

X-ray absorption and occultation in LS 5039

Anna Szostek & Guillaume Dubus

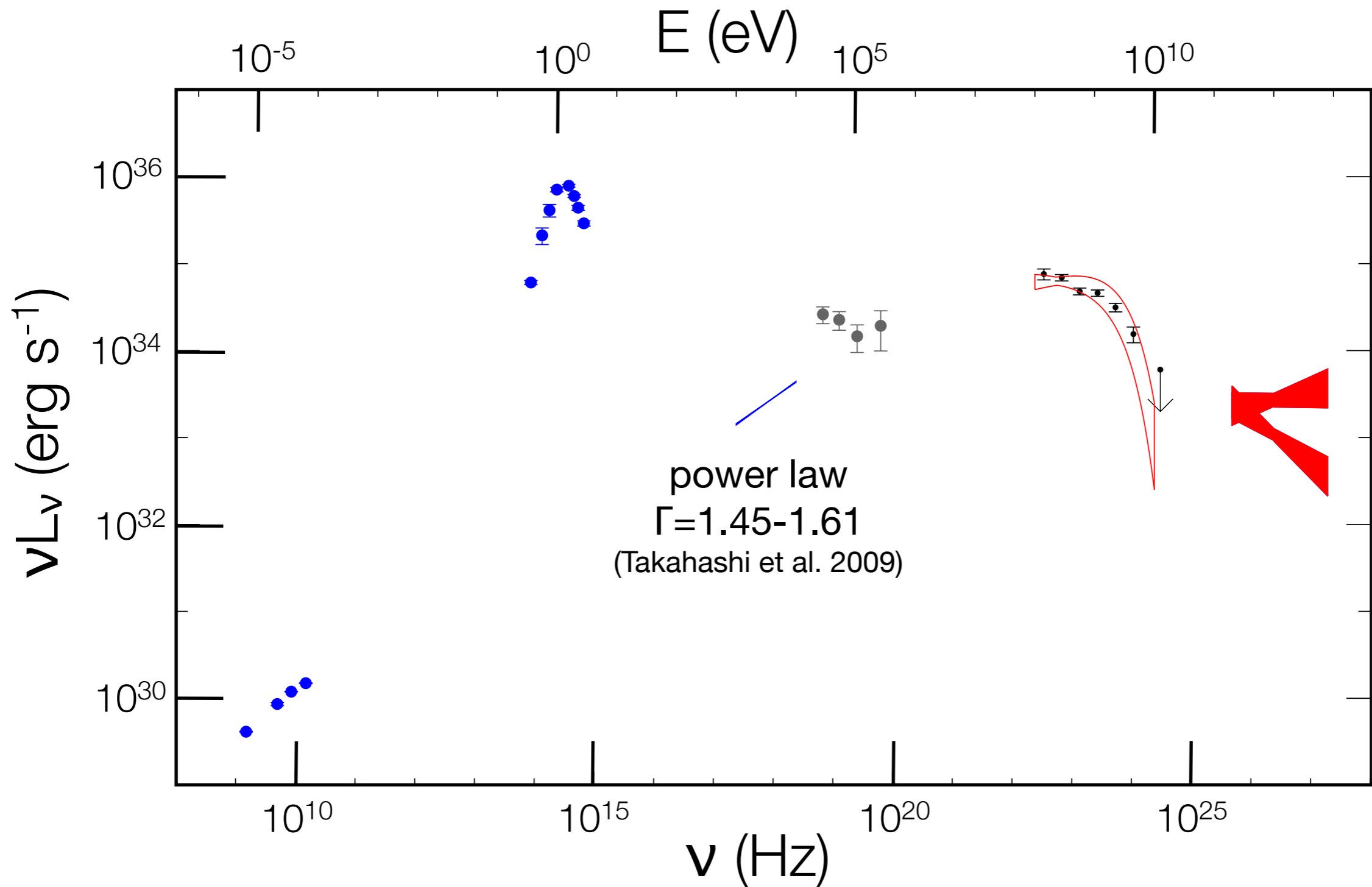
Laboratoire d'Astrophysique de l'Observatoire de Grenoble
Astronomical Observatory of the Jagiellonian University



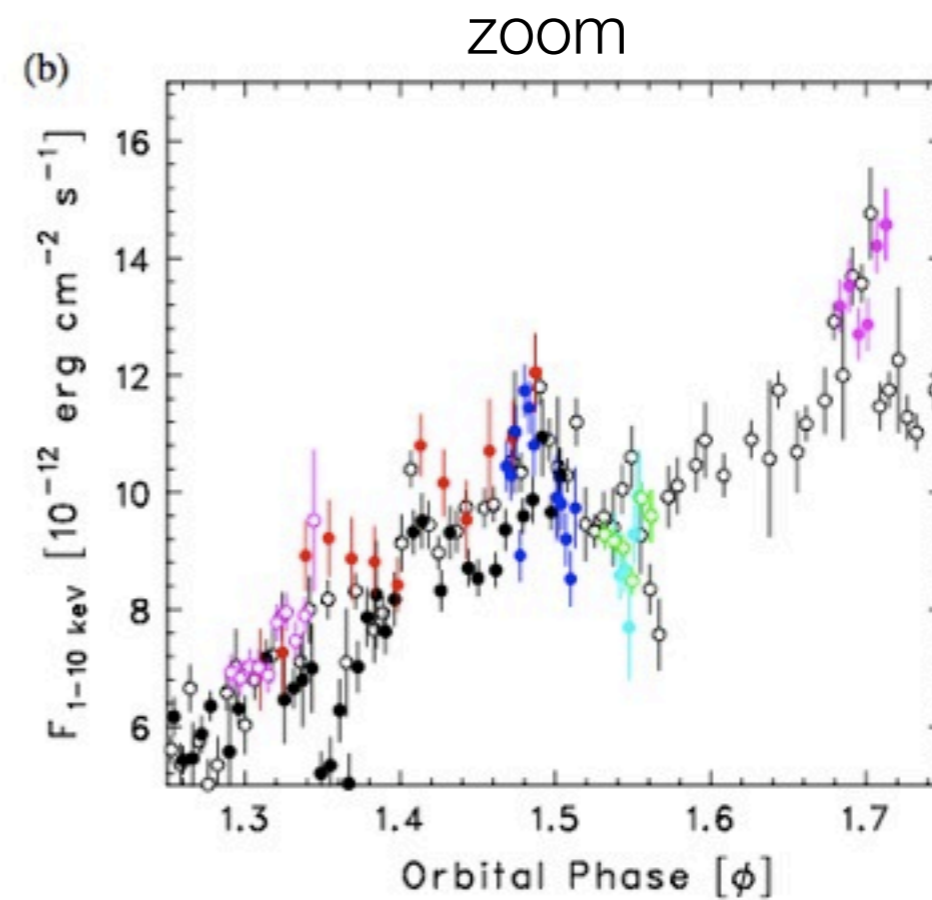
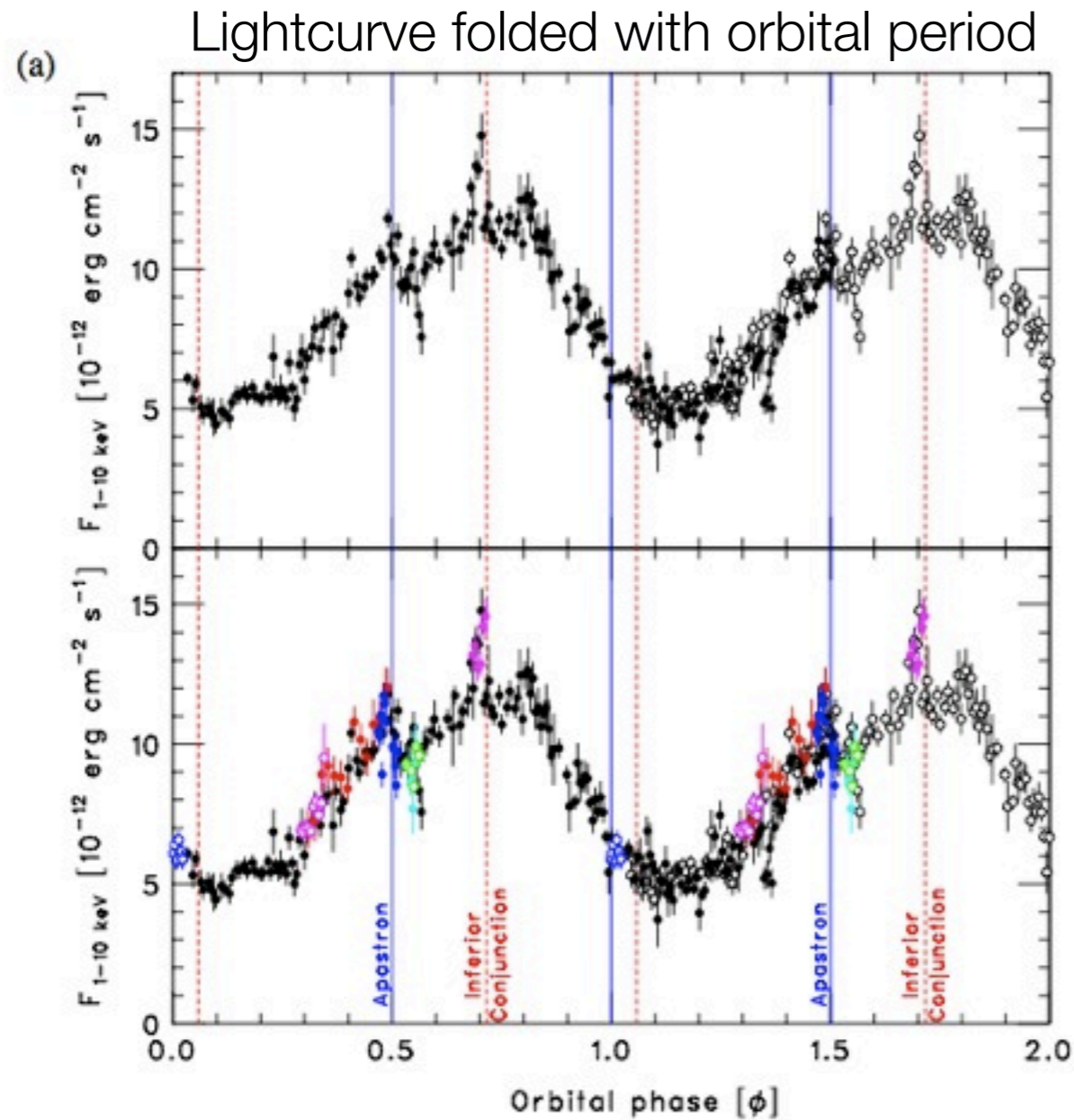
European Research Council



