

PREDICTIONS FOR AND IMPLICATIONS OF $0\nu\beta\beta$



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LAUNCH, 22/03/07

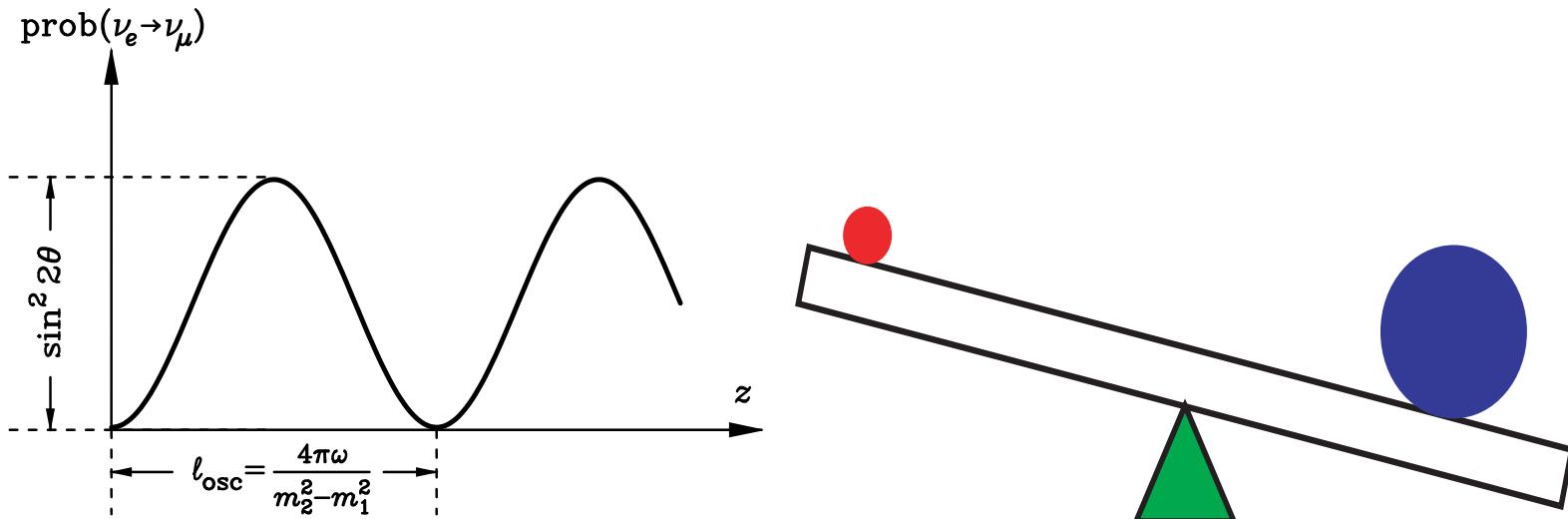


- Neutrino Mixing and $0\nu\beta\beta$
- $|m_{ee}|$ and NH vs. IH
- $|m_{ee}|$ and $|U_{e3}|$
- What's more to $0\nu\beta\beta$
- If we don't observe $0\nu\beta\beta$: vanishing effective mass
- If we observe $0\nu\beta\beta$...

STATUS AND GOAL OF NEUTRINO PHYSICS

understand form and origin of
fundamental object in low energy Lagrangian:

$$\mathcal{L} = \frac{1}{2} \overline{\nu_\alpha^c} (m_\nu)_{\alpha\beta} \nu_\beta + h.c. \text{ with } m_\nu = U^T \text{ diag}(m_1, m_2, m_3) U$$



$$P(\nu_e \rightarrow \nu_\mu) = \sin^2 2\theta \sin^2 \frac{\Delta m^2}{4E} L$$

Oscillations!

Masses?

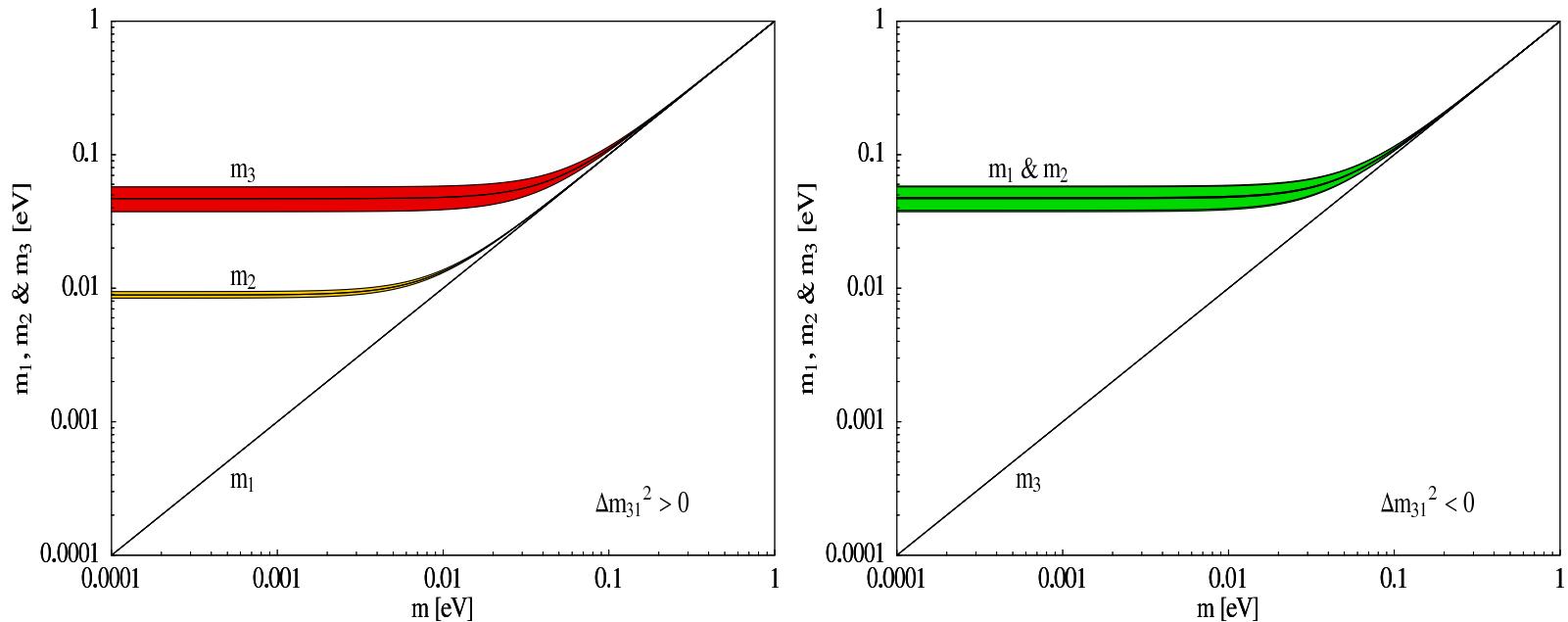
$$m_\nu = -m_D^T M_R^{-1} m_D$$

See-saw!

Lepton Number Violation?

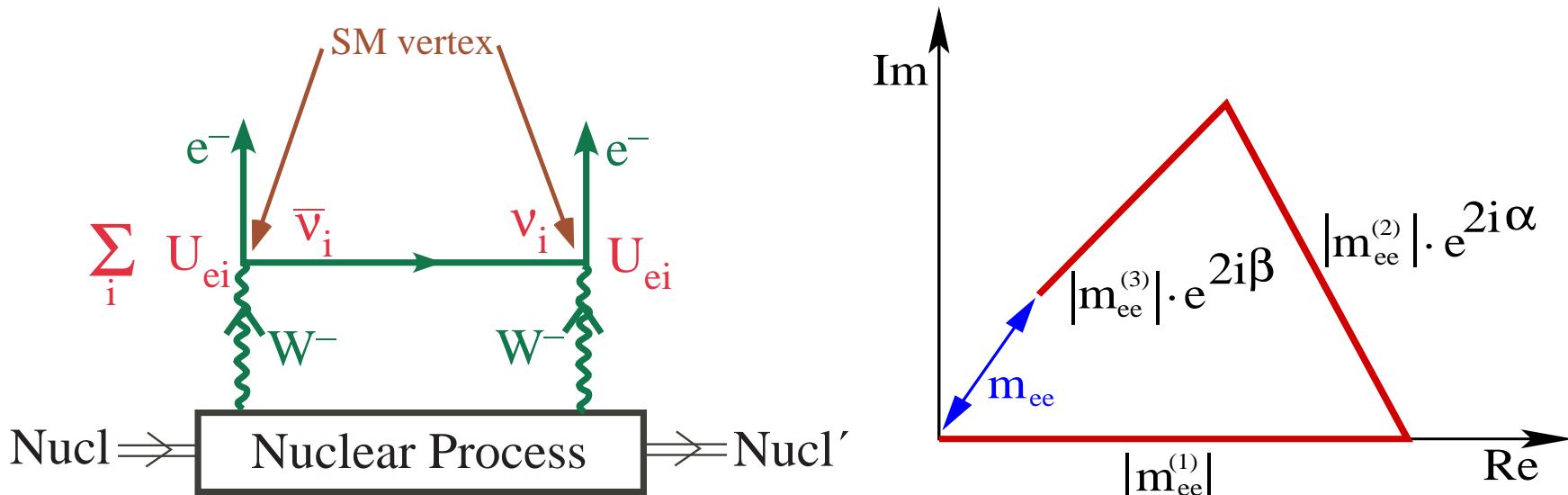
NEUTRINO MASSES

- 9 parameters in m_ν ; we only know θ_{12} and θ_{23}
- neutrino masses \leftrightarrow scale of their origin
- neutrino mass ordering \leftrightarrow form of m_ν



- $m_3^2 \simeq \Delta m_A^2 \gg m_2^2 \simeq \Delta m_\odot^2 \gg m_1^2$: normal hierarchy (NH)
- $m_2^2 \simeq |\Delta m_A^2| \simeq m_1^2 \gg m_3^2$: inverted hierarchy (IH)
- $m_3 \simeq m_2 \simeq m_1 \equiv m_0 \gg \sqrt{\Delta m_A^2}$: quasi-degeneracy (QD)

$\Delta L \neq 0$: NEUTRINOLESS DOUBLE BETA DECAY

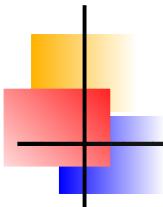


- only works when $\nu = \nu^c$ **and** $m_\nu \neq 0 \Leftrightarrow$ See-saw mechanism
- Nuclear Matrix Elements: Uncertainty $\zeta = \mathcal{O}(1)$!?

