

Running of Neutrino Mass Parameters

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mostly based on based on work
in collaboration with S. Antusch, J. Kersten, M. Lindner, M. Ratz [JHEP 0503:024]
and part of my diploma thesis [hep-ph/0703...]

Why do we need the renormalization group?

Measurements by
low-energy experiments



Why do we need the renormalization group?

Extra Dimensions

Predictions from
high-energy theories

SU(5)

SO(10)

String theory



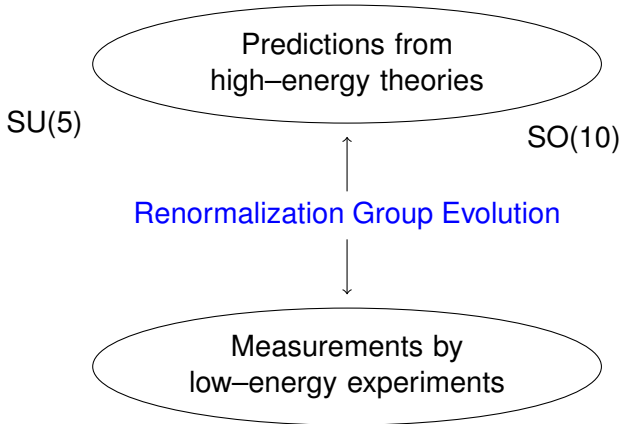
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Outline

1 Preliminaries

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Seesaw model

neutrino mass matrix

$$\begin{pmatrix} \nu \\ \bar{\nu} \end{pmatrix}^T \begin{pmatrix} 0 & m_D^T \\ m_D & M_R \end{pmatrix} \begin{pmatrix} \nu \\ \bar{\nu} \end{pmatrix}$$

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RH neutrinos decouple and give mass $(m_\nu)_{ij} = -\frac{v^2}{2} \frac{(Y_\nu)_{ki}(Y_\nu)_{kj}}{M_k}$ to light neutrinos by the seesaw mechanism [Minkowski; Yanagida; Glashow; Gell-Mann, Ramond, Slansky; Mohapatra, Senjanovic]

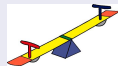


Seesaw model

neutrino mass matrix

$$\begin{pmatrix} \nu \\ \bar{\nu} \end{pmatrix}^T \begin{pmatrix} m_L & m_D^T \\ m_D & M_R \end{pmatrix} \begin{pmatrix} \nu \\ \bar{\nu} \end{pmatrix} \Rightarrow m_\nu = m_L - \frac{v^2}{2} Y_\nu^T M^{-1} Y_\nu$$

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Higgs triplet $\Delta \sim (\mathbf{3}, \mathbf{1})$:  $\Rightarrow m_L = \frac{v^2}{2} \frac{\Lambda_6}{M_\Delta^2} Y_\Delta$

[Maqq, Wetterich; Lazaridis, Shafi, Wetterich; Mohapatra, Senjanovic]

General structure of β -function

1 loop β -function:

$$P = P(Y_e, Y_\nu, Y_\Delta)$$

$$\alpha = \alpha(g_1, g_2, \text{Tr } Y^\dagger Y, \Lambda_i)$$

$$16\pi^2 \mu \frac{dm_\nu}{d\mu} = 16\pi^2 \beta_{m_\nu} = m_\nu P + P^T m_\nu + \alpha m_\nu$$

P.H. Chankowski, Z. Pluciennik (1993)

K.S. Babu, C.N. Leung, J. Pantaleone (1993)

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$$P = C_e Y_e^\dagger Y_e + C_\nu Y_\nu^\dagger Y_\nu + C_\Delta Y_\Delta^\dagger Y_\Delta$$

m_ν	SM			MSSM		
	C_e	C_ν	C_Δ	C_e	C_ν	C_Δ
$-\frac{v^2}{4} \kappa$	$-\frac{3}{2}$	$\frac{1}{2}$	$-\frac{3}{2}$	1	1	$\frac{3}{2}$
$-\frac{v^2}{2} Y_\nu^T M^{-1} Y_\nu$	$-\frac{3}{2}$	$\frac{1}{2}$	$-\frac{3}{2}$	1	1	$\frac{3}{2}$
$\frac{v^2}{2} \frac{\Lambda_6}{M_\Delta^2} Y_\Delta$	$\frac{1}{2}$	$\frac{1}{2}$	$-\frac{3}{2}$	1	1	$\frac{3}{2}$