Spectral energy distribution of LS 5039

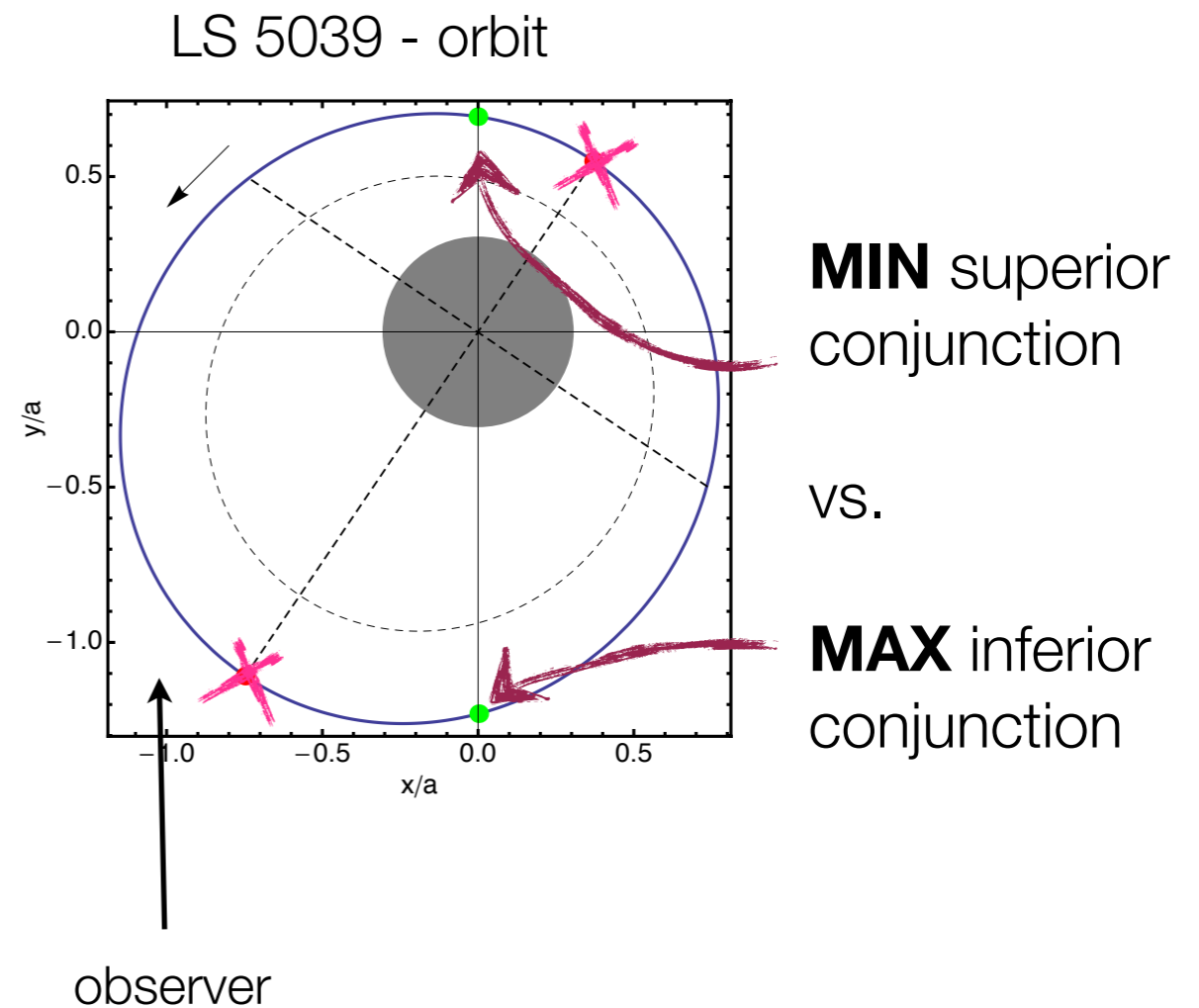
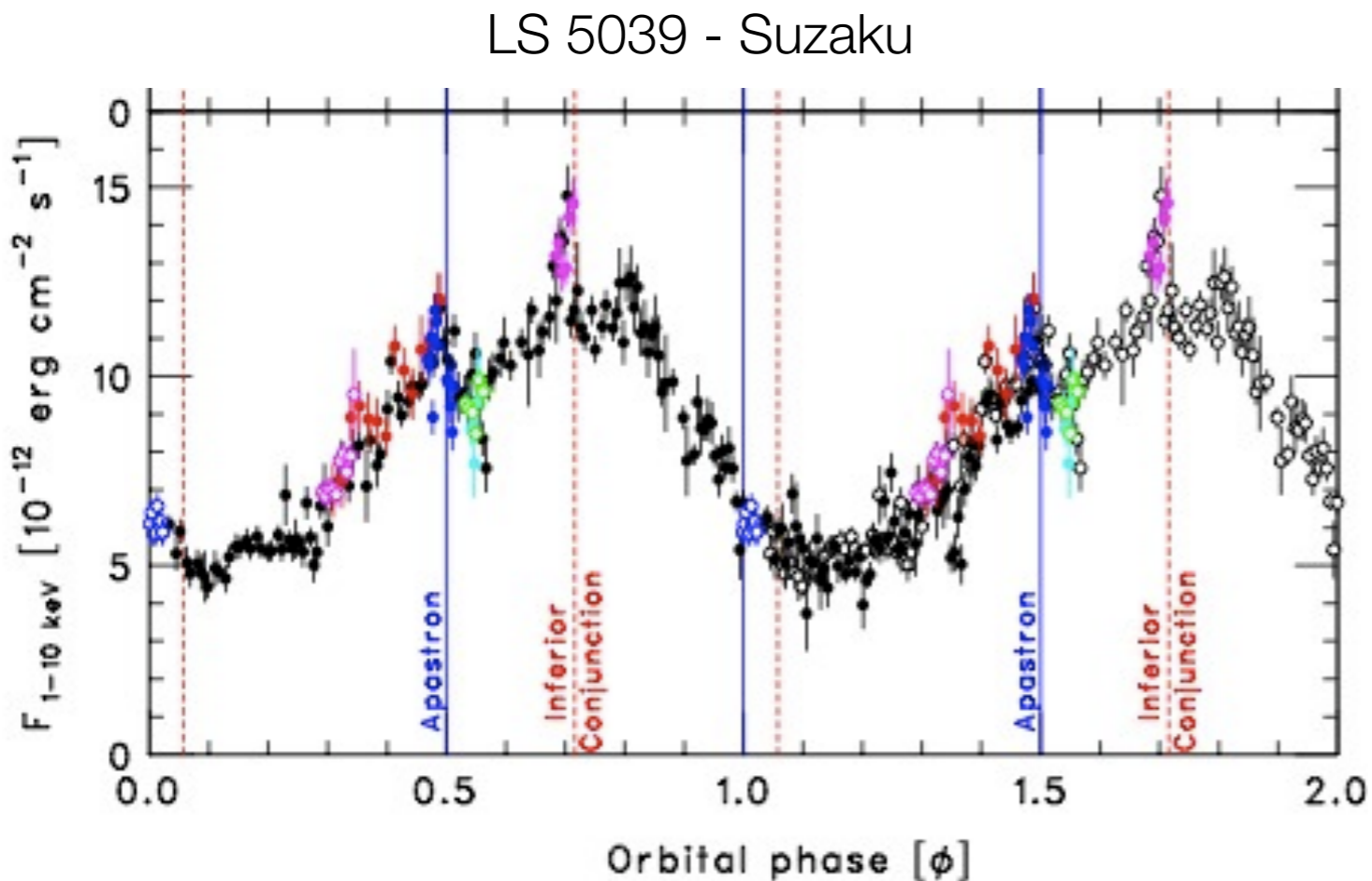


X-ray lightcurve of LS 5039 is very stable on long timescales.



1-10 keV, 1999-2007, Kishishita et al. 2009

X-ray lightcurve extrema occur during conjunctions, **NOT** during periastron or apastron.



Takahashi et al. 2009

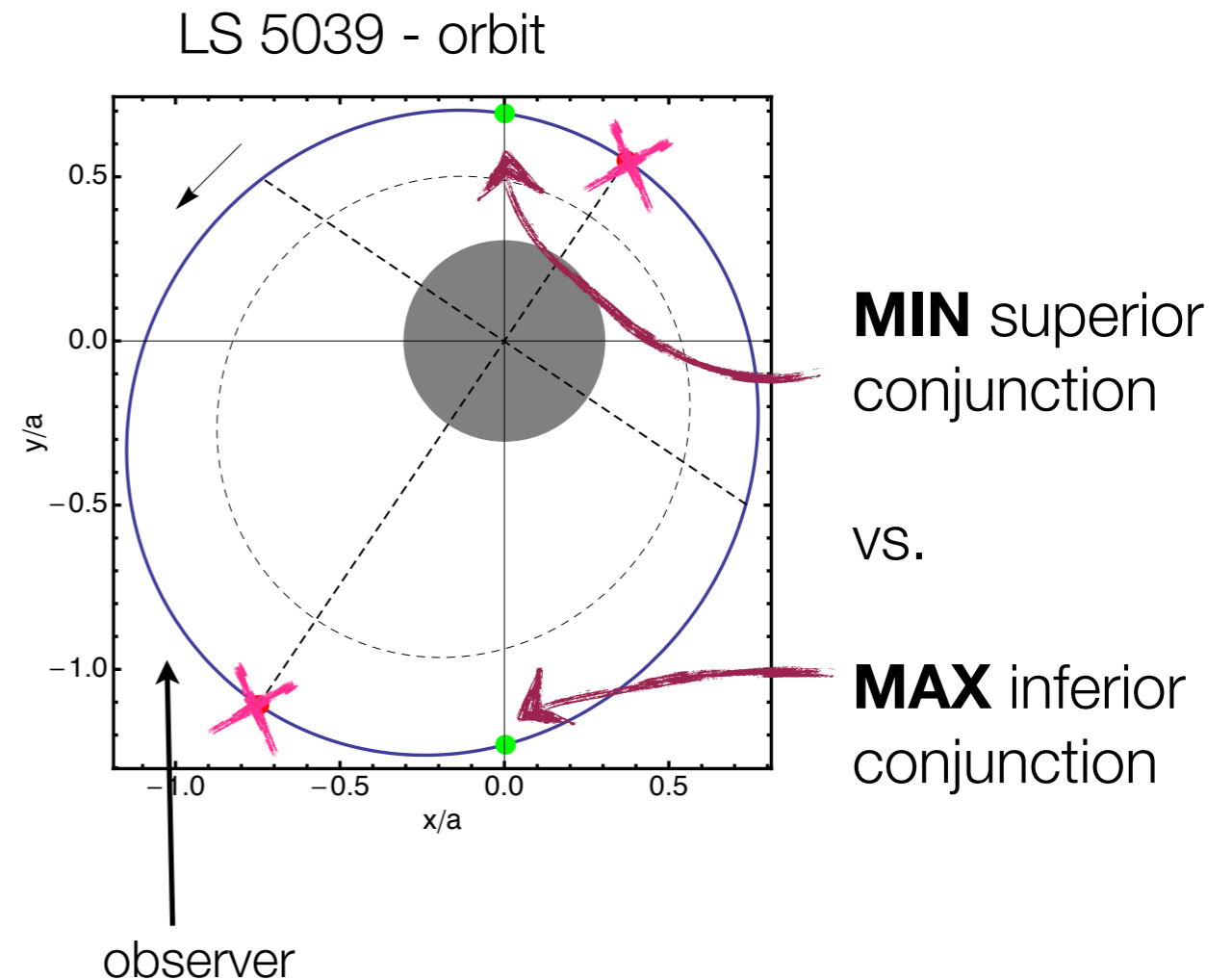
The modulation is a geometrical effect not related to physical changes when compact object travels on its orbit

Two example effects could be responsible for orbital modulation.

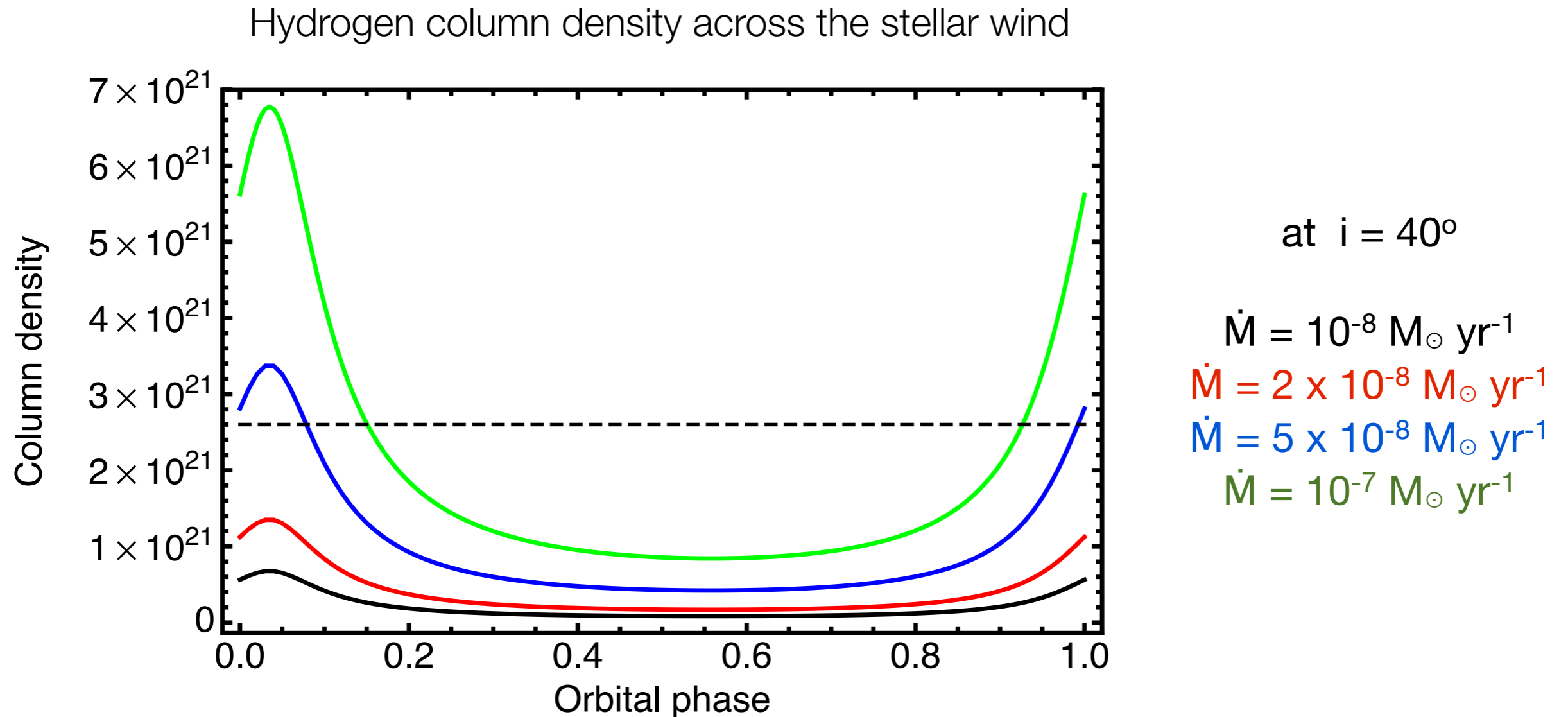
1. Photoelectric absorption of X-rays in the stellar wind