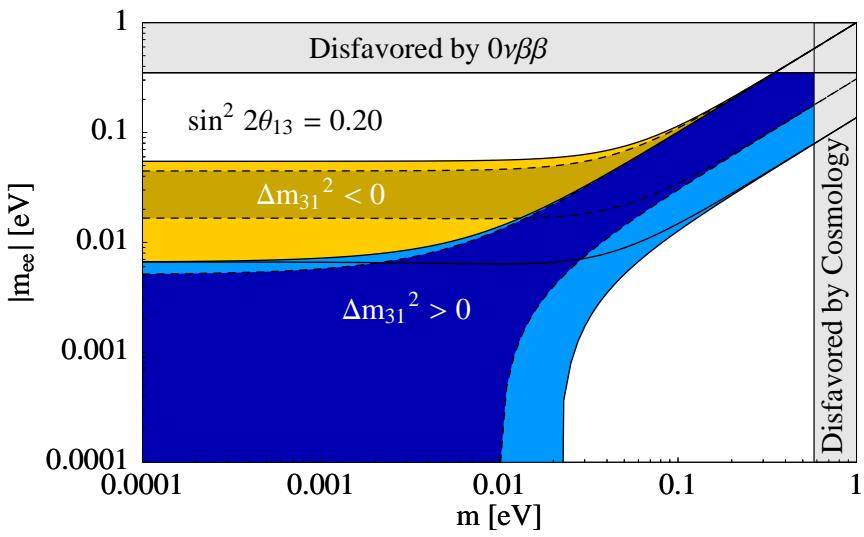
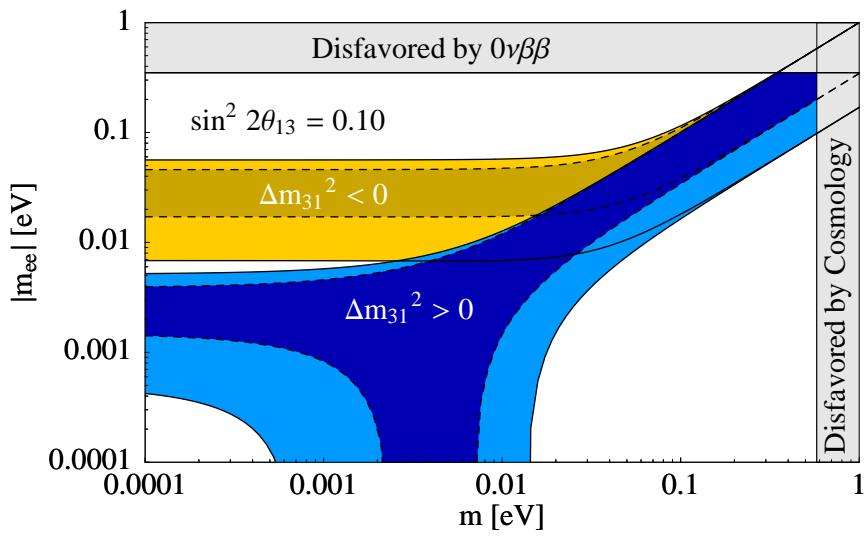
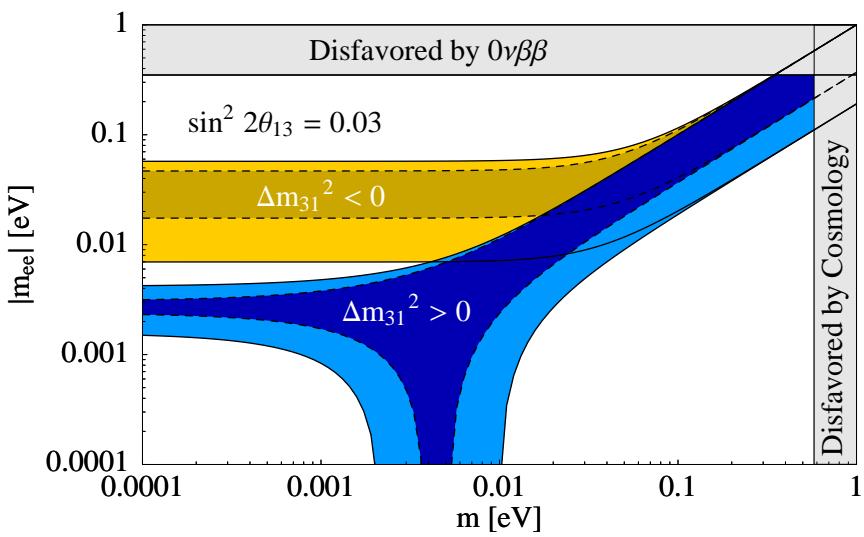
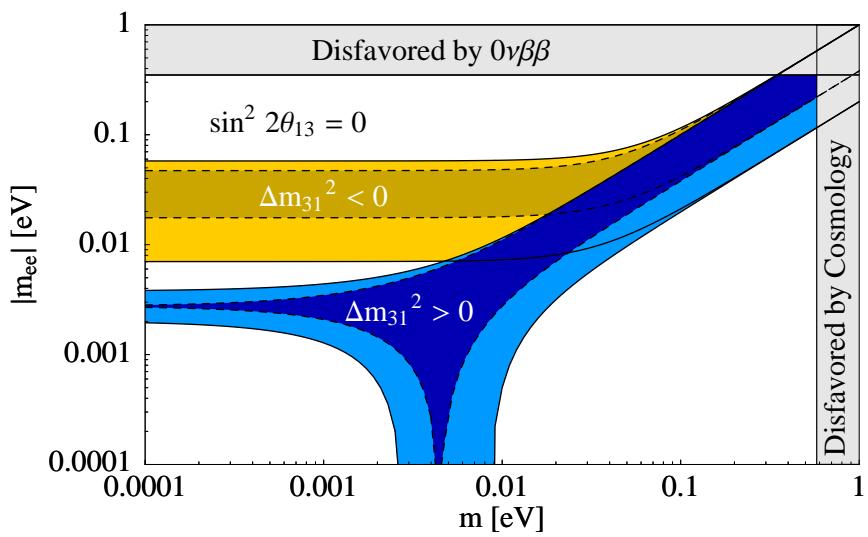
Amplitude proportional to coherent sum ("Effective mass" $|m_{ee}|$):

$$|m_{ee}| \equiv \left| \sum U_{ei}^2 m_i \right| = |c_{12}^2 c_{13}^2 m_1 + s_{12}^2 c_{13}^2 m_2 e^{2i\alpha} + s_{13}^2 m_3 e^{2i\beta}| \\ = \left| |m_{ee}^{(1)}| + |m_{ee}^{(3)}| e^{2i\alpha} + |m_{ee}^{(3)}| e^{2i\beta} \right| = f(\theta_{12}, m_i, |U_{e3}|, \text{sgn}(\Delta m_A^2), \alpha, \beta)$$

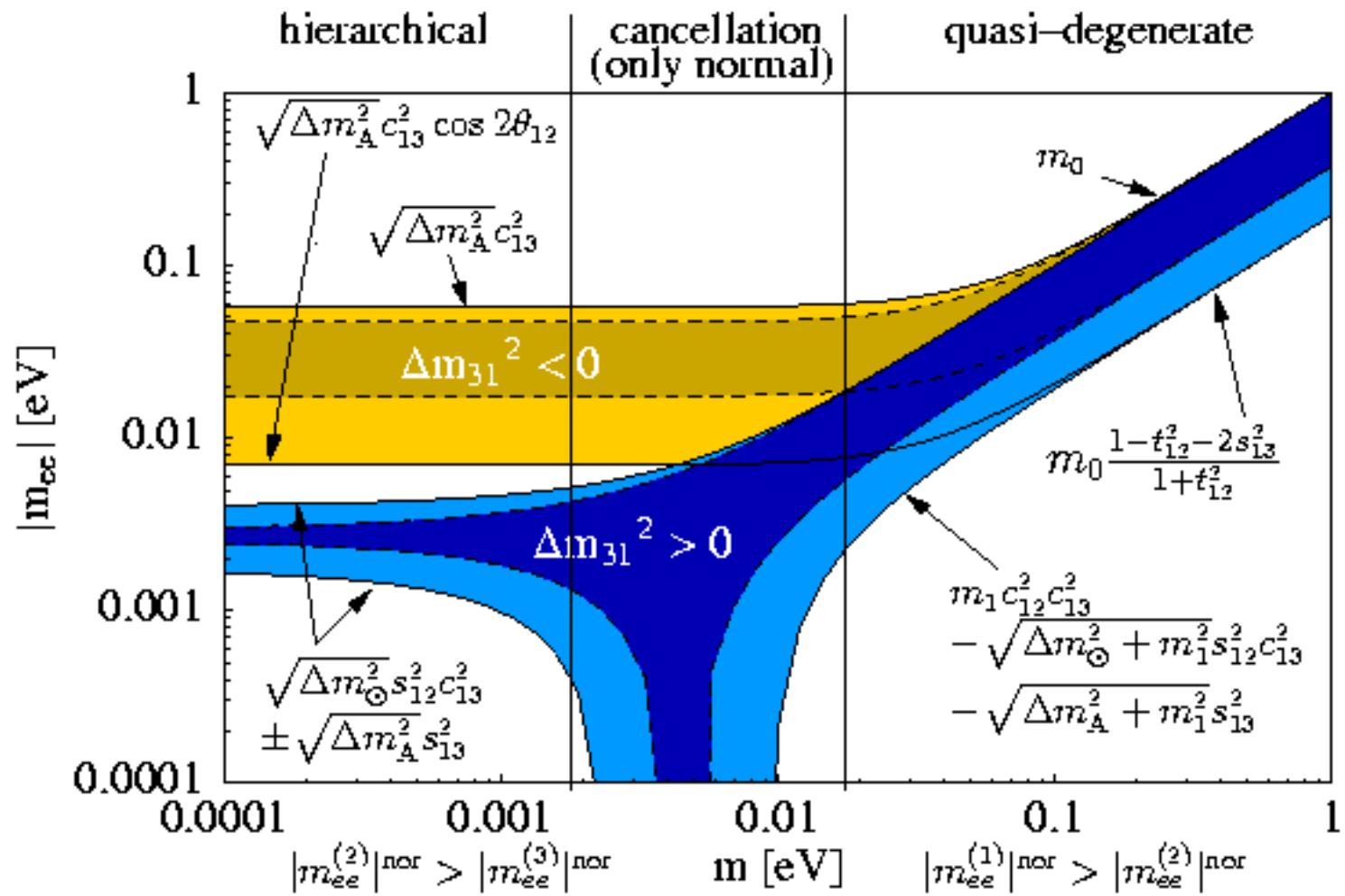
7 out of 9 parameters of m_ν ...



