

Three-flavour and matter effects in neutrino oscillations

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3 neutrino flavours exist: ν_e, ν_μ, ν_τ

But: 2f description – a good 1st approximation in many cases.

Reasons:

- Hierarchy of Δm^2 : $\Delta m_{\text{sol}}^2 \ll \Delta m_{\text{atm}}^2$
- Smallness of $|U_{e3}|$.

Exceptions: $P(\nu_\mu \leftrightarrow \nu_\tau)$, $P(\nu_\mu \rightarrow \nu_\mu)$ and $P(\nu_\tau \rightarrow \nu_\tau)$ when oscillations due to the solar frequency ($\sim \Delta m_{\text{sol}}^2$) are not frozen.

In any case, coorections due to 3-flavoriness can reach $\sim 10\%$

– cannot be ignored at present

Also: a number of pure 3f effects exist \Rightarrow

◇ 3f analyses are a must !

Two types of 3f effects

I. “Trivial” effects

- Existence of new physical channels – in addition to $\nu_e \leftrightarrow \nu_\mu$ there are $\nu_e \leftrightarrow \nu_\tau$ and $\nu_\mu \leftrightarrow \nu_\tau$; mutual influence of channels through unitarity (conservation of probability).
- New “parameter channels” for the same physical channel. E.g.: $\nu_e \leftrightarrow \nu_\mu$ oscill. can be governed by $(\Delta m_{21}^2, \theta_{12})$ and $(\Delta m_{31}^2, \theta_{13})$

Two types of 3f effects – contd.

II. Nontrivial effects

- Fundamental \mathcal{CP} and \mathcal{T}
- Matter-induced \mathcal{T}
- Interference of different “parameter channels” in $\nu_e \leftrightarrow \nu_{\mu,\tau}$ – specific contributions to oscillation probabilities
- Matter effects on $\nu_{\mu} \leftrightarrow \nu_{\tau}$ oscillations

Nontrivial 3f effects (except the last one): disappear if at least one mixing angle is 0 or 90°, or at least one

$$\Delta m_{ij}^2 = 0$$

3f neutrino mixing

Relates flavour eigenstates $\nu_f = (\nu_e, \nu_\mu, \nu_\tau)$ with mass eigenstates $\nu_m = (\nu_1, \nu_2, \nu_3)$:

$$\nu_f = \tilde{U} \nu_m$$

$$\tilde{U} = UK, \quad K = \text{diag}(1, e^{i\sigma_1}, e^{i\sigma_2})$$

Majorana-type phases $\sigma_{1,2}$ do not affect neutrino oscillations.

The relevant part of the mixing matrix:

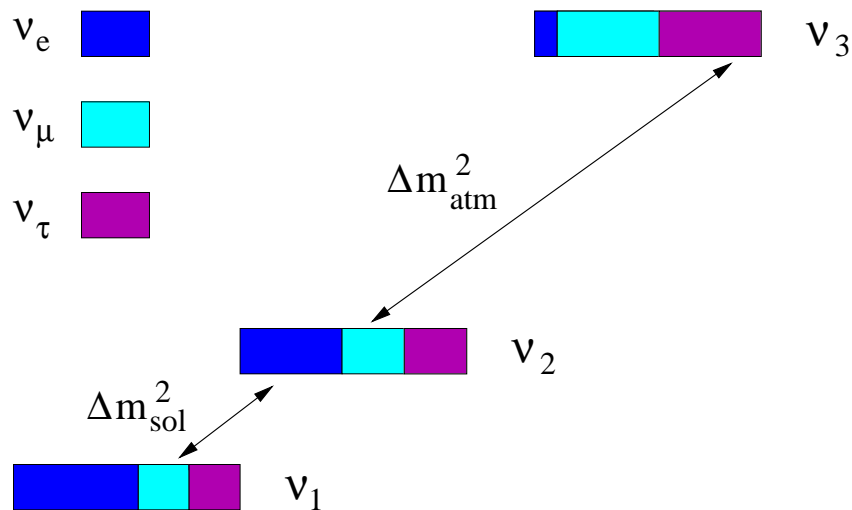
$$U = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta_{\text{CP}}} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta_{\text{CP}}} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$= O_{23} (\Gamma_\delta O_{13} \Gamma_\delta^\dagger) O_{12}, \quad \Gamma_\delta \equiv \text{diag}(1, 1, e^{i\delta_{\text{CP}}})$$

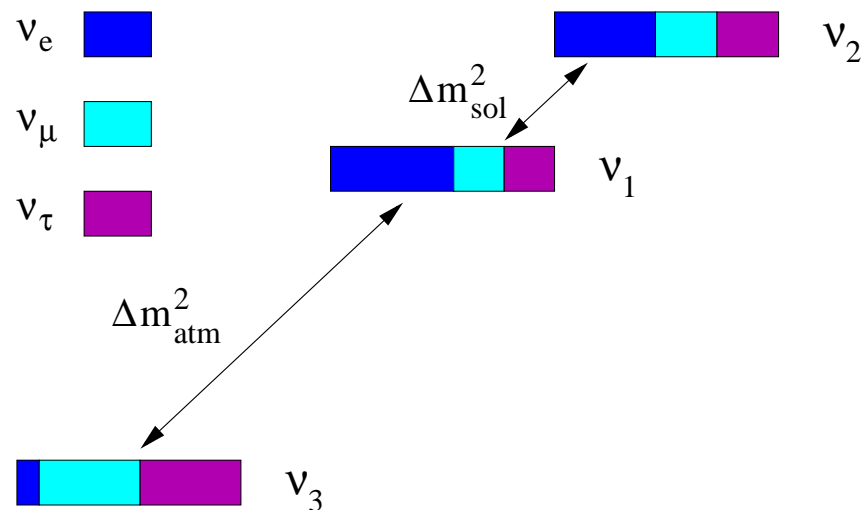
Leptonic mixing – contd.

$$U = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta_{CP}} \\ -s_{12}c_{23} - c_{12}s_{13}s_{23}e^{i\delta_{CP}} & c_{12}c_{23} - s_{12}s_{13}s_{23}e^{i\delta_{CP}} & c_{13}s_{23} \\ s_{12}s_{23} - c_{12}s_{13}c_{23}e^{i\delta_{CP}} & -c_{12}s_{23} - s_{12}s_{13}c_{23}e^{i\delta_{CP}} & c_{13}c_{23} \end{pmatrix}$$

Normal hierarchy:



Inverted hierarchy:



Genuine 3f effects

\mathcal{CP} and \mathcal{T} in ν oscillations in vacuum

- \mathcal{CP} : $P(\nu_a \rightarrow \nu_b) \neq P(\bar{\nu}_a \rightarrow \bar{\nu}_b)$
- \mathcal{T} : $P(\nu_a \rightarrow \nu_b) \neq P(\nu_b \rightarrow \nu_a)$

CPT invariance: $\diamond P(\nu_a \rightarrow \nu_b) = P(\bar{\nu}_b \rightarrow \bar{\nu}_a)$

$\mathcal{CP} \Leftrightarrow \mathcal{T}$ – consequence of CPT

Measures of \mathcal{CP} and \mathcal{T} – probability differences:

$$\Delta P_{ab}^{\text{CP}} \equiv P(\nu_a \rightarrow \nu_b) - P(\bar{\nu}_a \rightarrow \bar{\nu}_b)$$

$$\Delta P_{ab}^{\text{T}} \equiv P(\nu_a \rightarrow \nu_b) - P(\nu_b \rightarrow \nu_a)$$

From CPT:

$$\diamond \Delta P_{ab}^{\text{CP}} = \Delta P_{ab}^{\text{T}} ; \quad \Delta P_{aa}^{\text{CP}} = 0$$

\mathcal{CP} and \mathcal{T} in ν oscillations in matter

Normal matter [(# of particles) \neq (# of anti-particles)]:

The very presence of matter violates C, CP and CPT

\Rightarrow Fake (extrinsic) \mathcal{CP} . Exists even in 2f case. May complicate study of fundamental (intrinsic) \mathcal{CP}

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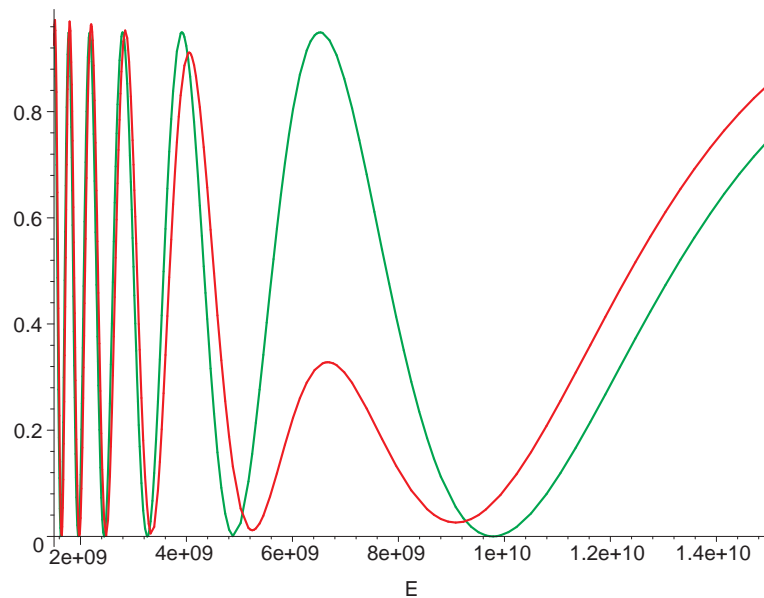
Induced \mathcal{T} : absent when either $U_{e3} = 0$ or $\Delta m_{\text{sol}}^2 = 0$ (2f limits)

\Rightarrow Doubly suppressed by both these small parameters – effects in terrestrial experiments are small

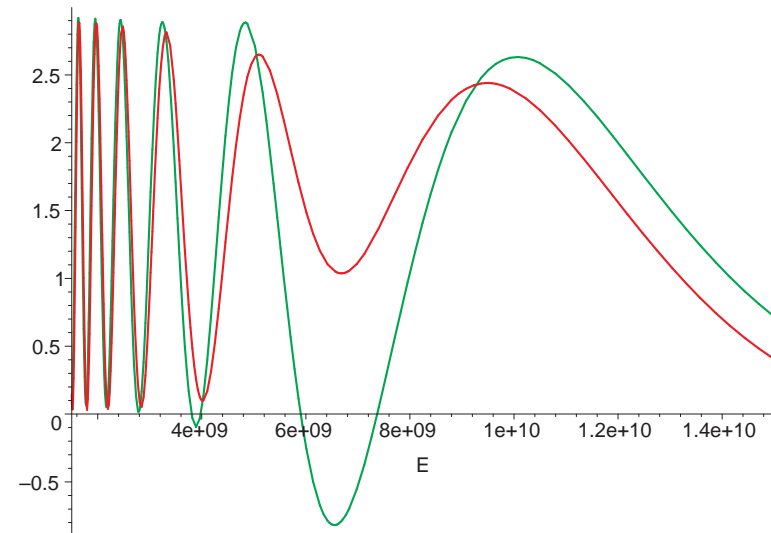
Matter effects on $\nu_\mu \leftrightarrow \nu_\tau$ oscillations

In 2f approximation: no matter effects on $\nu_\mu \leftrightarrow \nu_\tau$ oscillations
[$V(\nu_\mu) = V(\nu_\tau)$ modulo tiny rad. corrections].