2. Occultation of X-ray emitting region by a massive star

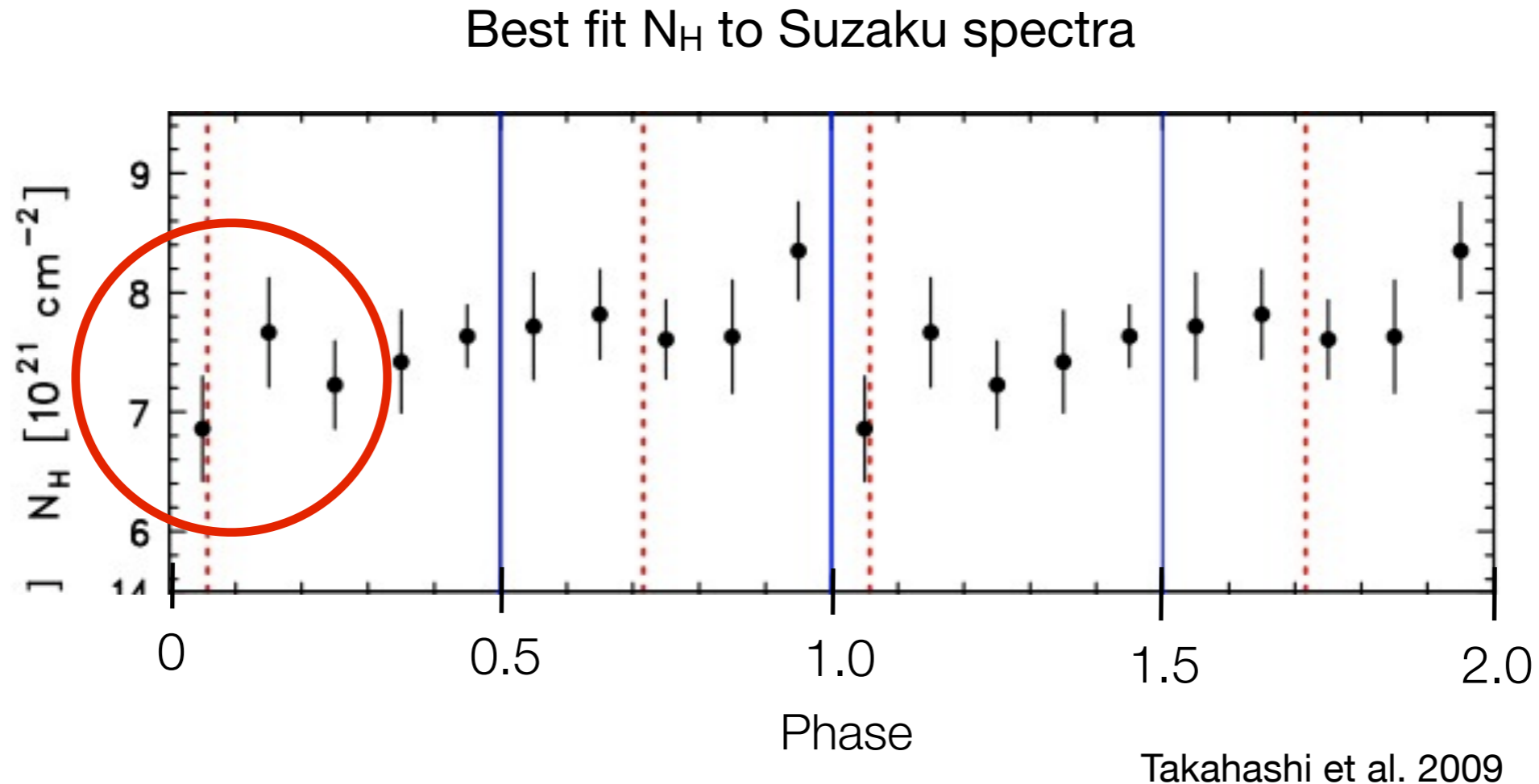
Both produce extrema at correct orbital phases



For X-ray point source at compact object, absorption effects can be seen only around superior conjunction.



In LS 5039 absorbing column density N_H does not change with orbital phase.

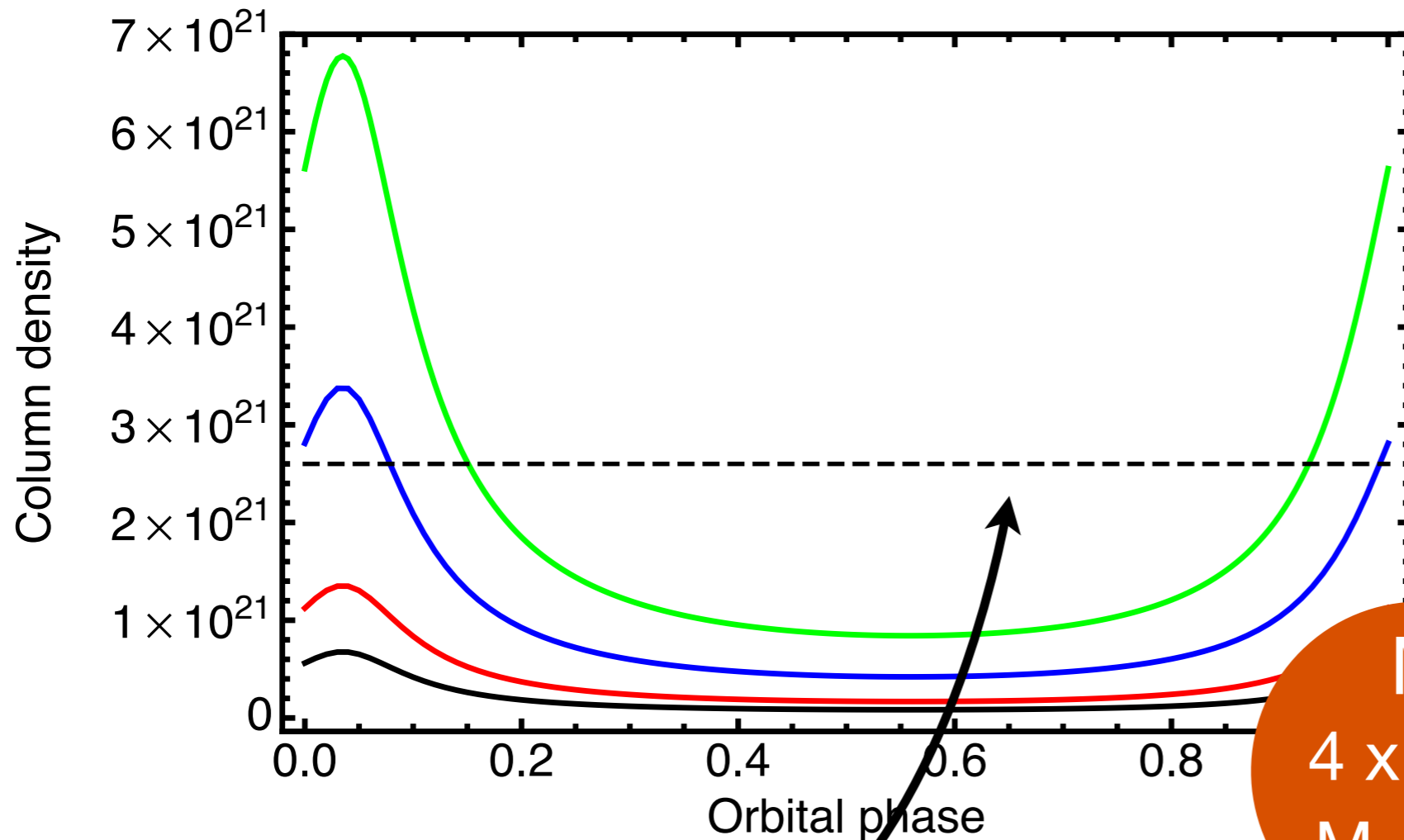


Taking into account absorption in ISM,
upper limit on intrinsic column density

$$N_H^{\text{max}} = 2.6 \times 10^{21} \text{ cm}^{-2}$$

(Bosch-Ramon et al. 2007, Takahashi et al. 2009)

For X-ray point source at compact object, mass-loss rate upper limit is low, $\sim 4 \times 10^{-8} M_{\odot} \text{ yr}^{-1}$.



at $i = 40^\circ$

$\dot{M} = 10^{-8} M_{\odot} \text{ yr}^{-1}$

$\dot{M} = 2 \times 10^{-8} M_{\odot} \text{ yr}^{-1}$

$\dot{M} = 5 \times 10^{-8} M_{\odot} \text{ yr}^{-1}$

$\dot{M} = 10^{-7} M_{\odot} \text{ yr}^{-1}$

\dot{M}
 4×10^{-8}
 $M_{\odot} \text{ yr}^{-1}$

\dot{M} at lower end for Galactic O stars
 $(\sim 10^{-8} - 10^{-6} M_{\odot} \text{ yr}^{-1})$
 (Mokiem et al. 2007)

Upper limit on intrinsic column density
 $N_{\text{H}}^{\text{max}} = 2.6 \times 10^{21} \text{ cm}^{-2}$
 (Bosch-Ramon et al. 2007, Takahashi et al. 2009)

\dot{M} for LS 5039 in literature
 $3 \times 10^{-8} - 7.5 \times 10^{-7} M_{\odot} \text{ yr}^{-1}$
 (McSwain & Gies 2002; Casares et al. 2005)