Kim, 1996; Minakata & Yasuda, 1996; Hirsch & Klapdor-Kleingrothaus, 1997; Bilenky, Giunti & Monteno, 1997; Fukuyama, Matsuda & Nishiura, 1997; Bilenky, Giunti, Kim & Monteno, 1998; Fukuyama, Matsuda & Nishiura, 1998; Vissani, 1999; Giunti, 1999; Bilenky, Giunti, Grimus, Kayser & Petcov, 1999; Ma, 1999; Wodecki & Kaminsky, 2000; Kalliomaki & Maalampi, 2000; Rodejohann, 2000; Matsuda, Takeda, Fukuyama & Nishiura, 2000; Klapdor-Kleingrothaus, Päs & Smirnov, 2001; Falcone & Tramontano, 2001; Bilenky, Pascoli & Petcov, 2001; Xing, 2001; Osland & Vigdel, 2001; Pascoli & Petcov, 2001; Barger, Glashow, Marfatia & Whisnant, 2002; Hambye, 2002; Minakata & Sugiyama, 2002; Klapdor-Kleingrothaus & Sarkar, 2002; Xing, 2002; Haba & Suzuki, 2002; Pakvasa & Roy, 2002; Rodejohann, 2002; Haba, Nakamura & Suzuki, 2002; Päs & Weiler, 2002; Barger, Glashow , Langacker, Marfatia, 2002; Civitarese & Suhonen, 2002; Pascoli, Petcov & Rodejohann, 2002; Sugiyama, 2002; Avignone & King, 2002; Minakata & Sugiyama, 2002; Cheung, 2003; Abada & Bhattacharyya, 2003; Giunti, 2003; Pascoli & Petcov, 2003; Elliott, 2003; Stoica, 2004; Brahmachari, 2004; Bilenky, Fäßler & Simkovic, 2004; Pascoli & Petcov, 2004; Deppisch, Päs & Suhonen, 2004; Joniec & Zralek, 2004; Pascoli & Petcov, 2005; Pascoli, Petcov & Schwetz, 2005; Goswami & Rodejohann, 2005; Choubey & Rodejohann, 2005; Bilenky, Fäßler, Gutsche, & Simkovic, 2005; Lindner, Merle & Rodejohann, 2005;



Our plots are blue and yellow...



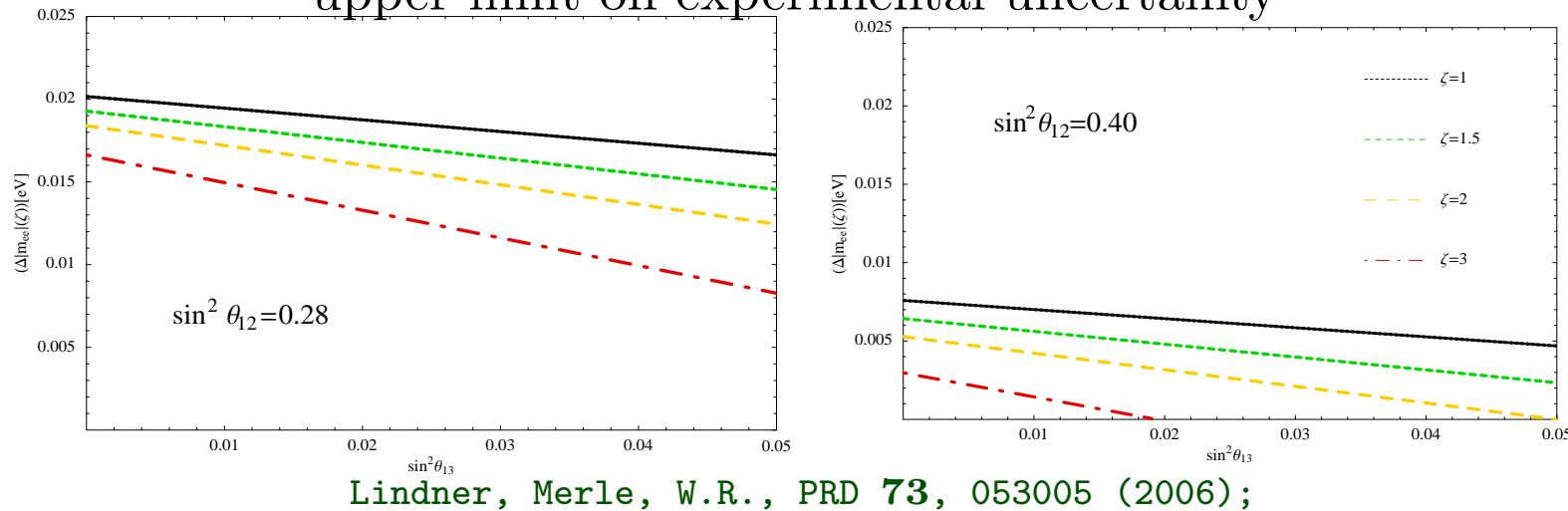
Lindner, Merle, W.R., Phys. Rev. D **73**, 053005 (2006)

DISTINGUISH NH FROM IH?!

NORMAL VS. INVERTED HIERARCHY VS. NUCLEAR PHYSICS

$$\Delta|m_{ee}| \equiv |m_{ee}|_{\text{MIN}}^{\text{IH}} - \zeta |m_{ee}|_{\text{MAX}}^{\text{NH}} \stackrel{!}{>} 0$$

upper limit on experimental uncertainty



Lindner, Merle, W.R., PRD **73**, 053005 (2006);

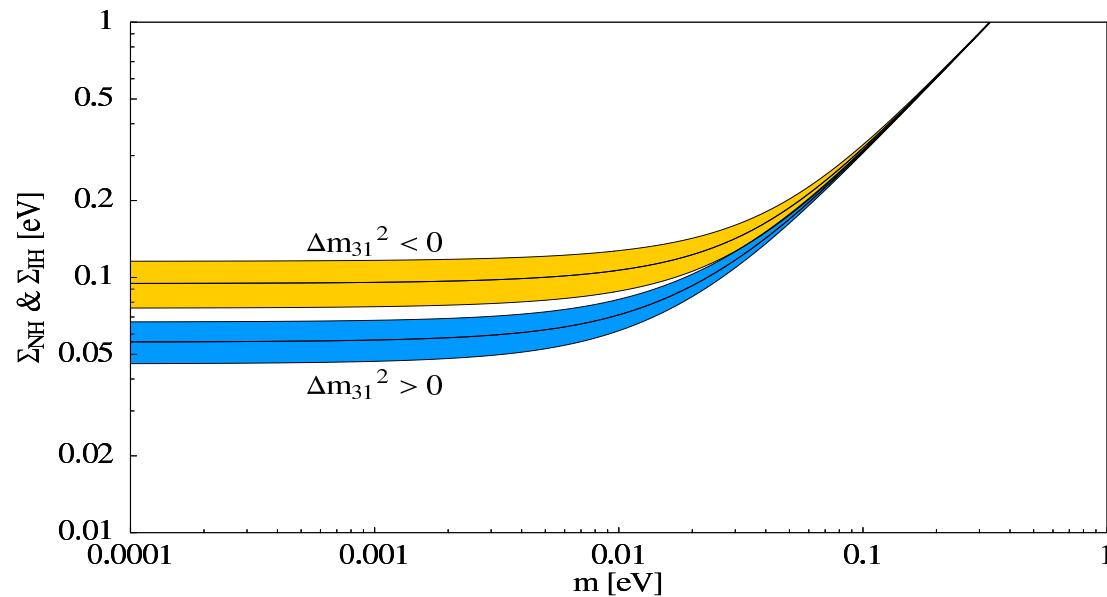
S. Choubey, W.R.; Pascoli, Petcov, Schwetz; de Gouvea, Jenkins

- likes $\zeta \lesssim 2$
- strong dependence on θ_{12} : likes small $\sin^2 \theta_{12} \lesssim 0.35$
- some dependence on U_{e3} : prefers small U_{e3} (\leftrightarrow oscillations)