Mixing parameters

MNS mixing matrix

$$U = \text{diag}(e^{i\delta_e}, e^{i\delta_\mu}, e^{i\delta_\tau}) \cdot V \cdot \text{diag}(e^{-i\varphi_1/2}, e^{-i\varphi_2/2}, 1)$$

where ($s_{ij} = \sin \theta_{ij}$, $c_{ij} = \cos \theta_{ij}$)

$$V = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -c_{23}s_{12} - s_{23}s_{13}c_{12}e^{i\delta} & c_{23}c_{12} - s_{23}s_{13}s_{12}e^{i\delta} & s_{23}c_{13} \\ s_{23}s_{12} - c_{23}s_{13}c_{12}e^{i\delta} & -s_{23}c_{12} - c_{23}s_{13}s_{12}e^{i\delta} & c_{23}c_{13} \end{pmatrix}$$

Experimental data

T. Schwetz [hep-ph/0606060]

Parameter	Best-fit	Allowed range (3σ)
$\sin^2 \theta_{12}$ [$^\circ$]	0.30	0.24 ... 0.40
$\sin^2 \theta_{23}$ [$^\circ$]	0.50	0.34 ... 0.68
$\sin^2 \theta_{13}$ [$^\circ$]	0.000	≤ 0.041
Δm_{21}^2 [10^{-5}eV^2]	7.9	7.1 ... 8.9
$ \Delta m_{31}^2 $ [10^{-3}eV^2]	2.5	1.9 ... 3.2

Effective field theory

$$P = C_e Y_e^\dagger Y_e$$

$$C_e^{\text{SM}} = -\frac{3}{2}, C_e^{\text{MSSM}} = 1$$

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General structure (also applicable for phases):

Renormalization scale

$$\mu \frac{d\theta_{ij}}{d\mu} \propto \frac{f(m_l, \delta, \varphi_1, \varphi_2)}{m_j^2 - m_i^2} \times F^{(ij)}(y_k, \theta_{lm})$$

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- The more **degenerate** the **spectrum** is, the larger is the running.
- **Phases** can damp the running.

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- $F^{(ij)} \propto$ Yukawa coupling
 $\Rightarrow \tan \beta$ dependence
- $F^{(ij)} = 0$ for **zero mixing**
 (fixed point)

S. Antusch, J. Kersten, M. Lindner, M. Ratz [hep-ph/0305273]



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Full theory

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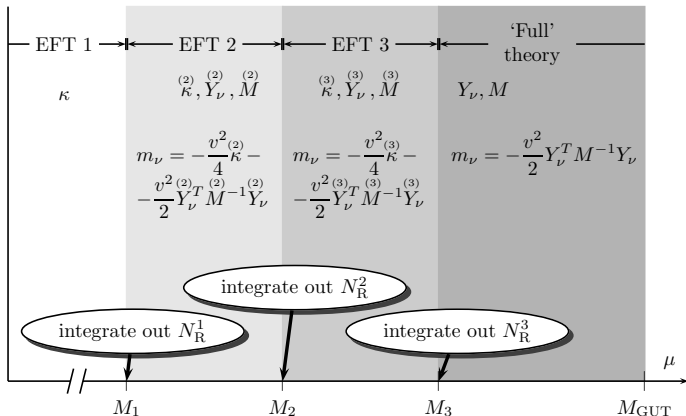
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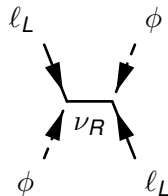
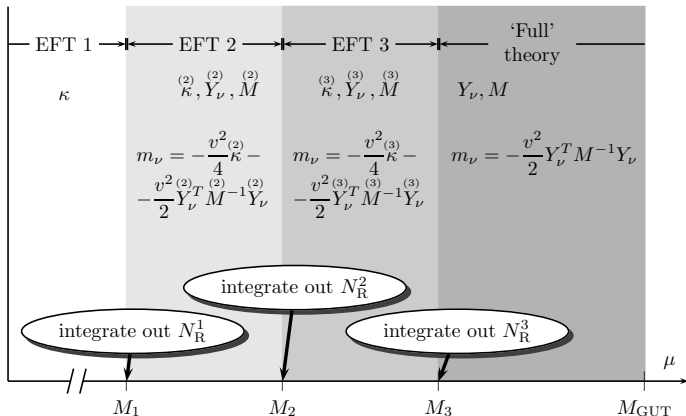
$$\mu \frac{d\theta_{ij}}{d\mu} \propto \sum_{kl} \frac{f_{kl}(m_l, \delta, \varphi_1, \varphi_2)}{m_j^2 - m_i^2} \times F_{kl}^{(ij)}(\theta_{lm}) \times P_{kl}(Y_e, Y_\nu) \\ + F_e^{(ij)}(Y_e, Y_\nu)$$