Not true in the full 3f framework! (E.A., 2002; Gandhi et al., 2004)



$P_{\mu\tau}$



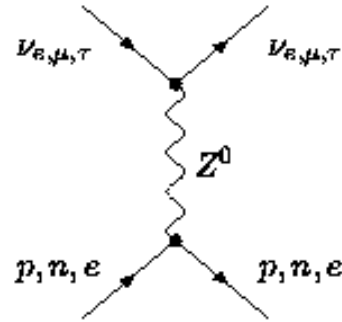
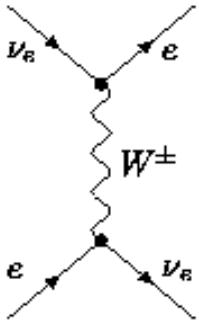
Oscillated flux of atm. ν_μ

$$\Delta m_{31}^2 = 2.5 \times 10^{-3} \text{ eV}^2, \quad \sin^2 \theta_{13} = 0.026, \quad \theta_{23} = \pi/4, \quad \Delta m_{21}^2 = 0, \quad L = 9400 \text{ km}$$

Red curves – w/ matter effects, green curves – w/o matter effects on $P_{\mu\tau}$

Neutrino oscillations in matter

Coherent forward scattering on the particles in matter



$$V_e^{\text{CC}} \equiv V = \sqrt{2} G_F N_e$$

2f neutrino evolution equation:

$$\diamond \quad i \frac{d}{dt} \begin{pmatrix} \nu_e \\ \nu_\mu \end{pmatrix} = \begin{pmatrix} -\frac{\Delta m^2}{4E} \cos 2\theta + V & \frac{\Delta m^2}{4E} \sin 2\theta \\ \frac{\Delta m^2}{4E} \sin 2\theta & \frac{\Delta m^2}{4E} \cos 2\theta \end{pmatrix} \begin{pmatrix} \nu_e \\ \nu_\mu \end{pmatrix}$$

Mixing in matter

$$\diamond \sin^2 2\theta_m = \frac{\sin^2 2\theta \cdot \left(\frac{\Delta m^2}{2E}\right)^2}{\left[\frac{\Delta m^2}{2E} \cos 2\theta - \sqrt{2}G_F N_e\right]^2 + \left(\frac{\Delta m^2}{2E}\right)^2 \sin^2 2\theta}$$

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Mikheyev - Smirnov - Wolfenstein (MSW) resonance:

$$\diamond \sqrt{2}G_F N_e = \frac{\Delta m^2}{2E} \cos 2\theta$$

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$$|\nu_\mu\rangle = -\sin \theta_m |\nu_{1m}\rangle + \cos \theta_m |\nu_{2m}\rangle$$

$|\nu_{1m}\rangle, |\nu_{2m}\rangle$ – eigenstates of H in matter (matter eigenstates)

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$$|\nu_e\rangle = \cos \theta_m |\nu_{1m}\rangle + \sin \theta_m |\nu_{2m}\rangle \quad N_e \gg (N_e)_{\text{res}} : \theta_m \approx 90^\circ$$

$$|\nu_\mu\rangle = -\sin \theta_m |\nu_{1m}\rangle + \cos \theta_m |\nu_{2m}\rangle \quad N_e = (N_e)_{\text{res}} : \theta_m = 45^\circ$$

$$N_e \ll (N_e)_{\text{res}} : \theta_m \approx \theta$$

$|\nu_{1m}\rangle, |\nu_{2m}\rangle$ – eigenstates of H in matter (matter eigenstates)

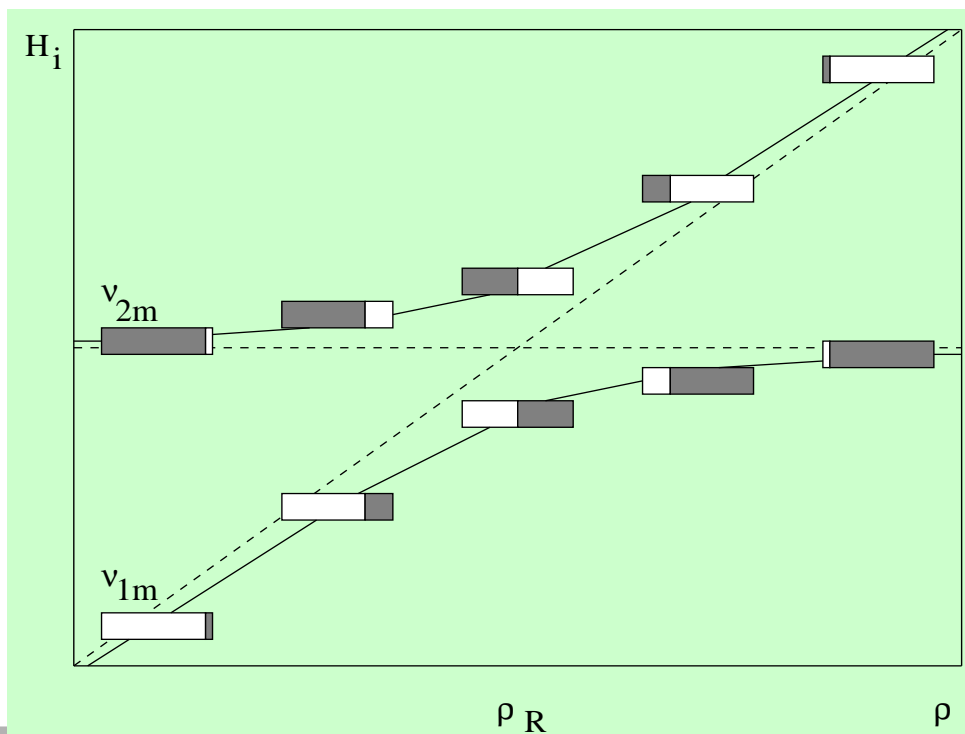
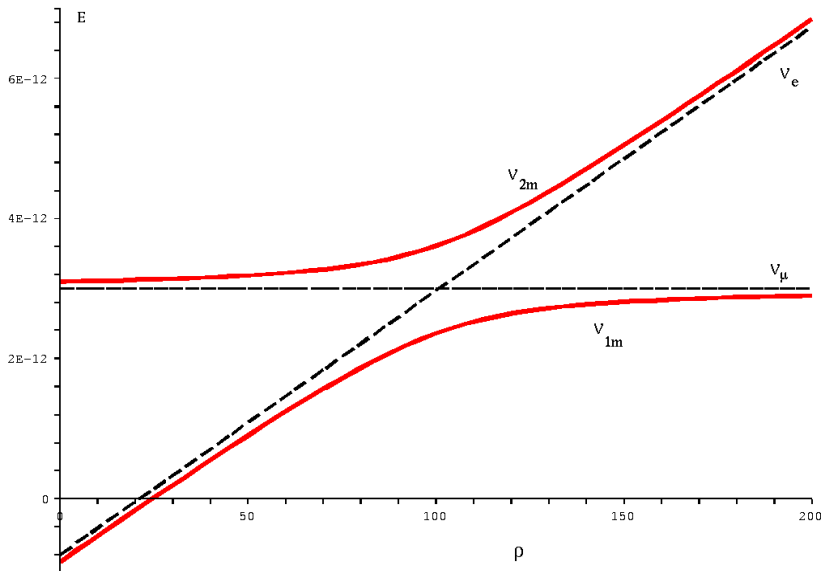
Adiabatic flavour conversion

Adiabaticity: slow density change along the neutrino path

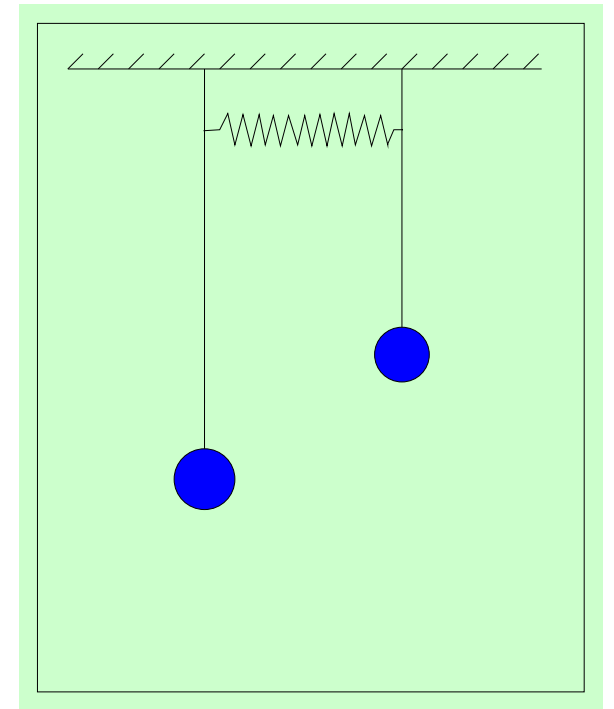
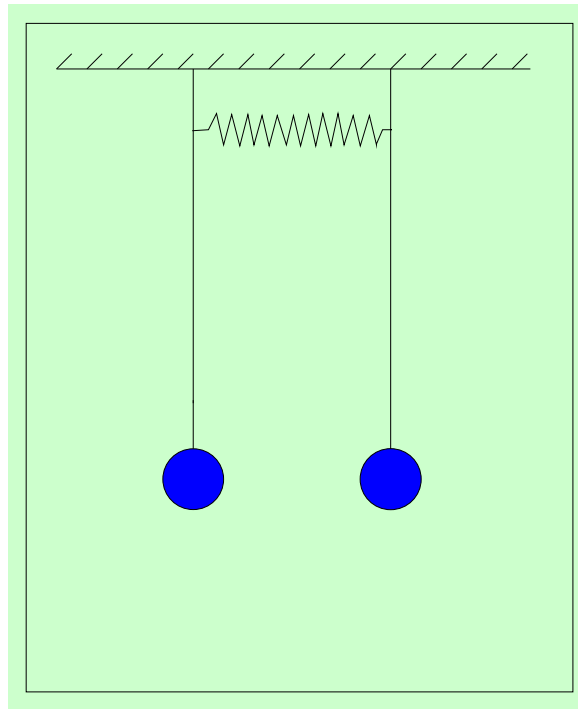
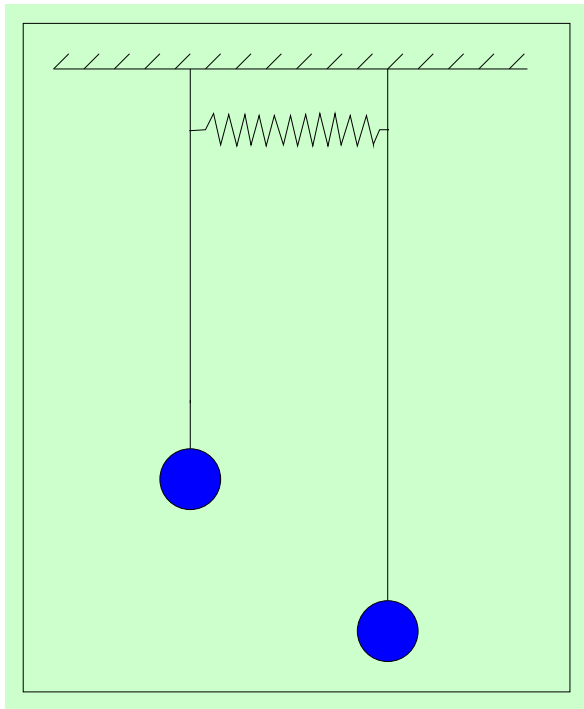
$$\frac{\sin^2 2\theta}{\cos 2\theta} \frac{\Delta m^2}{2E} L_\rho \gg 1$$

L_ρ – electron density scale height:

$$L_\rho = \left| \frac{1}{N_e} \frac{dN_e}{dx} \right|^{-1}$$

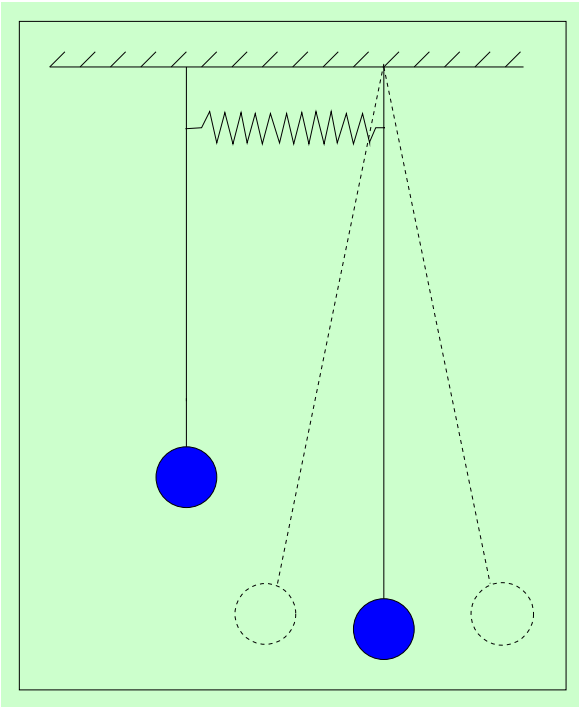


Analogy: Two coupled pendula



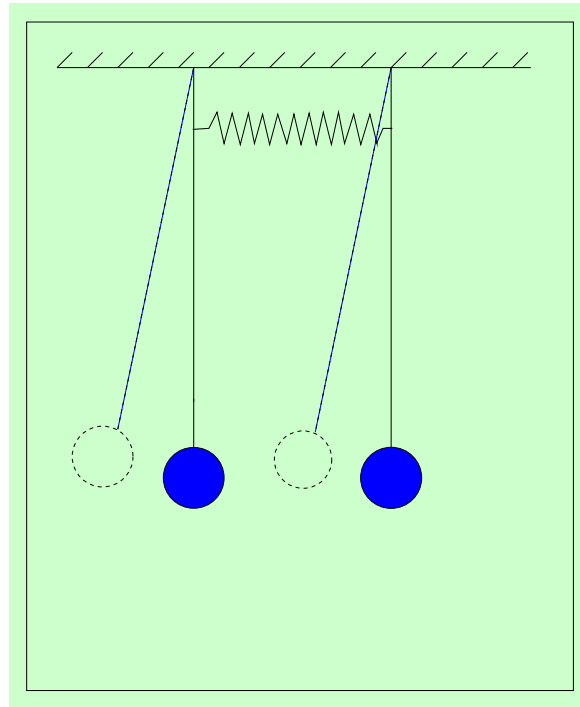
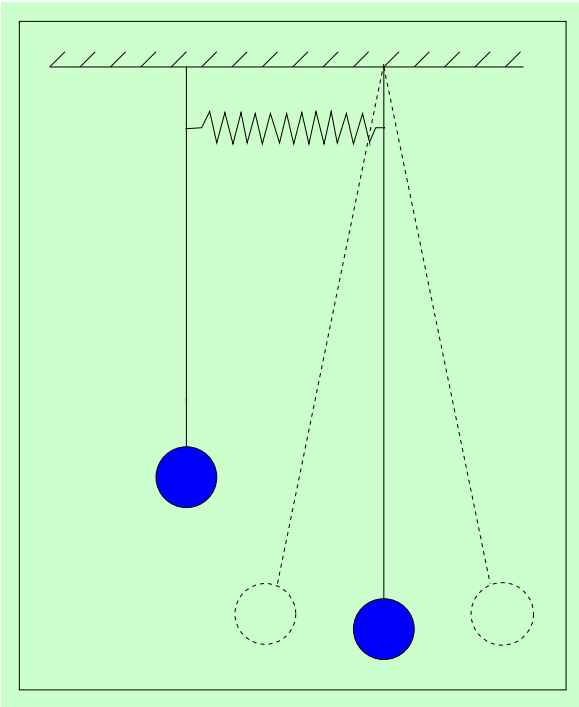
Mechanical model of the MSW effect

