## Experimental review on neutrinoless double beta decay

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#### Neutrinos and masses of elementary particles

• Neutrinos: much lighter than other known particles



#### Double beta decay

- The decay with emission of 2 neutrinos was observed in more than 10 different nuclei: <sup>48</sup>Ca, <sup>76</sup>Ge, <sup>82</sup>Se, <sup>96</sup>Zr, <sup>100</sup>Mo, <sup>116</sup>Cd, <sup>128</sup>Te, <sup>130</sup>Te, <sup>136</sup>Xe, <sup>150</sup>Nd, <sup>238</sup>U
- The observed energy spectrum of the two electrons is continuous, up to the Q-value

$$\Gamma^{2\nu} = \frac{1}{T_{1/2}^{2\nu}} = G^{2\nu}(Q,Z)|M^{2\nu}|^2 \qquad G^{2\nu} \propto (G_F \cos\theta_C)^4 Q^7 \left(1 + \frac{Q}{2} + \frac{Q^2}{9} + \frac{Q^3}{90} + \frac{Q^4}{1980}\right)$$



$$Q = E_{e1} + E_{e2} + E_{v1} + E_{v2} - 2m_e$$



#### Neutrinoless double beta decay

• More interesting: the decay mode without emission of neutrinos ("forbidden" in the SM, since  $\Delta L = 2$ )

$$\Gamma^{0\nu} = \frac{1}{T_{1/2}^{0\nu}} = G^{0\nu}(Q, Z) |M^{0\nu}|^2 \frac{|m_{\beta\beta}|^2}{m_e^2} \qquad n \qquad p \\ G^{0\nu} \propto (G_F \cos\theta_c)^4 \cdot \left[\frac{Q^5}{30} - \frac{2Q^2}{3} + Q - \frac{2}{5}\right] \propto (G_F \cos\theta_c)^4 \cdot Q^5 \qquad n \qquad p \\ for each or each of the decay \\ for each or each of the decay \\ g = E_{e1} + E_{e2} - 2m_e \\ Energy [keV]$$





#### Matrix elements











FIG. 7: (Color online) Neutrinoless double beta decay transition matrix elements for the different approaches: QRPA [5, 6], the SM [8–10], the projected HFB method [14] and the IBM [15]. The error bars for the QRPA are calculated as the highest and the lowest values for three different single nucleon basis sets, two different axial charges  $g_A = 1.25$  and the quenched value  $g_A = 1.00$  and two different treatments of short range correlations (Jastrow-like [25] and the Unitary Correlator Operator Method (UCOM) [26]). The radius parameter is as in this whole work  $r_0 = 1.2$  fm.

arXiv:1001.3519

#### Matrix elements: vary by a factor of 2-3 for a given A

#### Effective Majorana neutrino mass

•  $|m_{\beta\beta}|$  is a mixture of m<sub>1</sub>, m<sub>2</sub>, m<sub>3</sub>, proportional to the U<sub>ei</sub><sup>2</sup>, where U<sub>ei</sub> are complex entries

$$|m_{\beta\beta}| = |m_1|U_{e1}|^2 + m_2|U_{e2}|^2 e^{i(\alpha_1 - \alpha_2)} + m_3|U_{e3}|^2 e^{i(-\alpha_1 - 2\delta)}$$

• where U = neutrino mixing matrix,  $c_{ij} = cos\theta_{ij}$ ,  $s_{ij} = sin\theta_{ij}$ ,  $\alpha_1, \alpha_2 = Majorana$  phases



Remark: here the exchange of a light neutrinos is considered; many other contributions are possible (Majoron, heavy Majorana neutrino exchange, right-handed currents, SUSY, etc)

For a recent review, see: http://xxx.lanl.gov/pdf/ 1205.0649.pdf

#### Experimental sensitivity

• Experiments observe:

$$N^{0\nu}_{\beta\beta} = \frac{a \cdot M \cdot N_A}{A} \frac{\ln 2}{T^{0\nu}_{1/2}} \cdot \epsilon \cdot t$$

