Modelling of blazar SEDs with the nonlinear SSC cooling process

Michael Zacharias & Reinhard Schlickeiser

Institut für Theoretische Physik, Lehrstuhl IV: Weltraum- und Astrophysik, Ruhr-Universität Bochum, Germany



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Figure 1: SED of 3C 454.3 during a large outburst in November 2010 (Vercellone et al. 2011)

- Numerical calculations of SEDs implement time-dependent nature of SSC cooling
- Most modeling attempts fail to recognize the important effects



• Electron kinetic equation:

$$\frac{\partial n(\gamma, t)}{\partial t} - \frac{\partial}{\partial \gamma} \left[|\dot{\gamma}|_{\text{TOT}} n(\gamma, t) \right] = Q(\gamma, t)$$

• Source term:

$$Q(\gamma, t) = q_0 \delta(\gamma - \gamma_0) \delta(t)$$

• Description of flares, i.e. single injection of relativistic $(\gamma_0\gg 1)$ particles

Not useful for steady-state sources



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• Cooling rate:

$$|\dot{\gamma}|_{\text{TOT}} = |\dot{\gamma}|_{\text{SYN}} + |\dot{\gamma}|_{\text{SSC}} = D_0 \gamma^2 + A_0 \gamma^2 \int_0^\infty d\gamma \, \gamma^2 n(\gamma, t)$$

- SSC cooling several orders of magnitude quicker than synchrotron cooling
- SSC cooling depends on electron distribution (i.e. time-dependent) \Rightarrow synchrotron cooling dominates after some time



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Injection parameter α



• Defined as the ratio of the cooling terms at time of injection (t = 0)

$$\alpha = \left[\frac{|\dot{\gamma}(t=0)|_{\rm SSC}}{|\dot{\gamma}|_{\rm SYN}}\right]^{1/2} = \gamma_0 \left(\frac{q_0 A_0}{D_0}\right)^{1/2}$$

• $\alpha \ll 1$: linear cooling \Rightarrow dominance of the synchrotron peak • $\alpha \gg 1$: initial nonlinear cooling \Rightarrow dominance of the IC peak



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Figure 2: Model-SED for $\alpha \ll 1$ in the Thomson-limit





Figure 3: Model-SED for $\alpha \gg 1$ in the Thomson-limit

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- The synchrotron peak shows a characteristic feature for $\alpha \gg 1$ (without the need for fancy electron distributions)
- Different electron distributions affect only the high-energy end of the peaks (MZ & RS 2010)





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Standard scenario:

• Efficient accretion disk

- $\bullet \Rightarrow \mathsf{Strong} \ \mathsf{BLR} \ \mathsf{and} \ \mathsf{torus}$
- $\bullet \Rightarrow \mathsf{Lots} \text{ of seed photons}$
- \Rightarrow Strong EC cooling of the electrons
- ⇒ Dominance of the IC peak



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- \Rightarrow High electron density in the jet
- ⇒ Strong nonlinear SSC cooling of the electrons
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- \Rightarrow Dominance of the IC peak



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Standard scenario:

New scenario:

- Efficient accretion disk
- $\bullet \ \Rightarrow \mathsf{Strong} \ \mathsf{BLR} \ \mathsf{and} \ \mathsf{torus}$
- \Rightarrow Lots of seed photons
- \Rightarrow Strong EC cooling of the electrons
- ⇒ Dominance of the IC peak

- Efficient accretion disk
- $\bullet \ \Rightarrow \ {\rm High \ electron \ density \ in} \\ the \ jet$
- → Strong nonlinear SSC cooling of the electrons
- ⇒ Dominance of the IC peak

\Rightarrow Potentially both processes are equally important

External Compton (first results)





Figure 4: Model-SED for $\alpha_{ec} \ll 1$ and $l_{ec} \gg 1$ in the Thomson-limit

- Inclusion of external Compton requires new parameters:
 - Relative strengths $l_{ec} = u_{ec} \Gamma_b^2 / u_B$
 - Injection parameter $\alpha_{ec}^2 = \gamma_0^2 q_0 A_0 / D_0 (1 + l_{ec})$
- For $\alpha_{ec} = \alpha$
 - \Rightarrow q_0 increases
 - $\Rightarrow \mathsf{SSC}\ \mathsf{luminosity}$
 - increases with $(1 + I_{ec})$

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Figure 5: Model-SED for $\alpha_{ec} \gg 1$ and $l_{ec} \ll 1$ in the Thomson-limit

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Figure 6: Model-SED for $\alpha_{ec} \gg 1$ and $l_{ec} \gg 1$ in the Thomson-limit

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- Dominating IC peak can be achieved by time-dependent SSC scenario
- Comparable with EC depending on parameters
- Follow-ups: Lightcurves and optical depth



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Thank you very much!

- Zacharias & Schlickeiser 2012, in prep.
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