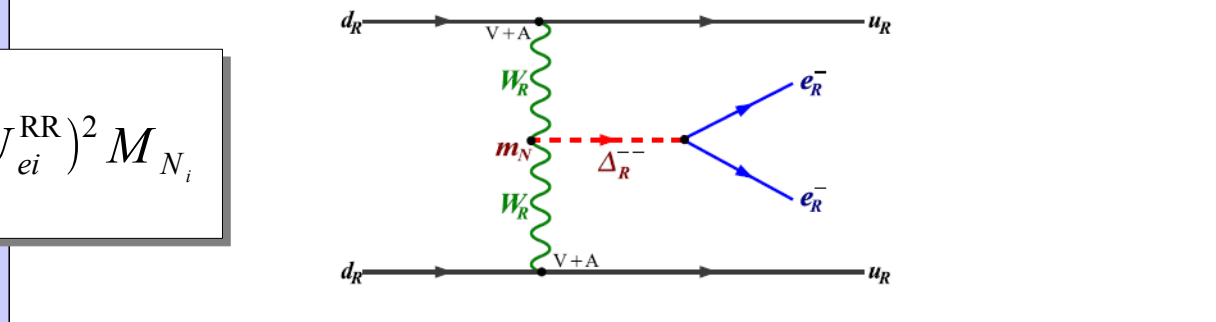
$$\frac{M_{W_L}^4}{M_{W_R}^4} \sum_i \frac{(U_{ei}^{RR})^2}{M_{N_i}}$$

$$\left(\frac{M_{W_L}}{M_{W_R}} \right)^2 \sum_i U_{ei}^{LL} U_{ei}^{LR}$$



$$\sin^2 \zeta \sum_i U_{ei}^{LL} U_{ei}^{LR}$$

$$\frac{M_{W_L}^4}{M_{W_R}^4} \frac{m_p}{M_{\Delta_R^{--}}^2} \sum_i (U_{ei}^{RR})^2 M_{N_i}$$



Charged Lepton Flavour Violation

- Lepton flavour practically conserved in the Standard Model

$$Br(\mu \rightarrow e \gamma) = \frac{3\alpha}{32\pi} \left| \sum_i U_{\mu i}^* U_{ei} \frac{\Delta m_{1i}^2}{m_W^2} \right|^2 \approx 10^{-54}$$

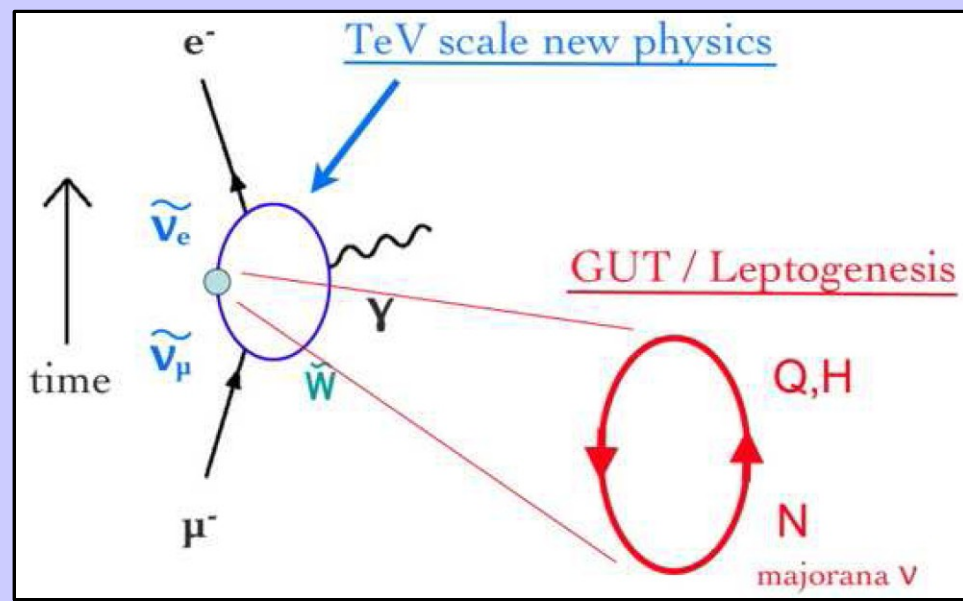
LFV is clear sign for BSM physics

- Flavour violation in quark and neutrino sector

Strong case to look for charged LFV

- LFV can shed light on

- Grand Unification models
- Flavour symmetries
- Origin of flavour



Rare LFV Processes

- Mediated by right-handed neutrinos and doubly charged Higgs bosons

$$BR(\mu \rightarrow e\gamma) \approx 2 \times 10^{-9} \sin^2(2\phi) \left(\frac{\Delta m_{12}^2}{m_{W_R}^2}\right)^2 \left(\frac{2 \text{ TeV}}{m_{W_R}}\right)^4,$$

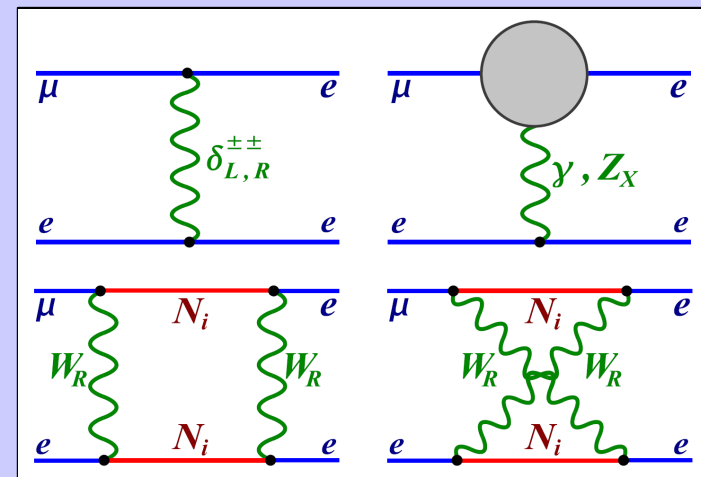
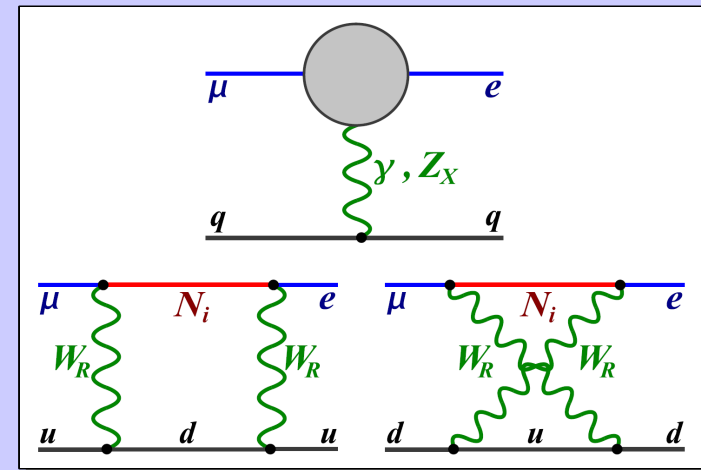
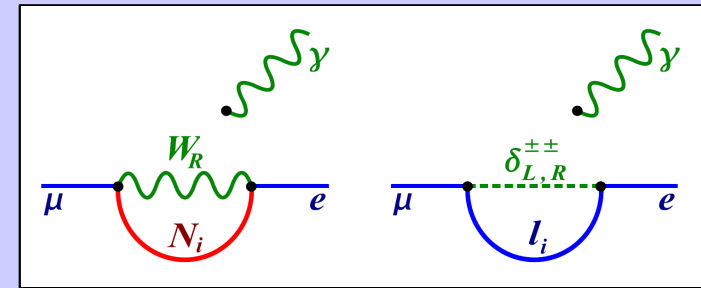
- μ -e conversion in nuclei enhanced via box diagrams