OTHER (NON-OSCILLATION) PROBES OF NH VS. IH

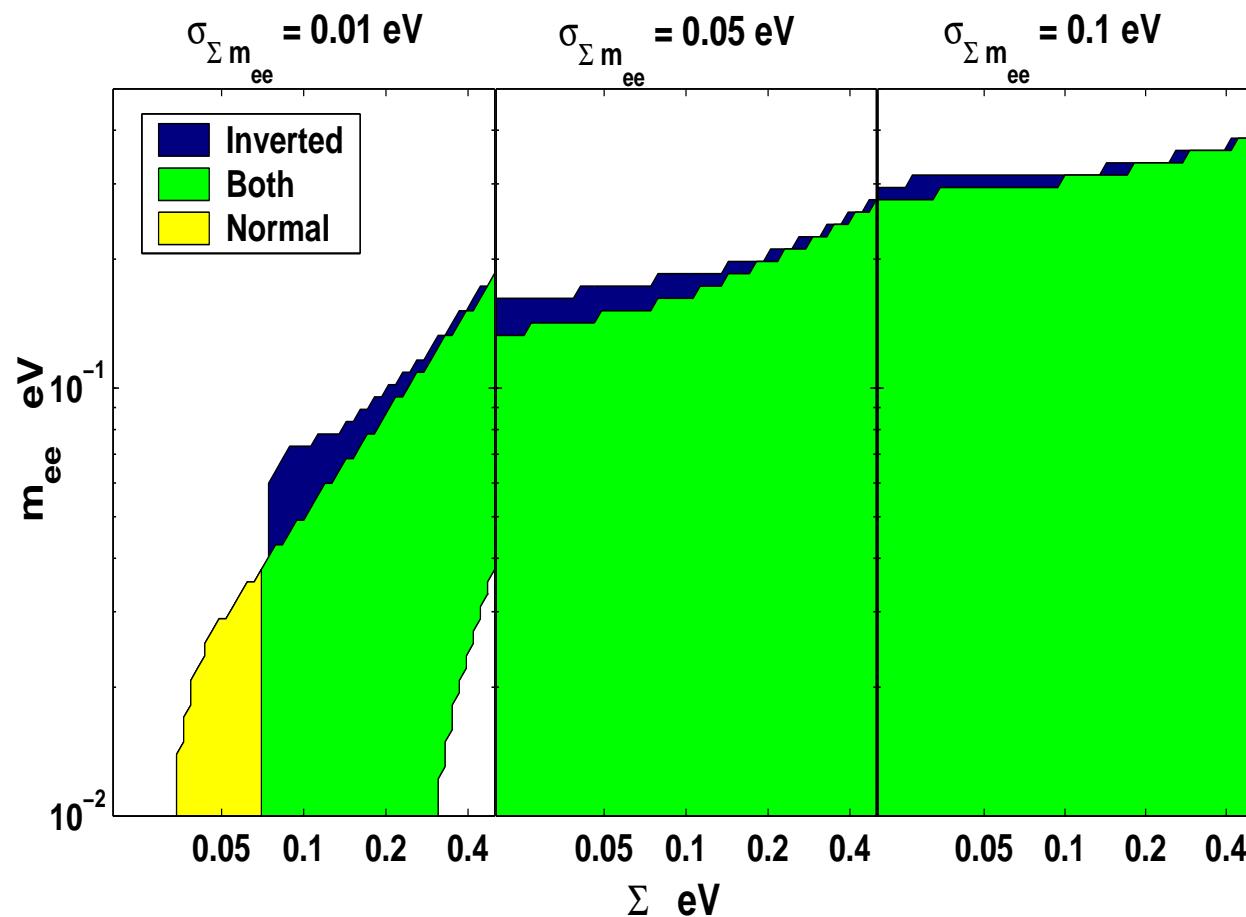
1) Cosmology: $\Sigma = \sum m_i$

$$\Sigma^{\text{NH}} \simeq \sqrt{\Delta m_A^2} < \Sigma^{\text{IH}} \simeq 2\sqrt{\Delta m_A^2}$$



- independent on mixing angles
- systematics?

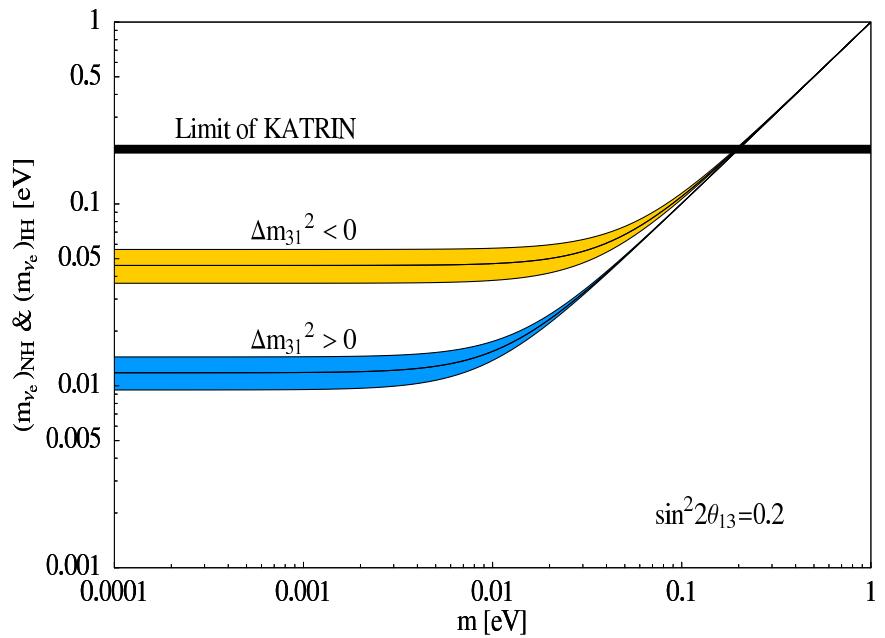
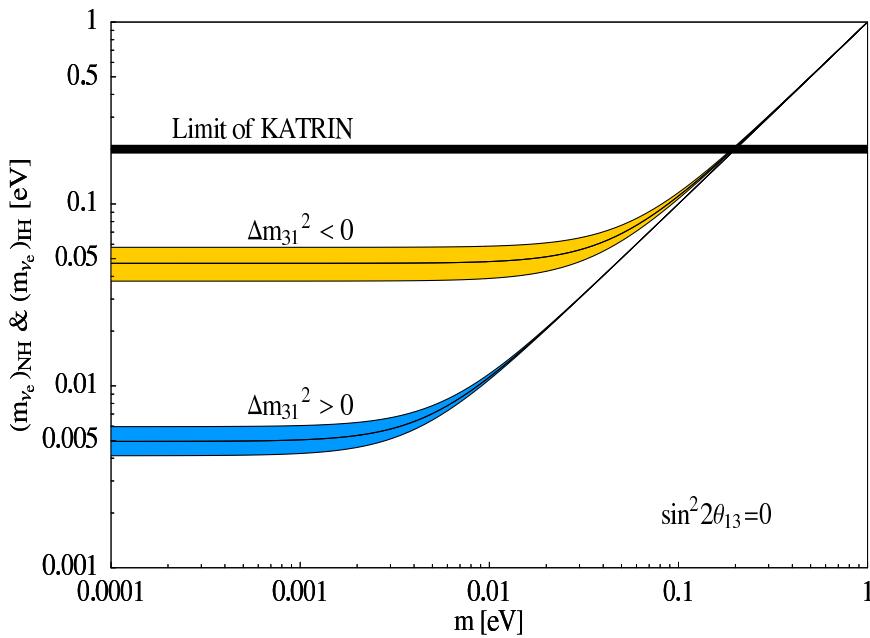
COSMOLOGY AND $0\nu\beta\beta$



Jenkins, de Gouvea, hep-ph/0507021

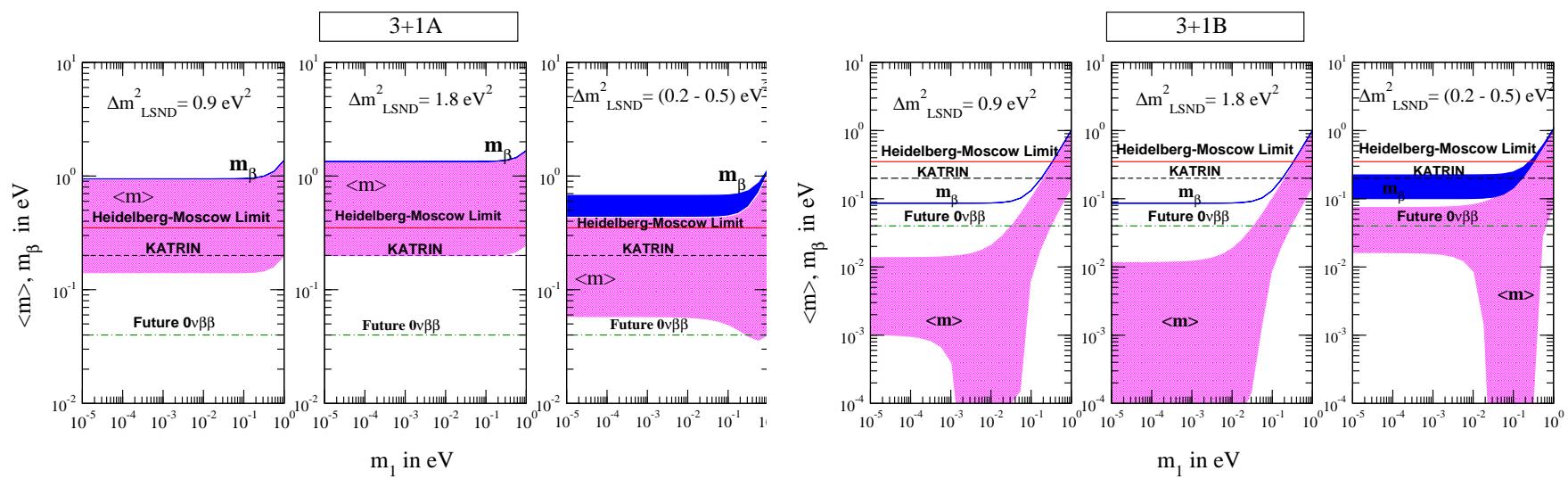
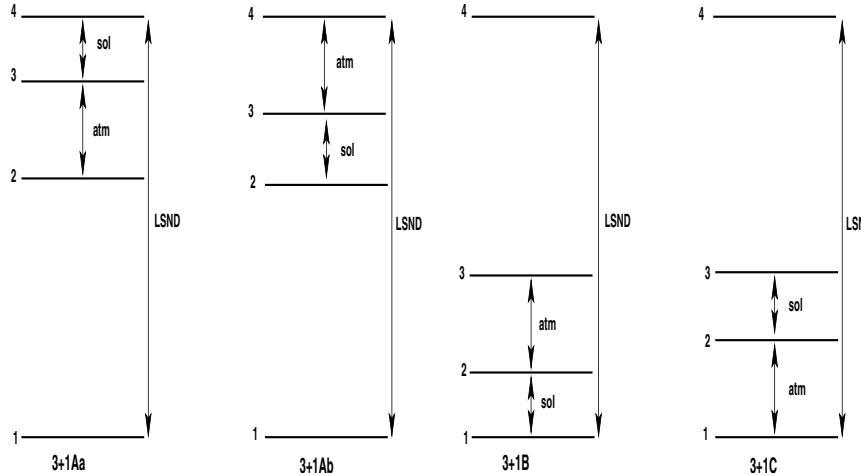
2) β -decay: $m_{\nu_e} = \sqrt{\sum |U_{ei}|^2 m_i^2}$

$$m_{\nu_e}^{\text{NH}} \simeq \sqrt{s_{12}^2 c_{13}^2 \Delta m_{\odot}^2 + s_{13}^2 \Delta m_{\text{A}}^2} \ll m_{\nu_e}^{\text{IH}} \simeq \sqrt{c_{13}^2 \Delta m_{\text{A}}^2}$$



- almost independent on mixing angles
- difference of normal and inverted shows up well below KATRIN limit
- different for sterile (LSND/MiniBoone!!) neutrinos

LSND AND $|m_{ee}|$ (Goswami, W.R., Phys. Rev. D 73, 113003 (2006))