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- off-diagonal terms \rightarrow mixing can be generated
- GUT: Y_ν strongly hierarchical $\rightarrow P_{33}$ dominates

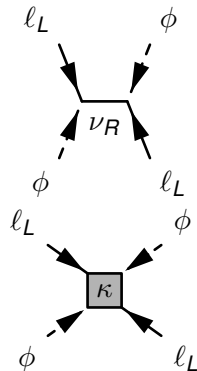
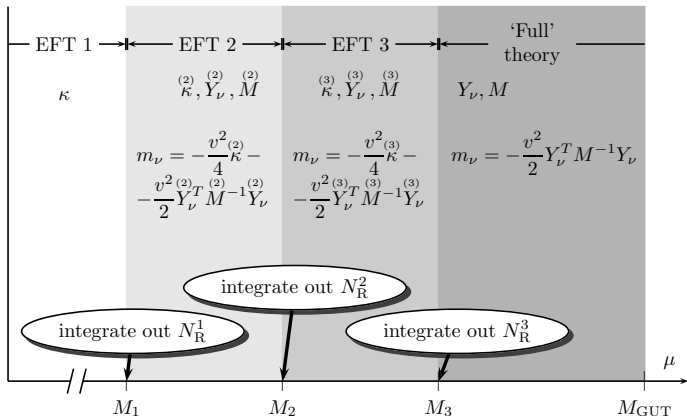
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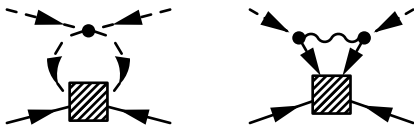
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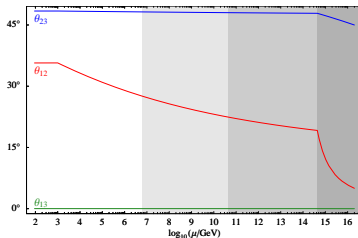
- rescaling of right-handed neutrino masses

M. Lindner, MS, A. Smirnov, hep-ph/0505162

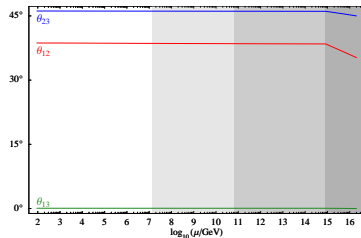
$$m_\nu = Z_{\text{ext}}^T Y_\nu^T X M^{-1} Y_\nu Z_{\text{ext}}$$