$$R(\mu \rightarrow e) \approx Br(\mu \rightarrow e\gamma)$$

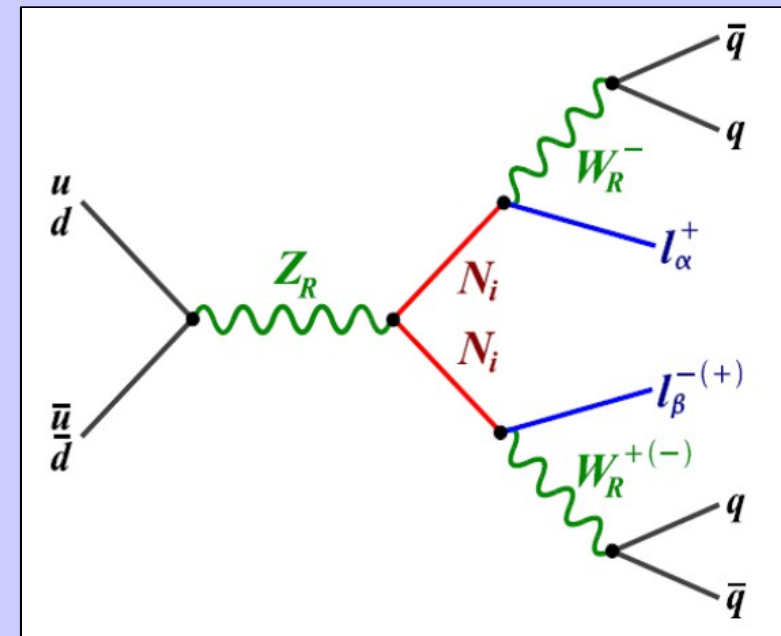
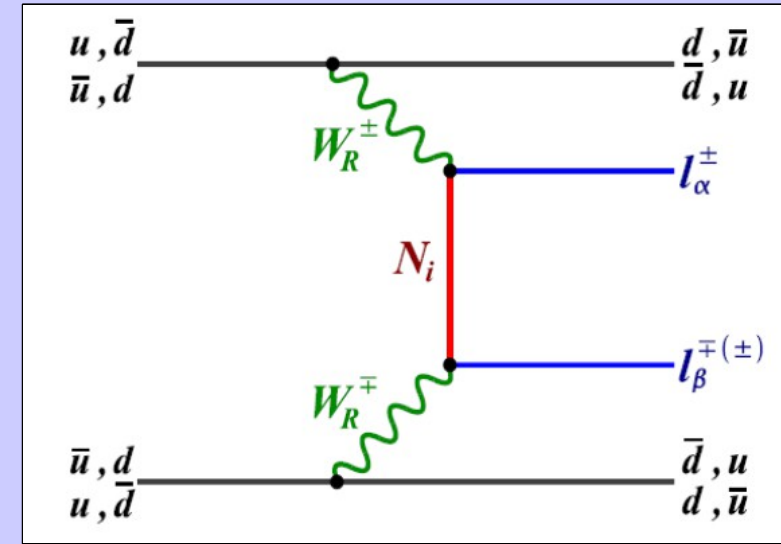
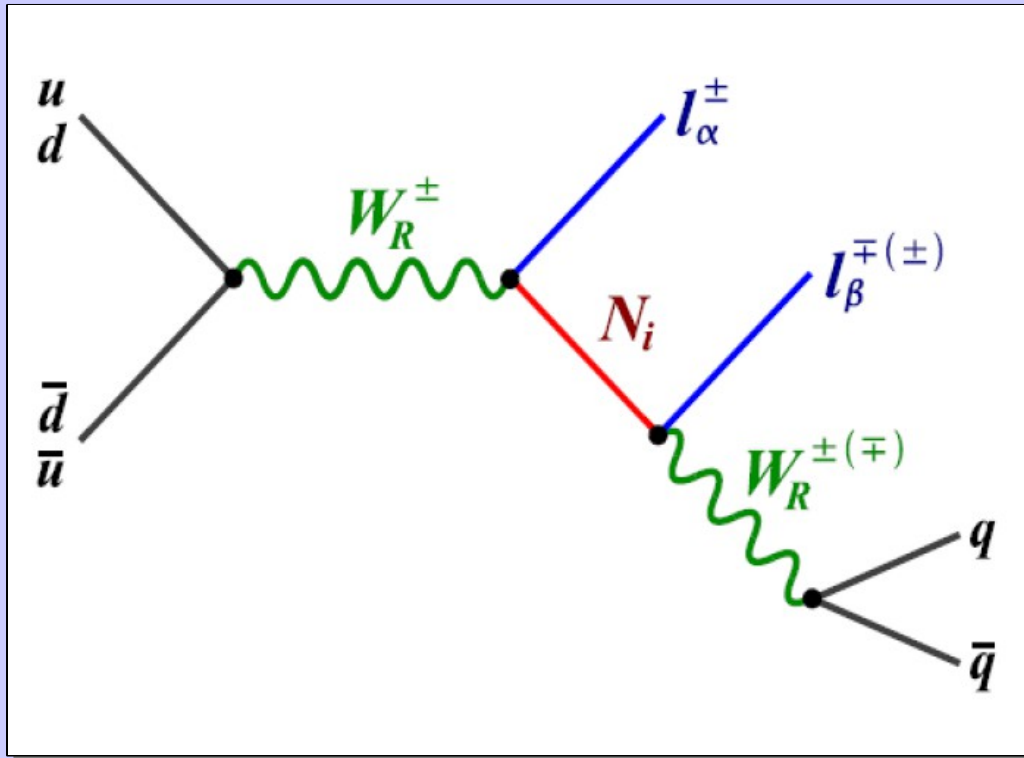
- $\mu \rightarrow eee$ strongly enhanced due to tree level contribution