• with a non-zero number of background events:

$$N_{bg} = M \cdot t \cdot B \cdot \Delta E$$

a = enrichment

 $\epsilon$  = detector efficiency

M = total mass

- t = measuring time
- $\Delta E$  = energy resolution
- B = background index

n<sub>sigma</sub>= confidence level in units of sigma

• The experimental sensitivity is thus:

$$T_{1/2}^{0\nu}(n_{\sigma}) = \frac{N_A \ln 2}{\sqrt{2}n_{\sigma}} \frac{a \cdot \epsilon}{A} \sqrt{\frac{M \cdot t}{B \cdot \Delta E}}$$

#### Experimental requirements

 $\bullet$  Experiments thus measure the half life of the decay,  $T_{1/2}$ 

$$T_{1/2}^{0\nu} \propto a \cdot \epsilon \cdot \sqrt{\frac{M \cdot t}{B \cdot \Delta E}}$$



Minimal requirements:

large detector masses enriched materials ultra-low background noise excellent energy resolution



Additional tools to distinguish signal from background:

angular distribution decay to excited states (gamma-rays) identification of daugther nucleus

#### Backgrounds for double beta experiments

✤ primordial radionuclides (<sup>238</sup>U, <sup>232</sup>Th, <sup>40</sup>K) in the detector materials, in the shielding and the concrete/rock (alpha, beta, gamma and neutrons)

✤ cosmic activation of detector materials (<sup>60</sup>Co, <sup>54</sup>Mn, <sup>65</sup>Zn,...)

✤ cosmic rays - muons - and secondary particles

✤ radon in air, radon emanation of materials,....

✤ anthropogens (<sup>85</sup>Kr, <sup>137</sup>Cs, <sup>207</sup>Bi,...)

2vββ-events: irreducible background an excellent energy resolution of the detector is crucial



F. Piquemal, Neutrino2012, Kyoto

#### Experiments: Main Approaches

Source ≠ Detector



Source as thin foil Electrons detected with: scintillator, TPC, drift chamber, semiconductor detectors Event topology Low energy resolution and detection efficiency

#### Source = Detector (calorimeters)



The sum of the energy of the two electrons is measured Signature: peak at the Q-value of the decay Scintillators, semiconductors, bolometers High resolution + detection efficiency No event topology (unless pixellized)



#### Source = Detector = Tracker

Source is - for example - the (high-pressure) gas of a TPC Charge and light detected with electron multipliers and/or photosensors Good energy and position resolution, high efficiency

electrons which drift to the TPC anode and generate EL light (or secondary scintillaten potential signal entering the region of interse field (Efp  $\approx 3$  kV/cm.bar) between the transparent EL grids. This light is recorded by an array of silicon photomultipliers (SiPM) located right behind the EL grids and used for tracking measurement. It is also recorded in the PMT plane behind the cathode for

# Existing experimental limits on T<sub>1/2</sub> and the effective Majorana neutrino mass

#### **Current best sensitivities are around a few 100 meV**

Table 1. A list of recent  $0\nu\beta\beta$  experiments and their 90% confidence level (except as noted) limits on  $T_{1/2}^{0\nu}$ . The  $\langle m_{\beta\beta} \rangle$  limits are those quoted by the authors using the  $M_{0\nu}$  of their choice.