There is also no evidence for occultation in LS 5039

From the absence of X-ray eclipses
Upper limit on binary inclination is

65°

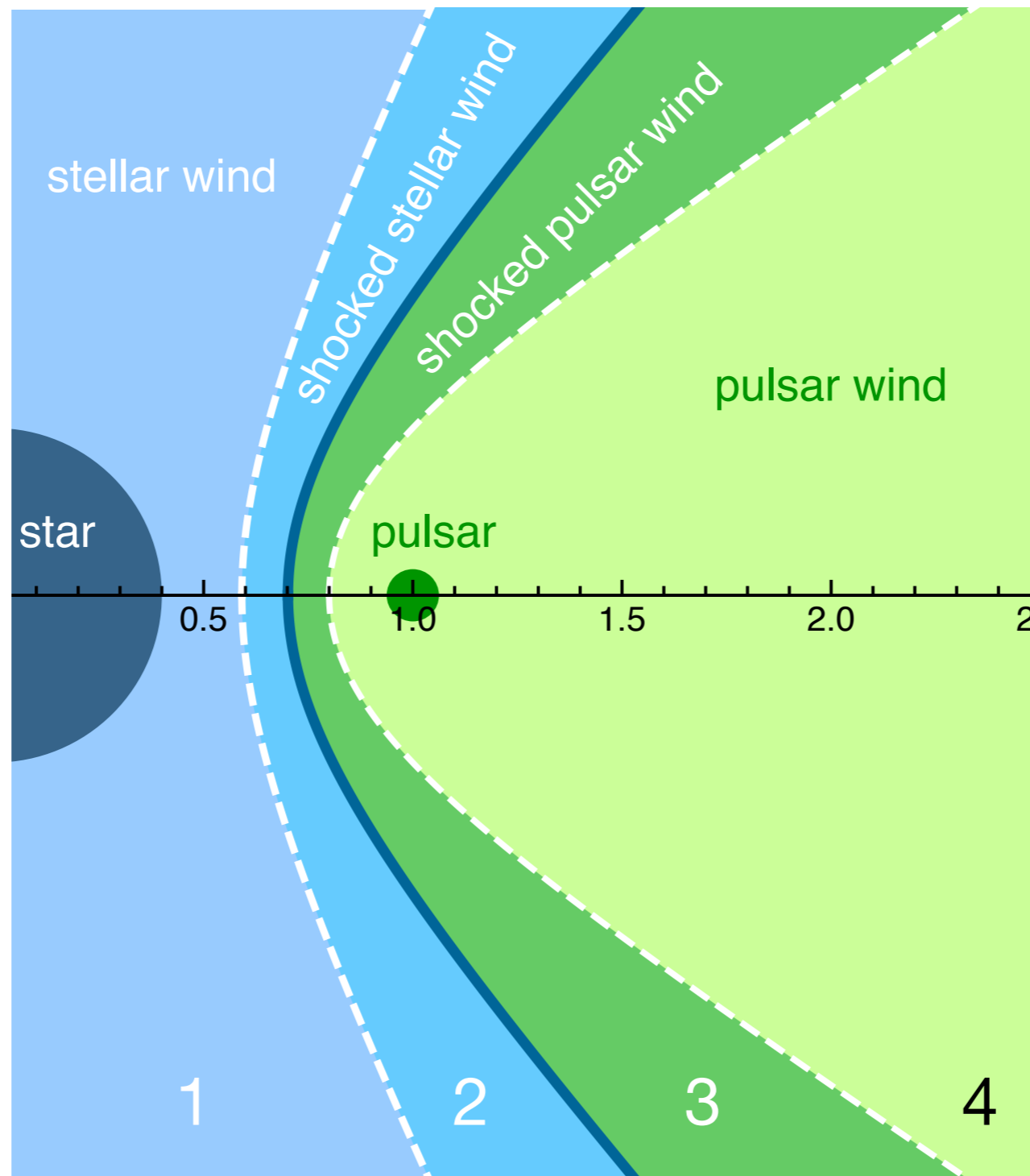
(Casares et al. 2005)

If the X-ray source is extended, the strict limits on \dot{M} and binary inclination are not valid anymore.

Motivation

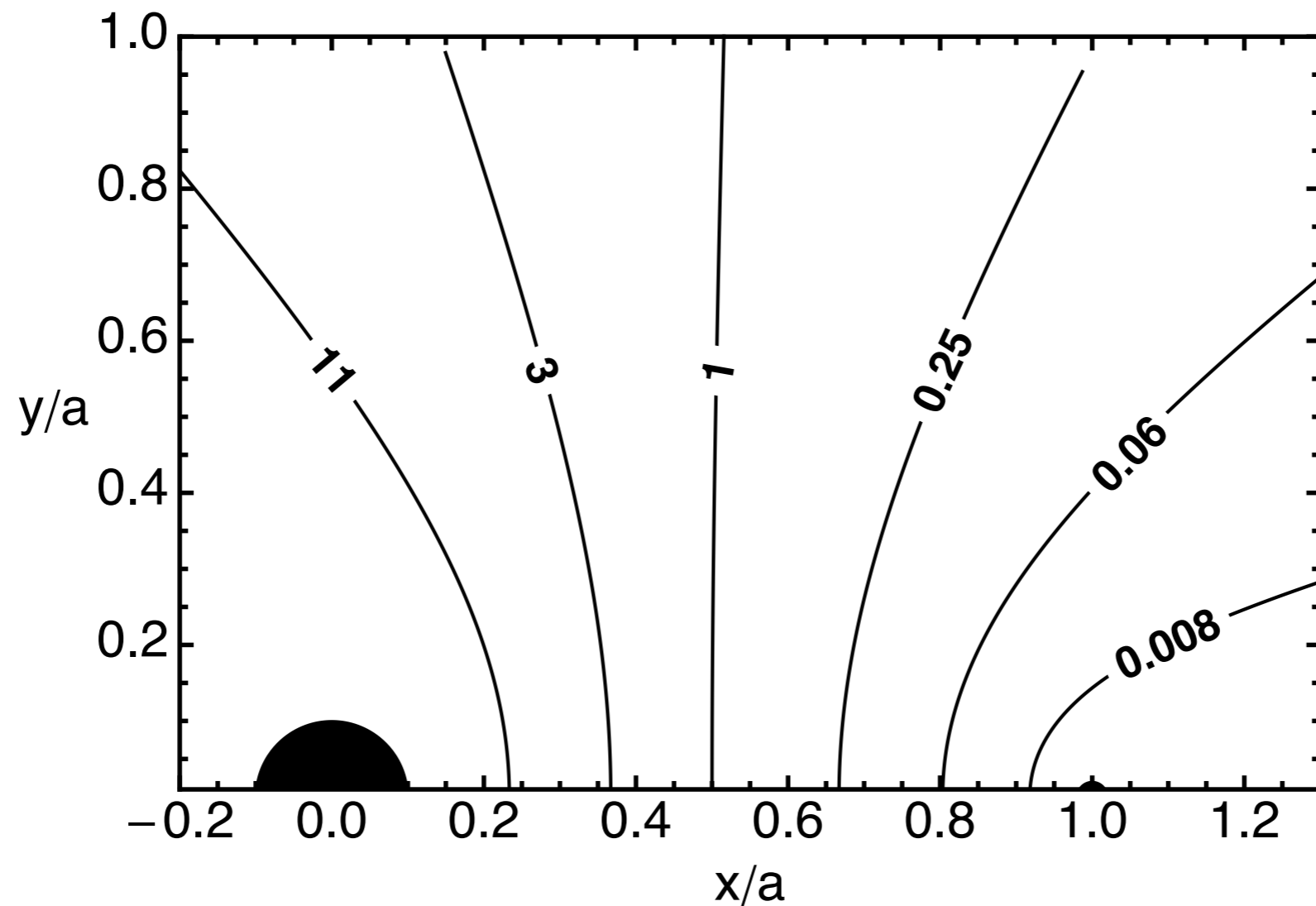
To derive constraints on the binary parameters from the lack of absorption and occultation signatures in the observed lightcurves and assuming the X-ray source is extended.

We created a toy model of LS 5039 in the framework of pulsar wind scenario.



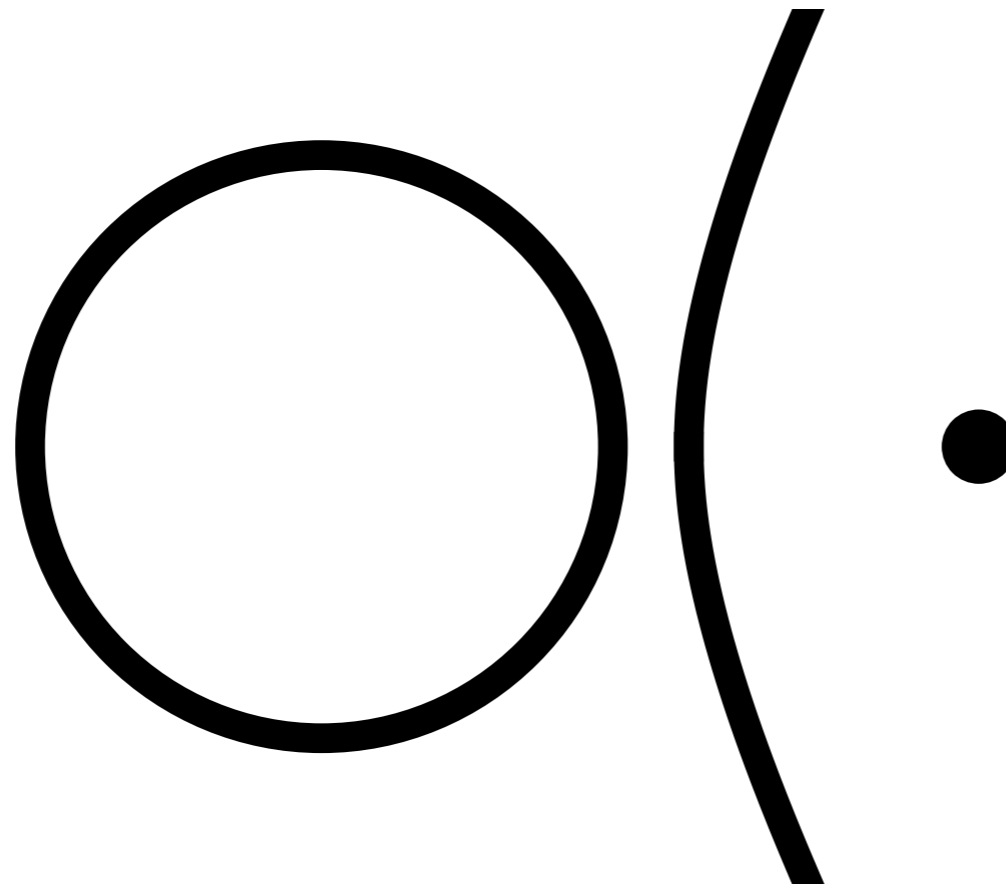
In general,
shape of the CD is parametrized by parameter η_0

$$\eta_0 = \eta(R_s, a - R_s) = \frac{\dot{E}}{\dot{M}v(R_s)c} = \frac{(a - R_s)^2}{R_s^2}$$



If \dot{E} is too large, the pulsar wind impacts the star surface, just like in black widow pulsars.

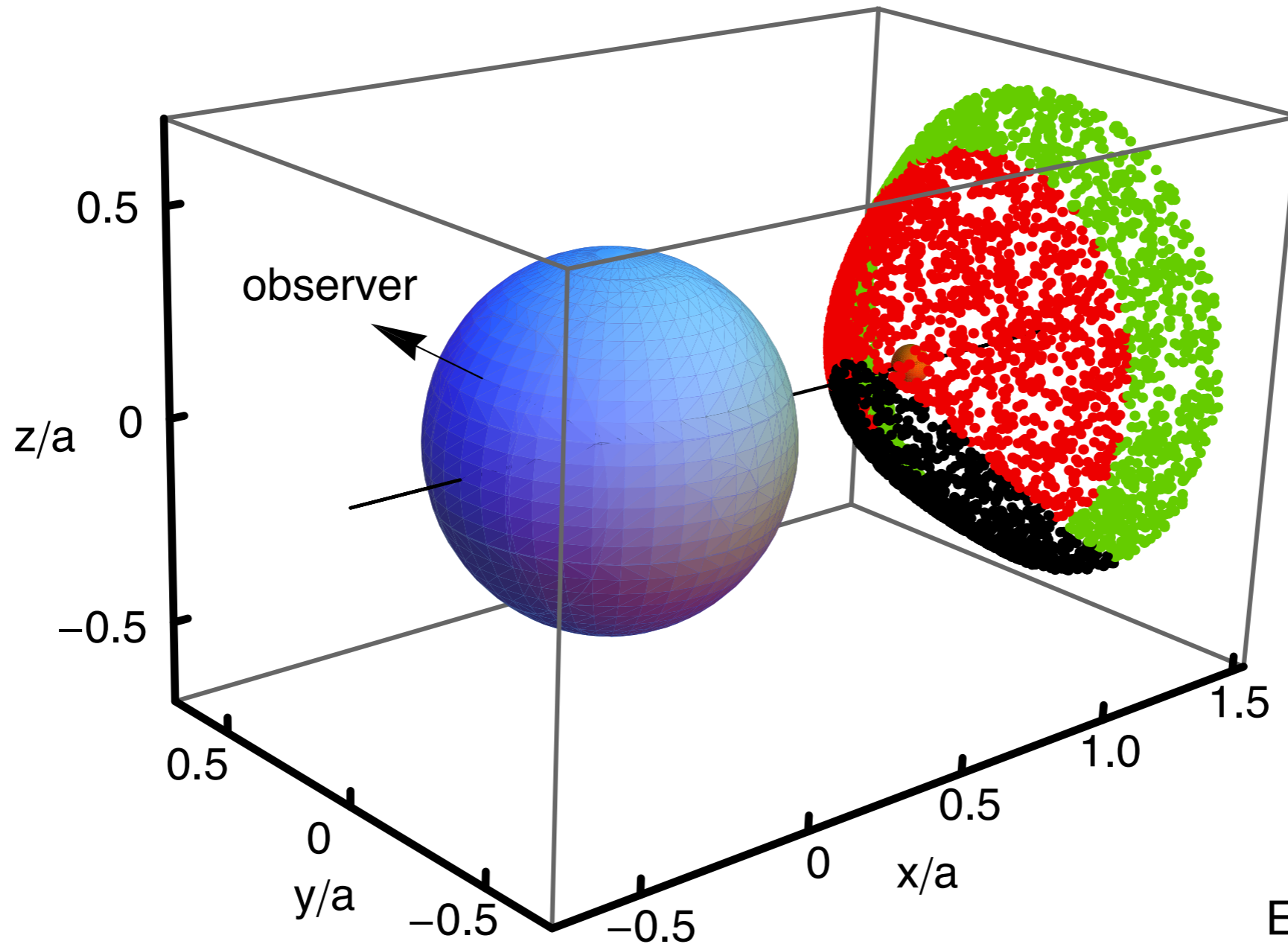
Assuming it does not happen: in LS 5039 η_0 never exceeds 0.6. CD opening angle $< 62^\circ$



For a range of \dot{M} ,

$$\dot{E}_{\max} = 1.5 \times 10^{36} \text{ erg s}^{-1} - 3.8 \times 10^{37} \text{ erg s}^{-1}$$

A 3D model of an extended X-ray source.