$$Br(\mu \rightarrow eee) \approx 10^2 \times R(\mu \rightarrow e)$$

- BSM Flavour Problem**
Small mixing and / or mass differences required

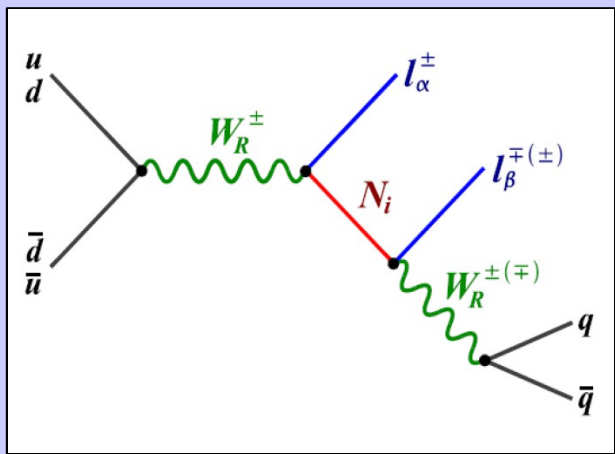


Right-handed Neutrino Production at the LHC



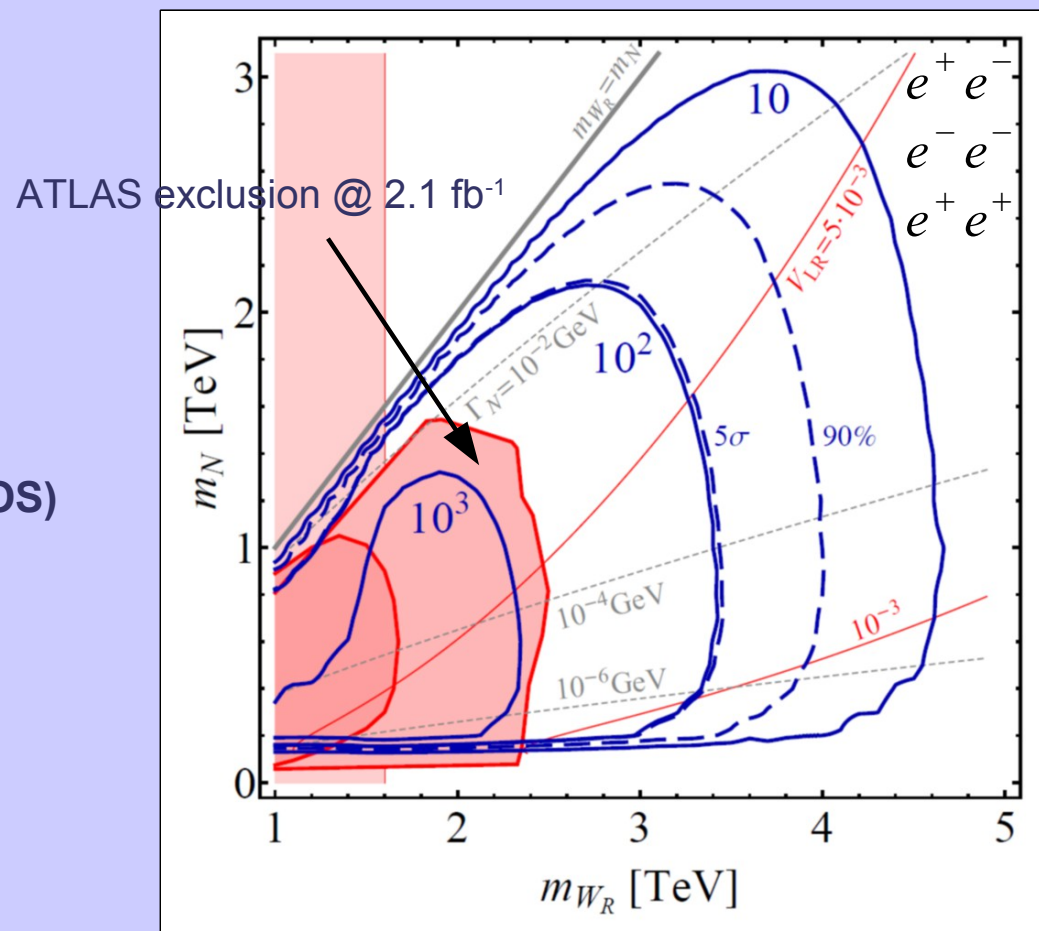
Single Neutrino Production

Sensitivity Reach



- Monte Carlo Simulation (PROTOS)
- Background $t\bar{t}$, Z + jets (Pythia, Alpgen)
- Fast Detector Simulation (AcerDET)
- Selection Criteria

number of jets	$N_j \geq 2$
number of isolated leptons	$N_\ell = 2$
invariant dilepton mass	$m_{\ell\ell} > 300 \text{ GeV}$
total invariant mass	$m_{\ell\ell jj} > 1.5 \text{ TeV}$

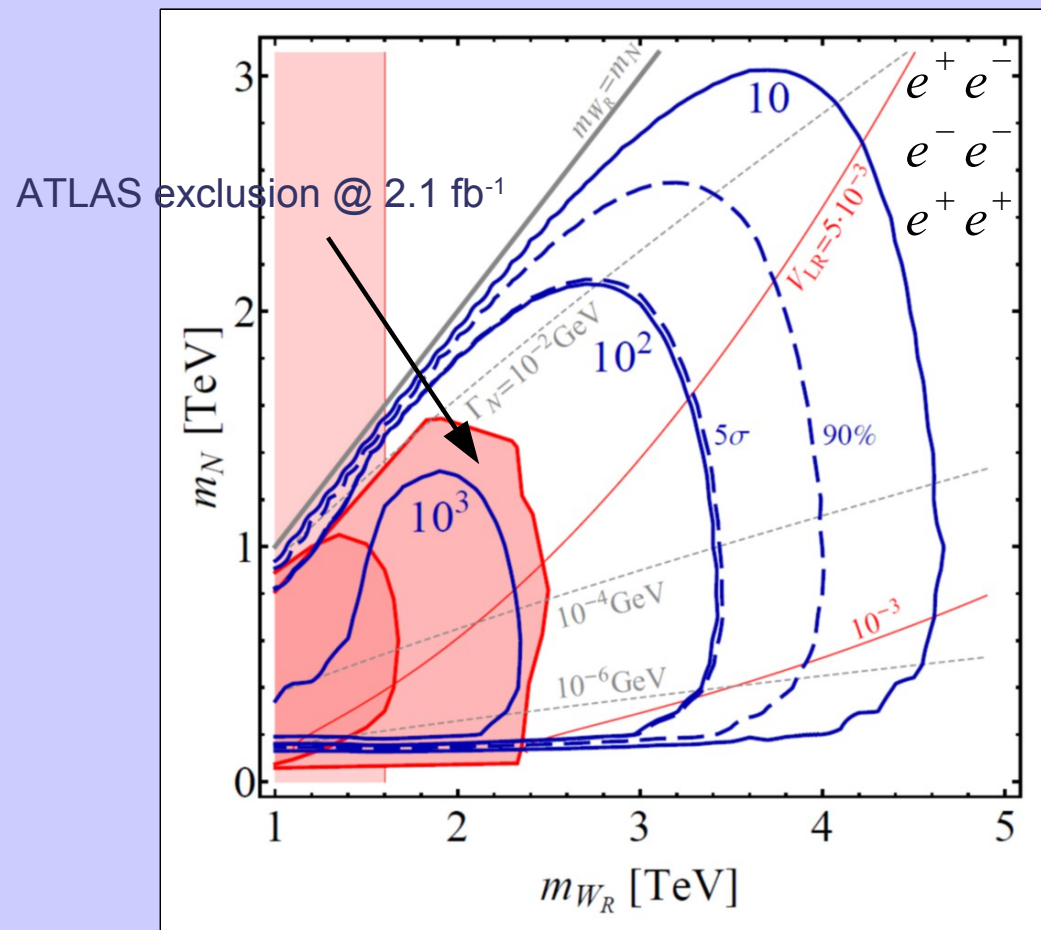
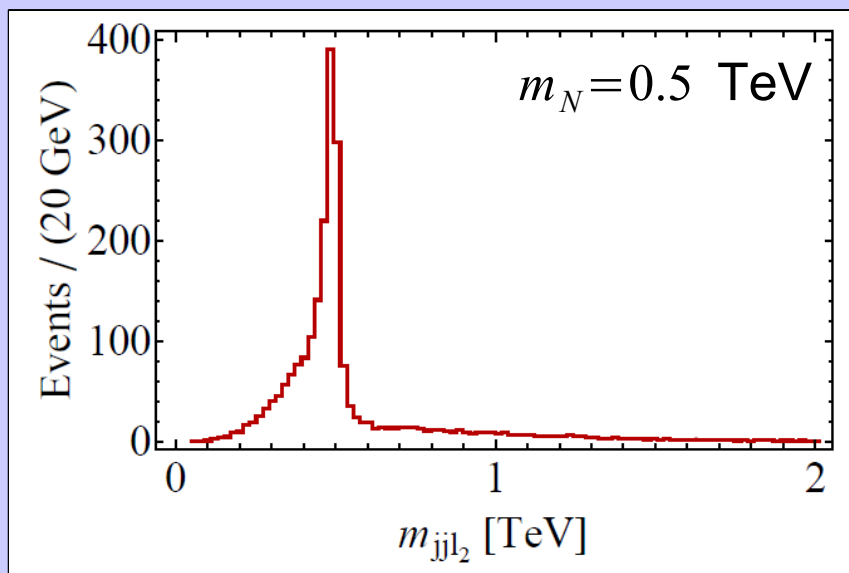
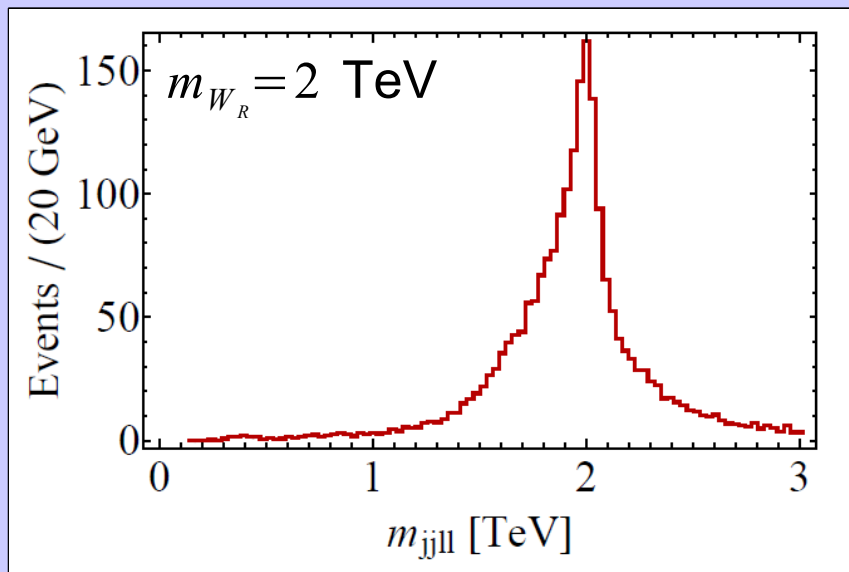


