

# phase space and NME

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A. Smolnikov & P. Grabmayr, PRC81 (2010) 28502

PHYSICAL REVIEW C 81, 028502 (2010)

**Conversion of experimental half-life to effective electron neutrino mass in  $0\nu\beta\beta$  decay**

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(Received 18 December 2009; published 23 February 2010)*

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arXiv:1002.1608v1 [nucl-th] 8 Feb 2010

# some formulas



$$\frac{1}{T_{1/2}} = G^{0\nu} |\mathcal{M}^{0\nu}|^2 \left| \frac{\langle m_{\beta\beta} \rangle}{m_e} \right|^2 = F^{0\nu} |\mathcal{M}^{0\nu}|^2 m_{\beta\beta}^2$$

Phase space factor

$$G^{0\nu} = G_{01} = \frac{a_{0\nu}}{(m_e R_A)^2 \ln 2} \int d\Omega_{0\nu} b(\varepsilon_1, \varepsilon_2)$$

$$a_{0\nu} = \frac{(G g_A)^4 m_e^9}{64\pi^5},$$

$$b(\varepsilon_1, \varepsilon_2) = F_0(Z, \varepsilon_1) F_0(Z, \varepsilon_2)$$

$$d\Omega_{0\nu} = \frac{p_1 \varepsilon_1 p_2 \varepsilon_2}{m_e^5} \delta(\varepsilon_1 + \varepsilon_2 + M_f - M_i) d\varepsilon_1 d\varepsilon_2 d \cos \theta$$

$$G^{0\nu} \sim g_A^4 \frac{1}{R_A^2}$$

# radius

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$$R = 3,108 \times 10^{-3} A^{1/3} / m_e$$

Doi

$$R = 1.1 A^{1/3}$$

Faessler

$$R = 1.2 A^{1/3}$$

Suhonen, Iachello,  
Poves



Author	Ref.	$r_0$ [fm]	$G^{0\nu} \times 10^{15}$ [ $y^{-1}$ ]	$\mathcal{M}^{0\nu}$	$s \times 10^{25}$ [ $eV^2 \times y$ ] $^{-1}$	Comments
Claim	[3]	1.2	6.31	4.22	4.30	
Pantis	[11]	1.1	7.93	1.34	0.55	np pairing
Simkovic	[12]	1.1	7.93	2.80	2.38	RQRPA
Simkovic	[12]	1.1	7.93	3.60	3.94	RQRPA
Rodin	[13]	1.1	7.93	3.92	4.67	RQRPA
Simkovic	[14]	1.1	7.93	3.33	3.37	Jastrow <sup>a</sup>
Simkovic	[14]	1.1	7.93	4.68	6.65	Jastrow <sup>b</sup>
Simkovic	[14]	1.1	7.93	3.92	4.67	UCOM <sup>a</sup>
Simkovic	[14]	1.1	7.93	5.73	9.97	UCOM <sup>b</sup>
Caurier	[15]	1.2	6.31	2.22	1.19	SM
Barea	[16]	1.2	6.31	5.47	7.23	IBM2-I
Barea	[16]	1.2	6.31	4.64	5.20	IBM2-II
Suhonen	[18]	1.2	6.31	2.78	1.87	Jastrow <sup>c</sup>
Suhonen	[18]	1.2	6.31	2.28	1.26	Jastrow <sup>d</sup>
Suhonen	[18]	1.2	6.31	4.11	4.08	UCOM <sup>c</sup>
Suhonen	[18]	1.2	6.31	3.23	2.52	UCOM <sup>d</sup>
Menendez	[20]	1.2	6.31	3.00	2.17	SM gcn <sup>e</sup>
Menendez	[20]	1.2	6.31	3.52	2.99	SM rg <sup>e</sup>
Simkovic	[21]	1.1	7.93	5.44	8.99	RQRPA <sup>e</sup>
Simkovic	[21]	1.1	7.93	4.07	5.03	RQRPA <sup>ae</sup>
Simkovic	[21]	1.1	7.93	6.64	13.39	RQRPA <sup>be</sup>