Example
SUPC, $i=45^\circ$, $\eta_0=0.1$

Absorption study

In the model absorption depends on

\dot{M} , i , η_0 , R_{out} (size of X-ray emission region)

Assumption: the observational data (point vs. extended source) are not distinguished by spectral fitting (e.g. xspec) if equal at 1 keV

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$$\frac{1}{N} \sum_{j=1}^N \exp[-\sigma(1\text{keV})N_{\text{H}j}(i, \phi)] = \exp[-\sigma(1\text{keV})N_{\text{H}}^{\text{max}}]$$

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Absorption of 1 keV photons from a point source by the not ionized homogeneous cloud with $N_{\text{H}}^{\text{max}} = 2.6 \times 10^{21} \text{ cm}^{-2}$

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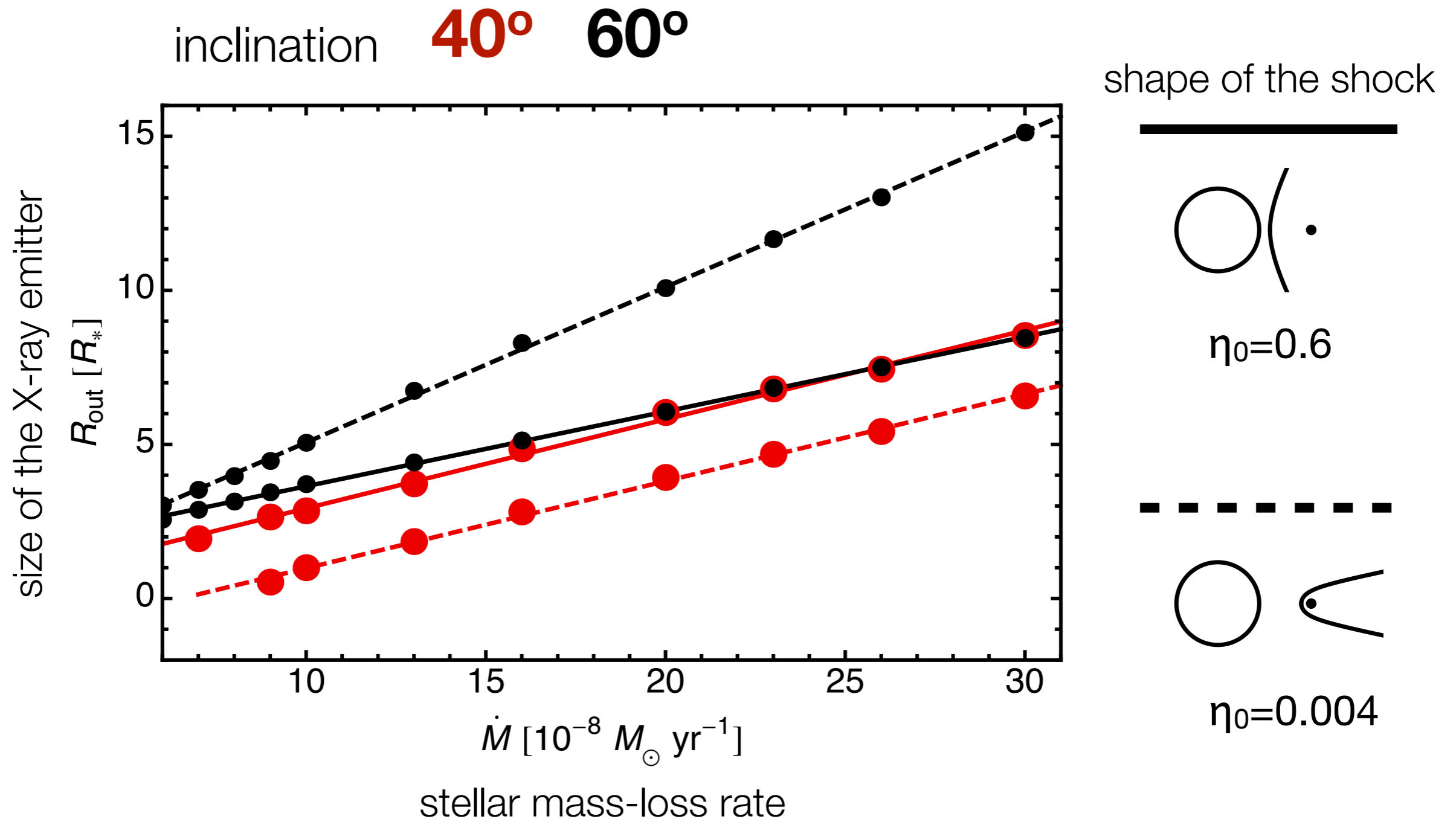
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Exact absorption of a 3D source in the not ionized stellar wind at SUPC at 1 keV

Absorption of 1 keV photons from a point source by the not ionized homogeneous cloud with $N_{\text{H}}^{\text{max}} = 2.6 \times 10^{21} \text{ cm}^{-2}$

Estimate based on absorption limit:

The higher the \dot{M} the larger the emitter size

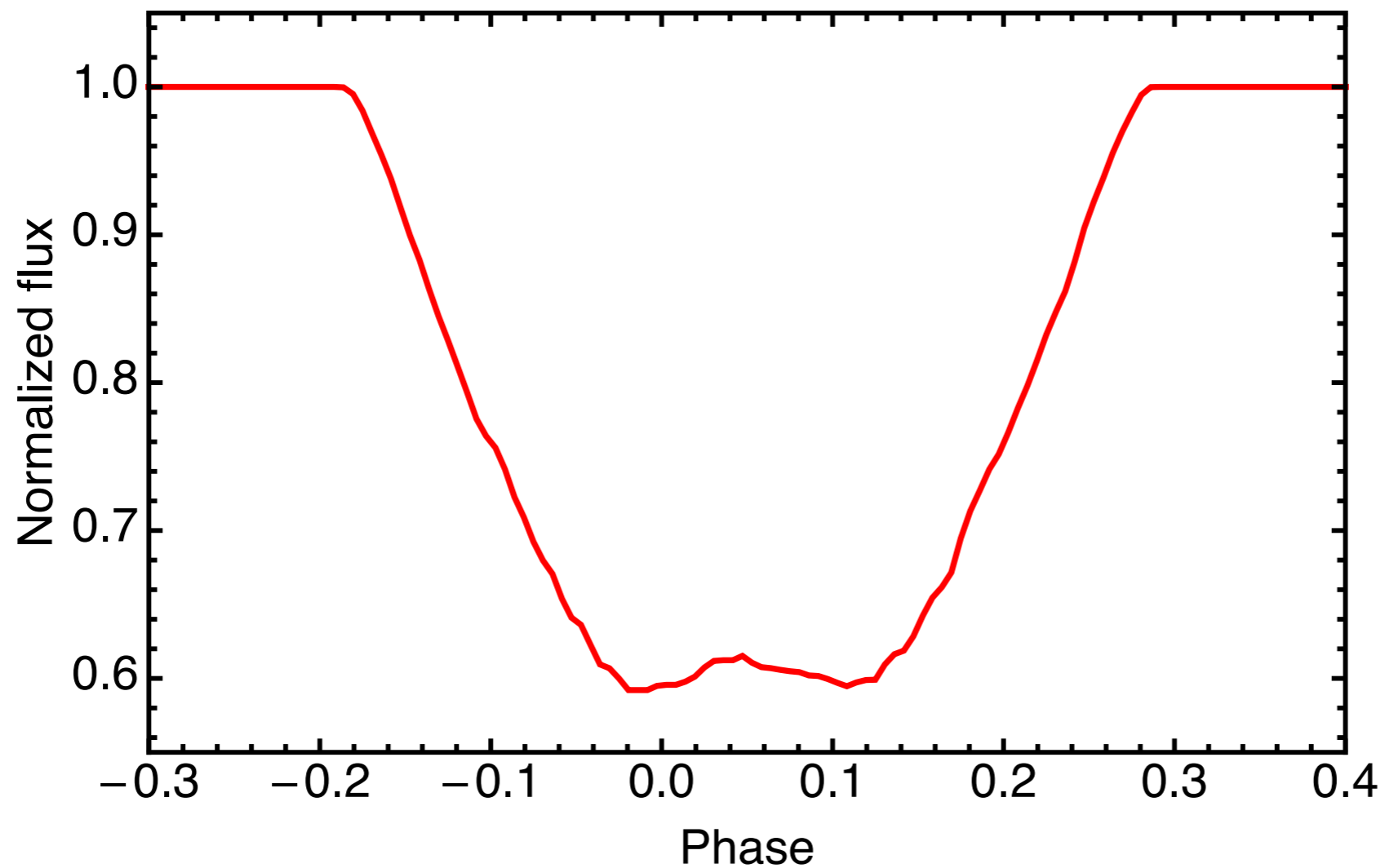


Occultation study

In the model occultation depends on:

i , η_0 , R_{out}

For example: circular orbit, $\eta_0 > 1$

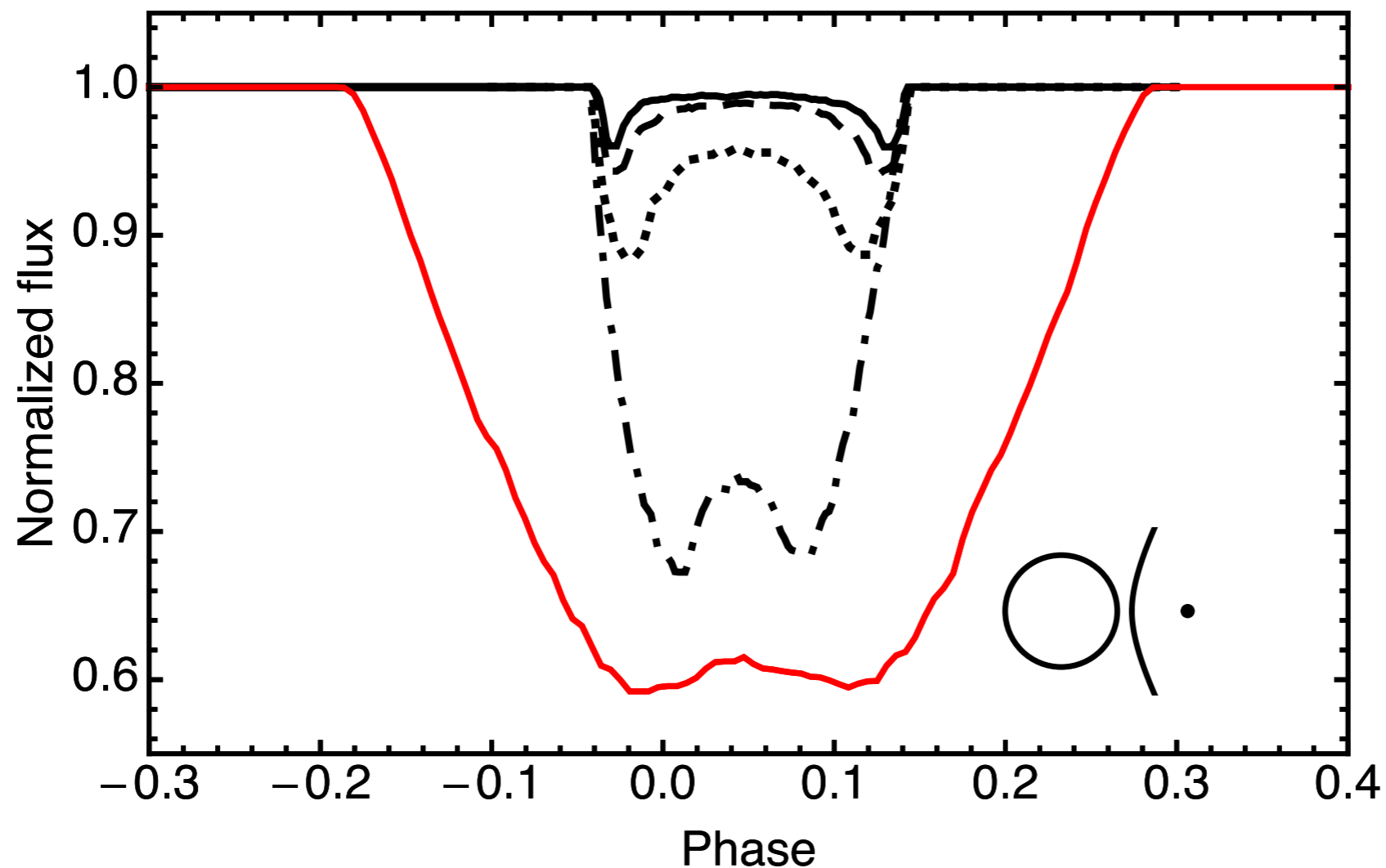


In the model occultation depends on:

i , η_0 , R_{out}

Occultation in LS 5039, in the most favorable condition

$\eta_0=0.6$, $i=60^\circ$, $R_{\text{out}} = 2R_\star, 5R_\star, 10R_\star, 15R_\star$