Opposite Sign + Same Sign Leptons
LHC reach @ 14 TeV, 30 fb⁻¹

Single Neutrino Production

Sensitivity Reach

- Reconstruction of W_R and N

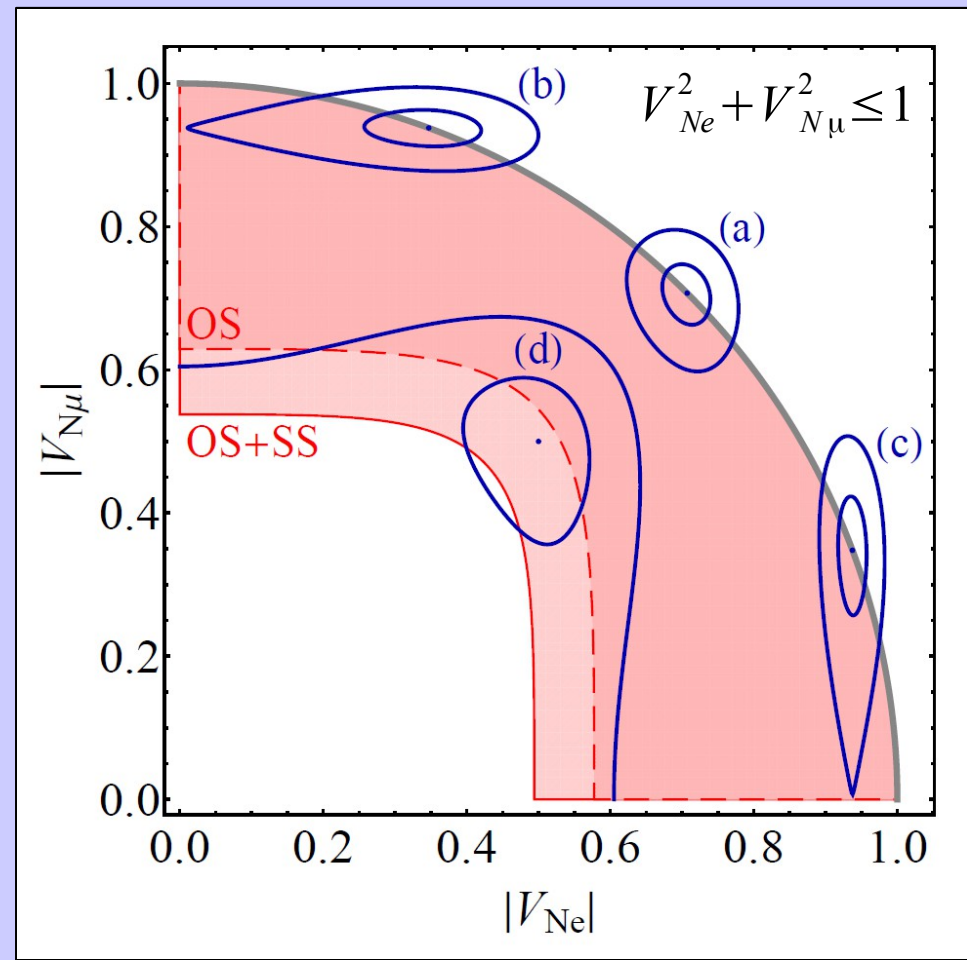
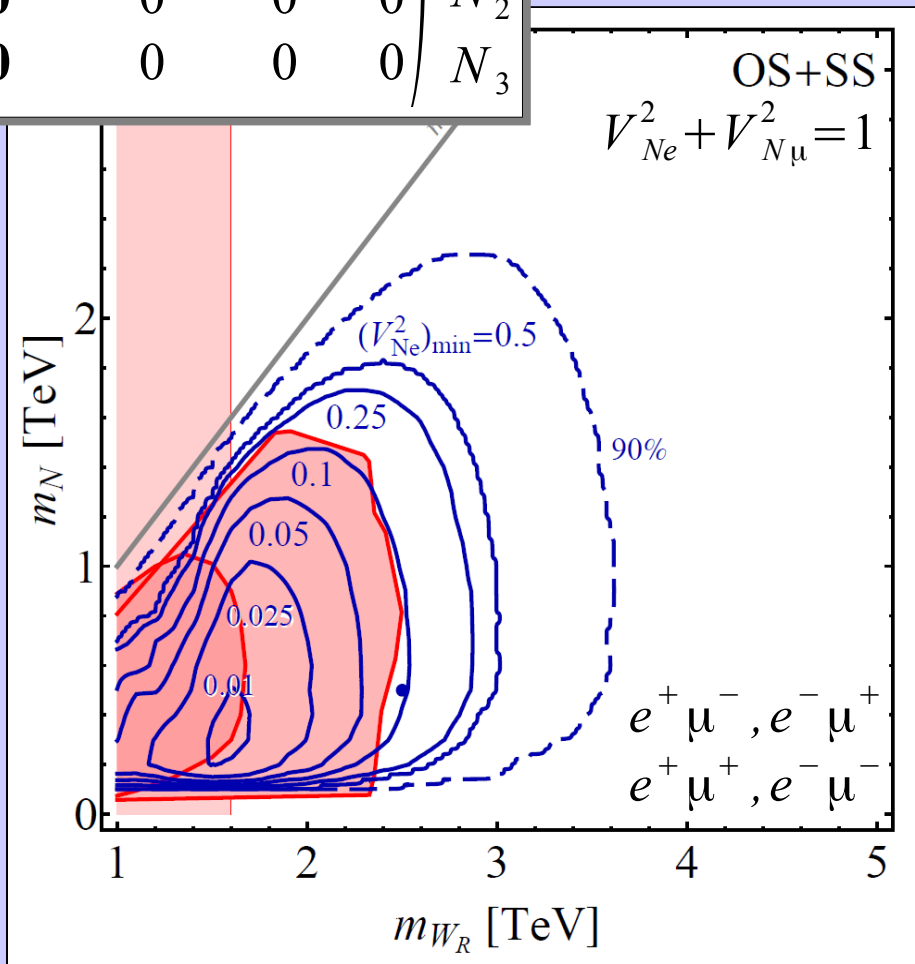