additional vertex corrections

Example for RG evolution



(a) $\tan \beta = 20$, $m_1 = 0.05 \text{ eV}$

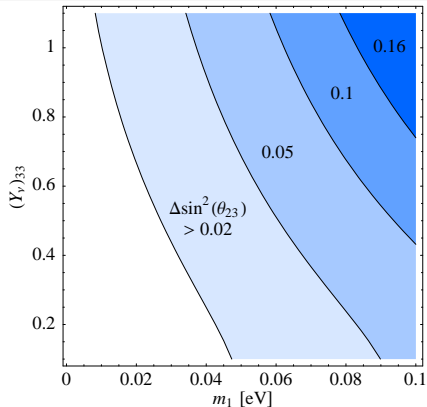


(b) $\tan \beta = 10$, $m_1 = 0.015 \text{ eV}$

$$\delta = \varphi_1 = \varphi_2 = 0$$

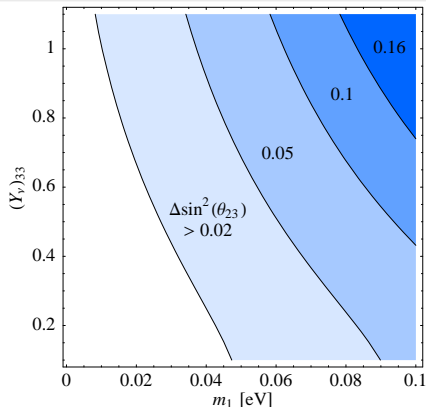
θ_{23}

Current	Beams	T2K+NuMI	JPARC-HK	NuFact-II
0.16	0.1	0.050	0.020	0.055

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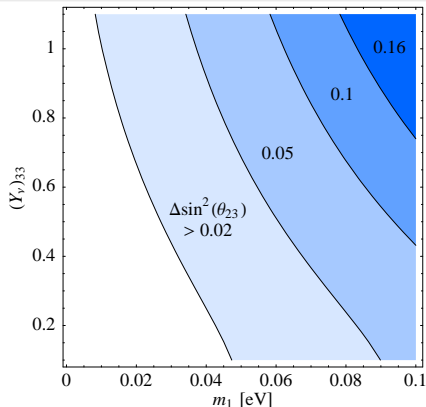
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$$\bullet \quad |0.5 - \sin^2 \theta_{23}| \leq 0.16$$

T. Schwetz [hep-ph/0606060]

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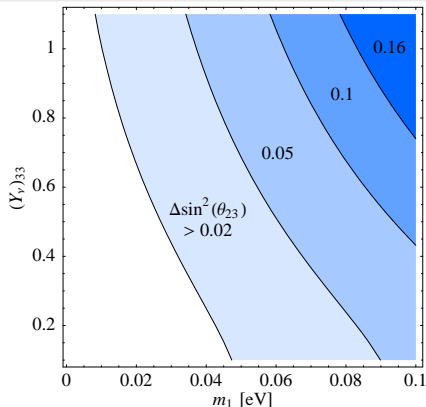
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T. Schwetz [hep-ph/0606060]
- small deviations from maximal mixing

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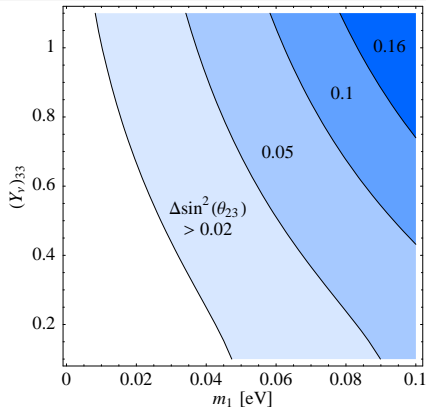
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T. Schwetz [hep-ph/0606060]

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- suppression by phases possible

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convenient basis: $Y_{\Delta} = \text{diag}(y_1, y_2, y_3)$ diagonal

$$16\pi^2 \dot{y}_i = 2C_{\Delta} y_i^3 + \alpha y_i$$

$$16\pi^2 \left(\dot{Y}_e \right)_{ij} = (Y_e)_{ij} D_{\Delta} y_j^2 + \alpha_e (Y_e)_{ij}$$

Chao,Zhang [hep-ph/0611323]

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Chao,Zhang [hep-ph/0611323]

- Neutrino mass matrix stays diagonal
- Running of angles and phases described by charged leptons