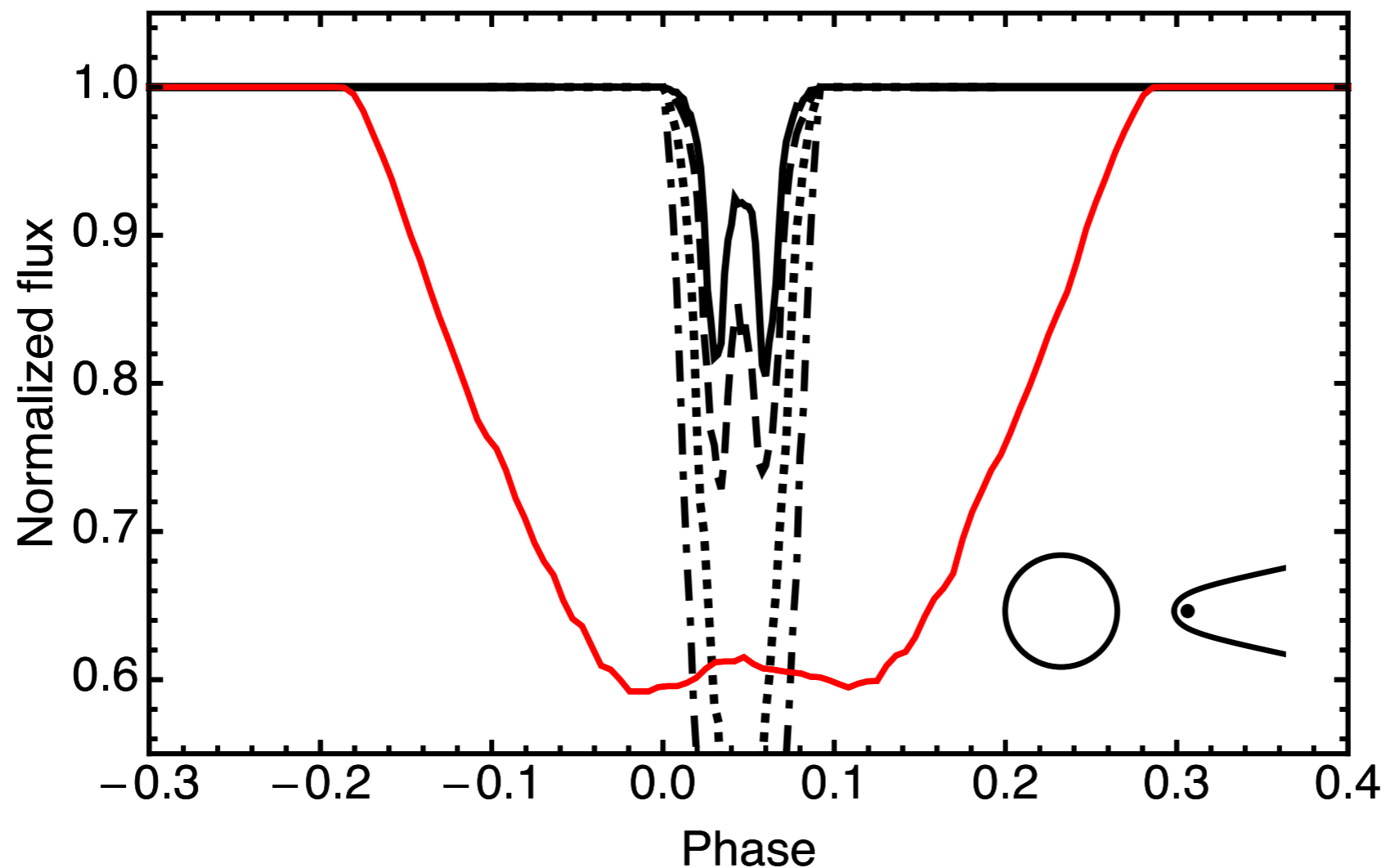


In the model occultation depends on:

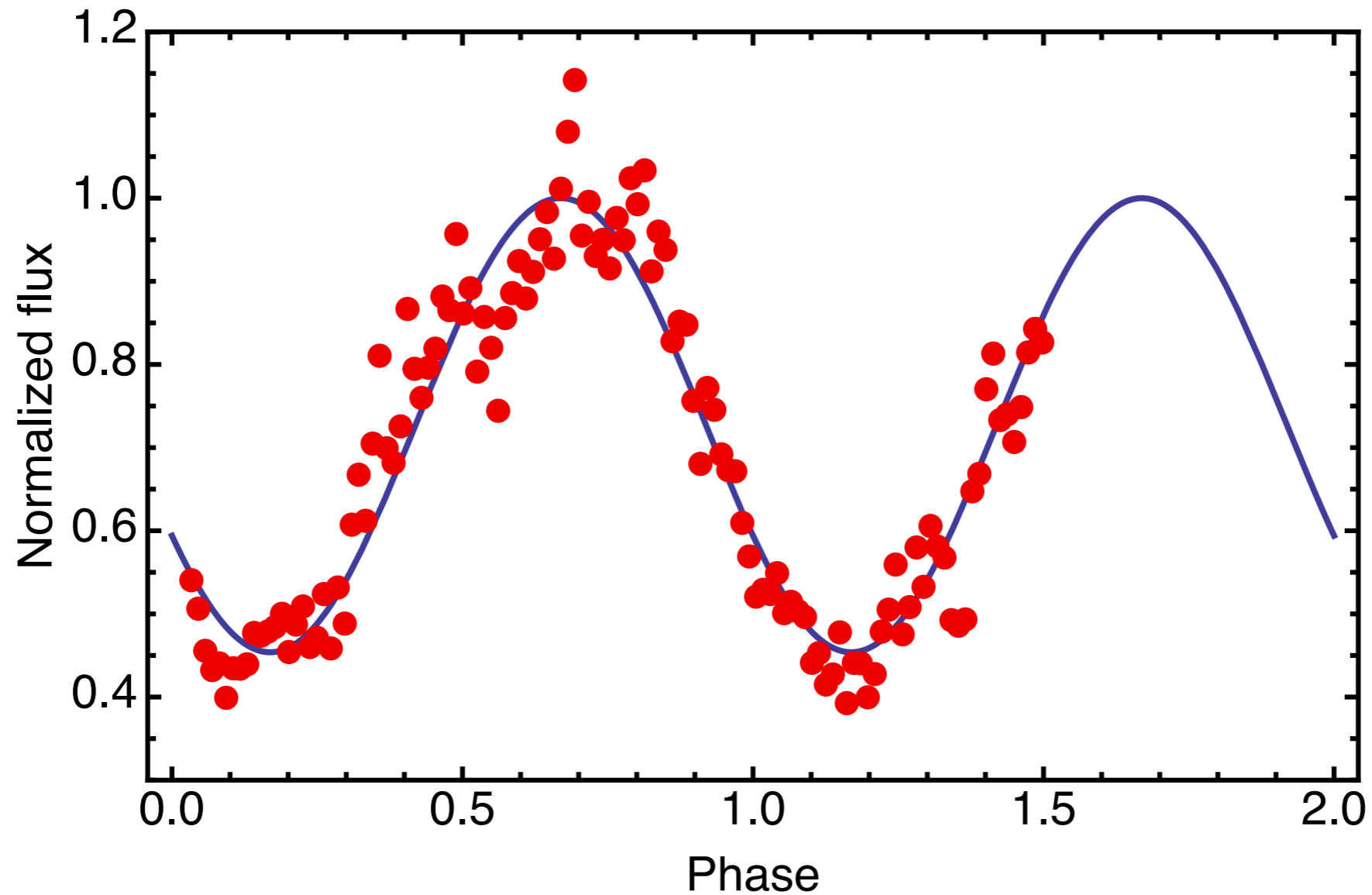
i , η_0 , R_{out}

Occultation in LS 5039, the less favorable condition

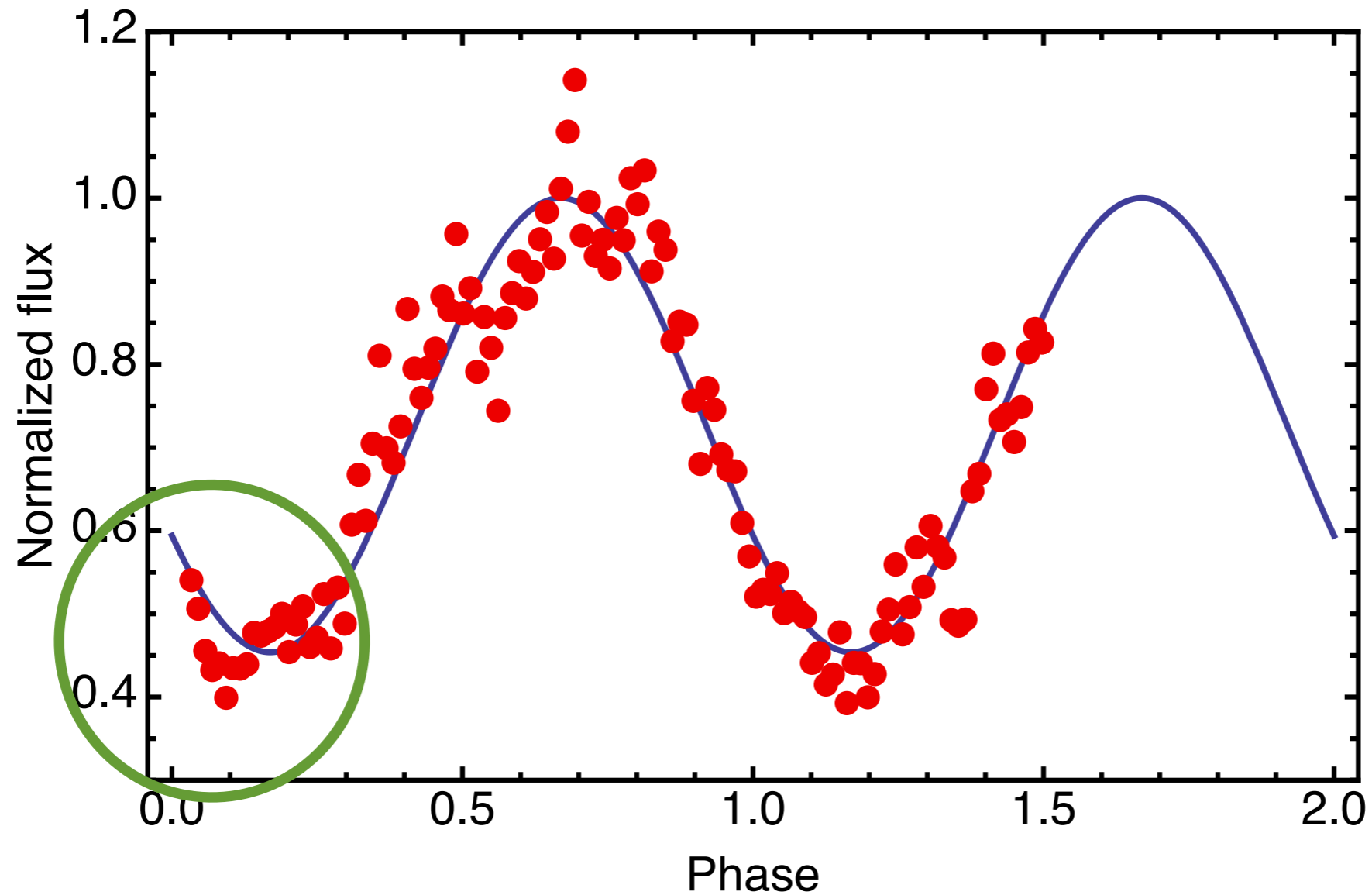
$\eta_0=0.004$, $i=90^\circ$, $R_{\text{out}} = 2R_\star, 5R_\star, 10R_\star, 15R_\star$