# Auxiliary variable $s$

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$$\frac{1}{T_{1/2}} = G^{0\nu} |\mathcal{M}^{0\nu}|^2 \left| \frac{\langle m_{\beta\beta} \rangle}{m_e} \right|^2 = F^{0\nu} |\mathcal{M}^{0\nu}|^2 m_{\beta\beta}^2$$

$$s = G^{0\nu} |\mathcal{M}^{0\nu}|^2 m_e^{-2} \quad [\text{eV}^2 \text{y}]^{-1}$$

$$\langle m_{\beta\beta} \rangle = (s T_{1/2})^{-1/2}$$

# Masses in [meV]



$$s = G^{0\nu} |\mathcal{M}'^{0\nu}|^2 m_e^{-2} \quad [\text{eV}^2 \text{y}]^{-1}$$

Author	Ref.	$s \times 10^{25}$ [eV <sup>2</sup> × y] <sup>-1</sup>	$T_{1/2} \times 10^{-25}$ [y]				
			1.2	2.2	3	15	20
			Claim	Phase I		Phase II	
Menendez	[20]	2.99	528	390	334	149	129
Suhonen	[18]	4.08	452	334	286	128	111
Rodin	[13]	4.67	422	312	267	119	103
Barea	[16]	7.23	340	251	215	96	83
Simkovic	[21]	8.99	304	225	193	86	75

# recommendation



Experimental sensitivity	Ref.	$T_{1/2}$ [ $10^{-25}$ y]	$\langle m_{\beta\beta} \rangle$ [eV]
Claim	[3]	1.2	0.30–0.53
GERDA Phase I	[1]	2.2	0.23–0.39
GERDA Phase II	[1]	15.0	0.09–0.15