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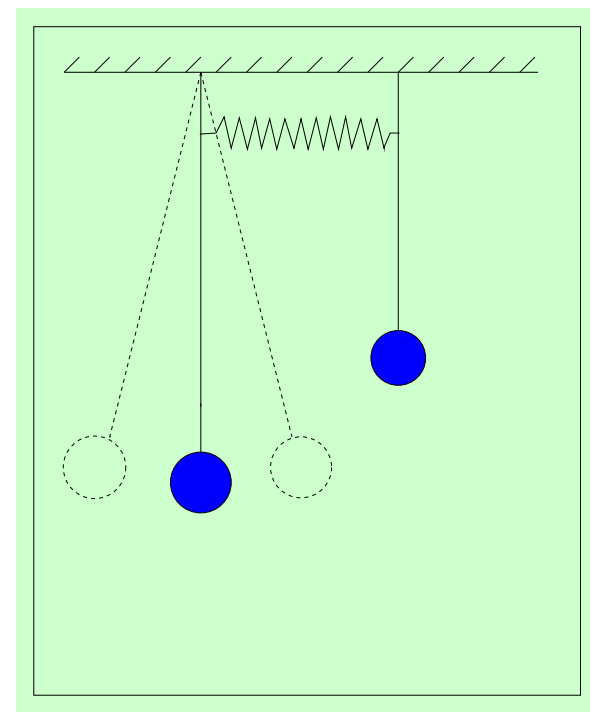
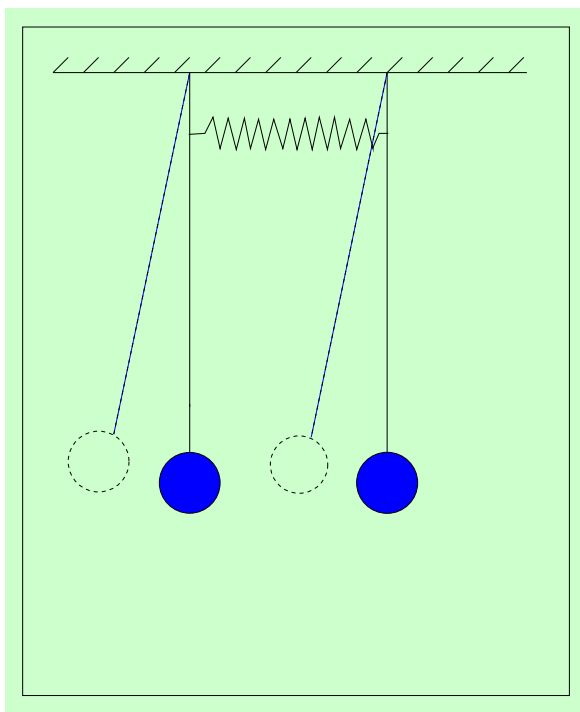
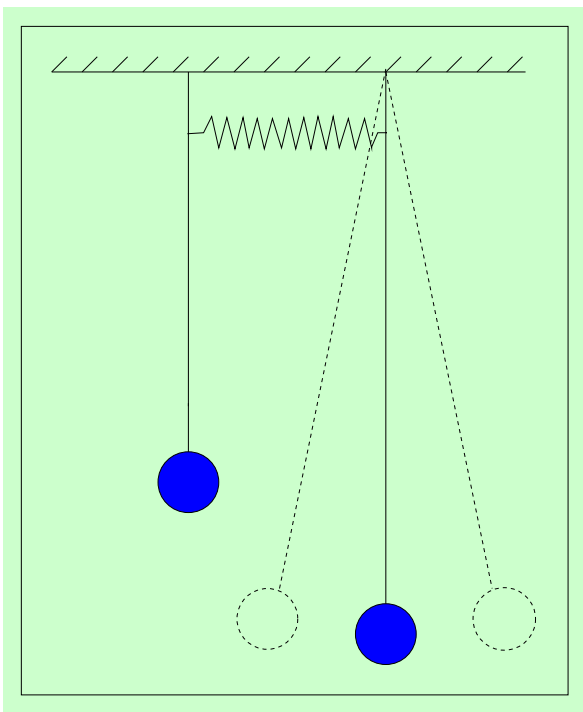
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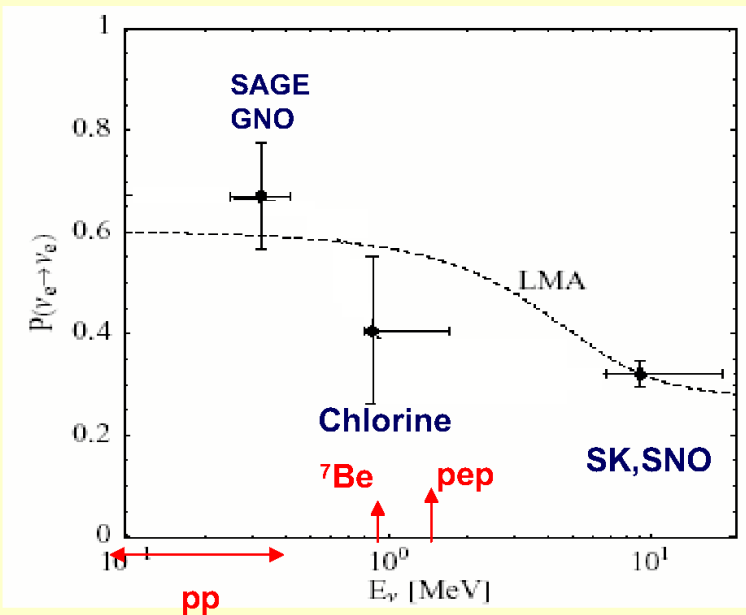


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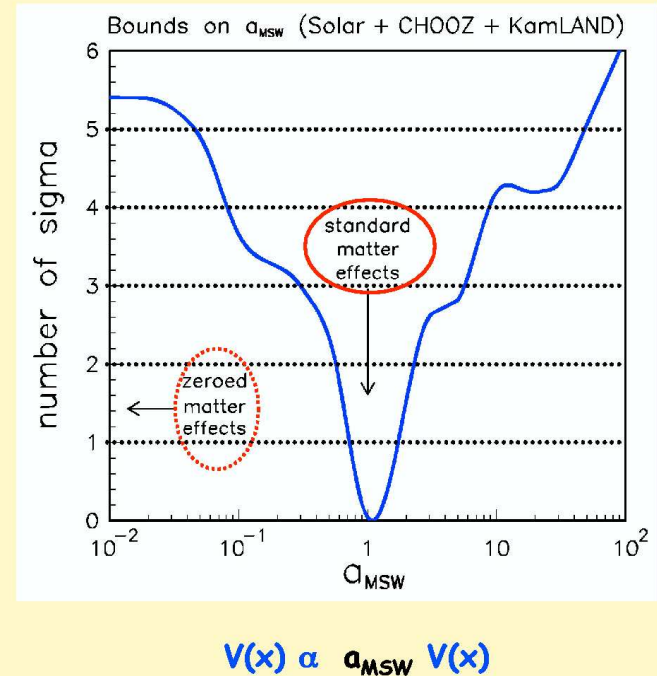
Evidence for the MSW effect

Matter Interaction Effect:LMA

Current Data for ν_e Survival



matter effects with standard size ($V = \alpha^2 G_F N_e$) confirmed



$V(x) \Rightarrow a_{\text{MSW}} V(x); a_{\text{MSW}} = 1$ strongly favoured

(Fogli et al. 2003, 2004; Fogli & Lisi 2004)

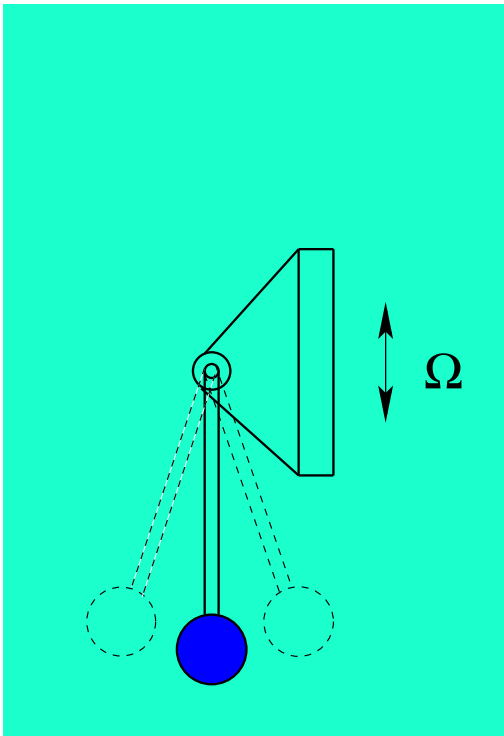
Another possible matter effect

Parametric resonance in neutrino oscillations

Parametric resonance in oscillating systems with varying parameters: occurs when the rate of the parameter change is correlated in a certain way with the values of the parameters themselves