General structure of RG equations I

$$\begin{aligned}
 16\pi^2 \dot{\theta}_{12} &= -\frac{D_\Delta}{2} y_{21}^2 \sin 2\theta_{12} \\
 16\pi^2 \dot{\theta}_{13} &= -\frac{D_\Delta}{4} \left[y_{31}^2 + y_{32}^2 + y_{21}^2 \cos 2\theta_{12} \right] \sin 2\theta_{13} \\
 16\pi^2 \dot{\theta}_{23} &= -\frac{D_\Delta}{2} \left[\left(y_{32}^2 c_{12}^2 + y_{31}^2 s_{12}^2 \right) \sin 2\theta_{23} \right. \\
 &\quad \left. + y_{21}^2 \cos \delta \sin 2\theta_{12} c_{23}^2 \theta_{13} \right] + \mathcal{O}(\theta_{13}^2)
 \end{aligned}$$

$$y_{ji}^2 = y_j^2 - y_i^2, \quad y_i \in \{y_1, y_2, y_3, y_e, y_\mu, y_\tau\}$$

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$$16\pi^2 \dot{\theta}_{23} = -\frac{D_\Delta}{2} \left[\left(y_{32}^2 c_{12}^2 + y_{31}^2 s_{12}^2 \right) \sin 2\theta_{23} \right. \\ \left. + y_{21}^2 \cos \delta \sin 2\theta_{12} c_{23}^2 \theta_{13} \right] + \mathcal{O}(\theta_{13}^2)$$

$$16\pi^2 \dot{\theta}_{ij} \approx -\frac{D_\Delta}{2} (y_j^2 - y_i^2) \sin 2\theta_{ij}$$

$$y_{ji}^2 = y_j^2 - y_i^2, y_i \in \{y_1, y_2, y_3, y_e, y_\mu, y_\tau\}$$

General structure of RG equations II

$$16\pi^2 \dot{\theta}_{12} = -\frac{1}{2} \left[D_\Delta y_{21}^2 + C_e y_\tau^2 \frac{y_2 + y_1}{y_2 - y_1} \sin \theta_{23} \right] \sin 2\theta_{12} + \mathcal{O}(\theta_{13})$$

$$16\pi^2 \dot{\theta}_{13} = -\frac{C_e}{2} \frac{(y_2 - y_1)y_3}{(y_3 - y_1)(y_3 - y_2)} y_\tau^2 \cos \delta \sin 2\theta_{12} \sin 2\theta_{23} + \mathcal{O}(\theta_{13})$$

$$16\pi^2 \dot{\theta}_{23} = -\frac{1}{2} \left[D_\Delta \left(y_3^2 - y_1^2 \sin^2 \theta_{12} - y_2^2 \cos^2 \theta_{12} \right) + C_e \frac{y_3^2 - y_1 y_2 + (y_2 - y_1) \cos 2\theta_{12}}{(y_3 - y_2)(y_3 - y_1)} y_\tau^2 \right] \sin 2\theta_{23} + \mathcal{O}(\theta_{13})$$

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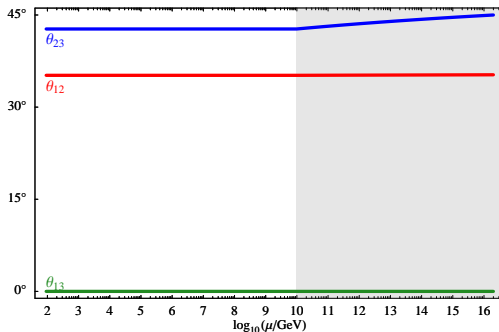
$$16\pi^2 \dot{\theta}_{23} = -\frac{1}{2} \left[D_\Delta \left(y_3^2 - y_1^2 \sin^2 \theta_{12} - y_2^2 \cos^2 \theta_{12} \right) + C_e \frac{y_3^2 - y_1 y_2 + (y_2 - y_1) \cos 2\theta_{12}}{(y_3 - y_2)(y_3 - y_1)} y_\tau^2 \right] \sin 2\theta_{23} + \mathcal{O}(\theta_{13})$$

$$16\pi^2 \dot{\delta} = \frac{C_e}{2} \frac{(y_2 - y_1)y_3}{(y_3 - y_1)(y_3 - y_2)} y_\tau^2 \sin \delta \sin 2\theta_{12} \sin 2\theta_{23} \theta_{13}^{-1} + \mathcal{O}(\theta_{13})$$