Isotope	Technique	$T^{0\nu}_{1/2}$	$\langle m_{\beta\beta} \rangle ~(\mathrm{eV})$	Reference
$^{48}$ Ca	$CaF_2$ scint. crystals	$> 1.4 \times 10^{22} \text{ y}$	<7.2-44.7	14
$^{76}\mathrm{Ge}$	enrGe det.	$> 1.9 \times 10^{25} \text{ y}$	< 0.35	15
$^{76}\mathrm{Ge}$	$^{enr}$ Ge det.	$(1.19^{+2.99}_{-0.50}) \times 10^{25} \text{ y} (3\sigma)$	0.24 - 0.58	16
$^{76}\mathrm{Ge}$	$^{enr}$ Ge det.	$> 1.57 \times 10^{25} \text{ y}$	< (0.33 - 1.35)	17
$^{82}\mathrm{Se}$	Thin metal foils and tracking	$> 3.6 \times 10^{23} \text{ y}$	<(0.89-2.54)	18
$^{96}\mathrm{Zr}$	Thin metal foils and tracking	$> 9.2 \times 10^{21} \text{ y}$	<(7.2-19.5)	19
$^{100}\mathrm{Mo}$	Thin metal foils and tracking	$> 1.1 \times 10^{24} \text{ y}$	<(0.45-0.93)	18
$^{116}\mathrm{Cd}$	$^{116}$ CdWO <sub>4</sub> scint. crystals	$> 1.7 \times 10^{23} \text{ y}$	< 1.7	20
$^{128}\mathrm{Te}$	geochemical	$> 7.7 \times 10^{24} \text{ y}$	<(1.1-1.5)	21
$^{130}\mathrm{Te}$	$TeO_2$ bolometers	$> 2.8 \times 10^{24} \text{ y}$	< (0.3-0.7)	22
$^{136}\mathrm{Xe}$	Xe disolved in liq. scint.	$> 5.7 \times 10^{24} \text{ y}$	< (0.3-0.6)	23
$^{150}\mathrm{Ne}$	Thin metal foil within TPC	$> 1.8 \times 10^{22} \text{ y}$	N.A.	24

S. Elliott, http://arxiv.org/pdf/1203.1070v1.pdf



#### Current, near-future, future experiments



#### Existing and proposed experiments

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Experiment	Isotope	Mass	Technique	Present Status	Location
AMoRE <sup>89,90</sup>	$^{100}\mathrm{Mo}$	50 kg	CaMoO <sub>4</sub> scint. bolometer crystals	Development	Yangyang
CANDLES <sup>91</sup>	$^{48}$ Ca	$0.35 \ \mathrm{kg}$	$CaF_2$ scint. crystals	Prototype	Kamioka
$CARVEL^{92}$	$^{48}$ Ca	1  ton	$CaF_2$ scint. crystals	Development	Solotvina
COBRA <sup>93</sup>	$^{116}\mathrm{Cd}$	183 kg	<sup>enr</sup> Cd CZT semicond. det.	Prototype	Gran Sasso
$CUORE-0^{69}$	$^{130}\mathrm{Te}$	11 kg	$TeO_2$ bolometers	Construction - 2012	Gran Sasso
$CUORE^{69}$	$^{130}\mathrm{Te}$	203  kg	$TeO_2$ bolometers	Construction - 2013	Gran Sasso
$DCBA^{94}$	$^{150}\mathrm{Ne}$	20  kg	<sup>enr</sup> Nd foils and tracking	Development	Kamioka
$EXO-200^{57}$	$^{136}$ Xe	160 kg	Liq. <sup>enr</sup> Xe TPC/scint.	Operating - 2011	WIPP
$\mathrm{EXO}^{70}$	$^{136}$ Xe	1-10 t	Liq. <sup>enr</sup> Xe TPC/scint.	Proposal	SURF
$GERDA^{71}$	$^{76}\mathrm{Ge}$	$\approx 35 \text{ kg}$	enrGe semicond. det.	Operating - 2011	Gran Sasso
$\mathrm{GSO}^{95}$	$^{160}\mathrm{Gd}$	$2 \tan$	$Gd_2SiO_5:Ce$ crys. scint. in liq. scint.	Development	
KamLAND-Zen <sup>96</sup>	$^{136}$ Xe	400  kg	$^{enr}$ Xe disolved in liq. scint.	Operating - 2011	Kamioka
LUCIFER <sup>97,98</sup>	$^{82}$ Se	18 kg	ZnSe scint. bolometer crystals	Development	Gran Sasso
Majorana <sup>77,78,79</sup>	$^{76}\mathrm{Ge}$	26  kg	enrGe semicond. det.	Construction - 2013	SURF
MOON <sup>99</sup>	$^{100}\mathrm{Mo}$	1 t	<sup>enr</sup> Mofoils/scint.	Development	
SuperNEMO-Dem <sup>87</sup>	$^{82}$ Se	$7 \mathrm{kg}$	enrSe foils/tracking	Construction - 2014	Fréjus
$SuperNEMO^{87}$	$^{82}$ Se	100 kg	enrSe foils/tracking		
NEXT <sup>82,83</sup>	$^{136}$ Xe	100 kg	gas TPC		
SNO + 84,85	$^{150}$ Nd	55  kg	Nd loaded liq. scint.		