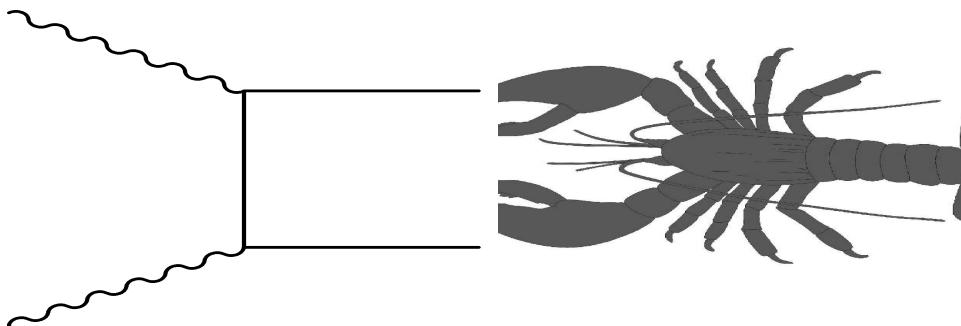


KATRIN AND $0\nu\beta\beta$

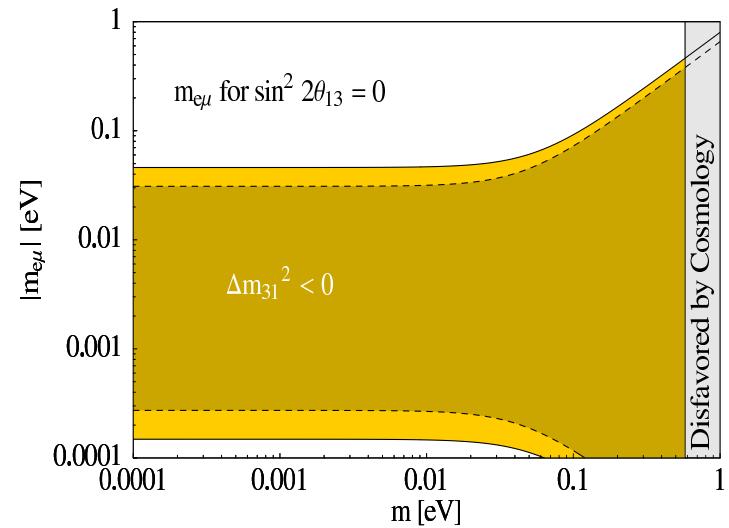
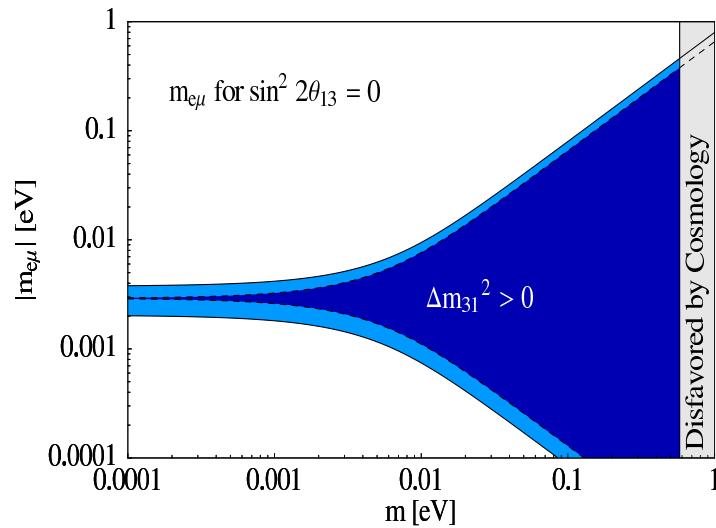
$0\nu\beta\beta$	Δm_A^2	KATRIN	Conclusion
yes	> 0	yes	QD, Majorana
yes	> 0	no	QD, Majorana or NH, Majorana + heavy particles
yes	< 0	no	IH, Majorana
yes	< 0	yes	QD, Majorana
no	> 0	no	NH, Dirac or Majorana
no	< 0	no	Dirac
no	< 0	yes	Dirac
no	> 0	yes	Dirac

OTHER (NON-OSCILLATION) PROBES OF NH VS. IH

3) Other elements of m_ν : “the lobster”

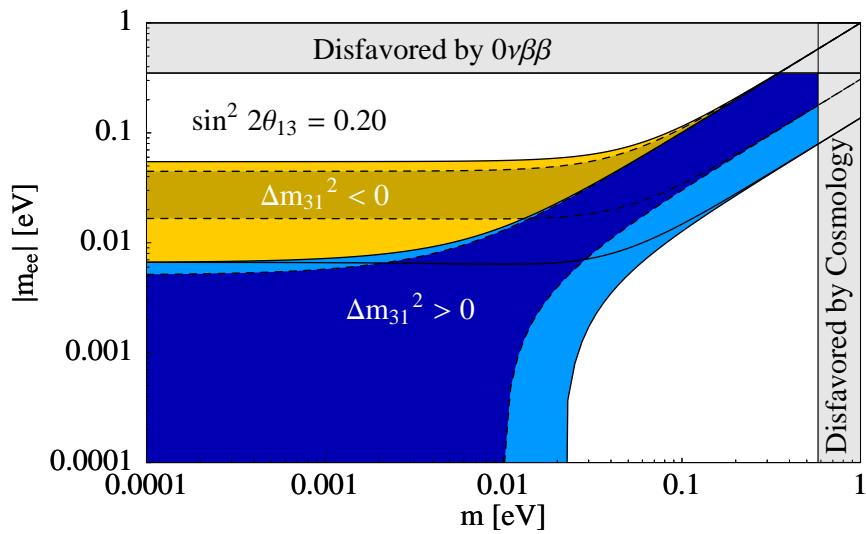
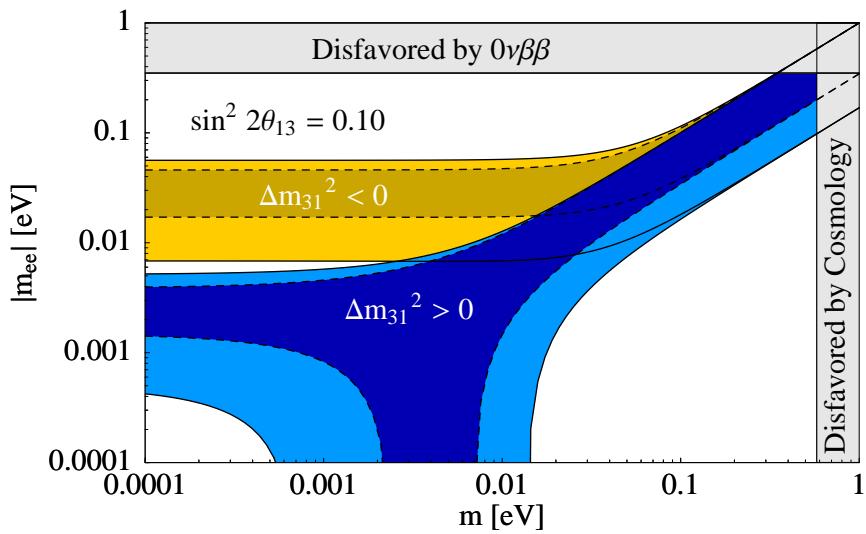
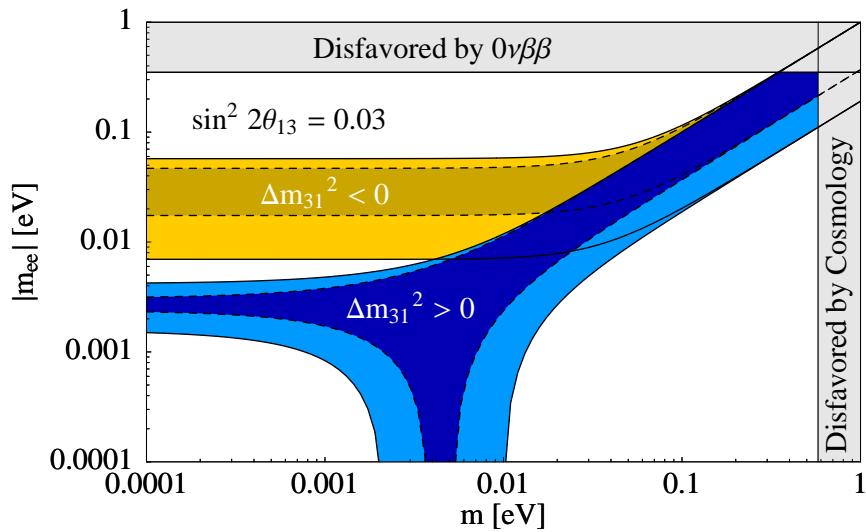
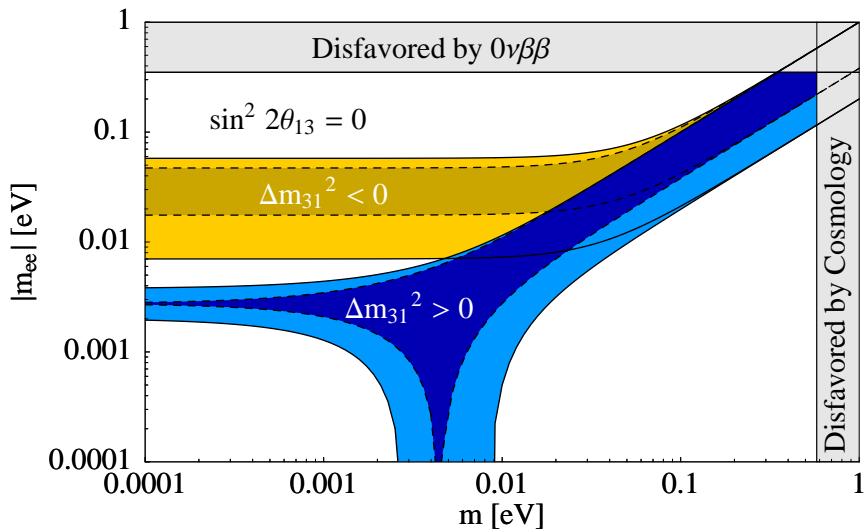


$$\text{BR}(K^+ \rightarrow \pi^- e^+ \mu^+) \propto |m_{e\mu}|^2 = \left| \sum U_{ei} U_{\mu i} m_i \right|^2 \sim 10^{-30} \left(\frac{|m_{e\mu}|}{\text{eV}} \right)^2$$