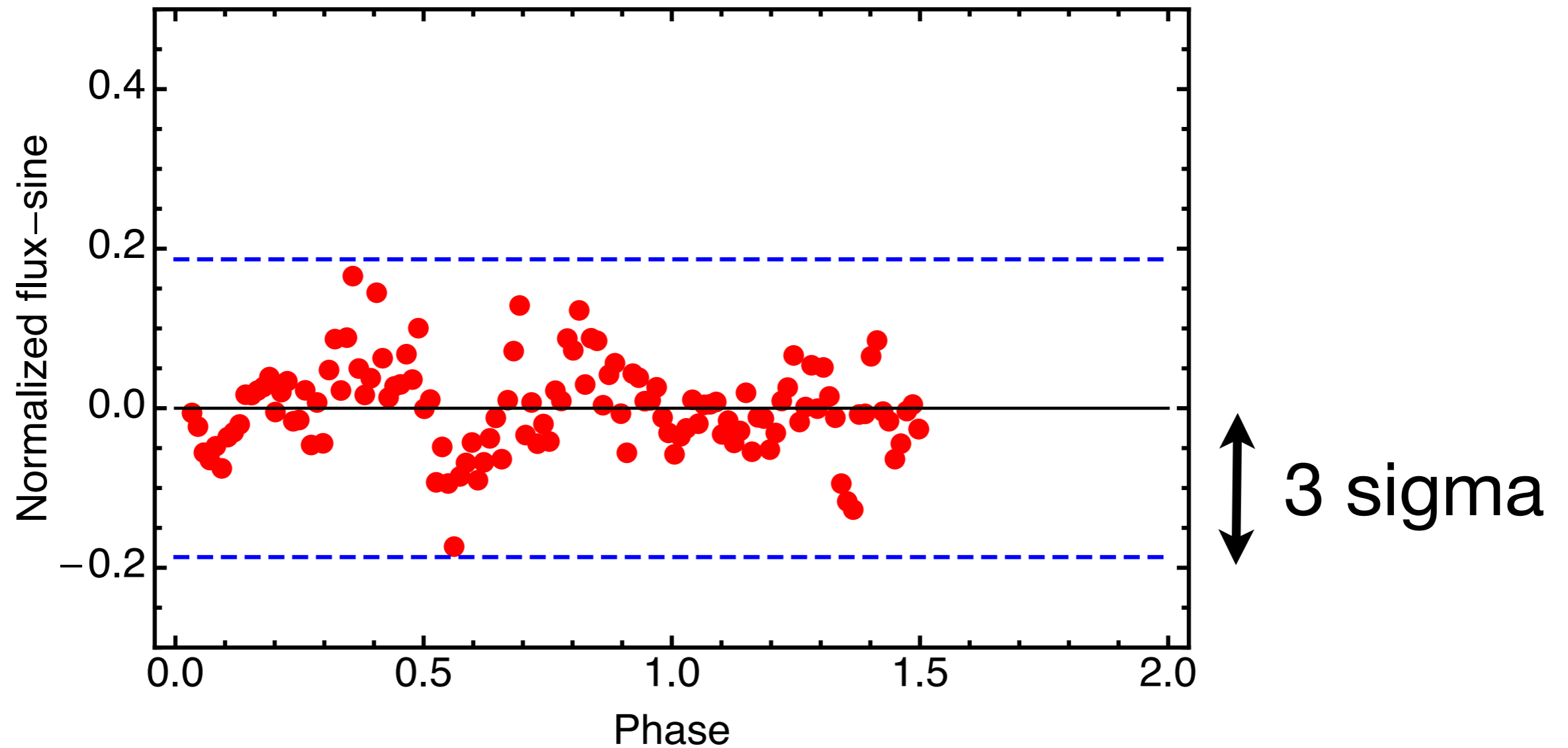
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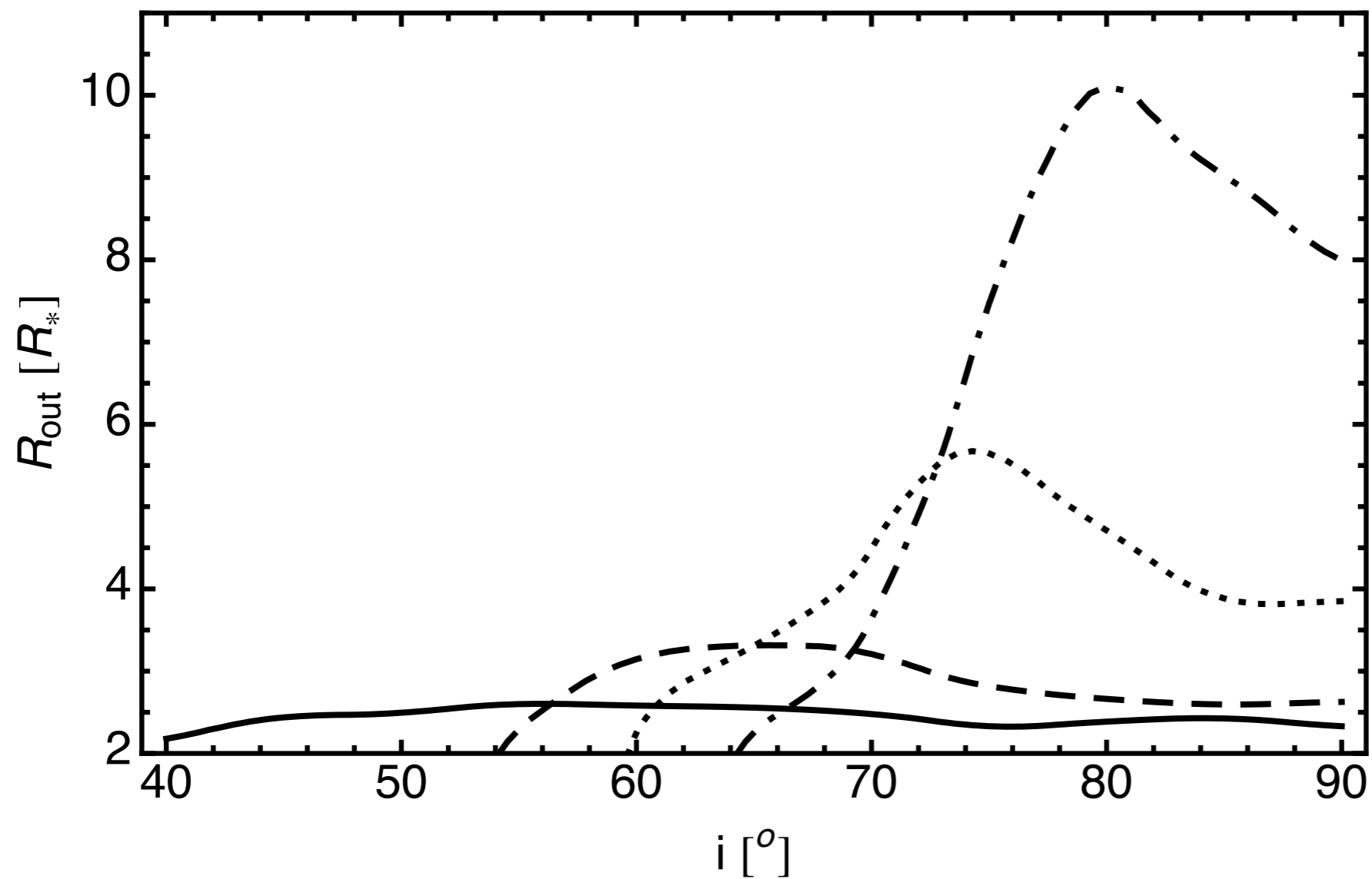


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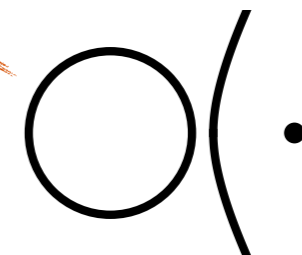
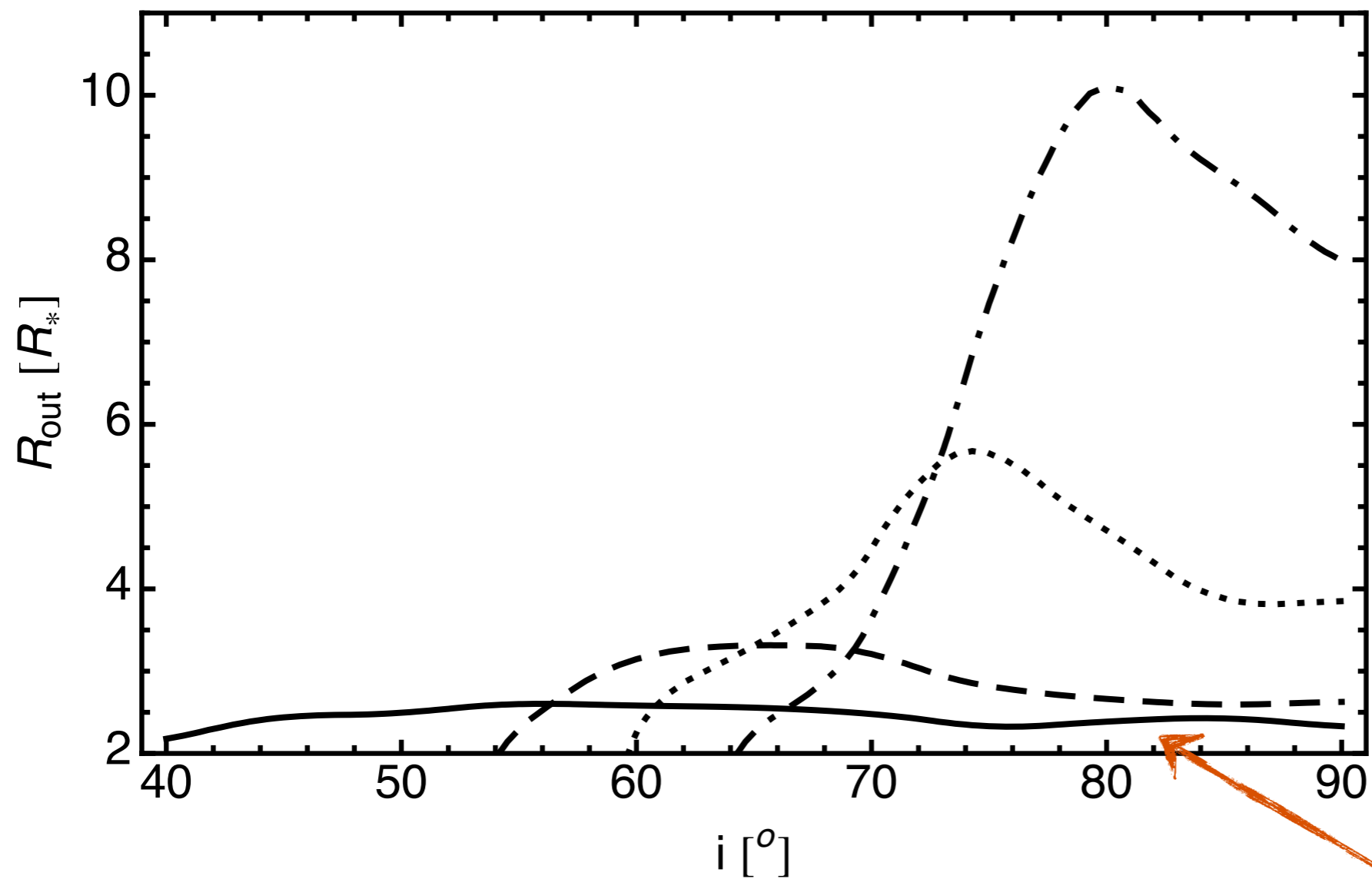
Occultation for different η_0

