

# The GERDA Neutrinoless Double Beta-Decay Experiment

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**Abstract.** Neutrinoless double beta ( $0\nu\beta\beta$ )-decay is the key process to gain understanding of the nature of neutrinos. The GERmanium Detector Array (GERDA) is designed to search for  $0\nu\beta\beta$ -decay of the isotope  $^{76}\text{Ge}$ . Germanium crystals enriched in  $^{76}\text{Ge}$ , acting as source and detector simultaneously, will be submerged directly into an ultra pure cooling medium that also serves as a radiation shield. This concept will allow for a reduction of the background by up to two orders of magnitudes with respect to earlier experiments.

**Keywords:** Neutrinoless Double Beta-Decay, Germanium Detectors

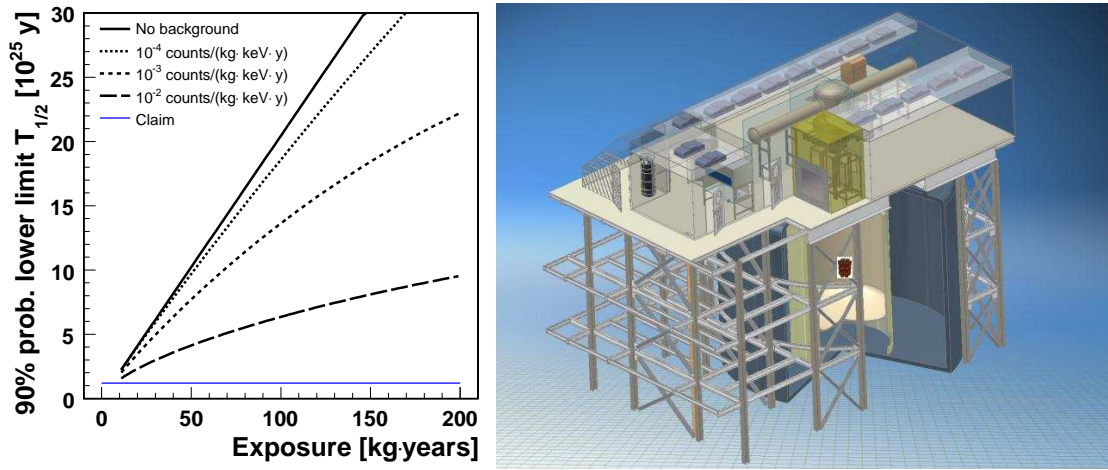
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## INTRODUCTION

Neutrino accompanied double beta-decay is an allowed second order weak process. Only if the neutrino is a massive Majorana-particle the decay can also occur without the emission of a neutrino [1]. The GERmanium Detector Array, GERDA [2], is designed to search for the  $0\nu\beta\beta$ -decay of  $^{76}\text{Ge}$ . The importance of such a search is emphasized by the observation of a non-zero neutrino mass of yet unknown scale from flavor oscillations [3]. The observation of neutrino-oscillation do allow for two different mass scenarios with a normal or an inverted hierarchy [4]. These can be disentangled by  $0\nu\beta\beta$ -decay for a Majorana neutrino if a sensitivity for the effective Majorana neutrino-mass of  $\approx 10$  meV can be achieved [5].

The most sensitive  $0\nu\beta\beta$  experiments so far are based on High-Purity-Germanium, HPGe, detector technology. This is due to the combination of a very good energy resolution of the detectors at the  $Q_{\beta\beta}$ -value of  $^{76}\text{Ge}$ , the very high purity of the detectors (very low intrinsic background) and the high signal detection efficiency of an experiment with the detector being equal the source.

Currently the Heidelberg-Moscow (HdMo) and IGEX experiments give limits on the lifetime of  $0\nu\beta\beta$ -decay of  $1.9 \cdot 10^{25}$  y and  $1.6 \cdot 10^{25}$  y, respectively (90% C.L.) [6, 7]. These lower limits can be translated into upper limits for the effective Majorana neutrino-mass in the range between 0.35 eV-1.2 eV and 0.3 eV-1.5 eV for the two experiments for different matrix elements (see [6, 7] and references therein). A part of the HdMo collaboration claims to have observed a peak at  $Q_{\beta\beta}$  with  $4.2\sigma$  confidence level which can be attributed to  $0\nu\beta\beta$ -decay with a half-life of  $T_{1/2}^{0\nu\beta\beta} = 1.19^{+2.99}_{-0.50} \cdot 10^{25}$  y (3  $\sigma$  range) [8].