Opposite Sign + Same Sign Leptons
LHC reach @ 14 TeV, 30 fb^{-1}

Single Neutrino Production

General e- μ Mixing

$$\begin{pmatrix} l_{jL} & e_R & \mu_R & \tau_R \\ \mathbf{0} & V_{Ne} & V_{N\mu} & \mathbf{0} \\ \mathbf{0} & \mathbf{0} & \mathbf{0} & \mathbf{0} \\ \mathbf{0} & \mathbf{0} & \mathbf{0} & \mathbf{0} \end{pmatrix} \begin{matrix} \nu_i \\ N_1 \\ N_2 \\ N_3 \end{matrix}$$



$$m_{W_R} = 2.5 \text{ TeV}, m_N = 0.5 \text{ TeV}$$

LHC reach @ 14 TeV, 30 fb⁻¹

Two Neutrino Oscillations

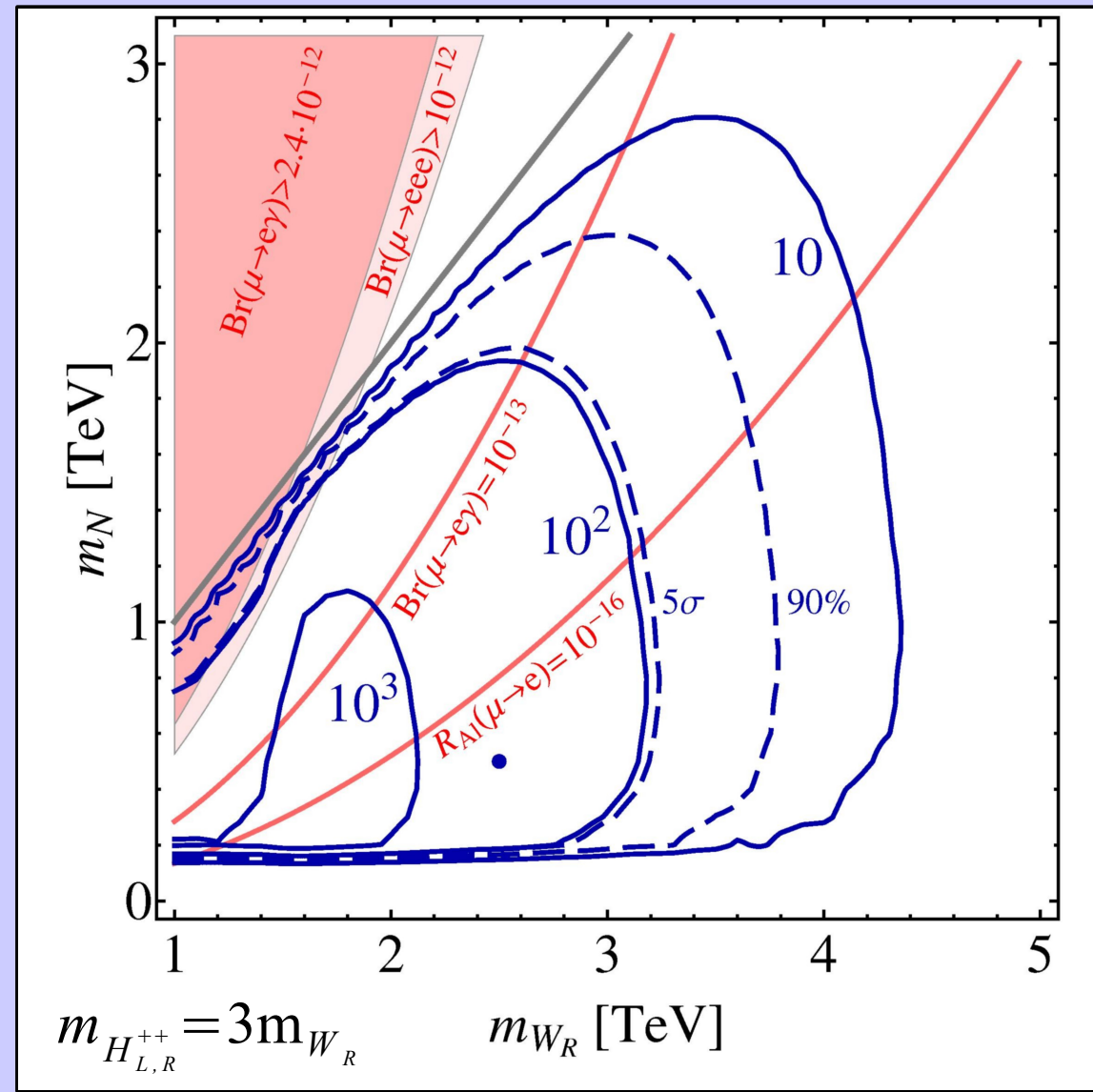
Maximal Lepton Flavour Violation

- Two neutrinos exchanged with maximal mixing and 1% mass splitting

$$\begin{pmatrix} l_{jL} & e_R & \mu_R & \tau_R \\ \mathbf{0} & \cos \phi & \sin \phi & 0 \\ \mathbf{0} & -\sin \phi & \cos \phi & 0 \\ \mathbf{0} & 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} \nu_i \\ N_1 \\ N_2 \\ N_3 \end{pmatrix}$$