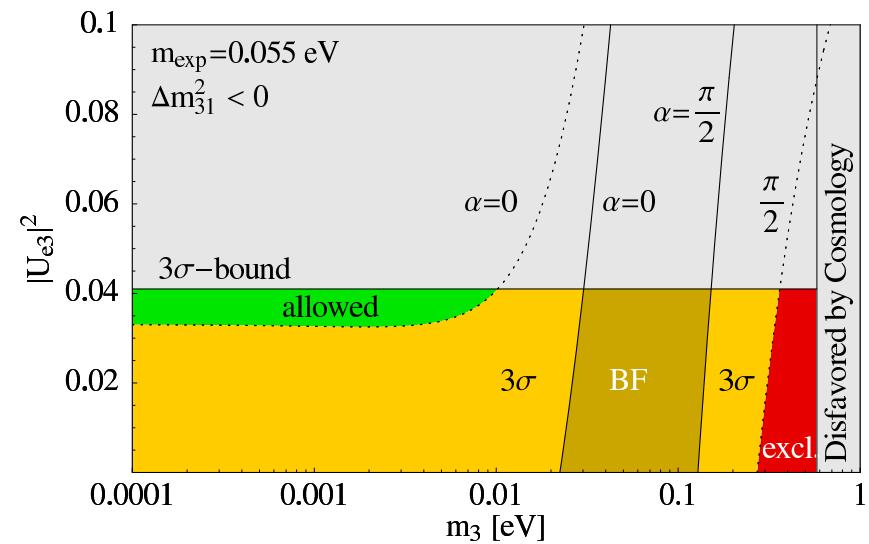
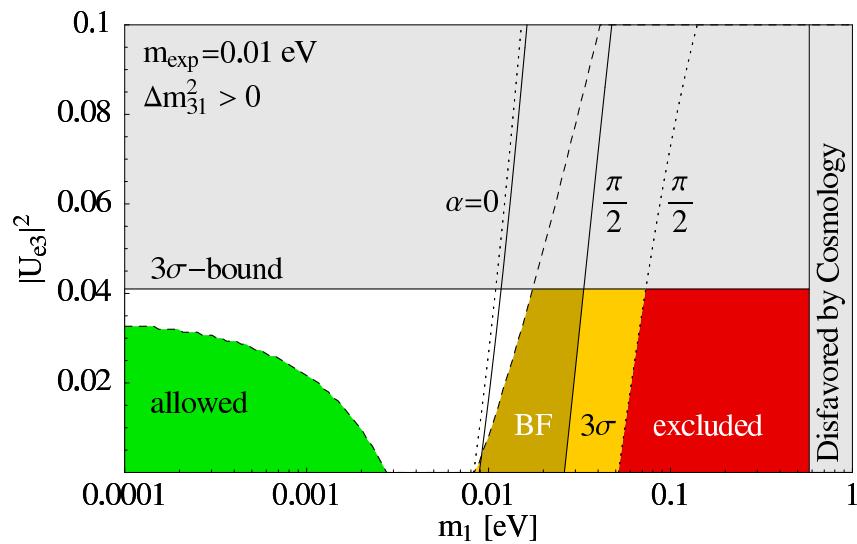
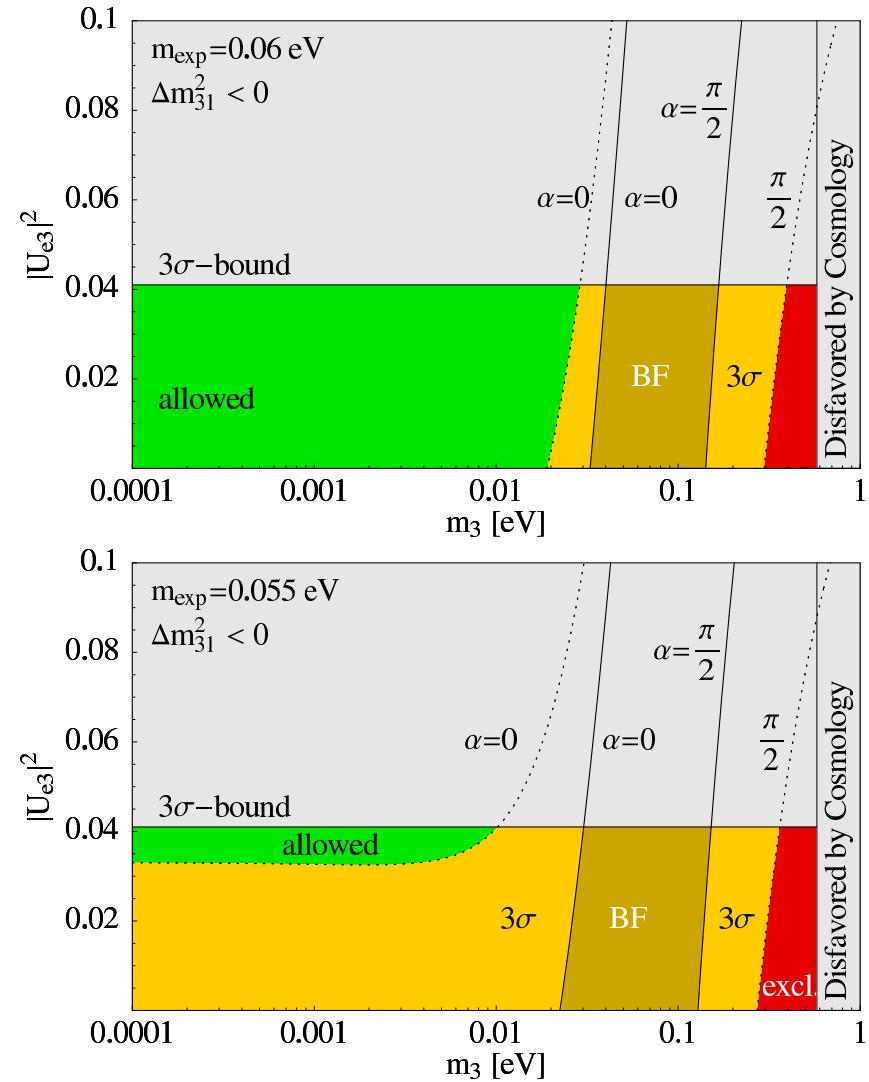
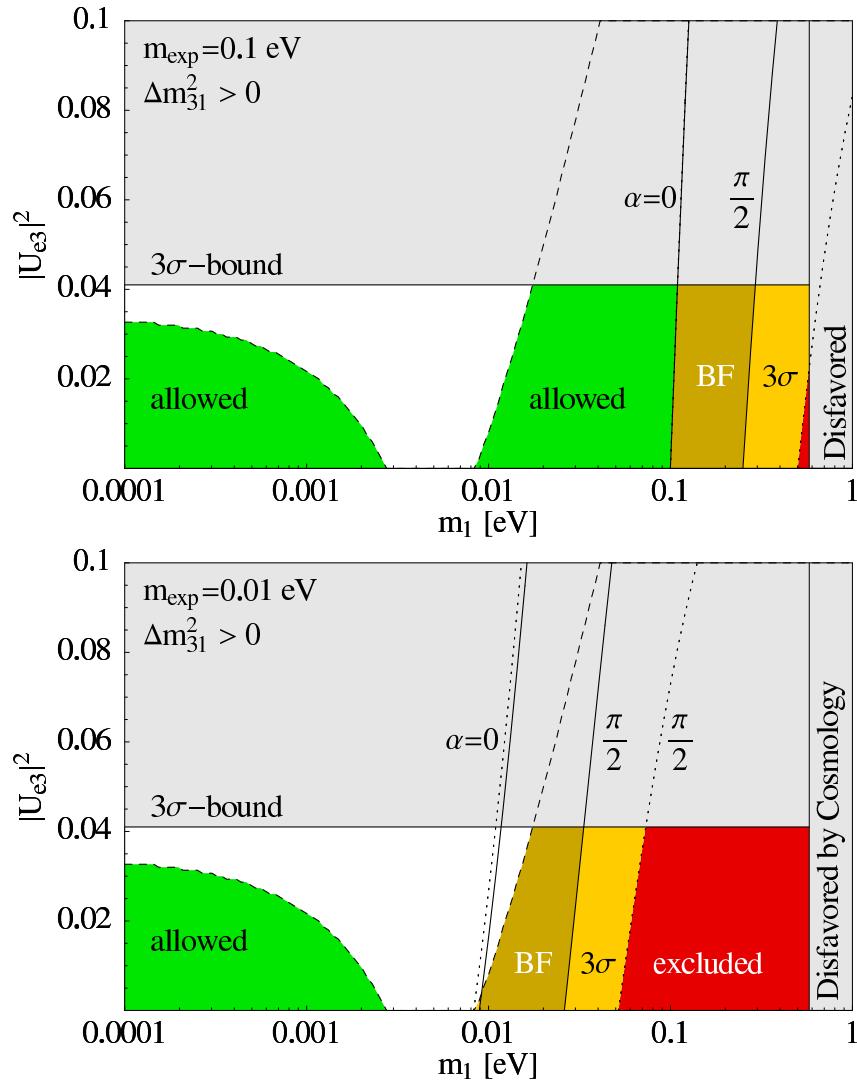
Merle, W.R., Phys. Rev. D 74, 017701 (2006)