**FIGURE 1.** Left: Expected 90% probability lower limit on  $T_{1/2}^{0\nu\beta\beta}$  as a function of exposure for different background indices (taken from [9]). Right: Schematic view of the GERDA setup. The detector array is sitting in the center of the cryostat. The cryo tank is surrounded by a water buffer serving as additional shield and as a muon-Cherenkov veto system. The detector array is loaded through a lock system located in a clean-room on the top of the tank.

## SENSITIVITY AND THE PRINCIPLES OF GERDA

The sensitivity obtainable in double beta experiments with a given exposure and background index is displayed in the left panel of Fig. 1. The expected 90 % probability lower limit as a function of exposure for given background indices are shown. The values were calculated using Monte-Carlo ensemble test on the basis of Bayesian statistics [9].

For a given exposure the background limits the sensitivity. Therefore the goal is to minimize the background. This can be achieved by using an ultra-pure cryogenic liquid as the cooling medium and as a shield against gamma radiation simultaneously [10]. The cryostat will be made out of carefully selected stainless steel. Additionally it will contain a low-background copper inlet as a radiation shield against the steel of the cryo-tank. The cryogenic volume is surrounded by a buffer of ultra-pure water acting as an additional gamma and neutron shield. The water buffer is additionally used as a muon-Cherenkov veto. The setup is schematically depicted in Fig. 1.

GERDA will be installed in Hall A of the Gran Sasso underground laboratory (LNGS), Italy. The experiment is foreseen to proceed in two phases. In the first phase, enriched detectors which were previously operated by the HdMo and IGEX collaborations will be re-deployed. If the HdMo claim is correct,  $6.0 \pm 1.4$  events above a background of 0.5 events after one year of measurement (15 kg · y) are expected. This will be enough to either confirm or refute the claim from the HdMo experiment (see Fig. 1). In the second phase custom made segmented detectors will be installed which have a true coaxial geometry. An exposure of 100 kg·y with a background level of  $10^{-3}$  counts/(kg·keV·y) is foreseen. A lower limit on the half-life of  $0\nu\beta\beta$ -decay of higher than  $1.5 \cdot 10^{26}$  y corresponding to an upper limit of  $\approx 140$  meV (using the corrected matrix elements from [11]) can be set with 90 % probability if no signal is observed.

## MAIN BACKGROUND SOURCES

Monte Carlo studies using the GEANT 4 application MaGe [12] were performed to estimate the background expected for Phase II of the GERDA experiment. The detector geometries assumed reflect the design of the first segmented prototype detector that has been developed and successfully operated [13].

A background of  $\approx 5 \cdot 10^{-4}$  counts/(kg·keV·y) is expected from internal contaminations of the detector (after two years of storage underground). The origin of this background is mainly cosmogenic  $^{68}\text{Ge}$  and  $^{60}\text{Co}$ . This background can be reduced by minimizing the exposure of the germanium to cosmic rays. Since  $^{68}\text{Ge}$  decays with a half-life of  $T_{1/2}=271$  d, this component decreases by a factor of 2.54 per year.

A contribution of  $\approx 5 \cdot 10^{-4}$  counts/(kg·keV·y) is expected from the detector infrastructure, i.e. detector support, cabling and electronics. Presently for the cabling and holder system 2.5 g of copper on Kapton cable, 31g of copper and 7 g of PTFE are used per detector in the direct vicinity of the crystals. Material selection and material minimization that could lead to a considerable decrease of this background component are under way.

Roughly  $3 \cdot 10^{-5}$  counts/ (kg·keV·y) are expected from the external infrastructure (cryo-tank, water shield, etc.) from  $^{208}\text{Tl}$  and  $^{214}\text{Bi}$  photons being negligible compared to the dominant background sources.

Neutrons and muons are expected to yield roughly  $2 \cdot 10^{-4}$  counts/(kg·keV·y) with an important contribution being due to the delayed decay of  $^{77}\text{Ge}$  produced by the neutron capture on  $^{76}\text{Ge}$  [14].

During the lifetime of the experiment the internal background will further decrease. In connection with the ongoing work on material selection this will ensure that the goal of a total background index of less than  $10^{-3}$  counts/(kg·keV·y) can be reached.