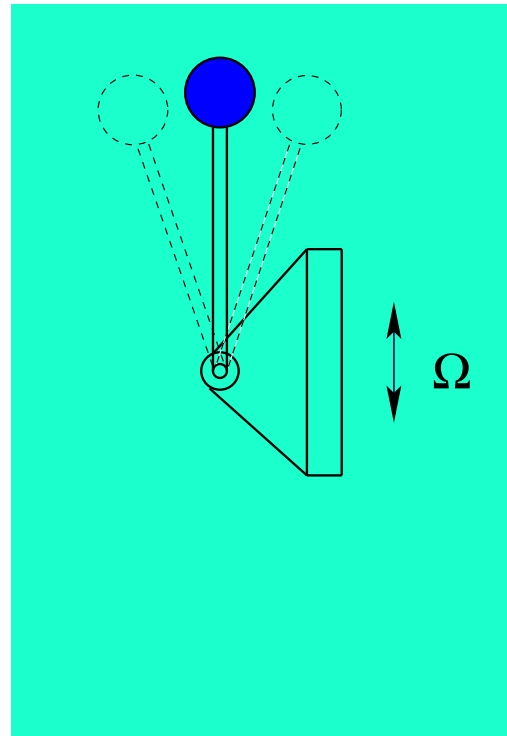
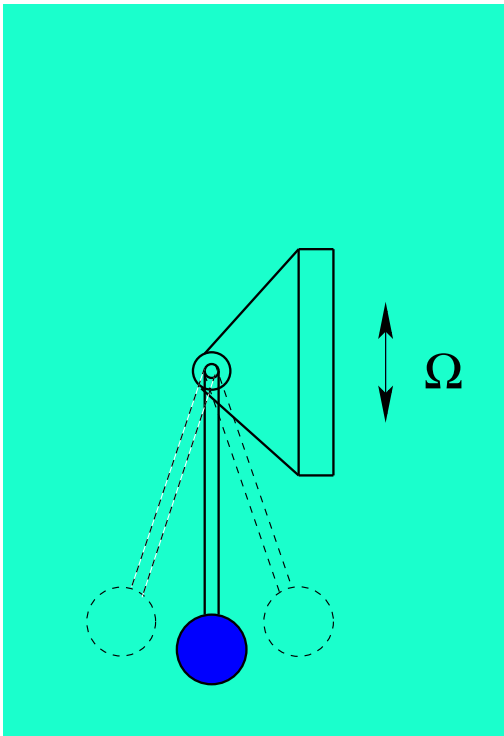
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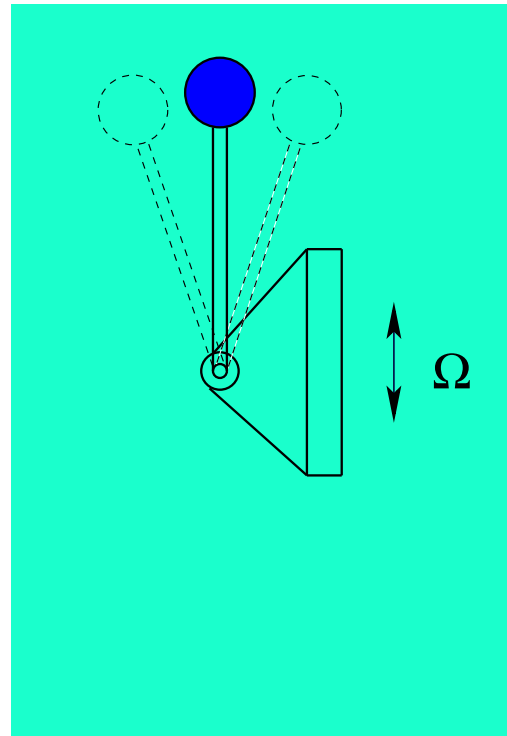
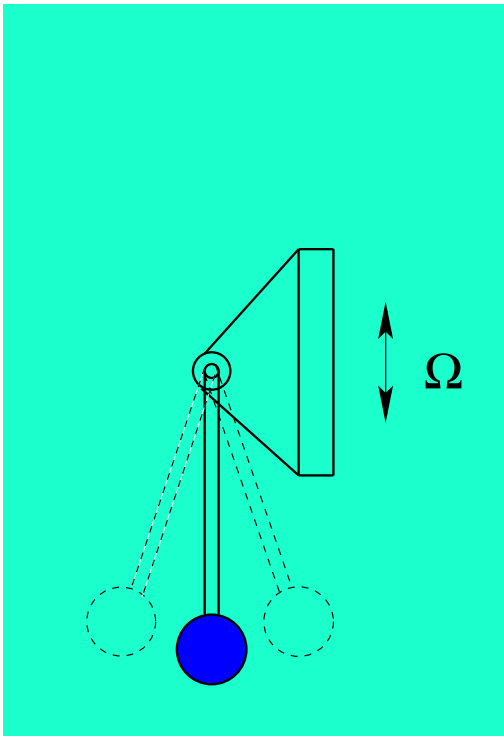
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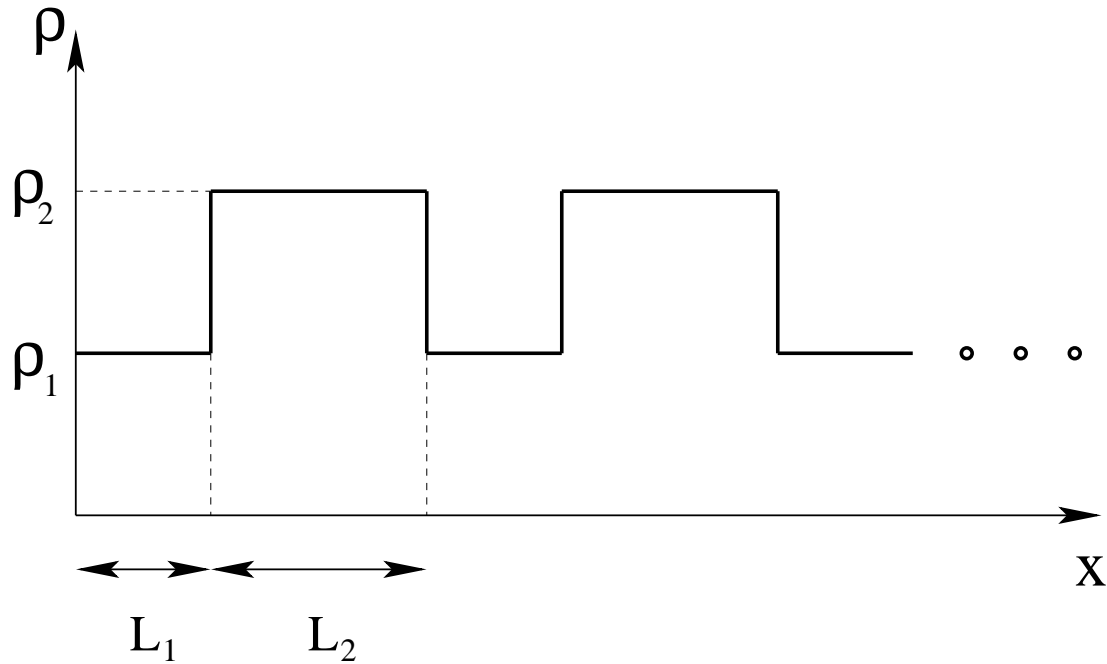
For small-ampl. osc.:

$$\Omega_{\text{res}} = \frac{2\omega}{n}$$

$$n = 1, 2, 3 \dots$$

Different from MSW eff. – no level crossing !

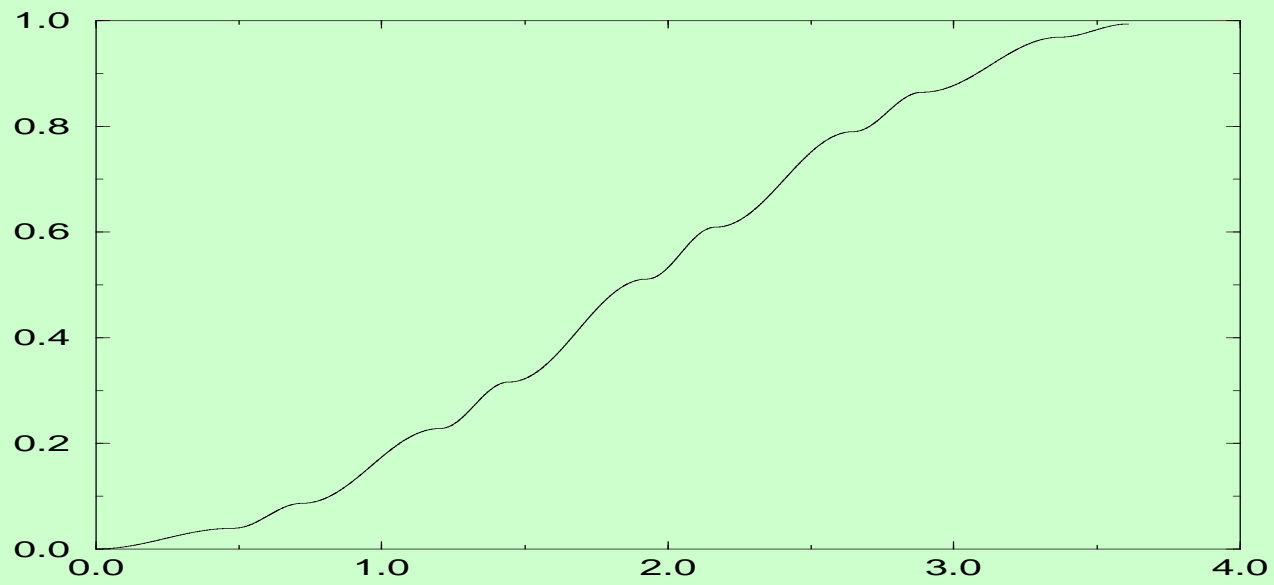
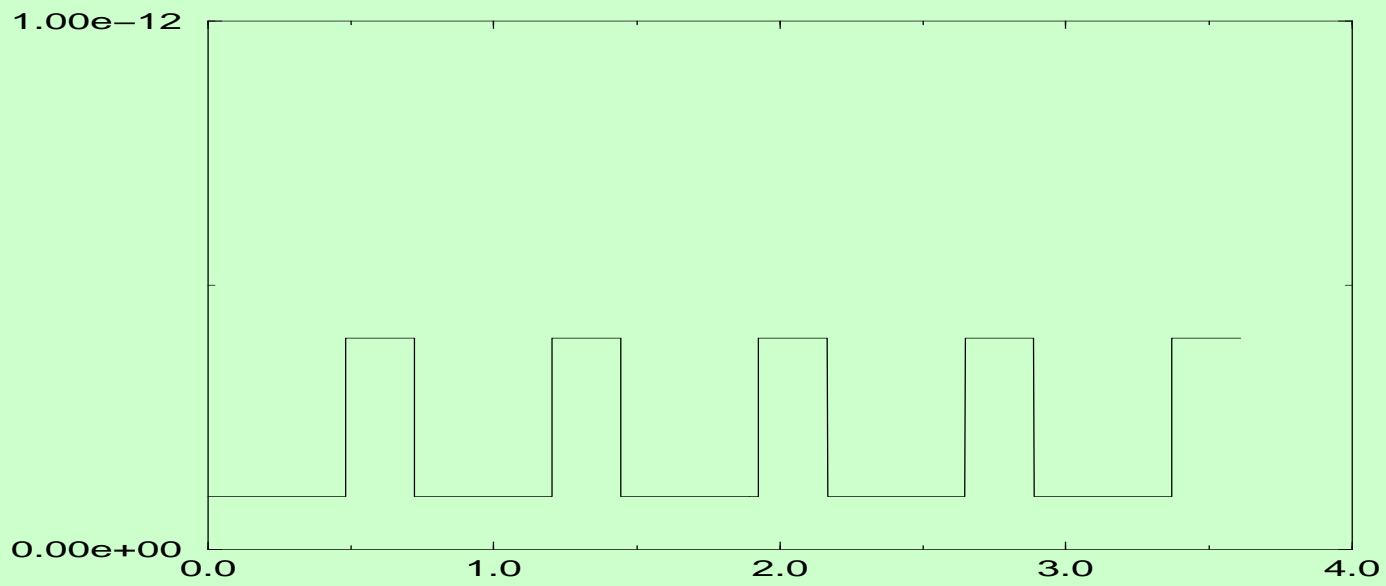
An example admitting an exact analytic solution – “castle wall” density profile (E.A., 1987, 1998):



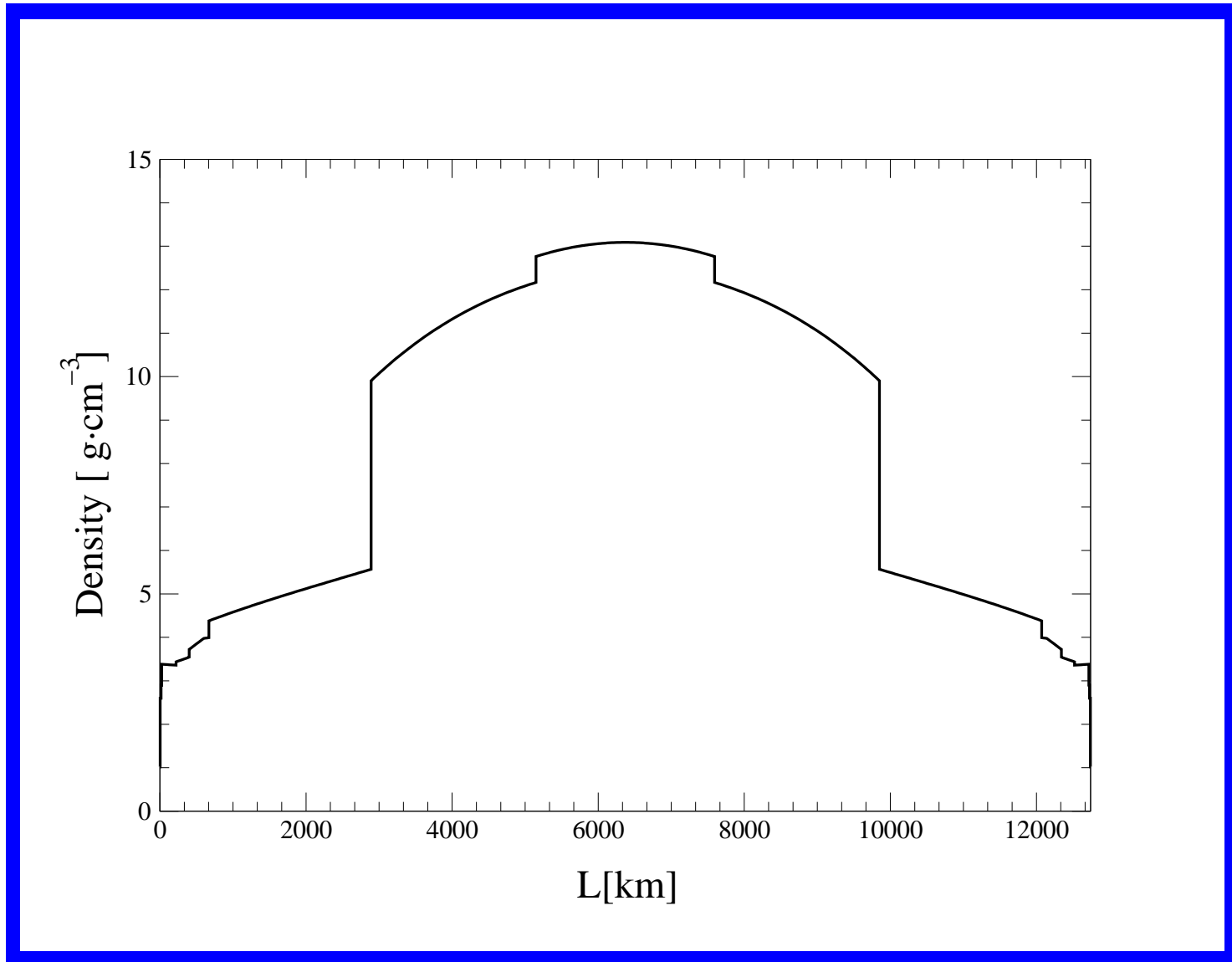
Resonance condition:

$$X_3 \equiv -(\sin \phi_1 \cos \phi_2 \cos 2\theta_{1m} + \cos \phi_1 \sin \phi_2 \cos 2\theta_{2m}) = 0$$