Allowed emitter size is above each curve



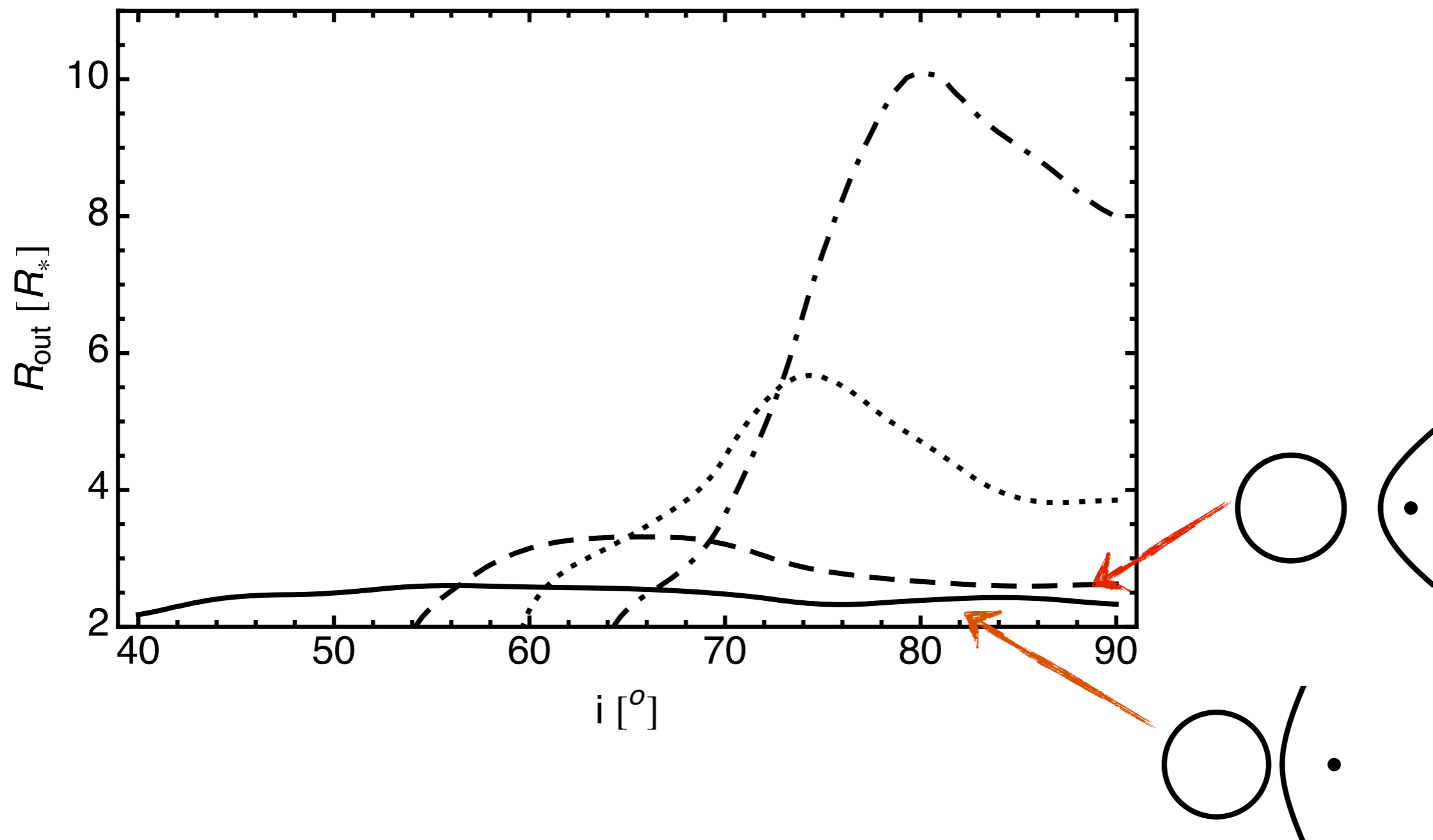
Occultation for different η_0

Allowed emitter size is above each curve



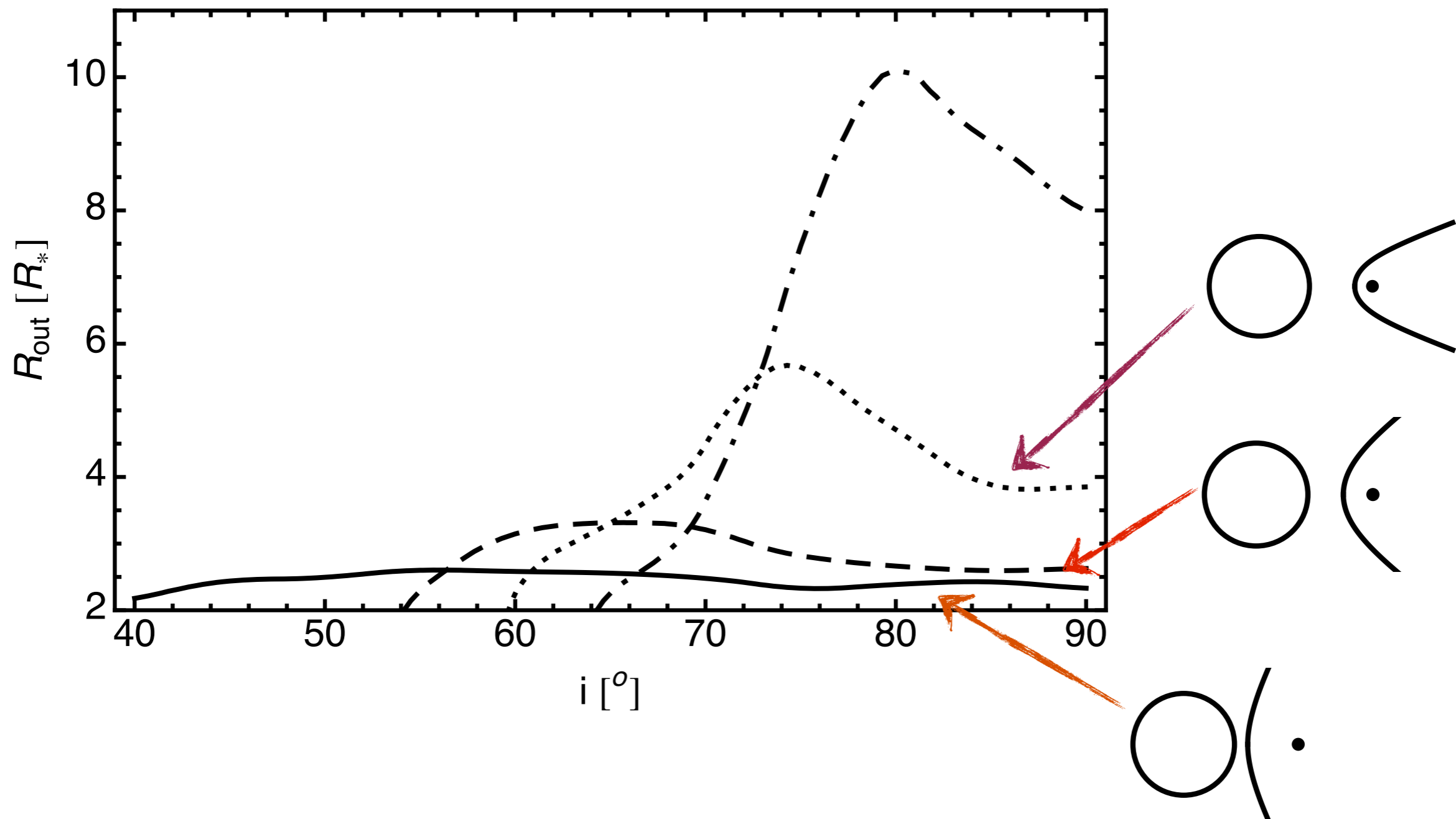
Occultation for different η_0

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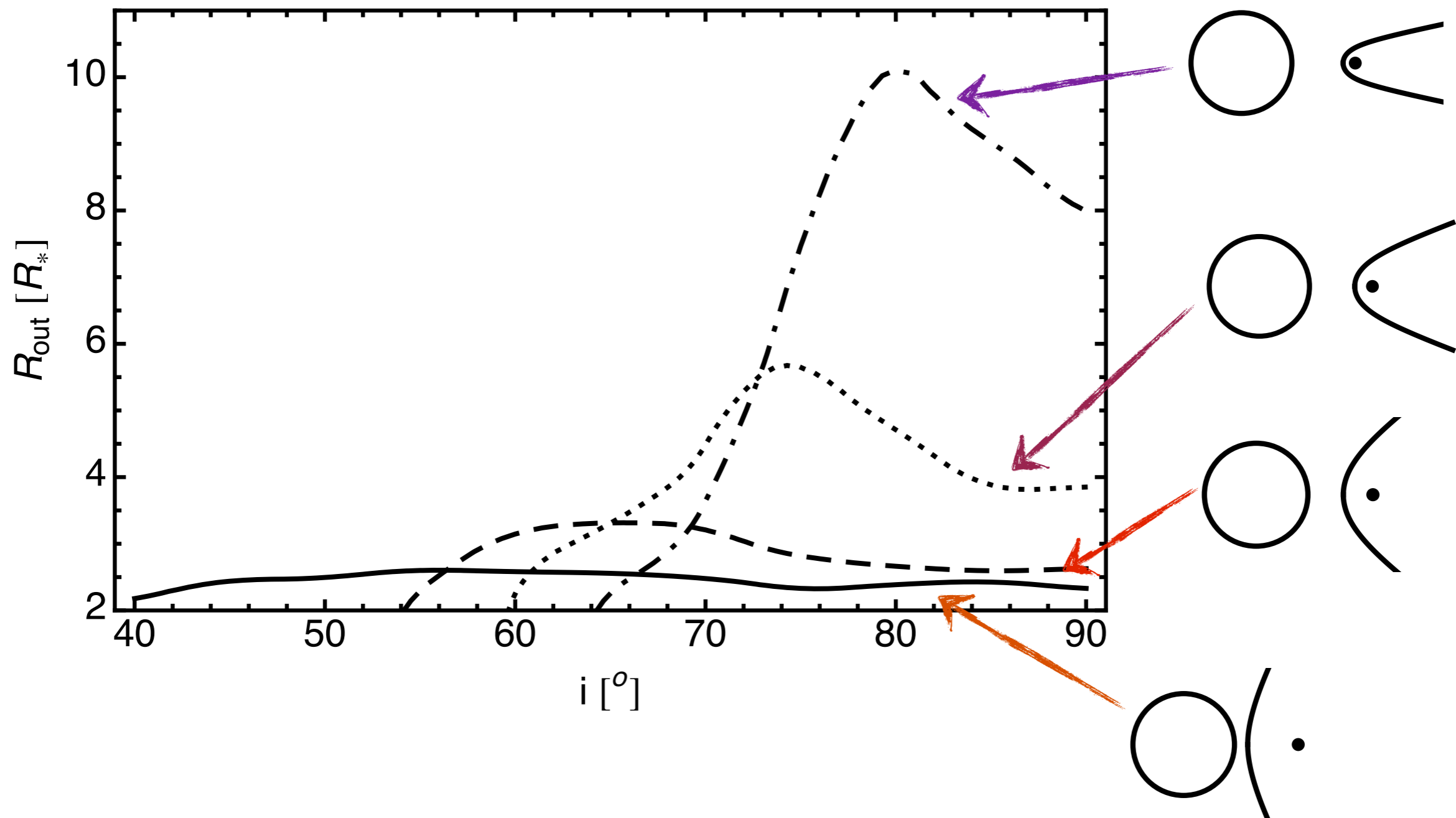
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Summary

An extended X-ray source appears necessary.

- Minimum size of X-ray emitting region, varies between 3-10 R_{\star}
- Inclination limit of 60° is not valid if $R_{\text{out}} > 4 R_{\star}$
- Order of magnitude higher stellar mass-loss rates are possible
- The extended X-ray source is expected in the pulsar wind scenario

This work was published here:

MNRAS in press
arXiv:1009.1980

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