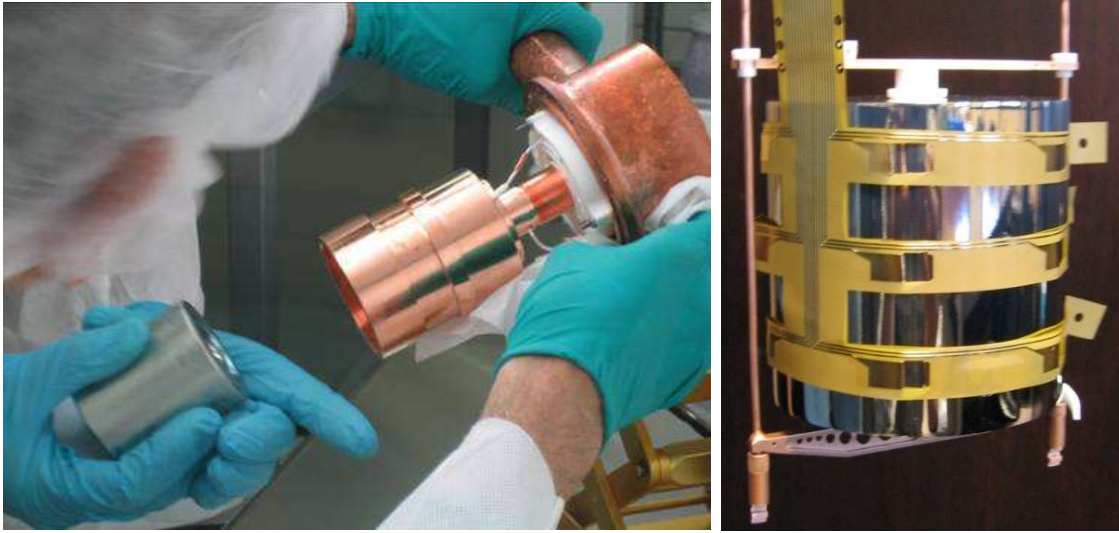
## PHASE I AND PHASE II STATUS

All IGEX and HdMo detectors are presently stored underground. They have been removed from their cryostats without any technical problems. The left panel of Fig. 2 shows the dismantling of the first enriched HdMo detector. The detectors are presently being refurbished and will then be installed into the phase I suspension. A prototype detector has been operated in liquid argon in the holder at LNGS since beginning of 2006. It has gone through more than twenty cooling and warming up cycles.

The enriched materials for phase II of the experiment, 35.5 kg of germanium enriched to 87%-88% in  $^{76}\text{Ge}$  in form of  $\text{GeO}_2$ , have been procured and were transported from Siberia to Germany in a steel cylinder designed to reduce cosmogenic activation. The material was unloaded and weighed and is now stored in a 500 mwe underground site.

The phase II detectors will be segmented true coaxial n-type, 18-fold in the presently working prototype design. The segmentation will help to identify multiple Compton-scattering events in the region of interest. As shown in [15] the Compton background can be identified with high efficiency depending on its source and location.

A prototype detector (see right panel of Fig. 2) was produced. The 18 segments are

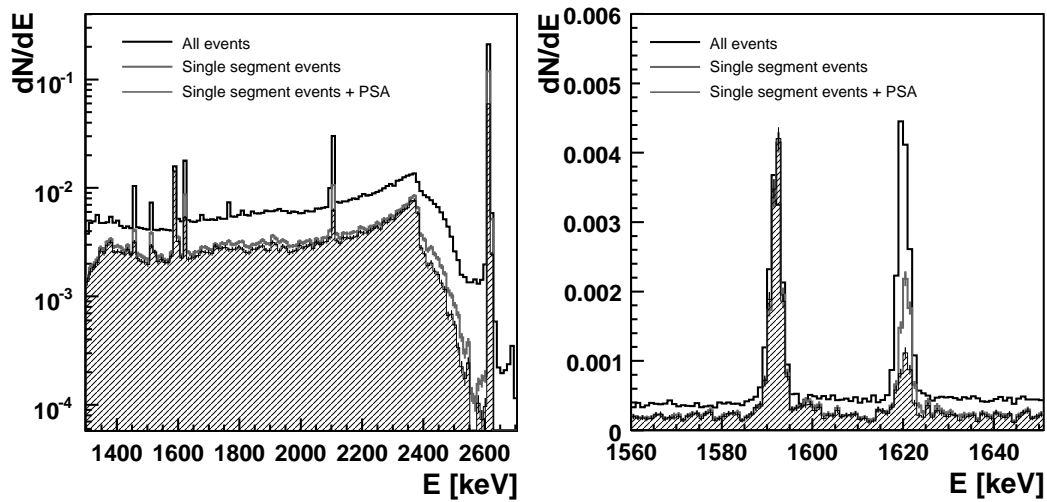


**FIGURE 2.** Left: Dismounting of first enriched detector of the HdMo experiment from its copper vacuum-cryostat. Right: Germanium crystal with the new contacting scheme in its low mass holder. The 18 segments are contacted by pressing the contact pads of Kapton printed circuit boards onto the segment contact area.