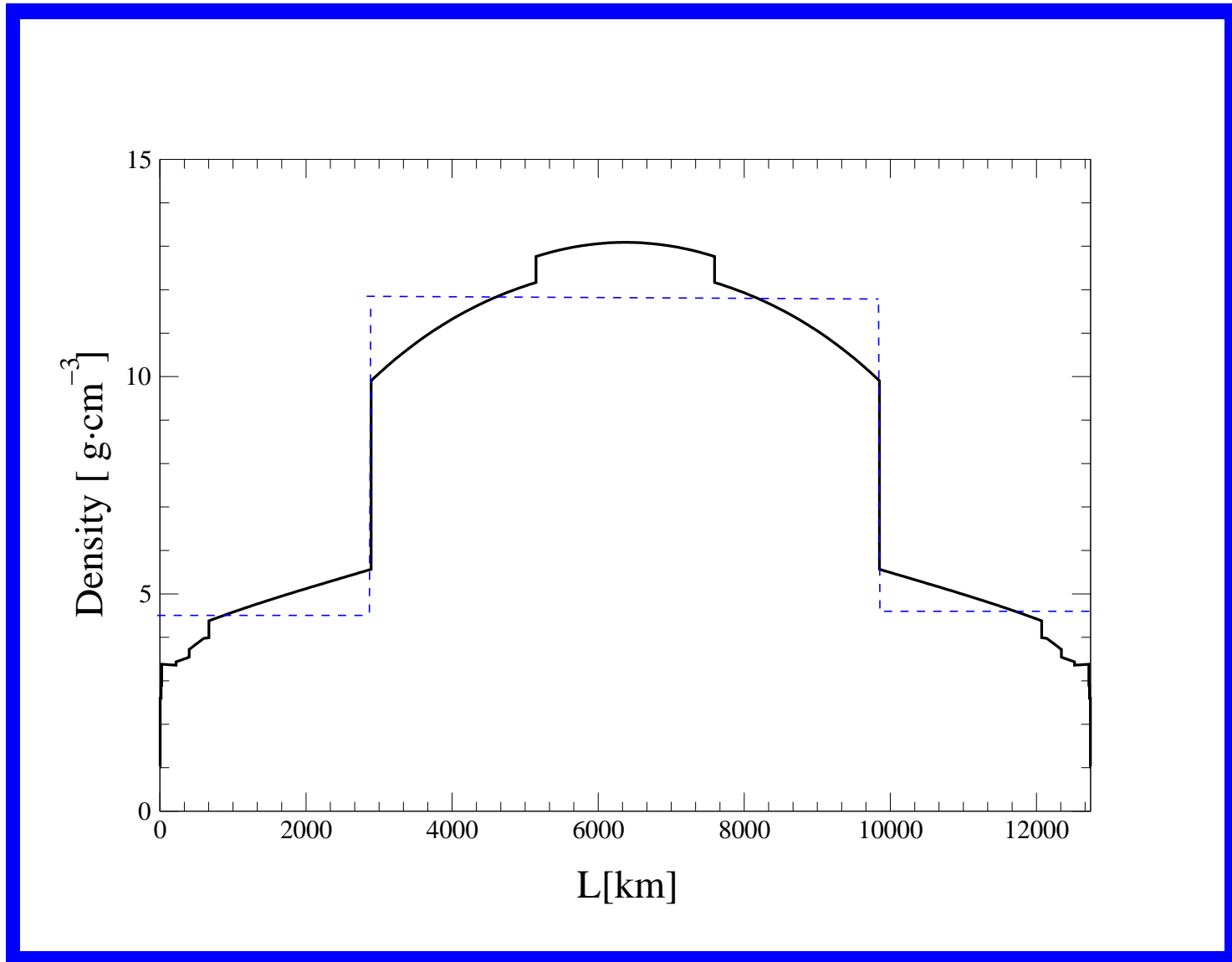
$\phi_{1,2}$ – oscillation phases acquired in layers 1, 2



Earth's density profile (PREM model) :

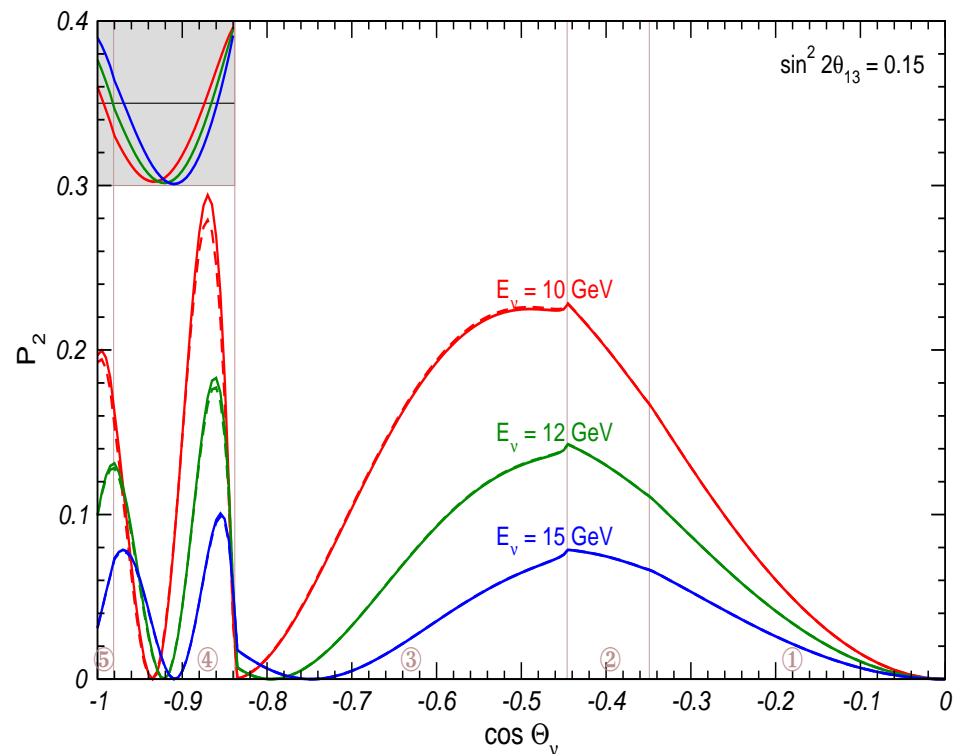
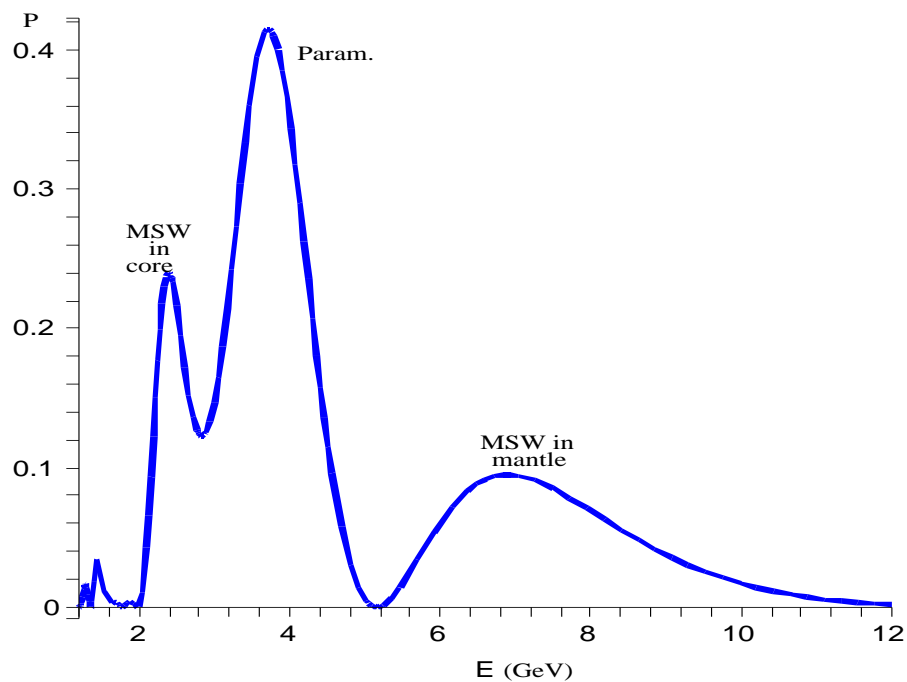


Earth's density profile (PREM model) :



Param. res. condition: $(l_{\text{osc}})_{\text{matt}} \simeq l_{\text{density mod.}}$

Fulfilled for $\nu_e \leftrightarrow \nu_{\mu,\tau}$ oscillations of core-crossing ν 's in the Earth for a wide range of energies and zenith angles !



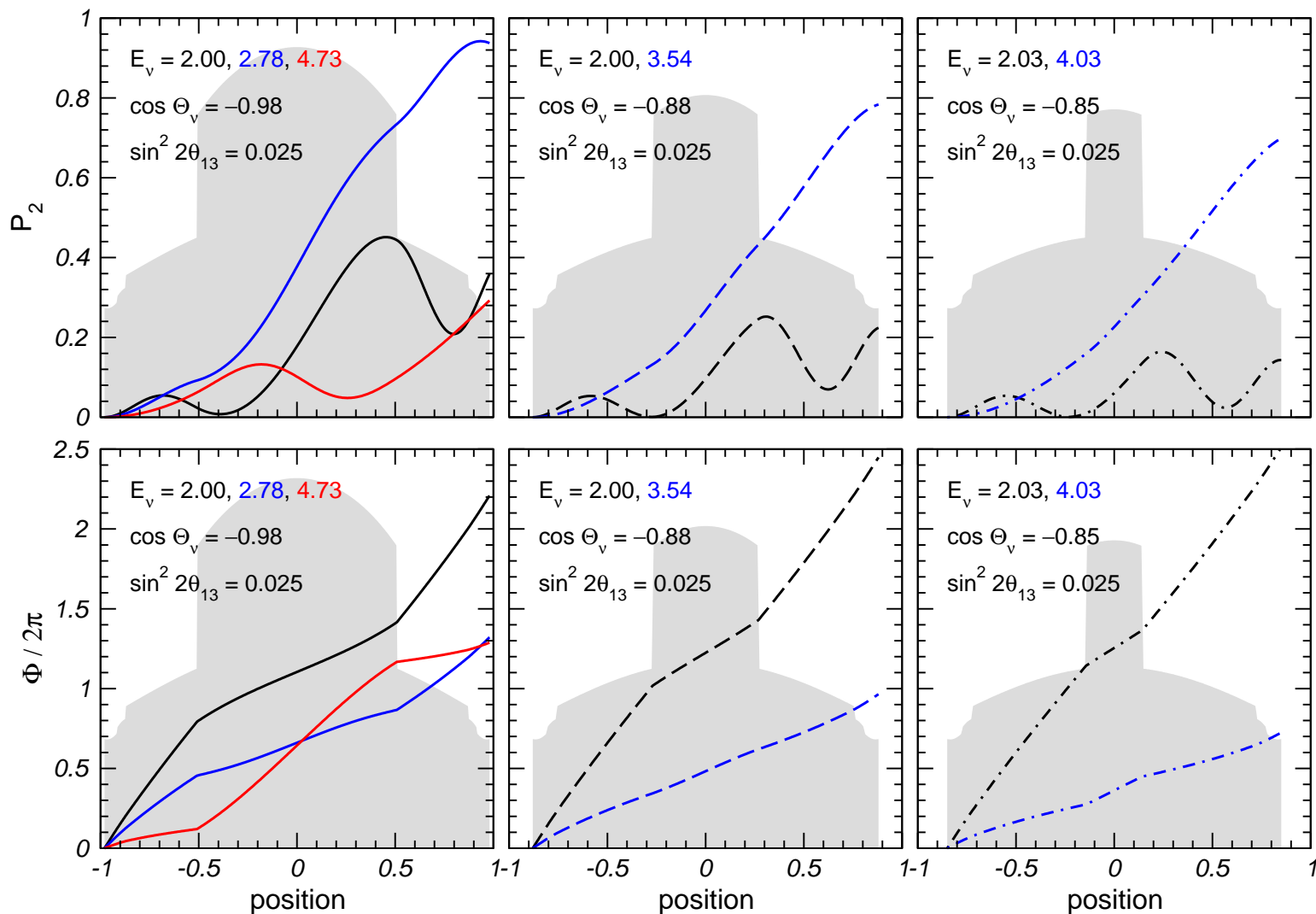
Intermed. energies

$$\cos \Theta = -0.89 \quad \sin^2 2\theta_{13} = 0.01$$

(Liu, Smirnov, 1998; Petcov, 1998; E.A. 1998)

High energies, $\cos \Theta$ - dependence

(E.A., Maltoni & Smirnov, 2005)



◇ Parametric resonance of ν oscillations in the Earth:
 can be observed in future atmospheric or accelerator
 experiments if θ_{13} is not much below its current upper limit

Neutrino oscillations in the Earth

A coherent description in terms of different realizations of just 2 conditions
– amplitude and phase conditions

Matter with $N_e = \text{const}$:

- amplitude condition = MSW resonance condition
- phase condition: $\phi = \pi/2 + \pi n$

3 layers of constant densities (or the “castle wall” density profile):

- amplitude condition = parametric resonance condition ($X_3 = 0$)
- phase condition: $\Phi \equiv \arccos Y = \pi/2 + \pi n$

Evolution matrix: $\nu(t) = U(t, t_0) \nu(0)$. For 2 layers:

$$U^{(2)}(t, t_0) = \begin{pmatrix} Y - iX_3 & -i(X_1 - iX_2) \\ -i(X_1 + iX_2) & Y + iX_3 \end{pmatrix}, \quad Y^2 + \mathbf{X}^2 = 1$$

The meaning of the amplitude condition

Alignment of the transitions amplitudes in different layers.