# open problem

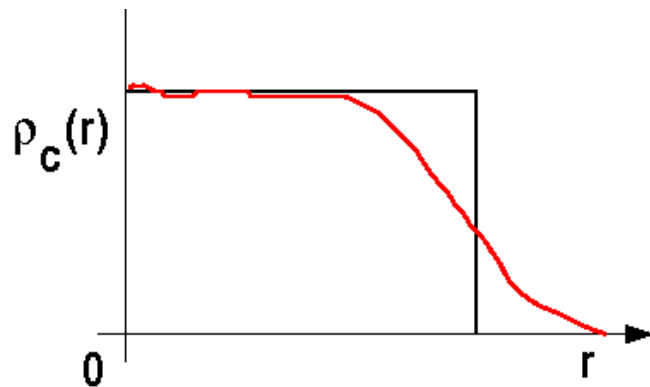


Fermi function  $F_0$  & integrand  $b(\varepsilon_1, \varepsilon_2)$

$$b(\varepsilon_1, \varepsilon_2) = F_0(Z, \varepsilon_1)F_0(Z, \varepsilon_2)$$

assumes: 1. 2e emission  
2. distortion by same  
Z-distribution

$F_0$  calculated from hard sphere distribution



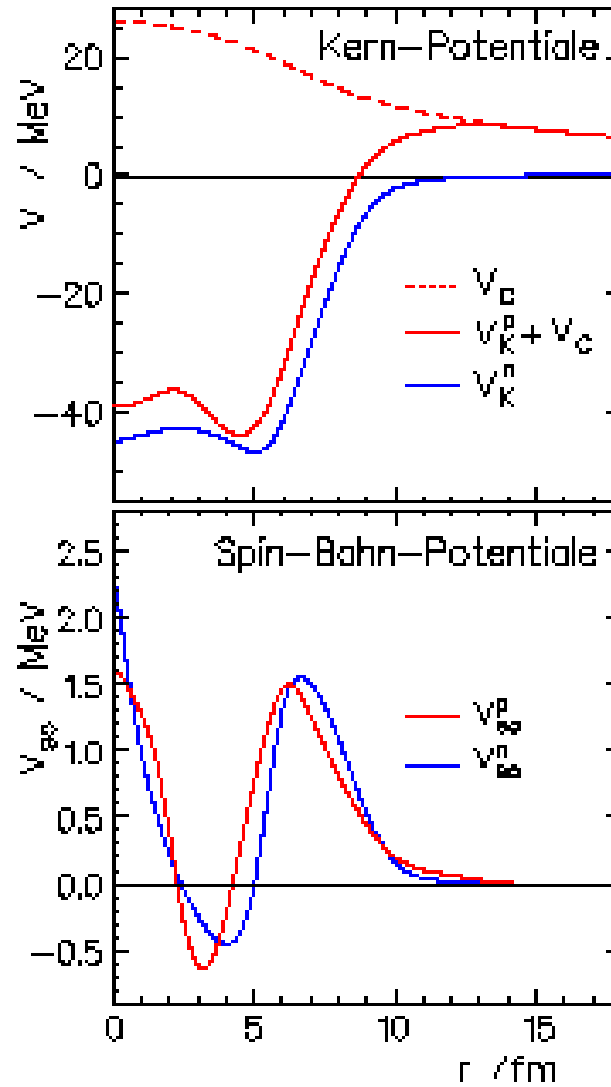
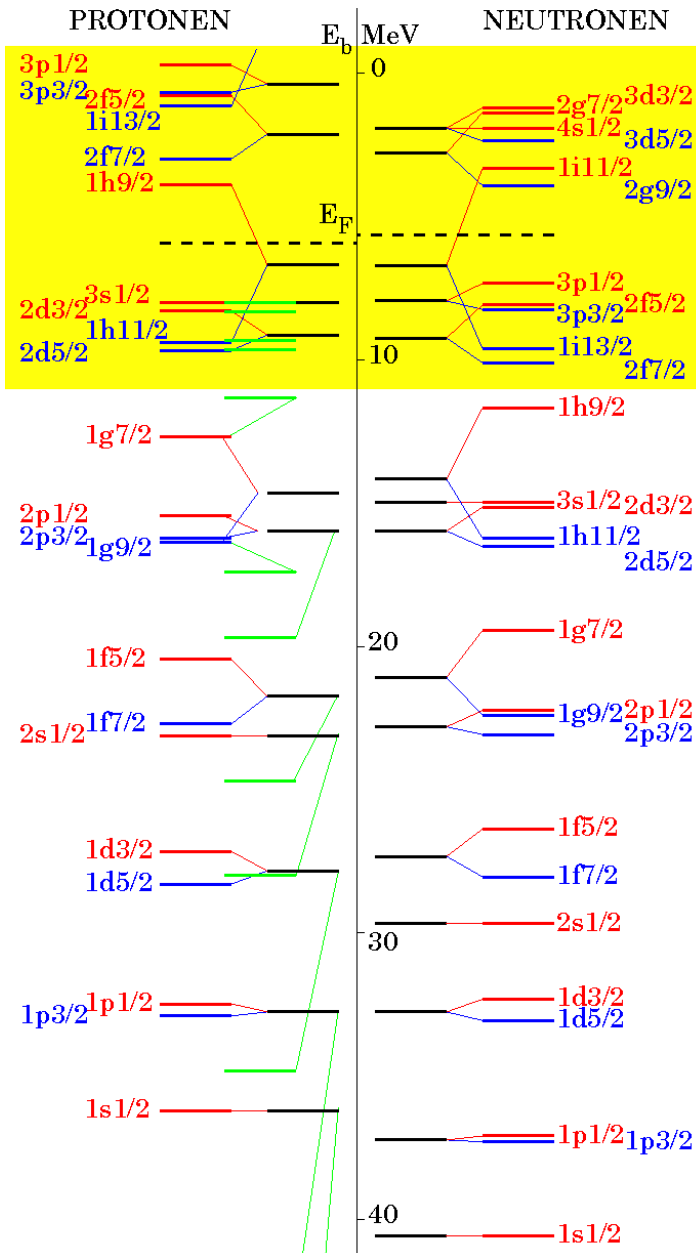
solve Dirac eq. in potential  
due to  $\rho_c(r)$

argued by Doi etc: not important (few %) (comp  $^{12}\text{B}$ )