read out using a novel contacting scheme described in [13]. The prototype detector was extensively tested in a conventional test cryostat. The energy resolution of all signals and the core were around 3 keV at 1.3 MeV. Details of the setup can be found in [13]. Fig. 3 shows a spectrum taken with a  $^{228}\text{Th}$  calibration source. Three spectra are shown: The full spectrum as measured by the core electrode, the spectrum of events in which an energy deposition  $\geq 20$  keV was detected in a single segment only (single segment events) and the spectrum with single segment events after application of pulse shape analysis (PSA). The left panel shows the energy spectrum up to 2.7 MeV, the right panel shows the energy region around the double escape peak of the 2.615 MeV gamma and the 1620 keV Compton peak. The Compton background discrimination works well and was a factor of three around 2 MeV for this configuration. The data could be reproduced very well by Monte Carlo simulations [15]. The double escape peak is only marginally reduced by application of the single segment requirement and remains basically unchanged after further application of PSA. The Compton peak is suppressed by requiring single segment events. PSA further reduces the peak by removing events close to the segment boundaries [16].

## CONCLUSIONS

The Gerda neutrinoless double beta-decay experiment will be installed in the LNGS underground laboratory. In the first phase an exposure of  $\geq 15$  kg·y with a background index of  $10^{-2}$  counts/(kg·keV·y) are planned. This will allow to check the claim of a positive evidence based on the HdMo experiment data. The goal of the second phase of the experiment is to collect 100 kg·y of data with a background index of not more than



**FIGURE 3.**  $^{228}\text{Th}$  calibration measurement of the 18-fold segmented prototype detector. The left panel shows the full spectrum up to 2.7 MeV, the right panel shows the energy region around the double escape peak of the 2.615 MeV gamma and the 1620 keV Compton peak (taken from [16]). See text for more details.

$10^{-3}$  counts/(kg·keV·y). This will allow to set a lower limit for the  $0\nu\beta\beta$  half life of  $1.5 \cdot 10^{26}$  y (90% C.L.) corresponding to an upper limit of the effective Majorana neutrino-mass of 140 meV. According to the Monte Carlo calculations and the material screening done so far this goal can be achieved. The successfully dismantled IGEX and HdMo detectors for phase I are being refurbished. A first prototype 18-fold segmented n-type HPGe detector has been produced and checked. Its performance fulfills all expectations.

## REFERENCES

1. S. Elliott and P. Vogel, *Ann. Rev. Nucl. Part. Sci* **52**(2002)115
2. I. Abt et al., hep-exp/0404039
3. S. Fakuda et al., *Phys. Rev. Lett.* **82**(1999)2644 and C. Ahmad et al., *Phys. Rev. Lett* **92**(2004)181301
4. S. Bilenky et al., *Phys. Rev. C* **64**(2001)053010
5. F. Feruglio et al., *Nucl. Phys. B* **637**(2002)345
6. L. Baudis et al., *Phys. Rev. Lett.* **83**(1999)41
7. C. Aalseth et al., *Phys. Rev. D* **65**(2002)092007
8. H.V. Klapdor-Kleingrothaus et al., *Phys. Lett. B* **586**(2004)198
9. A. Caldwell and K. Kröniger, *Phys. Rev. D* **74**(2006)092003
10. G. Heusser, *Ann. Rev. Nucl. Part. Sci.* **45** (1995)543
11. V. Rodin et al., *Nucl. Phys. A* **766**(2006)107 and V. Rodin, et al., arXiv:0704.4304
12. M. Bauer et al., *J. Phys. Conf. Series* **39**(2006)362
13. I. Abt et al., *Nucl. Instr. Meth. A* **577**(2007)574
14. L. Pandola et al., *Nucl. Instr. Meth. A* **570**(2007)149
15. I. Abt et al., *Nucl. Instr. Meth. A* (2007)479
16. I. Abt et al., arXiv:0704.3016, accepted for publication by European Physical Journal C