Evolution matrices for individual layers:

$$U_i(t, t_0) = \begin{pmatrix} \alpha_i & \beta_i \\ -\beta_i^* & \alpha_i^* \end{pmatrix}, \quad |\alpha_i|^2 + |\beta_i|^2 = 1, \quad i = 1, 2, 3$$

For 2 layers: $U^{(2)} = U_2 U_1$,

$$\beta^{(2)} = \alpha_2 \beta_1 + \beta_2 \alpha_1^*$$

Alignment (collinearity) condition:

$$\arg(\alpha_2 \beta_1) = \arg(\beta_2 \alpha_1^*) \quad \text{mod } (\pi)$$

– potentially leads to maximal trans. probability.

For 2 layers of const. densities: **align. cond.** $\Leftrightarrow s_1 s_2 X_3 = 0$

How about 3 layers?

$$U^{(3)} = U_3 U_2 U_1. \quad \text{For the Earth, } U^{(3)} = U_1^T U_2 U_1.$$

Transition amplitude:

$$\beta^{(3)} = \alpha_1 \alpha_2 \beta_1 - \alpha_1^* \alpha_2^* \beta_1^* + |\alpha_1|^2 \beta_2 + |\beta_1|^2 \beta_2^*$$

⇒ If the 2-layer align. cond. is satisfied, so is the 3-layer one !

A consequence of

- The symmetry of the core density profile
- The symmetry of the overall density profile of the Earth (3rd layer's profile is the reverse of the 1st layer's one)

⇒ The generalized amplitude condition is the alignment condition in the case of non-constant density layers

Generalized phase condition

For constant density matter: $\phi = \pi/2 + \pi n \Leftrightarrow \text{Im } \alpha^{(1)} \beta^{(1)*} = 0.$

\Rightarrow Generalize to an arbitrary density profile:

$$\text{Im } \alpha \beta^* = 0 \Leftrightarrow \frac{dP_{\text{tr}}}{dL} = 0$$

The whole complex oscillation pattern:

- MSW resonances
- parametric resonances
- saddle points
- local maxima and minima
- absolute maxima and minima

can be understood in terms of the generalized amplitude and phase conditions ! (E.A., Maltoni & Smirnov, 2007)

Neutrino oscillograms of the Earth

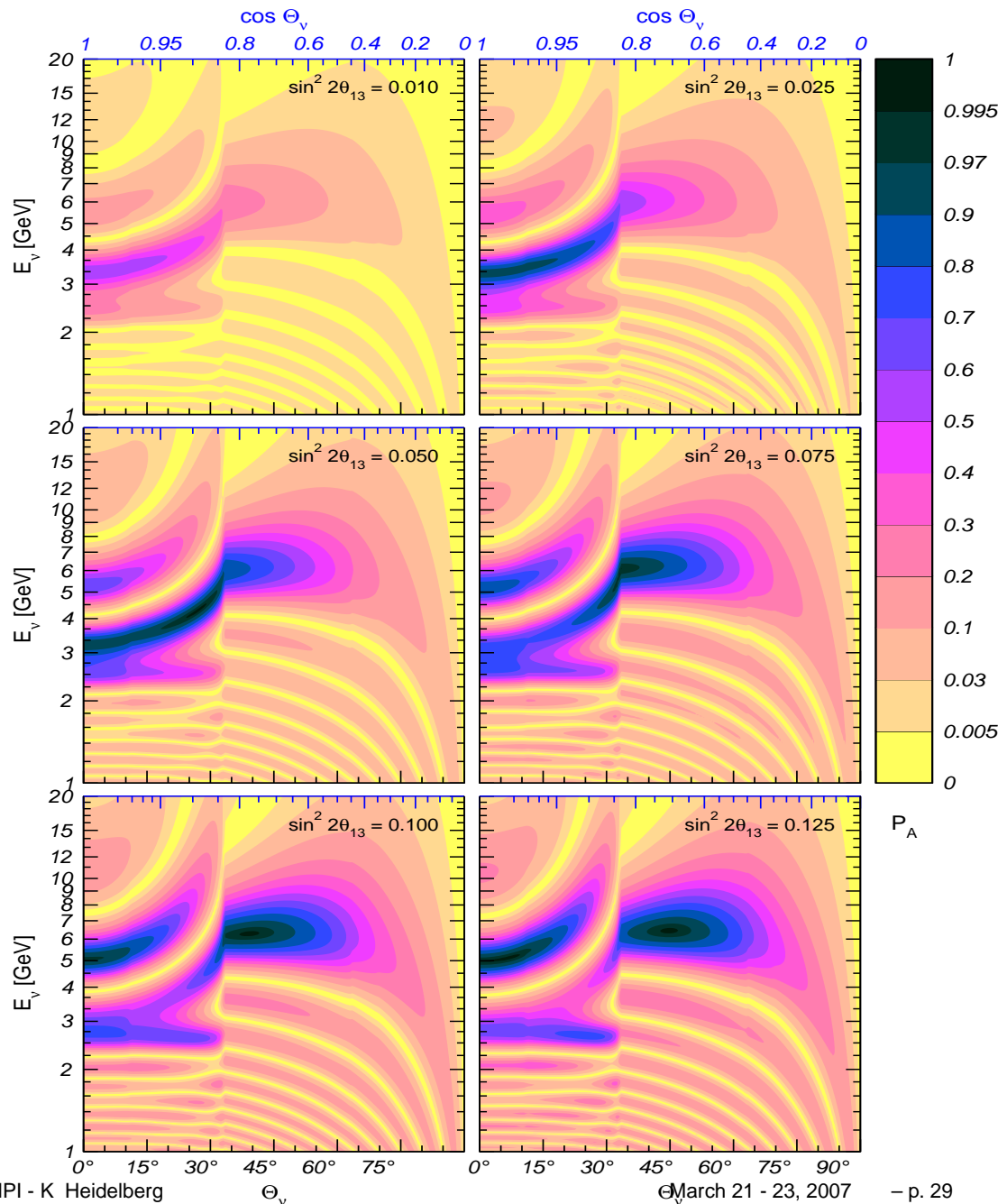
Contours of equal osc. probabilities in (Θ_ν, E_ν) plane

Θ_{13} - dependence of $P_A \Rightarrow$




P_A – effective 2f transition probability ($\Delta m_{\text{sol}}^2 \rightarrow 0$)

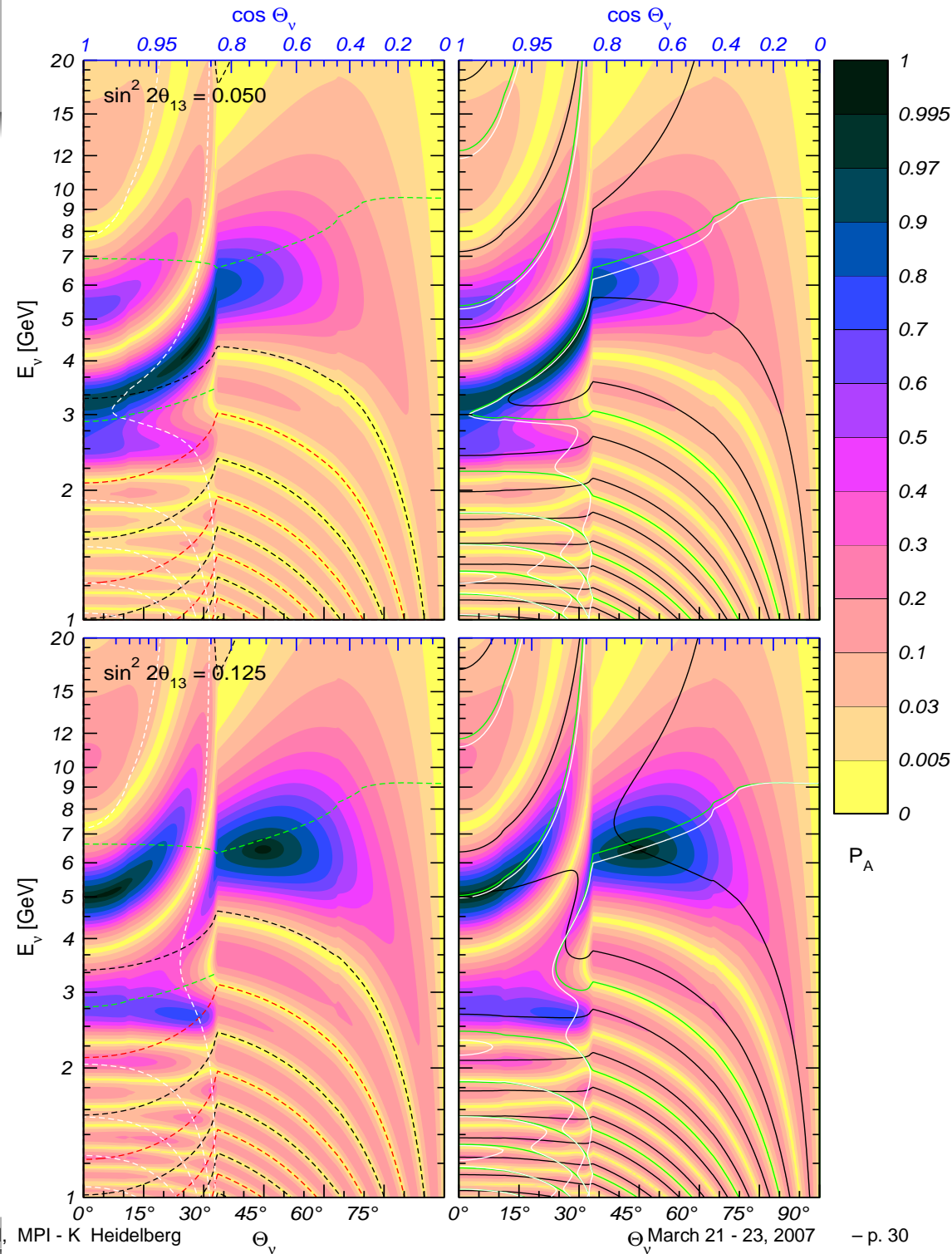
$$P_{e\mu} = s_{23}^2 P_A$$

$$P_{e\tau} = c_{23}^2 P_A$$



In the right panels:

-  — alignment
(collin.) cond.
-  — generalized
res. cond.
($\text{Im}\alpha^{(2)} = 0$)
-  — generalized
phase cond.



Including the effects of Δm_{sol}^2

Fundamental \mathcal{CP} and \mathcal{T} ; dependence of P_{ab} on δ_{CP}
(also in CP - and T - even terms) \Rightarrow
parameter correlations and degeneracies (e.g. θ_{13} and δ_{CP})

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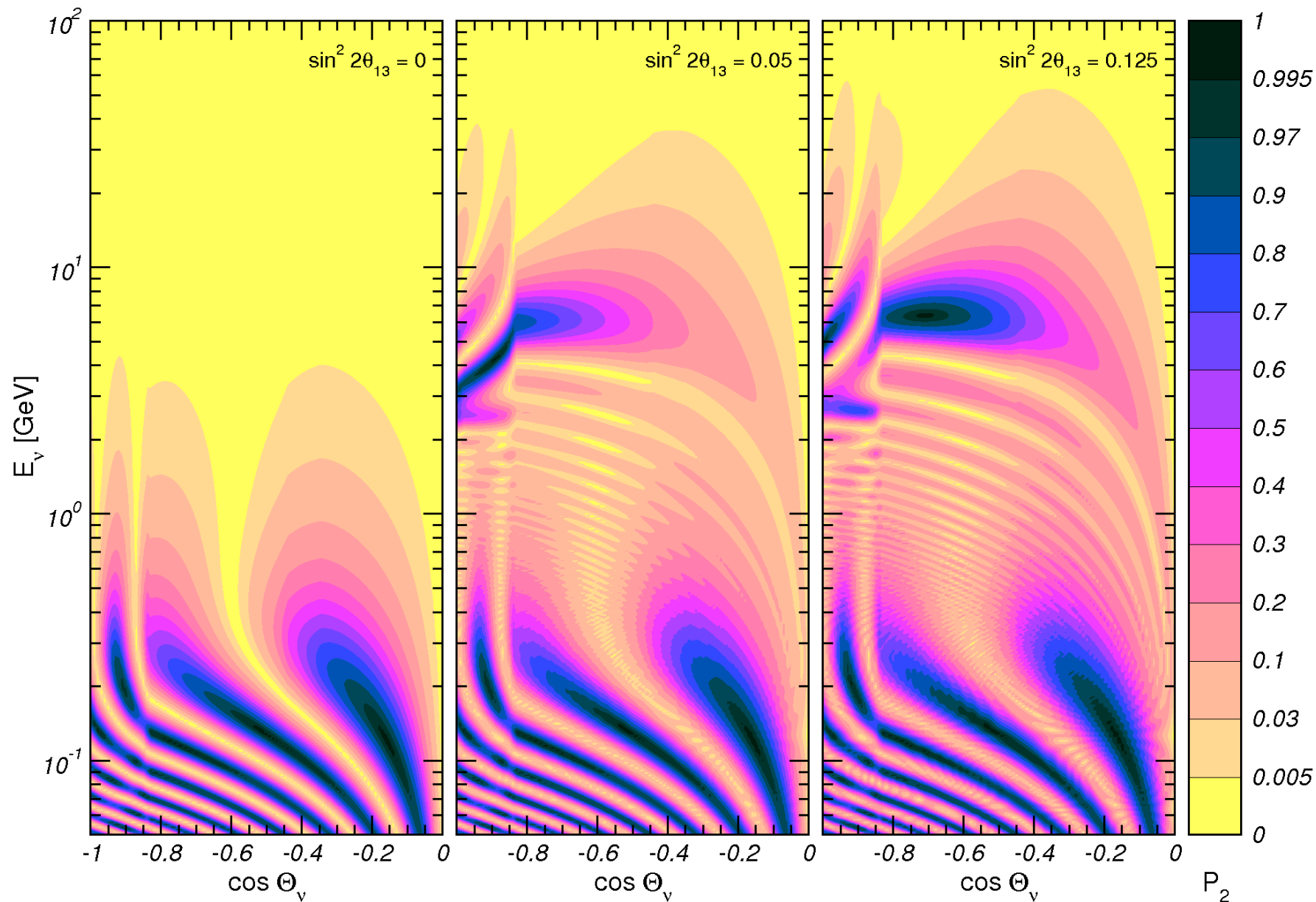
$\Rightarrow L_{\text{magic}} : \quad \phi = \pi n$