$$\phi = \pi/4$$

- Correlation with low energy LFV processes



LHC reach @ 14 TeV, 30 fb⁻¹

Two Neutrino Oscillations

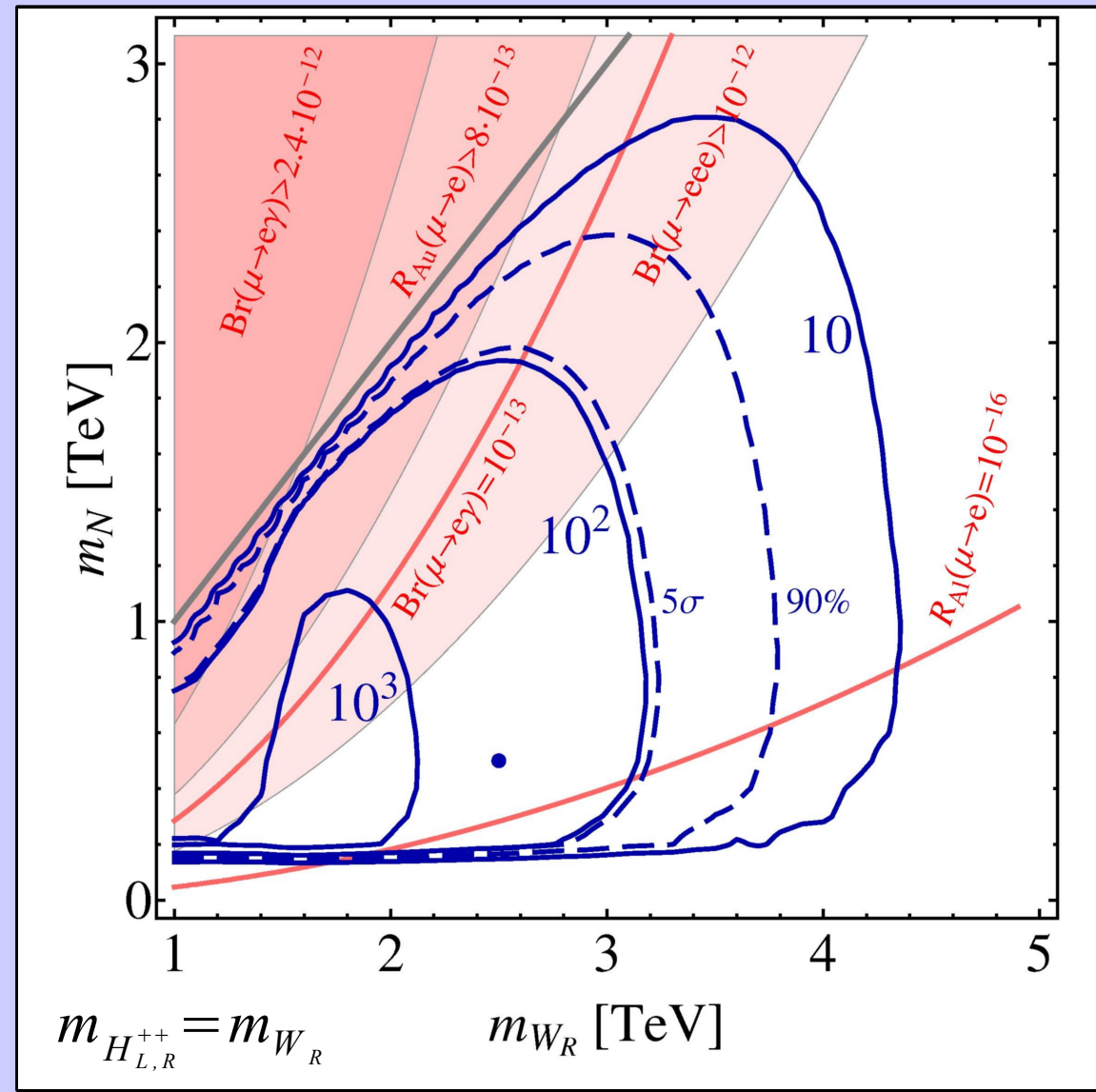
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LHC reach @ 14 TeV, 30 fb⁻¹

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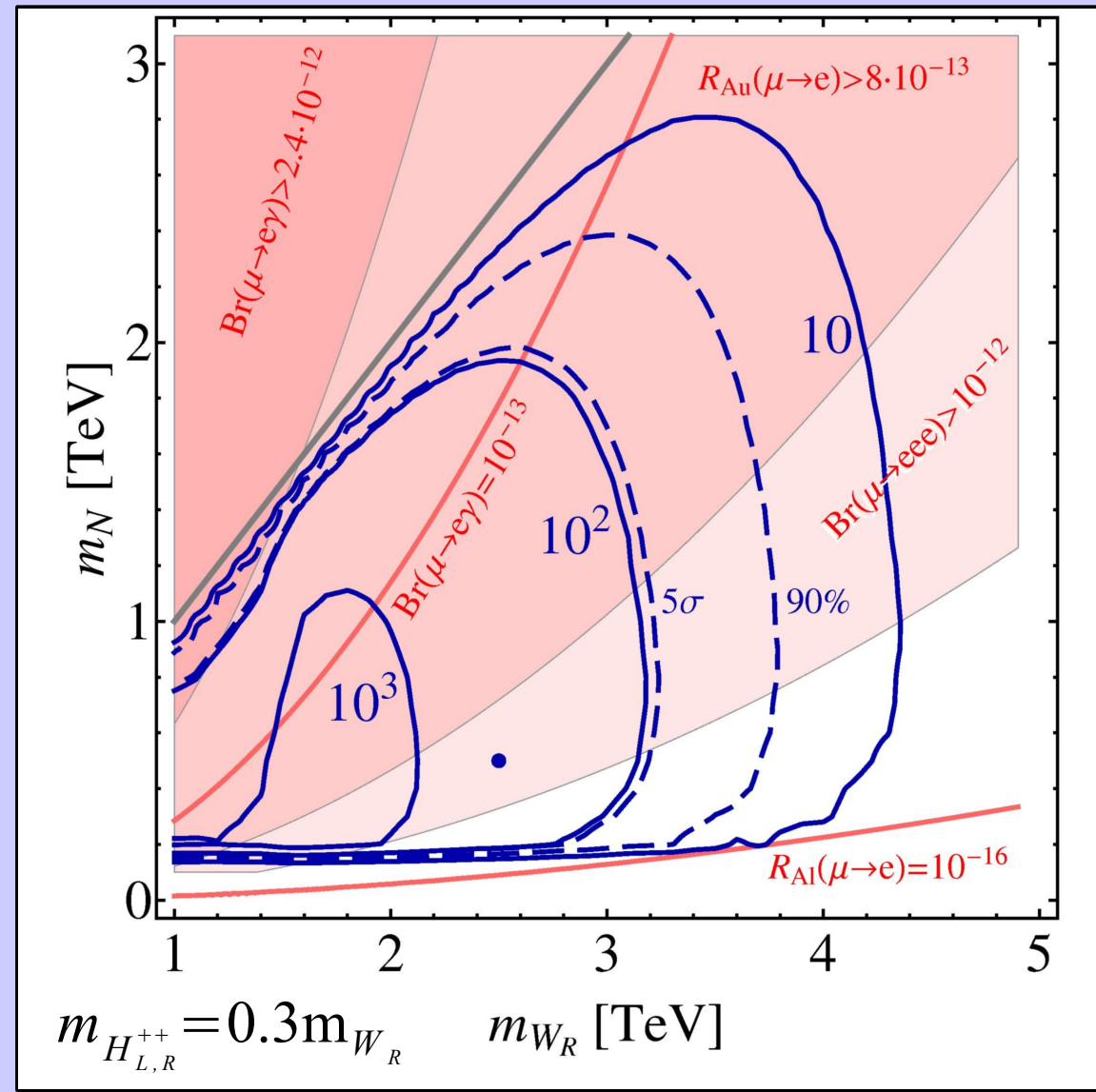
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$$\phi = \pi/4$$