# realistic potentials

# 208Pb



G.Mairle & P. Grabmayr  
EPJA 9 (2000) 313

# difference in occupancies



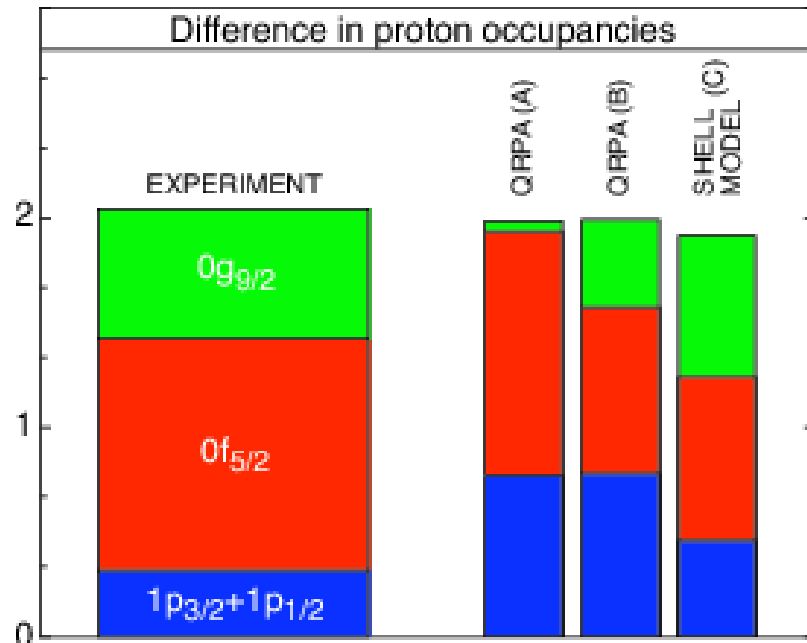
JP Schiffer et al., PRL 100 (2008) 112501

$$\Delta n = -2 \quad \Delta p = +2$$

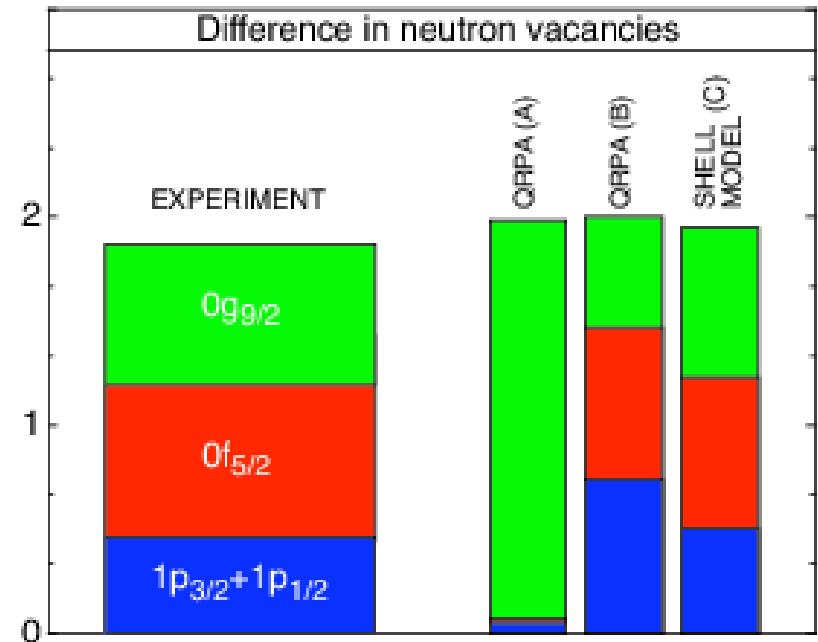
PHYSICAL REVIEW C 79, 021301(R) (2009)

**Nuclear structure relevant to neutrinoless double  $\beta$  decay: The valence protons in  $^{76}\text{Ge}$  and  $^{76}\text{Se}$**

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protons



neutrons

# Summary

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There are more uncertainties in the theoretical results than just the nuclear matrix elements.

Importance of support experiments.

## Deformation of Ge & Se