At high energies:

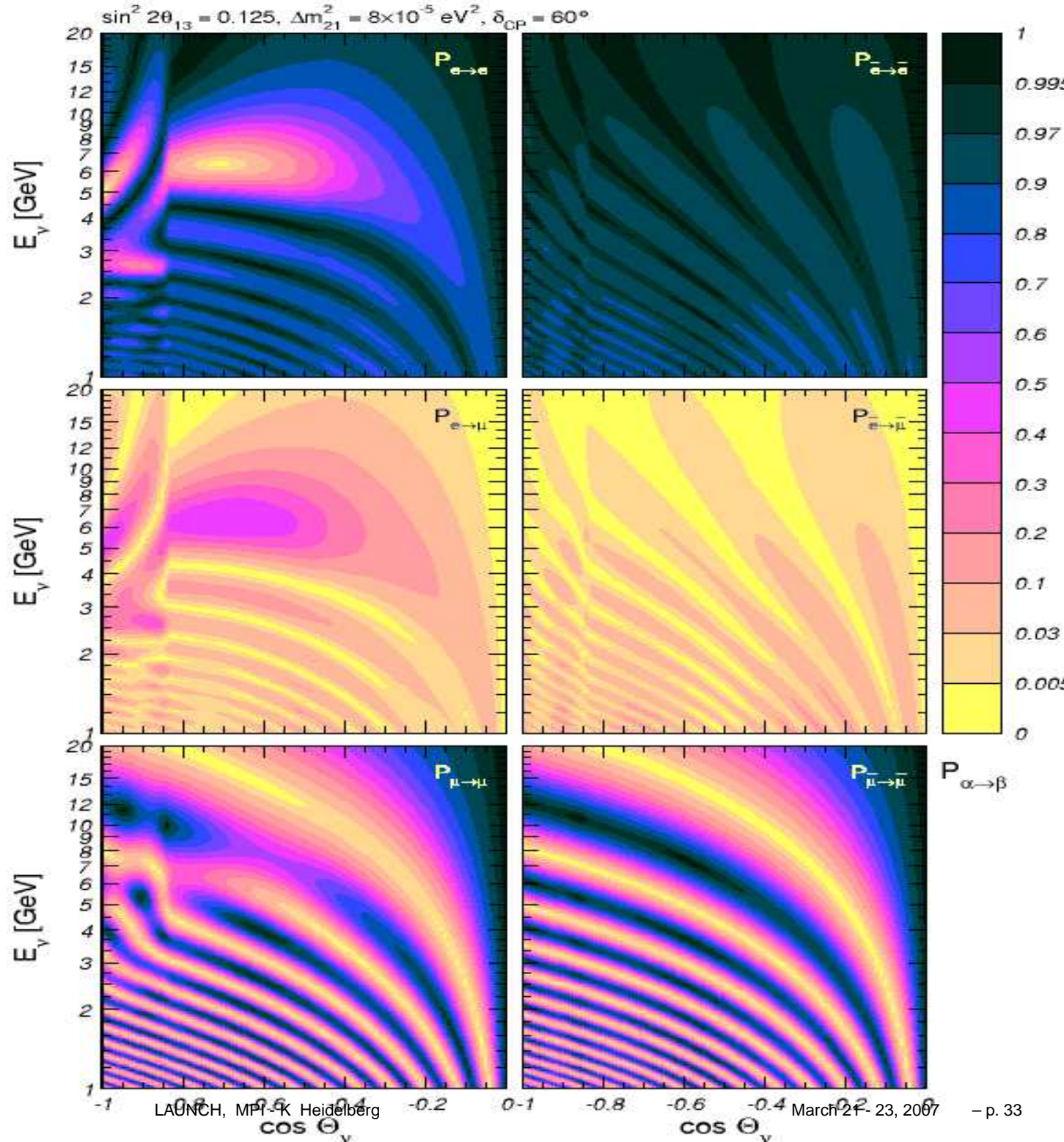
$$VL = 2\pi n$$

$$\Leftrightarrow L = l_{\text{refr.}} n$$

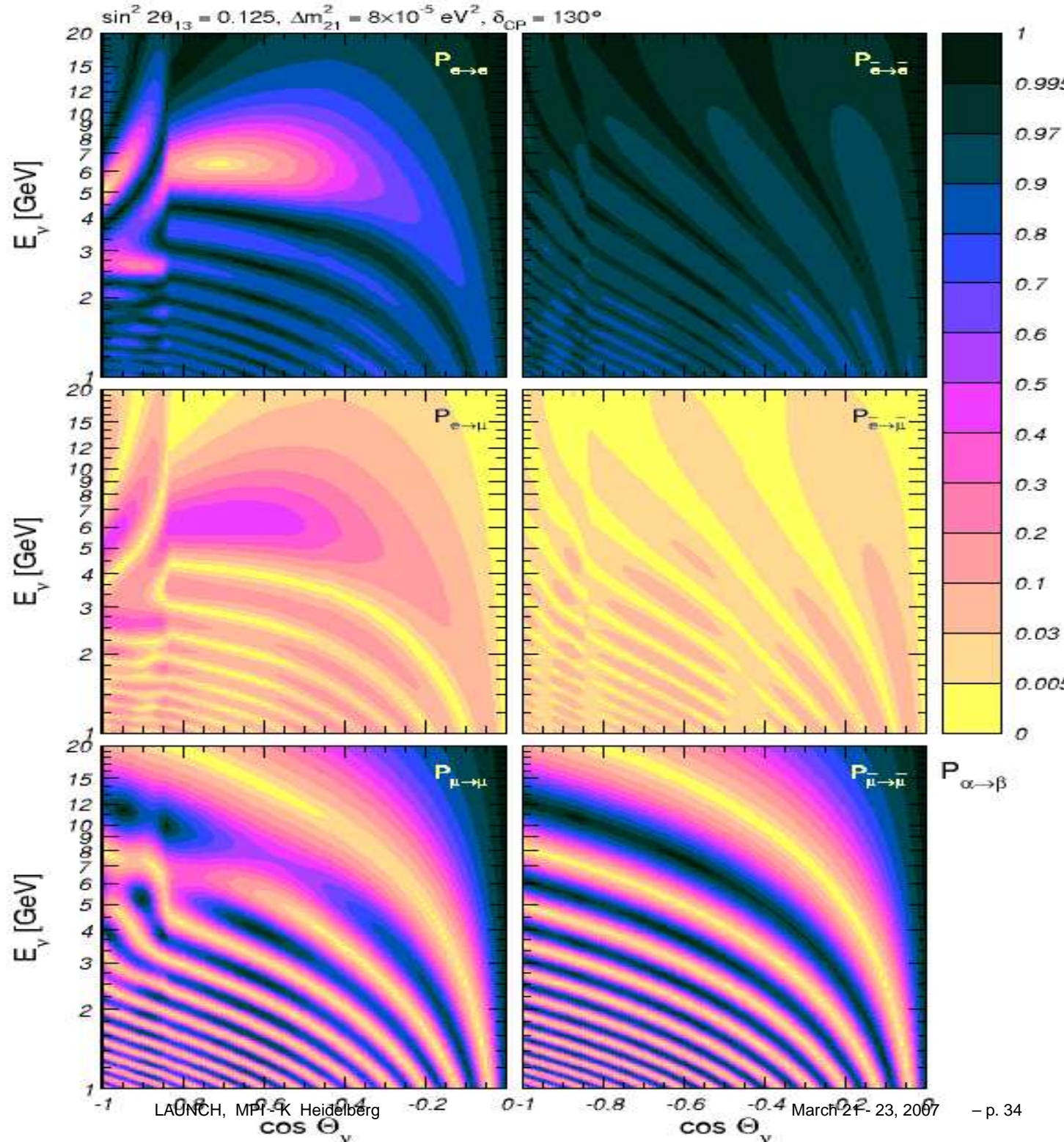
Including the effects of $\Delta m_{\text{sol}}^2 : (1 - P_{ee})$