- Correlation with low energy LFV processes



LHC reach @ 14 TeV, 30 fb⁻¹

Two Neutrino Oscillations

Mixing Angle and Mass Difference

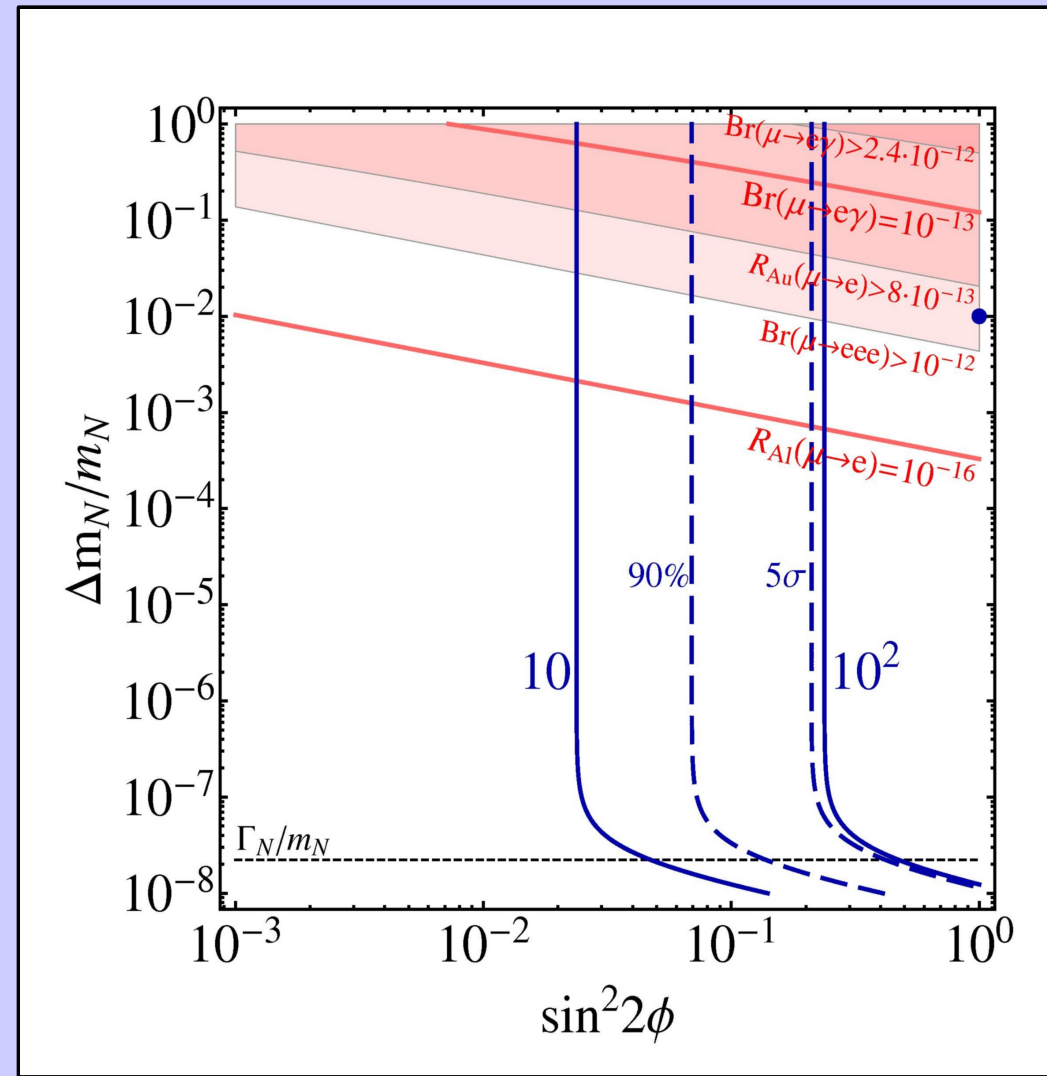
- Two neutrinos exchanged with mixing angle ϕ and mass diff. Δm_N

$$\begin{pmatrix} l_{jL} & e_R & \mu_R & \tau_R \\ \mathbf{0} & \cos \phi & \sin \phi & 0 \\ \mathbf{0} & -\sin \phi & \cos \phi & 0 \\ \mathbf{0} & 0 & 0 & 0 \end{pmatrix} U_{\text{PMNS}} \begin{pmatrix} \nu_i \\ N_1 \\ N_2 \\ N_3 \end{pmatrix}$$

suppressed as $\Delta m_N^2 / (m_N \Gamma_N)$

- Correlation with low energy LFV processes

suppressed as $\Delta m_N^2 / m_{W_R}^2$



$$m_{W_R} = 2.5 \text{ TeV}, m_N = 0.5 \text{ TeV}$$

LHC reach @ 14 TeV, 30 fb⁻¹

Single Neutrino Production

Unitary e - μ - τ Mixing

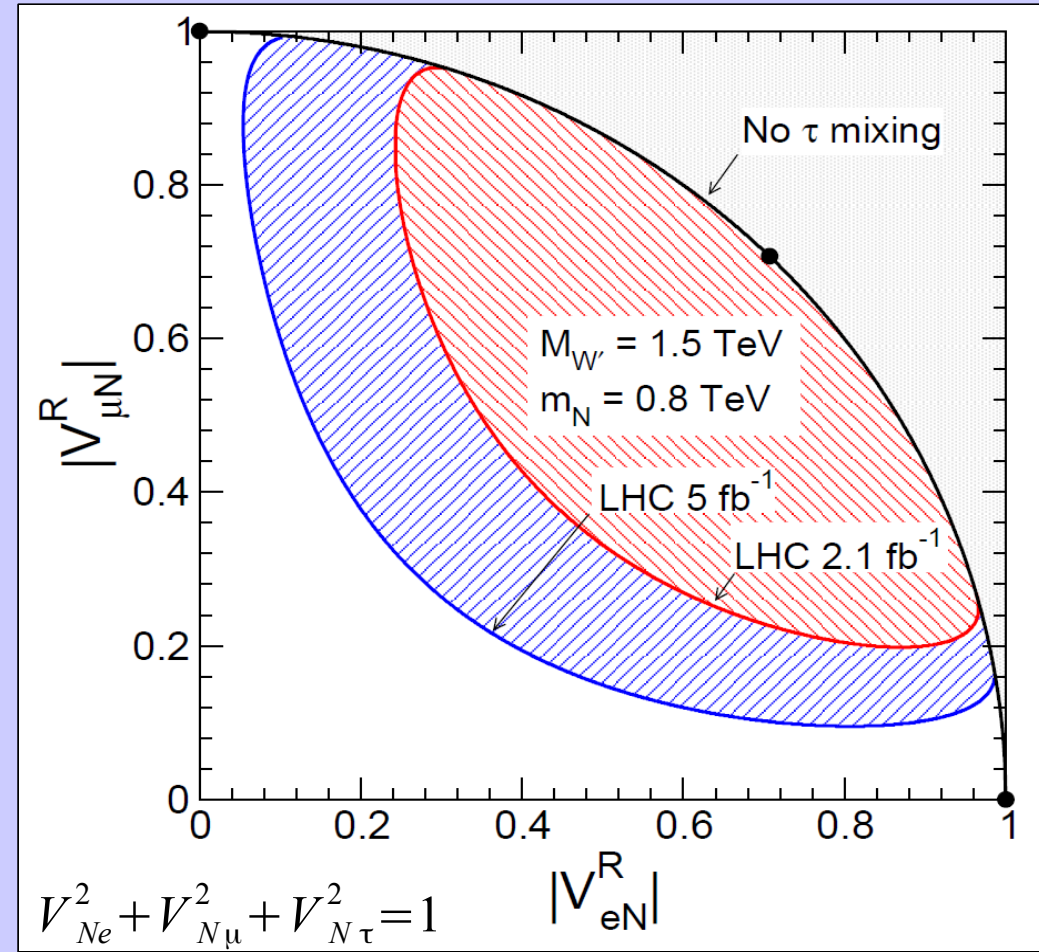
- Including coupling to taus

$$\begin{pmatrix} l_{jL} & e_R & \mu_R & \tau_R \\ \mathbf{U}_{\text{PMNS}} & \mathbf{0} & \mathbf{0} & \mathbf{0} \\ \mathbf{0} & V_{Ne} & V_{N\mu} & V_{N\tau} \\ \mathbf{0} & 0 & 0 & 0 \\ \mathbf{0} & 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} \nu_i \\ N_1 \\ N_2 \\ N_3 \end{pmatrix}$$