$$y_{ji}^2 = y_j^2 - y_i^2, y_i \in \{y_1, y_2, y_3, y_e, y_\mu, y_\tau\}$$

RG evolution of θ_{23}

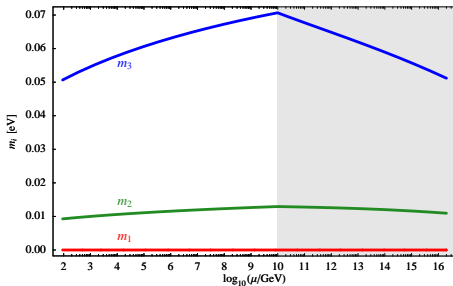
$$16\pi^2 \dot{\theta}_{23} \approx -\frac{D_\Delta}{2} \left(y_3^2 - y_1^2 \sin^2 \theta_{12} - y_2^2 \cos^2 \theta_{12} \right) \sin 2\theta_{23}$$



n.h., $m_1 = 0$ eV, $\langle \Delta \rangle \sim 0.1$ eV, $Y_\Delta \sim \mathcal{O}(0.1 - 1)$, $D_\Delta = -\frac{3}{2} \frac{c}{s}$

Masses

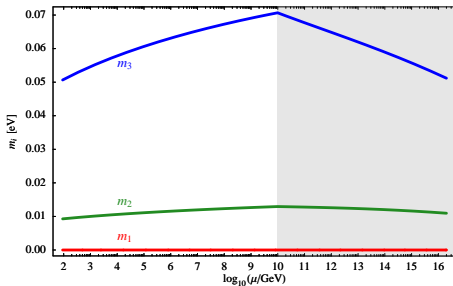
$$16\pi^2 \dot{y}_i = 2C_\Delta y_i^3 + \alpha y_i$$



n.h., $m_1 = 0 \text{ eV}$, $\langle \Delta \rangle \sim 0.1 \text{ eV}$, $Y_\Delta \sim \mathcal{O}(0.1 - 1)$, $C_\Delta = -\frac{3}{2}$, $\alpha(y_i, g_2, \Lambda_i) < 0$

Masses

$$16\pi^2 \dot{y}_i = 2C_\Delta y_i^3 + \alpha y_i$$



n.h., $m_1 = 0 \text{ eV}$, $\langle \Delta \rangle \sim 0.1 \text{ eV}$, $Y_\Delta \sim \mathcal{O}(0.1 - 1)$, $C_\Delta = -\frac{3}{2}$, $\alpha(y_i, g_2, \Lambda_i) < 0$

REAP/MixingParameterTools: <http://www.ph.tum.de/~rge>

Outline

- 1 Preliminaries
- 2 RG evolution in standard seesaw model
- 3 RG evolution in type-II seesaw model
- 4 Summary**

Summary

Standard seesaw

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- Large renormalization group effects above and between thresholds possible. → High-energy symmetries can be destroyed by RG effects.

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- RG effects become comparable to sensitivity of precision experiments.

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Type-II seesaw

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- Large renormalization group effects above and between thresholds possible. → High-energy symmetries can be destroyed by RG effects.
- RG effects become comparable to sensitivity of precision experiments.

Type-II seesaw

- No enhancement factor, but RG effect for considerable Y_{Δ}

Summary

Standard seesaw

- Large renormalization group effects above and between thresholds possible. → High-energy symmetries can be destroyed by RG effects.
- RG effects become comparable to sensitivity of precision experiments.

Type-II seesaw

- No enhancement factor, but RG effect for considerable Y_{Δ}
- RG effect proportional to mass squared difference

Summary

Standard seesaw

- Large renormalization group effects above and between thresholds possible. → High-energy symmetries can be destroyed by RG effects.
- RG effects become comparable to sensitivity of precision experiments.

Type-II seesaw

- No enhancement factor, but RG effect for considerable Y_{Δ}
- RG effect proportional to mass squared difference
- Sizable RG effect for θ_{23}