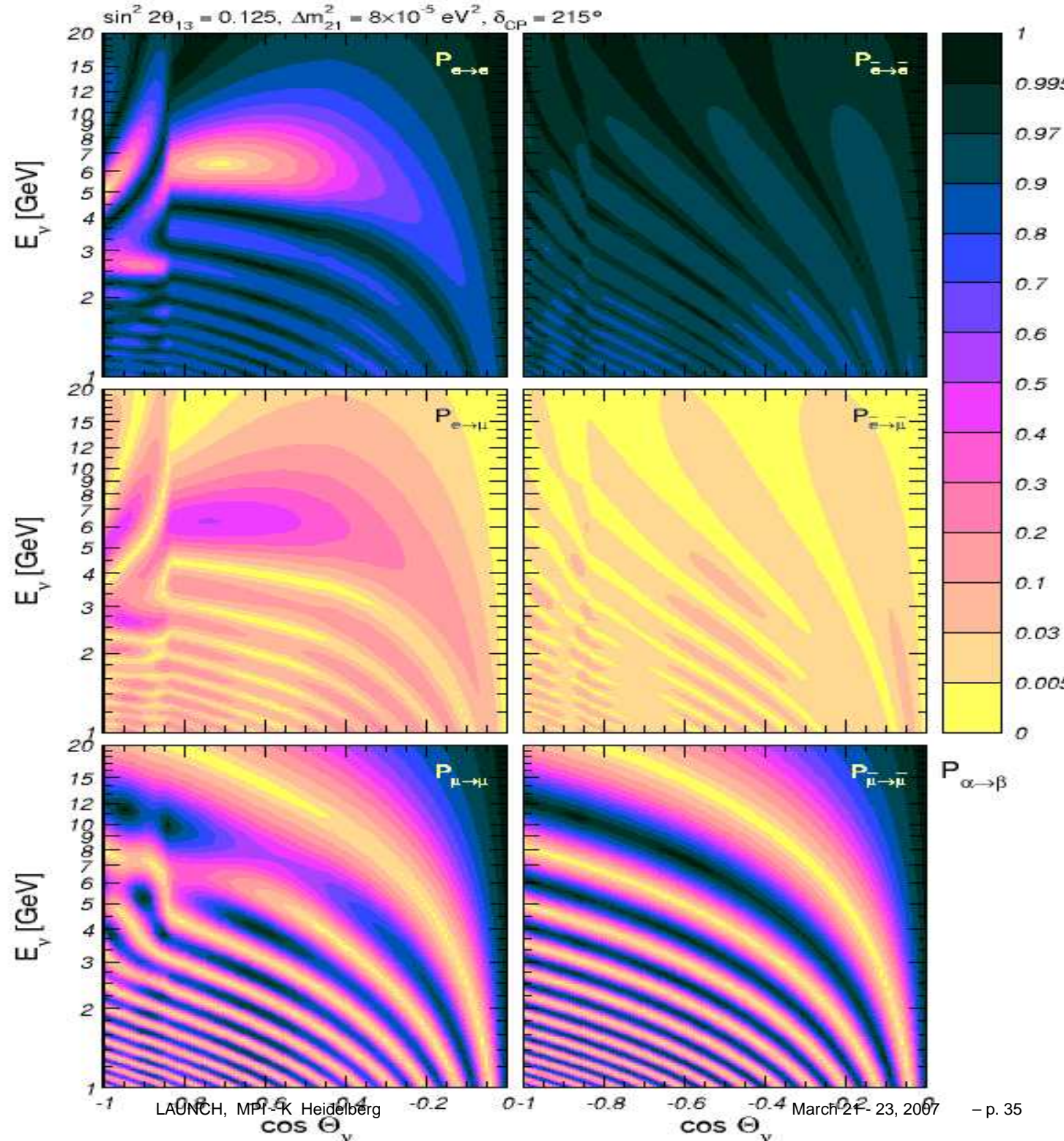
Dependence on δ_{CP}



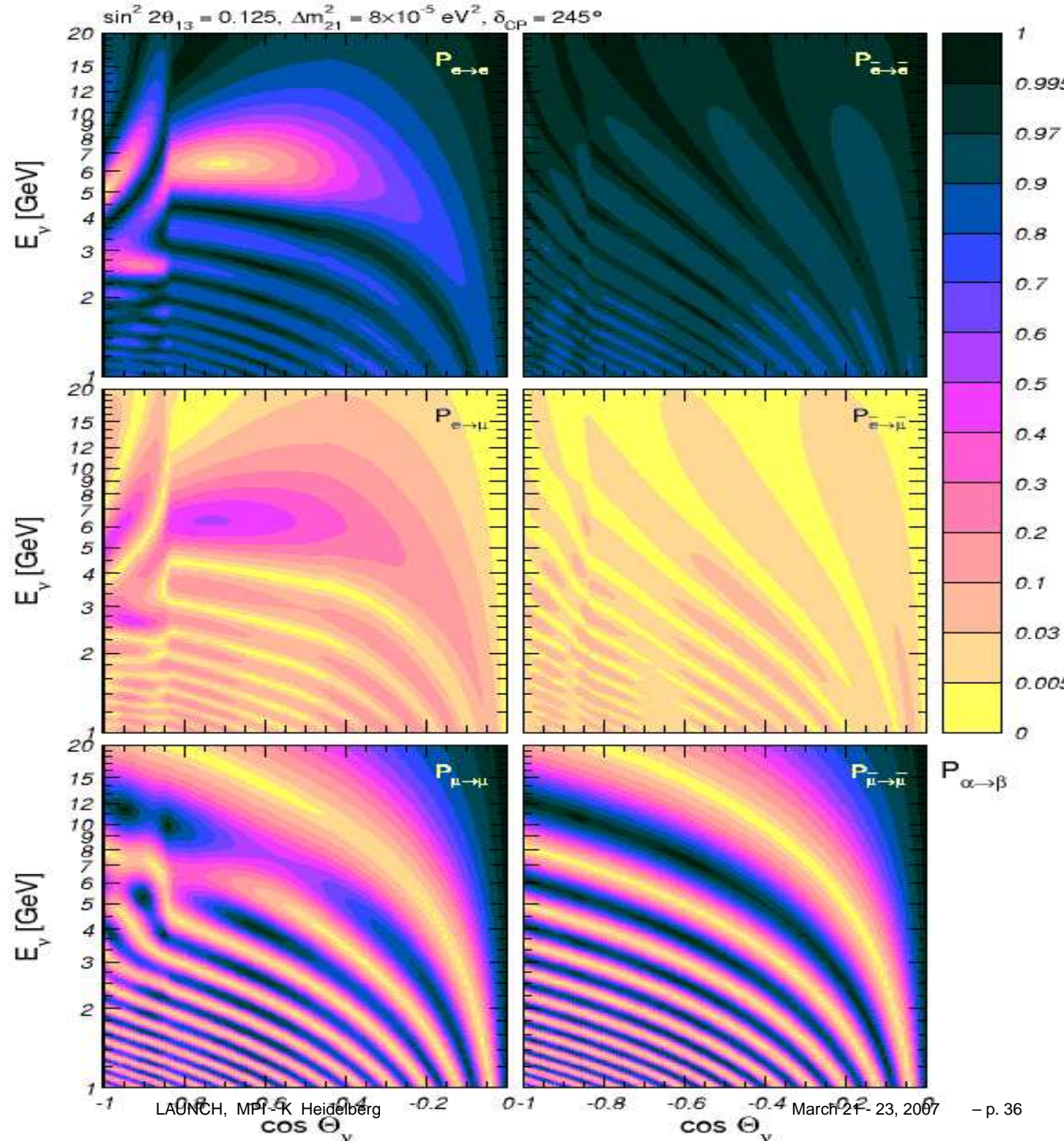
Dependence on δ_{CP}



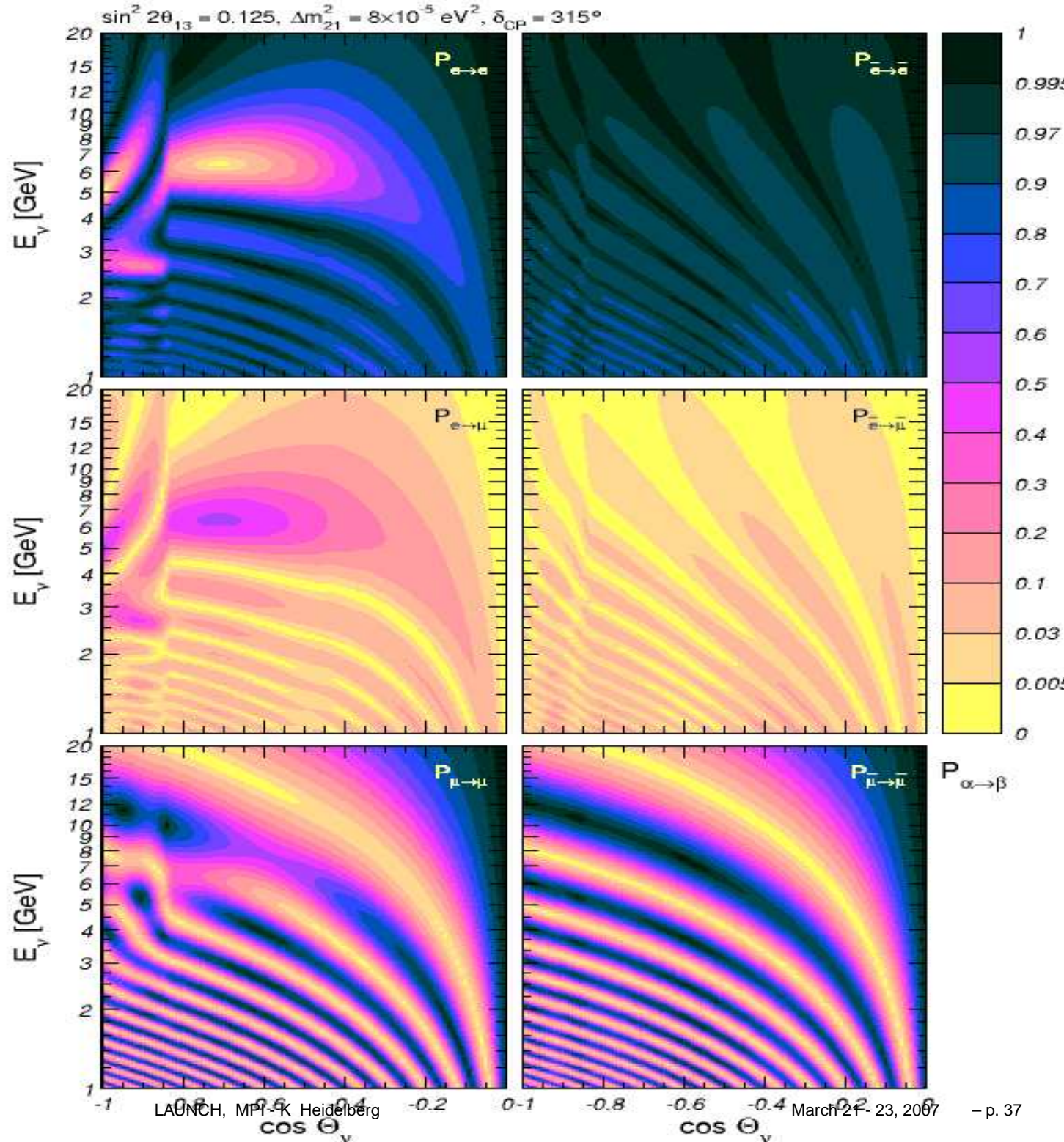
Dependence on
 δ_{CP}



Dependence on δ_{CP}



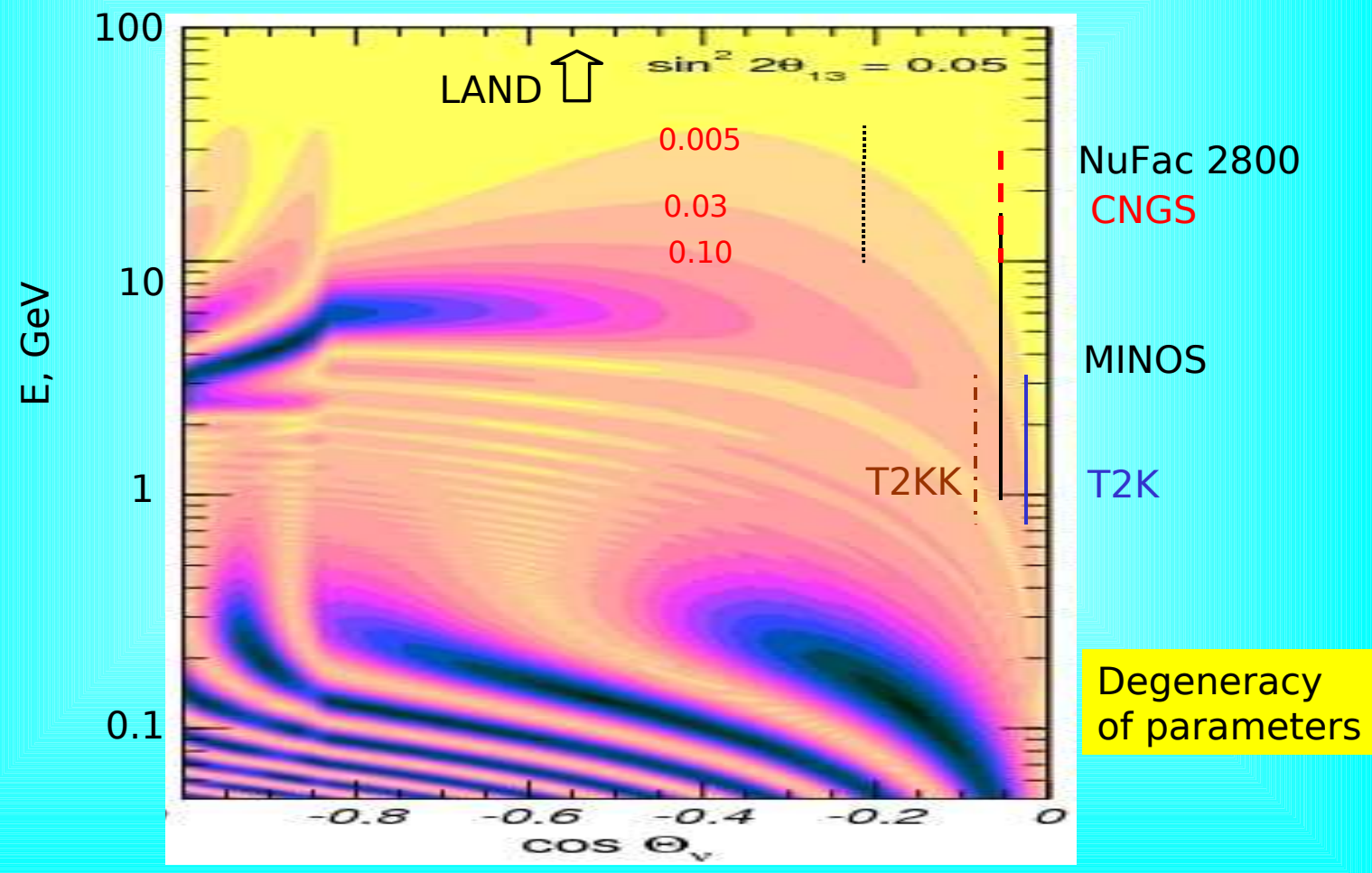
Dependence on δ_{CP}



Producing the oscillograms

Accelerators

Large atmospheric neutrino detectors



A. Smirnov, UCLA seminar

Producing the oscillograms – contd.

Huge atmospheric neutrino detectors may be necessary!

Would require :

- Very good energy and angle resolution
- Low threshold ($E_{\text{thr.}} \sim 3 \text{ GeV}$)
- Charge discrimination (?)
- High statistics

Very ambitious, but the gain may be overwhelming \Rightarrow

It is worth studying the oscillograms with Huge Atmospheric Neutrino Detectors !

Conclusions

- 3f corrections to 2f oscillation probabilities can reach $\sim 10\%$
 - at the level of current experimental sensitivity. Depend on $|U_{e3}| = \sin \theta_{13}$
- A number of interesting pure 3f effects exist – fundamental CP and T-violation, matter - induced T violation, matter effects on $\nu_{\mu} \leftrightarrow \nu_{\tau}$ oscillations
- Matter can strongly affect ν oscillations inside the Earth through the MSW and parametric resonance effects
- Neutrino oscillograms of the Earth carry a wealth of information both on neutrinos and the Earth:

Conclusions

They :

- Depend strongly on the neutrino mass hierarchy and the value of θ_{13}
- Depend sensitively on the \mathcal{CP} phase $\delta_{\mathcal{CP}}$ and on the Earth density profile
- Their specific structures (MSW resonances, parametric ridges, local and global extrema, saddle points) and their dependence on ν parameters can be fully described in terms of the amplitude and phase conditions
- This can be used for looking for best strategies for future ν oscillations experiments