$|U_{e3}|$ AND $0\nu\beta\beta$

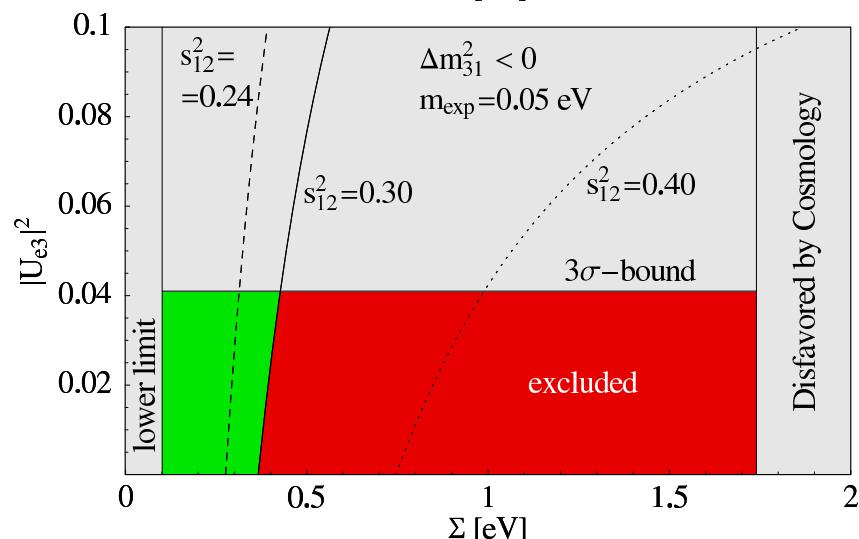
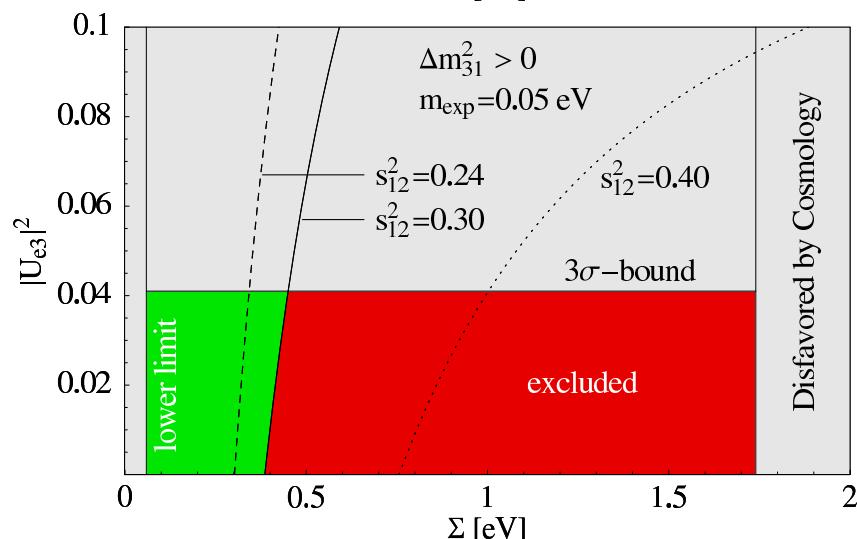
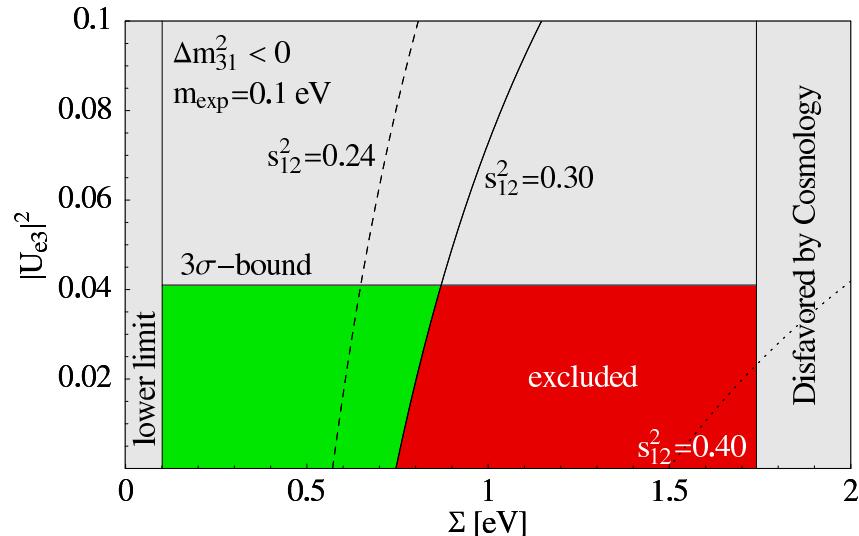
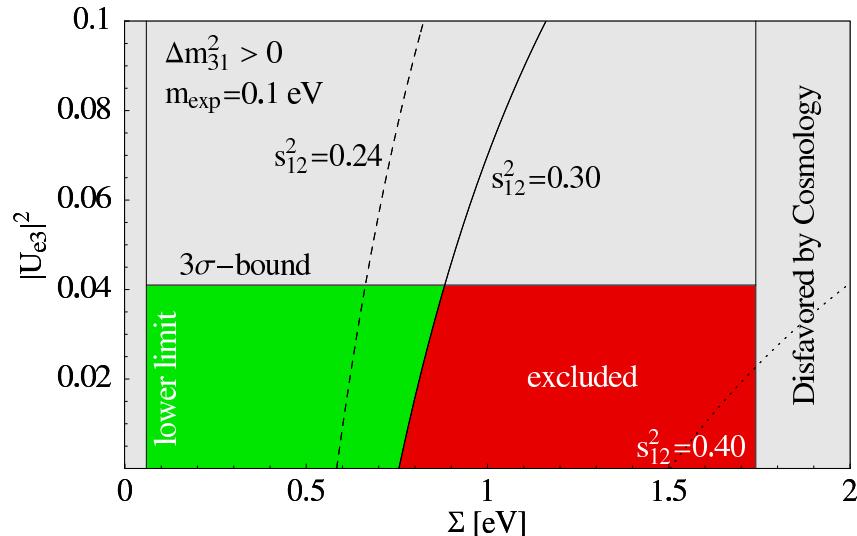


⇒ Value of U_{e3} influences $|m_{ee}|$, in particular NH vs. IH

LOWER LIMIT ON $|U_{e3}|^2$



COSMOLOGY AND $|U_{e3}|^2$

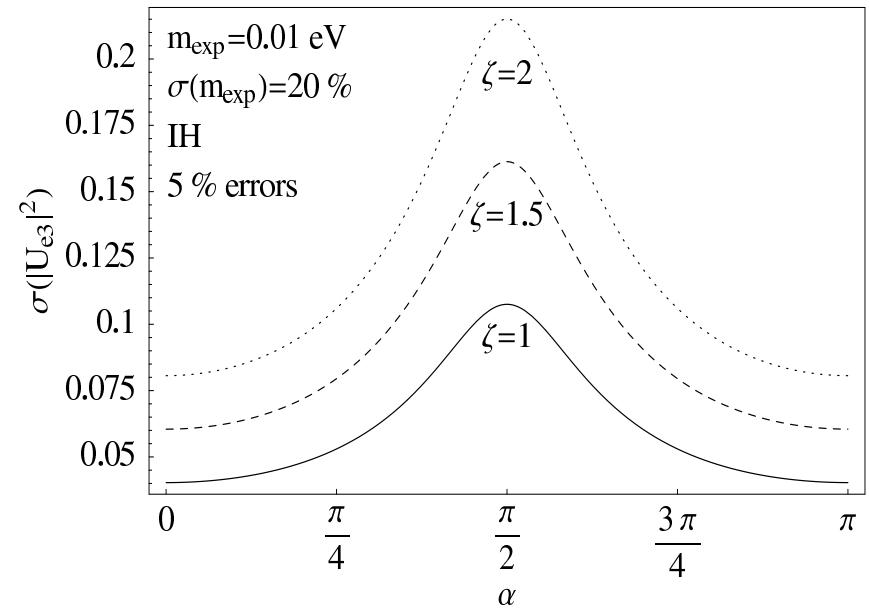
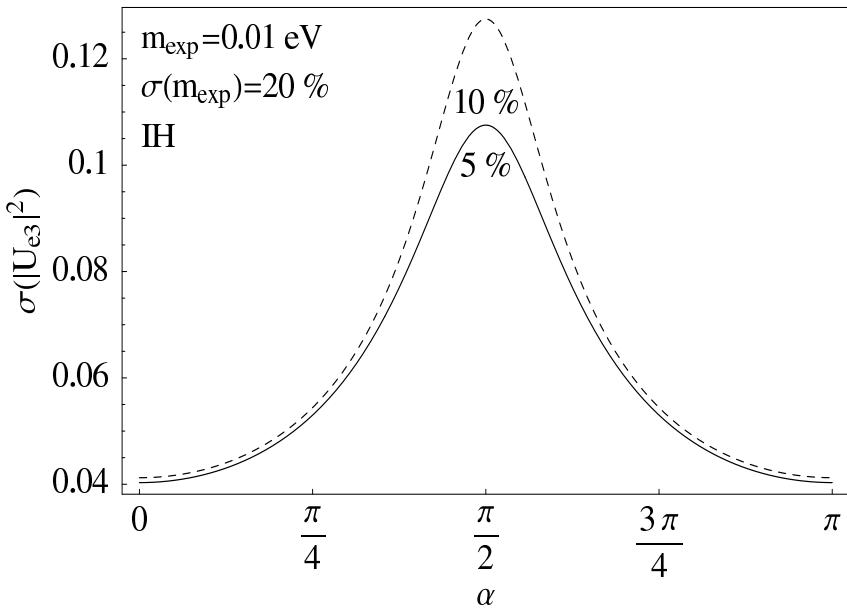


Merle, W.R., hep-ph/0703135

MEASURING $|U_{e3}|^2$ WITH NEUTRINO-LESS DOUBLE BETA DECAY?

Example inverted hierarchy:

$$|U_{e3}|^2 = 1 - \frac{|m_{ee}|}{\sqrt{\Delta m_A^2} \sqrt{1 - \sin^2 2\theta_{12} \sin^2 \alpha}}$$



WHAT'S MORE TO $0\nu\beta\beta$?

* Mass scale: consider QD spectrum

$$m_0 \leq \frac{1 + \tan^2 \theta_{12}}{1 - \tan^2 \theta_{12} - 2 |U_{e3}|^2} |m_{ee}|_{\text{exp}} \lesssim 5 \text{ eV}$$

* Majorana phases: consider IH spectrum

$$\sin^2 \alpha = \left(1 - \frac{|m_{ee}|}{\sqrt{|\Delta m_A^2|} (1 - |U_{e3}|^2)} \right)^2 \frac{1}{\sin^2 2\theta_{12}}$$