Table 2. A summary list of the  $0\nu\beta\beta$  proposals and experiments.

Steve Elliott: http://arxiv.org/pdf/1203.1070v1.pdf



## Recent results



- Liquid xenon TPC: 175 kg LXe, 80.6% enriched in 136Xe
- Charge and light readout (triplet wire channels and large area avalanche photodiodes)
- Drift field: 376 V/cm







- So far, 2 data taking phases
- First measurement of <sup>136</sup>Xe 2-neutrino half life; limit on the 0-neutrino mode

#### Data taking phases and Xenon Purity



## EXO-200: resolution and calibration

 $\sigma_{Tot}^2 = p_0^2 E + p_1^2 + p_2^2 E^2$ 

• Good energy resolution by linear combination of scintillation and charge signals









## KAMLAND-Zen

- Scintillator loaded with xenon
- 320 kg 90% enriched <sup>136</sup>Xe so far (more than 600 kg in the Kamioka mine)
- Advantages: huge and clean (U: 3.5e-18 g/g, Th: 5.2e-17 g/g) running detector
- Xe-LS can be purified, and is highly scalable
- No escape or invisible energy from gammas and beta: good background identification
- Disadvantage: relatively poor energy resolution
- no beta/gamma discrimination
- limited LS composition



#### KamLAND-Zen: installation



#### balloon and corrugated tube deployment







mini-balloon inflated with dummy LS and then replaced with Xe-loaded LS density tuning finished and tubes to be extracted

#### KamLAND-Zen



## KamLAND-Zen: energy calibration and lowbackground spectrum





Peak around the Q-value; however, peak position is different







#### Ge detectors in inquid argon (U/Th in LAr < 7x10<sup>-4</sup> µBq/kg)



• Physics run started on November 9, 2011



#### **GERDA** Calibration

Energy resolution: ~ 4.5 - 5 keV (FWHM) at 2.6 MeV



## GERDA low-background spectrum

- Background goal of ~ 10<sup>-2</sup> events/(kg yr keV) was reached
- Phase II (BEGe) detectors in production and testing
- LAr instrumentation (PMTs or SiPM & scintillating fibers) in development
- End of phase I and start of phase





#### GERDA low-background spectrum

• Analysis of 2-neutrino decay mode is in progress



## Summary

- Two-neutrino decay mode was measured for the first time in <sup>136</sup>Xe
- Xenon experiments provide competitive limits to germanium for the neutrinoless mode
- Several experiments are taking data, new results are expected soon
- Experiments under construction (or phase II of existing experiments) should achieve a sensitivity of 50 - 100 meV
- To go beyond, much lower backgrounds and larger masses are needed
- Tracking will be important to confirm a potential signal



#### Let us hope that...

• this prediction is true - it could be probed with future double beta experiments!



## End

#### Double beta decay

 If simple β<sup>-</sup> or β<sup>+</sup>-decay is forbidden on energetic grounds a nucleus can decay through a double beta mode:

$$^{106}_{48}Cd \rightarrow ^{106}_{46}Pd + 2e^+ + 2v_e$$

• The probability for a decay is very small, the mean lifetime of a nucleus is much larger than the age of the universe  $(\tau_{U} \sim 1.4 \times 10^{10} \text{ a})$ 

$$\tau_{2v} \approx 10^{20} a$$

- This is indeed a very rare process (as for instance proton decay, which was not yet observed)
- Nonetheless if one uses a large amount of nuclei, the process can be observed experimentally



Nuclear charge Z