- Tau reconstruction efficiency reduced by

$$Br(\tau \rightarrow e(\mu) \nu \bar{\nu}) \approx 1/3$$

- Highly boosted secondary leptons
- No cut on missing p_T



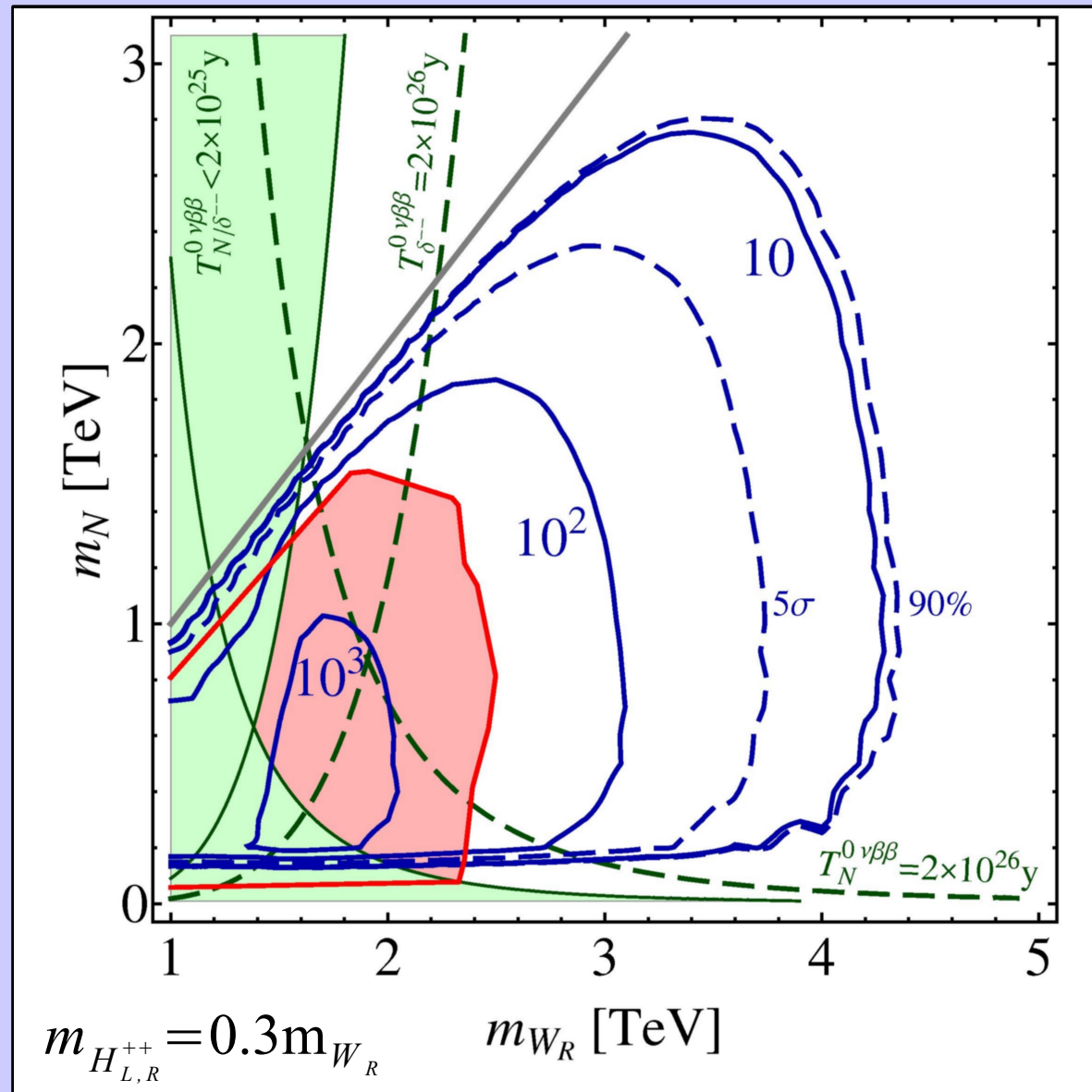
Comparison to $0\nu\beta\beta$

- Consider contributions to $0\nu\beta\beta$ from triplet Higgs

$$\frac{M_{W_L}^4}{M_{W_R}^4} \frac{m_p}{M_{\Delta_R^{--}}} \sum_i (U_{ei}^{RR})^2 M_{N_i}$$

- and heavy neutrinos

$$\frac{M_{W_L}^4}{M_{W_R}^4} \sum_i \frac{(U_{ei}^{RR})^2}{M_{N_i}}$$



LHC reach @ 14 TeV, 30 fb^{-1}

Conclusion

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Strong experimental program to probe absolute mass

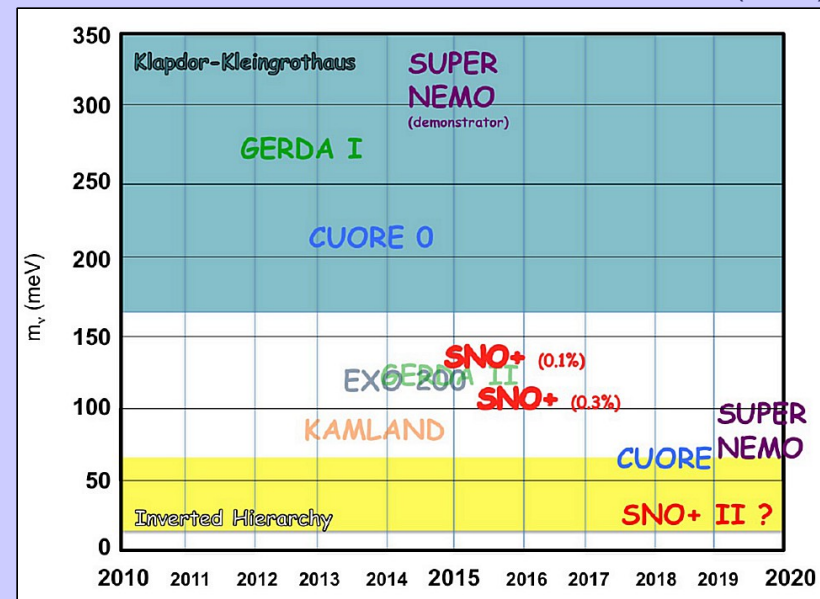
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Lefeuvre (2011)



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- **Rich phenomenology in models of neutrino mass generation**
 - Charged lepton flavour violation
 - LFV and LNV processes at the LHC
 - Connection to Leptogenesis?

Lefeuvre (2011)