measurement requires $\zeta \lesssim 1.5$ and

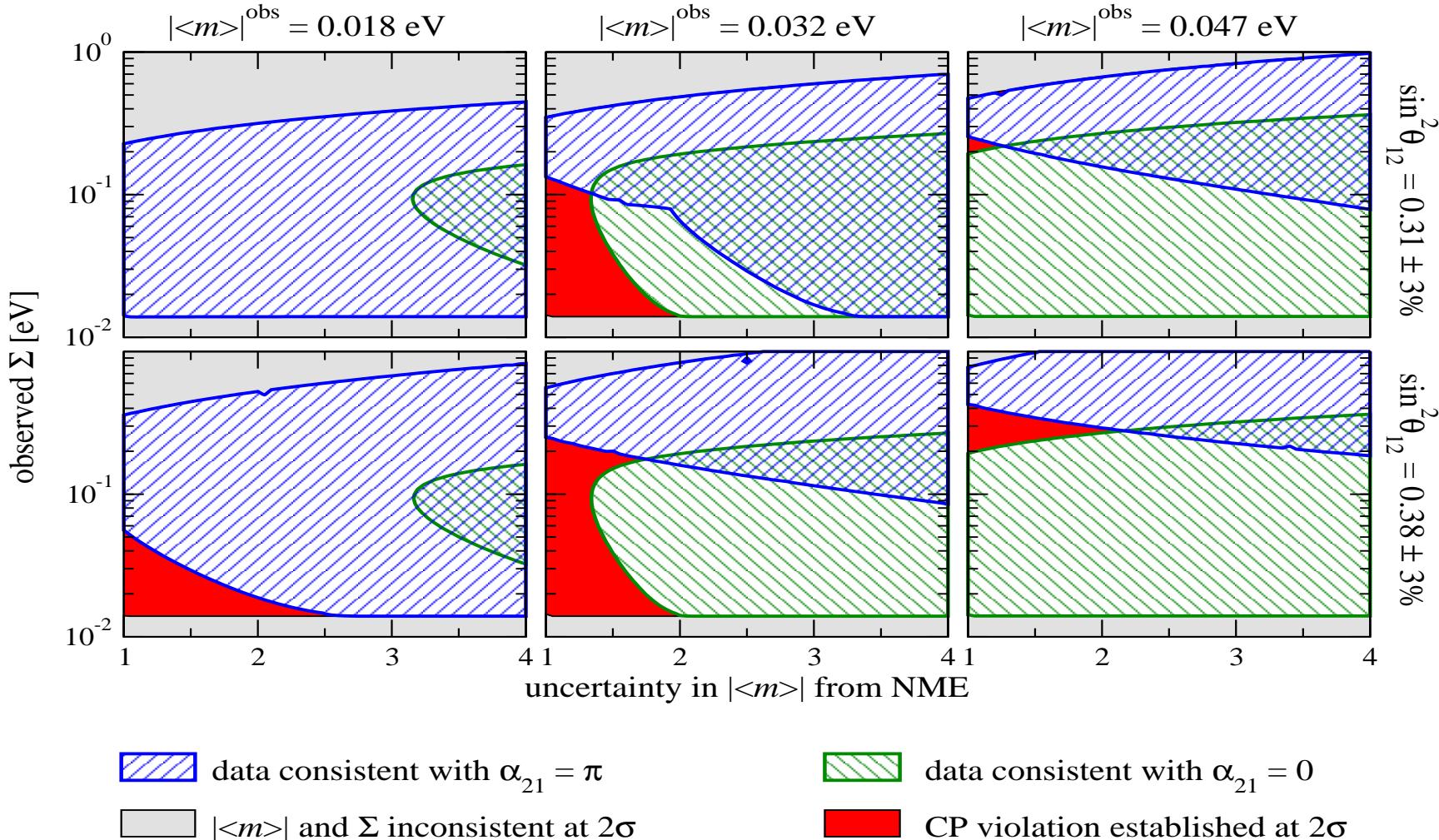
- $\sigma(|m_{ee}|) \lesssim 15\%$
- $\sigma(\Delta m_A^2) \lesssim 10\%$ (IH) or $\sigma(m_0) \lesssim 10\%$ (QD)
- $\sin^2 \theta_{12} \gtrsim 0.29$
- $2\alpha \in [\pi/4, 3\pi/4]$ or $[5\pi/4, 7\pi/4]$

Pascoli, Petcov, W.R., Phys. Lett. B 549, 177 (2002)

No to “no-go” from Barger *et al.*, Phys. Lett. B 540, 247 (2002)

MAJORANA PHASES

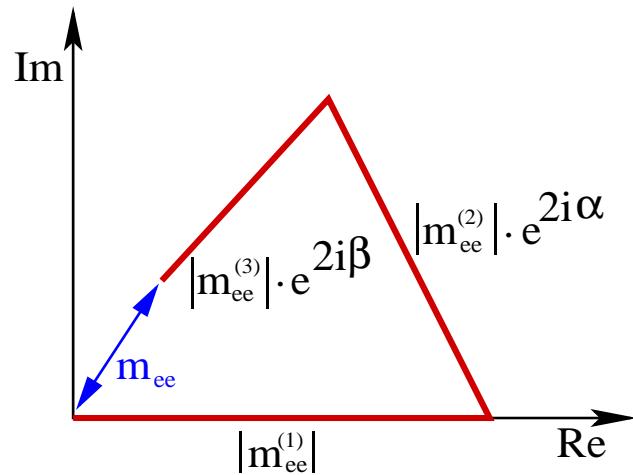
Pascoli, Petcov, Schwetz, Nucl. Phys. B 734, 24 (2006)



$$\sin^2 \theta_{13} = 0 \pm 0.002, \Delta m_{21}^2 = 8 \times 10^{-5} \pm 2\%, \Delta m_{31}^2 = -2.2 \times 10^{-3} \pm 3\%, \sigma_{\beta\beta} = 0.004 \text{ eV}, \sigma_\Sigma = 0.04 \text{ eV}$$

IF WE DON'T OBSERVE $0\nu\beta\beta$: VANISHING $|m_{ee}|$

- a triangle can be formed! $\left| |m_{ee}^{(1)}| + |m_{ee}^{(3)}| e^{2i\alpha} + |m_{ee}^{(3)}| e^{2i\beta} \right|$



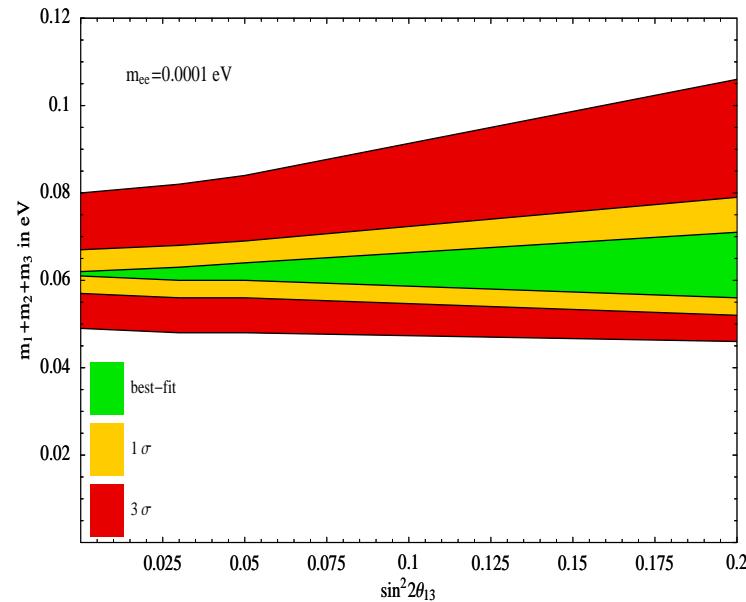
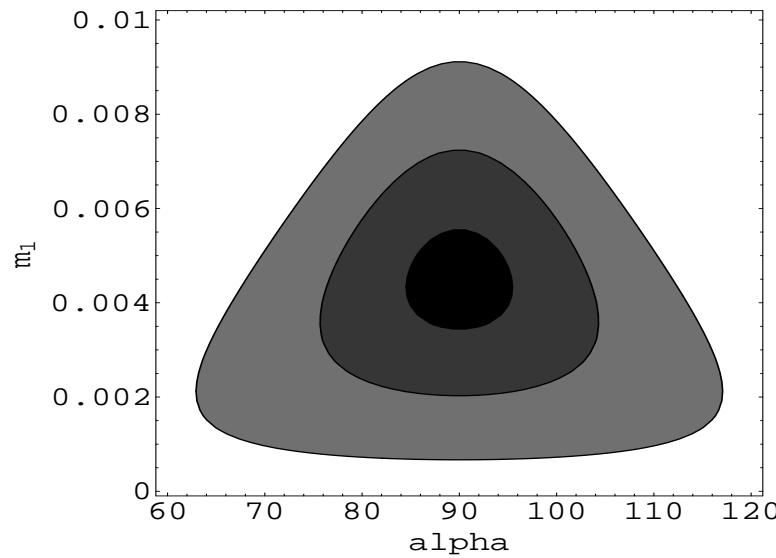
- texture zero in charged lepton basis!

$$\cos 2\alpha = \frac{|m_{ee}^{(1)}|^2 + |m_{ee}^{(2)}|^2 - |m_{ee}^{(3)}|^2}{2|m_{ee}^{(1)}| |m_{ee}^{(2)}|} =$$

$$\frac{m_1^2 \left(c_{13}^4 (s_{12}^4 + c_{12}^4) - s_{13}^4 \right) + \Delta m_\odot^2 s_{12}^4 c_{13}^4 - \Delta m_A^2 s_{13}^4}{2m_1 \sqrt{m_1^2 + \Delta m_\odot^2} s_{12}^2 c_{12}^2 c_{13}^4}$$

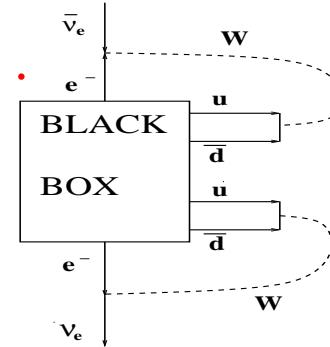
IF WE DON'T OBSERVE $0\nu\beta\beta$: VANISHING $|m_{ee}|$

- only possible for NH
- if $\theta_{13} = 0$: $m_1 = \sin^2 \theta_{12} \sqrt{\frac{\Delta m_\odot^2}{\cos 2\theta_{12}}} \simeq 4.5$ ($2.8 \div 8.4$) meV
- if $m_1 = 0$: $\sin^2 2\theta_{13} \simeq 4 \sin^2 \theta_{12} \sqrt{\frac{\Delta m_\odot^2}{\Delta m_A^2}} \simeq 0.24$ ($0.14 \div 0.40$)



IF WE OBSERVE $0\nu\beta\beta\dots$

- Neutrinos are Majorana (**Schechter-Valle**)



- we still need to identify the mechanism of $0\nu\beta\beta$: SUSY, RH currents, heavy Majorana neutrinos,...
solution: angular distribution, measure $0\nu\beta\beta$ in many different nuclei...
- reduce NME uncertainty: measure $0\nu\beta\beta$ in many different nuclei...
- if non-zero, Majorana phases **do contribute** to Y_B in leptogenesis (flavor effects)
- in general, we will believe much more firmly in leptogenesis
- $|m_{ee}|$ with precision experiments of oscillation parameters fixes structure of m_ν

SUMMARY

- – Oscillations $\Rightarrow m_\nu \neq 0$
 - almost all models explain $m_\nu \neq 0$ in connection with Lepton Number Violation $\Rightarrow 0\nu\beta\beta$
- $0\nu\beta\beta$ can
 - distinguish NH from IH
 - Mass scale (consistency with KATRIN and cosmology)
 - consistency with U_{e3}
 - Majorana phases (Leptogenesis?)
- identify structure of mass matrix

$|m_{ee}| + \text{precision data will help}$